

Some integrals with Gaussian wavepackets

1. Definition

According to Heller's notation, the normalized Gaussian wavepacket is defined as:

$$G = \left(\frac{2a}{\pi\hbar} \right)^{1/4} \exp \left(-\frac{a+ib}{\hbar} (x-x_0)^2 + \frac{ip_0}{\hbar} (x-x_0) + \frac{i}{\hbar} \varphi \right), \quad (1)$$

$$\alpha = 2(a+ib)$$

Here, a is the width parameter, b is the position-momentum correlation, x_0 is the center of the wavepacket, p_0 is the momentum of the wavepacket, φ is the phase.

2. Derivatives

Now, assume all the parameters a , b , x_0 , p_0 , and φ depend on time explicitly.

Compute the time derivative of the Gaussian:

$$\begin{aligned} \frac{\partial}{\partial t} G &= \left(\frac{2a}{\pi\hbar} \right)^{1/4} \exp \left(-\frac{\alpha}{2\hbar} (x-x_0)^2 + \frac{ip_0}{\hbar} (x-x_0) + \frac{i}{\hbar} \varphi \right) = \\ &= \frac{2\dot{a}}{\pi\hbar} \frac{1}{4} \left(\frac{2a}{\pi\hbar} \right)^{-3/4} \exp \left(-\frac{\alpha}{2\hbar} (x-x_0)^2 + \frac{ip_0}{\hbar} (x-x_0) + \frac{i}{\hbar} \varphi \right) + \\ &+ \left\{ -\frac{\dot{\alpha}}{2\hbar} (x-x_0)^2 + \frac{\alpha}{\hbar} (x-x_0) \dot{x}_0 + \frac{i\dot{p}_0}{\hbar} (x-x_0) - \frac{ip_0}{\hbar} \dot{x}_0 + \frac{i}{\hbar} \dot{\varphi} \right\} G = \\ &= \frac{2\dot{a}}{\pi\hbar} \frac{1}{4} \left(\frac{2a}{\pi\hbar} \right)^{-1} G + \left\{ -\frac{\dot{\alpha}}{2\hbar} (x-x_0)^2 + \frac{\alpha}{\hbar} (x-x_0) \dot{x}_0 + \frac{i\dot{p}_0}{\hbar} (x-x_0) - \frac{ip_0}{\hbar} \dot{x}_0 + \frac{i}{\hbar} \dot{\varphi} \right\} G = \\ &= \left\{ \frac{\dot{a}}{4a} - \frac{\dot{\alpha}}{2\hbar} (x-x_0)^2 + \frac{\alpha}{\hbar} (x-x_0) \dot{x}_0 + \frac{i\dot{p}_0}{\hbar} (x-x_0) - \frac{ip_0}{\hbar} \dot{x}_0 + \frac{i}{\hbar} \dot{\varphi} \right\} G \end{aligned} \quad (2)$$

$$\nabla G = \left(-\frac{\alpha}{\hbar} (x-x_0) + \frac{ip_0}{\hbar} \right) G, \quad (3)$$

$$\nabla^2 G = \left\{ \left(-\frac{\alpha}{\hbar} (x-x_0) + \frac{ip_0}{\hbar} \right)^2 - \frac{\alpha}{\hbar} \right\} G, \quad (4)$$

3. Time-dependent Schrodinger equation

$$i\hbar \frac{\partial G}{\partial t} = \left(-\frac{\hbar^2}{2M} \nabla^2 + V(x) \right) G \quad (5)$$

Substitute derivatives and simplify

$$\begin{aligned}
 i\hbar \left\{ \frac{\dot{a}}{4a} - \frac{\dot{\alpha}}{2\hbar} (x-x_0)^2 + \frac{\alpha}{\hbar} (x-x_0)\dot{x}_0 + \frac{i\dot{p}_0}{\hbar} (x-x_0) - \frac{ip_0}{\hbar} \dot{x}_0 + \frac{i}{\hbar} \dot{\phi} \right\} G = \\
 = -\frac{\hbar^2}{2M} \left\{ \left(-\frac{\alpha}{\hbar} (x-x_0) + \frac{ip_0}{\hbar} \right)^2 - \frac{\alpha}{\hbar} \right\} G + V(x)G
 \end{aligned} \tag{6a}$$

or

$$\begin{aligned}
 i\hbar \left\{ \frac{\dot{a}}{4a} - \frac{\dot{\alpha}}{2\hbar} (x-x_0)^2 + \frac{\alpha}{\hbar} (x-x_0)\dot{x}_0 + \frac{i\dot{p}_0}{\hbar} (x-x_0) - \frac{ip_0}{\hbar} \dot{x}_0 + \frac{i}{\hbar} \dot{\phi} \right\} = \\
 = -\frac{\hbar^2}{2M} \left\{ \left(-\frac{\alpha}{\hbar} (x-x_0) + \frac{ip_0}{\hbar} \right)^2 - \frac{\alpha}{\hbar} \right\} + V(x)
 \end{aligned} \tag{6b}$$

This is exact propagation!!!

Simplify:

$$\begin{aligned}
 \left\{ i\hbar \frac{\dot{a}}{4a} - \frac{i\hbar\dot{a} - \hbar\dot{b}}{\hbar} (x-x_0)^2 + 2(ia-b)(x-x_0)\dot{x}_0 - \dot{p}_0(x-x_0) + p_0\dot{x}_0 - \dot{\phi} \right\} = \\
 = -\frac{\hbar^2}{2M} \left\{ 4\frac{a^2-b^2+2iab}{\hbar^2} (x-x_0)^2 - \frac{p_0^2}{\hbar^2} - 4\frac{(ia-b)}{\hbar} (x-x_0)\frac{p_0}{\hbar} \right\} - \frac{\alpha}{\hbar} + V(x) = \tag{7} \\
 = \frac{p_0^2}{2M} + \frac{2(ia-b)}{M} (x-x_0)p_0 - 2\frac{a^2-b^2+2iab}{M} (x-x_0)^2 + \frac{a+ib}{M} \hbar + V(x)
 \end{aligned}$$

Separate real and imaginary parts:

Real:

$$\begin{aligned}
 \dot{b}(x-x_0)^2 - 2b(x-x_0)\dot{x}_0 - \dot{p}_0(x-x_0) + p_0\dot{x}_0 - \dot{\phi} = \\
 = \frac{p_0^2}{2M} - 2b\frac{p_0}{M} (x-x_0) - 2\frac{a^2-b^2}{M} (x-x_0)^2 + \frac{a}{M} \hbar + V(x)
 \end{aligned} \tag{8a}$$

Imaginary:

$$\left\{ \hbar \frac{\dot{a}}{4a} - \dot{a}(x-x_0)^2 + 2a(x-x_0)\dot{x}_0 \right\} = \frac{2a}{M} (x-x_0)p_0 - 4\frac{ab}{M} (x-x_0)^2 + \frac{b}{M} \hbar \tag{8b}$$

From Eq. 8b, comparing the powers of $(x-x_0)$, we get:

$$\dot{x}_0 = \frac{p_0}{M} \tag{9a}$$

$$\dot{a} = \frac{4ab}{M} \tag{9b}$$

Eq. 8a then simplifies:

$$\dot{b}(x-x_0)^2 - \dot{p}_0(x-x_0) - \dot{\phi} = -\frac{p_0^2}{2M} - 2\frac{a^2-b^2}{M} (x-x_0)^2 + \frac{a}{M} \hbar + V(x) \tag{9c}$$

Lets use the 2-nd order expansion of $V(x)$:

$$V(x) = V(x_0) + V'(x_0)(x - x_0) + \frac{1}{2} V''(x_0)(x - x_0)^2 \quad (10)$$

Then

$$\begin{aligned} \dot{b}(x - x_0)^2 - \dot{p}_0(x - x_0) - \dot{\phi} = & -\frac{p_0^2}{2M} - 2\frac{a^2 - b^2}{M}(x - x_0)^2 + \frac{a}{M}\hbar + \\ & + V(x_0) + V'(x_0)(x - x_0) + \frac{1}{2} V''(x_0)(x - x_0)^2 \end{aligned} \quad (11)$$

Comparing the same powers of $(x - x_0)$ yields:

$$\dot{b} = -2\frac{a^2 - b^2}{M} + \frac{1}{2} V''(x_0), \quad (12a)$$

$$-\dot{p}_0 = V'(x_0), \quad (12b)$$

$$-\dot{\phi} = -\frac{p_0^2}{2M} + \frac{a}{M}\hbar + V(x_0). \quad (12c)$$

4. Equations of motion

$$\dot{p}_0 = -V'(x_0), \quad (13a)$$

$$\dot{x}_0 = \frac{p_0}{M}, \quad (13b)$$

$$\dot{\phi} = \frac{p_0^2}{2M} - V(x_0) - \frac{a}{2M}\hbar = L(x_0, p_0) - \frac{a}{2M}\hbar, \quad (13c)$$

$$\dot{a} = \frac{4ab}{M}, \quad (13d)$$

$$\dot{b} = \frac{1}{2} V''(x_0) - 2\frac{a^2 - b^2}{M}. \quad (13e)$$