

# Green's functions

Green's functions:

- connect integral and differential operators via boundary conditions

- connection to QM: propagator

In general

$$\mathcal{O}(x)y(x) = f(x)$$

Simple equation

$$-i \frac{dy}{dx} = f(x)$$

, if  $y(a) = y_0$

Can write  $\mathcal{O}(x)y = f(x)$ ,  $\mathcal{O}(x) = -i \frac{d}{dx}$

$$y(x) = y_0 + i \int_a^x dx' f(x')$$

notice  $(-i \frac{d}{dx} \sim p_x)$

- trivial, any integral can be considered the solution to a 1<sup>st</sup> order differential eq.

Consider  $x \in [a, b]$

$$\Rightarrow y(x) = y_0 + i \int_a^b \Theta(x-x') f(x') dx'$$

$\hookrightarrow$  cuts of integral at  $x = x'$

$$\Rightarrow y(x) = y_0 + i \int_a^x dx' f(x') = y_0 + i \int_a^b \Theta(x-x') f(x') dx'$$

We can write,  $y(x) = y_0 + K f(x)$ ,

$$K f(x) \equiv i \int_a^b \Theta(x-x') f(x') dx'$$

We say  $i\theta(x-x')$  is the kernel of the integral operator  $K$ , and when the kernel comes from the solution of a differential eq it is often referred to as the Green's function for that differential operator for the relevant boundary conditions

i.e. For  $i\frac{d}{dx}$  and  $y(a) = y_0$ ,

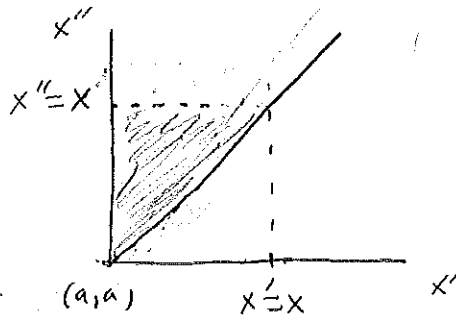
$$G_1(x, x') = i\theta(x-x')$$

$$y(x) = y_0 + \int_a^b G(x, x') f(x') dx'$$

$$\frac{d^2 y}{dx^2} = f(x), \quad y(a) = y_0, \quad y'(a) = \bar{y}_0$$

$$\Rightarrow \frac{dy}{dx} = \bar{y}_0 + \int_a^x f(x') dx'$$

$$\Rightarrow y(x) = y_0 + \bar{y}_0(x-a) + \int_a^x dx'' \int_a^{x''} dx' f(x')$$



We are integrating  $x'$  over  $[a, x'']$  then  $x''$  over  $[a, x]$

Equivalent to integrating  $x''$  over  $[x', x]$  then  $x'$  over  $[a, x]$

$$\Rightarrow y(x) = y_0 + \bar{y}_0(x-a) + \int_a^x dx' f(x') \int_{x'}^x dx''$$

$$= y_0 + \bar{y}_0(x-a) + \int_a^x (x-x') f(x') dx' = y_0 + \bar{y}_0(x-a) + \int_a^b (x-x') \theta(x-x') f(x') dx'$$

$$\Rightarrow G_2(x, x') = (x - x') \Theta(x - x')$$

$G_2$  is continuous while  $G_1$  is discontinuous at  $x = x'$

$y_0, \bar{y}_0$  correspond to the homogeneous solution,

In general

$$y(x) = \alpha + \beta x + \int_a^b (x - x') \Theta(x - x') f(x') dx'$$

$\alpha, \beta$  determined by boundary conditions

Now consider  $\Theta(x) = -i \frac{d}{dx}$  on  $[0, 1]$ ,  $y(1) = C y(0)$

In general

$$y(x) = A + i \int_0^1 \Theta(x - x') f(x') dx'$$

$$y(1) = A + i \int_0^1 f(x') dx = C y(0) = C A$$

$$\Rightarrow A = -\frac{i}{1-C} \int_0^1 f(x') dx$$

$$\begin{aligned} \Rightarrow y(x) &= -\frac{i}{1-C} \int_0^1 f(x') dx + i \int_0^1 \Theta(x - x') f(x') dx \\ &= \int_0^1 G_4(x, x') f(x') dx' \end{aligned}$$

$$G_q(x, x') = -i \left[ \frac{1}{1-c} - \theta(x-x') \right]$$

$$= \begin{cases} \frac{-i c}{1-c} & x > x' \\ \frac{-i}{1-c} & x \leq x' \end{cases}$$

$$\text{If } |c|=1 \Rightarrow c = 1/c^*$$

$$\Rightarrow G_q(x, x') = \begin{cases} \frac{-i}{c^* - 1} & x > x' \\ \frac{+i}{c - 1} & x \leq x' \end{cases}$$

$$\Rightarrow G_q(x, x') = G_q^*(x', x) \quad \text{"symmetric"}$$

$$\sim a_{ij} = a_{ji}^*$$

- Hermitian operator

- Boundary conditions can affect symmetry ( $i\theta(x-x')$  not symmetric)

$$p = -i \frac{d}{dx}$$

, symmetric in the sense  $\langle f, pg \rangle = \langle pf, g \rangle$

"inner product"

$$\langle f, g \rangle = \int dx \, w(x) f(x) g(x)$$

( $w(x) = 1$  often)

Consider

$$-i \frac{dy}{dx} = f(x, y) \quad \text{on } [0, 1]$$

$$\Rightarrow y(x) = \alpha + \int_0^1 G(x, x') f[x', y(x')] dx'$$

$$\text{If } y(0) = y_0 \Rightarrow G(x, x') = -i \Theta(x - x') \quad , \quad \alpha = y_0$$

$$\text{If } y(1) = cy(0), \Rightarrow \alpha = 0, \quad G(x, x') = -i \left[ \frac{1}{1-c} - \Theta(x - x') \right]$$

$|c| = 1$

Consider now,

$$\frac{d^2 y}{dx^2} = f(x, y) \quad y(0) = y_0, \quad y(1) = y_1$$

$$\Rightarrow y(x) = y_0 + (y_1 - y_0)x + \int_0^1 G(x, x') f[x', y(x')] dx'$$

$$G(x, x') = xx' - x' \Theta(x - x') - x \Theta(x' - x)$$

Ex

$$\frac{d^2 y}{dx^2} + \lambda y = 0 \quad , \quad y(0) = 0 = y(1)$$

$$\Rightarrow y(x) = -\lambda \int_0^1 G(x, x') y(x') dx'$$

- transformed into an integral equation for  $y(x)$



## Green's functions

- Connect differential and integral operators via boundary conditions
- propagator in QM, QFT

$$\psi(\vec{x}'', t) = \int d^3x' \underbrace{K(\vec{x}'', t; \vec{x}', t_0)}_{\text{propagator}} \psi(\vec{x}', t_0)$$

7.2  $K(\vec{x}'', t; \vec{x}', t_0) = \langle \vec{x}'' | \mathcal{U}(t, t_0) | \vec{x}' \rangle$

Consider,

$$Ly = f, \quad L: \text{linear ordinary differential operator}$$

Suppose  $L$  has complete set of eigenfunctions  $\{\phi_n(x)\}$

such that,  $\uparrow \langle \phi_n, \phi_m \rangle = \int dx \, w(x) \phi_n^*(x) \phi_m(x) \uparrow$   
 $= \delta_{nm}$

$$L\phi_n(x) = \lambda_n \phi_n(x)$$

"Basis of a Hilbert space"

$$\Rightarrow y(x) = \sum_{n=1}^{\infty} \alpha_n \phi_n(x), \quad f(x) = \sum_{n=1}^{\infty} \beta_n \phi_n(x)$$

$$Ly = L \sum_{n=1}^{\infty} \alpha_n \phi_n(x) = \sum_{n=1}^{\infty} \alpha_n \lambda_n \phi_n(x)$$

$$= f = \sum_{n=1}^{\infty} \beta_n \phi_n(x)$$

$$\Rightarrow \alpha_n = \frac{\beta_n}{\lambda_n}$$

$$\Rightarrow y(x) = \sum_{n=1}^{\infty} \frac{1}{\lambda_n} \beta_n \phi_n(x), \quad \beta_n = \underbrace{\langle \phi_n, f \rangle}_{\text{inner product of space}}$$

If  $\lambda_n = 0 \Rightarrow \exists \beta_n \neq 0 \Rightarrow$  non unique sol, assume  $\lambda_n \neq 0$

$$\begin{aligned}
y(x) &= \sum_{n=1}^{\infty} \frac{1}{\lambda_n} \phi_n(x) \langle \phi_n, f \rangle \\
&= \sum_{n=1}^{\infty} \frac{1}{\lambda_n} \phi_n(x) \int dx' \phi_n^*(x') f(x') \\
&= \int \sum_{n=1}^{\infty} \frac{\phi_n(x) \phi_n^*(x')}{\lambda_n} f(x') dx' \\
&= \int G(x, x') f(x') dx'
\end{aligned}$$

$$\begin{aligned}
\Rightarrow G(x, x') &= \sum_{n=1}^{\infty} \frac{\phi_n(x) \phi_n^*(x')}{\lambda_n} \\
&= G(x', x)^*
\end{aligned}$$

Ex:

$$L = \frac{d^2}{dx^2} \quad \text{on interval } [0, 1]$$

$$\Rightarrow \phi_n(x) = \sqrt{2} \sin(n\pi x) \quad \text{"particle in a box"}$$

$$\Rightarrow G(x, x') = -\frac{2}{\pi^2} \sum_{n=1}^{\infty} \frac{\sin(n\pi x) \sin(n\pi x')}{n^2}$$

Consider,

$$L G(x, x') = \sum_{n=1}^{\infty} L \frac{\phi_n(x) \phi_n^*(x')}{\lambda_n} = \sum_{n=1}^{\infty} \phi_n(x) \phi_n^*(x')$$

$$\Rightarrow \int dx' f(x') \sum_{n=1}^{\infty} \phi_n(x) \phi_n^*(x') = \sum_{n=1}^{\infty} \phi_n(x) \int dx' \phi_n^*(x') f(x')$$

$$= \sum_{n=1}^{\infty} \beta_n \phi_n(x)$$

$$= f(x)$$

$$\Rightarrow \sum_{n=1}^{\infty} \phi_n(x) \phi_n^*(x') = \delta(x - x') \quad , \quad \omega(x) = 1$$



$$\Rightarrow L G(x, x') = \delta(x - x')$$

$$y = L^{-1} f = K f$$

$$, L K = \delta(x - x')$$

Ex

$$L = -i \frac{d}{dx}$$

$$\Rightarrow -i \frac{d}{dx} G(x, x') = \delta(x - x')$$

$$\Rightarrow G(x, x') = i \int_{-\infty}^x dx'' \delta(x' - x'')$$

$$= i \Theta(x - x')$$

$$-i \frac{dy}{dx} = f(x), y(a) = y_0$$

$$\Rightarrow y(x) = y_0 + i \int_a^x dx' f(x')$$

$$= y_0 + i \int_a^b dx' \Theta(x - x') f(x')$$

$$\Rightarrow G(x, x') = i \Theta(x - x')$$

Example

$$\ddot{x}(t) + 2\gamma \dot{x}(t) + \omega_0^2 x(t) = F(t)$$

$$\left( \frac{d^2}{dt^2} + 2\gamma \frac{d}{dt} + \omega_0^2 \right) x(t) = F(t)$$

$$f(t) = \int_{-\infty}^{\infty} dk \hat{f}(k) e^{ikt}$$

$$dk = \frac{dK}{2\pi}$$

$$\Rightarrow \int dk \left( \frac{d^2}{dt^2} + 2\gamma \frac{d}{dt} + \omega_0^2 \right) e^{ikt} \hat{x}(k) = \int dk \hat{F}(k)$$

$$= \int dk \left( -k^2 + 2i\gamma k + \omega_0^2 \right) e^{ikt} \hat{x}(k)$$

$$\Rightarrow \hat{x}(k) = \frac{\hat{F}(k)}{-k^2 + 2i\gamma k + \omega_0^2}$$

$$\hat{F}(k) = \int_{-\infty}^{\infty} dt F(t) e^{-ikt}$$

$$\Rightarrow x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk \frac{\tilde{F}(k)}{-k^2 + 2i\gamma k + \omega_0^2} e^{ik t}$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} dt' \int_{-\infty}^{\infty} dk \frac{F(t') e^{ik(t-t')}}{-k^2 + 2i\gamma k + \omega_0^2}$$

For completeness,

$$x(t) = A x_1(t) + B x_2(t) + \int G(t, t') F(t') dt'$$

$$G(t, t') = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk \frac{e^{ik(t-t')}}{-k^2 + 2i\gamma k + \omega_0^2} \leftarrow \text{contour integral}$$

Can solve homogeneous eq to find

$$x_1(t) = e^{-\gamma t} \sin \Omega t$$

$$x_2(t) = e^{-\gamma t} \cos \Omega t$$

$$\Omega = \sqrt{\omega_0^2 - \gamma^2}$$

\* Flip will go over how to solve contour integrals

$$G(t, t') = \theta(t - t') \frac{e^{-\gamma(t-t')} \sin[\Omega(t-t')]}{\Omega}$$

$$x(t) = A x_1(t) + B x_2(t) + \int_{t_0}^t dt' \frac{e^{-\gamma(t-t')}}{\Omega} \sin[\Omega(t-t')] F(t')$$

Let  $F(t) = F_0 e^{-\alpha t}$ ,  $t_0 = 0$

Equilibrium at  $t=0 \Rightarrow A=B=0$

$$x(t) = \frac{F_0}{\Omega} e^{-\gamma t} \int_0^t dt' \sin[\Omega(t-t')] e^{-(\alpha-\gamma)t'}$$

$$= \frac{F_0}{\sqrt{\omega_0^2 - \gamma^2}} \frac{\sin[\sqrt{\omega_0^2 - \gamma^2} t - \delta]}{\sqrt{\omega_0^2 + \alpha^2 - 2\alpha\gamma}} e^{-\gamma t} + \frac{F_0}{\omega_0^2 + \alpha^2 - 2\alpha\gamma} e^{-\alpha t}$$

$$\tan \delta = \sqrt{\omega_0^2 - \gamma^2} / (\alpha - \gamma)$$

For  $\gamma=0$ ,  $t \rightarrow \infty$

$$x(t) = \frac{F_0}{\omega_0} \frac{\sin(\omega_0 t - \delta)}{\sqrt{\omega_0^2 + \alpha^2}}$$

$$E = \frac{1}{2} \dot{x}^2 + \frac{1}{2} \omega_0^2 x^2 = \frac{F_0^2}{2(\omega_0^2 + \alpha^2)}$$

$G(t, t')$  closely resembles homogeneous solution,  
It is a linear combination of  $x_1$  and  $x_2$

$$G(t, t') = A(t') e^{-\gamma t} \sin \Omega t + B(t') e^{-\gamma t} \cos \Omega t$$

$$A(t') = \frac{e^{\gamma t'}}{\Omega} \cos \Omega t' \quad B(t') = -\frac{e^{\gamma t'} \sin \Omega t'}{\Omega}$$

obvious because

$$\left( \frac{d^2}{dt^2} + 2\gamma \frac{d}{dt} + \omega_0^2 \right) G(t, t') = \delta(t - t')$$

$\Rightarrow$  for  $t < t'$  and  $t > t'$

$G(t, t')$  satisfies the same as equation as  $X(t)$

$\Rightarrow G(t, t')$  is a linear combination of  $x_1(t)$ ,  $x_2(t)$   
w/ coefficients dependant on  $t'$

In general

$$L \equiv f_0(t) \frac{d^2}{dt^2} + f_1(t) \frac{d}{dt} + f_2(t)$$

$x_1(t)$  and  $x_2(t)$  linearly independent solutions to

$$LX(t) = 0 \quad \Rightarrow \quad LG(t, t') = \delta(t - t')$$