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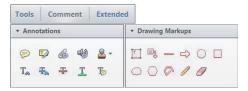
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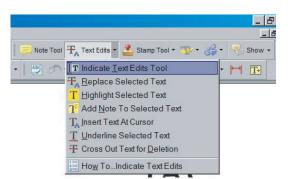
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Results of two multichord stellar occultations by dwarf planet (1) Ceres

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ABSTRACT

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We report the results of two multichord stellar occultations by the dwarf planet (1) Ceres that were observed from Brazil on 2010 August 17, and from the USA on 2013 October 25. Four positive detections were obtained for the 2010 occultation, and nine for the 2013 occultation. Elliptical models were adjusted to the observed chords to obtain Ceres' size and shape. Two limb-fitting solutions were studied for each event. The first one is a nominal solution with an indeterminate polar aspect angle. The second one was constrained by the pole coordinates as given by Drummond et al. Assuming a Maclaurin spheroid, we determine an equatorial diameter of 972 ± 6 km and an apparent oblateness of 0.08 ± 0.03 as our best solution. These results are compared to all available size and shape determinations for Ceres made so far, and shall be confirmed by the NASA's Dawn space mission.

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Key words: occultations – minor planets, asteroids: individual (1, Ceres) – planets and satellites: fundamental parameters.

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1 INTRODUCTION

Ceres is the sole example of a dwarf planet in the inner Solar system. Far from being mere taxonomic information, this suggests

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the great impact its study can have on the understanding of planetary formation and evolution of the Solar system. Indeed, it was proposed that Ceres' origin could have been as a trans-Neptunian object (McKinnon 2012), later scattered to the Main Belt due to the giant planets' migration predicted by the 'Nice Model' (Gomes et al. 2005). Even if it was formed close to its current location, the dynamical history of the Solar system must have left its signatures on Ceres. These could include not only the late heavy bombardment features that might exist on its surface, but also the makeup of its volatiles, which could have been transported from the outer region.

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Table 1. Circumstances of observation for all observing stations of the 2010 event.

| Site | Longitude Latitude Height | Telescope Aperture f-ratio | Exposure Cycle time Camera | Result Ingress Egress | Observer |
|--------------------------------|--|----------------------------|----------------------------------|--|---|
| Belo Horizonte CEAMIG-REA | 43°59′51″.1 W 19°49′49″.0 S 825 m | LX200 31 cm f/10 | 5 s 12 s SBIG ST10 | Positive 22:39:03.9 \pm 0.6 s 22:40:20 \pm 5 s | C. Jacques E. Pimentel |
| Pico dos Dias LNA | 45°34′45″.1 W 22°32′03″.7 S 1864 m | Zeiss 60 cm f/12.5 | 1 s 2 s Andor Ikon | Positive $22:37:30.3 \pm 0.6 \text{ s}$ $22:41:55.3 \pm 0.7 \text{ s}$ | J. I. B. Camargo G. B. Rossi |
| São José dos Campos INPE | 45°51′44″.0 W 23°12′33″.0 S 620 m | C11 28 cm f/6.3 | 2 s 5 s SBIG ST7 | Egress only Start obs.: 22:39:44 $22:42:03.0 \pm 0.2 \text{ s}$ | A. C. Milone T. Maldonado M. Okada |
| Ponta Grossa UEPG | 50°05′56″.0 W 25°05′22″.2 S 910 m | RCX 400 40 cm f/8 | 30 s 52 s SBIG STL6E | Positive $22:37:17 \pm 13 \text{ s}$ $22:39:56 \pm 13 \text{ s}$ | M. Emilio L. Mehret |
| Florianópolis UFSC | 48°31′20″.5 W 27°36′12″.3 S 20 m | C11 28 cm f/10 | 3 s 6 s SBIG ST7 | No occultation Start obs.: 21:49:27 End obs.: 22:51:21 | W. Schoenell A. J. T. Mello F. R. Herpich |

Since the 1970s it has been speculated that Ceres could contain water i hich was recently verified (Küppers et al. 2014). Although the water regime on this object is still unknown, some internal structure models suggest the existence of a water vapour - or even a liquid water - layer (Castillo-Rogez 2011). Yet the very question of whether Ceres underwent differentiation is open and, on the assumption of an affirmative answer, it is natural to ask if it ever had tectonics, what its geological evolution was, and if it is still active. Inarguably NASA's *Dawn* mission (Russell et al. 2004), which is currently orbiting the dwarf planet, will shed light on several open issues concerning Ceres.

Containing approximately one-fifth of the whole Main Belt's mass, Ceres is expected to have an equilibrium figure, i.e. a Maclaurin or a Jacobi ellipsoid. In fact, direct observations of Ceres by means of adaptive optics indicate that it is an oblate spheroid (Drummond et al. 2014). The precise knowledge of its size and shape is of utmost importance, for the models of density, internal structure, and differentiation.

The best ground-based technique for determining shape and size of a faraway object is the study of its shadow, cast by a star during an occultation. Since the 1960s occultations have provided measurements of hundreds of asteroids, thanks partially to the fruitful professional-amateur collaboration on the field. More recently, this technique has been applied to objects of the outer Solar system and has unveiled outstanding features of distant bodies, e.g. the ring system around the Centaur (10199) Chariklo (Braga-Ribas et al. 2014).

The first stellar occultation by Ceres was observed in 1984 (Millis et al. 1987) and led to the determination of its size to the precision of some kilometres, at a time when the uncertainties were often ten times larger. The high apparent brightness of Ceres, as compared to most asteroids, imposes a somewhat strong constraint on the stars capable of causing a detectable magnitude drop when occulted. For instance, after the 1984 event, to our knowledge, only four stellar occultations have been observed (Dunham et al. 2014). Two of them had only two chords each, thus not sufficient for providing accurate results. The two remaining events, which occurred on 2010 August 17 and 2013 October 25, are reported in the present work and are the first ones that used charged-couple devices (CCD) as recording systems. Throughout the paper we shall refer to these events as the '2010 occultation' and the '2013 occultation'.

Both events were predicted by Steve Preston² on behalf of the International Occultation Timing Association, during routine prediction of asteroidal occultations of bright stars.

This work is organized as follows. In Sections 2 and 3 we analyse the 2010 and the 2013 occultations, respectively. In Section 4 we give the geocentric positions of Ceres derived from the occultations. The comparison of our results to those in the literature is carried out in Section 5.

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2 THE 2010 OCCULTATION

2.1 Observations

As predicted, on 2010 August 17 Ceres occulted the star TYC 6833-163-1 (UCAC4 313-111823), which has magnitude V = 11.55 and has the ICRS position, based on UCAC4 catalogue (Zacharias et al. 2013), to the date of occultation:

$$\begin{cases} \alpha = 17^{h}18^{m}29^{s}.0085 \\ \delta = -27^{\circ}26'38''.867. \end{cases}$$
 (1)

Observations were carried out in Brazil from five different sites as displayed in Table 1 and Fig. 1. The occultation was detected from Q7 four of them. The southernmost one (UFSC) had a negative chord. From the positive sites, the one named INPE started observing after the event was already in progress, due to technical difficulties, thus providing only the star's reappearance time; the other three recorded the whole phenomenon.

A remarkable circumstance of this event was the low velocity of Ceres: only $3.9 \,\mathrm{km}\,\mathrm{s}^{-1}$ in the plane of the sky. Therefore, even exposures of a few seconds would translate in a relevant spatial

All the observations were made through a sequence of images obtained with CCDs. The times of each exposure were available on the header of each image with an internal accuracy of a few

¹ These events took place on 1994 August 22 and 2010 October 30.

² Predictions are published at http://asteroidoccultation.com

Figure 1. Post-occultation reconstruction of Ceres' shadow path on Earth for the 2010 August 17 event. The big red dot is the geocentric closest approach at 22:40:25 UT. The small red ones represent the centre of the shadow separated by one minute, shadow moves from the left to the right. Blue dots are the sites that have observed the event. As described in text, UFSC had a negative chord.

hundredths of a second. CEAMIG had only the integer part of the second available, due to the acquisition software used. The fraction of a second could be retrieved as described on the section 3.1 of Braga-Ribas et al. (2013, which shall be referred as BR13 hereafter). Cycle times (exposure plus read-out) varied from 2 to 52 s, as can be verified in Table 1, making it a heterogeneous set of observations and imposing an error of a few seconds for the ingress/egress times of some sites.

2.2 Light curves

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The flux of the star in the five occultation chords was obtained from the FITS images with the Platform for Reduction of Astronomical Q8 Images Automatically (PRAIA) (Assafin et al. 2011). The light curves were normalized to the flux of the star plus Ceres, as they were merged right before and after the occultation. Additionally, they were normalized by fitting a polynomial curve (of first or second order) outside the flux drop, so that the flux ratio was set to 1 outside the occultation.

The ingress (disappearance) and egress (reappearance) instants of the occultation were obtained for each light curve by fitting a square-well model convoluted with the Fresnel diffraction, the CCD bandwidth, the stellar apparent diameter, and the applied finite exposure time; see Widemann et al. (2009) and BR13. The smallest integration time used in the positive observations was 1.0 s, which translates to almost 3.9 km in the celestial plane. Therefore, the error on the time determination of the ingress and egress is largely dominated by the integration times, not by Fresnel diffraction or star diameter, which are both of the order of a few hundred metres for this event. The occultation data fit consists in minimizing a classical χ^2 function for each light curve, as described in Sicardy et al. (2011) and BR13. The free parameter to adjust is the ingress or egress instant, which provides the minimum value of χ^2 , denoted as χ^2_{min} . The best fittings to the 2010 occultation light curves are shown in Fig. 2, and the derived occultation times are listed in Table 1.

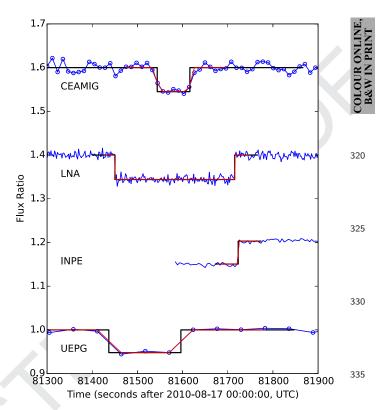


Figure 2. The four occultation light curves normalized and vertically shifted by a factor of 0.2 for better visualization. The solid black lines are the best fit of the square-well model to the data. Red lines are the square-well model convoluted with the Fresnel diffraction, the star diameter, and the applied exposure time. The mid-times of the occultations do not coincide due to the propagation delays of the shadow due to the distinct longitude of the sites. Exposures at INPE started after the immersion, as explained in the text

2.3 Limb-fitting methodology

The methodology used to analyse Ceres' profile from the observations is the same described by Sicardy et al. (2011) and BR13. Each combination of site position and recorded ingress/egress instant, together with star coordinates and Ceres' ephemeris, corresponds to a point in the plane of the sky. The collection of all these points ideally determines the apparent limb of Ceres.

We adopt an elliptic model for the limb profile, resulting from the projection of an oblate spheroid on to the sky plane. This choice is supported by the work of Drummond et al. (2014), by means of direct imaging of Ceres. Hence, we have N = 7 chord extremities to adjust the M=5 parameters which define an ellipse: apparent semimajor and semiminor axis (a' and b', respectively), position angle P of its semiminor axis and the position (f_c, g_c) of its centre with respect to the occulted star. Of course, the apparent semimajor axis a' is equivalent to the equatorial radius R_{equa} of the ellipsoid. The coordinates f_c and g_c , in kilometres, are calculated using the JPL#33 Ceres' ephemeris (Giorgini et al. 1996) and the occulted star's position. They are positive towards the local celestial east and north, respectively. The position angle P is counted positively from the direction of local celestial north to celestial east. The apparent oblateness can be defined by $\epsilon' = 1 - (b'/a')$. The bestfitting solution is obtained minimizing a reduced χ_r^2 function, where we define the number of degrees of freedom of the problem as $\mathcal{N} \equiv N - M$. All the procedures that allow the determination of the error bars of the physical parameters can be found in BR13.

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Table 2. Results of limb fitting to the data of the 2010 and 2013 events.

| Solution | 2010/Nominal | 2010/Pole-constrained | 2013/Nominal | 2013/Pole-constrained |
|--------------------------|-----------------|-----------------------|-----------------|-----------------------|
| Equatorial diameter (km) | 982 ± 14 | 972 ± 6 | 971 ± 7 | 971 ± 7 |
| True oblateness | 0.08 ± 0.03 | 0.08 ± 0.03 | 0.08 ± 0.04 | 0.08 ± 0.04 |
| Position angle (deg) | 5 ± 10 | $12 \pm 3 (*)$ | 22 ± 5 | $25 \pm 3 (*)$ |
| f_c (km) | 97 ± 9 | 102 ± 5 | 77 ± 6 | 78 ± 6 |
| g_c (km) | 16 ± 15 | 21 ± 11 | 13 ± 16 | 13 ± 16 |
| $\chi^2_{r,\min}$ | 0.24 | 0.42 | 1.27 | 1.27 |

Notes. In bold we highlight our best solution. Error bars are at 1σ level. The polar diameter (D_{pol}) can be easily calculated from $D_{\text{pol}} = D_{\text{equa}}(1 - \epsilon)$. (*) Position angle derived from Ceres' rotational pole coordinates determined by Drummond et al. (2014).

2.4 Limb-fitting solutions

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Two possible solutions were considered for the limb fitting. The first, which we call nominal solution, consists of determining the five parameters that characterize an ellipse from the seven observed contacts. As we shortly show, it led to a rather large uncertainty on the position angle. Furthermore, the nominal solution alone is not capable of returning the true oblateness, which can be evaluated through equation 2 of BR13 provided that the aspect angle ζ is known

It is possible to derive the pole position (ζ, P) from the coordinates of Ceres' pole (α_p, δ_p) and its ephemeris (α, δ) via

$$\begin{cases}
\cos(\zeta) = -\sin(\delta_p)\sin(\delta) - \cos(\delta_p)\cos(\delta)\cos(\alpha_p - \alpha) \\
\sin(P) = \frac{\sin(\alpha_p - \alpha)}{\tan(\delta_p)\cos(\delta) - \sin(\delta)\cos(\alpha_p - \alpha)},
\end{cases} (2)$$

where $\zeta=0^\circ$ and $\zeta=90^\circ$, respectively, correspond to a pole-on and an equator-on geometry.

From observations with adaptive optics spanning a 10-year period, Drummond et al. (2014) determined the position of Ceres' polar axis within an error of only 3°:

$$\alpha_p = (287 \pm 3)^\circ, \quad \delta_p = (+64 \pm 3)^\circ,$$
 (3)

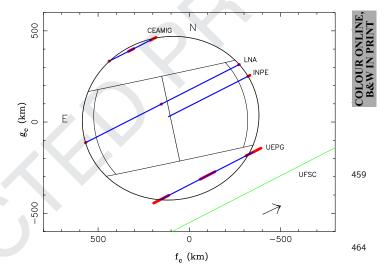
in equatorial J2000 coordinates. Together with Ceres' ephemeris at the moment of the occultation, this corresponds to the polar aspect angle $\zeta = 86$?1, which is very close to an equator-on geometry. Hence, we expect true figures to be similar to apparent ones.

The knowledge of Ceres' pole not only allows the determination of its polar aspect, it suffices to set its position angle. Therefore equations (2) and (3) may act as a constraint for P, and a second solution can be obtained by probing the parameter space with the restriction that the position angle is confined to the range that follows from equation (3). We call this the 'pole-constrained solution'.

2.4.1 Nominal solution

With the seven observed contacts it is possible to adjust the five parameters which define an ellipse. For the best-fitting solution we find $\chi^2_{r, \min} = 0.24$, which could be interpreted as a slightly overestimation of the error bars with regard to the good quality of the fit. However, inasmuch as the problem has only two degrees of freedom, it is far from the statistical realm and relatively small $\chi^2_{r, \min}$ are acceptable.

The resulting values of equatorial diameter, oblateness, position angle, and centre coordinates are presented in the second column of Table 2.



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Figure 3. The best elliptical fit for the occultation chords for the event of 2010 using the times from Table 1 and the pole-constrained solution. The arrow indicates the direction of motion, blue lines are the observed chords, the red segments are the ingress, egress and mid-occultation error bars at 1σ

Already mentioned, the parameter with the largest uncertainty is the position angle: spanning on a 20° interval, its determination has a relative precision worse than 10 per cent. Clearly, the coordinates of Ceres' pole (equation 3) can impose a strong constraint on the position angle, as the next solution shows.

Finally, the correction to the oblateness due to Ceres' polar aspect angle lies within the 1σ error bar and has no statistical relevance; hence $\epsilon=0.08\pm0.03$.

2.4.2 Pole-constrained solution

At the occultation, the coordinates (equation 3) of Ceres' rotational pole correspond to the position angle $P=(12\pm3)^\circ$. Exploration of the parameter space, restricted to ellipses with position angles laying within this range, results in the pole-constrained solution. The related physical parameters are displayed on the third column of Table 2 in boldface, while the best-fitting solution is depicted in Fig. 3.

We notice that the constraint corresponds to the upper limit of nominal solution's 1σ error bar for P. On the other hand, it selects the smallest values of semimajor axis, improving its determination by a factor of about 2. Notwithstanding, oblateness' figures remain the same.

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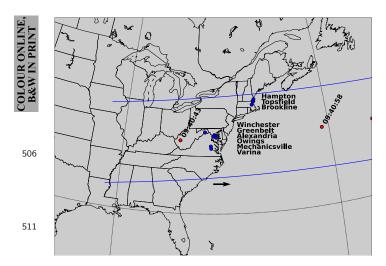


Figure 4. Post-occultation reconstruction of Ceres' shadow path on Earth for the 2013 October 25 event at the east coast of USA. Upper view of the 516 occultation over the sites that observed the event (blue dots). Red points are the centre of the shadow separated by 15 s.

3 THE 2013 OCCULTATION

3.1 Observations

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The event which took place on 2013 October 25 involved the star TYC 865-911-1 (UCAC4 496-058191), of magnitude $V = 10.\overline{05}$. Based on UCAC4 (Zacharias et al. 2013), the ICRS position to the date of occultation is

$$\begin{cases} \alpha = 11^{\text{h}}57^{\text{m}}52^{\text{s}}.7641\\ \delta = +09^{\circ}07'49'.835. \end{cases}$$
 (4)

The occultation could only be visible in the United States, before dawn, as depicted in Fig. 4.

Nine positive chords were obtained by the variety of instruments listed in Table 3. Each station was equipped with a video camera with negligible readout times. This was of particular importance, since in this event Ceres' shadow speed was 42.6 km s⁻¹, much faster than in the 2010 event.

Three different timing synchronization procedures were adopted among the set of observing stations. At Greenbelt and Owings, the 1PPS signal of a GPS unit was used to calibrate time stamps which were inserted at each frame of the video. Time extraction is thus straightforward, after taking camera delays into account. On the other hand, at Brookline the clock would be synchronized by an internet server. A lack of connection, however, resulted in spurious times. In fact, comparison between the times obtained at this station and the others suggests that the former have a delay of about 64 s. Therefore, we do not use Brookline's times in the analysis that follows. Finally, at the remaining six stations the videos were recorded by camcorders on digital tapes. The timing method consisted in the comparison of the camcorder internal clock to a 1PPS GPS signal, before and after the recording of the occultation. Absolute timing errors of this procedure are expected to be less than $0.1 \, \mathrm{s}.$

Table 3. Circumstances of observation for the observing stations of the 2013 event.

| Site | Longitude Latitude Height | Telescope: Aperture f-ratio | Camera Cycle time | Result Ingress Egress | Observer | |
|----------------|---|-----------------------------------|-------------------------|--|--------------------------|--|
| Hampton | 70°48′59″.7 W 42°53′52″.8 N 7 m | 12 cm f/5 | PC164C-EX2 0.033 s | Positive $09:40:46.9 \pm 0.1 \text{ s}$ $09:40:57.26 \pm 0.08 \text{ s}$ | T. Blank | |
| Topsfield | 70°55′16″.6 W 42°37′55″.9 N 45 m | 12 cm f/5 | PC164C-EX2 0.033 s | Positive $09:40:45.4 \pm 0.1 \text{ s}$ $09:40:58.0 \pm 0.1 \text{ s}$ | T. Blank | |
| Brookline | 71°08′14″.5 W 42°18′27″.4 N 109 m | 64 cm f/9.6 | Infinity2-1R 0.015 s | Positive $09:41:48.00 \pm 0.01 \text{ s}$ $09:42:02.93 \pm 0.01 \text{ s}$ | N. Weber R. Dantowitz | |
| Winchester | 78°14′39″.6 W 39°16′21″.5 N 211 m | 36 cm f/5 | PC164C 0.033 s | Positive $09:40:33.26 \pm 0.08 \text{ s}$ $09:40:55.86 \pm 0.09 \text{ s}$ | J. Brooks | |
| Greenbelt | 76°52′09″.4 W 38°59′12″.1 N 52 m | 12 cm f/2.5 | PC164C-EX2 0.033 s | Positive $09:40:33.6 \pm 0.1 \text{ s}$ $09:40:56.4 \pm 0.1 \text{ s}$ | J. Dunham D. Dunham | |
| Alexandria | 77°02′28″3 W 38°49′19″.1 N 8 m | 7 cm f/10 | Watec120N 0.067 s | Positive $09:40:33.3 \pm 0.1 \text{ s}$ $09:40:56.1 \pm 0.1 \text{ s}$ | P. Maley | |
| Owings | 76°38′06″3 W 38°41′26″5 N 38 m | 25 cm f/3.3 | PC164C 0.033 s | Positive $09:40:34.27 \pm 0.05 \text{ s}$ $09:40:56.0 \pm 0.2 \text{ s}$ | C. Ellington | |
| Mechanicsville | 77°23′06″.7 W 37°41′26″.1 N 60 m | 12 cm f/2.5 | PC164C-EX2 0.033 s | Positive $09:40:33.0 \pm 0.1 \text{ s}$ $09:40:54.8 \pm 0.1 \text{ s}$ | D. Dunham | |
| Varina | 77°19′49′.3 W 37°25′58′.6 N 19 m | 12 cm f/2.5 | PC164C-EX2 0.033 s | Positive $09:40:32.4 \pm 0.3 \text{ s}$ $09:40:53.1 \pm 0.2 \text{ s}$ | D. Dunham | |

During the event Ceres was low in the horizon, with altitudes between 15° (Winchester) and 20° (Hampton). Strong scintillation is expected in such a scenario which, combined to short integration times and the low magnitude drop of the event, resulted in rather noisy light curves and thus larger uncertainties in the time of the contacts, as is shown in the next section.

3.2 Light curves

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All videos were converted to FITS images and the photometry of the target was obtained via PRAIA (Assafin et al. 2011). The light curves were normalized to the flux of a reference star when available on the field

To reduce the noise, the data were binned by groups of five images – with the exception of Greenbelt, which was averaged in sets of 10. This procedure enlarges the effective integration times reported in Table 3 by a factor of 5 (or 10). As in the 2010 event, an additional normalization by a polynomial curve was applied.

The ingress and egress instants of the occultation were obtained by the same procedure explained in Section 2.2. Since the typical effective integration time used (0.17 s) translates to about 7 km in the celestial plane, and the Fresnel scale and star diameter are again of the order of a few hundreds of metres, the theoretical occultation light curves are largely dominated by the integration times, as for the 2010 event.

The best-fittings to the occultation light curves are shown in Fig. 5, and the derived occultation times are listed in Table 3.

3.3 Limb-fitting solutions

Elliptic limb profiles were adjusted to all the available³ chords by the same procedure described in Section 2.3. This yielded $\chi^2_{r, \min} = 13$, suggesting that an elliptic model is not satisfactory to the data. Indeed, a quick glance at a plot of the observed chords, such as in Fig. 6, shows that the Varina chord seems to be somewhat advanced with respect to the others. Taking into account that in this station time stamps were not inserted on the video frames, it is fairly possible to attribute this advance to an eventual problem on the correspondence between camcorder's and GPS' times. This could be caused, for example, if the camcorder delayed to start the recording and, since this was an unattended pre-pointed station, this fact would not be noticed by the analysis of the video itself.

The immersion recorded at Owings also seems to be shifted (delayed) with respect to the nearby chords (see Fig. 6). This chord, actually, has roughly the same length of Mechanicsville's, despite the fact that they are separated by about 100 km. Differently from Varina, though, this station had time stamps inserted in each frame of the video, which makes it more unlikely to justify an eventual timing issue. Another possible explanation for the apparent problem of the times of this chord is the determination of the ingress and egress instants in the light-curve analysis, which could have been affected by noise. Finally, the delay during the ingress could be caused by a relief feature in Ceres; we shall soon return to this hypothesis.

In a second limb fitting, thus, we did not consider Brookline, Varina, and Owings chords. The adjustment of the five parameters which define an ellipse to the 12 contacts then resulted in $\chi^2_{r,\min} = 1.27$, indicating that the fitting is in good agreement with

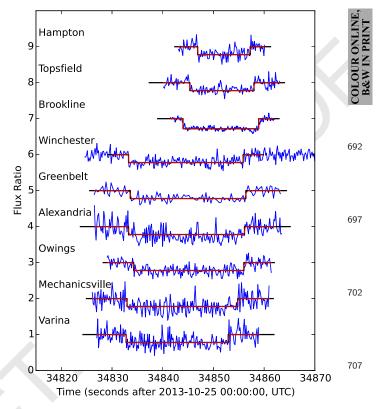


Figure 5. The nine occultation light curves normalized and vertically shifted by a factor of 1.0 for better viewing, see Fig. 2 for the explanation of the graph. The light curve of Brookline is shifted by -64 s as explained in the text.

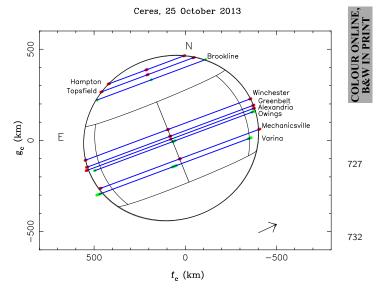


Figure 6. The best elliptical fit for the occultation chords for the event of 2013 using timing from Table 3 and the pole-constrained solution. The arrow indicates the direction of motion, blue lines are the observed chords, the red and green segments are the ingress, egress and mid-occultation error bars at 1σ level. The chords with green error bars were not used during the limb-fitting process. The chord of Brookline if shifted by -64 s as explained in the text.

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³ Brookline's chord was not used for limb fitting since it had an inaccurate absolute time.

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the observed data within the error bars. This is the solution depicted in Fig. 6, where we see that the chord length measured in Brookline is compatible to the model. The associated physical parameters are presented in Table 2 as the nominal solution.

For this event the polar aspect angle is $\zeta = 90^{\circ}.7$, which makes the true oblateness equal to the apparent one, within the error bars.

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The position angle of this nominal solution is better constrained than the 2010 occultation. Actually, the uncertainty in the former is of 5° in contrast to 10° of the latter. This suggests that the poleconstrained solution obtained via the position of Ceres' pole (equation 3) would not be significantly different from the nominal one.

This assumption was confirmed when we carried out the limb fitting with the constraint $P = (25 \pm 3)^{\circ}$, the position angle at the occultation which follows from equation (3). The physical parameters related to this pole-constrained solution are presented in the last column of Table 2, and are essentially the same of the nominal solution.

In the hypothesis that the delay observed in the immersion at Owings could be associated with a limb topography feature, the recorded contact would then correspond to a negative elevation of 31 ± 4 km with respect to the best-fitting ellipse. Theoretical models, however, predict that reliefs in Ceres should not be higher than about $10-20\,\mathrm{km}$ (Johnson & McGetchin 1973), while published observational data sets the bound of $18\,\mathrm{km}$ (Carry et al. 2008). More recently, images by the probe Dawn also reveal an even smoother surface. Therefore, the association of Owings first contact to a relief is improbable.

4 ASTROMETRY FROM OCCULTATIONS

The objects' geocentric positions derived from stellar occultations are most valuable for improving their orbits (Desmars et al. 2015). Usually the precision of the astrometric positions are limited by the accuracy of the occulted star position, not by the limb fit. Ceres' geocentric J2000 positions at the time of each occultation are displayed in equations (5) and (6). The errors of the positions come from the errors of the star positions, taken from the catalogues and from the errors of the relative apparent distances between star and Ceres, derived from the limb fit (which are displayed as $\Delta\alpha\cos\delta$, $\Delta\delta$).

2010 Aug 17
$$\begin{cases} \text{Time} = 22:40:00\\ \alpha = 17^{\text{h}}18^{\text{m}}29.0122 \pm 0.027\\ \delta = -27^{\circ}26.38.617 \pm 0.028\\ \Delta\alpha\cos\delta = 0.003; \ \Delta\delta = 0.007 \end{cases}$$
(5)

$$2013 \text{ Oct } 25 \begin{cases}
\text{Time} = 09 : 45 : 00 \\
\alpha = 11^{\text{h}} 57^{\text{m}} 52^{\text{s}} 9154 \pm 0.019 \\
\delta = +09^{\circ} 07' 49'.865 \pm 0.021 \\
\Delta \alpha \cos \delta = 0.002; \quad \Delta \delta = 0.007.
\end{cases} \tag{6}$$

5 DISCUSSION

A quick glance at Table 2 shows overall agreement between the physical parameters derived from both occultations, especially in the equatorial diameter. The differences of the solutions occur basically on the size of their error bars, and can be justified by the particularities of each set of data, as discussed below.

The 2010 event, for example, had only seven contacts; none the less, they were well distributed over Ceres' disc (see Fig. 3) acting as

Table 4. Ceres' equatorial diameter and oblateness.

| Eq. diameter (km) | Oblateness | Method | Ref. | |
|-------------------|-------------------|-------------|------|-----|
| 972 ± 6 | 0.08 ± 0.03 | Occultation | 1 | |
| 967 ± 10 | 0.078 ± 0.015 | Keck+VTL | 2 | 811 |
| 959 ± 5 | 0.074 ± 0.007 | Keck | 3 | |
| 975 ± 4 | 0.067 ± 0.005 | HST | 4 | |
| 959 ± 5 | 0.05 ± 0.01 | Occultation | 5 | |

References. 1: Present work. 2: Drummond et al. (2014). 3: Carry et al. (2008). 4: Thomas et al. (2005). 5: Millis et al. (1987).

a constraint to its shape. On the other hand, the 2013 event had five more exploitable contacts, but they were concentrated in certain regions of the body. In particular, the absence of chords close to Ceres' south pole made its oblateness less well determined here than in the 2010 event.

However, even our best measurement for the oblateness, $\epsilon=0.08\pm0.03$, has still a higher uncertainty with regard to other figures published in the literature, as Table 4 shows. A larger number of uniformly spaced chords would be necessary to offer a best constraint to the oblateness.

The few chords of the 2010 occultation could themselves only constrain the position angle of the object to a uncertainty of 10° . This uncertainty was reduced by a factor of 2 in the 2013 event, approaching – and verifying – the result predicted by the work of Drummond et al. (2014). As was shown, using the coordinates of Ceres' polar axis to limit the position angle was not an efficient procedure in the 2013 event, in the sense that it did not result in significant changes in the parameters obtained in the nominal solution (see Table 2).

On the other hand, constraining the position angle on the 2010 occultation was proved to reduce the error bars of the other parameters (disregarding oblateness). Moreover, this procedure resulted in excellent agreement between the equatorial radius figures of both events.

The 2013 occultation, therefore, offers not only an independent verification of the figures resulted from the 2010 event, but also validates the procedure carried out there which led to the best-constrained parameters in this work.

Comparison of Ceres' equatorial diameter as measured by different techniques is carried out in Table 4. We note overall agreement between our result to those obtained via direct imaging by the *Hubble Space Telescope (HST*; Thomas et al. 2005), the Keck Observatory and the ESO VLT (Drummond et al. 2014). The smaller figure reported by Carry et al. (2008) may be justified by the fact that in this study the effects of limb darkening were not taken into account, as pointed out by Drummond et al. (2014).

As mentioned in the Introduction, the 1984 event (Millis et al. 1987) is the only occultation data to which we can compare our result. The measured diameters do not agree, ours being larger by 2σ . It is difficult to state for sure the reasons for this divergence. One way to clarify the issue is to redetermine the immersion and emersion instants from the original light curves using the same methodology presented in our work. Note that in this occultation a variation of 0.5 s in the contact times correspond to almost 7 km on Ceres' limb, which is on the same order of the residuals of their elliptical fit. However, we could not do that, as we had no access to the original 1984 light-curve data. Moreover, we have no description of how the immersion and emersion instants of the

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chords were actually determined from these light curves in Millis et al. (1987).

NASA's *Dawn* mission shall bring light to these questions, which will be important not only for the knowledge of Ceres itself, but also for all the techniques used so far to study the physical properties of small Solar system objects, such as the stellar occultations.

ACKNOWLEDGEMENTS

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|---|------------------|
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| This paper has been typeset from a TEX/LATEX file prepared by the author. | 955 |
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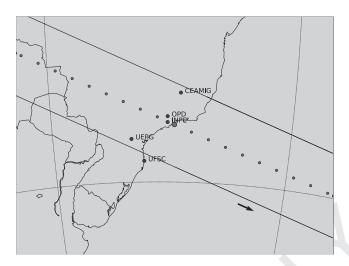


Figure 1. Post-occultation reconstruction of Ceres' shadow path on Earth for the 2010 August 17 event. The big red dot is the geocentric closest approach at 22:40:25 ut. The small red ones represent the centre of the shadow separated by one minute, shadow moves from the left to the right. Blue dots are the sites that have observed the event. As described in text, UFSC had a negative chord.

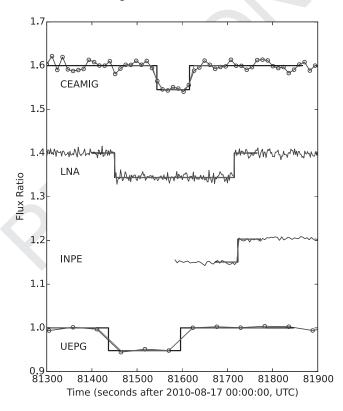


Figure 2. The four occultation light curves normalized and vertically shifted by a factor of 0.2 for better visualization. The solid black lines are the best fit of the square-well model to the data. Red lines are the square-well model convoluted with the Fresnel diffraction, the star diameter, and the applied exposure time. The mid-times of the occultations do not coincide due to the propagation delays of the shadow due to the distinct longitude of the sites. Exposures at INPE started after the immersion, as explained in the text.

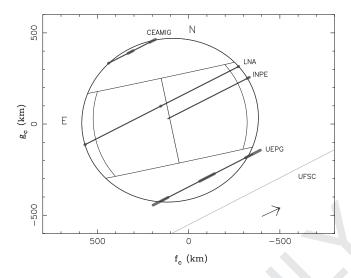


Figure 3. The best elliptical fit for the occultation chords for the event of 2010 using the times from Table 1 and the pole-constrained solution. The arrow indicates the direction of motion, blue lines are the observed chords, the red segments are the ingress, egress and mid-occultation error bars at 1σ level.

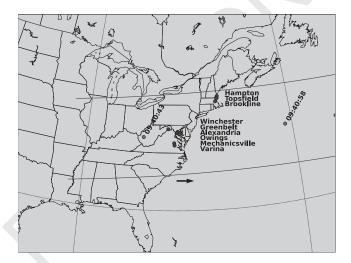


Figure 4. Post-occultation reconstruction of Ceres' shadow path on Earth for the 2013 October 25 event at the east coast of USA. Upper view of the occultation over the sites that observed the event (blue dots). Red points are the centre of the shadow separated by 15 s.

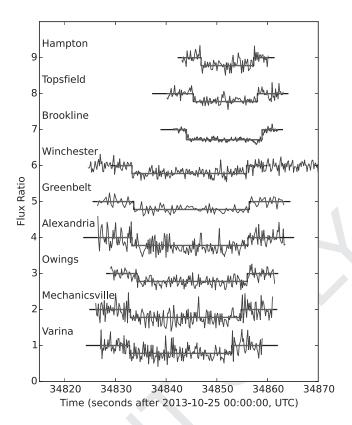


Figure 5. The nine occultation light curves normalized and vertically shifted by a factor of 1.0 for better viewing, see Fig. 2 for the explanation of the graph. The light curve of Brookline is shifted by -64 s as explained in the text.

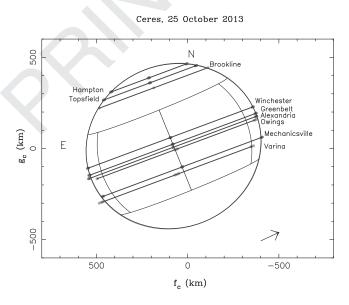


Figure 6. The best elliptical fit for the occultation chords for the event of 2013 using timing from Table 3 and the pole-constrained solution. The arrow indicates the direction of motion, blue lines are the observed chords, the red and green segments are the ingress, egress and mid-occultation error bars at 1σ level. The chords with green error bars were not used during the limb-fitting process. The chord of Brookline if shifted by -64 s as explained in the text.

List of astronomical key words (updated 2013 July)

This list is common to Monthly Notices of the Royal Astronomical Society, Astronomy and Astrophysics, and The Astrophysical Journal. In order to ease the search, the key words are subdivided into broad categories. No more than six subcategories altogether should be listed for a paper.

The subcategories in boldface containing the word 'individual' are intended for use with specific astronomical objects; these should never be used alone, but always in combination with the most common names for the astronomical objects in question. Note that each object counts as one subcategory within the allowed limit of six.

The parts of the key words in italics are for reference only and should be omitted when the key words are entered on the manuscript.

General

editorials, notices errata, addenda extraterrestrial intelligence history and philosophy of astronomy miscellaneous obituaries, biographies publications, bibliography sociology of astronomy standards

Physical data and processes

acceleration of particles accretion, accretion discs asteroseismology astrobiology astrochemistry astroparticle physics atomic data atomic processes black hole physics

chaos conduction convection dense matter diffusion dynamo elementary particles

equation of state

gravitation

gravitational lensing: strong gravitational lensing: weak gravitational lensing: micro gravitational waves hydrodynamics instabilities line: formation

line: identification line: profiles magnetic fields

magnetic reconnection

(magnetohydrodynamics) MHD

masers

molecular data

molecular processes

neutrinos

nuclear reactions, nucleosynthesis, abundances

opacity plasmas polarization

radiation: dynamics

radiation mechanisms:general radiation mechanisms: non-thermal radiation mechanisms: thermal

radiative transfer relativistic processes

scattering shock waves solid state: refractory

solid state: volatile turbulence

waves

Astronomical instrumentation, methods and techniques

atmospheric effects

balloons

instrumentation: adaptive optics instrumentation: detectors

instrumentation: high angular resolution instrumentation: interferometers

instrumentation: miscellaneous instrumentation: photometers instrumentation: polarimeters instrumentation: spectrographs

light pollution methods: analytical methods: data analysis methods: laboratory: atomic methods: laboratory: molecular

methods: laboratory: solid state methods: miscellaneous methods: numerical methods: observational methods: statistical

site testing space vehicles

space vehicles: instruments techniques: high angular resolution techniques: image processing techniques: imaging spectroscopy techniques: interferometric techniques: miscellaneous techniques: photometric techniques: polarimetric techniques: radar astronomy

techniques: radial velocities

techniques: spectroscopic

telescopes

Astronomical data bases

astronomical data bases: miscellaneous

atlases catalogues surveys

virtual observatory tools

Astrometry and celestial mechanics

astrometry

celestial mechanics

eclipses ephemerides occultations parallaxes proper motions reference systems

time

The Sun

Sun: abundances Sun: activity Sun: atmosphere Sun: chromosphere Sun: corona

Sun: coronal mass ejections (CMEs)

Sun: evolution
Sun: faculae, plages

Sun: filaments, prominences

Sun: flares

Sun: fundamental parameters

Sun: general
Sun: granulation
Sun: helioseismology
Sun: heliosphere
Sun: infrared
Sun: interior
Sun: magnetic fields

Sun: oscillations
Sun: particle emission
Sun: photosphere
Sun: radio radiation

Sun: rotation

(Sun:) solar-terrestrial relations

(Sun:) solar wind (Sun:) sunspots Sun: transition region Sun: UV radiation Sun: X-rays, gamma-rays

Planetary systems

comets: general

comets: individual: ...

Earth

interplanetary medium Kuiper belt: general

Kuiper belt objects: individual: ...

meteorites, meteors, meteoroids minor planets, asteroids: general

minor planets, asteroids: individual: ...

Moon Oort Cloud

planets and satellites: atmospheres planets and satellites: aurorae planets and satellites: composition planets and satellites: detection

planets and satellites: dynamical evolution and stability

planets and satellites: formation

planets and satellites: fundamental parameters

planets and satellites: gaseous planets planets and satellites: general planets and satellites: individual: ... planets and satellites: interiors

planets and satellites: interiors planets and satellites: magnetic fields planets and satellites: oceans

planets and satellites: physical evolution

planets and satellites: rings planets and satellites: surfaces planets and satellites: tectonics planets and satellites: terrestrial planets

planet–disc interactions`
planet–star interactions

zodiacal dust

protoplanetary discs

Stars

stars: abundances stars: activity

stars: AGB and post-AGB

stars: atmospheres

(stars:) binaries (including multiple): close

(stars:) binaries: eclipsing (stars:) binaries: general (stars:) binaries: spectroscopic (stars:) binaries: symbiotic (stars:) binaries: visual stars: black holes (stars:) blue stragglers (stars:) brown dwarfs

stars: carbon stars: chemically po

stars: chemically peculiar stars: chromospheres (stars:) circumstellar matter

stars: coronae stars: distances stars: dwarf novae stars: early-type stars: emission-line, Be stars: evolution

stars: flare stars: formation

stars: fundamental parameters (stars:) gamma-ray burst: general (stars:) gamma-ray burst: individual: ...

stars: general

(stars:) Hertzsprung-Russell and colour-magnitude diagrams

stars: horizontal branch stars: imaging stars: individual: ...

stars: interiors The Galaxy stars: jets Galaxy: abundances stars: kinematics and dynamics Galaxy: bulge stars: late-type Galaxy: centre stars: low-mass Galaxy: disc stars: luminosity function, mass function Galaxy: evolution stars: magnetars Galaxy: formation stars: magnetic field Galaxy: fundamental parameters stars: massive Galaxy: general stars: mass-loss (Galaxy:) globular clusters: general stars: neutron (Galaxy:) globular clusters: individual: ... (stars:) novae, cataclysmic variables Galaxy: halo stars: oscillations (including pulsations) Galaxy: kinematics and dynamics stars: peculiar (except chemically peculiar) (Galaxy:) local interstellar matter (stars:) planetary systems Galaxy: nucleus stars: Population II (Galaxy:) open clusters and associations: general stars: Population III (Galaxy:) open clusters and associations: individual: ... stars: pre-main-sequence (Galaxy:) solar neighbourhood stars: protostars Galaxy: stellar content (stars:) pulsars: general Galaxy: structure (stars:) pulsars: individual: ... stars: rotation Galaxies stars: solar-type (stars:) starspots galaxies: abundances galaxies: active stars: statistics (galaxies:) BL Lacertae objects: general (stars:) subdwarfs (galaxies:) BL Lacertae objects: individual: ... (stars:) supergiants galaxies: bulges (stars:) supernovae: general (stars:) supernovae: individual: ... galaxies: clusters: general galaxies: clusters: individual: ... stars: variables: Cepheids galaxies: clusters: intracluster medium stars: variables: δ Scuti galaxies: distances and redshifts stars: variables: general galaxies: dwarf stars: variables: RR Lyrae stars: variables: S Doradus galaxies: elliptical and lenticular, cD stars: variables: T Tauri, Herbig Ae/Be galaxies: evolution (stars:) white dwarfs galaxies: formation galaxies: fundamental parameters stars: winds, outflows galaxies: general stars: Wolf-Rayet galaxies: groups: general Interstellar medium (ISM), nebulae galaxies: groups: individual: ... ISM: abundances galaxies: haloes ISM: atoms galaxies: high-redshift galaxies: individual: ... ISM: bubbles ISM: clouds galaxies: interactions (galaxies:) intergalactic medium (ISM:) cosmic rays galaxies: irregular (ISM:) dust, extinction galaxies: ISM ISM: evolution galaxies: jets ISM: general (ISM:) HII regions galaxies: kinematics and dynamics (galaxies:) Local Group (ISM:) Herbig-Haro objects ISM: individual objects: ... galaxies: luminosity function, mass function (galaxies:) Magellanic Clouds (except planetary nebulae) galaxies: magnetic fields ISM: jets and outflows galaxies: nuclei ISM: kinematics and dynamics ISM: lines and bands galaxies: peculiar ISM: magnetic fields galaxies: photometry (galaxies:) quasars: absorption lines ISM: molecules (galaxies:) quasars: emission lines (ISM:) planetary nebulae: general (ISM:) planetary nebulae: individual: ... (galaxies:) quasars: general (ISM:) photodissociation region (PDR) (galaxies:) quasars: individual: ... ISM: structure (galaxies:) quasars: supermassive black holes ISM: supernova remnants galaxies: Seyfert

galaxies: spiral galaxies: starburst

galaxies: star clusters: general

galaxies: star clusters: individual: ...

galaxies: star formation galaxies: statistics galaxies: stellar content galaxies: structure

Cosmology

(cosmology:) cosmic background radiation (cosmology:) cosmological parameters

cosmology: miscellaneous cosmology: observations cosmology: theory

(cosmology:) dark ages, reionization, first stars

(cosmology:) dark energy (cosmology:) dark matter (cosmology:) diffuse radiation (cosmology:) distance scale (cosmology:) early Universe (cosmology:) inflation

(cosmology:) large-scale structure of Universe (cosmology:) primordial nucleosynthesis

Resolved and unresolved sources as a function of wavelength

gamma-rays: diffuse background

gamma-rays: galaxies

gamma-rays: galaxies: clusters

gamma-rays: general gamma-rays: ISM gamma-rays: stars

infrared: diffuse background

infrared: galaxies infrared: general infrared: ISM

infrared: planetary systems

infrared: stars

radio continuum: galaxies radio continuum: general radio continuum: ISM

radio continuum: planetary systems

radio continuum: stars radio lines: galaxies radio lines: general radio lines: ISM

radio lines: planetary systems

radio lines: stars

submillimetre: diffuse background

submillimetre: galaxies submillimetre: general submillimetre: ISM

submillimetre: planetary systems

submillimetre: stars ultraviolet: galaxies ultraviolet: general ultraviolet: ISM

ultraviolet: planetary systems

ultraviolet: stars X-rays: binaries X-rays: bursts

X-rays: diffuse background

X-rays: galaxies

X-rays: galaxies: clusters X-rays: general

X-rays: general
X-rays: individual: ...

X-rays: ISM X-rays: stars