Two state system

Consider the TD-SE for two level system:

$$i\hbar\dot{c}_{0} = E_{0}c_{0} - i\hbar d_{01}c_{1} i\hbar\dot{c}_{1} = -i\hbar d_{10}c_{0} + E_{1}c_{1}$$
(1)

Introduce new variables:

$$a_i = -i\frac{E_i}{\hbar}, (2)$$

$$b = -d_{01}. (3)$$

The Eqs. (1) then turn to:

$$\dot{c}_0 = a_0 c_0 + b c_1
\dot{c}_1 = -b c_0 + a_1 c_1$$
(4)

Assuming that a_i and b are constant (do not depend on time), by taking second time-derivative we obtain:

$$\ddot{c}_0 = a_0 \dot{c}_0 + b \dot{c}_1 = a_0 (a_0 c_0 + b c_1) + b(-b c_0 + a_1 c_1) =$$

$$= (a_0^2 - b^2) c_0 + (a_0 + a_1) b c_1 = (a_0^2 - b^2) c_0 + (a_0 + a_1) (\dot{c}_0 - a_0 c_0),$$
(5)

or

$$\ddot{c}_0 - (a_0 + a_1)\dot{c}_0 + \left[a_0(a_0 + a_1) - (a_0^2 - b^2)\right]c_0 = 0,$$
(6)

or, finally

$$\ddot{c}_0 - (a_0 + a_1)\dot{c}_0 + (a_0a_1 + b^2)c_0 = 0.$$
(7)

To solve the second-order ODE Eq. (7) we assume the following ansatz: $c_0 \sim \exp(\alpha t)$. (8)

Substitution of Eq. (8) into Eq. (7) gives the quadratic equation:

$$\alpha^2 - \alpha(a_0 + a_1) + (a_0 a_1 + b^2) = 0. (9)$$

with the roots

$$\alpha_{1,2} = \frac{(a_0 + a_1) \pm \sqrt{(a_0 + a_1)^2 - 4(a_0 a_1 + b^2)}}{2} = \frac{(a_0 + a_1) \pm \sqrt{(a_0 - a_1)^2 - 4b^2}}{2}.$$
 (10)

Each of the roots gives the specific solution of the ODE Eq. (7). The general solution is obtained by summing specific solutions with some coefficients:

$$c_0(t) = A \exp(\alpha_1 t) + B \exp(\alpha_2 t) = \exp\left(\frac{(a_0 + a_1)}{2}t\right) \left[A \exp(i\Omega t) + B \exp(-i\Omega t)\right],\tag{11}$$

where

$$\Omega = \frac{\sqrt{4b^2 - (a_0 - a_1)^2}}{2} = \sqrt{b^2 - \left(\frac{a_0 - a_1}{2}\right)^2} = \sqrt{b^2 + \left(\frac{E_0 - E_1}{2\hbar}\right)^2} \ . \tag{12}$$

The coefficients are determined by the initial conditions. For example If $c_0(t) = 0$, then B = -A and solution Eq. (11) takes the form:

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$$c_0(t) = A \exp\left(\frac{\left(a_0 + a_1\right)}{2}t\right) \left[\exp(i\Omega t) - \exp(-i\Omega t)\right] = 2iA \exp\left(\frac{\left(a_0 + a_1\right)}{2}t\right) \sin(\Omega t). \tag{13}$$

The population of the considered state (0) is then given by:

$$\left|c_0(t)\right|^2 = 4A^2 \sin^2(\Omega t). \tag{14}$$

The exponent disappears because a_i are the pure imaginary quantities.

The transition rate can be obtained as:

$$r_{1\to 0} = \frac{d}{dt} \left| c_0(t) \right|^2 = 8A^2 \Omega \sin(\Omega t) \cos(\Omega t) = 4\Omega A^2 \sin(2\Omega t). \tag{15}$$

This shows that the population oscillates with some frequency, depending on the energy difference of two levels $E_0 - E_1$ and the coupling between them $b = -d_{01}$. For the degenerate case $E_0 - E_1 = 0$ the rate is determined only by the coupling d_{01} . Alternatively, if the states are not coupled $d_{01} = 0$ the rate of population transfer is determined only by the energy difference $E_0 - E_1$ - this is a qualitatively incorrect result, because the derivation scheme does not work under such conditions – see the paradox and its solution.

Some asymptotic analysis: using $\sin x \approx x$ we can find that at short times $\Omega t \to 0$:

$$r_{1\to 0} = 4\Omega A^2 \sin(2\Omega t) \approx 8tA^2\Omega^2 = 8tA^2 \left(b^2 + \left(\frac{E_0 - E_1}{2\hbar}\right)^2\right).$$
 (16)

A paradox:

If we assume the states are not coupled ($d_{01} = 0$) in the very beginning of the derivation, the equation (1) will read:

$$\dot{c}_0 = a_0 c_0
\dot{c}_1 = a_1 c_1$$
(17)

leading straight to the solution:

$$c_0(t) = c_0(0) \exp(a_0 t) = c_0(0) \exp\left(-i\frac{E_0}{\hbar}t\right),\tag{18}$$

so the population is conserved over time:

$$\left|c_{0}(t)\right|^{2} = \left|c_{0}(0)\right|^{2}.$$
 (19)

Solving the paradox:

If $d_{01} = 0 \Rightarrow b = 0$, then the substitution in Eq. (5) is invalid, that is the $b\dot{c}_1$ and bc_1 should be set to zero, leading to

$$\ddot{c}_0 = a_0 \dot{c}_0 + b \dot{c}_1 = a_0 \dot{c}_0 = a_0^2 c_0. \tag{20}$$

The equation Eq. (20) has the same solution as Eqs. (17), what solves the paradox.