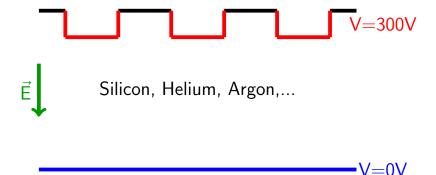


Summary

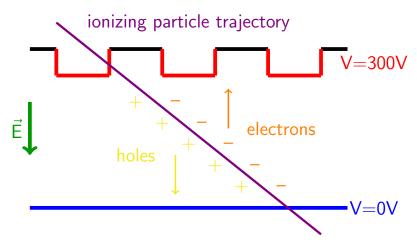
- Problem statement
 - Particle detector
 - Media ionization
 - Drift velocity
 - Shockley Ramo theorem
 - Townsend avalanche
- The software pdetect
 - Features and technical characteristicsGeometries and boundary conditions

 - Computing the potential in the detector
 - Computing the current induced by a particle
- Conclusion

The detector



Media ionization



Charges drift with velocity

$$|\vec{\mathbf{v}}| = \mu |\vec{\mathbf{E}}|$$

where μ is the mobility

- Electrons mobility is much higher than hole mobility
- mobility decreases when fields of 104 Vcm⁻¹
 and higher are applied due to saturation

Shockley Ramo theorem

Instantaneous current generated by one moving charge q at speed \vec{v}

$$i_{e,h} = -q\vec{v}\cdot\vec{E}_w$$

Weighting field $\vec{E}_w = \vec{E}$ computed when applying

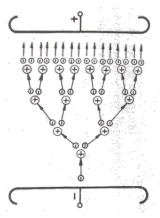
- 1V to the measurement electrode
- 0V to the other electrodes

Instantaneous current induced by the particle

$$i_{tot} = \sum_{holes} i_h + \sum_{electrons} i_e$$

Townsend avalanche

- e⁻ extract e⁻ from media molecules
- \bullet In gas detectors when $|\vec{E}|>10^6 {\rm Vm^{-1}}$



Summary: what do we need to compute the current?

$$egin{aligned} ec{i} &= -qec{v}\cdotec{E}_w & ec{v} &= rac{q}{|q|}\muec{E} \ &\Longrightarrow i &= -|q|\muec{E}\cdotec{E}_w \ &ec{E} &= -
abla V & ec{E}_w &= -
abla V_w \end{aligned}$$

 \implies to compute i you need V and V_w \implies Solve $\nabla^2 V = 0$ and $\nabla^2 V_w = 0$ for corresponding boundary conditions

- Simulates both silicon and gas detectors
- Computes \vec{E} and \vec{E}_w for three different 2D geometries
- Computes itot for any particle trajectory
 - Simulates Townsend avalanche and mobility saturation
- Outputs:
 - plots of $V(x, y), V_w(x, y), \vec{E}(x, y), \vec{E}_w(x, y)$
 - points $(t, i_{tot}(t))$ in a text file

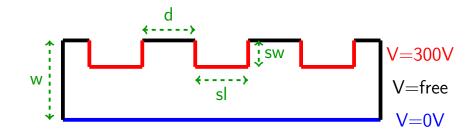
Developed in C++ (5209 lines of code)

Object oriented programming (inheritance, generic programming, composition,...) advantages:

- avoids code duplication
- add features to a module with few modifications to other modules
- improve code readability

Fast: multithreading, adaptive grid refinement,...

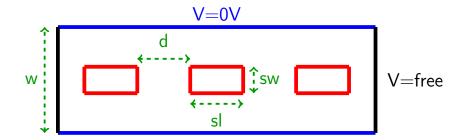
Detector geometry and boundary conditions (I)



- number of strips
- strip length sl
- strip width sw

- strip potential V
- detector width w
- inter-strips distance d

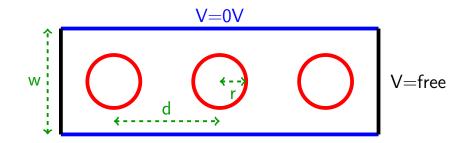
Detector geometry and boundary conditions (II)



- number of strips
- strip length sl
- strip width sw

- strip potential V
- detector width w
- inter-strips distance d

Detector geometry and boundary conditions (III)



- number of strips
- strip radius r
- detector width w

- strip potential V
- inter strips centers distance d



Finite element method solves $\nabla^2 V = 0$ numerically

- function domain is divided in subdomains called finite elements
- piecewise interpolation: one polynomial fits the restriction of the searched function on one finite element
- build a linear system with the polynomials coefficients
- solves the system thanks to constraints:
 - satisfy $\nabla^2 V = 0$
 - match imposed domain boundary conditions
 - boundary conditions between polynomials defined on adjacent finite elements

Use deal.ii to solve the Laplace eq. (fast thanks to multithreading)

Inputs:

- Detector geometry
- Boundary conditions
- Max relative error on V

Outputs for each FE:

- lacktriangledown a funct approximating V
- $oldsymbol{2}$ a bound on the error on V
- lacktriangle a funct approximating ∇V

Perform this process two times

- get $\vec{E} = -\nabla V$
- get $\vec{E}_w = -\nabla V_w$ with other boundary conditions

Computing the potential: adaptive grid refinement (I)

Electric field constant in most part of the detector

quick computation with large FE sufficient

But Electric field varies strongly close to the strips

much smaller FE required to get precise results

Uniform coarse grid composed of large FE

→ results not precise enough close to the strips

Uniform dense grid composed of small FE

⇒ waste of time

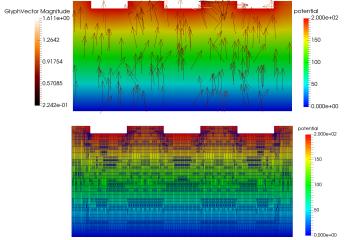
Iterative process

- Starts with the coarsest grid
- Compute V and error at each FE
- If max error > max tolerated error
 - refine only cells of the grid with highest errors
 - return to step 2

Final grid is composed of small cells only where the Electric field varies strongly

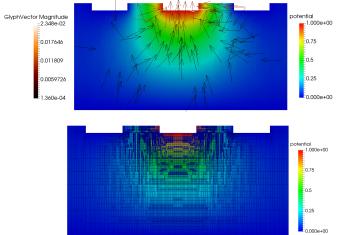
Serrated rectangle: potential and electric field

Strip potential = 200V, detector width = 300μ m, strip length = 100μ m, strip width = 20μ m, inter-strips distance = 100μ m, max relative error = 0.9%



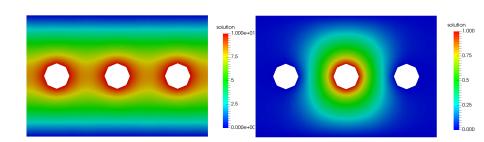
Serrated rectangle: weighting potential and weighting electric field

Strip potential = 200V, detector width = $300\mu m$, strip length = $100\mu m$, strip width = $20\mu m$, inter-strips distance = $100\mu m$, max relative error = 0.9%

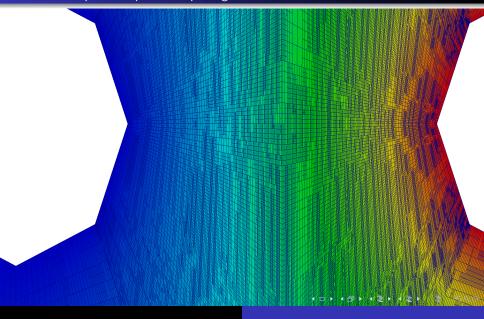


Circular strips: potential and weighting potential

Detector width $=200\mu\text{m}$, number of strips =3, radius $=21\mu\text{m}$, inter strips distance $=50\mu\text{m}$, potential =10V, max relative error 1%

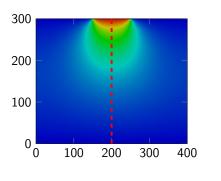


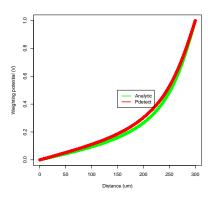
Circular strips: example of adaptive grid



Potential: numerical vs analytical solution

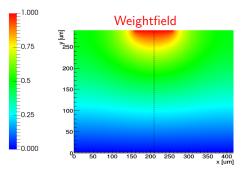
Potential along a vertical line (middle of the detector) Pitch = 400μ m, strip length = 100μ m, width = 300μ m, $V_{strip} = 1V$

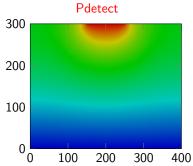




- The two curves perfectly match for a different strip width
- Strip width used to compute analytical solution is unknown

V=1V, 1 strip, pitch = 400μ m, strip length = 100μ m, width = 300μ m





Computing the current

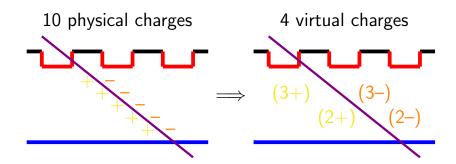
User must choose

- particle trajectory (all supported)
- a precision level affecting
 - the time interval Δt of the simulation
 - the number of punctual charges

Initial charges are grouped among punctual charges

- equidistant and uniformly distributed along the trajectory
- ullet punctual charge $= \Sigma$ physical charges at that point
- ullet fewer than real charges \Longrightarrow faster computation

Punctual / virtual charges



$$t=0$$
 while there are charges inside detector $i_{tot}=0$ for each punctual charge

- \bullet computes \vec{v} using \vec{E}
- ullet computes i using \vec{v} and \vec{E}_w
 - $\bullet \ i_{tot} = i_{tot} + i$
- move punctual charge

$$\bullet (x,y) = (x,y) + (\Delta t \times v_x, \Delta t \times v_y)$$

save
$$(t, i_{tot})$$

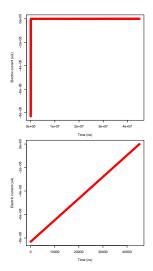
 $t = t + \Delta t$

In gas detectors holes slower than e^- (1000 times)

- small Δt: fine to capture electrons current but far too slow to compute holes current
- large Δt: capture quickly holes current but miss electrons current

Solution: at each iteration select Δt depending on max speed at previous iteration

$$\Delta t = \frac{\textit{width}}{\textit{max_speed_y} \times 2^{\textit{precision_level}}}$$



Current computation algorithm accesses $\vec{E}(x, y)$ and $\vec{E}_w(x, y)$ very often \implies access must be fast

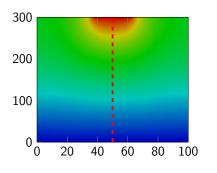
The problem: find the cell (FE) in which the point (x, y) belong

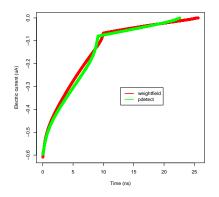
Running through each cell is too slow: O(n) time complexity (n= #cells)

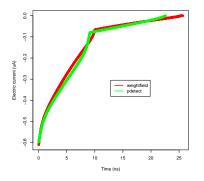
 \implies rtree data structure allows to access the right cell with $\mathcal{O}(\log(n))$ time complexity

Computing the current: Pdetect vs Weightfield

Current induce by particle passing vertically (middle of the detector) 1 strip, pitch = 100μ m, strip length = 25μ m, width = 300μ m, $V_{strip} = 100V$







Some differences, due to effect in Weightfield:

- Depletion voltage
- Gain
- Temperature
- Different number of hole/electron

Conclusion

- This work is a first step to the development of a properly designed software to simulate particle detectors
- Computed potentials match analytical solution and Weightfield results
- Computed current for silicon detectors match Weightfield results
- Computed current for gas detectors need to be validated
- Solving the Laplace equation is fast for 2D geometries
- The slowest phase is the current computation

Future work

- Validate results for gas detectors and Townsend avalanche
- Simulation of additional physical effects: non uniform energy deposition by the particle in the detector,...
- Additional 2D geometries: gaussian curves as corners
- 3D geometries
- Graphic user interface
- Clusters, GPUs
- Simulation of the electronic
- Dependance regarding pressure and temperature
- External magnetic field
- ...

Number of hole-electron pairs

$$N = \frac{dE/dx}{w} \times L$$

- dE/dx energy deposited per unit of length
- w energy to produce one hole-electron pair
- L distance covered in the detector

$$\mu_{s} = \mu \left(1 + \left(\frac{\mu |E_{y}|}{v_{sat}} \right)^{\beta} \right)^{-\frac{1}{\beta}} \tag{1}$$

- \bullet β is a constant
- $m{\mu}$ is the mobility when not considering saturation
- E_y is the component of the electric field along the axis orthogonal to the anode and cathode
- v_{sat} the saturation velocity, it depends on the media and the temperature.

During Δt the initial electric charge q_0 is multiplied:

$$q = q_0 e^{\alpha \Delta x} \tag{2}$$

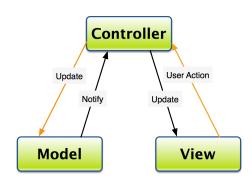
- Δx : distance covered by the e⁻ during Δt
- \bullet α is the first Townsend coefficient

$$\alpha = ap \ e^{-bp/E}$$

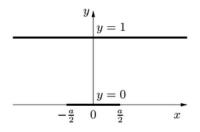
a, b: constants depending on the gas mediap: pressure in the detectorE: norm of the electric field at the charge position

Model-view-controller

- model manages the data, logic and rules of the application
- view: any output representation of information, multiple views of the same information are possible
- controller accepts input and converts it to commands for the model or view

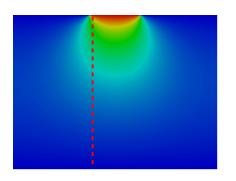


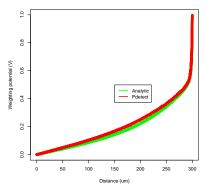
Potential analytical solution



$$\phi_w(x,y) = atan\left(\frac{\sin(\pi y)\sinh(\pi \frac{a}{2})}{\cosh(\pi x) - \cos(\pi y)\cosh(\pi \frac{a}{2})}\right)$$

Potential along a vertical line (edge of the strip) Pitch = 400μ m, strip length = 100μ m, width = 300μ m, $V_{strip} = 1V$

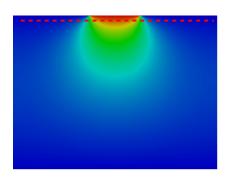


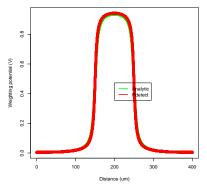


Appendix

Potential: numerical vs analytical solution (II)

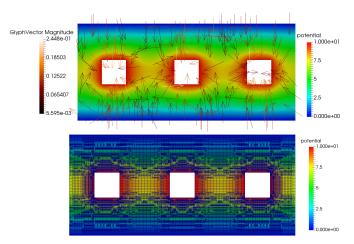
Potential along a horizontal line (below the strip at y=298 μ m) Pitch = 400 μ m, strip length = 100 μ m, width = 300 μ m, V_{strip} = 1V





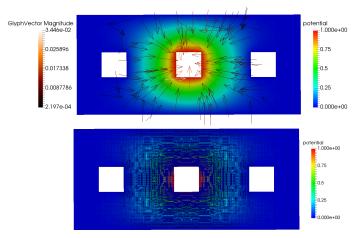
Rectangular strips: potential and electric field

Detector width = 200 μm , strip length = 50 μm , strip width = 50 μm , inter strips distance = 100 μm , number of strips = 3, potential = 10V, max relative error = 0.9%



Rectangular strips: weighting potential and weighting electric field

Detector width = 200 μm , strip length = 50 μm , strip width = 50 μm , inter strips distance = 100 μm , number of strips = 3, potential = 10V, max relative error = 0.9%



- Kelly error estimator
- Integration of the "jump" of gradient of the function on the sides.

$$\eta_K^2 = \sum_{F \in \delta K} c_F \int_{\delta K_F} \left(a \frac{\delta u_h}{\delta n} \right)^2 do$$

- Upper bound of the error
- Modified version of Kelly error system
- Reliable for Poisson equation

More information: see the deal.ii Kelly error estimator page