Some integrals with Gaussian wavepackets

1. Definition

The normalized Gaussian wavepacket is defined as:

$$G = \left(\frac{a}{\pi}\right)^{1/4} \exp\left(-\frac{a}{2}(x - x_0)^2 + ip_0(x - x_0)\right). \tag{1}$$

Here, a is the width parameter, x_0 is the center of wavepacket, p_0 is the momentum of wavepacket.

2. Overlap integral

$$S_{12} = \int G_{1}^{*} G_{2} dx = \left(\frac{a}{\pi}\right)^{1/2} \int \exp\left(-\frac{a}{2}(x - x_{0,1})^{2} - \frac{a}{2}(x - x_{0,2})^{2} - ip_{0,1}(x - x_{0,1}) + ip_{0,2}(x - x_{0,2})\right) dx =$$

$$= \left(\frac{a}{\pi}\right)^{1/2} \int \exp\left(\left[-ax^{2} + a(x_{0,1} + x_{0,2})x - \frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right] + ix(p_{0,2} - p_{0,1}) dx =$$

$$= \left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \int \exp\left(-ax^{2} + \left[a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})\right]x\right) dx =$$

$$= \exp\left(-\frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \exp\left(\frac{\left[a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})\right]^{2}}{4a}\right) =$$

$$= \exp\left(-\frac{a}{4}(x_{0,1} - x_{0,2})^{2} - \frac{1}{4a}(p_{0,2} - p_{0,1})^{2}\right) \exp\left(i\left(x_{0,1} - x_{0,2}) + \frac{(p_{0,1} + p_{0,2})}{2}\right)\right)$$

Here, we have used the integral

$$\int \exp(-ax^2 + bx + c)dx = \left(\frac{\pi}{a}\right)^{1/2} \exp\left(\frac{b^2}{4a} + c\right)$$
(3)

3. Position expectation value

$$\mu_{12} = \int G_{1}^{*}xG_{2}dx = \left(\frac{a}{\pi}\right)^{1/2} \int x \exp\left(-\frac{a}{2}(x-x_{0,1})^{2} - \frac{a}{2}(x-x_{0,2})^{2} - ip_{0,1}(x-x_{0,1}) + ip_{0,2}(x-x_{0,2})\right) dx =$$

$$= \left(\frac{a}{\pi}\right)^{1/2} \int x \exp\left(\left[-ax^{2} + a(x_{0,1} + x_{0,2})x - \frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right] + ix(p_{0,2} - p_{0,1}) dx =$$

$$= \left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \int x \exp\left(-ax^{2} + \left[a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})\right]x\right) dx =$$

$$= \left(\frac{b}{2a}\right) \exp\left(-\frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \exp\left(\frac{\left[a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})\right]^{2}}{4a}\right) =$$

$$= \left(\frac{b}{2a}\right) S_{12} = \frac{a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})}{2a} S_{12} = \left[\frac{(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})}{2a}\right] S_{12}$$

$$\int x \exp(-ax^2 + bx + c)dx = \int x \exp\left(-a\left(x^2 - 2\frac{b}{2a}x + \left(\frac{b}{2a}\right)^2 - \left(\frac{b}{2a}\right)^2\right) + c\right)dx =$$

$$= \int x \exp\left(-a\left(x - \frac{b}{2a}\right)^2 + \left(\frac{b^2}{4a} + c\right)\right)dx = \exp\left(\frac{b^2}{4a} + c\right)\int x \exp\left(-a\left(x - \frac{b}{2a}\right)^2\right)dx =$$

$$= |y = x - \frac{b}{2a}, dx = dy| = \exp\left(\frac{b^2}{4a} + c\right)\int \left(y + \frac{b}{2a}\right) \exp(-ay^2)dy =$$

$$= \exp\left(\frac{b^2}{4a} + c\right)\left[\frac{b}{2a}\left(\frac{\pi}{a}\right)^{1/2}\right]$$

4. Derivative coupling integral

$$D_{12} = \int G_{1}^{*} \nabla G_{2} dx = \left(\frac{a}{\pi}\right)^{1/2} \int \left[-a(x - x_{0,2}) + ip_{0,2}\right] \exp\left(-\frac{a}{2}(x - x_{0,1})^{2} - \frac{a}{2}(x - x_{0,2})^{2} - ip_{0,1}(x - x_{0,1}) + ip_{0,2}(x - x_{0,2})\right) dx =$$

$$= (ax_{0,2} + ip_{0,2})S - a\left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \int x \exp\left(-ax^{2} + \left[a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})\right]x\right) dx =$$

$$= (ax_{0,2} + ip_{0,2})S - a\left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \frac{b}{2a} \exp\left(\frac{b^{2}}{4a}\right) \left(\frac{\pi}{a}\right)^{1/2} =$$

$$= (ax_{0,2} + ip_{0,2})S - \frac{b}{2}S = \left\{(ax_{0,2} + ip_{0,2}) - \frac{\left[a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})\right]}{2}\right\}S = \frac{\left[a(x_{0,2} - x_{0,1}) + i(p_{0,2} + p_{0,1})\right]}{2}S$$

Check:

$$D_{21} = \int (G_2^* \nabla G_1) dx = \frac{\left[a(x_{0,1} - x_{0,2}) + i(p_{0,1} + p_{0,2})\right]}{2} S \Rightarrow$$

$$\left(\int (G_2^* \nabla G_1) dx\right)^* = \frac{\left[-a(x_{0,2} - x_{0,1}) - i(p_{0,1} + p_{0,2})\right]}{2} S = -D_{12}$$
So:
$$D_{12}^* = -D_{21}.$$
(5)

5. Kinetic energy integral

$$g = \exp\left(-\frac{a}{2}(x - x_0)^2 + ip_0(x - x_0)\right)$$

$$g'_x = (-a(x - x_0) + ip_0)G$$
(8)

(6)

$$g_{x}^{"} = -ag + (-a(x - x_{0}) + ip_{0})g' = -ag + (-a(x - x_{0}) + ip_{0})^{2}g =$$

$$= (a^{2}(x - x_{0})^{2} - 2iap_{0}(x - x_{0}) - a - p_{0}^{2})g$$
(9)

$$\int G_1^* \nabla^2 G_2 dx = \left(\frac{a}{\pi}\right)^{1/2} \int \left(a^2 \left(x - x_{0,2}\right)^2 - 2iap_{0,2} \left(x - x_{0,2}\right) - a - p_{0,2}^2\right) g_1^* g_2 dx =$$

$$= \left(-a - p_{0,2}^2\right) S_{12} + I_2 - I_1$$
(10)

$$I_{2} = \left(\frac{a}{\pi}\right)^{1/2} \int a^{2}(x - x_{0,2})^{2} g_{1}^{*} g_{2} dx =$$

$$= a^{2} \left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right)$$

$$\int (x - x_{0,2})^{2} \exp\left(-ax^{2} + [a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})]x) dx =$$

$$= C\int (x - x_{0,2})^{2} \exp\left(-ax^{2} + bx\right) dx = C\int (x - x_{0,2})^{2} \exp\left(-a\left(x - \frac{b}{2a}\right)^{2} + \frac{b^{2}}{4a}\right) dx = |x| = x - \frac{b}{2a}|x|$$

$$= C'\int \left(x' + \frac{b}{2a} - x_{0,2}\right)^{2} \exp\left(-ax'^{2}\right) dx' = C'\left[\frac{1}{2}\left(\frac{\pi}{a^{3}}\right)^{1/2} + \left(\frac{b}{2a} - x_{0,2}\right)^{2}\left(\frac{\pi}{a}\right)^{1/2}\right] =$$

$$= a^{2}\left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \exp\left(\frac{b^{2}}{4a}\left[\frac{1}{2}\left(\frac{\pi}{a^{3}}\right)^{1/2} + \left(\frac{b}{2a} - x_{0,2}\right)^{2}\left(\frac{\pi}{a}\right)^{1/2}\right] =$$

$$= a^{2} \exp\left(-\frac{a}{2}(x_{0,1}^{2} + x_{0,2}^{2}) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}) + \frac{b^{2}}{4a}\left[\frac{1}{2a} + \left(\frac{b}{2a} - x_{0,2}\right)^{2}\right] =$$

$$= S_{12}a^{2}\left[\frac{1}{2a} + \left(\frac{b}{2a} - x_{0,2}\right)^{2}\right] = S_{12}\left[\frac{1a}{2} + \left(\frac{b - 2ax_{0,2}}{2}\right)^{2}\right]$$
(11)

$$C = a^{2} \left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}\left(x_{0,1}^{2} + x_{0,2}^{2}\right) + i\left(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}\right)\right)$$
(12)

$$C' = C \exp\left(\frac{b^2}{4a}\right) \tag{13}$$

$$b = \left[a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1}) \right]$$
(14)

$$I_{1} = 2iap_{0,2} \left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}\left(x_{0,1}^{2} + x_{0,2}^{2}\right) + i\left(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}\right)\right)$$

$$\int (x - x_{0,2}) \exp\left(-ax^{2} + \left[a\left(x_{0,1} + x_{0,2}\right) + i\left(p_{0,2} - p_{0,1}\right)\right]x\right) dx =$$

$$= C\int (x - x_{0,2}) \exp\left(-ax^{2} + bx\right) dx = C\int (x - x_{0,2}) \exp\left(-a\left(x - \frac{b}{2a}\right)^{2} + \frac{b^{2}}{4a}\right) dx =$$

$$= C'\int (x - x_{0,2}) \exp\left(-a\left(x - \frac{b}{2a}\right)^{2}\right) dx = x' = x - \frac{b}{2a} = C'\int (x' + \frac{b}{2a} - x_{0,2}) \exp\left(-ax'^{2}\right) dx' =$$

$$= C'\left(\frac{b}{2a} - x_{0,2}\right)\left(\frac{\pi}{a}\right)^{1/2} = 2iap_{0,2}\left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}\left(x_{0,1}^{2} + x_{0,2}^{2}\right) + i\left(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}\right)\right) \exp\left(\frac{b^{2}}{4a}\right)\left(\frac{b}{2a} - x_{0,2}\right)\left(\frac{\pi}{a}\right)^{1/2} =$$

$$= 2iap_{0,2} \exp\left(-\frac{a}{2}\left(x_{0,1}^{2} + x_{0,2}^{2}\right) + i\left(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}\right) + \frac{b^{2}}{4a}\right)\left(\frac{b}{2a} - x_{0,2}\right) =$$

$$= S_{12}ip_{0,2}\left(b - 2ax_{0,2}\right)$$

$$C = 2iap_{0,2}\left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}\left(x_{0,1}^{2} + x_{0,2}^{2}\right) + i\left(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}\right) + i\left(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}\right)\right)$$

$$(16)$$

$$C' = C \exp\left(\frac{b^2}{4a}\right) \tag{17}$$

$$b = \left[a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1}) \right]$$
(18)

So:

$$\int G_1^* \nabla^2 G_2 dx = \left(-a - p_{0,2}^2\right) S_{12} + I_2 - I_1 =$$

$$= \left(-a - p_{0,2}^{2}\right)S_{12} + S_{12}\left[\frac{a}{2} + \left(\frac{b - 2ax_{0,2}}{2}\right)^{2}\right] - S_{12}ip_{0,2}(b - 2ax_{0,2}) =$$

$$= S_{12}\left\{-a - p_{0,2}^{2} + \frac{a}{2} + \frac{\left[a(x_{0,1} - x_{0,2}) + i(p_{0,2} - p_{0,1})\right]^{2}}{4} - ip_{0,2}\left[a(x_{0,1} - x_{0,2}) + i(p_{0,2} - p_{0,1})\right]\right\} =$$

$$= S_{12}\left\{-\frac{a}{2} - p_{0,2}^{2} + \frac{1}{4}a^{2}(x_{0,1} - x_{0,2})^{2} + \frac{1}{2}ai(x_{0,1} - x_{0,2})(p_{0,2} - p_{0,1}) - \frac{1}{4}(p_{0,2}^{2} - 2p_{0,2}p_{0,1} + p_{0,1}^{2})\right\} =$$

$$= S_{12}\left\{-\frac{a}{2} + a^{2}\left(\frac{x_{0,1} - x_{0,2}}{2}\right)^{2} - \left(\frac{p_{0,2} + p_{0,1}}{2}\right)^{2} - \frac{1}{2}ai(x_{0,1} - x_{0,2})(p_{0,2} + p_{0,1})\right\}$$

6. Summary

$$G = \left(\frac{a}{\pi}\right)^{1/4} \exp\left(-\frac{a}{2}(x - x_0)^2 + ip_0(x - x_0)\right),\tag{20}$$

$$S_{G,12} = \int G_1^* G_2 dx = \exp\left(-\frac{a}{4} \left(x_{0,1} - x_{0,2}\right)^2 - \frac{1}{4a} \left(p_{0,2} - p_{0,1}\right)^2\right) \exp\left(i\left(\left(x_{0,1} - x_{0,2}\right) \frac{\left(p_{0,1} + p_{0,2}\right)}{2}\right)\right) (21)$$

$$\mu_{G,12} = \int G_1^* x G_2 dx = \left[\frac{\left(x_{0,1} + x_{0,2} \right)}{2} + i \frac{\left(p_{0,2} - p_{0,1} \right)}{2a} \right] S_{G,12}$$

$$D_{G,12} = \int G_1^* \nabla G_2 dx = \frac{\left[a \left(x_{0,2} - x_{0,1} \right) + i \left(p_{0,2} + p_{0,1} \right) \right]}{2} S_{G,12}$$
(22)

$$T_{G,12} = \int G_1^* \nabla^2 G_2 dx = S_{G,12} \left\{ -\frac{a}{2} + a^2 \left(\frac{x_{0,1} - x_{0,2}}{2} \right)^2 - \left(\frac{p_{0,2} + p_{0,1}}{2} \right)^2 - \frac{1}{2} ai \left(x_{0,1} - x_{0,2} \right) \left(p_{0,2} + p_{0,1} \right) \right\}$$
(23)

6. Generalization

The form of the Gaussian, Eq. 1, is not general – the momentum term should have the \hbar denominator. In addition, the complex phase factor is often present. So the more appropriate form is

$$G = \left(\frac{a}{\pi}\right)^{1/4} \exp\left(-\frac{a}{2}(x - x_0)^2 + \frac{ip_0}{\hbar}(x - x_0) + \frac{i\gamma}{\hbar}\right)$$
 (24)

Let also denote $\alpha = \frac{a}{2}$, so:

$$G = \left(\frac{2\alpha}{\pi}\right)^{1/4} \exp\left(-\alpha(x-x_0)^2 + \frac{ip_0}{\hbar}(x-x_0) + \frac{i\gamma}{\hbar}\right)$$
 (25)

This form, for instance, is present in the paper of Makhov et al. [Makhov, D. V.; Glover, W. J.; Martinez, T. J.; Shalashilin, D. V. Ab Initio Multiple Cloning Algorithm for Quantum Nonadiabatic Molecular Dynamics. The Journal of Chemical Physics 2014, 141 (5), 54110.]

The results Eqs. 21-23 derived for the form Eq. 1 will transform to:

$$S_{G,12} = \int G_{1}^{*}G_{2}dx = \exp\left(-\frac{\alpha}{2}(x_{0,1} - x_{0,2})^{2} - \frac{1}{8\alpha\hbar^{2}}(p_{0,2} - p_{0,1})^{2} + i(x_{0,1} - x_{0,2})\frac{(p_{0,1} + p_{0,2})}{2\hbar} + \frac{i}{\hbar}(\gamma_{2} - \gamma_{1})\right)$$

$$(26)$$

$$\mu_{G,12} = \int G_{1}^{*}xG_{2}dx = \left[\frac{(x_{0,1} + x_{0,2})}{2} + i\frac{(p_{0,2} - p_{0,1})}{4\alpha\hbar}\right]S_{G,12}$$

$$D_{G,12} = \int G_{1}^{*}\nabla G_{2}dx = \left[\alpha(x_{0,2} - x_{0,1}) + i\frac{(p_{0,2} + p_{0,1})}{2\hbar}\right]S_{G,12}$$

$$T_{G,12} = \int G_{1}^{*}\nabla^{2}G_{2}dx = S_{G,12}\left\{-\alpha + \alpha^{2}(x_{0,1} - x_{0,2})^{2} - \frac{1}{\hbar^{2}}\left(\frac{p_{0,2} + p_{0,1}}{2}\right)^{2} - 2\alpha i(x_{0,1} - x_{0,2})\frac{(p_{0,2} + p_{0,1})}{2}\right\}$$

$$(28)$$

Introducing variables

$$\Delta X = x_{0,2} - x_{0,1}, \ \Delta P = p_{0,2} - p_{0,1}, \ \Delta \gamma = \gamma_2 - \gamma_1 \ \overline{P} = \frac{p_{0,1} + p_{0,2}}{2}, \ \overline{X} = \frac{x_{0,1} + x_{0,2}}{2}. \tag{29}$$

The above equations simplify to:

$$S_{G,12} = \int G_1^* G_2 dx = \exp\left(-\frac{\alpha}{2} (\Delta X)^2 - \frac{1}{8\alpha\hbar^2} (\Delta P)^2 - \frac{i\overline{P}\Delta X}{\hbar} + \frac{i\Delta\gamma}{\hbar}\right)$$
(30)

$$\mu_{G,12} = \int G_1^* x G_2 dx = \left[\overline{X} + \frac{\Delta P}{4\alpha \hbar} i \right] S_{G,12}$$

$$D_{G,12} = \int G_1^* \nabla G_2 dx = \left[\alpha (\Delta X) + \frac{i\overline{P}}{\hbar} \right] S_{G,12}$$
(31)

$$T_{G,12} = \int G_1^* \nabla^2 G_2 dx = S_{G,12} \left\{ -\alpha + \alpha^2 (\Delta X)^2 - \frac{\overline{P}^2}{\hbar^2} + 2\alpha i \Delta X \cdot \overline{P} \right\}$$
 (32)

The equations for the overlap and for the kinetic energy are consistent with those of Makhov et al. However, the equation for the derivative coupling is not – they attribute similar equation to a different type of the derivative coupling.

To show consistency of the overlap, consider

$$\begin{split} &\frac{i}{\hbar}\Bigg(p_{0,1}x_{0,1}-p_{0,2}x_{0,2}+\frac{x_{0,1}+x_{0,2}}{2}\Big(p_{0,2}-p_{0,1}\Big)\Bigg)=\\ &=\frac{i}{\hbar}\Bigg(p_{0,1}x_{0,1}-p_{0,2}x_{0,2}+\frac{x_{0,1}p_{0,2}}{2}-\frac{x_{0,1}p_{0,1}}{2}+\frac{x_{0,2}p_{0,2}}{2}-\frac{x_{0,2}p_{0,1}}{2}\Bigg)=\\ &=\frac{i}{\hbar}\Bigg(\frac{x_{0,1}p_{0,2}}{2}+\frac{x_{0,1}p_{0,1}}{2}-\frac{x_{0,2}p_{0,2}}{2}-\frac{x_{0,2}p_{0,1}}{2}\Bigg)=\frac{i}{\hbar}\bigg(\frac{\left(x_{0,1}-x_{0,2}\right)p_{0,2}}{2}+\frac{\left(x_{0,1}-x_{0,2}\right)p_{0,1}}{2}\bigg)=\\ &=-\frac{i\Delta X}{\hbar}\bigg(\frac{p_{0,2}+p_{0,1}}{2}\bigg)=-\frac{i\Delta X\cdot\overline{P}}{\hbar} \end{split}$$

So, their expression coincides with our term $-\frac{i\overline{P}\Delta X}{\hbar}$