

Sunny.jl: Overview and SU(N) Capabilities

ACNS Tutorial, 2024



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Su(n)ny



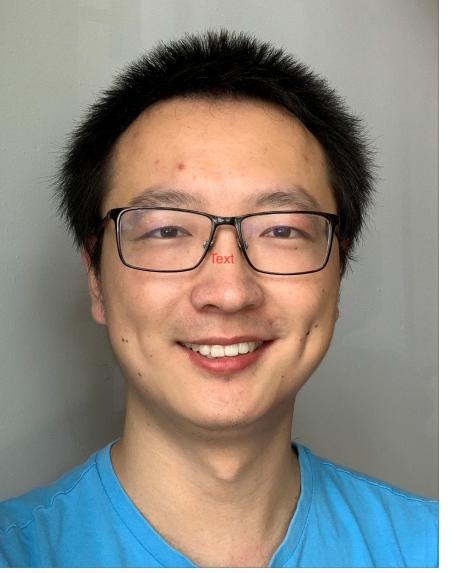
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Daniel Pajerovski
(ORNL)



M. Mourigal
(Georgia Tech)



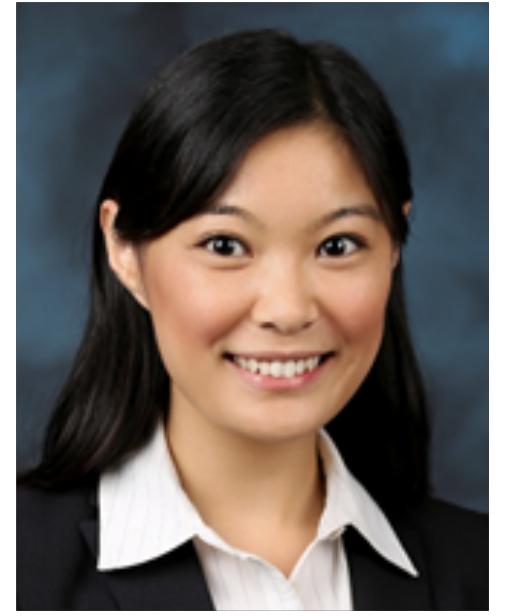
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David Dahlbom
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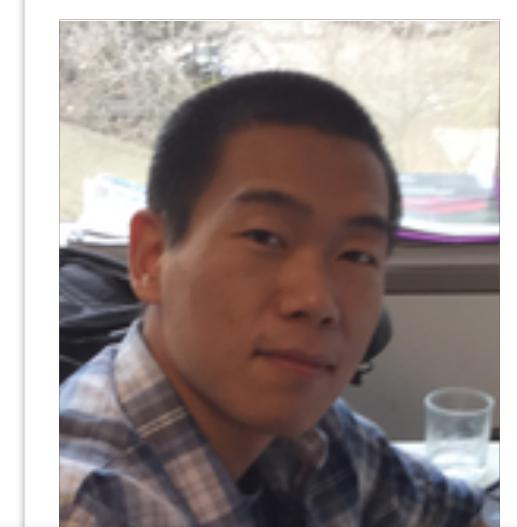
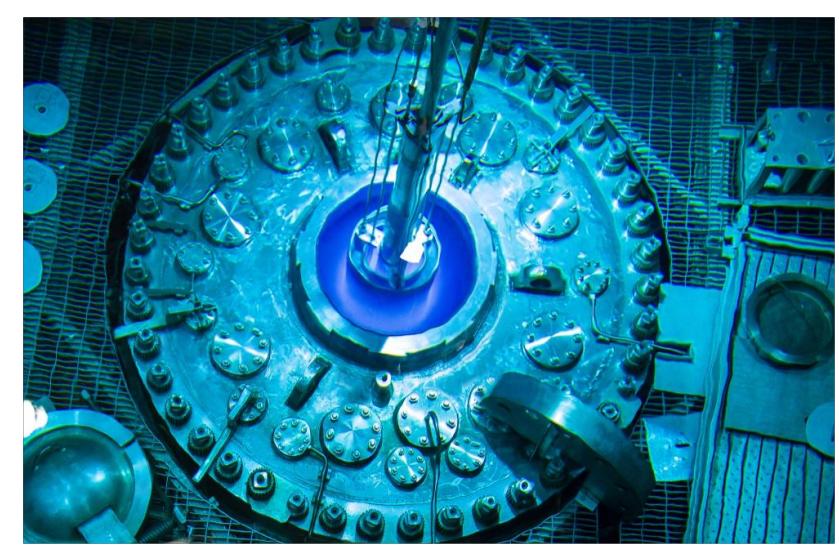
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Matt Wilson
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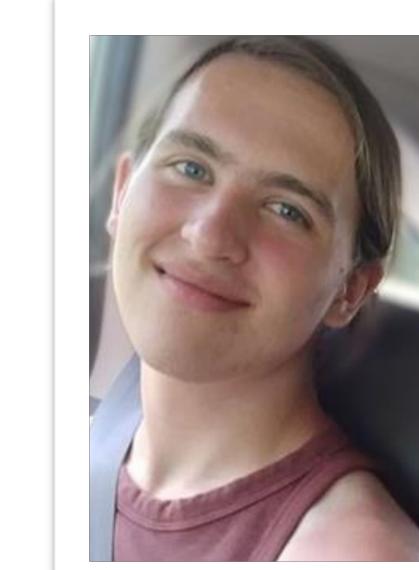
Hao Zhang
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Sakib Matin
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Het Mankad
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Sam Quinn
(Georgia Tech)



Harry Lane
(St. Andrews)



Alin Niraula
(LSU)



**Bhushan
Thipe**
(LSU)



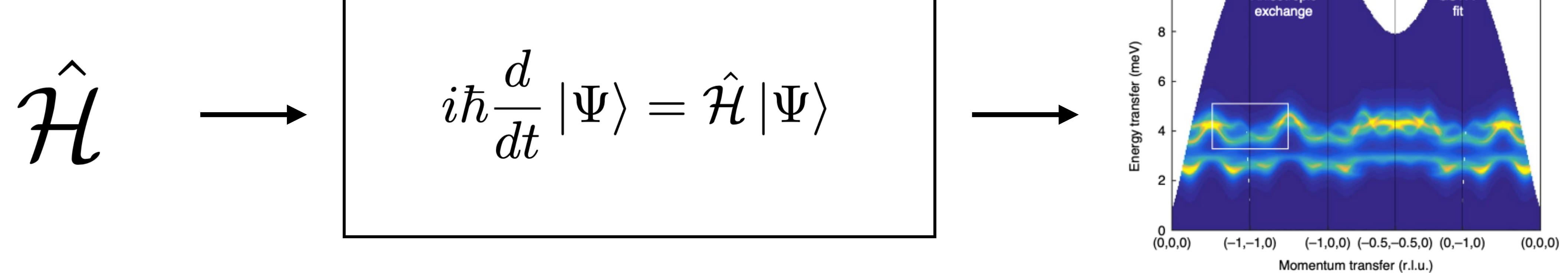
Cole Miles
(Kodiak)

Motivation for developing Sunny

Experiment



Theory



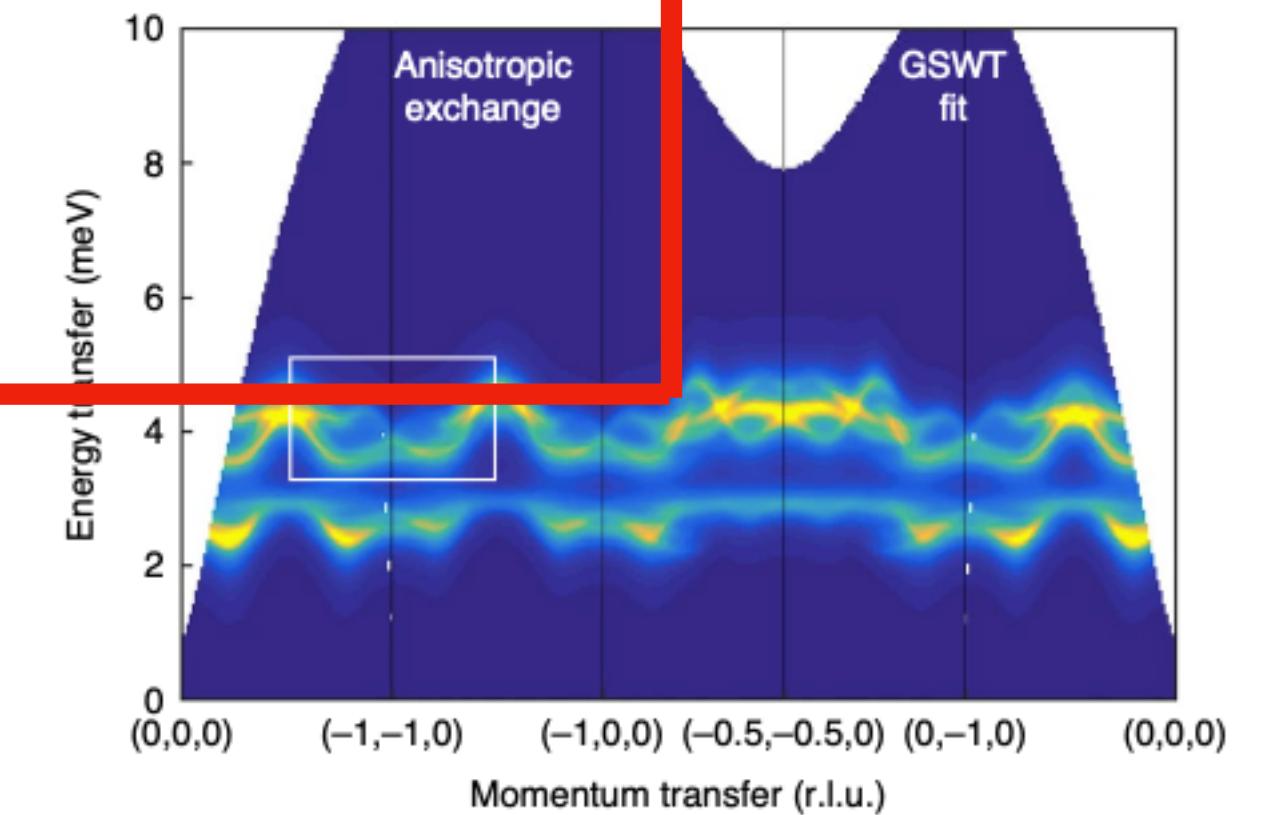
Motivation for developing Sunny

Experiment



Theory

$$i\hbar \frac{d}{dt} |\Psi\rangle - \hat{\mathcal{H}} |\Psi\rangle$$



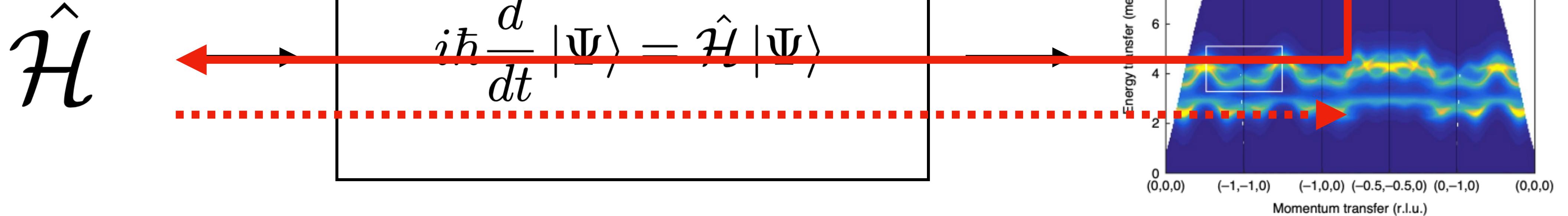
Motivation for developing Sunny

$$\propto \mathcal{S}^{\alpha\beta}(\mathbf{q}, \omega)$$

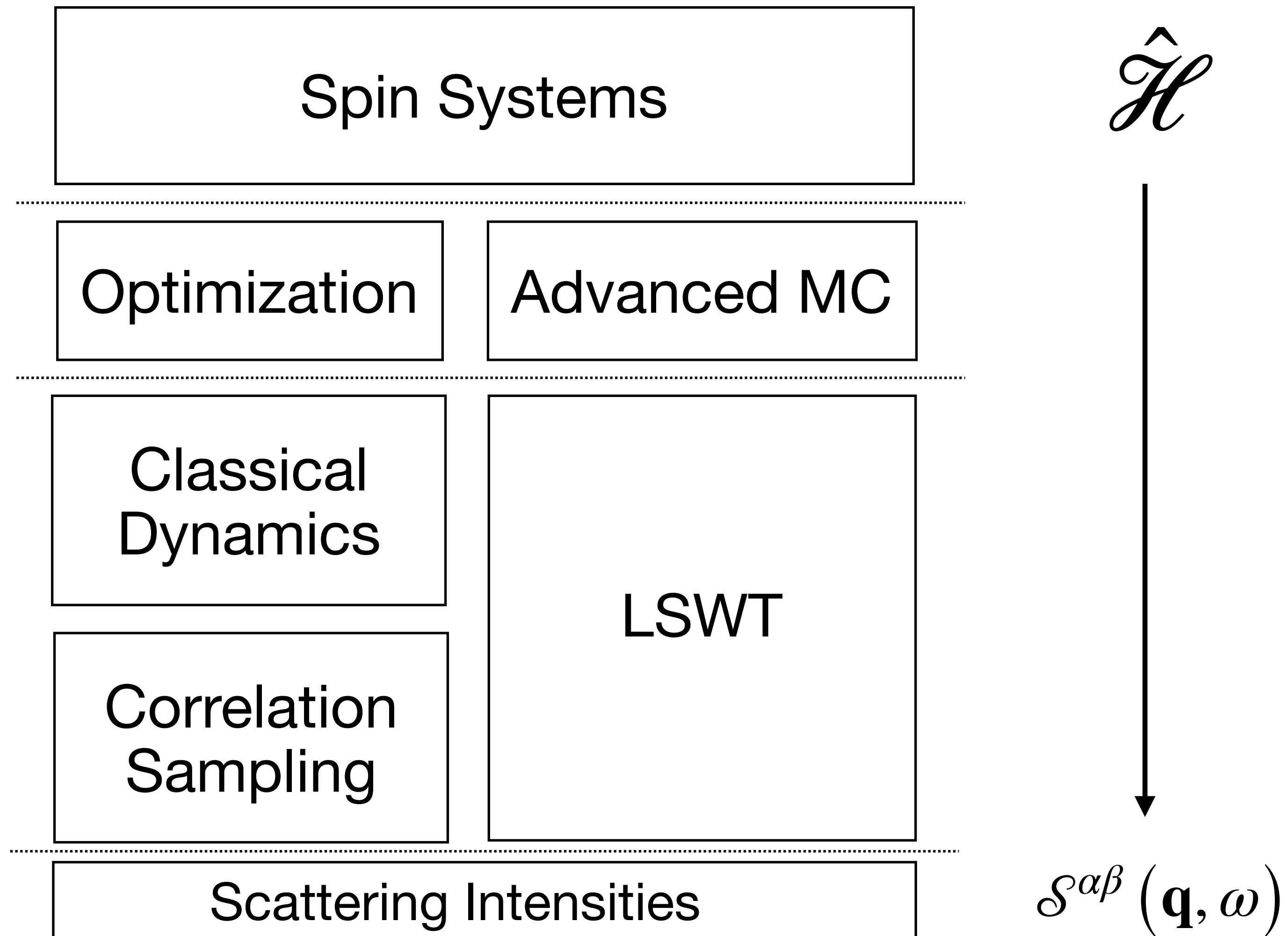
Experiment



Theory



Structure of Sunny



Structure of Sunny

Spin Systems

Optimization

Advanced MC

Classical
Dynamics

LSWT

Correlation
Sampling

Scattering Intensities

- Crystals and symmetry analysis
- Visualization
- Hamiltonian specification
 - General pair-wise interactions
 - Single-ion anisotropies (and renormalizations)
- Storage of product states (dipoles or $SU(N)$ coherent states)

Structure of Sunny

Spin Systems

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Scattering Intensities

- Symmetry analysis
 - `Crystal`
 - `print_symmetry_table`
 - `view_crystal`

Structure of Sunny

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Scattering Intensities

- System and Hamiltonian specification
 - **System**
 - **SpinInfo**
 - **set_exchange!**
 - **set_pair_coupling!**
 - **set_onsite_coupling!**

Structure of Sunny

Spin Systems

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Scattering Intensities

- Gradient-based optimizers to find ground states
 - Work for both dipoles and $SU(N)$ coherent states
- Advanced Monte Carlo classical spin systems
 - Parallel tempering
 - Wang-Landau

Structure of Sunny

Spin Systems

Optimization

Advanced MC

Classical
Dynamics

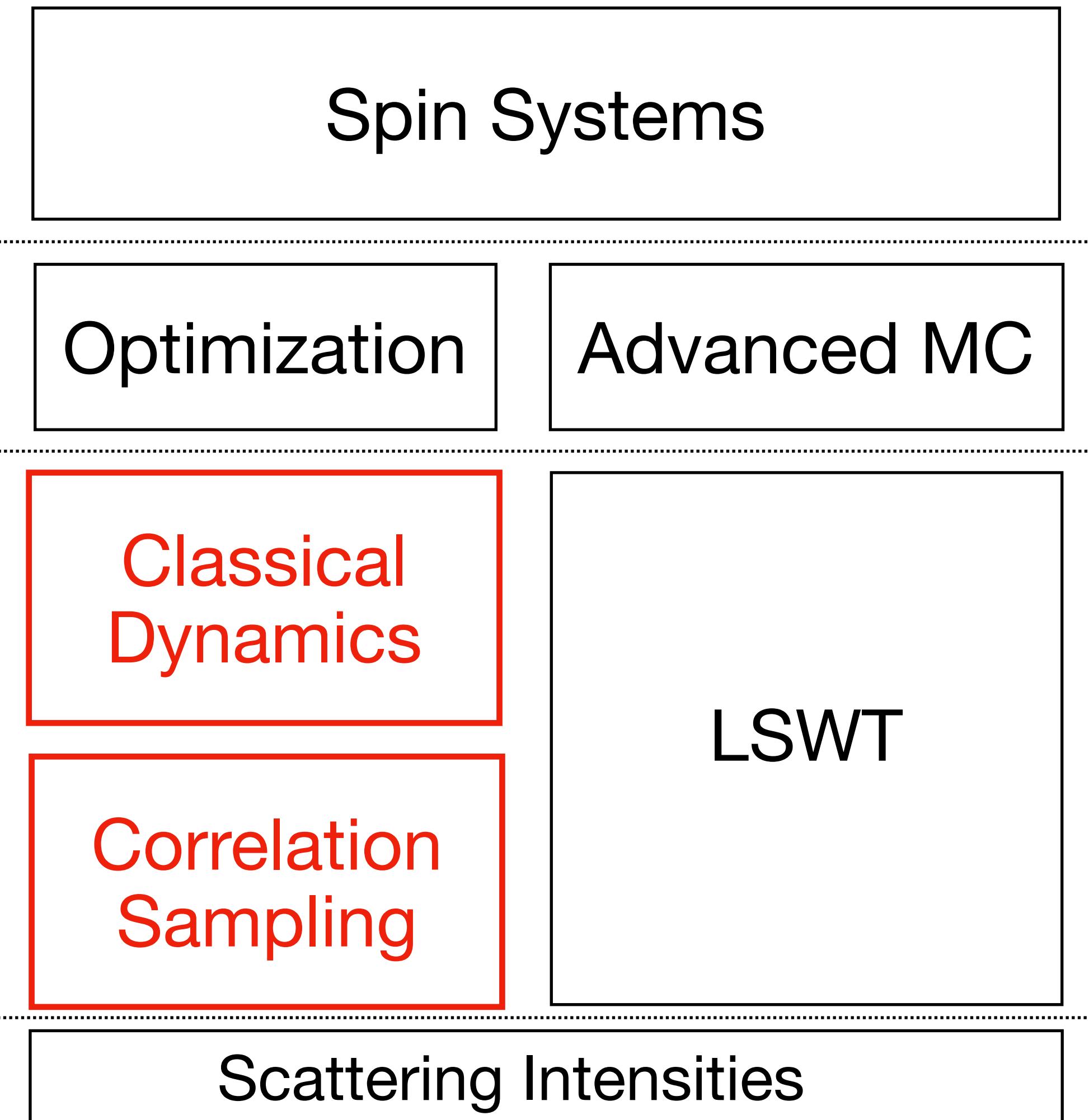
LSWT

Correlation
Sampling

Scattering Intensities

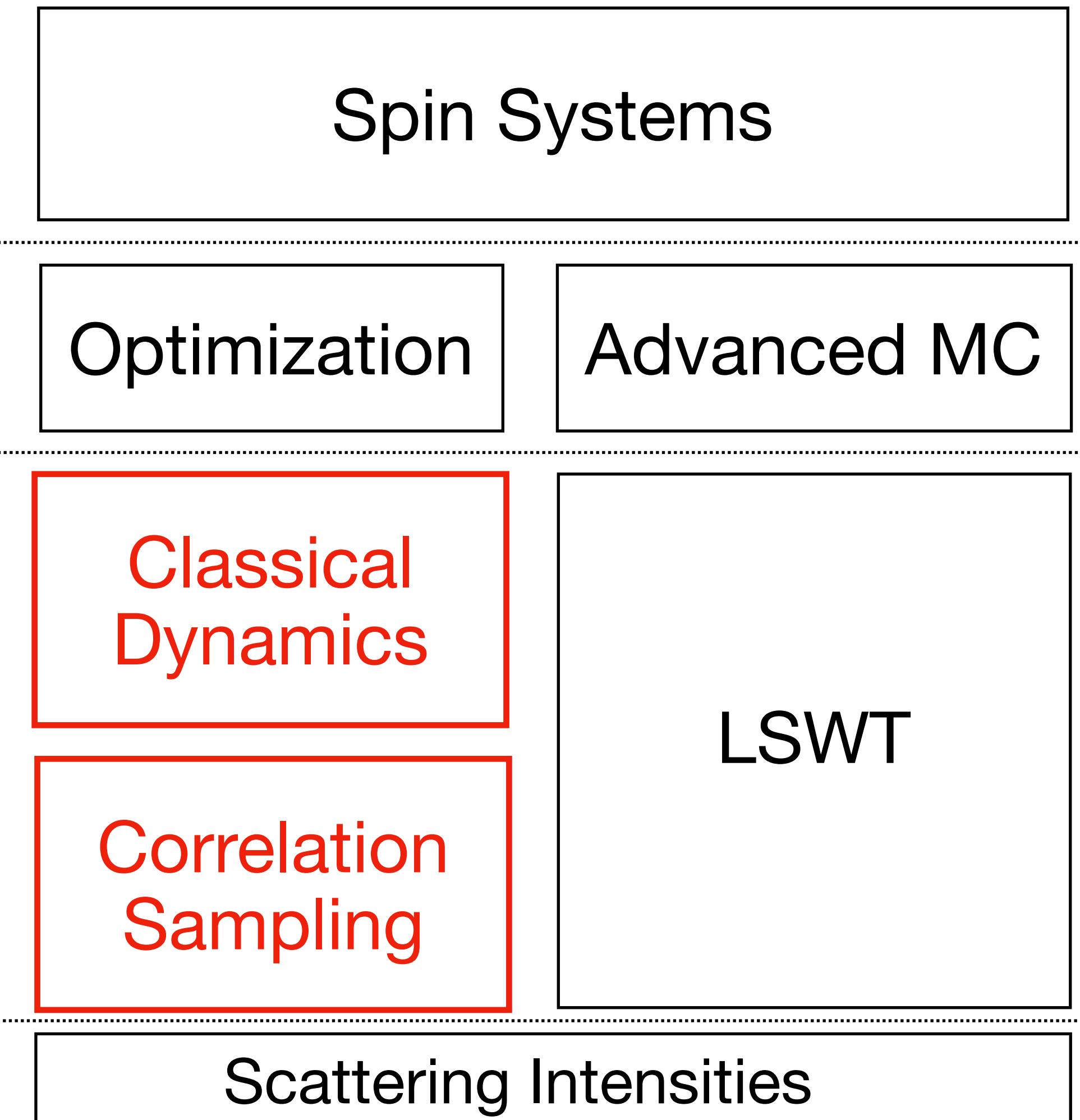
- Optimization
 - **randomize_spins!**
 - **minimize_energy!**
- Annealing via
 - **Langevin** and **step!**
 - **LocalSampler** and **step!**

Structure of Sunny



- Traditional Landau-Lifshitz and generalized $SU(N)$ classical dynamics
- Advanced integration techniques
 - Symplectic integrators for energy-conserving dynamics
- Dynamics with presence of thermal bath (both dipole and $SU(N)$).

Structure of Sunny



- Classical dynamics
 - **ImplicitMidpoint** and **step!**
 - **Langevin** and **step!**
 - **suggest_timestep**

Structure of Sunny

Spin Systems

Optimization

Advanced MC

Classical
Dynamics

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LSWT

Scattering Intensities

- Correlation sampling, $S(q,\omega)$
 - **dynamical_correlations**
 - **instant_correlations**
 - **Langevin** and **step!**
 - **add_sample!**

Structure of Sunny

Spin Systems

Optimization

Advanced MC

Classical
Dynamics

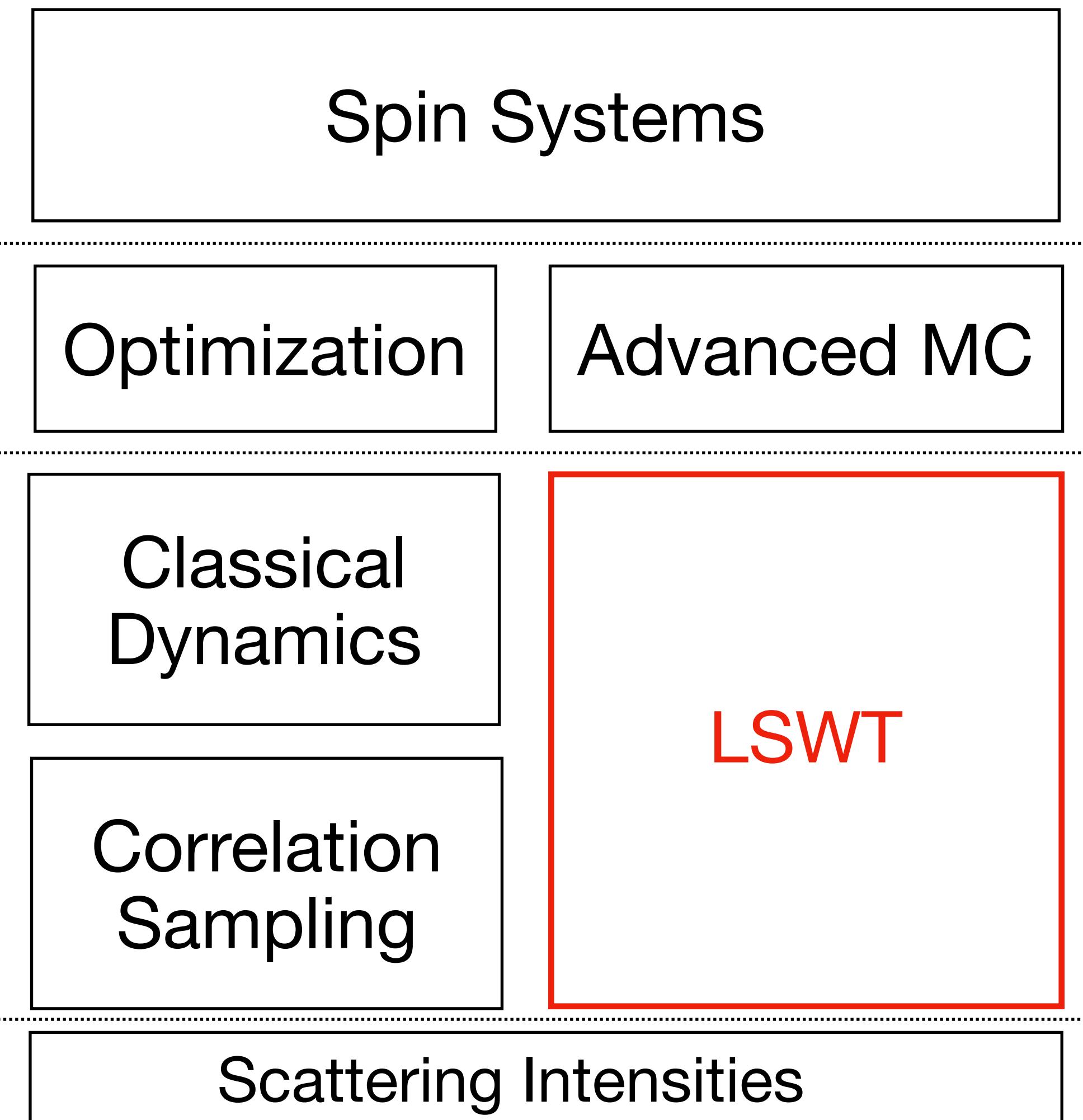
Correlation
Sampling

Scattering Intensities

- Construct and diagonalize LSWT Hamiltonian
- Both traditional S-based formalism and generalization based on $SU(N)$ coherent states.

LSWT

Structure of Sunny



- LSWT
 - **SpinWaveTheory**

Structure of Sunny

Spin Systems

Optimization

Advanced MC

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Dynamics

Correlation
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LSWT

Scattering Intensities

- Tools for extracting scattering intensities for comparison with experiment.
- Polarization, form factors, convolution kernels.

Structure of Sunny

Spin Systems

Optimization

Advanced MC

Classical
Dynamics

Correlation
Sampling

LSWT

Scattering Intensities

- Intensity calculations
 - **intensity_formula**
 - **intensities_bands**
 - **intensities_broadened**
 - **intensities_interpolated**

Typical workflow

1) Specify a system

Specify a crystal, supercell, local Hilbert space information, and interactions

- **Crystal**
- **view_crystal**
- **print_symmetry_table**
- **System**
- **SpinInfo**
- **set_exchange!**
- **set_onsite_coupling!**

Typical workflow

2) Find a ground state

Thermalize system or randomize spins, then find lowest energy configuration

- **randomize_spins!**
- **minimize_energy!**
- **Langevin** or **LocalSampler** and **step!** with an annealing schedule

Typical workflow

3a) Classical structure factors

Build a **SampledCorrelations** object that accumulates spin-spin correlation information.

- **dynamical_correlations**
- **Langevin** and **step!**
- **add_sample!**

Typical workflow

3b) LSWT structure factors

Specify a magnetic unit cell and build a **SpinWaveTheory**. This stores observables and information for constructing the spin wave Hamiltonian.

- **print_wrapped_intensities**
- **suggest_magnetic_supercell**
- **reshape_supercell**
- **SpinWaveTheory**

Typical workflow

4) Retrieve intensities

For both classical and LSWT calculations, specify a “formula”

- **intensity_formula**

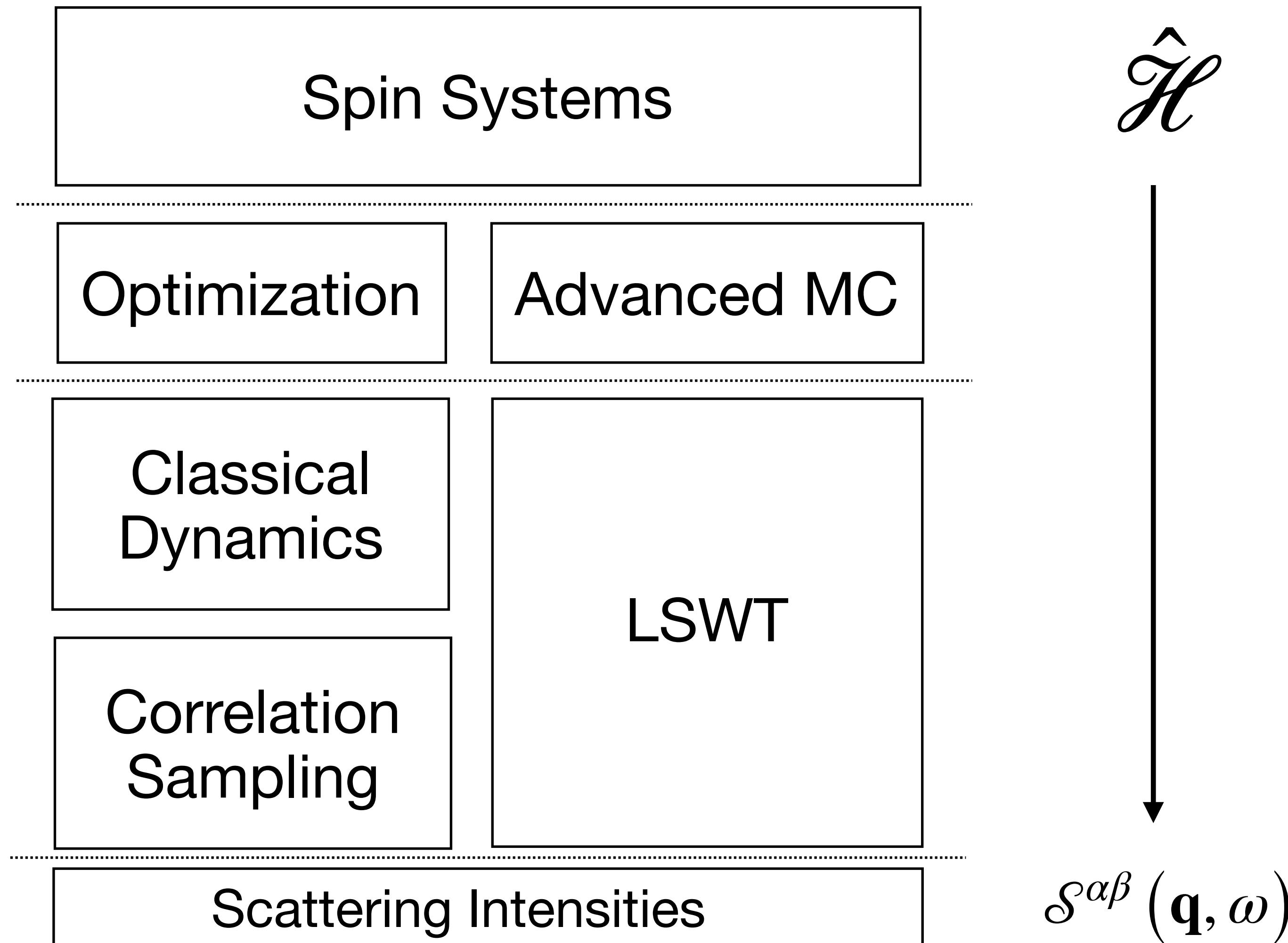
For classical calculations

- **intensities_interpolated**

For LSWT calculations

- **intensities_bands , intensities_broadened , dispersion**

Structure of Sunny



Advantages of Sunny

1. Provides tools for calculating **classical phase diagrams** in addition to spin wave calculations.
2. Can perform spin wave calculations both **classically** (at finite temperature) and **semi-classically** (SWT).
3. Generalizes both the classical and SWT calculations in corresponding ways with **SU(N) coherent states**.
4. Is written in a language with dynamic features (as easy as Python) that is nonetheless compiled (as fast as C++) – both **fast** and **interactive**.

Summary of Sunny “modes”

Modes when creating a system

- **:dipole_large_S**
 - Both LL and LSWT are performed using the traditional $S \rightarrow \infty$ limit
 - 1 mode per site
- **:SUN**
 - Both LL and LSWT are performed using the $SU(N)$ formalism
 - $(N - 1)$ modes per site
- **:dipole**
 - Use $SU(N)$ formalism but constrain to dipolar states, leading to renormalizations
 - 1 mode per site

Summary of Sunny “modes”

Guidance about when to use each mode

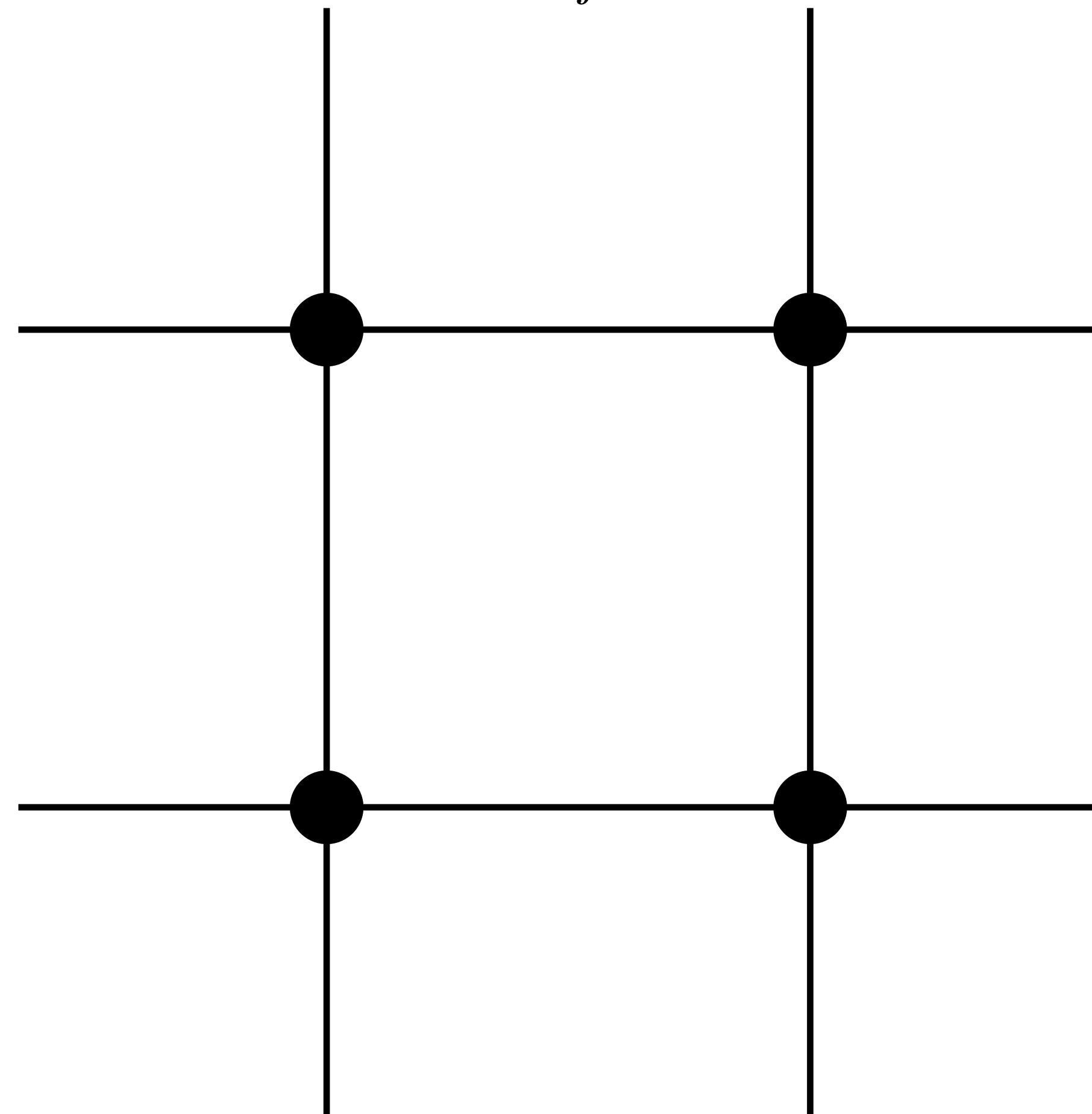
- **:dipole_large_S**
 - For reproducing existing results (verification)
 - Will differ from :dipole mode only when there are nonlinear terms
- **:SUN**
 - When you have large spins with nonlinear terms
 - When you want to model “localized” entanglement
- **:dipole**
 - For Hamiltonians that have a predominantly dipolar character
 - Nonlinear terms small compared to exchange terms.

General structure of SU(2) methods

Product State Assumption

- In traditional (large-S) theories, decompose into product of SU(2) coherent states.

$$|\Psi\rangle = \bigotimes_j |\Psi_j\rangle$$

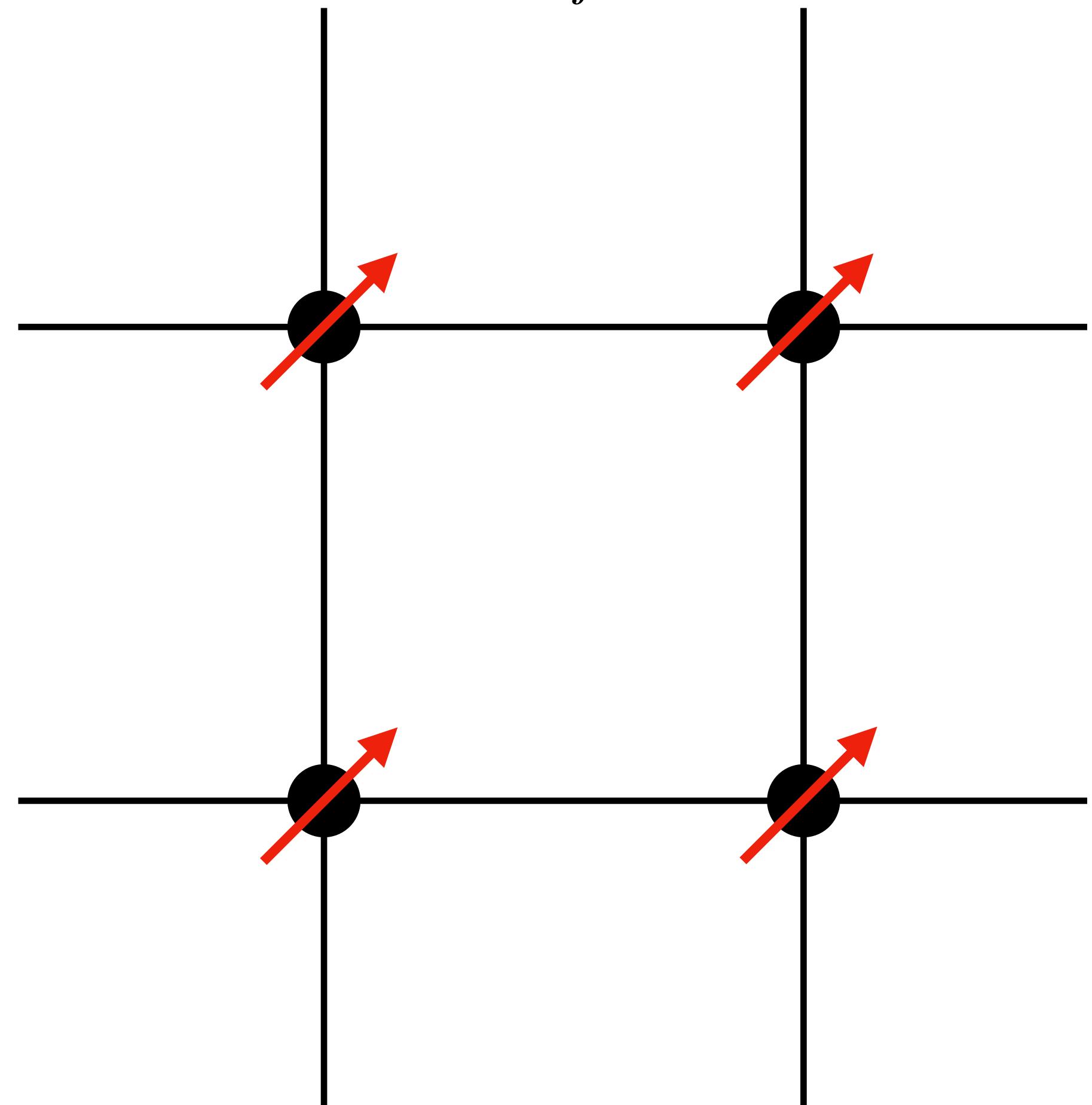


General structure of SU(2) methods

Product State Assumption

- In traditional (large-S) theories, decompose into product of SU(2) coherent states.
- SU(2) coherent states can always be put into one-to-one correspondence with points on a sphere

$$|\Psi\rangle = \bigotimes_j |\Psi_j\rangle$$

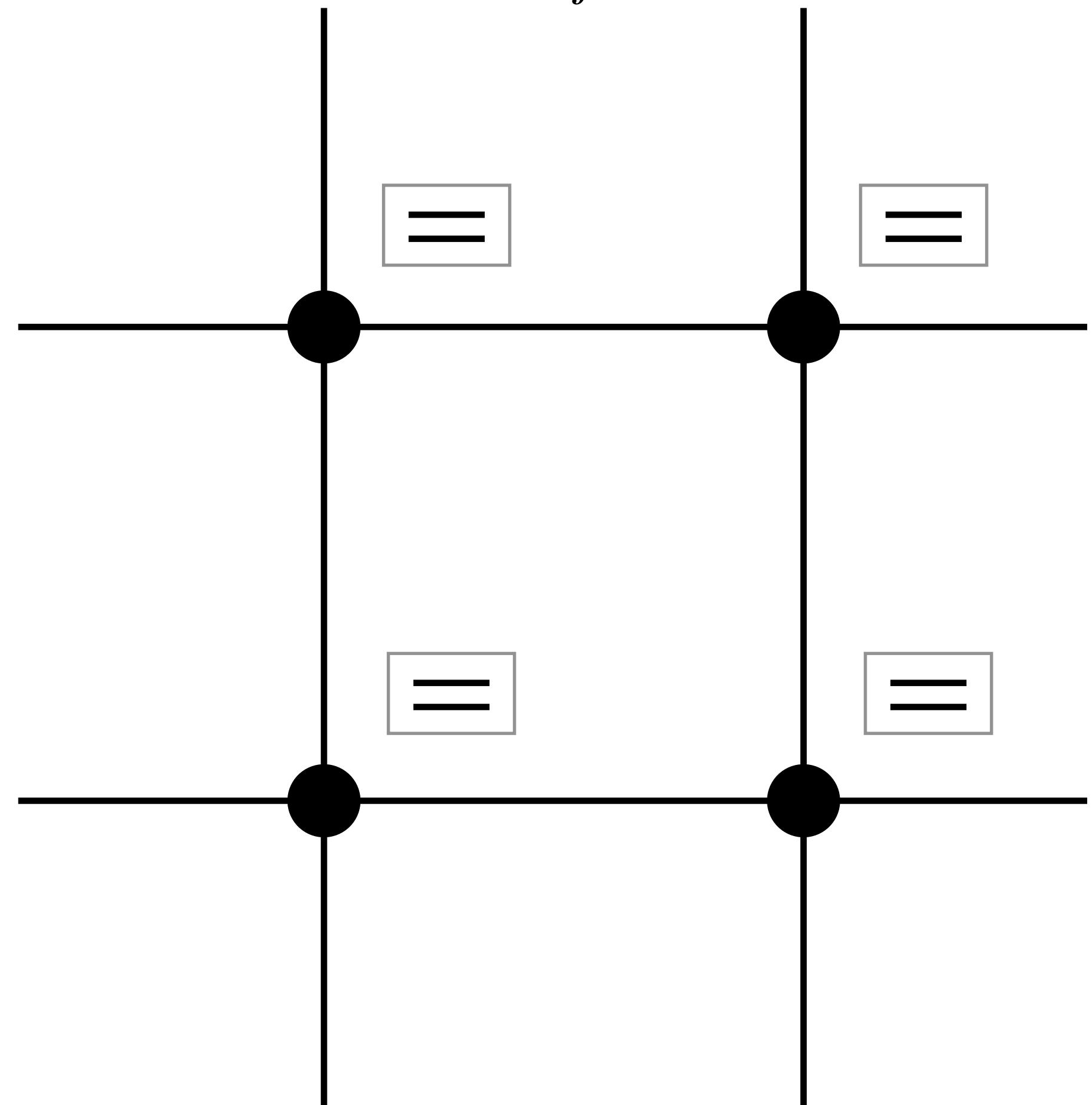


General structure of SU(2) methods

Product State Assumption

- In traditional (large-S) theories, decompose into product of SU(2) coherent states.
- SU(2) coherent states can always be put into one-to-one correspondence with points on a sphere
- This is isomorphic to the state space for a 2-level system

$$|\Psi\rangle = \bigotimes_j |\Psi_j\rangle$$



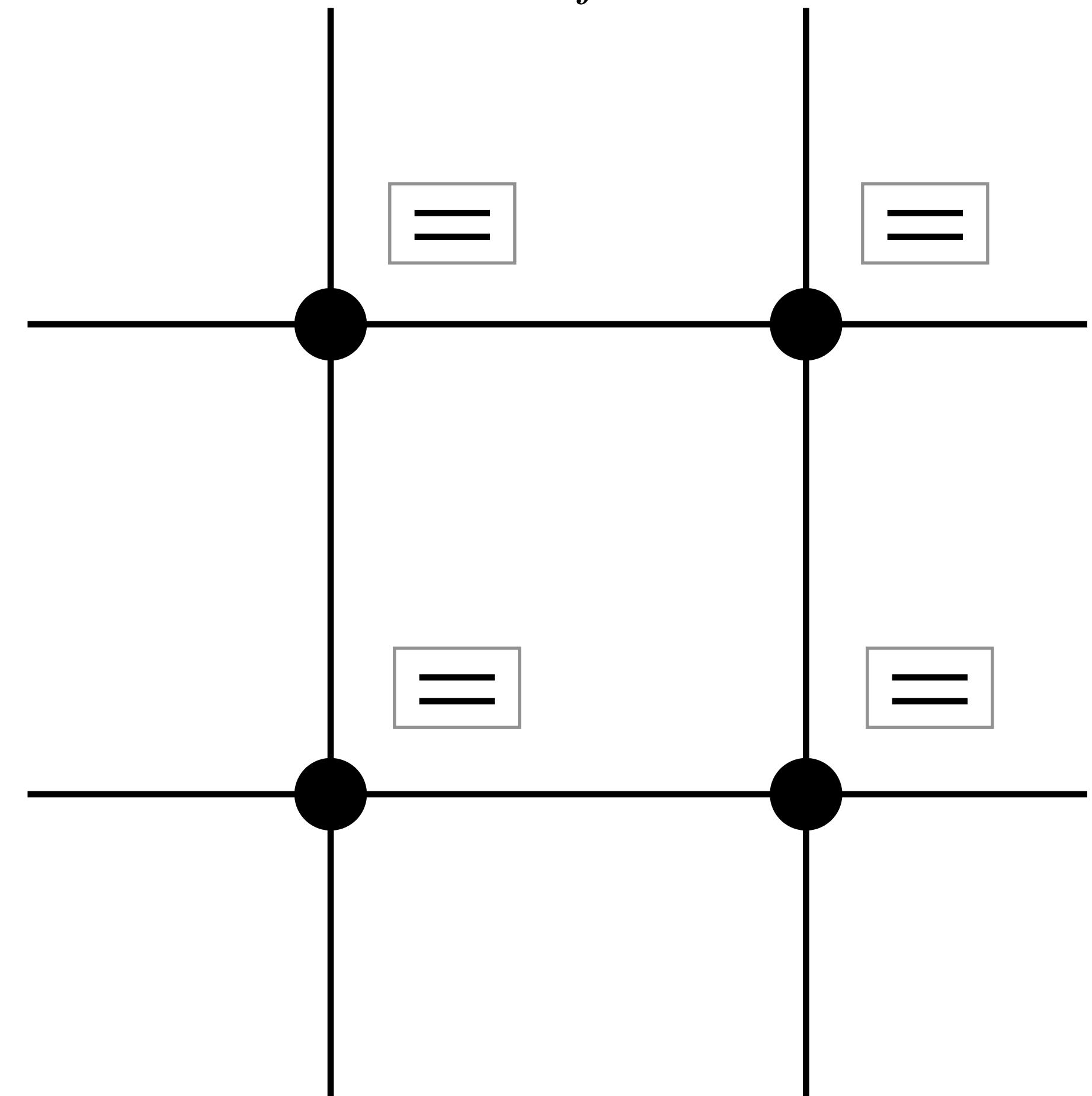
General structure of SU(2) methods

Classical dynamics

- Well defined procedures for deriving **classical dynamics** ($S \rightarrow \infty$)

$$\frac{d\mathbf{s}_j}{dt} = -\mathbf{s}_j \times \nabla_{\mathbf{s}_j} H_{SU(2)}$$

$$|\Psi\rangle = \bigotimes_j |\Psi_j\rangle$$



General structure of SU(2) methods

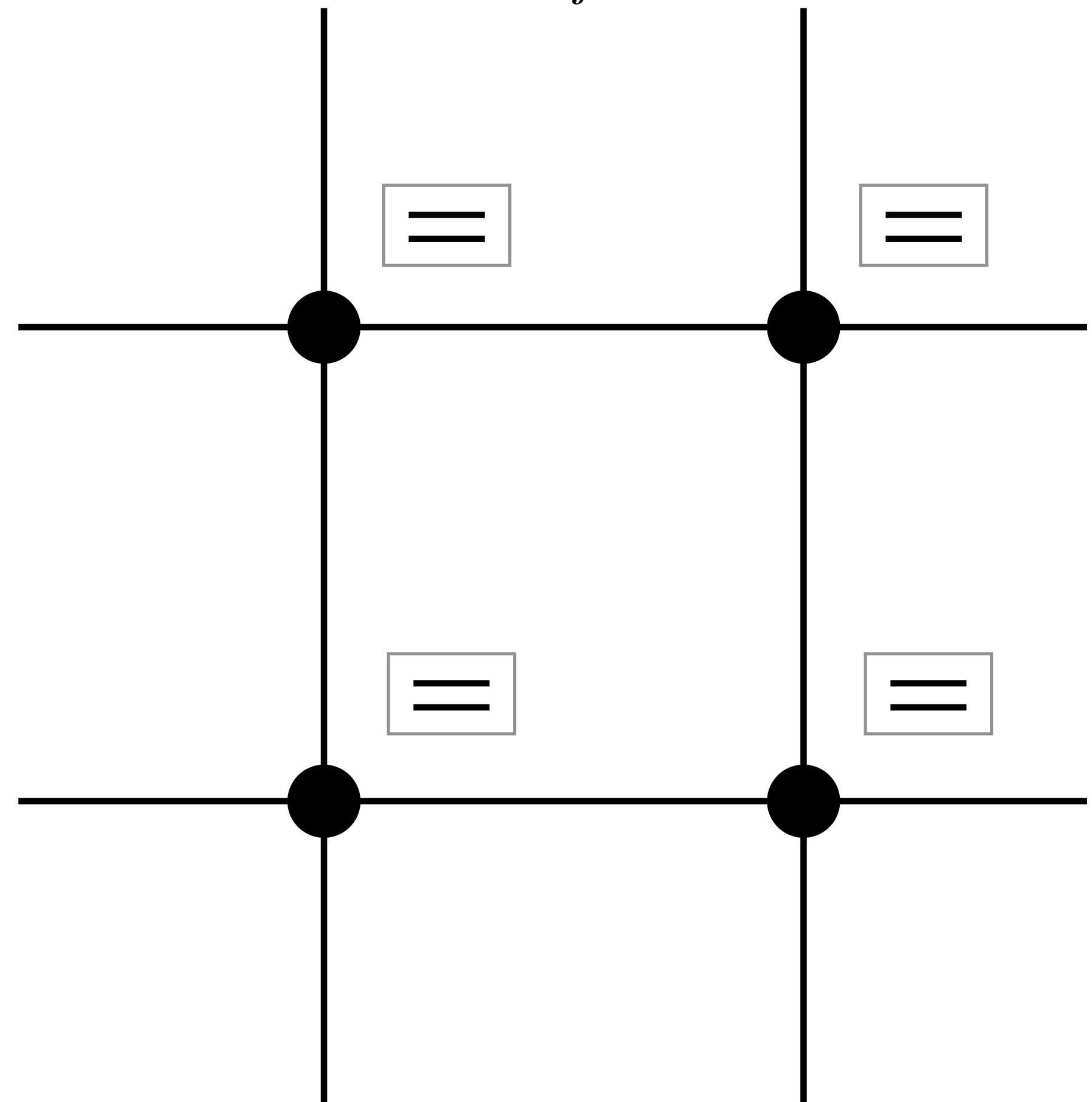
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General structure of SU(2) methods

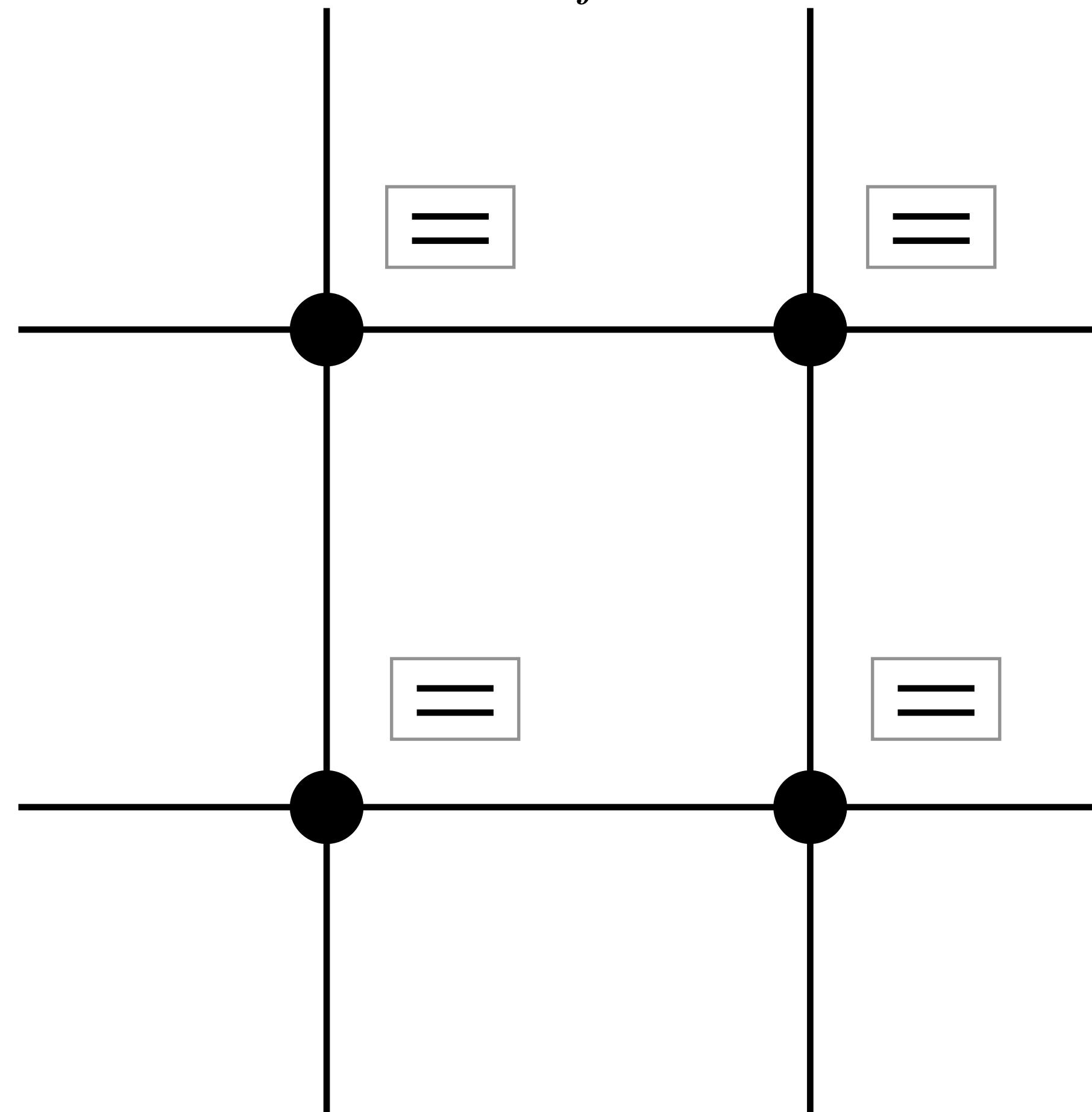
Semiclassical dynamics

- Well defined procedures **bosonizing a Hamiltonian** (expansion in $1/S$)

$$\{ b_{j,0}, b_{j,1} \} \quad b_{0,j}^\dagger b_{0,j} + b_{1,j}^\dagger b_{1,j} = 2S$$

$$b_{j,0}^\dagger = b_{j,0} = \sqrt{2S} \sqrt{1 - \frac{b_{j,1}^\dagger b_{j,1}}{2S}}$$

$$|\Psi\rangle = \bigotimes_j |\Psi_j\rangle$$

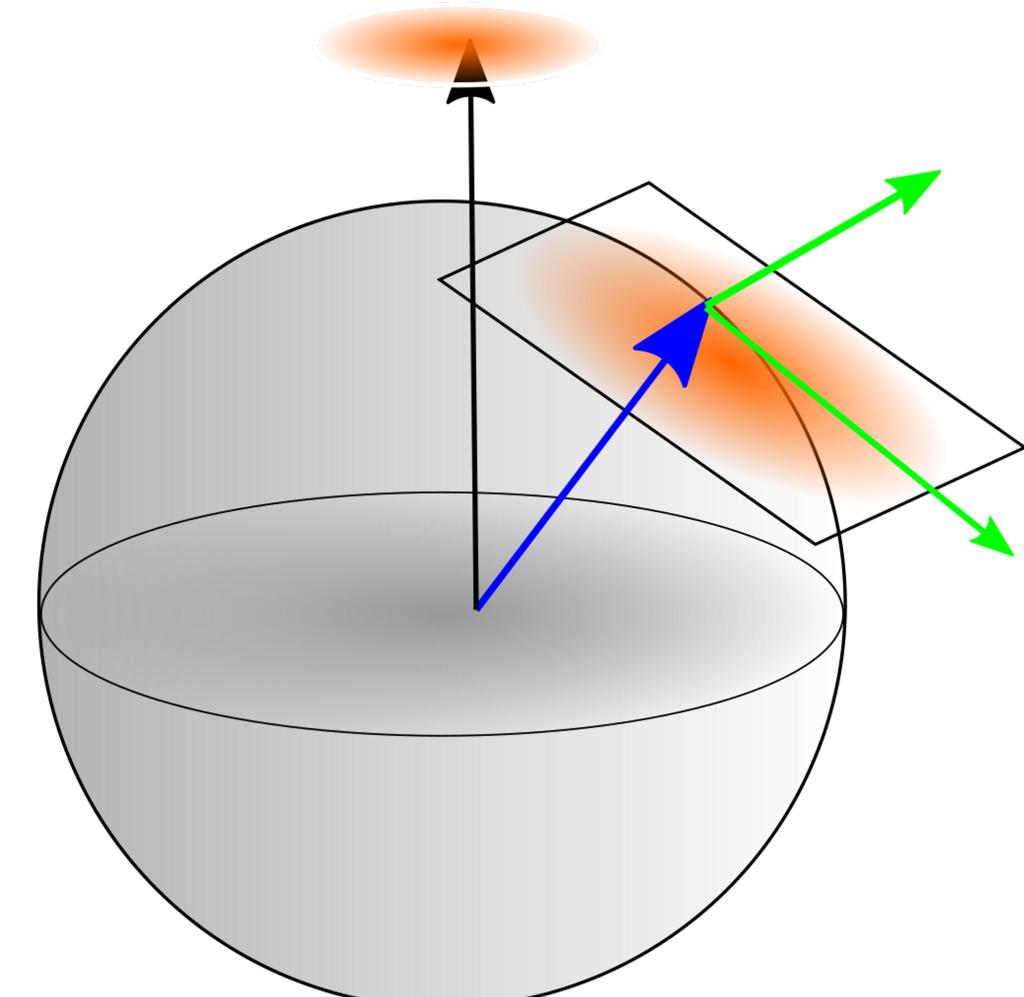


Correspondence between LL and LSWT

A few assertions

- LL and LSWT connected by SU(2) coherent states and the use of S as a control parameter.
- Linearized LL about the ground state yields same dispersion as LSWT
- Can recover quantum behavior by quantizing the normal modes of the classical solution

$$\mathbf{s}_j = \mathbf{s}_{0,j} + \delta\mathbf{s}_j$$



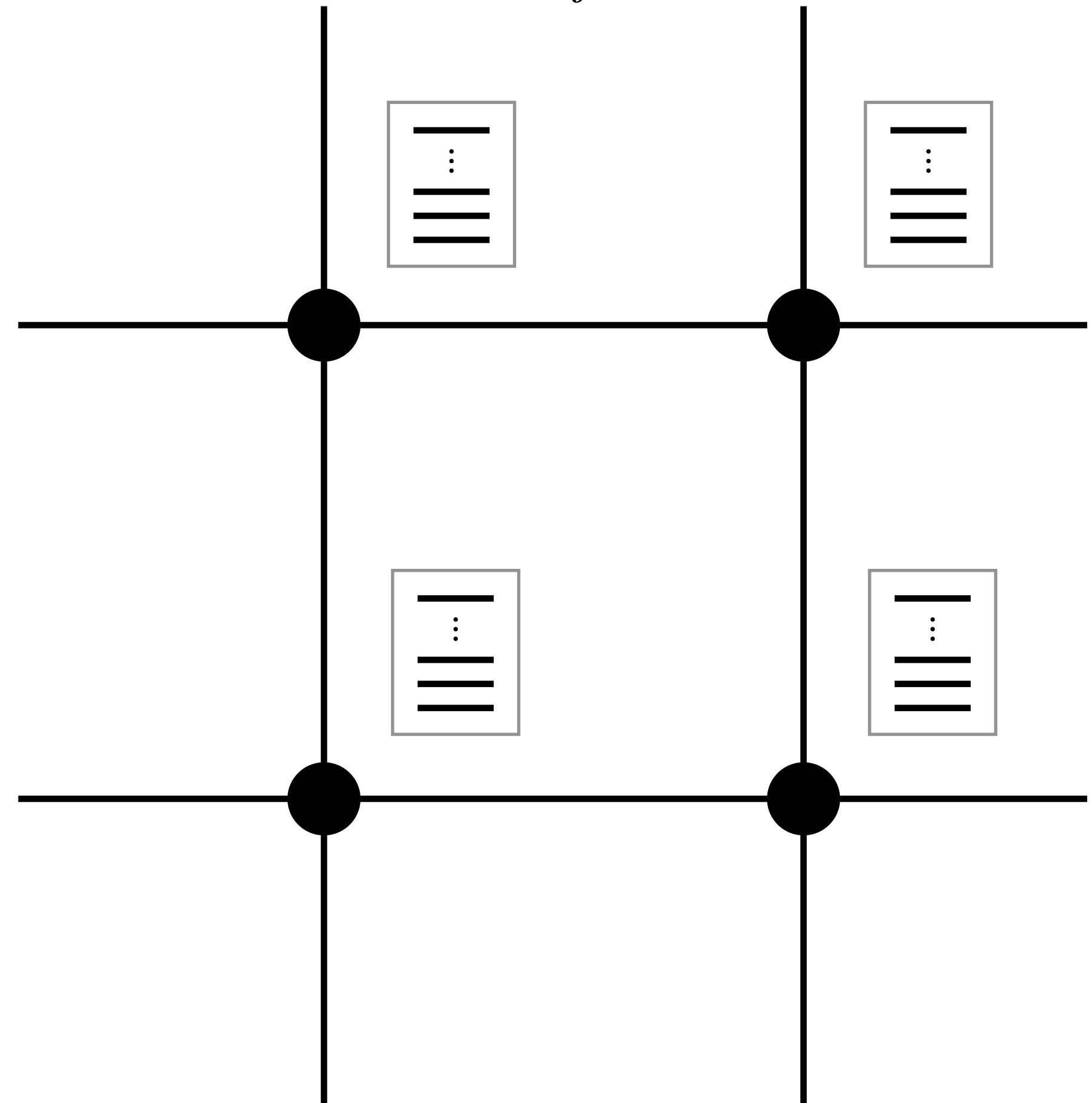
$$S_Q^H(\mathbf{q}, \omega) = \frac{\hbar\omega}{k_B T} [1 + n_B(\omega/T)] S_{cl}^H(\mathbf{q}, \omega)$$

General structure of $SU(N)$ methods

Product State Assumption

- Again decompose into a product state.
- This time put an $SU(N)$ coherent state on each site
- Retain all the richness of an N -level system – don't reduce to a dipole.

$$|\Psi\rangle = \bigotimes_j |\Psi_j\rangle$$



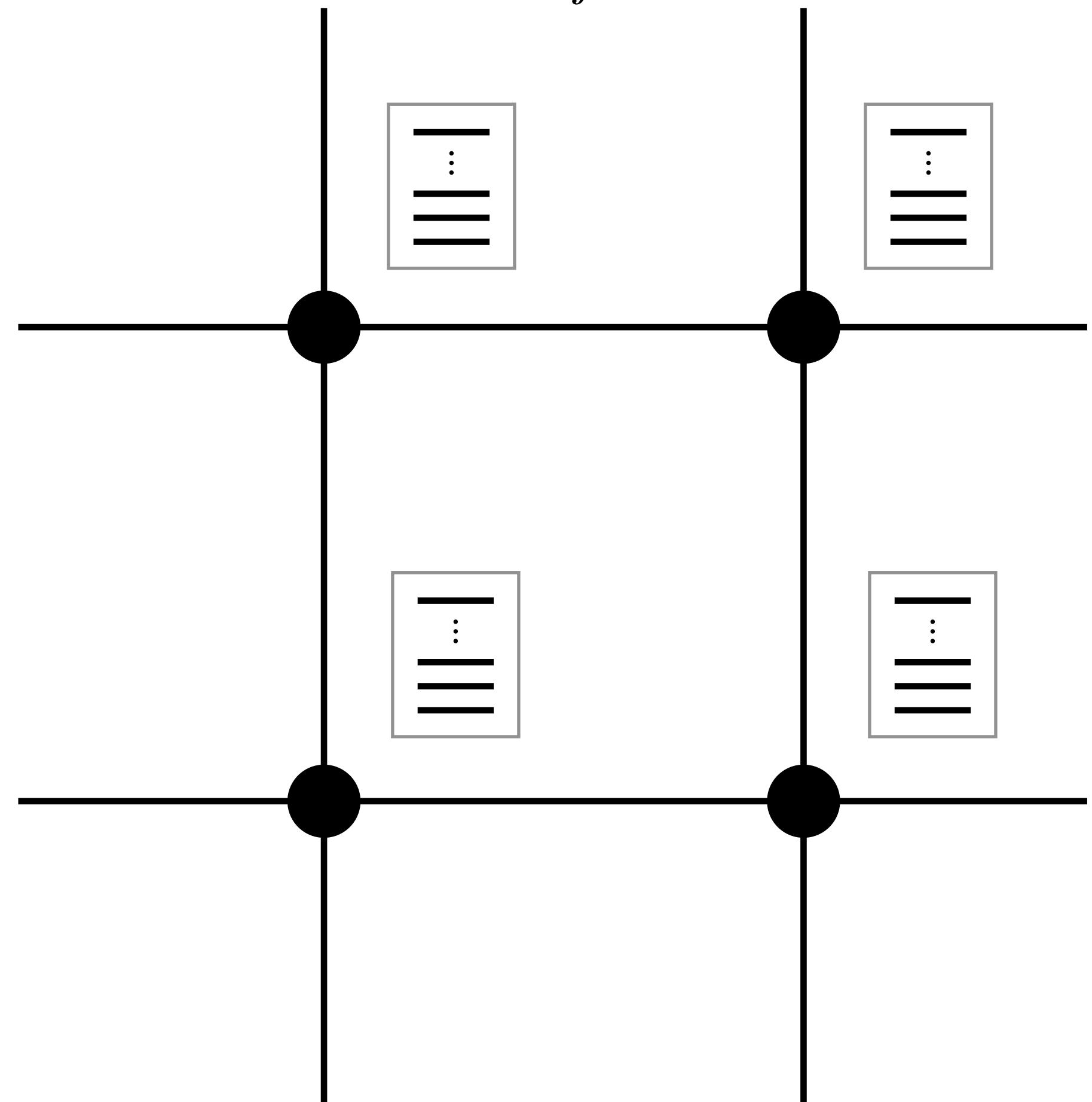
General structure of $SU(N)$ methods

Classical dynamics

- Well defined procedure for defining a classical dynamics.

$$\frac{d\mathbf{n}_j}{dt} = - \mathbf{n}_j \star \nabla_{\mathbf{n}_j} H_{SU(N)}$$

$$|\Psi\rangle = \bigotimes_j |\Psi_j\rangle$$



General structure of $SU(N)$ methods

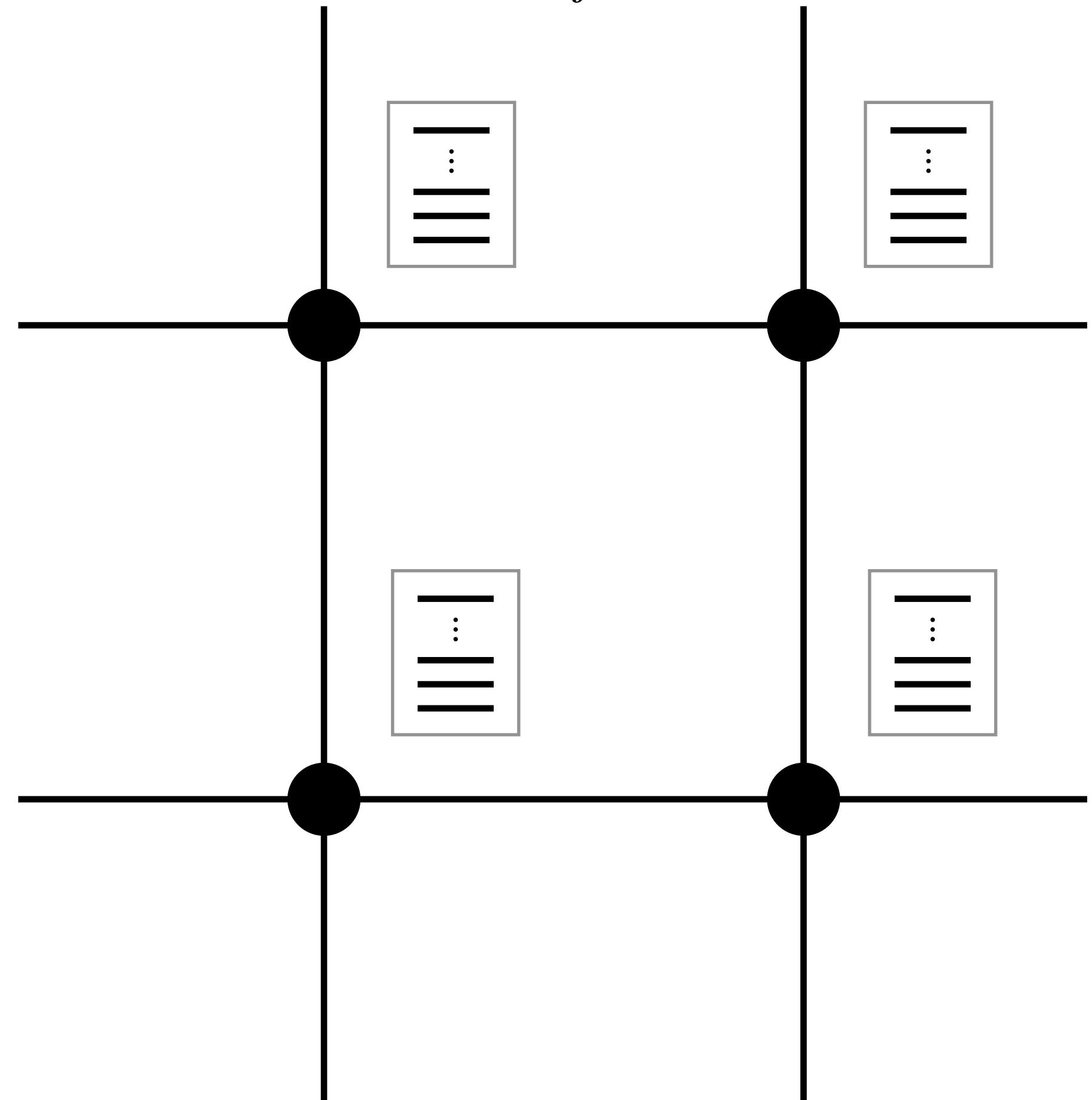
Classical dynamics

- Well defined procedure for bosonizing a Hamiltonian in N -flavors

$$\left\{ b_{j,0}, \dots, b_{j,N-1} \right\} \quad \sum_{m=0}^{N-1} b_{j,m}^\dagger b_{j,m} = M$$

$$b_{j,0}^\dagger = b_{j,0} = \sqrt{M} \sqrt{1 - \frac{\sum_{m=1}^{N-1} b_{j,m}^\dagger b_{j,m}}{M}}$$

$$|\Psi\rangle = \bigotimes_j |\Psi_j\rangle$$

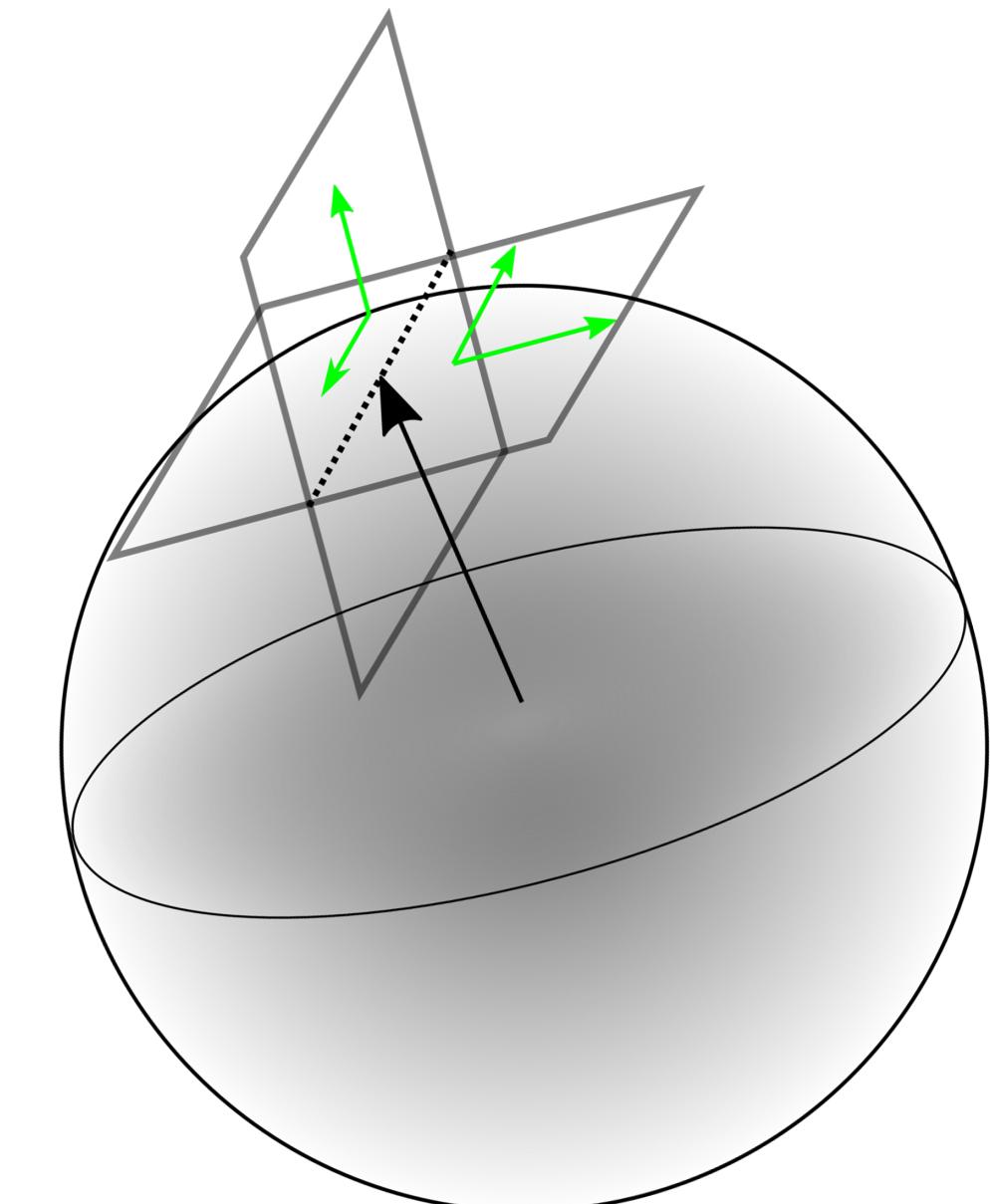


Correspondence between LL and LSWT

SU(N) generalization

- Generalized LL (GLL) and multi-flavor LSWT (GLSWT) are connected by SU(N) coherent states and the use of M as a control parameter.
- Linearized GLL about the ground state yields same dispersion as GLSWT
- The generalizations of both the classical and semiclassical theories yield $N - 1$ modes per site

$$\mathbf{n}_j = \mathbf{n}_{0,j} + \delta\mathbf{n}_j$$



SU(N) generalization

Motivation: Ground states

- Consider a simple single-site problem.

$$\begin{aligned}\mathcal{H} &= D (S^z)^2 \\ &= D \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}\end{aligned}$$

$$===== \langle \pm 1 | \mathcal{H} | \pm 1 \rangle = D$$

$$----- \langle 0 | \mathcal{H} | 0 \rangle = 0$$

SU(N) generalization

Motivation: Ground states

- The ground state is clearly $|0\rangle$, which is non-magnetic:

$$\langle 0 | S^x | 0 \rangle = \langle 0 | S^y | 0 \rangle = \langle 0 | S^z | 0 \rangle = 0$$

- But $|0\rangle$ is not an $SU(2)$ coherent state. Recall the $S = 1$ parameterization:

$$|\Omega(\theta, \phi)\rangle = e^{i\phi} \cos^2\left(\frac{\theta}{2}\right) |1\rangle + \sqrt{2} \cos\left(\frac{\theta}{2}\right) \sin\left(\frac{\theta}{2}\right) |0\rangle + e^{-i\phi} \sin^2\left(\frac{\theta}{2}\right) |-1\rangle$$

SU(N) generalization

Motivation: Dynamics

$$\hat{\mathcal{H}}_{SI} = D(\hat{S}^z)^2 - BS^z$$

Dipolar precession

$$\omega_D = (D + B)/\hbar$$

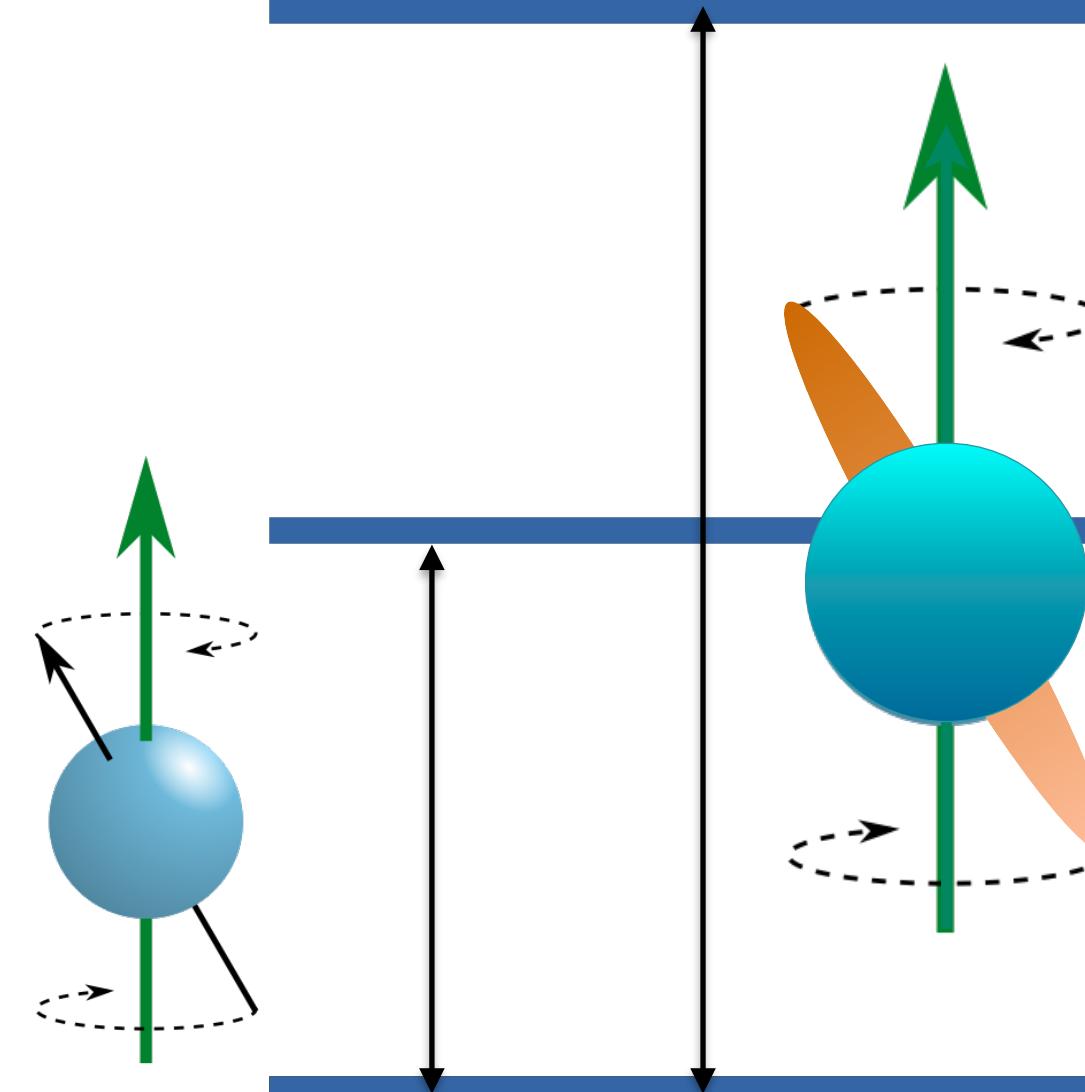
Quadrupolar precession

$$\omega_Q = 2B/\hbar$$

$| \bar{1} \rangle$

$| 0 \rangle$

$| 1 \rangle$



$$\epsilon = D + B$$

$$\epsilon = 0$$

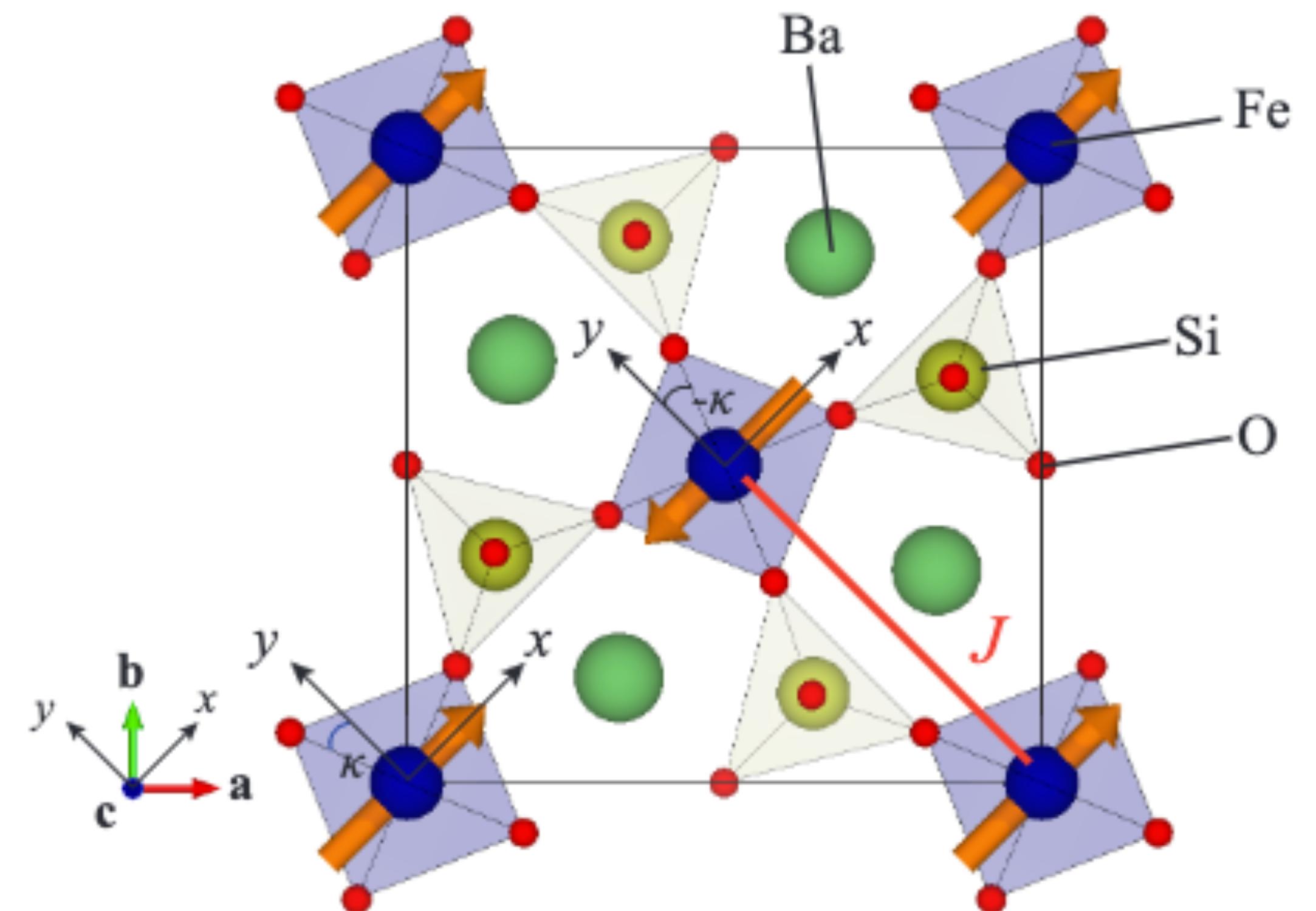
$$\epsilon = D - B$$

SU(N) use case: Ba₂FeSi₂O₇

Rich behavior of $S = 2$ spins

- AFM exchange and hard-axis anisotropy on a square lattice.
- Quasi-2D – weak AFM coupling between layers.
- Effective $S = 2$, so use SU(5) with local basis

$$|2\rangle, |1\rangle, |0\rangle, |-1\rangle, |-2\rangle$$



SU(N) use case: Ba₂FeSi₂O₇

Rich behavior of $S = 2$ spins

- AFM exchange and hard-axis anisotropy on a square lattice.
- Quasi-2D – weak AFM coupling between layers.
- Effective $S = 2$, so use SU(5) with local basis

$$|2\rangle, |1\rangle, |0\rangle, |-1\rangle, |-2\rangle$$

$$\begin{aligned}\mathcal{H} = & J \sum_{i,j} \mathbf{S}_i \cdot \mathbf{S}_j + J' \sum_{\langle\langle i,j \rangle\rangle} \mathbf{S}_i \cdot \mathbf{S}_j \\ & + D \sum_i (S_i^z)^2 - h \sum_i S_i^z \\ & + A \sum_{i \in A} \left[(\mathbf{S}_i \cdot \mathbf{n}_1^A)^4 + (\mathbf{S}_i \cdot \mathbf{n}_2^A)^4 \right] \\ & + A \sum_{j \in B} \left[(\mathbf{S}_j \cdot \mathbf{n}_1^B)^4 + (\mathbf{S}_j \cdot \mathbf{n}_2^B)^4 \right] \\ & + C \sum_i (S_i^z)^4,\end{aligned}$$

SU(N) use case: Ba₂FeSi₂O₇

Rich behavior of $S = 2$ spins

- Magnetization versus field

PRB **107** 144427 (2023)

- Generalized LSWT calculations

Nature Communications 12:5331 (2021)

- Finite temperature behavior

npj quantum materials 5 (2023)

Field-induced spin level crossings within a quasi-XY antiferromagnetic state in Ba₂FeSi₂O₇

Minseong Lee  ^{1,*}, Rico Schönenmann  ¹, Hao Zhang  ^{2,3}, David Dahlbom  ³, Tae-Hwan Jang  ⁴, Seung-Hwan Do  ², Andrew D. Christianson  ², Sang-Wook Cheong  ^{4,5}, Jae-Hoon Park  ^{4,6}, Eric Brosha  ⁷, Marcelo Jaime  ¹, Kipton Barros  ⁸, Cristian D. Batista  ^{3,†} and Vivien S. Zapf  ^{1,‡}

Decay and renormalization of a longitudinal mode in a quasi-two-dimensional antiferromagnet

Seung-Hwan Do  ^{1,9}, Hao Zhang  ^{1,2,9}, Travis J. Williams  ³, Tao Hong  ³, V. Ovidiu Garlea  ³, J. A. Rodriguez-Rivera  ^{4,5}, Tae-Hwan Jang ⁶, Sang-Wook Cheong  ^{6,7}, Jae-Hoon Park  ^{6,8}, Cristian D. Batista  ² & Andrew D. Christianson  ¹

Understanding temperature-dependent SU(3) spin dynamics in the $S=1$ antiferromagnet Ba₂FeSi₂O₇

Seung-Hwan Do  ^{1,8}, Hao Zhang  ^{1,2,8}, David A. Dahlbom  ², Travis J. Williams  ³, V. Ovidiu Garlea ³, Tao Hong  ³, Tae-Hwan Jang  ⁴, Sang-Wook Cheong  ^{4,5}, Jae-Hoon Park  ^{4,6}, Kipton Barros ⁷, Cristian D. Batista  ² and Andrew D. Christianson  ¹

SU(N) use case: Ba₂FeSi₂O₇

M and staggered magnetization vs H

- Zero field:

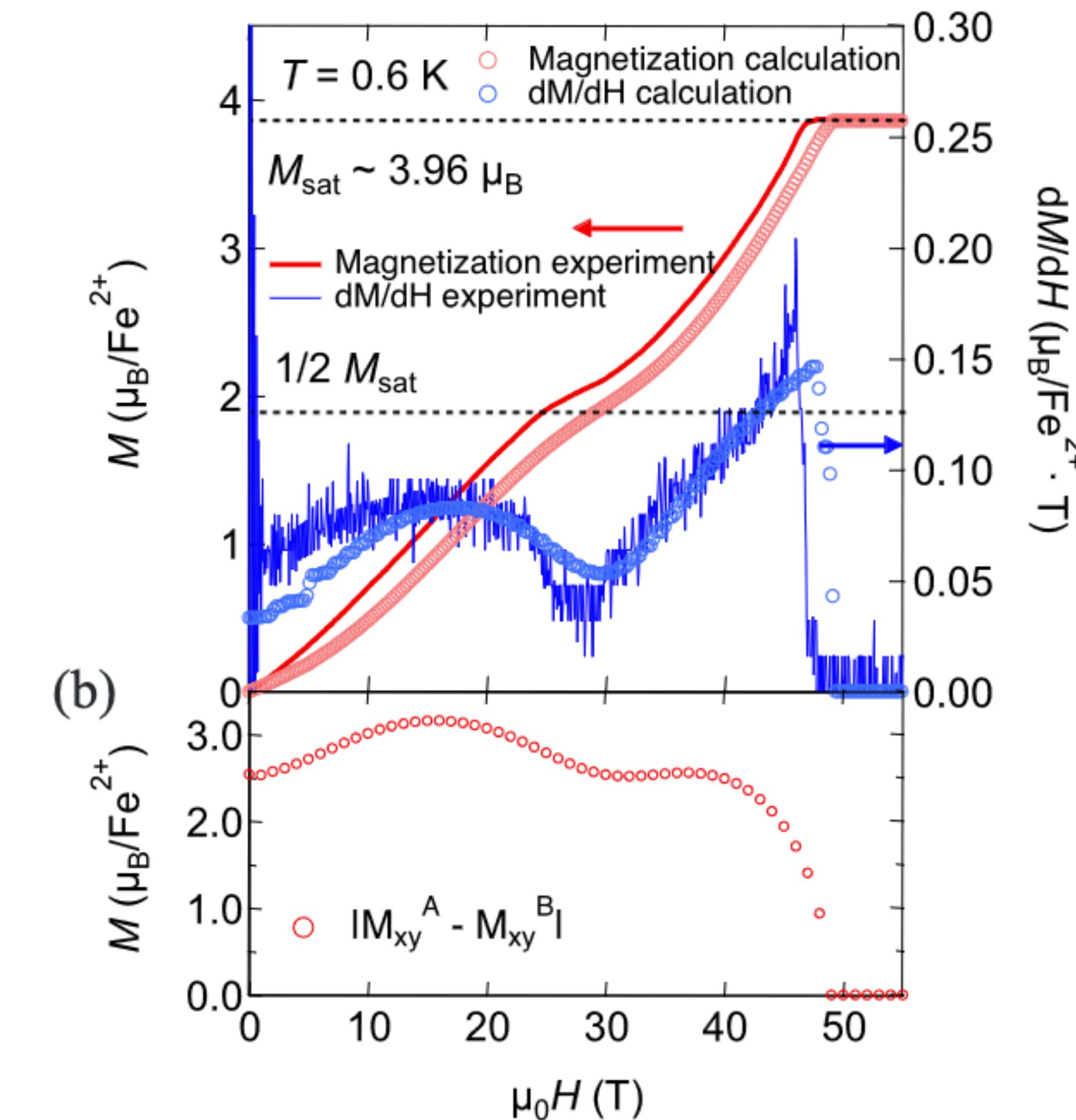
$|1\rangle, |0\rangle$

- Applied field:

$|2\rangle, |0\rangle$

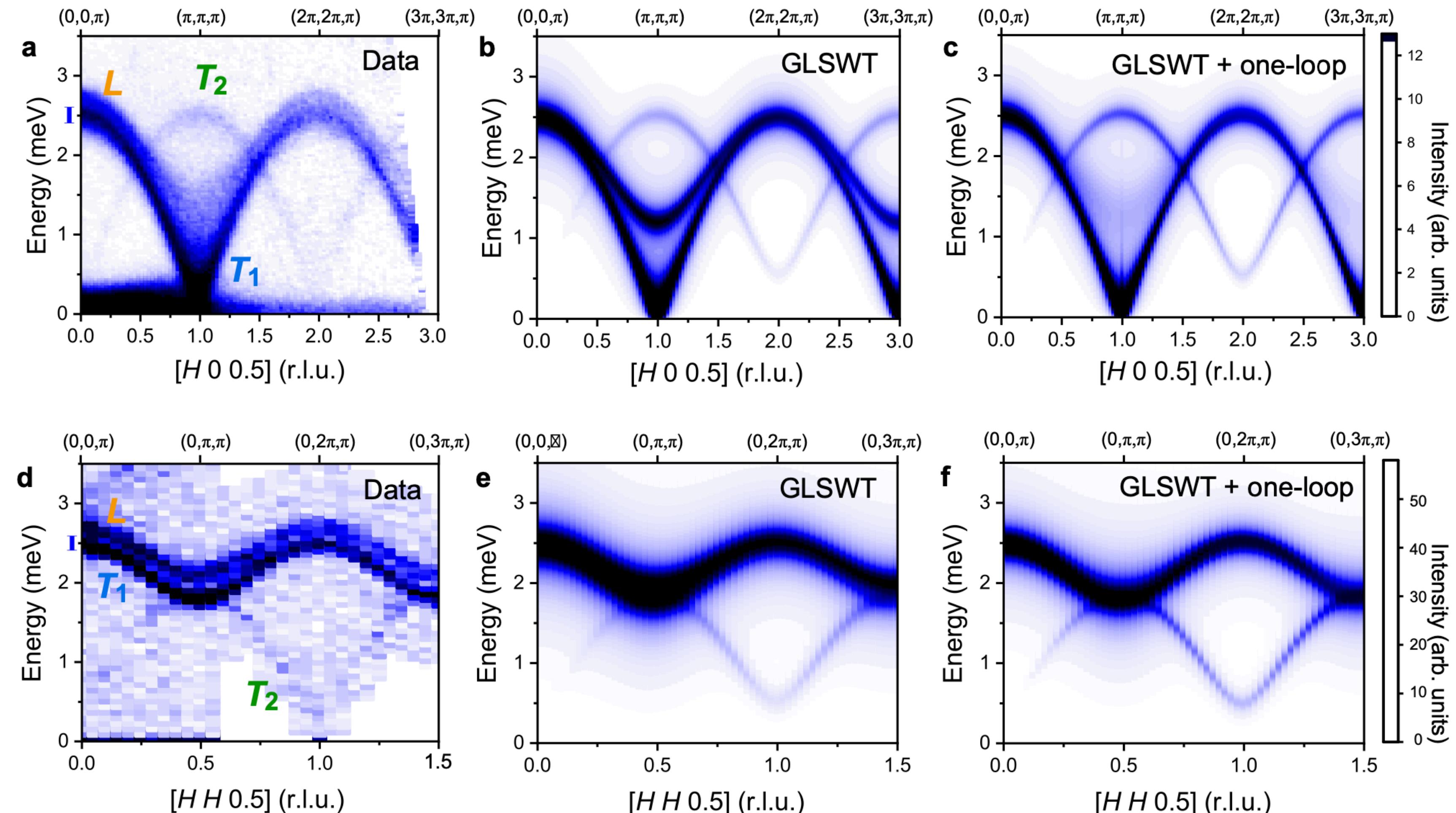
- Fully saturated:

$|2\rangle$



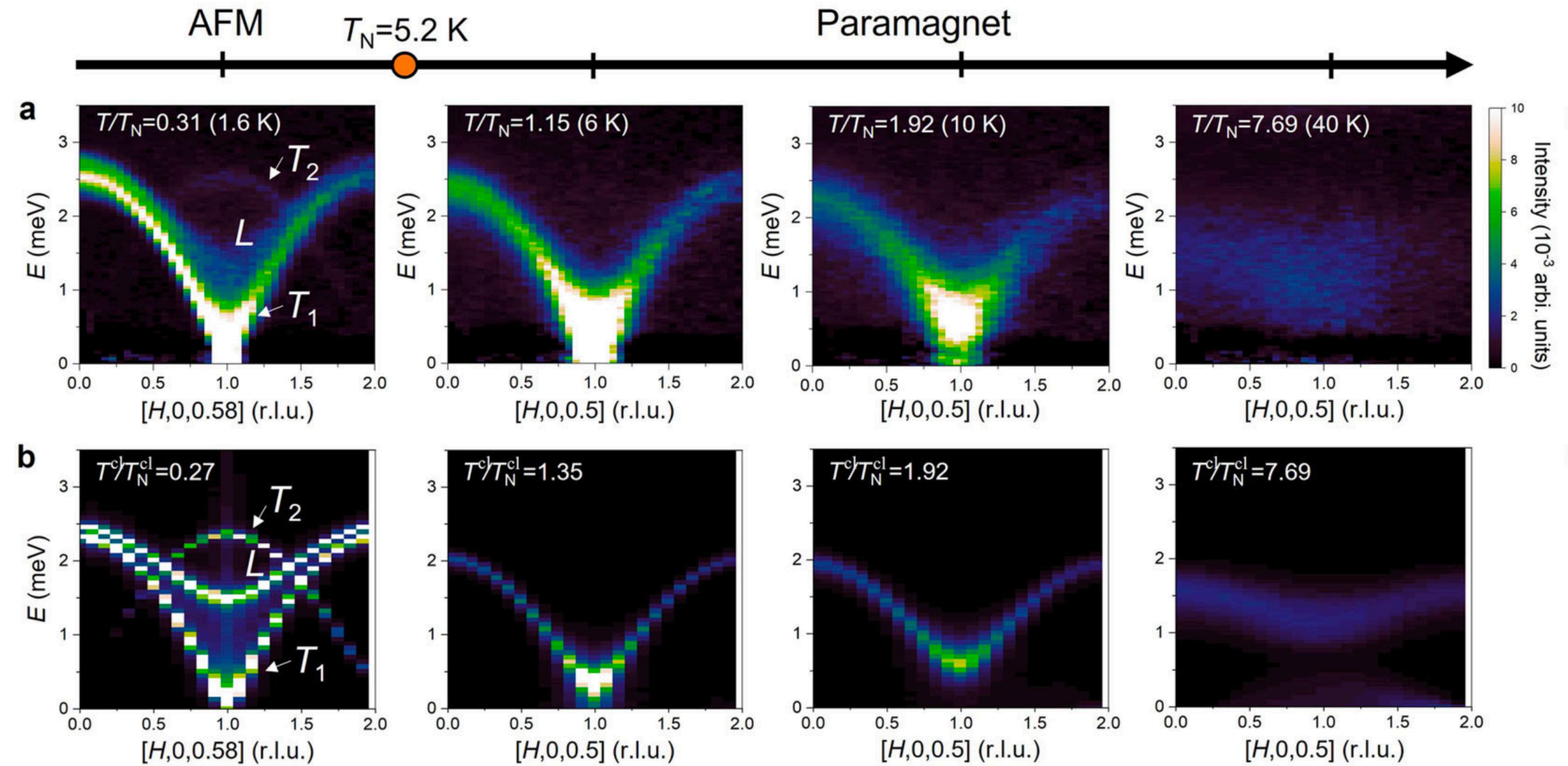
SU(N) use case: Ba₂FeSi₂O₇

Multi-flavor spin wave calculation



SU(N) use case: Ba₂FeSi₂O₇

Finite temperature



Takeaways

What Sunny is

- A collection of tools for
 - Performing **symmetry analysis** and specifying spin Hamiltonians

Takeaways

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 - Finding **classical ground states** and constructing phase diagrams

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Takeaways

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- A collection of tools for
 - Performing **symmetry analysis** and specifying spin Hamiltonians
 - Finding **classical ground states** and constructing phase diagrams
 - Calculating **structure factors** both classically and semiclassically
 - Generating **intensity output** that can be compared directly to neutron scattering data

Takeaways

What Sunny is

- A collection of tools for
 - Performing **symmetry analysis** and specifying spin Hamiltonians
 - Finding **classical ground states** and constructing phase diagrams
 - Calculating **structure factors** both classically and semiclassically
 - Generating **intensity output** that can be compared directly to neutron scattering data
- It is also a broad collaboration dedicated to making general, high-quality implementations of theoretical techniques available to neutron scatterers.

Takeaways

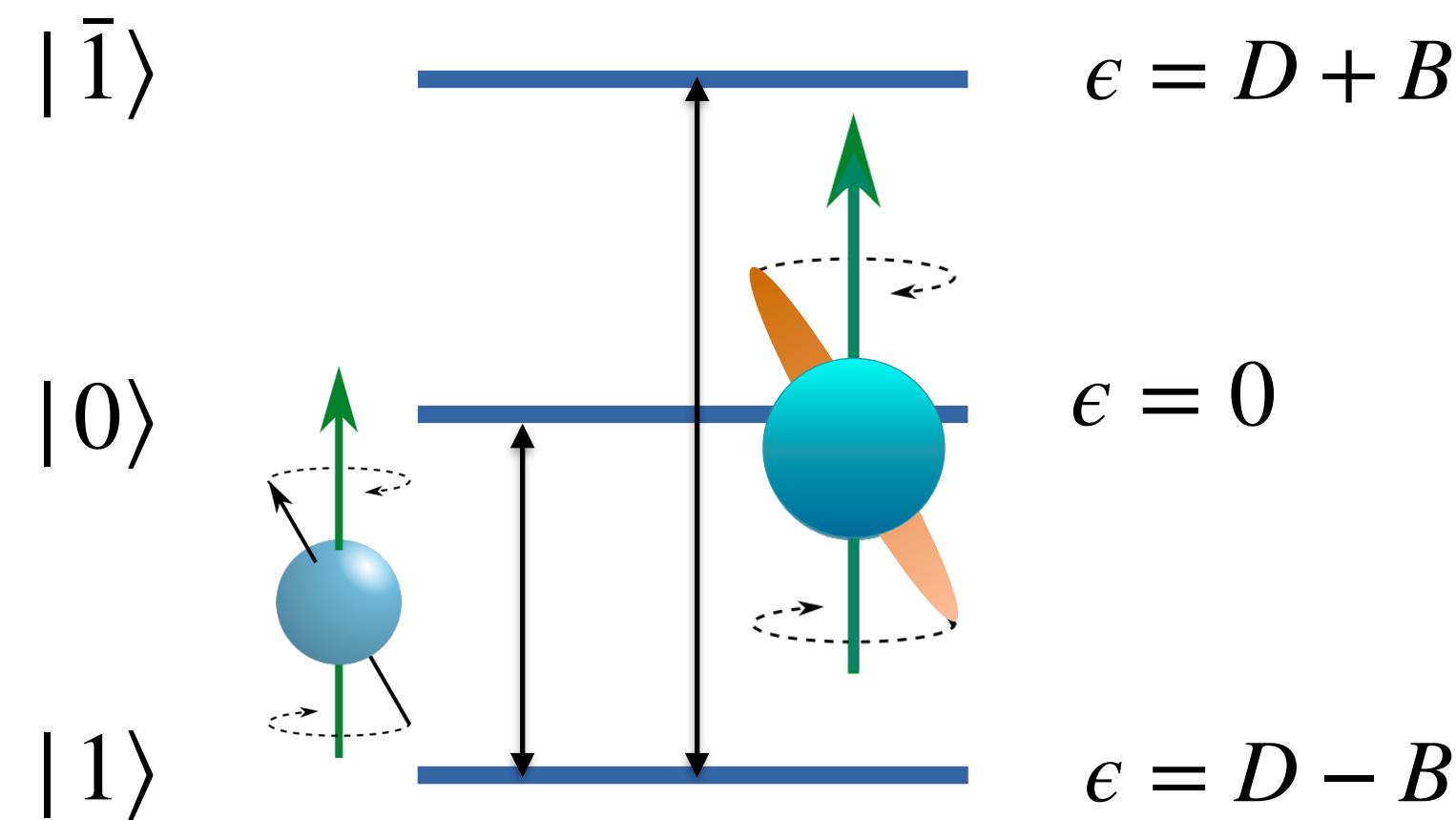
What $SU(N)$ is good for

- Capture more **local physics** at the linear level

Takeaways

What $SU(N)$ is good for

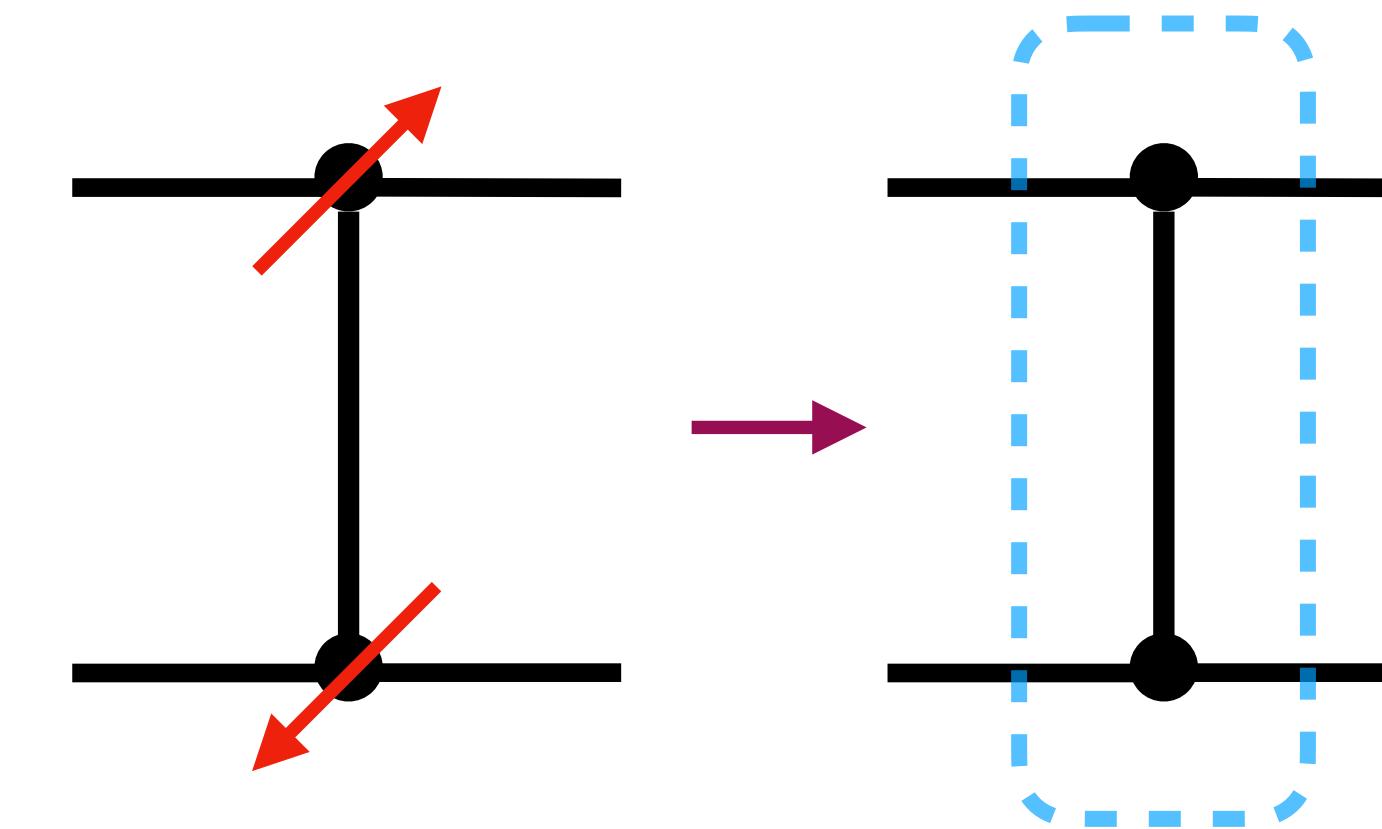
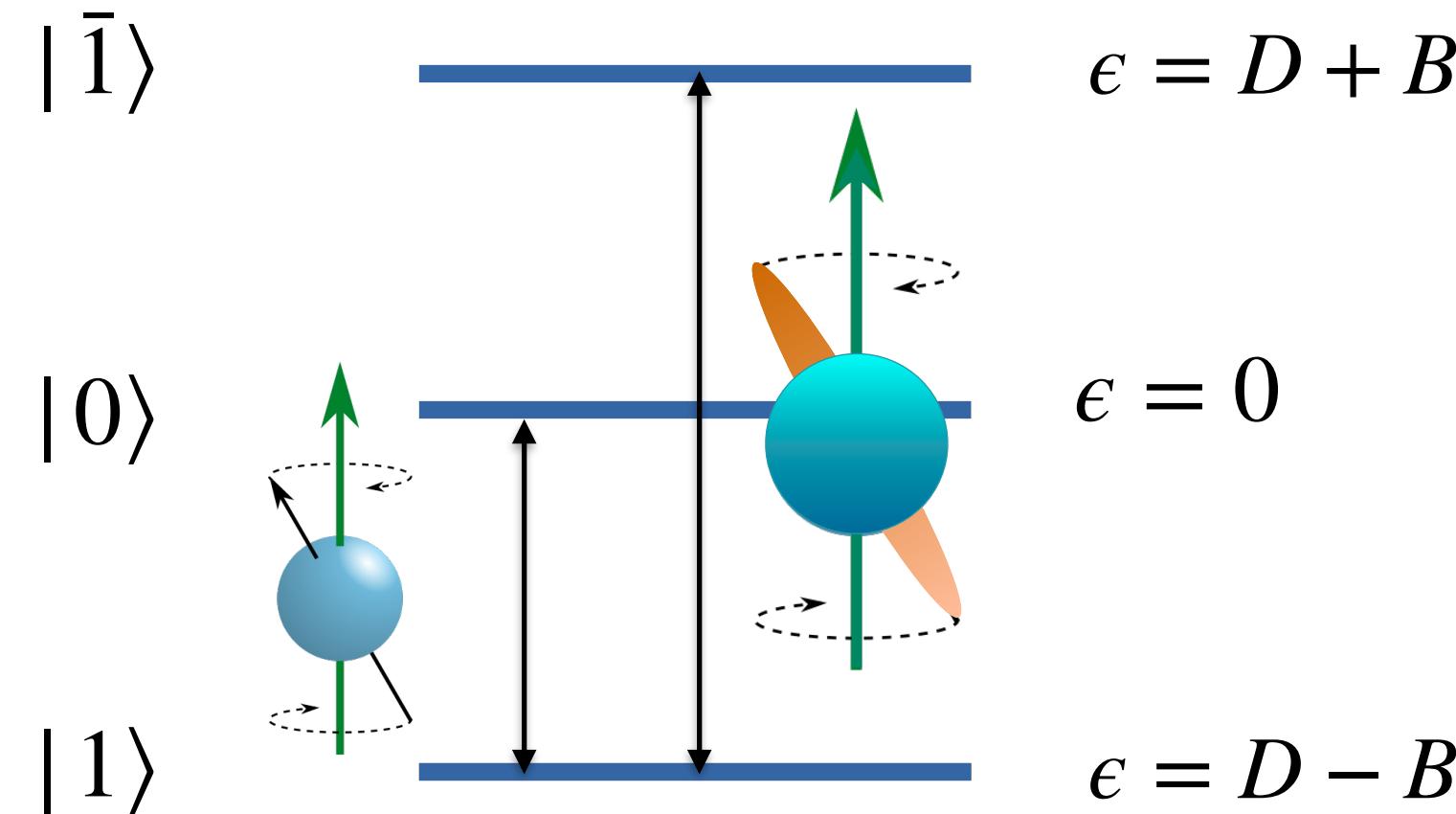
- Capture more **local physics** at the linear level
- Useful for **large spins** with “nonlinear” operators in your Hamiltonian (4d-5d, 4f-5f and some 3d materials).



Takeaways

What $SU(N)$ is good for

- Capture more **local physics** at the linear level
- Useful for **large spins** with “nonlinear” operators in your Hamiltonian (4d-5d, 4f-5f and some 3d materials).
- Model “localized” entanglement (dimers, trimers, tetramers, etc.).



Takeaways

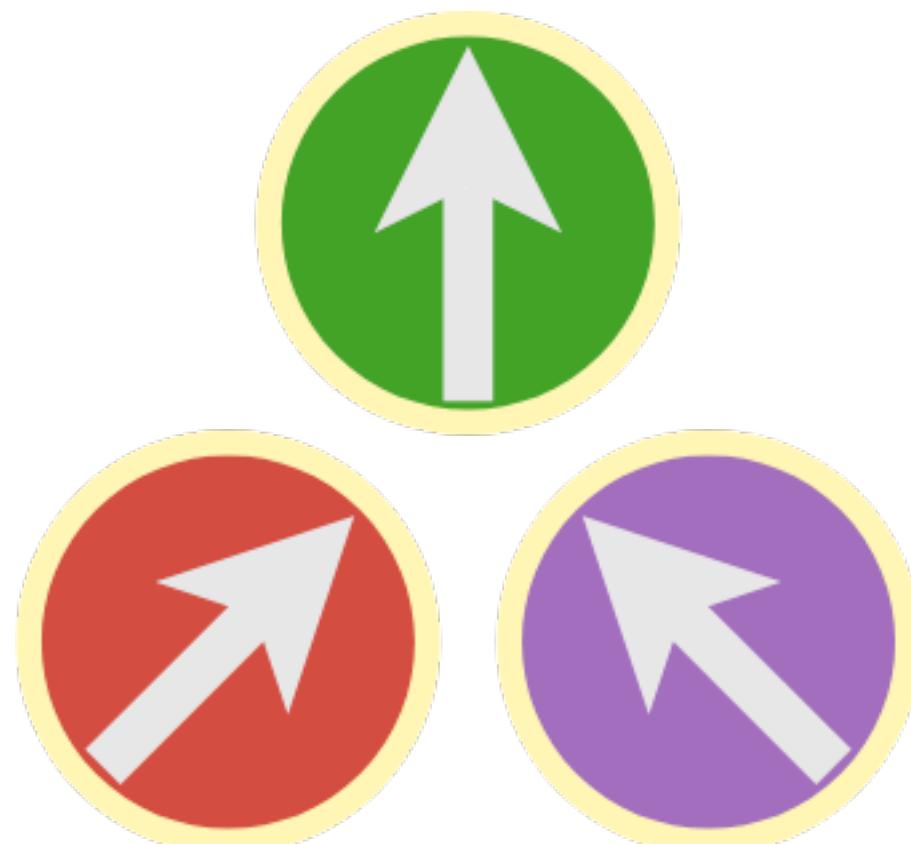
What makes Sunny unique

- Powerful set of tools for analyzing symmetries, visualizing crystals, and specifying a **broad class of spin Hamiltonians**.
- Easy access to both **classical and semiclassical techniques**, which are coupled together productively.
- Generalization of both classical and semiclassical dynamics to **SU(N) coherent states**.
- Written in **Julia**, which is both interactive and fast.

Takeaways

How to get involved

- Visit the GitHub site (<https://github.com/SunnySuite/Sunny.jl>)
- Join the Sunny Slack (link on bottom of GitHub page)
- Talk to one of us!



Su(n)ny