

# PHY 410

## Homework Assignment 10

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### Abstract

The goal of this assignment is to get more familiar with random process and method, as well as its application on random walk and statistical mechanics.

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# 1 Problem 1

## 1.1 Description

Modify `rwalk.cpp` or `rwalk.py` to simulate a random walk on a 2-dimensional square lattice and a 3-dimensional cubic lattice. Measure the diffusion constant  $D$  in 2 and 3 dimensions and compare with the 1-dimensional walk and theoretical expectations.

## 1.2 Numerical result and analysis

Theoretically, for example in one dimension, assume for each step the displacement will be  $\pm\delta$ , thus the position after  $N$  steps can be expressed as:

$$x^2(N) = [x(N-1) \pm \delta]^2 = x^2(N-1) \pm 2\delta x(N-1) + \delta^2$$

Thus the average will be:

$$\langle x^2(N) \rangle = \langle x^2(N-1) \rangle \pm 2\delta \langle x(N-1) \rangle + \delta^2$$

the middle term is zero, therefore we have:

$$\langle x^2(N) \rangle = \langle x^2(N-1) \rangle + \delta^2$$

consequently:

$$\langle x^2(N) \rangle = N\delta^2$$

In 2 and 3 dimensions it is nothing but the summation in each individual direction. In my case, for each step the displacement in x, y, z direction is 1, therefore the diffusion constants for 1, 2, 3 dimensions are 0.5, 1, 1.5.

Modify the code to simulate for 2 and 3 dimensions, for 2 dimensions, since the random number generator is not a perfect one and limited by the number of steps, when using 1000 ensembles and 1000 steps the result is  $1.017 \pm 0.001$  Fig. 1, while using 10000, 10000 setting, the result is  $0.9990 \pm 0.0001$  very close to theoretical value. Fig 2

For 3 dimensions: for 1000 walkers 1000 steps, the result is  $1.499 \pm 0.001$  Fig. 3

Since it is a random process, each time the result will be different and roughly fluctuating around the theoretical value, I was quite lucky in this trial, thus I stopped here.

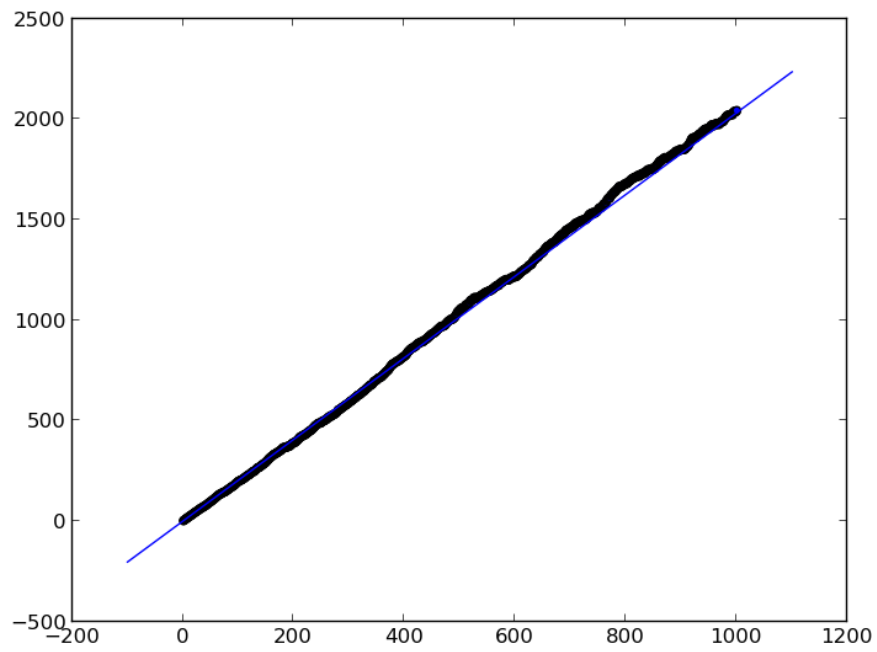


Figure 1: 2d average of  $displacement^2$  vs. steps, 1000 walkers

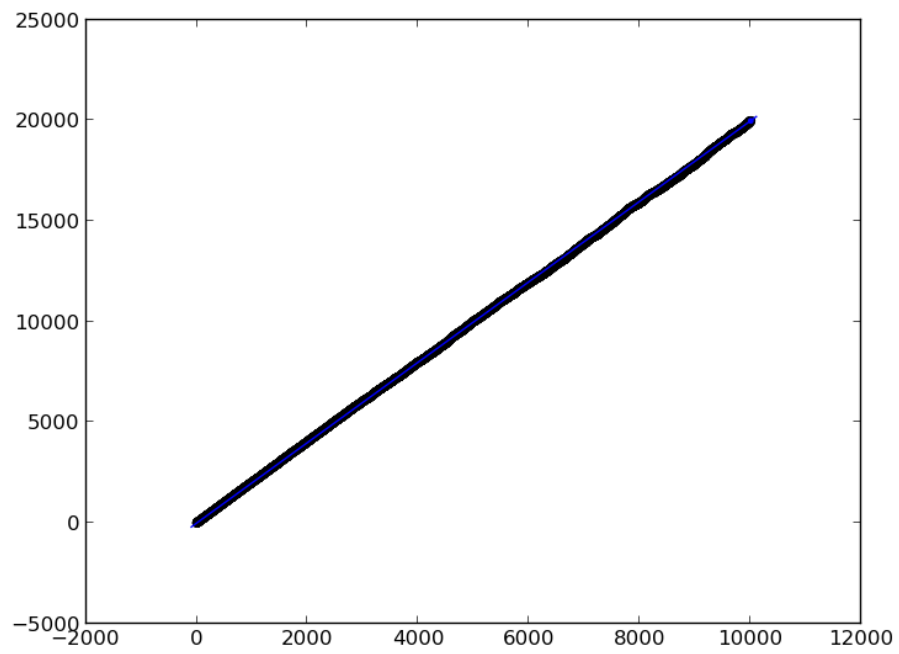


Figure 2: 2d average of  $displacement^2$  vs. steps, 10000 walkers

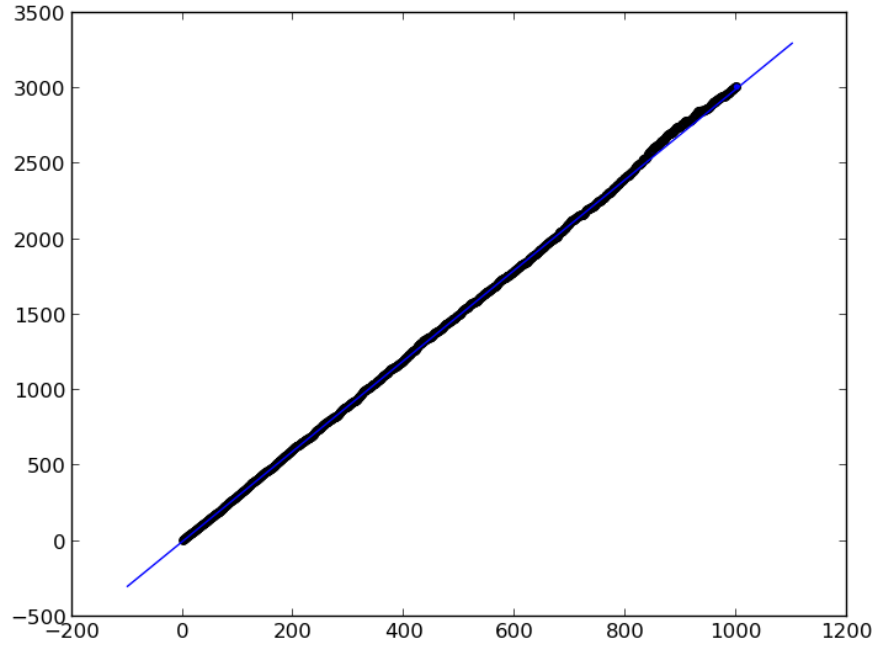


Figure 3: 3d average of  $displacement^2$  vs. steps, 1000 walkers

## 2 Problem 2

### 2.1 Description

Measure and plot the average magnetization  $\langle |M| \rangle$  and Heat Capacity  $C$  of the Ising model as a function of temperature  $T$  for three different lattice sizes. Use the data to identify the transition temperature  $T_C$ .

### 2.2 Result

I used the lattice of size 10x10, 50x50, 100x100, using 500 MC steps. (This is a fair number for 10x10 yet not sufficient for larger grids, however, limited by the performance even using c++, I chose this number, we can see below for large grids the result is not so obvious)

For 10x10: Fig. 4 For 50x50: Fig. 5 For 100x100: Fig. 5

We can see for larger lattice, since I didn't use enough number of steps, the pattern is not as obvious as that of 10x10, yet we can still see the shape of it. To better evaluate the transition temperature, I applied FFT and got rid of the high frequency waves, as you can see in this plot: 7 8 9.

After fft, we can see for 10x10, 50x50, 100x100, the  $T_c = 2.62, 2.22, 2.23$ . We can see later in the third problem, when the lattice is infinitely large, approximately the  $T_c = 2.27$ , my result here is really close, considering the limit on the size and the random nature of the algorithm. I am going to use chi-square fit later in problem 3

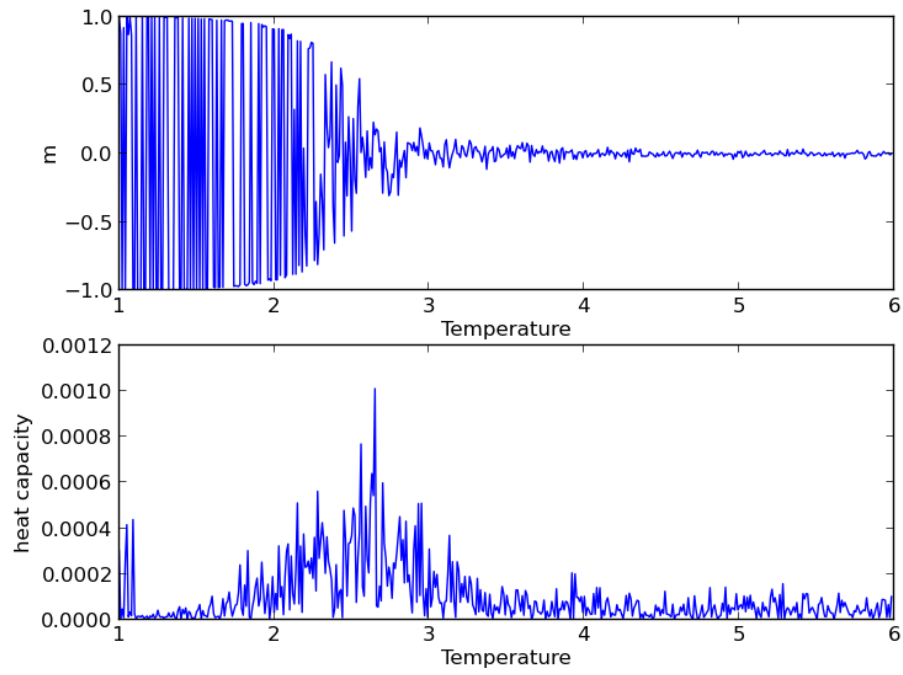


Figure 4: 10x10 lattice m and c vs. T

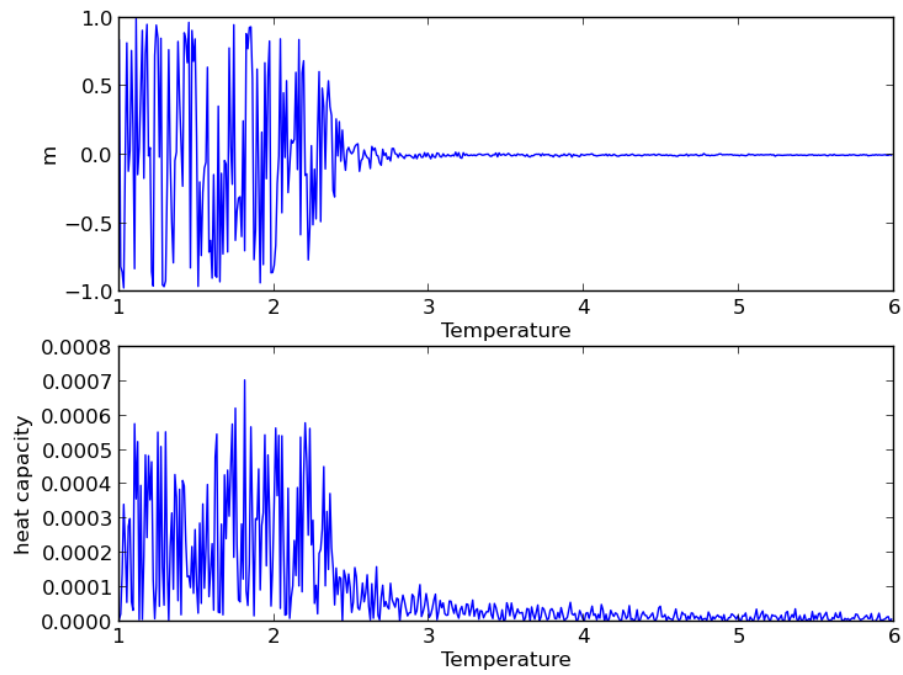


Figure 5: 50x50 lattice m and c vs. T

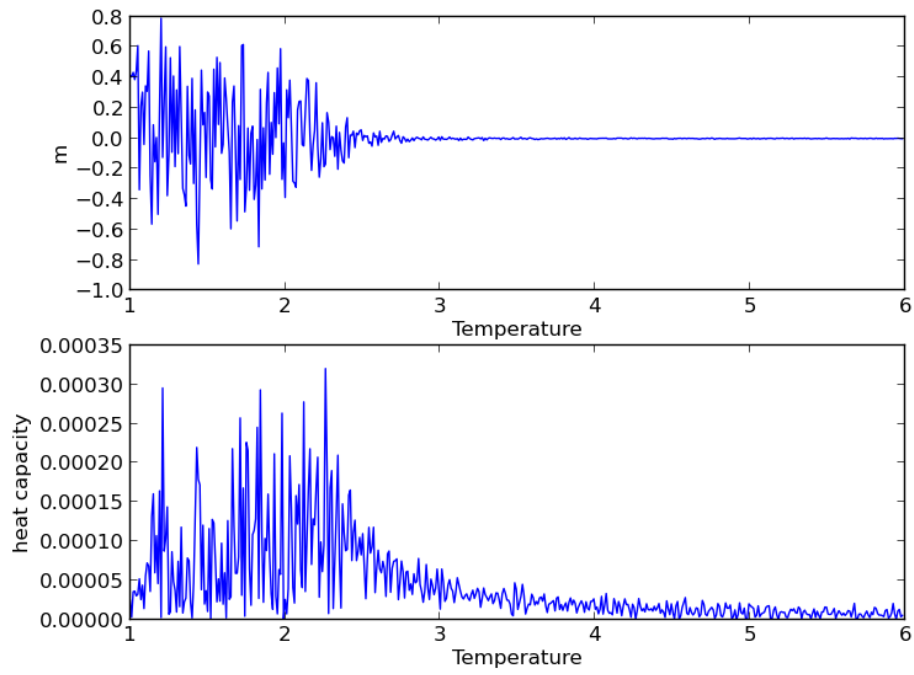


Figure 6: 100x100 lattice m and c vs. T

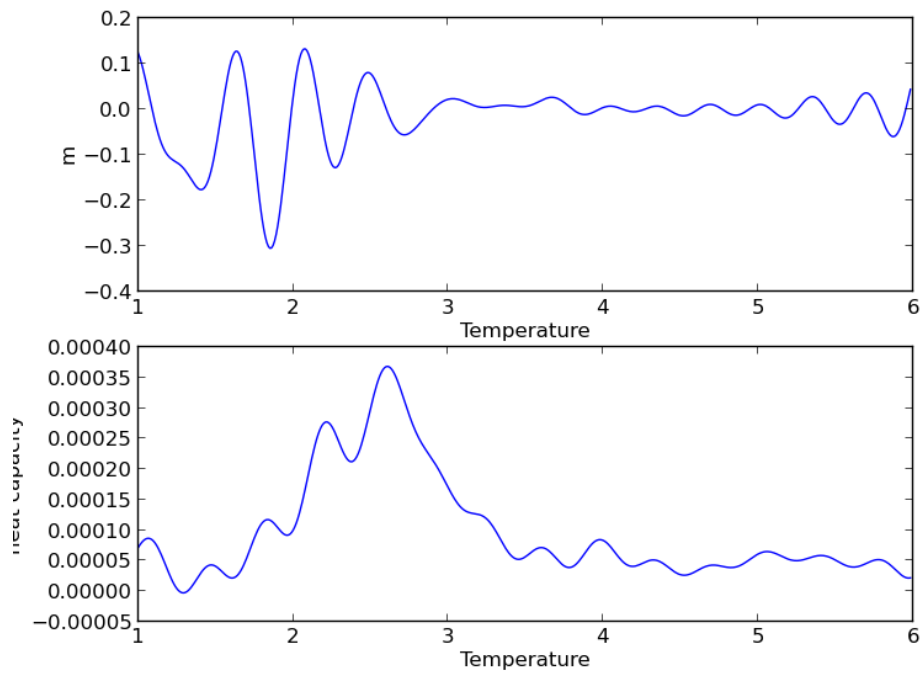
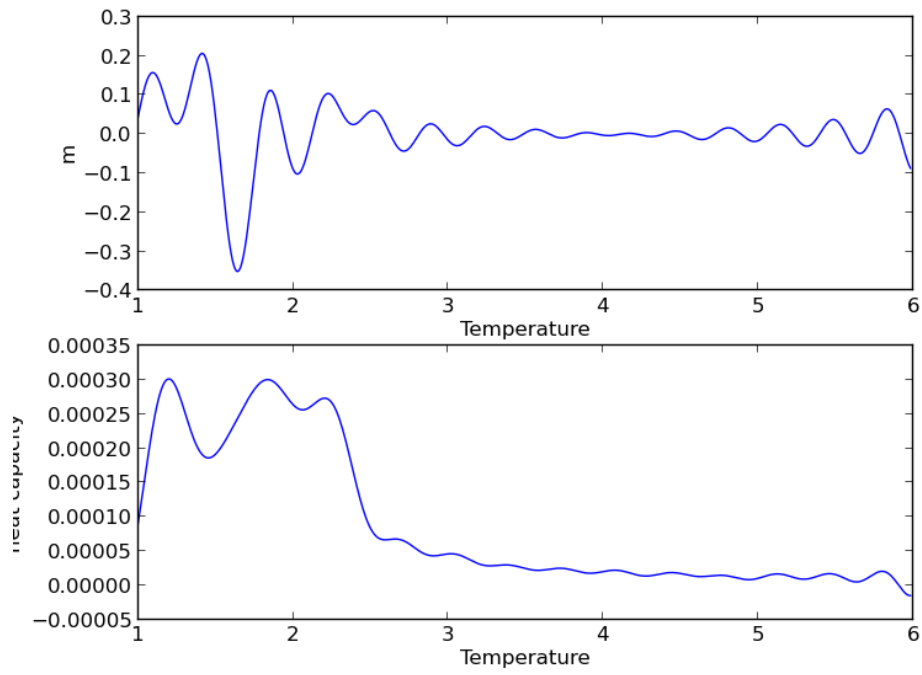
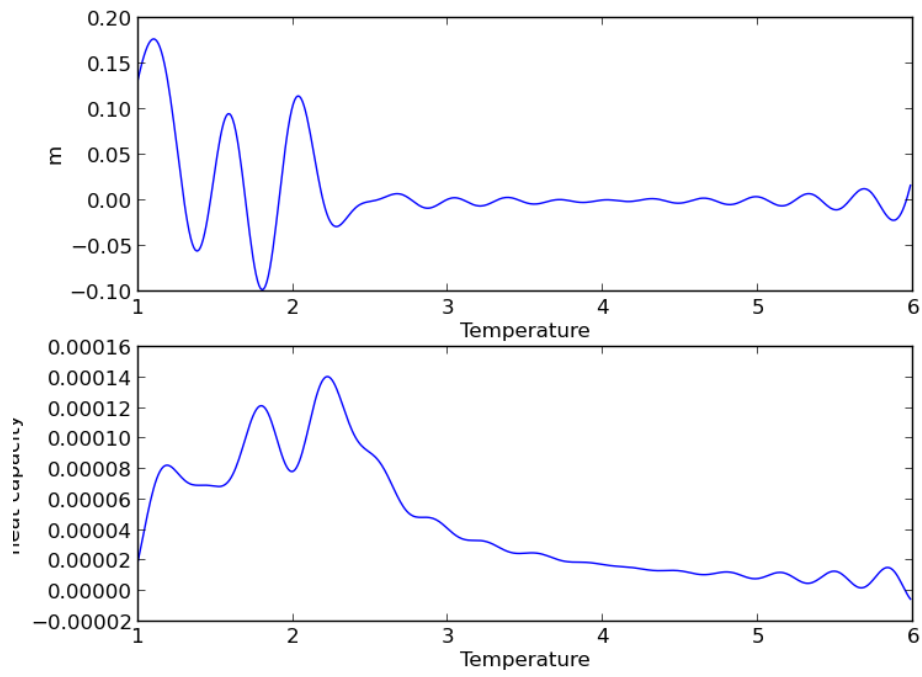


Figure 7: 10x10 fft lattice m and c vs. T

Figure 8: 50x50 ft lattice  $m$  and  $c$  vs.  $T$ Figure 9: 100x100 ft lattice  $m$  and  $c$  vs.  $T$

**Problem 3 (PHY 505 ONLY) :** Consider the Ising model for a 2-d simplified ferromagnet. In the thermodynamic limit  $N \rightarrow \infty$  and zero magnetic field  $H = 0$ , the magnetization per spin satisfies :

$$m = \lim_{N \rightarrow \infty} \frac{\langle \sum_i s_i \rangle}{N} = \begin{cases} \left[ 1 - \left\{ \sinh \left( \frac{2J}{k_B T} \right) \right\}^{-4} \right]^{1/8}, & \text{for } T \leq T_c \\ 0, & \text{for } T > T_c \end{cases}$$

Close to the critical temperature, the magnetization per spin satisfies :

$$m \sim (T_c - T)^\beta,$$

where  $\beta$  is a critical exponent. Onsager's formula shows that  $\beta = 1/8$  for the 2-D Ising model.

**a)** Recall that sudden reversals of the magnetization occur from time to time in systems of finite time. From the program "ising" in Lecture 32, what is a reasonable value of L to be chosen to maintain one single domain instead of flipping throughout? Show a plot of the average magnetization per spin.

**b)** Recall from class that the average magnetization per spin ( $m$ ) can be estimated using the above formula in the thermodynamic limit. Using the above value of L from Problem 1a, what is the minimum number of Monte Carlo steps that should be taken to ensure proper coverage in the "metropolis" algorithm? Why?

**c)** Using the values for L and number of steps from a and b respectively, run the "ising" program for  $H=0$ , and vary T from 2.0 to 2.5 to compute the average magnetization ( $\langle |M| \rangle$ ) as a function of temperature  $T$ . Use the data to identify the transition temperature  $T_c$  using the above relationship for  $m$ . Hint : you may find it useful to use one of the chi-squared fitting functions from earlier in the semester (Lecture 2).

Figure 10: p3 description, sorry I am a little lazy

### 3 Problem 3

#### 3.1 Description

#### 3.2 Result

##### 3.2.1 a

To theoretically determine the proper number of L require probability theory and using Gaussian distribution. Yet that might be too troublesome here, after several trials I found L=30 should be a reasonable number. Here is the plot: 11

##### 3.2.2 b

I firstly found all the Boltzmann factors are:0.026347980814448734

37.95357249735125

0.16232061118184818

6.160647084304639

1.0

1.0

6.160647084304639

0.16232061118184818

37.95357249735125

0.026347980814448734

According to the Metropolis algorithm, firstly we have fifty fifty chance to have a spin=1 or 0 ( $s_i$ ), for spin=0, the factor will always be 0 and this trial will not be used. For spin=1, after multiplying with Boltzmann factor, we can see six out of nine B-factors is larger than 1 and when times them, the trial will be effective. For other three, we can roughly think the chance the trial can be used is  $0.11(\frac{1}{3} \times 0.1623 \times 2 + \frac{1}{3} \times 0.026)$ . Thus, the chance one trial can be used is  $0.5 \times (\frac{2}{3} + 0.11) = 0.388$ . A



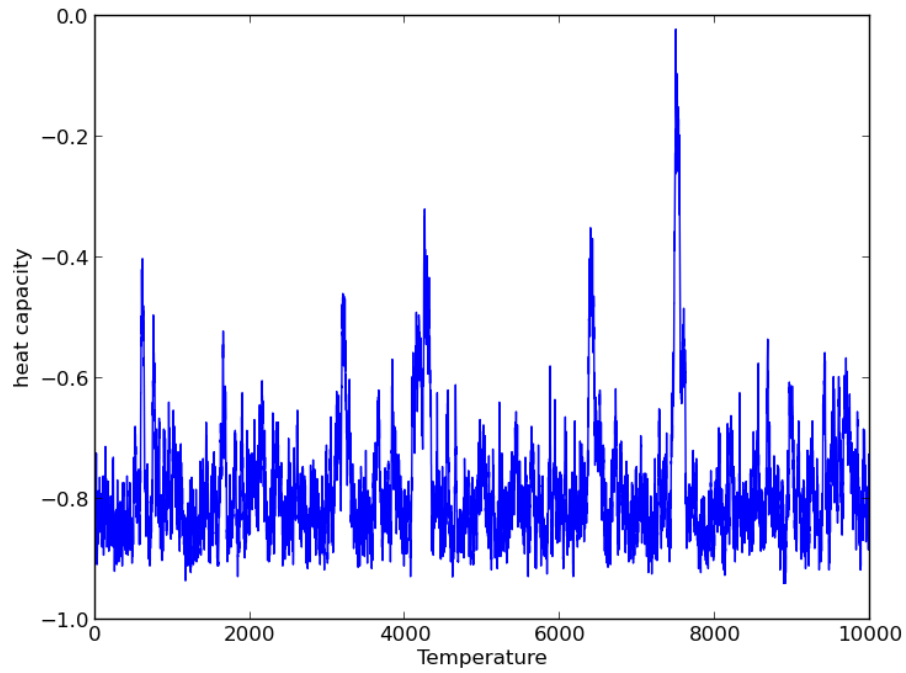


Figure 11: 30x30 average magnetization vs. steps

fair number of steps should satisfy that the expectation flipping number for individual spin greater than 1. Giving  $N > L * L / 0.388 = 2318$  I will say 2500 steps for simplicity.

### 3.2.3 c

Using the parameters above and I applied two chi-squared fits in two small ranges symmetric about a estimation of  $T_c$ , then find the intersect of these two lines. I found  $T_c = 2.272$ , very close to the predicted value. Fig. 12

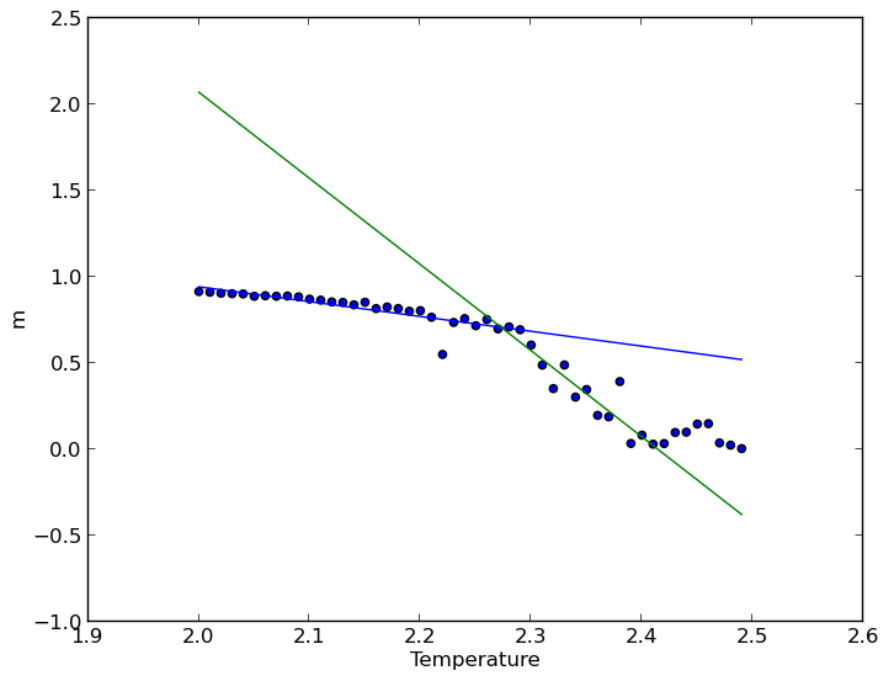


Figure 12: 30x30 average magnetization vs. temperature

## Acknowledgements

I discussed this assignment with my classmates and used material from the cited references, but this write-up is my own.

## References

- [1] PHY 410-505 Webpage, <http://www.physics.buffalo.edu/phy410-505>.

## A Appendix

### A.1 python code

The following python code was used to obtain the results in this report:

```
#include "cptstd.hpp"
#include "linalg.hpp"
#include "random.hpp"
#include <cmath>
using namespace cpt;

class Ising {
public :

    Ising(double iJ=1.0, int iL=10, int iN=100, double iT=2.0, double iH=0.0) :
        J(iJ), L(iL), N(iN), T(iT), H(iH), s(L,L), Lx(L), Ly(L)
    {
        s = Matrix<int,2>(Lx, Ly);
        for (int i = 0; i < Lx; i++)
            for (int j = 0; j < Ly; j++)
                s[i][j] = rng.rand() < 0.5 ? +1 : -1;    // hot start
        compute_boltzmann_factors();
        steps = 0;
    }

    void compute_boltzmann_factors()
    {
        for (int i = -8; i <= 8; i += 4) {
            w[i + 8][0] = exp( - (i * J - 2 * H) / T);
            w[i + 8][2] = exp( - (i * J + 2 * H) / T);
        }
    }

    bool metropolis_step()
    {
        // choose a random spin
        int i = int(Lx * rng.rand());
        int j = int(Ly * rng.rand());

        // find its neighbors using periodic boundary conditions
        int iPrev = i == 0 ? Lx-1 : i-1;
        int iNext = i == Lx-1 ? 0 : i+1;
        int jPrev = j == 0 ? Ly-1 : j-1;
        int jNext = j == Ly-1 ? 0 : j+1;
```

```

// find sum of neighbors
int sumNeighbors = s[iPrev][j] + s[iNext][j] + s[i][jPrev] + s[i][jNext];
int delta_ss = 2*s[i][j]*sumNeighbors;

// ratio of Boltzmann factors
double ratio = w[delta_ss+8][1+s[i][j]];
if (rng.rand() < ratio) {
    s[i][j] = -s[i][j];
    return true;
} else return false;
}

double acceptanceRatio;

void one_monte_carlo_step_per_spin ( ) {
    int accepts = 0;
    for (int i = 0; i < N; i++)
        if (metropolis_step())
            ++accepts;
    acceptanceRatio = accepts/double(N);
    ++steps;
}

double magnetizationPerSpin ( ) {
    int sSum = 0;
    for (int i = 0; i < Lx; i++)
        for (int j = 0; j < Ly; j++) {
            sSum += s[i][j];
        }
    return sSum / double(N);
}

double energyPerSpin ( ) {
    int sSum = 0, ssSum = 0;
    for (int i = 0; i < Lx; i++)
        for (int j = 0; j < Ly; j++) {
            sSum += s[i][j];
            int iNext = i == Lx-1 ? 0 : i+1;
            int jNext = j == Ly-1 ? 0 : j+1;
            ssSum += s[i][j]*(s[iNext][j] + s[i][jNext]);
        }
    return -(J*ssSum + H*sSum)/N;
}

protected :

Random rng; // random number generator

```

```

double J;                // ferromagnetic coupling
int L, Lx, Ly;           // number of spins in x and y
int N;                   // number of spins
Matrix<int,2> s;          // the spins
double T;                // temperature
double H;                // magnetic field

double w[17][3];         // Boltzmann factors

int steps;               // steps so far

};

int main (int argc, char *argv[]) {

    int Lx, Ly, N;
    double T,Th, H, c, e1,e2,e3,m1,m2;
    cout << "Two-dimensional Ising Model--Metropolis simulation\n"
    << "-----\n"
    << "Enter number of spins L in each direction: ";
    cin >> Lx;
    Ly = Lx;
    N = Lx*Ly;
    //      cout << " Enter temperature T: ";
    //      cin >> T;
    cout << "Enter magnetic field H: ";
    cin >> H;
    cout << "Enter number of Monte Carlo steps: ";
    int MCSteps;
    cin >> MCSteps;

    T=1.0;
    Th=0.01;
    int Nloop=500;
    e1=0.0;
    e2=0.0;
    e3=0.0;
    ofstream dataFile;
    dataFile.open("1isingcpp.data");
    for (int count=0; count<Nloop; count++ ) {

        Ising ising(1.0, Lx, N, T, H);

```

```

int thermSteps = int(0.2 * MCSteps);
cout << "Performing " << thermSteps
    << " steps to thermalize the system..." << flush;
for (int s = 0; s < thermSteps; s++)
    ising.one_monte_carlo_step_per_spin();

cout << "Done\nPerforming production steps..." << flush;
double mAv = 0, m2Av = 0, eAv = 0, e2Av = 0;

for (int s = 0; s < MCSteps; s++) {
    ising.one_monte_carlo_step_per_spin();
    double m = ising.magnetizationPerSpin();
    double e = ising.energyPerSpin();
    mAv += m; m2Av += m * m;
    eAv += e; e2Av += e * e;
    //dataFile << m << '\t' << e << '\n';
}
//dataFile.close();
mAv /= MCSteps; m2Av /= MCSteps;
eAv /= MCSteps; e2Av /= MCSteps;
e3=e2;
e2=e1;
e1=eAv;
m2=m1;
m1=mAv;

if (count==1) {
    c = abs(e2-e1)/2*Th;
    dataFile << T-Th << '\t' << m2 << '\t' << c << '\n';
}

else if(count>1 and count < Nloop-1) {
    c = abs(e1-e3)/2*Th;
    dataFile << T-Th << '\t' << m2 << '\t' << c << '\n';
    cout<< T << '\t' << m1 << '\t' << c << '\n';
}

else if(count==Nloop-1) {
    c = abs(e1-e3)/2*Th;
    dataFile << T-Th << '\t' << m2 << '\t' << c << '\n';
    c = abs(e1-e2)/Th;
    dataFile << T << '\t' << m1 << '\t' << c << '\n';
    cout<< T << '\t' << m1 << '\t' << c << '\n';
}

```

```

//          cout << " \n Magnetization and energy per spin written in file "//
//          //<< " \"ising.data\" " << endl;
//cout << " <m> = " << mAv << " +/- " << sqrt(m2Av - mAv*mAv) << endl;
//cout << " <e> = " << eAv << " +/- " << sqrt(e2Av - eAv*eAv) << endl;
T +=Th;
}

dataFile.close();
}

import math
import random
import matplotlib.pyplot as plt
from cpt import *

class Ising :
    def __init__(self, J=1.0, L=10, N=100, T=2., H=0.) :

        self.J = J                    # spin-spin coupling += ferro, -= antiferro
        self.L_x = L; self.L_y = L    # number of spins in x and y
        self.N = N                    # total number of spins
        self.s = []                   # L_x x L_y array of spin values
        self.T = T                     # Temperature
        self.H = H                     # magnetic field

        self.w = []                   # Boltzmann factors at fixed T and H
        self.steps = 0                 # Monte Carlo steps so far
        self.acceptance_ratio = 0      # accepted steps / total number of steps

        # create spin lattice and set spin randomly up or down (hot start)
        for i in range(self.L_x):
            self.s.append( [ ] )
            for j in range(self.L_y):
                self.s[i].append(random.choice( (-1, 1) ))
        self.compute_Boltzmann_factors()
        self.steps = 0

    def compute_Boltzmann_factors(self):
        self.w = []
        for m in range(5):
            self.w.append( [ ] )
            sum_of_neighbors = -4 + 2 * m
            for n in range(2):
                s_i = -1 + 2 * n
                factor = math.exp( -2.0 * (self.J * sum_of_neighbors + self.H) * s_i )
                self.w[m].append(factor)

```

```

def Metropolis_step_accepted(self):

    # choose a random spin
    i = random.randrange(self.L_x)
    j = random.randrange(self.L_y)

    # find the sum of neighbors assuming periodic boundary conditions
    sum_of_neighbors = ( self.s[(i-1)%self.L_x][j] + self.s[(i+1)%self.L_x][j]
                        self.s[i][(j-1)%self.L_y] + self.s[i][(j+1)%self.L_y])

    # access ratio of precomputed Boltzmann factors ,
    ratio = self.w[2 + int(sum_of_neighbors/2)][int((1 + self.s[i][j])/2)]

    # apply the Metropolis test
    if ratio > 1.0 or ratio > random.random():
        self.s[i][j] = -self.s[i][j]
        return True
    else:
        return False

def one_Monte_Carlo_step_per_spin(self):
    accepts = 0
    for n in range(self.N):
        if self.Metropolis_step_accepted():
            accepts += 1
    self.acceptance_ratio = accepts / float(self.N)
    self.steps += 1

def magnetization_per_spin(self):

    s_sum = 0.0
    for i in range(self.L_x):
        for j in range(self.L_y):
            s_sum += self.s[i][j]
    return s_sum / float(self.N)

def energy_per_spin(self):

    s_sum = 0.0
    ss_sum = 0.0
    for i in range(self.L_x):
        for j in range(self.L_y):
            s_sum += self.s[i][j]

```



```

        ss_sum += self.s[i][j] * (self.s[(i+1)%self.L_x][j] + self.s[i][(j-1)%self.L_y])
    return -(self.J * ss_sum + self.H * s_sum) / float(self.N)

print "Two-dimensional Ising Model--Metropolis simulation"
print "-----"
L = int(input("Enter number of spins L in each direction: "))
#T = float(input("Enter temperature T: "))
H = float(input("Enter magnetic field H: "))
MC_steps = int(input("Enter number of Monte Carlo steps: "))

def my_range(start, end, step):
    while start <= end:
        yield start
        start += step

Tout=[]
mav=[]
eav=[]
hcout=[]
T=2.0
Nl=50
Th=0.01
fp=open("p3data","w")
for count in range(Nl):
    print count
    ising = Ising(L=L, N=L*L, T=T, H=H)

    therm_steps = int(0.2 * MC_steps)

    print "Performing", therm_steps, "thermalization steps..."
    for i in range(therm_steps):
        ising.one_Monte_Carlo_step_per_spin()

    print "Done...Performing production steps..."
    m_av = 0.0; m2_av = 0.0; e_av = 0.0; e2_av = 0.0
    #data_file = open("isingT.data", "w")
    for i in range(MC_steps):
        ising.one_Monte_Carlo_step_per_spin()
        m = ising.magnetization_per_spin()
        e = ising.energy_per_spin()
        m_av += m
        m2_av += m**2
        e_av += e
        e2_av += e**2
        #data_file.write(repr(m) + "\t" + repr(e) + "\n")
    #data_file.close()

```

```

print "M/spin_and_E/spin_values_written_in_isingT.data"
m_av /= float(MC_steps)
m2_av /= float(MC_steps)
e_av /= float(MC_steps)
e2_av /= float(MC_steps)
mav.append(abs(m_av))
eav.append(e_av)
Tout.append(T)
s = '{0:8.4f}_{1:8.4f}\n'.format(T,abs(m_av))
fp.write(s)
T += Th
#   if count==1 :
#       c = (eav[1]-eav[0])/Th
#       print c
#       hcout.append(c)                                #hcout[0]
#   elif count>1 and count < Nl-1:
#       c = (eav[count]-eav[count-2])/Th*2
#       print c
#       hcout.append(c)                                #hcout[count-1]
#   elif count == Nl-1:
#       c = (eav[count]-eav[count-2])/Th*2
#       print c
#       hcout.append(c)
#       c = (eav[count]-eav[count-1])/Th
#       print c
#       hcout.append(c)                                #hcout[N-1] hcout[N-2]
#print " <m> =", m_av, "+/-", math.sqrt(m2_av - m_av**2)
#print " <e> =", e_av, "+/-", math.sqrt(e2_av - e_av**2)
#print ising.w

print Tout, len(Tout)
print mav, len(mav)
#print hcout, len(hcout)
#plt.subplot(2, 1, 1)
plt.plot(Tout, mav)
plt.xlabel('Temperature')
plt.ylabel('m')

#plt.subplot(2,1,2)
#plt.plot(Tout, hcout)
#plt.xlabel('Temperature')
#plt.ylabel('heat capacity')

plt.show()

import matplotlib.pyplot as plt
import math

```

```

from numpy import array , real
from fft import fft , fft_power , ifft

filename = "l100.data"
file = open(filename , "r")
lines = file.readlines()
T=[]
m=[]
c=[]

for line in lines:
    if line != '\n' :
        words = line.split()
        tv,mv,cv = [float(s) for s in words]
        T.append(tv)
        m.append(mv)
        c.append(cv)

N=len(m)
log2N = math.log(N,2)
if log2N - int(log2N) > 0.0 :
    print 'Padding with zeros!'
    pads = [0.0] * (pow(2, int(log2N)+1) - N)    #now paddle with 0s
    ma = m + pads
    ca = c + pads
    N = len(m)
    print len(m)
M = fft(ma)
C = fft(ca)

maxfreq = 15
# Now smooth the data
for i in range(maxfreq, len(M)-maxfreq) :
    M[i] = complex(0,0)
    C[i] = complex(0,0)

cp = ifft(C)
mp = ifft(M)
cpr = real(cp)
mpr = real(mp)

mf = mpr[0:len(m)]
cf = cpr[0:len(c)]

plt.subplot(2, 1, 1)
plt.plot(T,mf)

```

```
plt.xlabel('Temperature')
plt.ylabel('m')

plt.subplot(2,1,2)
plt.plot(T, cf)
plt.xlabel('Temperature')
plt.ylabel('heat_capacity')

plt.show()
```