

these notes overlap w lec 14.

NOTE ON HW

HOW DO YOU INTEGRATE A GREEN'S FUNCTION TWICE?

$$\int_{y-\epsilon}^{y+\epsilon} \left[\frac{d^2 G}{dx^2} = \delta(x) \right] dx$$

$$\frac{dG^>(y+\epsilon)}{d(\arg)} - \frac{dG^<(y-\epsilon)}{d(\arg)} = 1$$

\uparrow $x=y+\epsilon$ \uparrow $x=y-\epsilon$
 $dx = d\epsilon$ $dx = -d\epsilon$

treat this as: $G^>(x < y) = G^>(x > y) = 0$

NOW INTEGRATE WRT ϵ $\leftarrow y$ IS FIXED
 ϵ IS A PARAMETER

$$\int_{-\delta}^{\delta} \frac{dG^>(y+\epsilon)}{d(y+\epsilon)} d\epsilon - \int_{-\delta}^{\delta} \frac{dG^<(y-\epsilon)}{d(y-\epsilon)} d\epsilon = \delta$$

$\downarrow \delta \rightarrow 0$ $\delta \rightarrow 0$

$$G^>(y) - \int_{y+\delta}^{y-\delta} \frac{dG^<(u)}{du} (-du) = 0$$

$$- [G^<(y-\delta) - G^<(y+\delta)]$$

\downarrow
0

$$\boxed{G^>(y) - G^<(y) = 0}$$

so G IS CONTINUOUS

G' IS DISCONTINUOUS

G'' HAS A δ

HARMONIC OSCILLATOR

$$\left[\underbrace{\left(\frac{d}{dt}\right)^2 + \omega^2}_{\substack{\uparrow \\ \text{SPRING} \\ \text{CONST.} \\ \text{(res. freq.)}}} \right] \underbrace{q(t)}_{\substack{\downarrow \\ \text{POSITION} \\ \text{OF THING} \\ \text{THAT IS} \\ \text{OSCILLATING}}} = \underbrace{S(t)}_{\substack{\downarrow \\ \text{SOURCE} \\ \text{(DRIVING FORCE)}}$$

$q(t)$
↓

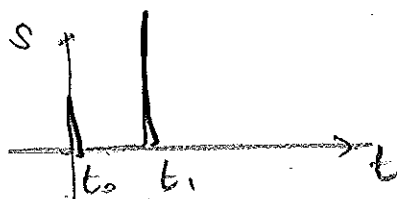
GREEN'S FUNCTION: $G(t, t_0)$ GIVES RESPONSE TO A UNIT DRIVING FORCE AT TIME t_0
 ↑
 WE'LL TAKE $= 0$
 LIKE PLUCKING A GUITAR STRING.

So the Green's function eq. is

$$\boxed{G''(t) + \omega^2 G(t) = \delta(t)}$$

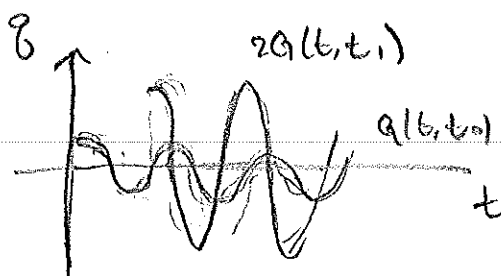
Why?

SUPPOSE WE PLUCK GUITAR STRING TWICE SECOND TIME IS TWICE AS HARD



then response is a superposition

$$q(t) = G(t, t_0) + 2G(t, t_1)$$



FOURIER: GIVEN "POSITION SPACE" $f(t)$,
 DEFINE MOMENTUM SPACE BY

$$\boxed{f(t) = \int \frac{dk}{2\pi} e^{-ikt} \tilde{f}(k)}$$

HWS } CHOICE: NEED TO PUT A $1/2\pi$ SIGN
 } NEED TO CHOOSE
 } (i "units")

just be consistent.

w/ these choices

$$\tilde{f}(k) = \int dx e^{+ikt} f(t)$$

PLUGGING IN FOURIER EXP.

$$\left(\frac{d}{dt}\right)^2 G(t) = \left(\frac{d}{dt}\right)^2 \int dk e^{-ikt} \tilde{G}(k)$$

$$= \int dk (-k^2) e^{-ikt} \tilde{G}(k)$$

$$\left[\left(\frac{d}{dt}\right)^2 + \omega^2\right] G(t) = \int dk e^{-ikt} (\omega^2 - k^2) \tilde{G}(k)$$

$$\delta(t) = \int dk e^{-ikt}$$

$$\downarrow$$

$$= 1$$

$$\Rightarrow \boxed{\tilde{G}(k) = \frac{1}{\omega^2 - k^2}}$$

done. up to integral back to t . var.

$$g(t) = \int_{-\infty}^{\infty} \frac{dk}{2\pi} \frac{-e^{-ikt}}{k^2 - \omega^2}$$

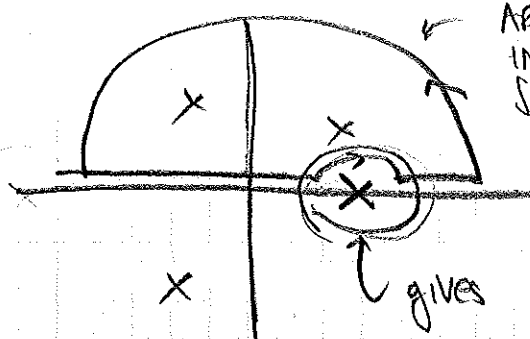
k is the var. that's promoted to a \mathbb{C} #.

SCOP: PHYSICAL SIGNIFICANCE IS CLEAR! RESONANCE

$(k+\omega)(k-\omega)$: POLES @ $\pm\omega$
 these are on the integration contour!!

LAST TIME: PRINCIPAL VALUE

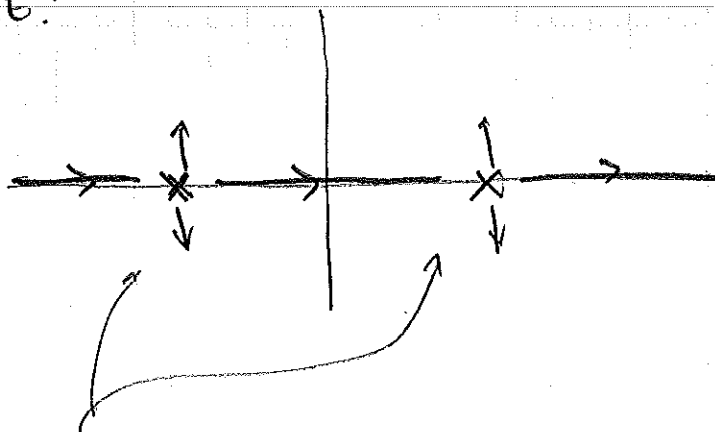
$$\frac{g(x)}{(x-x_0)}$$



ARC CHOSEN BY INTEGRAND: WANT $\int_{\text{arc}} f(z) dz \rightarrow 0$
 eg from e^{iz} factor

CAN GO OVER OR UNDER, GET SAME RESULT - choice doesn't matter!

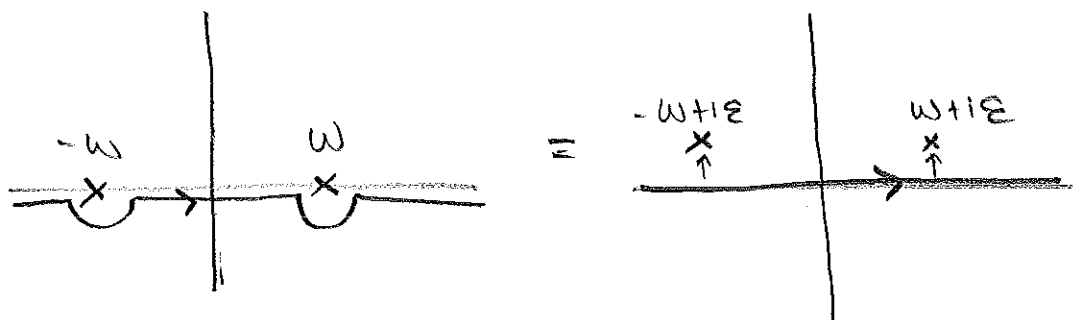
THIS TIME:



two choices: now it matters: relative to one pole, the other pole may or may not be in the contour!?

CHOICE MATTERS!

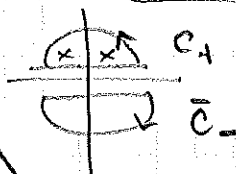
TRY: GO UNDER FOR BOTH POLES



nb: this is equivalent to moving pole an ϵ amount in the imaginary direction

- take $\epsilon \rightarrow 0$ limit @ end
- secretly: same as saying "P" asterisk defined as principal val.

How to close contour?



integrand: (pulling out 2π)

$$f(z) dk = \frac{-e^{-ikt}}{k^2 - w^2} dk \quad \boxed{e^{Rt \sin \theta}} \quad \overset{\text{PHASE}}{e^{-iRt \cos \theta}}$$

\uparrow $R^2 e^{2i\theta}$ \uparrow $iR e^{i\theta} d\theta$

WANT CONTOUR s.t. THIS GOES LIKE $e^{-\text{BIG}}$
... DEPENDS ON THE SIGN OF t !

If $t > 0$, need $\sin \theta < 0 \rightarrow \bar{C}$.

↳ BUT \bar{C} ENCIRCLES NO POLES ?!

$$Q^a(t > 0) = \int_{\bar{C}} f(z) dk = 0$$

no 'forward propagation of information'!!

WHAT ABOUT $t < 0$?

$$G(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk f(k)$$

$$= \frac{1}{2\pi} \int dz f(z) \quad \left\{ \begin{array}{l} \text{arc contributes} \\ \text{nothing} \end{array} \right.$$

$$= \frac{1}{2\pi} \times 2\pi i \times (\text{Res}_f w + \text{Res}_f -w)$$

take $\epsilon \rightarrow 0$ lim.

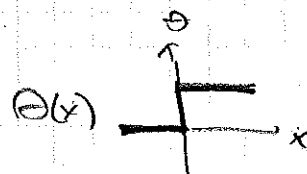
$$\frac{-e^{-i\omega t}}{2\omega}$$

$$\frac{-e^{i\omega t}}{-2\omega}$$

$$= \frac{-1}{\omega} \left[\frac{1}{2i} (e^{i\omega t} - e^{-i\omega t}) \right]$$

$\sin \omega t$

$$= \left[\frac{-1}{\omega} \sin \omega t \right] \quad \text{for } t < 0$$



OR:

$$G^a(t) = \frac{-1}{\omega} \sin(\omega t) \Theta(-t)$$

↑

↑ wrong sign!

ADVANCED PROPAGATOR (Green's function)

↑

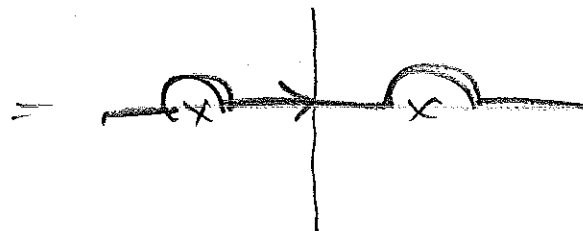
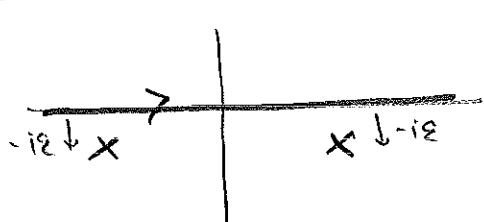
ACCAUSAL. no good.

BUT THIS IS A SOLUTION
MATHEMATICALLY

→ the dynamics are time reversal
symmetric!

INTERP: \exists SOME \sim IN ASYMPTOTIC PAST THAT
GETS CANCELED EXACTLY WHEN WE PLUCK
STRING @ $t=0$.

LET'S DO THE "RIGHT" THINGS:



PUSH BOTH POLES DOWN

equivalent to going over

CONVERGENCE STILL SET BY e^{-ikt}

$t > 0 \rightarrow C_-$

← has poles

$t < 0 \rightarrow C_+$

← no poles

$$\Rightarrow [G^r(t < 0) \equiv 0] \checkmark$$

first: $\oint_{C_-} = - \oint_{C_+}$ ← ORIENTATION OF CONTOUR

$$G^r(t > 0) = -\frac{1}{2\pi} \oint_{C_-} f(z) dz$$

$$= 2\pi i \times \frac{-1}{2\pi} (\text{Res}_f \omega + \text{Res}_f -\omega)$$

sign difference w/r-t adv.

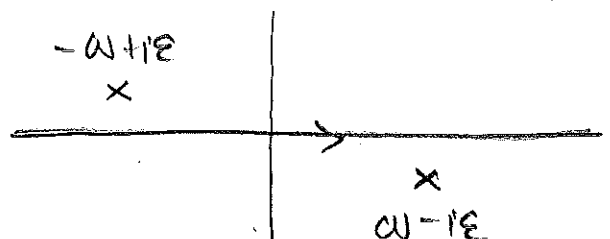
$$= \frac{1}{\omega} \sin \omega t$$

$$G^r(t > 0) = \frac{1}{\omega} \sin(\omega t) \Theta(t)$$

↑ retarded propagator

EFFECT COMES AFTER CAUSE.

other choices? eg. FEYNMAN PROPAGATOR
 ↑ you know it's important



one pole in C_+

one pole in C_-

$$G^F(t > 0) = \frac{i}{2\omega} e^{-i\omega t}$$

$$G^F(t < 0) = \frac{-i}{2\omega} e^{i\omega t} \leftarrow \text{related to } t > 0 \text{ by } \boxed{\omega \rightarrow -\omega}$$

ω
imaginary?

turns out: USEFUL IN RELATIVISTIC THEORIES

$t < 0$ looks like a NEGATIVE freq solution
moving forward in time

DIRAC: found negative energy states
 in theory of RELATIVISTIC ELECTRONS
 ENERGY \leftrightarrow FREQ.

so this was identified as

POSITIVE ENERGY position

moving fwd in time.

two
minus signs