Lecture 4. Linear transformations and operators

Aleks Avdiushenko

Neapolis University Paphos

May 31, 2023



Linear Transformations

A linear transformation T is a function that maps vectors from one vector space V to another vector space W and satisfies the following properties:

- Additivity: T(u+v) = T(u) + T(v) for all $u, v \in V$
- ► Homogeneity: T(cu) = cT(u) for all $c \in \mathbb{R}$ and $u \in V$

Linear Transformations

A linear transformation T is a function that maps vectors from one vector space V to another vector space W and satisfies the following properties:

- Additivity: T(u+v) = T(u) + T(v) for all $u, v \in V$
- ► Homogeneity: $T(c\mathbf{u}) = cT(\mathbf{u})$ for all $c \in \mathbb{R}$ and $\mathbf{u} \in V$

Remark

If $T: V \to W$ is bijective then T is an isomorphism.

Linear transformations can be represented as matrices. If $T: V \to W$ is a linear transformation and A is a matrix representing T, then T(x) = Ax.

Example

A rotation in \mathbb{R}^2 by an angle θ counterclockwise is a linear transformation. The matrix representation is:

$$A = \begin{bmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{bmatrix}$$

Linear transformations can be represented as matrices. If $T: V \to W$ is a linear transformation and A is a matrix representing T, then T(x) = Ax.

Example

A rotation in \mathbb{R}^2 by an angle θ counterclockwise is a linear transformation. The matrix representation is:

$$A = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

Lemma

Let e_1, \ldots, e_n be basis of V then

$$\forall w_1, \ldots, w_n \in W \exists ! T : T(e_i) = w_i$$

Task 1

Is there any linear transformation $T: \mathbb{R}^2 \to \mathbb{R}^2$ such that

$$T\begin{pmatrix} -1\\1\\v_1\end{pmatrix} = \begin{pmatrix} 1\\0\\\end{pmatrix}, T\begin{pmatrix} 2\\1\\v_2\end{pmatrix} = \begin{pmatrix} 1\\1\\\end{pmatrix}, T\begin{pmatrix} 1\\2\\v_3\end{pmatrix} = \begin{pmatrix} 0\\1\\\end{pmatrix}$$

Task 1

Is there any linear transformation $T: \mathbb{R}^2 \to \mathbb{R}^2$ such that

$$T\begin{pmatrix} -1\\1\\v_1 \end{pmatrix} = \begin{pmatrix} 1\\0 \end{pmatrix}, T\begin{pmatrix} 2\\1\\v_2 \end{pmatrix} = \begin{pmatrix} 1\\1 \end{pmatrix}, T\begin{pmatrix} 1\\2\\v_3 \end{pmatrix} = \begin{pmatrix} 0\\1 \end{pmatrix}$$

Solution

The first two vectors are linearly independent and also $v_3 = v_1 + v_2$ It means, that $T(v_3) = T(v_1) + T(v_2)$, but this is not the case.

Task 1

Is there any linear transformation $T: \mathbb{R}^2 \to \mathbb{R}^2$ such that

$$T\begin{pmatrix} -1\\1\\v_1 \end{pmatrix} = \begin{pmatrix} 1\\0 \end{pmatrix}, T\begin{pmatrix} 2\\1\\v_2 \end{pmatrix} = \begin{pmatrix} 1\\1 \end{pmatrix}, T\begin{pmatrix} 1\\2\\v_3 \end{pmatrix} = \begin{pmatrix} 0\\1 \end{pmatrix}$$

Solution

The first two vectors are linearly independent and also $v_3 = v_1 + v_2$ It means, that $T(v_3) = T(v_1) + T(v_2)$, but this is not the case. Also you could find the matrix of T directly from the equations:

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} -1 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 2 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

Useful formalism

Let $\vec{e}_1, ..., \vec{e}_n$ be basis of V and $\vec{f}_1, ..., \vec{f}_m$ be basis of W.

$$T(\vec{e}_i) = \vec{w}_i = a_{1i}\vec{f}_1 + \dots + a_{mi}\vec{f}_m$$
$$T(\vec{e}_i) = (\vec{f}_1, \dots, \vec{f}_m) \begin{pmatrix} a_{1i} \\ \vdots \\ a_{mi} \end{pmatrix}$$

Useful formalism

Let $\vec{e}_1, ..., \vec{e}_n$ be basis of V and $\vec{f}_1, ..., \vec{f}_m$ be basis of W.

$$T(\vec{e}_i) = \vec{w}_i = a_{1i}\vec{f}_1 + \dots + a_{mi}\vec{f}_m$$
$$T(\vec{e}_i) = (\vec{f}_1, \dots, \vec{f}_m) \begin{pmatrix} a_{1i} \\ \vdots \\ a_{mi} \end{pmatrix}$$

$$(T(\vec{e}_1),\ldots,T(\vec{e}_n))=(\vec{f}_1,\ldots,\vec{f}_m)\begin{pmatrix} a_{11}&\ldots&a_{1n}\\ \vdots&\ddots&\vdots\\ a_{m1}&\ldots&a_{mn} \end{pmatrix}$$

Useful formalism

Let $\vec{e}_1, ..., \vec{e}_n$ be basis of V and $\vec{f}_1, ..., \vec{f}_m$ be basis of W.

$$T(\vec{e}_i) = \vec{w}_i = a_{1i}\vec{f}_1 + \dots + a_{mi}\vec{f}_m$$

$$T(\vec{e}_i) = (\vec{f}_1, \dots, \vec{f}_m) \begin{pmatrix} a_{1i} \\ \vdots \\ a_{mi} \end{pmatrix}$$

$$(T(\vec{e}_1), \dots, T(\vec{e}_n)) = (\vec{f}_1, \dots, \vec{f}_m) \begin{pmatrix} a_{11} & \dots & a_{1n} \\ \vdots & \ddots & \vdots \\ a_{m1} & \dots & a_{mn} \end{pmatrix}$$

$$T(\vec{e}_1, \dots, \vec{e}_n) = (\vec{f}_1, \dots, \vec{f}_m)A$$

Changing the Linear Operator Matrix with Changing Basis

Let $T:V\to W$ be a linear transformation, and let A be the matrix representing T with respect to bases \mathcal{B}_V and \mathcal{B}_W for vector spaces V and W, respectively.

Changing the Linear Operator Matrix with Changing Basis

Let $T:V\to W$ be a linear transformation, and let A be the matrix representing T with respect to bases \mathcal{B}_V and \mathcal{B}_W for vector spaces V and W, respectively.

If we change the bases for V and W to \mathscr{B}'_V and \mathscr{B}'_W , respectively, we can find the new matrix representation A' of T with respect to these new bases using the following formula:

$$A' = P^{-1}AQ$$

where P is the change of basis matrix from \mathscr{B}_W to \mathscr{B}_W' and Q is the change of basis matrix from \mathscr{B}_V to \mathscr{B}_V'

Changing the Linear Operator Matrix with Changing Basis

Let $T:V\to W$ be a linear transformation, and let A be the matrix representing T with respect to bases \mathcal{B}_V and \mathcal{B}_W for vector spaces V and W, respectively.

If we change the bases for V and W to \mathscr{B}'_V and \mathscr{B}'_W , respectively, we can find the new matrix representation A' of T with respect to these new bases using the following formula:

$$A' = P^{-1}AQ$$

where P is the change of basis matrix from \mathscr{B}_W to \mathscr{B}_W' and Q is the change of basis matrix from \mathscr{B}_V to \mathscr{B}_V'

Note

If V = W and $\mathcal{B}_V = \mathcal{B}_W$, we have P = Q and the formula simplifies to $A' = P^{-1}AP$



Remarks

$$V \rightarrow W$$

On the coordinates level: x = Qx' and y = Py'

$$A' = P^{-1}AQ$$

Remarks

$$V \rightarrow W$$

On the coordinates level: x = Qx' and y = Py'

$$A' = P^{-1}AQ$$

Remark 1

$$x' \to x = Qx' \to y = Ax = AQx' \to y' = P^{-1}y = P^{-1}AQx' = A'x'$$

Remarks

 $V \rightarrow W$

On the coordinates level: x = Qx' and y = Py'

$$A' = P^{-1}AQ$$

Remark 1

$$x' \rightarrow x = Qx' \rightarrow y = Ax = AQx' \rightarrow y' = P^{-1}y = P^{-1}AQx' = A'x'$$

Remark 2

In fact, by changing the bases, one can bring the matrix of any linear transformation to a «block-identity» one:

	k	
k A' =	E	0
A =	0	0



Kernel and Image

Let $T: V \to W$ be a linear transformation. The *kernel* and *image* of T are defined as follows:

Kernel

The kernel (or null space) of T is the set of all vectors $v \in V$ such that $T(v) = \vec{0}_W$, where $\vec{0}_W$ is the zero vector in W:

$$\ker(T) = \{ v \in V : T(v) = \vec{0}_W \}$$

Image

The image (or range) of T is the set of all vectors in W that can be obtained by applying T to some vector in V:

$$im(T) = \{T(v) : v \in V\}$$



Dimension Theorem (rank-nullity)

The dimension theorem (or rank-nullity theorem) states that for a linear transformation $T: V \to W$:

$$\dim(\ker(T)) + \dim(\operatorname{im}(T)) = \dim(V)$$

Remarks

For $T: \mathbb{R}^n \to \mathbb{R}^m$ its matrix can be represented as set of columns:

$$A = [A_1 \mid \dots \mid A_n]$$

$$\ker(T) = \{x \in \mathbb{R}^n \mid Ax = 0\}$$

$$\operatorname{im}(T) = \{x_1 A_1 + \dots + x_n A_n \mid x_i \in \mathbb{R}\} = \operatorname{span}(A_1, \dots, A_n)$$

- 1. T is onto \Leftrightarrow im(T) = W
- 2. T is one-to-one $\Leftrightarrow \ker(T) = 0$



Theorem 2

(in some sense the inverse of the dimension theorem) Let $K \subseteq V, I \subseteq W$ and

$$\dim(K) + \dim(I) = \dim(V)$$

Than exists linear transformation $T: V \rightarrow W$ such that

$$\ker(T) = K$$
, $\operatorname{im}(T) = I$

Proof scheme for \mathbb{R} case

Theorem 2

(in some sense the inverse of the dimension theorem) Let $K \subseteq V, I \subseteq W$ and

$$\dim(K) + \dim(I) = \dim(V)$$

Than exists linear transformation $T: V \rightarrow W$ such that

$$\ker(T) = K$$
, $\operatorname{im}(T) = I$

Proof scheme for \mathbb{R} case

Suppose $V=\mathbb{R}^n$ and $W=\mathbb{R}^m$. We will find matrix A with linearly independent rows, that $K=\{x\in\mathbb{R}^n\mid Ax=0\}$ and vectors $w_1,\ldots,w_k\in\mathbb{R}^m$ that $I=\langle w_1,\ldots,w_k\rangle$

Theorem 2

(in some sense the inverse of the dimension theorem) Let $K \subseteq V, I \subseteq W$ and

$$\dim(K) + \dim(I) = \dim(V)$$

Than exists linear transformation $T: V \to W$ such that

$$\ker(T) = K$$
, $\operatorname{im}(T) = I$

Proof scheme for ℝ case

Suppose $V = \mathbb{R}^n$ and $W = \mathbb{R}^m$. We will find matrix A with linearly independent rows, that $K = \{x \in \mathbb{R}^n \mid Ax = 0\}$ and vectors $w_1, \dots, w_k \in \mathbb{R}^m$ that $I = \langle w_1, \dots, w_k \rangle$

Consider matrix $B = [w_1 \mid \cdots \mid w_k]$.

$$\#(\text{of rows in } A) = \text{rk}(A) = n - \text{dim}(K) = k$$

Than shapes of matrices A, B allow us to consider the product $BA \in M_{mn}(\mathbb{R})$, which is the matrix representing T (for $BCA \in M_{mn}(\mathbb{R}), \forall C \in M_k(\mathbb{R}), \det C \neq 0$ it is also true).

Linear Operators

Routhly speaking up to now we used only Guass algorithm for various problems.

Linear Operators

Routhly speaking up to now we used only Guass algorithm for various problems.

A *linear operator* is a linear transformation from a vector space V to itself, i.e., $T: V \rightarrow V$

Matrix Representation

Let $\mathscr{B} = \{v_1, v_2, ..., v_n\}$ be a basis of V. The matrix representation of a linear operator $T: V \to V$ with respect to \mathscr{B} is an $n \times n$ matrix $A = [a_{ij}]$, where:

$$T(v_j) = \sum_{i=1}^n a_{ij} v_i$$

So for basis vectors $v_1, ..., v_n \in V$

$$T(v_1,...,v_n) = (v_1,...,v_n)A$$



$$A' = P^{-1}AP$$

$$A' = P^{-1}AP$$

$$A' = P^{-1}AP$$

1.
$$tr(A') = tr(P^{-1}AP) = tr(APP^{-1}) = tr(A) = tr(T)$$

$$A' = P^{-1}AP$$

1.
$$tr(A') = tr(P^{-1}AP) = tr(APP^{-1}) = tr(A) = tr(T)$$

2.
$$det(A') = det(P^{-1}AP) = det(P^{-1})det(A)det(P) = det(A) = det(T)$$

$$A' = P^{-1}AP$$

- 1. $tr(A') = tr(P^{-1}AP) = tr(APP^{-1}) = tr(A) = tr(T)$
- 2. $det(A') = det(P^{-1}AP) = det(P^{-1})det(A)det(P) = det(A) = det(T)$
- 3. Characteristic Polynomial: $det(\lambda E A') = det(\lambda E A) = det(\lambda id T)$

$$A' = P^{-1}AP$$

- 1. $tr(A') = tr(P^{-1}AP) = tr(APP^{-1}) = tr(A) = tr(T)$
- 2. $det(A') = det(P^{-1}AP) = det(P^{-1})det(A)det(P) = det(A) = det(T)$
- 3. Characteristic Polynomial: $det(\lambda E A') = det(\lambda E A) = det(\lambda id T)$
- 4. Minimal annihilating polynomial

The main question

How to choose a basis in which the matrix of a linear operator will have the simplest form?

The main question

How to choose a basis in which the matrix of a linear operator will have the simplest form?

We can't find such basis without Eigenvalues and Eigenvectors.

The main question

How to choose a basis in which the matrix of a linear operator will have the simplest form?

We can't find such basis without Eigenvalues and Eigenvectors.

Definition

An eigenvalue λ of a linear operator $T: V \to V$ is a scalar such that there exists a non-zero vector $v \in V$, called an eigenvector of T, satisfying:

$$Tv = \lambda v$$

The main question

How to choose a basis in which the matrix of a linear operator will have the simplest form?

We can't find such basis without Eigenvalues and Eigenvectors.

Definition

An eigenvalue λ of a linear operator $T: V \to V$ is a scalar such that there exists a non-zero vector $v \in V$, called an eigenvector of T, satisfying:

$$Tv = \lambda v$$

 $V_{\lambda} = \{ v \in V \mid Tv = \lambda v \} = \ker(T - \lambda Id)$

The main question

How to choose a basis in which the matrix of a linear operator will have the simplest form?

We can't find such basis without Eigenvalues and Eigenvectors.

Definition

An eigenvalue λ of a linear operator $T: V \to V$ is a scalar such that there exists a non-zero vector $v \in V$, called an eigenvector of T, satisfying:

$$Tv = \lambda v$$

- $V_{\lambda} = \{ v \in V \mid Tv = \lambda v \} = \ker(T \lambda Id)$
- λ eigenvalue ⇔ ∃x ≠ 0 : Ax = λx ⇔ ∃x ≠ 0 : (A λE)x = 0⇔ (A - λE) is singular (irreversible) ⇔ det(A - λE) = 0

The main question

How to choose a basis in which the matrix of a linear operator will have the simplest form?

We can't find such basis without Eigenvalues and Eigenvectors.

Definition

An eigenvalue λ of a linear operator $T: V \to V$ is a scalar such that there exists a non-zero vector $v \in V$, called an eigenvector of T, satisfying:

$$Tv = \lambda v$$

- $V_{\lambda} = \{ v \in V \mid Tv = \lambda v \} = \ker(T \lambda Id)$
- ▶ λ eigenvalue $\Leftrightarrow \exists x \neq 0 : Ax = \lambda x \Leftrightarrow \exists x \neq 0 : (A \lambda E)x = 0$ $\Leftrightarrow (A - \lambda E)$ is singular (irreversible) \Leftrightarrow det $(A - \lambda E) = 0$

Eigenvalues of $A = \operatorname{Spec} A = \operatorname{roots}$ of χ_A



A rotation in \mathbb{R}^2 by an angle θ counterclockwise is a linear transformation. The matrix representation is:

$$A = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

A rotation in \mathbb{R}^2 by an angle θ counterclockwise is a linear transformation. The matrix representation is:

$$A = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

Also we know, that a complex number $z = r(\cos\theta + i\sin\theta)$ corresponds to a rotation by θ and dilation by r

A rotation in \mathbb{R}^2 by an angle θ counterclockwise is a linear transformation. The matrix representation is:

$$A = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

- Also we know, that a complex number $z = r(\cos \theta + i \sin \theta)$ corresponds to a rotation by θ and dilation by r
- So we can represent z by the matrix

$$Z = \begin{bmatrix} r\cos\theta & -r\sin\theta \\ r\sin\theta & r\cos\theta \end{bmatrix}$$

A rotation in \mathbb{R}^2 by an angle θ counterclockwise is a linear transformation. The matrix representation is:

$$A = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

- Also we know, that a complex number $z = r(\cos \theta + i \sin \theta)$ corresponds to a rotation by θ and dilation by r
- ► So we can represent z by the matrix

$$Z = \begin{bmatrix} r\cos\theta & -r\sin\theta \\ r\sin\theta & r\cos\theta \end{bmatrix}$$

► Having the alternative construction of complex numbers, which allow us to get square root from negatives numbers



Quaternions extend the concept of rotation from 2D to 3D and are used in 3D geometry, physics, and computer graphics

Quaternions extend the concept of rotation from 2D to 3D and are used in 3D geometry, physics, and computer graphics

 Quaternions, discovered by William Rowan Hamilton, are a number system that extends the complex numbers

Quaternions extend the concept of rotation from 2D to 3D and are used in 3D geometry, physics, and computer graphics

- Quaternions, discovered by William Rowan Hamilton, are a number system that extends the complex numbers
- ► They can be expressed in the form q = a + bi + cj + dk where a, b, c, d are real numbers, and i, j, k are the fundamental quaternion units with the relations: $i^2 = j^2 = k^2 = ijk = -1$

Quaternions extend the concept of rotation from 2D to 3D and are used in 3D geometry, physics, and computer graphics

- Quaternions, discovered by William Rowan Hamilton, are a number system that extends the complex numbers
- ▶ They can be expressed in the form q = a + bi + cj + dk where a, b, c, d are real numbers, and i, j, k are the fundamental quaternion units with the relations: $i^2 = j^2 = k^2 = ijk = -1$
- ▶ In 3D space, b, c, and d can be interpreted as coefficients for the i, j, and k directions, respectively, while a is the real part

Quaternions extend the concept of rotation from 2D to 3D and are used in 3D geometry, physics, and computer graphics

- Quaternions, discovered by William Rowan Hamilton, are a number system that extends the complex numbers
- ▶ They can be expressed in the form q = a + bi + cj + dk where a, b, c, d are real numbers, and i, j, k are the fundamental quaternion units with the relations: $i^2 = j^2 = k^2 = ijk = -1$
- ▶ In 3D space, b, c, and d can be interpreted as coefficients for the i, j, and k directions, respectively, while a is the real part

Example

Quaternion multiplication is non-commutative. For instance, if $q_1 = a_1 + b_1 i + c_1 j + d_1 k$ and $q_2 = a_2 + b_2 i + c_2 j + d_2 k$, then

$$q_1q_2 \neq q_2q_1$$



A quaternion q = a + bi + cj + dk can be represented as the following 4×4 real matrix:

$$\begin{bmatrix} a & -b & -c & -d \\ b & a & -d & c \\ c & d & a & -b \\ d & -c & b & a \end{bmatrix}$$

A quaternion q = a + bi + cj + dk can be represented as the following 4×4 real matrix:

$$\begin{bmatrix} a & -b & -c & -d \\ b & a & -d & c \\ c & d & a & -b \\ d & -c & b & a \end{bmatrix}$$

Similarly, it can be represented as a 2x2 complex matrix:

$$\begin{bmatrix} a+bi & c+di \\ -c+di & a-bi \end{bmatrix}$$

A quaternion q = a + bi + cj + dk can be represented as the following 4×4 real matrix:

$$\begin{bmatrix} a & -b & -c & -d \\ b & a & -d & c \\ c & d & a & -b \\ d & -c & b & a \end{bmatrix}$$

Similarly, it can be represented as a 2x2 complex matrix:

$$\begin{bmatrix} a+bi & c+di \\ -c+di & a-bi \end{bmatrix}$$

These representations allow us to apply linear algebra techniques to quaternions.

A quaternion q = a + bi + cj + dk can be represented as the following 4×4 real matrix:

$$\begin{bmatrix} a & -b & -c & -d \\ b & a & -d & c \\ c & d & a & -b \\ d & -c & b & a \end{bmatrix}$$

Similarly, it can be represented as a 2x2 complex matrix:

$$\begin{bmatrix} a+bi & c+di \\ -c+di & a-bi \end{bmatrix}$$

These representations allow us to apply linear algebra techniques to quaternions.

For example, quaternion multiplication corresponds to matrix multiplication in these representations.



Octonions, also known as Cayley numbers, are a non-associative extension of Quaternions. They form an 8-dimensional normed division algebra over the real numbers.

Octonions, also known as Cayley numbers, are a non-associative extension of Quaternions. They form an 8-dimensional normed division algebra over the real numbers.

Let
$$O = a_0 + a_1 e_1 + a_2 e_2 + \dots + a_7 e_7$$

Where:

- ▶ $a_0, a_1, ..., a_7 \in \mathbb{R}$
- $e_1, e_2, ..., e_7$ are the octonion units

Octonions, also known as Cayley numbers, are a non-associative extension of Quaternions. They form an 8-dimensional normed division algebra over the real numbers.

Let
$$O = a_0 + a_1 e_1 + a_2 e_2 + \dots + a_7 e_7$$

Where:

- ▶ $a_0, a_1, ..., a_7 \in \mathbb{R}$
- $e_1, e_2, ..., e_7$ are the octonion units

Multiplication of Octonions is not associative, i.e., $(ab)c \neq a(bc)$ for some Octonions a, b, c, so they can't have a matrix representation

Octonions, also known as Cayley numbers, are a non-associative extension of Quaternions. They form an 8-dimensional normed division algebra over the real numbers.

Let
$$O = a_0 + a_1 e_1 + a_2 e_2 + \dots + a_7 e_7$$

Where:

- ▶ $a_0, a_1, ..., a_7 \in \mathbb{R}$
- $ightharpoonup e_1, e_2, ..., e_7$ are the octonion units

Multiplication of Octonions is not associative, i.e., $(ab)c \neq a(bc)$ for some Octonions a,b,c, so they can't have a matrix representation Octonions find applications in different areas of theoretical physics, like string theory and quantum mechanics

Octonions, also known as Cayley numbers, are a non-associative extension of Quaternions. They form an 8-dimensional normed division algebra over the real numbers.

Let
$$O = a_0 + a_1 e_1 + a_2 e_2 + \dots + a_7 e_7$$

Where:

- ▶ $a_0, a_1, ..., a_7 \in \mathbb{R}$
- $ightharpoonup e_1, e_2, ..., e_7$ are the octonion units

Multiplication of Octonions is not associative, i.e., $(ab)c \neq a(bc)$ for some Octonions a, b, c, so they can't have a matrix representation Octonions find applications in different areas of theoretical physics, like string theory and quantum mechanics

Problem

Why we have something interesting in the 2, 4 and 8 dimensional spaces?

The dimensions 2, 4, and 8 are special due to the properties of the normed division algebras that exist in these dimensions. These are the real numbers (dimension 1), complex numbers (dimension 2), quaternions (dimension 4), and octonions (dimension 8).

The dimensions 2, 4, and 8 are special due to the properties of the normed division algebras that exist in these dimensions. These are the real numbers (dimension 1), complex numbers (dimension 2), quaternions (dimension 4), and octonions (dimension 8).

There are several reasons these dimensions are special:

The dimensions 2, 4, and 8 are special due to the properties of the normed division algebras that exist in these dimensions. These are the real numbers (dimension 1), complex numbers (dimension 2), quaternions (dimension 4), and octonions (dimension 8).

There are several reasons these dimensions are special:

Frobenius Theorem: According to the Frobenius theorem, these are the
only dimensions in which normed division algebras over the reals can
exist. In other words, there are no other dimensions where we can define a
multiplication operation that behaves like the one we're used to from real
and complex numbers.

The dimensions 2, 4, and 8 are special due to the properties of the normed division algebras that exist in these dimensions. These are the real numbers (dimension 1), complex numbers (dimension 2), quaternions (dimension 4), and octonions (dimension 8).

There are several reasons these dimensions are special:

- Frobenius Theorem: According to the Frobenius theorem, these are the
 only dimensions in which normed division algebras over the reals can
 exist. In other words, there are no other dimensions where we can define a
 multiplication operation that behaves like the one we're used to from real
 and complex numbers.
- 2. **Algebraic Structure**: As we increase the dimension, the algebraic structure gets progressively «weaker». Real numbers form a field, as do complex numbers. Quaternions, however, are non-commutative, meaning that the order in which you multiply them matters. Octonions go a step further and are non-associative, meaning that even the property (ab)c = a(bc) doesn't hold.

So, while you might initially expect every dimension to have its own «numbers», algebraic properties restrict the possibilities, leaving 1, 2, 4, and 8 as the special dimensions.

▶ Dimension 16 corresponds to the «Sedenions», which are an extension of the Octonions. However, the Sedenions are not a division algebra because they contain zero divisors, i.e., non-zero elements whose product is zero.

- ▶ Dimension 16 corresponds to the «Sedenions», which are an extension of the Octonions. However, the Sedenions are not a division algebra because they contain zero divisors, i.e., non-zero elements whose product is zero.
- Just like how the Quaternions lose commutativity and the Octonions lose associativity, the Sedenions lose the property of being an «alternative algebra», which means they no longer satisfy the property (xx)y = x(xy) for all x, y in the algebra. This leads to the existence of zero divisors, making them much less useful for many mathematical and physical applications compared to the lower-dimensional algebras.

- ▶ Dimension 16 corresponds to the «Sedenions», which are an extension of the Octonions. However, the Sedenions are not a division algebra because they contain zero divisors, i.e., non-zero elements whose product is zero.
- Just like how the Quaternions lose commutativity and the Octonions lose associativity, the Sedenions lose the property of being an «alternative algebra», which means they no longer satisfy the property (xx)y = x(xy) for all x, y in the algebra. This leads to the existence of zero divisors, making them much less useful for many mathematical and physical applications compared to the lower-dimensional algebras.
- ▶ The pattern of losing properties as we double dimensions stops here; there are no interesting algebras in dimension 32. Sedenions are generally seen as the end of the line for the Cayley-Dickson construction, the process by which each algebra is formed by doubling the dimension of the previous one.

- ▶ Dimension 16 corresponds to the «Sedenions», which are an extension of the Octonions. However, the Sedenions are not a division algebra because they contain zero divisors, i.e., non-zero elements whose product is zero.
- Just like how the Quaternions lose commutativity and the Octonions lose associativity, the Sedenions lose the property of being an «alternative algebra», which means they no longer satisfy the property (xx)y = x(xy) for all x, y in the algebra. This leads to the existence of zero divisors, making them much less useful for many mathematical and physical applications compared to the lower-dimensional algebras.
- ▶ The pattern of losing properties as we double dimensions stops here; there are no interesting algebras in dimension 32. Sedenions are generally seen as the end of the line for the Cayley-Dickson construction, the process by which each algebra is formed by doubling the dimension of the previous one.
- ▶ So while there are «numbers» in dimension 16, they lose many of the nice properties we have in lower dimensions. This is why we often focus more on dimensions 1, 2, 4, and 8.

Diagonalizable Operators

A linear operator $T: V \to V$ is called *diagonalizable* if there exists an invertible matrix P and a diagonal matrix D such that:

$$P^{-1}AP = D$$

where A is the matrix representation of T with respect to some basis of V.

Diagonalizable Operators

A linear operator $T: V \to V$ is called *diagonalizable* if there exists an invertible matrix P and a diagonal matrix D such that:

$$P^{-1}AP = D$$

where A is the matrix representation of T with respect to some basis of V.

For a diagonalizable operator T, the diagonal entries of the matrix D are the eigenvalues of T, and the columns of the matrix P are the corresponding eigenvectors.

Diagonalizable Operators

A linear operator $T: V \to V$ is called *diagonalizable* if there exists an invertible matrix P and a diagonal matrix D such that:

$$P^{-1}AP = D$$

where A is the matrix representation of T with respect to some basis of V.

For a diagonalizable operator T, the diagonal entries of the matrix D are the eigenvalues of T, and the columns of the matrix P are the corresponding eigenvectors.

Conditions for Diagonalizability

A linear operator T is diagonalizable iff the following conditions hold:

- ▶ The sum of the dimensions of the eigenspaces is equal to *n*
- ► There exist *n* linearly independent eigenvectors of *T*



Diagonalization Algorithm

To diagonalize a linear operator T with matrix representation A, follow these steps:

- 1. Find the eigenvalues $\lambda_1, \lambda_2, ..., \lambda_r$ of A
- 2. Find a basis for each eigenspace V_{λ_i}
- 3. Form the matrix P with the eigenvectors as columns
- 4. Compute $P^{-1}AP$, which should be the diagonal matrix D

Condition, sufficient for diagonalization

Let
$$T: \mathbb{R}^n \to \mathbb{R}^n$$
, $x \mapsto Ax$.

- 1. Suppose $p(T) = 0 \Leftrightarrow p(A) = 0$ for some polynomial p
- 2. If $p(x) = (x \lambda_1)...(x \lambda_k)$ than T is diagonalizable
- 3. Common case is to use χ_A as p

Task 2

Let $E \in M_n(\mathbb{R})$ — identity matrix and $X \in M_n(\mathbb{R})$. Solve equation

$$X^2 = E$$

Task 2

Let $E \in M_n(\mathbb{R})$ — identity matrix and $X \in M_n(\mathbb{R})$. Solve equation

$$X^2 = E$$

Solution

Than $p(t) = t^2 - 1 = (t - 1)(t + 1)$ is annihilating polynomial for X It means, that X is diagonalizable (there is no multiple roots) and

$$X = PDP^{-1}$$

and D has only ± 1 on the diagonal. Vice versa, for any $X = PDP^{-1}$ with such matrix D we have $X^2 = E$.

Task 2

Let $E \in M_n(\mathbb{R})$ — identity matrix and $X \in M_n(\mathbb{R})$. Solve equation

$$X^2 = E$$

Solution

Than $p(t) = t^2 - 1 = (t - 1)(t + 1)$ is annihilating polynomial for X It means, that X is diagonalizable (there is no multiple roots) and

$$X = PDP^{-1}$$

and D has only ± 1 on the diagonal. Vice versa, for any $X = PDP^{-1}$ with such matrix D we have $X^2 = E$.

Remark

For recurrent formula $x_n = Ax_{n-1} = \cdots = A^{n-1}x_1$ If $A = PDP^{-1}$, then $A^n = PD^nP^{-1}$.



Projectors

Let V be vector space and $U, W \subseteq V$:

- 1. $U \cap W = 0$
- 2. $\langle U, W \rangle = V \Leftrightarrow \forall v \in V \ v = u + w$

Projectors

Let V be vector space and $U, W \subseteq V$:

- 1. $U \cap W = 0$
- 2. $\langle U, W \rangle = V \Leftrightarrow \forall v \in V \ v = u + w$

Projection $P: V \rightarrow V$ "on U along W":

$$v = u + w \mapsto u$$

Projector Matrix: Geometric Approach

- A projector matrix P is a square matrix that satisfies the property: $P^2 = P$
- Geometrically, a projector matrix represents a linear transformation that projects vectors onto a subspace U of the original vector space, U = Im P, W = ker P
- ► For a given subspace U, the projector matrix P maps each vector v in the vector space to its orthogonal projection Pv on U

Projector Matrix: Geometric Approach

- A projector matrix P is a square matrix that satisfies the property: $P^2 = P$
- Geometrically, a projector matrix represents a linear transformation that projects vectors onto a subspace U of the original vector space, U = Im P, W = ker P
- For a given subspace U, the projector matrix P maps each vector v in the vector space to its orthogonal projection Pv on U
- Pv is the closest point to v in the subspace U

Projector Matrix: Geometric Approach

- A projector matrix P is a square matrix that satisfies the property: $P^2 = P$
- Geometrically, a projector matrix represents a linear transformation that projects vectors onto a subspace U of the original vector space, U = Im P, W = ker P
- For a given subspace U, the projector matrix P maps each vector v in the vector space to its orthogonal projection Pv on U
- ightharpoonup Pv is the closest point to v in the subspace U
- ▶ If *U* is spanned by an orthonormal basis $\{u_1, u_2, ..., u_k\}$, then the projector matrix is given by:

$$P = u_1 u_1^T + u_2 u_2^T + \dots + u_k u_k^T$$



Projector Matrix: Algebraic Approach

- Algebraically, a projector matrix can be derived using the Gram-Schmidt orthogonalization process or by solving a system of linear equations
- Suppose we have a basis $\{u_1, u_2, ..., u_k\}$ for subspace U. We can form a matrix B whose columns are the basis vectors, i.e., $B = [u_1 \ u_2 \cdots u_k]$
- ▶ If B has linearly independent columns, the projector matrix onto the column space of B (which is U) can be computed as:

$$P = B(B^T B)^{-1} B^T$$

► If the basis vectors are already orthogonal, the projector matrix simplifies to:

$$P = B(B^TB)^{-1}B^T = BB^T$$



Question

What is the diagonal form of projector $P = C^{-1}DC$?

Question

What is the diagonal form of projector $P = C^{-1}DC$?

Example

 $U, W \subset \mathbb{R}^n$, $U \cap W = 0$, dim $U + \dim W = n$ How to express the projector on U?

Suppose we have a basis $\{u_1, u_2, ..., u_k\}$ for subspace U. We can form a matrix B whose columns are the basis vectors, i.e.,

$$B = [u_1 \ u_2 \cdots u_k]$$

And let $W = \{y \in \mathbb{R}^n \mid Ay = 0\}$, then

$$P_U = B(AB)^{-1}A$$

Remark

In computation use this formula carefully: $Pv = B((AB)^{-1}(Av))$



Jordan Normal Form

- The Jordan normal form is a canonical representation of a linear operator or matrix
- It's a decomposition of a matrix into a block diagonal matrix with Jordan blocks
- ► A Jordan block is a square matrix with the following structure:

$$J_k(\lambda) = \begin{bmatrix} \lambda & 1 & & & \\ & \lambda & \ddots & & \\ & & \ddots & 1 \\ & & & \lambda \end{bmatrix}$$

- ▶ Not all matrices have a Jordan normal form over the real numbers, but they do have one over the complex numbers
- For a given matrix A, there exists an invertible matrix P such that:

$$A = PJP^{-1}$$

where J is the Jordan normal form of A