

# Nonlinear Dynamics and Chaos I

## Problem Set 3 - Questions

### Question 1

Consider a ball of mass  $m$  that slides on a rotating hoop (see Fig. 1).

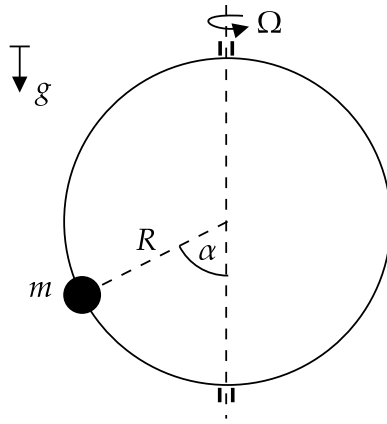


Figure 1: Mass on a hoop

The angular velocity of the hoop is  $\Omega$ , the viscous friction coefficient between the hoop and the ball is  $b$ , and the constant of gravity is  $g$ . The equation of motion for the sliding ball is given by

$$mR^2\ddot{\alpha} + bR^2\dot{\alpha} + mR^2(g/R - \Omega^2 \cos(\alpha)) \sin(\alpha) = 0.$$

- Plot the location of equilibria of the ball as a function of the non-dimensionalized rotation parameter  $\nu = R\Omega^2/g$ .
- Using linearization, determine the stability type of the different equilibrium branches on the plot. Identify the critical angular velocity at which a bifurcation of equilibria occurs.

### Question 2

Consider a discrete dynamical system given by the iterated mapping

$$x_{k+1} = f(x_k), \quad f: \mathbb{R}^n \rightarrow \mathbb{R}^n, \quad x_k \in \mathbb{R}^n.$$

Assume that  $x = p$  is a fixed point for the mapping, i.e.,  $p = f(p)$ .

- Derive a linearized mapping of the form

$$y_{k+1} = Ay_k \tag{1}$$

to describe the discrete dynamics in the vicinity of the fixed point.

- Assume that  $A$  has eigenvalues  $\lambda_1, \dots, \lambda_n \in \mathbb{C}$  with corresponding  $n$  linearly independent eigenvectors  $s_1, \dots, s_n \in \mathbb{C}^n$ . Show that the general solution of (1) is of the form

$$y_k = c_1\varphi_1(k) + \dots + c_n\varphi_n(k), \quad \varphi_i(k) = \lambda_i^k s_i. \tag{2}$$

- Formulate a definition of stability, asymptotic stability, and instability for the  $y = 0$  fixed point of (1).
- Using (2), find a sufficient and necessary condition for the asymptotic stability you have defined in (c).

### Question 3

The first three modes of a convecing fluid motion in a two-dimensional layer heated from below are described by the famous *Lorenz equations*

$$\begin{aligned}\dot{x} &= a(y - x), \\ \dot{y} &= bx - y - xz, \\ \dot{z} &= xy - cz.\end{aligned}$$

Here  $a > 0$  denotes the Prandtl number,  $b > 0$  is the Rayleigh number, and  $c > 0$  is the aspect ratio. Lorenz's original assumption is that  $a > 1 + c$ .

The above equations describe the evolution of the amplitudes of one velocity mode and two temperature modes. As a paradigm of chaotic dynamics, this system has inspired much of the development of the modern geometric theory of dynamical systems.

- (a) Complicated dynamics arise when all possible equilibria of the system become unstable, and hence solutions cannot settle down to any simple steady state. Show that this is the case when

$$b > \frac{a(3 + a + c)}{a - c - 1}$$

*Note:* Note: A negative sign for one of the Routh-Hurwitz determinants actually implies instability, not just the lack of asymptotic stability.

- (b) Solve the Lorenz equations numerically for  $a = 10$ ,  $b = 28$ , and  $c = 8/3$ , choosing an initial condition close to  $x = y = z = 0$ . Plot the trajectory in three dimensions to show that it converges to a complicated surface, a *chaotic attractor*.

### Question 4

Recall from Question 1 that a ball of mass  $m$  sliding on a hoop rotating with angular velocity  $\Omega$  satisfies the differential equation

$$mR^2\ddot{\alpha} + bR^2\dot{\alpha} + mR^2(g/R - \Omega^2 \cos(\alpha)) \sin(\alpha) = 0 \quad (3)$$

if there is no friction between the hoop and the mass. Assume that the parameter values are such that the lower equilibrium position of the ball is stable in linear approximation.

- (a) Show that in this case, the equilibrium is also nonlinearly stable.

*Hint:* Note that system (3) is not conservative: an external torque is required to maintain the constant rotation of the hoop. As a result, the total mechanical energy of the system is not expected to work as a Lyapunov function. To find another candidate for a Lyapunov function, find a quantity that *is* conserved along trajectories. Such a quantity can be found by multiplying equation (3) by  $\dot{\alpha}$  and integrating in time.

- (b) Prove that the equilibrium *cannot* be asymptotically stable for the nonlinear system.

*Hint:* use the Lyapunov function you have found in (a).

### Question 5

Consider the damped pendulum equation

$$\ddot{x} + c\dot{x} + \sin(x) = 0. \quad (4)$$

- (a) Using the energy of the pendulum as a Lyapunov function, can we conclude the nonlinear asymptotic stability of the  $x = 0$  equilibrium? Give detailed reasoning why.
- (b) A theorem due to Krasovski states the following: Assume that  $x = 0$  is a fixed point for the  $n$ -dimensional dynamical system  $\dot{x} = f(x)$ . Assume that there exists a smooth scalar function  $V(x)$  such that
- $V(x)$  is positive definite on an open neighborhood  $U$  of  $x = 0$
  - $\dot{V}$  is negative semi-definite on the same neighborhood
  - the only trajectory lying *completely* in the set  $S = \{x \in U : \dot{V} = 0\}$  is the fixed point  $x = 0$ . Then  $x = 0$  is asymptotically stable.

Use Krasovski's theorem to conclude the asymptotic stability of the origin for (4).

## Question 6

Consider an  $n$ -degree-of-freedom holonomic mechanical system (i.e. one that has only position-dependent constraints) with generalized coordinates  $q \in \mathbb{R}^n$  and generalized velocities  $\dot{q} \in \mathbb{R}^n$ . The total energy of the system is of the form

$$E(q, \dot{q}) = \frac{1}{2} \dot{q}^T M(q) \dot{q} + V(q),$$

where  $M \in \mathbb{R}^{n \times n}$  is the mass matrix (symmetric and positive definite), and  $V(q)$  is the potential energy. In the absence of external forces, the associated Lagrangian equations of motion are

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} = 0,$$

where  $L(q, \dot{q}) = \frac{1}{2} \dot{q}^T M(q) \dot{q} + V(q)$  is the Lagrangian of the mechanical system.

Show that if  $V(q)$  admits a strict local minimum at a point  $q_0$ , then  $q_0$  is a (nonlinearly) stable equilibrium for the mechanical system. (This result is also known as *Dirichlet's Theorem* in classical mechanics).

## Question 7

For what values of  $a$  and  $b$  does the dynamical system  $\dot{x} = Ax$  have a center fixed point at the origin ?

$$A = \begin{pmatrix} a & -b \\ b & 2 \end{pmatrix}$$

- (a)  $a = -2, b = 2$
- (b)  $a = \frac{1}{2}, b = 2$
- (c)  $a = -2, b = -4$
- (d)  $a = \frac{1}{2}, b = -4$

## Question 8

Consider the dynamical system

$$\ddot{\varphi} + a\dot{\varphi} + (b - \omega_0^2 \cos(\varphi)) \sin(\varphi) = 0$$

where  $a, b, \omega_0$  are positive constants. Which of the following is not correct?

- (a) If  $\omega_0^2 < b$ , then  $(0, 0)$  is an asymptotically stable fixed point.
- (b) If  $\omega_0^2 > b$ , then  $(0, 0)$  is a (Lyapunov) unstable fixed point.
- (c) For any  $a, b, \omega_0$ ,  $(\pi, 0)$  is an unstable fixed point.
- (d) For any  $a, b, \omega_0$ ,  $(0, 0)$  is a (Lyapunov) stable fixed point.

## Question 9

Consider a dynamical system of the form

$$\ddot{y} + p(t)\dot{y} + q(t)y = 0 \quad (5)$$

with  $C^1$  functions  $p(t)$  and  $q(t)$ . The transformation

$$x(t) = y(t)e^{-\frac{1}{2} \int_{t_0}^t p(s) \, ds}$$

puts this system in the form

$$\ddot{x} + \omega(t)x = 0 \quad (6)$$

For an appropriate  $C^1$  function  $\omega(t)$ . Does the Lyapunov stability of  $x = 0$  in (6) imply the Lyapunov stability of  $y = 0$  in (5)?

- (a) Yes, the two systems are topologically equivalent.
- (b) Only for a large enough  $t$ .
- (c) Yes, when  $p(t) < 0$  for all  $t \in \mathbb{R}$ .
- (d) None of the above