Spectral Graph Theory – Electric Flow

Zhiwei Zhang

March 18, 2019

1 Introduction

In this notes, we are going to discuss the interestingproperties when turning a graph into a network of resistors.

1.1 Electrical Laws

First recall some E&M Laws:

$$I = \frac{U}{R} \qquad \qquad \text{Ohm's law}$$
$$E = I^2 R \qquad \qquad \text{Energy formula}$$
$$|I_{v,in}| = |I_{v,out}| \qquad \qquad \text{conservation of flow}$$

Note that the last law only holds for nodes that are not source or sink.

1.2 Matrices

Also recall the laplacian of a graph:

$$L_{i,j} := \begin{cases} \deg(v_i) & \text{if } i = j \\ -1 & \text{if } i \neq j \text{ and } v_i \text{ is adjacent to } v_j \\ 0 & \text{otherwise} \end{cases}$$

Weighted laplacian of a graph:

$$L_{ij} = \begin{cases} -w_{ij} & \text{if } i \sim j \\ w_i & \text{if } i = j \\ 0 & \text{otherwise} \end{cases}$$

where $w_i = \sum_{j \sim i} w_{ij}$ is the sum of the weights of edges incident on vertex i.

Pseudoinverse of laplacian: Let $0 = \lambda_1 < \lambda_2 \leqslant \ldots \leq \lambda_n$ be the eigenvalues of L_G with associated eigenvectors $u_1, u_2, \ldots, u_n$. Then

$$L_G = \sum_{i=1}^{n} \lambda_i u_i u_i^T$$

. The pseudo-inverse of L_G is

$$L_G^+ = \sum_{i=2}^{n} \frac{1}{\lambda_i} u_i n_i^T$$

.

1.3   Formation on Graph

We will write a matrix formation of the problem.

- Let $G(V, E)$ be an undirected graph with $|V| = n, |E| = m$.

- Let $v \in \mathbb{R}^n$ be the vector representing the potentials of vertices.

- Edges represent the resistors, and $\forall e(u, v) \in E$. Edge e has resistance r_e.

- Let $f \in \mathbb{R}^m$ representing the flow of all edges, where $f(a, b)$ represents the flow from a to b with. Since $f(a, b)$ is directed, we have $f(a, b) = -f(b, a)$.

- Let weight $w_e = \frac{1}{r_e}$, or the "conductance" of e.

2   Matrix Formation

We also define $f_{ext}(a) = \sum_{b=(a,b)\in E} f(a, b)$, and $f_{ex}t(a)$ basically denotes the external current on a, which is **positive number if a is souce, negative number with equal magnitude if a is source, and zero otherwise**. So f_{ext} is a very sparse vector.

Ohm's law directly states that $f(a, b) = \frac{v(a)-v(b)}{r_{a,b}} = w_{a,b}(v(a) - v(b))$, therefore

$$\sum_{b:(a,b)\in E} f(a, b) = \sum_{b:(a,b)\in E} w_{a,b}(v(a) - v(b)) = d(a)v(a) - \sum_{b:(a,b)\in E} w_{a,b}v(b)$$

where $d(a) = \sum_{b:(a,b)\in E} w_{a,b}$, the weighted degree of a.

Notice that $d(a), w_{a,b}$ are entries of the weighted laplacian L_G, and through simple verification, we can show that the equation above the equivalent to $L_G v = f_{ext}$

3   Computing Voltages

Since we know that $Nul(L_G) = \vec{1}$, it's trivial to see that for any vector x, $L_G x$ is perpendicular to $\vec{1}$, which implies there's a solution to $L_G v = f_{ext}$ iff f_{ext} is perpendicular to 1. This is also simply true since the two non-zero entries of f_{ext} have the same magnitude with opposite sign.

Therefore $v = L_G^+ f_{ext}$ is the only solution with $v \perp \vec{1}$, and the whole set of solution is $\{v + c\vec{1} | c \in \mathbb{R}\}$

This also makes sense in the physical way, as if we increase the potential of all nodes, the physical flow or energy will not change as electrical potentials are only significant when taking differences.

4   Computing Currents

First we reintroduce the incidence matrix B of dimension $m \times n$, and the row corresponding to the edge $e = (a, b)$ where $a < b$ is $(x_a - x_b)^T$, where x_a is the characteristic with the only non-zero entry at $a - th$ entry of value 1.

Let w be the mxm diagonal matrix where $W_{e,e} = w_e$ is the weight of edge e.

Notice that Bv gives the potential difference on each edge, and therefore $f = WBv$.

It's also true that $L_G = \sum_{e:(a,b)} w_e (x_a - x_b) (x_a - x_b)^T = B^T W B$.

Therefore $f_{ext} = L_G v = B^T W B v = B^T f$

5   Random Walk and Effective Resistance

Random walk is largely studied on graphs, and it turns out that the concept of effective resistance is very useful in this area. Before we get into the theorems, we need to specify some restrictions on the graph we are using and define some terms.

5.1   Specification and Definitions

When modeling random walks with electric flows, we general set the resistance of each edge to be 1, and $\forall e \in E, w_e = w/r_e = 1$.

In addition:

- We define hitting time $h_{uv} :=$ expected number of steps to reach v from u

- We define commutime $C_{u,v} :=$ expected number of steps to start from u, travelling to v, and come back to u. More specifically, $C_{u,v} = h_{u,v} + h_{v,u} = C_{v,u}$. Notice that $h_{u,v}$ not necessarily $= h_{v,u}$.

5.2   Theorems and Corollaries

Theorem 1. $C_{s,t} = 2mR_{eff}(s,t)$

Proof. First using recursive definition, we can develop the equation:

$$\forall v \in V - t, h_{vt} = \sum_{w:vw \in E} \frac{1}{d(v)}(1 + h_{wt}).$$

Notice that this only holds for all vertices except t, which is very important since $h_{tt} = 0$, and there'd also be no reason to walk away from t and walk back. We will reemphasize this point later again.

The equation is equivalent to

$$d(v) = d(v)h_{vt} - \sum_{w:vw \in E} h_{wt} = \sum_{w:vw \in E} (h_{vt} - h_{wt})$$

Notice that this form is similar to the Laplacian system of linear equations (reacall $\sum_{b:(a,b) \in E} f(a,b) = \sum_{b:(a,b) \in E} w_{a,b}(v(a) - v(b)) = d(a)v(a) - \sum_{b:(a,b) \in E} w_{a,b}v(b)$)

Next, we denote ϕ_{vt} be the potential at v with $\phi_{tt=0}$. Now if for all $v \in V - t$, we inject $d(v)$ units of current (so $|V| - 1$ sources). Totally we inject $2m - d(t)$ units of current, and we have to remove the same amount from t.

Since $\forall v \in V - t$, $d(v)$ units of current are injected, its external current is therefore $d(v)$, which dissipates through its neighbors. Therefore

$$d(v) = \sum_{w:vw \in E} (\phi_{vt} - \phi_{wt})/r_{(v,w)} = \sum_{w:vw \in E} (\phi_{vt} - \phi_{wt}).$$

4

Let f_t be the external flow vector if t is the only sink. In this situation, $\forall v \in V - t, f_t(v) = d(v, f_t(t) = d(t) - 2m)$

Therefore we get $\forall v \in V - t, d(v) = \sum_{w:vw \in E}(\phi_{vt} - \phi_{wt}) = \sum_{w:vw \in E}(h_{vt} - h_{wt})$, so $\vec{\phi_t}$ and $\vec{h_t}$ satisfy the same equation: $L_G x = f_t$. Since we have a set of possible solutions, but also since $h_{tt} = \phi_{tt} = 0$, $\vec{h_t} = \vec{\phi_t}$ at all points. We now reamphasize the importance of $d(v) = \sum_{w:vw \in E}(h_{vt} - h_{wt})$ only satisfies for $v \in V - t$: we can see that if we extend this definition to t, they will not satisfy the laplacian equation.

Now, we define f_s similar to f_t (now s is the only sink, all other vertices are sources).

Therefore we have $L_G\left(\vec{h_t} - h_S\right) = f_t - f_S = 2m\left(x_s - x_t\right)$, and so $\left(\vec{h_t} - \vec{h_s}\right)/2m = L_G^\dagger\left(x_s - x_t\right)$.

$$R_{eff}(s,t) = (x_s - x_t)^T LG^+(x_s - x_t)$$
$$= (x_s - x_t)^T \left(\frac{1}{2m}\left(\vec{h_t} - \vec{h_s}\right)\right)$$
$$= \frac{1}{2m}\left(h_t(s) + \vec{h_s}(t)\right)$$
$$= \frac{1}{2m}(h_{st} + h_{ts})$$
$$= \frac{1}{2m}C_{s,t}$$

$\square$