# PHY 410 Homework Assignment 8

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# November 9, 2014

#### Abstract

The goal of this assignment is to get more familiar with RK-4 method in ODE, as well as other methods, examples of celestial orbits, quantum harmonic oscillator and Kronig-Penney model will be presented.

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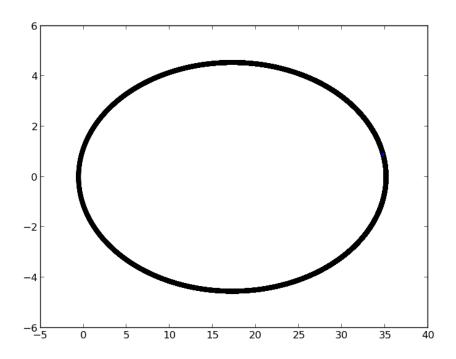


Figure 1: Halley Comet orbit with RK4 method

## 1 Problem 1

#### 1.1 Description

Generate the orbits of (1) Halley's Comet and at least one planetary orbit using fixed-time step Runge-Kutta and adaptive time step Runge-Kutta and compare the efficiency of the two methods (you can estimate errors by measuring the period for example); and (2) a few interesting orbits near the Lagrangian points of the restricted 3-body problem, for example the Halo orbit selected by NASA for the James Webb telescope described by Dr. Mather earlier this year, or a Lissajous orbit.

## 1.2 Numerical Analysis

#### 1.2.1 part a

Using the program, combined with the data from wikipedia [2], when using RK-4 method: Enter aphelion distance in AU: 35.1

Enter eccentricity: 0.967

Semimajor axis a = 17.8444331469 AU

Period T = 75.3796531592 yr $v_u(0) = 0.192656328719 \text{ AU/yr}$ 

With the time step 0.001, the result is shown below: Fig. 1

When using RK-4 adaptive method: Enter aphelion distance in AU: 35.1

Enter eccentricity: 0.967

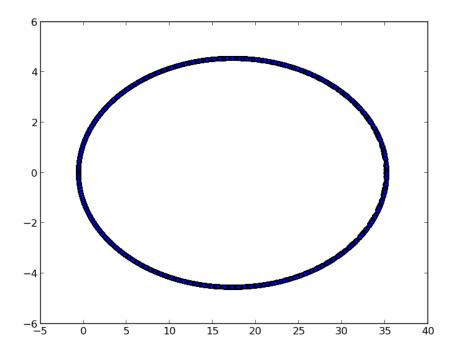


Figure 2: Halley Comet orbit with RK4-adaptive method

Semimajor axis a = 17.8444331469 AU

 ${\rm Period}\ {\rm T} = 75.3796531592\ {\rm yr}$ 

 $v_u(0) = 0.192656328719 \text{ AU/yr}$ 

Enter step size dt: 0.001

Enter desired accuracy for adaptive integration: 0.000001

The result is shown below: Fig 2

I found in this case the regular RK-4 method is more efficient.

#### 1.2.2 part b

With the parameters included in here [3], around point L2, with the proper initial condition, I generate the halo orbit shown here: Fig. 3. We can see it's fairly a good plot, although since the data is not precise enough, a little "Lissajous orbit" style will be mixed into the simulation plot.

1 PROBLEM 1 4

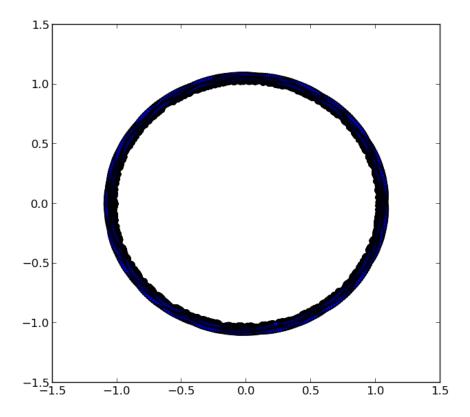


Figure 3: Halo orbit

2 PROBLEM 2 5

#### 2 Problem 2

#### 2.1 Description

Modify the Schroedinger code to add (1) a cubic  $x^3$  term and (2) a quartic  $x^4$  term, to the harmonic oscillator potential. Discuss how each type of perturbation modifies the energy levels and eigenfunctions. Be sure to increase the left and right boundaries if needed!

#### 2.2 Result

To be consistent with the perturbation theory, I added a parameter  $\lambda$  before the cubic or the quartic term, and making  $\lambda$  a rather small number. I chose E=5 in our case so that the number of eigenvalues is appropriate. Without the perturbation, the result is: Level Energy Simple Steps Secant Steps

```
-1.001250 0.500076
1.001250 0.500076
-1.732500 1.500095
1.732500 1.500095
-2.236250 2.500113
```

-2.646250 3.500115 2.646250 3.500115

 $2.236250\ 2.500113$ 

 $-3.001250 \ 4.500120$  $3.001250 \ 4.500120$ 

-3.317500 5.500139 3.317500 5.500139 With the diagram: Fig 4

When it comes to cubic perturbation, with  $\lambda = 0.05$ :

Level Energy Simple Steps Secant Steps

```
-9.903750 0.475473
0.933750 0.475473
-9.697500 1.422769
```

-9.482500 2.329516 1.973750 2.329516

 $1.568750 \ 1.422769$ 

-9.262500 3.166497 2.272500 3.166497

-9.041250 3.921575

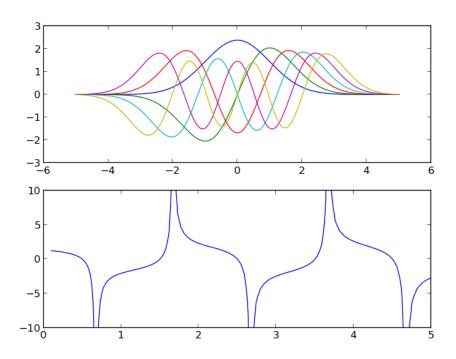


Figure 4: Harmonic oscillation without perturbation

 $2.505000\ 3.921575$ 

-8.811250 4.618098

 $2.697500\ 4.618098$ 

 $\begin{array}{c} -8.198750 \ 6.054249 \\ 3.047500 \ 6.054249 \end{array}$ 

And the plot: 5 When quadric perturbation, with  $\lambda=0.05$ : Level Energy Simple Steps Secant Steps

 $\begin{array}{c} -0.986250 \ 0.532749 \\ 0.986250 \ 0.532749 \end{array}$ 

-1.620000 1.653598

1.620000 1.653598

 $\begin{array}{c} -2.021250 \ 2.874194 \\ 2.021250 \ 2.874194 \end{array}$ 

 $-2.328750\ 4.176595$ 

2 PROBLEM 2 7

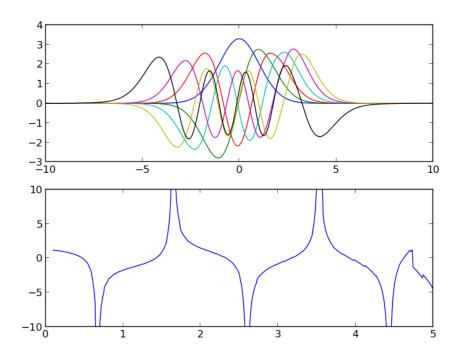


Figure 5: Harmonic oscillation with cubic perturbation

 $2.328750\ 4.176595$ 

 $\begin{array}{c} -2.581250\ 5.549612 \\ 2.581250\ 5.549612 \end{array}$ 

And the plot: 6

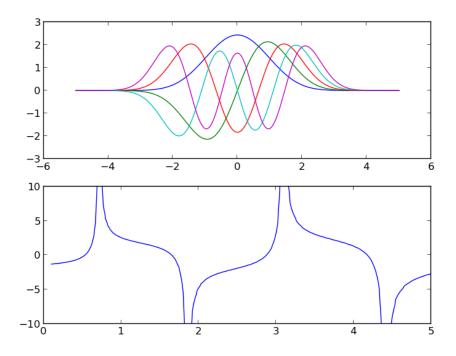


Figure 6: Harmonic oscillation with quadric perturbation

We can see, with the cubic perturbation, the number of energy level increase from 5 to 6, the whole numbers of level shifted and the values of energy for the first 5 levels decrease a little, while with the quadric perturbation, the number remains the same and the values of energy increase a little. Those property origin from the form of the perturbation, namely the cubic form is antisymmetric and the quadric one is symmetric.

## 3 Problem 3

#### 3.1 Description

Study the dependence of the Kronig-Penney model band structure on the width and height of the potential step, and describe any systematic trends you notice.

## 3.2 Result

Guided by the theoretical result as well as the symmetric property of the system. Based on the result of different trials, I found when fix the width and increase the height, the band gap increase Fig 7. When fix the height and change the width, the more close the width is to half of the size of the unit cell, the bigger the gap. When the width is approaching 0 or the size of the unit cell, the gap is vanishing. 8

Here is the plot with width 0.2 height -5.0 Fig 9

Here is the plot with width 0.5 height -5.0 Fig 10

Here is the plot with width 0.2 height -10.0 Fig 11

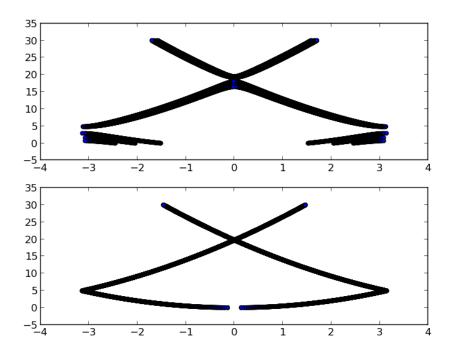


Figure 7: Kronig Penney with height -5 -8 -10 and width 0.2

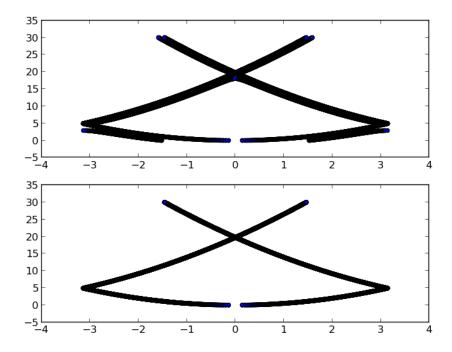


Figure 8: Kronig Penney with width 0.2 0.5 1.0 and height -5.0  $\,$ 

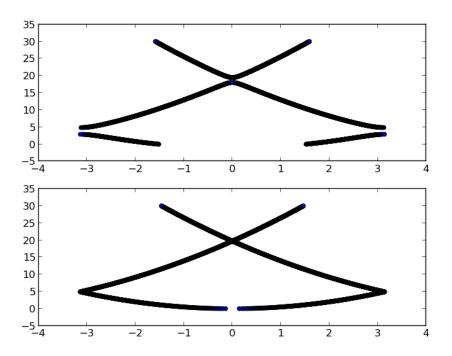


Figure 9: Kronig Penney with width 0.2 height -5.0

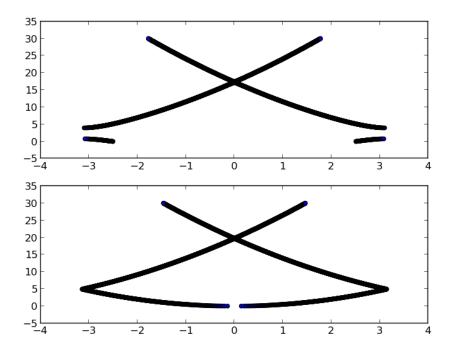


Figure 10: Kronig Penney with width 0.5 height -5.0  $\,$ 

REFERENCES 11

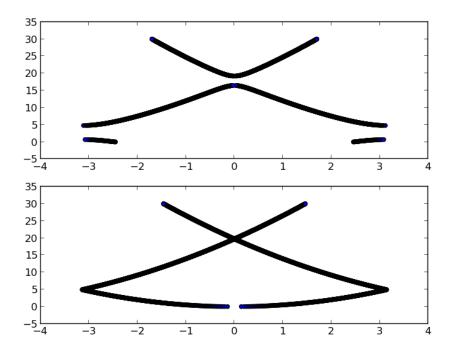


Figure 11: Kronig Penney with width 0.2 height -10.0

# Acknowledgements

I discussed this assignment with my classmates and used material from the cited references, but this write-up is my own.

## References

- [1] PHY 410-505 Webpage, http://www.physics.buffalo.edu/phy410-505.
- [2] Halley Comet Wikipedia http://en.wikipedia.org/wiki/Halley%27s\_Comet
- [3] Lagrange point http://www.physics.montana.edu/faculty/cornish/lagrange.pdf

 $A \quad APPENDIX$  12

## A Appendix

## A.1 python code

```
The following python code was used to obtain the results in this report:
\# pendul - Program to compute the motion of a simple pendulum
# using the Euler or Verlet method
from cpt import *
from math import *
from numpy import array
class Kepler:
    G_m1_plus_m2 = 0.
    step_using_y = False
    def __init__( self , G_m1_plus_m2 , step_using_y ):
         self.G_m1_plus_m2 = G_m1_plus_m2
         self.step_using_y = step_using_y
    \mathbf{def} = \mathbf{call} = (\mathbf{self}, \mathbf{p}) :
        t = p[0]
        x = p[1]
        y = p[2]
        vx = p[3]
        vy = p[4]
        r = math.sqrt(x**2 + y**2)
        ax = - self.G_m1_plus_m2 * x / r**3
        ay = - self.G_m1_plus_m2 * y / r**3
        flow = [1, vx, vy, ax, ay]
        if self.step_using_y:
                                            # change independent variable from t to y
             for i in range (5):
                 flow[i] /= vy
        return flow
def main () :
    kepler = Kepler (G_m1_plus_m2 = 4 * math.pi**2, step_using_y=False
)
    #* Select the numerical method to use: Euler or Verlet
    method = input( "Choose_a_numerical_method_: RK4_{-}(0), or_Adaptive_RK4_{-}(1): _")
    print "_Kepler_orbit_using_fixed_and_then_adaptive_Runge-Kutta"
    r_aphelion = float (input ("_Enter_aphelion_distance_in_AU:_"))
    eccentricity = float (input ("_Enter_eccentricity:_"))
    a = r_aphelion / (1 + eccentricity)
```

```
T = a **1.5
vy0 = math. sqrt(kepler. G_m1_plus_m2 * (2 / r_aphelion - 1 / a))
print "_Semimajor_axis_a_=_", a, "_AU"
print "_Period_T_=_", T, "_yr"
print "_v_y(0) == ", vy0, "_AU/yr"
tau = float (input ("_Enter_step_size_dt:_"))
accuracy = float(input("_Enter_desired_accuracy_for_adaptive_integration:_"))
# Initialize vector
xv = [0.0] * 5
xv[0] = 0.0
xv[1] = r_aphelion
xv[2] = 0.0
xv[3] = 0.0
xv[4] = vy0
#* Loop over desired number of steps with given time step
# and numerical method
nStep = input ( "Enter_number_of_time_steps:_" )
t_plot = []
x_plot = []
y_plot = []
dt_{-}min = tau
dt_{-}max = tau
for iStep in xrange(nStep):
    \#if (xv/0) > T :
    \# break
    #* Record angle and time for plotting
    t_{plot.append}(xv[0])
    x_{plot.append}(xv[1])
    y_plot.append(xv[2])
       method = 0:
        RK4_step(xv, tau, kepler)
    elif method == 1 :
        tau = RK4_adaptive_step(xv, tau, kepler, accuracy)
        dt_{-min} = min(tau, dt_{-min})
        dt_{max} = max(tau, dt_{max})
import matplotlib
matplotlib.rcParams['legend.fancybox'] = True
```

```
import matplotlib.pyplot as plt
    plt.scatter(x_plot, y_plot)
    plt.show()
if _-name_- = "_-main_-" :
    main()
import math
from cpt import *
\# the restricted circular planar 3-body problem has one parameter
\# represent a point in the extended phase space by a 5-component vector
\# trv = [t, r, v] = [t, x, y, vx, vy]
class Rcp3Body :
    \# alpha = m2/(m1+m2) in the webnotes
    a = 0.0
    # switch to zero in on Poincare section point
    # use y instead of t as independent variable
    step_using_y = False
    def __init__(self , a , step_using_y ) :
        self.a = a
        self.step_using_y=step_using_y
    def set_step_using_y( self , step_using_y ) :
        self.step\_using\_y = step\_using\_y
    def __call__(self, trv):
                                     # equations in co-rotating frame
        t = trv[0]
        x = trv[1]; y = trv[2]; vx = trv[3]; vy = trv[4]
        d1 = \text{math.pow}((x - \text{self.a}) **2 + y **2, 1.5)
        d2 = \text{math.pow}((x + 1 - \text{self.a}) **2 + y **2, 1.5)
        ax = -(1 - self.a) * (x - self.a) / d1 - self.a * (x + 1 - self.a) / d2 + x
        ay = -(1 - self.a) * y / d1 - self.a * y / d2 + y - 2 * vx
        flow = [1.0, vx, vy, ax, ay]
        if self.step_using_y:
                                     \# change integration variable from t to y
```

```
for i in range(len(flow)):
                   flow[i] /= vy
         return flow
    def Jacobi (self, trv): # Jacobi Integral
         t = trv[0]
         x = trv[1]; y = trv[2]; vx = trv[3]; vy = trv[4]
         r1 = math. sqrt((x - self.a)**2 + y**2)
         r2 = math. sqrt((x + 1 - self.a)**2 + y**2)
         return x**2 + y**2 + 2 * (1 - self.a) / r1 + 2 * self.a / r2 - vx**2 - vy**
    \mathbf{def} \ \mathbf{f}_{-\mathbf{x}}(\mathbf{self}, \mathbf{x}): \quad \# \ effective \ x \ component \ of \ force \ on \ the \ x-axis
         return (x - (1 - self.a) * (x - self.a) / abs(x - self.a) / (x - self.a) **
                      - self.a * (x + 1 - self.a) / abs(x + 1 - self.a) / (x + 1 - self.a)
    \mathbf{def} \ \mathbf{zero} \ (\mathbf{self}, \mathbf{f}, \mathbf{x}_{\mathbf{lower}}, \mathbf{x}_{\mathbf{upper}}, \mathbf{accuracy} = 1.0 \mathbf{e} - 6, \mathbf{max}_{\mathbf{steps}} = 1000):
         \# use bisection search to solve f(x)=0 in interval [x_lower, x_upper]
         x_{mid} = (x_{upper} + x_{lower}) / 2
         dx = x_{-upper} - x_{-lower}
         f_{lower} = f(x_{lower})
         step = 0
         while abs(dx) > accuracy:
              f_{-mid} = f(x_{-mid})
              \mathbf{if} \quad f_{-} \text{mid} == 0:
                   dx = 0
              else:
                   if f_{\text{lower}} * f_{\text{mid}} > 0:
                        x_{lower} = x_{mid}
                        f_{\text{-lower}} = f_{\text{-mid}}
                   else:
                        x_{upper} = x_{mid}
                   x_{mid} = (x_{upper} + x_{lower}) / 2
                   dx = x_{upper} - x_{lower}
              step += 1
               assert step < max_steps, "_zero(f)_too_many_steps_" + str(max_steps)
         return x_mid
# get parameters from user
print "_Restricted_circular_planar_3-body_problem"
while True:
    a = input(" \_Enter \_alpha \_= \_m2/(m1+m2) \_> \_0 \_and \_< \_0.5: \_")
     if a \le 0 or a > 0.5:
```

```
print "_Bad_alpha , _please_try_again"
                     continue
          break
rcp3Body = Rcp3Body (a, False)
# find Lagrangian points for this alpha
eps = 1e-6
print "_Lagrangian_points:"
print "L1_: Lx_=_", rcp3Body.zero(rcp3Body.f_x, rcp3Body.a - 1 + eps, rcp3Body.a -
\mathbf{print} \ \text{``} \bot L3 \bot : \bot x \bot = \bot \text{''}, \ \operatorname{rcp3Body}. \ \operatorname{zero}\left(\operatorname{rcp3Body}. \ \operatorname{f\_x}, \ \operatorname{rcp3Body}. \ \operatorname{a} + \ \operatorname{eps}, \ 1.5\right), \ \text{``} \bot y \bot = \bot \text{''}, \ 0
print "L5_: Lx_=_", rcp3Body.a - 0.5, "_y_=_", math.sqrt(3.0) / 2
# get initial values
one = input ("\_Enter\_0\_to\_specify\_[x,y,vx,vy]\_or\_1\_to\_specify\_C\_and\_[x,y,vx]:\_")
if one:
          while True:
                     C = input("_Enter_value_of_the_Jacobi_integral_C:_")
                     x, y, vx = input("\_Enter\_x, \_y, \_vx: \_")
                     r1 = math. sqrt((x - a)**2 + y**2)
                     r2 = math. sqrt ((x + 1 - rcp3Body.a)**2 + y**2)
                     vy_sqd = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / r1 + 2 * rcp3Body.a / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + 2 * (1 - rcp3Body.a) / range = -C + x**2 + y**2 + y
                     if vy_sqd < 0:
                                print "_Sorry_C_too_large,_cannot_solve_for_vy"
                     else:
                                vy = math.sqrt(vy\_sqd)
                                print "_vy_=_", vy
                                break
else:
          x, y, vx, vy = input("\_Enter\_x, \_y, \_vx, \_vy: \_")
          print "_Jacobi_Integral_C_=_", rcp3Body.Jacobi([0, x, y, vx, vy])
t_max = input("_Enter_maximum_integration_time_t_max:_")
crossing = 0
dt = 0.01
t = 0
trv = [t, x, y, vx, vy]
x_{plot} = []
y_plot = []
while t < t_max:
```

```
# write trajectory point
    t = trv[0]
    x = trv[1]; y = trv[2]; vx = trv[3]; vy = trv[4]
    y_save = y
                   # remember y to check section crossing
    x_{plot.append}(x)
    y_plot.append( y )
    \# use adaptive Runge-Kutta with default accuracy
    dt = RK4_adaptive_step(trv, dt, rcp3Body)
    \# Poincare section at y = 0 and vy positive
    x = trv[1]; y = trv[2]; vx = trv[3]; vy = trv[4]
    if y_save < 0 and y >= 0 and vy >= 0:
        rcp3Body.set_step_using_y(True)
        dy = -y
        RK4_step(trv, dy, rcp3Body)
        t = trv[0]; x = trv[1]; y = trv[2]; vx = trv[3]; vy = trv[4]
        crossing += 1
        print "_Crossing_No.", crossing, "_at_t_=_", t, "_C_=_", rcp3Body.Jacobi(ti
        rcp3Body.set_step_using_y(False)
import matplotlib
matplotlib.rcParams['legend.fancybox'] = True
import matplotlib.pyplot as plt
plt.scatter(x_plot, y_plot, c='blue')
plt.show()
import math
import cpt
import matplotlib.pyplot as plt
class Schroedinger:
    \mathbf{def} __init__(self) :
        self.hbar = 1.0
                                         # Planck's constant / 2pi
        self.m = 1.0
                                         # particle mass
        self.omega = 1.0
                                         # oscillator frequency
        self.E = 0.0
                                         # current energy in search
        self.N = 500
                                         \# number of lattice points = N+1
        self.x_left = -5.0
                                         # left boundary
        self.x_right = 5.0
                                         # right boundary
        self.h = (self.x_right - self.x_left) / self.N # grid spacing
        self.phi_left = [0.0]*(self.N+1) # wave function integrating from left
```

```
self.phi_right = [0.0]*(self.N+1)
                                          \# wave function integrating from right
    self.phi = [0.0]*(self.N+1)
                                            # whole wave function
    self.sign = 1
                                       \# current sign used to make F(E) continuous
    self.nodes = 0
                                       # current number of nodes in wavefunction
                                    \# \ harmonic \ oscillator \ potential
\mathbf{def} \ V(self, x):
    return 0.5 * self.m * self.omega**2 * x**2 + 0.05*x**4
                                    # Sturm-Liouville q function
\mathbf{def} \ \mathbf{q}(\mathbf{self}, \mathbf{x}):
    return 2 * self.m / self.hbar**2 * (self.E - self.V(x))
                                   \# eigenvalue \ at \ F(E) = 0
def F(self, energy):
    # set energy needed by the q(x) function
    self.E = energy
    # find the right turning point
    i_match = self.N
    x = self.x_right
                                  \# start at right boundary
    while self.V(x) > self.E:
                                       # in forbidden region
        i_match = 1
        x = self.h
        if i_{\text{match}} < 0:
            raise Exception ("can't_find_right_turning_point")
    \# integrate self. phi\_left using Numerov algorithm
    self.phi_left[0] = 0.0
    self.phi_left[1] = 1.0e-10
                                   # constant in Numerov formula
    c = self.h**2 / 12.0
    for i in range (1, i_{-}match+1):
        x = self.x_left + i * self.h
        self.phi_left[i+1] = 2 * (1 - 5 * c * self.q(x)) * self.phi_left[i]
        self.phi_left[i+1] = (1 + c * self.q(x - self.h)) * self.phi_left[i-1]
        self.phi_left[i+1] /= 1 + c * self.q(x + self.h)
    \# integrate self.phi_right
    self.phi[self.N] = self.phi_right[self.N] = 0.0
    self.phi[self.N-1] = self.phi_right[self.N-1] = 1.0e-10
    for i in range (self.N -1, i_match -1, -1):
        x = self.x_right - i * self.h
        self.phi_right[i-1] = 2 * (1 - 5 * c * self.q(x)) * self.phi_right[i]
        self.phi_right[i-1] = (1 + c * self.q(x + self.h)) * self.phi_right[i-1]
        self.phi_right[i-1] /= 1 + c * self.q(x - self.h)
        self.phi[i-1] = self.phi_right[i-1]
```

 $\# rescale self.phi_left$ 

```
scale = self.phi_right[i_match] / self.phi_left[i_match]
        for i in range(i_match + 2):
             self.phi_left[i] *= scale
             self.phi[i] = self.phi_left[i]
        \# make F(E) continuous
        # count number of nodes in self.phi_left
        n = 0
        for i in range (1, i_{-}match+1):
            if self.phi_left[i-1] * self.phi_left[i] < 0.0:
        # flip its sign when a new node develops
        if n != self.nodes:
            self.nodes = n
             self.sign = -self.sign
        return ( self.sign *
                  (self.phi_right[i_match-1] - self.phi_right[i_match+1] -
                    self.phi_left [i_match-1] + self.phi_left[i_match+1]) /
                 (2 * self.h * self.phi_right[i_match]) )
    def normalize (self):
        norm = 0.0
        for i in range (self.N):
            norm += self.phi[i]**2
        norm /= self.N
        norm = math.sqrt(norm)
        for i in range (self.N):
             self.phi[i] /= norm
print "_Eigenvalues_of_the_Schroedinger_equation"
print "_for_the_harmonic_oscillator_V(x) = 0.5 x^2"
E_max = input("\_Enter\_maximum\_energy\_E:\_")
phi_file = open("phi.data", "w")
level = 0
                             \# level number
E_{\text{old}} = 0.0
                             # previous energy eigenvalue
schroedinger = Schroedinger()
schroedinger.E = 0.1 \# guess \ and \ E \ below \ the \ ground \ state
```

```
\# draw the potential
for i in range (schroedinger.N+1):
        x = schroedinger.x_left + i * schroedinger.h
swrite = (0:12.6 f) \{1:12.6 f\}'.format(x, schroedinger.V(x))
print swrite
print ''
# find the energy levels
print "Level____Energy____Simple_Steps___Secant_Steps"
                            print "_-----
x_data = []
phi_data = []
                              # loop over levels
while True:
        \# estimate next E and dE
        dE = 0.5 * (schroedinger.E - E_old)
        E_{\text{old}} = \text{schroedinger.} E
        schroedinger.E += dE
        \# use simple search to locate root with relatively low accuracy
        accuracy = 0.01
        schroedinger.E = cpt.root_simple(schroedinger.F, schroedinger.E, dE, accura
        \#simple\_steps = cpt.root\_steps()
        # use secant search with relatively high accuracy
        accuracy = 1.0e-6
        \mathrm{E1} = \mathrm{schroedinger.E} + 100 * \mathrm{accuracy} \ \# \ \mathit{guess} \ \mathit{second} \ \mathit{point} \ \mathit{required}
        schroedinger.E = cpt.root_secant(schroedinger.F, schroedinger.E, E1, accura
        \#secant\_steps = cpt.root\_steps()
        \#ans = "" + repr(level).rjust(3) + ""*5 + repr(E).ljust(20) + ""*6
        \#ans += repr(simple\_steps). rjust(3) + ""*11 + repr(secant\_steps). rjust(3)
        \#print ans
        level += 1
        accuracy = 0.001
        x = cpt.root_simple(schroedinger.q, schroedinger.x_left, schroedinger.h, ac
        swrite = {}^{\prime} \{0:12.6f\} \cup \{1:12.6f\} {}^{\prime}.format(x, schroedinger.E)
        print swrite
        x = cpt.root_simple(schroedinger.q, schroedinger.x_right, -schroedinger.h,
        swrite = \{0:12.6f\} \cup \{1:12.6f\}'.format(x, schroedinger.E)
        print swrite
        print ''
```

```
schroedinger.normalize()
        x_data.append([])
        phi_data.append([])
        iphi = len(x_data) - 1
        for i in range (schroedinger.N+1):
                x = schroedinger.x_left + i * schroedinger.h
                x_{data}[iphi].append(x)
                phi_data[iphi].append( schroedinger.phi[i] )
        if schroedinger.E >= E_max:
                                              # we are done
                break
# print the search function
schroedinger.E = 0.1
dE = 0.01
E_{data} = []
F_{data} = []
while schroedinger. E < E_max:
        E_data.append( schroedinger.E )
        F_data.append( schroedinger.F(schroedinger.E) )
        schroedinger.E += dE
s1 = plt.subplot(2,1,1)
for i in range(len(x_data)) :
        plt.plot(x_data[i], phi_data[i])
s2 = plt.subplot(2,1,2)
plt.plot(E_data, F_data)
plt.ylim ([-10, 10])
plt.show()
import math
import cmath
import matplotlib.pyplot as plt
class KronigPenney:
        \mathbf{def} __init__(self) :
                self.a = 1.0
                                                           \# size of unit cell - late
                self.V_0 = -5.0
                                                           # height of potential bar
                self.Delta = 0.2
                                                           # width of potential barr
```

```
def solve_for_E (self, E, k):
                                                        # to solve 2x2 eigenvalue problem
             # E is the desired enegy (input)
             # k is a list of the two solutions
             q = math. sqrt (2 * E)
             kappa = math.sqrt(2 * (E - self.V_0))
             i = 1.0 j
             T11 = ( cmath.exp(i * q * (self.a - self.Delta)) / (4 * q * kappa) *
                       (\text{cmath.exp}(i * \text{kappa} * \text{self.Delta}) * (q + \text{kappa}) * * 2 -
                         \operatorname{cmath.exp}(-i * \operatorname{kappa} * \operatorname{self.Delta}) * (q - \operatorname{kappa}) * * 2
) )
             T22 = T11.conjugate()
             T12 = (-i * cmath.exp(i * q * (self.a - self.Delta)) / (2 * q * kappa)
                      (q**2 - kappa**2) * math.sin(kappa * self.Delta)
             T21 = T12.conjugate()
             # solve quadratic determinantal equation
             b = - (T11 + T22)
             c = (T11 * T22 - T12 * T21)
             k[0] = (-b + cmath.sqrt(b**2 - 4*c)) / 2.0
             k[1] = (-b - cmath.sqrt(b**2 - 4*c)) / 2.0
             for j in range (2):
                  k[j] = \operatorname{cmath.log}(k[j]) / (i * self.a)
         def compute_bands(self, dE, steps, band_file_name):
             \# dE = step \ size \ in \ E \ for \ search
             \# steps = number of steps
             E_{\text{values}} = []
             rq_values = []
             file = open(band_file_name, "w")
             E = dE
             for step in range(steps):
                  q = [0.0 + 0.0j, 0.0 + 0.0j]
                  self.solve_for_E(E, q)
                  for j in range (2):
                      rq = q[j]. real
                      if rq > 0.0 and rq < math.pi / self.a:
                           rq_values.append( rq )
                           rq_values.append( -rq )
                           E_values.append(E)
                           E_values.append(E)
                  E += dE
```

return rq\_values, E\_values

```
kp = KronigPenney()
print "_Kronig-Penney_Model"
dE = 0.01
steps = 3000
ax1 = plt.subplot(2,1,1)
\mathbf{print} "_{-}V_{-}0_{-}=", \mathrm{kp.}V_{-}0, "_{-}\mathrm{Delta}_{-}=", \mathrm{kp.}\mathrm{Delta}_{-}
rq_values1, E_values1 = kp.compute_bands(dE, steps, "band.data")
kp.V_0 = 0.0
\mathbf{print} "_{-}V_{-}0_{-}=", \mathrm{kp.}V_{-}0, "_{-}\mathrm{Delta}_{-}=", \mathrm{kp.}\mathrm{Delta}_{-}
rq_values2, E_values2 = kp.compute_bands(dE, steps, "band0.data")
plt.scatter( rq_values1, E_values1)
kp.Delta=0.5
\mathbf{print} "_V_0_=", kp.V_0, "_Delta_=", kp.Delta
rq_values1, E_values1 = kp.compute_bands(dE, steps, "band2.data")
kp.V_0 = 0.0
print "_V_0_=", kp.V_0, "_Delta_=", kp.Delta
rq2_values2, E2_values2 = kp.compute_bands(dE, steps, "band20.data")
plt.scatter(rq_values1, E_values1)
kp. Delta= 1.0
print "_V_0_=", kp.V_0, "_Delta_=", kp.Delta
rq_values1, E_values1 = kp.compute_bands(dE, steps, "band3.data")
kp.V_0 = 0.0
rq3_values2, E3_values2 = kp.compute_bands(dE, steps, "band30.data")
\#ax1 = plt.subplot(2,1,1)
plt.scatter( rq_values1, E_values1)
\#plt.scatter(rq3\_values1, E2\_values1)
\#plt.scatter(rq2\_values1, E3\_values1)
ax2 = plt.subplot(2,1,2)
plt.scatter( rq_values2, E_values2)
plt.scatter(rg2_values2, E2_values2)
plt.scatter( rq3_values2, E3_values2)
plt.show()
```