

Gold nanoparticles as absolute nano-thermometers

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Abstract

Nano-thermometry is a challenging field that can open the door to intriguing questions ranging from biology and medicine to material sciences. Gold nanorods are excellent candidates to act as nanoprobe because they are reasonably bright emitters upon excitation with a monochromatic source. Gold nanoparticles are commonly used in photothermal therapy as efficient transducers of electromagnetic radiation into heat. In this work we use the spectrum of the anti-Stokes emission from gold nanorods irradiated in resonance to measure the *absolute temperature* of the nanoparticles and their surrounding medium *without* the need for a previous calibration. We show a 4K accuracy in the determination of the temperature of the medium with spectral measurements of 180 s integration time. This procedure can be easily implemented in any microscope capable of acquiring emission spectra and it is not limited to any specific shape of nanoparticles.

Keywords

Gold nanorods, Plasmon, Anti-Stokes, Sensing, Temperature

Most physical, chemical and biological processes depend on temperature. Together with the miniaturization of devices and the advent of nanotechnology the need for measuring temperature with high spatial accuracy started to emerge. Notably in biology^{1,2} and medicine³ measuring and controlling temperature at a sub-cellular scale are the challenges that need to be overcome to achieve better understanding and control of the mechanisms involved in new therapies such as photothermal tumor ablation⁴ or controlled drug delivery.^{5,6}

Nanometer-size probes with distinctive spectral features are ideal candidates for temperature measurements since they provide high spatial accuracy while far-field optics allow a non-contact readout. Some of the proposed strategies include structures that undergo a conformational change upon an increase in temperature,⁷ thus inducing variations in fluorescence intensity of a dye molecule embedded in them.

Also cleverly designed lanthanide-based fluorescent probes in which the ratio of particular emission peaks depends on temperature provide a high accuracy and can be used as nanothermometers⁸ even in biological samples.⁹ Photobleaching is often an important limitation of these approaches. Recently, Surface Enhanced Raman Spectroscopy (SERS) allowed to measure spectral changes induced by temperature down to single molecules,¹⁰ but a careful calibration of the measurements is crucial.

Gold nanoparticles continue to receive a fair amount of attention because of their unique optical properties.¹¹ The collective oscillation of conduction electrons, also known as plasmon, shows a resonance in the visible to near infra-red wavelengths. This resonance can be tuned by changing the shape of the particles¹² and is responsible for a large absorption and scattering cross section at the resonance wavelength. These cross sections can be calculated by solving the Maxwell equations numerically employing different computer packages,^{13–15} providing a good agreement between calculations and what is experimentally achievable.

Thanks to their high absorption and scattering cross section (several times higher than their geometrical cross section) it is relatively simple to detect nanoparticles in a dark-field scattering configuration¹⁶ or via photothermal imaging.^{17,18} Alternatively, detecting gold

nanoparticles through their luminescence¹⁹ is also possible; their low quantum yield,^{20–23} around 10^{-6} , is compensated by the enhanced cross section at the surface plasmon resonance (SPR). The luminescence signal is stable over time; gold nanoparticles do not blink nor bleach, therefore are useful labeling agents for processes that require extended periods of observation.²⁴

Different metallic nano-objects are being introduced as agents for photothermal therapy^{5,25} or drug delivery.²⁶ One of the advantages of gold nanoparticles is the possibility of tuning their resonance to the near infra-red range, where the penetration of light into tissues can be of several centimeters.^{3–5,25,27,28} Moreover the particles can be used not only for treatment, but also for imaging.^{5,29} In the case of photothermal therapy, nanoparticles are used as heat sources^{4,27} to locally increase the temperature in order to induce the death of specific cells in a tissue.^{5,25} However, the temperatures reached³⁰ can only be estimated from models²⁹ or from an ad-hoc calibration. Therefore a method to simultaneously increase and monitor the local temperature will be of great interest in a broad range of fields. In this paper we show that the anti-Stokes luminescence of *single* gold nanorods can be used to measure their temperature *upon resonant CW irradiation*. *The temperature error for a single measurement is less than 10 K with an acquisition time of a few minutes; with a set of such measurements, the temperature of the surrounding medium can be determined with 4 K accuracy or better.*

Luminescence of metallic nanoparticles has been the subject of extensive study in recent years. Since the first observation of luminescence from bulk gold,³¹ different groups have tried to quantitatively describe the observed properties^{32,33} *34–36*, such as the quantum yield^{20–23,37} and the emission spectrum.³⁸ In particular, gold nanorods present two distinct resonance energies, namely the transverse and the longitudinal plasmon resonances. These particles can therefore be excited efficiently at one of those energies; the transverse resonance corresponds to a wavelength of about 530 nm and will give rise to a broad luminescence emission with a peak at the longitudinal plasmon energy. Conversely it is possible to excite the

particles with a wavelength matching the longitudinal plasmon resonance. In this case the excitation benefits from an enhanced absorption cross section, but the emission that overlaps the plasmon resonance will be mostly blocked by the filters needed to prevent direct excitation light from reaching the detectors.

In this work, we call “photoluminescence” any secondary light emission[?] at energies different from the excitation laser energy, $\hbar\omega_L$. After (virtual or real)¹ absorption of an excitation photon, the excited electronic state^{31,37} may interact and exchange energy with the phonon bath or, in the case of metals, with the bath of thermally excited charge carriers around the Fermi level. After a number of interactions, the excited electronic state will re-emit a photon which can possess a lower or higher energy than that of the excitation photon.^{39–41} For a non-resonant excitation, the probability of more than one interaction is negligible and the main contribution to secondary emission is Raman scattering.³⁴ This is the case, for example, of insulators excited well below their electronic absorption edge. For resonant excitation, a relatively long-lived excited state is prepared. It will have enough time to interact repeatedly with thermal baths, particularly with phonons. This is the case of organic dye molecules or semiconductors, in which relaxed fluorescence is observed. We also note that fluorescence always presents hot bands on the anti-Stokes side of the excitation laser. In most fluorescence detection schemes, however, these hot bands are ignored, but they are far from negligible in heavily doped samples.⁴²

Metal nanoparticles fall between those two extremes because the excited electronic state, an electron-hole pair, relaxes very rapidly by interacting with other charge carriers and with phonons. The photoluminescence lifetimes are on the order of tens of femtoseconds⁴³ and therefore there is not enough time to obtain a fully relaxed luminescence. In other words, the photoluminescence is always “hot”. It is worth noting that Raman scattering, corresponding to the lowest order of interaction with baths, will be an important contribution to photoluminescence.^{34,36} However, second and higher orders may also contribute significantly. Because

¹we do not specify whether the absorption is real or virtual.

all these processes obey a Boltzman-type of relation between anti-Stokes and Stokes emission, they cannot be easily distinguished from each other on the basis of their temperature dependence.

The anti-Stokes emission is highly sensitive to temperature and thus it can be used for thermometry.⁴⁴ In this letter we present a simple procedure to extract the absolute temperature from the anti-Stokes photoluminescence spectrum of individual gold nanorods without the need of any previous temperature calibration. We show that we can determine the particle temperature *in-situ* with an accuracy of 6% by recording a single anti-Stokes spectrum (with an acquisition time of 3 minutes). Moreover, by performing this measurement at different excitation powers we can obtain the temperature of the surrounding medium with an accuracy better than 2%.

Phenomenological model for the luminescence emission. In a nutshell, we consider the luminescence emission as radiative recombination of electron-hole pairs created by the decay of the plasmon and safter their interaction with thermal baths. Before the recombination, carriers may interact with the baths one or more times, leading to secondary light emission with an energy different from the initial internal energy of the pair. The anti-Stokes spectral contribution arises from interactions that increase the energy of the pair, whereas the Stokes emission to a decrease in energy. The emission process will be enhanced by the plasmon; therefore the luminescence spectrum will be modulated with the plasmon shape. A scheme of these ideas is shown in the Supporting Information.

Exciting a gold nanorod with a monochromatic beam at its resonance frequency, ω_{SPR} , generates a collective oscillation of the gas of conduction electrons, also called a plasmon. The plasmon decays by forming a pair of hot electron and hole with an internal energy equal to the exciting photon energy,⁴⁵⁻⁴⁷ i.e. $E_{e-h} = \hbar\omega_L$.

This hot electron and hole have a small probability of recombining radiatively, i.e. of re-emitting their high electronic energy as a photoluminescence photon. If they have interacted only with static surfaces or defects, their energy will be the same and therefore the emitted

photon will have the same energy as the incoming photon, and will not contribute to the measured photoluminescence. It will be blocked by the notch filter used to remove the exciting laser from detection. If, on the other hand, they have interacted with a phonon or a thermally excited electron or hole, they may have lost or acquired energy. In both cases the energy available upon recombination cannot much exceed $\hbar\omega_L + k_B T$, where k_B represents Boltzmann's constant and T the absolute temperature.

Radiative recombination gives rise to emission spectrally and spatially distributed throughout the particle over a broad frequency band with an exponential cutoff at $\hbar\omega_L + k_B T$. The weak recombination emission can be greatly enhanced by the surface plasmon resonance, acting as an antenna. With this model the following predictions can be made. Firstly the emission spectrum must follow the plasmon spectrum if the excitation laser is well above the plasmon resonance as shown in Figure 1, green line. If the excitation falls within the plasmon resonance, the spectrum is expected to follow the plasmon spectrum multiplied by a Bose-Einstein statistics factor arising from phonon population. Here we assume that the coupling to the phonons dominates the process, while references^{34,36} assume that carrier-carrier interactions dominate. Thus, under our assumption, the emission should be proportional to the phonon occupation number \bar{n} for anti-Stokes and $\bar{n} + 1$ for Stokes processes, with

$$\bar{n} = \left(\exp \frac{\hbar\Omega}{k_B T} - 1 \right)^{-1}, \quad (1)$$

where $\hbar\Omega$ is the phonon energy. It is possible to consider carrier-carrier interactions obeying Fermi statistics, i.e. $n = \left(\exp(\frac{\hbar\Omega}{k_B T}) + 1 \right)^{-1}$.^{34,36}

With this model, we can also predict that the emission should be polarized. For the strong longitudinal plasmon of gold nanorods this polarization coincides with the longitudinal axis of the particle.⁴⁸ Moreover, the lifetime should be determined by the lifetime of hot electrons and holes and should be significantly shorter than the thermalization time of the carriers. Indeed, a few interactions would suffice to reduce the carriers' energy significantly and therefore the electron and hole wouldn't have the energy required to produce an optical

photon. One important assumption for this model is that the emission spectrum of radiative recombination is much broader than the plasmon. Therefore excitation just above the plasmon resonance should excite the electron-hole pairs with nearly the same efficiency as well above the plasmon resonance.^{23,36}

Application to nanothermometry. According to the model just described, the anti-Stokes emission spectrum follows the form,

$$I(\omega) = I_{\text{SPR}}(\omega) \cdot \left(\exp \frac{\hbar(\omega - \omega_L)}{k_B T} - 1 \right)^{-1} \quad (2)$$

where $I(\omega)$ is the emitted intensity, ω is the angular frequency of the photons and ω_L is the frequency of the exciting laser. $I_{\text{SPR}}(\omega)$ is the surface plasmon resonance profile that will be obtained by exciting the particle at energy much higher than the resonance. The only remaining free parameter is the temperature T (and a normalization constant not included in equation 2). Fitting measured anti-Stokes emission spectra with this equation will yield the absolute temperature T of the particle under study.

The procedure we propose to obtain the absolute temperature of gold nanorods from the anti-Stokes luminescence emission and without the need of any previous temperature calibration involves the following steps.

1. Obtain the plasmon resonance spectrum of the particle. This is usually expressed as a Lorentzian function,¹¹ i.e.

$$I_{\text{SPR}}(\omega) = \frac{(\Gamma/2)^2}{(\omega - \omega_{\text{SPR}})^2 + (\Gamma/2)^2}$$

where ω is the photon energy frequency, ω_{SPR} is the resonance frequency and Γ is the width of the surface plasmon resonance. In our case, we detect the spectrum of photoluminescence, excited at 532 nm to extract ω_{SPR} and Γ , as explained in the Supporting Information, along with an explanation of the experimental errors.

2. Excite near the longitudinal plasmonic resonance and detect the blue-shifted anti-Stokes emission spectra. For this we employed a 633 nm laser as a source.
3. Fit the high-energy part of the spectrum using equation 2 with T as the *only* free parameter.

We emphasize that we cannot simply use the anti-Stokes to Stokes intensity ratio to obtain the temperature of the particle, as is commonly done with Raman lines of molecules,^{49,50} due to the presence of the strong plasmonic enhancement of the emission that must be considered in addition to the Boltzmann factor.

Experimental method. All the measurements in this work were performed with a home-built confocal microscope equipped a spectrometer (Acton 500i) in the emission path. We focused our lasers to a diffraction-limited spot using a 60 \times , NA 1.4 oil immersion objective (Olympus) and collected the emitted photons through the same objective. This provided high excitation and collection efficiency. We employed a 532 nm (CNI) laser for characterizing the nanorods' plasmon and a 633 nm HeNe (Thorlabs) to excite the nanorods in resonance. We give more experimental details in the Supporting Information.

Wet chemically synthesized nanorods⁵¹ with average dimensions of 21 nm \times 50 nm and a plasmon resonance around 650 nm were spin-coated onto clean coverslips, controlling the superficial concentration to separate individual nanorods within the diffraction-limited spot.¹¹ In addition, the samples were mounted in a flow cell that allowed us to increase the temperature of the medium up to 60 °C and to monitor it through a Pt100 resistance thermometer placed 1 mm away from the observation area. More details about the experimental setup are given in the Supporting Information.

To compensate for the drift of the setup while increasing the temperature, we developed a computer program to continuously track a reference particle. The same program was responsible for recording the temperature and triggering the spectrometer. In this way complete data sets were acquired at different temperatures, with excitation at 532 nm and 633 nm, at different laser intensities. A spectrum with 532 nm laser excitation was taken after

every cycle to ensure that the particle under study had not reshaped due to high excitation power.

The intensity of the laser was controlled via an acousto-optic modulator. Six accumulations of each spectrum were recorded with an exposure time of 30 s. This not only allowed us to lower the noise of the measurement because of a longer exposure time (3 minutes in total), but also allowed us to remove bright pixels generated by cosmic rays. Having several accumulations is also useful to monitor changes in the intensity of the spectra during the acquisition itself. These changes can be due to a drift of the setup while measuring or to a reshaping of the particle. If the reshaping was confirmed by comparing the spectra acquired with the 532 nm laser,⁵² the measurements were rejected. If the changes in the observed emission spectra were due to drift of the setup, the particular data set was not taken into account. For the purposes of this work the excitation intensity is crucial for characterizing the method; if the particle is not in focus it would result in an overestimation of the excitation power.

Experimental Results. The proposed model for the anti-Stokes emission requires the plasmon spectrum ($I_{\text{SPR}}(\omega)$ in equation 2) in order to fit the emission at shorter wavelengths and extract the particle temperature. It has been shown that both scattering and luminescence spectra roughly overlap over a broad range of wavelengths.²² Therefore exciting gold nanorods with 532 nm allows us to record the longitudinal plasmon spectra, as shown in the green solid curve of Figure 1. It has to be recalled that the luminescence spectrum is not a perfect Lorentzian since there is a broadband contribution also observed in bulk gold.³¹ The procedure to extract the SPR profile from such a measurement is explained in the Supporting Information. We show in figure 1 the extracted surface plasmon profile in the red solid curve.

The other curves in Figure 1 show the luminescence emission of the same nanorod with irradiation at 633 nm at different powers, ranging from 25 μW to 75 μW at the back aperture of the objective. The vertical black line shows the wavelength of the laser. The Stokes

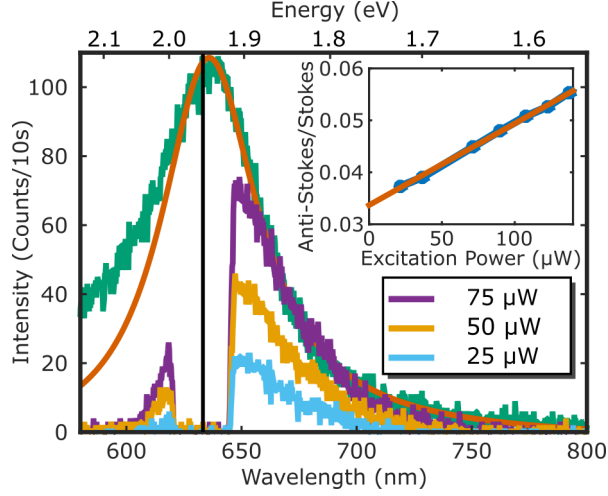


Figure 1: **Luminescence emission spectra of a single gold nanorod.** The green curve is the measured luminescence emission under 532 nm excitation and the red curve shows the extracted $I_{\text{SPR}}(\omega)$ from this spectra. The other curves are the emission of the same particle under 633 nm irradiation at three different powers indicated in the legend. The inset shows the anti-Stokes-to-Stokes ratio as a function of the excitation power, overlapped with a linear fit in red. The dip centered on the laser wavelength is caused by the notch filter used to prevent the excitation laser from reaching the detectors.

part of the spectrum at longer wavelengths than the excitation shows the same shape as the plasmon emission observed under 532 nm excitation, apart from a normalization factor. From the figure it can readily be seen that the shape of the anti-Stokes emission, at shorter wavelengths than excitation, is exponential-like and doesn't follow the Lorentzian shape of the Stokes emission. The dip between Stokes and anti-Stokes is caused by the notch filter that prevents direct excitation light from reaching the detectors.

The inset of Figure 1 shows the anti-Stokes-to-Stokes ratio of the integrated luminescence for different laser excitation intensities. It is possible to see that even though the photoluminescence process is linear with a linear behavior, the anti-Stokes intensity increases slightly more rapidly than the Stokes emission. We already exploited this phenomenon to image gold nanorods in high-background conditions.⁴² For more information on the power dependence of both the anti-Stokes and Stokes luminescence, please refer to the Supporting Information.

To further characterize the anti-Stokes emission in gold nanorods, we measured the emission for 105 nanorods with different plasmon resonances under the same 633 nm excitation

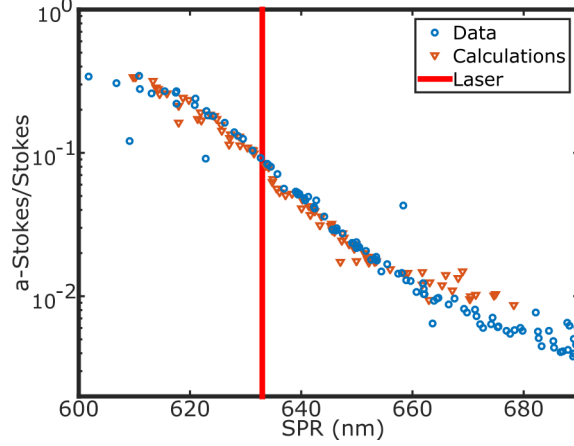


Figure 2: **Characterization of anti-Stokes emission for different surface plasmon resonance.** Ratio of the anti-Stokes to Stokes emission under 633 nm excitation as a function of the resonance wavelength of the particle. The blue circles are experimental results (105 different nanorods), while the red triangles are the results of numerical simulations with equation 1 (82 nanorods with different aspect ratio). There is a very good agreement between experiment and calculations. Particles with a resonance to the blue of the laser (indicated by the vertical red line) have an increased anti-Stokes emission.

and calculated the ratio of integrated anti-Stokes to Stokes emissions. Figure 2 shows the experimental ratios as blue circles, where as a function of the surface plasmon resonance (SPR) of the particle. The vertical red line marks the laser wavelength. The particles measured had resonances between 600 nm and 690 nm; the ones showing the maximum ratio of anti-Stokes to Stokes are those with a resonance to the blue of the laser. For these particles the longitudinal plasmon is enhancing preferably the anti-Stokes emission. For particles with a resonance at the laser wavelength the anti-Stokes and the Stokes emission have similar enhancement and show a ratio close to 10%.

Figure 2 also shows as red triangles the results of numerical calculations following the protocol presented before. An excellent overlap between the measured and the calculated data can be observed. The absorption cross section of 82 particles was calculated numerically with the ADDA package¹⁴ using a fixed width of 23 nm and different lengths to achieve different SPR wavelengths. In the Supporting Material we show TEM images from which we extracted the size of the nanorods. Each calculated absorption spectrum was fitted by a Lorentzian and used as $I_{\text{SPR}}(\omega)$ in eqn. 2. Assuming a diffraction-limited laser spot and

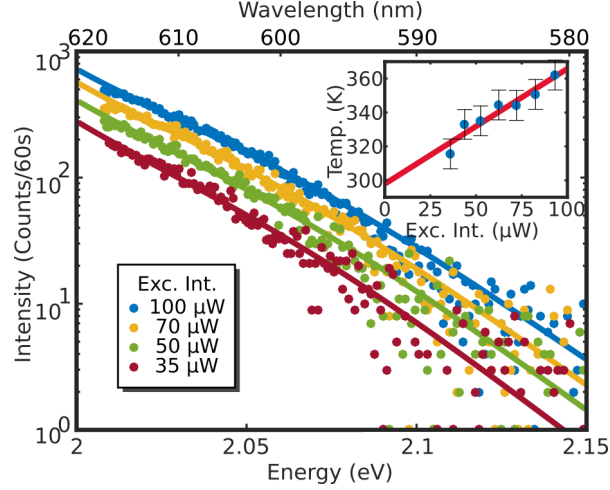


Figure 3: **Anti-Stokes emission of a single nanorod at different irradiation powers.** We used the model from equation 2 to fit the experimental data. There is an excellent agreement between data and model. The inset shows the extracted temperature at each power (blue dots) and a linear extrapolation of the data to $0 \mu\text{W}$ excitation power. The value obtained for room temperature was 293 K while the measured value was 296 K.

using the calculated absorption cross section we calculated the temperature of the particle. This value was used in equation 2 to compute the anti-Stokes emission spectrum. The Stokes emission was set proportional to the excitation power with a shape given by the calculated absorption spectra. Since both anti-Stokes and Stokes emissions are proportional to the excitation power, this term cancels out when computing the ratio. The laser power therefore only enters into the equation when calculating the temperature of the particles. It is remarkable that the agreement between data and calculations was achieved without free parameters, solely taking into account the transmission spectra of the filters.

Furthermore, by fitting the anti-Stokes part of the spectra shown in Figure 1 with equation 2 it is possible to extract the temperature of the particle at each excitation power. Figure 3 shows the results of this procedure. The spectra shown were recorded at 4 different excitation intensities while the full lines are the fits; again, there is an excellent agreement between data and model. For every anti-Stokes measurement we have also acquired the full plasmon spectrum exciting with a 532 nm laser before and after the temperature extraction. The full plasmon spectrum is necessary to calculate the parameters of $I_{\text{SPR}}(\omega)$ from equation

2 and also to verify that the particle did not reshape while being excited at resonance.

The inset in Figure 3 shows the temperatures resulting from the fits at different irradiation intensities (blue dots). Note that the absolute temperatures of the particle at each excitation power were calculated without any calibration. As expected, the temperature of the nanorod varied linearly with excitation intensity, or equivalently with the absorbed energy. Thus, this method provides an *in situ* way to measure the temperature reached by nanoparticles when they are excited with resonant monochromatic light. Additionally, from these data sets it is also possible to calculate the temperature at 0 μW excitation power, i.e. room temperature, by extrapolating the results with a linear fit. The value we obtained in this case is $293 \pm 6 \text{ K}$, while room temperature was 296 K, a 2% accuracy.

The accuracy of the obtained temperature depends on an accurate modeling of the photoluminescence. The first step in the protocol is the determination of the surface plasmon spectral profile, $I_{\text{SPR}}(\omega)$ in equation 2. In this paper we obtained this term by fitting an exponential background plus a Lorentzian to the spectra obtained at 532 nm excitation, as explained in detail in the Supporting Information. We note that this choice was made for experimental convenience in our setup, but other options to obtain the SPR profile, are suited for the procedure as well. The error bar in the inset of figure 3 and in the following figures is the result of the estimated variance in the fit parameters for the anti-Stokes spectra using equation ??, step 3. in our protocol.

As expected from the model, the anti-Stokes emission depends not only on the particle's intrinsic properties but also on the temperature of the surrounding medium.⁵³ In order to further test this point, we changed the temperature of the sample in a controlled manner and recorded the luminescence emitted by a single nanorod.

For this set of experiments we employed an air objective (60 \times , NA 0.9, Olympus) to avoid the presence of a heat sink directly in contact with the observed area. We employed higher laser powers to compensate for the lower excitation efficiency. At each temperature, 7 spectra were acquired at different 633 nm excitation powers and also a spectrum of the

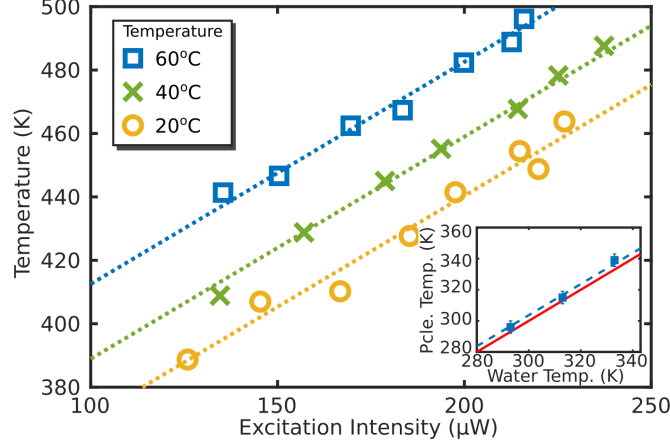


Figure 4: **Calibration-free temperature measurement.** Extracted temperatures from the anti-Stokes-luminescence emission of an individual nanorod at different excitation powers and at different sample temperatures. The dashed lines are fits with the same slope for the three temperatures. The **squares** in the inset plot show the local temperature of the sample obtained by extrapolating **the** temperature at zero excitation power as a function of the water temperature. The red line represents the expected curve if both temperatures are identical (equal). The dashed blue line is a fit to the data points with unit slope that shows a systematic offset of 3.8 K, a 1.2% difference.

plasmon before and after each measurement in order to monitor any possible reshaping of the particles during the experiment.

Figure 4 shows the extracted temperature of a particle at varying excitation powers and at different water temperatures. The blue squares are the results of the measurement at 20 °C, while the green crosses are measured at 40 °C and the yellow circles at 60 °C. The full lines are the calculated temperatures for a particle with plasmon overlapping the measured one and assuming a diffraction-limited focus spot. For the dimensions of the particle, the mean values from TEM images were used and the length was adjusted to obtain the measured resonance. There is a remarkable agreement between the calculation and the measured values. Moreover it is possible to extrapolate the temperature at zero excitation power for each case as was explained earlier. The results are shown in the inset of the figure for each temperature. The red line with slope 1 is a guide to the eye.

Figure 4 clearly shows that the extracted temperature varies with the temperature of the surrounding medium. More strikingly the method does not require any previous calibration

nor adjustment. The values obtained with the extrapolation to $0\text{ }\mu\text{W}$ excitation power were $296 \pm 4\text{ K}$, $315 \pm 4\text{ K}$ and $339 \pm 4\text{ K}$ for water temperatures of 293 K , 313 K and 333 K respectively. The inset plot in figure 4 presents these points and a red solid line with the expected curve if both temperatures are identical. The dashed line shows a fit of the data that evidences a small systematic offset of 3.8 K . This represents an inaccuracy of 1.2% which is a good result for a calibration-free method. Notably, the presented calibration-free procedure would allow us to perform the same measurements in any other setup and could act as a reference for calibration of other nano-thermometers.

Being able to control and monitor temperature at the nanoscale is of utmost importance in different fields ranging from photothermal therapy⁵ to nano fabrication.⁵⁴ In this work we have shown a simple procedure that allows us to measure the temperature of single gold nanorods irradiated by a monochromatic continuous laser and without any previous calibration. The level of accuracy of the temperature measurement depends on several factors, but for nanorods it can be estimated to be better than 6 K with an integration time of 3 minutes without any previous calibration.

The model employed for describing the anti-Stokes emission takes into account the surface plasmon resonance of the particles under study, which is responsible for enhancing the emission, as well the electron-hole pairs interaction with the thermal baths. Particles with a resonance to the red of the excitation wavelength would be more reliable in the temperature extraction procedure, but would also exhibit a lower emission towards shorter wavelengths. The trade-off between both effects and the possibility to fully characterize the plasmon resonance, will determine the specific particles that are better suited for each application.

A possible improvement of this technique would be the use nanostructures with a narrow shape distribution such as gold bipyramids.⁵⁵ Such structures would be ideal candidates for temperature extraction since they present negligible size dispersion and thus their plasmon can be measured in bulk or determined from theory, avoiding the need of a second excitation source. This would reduce the main source of inaccuracies for the method.

The proposed method does not require any temperature calibration, since the only free parameter of the model is the absolute temperature of the nanoparticle under study. Moreover the recording of the anti-Stokes spectrum is readily achievable in any confocal microscope with a coupled spectrometer. A 6 K accuracy may suffice for several applications; it is important to point out that this value can be improved in different ways: by carefully selecting the particles that show the most favorable plasmon resonance; by determining the plasmon resonance through white-light scattering, reducing the uncertainty in the fit; by increasing the exposure times to increase the signal-to-noise ratio. [A cheaper alternative implementation would be to have point two detectors and use filters to detect the anti-Stokes emission in one and the Stokes emission in the other, but this approach would require a temperature calibration.](#)

The authors declare no competing financial interest.

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Supporting Information Available

The Supporting Information is available free of charge on the ACS Publications website with the following sections: Anti-Stokes emission from gold nanorods, Experimental setup, Gold nanorods sample characterization, Determination of the error in the temperature extraction, Gold Nanorod temperature numerical calculations.

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Graphical TOC Entry

