

Meyer-Miller-Thoss-Stock (MMTS) formulation revised

1. Statement 1

In variables $q_i = \text{Re}(c_i)$ and $p_i = \text{Im}(c_i)$, the TD-SE:

$$i\hbar \frac{\partial c_i(t)}{\partial t} = \sum_j H_{ij} c_j(t), \quad (1.1)$$

can be written as:

$$\begin{aligned} \frac{dq_k}{dt} &= \sum_j [\text{Re}(\Omega_{kj}) p_j + \text{Im}(\Omega_{kj}) q_j] = \frac{\partial H_{eff}}{\partial p_k} \\ \frac{dp_k}{dt} &= -\sum_j [\text{Re}(\Omega_{kj}) q_j - \text{Im}(\Omega_{kj}) p_j] = -\frac{\partial H_{eff}}{\partial q_k}, \end{aligned} \quad (1.2)$$

with

$$\Omega = \frac{H}{\hbar}. \quad (1.3)$$

Proof:

$$\begin{aligned} i\hbar \frac{d \text{Re}(c_k) + i \text{Im}(c_k)}{dt} &= \sum_j H_{kj} [\text{Re}(c_j) + i \text{Im}(c_j)] = \sum_j [\text{Re}(H_{kj}) + i \text{Im}(H_{kj})] [\text{Re}(c_j) + i \text{Im}(c_j)] = \\ &= \sum_j [\text{Re}(H_{kj}) \text{Re}(c_j) - \text{Im}(H_{kj}) \text{Im}(c_j)] + i [\text{Re}(H_{kj}) \text{Im}(c_j) + \text{Im}(H_{kj}) \text{Re}(c_j)] \Leftrightarrow \\ \frac{d \text{Re}(c_k)}{dt} &= \sum_j [\text{Re}(\Omega_{kj}) \text{Im}(c_j) + \text{Im}(\Omega_{kj}) \text{Re}(c_j)] \\ \frac{d \text{Im}(c_k)}{dt} &= -\sum_j [\text{Re}(\Omega_{kj}) \text{Re}(c_j) - \text{Im}(\Omega_{kj}) \text{Im}(c_j)] \end{aligned} \quad (1.4)$$

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Note that the leftmost equalities in Eqs. 1.2 are exact and true for a general Hamiltonian, Hermitian or non-Hermitian. The rightmost equalities in Eq. 1.2 are desired, but are not always possible – they are only possible if the Hamiltonian H is Hermitian (classical analog: Hamiltonian dynamics). If H is non-Hermitian, the rightmost

equalities can not be met, additional terms are needed (classical analog: non-Hamiltonian dynamics).

2. Statement 2

For a Hermitian Hamiltonian, $H_{ij} = H_{ji}^*$, Eqs. 1.2 are Hamiltonian (proof of the rightmost equalities), with the effective Hamiltonian, H_{eff} , given by:

$$H_{eff} = \sum_{i,j} \text{Re} \left(\frac{H_{ij}}{2\hbar} \right) (q_i q_j + p_i p_j) + \sum_{i,j} \text{Im} \left(\frac{H_{ij}}{\hbar} \right) p_i q_j. \quad (2.1)$$

Proof:

First, assume a very general quadratic form of the effective Hamiltonian:

$$H_{eff} = \sum_{i,j} B_{ij} q_i q_j + \sum_{i,j} C_{ij} p_i p_j + \sum_{i,j} D_{ij} p_i q_j. \quad (2.2)$$

Then

$$\begin{aligned} \frac{\partial H_{eff}}{\partial p_k} &= \sum_j (C_{kj} + C_{jk}) p_j + \sum_j D_{kj} q_j \\ \frac{\partial H_{eff}}{\partial q_k} &= \sum_j (B_{kj} + B_{jk}) q_j + \sum_j D_{jk} p_j \end{aligned}, \quad (2.3)$$

and

$$\begin{aligned} \frac{dq_k}{dt} &= \sum_j [\text{Re}(\Omega_{kj}) p_j + \text{Im}(\Omega_{kj}) q_j] = \frac{\partial H_{eff}}{\partial p_k} \\ \frac{dp_k}{dt} &= -\sum_j [\text{Re}(\Omega_{kj}) q_j - \text{Im}(\Omega_{kj}) p_j] = -\frac{\partial H_{eff}}{\partial q_k} \end{aligned}. \quad (2.4)$$

We want the rightmost sides of Eq. 2.4 be satisfied:

$$\begin{aligned} \frac{\partial H_{eff}}{\partial p_k} &= \sum_j (C_{kj} + C_{jk}) p_j + \sum_j D_{kj} q_j = \dot{q}_k = \sum_j \text{Re}(\Omega_{kj}) p_j + \sum_j \text{Im}(\Omega_{kj}) q_j \\ -\frac{\partial H_{eff}}{\partial q_k} &= \sum_j (-B_{kj} - B_{jk}) q_j + \sum_j -D_{jk} p_j = \dot{p}_k = -\sum_j \text{Re}(\Omega_{kj}) q_j + \sum_j \text{Im}(\Omega_{kj}) p_j \end{aligned}. \quad (2.5)$$

These equations are equivalent to:

$$\text{Im}(\Omega_{kj}) = D_{kj}, \quad (2.6a)$$

$$\text{Im}(\Omega_{jk}) = -D_{jk}. \quad (2.6b)$$

Equations 2.6a and 2.6b imply:

$$\text{Im}(\Omega_{jk}) = -D_{kj} = -\text{Im}(\Omega_{kj}). \quad (2.7c)$$

From Eq. 2.5, we find that:

$$\text{Re}(\Omega_{kj}) = (C_{kj} + C_{jk}), \quad (2.8a)$$

$$\text{Re}(\Omega_{jk}) = (B_{kj} + B_{jk}). \quad (2.8b)$$

Equations 2.8a and 2.9b imply:

$$\text{Re}(\Omega_{kj}) = \text{Re}(\Omega_{jk}). \quad (2.8c)$$

To see that:

$$\text{Re}(\Omega_{kj}) = (C_{kj} + C_{jk}) = (C_{jk} + C_{kj}) = \text{Re}(\Omega_{jk}). \quad (2.9)$$

The matrices C and B can be related to each other as $C = B$ or $C = B^T$. These two conditions can be satisfied simultaneously, if the matrices are symmetric. There is little physical meaning in requiring either of matrices C or B to be non-symmetric, so this condition is met practically always.

In summary, we have shown that:

$$\text{Re}(\Omega_{kj}) = \text{Re}(\Omega_{jk}), \quad (2.10a)$$

$$\text{Im}(\Omega_{jk}) = -\text{Im}(\Omega_{kj}). \quad (2.10b)$$

In other words, the Hamiltonian equations of motion, Eq. 2.4, are held only when the quantum Hamiltonian, H , is Hermitian. The form of the effective Hamiltonian can be derived as follows.

Using the relationship, $C = B$, we can simplify Eq. 2.2 to:

$$H_{eff} = \sum_{i,j} B_{ij} (q_i q_j + p_i p_j) + \sum_{i,j} D_{ij} p_i q_j. \quad (2.11)$$

Using the fact that $\text{Re}(\Omega_{kj}) = (B_{kj} + B_{jk})$, and $B = B^T$, we obtain:

$$B = \frac{1}{2} \text{Re}(\Omega) = \text{Re}\left(\frac{H}{2\hbar}\right). \quad (2.12)$$

From the above definitions, it follows that:

$$D = \text{Im}(\Omega) = \text{Im}\left(\frac{H}{\hbar}\right). \quad (2.13)$$

Combining Eqs. 2.11-2.13, we finally obtain:

$$H_{eff} = \sum_{i,j} \text{Re}\left(\frac{H_{ij}}{2\hbar}\right) (q_i q_j + p_i p_j) + \sum_{i,j} \text{Im}\left(\frac{H_{ij}}{\hbar}\right) p_i q_j. \quad (2.14)$$

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3. Statement 3

If the electronic Hamiltonian is Hermitian, then the total electronic energy, E , is related to the effective mapped Hamiltonian, H_{eff} , as:

$$E = 2\hbar H_{eff}. \quad (3.1)$$

Proof:

$$\begin{aligned} E &= \langle \Psi | H | \Psi \rangle = \sum_{i,j} c_i^* c_j H_{ij} = \sum_{i,j} H_{ij} [(q_i q_j + p_i p_j) + i(q_i p_j - p_i q_j)] = \\ &= \sum_{i,j} [\text{Re}(H_{ij}) + i \text{Im}(H_{ij})] [(q_i q_j + p_i p_j) + i(q_i p_j - p_i q_j)] = \\ &= \sum_{i,j} [\text{Re}(H_{ij})(q_i q_j + p_i p_j) - \text{Im}(H_{ij})(q_i p_j - p_i q_j)] + \\ &+ i \sum_{i,j} \text{Re}(H_{ij})(q_i p_j - p_i q_j) + i \sum_{i,j} \text{Im}(H_{ij})(q_i q_j + p_i p_j) \end{aligned} \quad (3.2)$$

Then, using the hermiticity of the vibronic Hamiltonian, H :

$$\text{Re}(H_{ij}) = \text{Re}(H_{ji}), \quad (3.2a)$$

$$\text{Im}(H_{ij}) = -\text{Im}(H_{ji}), \quad (3.2b)$$

and re-summing the resulting terms, we obtain:

$$\begin{aligned} \sum_{i,j} \text{Re}(H_{ij})(q_i p_j - p_i q_j) &= \sum_{i < j} \text{Re}(H_{ij})(q_i p_j - p_i q_j) + \sum_{i > j} \text{Re}(H_{ij})(q_i p_j - p_i q_j) = \\ &= \sum_{i < j} \text{Re}(H_{ij})(q_i p_j - p_i q_j) + \sum_{i < j} \text{Re}(H_{ji})(q_j p_i - p_j q_i) = 0 \end{aligned} \quad (3.3)$$

and

$$\begin{aligned} \sum_{i,j} \text{Im}(H_{ij})(q_i q_j + p_i p_j) &= \sum_{i < j} \text{Im}(H_{ij})(q_i q_j + p_i p_j) + \sum_{i > j} \text{Im}(H_{ij})(q_i q_j + p_i p_j) = \\ &= \sum_{i < j} \text{Im}(H_{ij})(q_i q_j + p_i p_j) + \sum_{i < j} \text{Im}(H_{ji})(q_j q_i + p_j p_i) = 0 \end{aligned} \quad (3.4)$$

So, the imaginary term is zero, as it should be, and the energy is:

$$E = \sum_{i,j} [\text{Re}(H_{ij})(q_i q_j + p_i p_j) - \text{Im}(H_{ij})(q_i p_j - p_i q_j)]. \quad (3.5)$$

The last term is:

$$\begin{aligned} \sum_{i,j} [-\text{Im}(H_{ij})(q_i p_j - p_i q_j)] &= \sum_{i,j} \text{Im}(H_{ij})(p_i q_j - q_i p_j) = \sum_{i,j} \text{Im}(H_{ij}) p_i q_j - \sum_{i,j} \text{Im}(H_{ij}) q_i p_j = \\ &= \sum_{i,j} \text{Im}(H_{ij}) p_i q_j + \sum_{i,j} \text{Im}(H_{ji}) q_i p_j = \sum_{i,j} \text{Im}(H_{ij}) p_i q_j + \sum_{i,j} \text{Im}(H_{ij}) q_j p_i = 2 \sum_{i,j} \text{Im}(H_{ij}) p_i q_j \end{aligned} \quad (3.6)$$

Thus, the energy is:

$$E = \sum_{i,j} \text{Re}(H_{ij})(q_i q_j + p_i p_j) + 2 \sum_{i,j} \text{Im}(H_{ij}) p_i q_j = 2\hbar H_{\text{eff}}. \quad (3.7)$$

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