Restoration from Multispectral Blurred Data with Non-Stationary Instrument Response

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Abstract—In this paper we propose an approach of image restoration from multispectral data provided by an imaging system. We specifically address two topics: (i) Development of a multi-wavelength direct model for non-stationary instrument response that includes a spatial convolution and a spectral integration, (ii) Implementation of multispectral image restoration using a regularized least-square, based on a quadratic criterion and minimized by a gradient algorithm. We test our approach on simulated data of the Mid-InfraRed Instrument IMager (MIRIM) of the James Webb Space Telescope (JWST). Our method shows a clear increase of spatial resolution compare to conventional methods.

Index Terms—Direct Model, Multispectral Imaging, Inverse **Problems, Image Restoration**

I. Introduction

Multispectral imaging instrument are used in many fields, e.g. remote sensing [1], medicine [2], astrophysics [3]. Imaging system on board space telescopes are specific subject of this work, especially the Mid-InfraRed Instrument IMager (MIRIM) [4] on board the James Webb Space Telescope (JWST)¹ which is the next space telescope of NASA, ESA and the Canadian Space Agency (CSA) to be launched at the end of 2018. Such a system is mainly composed of an optical system (or optic) that gathers source's light (or object) from space and provides it to the imaging instrument, which contains a spectral filter (or photometric band) to select the wavelength band of interest and an infrared detector that integrates and discretizes the light, providing a 2D discrete image data.

Due to its passage through the optics, photon flux of the object of interest is diffracted on the focal plane of the telescope. The result is a spatial 2D-convolution of the object with the optic response known as Point Spread Function (PSF) [5]. On the other hand, the PSF vary accordingly to the wavelength, blurring the object of interest and limiting its spatial resolution. The second problem occurs during the integration of the object by the detector, where the continuous $2D+\lambda$ object is integrated spectrally over the filter bandpass, the 2D discrete image data has thus no spectral dimension.

Conventional approaches consider measured PSFs [6], [4] but only PSFs at specific wavelengths are measured. Other

1https://jwst.nasa.gov/

approaches assume a broadband PSF, by averaging available monochromatic PSFs, weighted by the filter+detector spectral response and the object spectrum to observe [7]. Broadband PSF smooths the structure of the PSF and makes it dependent on the object spectrum to observe, especially for filters with wide bands. In addition, a monochromatic object spectrum is assumed for unknown sources, which results in a stationary PSF and leads to inaccurate optics response and direct model. Spatially-variant PSF has also been studied in [8] which approximates the spatial-variant PSF using PSF-interpolation. The same idea was used in [9] to approximate the spectralvariant PSF. More recently, [10] reviews and provides models of spatial-variant PSF. Other works treated deconvolution problem with unknown parameter of the PSF or unknown PSF, known as myope or blind deconvolution [11]. In the other hand, the spectral integration is mostly approached by a spectral convolution, such as in [12] where the 3D-PSF is separated into spectral and spatial-invariant PSF, and [9] by using PSF-interpolation so that the direct model is a 2D+ λ convolution. In this paper we deal with the continuous spectral integration by taking into account informations from available multispectral data of multi-filter, and preserve the spectral non-stationarity of PSF using PSF simulator. This work aims to restore a discrete version of the original spatio-spectral object from available multispectral discrete data. Our main contribution is the development of an instrument model and a direct model for an imaging system with a non-stationary instrument response.

The paper is organized as follows. In Sec. II, we develop the instrument model. Then we present the direct model along with the restoration in Sec. III. Experimental results on simulated data with an application to the instrument MIRIM are presented and discussed in Sec. IV. Finally, we conclude our work and provide perspectives in Sec. V.

II. INSTRUMENT MODEL

In this section we develop the instrument model for the imager. The block diagram of the instrument is represented in Fig. 1, and is composed of an optic response, spectral filtering and detector integration. The object of interest at the entrance of the acquisition system is a 3D object, $\phi(\alpha, \beta, \lambda) : \mathbb{R}^3 \to \mathbb{R}$, having two spatial dimensions $(\alpha, \beta) \in \mathbb{R}^2$ (angles in radian)

and one spectral dimension $\lambda \in \mathbb{R}^+$ (in microns). During the observation process with the filter f ($f \in [1, n_f]$ where n_f stands for total number of filters), the object is modified by the instrument components, providing 2D-discrete data $y^{(f)}$.

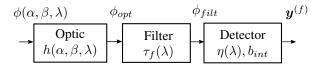


Fig. 1. Block diagram of the instrument model of the imager.

A. Instrument Response

1) Optic: The effect of the optic is carried out by the (α, β) -convolution of the input object $\phi(\alpha, \beta, \lambda)$ with the spectral-variant PSF $h(\alpha, \beta, \lambda)$ as follows

$$\phi_{opt}(\alpha, \beta, \lambda) = \phi(\alpha, \beta, \lambda) *_{(\alpha, \beta)} h(\alpha, \beta, \lambda)$$

$$= \iint_{\mathbb{R}^2} \phi(\alpha', \beta', \lambda) h(\alpha - \alpha', \beta - \beta', \lambda) d\alpha' d\beta',$$
(1)

where $\underset{(\alpha,\beta)}{*}$ stands for 2D spatial convolution.

Fig. 2 illustrates the non-stationarity of the PSF. We display two simulated PSFs of the instrument JWST/MIRIM at different wavelengths from the mid-infrared range, $\lambda_1 = 7.7 \mu m$ and $\lambda_2 = 25.25 \mu m$, PSF $_{\lambda_2}$ is clearly larger than PSF $_{\lambda_1}$.

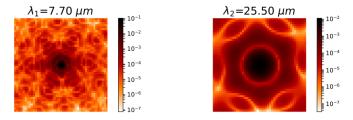


Fig. 2. Monochromatic PSFs simulated at $\lambda_1=7.7\,\mu m$ and $\lambda_2=25.25\,\mu m$, using the simulation tool WebbPSF [13] for the instrument MIRIM on board the JWST.

2) Filter: The object passes through the filter f in order to select the wavelength range of interest. This operation is performed through the filter response $\tau_f(\lambda)$ as follows

$$\phi_{filt}^{(f)}(\alpha, \beta, \lambda) = \tau_f(\lambda) \,\phi_{opt}(\alpha, \beta, \lambda). \tag{2}$$

- 3) Detector: The detector integrates the filtered object and performs sampling. We model its response in two steps:
 - The spectral integration is weighted by the spectral response of the detector (or quantum efficiency) $\eta(\lambda)$,

$$y^{(f)}(\alpha,\beta) = \int_{\mathbb{R}_+} \eta(\lambda) \phi_{filt}^{(f)}(\alpha,\beta,\lambda) d\lambda.$$
 (3)

• We introduce a basis function b_{int} to perform spatial sampling within spatial integration, e.g. indicator function

of a sensitive area of the pixel Ω_{pix} . The discrete data of the filter f and pixel (i, j) is

$$y_{i,j}^{(f)} = \iint_{\Omega_{pix}} y^{(f)}(\alpha, \beta) b_{\text{int}} \left(\alpha - \alpha_{i,j}, \beta - \beta_{i,j}\right) d\alpha d\beta.$$
(4)

with $\alpha_{i,j}$ and $\beta_{i,j}$ the angular directions of the pixel position (i,j) defined on the detector grid. We denote N the total number of pixels.

B. Complete Formulation

The complete formulation of the instrument model is obtained by substituting Eqs.(1)-(2) and (3) in Eq.(4). This yields

$$y_{i,j}^{(f)} = \int_{\mathbb{R}_{+}} \eta(\lambda) \tau_{f}(\lambda) \left(\iint_{\Omega_{pix}} \phi(\alpha, \beta, \lambda) *_{(\alpha, \beta)} h(\alpha, \beta, \lambda) \right) d\alpha d\beta d\lambda.$$

$$b_{\text{int}}(\alpha - \alpha_{i,j}, \beta - \beta_{i,j}) d\alpha d\beta d\lambda.$$
(5)

This model links the continuous 3D object $\phi(\alpha, \beta, \lambda)$ at the entrance of the imaging system to the 2D discrete data $y^{(f)}$ through the instrument response, which includes spectral windowing and five integrations, two for spatial 2D-convolution, two for spatial integration and one for spectral integration. The developed instrument model takes into account the non-stationarity of the instrument response through the spectral-variant PSF. Moreover, the above model does not include non-ideal characteristics of the detector [14]. All these effects are assumed to be corrected through the pipeline stages of the data reduction plan [15]. In this paper, we consider the same sampling grid of the object as for the data, and that all pixels are regularly disposed on the detector grid, having the same area Ω_{vix} .

In order to simplify the instrument model, we consider a constant object over pixel area. Thus, we define $b_{\rm int}$ as a rectangular function over the pixel area. The impact of this approximation is slightly important for detectors with good resolution, such as MIRIM detector [16]. Thus, Eq.(5) becomes

$$y_{i,j}^{(f)} = \Omega_{pix} \int_{\mathbb{R}_{+}} \eta(\lambda) \tau_{f}(\lambda)$$

$$\phi\left(\alpha_{i,j}, \beta_{i,j}, \lambda\right) \underset{\left(\alpha_{i,j}, \beta_{i,j}\right)}{*} h\left(\alpha_{i,j}, \beta_{i,j}, \lambda\right) d\lambda \qquad (6)$$

where $*_{(\alpha_{i,j},\beta_{i,j})}$ stands for discrete 2D-convolution.

III. DIRECT MODEL AND INVERSE PROBLEM

In this section we first model the object spectrum, then we develop the direct model, thereafter we tackle the inverse problem of the object.

A. Continuous Object Spectrum

In this paper we model the spectrum object with a continuous piecewise linear function, as shown in Fig.(3). This choice allows us to obtain a simple model that preserves the spectral distribution of the object with less complexity, whereas conventional approaches do not. They generally consider a

model with broadband PSF, defined as a spectral integration of monochromatic PSF weighted by the filter + detector + object spectrum window [7]. Hence, one can only attempt to restore a spatial distribution with such a model [17]. The object is modeled as follows

$$\phi(\alpha_{i,j}, \beta_{i,j}, \lambda) = \sum_{b=1}^{n_b} \left(x^{(b)} \left(\alpha_{i,j}, \beta_{i,j} \right) g_+^{(b)}(\lambda) + x^{(b-1)} \left(\alpha_{i,j}, \beta_{i,j} \right) g_-^{(b)}(\lambda) \right) \mathbb{1}_{\left[\lambda^{(b-1)}, \lambda^{(b)} \right]}(\lambda),$$
(7)

where $b \in [1,...,n_b]$ is the index of the bandwidth, $\mathbbm{1}(\lambda)$ the indicator function, which is equal to 1 for $\lambda \in [\lambda^{(b-1)},\lambda^{(b)}]$ and 0 otherwise. The discrete parameters $x_{i,j}^{(b-1)}$ and $x_{i,j}^{(b)}$ are intensities of the object at $\lambda^{(b-1)}$ and $\lambda^{(b)}$, respectively, $g_-^{(b)}(\lambda) = \frac{1}{2} - \frac{\lambda - \lambda_c^{(b)}}{\lambda^{(b)} - \lambda^{(b-1)}}$ and $g_+^{(b)}(\lambda) = \frac{1}{2} + \frac{\lambda - \lambda_c^{(b)}}{\lambda^{(b)} - \lambda^{(b-1)}}$ express the linear slope of the object spectrum.

Thanks to the parameterization in Eq. (7), the object spectrum is continuous in λ and the positivity constraint is naturally fulfilled.

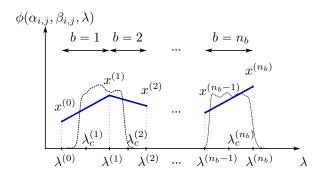


Fig. 3. Representation of a piecewise linear spectrum, for a single position over a wavelength range of $[\lambda^{(0)}, \lambda^{(n_b)}]$. (in dots) Filters transmission to illustrate the multi-filter analysis.

B. Direct model

The discrete data $y^{(f)}$ contains information about the object of interest within the filter bandpass. In fact, we use a multifilter analysis to develop a direct model that accounts for informations of the whole wavelength range of the instrument. By substituting Eq.(7) in Eq.(6) and rearranging similar terms, only wavelength-dependent terms are left in the spectral integration without the object parameter x. This yields

$$y_{i,j}^{(f)} = \sum_{b=0}^{n_b} h_{\text{int}}^{(f,b)}(\alpha_{i,j}, \beta_{i,j}) \underset{(\alpha_{i,j}, \beta_{i,j})}{*} x^{(b)}(\alpha_{i,j}, \beta_{i,j})$$
(8)

where $h_{\text{int}}^{(f,b)}$ is a 2D integrated PSF

$$h_{\text{int}}^{(f,b)}(\alpha_{i,j}, \beta_{i,j}) = \Omega_{pix} \int_{\mathbb{R}_+} \eta(\lambda) \tau_f(\lambda) \left[g_-^{(b+1)}(\lambda) + g_+^{(b)}(\lambda) \right] h(\alpha_{i,j}, \beta_{i,j}, \lambda) d\lambda,$$

with the boundary conditions $g_{-}^{(n_b+1)}(\lambda) = g_{+}^{(0)}(\lambda) = 0$.

Eq. (8) can be formulated in a vector form as follows

$$\mathbf{y}^{(f)} = \sum_{b=0}^{n_b} \mathbf{H}_{int}^{(f,b)} \mathbf{x}^{(b)} + \epsilon^{(f)},$$
 (10)

where $\boldsymbol{H}_{\text{int}}^{(f,b)} \in \mathbb{R}^{N \times N}$ is a convolution matrix. $\boldsymbol{x}^{(b)} \in \mathbb{R}^{N}$ and $\boldsymbol{y}^{(f)} \in \mathbb{R}^{N}$ are vector representations of $y^{(f)}$ and $x^{(b)}$. $\boldsymbol{\epsilon}^{(f)} \in \mathbb{R}^{N}$ corresponds to readout noise and modeling error.

The linear direct model in Eq. (10) is a sum of $(n_b + 1)$ spatial convolutions, between the 2D object parameters $\boldsymbol{x}^{(b)}$ and 2D system response $\boldsymbol{H}_{\text{int}}^{(f,b)}$. Moreover, the observation $\boldsymbol{y}^{(f)}$ contains information about the unknown object over the whole spectral range, weighted by the spectral windows $\eta(\lambda)\tau_f(\lambda)$. By combining all available multispectral data in one vector \boldsymbol{y} , Eq. (10) becomes

$$y = Hx + \epsilon, \tag{11}$$

where $\boldsymbol{H} = \left\{\boldsymbol{H}_{\mathrm{int}}^{(f,b)}\right\}_{f=1:n_f}$ is a block matrix containing $n_f \times (n_b+1)$ convolution matrices corresponding to the integrated PSFs $h_{\mathrm{int}}^{(f,b)}$. We denote by $\boldsymbol{x} = \left[\boldsymbol{x}^{(0)} \dots \boldsymbol{x}^{(n_b)}\right]^t$ the vector of (n_b+1) unknowns object parameters, $\boldsymbol{y} = \left[\boldsymbol{y}^{(1)} \dots \boldsymbol{y}^{(n_f)}\right]^t$ vector of all available data, and $\boldsymbol{\epsilon} = \left[\boldsymbol{\epsilon}^{(1)} \dots \boldsymbol{\epsilon}^{(n_f)}\right]^t$ vector of additive noise associated to data.

C. Restoration

Restoration of x relies on regularized least-square approach in inverse problems [18], where the solution is obtained by minimizing a criterion \mathcal{J} as follows

$$\hat{x} = \underset{\sim}{\operatorname{argmin}} \mathcal{J}(x).$$
 (12)

The criterion of least-square (called data fidelity) is $\| \boldsymbol{y} - \boldsymbol{H} \boldsymbol{x} \|_2^2$. It enforces agreement of the solution $\hat{\boldsymbol{x}}$ with the data \boldsymbol{y} . Least-square problem is ill-posed because of the ill-conditioning of \boldsymbol{H} . In fact, regularized least-square estimator aims of correcting this ill-posedness by adding a regularization term $\| \boldsymbol{C} \boldsymbol{x} \|_2^2$ to the criterion, with \boldsymbol{C} a well-conditioned matrix. We are particularly interested in a quadratic-regularization (l_2 -norm) in order to obtain a differentiable criterion. This allow us to implement a fast calculation of the solution. Therefore, the regularized least-square criterion is

$$\mathcal{J}(x) = \|y - Hx\|_{2}^{2} + \|Cx\|_{2}^{2}, \tag{13}$$

the matrix $C = \operatorname{diag} \left\{ \sqrt{\mu_0} D_{\alpha,\beta}, \ldots, \sqrt{\mu_{n_b}} D_{\alpha,\beta} \right\}$, where $D_{\alpha,\beta}$ is a 2D spatial constraint operator, e.g. discrete Laplacian operator in case of a spatially smooth object. Moreover, regularization parameters μ_0, \ldots, μ_{n_b} are set to compromise between the two terms of the criterion, data fidelity and regularization.

As \mathcal{J} is a quadratic form, solution of the problem \hat{x} is explicit and obtained by canceling the gradient of the criterion defined in Eq.(13):

$$\hat{\boldsymbol{x}} = (\underbrace{\boldsymbol{H}^t \boldsymbol{H} + \boldsymbol{C}^t \boldsymbol{C}}_{\boldsymbol{Q}})^{-1} \boldsymbol{H}^t \boldsymbol{y}$$
 (14)

where Q is a Hessian matrix of size $(n_b+1)N\times(n_b+1)N$. We first attempt to compute the solution by inverting Q using diagonalization of circular matrices in Fourier space, therefore, computing Q^{-1} means inverting N square matrices of size $(n_b+1)\times(n_b+1)$. This calculation turns out heavy for $n_b>4$. We propose instead to compute the solution without inverting Q, i.e. by solving a linear system $(H^tH+C^tC)\hat{x}=H^ty$ through an iterative scheme, e.g. using numerical optimization algorithm such as the conjugated gradient (CG) algorithm [19]. Moreover, we consider circular convolution and compute the solution in Fourier space for efficient computation.

IV. SIMULATION RESULTS

A. Application to JWST/MIRIM

We apply the proposed approach to the instrument MIRIM on board the space telescope JWST. The optical system of the telescope is mainly composed of a 6.5-meters primary mirror, made up of 18-hexagonal segmented mirror. The imaging instrument MIRIM has nine photometric bands that cover the mid-infrared range of 5 to $28\mu m$. Their transmission profiles are given in [4]. The object of interest is guided through mirrors to the MIRIM detector [16], where it will be integrated and sampled. MIRIM detector has a pixel scale of $0.11~{\rm arcsec/pixel}~(\Omega_{pix}=0.11^2~{\rm arcsecond^2})$ and its spectral window $\eta(\lambda)\tau_f(\lambda)$ is known as Photon Conversion Efficiency (PCE) [20], see Fig. 4. Moreover, non-stationary PSF is a set of monochromatic PSFs simulated on a discrete wavelength grid that covers the whole spectral range of the instrument, simulated with WebbPSF [13] the official PSF simulator of mission JWST.

B. Setup

In order to validate the proposed approach, all tests were done on simulated data, where we simulated the original object as a cube of size $64 \times 64 \times 9$ $(N = 64, n_b = 8)$ with a Gaussian spatial distribution ($\sigma = 3$) and a linear spectrum over the instrument spectral range. We implemented the direct model in Eq.(11) by computing the nine integrated PSFs and performing convolutions with original object parameters. As a result, nine multispectral data are simulated corresponding to nine MIRIM photometric bands. Signal-to-noise ratio (SNR) [21] is used to evaluate the level of the white Gaussian zero-mean noise added to the data. We implement the conjugate gradient algorithm to restore the object parameter \hat{x} and compare it to the conventional approach given in Sec. III-A, referenced as as "standard". Object restoration in case of the standard approach consists of deconvolution of nine data by considering computed broadband PSFs. We use the unsupervised deconvolution method of [22], with a quadratic-regularization and same prior. We obtain deconvolved observations as well as estimation of regularization parameters, these latter are used for both approaches, proposed and standard ones. Finally, the restored object of interest $\hat{\phi}$ is deducted from the restored object parameter \hat{x} using Eq.(7). Restoration quality is measured with a relative error $\left\|\phi_{\text{original}} - \hat{\phi}\right\|_2 / \left\|\phi_{\text{original}}\right\|_2$ between the original object of interest ϕ_{original} and the restored object $\hat{\phi}$.

C. Discussion

Fig. 4 illustrates restoration results in a spectral dimension, where we display the spectrum of the central pixel of the original object, proposed and standard restorations. All data are corrupted by an additive zero-mean Gaussian noise of 30 dB. We observe that the restored spectrum using the proposed approach fits very well the original spectrum along the whole wavelength range. However, the standard approach fails to do so, because of the inaccuracy of the instrument response considered in the standard approach. In addition, Fig. 5 shows observed data using filter f = 8 (named F2100W), the original object at $\lambda = 18.7 \mu m$ and the spatial restoration with both approaches. The restoration result using our approach shows a good restoration of the dynamic and the spatial details, the blur caused by the PSF is also removed, with an error of 6.27% compared to 21.29% obtained with the standard approach. Moreover, we notice that the wider the filter band, the better the results of our approach compared to the conventional approach, which is the case for the instrument MIRIM.

V. CONCLUSION

In this paper we present the problem of image restoration from multispectral data acquired by an imager, where we restored a spatio-spectral object. An instrument model with a complex non-stationary response is detailed, including a spatial convolution with a spectral-variant PSF and a spectral integration. Then, we developed a multi-wavelength direct model, which is the sum of spatial 2D-convolution of the object parameter with an integrated PSF. The idea relies on using a multi-filter processing and an approximation of a continuous piecewise linear spectrum. Multispectral restoration is implemented using regularized least-square based on quadratic criterion. Preliminary restoration results are provided on simulated data, where we obtained a clear improvement of restoration quality compared to classical approach.

Several aspects are to improve in our approach, such as the estimation of the (n_b+1) regularization parameters. As future work, generalization of the direct model by introducing a set of basis functions of the object discretization, and increasing the spectral sampling of the sought object by adding spectral prior of the object. Finally, validation of direct model and restoration of real astrophysical object.

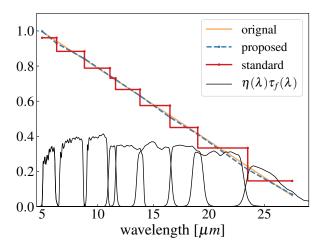


Fig. 4. Illustration of the spectral restoration. All intensities were normalized to 1 for better visualization and comparison.

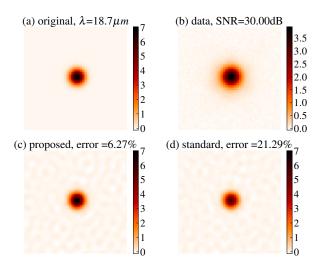


Fig. 5. Illustration of the spatial restoration : (a) original object at $\lambda=18.7\mu m$. (b) data for the filter F2100W. (c) proposed restoration, (d) standard restoration. Parameters used are : SNR = 30 dB, $\sigma_{\rm Gauss}=3$

REFERENCES

- D. A. Landgrebe, Signal theory methods in multispectral remote sensing. John Wiley & Sons, 2005, vol. 29.
- [2] S. Ozer, D. L. Langer, X. Liu, M. A. Haider, T. H. van der Kwast, A. J. Evans, Y. Yang, M. N. Wernick, and I. S. Yetik, "Supervised and unsupervised methods for prostate cancer segmentation with multispectral mri," *Medical physics*, vol. 37, no. 4, pp. 1873–1883, 2010.
- [3] R. Lionello, J. A. Linker, and Z. Mikić, "Multispectral emission of the sun during the first whole sun month: Magnetohydrodynamic simulations," *The Astrophysical Journal*, vol. 690, no. 1, p. 902, 2008.
- [4] P. Bouchet, M. García-Marín, P.-O. Lagage, J. Amiaux, J.-L. Auguéres, E. Bauwens, J. Blommaert, C. Chen, Ö. Detre, D. Dicken et al., "The Mid-Infrared Instrument for the James Webb Space Telescope, III: MIRIM, The MIRI Imager," Publications of the Astronomical Society of the Pacific, vol. 127, no. 953, p. 612, 2015.
- [5] J. W. Goodman, Introduction to Fourier optics. Roberts and Company Publishers, 2005.
- [6] P. Guillard, T. Rodet, S. Ronayette, J. Amiaux, A. Abergel, V. Moreau, J. Augueres, A. Bensalem, T. Orduna, C. Nehmé et al., "Optical performance of the JWST/MIRI flight model: characterization of the point spread function at high resolution," in SPIE Astronomical Telescopes+

- *Instrumentation*. International Society for Optics and Photonics, 2010, pp. 77310J–77310J.
- [7] N. Geis and D. Lutzn, "Herschel/PACS modelled point-spread functions."
- [8] L. Denis, É. Thiébaut, and F. Soulez, "Fast model of space-variant blurring and its application to deconvolution in astronomy," in *Image Processing (ICIP)*, 2011 18th IEEE International Conference on. IEEE, 2011, pp. 2817–2820.
- [9] F. Soulez, E. Thiébaut, and L. Denis, "Restoration of hyperspectral astronomical data with spectrally varying blur," *EAS Publications Series*, vol. 59, pp. 403–416, 2013.
- [10] É. Thiébaut, L. Dénis, F. Soulez, and R. Mourya, "Spatially variant psf modeling and image deblurring," in SPIE Astronomical Telescopes+ Instrumentation. International Society for Optics and Photonics, 2016, pp. 99 097N–99 097N.
- [11] R. Molina, J. Mateos, and A. K. Katsaggelos, "Blind deconvolution using a variational approach to parameter, image, and blur estimation," *IEEE Transactions on Image Processing*, vol. 15, no. 12, pp. 3715–3727, 2006.
- [12] E. Villeneuve and H. Carfantan, "Nonlinear deconvolution of hyperspectral data with meme for studying the kinematics of galaxies," *IEEE Transactions on Image Processing*.
- [13] M. D. Perrin, R. Soummer, E. M. Elliott, M. D. Lallo, and A. Sivara-makrishnan, "Simulating point spread functions for the James Webb Space Telescope with WebbPSF," in SPIE Astronomical Telescopes+Instrumentation. International Society for Optics and Photonics, 2012, pp. 84 423D–84 423D.
- [14] G. Wright, D. Wright, G. Goodson, G. H. Rieke, G. Aitink-Kroes, J. Amiaux, A. Aricha-Yanguas, R. Azzollini, K. Banks, D. Barrado-Navascues et al., "The mid-infrared instrument for the james webb space telescope, ii: Design and build," *Publications of the Astronomical Society of the Pacific*, vol. 127, no. 953, p. 595, 2015.
- [15] K. D. Gordon, C. Chen, R. E. Anderson, R. Azzollini, L. Bergeron, P. Bouchet, J. Bouwman, M. Cracraft, S. Fischer, S. D. Friedman et al., "The Mid-Infrared Instrument for the James Webb Space Telescope, X: Operations and Data Reduction," Publications of the Astronomical Society of the Pacific, vol. 127, no. 953, p. 696, 2015.
- [16] G. Rieke, M. Ressler, J. E. Morrison, L. Bergeron, P. Bouchet, M. García-Marín, T. Greene, M. Regan, K. Sukhatme, and H. Walker, "The Mid-Infrared Instrument for the James Webb Space Telescope, VII: The MIRI Detectors," *Publications of the Astronomical Society of the Pacific*, vol. 127, no. 953, p. 665, 2015.
- [17] G. Aniano, B. Draine, K. Gordon, and K. Sandstrom, "Commonresolution convolution kernels for space-and ground-based telescopes," *Publications of the Astronomical Society of the Pacific*, vol. 123, no. 908, p. 1218, 2011.
- [18] J. Idier, Bayesian approach to inverse problems.
- [19] J. R. Shewchuk, "An introduction to the conjugate gradient method without the agonizing pain," 1994.
- [20] A. Glasse, G. Rieke, E. Bauwens, M. García-Marín, M. Ressler, S. Rost, T. V. Tikkanen, B. Vandenbussche, and G. Wright, "The Mid-Infrared Instrument for the James Webb Space Telescope, IX: Predicted Sensitivity," *Publications of the Astronomical Society of the Pacific*, vol. 127, no. 953, p. 686, 2015.
- [21] R. C. Gonzalez and R. E. Woods, "Image processing," *Digital image processing*, vol. 2, 2007.
- [22] F. Orieux, J.-F. Giovannelli, and T. Rodet, "Bayesian estimation of regularization and point spread function parameters for wiener-hunt deconvolution," *JOSA A*, vol. 27, no. 7, pp. 1593–1607, 2010.