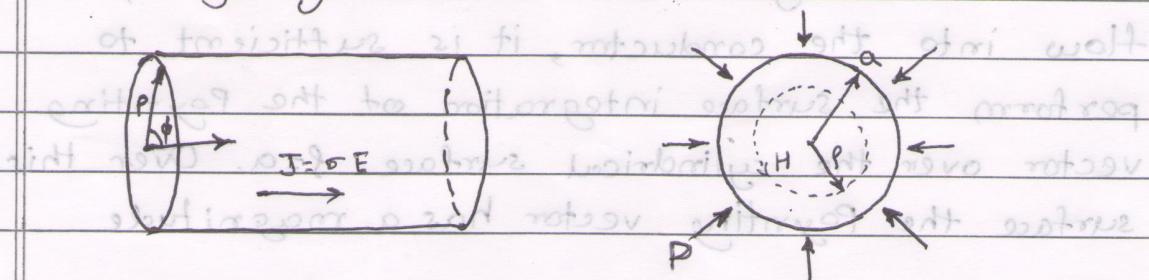


(1)

orb ziloyot si rotayot to osibisib ott
The Poynting vector is $\bar{P} = \bar{E} \times \bar{H}$ ott to zixi



$$\text{We know that, } \bar{J} = \sigma \bar{E} \quad \therefore \bar{E} = \frac{\bar{J}}{\sigma}$$

$$\therefore \bar{P} = \bar{J}_z \times \bar{H} \text{ ides to longotai ziloyot ott}$$

surface to zero $\times \frac{\partial \bar{J}_z}{\partial r} = 0$

Assume the current density \bar{J}_z is \hat{a}_z direction then,

$$\bar{J} = J_z \hat{a}_z \quad \text{and} \quad \bar{H} = H_\phi \hat{a}_\phi$$

$$(\text{corr}) (\hat{a}_z) =$$

$$\bar{P} = \frac{J_z}{\sigma} \hat{a}_z \times H_\phi \hat{a}_\phi = - \frac{J_z}{\sigma} H_\phi \hat{a}_\rho \quad \dots [\because \hat{a}_z \times \hat{a}_\phi = -\hat{a}_\rho]$$

Now, the magnetic field intensity H_ϕ at any radius ρ is

$$\text{introduced ott to } I = \frac{J_z \times \pi \rho^2}{2\pi \rho} = \frac{J_z \rho}{2}$$

$$\dots [\because I = J_z \times \text{area}]$$

$$\bar{P} \text{ at any radius} = - \frac{J_z}{\sigma} \times \frac{J_z \rho}{2} \hat{a}_\rho = - \frac{J_z^2 \rho}{2 \sigma} \hat{a}_\rho$$

~~11~~

(2) (i) We have $a = 1\text{cm}$, $\sigma = 2 \times 10^7 \text{ S/m}$, $I = 100 \text{ A}$, $l = 1\text{m}$

$$R = \frac{\rho l}{A} = \frac{l}{\sigma A}$$

$$= \frac{1}{2 \times 10^7 \times \pi \times (1 \times 10^{-2})^2}$$

$$= 1.59 \times 10^{-4} \Omega$$

$$P = I^2 R$$

$$= (100)^2 \times 1.59 \times 10^{-4}$$

$$= 1.59 \text{ W}$$

(ii) Average J is $\bar{J} = \frac{I}{\pi r^2}$ A/m^2 for part (iii)

$$= 318.31 \times 10^3 \text{ A/m}^2$$

$$\therefore \bar{J} = 318.31 \times 10^3 \text{ A/m}^2$$

We know that $\bar{J} = \sigma \bar{E}$

$$\sigma \bar{E} = \bar{J}$$

$$2 \times 10^7 \times \bar{E} = 318.31 \times 10^3$$

$$\therefore \bar{E} = 15.92 \times 10^{-3} \text{ V/m}$$

Magnetic field within the conductor is,

$$\bar{H} = I \left(\frac{1}{2\pi a^2} \right) \bar{a}_\phi$$

where, a is the radius of the conductor and δ is the radius of circular path inside, then

$$\bar{H} = 100 \times \frac{1}{2\pi (10^{-2})^2} \delta \bar{a}_\phi$$

$$= 159.2 \times 10^3 \delta \bar{a}_\phi \text{ A/m}$$

The Poynting vector is

$$\bar{P} = \bar{E} \times \bar{H}$$

$$= \begin{vmatrix} \bar{a}_r & \bar{a}_\theta & \bar{a}_z \\ 0 & 0 & 15.92 \times 10^{-3} \\ 0 & 15.92 \times 10^3 & 0 \end{vmatrix}$$

$$= (-15.92 \times 10^{-3} \times 15.92 \times 10^3) \cdot \bar{a}_\theta$$

$$\bar{P} = -2534.1 \bar{a}_\theta$$

(iii) Integrating \bar{P} over the cylindrical surface, enclosing one meter length of conductor, we get,

$$\begin{aligned} \int_S \bar{P} \cdot d\bar{s} &= \int_S \bar{P} \cdot (\hat{s} d\phi dz \bar{a}_\theta) \\ &= \iint_{0 \rightarrow 1}^{0 \rightarrow 2\pi} -2534.1 \rho^2 d\phi dz \\ &= -2534.1 \times [\phi]_0^{2\pi} [z]_0^1 \times (1 \times 10^{-2})^2 \\ &= -1.59 \text{ W} \end{aligned}$$

and $|\bar{P} \cdot d\bar{s}| = 1.59 \text{ W}$ which is same as in part (i)

$$\bar{a}_\theta \left(\frac{\partial}{\partial r} \right) I = \bar{H}$$

$$\bar{a}_\theta \left(\frac{\partial}{\partial r} \right) I = \bar{H}$$

$$(m/A) \bar{a}_\theta \left(\frac{\partial}{\partial r} \right) I =$$

(B)

Given: $H_0 = 0.2$ and for free space, $\eta = 120\pi (S2)$

As the expression for fields says that the wave is travelling in $+x$ direction,

$$\bar{P}_{avg} = \frac{1}{2\eta} E_0^2 \bar{a}_x = \frac{1}{2} \eta H_0^2 \bar{a}_x$$

$$\bar{P}_{avg} = \frac{1}{2} (120\pi) (0.2)^2 \bar{a}_x$$

Normal to plane with equation $Ax + By + Cz + D = 0$

$$\bar{a}_n = \pm \frac{A\bar{a}_x + B\bar{a}_y + C\bar{a}_z}{\sqrt{A^2 + B^2 + C^2}}$$

For the given plane, $\bar{a}_n = \frac{\bar{a}_x + \bar{a}_y}{\sqrt{2}} = \frac{1}{\sqrt{2}}$

The differential area for the given plane,

$$ds = d\bar{s} \bar{a}_n$$

Power crossing the given plane is,

$$\bar{P} = \int_S \bar{P}_{avg} \cdot \bar{ds} = \int_S \frac{1}{2} (120\pi) (0.2)^2 \bar{a}_x \cdot ds \left(\frac{\bar{a}_x + \bar{a}_y}{\sqrt{2}} \right)$$

$$= \frac{1}{2} (120\pi) (0.2)^2 \frac{1}{\sqrt{2}} \int_S ds = \frac{1}{2} (120\pi) (0.2)^2 \frac{1}{\sqrt{2}} S$$

$$= \frac{1}{2} (120\pi) (0.2)^2 \times \frac{1}{\sqrt{2}} \times (0.1)^2$$

$$= 53.31 \text{ mW}$$

(ii) The given plane is $x=1$ = constant,

(2) ITOS135 ~~so it's not true~~ $S \cdot 0 = 0$ ~~it's not~~

$$\bar{a}_n = \bar{a}_x$$

~~if surface area $dS = d\bar{s} \bar{a}_n = d\bar{s} \bar{a}_x$ then $\int d\bar{s} \bar{a}_x$ is the power~~

Power crossing the circular disc of radius
5 cm

$$P = \int_S \bar{P}_{avg} \cdot \bar{a}_x(0,0)(\text{ITOS1}) \frac{1}{2} = \text{avg}$$

$$= \int_0^{\pi} \frac{1}{2} (120\pi) (0.2)^2 \bar{a}_x \cdot d\bar{s} \bar{a}_x$$

$$= \frac{1}{2} (120\pi) (0.2)^2 s$$

$$= \frac{1}{2} (120\pi) (0.2)^2 \times \pi (0.05)^2$$

$$P = \underline{\underline{59.23 \text{ mW}}} \quad \text{avg over not}$$

$$\text{avg over not even though } \bar{a}_x = \bar{a}_y$$

~~avg over not even though $\bar{a}_x = \bar{a}_y$~~

$$\left(\frac{\sqrt{0} + \sqrt{0}}{2} \right) 2b \cdot \bar{x}^2 (0,0)(\text{ITOS1}) \frac{1}{2} = 2b \cdot \bar{x}^2 \frac{1}{2} = 0$$

$$\frac{1}{2} \left[\bar{x}^2 (0,0)(\text{ITOS1}) \frac{1}{2} - 2b \right] \frac{1}{2} \left[\bar{x}^2 (0,0)(\text{ITOS1}) \frac{1}{2} - 2b \right]$$

$$= \frac{1}{2} (1.0) \times 1 \times \bar{x}^2 (0,0)(\text{ITOS1}) \frac{1}{2} =$$

$$\bar{x} = 18.25$$

(4) Given electric field intensity in phasor form

$$\bar{E}_s = E_0 (\bar{a}_y - j \bar{a}_z) e^{-j \beta z}$$

Instantaneous expression of electric field intensity will be,

$$\bar{E} = \operatorname{Re} \left\{ E_0 (\bar{a}_y - j \bar{a}_z) e^{j(\omega t - \beta z)} \right\}$$

$$= \operatorname{Re} \left\{ E_0 (\bar{a}_y - j \bar{a}_z) [\cos(\omega t) + j \sin(\omega t)] \right\} e^{-j \beta z}$$

$$= E_0 (\bar{a}_y \cos(\omega t) + \bar{a}_z \sin(\omega t)) e^{-j \beta z}$$

Therefore, the magnitude of the field is

$$|\bar{E}| = \sqrt{(E_0 \cos \omega t)^2 + (E_0 \sin \omega t)^2}$$

$$\text{or, } |E_x|^2 + |E_y|^2 = E_0^2$$

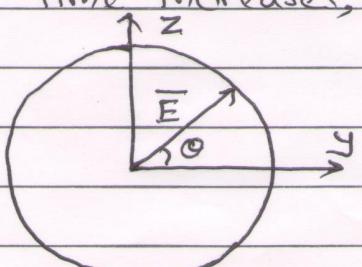
which is a circular equation i.e. the wave is circularly polarized.

Now, the instantaneous angle Θ that the field \bar{E} makes with y -axis is given as

$$\tan \Theta = \frac{E_0 \sin \omega t}{E_0 \cos \omega t}$$

$$\text{or, } \Theta = \omega t$$

\therefore As time increases, \bar{E} rotates from y to z as shown and since the direction of wave propagation is in $+\hat{a}_z$ direction so, the rotation from y to z obeys the right hand rule.



\therefore The field is right hand circularly polarized.

(5) Given, the phasor form of electric field intensity,

$$\bar{E}_s = 4(\bar{a}_z - j\bar{a}_x)e^{-jBy}$$

Show to write about wave

So, the electric field intensity of the reflected wave will be

at a more distance of ωz

$$\bar{E}_{rs} = R[4(\bar{a}_z - j\bar{a}_x)]e^{jBy}$$

at first distance ωz

where R is the reflection coefficient at the interface.

For perfect conductor $R = -1$,

$$\therefore \bar{E}_{rs} = 4(-\bar{a}_z + j\bar{a}_x)e^{jBy}$$

∴ Instantaneous expression of the electric field of reflected wave will be

$$\begin{aligned}\bar{E} &= \operatorname{Re} \{ 4(-\bar{a}_z + j\bar{a}_x)(\cos \omega t + j \sin \omega t) \} e^{jBy} \\ &= 4(-\cos(\omega t)\bar{a}_z - \sin(\omega t)\bar{a}_x)e^{jBy}\end{aligned}$$

Therefore, the magnitude of the reflected field is

$$|E| = \sqrt{(4\cos \omega t)^2 + (4\sin \omega t)^2}$$

$$\text{or, } |E_1|^2 + |E_2|^2 = 4$$

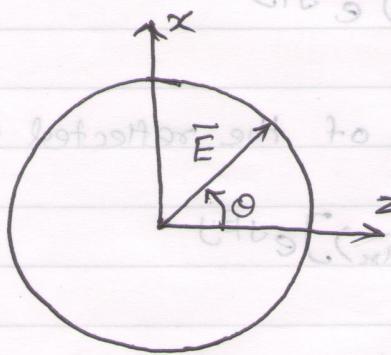
which is a circular equation i.e. the wave is circularly polarized.

Now, the instantaneous angle θ that \bar{E} makes with z -axis is given as,

$$\tan \theta = \frac{-4 \sin \omega t}{-4 \cos \omega t}$$

$$\therefore \theta = \omega t$$

So, as time increases, electric field \vec{E} rotates from z to x as shown in the figure below



Since the direction of wave

propagation is along $-\hat{x}$,

so, the rotation from z to x follows left hand rule.

Thus, we conclude that the

EM wave is LHC (left hand circularly) polarized.

$$E = E_0 (\cos(\omega t + \phi_0) \hat{i} + \sin(\omega t + \phi_0) \hat{j})$$

better to write vectors in terms of unit vectors \hat{i} and \hat{j} .

as lines show

$$E = E_0 (\cos(\omega t + \phi_0) \hat{i} + \sin(\omega t + \phi_0) \hat{j})$$

$$E = E_0 (\cos(\omega t) \hat{i} - \sin(\omega t) \hat{j})$$

if we take $\phi_0 = 0$ then

$$E = E_0 (\cos(\omega t) \hat{i} + \sin(\omega t) \hat{j})$$

$$E = E_0 \hat{i} + E_0 \hat{j}$$

which shows E always remains in xy plane

logically

E is not zero along z axis so it must be zero

now E is zero at $t = 0$

$$\omega t = 0 \quad \therefore$$

$$\cos(\omega t) = 1 \quad \sin(\omega t) = 0$$

$$\cos(\omega t) = 1 \quad \sin(\omega t) = 0$$

(6)

Given, the electric field intensity of the propagating wave,

$$\vec{E} = \bar{a}_x \sin(\omega t - \beta z) + \bar{a}_y \sin(\omega t - \beta z + \pi/2)$$

So, we conclude that the wave is propagating along \bar{a}_z direction and the field components are along \bar{a}_x and \bar{a}_y and equal.

i.e. $E_x = E_y$

\therefore The wave is circularly polarized. Now we will determine the field is either right circular or left circular. The angle between the electric field \vec{E} and x -axis is given as,

$$\theta = \tan^{-1} \left(\frac{\cos \omega t}{\sin \omega t} \right) = \frac{\pi}{2} - \omega t$$

So, with increase in time the tip of the field intensity moves from y to x -axis and as the wave is propagating in \bar{a}_z direction therefore, the wave is left hand circularly polarized.

$\omega t = \text{time}$ $\theta = \text{phase}$

(8) Linearly polarized wave A in the x direction can be
(7) written as, $E_A(z) = \bar{a}_x E_0 e^{-j\beta z} [V/m]$

$$E_A(z) = \bar{a}_x E_0 e^{-j\beta z} [V/m]$$

Wave B has equal components in x and y direction,

$$\therefore E_B(z) = (\bar{a}_x + \bar{a}_y) E_0 \frac{e^{-j\beta z}}{\sqrt{2}} e^{-j\theta} [V/m]$$

The factor $\sqrt{2}$ ensures equal amplitudes of wave A and B. $e^{-j\theta}$ term in the equation for wave B implies the lag compared to wave A. Superposition gives,

$$E_A + E_B = \bar{a}_x \left(E_0 e^{-j\beta z} + \frac{E_0}{\sqrt{2}} e^{-j\beta z} e^{-j\theta} \right) + \bar{a}_y \frac{E_0}{\sqrt{2}} e^{-j\beta z} e^{-j\theta} [V/m]$$

In time domain,

$$E_A + E_B = \bar{a}_x \left(E_0 \cos(\omega t - \beta z) + \frac{E_0}{\sqrt{2}} \cos(\omega t - \beta z - \theta) \right)$$

$$+ \bar{a}_y \frac{E_0}{\sqrt{2}} \cos(\omega t - \beta z - \theta) [V/m]$$

when $z=0$ and $\omega t=0$:

$$\left[E_A + E_B \right]_{z=0, \omega t=0} = \hat{a}_x \left[E_0 + \frac{E_0}{\sqrt{2}} \cos(\theta) \right] + \hat{a}_y \frac{E_0}{\sqrt{2}} \cos \theta [V/m]$$

- (1)

when $z=0$ and $\omega t = \pi/2$

$$[E_A + E_B] = \bar{a}_x \left[\frac{E_0}{\sqrt{2}} \sin \theta \right] + \bar{a}_y \left[\frac{E_0}{\sqrt{2}} \sin \theta \right] \text{ [V/m]}$$

$\therefore z=0, \omega t=\pi/2 \quad \Rightarrow \quad [E_A + E_B] = (\vec{s})_A$

(2)

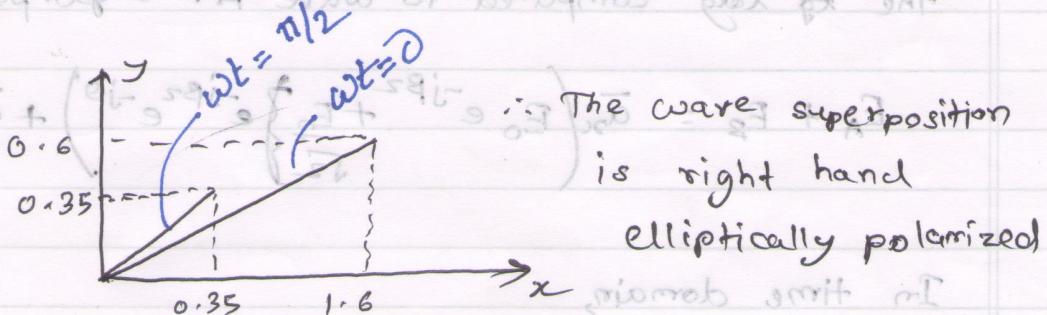
Let us draw x vs z at two times t and $\omega t = \pi/2$

Substituting $\theta = 30^\circ$ in (1) and (2), we get,

$$[E_A + E_B] \Big|_{z=0, \omega t=0, \theta=30} = \bar{a}_x 1.612 E_0 + \bar{a}_y 0.612 E_0$$

Now A shows to elliptical loops towards the right side

$$[E_A + E_B] \Big|_{z=0, \omega t=\pi/2, \theta=30} = \bar{a}_x 0.354 E_0 + \bar{a}_y 0.354 E_0$$



$$\left((0 - \omega q - \omega) \frac{\partial}{\partial t} \hat{x} + (\omega q - \omega) \frac{\partial}{\partial t} \hat{y} \right) \vec{v} = \hat{x} + \hat{y}$$

$$[m/v] (0 - \omega q - \omega) \frac{\partial}{\partial t} \hat{x} +$$

$$; \omega = \omega_0 \text{ and } \omega = s \text{ rad/s}$$

$$[m/v] 0 \frac{\partial}{\partial t} \hat{x} + \left[(\omega) \frac{\partial}{\partial t} \hat{x} + \hat{y} \right] \frac{\partial}{\partial t} \hat{y} = [\hat{x} + \hat{y}]$$

extreme

(1) →