

## Some integrals with Gaussian wavepackets

### 1. Definition

The normalized Gaussian wavepacket is defined as:

$$G = \left(\frac{a}{\pi}\right)^{1/4} \exp\left(-\frac{a}{2}(x-x_0)^2 + ip_0(x-x_0)\right). \quad (1)$$

Here,  $a$  is the width parameter,  $x_0$  is the center of wavepacket,  $p_0$  is the momentum of wavepacket.

### 2. Overlap integral

$$\begin{aligned} S_{12} &= \int G_1^* G_2 dx = \left(\frac{a}{\pi}\right)^{1/2} \int \exp\left(-\frac{a}{2}(x-x_{0,1})^2 - \frac{a}{2}(x-x_{0,2})^2 - ip_{0,1}(x-x_{0,1}) + ip_{0,2}(x-x_{0,2})\right) dx = \\ &= \left(\frac{a}{\pi}\right)^{1/2} \int \exp\left(\left[-ax^2 + a(x_{0,1}+x_{0,2})x - \frac{a}{2}(x_{0,1}^2+x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right] + ix(p_{0,2} - p_{0,1})\right) dx = \\ &= \left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^2+x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \int \exp(-ax^2 + [a(x_{0,1}+x_{0,2}) + i(p_{0,2} - p_{0,1})]x) dx = \\ &= \exp\left(-\frac{a}{2}(x_{0,1}^2+x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \exp\left(\frac{[a(x_{0,1}+x_{0,2}) + i(p_{0,2} - p_{0,1})]^2}{4a}\right) = \\ &= \exp\left(-\frac{a}{4}(x_{0,1}-x_{0,2})^2 - \frac{1}{4a}(p_{0,2} - p_{0,1})^2\right) \exp\left(i\left((x_{0,1}-x_{0,2})\frac{(p_{0,1}+p_{0,2})}{2}\right)\right) \end{aligned} \quad (2)$$

Here, we have used the integral

$$\int \exp(-ax^2 + bx + c) dx = \left(\frac{\pi}{a}\right)^{1/2} \exp\left(\frac{b^2}{4a} + c\right) \quad (3)$$

### 3. Derivative coupling integral

$$\begin{aligned} D_{12} &= \int G_1^* \nabla G_2 dx = \left(\frac{a}{\pi}\right)^{1/2} \int [-a(x-x_{0,2}) + ip_{0,2}] \exp\left(-\frac{a}{2}(x-x_{0,1})^2 - \frac{a}{2}(x-x_{0,2})^2 - ip_{0,1}(x-x_{0,1}) + ip_{0,2}(x-x_{0,2})\right) dx = \\ &= (ax_{0,2} + ip_{0,2})S - a\left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^2+x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \int x \exp(-ax^2 + [a(x_{0,1}+x_{0,2}) + i(p_{0,2} - p_{0,1})]x) dx = \\ &= (ax_{0,2} + ip_{0,2})S - a\left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^2+x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \frac{b}{2a} \exp\left(\frac{b^2}{4a}\right) \left(\frac{\pi}{a}\right)^{1/2} = \\ &= (ax_{0,2} + ip_{0,2})S - \frac{b}{2}S = \left\{(ax_{0,2} + ip_{0,2}) - \frac{[a(x_{0,1}+x_{0,2}) + i(p_{0,2} - p_{0,1})]}{2}\right\}S = \frac{[a(x_{0,2}-x_{0,1}) + i(p_{0,2} + p_{0,1})]}{2}S \end{aligned} \quad (4)$$

Check:

$$\begin{aligned} D_{21} &= \int (G_2^* \nabla G_1) dx = \frac{[a(x_{0,1}-x_{0,2}) + i(p_{0,1} + p_{0,2})]}{2}S \Rightarrow \\ (\int (G_2^* \nabla G_1) dx)^* &= \frac{[-a(x_{0,2}-x_{0,1}) - i(p_{0,1} + p_{0,2})]}{2}S = -D_{12} \end{aligned} \quad (5)$$

So:

$$D_{12}^* = -D_{21}. \quad (6)$$

#### 4. Kinetic energy integral

$$g = \exp\left(-\frac{a}{2}(x-x_0)^2 + ip_0(x-x_0)\right) \quad (7)$$

$$g'_x = (-a(x-x_0) + ip_0)G \quad (8)$$

$$\begin{aligned} g''_x &= -ag + (-a(x-x_0) + ip_0)g' = -ag + (-a(x-x_0) + ip_0)^2 g = \\ &= (a^2(x-x_0)^2 - 2iap_0(x-x_0) - a - p_0^2)g \end{aligned} \quad (9)$$

$$\begin{aligned} \int G_1^* \nabla^2 G_2 dx &= \left(\frac{a}{\pi}\right)^{1/2} \int (a^2(x-x_{0,2})^2 - 2iap_{0,2}(x-x_{0,2}) - a - p_{0,2}^2) g_1^* g_2 dx = \\ &= (-a - p_{0,2}^2) S_{12} + I_2 - I_1 \end{aligned} \quad (10)$$

$$\begin{aligned} I_2 &= \left(\frac{a}{\pi}\right)^{1/2} \int a^2(x-x_{0,2})^2 g_1^* g_2 dx = \\ &= a^2 \left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^2 + x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \\ &\int (x-x_{0,2})^2 \exp(-ax^2 + [a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})]x) dx = \\ &= C \int (x-x_{0,2})^2 \exp(-ax^2 + bx) dx = C \int (x-x_{0,2})^2 \exp\left(-a\left(x - \frac{b}{2a}\right)^2 + \frac{b^2}{4a}\right) dx = \int x' = x - \frac{b}{2a} \int \\ &= C' \int \left(x' + \frac{b}{2a} - x_{0,2}\right)^2 \exp(-ax'^2) dx' = C' \left[ \frac{1}{2} \left(\frac{\pi}{a^3}\right)^{1/2} + \left(\frac{b}{2a} - x_{0,2}\right)^2 \left(\frac{\pi}{a}\right)^{1/2} \right] = \\ &= a^2 \left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^2 + x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \exp\left(\frac{b^2}{4a}\right) \left[ \frac{1}{2} \left(\frac{\pi}{a^3}\right)^{1/2} + \left(\frac{b}{2a} - x_{0,2}\right)^2 \left(\frac{\pi}{a}\right)^{1/2} \right] = \\ &= a^2 \exp\left(-\frac{a}{2}(x_{0,1}^2 + x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}) + \frac{b^2}{4a}\right) \left[ \frac{1}{2a} + \left(\frac{b}{2a} - x_{0,2}\right)^2 \right] = \\ &= S_{12} a^2 \left[ \frac{1}{2a} + \left(\frac{b}{2a} - x_{0,2}\right)^2 \right] = S_{12} \left[ \frac{1a}{2} + \left(\frac{b-2ax_{0,2}}{2}\right)^2 \right] \end{aligned} \quad (11)$$

$$C = a^2 \left(\frac{a}{\pi}\right)^{1/2} \exp\left(-\frac{a}{2}(x_{0,1}^2 + x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2})\right) \quad (12)$$

$$C' = C \exp\left(\frac{b^2}{4a}\right) \quad (13)$$

$$b = [a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})] \quad (14)$$

$$\begin{aligned}
 I_1 &= 2iap_{0,2} \left( \frac{a}{\pi} \right)^{1/2} \exp \left( -\frac{a}{2} (x_{0,1}^2 + x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}) \right) \\
 &\int (x - x_{0,2}) \exp(-ax^2 + [a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})]x) dx = \\
 &= C \int (x - x_{0,2}) \exp(-ax^2 + bx) dx = C \int (x - x_{0,2}) \exp \left( -a \left( x - \frac{b}{2a} \right)^2 + \frac{b^2}{4a} \right) dx = \\
 &= C' \int (x - x_{0,2}) \exp \left( -a \left( x - \frac{b}{2a} \right)^2 \right) dx \stackrel{x' = x - \frac{b}{2a}}{=} C' \int \left( x' + \frac{b}{2a} - x_{0,2} \right) \exp(-ax'^2) dx' = \\
 &= C' \left( \frac{b}{2a} - x_{0,2} \right) \left( \frac{\pi}{a} \right)^{1/2} = 2iap_{0,2} \left( \frac{a}{\pi} \right)^{1/2} \exp \left( -\frac{a}{2} (x_{0,1}^2 + x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}) \right) \exp \left( \frac{b^2}{4a} \right) \left( \frac{b}{2a} - x_{0,2} \right) \left( \frac{\pi}{a} \right)^{1/2} = \\
 &= 2iap_{0,2} \exp \left( -\frac{a}{2} (x_{0,1}^2 + x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}) + \frac{b^2}{4a} \right) \left( \frac{b}{2a} - x_{0,2} \right) = \\
 &= S_{12} ip_{0,2} (b - 2ax_{0,2})
 \end{aligned} \tag{15}$$

$$C = 2iap_{0,2} \left( \frac{a}{\pi} \right)^{1/2} \exp \left( -\frac{a}{2} (x_{0,1}^2 + x_{0,2}^2) + i(p_{0,1}x_{0,1} - p_{0,2}x_{0,2}) \right) \tag{16}$$

$$C' = C \exp \left( \frac{b^2}{4a} \right) \tag{17}$$

$$b = [a(x_{0,1} + x_{0,2}) + i(p_{0,2} - p_{0,1})] \tag{18}$$

So:

$$\begin{aligned}
 \int G_1^* \nabla^2 G_2 dx &= (-a - p_{0,2}^2) S_{12} + I_2 - I_1 = \\
 &= (-a - p_{0,2}^2) S_{12} + S_{12} \left[ \frac{a}{2} + \left( \frac{b - 2ax_{0,2}}{2} \right)^2 \right] - S_{12} ip_{0,2} (b - 2ax_{0,2}) = \\
 &= S_{12} \left\{ -a - p_{0,2}^2 + \frac{a}{2} + \frac{[a(x_{0,1} - x_{0,2}) + i(p_{0,2} - p_{0,1})]^2}{4} - ip_{0,2} [a(x_{0,1} - x_{0,2}) + i(p_{0,2} - p_{0,1})] \right\} = \\
 &= S_{12} \left\{ -\frac{a}{2} - p_{0,2}^2 + \frac{1}{4} a^2 (x_{0,1} - x_{0,2})^2 + \frac{1}{2} ai(x_{0,1} - x_{0,2})(p_{0,2} - p_{0,1}) - \frac{1}{4} (p_{0,2}^2 - 2p_{0,2}p_{0,1} + p_{0,1}^2) \right. \\
 &\quad \left. - ai(x_{0,1} - x_{0,2})p_{0,2} + p_{0,2}^2 - p_{0,2}p_{0,1} \right\} = \\
 &= S_{12} \left\{ -\frac{a}{2} + a^2 \left( \frac{x_{0,1} - x_{0,2}}{2} \right)^2 - \left( \frac{p_{0,2} + p_{0,1}}{2} \right)^2 - \frac{1}{2} ai(x_{0,1} - x_{0,2})(p_{0,2} + p_{0,1}) \right\}
 \end{aligned} \tag{19}$$

## 5. Summary

$$G = \left( \frac{a}{\pi} \right)^{1/4} \exp \left( -\frac{a}{2} (x - x_0)^2 + ip_0(x - x_0) \right), \tag{20}$$

$$S_{G,12} = \int G_1^* G_2 dx = \exp \left( -\frac{a}{4} (x_{0,1} - x_{0,2})^2 - \frac{1}{4a} (p_{0,2} - p_{0,1})^2 \right) \exp \left( i \left( (x_{0,1} - x_{0,2}) \frac{(p_{0,1} + p_{0,2})}{2} \right) \right) \tag{21}$$

$$D_{G,12} = \int G_1^* \nabla G_2 dx = \frac{[a(x_{0,2} - x_{0,1}) + i(p_{0,2} + p_{0,1})]}{2} S_{G,12} \tag{22}$$

$$T_{G,12} = \int G_1^* \nabla^2 G_2 dx = S_{G,12} \left\{ -\frac{a}{2} + a^2 \left( \frac{x_{0,1} - x_{0,2}}{2} \right)^2 - \left( \frac{p_{0,2} + p_{0,1}}{2} \right)^2 - \frac{1}{2} ai(x_{0,1} - x_{0,2})(p_{0,2} + p_{0,1}) \right\} \tag{23}$$

## 6. Generalization

The form of the Gaussian, Eq. 1, is not general – the momentum term should have the  $\hbar$  denominator. In addition, the complex phase factor is often present. So the more appropriate form is

$$G = \left(\frac{a}{\pi}\right)^{1/4} \exp\left(-\frac{a}{2}(x-x_0)^2 + \frac{ip_0}{\hbar}(x-x_0) + \frac{i\gamma}{\hbar}\right) \quad (24)$$

Let also denote  $\alpha = \frac{a}{2}$ , so:

$$G = \left(\frac{2\alpha}{\pi}\right)^{1/4} \exp\left(-\alpha(x-x_0)^2 + \frac{ip_0}{\hbar}(x-x_0) + \frac{i\gamma}{\hbar}\right) \quad (25)$$

This form, for instance, is present in the paper of Makhov et al. [ Makhov, D. V.; Glover, W. J.; Martinez, T. J.; Shalashilin, D. V. Ab Initio Multiple Cloning Algorithm for Quantum Nonadiabatic Molecular Dynamics. The Journal of Chemical Physics 2014, 141 (5), 54110.]

The results Eqs. 21-23 derived for the form Eq. 1 will transform to:

$$S_{G,12} = \int G_1^* G_2 dx = \exp\left(-\frac{\alpha}{2}(x_{0,1} - x_{0,2})^2 - \frac{1}{8\alpha\hbar^2}(p_{0,2} - p_{0,1})^2 + i(x_{0,1} - x_{0,2})\frac{(p_{0,1} + p_{0,2})}{2\hbar} + \frac{i}{\hbar}(\gamma_2 - \gamma_1)\right) \quad (26)$$

$$D_{G,12} = \int G_1^* \nabla G_2 dx = \left[\alpha(x_{0,2} - x_{0,1}) + i\frac{(p_{0,2} + p_{0,1})}{2\hbar}\right] S_{G,12} \quad (27)$$

$$T_{G,12} = \int G_1^* \nabla^2 G_2 dx = S_{G,12} \left\{-\alpha + \alpha^2(x_{0,1} - x_{0,2})^2 - \frac{1}{\hbar^2}\left(\frac{p_{0,2} + p_{0,1}}{2}\right)^2 - 2\alpha i(x_{0,1} - x_{0,2})\frac{(p_{0,2} + p_{0,1})}{2}\right\} \quad (28)$$

Introducing variables

$$\Delta X = x_{0,2} - x_{0,1}, \Delta P = p_{0,2} - p_{0,1}, \Delta\gamma = \gamma_2 - \gamma_1, \bar{P} = \frac{p_{0,1} + p_{0,2}}{2}. \quad (29)$$

The above equations simplify to:

$$S_{G,12} = \int G_1^* G_2 dx = \exp\left(-\frac{\alpha}{2}(\Delta X)^2 - \frac{1}{8\alpha\hbar^2}(\Delta P)^2 - \frac{i\bar{P}\Delta X}{\hbar} + \frac{i\Delta\gamma}{\hbar}\right) \quad (30)$$

$$D_{G,12} = \int G_1^* \nabla G_2 dx = \left[\alpha(\Delta X) + \frac{i\bar{P}}{\hbar}\right] S_{G,12} \quad (31)$$

$$T_{G,12} = \int G_1^* \nabla^2 G_2 dx = S_{G,12} \left\{-\alpha + \alpha^2(\Delta X)^2 - \frac{\bar{P}^2}{\hbar^2} + 2\alpha i\Delta X \cdot \bar{P}\right\} \quad (32)$$

The equations for the overlap and for the kinetic energy are consistent with those of Makhov et al. However, the equation for the derivative coupling is not – they attribute similar equation to a different type of the derivative coupling.

To show consistency of the overlap, consider

$$\begin{aligned}
 & \frac{i}{\hbar} \left( p_{0,1} x_{0,1} - p_{0,2} x_{0,2} + \frac{x_{0,1} + x_{0,2}}{2} (p_{0,2} - p_{0,1}) \right) = \\
 & = \frac{i}{\hbar} \left( p_{0,1} x_{0,1} - p_{0,2} x_{0,2} + \frac{x_{0,1} p_{0,2}}{2} - \frac{x_{0,1} p_{0,1}}{2} + \frac{x_{0,2} p_{0,2}}{2} - \frac{x_{0,2} p_{0,1}}{2} \right) = \\
 & = \frac{i}{\hbar} \left( \frac{x_{0,1} p_{0,2}}{2} + \frac{x_{0,1} p_{0,1}}{2} - \frac{x_{0,2} p_{0,2}}{2} - \frac{x_{0,2} p_{0,1}}{2} \right) = \frac{i}{\hbar} \left( \frac{(x_{0,1} - x_{0,2}) p_{0,2}}{2} + \frac{(x_{0,1} - x_{0,2}) p_{0,1}}{2} \right) = \\
 & = -\frac{i \Delta X}{\hbar} \left( \frac{p_{0,2} + p_{0,1}}{2} \right) = -\frac{i \Delta X \cdot \bar{P}}{\hbar}
 \end{aligned}$$

So, their expression coincides with our term  $-\frac{i \bar{P} \Delta X}{\hbar}$