(ADAPTIVE) GROVER FIXED-POINT SEARCH FOR BINARY OPTIMIZATION PROBLEMS

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ABSTRACT. *to be completed later...*

1. INTRODUCTION

to be completed later...

Organization of the paper:

to be completed later...

2. (POLYNOMIAL UNCONSTRAINED) BINARY OPTIMIZATION AND QUANTUM DICTIONARIES

For the rest of the paper, $\mathbb{Z}_2^n$ denotes the space of length-n bitstrings. Given a function $f : \mathbb{Z}_2^n \to \mathbb{R}$, the associated (Unconstrained) Binary Optimization problem is the task of finding an element $x \in \mathbb{Z}_2^n$ such that $f(x)$ is maximal. Note that every binary function is polynomial, which can be seen by simple dimension count.

Many interesting Binary Optimization problems, such as finding maximal graph cuts or the Max 2-SAT problems are quadratic, and most of the contemporary research centers around Quadratic Unconstrained Binary Optimization (QUBO) problems.

The first main contribution of the paper is an oracle design for QUBO problems. More concretely, we construct encoding operators of *quantum dictionaries*, as introduced in [1]. Such oracles have applications, for example, in Grover type algorithms and threshold QAOA [3]. While such designs have already existed, cf. [2], ours has better circuit depth, gate count, and CNOT count. Thus, they are simultaneously faster and more noise-resistant.

Briefly, the quantum dictionary, corresponding to a function (thought of as a classical dictionary), $F : \mathrm{dom}(F) \to \mathbb{Z}_2^d$, where $\mathrm{dom}(F) \subseteq \mathbb{Z}_2^n$, is the following quantum state on $n+d$ qubits:

$$|\mathrm{QDICT}(F)\rangle := \frac{1}{\sqrt{|\mathrm{dom}(F)|}} \sum_{x \in \mathrm{dom}(F)} |x\rangle_n |F(x)\rangle_d .$$

An integer-valued function $f : \mathbb{Z}_n^2 \to \mathbb{Z}$ canonically determines a quantum dictionary via first defining $F(x)$ to be the digits of $f(x)$, then setting, by a slight abuse of notation, $|\mathrm{QDICT}(f)\rangle = |\mathrm{QDICT}(F)\rangle$. Let us handle signs via the "Two's complement" convention, in particular, a binary

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number $y_0 y_1 \dots y_{d-1}$ is negative exactly when $y_0 = 1$. In fact, every quantum dictionary can be realized in this a way.

We construct the above-mentioned encoding oracles in two steps. First, we outline a modified version of the encoding operator given in [2] that is convenient to encode quadratic polynomials. Then we show that quadratic polynomial can be expressed in a basis of functions that can be more efficiently encoded.

2.1. **Quadratic encoder.** Let $I \subseteq \{1,2,\dots,n-1\}$ and $x_I := x_{i_1} x_{i_2} \cdots x_{i_j}$ be an arbitrary monomial and consider a quantum circuit with $n+d$ qubits. Following [2], we construct an oracle that sends $|x\rangle_n |0\rangle_d$ to $|x\rangle_n |x_I\rangle_d$, for any $x \in \mathbb{Z}_2^n$.

Let us make two definitions: Let QFT_d be the Quantum Fourier Transform on m qubits, that is for any $-2^{d-1} \leqslant y < 2^{d-1}$, we have

$$\mathrm{QFT}_d |y\rangle_d = 2^{-\frac{m}{2}} \sum_{z=-2^{d-1}}^{2^{d-1}-1} e^{\frac{2\pi yz}{2^d} i} |z\rangle_d .$$

Then

$$\mathrm{QFT}_d^\dagger |z\rangle_d = 2^{-\frac{d}{2}} \sum_{y'=-2^{d-1}}^{2^{d-1}-1} e^{-\frac{2\pi y'z}{2^d} i} |y'\rangle_d .$$

Now let $\mathscr{P}_d(k)$ be the following m-qubit gate

$$\begin{aligned}
|z_0\rangle &\text{ —— } \boxed{\mathrm{PHASE}(\pi k)} \text{ —— } e^{\frac{2\pi k z_0 2^{d-1}}{2^d} i} |z_0\rangle \\
&\vdots \\
|z_j\rangle &\text{ —— } \boxed{\mathrm{PHASE}\left(\frac{2\pi k}{2^{j+1}}\right)} \text{ —— } e^{\frac{2\pi z_j 2^{m-j-1}}{2^d}} |z_j\rangle \\
&\vdots \\
|z_{m-1}\rangle &\text{ —— } \boxed{\mathrm{PHASE}\left(\frac{2\pi k}{2^d}\right)} \text{ —— } e^{\frac{2\pi k z_{m-1}}{2^d} i} |z_d\rangle
\end{aligned}$$

Thus $\mathscr{P}_d(k) |z\rangle_d = e^{\frac{2\pi kz}{2^d} i} |z\rangle_d$.

Now we can prove a well-known lemma. citation needed

Lemma 2.1. *For any $-2^{d-1} \leqslant y < 2^{d-1}$ and $k \in \mathbb{Z}$ we have*

$$\mathrm{QFT}_d^\dagger \circ \mathscr{P}_d(k) \circ \mathrm{QFT}_d |y\rangle_d = |y + k \mod 2^{d-1}\rangle .$$

Proof. First we compute

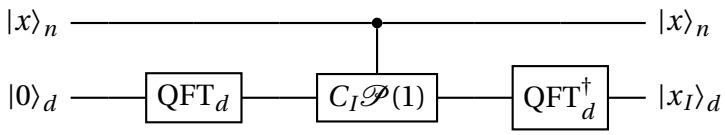
$$\begin{aligned}
\mathscr{P}_d(k) \circ \mathrm{QFT}_d \, |y\rangle_d &= \mathscr{P}(k) \left(2^{-\frac{m}{2}} \sum_{z=-2^{d-1}}^{2^{d-1}-1} e^{\frac{2\pi yz}{2^d} i} \, |z\rangle_d \right) \\
&= 2^{-\frac{d}{2}} \sum_{z=-2^{d-1}}^{2^{d-1}-1} e^{\frac{2\pi yz}{2^d} i} \mathscr{P}(k) \, |z\rangle_d \\
&= 2^{-\frac{d}{2}} \sum_{z=-2^{d-1}}^{2^{d-1}-1} e^{\frac{2\pi (y+k)z}{2^d} i} \, |z\rangle_d \\
&= \mathrm{QFT}_d \, |y + k \mod 2^{d-1}\rangle ,
\end{aligned}$$

hence

$$\mathrm{QFT}_d^\dagger \circ \mathscr{P}(k) \circ \mathrm{QFT}_d \, |y\rangle_d = |y + k \mod 2^{d-1}\rangle .$$

□

By Lemma 2.1 it is immediate that



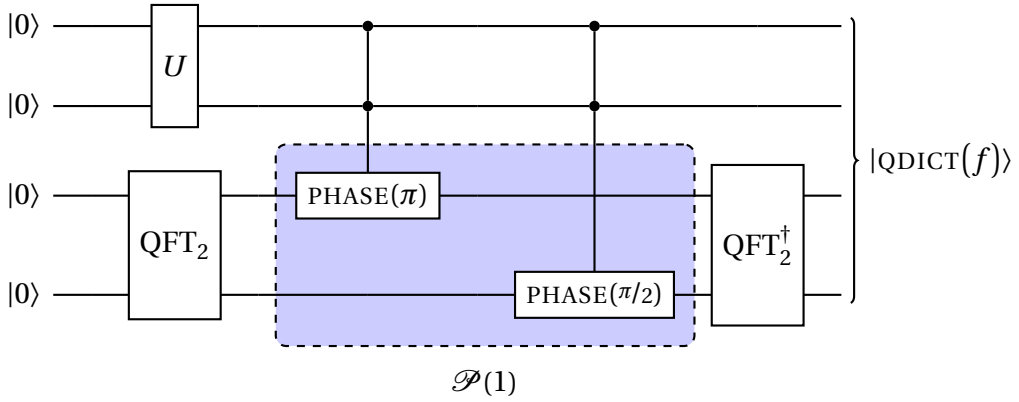
is the desired oracle, where C_I means control by the qubits $i_1, i_2, \ldots i_j$.

Finally, to create the full quantum dictionary, $|\mathrm{QDICT}(f)\rangle$, we need to pre-compose an oracle, call U, for which we have

$$U \, |0\rangle_n = \frac{1}{\sqrt{|\mathrm{dom}(f)|}} \sum_{x \in \mathrm{dom}(f)} |x\rangle_n .$$

If $\mathrm{dom}(f) = \mathbb{Z}_2^n$, then $U = H^{\otimes n}$.

Example 2.2. *Let $n = d = 2$ and $f(x) = x_0 x_1$. Now the encoder oracle takes the form*

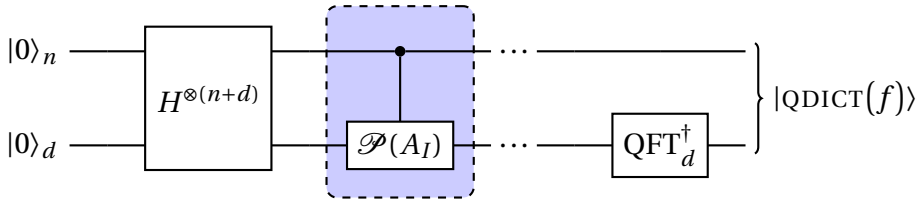


For the rest of the paper we assume that $\mathrm{dom}(f) = \mathbb{Z}_2^n$ and use $U = H^{\otimes n}$. Furthermore, since $\mathrm{QFT}_d |0\rangle_d = H^{\otimes m} |0\rangle_d$, we can replace QFT_d with $H^{\otimes d}$ in the oracle, as the latter has depth 1 and uses only single-qubit gates.

Let $\mathbb{P}_n$ be the power set of $\{0, 1, \ldots, n-1\}$. Now, for an arbitrary polynomial,

$$f(x) = \sum_{I \in \mathbb{P}_n} A_I x^I,$$

and $d \in \mathbb{N}_+$ large enough so that all values of f can be digitized on m bits, we have that



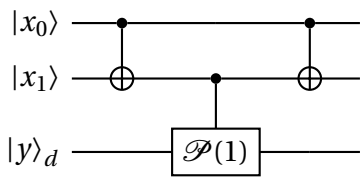
one for each $I = \{i_1, i_2, \ldots, i_{|I|}\} \in \mathbb{P}_n$, controlled by qubits $x_{i_1}, x_{i_2}, \ldots, x_{i_{|I|}}$

2.2. **New basis for quadratic polynomials.** We motivate the idea of the new basis by outlining it in the $n = 2$ case.

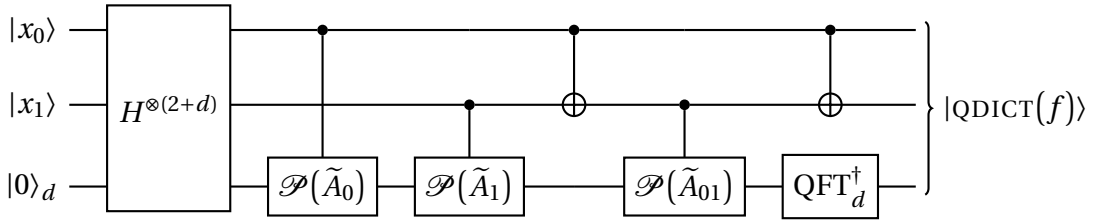
Note first that, since $x_i^2 = x_i$ for binary variables, we have that $x_0 x_1 = \frac{1}{2}\big(x_0 + x_1 - (x_0 - x_1)^2\big)$. Now $(x_0 - x_1)^2$ is also a binary variable, in fact, $(x_0 - x_1)^2 = x_0$ XOR x_1. Now let f be a generic polynomial, $f(x_0, x_1) = A_\emptyset + A_0 x_0 + A_1 x_1 + A_{01} x_0 x_1$. Since we are interested in finding the maximum of f, we can assume, without any loss of generality, that $f(0,0) = A_\emptyset = 0$. By generic, we mean that $0 \notin \{A_0, A_1, A_{01}\}$. Using d digits, we need $2d$ CNOT gates for each linear terms and $6d$ CNOT gates for the quadratic term, thus a total of $10d$ CNOT gates (not counting the CNOT gates in $\mathrm{QFT}_d^\dagger$). However, we can rewrite f as

$$\begin{aligned} f(x_0, x_1) &= A_0 x_0 + A_1 x_1 + A_{01} x_0 x_1 \\ &= \overbrace{\big(A_0 + \tfrac{1}{2} A_{01}\big)}^{\widetilde{A}_0 :=} x_0 + \overbrace{\big(A_1 + \tfrac{1}{2} A_{01}\big)}^{\widetilde{A}_1 :=} x_1 + \overbrace{\big(-\tfrac{1}{2} A_{01}\big)}^{\widetilde{A}_{01} :=} (x_0 - x_1)^2, \end{aligned}$$

and note that the last term can be implemented as



which now has CNOT count only $2+2d$. Thus the whole oracle (for arbitrary f) can be realized as



making the new CNOT count for the whole oracle to be (at most) $2+6d$, again, not counting the CNOT gates in $\mathrm{QFT}_d^\dagger$. In fact, the only time the two counts equal is when $A_0 = A_1 = 0$, $A_{01} \neq 0$, and $d = 1$, that is, when middle of the oracle is a single doubly-controlled phase gate, in which case this construction recovers the well-known one; cf. [4, Figure 4.8].

For a general, n-bit, quadratic polynomial, given by a symmetric, real, n-by-n matrix, Q via

$$f(x_0, x_1, \ldots, x_{n-1}) = x^T Q x = \sum_{i=0}^{n-1} Q_{ii} x_i + 2 \sum_{i=0}^{n-2} \sum_{j=i+1}^{n-1} Q_{ij} x_i x_j,$$

we have that if $q_i = \sum_{j=0}^{n-1} Q_{ij}$ is the sum of the i^{th} row, then

$$f(x_0, x_1, \ldots, x_{n-1}) = \sum_{i=0}^{n-1} q_i x_i - \sum_{i=0}^{n-2} \sum_{j=i+1}^{n-1} Q_{ij} \big(x_i - x_j\big)^2.$$

Since 1-controlled $\mathscr{P}$-gates require $2d$ and 2-controlled $\mathscr{P}$-gates require $8d$ CNOT gates, the oracle design of [2] requires $n(2d) + \frac{n(n-1)}{2}(8d) = \big(4n^2 - 2n\big)d$ CNOT gates. Our construction however requires only $2+2d$ CNOT gates for each quadratic term, so the total CNOT count is $n(2d) + \frac{n(n-1)}{2}(2+2d) = (d+1)n^2 + (d-1)n$, which is an approximately threefold improvement in the generic case.

Remark 2.3. *In the case a single, quadratic monomial and $d = 1$, there is no advantage; in fact, in this case, our construction is just the well-known 2-controlled phase gate from* [4, Figure 4.8]. *Using this observation and a bit more work, one can also show that our construction is never worse (in terms of gate count, CNOT count, or gate circuit depth), than that of* [2].

A further generalization of this construction to higher degree polynomials is also currently being prepared by the authors.

3. Grover Fixed-point search for QUBO

Grover Fixed-Point Search (GFPS) [5] is a variant of Grover's search algorithm that retains the original version query complexity while does not suffer from the soufflé problem (more

on these below). In this section we introduce the algorithm, show how our oracle design from Section 2.2 can be used to implement GFPS for QUBO problems, and argue that this method is better suited for an adaptive optimizer algorithm, then the original.

As opposed to Grover's algorithm, where the only input is the set of *good* (or *target*) configurations, $T \subset \mathbb{Z}_n^2$, the GFPS algorithm requires an additional one, that can be chosen to be either the target probability/amplitude or the query complexity. The target probability is the probability of finding the system in a good state after running the circuit. The price of this flexibility (and of the elimination of half of the soufflé problem) is that on needs to implement not only a pair of oracles, but two families of them, call $S_s(\alpha)$ and $S_t\big(\beta\big)$ (where the subscripts refer to the *start* and *target* states, and α, β are real parameters), with the following properties: let fix gauge so that, for some $\lambda \in (0,1)$, we can write

$$|s\rangle = \sqrt{\lambda}\,|t\rangle + \sqrt{1-\lambda}\,|\overline{t}\rangle\,,$$
$$|t\rangle = \sqrt{\lambda}\,|s\rangle + \sqrt{1-\lambda}\,|\overline{s}\rangle\,,$$

where $\langle t|\overline{t}\rangle = \langle s|\overline{s}\rangle = 0$. Then

$$S_s(\alpha)\big(A\,|s\rangle + B\,|\overline{s}\rangle\big) = e^{i\alpha}A\,|s\rangle + B\,|\overline{s}\rangle\,,$$
$$S_t\big(\beta\big)\big(D\,|t\rangle + C\,|\overline{t}\rangle\big) = e^{i\beta}C\,|s\rangle + D\,|\overline{t}\rangle\,.$$

Let $G\big(\alpha,\beta\big) := S_s(\alpha)S_t\big(\beta\big)$. Once in possession of such oracles and a target probability $P \in (0,1)$, the main result of [5] can be summarized as follows: for any $\lambda' \in (0,\lambda)$ large enough, there exists $l = l(P,\lambda) \in \mathbb{N}_+$ and $\delta = \delta_{P,\lambda} \in (0,1)$, such that if

$$\forall\, j \in \{1,\dots,l\} : \alpha_j = \beta_{l-j+1} = 2\cot^{-1}\left(\tan\left(\frac{2\pi j}{2l+1}\right)\tanh\left(\frac{\cosh^{-1}\left(\frac{1}{\delta}\right)}{2l+1}\right)\right),$$

then

$$P_{\text{success}} := \Big|\big\langle t\big|G\big(\alpha_1,\beta_1\big)\cdots G\big(\alpha_1,\beta_1\big)\big|s\big\rangle\Big|^2 \geqslant P.$$

Moreover, l, δ can be explicitly computed (see next section).

We implement $S_s(\alpha)$ and $S_t\big(\beta\big)$ in straightforward ways. If U_s is the state preparation oracle, that is, $|s\rangle = U_s\,|0\rangle$, and $\mathrm{MCP}_n(\alpha)$ is the $(n-1)$-controlled phase gate on n qubits, then

$$S_s(\alpha) = U_s\mathrm{MCP}_n(\alpha)U_s^\dagger,$$

works.

In general, the implementation of $S_t\big(\beta\big)$ is case specific. In the case of QUBO problems, we can use constructions of Sections 2.1 and 2.2, to get an oracle, U_t on $n+1$ qubits, such that

for all $x \in \mathbb{Z}_2^n$ and $y \in \mathbb{Z}_2$, we have

$$U|x\rangle|y\rangle = \begin{cases} |x\rangle|y \oplus 1\rangle, & \text{if } x \in T, \\ |x\rangle|y\rangle, & \text{if } x \notin T. \end{cases}$$

Then

$$S_t(\beta) = U_t \mathrm{PHASE}_{n+1}(\alpha) U_t,$$

works, if the $(n+1)^{\mathrm{st}}$ qubits is a clean ancilla kept in $|0\rangle$ before and after the usage of the oracle.

Everything below is old stuff that I might or might not want to include.

Input: A symmetric, integer-valued, n-by-n matrix, Q and a constant $c \in \mathbb{Z}$, or, equivalently, a quadratic function on $x \in \{0,1\}^n$ given by

$$f(x) := x^T Q x + c. \tag{3.1}$$

(Note that since $x_i^2 = x_i$, we can move linear terms into the diagonal of Q.)

Output: An estimate for the value

$$M := \max\left(\left\{ f(x) \;\middle|\; x \in \{0,1\}^n \right\}\right).$$

Example 3.1 (Maximal Graph Cuts)**.** *Given a simple, undirected graph, $G = (V,E)$, let Q be its graph Laplacian, defined as*

$$Q_{i,j} = \begin{cases} \deg(v_i), & \text{if } i = j, \\ -1, & \text{if } \{v_i, v_j\} \in E, \\ 0, & \text{otherwise,} \end{cases}$$

$b = 0$ and $c = 0$. Then $V = V^+ \coprod V^-$ is a maximal exactly when $\mathrm{MaxCut}(G) = f(x) = M$, where $x \in \{0,1\}^n$ is defined as $x_i = 1$ if $v_i \in V^+$ and zero otherwise.

The Edwards–Erdős bound yields

$$\mathrm{MaxCut}(G) \geqslant B_G := \begin{cases} \frac{2|V|+|E|-1}{4}, & \text{if (we know that) } G \text{ is connected,} \\ \frac{|V|}{2} + \sqrt{\frac{|V|}{8} + \frac{1}{64}} - \frac{1}{8}, & \text{otherwise.} \end{cases}$$

4. The oracles:

An element $x = (x_1, x_2, \ldots, x_n) \in \{0,1\}^n$ is also regarded as a binary number via $x \sim \overline{x_1 x_2 \ldots x_n} := \sum_i x_i 2^{n-i}$ and as an element of the computational basis via

$$|x\rangle_n := |x_1\rangle \ldots |x_{n-1}\rangle |x_n\rangle,$$

Given a function as in equation (3.1), let us pick $m \gg \log_2(M)$ (in fact, $m = \lceil \log_2(\mathrm{tr}(Q)) \rceil + 1$ works for our purposes). We use the binary 2s complement convention when digitizing integers and we with that in mind, we construct a oracle on $(n+m)$-qubits, U_f, so that

$$U_f |x\rangle_n |y\rangle_d = |x\rangle_n |y - f(x)\rangle_d.$$

Note that the $(n+1)^{\text{th}}$ register of $U_f |x\rangle_n |y\rangle_d$ is $|1\rangle$ exactly when $y < f(x)$.

5. Application to Grover fixed-point search and state preparation

Fix $\delta \in (0,1)$ and y. Let $\lambda := \frac{|C_y|}{2^n}$, where $C_y := \{x \in \{0,1\}^n \,|\, f(x) \geqslant y\}$. Finally let $l := \left\lceil \frac{\log_2\left(\frac{2}{\delta}\right)}{2\sqrt{\lambda}} - \frac{1}{2} \right\rceil$. Then, following [5], we can construct a Quantum circuit (using U_f from the previous section), that results in a state $S_l |0\rangle_n |y\rangle_d$ with the following significance: When the first n qubits are measured in the computational basis, then

$$P\big(x \in C_y\big) = \sum_{x \in C_y} |\langle x| S_l |0\rangle|^2 \geqslant 1 - \delta^2.$$

Let us make the following definitions:

$$\begin{aligned} U_S &:= H^{\otimes n} \otimes \mathbb{1}^{\otimes m}, \\ R_0(\alpha) &:= \mathbb{1}^{\otimes(n+m)} + \left(1 - e^{i\alpha}\right) |0\rangle_n \langle 0|_n \otimes \mathbb{1}^{\otimes(1+m)}, \\ R_T\big(\beta\big) &:= U_f^\dagger P_{n+1}\big(\beta\big) U_f, \\ G\big(\alpha, \beta\big) &:= -U_S R_0(\alpha) U_S^\dagger R_T\big(\beta\big). \end{aligned}$$

Let $\big(\boldsymbol{\alpha}, \boldsymbol{\beta}\big) = \big(\alpha_1, \beta_1, \cdots, \alpha_l, \beta_l\big)$ be given by

$$\forall\, j \in \{1, \ldots, l\}: \quad \alpha_j := -\beta_{l-j+1} = 2\cot^{-1}\left(\tan\left(\tfrac{2\pi j}{2l+1}\right)\sqrt{1-\gamma^2}\right),$$

where $\gamma := \big(T_{1/(2l+1)}\big(\delta^{-1}\big)\big)^{-1}$ and let

$$S_l\big(\boldsymbol{\alpha}, \boldsymbol{\beta}\big) = G\big(\alpha_l, \beta_l\big) G\big(\alpha_{l-1}, \beta_{l-1}\big) \cdots G\big(\alpha_1, \beta_1\big) U_S.$$

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