



Workshop on R&D on Superconducting Proton Linac

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Cavity mechanics and vibrations under Lorentz forces excitation

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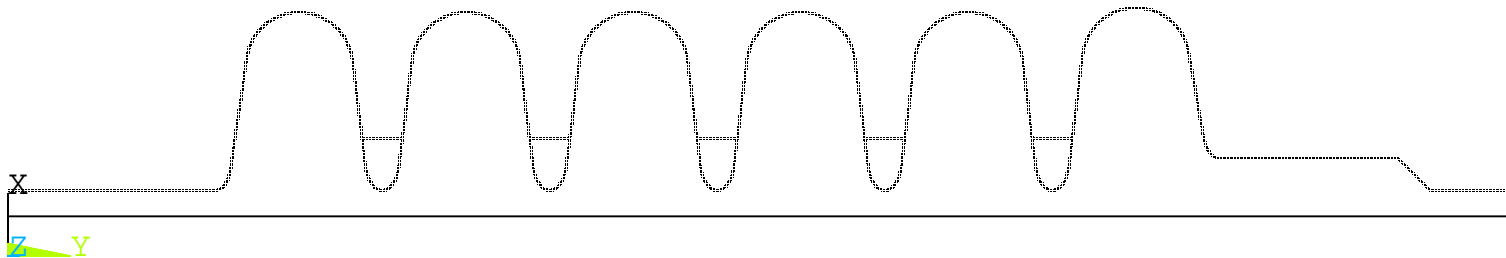
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Overview of the analysis presented here

- The behavior of the real system (cavity + environment) **cannot be precisely determined** until all details are fixed
- A “bare” multicell cavity is taken as a **reference** and **excited near its fundamental resonance** applying a pulsed Lorentz pressure load, to determine the induced RF phase shift on the “Qloaded” curve

Reference Parameters for calculation

- Cavity is the SNS Beta 061 (627 in BuildCavity Database)
 - Stiffening radius: 80mm
 - Niobium Thickness: 3.8mm
 - Young Modulus: 103 GPa
 - Poisson Ratio: 0.3
- Operating conditions:
 - $E_{acc} = 10 \text{ MV/m}$
 - $Q_{ext} = 1 \cdot 10^6$
 - Macro bunch length= 1ms
 - Repetition rate= 60Hz (but I will force resonance excitation of the modes!)



Lorentz Forces and Microphonic Vibrations: Some Considerations

The system under study is **ONLY** the “bare” cavity:

- **No ancillaries components**
 - No external mass added
 - No external position constrain
- No external load: **cavity length is fixed**
- Cavity has **not been pre-stressed**
 - No tuning compression
 - No residual stress influence

The results do not have absolute value but they represent a good exercise to understand the problem

Cavity Eigenfrequencies

A longitudinally axi-symmetric geometry has been studied.

Modal shapes have been computed for the **unstiffened** cavity

Hz
81.9
161.8
237.4
304.5
356.9
381.9
680.2
717.0
761.5
804.7
837.9

Two possible kinds of modal shape:

“**Bellow**” shape:

- Lower frequencies
- Produce “semirigid” displacement of cells

“**Cell**” shape

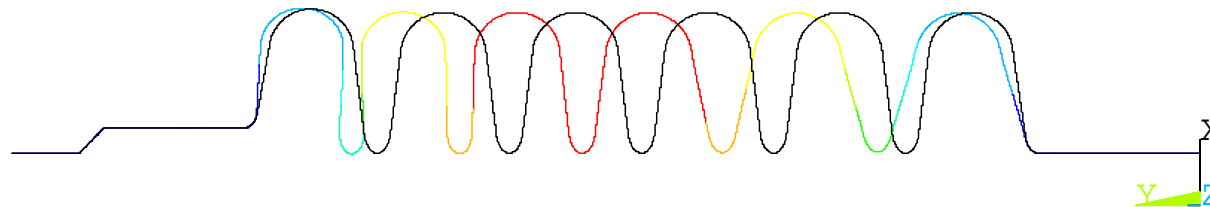
- Higher frequencies (>1 kHz)
- Cell are “independent”

Static Deformation is “similar” to both but...

Bellow shape

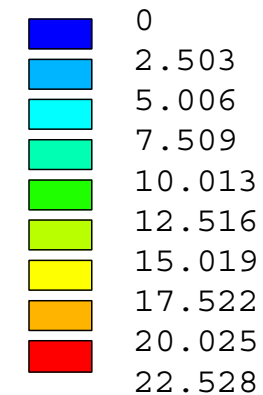
1

$f=81.9 \text{ Hz}$



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OCT 17 2000
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RSYS=0
PowerGraphics
EFACET=1
AVRES=Mat
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SMX =22.528

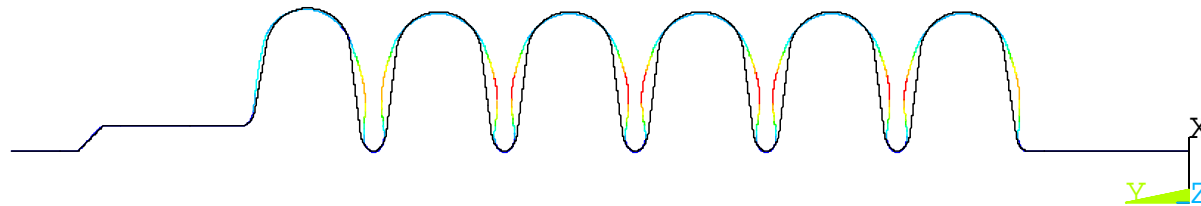
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XF =106.989
YF =510.8
A-ZS=90
Z-BUFFER



Cell shape

1

$f=1241$ Hz



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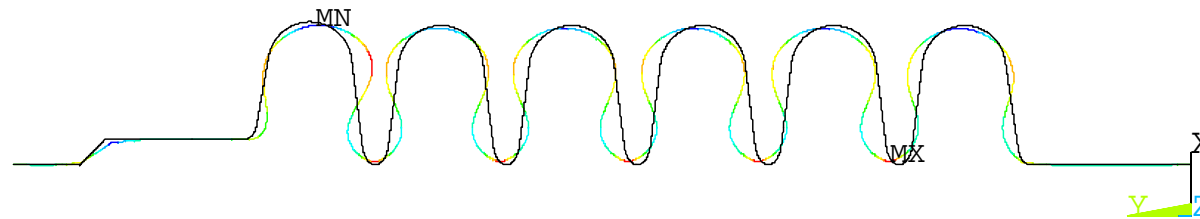
ZV =1
DIST=561.88
XF =105.817
YF =510.8
A-ZS=90
Z-BUFFER



Steady State deformation by Lorentz Forces

1

Steady state $\Delta f=1.7$ Hz @ 1 MV/m



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NODAL SOLUTION
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TIME=1
UX (AVG)
RSYS=0
PowerGraphics
EFACET=1
AVRES=Mat
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SMN =-.279E-06
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-.147E-06
-.814E-07
-.157E-07
.500E-07
.116E-06
.182E-06
.247E-06
.313E-06

Displacements (mm)
(@ 1 MV/m Eacc)

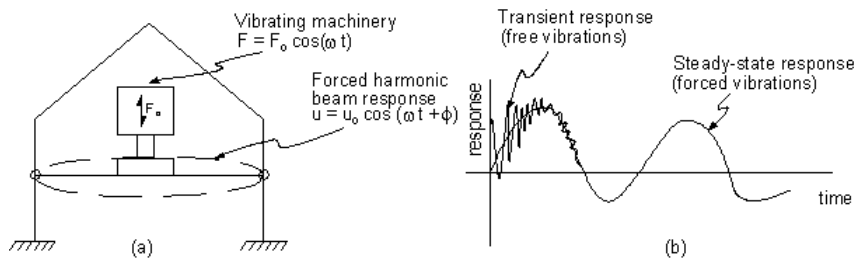
Steady State

Deterministic Harmonic Analysis

Only the **shape** information can be extracted from **modal** analysis.

The load is the Lorentz pressure distribution (calculated from Cavity SFO file)

Suppose that **load is varying harmonically** at different frequencies and forget the transitory effects



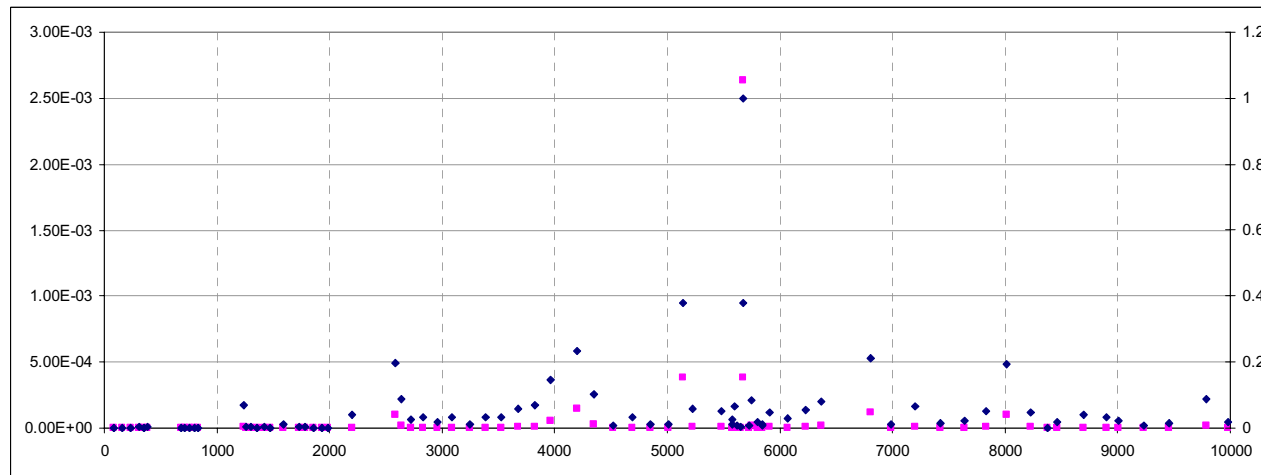
Solution can be expressed by **superimposition** of natural modes.

$$Y(x, \Omega) = \sum \frac{f_n}{(\omega_n^2 - \Omega^2) + i \cdot (2 \cdot \omega_n \cdot \Omega \cdot \xi_n)} \cdot \phi_n(x)$$

f_n is the participation factor and gives information on the “compatibility” between a mode shape and the load .

Partecipation Factor Effects

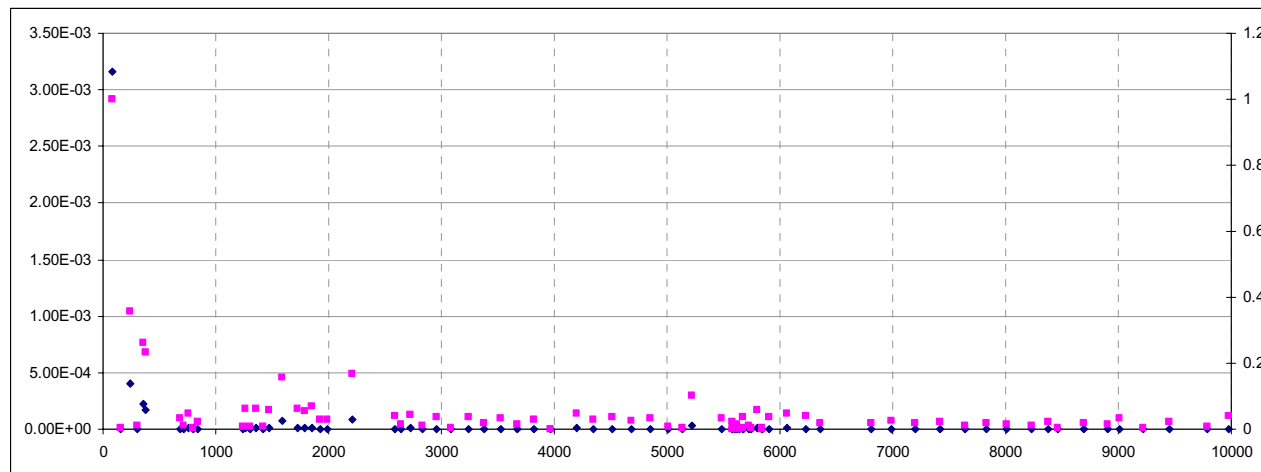
R direction



Formally the participation factor is the dot product between the modal shape and the load profile.

Modal shapes are also normalized to a proper unity vector so also normalization has to be kept into account

Z direction



Low frequency modes are well coupled to transversal components while higher frequencies are coupled to radial components

Modal analysis completed

Hz	Rcoupling	Zcoupling
81.9	0.0001	1.0000
161.8	0.0010	0.0021
237.4	0.0001	0.3554
304.5	0.0019	0.0127
356.9	0.0010	0.2630
381.9	0.0029	0.2344
680.2	0.0002	0.0336
717.0	0.0014	0.0108
761.5	0.0002	0.0480
804.7	0.0013	0.0051
837.9	0.0000	0.0206
	(...)	
5619.4	0.0075	0.0145
5647.3	0.0048	0.0011
5670.2	0.3805	0.0357
5672.9	1.0000	0.0034
5731.0	0.0085	0.0112
5745.3	0.0842	0.0033
5803.3	0.0188	0.0590
5840.2	0.0063	0.0042
5844.1	0.0116	0.0002

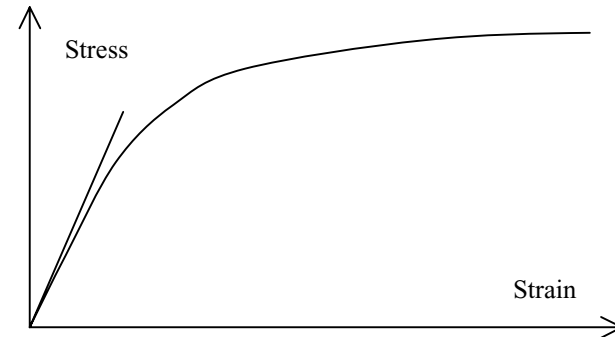
Once the load shape is defined the mode table can be completed
(i.e. loads can be projected on each mode shape)

Each mode is characterized by its influence in the response of the system to a specified load

Damping

We consider only material inner damping:

- No damping coming from supports
- No damping from ancillary components



Material damping is almost due to **plasticity effects during the oscillation.**

Material damping depends on the amplitude of the oscillation.

Analytical energy consideration and measurements pointed out that the energy lost during an oscillation with respect to the elastic energy can be generally expressed by a simplified relation that is a reasonable approximation for all metals

$$\psi = \frac{\epsilon}{\epsilon_{pl}} \cdot 8 \rightarrow Q = \frac{1}{\psi}$$

As reference I take **Q=50**, that corresponds to **0.25 % of plasticity effects.**

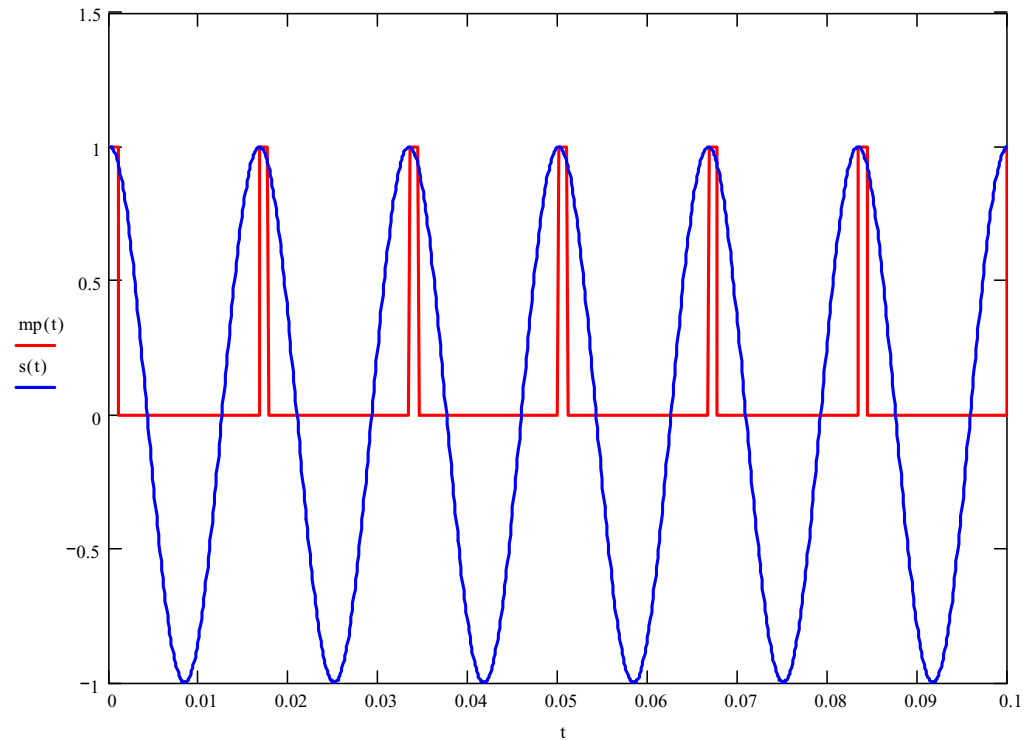
'Real' pulse description

Lorentz excitation is not harmonic, but the correspondence can be analytically corrected by means of Laplace formalism

$$\text{Out}(s) := \sum_{n=0}^{\infty} \frac{(1 - e^{-s \cdot \lambda})}{s} \cdot e^{\frac{-n \cdot s}{f}} \cdot H(s)$$

$$\text{End}(s) := \lim_{s \rightarrow 0} s \cdot \text{Out}(s)$$

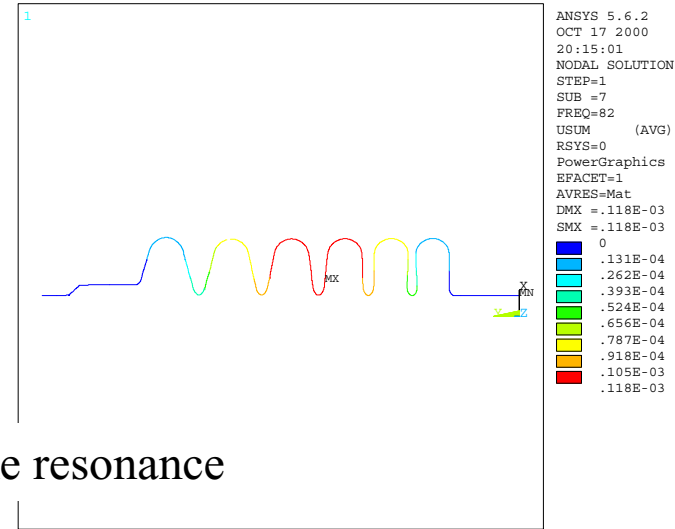
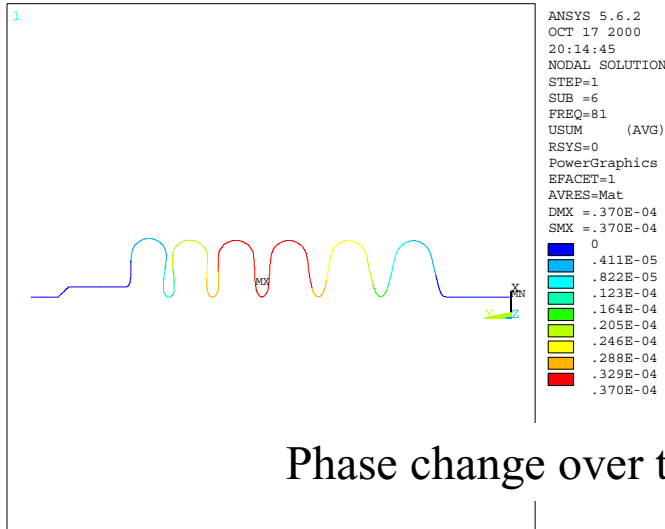
The results of the harmonic analysis give the $H(s)$ function



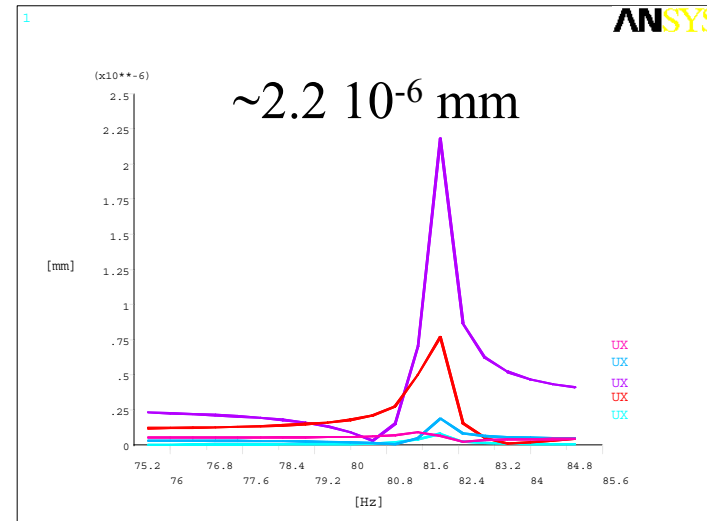
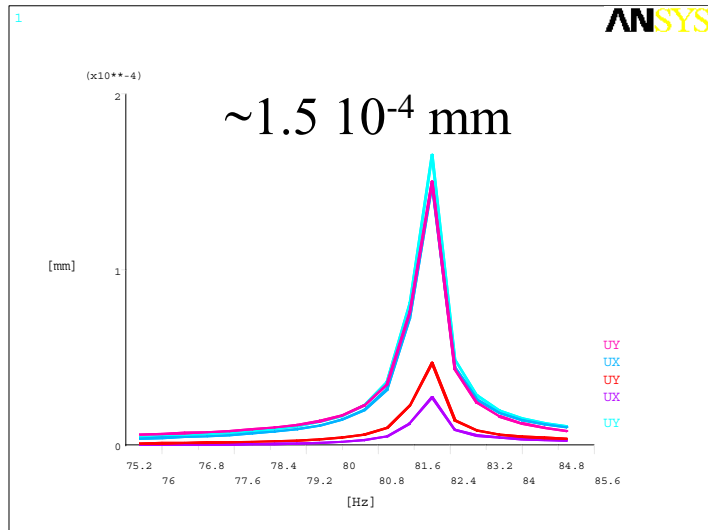
λ is the pulse length (1 ms)

$$\text{Amplitude}_{\text{Norm}} = \text{Amplitude} \cdot (\lambda \cdot f)$$

Unstiffened Cavity Analysis



Phase change over the resonance



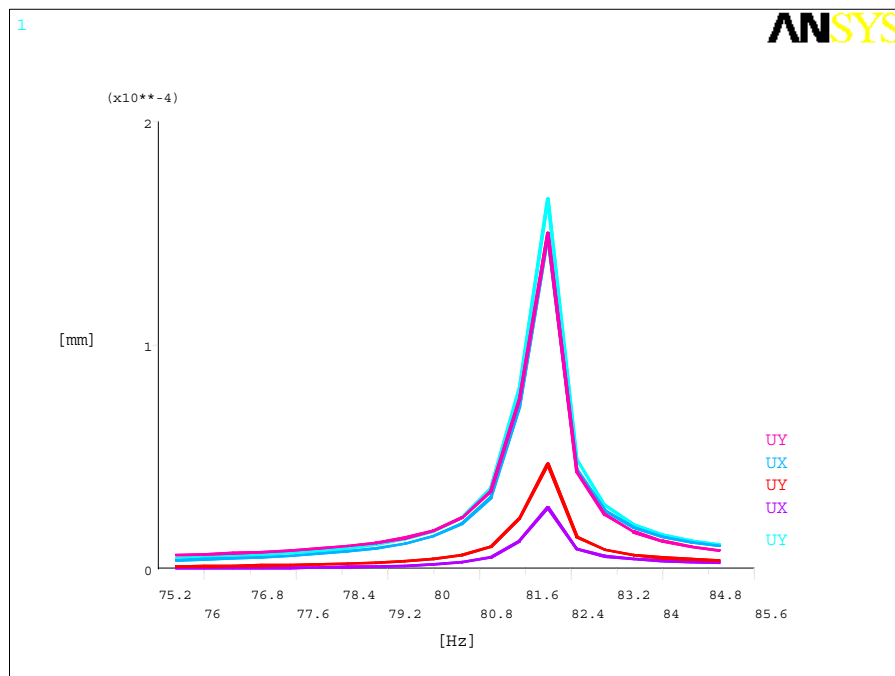
Modulus of displacement in Z and R direction

Unstiffened Cavity Analysis

The **real and imaginary part** of oscillation displacement **near the resonance** are extracted and used to calculate the **frequency shift by SLATER theorem**

Using the Q_{loaded} curve: **the frequency shift is moved to RF phase shift**

$$\tan(\phi) = \frac{\omega \cdot \omega_0}{Q \cdot (\omega_0^2 - \omega^2)} \longrightarrow \tan(\phi) = \frac{F}{2 \cdot Q \cdot \Delta F}$$



$$\Delta\phi = 2.9^\circ @ 81\text{Hz}$$

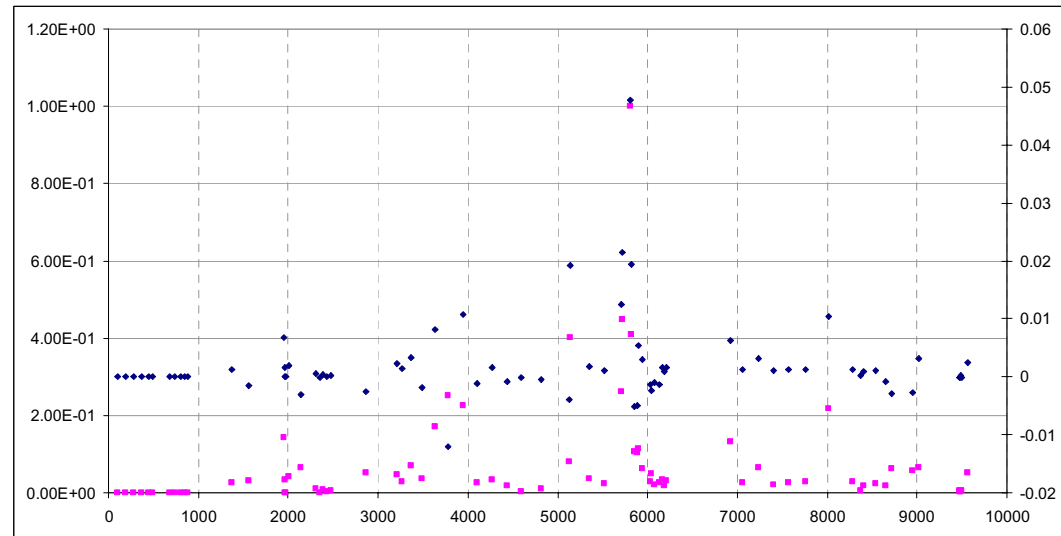
$$\Delta\phi = 6.1^\circ @ 81.9\text{Hz}$$

$$\Delta\phi = 0.5^\circ @ 83\text{Hz}$$

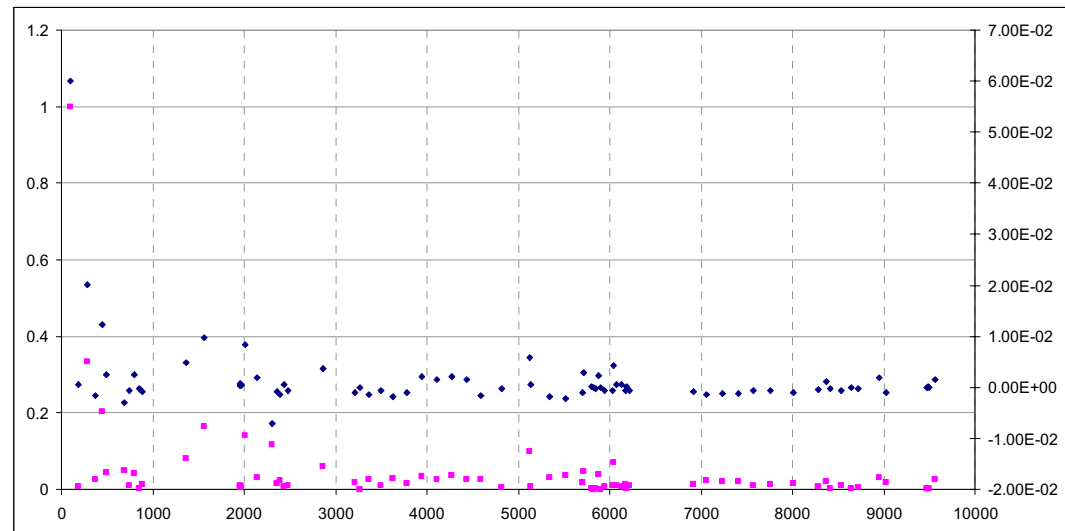
Also at the exact resonance the Lorentz load is able to transfer a very small amount of energy into the mode!

Stiffened Cavity Analysis

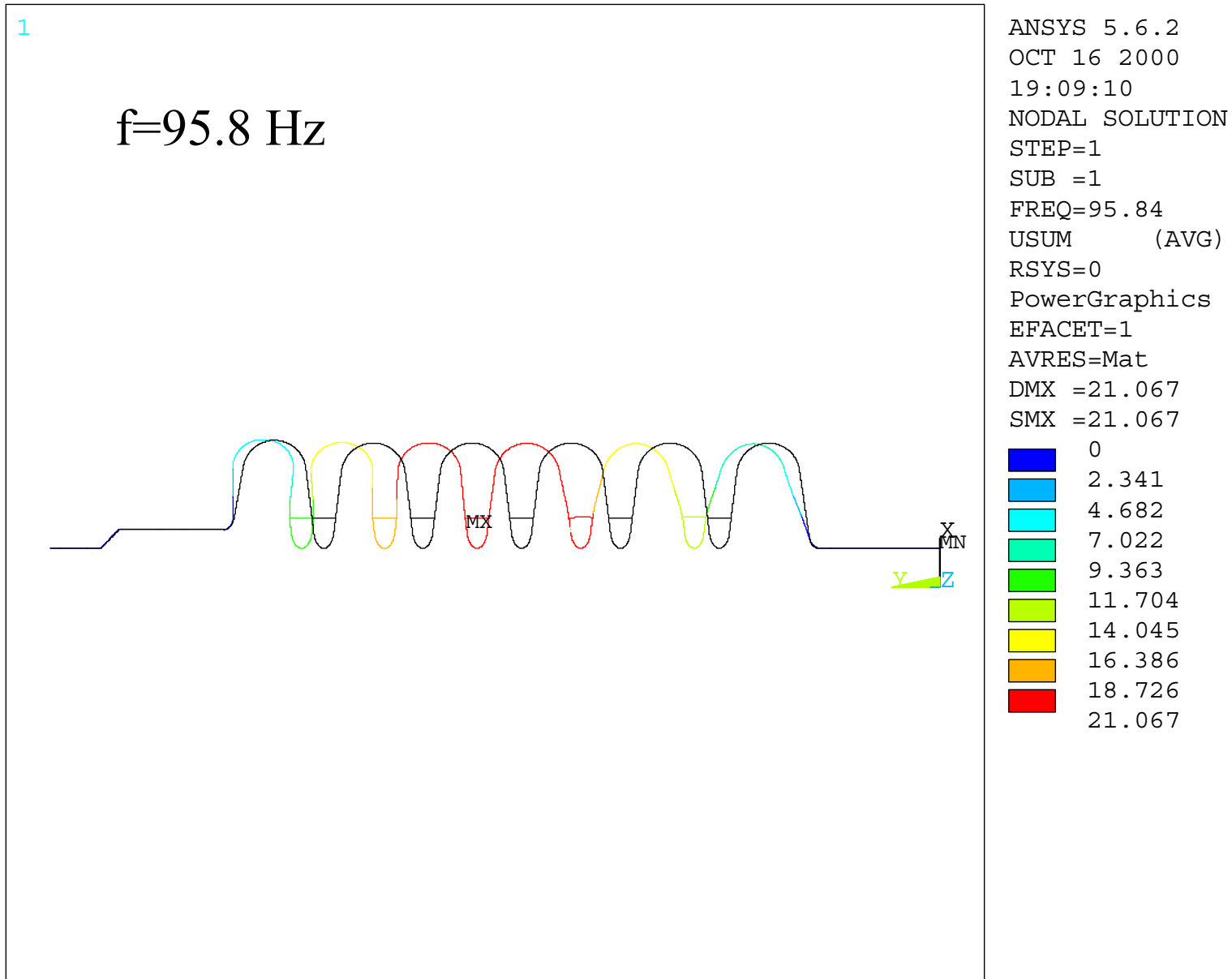
Hz	Rcoupling	Zcoupling
95.8	0.0001	1.0000
190.7	0.0003	0.0088
282.9	0.0004	0.3334
369.6	0.0004	0.0263
444.4	0.0007	0.2038
491.7	0.0003	0.0431
685.4	0.0004	0.0496
740.6	0.0004	0.0104
798.9	0.0005	0.0425
849.3	0.0006	0.0028
883.7	0.0002	0.0141
	(...)	
5946.4	0.0239	0.0374
6029.8	0.2615	0.0176
6041.5	0.4489	0.0473
6080.6	1.0000	0.0015
6132.0	0.4084	0.0013
6170.3	0.1064	0.0035
6189.1	0.1039	0.0395
6214.6	0.1130	0.0012
6922.9	0.0610	0.0090



Same kind of modes and same peculiarities



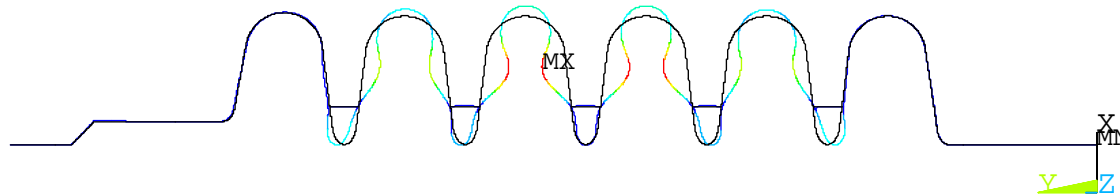
Bellow mode shape



Cell Shape

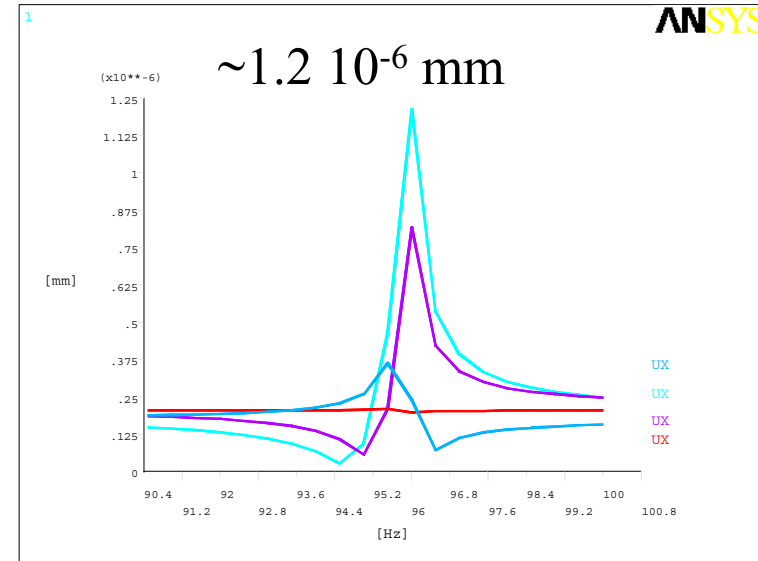
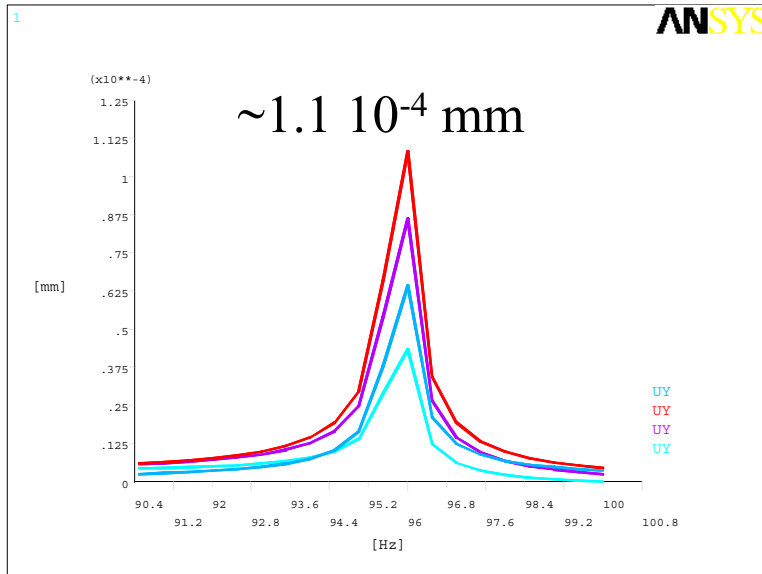
1

$f=1953$ Hz



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19:10:20
NODAL SOLUTION
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FREQ=1953
USUM (AVG)
RSYS=0
PowerGraphics
EFACET=1
AVRES=Mat
DMX =41.681
SMX =41.681
0
4.631
9.263
13.894
18.525
23.156
27.788
32.419
37.05
41.681

Stiffened Cavity Analysis

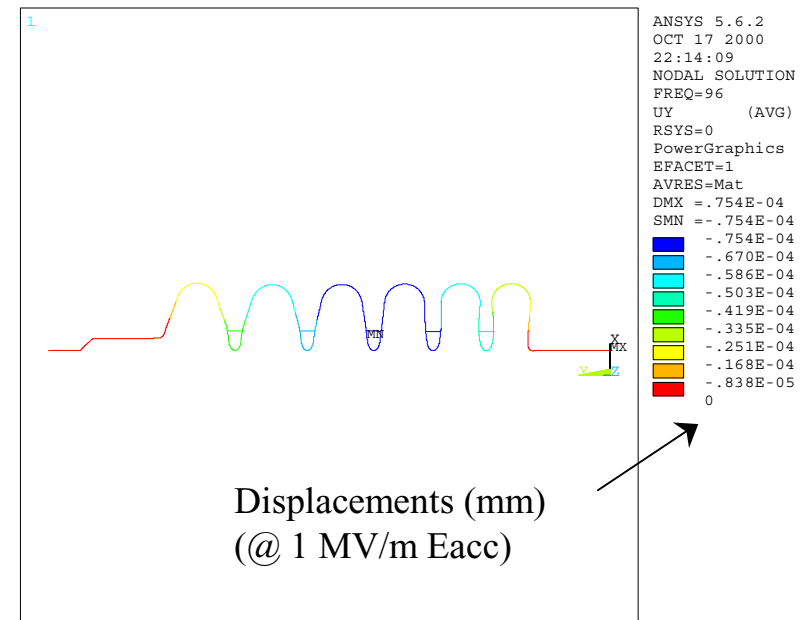


Deformations and frequency shifts are lower by $\sim 25\%$

$$\Delta\phi = 0.8^\circ @ 95\text{Hz}$$

$$\Delta\phi = 4.5^\circ @ 95.8\text{Hz}$$

$$\Delta\phi = 1.5^\circ @ 97\text{Hz}$$



Conclusion and plans

- Results are preliminary and must be checked.
 - It is, however, a pessimistic scenario, since modes are always excited near the resonance condition, independently from the 60 Hz “clock”
 - Only material damping
- Lorentz contribute **in full resonance condition** looks “small” with respect to static Lorentz contribution and:
 - Real system damping has to be considered
 - Transversal modes have not been checked but due to the axi-simmetry of the Lorentz load they should be weakly coupled and also Slater contributions should be self compensating
 - Transient effects could be significant during the pulsing and have to be studied (overelongation)
- Different harmonic loads could be checked:
 - Uniform pressure
 - Concentrated loads on the supports
- Prestress effects must be considered in the calculation for the real system
 - Change the frequency and the damping