# NE585 NUCLEAR FUEL CYCLES Interaction of radiation with matter

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#### **Learning objectives**

Explain the different ways radiation interacts with matter

Apply the concept of cross sections to neutron interactions with matter

Derive neutron scattering relationships

Book – Chapter 3

#### **Learning nodes**

#### Gamma ray interactions

Photoelectric effect Pair production Compton effect

#### **Fission**

#### **Neutron interactions**

Elastic scattering
Lethargy
Inelastic scattering
Radiative capture
Charged particle emission
Bremsstrahlung radiation

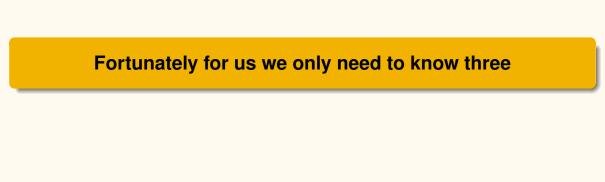
**Cross sections** 

**Neutron density** 

**Neutron fluence** 

**Neutron flux** 







### The photoelectric effect occurs when a high energy photon knocks out an electron

Interacts with entire atom

Electron ejected

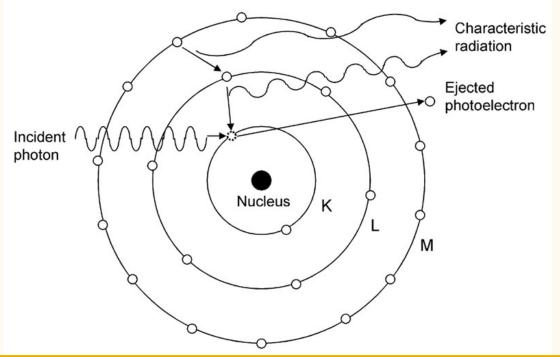
Leaves a positively charged ion

Usually an x ray emission

Atom recoil carries little Kinetic Energy (KE)

Highly probable with heavier atoms

Why?





# For pair production, a high energy photon creates a negatron and positron

Only occurs with photon above 1.022 MeV

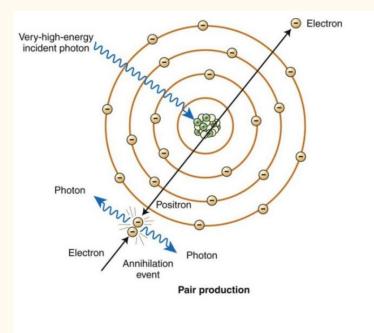
And near the nucleus

Why?

Then they react with other negatrons/positrons

Give off characteristic 0.511 MeV photons

Only in vicinity of Coulomb field





### The Compton effect is elastic scattering of a photon by an electron

Energy and momentum are conserved

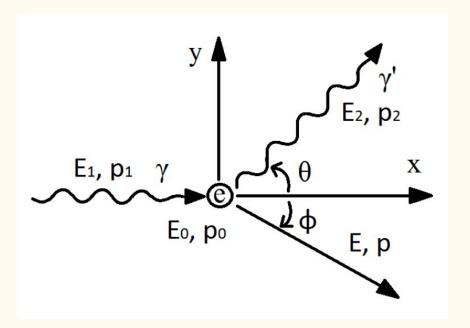
All angles of scattering possible

$$E_2 = \frac{E_1 E_{\epsilon}}{E_1 (1 - \cos \theta) + E_{\epsilon}} \tag{1}$$

Should be able to derive maximum energy from this equation

Why would the Compton effect cause problems with shielding? 1

<sup>(</sup>Talk about NE class experiment at WPI)





# What is the energy of fission from thorium and uranium fuel cycles?

$${}_{0}^{1}n + {}_{92}^{235}U \rightarrow {}_{36}^{92}Kr + {}_{56}^{141}Ba + {}_{0}^{1}n + {}_{0}^{1}n + {}_{0}^{1}n$$
 (2)

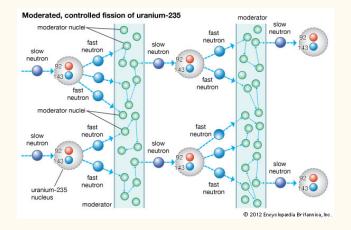
$${}_{0}^{1}n + {}_{92}^{233}U \rightarrow {}_{36}^{92}Kr + {}_{56}^{141}Ba + {}_{0}^{1}n$$
 (3)

$$Q = 173.3 \; MeV, \; 185.4 \; MeV$$
 (4)

For comparison, burning a carbon atom releases 4 eV

Some super heavy nuclei fission spontaneously

### Fission is the splitting of atoms, which releases a lot of energy





## What is the energy of fission from thorium and uranium fuel cycles?

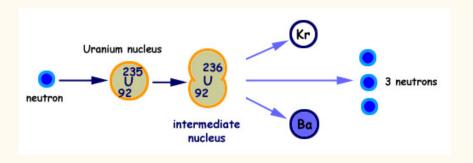
Why?

$$^{1}_{0}n + ^{233}_{92}U \rightarrow ^{236}_{92}U^{*} \rightarrow$$
 (5)

Sometimes the compound nucleus is left in an excited state, before the reaction proceeds, but this can be a very short time

What happens is highly dependent on neutron energy/temp

So you do not always write the compound nucleus in the reaction expression



### There are several ways neutrons interact with nuclei

- (1) Elastic scattering (probably most important)
- (2) Inelastic scattering
- (3) Radiative capture/neutron absorption
- (4) Charged-particle reactions
- (5) Neutron producing reactions

### What is the energy of fission from thorium and uranium fuel cycles?

$${}_{0}^{1}n + {}_{92}^{233}U \rightarrow {}_{92}^{236}U^{*} \rightarrow {}_{92}^{233}U + {}_{0}^{1}n$$
 (6)

$$^{1}_{0}n + ^{233}_{92}U \rightarrow ^{236}_{92}U^{*} \rightarrow ^{233}_{92}U + ^{1}_{0}n'$$

$$^{1}_{0}n + ^{233}_{92}U \rightarrow ^{236}_{92}U^{*} \rightarrow ^{233}_{92}U + \gamma$$

$$^{1}_{0}n + ^{233}_{92}U \rightarrow ^{236}_{92}U^{*} \rightarrow ^{233}_{92}U + ^{4}_{2}\alpha$$

$$^{1}_{0}n + ^{233}_{92}U \rightarrow ^{236}_{92}U^{*} \rightarrow ^{233}_{92}U + 2^{1}_{0}n$$

(10)

(7)

(8)

(9)



### With elastic scattering, the neutron strikes the nucleus, and they scatter

If you've ever played pool (billiards), it's that sort of

But the neutron is a lot smaller than the nucleus

Kinetic energy and momentum is conserved with nucleus and neutron

Nucleus can be at rest or moving, but neutron velocity is really faster

So basically the nucleus is 'at rest'

### With elastic scattering, the neutron strikes the nucleus, and they scatter

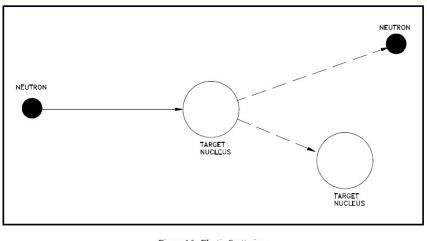
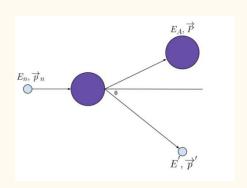


Figure 16 Elastic Scattering

# The neutron energy after the collision can be obtained by conservation principles



$$E = E' + E_A \tag{11}$$

$$\vec{p} = \vec{p}' + \vec{P} \tag{12}$$

$$E' = \frac{E_n}{(A+1)^2} \cdot [\cos\theta \sqrt{A^2 - \sin^2\theta}]^2$$
 (13)

$$\alpha \equiv (\frac{A-1}{A+1})^2 \tag{14}$$

(p58-9) - but in any edition



### The concept of lethargy was invented for neutron moderation

They aren't good at naming things

$$u \equiv ln \frac{E_M}{E} \tag{15}$$

 $E_M$  is the highest neutron energy in the system

What is happening mathematically?

#### The concept of lethargy was invented for neutron moderation

So, lethargy increases when neutron slows down

Log scale is used because neutrons have huge energy range

### What we really want is average log energy loss per elastic scatter/collision

$$\xi \equiv 1 + \frac{\alpha}{1 - \alpha} \ln \alpha \tag{16}$$

$$\alpha \equiv (\frac{A-1}{A+1})^2 \tag{17}$$

Average change in lethargy

What does this tell us?

In terms of thermal reactor design?

### Because elastic scattering can be applied to find out how a neutron will slow down

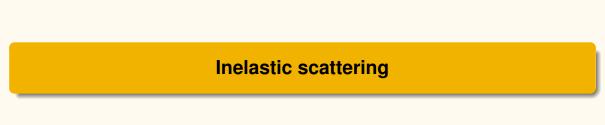
$$\overline{n}_{COL} \equiv \frac{1}{\xi} ln \frac{E_1}{E_2} \tag{18}$$

What is the physical interpretation?

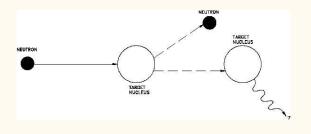
Moderation = getting high energy neutrons to low energy

This means from fast moving to slow

Why? This is a very fundamental reactor design constraint



### With inelastic scattering, it is the same as elastic, but the nucleus is left in excited state



Energy is lower for heavier nuclei

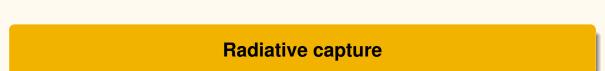
Kinetic energy and momentum not conserved

Compound nucleus emits lower energy neutron

Then the nucleus emits gamma rays

Occurs above a certain neutron energy for each nucleus

A neutron with same energy may not scatter with different nuclei



### Because elastic scattering can be applied to find out how a neutron will slow down

This is important with reactor design (fission and fusion)

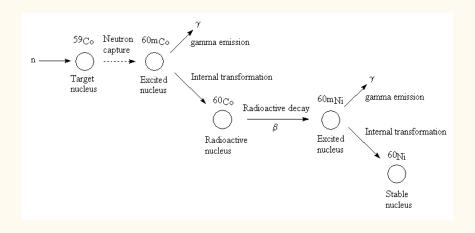
Take a guess why?

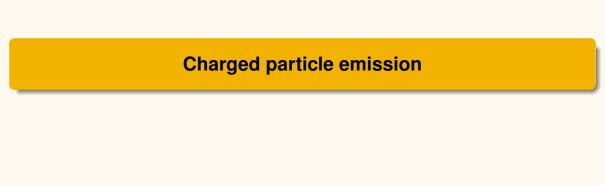
We also call this neutron absorption in this context

Usually a low energy neutron reaction (why?)

Depends on neutron energy and material

Other applications include NAA, isotope production





# The nucleus can also absorb a neutron and emit a charged particle

This can be used to design neutron detectors

Part of fusion reactions

Boron is used as neutron 'poison' and can absorb low energy

Materials testing for fuel cladding (fusion and fission)

Charged particle production could be adverse

Neutron depth profiling; surface of semiconductors, etc.

$$^{1}_{0}n + ^{10}_{5}B \rightarrow 11_{5}B^{*} \rightarrow ^{7}_{3}Li + \alpha + \gamma$$



# A consequence of charged particle reactions is the Bremsstrahlung radiation

Deceleration of the charged particle through matter

A photon is emitted from interaction with the electron cloud or nucleus with the charged particle

The moving particle loses energy and this is in the form of a photon

aka 'braking radiation'

Continuous radiation distribution

Can be used to identify galaxy clusters

## **Cross sections**

# A cross section tells us basically the probability of an interaction (with something)

Neutrons for us

Cross sections for other events like Compton, etc.

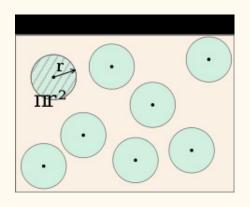
People dedicate careers to establishing precise cross sections (and good for them)

Cross sections are given as a hypothetical area around the target nucleus (big area, high probability)

$$1 b = 10^{-24} cm^2 (20)$$

Dependent on speed/energy/temperature

#### Try to think of this as the likelihood of interaction

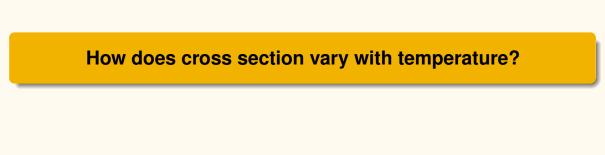


Since there is a lot of space in between nuclei, some neutrons just fly by

Radius of a nucleus  $\sim 10^{-12}~cm^2$ , so cross section is  $\sim 10^{-24}~cm^2$ 

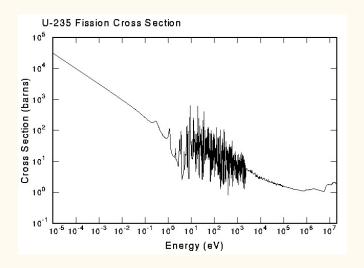
Which is where they derived barns

Probability of interaction is the ratio of the total surface area of the atoms to the total area of the medium

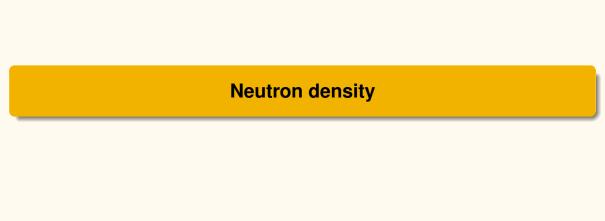




### Fission cross sections have this general trend



If fission neutrons start at 2 MeV, where would you want them for a reactor?



## Neutron density in the thermal region follows Maxwellian distribution

As fast neutrons slow down, they eventually come into thermal equilibrium with the thermal motion of the atoms in the medium through which they are moving

Also call the  $\frac{1}{v}$  region

Resonances occur if the compound nucleus is produced in one of its characteristic excited states

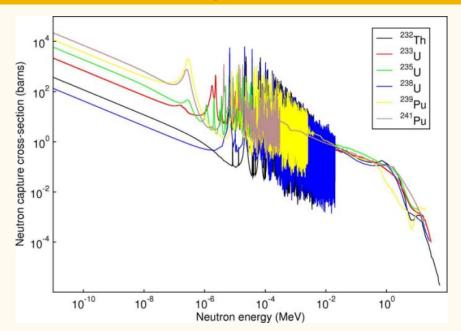
And the fast end is a smooth region

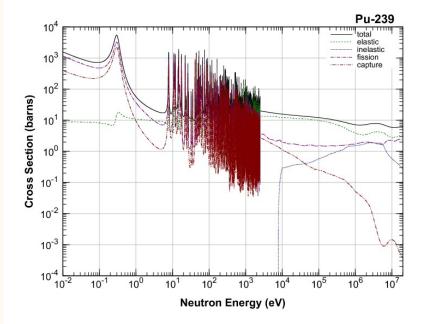
## What we really want is average log energy loss per elastic scatter/collision

$$n(E) = n_0 \frac{2\pi\sqrt{E}}{(\pi kT)^{\frac{3}{2}}} \cdot e^{-\frac{E}{kT}}$$
 (21)

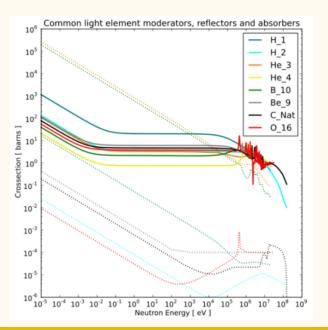
$$n(v) = n_0 \frac{4\pi v^2}{(\frac{2\pi kT}{m})^{\frac{3}{2}}} \cdot e^{-\frac{mv^2}{2kT}}$$
 (22)

## How does this affect reactor operation?





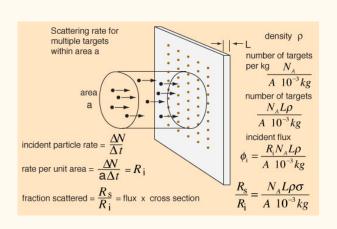
### How do you choose materials?



solid = scatter dotted = absorption



# How many neutrons/s that strike a target actually interact with target nuclei?



$$I\left[\frac{n}{cm^2 \cdot s}\right] = \left[\frac{n}{cm^3}\right] \cdot v\left[\frac{cm}{s}\right]$$
 (23)

$$\left[\frac{n}{cm^3}\right] \cdot v\left[\frac{cm}{s}\right] \cdot a[cm^2] = \frac{n}{s} \qquad (24)$$

$$nuclei_{tar} \equiv N[\frac{nuclei}{cm^3}] \cdot a[cm^2] \cdot L[cm]$$
 (25)

$$\frac{\textit{interactions}}{s} \equiv (\sigma I)(\textit{NaL}) \tag{26}$$

The cross section is the probability the neutron interacts with the nucleus

Section 3.2, example 3.1

#### Then we can define the mean free path

First, macroscopic cross section is -

$$\Sigma[cm^{-1}]N\sigma = \frac{\rho N_A}{A} \tag{27}$$

Then mean free path is -

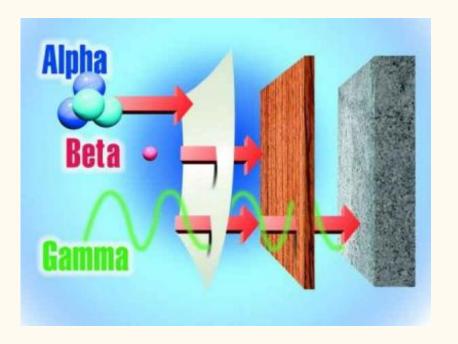
$$\lambda \equiv \Sigma^{-1} [cm] \tag{28}$$

Expected value of distance between interactions

Then the interaction rate is -

$$R_X = \phi \Sigma_X \tag{29}$$





Which flux is larger?
Point source
Isotropic

### Uncollided flux decreases exponentially with distance

$$-d\phi = \Sigma_t \phi dx \tag{30}$$

$$-\frac{d\phi}{\phi} = \Sigma_t dx \tag{31}$$

Probability a neutron interacts in dx

$$\phi(\mathbf{x}) = \phi_0 \mathbf{e}^{-\Sigma_t \mathbf{x}} \tag{32}$$

Probability a neutron does not interact in x

### It's the same for gamma rays

Shielding is material density dependent

$$\phi(\mathbf{x}) = \phi_0 \mathbf{e}^{-\mu \mathbf{x}} \tag{33}$$

Look up mass attenuation coefficient

$$\frac{\mu}{\rho} \left[ \frac{cm^2}{g} \right] \tag{34}$$

Then just multiply mass attentuation by material density

