

HOMEWORK 4: Operators on Function Space

COURSE: Physics 231, *Methods of Theoretical Physics* (2017)

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DUE BY: Friday, October 27

Reading: For this week, please read your favorite references to brush up on complex analysis. Our main goal is to understand contour integrals, though we will take a detour and see a nice application to causality. Please review the following key topics:

- The Cauchy–Riemann equations
- The meaning of an analytic function
- Laurent series
- Cauchy’s theorem
- The residue theorem.

Some possible references include: Byron & Fuller chapter 6, Boas chapter 14, Matthews & Walker chapter 3.3 and appendix A, Stone & Goldbart chapter 17, Appel chapter 4.1–4.6, Cahill chapter 5.1–5.14.

1 Sturm–Liouville Operator

This problem follows Lecture 9. Consider the differential operator

$$L = p_2(x) \left(\frac{d}{dx} \right)^2 + p_1(x) \frac{d}{dx} + p_0(x). \quad (1.1)$$

Assume $w(x) = 1$ and some domain $x \in [a, b]$.

REFERENCE: Refer to Stone & Goldbart chapter 4 for a discussion of formal versus concrete (what we called ‘proper’) linear differential operators. We use a slightly different notation, but the question of a formal adjoint is explored in 4.2 with the Sturm–Liouville operator identified in equation (4.28). A more general discussion that follows more of the spirit of the lecture is in Matthews & Walker Chapter 9-2.

1.1 Formal adjoint

Using integration by parts, show that the formal adjoint of this operator is

$$L^\dagger = \left(\frac{d}{dx} \right)^2 p_2(x) - \left(\frac{d}{dx} \right) p_1(x) + p_0(x), \quad (1.2)$$

where we use the notation

$$\left[\left(\frac{d}{dx} \right)^n p_n(x) \right] f(x) \equiv \left(\frac{d}{dx} \right)^n [p_n(x) f(x)] . \quad (1.3)$$

1.2 Self-Adjoint Case

Show that in order for L to be formally self-adjoint, that is $L = L^\dagger$ (without considering boundary conditions), then the following conditions must hold:

$$p_1(x) = 2p_2'(x) - p_1(x) \quad (1.4)$$

$$p_0(x) = p_x''(x) - p_1(x) + p_0(x) . \quad (1.5)$$

Thus L is self-adjoint (Hermitian) if $p_1(x) = p_2'(x)$. L can be written succinctly as:

$$L = \frac{d}{dx} \left(p_2(x) \frac{d}{dx} \right) + p_0(x) . \quad (1.6)$$

2 A Green's function by completeness

Three formulaic ways of solving for Green's functions are:

1. Fourier transforming to turn the differential operator into an algebraic one, and then doing a contour integral to go back to position space.
2. Solving the homogeneous equation and patching together two solutions over the δ -function.
3. Using the completeness of eigenfunctions.

We will spend most of this course highlighting method 1. Next week you will use method 2—it will be familiar from electrodynamics. This week, we will use the third method, which is based on the eigenfunction completeness relation,

$$\sum_n e_n^*(y) e_n(x) = \delta(x - y) . \quad (2.1)$$

Thus if you know the eigenvalues of the linear differential operator, L_x , then you can use this relation to hack together the Green's function G because $L_x G(x, y) = \delta(x - y)$. The subscript x in L_x means that the operator is a function of x and derivative operators with respect to x .

REFERENCES: This method is discussed in Matthews & Walker chapter 9-4, Stone & Goldbart chapter 5.4, Cahill chapter 6.38, and Byron & Fuller chapter 7.2.

In this problem, we consider a second order differential equation acting on a state $\psi(x)$ with a source $s(x)$.

$$\left(\frac{d^2}{dx^2} + k^2 \right) \psi(x) = s(x) .$$

We consider the domain $x \in [0, 1]$ with boundary conditions $\psi(0) = \psi(1) = 0$.

2.1 Eigenfunctions

What are the eigenfunctions $e_n(x)$ and the corresponding eigenvalues λ_n ? Recall that an eigenfunction of a linear operator L_x satisfies $L_x e_n(x) = \lambda_n e_n(x)$.

HINT: The eigenfunctions of d^2/dx^2 are $\psi_n(x) = \sqrt{2} \sin(n\pi x)$ with eigenvalues $n^2\pi^2$. What changes when the operator L_x includes an additive constant?

2.2 Completeness

Write down the Green's function $G(x, y)$ using the completeness trick with respect to eigenfunctions e_n with eigenvalues λ_n :

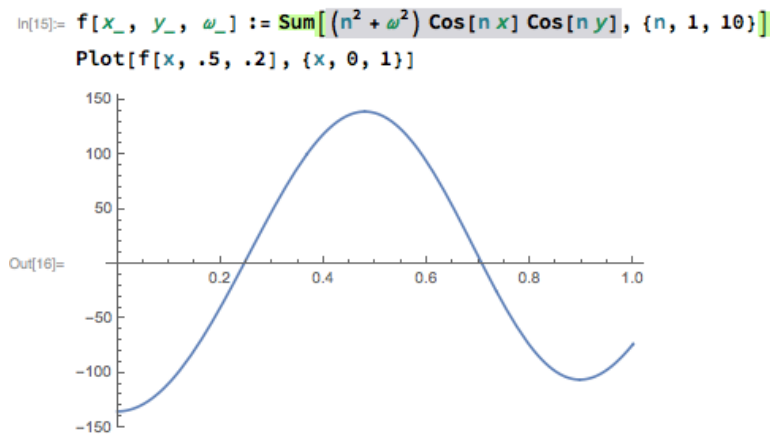
$$G(x, y) = \sum_n \frac{e_n^*(y)e_n(x)}{\lambda_n} . \quad (2.2)$$

HINT: Once you did the previous part, this step takes about twenty seconds depending on how quickly you write.

2.3 Hands on

You have a free license for several types of scientific software through UCR through MySoftware¹. This problem requires some plotting. I suggest using *Mathematica*².

Plot the solution to the Green's function $G(x, y)$ for $y = 0.5$ and $k = 0.2$, summing from $n = 1$ to $n = 10$. Here's the general way to format it in *Mathematica*:



The highlighted piece and the example plot are *completely wrong*! Fill it in with the correct $G(x, y)$ and plot it. It may help to use the ‘Basic Math Input’ palette if you’re unfamiliar with this. If this task is very painful, please ask a friend. If this is still very painful, then you may use any other plotting program that you wish.

Does this shape make sense? What do you expect will happen as n becomes larger? Try it for $n = 100$. Explore what happens as you change k and y . (Make sure y is in the domain of the function space!) Take a moment to have a nice warm cup of tea, enjoy the agreeable weather outside, and meditate upon why this shape of the Green's function makes sense and the kinds of physical problem where Green's function may be relevant. You do not have to write up these meditations, but I do strongly suggest the tea and outdoors³.

¹<http://cnc.ucr.edu/mysoftware/>

²It sucks that this is not open source software. The closest open source alternative is SciPy with a Jupyter notebook. You can use this instead if you prefer.

³If you are looking for nice tea to bring home, I suggest the *British Emporium* in Canyon Crest.

Extra Credit

These problems are not graded and are for your edification. You are strongly encouraged to explore and discuss these topics, especially if they are in a field of interest to you.

A Dimensional Analysis: Nature is Approximately Second Order

Most differential equations in physics are second order (in derivatives) or less. We observed that this can be understood from the mantra that *physics is local*; higher derivative terms encode non-locality. Here we give a toy example for why nature appears to stop at second order.

Often our favorite differential equations come from a least action principle, $\delta S = 0$. In natural units, the action is dimensionless, $[S] = 0$. We'll take the example of the dynamics of a scalar field, $h(x)$. This is a function defined over all of space and time. Quantum excitations of this field correspond to particles.

- (a) The action is written as an integral of a Lagrangian density, \mathcal{L} over spacetime: $S = \int d^4x \mathcal{L}$. To match onto something familiar, this is just saying that $S = \int dt L$ with $L = \int d^3\mathbf{x} \mathcal{L}$. In natural units, what is the mass dimension of \mathcal{L} ?
- (b) The Lagrangian density is a function of the field $h(x)$ and derivatives, $\partial_\mu = \partial/\partial x^\mu$. A dynamical particle necessarily has at least the following term in the Lagrangian density,

$$\mathcal{L} \supset \frac{1}{2} h(x) \partial_\mu \partial^\mu h(x) = \frac{1}{2} h(x) (\partial_t^2 - \partial_x^2 - \partial_y^2 - \partial_z^2) h(x).$$

This corresponds to the kinetic energy term in the Lagrangian. In natural units, what is the mass dimension of a derivative, ∂_μ ? (Does it matter whether it's a time derivative or a space derivative?) What is the mass dimension of the field $h(x)$? Observe that Lorentz invariance requires that the space derivatives come along with the time derivatives.

- (c) We can ignore terms in the Lagrangian density with fewer than two powers⁴ of $h(x)$. Lorentz invariance says that you cannot have a term with an odd number of derivatives, so let's imagine what would happen to a term in \mathcal{L} that had two powers of $h(x)$ and four derivatives, for example:

$$\mathcal{L} \supset \Lambda^n h(x) \partial^4 h(x).$$

We've included an overall prefactor, Λ^n . Assume that Λ has mass dimension of one. Given your answers to the previous question, what is the value of n for which the $\mathcal{O}(\partial^4)$ term has the correct dimension for a Lagrangian density term?

- (d) The answer to the above question is $n = -2$. One of the key ideas in theoretical physics is the *renormalization group* (related ideas: effective theories, anomalous scaling.). I hope you have the chance to study this in the many ways it shows up in different fields, but one of the punchlines of this idea is that the prefactor Λ^n in operators like the one above should be identified with a *cutoff* for the theory. That is: this is an energy scale (inverse length scale) at which the model you're using to describe the system breaks down because you're not including

⁴Extra credit: think about such terms would mean for the differential equation that comes from $\delta S = 0$.

micro-physics. For example: if you had an action for the behavior of water waves where the water is treated as a continuous medium, then the theory breaks down at the scale where individual water molecules are resolved. Is Λ a ‘big’ number or a ‘small’ number? Specify in what sense you mean this—what units are you using, and to what are you comparing? HINT: As you know from Fourier analysis, derivatives bring down powers of momentum. So a more concrete question is: I have a theory of the scalar field $h(x)$ which I know is valid at some energy scale, E —if all my momenta are of this characteristic energy scale, $\partial \rightarrow k \approx E$, then is the coefficient of the $\mathcal{O}(\partial^4)$ term in (c) bigger or smaller than the ordinary kinetic term in (b)?

B Self-adjointness and Variational Principles

Kuntal brought up a great point on Wednesday: there is a relation between the self-adjointness of a differential operator appearing in an equation of motion whether or not that equation could come from a variational principle. There is some discussion online⁵, but I’ve yet to find a satisfactory physically intuitive explanation (to the extent to which the statement is even true). However, I suspect such an explanation relates reversibility and conservation of probability in quantum mechanics. If you have thoughts on this, I would be happy to hear them.

⁵<https://physics.stackexchange.com/q/291697>