# 1 Cooper-Frye integral for decay products

The spectrum of final decay products is given by

$$E_{\mathbf{p}}\frac{dN}{d^{3}\mathbf{p}} = \frac{\nu}{(2\pi)^{3}} \int_{\sigma} d\sigma_{\mu} g^{\mu} \tag{1}$$

where  $g^{\mu}$  is the vector distribution function of decayed particles.  $g^{\mu}$  is typically given as a sum of a couple Lorentz invariant weight functions and some explicit Lorentz vectors, which depends on what is the initial primary resonance distribution function. Below we document, which distribution components are calculated in FastReso. See the paper and code for further details of how the the decays are implemented.

## 2 Initial distribution functions

Code stores distribution components in an array TParticle::fFj\_arr, which is of size [grid\_params:fN\_f]  $\times$  [grid\_params:fN\_pbar], where  $N_f = 12$  is the total number of irreducible components and  $N_{\bar{\mathbf{p}}} = 201$  is the number of (un-equally spaced) momentum grid points. Note that some momentum independent terms are factored out to make the stored values simple and dimensionless. As of time of writing (2018.09.27) FastReso calculates the decay components of the following initial distribution functions.

## 2.1 Equilibrium distribution

Initial distribution is given by Fermi-Dirac or Bose-Einstein distribution (there is also an option for Boltzmann distribution which is never used)

$$f_{\text{eq}}(\bar{E}_{\mathbf{p}}, T, \mu) = \left(e^{(\bar{E}_{\mathbf{p}} - \mu)/T} \mp 1\right)^{-1},$$
 (2)

where  $\mu = Q_B \mu_B + Q_{I_3} \mu_{I_3} + Q_S \mu_S + Q_C \mu_C$  where the charges and corresponding chemical potentials are that of baryon number, third component of the isospint, strangeness and charm. The decayed vector distribution function is given by

$$g_b^{\mu}(\bar{E}_{\mathbf{p}}) = f_1^{\text{eq}}(\bar{E}_{\mathbf{p}}) \left( p^{\mu} - \bar{E}_{\mathbf{p}} u^{\mu} \right) + f_2^{\text{eq}}(\bar{E}_{\mathbf{p}}) \bar{E}_{\mathbf{p}} u^{\mu}.$$
 (3)

The actual data stored and printed is

$$f_1^{\text{eq}}(\bar{E}_{\mathbf{p}}) = \text{fFj\_arr[0]}, \tag{4}$$

$$f_2^{\text{eq}}(\bar{E}_{\mathbf{p}}) = \text{fFj\_arr[1]} \tag{5}$$

which is initialized by

$$fFj\_arr[0] = f_{eq}(\bar{E}_{\mathbf{p}}, T, \mu), \tag{6}$$

$$fFj_arr[1] = f_{eq}(\bar{E}_{\mathbf{p}}, T, \mu) \tag{7}$$

#### 2.2 Viscous bulk and shear distributions

Initial perturbations are given by (taken from [arXiv:1509.06738])

$$\delta f^{\text{bulk}}(x,p) = f_{\text{eq}}(1 \pm f_{\text{eq}}) \left[ \frac{\bar{E}_p}{T} \left( \frac{1}{3} - c_s^2 \right) - \frac{1}{3} \frac{m^2}{T\bar{E}_p} \right] \times \frac{\tau_{\text{II}}\Pi}{\zeta}, \tag{8}$$

$$\delta f^{\text{shear}}(x,p) = f_{\text{eq}}(1 \pm f_{\text{eq}}) \times \pi_{\rho\nu} p^{\rho} p^{\nu} \times \frac{1}{2(e+p)T^2}.$$
 (9)

where  $c_s$  is the speed of sound on the freeze-out surface (given as a parameter) and m is a primary resonance mass. The decayed vectro distribution function is given by

$$g_{\text{bulk}}^{\mu}\Pi = \left(p^{\mu} - \bar{E}_{\mathbf{p}}u^{\mu}\right)f_1^{\text{bulk}}(\bar{E}_{\mathbf{p}}) + \bar{E}_{\mathbf{p}}u^{\mu}f_2^{\text{bulk}}(\bar{E}_{\mathbf{p}}),\tag{10}$$

and

$$g_{\text{shear}}^{\mu\nu\rho}\pi_{\nu\rho} = [p^{\nu}\pi_{\nu\rho}p^{\rho}(p^{\mu} - \bar{E}_{\mathbf{p}}u^{\mu}) - \frac{2}{5}|\bar{\mathbf{p}}|^{2}p_{\nu}\pi^{\nu\mu}]f_{1}^{\text{shear}}(\bar{E}_{\mathbf{p}}) + \frac{2}{5}|\bar{\mathbf{p}}|^{2}p_{\nu}\pi^{\nu\mu}f_{2}^{\text{shear}}(\bar{E}_{\mathbf{p}}) + \pi_{\nu\rho}p^{\nu}p^{\rho}\bar{E}_{\mathbf{p}}u^{\mu}f_{3}^{\text{shear}}(\bar{E}_{\mathbf{p}}).$$
(11)

The actual data stored and printed is

$$f_1^{\text{shear}}(\bar{E}_{\mathbf{p}}) = \text{fFj\_arr[2]} \times \frac{1}{2(e+p)T^2}, \tag{12}$$

$$f_2^{\text{shear}}(\bar{E}_{\mathbf{p}}) = \text{fFj\_arr[3]} \times \frac{1}{2(e+p)T^2}, \tag{13}$$

$$f_3^{\text{shear}}(\bar{E}_{\mathbf{p}}) = \text{fFj\_arr[4]} \times \frac{1}{2(e+p)T^2},\tag{14}$$

$$f_1^{\mathrm{bulk}}(\bar{E}_{\mathbf{p}}) = \mathtt{fFj\_arr[5]} \times \frac{-\tau_{\Pi}\Pi}{\zeta},$$
 (15)

$$f_2^{\mathrm{bulk}}(\bar{E}_{\mathbf{p}}) = \mathtt{fFj\_arr[6]} \times \frac{-\tau_{\Pi}\Pi}{\zeta},$$
 (16)

(17)

which is initialized by

$$fFj_{arr}[2] = f_{eq}(\bar{E}_{p}, T, \mu)(1 \pm f_{eq}(\bar{E}_{p}, T, \mu)), \tag{18}$$

$$fFj_{arr}[3] = f_{eq}(\bar{E}_{p}, T, \mu)(1 \pm f_{eq}(\bar{E}_{p}, T, \mu)), \tag{19}$$

$$fFj_{arr}[4] = f_{eq}(\bar{E}_{p}, T, \mu)(1 \pm f_{eq}(\bar{E}_{p}, T, \mu)), \tag{20}$$

$$fFj_{arr}[5] = f_{eq}(\bar{E}_{\mathbf{p}}, T, \mu)(1 \pm f_{eq}(\bar{E}_{\mathbf{p}}, T, \mu)) \left[ \frac{1}{3} \frac{m^2}{T\bar{E}_p} - \frac{\bar{E}_p}{T} \left( \frac{1}{3} - c_s^2 \right) \right], \tag{21}$$

$$\texttt{fFj\_arr[6]} = f_{eq}(\bar{E}_{\mathbf{p}}, T, \mu) (1 \pm f_{eq}(\bar{E}_{\mathbf{p}}, T, \mu)) \left[ \frac{1}{3} \frac{m^2}{T\bar{E}_p} - \frac{\bar{E}_p}{T} \left( \frac{1}{3} - c_s^2 \right) \right] \tag{22}$$

Note the sign in the bulk perturbation. Also, because of the explicit shear components carry  $|\bar{p}|^2$ , the irreducible weight function after decays are actually diverging at the origin, but no more than quadratically.

## 2.3 Temperature and velocity perturbations

Linear perturbations in temperature and velocity

$$\delta f^{\text{temp.}}(x,p) = f_{\text{eq}}(1 \pm f_{\text{eq}}) \frac{\bar{E}_{\mathbf{p}}}{\bar{T}} \frac{\delta T}{\bar{T}}, \tag{23}$$

$$\delta f^{\text{velocity}}(x,p) = f_{\text{eq}}(1 \pm f_{\text{eq}}) \frac{\delta u_{\nu} p^{\nu}}{\bar{T}}.$$
 (24)

Vector distribution function is then given by

$$g_{\text{temp}}^{\mu} \delta T = f_1^{\text{temp.}}(\bar{E}_{\mathbf{p}}) \left( p^{\mu} - \bar{E}_{\mathbf{p}} u^{\mu} \right) + f_2^{\text{temp.}}(\bar{E}_{\mathbf{p}}) \bar{E}_{\mathbf{p}} u^{\mu}. \tag{25}$$

and

$$g_{\text{velocity}}^{\mu\nu} \delta u_{\nu} = \left[\delta u_{\nu} p^{\nu} (p^{\mu} - \bar{E}_{\mathbf{p}} u^{\mu}) - \frac{1}{3} |\bar{\mathbf{p}}|^{2} \delta u^{\mu}\right] f_{1}^{\text{velocity}}(\bar{E}_{\mathbf{p}})$$

$$+ \frac{1}{3} |\bar{\mathbf{p}}|^{2} \delta u^{\mu} f_{2}^{\text{velocity}}(\bar{E}_{\mathbf{p}}) + \delta u_{\nu} p^{\nu} \bar{E}_{\mathbf{p}} u^{\mu} f_{3}^{\text{velocity}}(\bar{E}_{\mathbf{p}}).$$

$$(26)$$

The actual data stored and printed is

$$f_1^{\text{temp.}}(\bar{E}_{\mathbf{p}}) = \text{fFj\_arr[7]} \times \frac{\delta T}{T},$$
 (27)

$$f_2^{\text{temp.}}(\bar{E}_{\mathbf{p}}) = \text{fFj\_arr[8]} \times \frac{\delta T}{T},$$
 (28)

$$f_1^{\rm velocity}(\bar{E}_{\mathbf{p}}) = \mathtt{fFj\_arr[9]} \times \frac{1}{T}, \tag{29}$$

$$f_2^{\mathrm{velocity}}(\bar{E}_{\mathbf{p}}) = \mathtt{fFj\_arr[10]} \times \frac{1}{T},$$
 (30)

$$f_3^{\mathrm{velocity}}(\bar{E}_{\mathbf{p}}) = \mathtt{fFj\_arr[11]} \times \frac{1}{T},$$
 (31)

(32)

Note that irreducible components for velocity perturbation can also be divergent, but the explicit momentum prefactors will make it finite.

Initialization is given by

$$\texttt{fFj\_arr[7]} = f_{\text{eq}}(\bar{E}_{\mathbf{p}}, T, \mu) (1 \pm f_{\text{eq}}(\bar{E}_{\mathbf{p}}, T, \mu)) \frac{\bar{E}_{\mathbf{p}}}{T}, \tag{33}$$

$$fFj_{arr}[8] = f_{eq}(\bar{E}_{p}, T, \mu)(1 \pm f_{eq}(\bar{E}_{p}, T, \mu))\frac{E_{p}}{T},$$
(34)

$$fFj_{arr}[9] = f_{eq}(\bar{E}_{p}, T, \mu)(1 \pm f_{eq}(\bar{E}_{p}, T, \mu)), \tag{35}$$

$$fFj_{arr}[10] = f_{eq}(\bar{E}_{p}, T, \mu)(1 \pm f_{eq}(\bar{E}_{p}, T, \mu)), \tag{36}$$

$$fFj_{arr}[11] = f_{eq}(\bar{E}_{p}, T, \mu)(1 \pm f_{eq}(\bar{E}_{p}, T, \mu))$$
(37)

# 3 Header of output files

```
#
     Name
              Mass
                      Gamma Spin Isospin I3 Nq Ns Naq Nas Nc
Nac
         MC
               Nyield
# pi0139plu 0.139570 0.00000e+00 0.0 1.0 1.0 1 0 1 0 00
                                                          211 0.000701
# 1:pbar [GeV] 2:m [GeV]
                               3:feq 1
                                             4:feq 2 5:fshear 1
                                                                       6:fshear 2
  7:fshear 3
             8:fbulk 1
                              9:fbulk 2 10:ftemperature 1 11:ftemperature 2
12:fveloc
ity 1 13:fvelocity 2 14:fvelocity 3
```