```
Relation between fluctuation of total electric dipole moment
                                                                                                               by Takeshi Nishimatsu
   and electric susceptibility
          V = L_z L_z L_z a_o^3
                                                                           volume of the supercell
                                                                          external electric field
          p = \sum_{R} Z^* u(R)
                                                                          total electric dipole moment in the supercell
                                                                           P = (P_x, P_y, P_z)
                                                                          Hamiltonian of the supercell, s is a state
       74(5)=7,(5)-E.1f(5):
    Thermal average (here "average" means average-in-time) of P_{\alpha} and P_{\alpha} P_{\sigma} (\alpha, \sigma = x, y, z) can be calculated with partition function Z(\beta, E) = \sum_{s} e^{-\beta \mathcal{H}(s)}, where \beta = 1/k_s T, as
\langle \mathcal{P}_{\alpha} \rangle = \frac{1}{Z(\beta, \mathbf{E})} \sum_{s} e^{-\beta \left[\mathcal{H}_{o}(s) - \mathbf{E} \cdot \mathbf{P}(s)\right]} P_{\alpha}(s) = \frac{1}{\beta} \frac{\partial \log Z(\beta, \mathbf{E})}{\partial \mathcal{E}_{\alpha}}
\langle \mathcal{P}_{\alpha} \mathcal{P}_{\delta} \rangle = \frac{1}{Z(\beta, \mathbf{E})} \sum_{s} e^{-\beta \left[\mathcal{H}_{o}(s) - \mathbf{E} \cdot \mathbf{P}(s)\right]} P_{\alpha}(s) P_{\beta}(s) = \frac{1}{\beta^{2}} \frac{\partial \log Z(\beta, \mathbf{E})}{\partial \mathcal{E}_{\alpha} \partial \mathcal{E}_{\beta}} \frac{\partial^{2} Z(\beta, \mathbf{E})}{\partial \mathcal{E}_{\alpha} \partial \mathcal{E}_{\beta}}
    Now, fluctuation can be calculated as
      \langle P_{\alpha} - \langle P_{\alpha} \rangle \rangle \langle P_{\beta} - \langle P_{\beta} \rangle \rangle = \langle P_{\alpha} P_{\beta} \rangle - \langle P_{\alpha} \rangle \langle P_{\beta} \rangle
         \frac{1}{\beta^2} \frac{1}{Z(\beta, \mathfrak{E})} \frac{\partial^2 Z(\beta, \mathfrak{E})}{\partial \mathcal{E}_{\mathcal{X}} \partial \mathcal{E}_{\mathcal{S}}} = \frac{1}{\beta^2} \frac{\partial \log Z(\beta, \mathfrak{E})}{\partial \mathcal{E}_{\mathcal{X}}} \frac{\partial \log Z(\beta, \mathfrak{E})}{\partial \mathcal{E}_{\mathcal{S}}}
 = \frac{1}{\beta^2} \frac{\partial^2 \log Z(\beta, \mathcal{E})}{\partial \mathcal{E}_{\mathcal{A}} \partial \mathcal{E}_{\mathcal{S}}} = \frac{1}{\beta} \frac{\partial \langle \mathcal{P}_{\mathcal{A}} \rangle}{\partial \mathcal{E}_{\mathcal{S}}}
   Here, electric susceptibility per unit volume is,
    in CGS
 \chi_{dr} = \frac{4\pi}{V} \frac{Q\langle R_d \rangle}{\partial E_{r}} = 4\pi \frac{P}{V} \left( \langle R_d R_{r} \rangle - \langle R_d \rangle \langle R_{r} \rangle \right) = \frac{4\pi}{V k_B T} \left( \langle R_d R_{r} \rangle - \langle R_d \rangle \langle R_{r} \rangle \right)
in SI relative susceptibility (no dimension)
                          \frac{\chi_{\chi r} = \frac{1}{E_0 V k_B T} \left( \langle P_\alpha P_\sigma \rangle - \langle P_\alpha \rangle \langle P_r \rangle \right)}{\frac{F}{m} = \frac{C^2}{Jm} \frac{\sqrt{m^3}}{m^3} J \left( \frac{Cm}{m} \right)^2}
```