

# Supplementary Material (SM)

Throughout this supplement, we set use dot notation for time derivatives so that  $\dot{f}(x, t) = \frac{\partial}{\partial t} f(x, t)$  and set  $\Delta = \sum_{i=1}^d \frac{\partial^2}{\partial x_i^2}$  the Laplace operator on  $\mathbb{R}^d$ , except in §10.1.3 where  $\Delta$  represents a random variable.

## 1 Sufficient conditions for finite mean, variance and total abundance in the deterministic case

Recall that  $m(\nu, x)$  is shorthand for  $m((K\nu)(x, t), x)$ . That is,  $m : [0, +\infty) \times \mathbb{R} \rightarrow \mathbb{R}$ . Following our assumptions of the main text, we have that  $m(h, x)$  is differentiable with respect to both  $x$  and  $h$  and there exists  $r \in \mathbb{R}$  such that  $m(h, x) \leq r$  across all  $h \geq 0$  and  $x \in \mathbb{R}$ . As in the main text, we also assume the initial condition  $u(x)$  is continuous and integrable in  $x$  and satisfies

$$0 \leq \int_{\mathbb{R}} (|x| + x^2) u(x) dx < +\infty \quad (1)$$

and consider the Cauchy problem

$$\begin{cases} \dot{\nu}(x, t) = m(\nu, x)\nu(x, t) + \frac{\mu}{2}\Delta\nu(x, t) & t > 0 \\ \nu(x, 0) = u(x) & t = 0. \end{cases} \quad (2)$$

According to Theorem 2.5.6 of Zheng (2004), if the operator  $F$  defined by  $\nu(x, t) \rightarrow m(\nu, x)\nu(x, t)$  is locally Lipschitz, corresponding to equation (8) of the main text, then for some maximal  $T > 0$  problem (2) admits a unique local classical solution  $\nu(x, t)$  for  $t \in [0, T)$ . Furthermore, either  $T = +\infty$  or  $T < +\infty$  and  $\lim_{t \uparrow T} \int_{\mathbb{R}} |\nu(x, t)| dx = +\infty$ . In this section we show that our assumption  $m(h, x) \leq r$  for all  $h \geq 0$  and  $x \in \mathbb{R}$  implies  $T = +\infty$ . Replacing  $m$  with its upper bound  $r \in \mathbb{R}$ , equation (2) reduces to a simple parabolic equation that can be solved using elementary techniques (Farlow 1993). In particular, when  $m(\nu, x) \equiv 0$  denote the solution to (2) by  $\nu_0(x, t)$ . Then, denoting

$$\Phi(x, t) = \frac{\exp(-x^2/2\mu t)}{\sqrt{2\pi\mu t}}, \quad (3)$$

we have

$$\nu_0(x, t) = \int_{\mathbb{R}} \Phi(x - y, t) \nu(y, 0) dy. \quad (4)$$

In the more general case, when  $m(\nu, x) \equiv r \in \mathbb{R}$ , equation (2) has the solution  $\nu_r(x, t) = e^{rt} \nu_0(x, t)$ . Hence,  $\nu_r(x, t) \geq 0$  for all  $x \in \mathbb{R}$  and  $\int_{\mathbb{R}} \nu_r(x, t) dx = e^{rt} N(0) < +\infty$  for all  $t \geq 0$ . Furthermore, denoting

$$\begin{aligned} N_r(t) &= \int_{\mathbb{R}} \nu_r(x, t) dx, \\ p_r(x, t) &= \nu_r(x, t) / N_r(t), \\ \bar{x}_r(t) &= \int_{\mathbb{R}} x p_r(x, t) dx, \\ \sigma_r^2(t) &= \int_{\mathbb{R}} (x - \bar{x}_r(t))^2 p_r(x, t) dx, \end{aligned} \quad (5)$$

we have

$$\bar{x}_r(t) = \int_{\mathbb{R}} x \int_{\mathbb{R}} \Phi(x - y, t) p(y, 0) dy dx = \int_{\mathbb{R}} y p(y, 0) dy = \bar{x}(0), \quad (6)$$

$$\sigma_r^2(t) = \int_{\mathbb{R}} (x - \bar{x}_r(t))^2 \int_{\mathbb{R}} \Phi(x - y, t) p(y, 0) dy dx = \int_{\mathbb{R}} ((y - \bar{x}(0))^2 + \mu t) p(y, 0) dy = \sigma^2(0) + \mu t. \quad (7)$$

Hence,  $|\bar{x}_r(t)|, \sigma_r^2(t) < +\infty$  for all  $t \geq 0$ . For the sake of contradiction, suppose there exists  $x \in \mathbb{R}$  and  $t \geq 0$  such that  $\nu(x, t) > \nu_r(x, t)$ . Then

$$\nu(x, t) - \nu(x, 0) = \int_0^t m(\nu, x) \nu(x, s) ds + \frac{\mu}{2} \Delta \nu(x, s) ds > \int_0^t r \nu_r(x, s) + \frac{\mu}{2} \Delta \nu_r(x, s) ds = \nu_r(x, t) - \nu(x, 0) \quad (8)$$

which implies there exists  $h \geq 0$  and  $x \in \mathbb{R}$  such that  $m(h, x) > r$ . But this contradicts our assumption  $m(h, x) \leq r$  for all  $h \geq 0$  and  $x \in \mathbb{R}$ . So we have  $\nu(x, t) \leq \nu_r(x, t)$  for each  $x \in \mathbb{R}$  and  $t \geq 0$ . This implies that  $N(t) = \int_{\mathbb{R}} \nu(x, t) dx < +\infty$ ,

$$0 \leq \int_{\mathbb{R}} x^2 \nu(x, t) dx \leq \int_{\mathbb{R}} x^2 \nu_r(x, t) dx < +\infty \quad (9)$$

and in particular

$$0 \leq \sigma^2(t) + \bar{x}^2(t) = \frac{1}{N(t)} \int_{\mathbb{R}} x^2 \nu(x, t) dx < +\infty \quad (10)$$

for each  $t \geq 0$ .

## 2 Equilibrium moments for a deterministic population experiencing logistic growth and stabilizing selection

Here we set out to show, given the initial conditions  $\nu(\cdot, 0) \in L^1(\mathbb{R})$  such that  $\bar{x}(0) \in \mathbb{R}$ ,  $\sigma^2(0) \in [0, +\infty)$  and  $N(0) \in (0, +\infty)$ , and given the growth rate

$$m(\nu, x) = r - \frac{a}{2}(\theta - x)^2 - c \int_{\mathbb{R}} \nu(y, t) dy \quad (11)$$

such that  $\theta \in \mathbb{R}$ ,  $a, c, \mu > 0$  and  $r > \frac{1}{2}\sqrt{\mu a}$ , we have the following stable equilibrium  $N_{\infty} = \lim_{t \rightarrow \infty} N(t) = \frac{1}{c}(r - \frac{1}{2}\sqrt{a\mu})$ ,  $\bar{x}_{\infty} = \lim_{t \rightarrow \infty} \bar{x}(t) = \theta$  and  $\sigma_{\infty}^2 = \lim_{t \rightarrow \infty} \sigma^2(t) = \sqrt{\frac{\mu}{a}}$ .

The mean fitness can be calculated as

$$\bar{m}(t) = r - \frac{a}{2}((\theta - \bar{x}(t))^2 + \sigma^2(t)) - cN(t). \quad (12)$$

Recalling the ODE derived for  $N(t)$ ,

$$\frac{d}{dt}N(t) = \bar{m}(t)N(t) = (r - \frac{a}{2}((\theta - \bar{x}(t))^2 + \sigma^2(t)) - cN(t))N(t), \quad (13)$$

solving for equilibrium total abundance  $\hat{N}$  amounts to setting  $\frac{d}{dt}N(t) = 0$  and solving for  $N(t)$ . Ignoring the equilibrium  $N(t) = 0$ , this reduces to solving  $\bar{m}(t) = 0$  for  $N(t)$ , which returns

$$\hat{N} = \frac{1}{c} \left( r - \frac{a}{2}((\theta - \hat{x})^2 - \hat{\sigma}^2) \right). \quad (14)$$

Unfortunately, deriving ODE for  $\bar{x}(t)$  and  $\sigma^2(t)$  leads to expressions involving higher moments and finding ODE for these higher moments will lead to expressions involving yet even higher moments. To avoid this infinite regression, we find the equilibrium abundance density  $\hat{\nu}(x)$  by solving  $\frac{\partial}{\partial t}\nu(x, t) = 0$  for  $\nu(x, t)$ . This implies the following ordinary differential equation

$$\frac{d^2}{dx^2}\hat{\nu}(x) = \left( \frac{2c}{\mu}\hat{N} + \frac{a}{\mu}(\theta - x)^2 - \frac{2r}{\mu} \right) \hat{\nu}(x) \quad (15)$$

which has the solution

$$\hat{\nu}(x) = \frac{\hat{N}}{\sqrt{2\pi}} \left( \frac{a}{\mu} \right)^{\frac{1}{4}} \exp \left( -\sqrt{\frac{a}{\mu}} \frac{(\theta - x)^2}{2} \right). \quad (16)$$

From this expression we infer  $\hat{x} = \theta$  and  $\hat{\sigma}^2 = \sqrt{\frac{\mu}{a}}$ . Hence  $\hat{N} = \frac{1}{c} (r - \frac{1}{2}\sqrt{a\mu})$ . To show that  $N_\infty = \hat{N}$ ,  $\bar{x}_\infty = \hat{x}$  and  $\sigma_\infty^2 = \hat{\sigma}^2$ , we linearize ODE for  $N(t)$ ,  $\bar{x}(t)$  and  $\sigma^2(t)$  using the equilibrium  $\hat{\nu}(x)$ . In particular, since  $\hat{\nu}(x)$  is Gaussian, we do not run into the same issue with higher moments as above. Assuming the equilibrium  $\hat{\nu}(x)$  does not change the ODE for  $N(t)$ , but the ODE for  $\bar{x}(t)$  and  $\sigma^2(t)$  can now be expressed as

$$\frac{d}{dt} \bar{x}(t) = \sigma^2(t) \left( \frac{\partial \bar{m}(t)}{\partial \bar{x}(t)} - \frac{\partial \bar{m}(t)}{\partial \bar{x}(t)} \right) = a\sigma^2(t)(\theta - \bar{x}(t)), \quad (17)$$

$$\frac{d}{dt} \sigma^2(t) = 2\sigma^4(t) \left( \frac{\partial \bar{m}(t)}{\partial \sigma^2(t)} - \frac{\partial \bar{m}(t)}{\partial \sigma^2(t)} \right) + \mu = \mu - a\sigma^4(t), \quad (18)$$

where  $\hat{\rho}(x) = \hat{\nu}(x)/\hat{N}$ . These expressions confirm our findings that  $\hat{x} = \theta$  and  $\hat{\sigma}^2 = \sqrt{\frac{\mu}{a}}$ . Furthermore, calculating

$$\frac{\partial}{\partial \sigma^2(t)} \frac{d}{dt} \sigma^2(t) = -2a\sigma^2(t) \quad (19)$$

and evaluating at  $\sigma^2(t) = \hat{\sigma}^2$  demonstrates the equilibrium phenotypic variance is stable when  $a, \mu > 0$ . Hence, calculating

$$\frac{\partial}{\partial \bar{x}(t)} \frac{d}{dt} \bar{x}(t) = -a\sigma^2(t) \quad (20)$$

and evaluating at  $\sigma^2(t) = \hat{\sigma}^2$  and  $\bar{x}(t) = \hat{x}$  demonstrates the equilibrium phenotypic mean is stable when  $a, \mu > 0$ . Finally, calculating

$$\frac{\partial}{\partial N(t)} \frac{d}{dt} N(t) = r - \frac{a}{2}((\theta - \bar{x}(t))^2 + \sigma^2(t)) - 2cN(t) \quad (21)$$

and evaluating at  $\sigma^2(t) = \hat{\sigma}^2$ ,  $\bar{x}(t) = \hat{x}$  and  $N(t) = \hat{N}$  demonstrates the equilibrium total abundance is stable when  $a, c, \mu > 0$ , and  $r > \frac{1}{2}\sqrt{a\mu}$ .

### 3 The relation between diffusion and convolution with a Gaussian kernel

For continuous  $g : \mathbb{R}^d \rightarrow \mathbb{R}$ , consider the deterministic Cauchy problem

$$\begin{cases} \dot{f}(x, t) = \Delta f(x, t), & (x, t) \in \mathbb{R}^d \times (0, \infty) \\ f(x, t) = g(x), & (x, t) \in \mathbb{R}^d \times \{0\}. \end{cases} \quad (22)$$

According to Evans (2010), the fundamental solution of (22) is

$$\Phi(x, t) = \frac{1}{(4\pi t)^{d/2}} \exp \left( -\frac{|x|^2}{4t} \right), \quad (x, t) \in (0, \infty) \times \mathbb{R}^d, \quad (23)$$

where  $|x| = \sqrt{\sum_i x_i^2}$ . The solution  $f(x, t)$  of PDE (22) is then given by the convolution

$$f(x, t) = \int_{\mathbb{R}^d} \Phi(x - y, t) g(y) dy, \quad (x, t) \in (0, \infty) \times \mathbb{R}^d. \quad (24)$$

Hence, by the fundamental theorem of calculus,

$$\begin{aligned}
f(x, t) + \int_t^{t+1} \dot{f}(x, s) ds &= f(x, t+1) \\
&= \int_{\mathbb{R}^d} \Phi(x-y, t+1) g(y) dy = \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \Phi(x-y, 1) \Phi(y-z, t) g(z) dz dy \\
&= \int_{\mathbb{R}^d} \Phi(x-y, 1) f(t, y) dy. \quad (25)
\end{aligned}$$

66 In particular,

$$f(x, t) + \int_t^{t+1} \Delta f(x, s) ds = \int_{\mathbb{R}^d} \Phi(1, x-y) f(y, t) dy. \quad (26)$$

## 67 4 Deterministic dynamics of $\sigma^2(t)$

68 Picking up from the main text §2.1,

$$\begin{aligned}
\frac{d}{dt} \sigma^2(t) &= \frac{d}{dt} \int_{\mathbb{R}} (x - \bar{x}(t))^2 p(x, t) dx = \int_{\mathbb{R}} 2(x - \bar{x}(t)) \frac{d}{dt} \bar{x}(t) + (x - \bar{x}(t))^2 \frac{\partial}{\partial t} p(x, t) dx \\
&= \int_{\mathbb{R}} (x - \bar{x}(t))^2 \left( (m(\nu, x) - \bar{m}(t)) p(x, t) + \frac{\mu}{2} \frac{\partial^2}{\partial x^2} p(x, t) \right) dx \\
&= \int_{\mathbb{R}} ((x - \bar{x}(t))^2 - \sigma^2(t) + \sigma^2(t)) (m(\nu, x) - \bar{m}(t)) p(x, t) + (x - \bar{x}(t))^2 \frac{\mu}{2} \frac{\partial^2}{\partial x^2} p(x, t) dx \\
&= \text{Cov}_t((x - \bar{x}(t))^2, m(\nu, x)) + \frac{\mu}{2} \int_{\mathbb{R}} (x - \bar{x}(t))^2 \frac{\partial^2}{\partial x^2} p(x, t) dx. \quad (27)
\end{aligned}$$

69 Applying integration by parts twice yields

$$\int_{-\infty}^{+\infty} (x - \bar{x}(t))^2 \frac{\partial^2}{\partial x^2} p(x, t) dx = 2. \quad (28)$$

### 70 4.1 Simplifying covariances with fitness under the assumption of a Gaussian density

72 By assuming

$$p(x, t) = \frac{\exp\left(-\frac{(x - \bar{x}(t))^2}{2\sigma^2(t)}\right)}{\sqrt{2\pi\sigma^2(t)}} \quad (29)$$

73 we have

$$\begin{aligned}
\sigma^2 \left( \frac{\partial \bar{m}}{\partial \bar{x}} - \overline{\frac{\partial m}{\partial \bar{x}}} \right) &= \sigma^2 \left( \frac{\partial}{\partial \bar{x}} \int_{\mathbb{R}} m(\nu, x) p(x, t) dx - \int_{\mathbb{R}} p(x, t) \frac{\partial}{\partial \bar{x}} m(\nu, x) dx \right) \\
&= \sigma^2 \int_{\mathbb{R}} m(\nu, x) \frac{\partial}{\partial \bar{x}} p(x, t) dx = \sigma^2 \int_{\mathbb{R}} \frac{x - \bar{x}(t)}{\sigma^2} m(\nu, x) p(x, t) dx \\
&= \int_{\mathbb{R}} (x - \bar{x})(m(\nu, x) - \bar{m}) p(x, t) dx = \text{Cov}_t(m, x), \quad (30)
\end{aligned}$$

74 and

$$\begin{aligned}
2\sigma^4 \left( \frac{\partial \bar{m}}{\partial \sigma^2} - \overline{\frac{\partial m}{\partial \sigma^2}} \right) &= 2\sigma^4 \left( \frac{\partial}{\partial \sigma^2} \int_{\mathbb{R}} m(\nu, x) p(x, t) dx - \int_{\mathbb{R}} p(x, t) \frac{\partial}{\partial \sigma^2} m(\nu, x) dx \right) \\
&= 2\sigma^4 \int_{\mathbb{R}} \frac{(x - \bar{x})^2 - \sigma^2}{2\sigma^4} m(\nu, x) p(x, t) dx = \int_{\mathbb{R}} ((x - \bar{x})^2 - \sigma^2) (m(\nu, x) - \bar{m}) p(x, t) dx \\
&= \text{Cov}_t((x - \bar{x})^2, m). \quad (31)
\end{aligned}$$

## 75 5 Comparing our treatment of white noise to Da Prato and 76 Zabczyk (2014)

77 Our approach in the main text is inspired by the treatment provided in §4.2 of Da Prato and Zabczyk  
78 (2014). Here the authors develop a stochastic integral of operator-valued processes. In particular, they  
79 consider processes indexed by time  $t \geq 0$  valued as Hilbert-Schmidt operators  $\Phi(t)$  and define the norm

$$\|\Phi\|_t = \sqrt{\mathbb{E} \left( \int_0^t \text{Tr}[\Phi(s)\Phi^*(s)] ds \right)}, \quad t \geq 0. \quad (32)$$

80 In our case we only consider the so-called multiplication operators. That is, processes that consist of operators  
81  $\Phi(t)$  having the form  $\Phi(t)g(x) = \varphi(x, t)g(x)$  such that  $\varphi(\cdot, t) \in L^2(\mathbb{R})$  a.s. for each  $t \geq 0$ . In this case  
82  $\Phi(t) = \Phi^*(t)$  and

$$\|\Phi\|_t = \|\varphi\|_t = \sqrt{\mathbb{E} \left( \int_0^t \int_{\mathbb{R}} \varphi^2(x, s) dx ds \right)}, \quad t \geq 0. \quad (33)$$

83 Da Prato and Zabczyk (2014) form the space  $\mathcal{N}_W^2(0, T)$  of Hilbert-Schmidt operator-valued predictable  
84 processes  $\Phi(t)$  that satisfy  $\|\Phi\|_T < +\infty$  for some  $T > 0$ . This corresponds to our more specialized space  $\mathcal{N}_2$   
85 that consists of  $L^2(\mathbb{R})$  valued processes  $\varphi(x, t)$  such that  $\|\varphi\|_t < +\infty$  for all  $t \geq 0$ . In their treatment,  $W(t)$   
86 plays a similar role to our generalized process  $\mathbf{W}_t$ . For  $\Phi \in \mathcal{N}_W^2(0, T)$ , they denote the stochastic integral  
87 for  $t \in [0, T]$  by  $\Phi \cdot W(t)$ . Hence, for  $\Phi(t)g(x) = \varphi(x, t)g(x)$  as above,  $\mathbf{W}_t(\varphi) = \Phi \cdot W(t)$ . The authors then  
88 prove the following:

89 **Proposition 4.28** *Assume that  $\Phi_1, \Phi_2 \in \mathcal{N}_W^2(0, T)$ . Then*

$$\mathbb{E}(\Phi_i \cdot W(t)) = 0, \quad \mathbb{E}(\|\Phi_i \cdot W(t)\|^2) < +\infty, \quad \forall t \in [0, T].$$

90 **Corollary 4.29** *Under the same assumptions as Proposition 4.28,*

$$\mathbb{C}(\Phi_1 \cdot W(t), \Phi_2 \cdot W(s)) = \mathbb{E} \left( \int_0^{t \wedge s} \text{Tr}[\Phi_2(r)\Phi_1^*(r)] dr \right), \quad \forall t, s \in [0, T].$$

91 Simplifying these expressions for the multiplication operators described above returns equations (23) and  
92 (24) of the main text.

## 93 6 Simulating the rescaled process

94 Here we provide a detailed description of the branching random walk and how we have chosen to rescale it.  
95 We focus on the case

$$m(\nu, x) = r - \frac{a}{2}(\theta - x)^2 - c \int_{\mathbb{R}} \nu(x, t) dx. \quad (34)$$

## 6.1 Description of simulation

We begin by describing the branching random walk before introducing our scheme to rescale it. Our branching random walk follows closely the description of branching Brownian motion in the main text. However, we replace exponentially distributed lifetimes with deterministic unit time steps for easier implementation. Hence, we restrict time to  $t = 0, 1, 2, \dots$ , and so on. Furthermore, we allow individual fitness to depend on both trait value and the state of the entire population. For time  $t$  we write  $\{\xi_1(t), \dots, \xi_{N(t)}(t)\}$  as the set of trait values across all  $N(t)$  individuals alive in the population. Since our simulation follows discrete individuals instead of continuous distributions of trait values, we can write

$$\nu(x, t) = \sum_{i=1}^{N(t)} \delta(x - \xi_i(t)), \quad (35)$$

where  $\delta(x - \xi_i(t))$  denotes the point mass located at  $\xi_i(t)$ . For simplicity we assume perfect heritability. At each iteration we draw, for each individual, a random number of offspring from a Negative-Binomial distribution. Recall the Negative-Binomial distribution models the number of failed Bernoulli trials that occur before a given number of successful trials. Denoting  $q$  the probability of success for each trial and  $s$  the number of successes, the mean and variance is given respectively by

$$\frac{s(1-q)}{q}, \quad \frac{s(1-q)}{q^2}. \quad (36)$$

Then if we require the  $i$ th individual to have mean number offspring  $\mathcal{W}(\nu, \xi_i)$  and variance equal to  $V$ , the parameters of the associated Negative-Binomial distribution become

$$q(\nu, \xi_i) = \frac{\mathcal{W}(\nu, \xi_i)}{V}, \quad s(\nu, \xi_i) = \frac{\mathcal{W}^2(\nu, \xi_i)}{V - \mathcal{W}(\nu, \xi_i)}. \quad (37)$$

This imposes the restriction  $V > \mathcal{W}(\nu, \xi_i)$ . For each offspring produced by the individual with trait value  $\xi_i(t)$ , we assign indepently drawn trait values normally distributed around  $\xi_i(t)$  with variance  $\mu$ . This summarizes the basic structure of our simulation. To impose selection and density dependent growth rates, we set

$$\mathcal{W}(\nu, \xi_i) = \exp \left( r - \frac{a}{2}(\theta - \xi_i)^2 - c \int_{\mathbb{R}} \nu(x, t) dx \right), \quad (38)$$

where the above integral becomes  $\int_{\mathbb{R}} \nu(x, t) dx = \sum_{i=1}^{N(t)} 1 = N(t)$ .

## 6.2 Rescaling

To rescale the branching random walk by a positive integer  $n$ , we rescale mutational variance by  $\mu \rightarrow \mu/n$ , time by  $t \rightarrow t/n$  and the reproductive law by  $V \rightarrow V$  and

$$\mathcal{W}(\nu, \xi_i) \rightarrow \mathcal{W}^{(n)}(\nu, \xi_i) = \exp \left( \frac{r}{n} - \frac{a}{2n}(\theta - \xi_i)^2 - \frac{c}{n^2}N(t) \right) = \exp \left( \frac{r}{n} - \frac{a}{2n}(\theta - \xi_i)^2 - \frac{c}{n}N^{(n)}(t) \right). \quad (39)$$

We also replace individual mass with  $\frac{1}{n}$  and write rescaled abundance as  $N^{(n)}(t) = \frac{1}{n}N(nt)$ . When it exists, we denote by  $N^{(\infty)}(t)$  the limiting process of  $N^{(n)}(t)$ . Then

$$\lim_{n \rightarrow \infty} n \left( \mathcal{W}^{(n)}(\nu, \xi_i) - 1 \right) = r - \frac{a}{2}(\theta - \xi_i)^2 - cN^{(\infty)}(t). \quad (40)$$

Note that, since the limiting fitness function is bounded above and decreases linearly with respect to  $N^{(\infty)}(t)$ , it satisfies the hypotheses of Champagnat, Ferrière and Méléard (2006). We have implemented this simulation in the programming language Julia. A copy can be found at the url:

<https://github.com/bobweek/branching.brownian.motion.and.spde>

For the sake of illustration, we simulated the unscaled process ( $n = 1$ ) and the rescaled process with  $n = 5$  and  $n = 20$  for 50 units of time. Results are shown in Figure 2 of the main text.

## 7 Derivation of SDE for $\bar{x}$ and $\sigma^2$

Picking up from §2.2.3 of the main text, we have

$$\tilde{x}(t) = \int_{\mathbb{R}} x\nu(x, t)dx, \quad \tilde{\sigma}^2(t) = \int_{\mathbb{R}} x^2\nu(x, t)dx \quad (41)$$

and

$$\tilde{x}(t) = \tilde{x}(0) + \int_0^t \int_{\mathbb{R}} \nu(x, s)m(\nu, x)x + x\sqrt{V\nu(x, s)}\dot{W}(x, s)dx ds, \quad (42)$$

$$\tilde{\sigma}^2(t) = \tilde{\sigma}^2(0) + \int_0^t \int_{\mathbb{R}} \nu(x, s) (m(\nu, x)x^2 + \mu) + x^2\sqrt{V\nu(x, s)}\dot{W}(x, s)dx ds. \quad (43)$$

### 7.1 Derivation for trait mean

We make use of the notation

$$\begin{cases} \|N\|_2 = \sqrt{V \int_{\mathbb{R}} \nu(x, t)dx} = \sqrt{VN} \\ \|\tilde{x}\|_2 = \sqrt{V \int_{\mathbb{R}} x^2\nu(x, t)dx} \\ \langle \tilde{x}, N \rangle = V \int_{\mathbb{R}} x\nu(x, t)dx = \bar{x}VN. \end{cases} \quad (44)$$

Rewriting formula (42) as an SDE provides

$$d\tilde{x} = \left( \bar{x}mN + \frac{\mu}{2} \int_{\mathbb{R}} x\Delta\nu(x, t)dx \right) dt + \|\tilde{x}\|_2 d\tilde{W}_2, \quad (45)$$

where,

$$d\tilde{W}_2 = d\hat{\mathbf{W}}_t(\sqrt{Vx^2\nu}) = \frac{1}{\|\tilde{x}\|_2} \int_{\mathbb{R}} x\sqrt{V\nu(x, t)}\dot{W}(x, t)dx dt. \quad (46)$$

Using Itô's quotient rule on  $\bar{x} = \tilde{x}/N$ , we obtain

$$d\bar{x} = d\left(\frac{\tilde{x}}{N}\right) = \frac{\tilde{x}}{N} \left( \frac{d\tilde{x}}{\tilde{x}} - \frac{dN}{N} - \frac{d\tilde{x}}{\tilde{x}} \frac{dN}{N} + \left(\frac{dN}{N}\right)^2 \right) = \frac{d\tilde{x}}{N} - \bar{x} \frac{dN}{N} - \frac{d\tilde{x}}{N} \frac{dN}{N} + \bar{x} \left(\frac{dN}{N}\right)^2. \quad (47)$$

From Table 1 of the main text  $d\tilde{x}dN = \langle \tilde{x}, N \rangle$  and  $dN^2 = \|N\|_2^2$ . Hence,

$$\begin{aligned} d\bar{x} &= \bar{x}m dt + \frac{\|\tilde{x}\|_2}{N} d\tilde{W}_2 - \bar{x} \left( \bar{m} dt + \sqrt{\frac{V}{N}} dW_1 \right) - \frac{\langle \tilde{x}, N \rangle}{N^2} dt + \bar{x} \frac{\|N\|_2^2}{N^2} dt \\ &= (\bar{x}m - \bar{x}\bar{m}) dt + \frac{\|\tilde{x}\|_2}{N} d\tilde{W}_2 - \bar{x} \sqrt{\frac{V}{N}} dW_1 - V \frac{\bar{x}}{N} dt + V \frac{\bar{x}}{N} dt \\ &= \text{Cov}_t(x, m) + \frac{\|\tilde{x}\|_2}{N} d\tilde{W}_2 - \bar{x} \sqrt{\frac{V}{N}} dW_1. \end{aligned} \quad (48)$$

Note that

$$\begin{aligned}
\frac{\|\tilde{x}\|_2}{N} d\tilde{W}_2 - \bar{x} \sqrt{\frac{V}{N}} dW_1 &= \frac{1}{N} \int_{\mathbb{R}} x \sqrt{V\nu(x,t)} \dot{W}(x,t) dx - \frac{\bar{x}}{N} \int_{\mathbb{R}} \sqrt{V\nu(x,t)} \dot{W}(x,t) dx \\
&= \int_{\mathbb{R}} \frac{x - \bar{x}}{N} \sqrt{V\nu(x,t)} \dot{W}(x,t) dx
\end{aligned} \tag{49}$$

137 and

$$\mathbb{V} \left( \int_{\mathbb{R}} \frac{x - \bar{x}}{N} \sqrt{V\nu(x,t)} \dot{W}(x,t) dx \right) = \frac{V}{N} \int_{\mathbb{R}} (x - \bar{x})^2 p(x,t) dx = V \frac{\sigma^2}{N}. \tag{50}$$

138 Hence, by setting

$$dW_2 = \frac{\int_{\mathbb{R}} \frac{(x - \bar{x})}{N} \sqrt{V\nu(x,t)} \dot{W}(x,t) dx}{\sqrt{V\sigma^2/N}} \tag{51}$$

139 we can write

$$d\bar{x} = \text{Cov}_t(x, m) dt + \sqrt{V \frac{\sigma^2}{N}} dW_2. \tag{52}$$

## 140 7.2 Derivation for trait variance

141 We make use of the notation

$$\begin{cases} \|\tilde{\sigma}^2\|_2 = \sqrt{V \int_{\mathbb{R}} x^4 \nu(x,t) dx} \\ \langle \tilde{\sigma}^2, N \rangle = V \int_{\mathbb{R}} x^2 \nu(x,t) dx = \overline{x^2} V N. \end{cases} \tag{53}$$

142 Applying formula (43) provides

$$d\tilde{\sigma}^2 = \left( \overline{x^2 m} N + \mu N \right) dt + \|\tilde{\sigma}^2\|_2 d\tilde{W}_3 \tag{54}$$

143 where

$$d\tilde{W}_3 = d\hat{\mathbf{W}}_t(\sqrt{V x^4 \nu}) = \frac{1}{\|\tilde{\sigma}^2\|_2} \int_{\mathbb{R}} x^2 \sqrt{V\nu(x,t)} \dot{W}(x,t) dx. \tag{55}$$

144 Using Itô's quotient rule on  $\overline{x^2} = \tilde{\sigma}^2/N$ , we obtain

$$d\overline{x^2} = d \left( \frac{\tilde{\sigma}^2}{N} \right) = \frac{\tilde{\sigma}^2}{N} \left( \frac{d\tilde{\sigma}^2}{\tilde{\sigma}^2} - \frac{dN}{N} - \frac{d\tilde{\sigma}^2}{\tilde{\sigma}^2} \frac{dN}{N} + \left( \frac{dN}{N} \right)^2 \right) = \frac{d\tilde{\sigma}^2}{N} - \overline{x^2} \frac{dN}{N} - \frac{d\tilde{\sigma}^2}{N} \frac{dN}{N} + \overline{x^2} \left( \frac{dN}{N} \right)^2. \tag{56}$$

145 Table 1 of the main text implies  $d\tilde{W}_3 dW_1 = \langle \tilde{\sigma}^2, N \rangle$  and hence



$$\begin{aligned}
d\bar{x}^2 &= \left( \overline{x^2 m} + \mu \right) dt + \frac{\|\tilde{\sigma}^2\|_2}{N} d\tilde{W}_3 - \bar{x}^2 \left( \bar{m} dt + \sqrt{\frac{V}{N}} dW_1 \right) - \frac{\langle \tilde{\sigma}^2, N \rangle}{N^2} dt + \bar{x}^2 \frac{\|N\|_2^2}{N^2} dt \\
&= \left( \overline{x^2 m} - \bar{x}^2 \bar{m} dt + \mu \right) dt + \frac{\|\tilde{\sigma}^2\|_2}{N} d\tilde{W}_3 - \bar{x}^2 \sqrt{\frac{V}{N}} dW_1 - \bar{x}^2 \frac{V}{N} dt + \bar{x}^2 \frac{V}{N} dt \\
&= \left( \text{Cov}_t(x^2, m) + \mu \right) dt + \frac{\|\tilde{\sigma}^2\|_2}{N} d\tilde{W}_3 - \bar{x}^2 \sqrt{\frac{V}{N}} dW_1. \quad (57)
\end{aligned}$$

<sup>146</sup> Setting  $F(y, z) = y - z^2$ , use Itô's formula on  $\sigma^2 = F(\bar{x}^2, \bar{x}) = \bar{x}^2 - \bar{x}^2$  to obtain:

$$\begin{aligned}
d\sigma^2 &= d\bar{x}^2 - 2\bar{x}d\bar{x} - (d\bar{x})^2 = \left( \text{Cov}_t(x^2, m) + \mu \right) dt + \frac{\|\tilde{\sigma}^2\|_2}{N} d\tilde{W}_3 - \bar{x}^2 \sqrt{\frac{V}{N}} dW_1 \\
&\quad - 2\bar{x} \left( \text{Cov}_t(x, m) + \mu dt + \sqrt{\frac{V\sigma^2}{N}} dW_2 \right) - \left( \text{Cov}_t(x, m) dt + \mu dt + \sqrt{\frac{V\sigma^2}{N}} dW_2 \right)^2 \\
&= \left( \text{Cov}_t(x^2 - 2\bar{x}x, m) + \mu \right) dt + \frac{\|\tilde{\sigma}^2\|_2}{N} d\tilde{W}_3 - \bar{x}^2 \sqrt{\frac{V}{N}} dW_1 - 2\bar{x} \sqrt{\frac{V\sigma^2}{N}} dW_2 - \left( \frac{V\sigma^2}{N} \right) dt \\
&= \left( \text{Cov}_t(x - \bar{x})^2, m \right) + \mu - \frac{V\sigma^2}{N} dt + \frac{\|\tilde{\sigma}^2\|_2}{N} d\tilde{W}_3 - \bar{x}^2 \sqrt{\frac{V}{N}} dW_1 - 2\bar{x} \sqrt{\frac{V\sigma^2}{N}} dW_2. \quad (58)
\end{aligned}$$

<sup>147</sup> In light of

$$\begin{aligned}
\frac{\|\tilde{\sigma}^2\|_2}{N} d\tilde{W}_3 - \bar{x}^2 \sqrt{\frac{V}{N}} dW_1 - 2\bar{x} \sqrt{\frac{V\sigma^2}{N}} dW_2 &= \frac{1}{N} \int_{\mathbb{R}} (x^2 - \bar{\sigma}^2 - 2\bar{x}(x - \bar{x})) \sqrt{V\nu(x, t)} \dot{W}(x, t) dx \\
&= \frac{1}{N} \int_{\mathbb{R}} ((x - \bar{x})^2 - \sigma^2) \sqrt{V\nu(x, t)} \dot{W}(x, t) dx \quad (59)
\end{aligned}$$

<sup>148</sup> and

$$\begin{aligned}
\frac{1}{N} \int_{\mathbb{R}} \left( ((x - \bar{x})^2 - \sigma^2) \sqrt{V\nu(x, s)} \right)^2 dx &= \frac{V}{N} \left( \int_{\mathbb{R}} ((x - \bar{x})^4 - 2(x - \bar{x})^2 \sigma^2 + \sigma^4) p(x, t) dx \right) \\
&= \frac{V}{N} \left( \overline{(x - \bar{x})^4} - \sigma^4 \right) \quad (60)
\end{aligned}$$

<sup>149</sup> we set

$$dW_3 = \frac{\int_{\mathbb{R}} ((x - \bar{x})^2 - \sigma^2) \sqrt{V\nu(x, t)} \dot{W}(x, t) dx}{V \left( \overline{(x - \bar{x})^4} - \sigma^4 \right)} \quad (61)$$

150 so that

$$d\sigma^2 = \text{Cov}_t\left((x - \bar{x})^2, m\right)dt + \left(\mu - V\frac{\sigma^2}{N}\right)dt + \sqrt{V\frac{(x - \bar{x})^4 - \sigma^4}{N}}dW_3. \quad (62)$$

151 Table 1 of the main text implies

$$dW_1dW_2 = \frac{\int_{\mathbb{R}}(x - \bar{x})\nu(x, t)dx}{\sqrt{N\sigma^2}}dt = 0, \quad (63)$$

$$dW_1dW_3 = \frac{\int_{\mathbb{R}}((x - \bar{x})^2 - \sigma^2)\nu(x, t)dx}{\sqrt{(x - \bar{x})^4 - \sigma^4}}dt = 0, \quad (64)$$

$$dW_2dW_3 = \frac{\int_{\mathbb{R}}(x - \bar{x})((x - \bar{x})^2 - \sigma^2)p(x, t)dx}{\sqrt{\sigma^2((x - \bar{x})^4 - \sigma^4)}}dt = \frac{N\overline{(x - \bar{x})^3}}{\sqrt{\sigma^2((x - \bar{x})^4 - \sigma^4)}}dt. \quad (65)$$

152 In particular, when  $p$  is a Gaussian curve  $dW_2dW_3 = 0$ .

## 153 8 Relating fitness of expressed traits to fitness of breeding values

154

155 Following §2.3.2 of the main text, we have  $\sigma^2 = G + \eta$  and

$$m^*(\rho, g) = \int_{\mathbb{R}} m(\nu, x)\psi(x, g)dx. \quad (66)$$

156 Hence,

$$\begin{aligned} \overline{\frac{\partial m^*}{\partial \bar{x}}} &= \int_{\mathbb{R}} \frac{\rho(g, t)}{N(t)} \frac{\partial}{\partial \bar{x}} \int_{\mathbb{R}} m(\nu, x)\psi(x, g)dx dg = \\ &= \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{\rho(g, t)}{N(t)} \psi(x, g) dg \frac{\partial}{\partial \bar{x}} m(\nu, x) dx = \int_{\mathbb{R}} p(x, t) \frac{\partial}{\partial \bar{x}} m(\nu, x) dx = \overline{\frac{\partial m}{\partial \bar{x}}} \end{aligned} \quad (67)$$

157 and

$$\begin{aligned} \overline{\frac{\partial m^*}{\partial G}} &= \int_{\mathbb{R}} \frac{\rho(g, t)}{N(t)} \frac{\partial}{\partial G} \int_{\mathbb{R}} m(\nu, x)\psi(x, g)dx dg = \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{\rho(g, t)}{N(t)} \psi(x, g) dg \frac{\partial m}{\partial G} dx = \\ &= \int_{\mathbb{R}} p(x, t) \frac{\partial m}{\partial \sigma^2} \frac{\partial \sigma^2}{\partial G} dx = \overline{\frac{\partial m}{\partial \sigma^2}}. \end{aligned} \quad (68)$$

## 9 Derivation of diffuse coevolution model

In this section we provide a derivation of our model of diffuse coevolution driven by competition. Since most of the work in this derivation has already been completed in Supplementary Material §7, we focus here on deriving the Malthusian fitness  $m$  as a function of trait value  $x$ . We begin with discrete populations of individuals. In particular, we begin by assuming population size  $n_i$  is an integer for each species  $i = 1, \dots, S$  before passing to the large population size limit.

The reduction in fitness for an individual of species  $i$  caused by competition is captured multiplicatively by  $0 < C_i \leq 1$ . Biologically this assumes all competitors affect individuals of a given species equally by consuming the same amount of resources. This is a mean-field interaction since any individual that consumes resources has an effect on the fitness of all other individuals competing for the same resources. Denote by  $x_{ij}$  the trait value of the  $j$ -th individual belonging to species  $i$ . The set of trait values across all individuals in the community at time  $t \geq 0$  is written  $X = \{x_{ij}\}$ . We denote by  $\mathcal{B}_{ij}$  a function that maps  $X$  to the cumulative effect of all competitive interactions on the fitness of the  $j$ -th individual in species  $i$ . Since individuals do not compete with themselves the net multiplicative effects on fitness of both interspecific and intraspecific competition on the  $j$ -th individual in species  $i$  can be summarized by

$$\mathcal{B}_{ij}(X) = C_i^{\sum_{l \neq j} \mathcal{O}_{ii}(x_{ij} - x_{il}) + \sum_{k \neq i} \sum_{l=1}^{n_k} \mathcal{O}_{ik}(x_{ij} - x_{kl})}, \quad (69)$$

where  $\mathcal{O}_{ij}$ , defined in the main text, measures the overlap in resource use between individuals of species  $i$  and  $j$  as a function of their niche-centers. Writing  $\mathcal{W}_{ij}(X)$  as the average number of offspring left by the  $j$ -th individual of species  $i$ , we have

$$\mathcal{W}_{ij}(X) = \mathcal{A}_i(x_{ij}) \mathcal{B}_{ij}(X), \quad (70)$$

where  $\mathcal{A}_i(x) = \int_{\mathbb{R}} e_i(\zeta) u_i(\zeta, x) d\zeta$  accounts for abiotic selection and  $e_i$  has been defined in the main text.

We now turn to a diffusion limit. Since we have more than one population, we take the diffusion limit for each population one at a time starting with population 1. We write  $\mathbf{n} = (n_1, \dots, n_S)$ . Following Méléard and Roelly (1993, 1992) we rescale generation time and individual mass to  $\frac{1}{n_1}$  and mean of the reproductive law to

$$\mathcal{W}_{1j}^{(\mathbf{n})}(X) = \mathcal{A}_1(x_{1j})^{1/n_1} \exp \left( \frac{\ln C_1}{n_1^2} \sum_{l \neq j} \mathcal{O}_{11}(x_{1j} - x_{1l}) + \frac{\ln C_1}{n_1} \sum_{k \neq 1} \frac{1}{n_k} \sum_{l=1}^{n_k} \mathcal{O}_{1k}(x_{1j} - x_{kl}) \right). \quad (71)$$

For large  $n_1$ , we have the approximation

$$\mathcal{W}_{1j}^{(\mathbf{n})}(X) \approx \mathcal{A}_1(x_{1j})^{1/n_1} \left( 1 + \frac{\ln C_1}{n_1^2} \sum_{l \neq j} \mathcal{O}_{11}(x_{1j} - x_{1l}) + \frac{\ln C_1}{n_1} \sum_{k \neq 1} \frac{1}{n_k} \sum_{l=1}^{n_k} \mathcal{O}_{1k}(x_{1j} - x_{kl}) \right). \quad (72)$$

Hence

$$\lim_{n_1 \rightarrow \infty} n_1 \left( \mathcal{W}_{1j}^{(\mathbf{n})}(X) - 1 \right) = \ln \mathcal{A}_1(x_{1j}) + \left( \int_{\mathbb{R}} \mathcal{O}_{11}(x_{1j} - y) \nu_1(y, t) dy + \sum_{k \neq 1} \frac{1}{n_k} \sum_{l=1}^{n_k} \mathcal{O}_{1k}(x_{1j} - x_{kl}) \right) \ln C_1. \quad (73)$$

We write  $\lim_{\mathbf{n} \rightarrow \infty}$  for the iterated limit  $\lim_{n_S \rightarrow \infty} \dots \lim_{n_1 \rightarrow \infty}$  and, assuming  $\nu_i(\cdot, t) \in C_1^+(\mathbb{R})$  for  $i = 1, \dots, S$  and  $t \in [0, \infty)$ , we set  $\boldsymbol{\nu} = (\nu_1, \dots, \nu_S)$ . Then, for any  $\boldsymbol{\nu}$ , the diffusion limits for the remaining populations provides the Malthusian parameter for individuals in species  $i$  with trait value  $x_{1j}$  as

$$m_1(\boldsymbol{\nu}, x_{1j}) := \lim_{\mathbf{n} \rightarrow \infty} n_1 \left( \mathcal{W}_{1j}^{(\mathbf{n})}(X) - 1 \right) = \ln \mathcal{A}_1(x) + \left( \sum_{k=1}^S \int_{\mathbb{R}} \mathcal{O}_{1k}(x_{1j} - y) \nu_k(y, t) dy \right) \ln C_1. \quad (74)$$

186 We compute the average niche overlap of an individual in species  $i$  with niche location  $x$  across all individuals  
187 in species  $j$  as

$$\bar{\mathcal{O}}_{ij}(x, t) = \frac{\int_{\mathbb{R}} \mathcal{O}_{ij}(x - y) \nu_j(y, t) dy}{\int_{\mathbb{R}} \nu_j(y, t) dy}. \quad (75)$$

188 We now assume the resource utilization curves  $u_i(\zeta)$  and phenotypic densities  $\nu_i(x, t)$  are Gaussian curves  
189 for  $i = 1, \dots, S$ . In this case  $\bar{\mathcal{O}}_{ij}(x, t)$  simplifies to

$$\bar{\mathcal{O}}_{ij}(x, t) = \frac{\int_{\mathbb{R}} \mathcal{O}_{ij}(x - y) \nu_j(y, t) dy}{\int_{\mathbb{R}} \nu_j(y, t) dy} = \frac{U_i U_j}{\sqrt{2\pi(w_i + w_j + \sigma_j^2(t))}} \exp\left(-\frac{(x - \bar{x}_j(t))^2}{2(w_i + w_j + \sigma_j^2(t))}\right). \quad (76)$$

190 Setting

$$\sigma_i^2(t) = G_i(t) + \eta_i, \quad (77a)$$

$$R_i = \ln\left(\frac{Q_i U_i}{\sqrt{1 + A_i w_i}}\right), \quad (77b)$$

$$a_i = \frac{A_i}{1 + A_i w_i}, \quad (77c)$$

$$\tilde{b}_{ij}(t) = \frac{1}{w_i + w_j + \sigma_j^2(t)}, \quad (77d)$$

$$c_i = -\ln C_i, \quad (77e)$$

191 we get

$$m_i(\boldsymbol{\nu}, x) = R_i - \frac{a_i}{2} (x - \theta_i)^2 - c_i \sum_{j=1}^S N_j(t) U_i U_j \sqrt{\frac{\tilde{b}_{ij}(t)}{2\pi}} e^{-\frac{\tilde{b}_{ij}(t)}{2} (x - \bar{x}_j(t))^2}. \quad (78)$$

192 Hence, our fitness function is bounded above and decreases linearly with total abundance, as required by the  
193 assumptions of Champagnat, Ferrière and Méléard (2006). For the remainder of the derivation we suppress  
194 notation indicating dependency on  $\boldsymbol{\nu}$ ,  $x$  and  $t$ . From (78) we calculate

$$\frac{\partial m_i}{\partial \bar{x}_i} = c_i N_i U_i^2 \tilde{b}_{ii} (x - \bar{x}_i) \sqrt{\frac{\tilde{b}_{ii}}{2\pi}} e^{-\frac{\tilde{b}_{ii}}{2} (x - \bar{x}_i)^2} \quad (79)$$

$$\begin{aligned} \frac{\partial m_i}{\partial G_i} &= \frac{c_i N_i U_i^2}{2} \left( \frac{(x - \bar{x}_i)^2 - G_i - \eta_i - 2w_i}{(G_i + \eta_i + 2w_i)^2} \right) \sqrt{\frac{\tilde{b}_{ii}}{2\pi}} e^{-\frac{\tilde{b}_{ii}}{2} (x - \bar{x}_i)^2} \\ &= \frac{c_i N_i U_i^2 \tilde{b}_{ii}^2}{2} ((x - \bar{x}_i)^2 - \sigma_i^2 - 2w_i) \sqrt{\frac{\tilde{b}_{ii}}{2\pi}} e^{-\frac{\tilde{b}_{ii}}{2} (x - \bar{x}_i)^2}. \end{aligned} \quad (80)$$

195 Note that

$$\begin{aligned}
& \sqrt{\frac{\tilde{b}_{ii}}{2\pi}} \exp\left(-\frac{\tilde{b}_{ii}}{2}(x - \bar{x}_i)^2\right) \sqrt{\frac{1}{2\pi\sigma_i^2}} \exp\left(-\frac{(x - \bar{x}_i)^2}{2\sigma_i^2}\right) \\
&= \sqrt{\frac{1}{2\pi(\sigma_i^2 + 1/\tilde{b}_{ii})}} \sqrt{\frac{\sigma_i^2 + 1/\tilde{b}_{ii}}{2\pi\sigma_i^2/\tilde{b}_{ii}}} \exp\left(-\frac{\sigma_i^2 + 1/\tilde{b}_{ii}}{2\sigma_i^2/\tilde{b}_{ii}}(x - \bar{x}_i)^2\right) \\
&= \sqrt{\frac{1}{4\pi(\sigma_i^2 + w_i)}} \sqrt{\frac{2(\sigma_i^2 + w_i)}{2\pi\sigma_i^2(\sigma_i^2 + 2w_i)}} \exp\left(-\frac{\sigma_i^2(\sigma_i^2 + 2w_i)}{4(\sigma_i^2 + w_i)}(x - \bar{x}_i)^2\right). \quad (81)
\end{aligned}$$

Hence,

$$\frac{\partial \bar{m}_i}{\partial \bar{x}_i} = 0, \quad (82)$$

$$\begin{aligned}
\frac{\partial \bar{m}_i}{\partial G_i} &= \frac{c_i N_i U_i^2}{2(\sigma_i^2 + 2w_i)^2} \left( \frac{(\sigma_i^2 + 2w_i)\sigma_i^2}{2(w_i + \sigma_i^2)} - \sigma_i^2 - 2w_i \right) \sqrt{\frac{b_{ii}}{2\pi}} \\
&= \frac{c_i N_i U_i^2}{2(\sigma_i^2 + 2w_i)} \left( \frac{\sigma_i^2}{2(\sigma_i^2 + w_i)} - 1 \right) \sqrt{\frac{b_{ii}}{2\pi}} = -\frac{c_i N_i U_i^2 b_{ii}}{2} \sqrt{\frac{b_{ii}}{2\pi}}, \quad (83)
\end{aligned}$$

where

$$b_{ij} = \frac{1}{w_i + w_j + \sigma_i^2 + \sigma_j^2}. \quad (84)$$

The average fitness for species  $i$  is

$$\bar{m}_i = R_i - \frac{a_i}{2} \left( (\bar{x}_i - \theta_i)^2 + G_i + \eta_i \right) - c_i \sum_{j=1}^S N_j U_i U_j \sqrt{\frac{b_{ij}}{2\pi}} e^{-\frac{b_{ij}}{2}(\bar{x}_i - \bar{x}_j)^2}. \quad (85)$$

Thus,

$$\frac{\partial \bar{m}_i}{\partial \bar{x}_i} = a_i(\theta_i - \bar{x}_i) - c_i \sum_j N_j U_i U_j b_{ij}(\bar{x}_j - \bar{x}_i) \sqrt{\frac{b_{ij}}{2\pi}} e^{-\frac{b_{ij}}{2}(\bar{x}_i - \bar{x}_j)^2}, \quad (86)$$

$$\frac{\partial \bar{m}_i}{\partial G_i} = -\frac{a_i}{2} + \frac{c_i}{2} \sum_{j=1}^S N_j U_i U_j b_{ij} (1 - b_{ij}(\bar{x}_i - \bar{x}_j)^2) \sqrt{\frac{b_{ij}}{2\pi}} e^{-\frac{b_{ij}}{2}(\bar{x}_i - \bar{x}_j)^2}. \quad (87)$$

In particular

$$\frac{\partial \bar{m}_i}{\partial G_i} - \frac{\partial \bar{m}_i}{\partial G_i} = -\frac{a_i}{2} + \frac{c_i}{2} \left( N_i U_i^2 b_{ii} \sqrt{\frac{b_{ii}}{2\pi}} + \sum_{j=1}^S N_j U_i U_j b_{ij} (1 - b_{ij}(\bar{x}_i - \bar{x}_j)^2) \sqrt{\frac{b_{ij}}{2\pi}} e^{-\frac{b_{ij}}{2}(\bar{x}_i - \bar{x}_j)^2} \right). \quad (88)$$

Applying equations (37a), (47a) and (47b) of the main text recovers system (52) of the main text.

## 10 The relation between competition coefficients and selection

### 10.1 Derivation of analytical approximations

Just as with most calculations in this work, the derivations are straightforward applications of Gaussian products. That is, if

$$f_1(x) = \frac{1}{\sqrt{2\pi\sigma_1^2}} \exp\left(-\frac{(\mu_1 - x)^2}{2\sigma_1^2}\right), \quad f_2(x) = \frac{1}{\sqrt{2\pi\sigma_2^2}} \exp\left(-\frac{(\mu_2 - x)^2}{2\sigma_2^2}\right), \quad (89)$$

then

$$f_1(x)f_2(x) = \frac{1}{\sqrt{2\pi(\sigma_1^2 + \sigma_2^2)}} \exp\left(-\frac{(\mu_1 - \mu_2)^2}{2(\sigma_1^2 + \sigma_2^2)}\right) \frac{1}{\sqrt{2\pi\tilde{\sigma}^2}} \exp\left(-\frac{(\tilde{\mu} - x)^2}{2\tilde{\sigma}^2}\right), \quad (90)$$

where

$$\tilde{\mu} = \frac{\sigma_2^2\mu_1 + \sigma_1^2\mu_2}{\sigma_1^2 + \sigma_2^2}, \quad \tilde{\sigma}^2 = \frac{\sigma_1^2\sigma_2^2}{\sigma_1^2 + \sigma_2^2}. \quad (91)$$

#### 10.1.1 Calculating $\text{Cov}_{f_{\bar{X}}}(\alpha, \gamma)$

Recalling

$$\alpha(\bar{x}_i, \bar{x}_j) = \frac{c}{\bar{r}} \sqrt{\frac{b}{2\pi}} \exp\left(-\frac{b}{2}(\bar{x}_i - \bar{x}_j)^2\right), \quad (92)$$

$$\gamma(\bar{x}_i, \bar{x}_j) = cNb \left(1 - b(\bar{x}_i - \bar{x}_j)^2\right) \sqrt{\frac{b}{2\pi}} \exp\left(-\frac{b}{2}(\bar{x}_i - \bar{x}_j)^2\right), \quad (93)$$

we have

$$\begin{aligned} \bar{\alpha} &= \int_{\mathbb{R}} \int_{\mathbb{R}} \alpha(\bar{x}_i, \bar{x}_j) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\ &= \frac{c}{\bar{r}} \int_{\mathbb{R}} \frac{1}{\sqrt{2\pi(b^{-1} + V_{\bar{X}})}} \exp\left(-\frac{(\bar{x} - \bar{x}_j)^2}{2(b^{-1} + V_{\bar{X}})}\right) f_{\bar{X}}(\bar{x}_j) d\bar{x}_j = \frac{c/\bar{r}}{\sqrt{2\pi(b^{-1} + 2V_{\bar{X}})}}, \end{aligned} \quad (94)$$

$$\begin{aligned}
\bar{\gamma} &= \int_{\mathbb{R}} \int_{\mathbb{R}} \gamma(\bar{x}_i, \bar{x}_j) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\
&= cNb \int_{\mathbb{R}} \left\{ 1 - \left[ \left( \frac{\bar{x} + bV_{\bar{X}}\bar{x}_j}{1 + bV_{\bar{X}}} - \bar{x}_j \right)^2 + \frac{V_{\bar{X}}}{1 + bV_{\bar{X}}} \right] \right\} \frac{1}{\sqrt{2\pi(b^{-1} + V_{\bar{X}})}} \exp\left(-\frac{(\bar{x} - \bar{x}_j)^2}{2(b^{-1} + V_{\bar{X}})}\right) f_{\bar{X}}(\bar{x}_j) d\bar{x}_j \\
&= cNb \int_{\mathbb{R}} \left\{ 1 - \left[ \left( \frac{\bar{x} - \bar{x}_j}{1 + bV_{\bar{X}}} \right)^2 + \frac{V_{\bar{X}}}{1 + bV_{\bar{X}}} \right] \right\} \frac{1}{\sqrt{2\pi(b^{-1} + V_{\bar{X}})}} \exp\left(-\frac{(\bar{x} - \bar{x}_j)^2}{2(b^{-1} + V_{\bar{X}})}\right) f_{\bar{X}}(\bar{x}_j) d\bar{x}_j \\
&= cNb \left( 1 - \frac{(1 + bV_{\bar{X}})V_{\bar{X}}}{1 + 2bV_{\bar{X}}} \frac{1}{(1 + bV_{\bar{X}})^2} - \frac{V_{\bar{X}}}{1 + bV_{\bar{X}}} \right) \frac{1}{\sqrt{2\pi(b^{-1} + 2V_{\bar{X}})}} \\
&= cNb \left[ 1 - \left( \frac{1}{1 + 2bV_{\bar{X}}} + 1 \right) \frac{V_{\bar{X}}}{1 + bV_{\bar{X}}} \right] \frac{1}{\sqrt{2\pi(b^{-1} + 2V_{\bar{X}})}} \\
&= cNb \left( 1 - \frac{2V_{\bar{X}}}{1 + 2bV_{\bar{X}}} \right) \sqrt{\frac{b}{2\pi(1 + 2bV_{\bar{X}})}}, \quad (95)
\end{aligned}$$

$$\begin{aligned}
\text{Var}_{f_{\bar{X}}}(\alpha) &= \int_{\mathbb{R}} \int_{\mathbb{R}} (\bar{\alpha} - \alpha(\bar{x}_i, \bar{x}_j))^2 f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\
&= \frac{c^2}{\bar{r}^2} \left( \sqrt{\frac{b}{4\pi}} \int_{\mathbb{R}} \int_{\mathbb{R}} \sqrt{\frac{b}{\pi}} \exp(-b(\bar{x}_i - \bar{x}_j)^2) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j - \frac{1}{2\pi(b^{-1} + 2V_{\bar{X}})} \right) \\
&= \frac{c^2}{\bar{r}^2} \left( \sqrt{\frac{b}{4\pi}} \int_{\mathbb{R}} \sqrt{\frac{1}{2\pi(\frac{1}{2b} + V_{\bar{X}})}} \exp(-b(\bar{x} - \bar{x}_j)^2) f_{\bar{X}}(\bar{x}_j) d\bar{x}_j - \frac{1}{2\pi(b^{-1} + 2V_{\bar{X}})} \right) \\
&= \frac{c^2}{\bar{r}^2} \left( \sqrt{\frac{b}{4\pi}} \sqrt{\frac{1}{2\pi(\frac{1}{2b} + 2V_{\bar{X}})}} - \frac{1}{2\pi(b^{-1} + 2V_{\bar{X}})} \right) = \frac{c^2 b}{2\pi \bar{r}^2} \left( \frac{1}{\sqrt{1 + 4bV_{\bar{X}}}} - \frac{1}{1 + 2bV_{\bar{X}}} \right), \quad (96)
\end{aligned}$$

$$\begin{aligned}
\text{Cov}_{f_{\bar{X}}}(\alpha, \gamma) &= \int_{\mathbb{R}} \int_{\mathbb{R}} (\bar{\alpha} - \alpha(\bar{x}_i, \bar{x}_j)) (\bar{\gamma} - \gamma(\bar{x}_i, \bar{x}_j)) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\
&= \frac{c^2 Nb}{2\bar{r}} \sqrt{\frac{b}{\pi}} \int_{\mathbb{R}} \int_{\mathbb{R}} (1 - b(\bar{x}_i - \bar{x}_j)^2) \sqrt{\frac{b}{\pi}} \exp(-b(\bar{x}_i - \bar{x}_j)^2) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j - \bar{\alpha} \bar{\gamma} \\
&= \frac{c^2 Nb}{2\bar{r}} \sqrt{\frac{b}{\pi}} \frac{1 - 2bV_{\bar{X}}}{\sqrt{2\pi((2b)^{-1} + 2V_{\bar{X}})}} - \frac{c^2 Nb}{\bar{r}} \frac{1 - 2bV_{\bar{X}}}{2\pi(b^{-1} + 2V_{\bar{X}})} \\
&= \frac{c^2 b^2 N}{2\pi \bar{r}} (1 - 2bV_{\bar{X}}) \left( \frac{1}{\sqrt{1 + 4bV_{\bar{X}}}} - \frac{1}{1 + 2bV_{\bar{X}}} \right). \quad (97)
\end{aligned}$$

### 211 10.1.2 Calculating $\text{Cov}_{f_{\bar{X}}}(\alpha, |\beta|)$

212 To calculate moments of  $|\beta|$  we note that, as a random variable,  $|\beta|$  takes a folded normal distribution.  
213 Setting  $\Phi(x)$  equal to the cumulative density function of the standard normal distribution and using the  
214 properties of the folded normal distribution, we find

$$|\bar{\beta}| = \sqrt{\frac{2\text{Var}_{f_{\bar{X}}}(\beta)}{\pi}} \exp\left(-\frac{\bar{\beta}^2}{2\text{Var}_{f_{\bar{X}}}(\beta)}\right) - \bar{\beta} \left[ 1 - 2\Phi\left(\frac{\bar{\beta}}{\sqrt{\text{Var}_{f_{\bar{X}}}(\beta)}}\right) \right] \quad (98)$$

$$\text{Var}_{f_{\bar{X}}}(|\beta|) = \bar{\beta}^2 + \text{Var}_{f_{\bar{X}}}(\beta) - |\bar{\beta}|^2. \quad (99)$$

215 Recall that

$$\beta(\bar{x}_i, \bar{x}_j) = cNb(\bar{x}_i - \bar{x}_j) \sqrt{\frac{b}{2\pi}} \exp\left(-\frac{b}{2}(\bar{x}_i - \bar{x}_j)^2\right) \quad (100)$$

216 and hence

$$\begin{aligned} \bar{\beta} &= \int_{\mathbb{R}} \int_{\mathbb{R}} \beta(\bar{x}_i, \bar{x}_j) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\ &= cNb \int_{\mathbb{R}} (\bar{x} - \bar{x}_j) \frac{1}{\sqrt{2\pi(b^{-1} + V_{\bar{X}})}} \exp\left(-\frac{(\bar{x} - \bar{x}_j)^2}{2(b^{-1} + V_{\bar{X}})}\right) f_{\bar{X}}(\bar{x}_j) d\bar{x}_j = 0, \end{aligned} \quad (101)$$

$$\begin{aligned} \text{Var}_{f_{\bar{X}}}(\beta) &= \int_{\mathbb{R}} \int_{\mathbb{R}} (\bar{\beta} - \beta(\bar{x}_i, \bar{x}_j))^2 f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\ &= \int_{\mathbb{R}} \int_{\mathbb{R}} c^2 N^2 b^2 (\bar{x}_i - \bar{x}_j)^2 \frac{b}{2\pi} \exp(-b(\bar{x}_i - \bar{x}_j)^2) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\ &= \sqrt{\frac{b}{4\pi}} c^2 N^2 b^2 \int_{\mathbb{R}} \left[ \left( \frac{\bar{x} + 2bV_{\bar{X}}\bar{x}_j}{1 + 2bV_{\bar{X}}} - \bar{x}_j \right)^2 + \frac{V_{\bar{X}}}{1 + 2bV_{\bar{X}}} \right] \frac{\exp\left(-\frac{(\bar{x} - \bar{x}_j)^2}{2(\frac{1}{2b} + V_{\bar{X}})}\right)}{\sqrt{2\pi(\frac{1}{2b} + V_{\bar{X}})}} f_{\bar{X}}(\bar{x}_j) d\bar{x}_j \\ &= \sqrt{\frac{b}{4\pi}} c^2 N^2 b^2 \int_{\mathbb{R}} \left[ \frac{(\bar{x} - \bar{x}_j)^2}{(1 + 2bV_{\bar{X}})^2} + \frac{V_{\bar{X}}}{1 + 2bV_{\bar{X}}} \right] \frac{\exp\left(-\frac{(\bar{x} - \bar{x}_j)^2}{2(\frac{1}{2b} + V_{\bar{X}})}\right)}{\sqrt{2\pi(\frac{1}{2b} + V_{\bar{X}})}} f_{\bar{X}}(\bar{x}_j) d\bar{x}_j \\ &= \sqrt{\frac{b}{4\pi}} c^2 N^2 b^2 \left[ \frac{(1 + 2bV_{\bar{X}})V_{\bar{X}}}{1 + 4bV_{\bar{X}}} \frac{1}{(1 + 2bV_{\bar{X}})^2} + \frac{V_{\bar{X}}}{1 + 2bV_{\bar{X}}} \right] \frac{1}{\sqrt{2\pi(\frac{1}{2b} + 2V_{\bar{X}})}} \\ &= \frac{b}{\pi} \frac{c^2 N^2 b^2}{\sqrt{1 + 4bV_{\bar{X}}}} \frac{V_{\bar{X}}}{1 + 2bV_{\bar{X}}} \left( \frac{1}{1 + 4bV_{\bar{X}}} + 1 \right) = \frac{2c^2 N^2 b^3 V_{\bar{X}}}{\pi(1 + 4bV_{\bar{X}})^{3/2}}. \end{aligned} \quad (102)$$

217 Thus, using properties of the folded normal distribution, we find

$$|\beta| = \sqrt{\frac{2}{\pi}} \frac{cNb^{3/2}}{(1 + 4bV_{\bar{X}})^{3/4}} \sqrt{\frac{2V_{\bar{X}}}{\pi}} = \frac{2}{\pi} \frac{cNb^{3/2}}{(1 + 4bV_{\bar{X}})^{3/4}} \sqrt{V_{\bar{X}}}, \quad (103)$$

$$\text{Var}_{f_{\bar{X}}}(|\beta|) = \frac{c^2 N^2 b^3}{(1 + 4bV_{\bar{X}})^{3/2}} \frac{2V_{\bar{X}}}{\pi} \left( 1 - \frac{2}{\pi} \right). \quad (104)$$

218 We also calculate

$$\begin{aligned} \text{Cov}_{f_{\bar{X}}}(\alpha, \beta) &= \int_{\mathbb{R}} \int_{\mathbb{R}} (\bar{\alpha} - \alpha(\bar{x}_i, \bar{x}_j)) (\bar{\beta} - \beta(\bar{x}_i, \bar{x}_j)) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\ &= \frac{c^2 Nb}{2\bar{r}} \sqrt{\frac{b}{\pi}} \int_{\mathbb{R}} \int_{\mathbb{R}} (\bar{x}_i - \bar{x}_j) \sqrt{\frac{b}{\pi}} \exp(-b(\bar{x}_i - \bar{x}_j)^2) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j = 0. \end{aligned} \quad (105)$$

219 In attempt to calculate  $\text{Cov}_{f_{\bar{X}}}(\alpha, |\beta|)$  we find



$$\begin{aligned}
\text{Cov} f_{\bar{X}}(\alpha, |\beta|) &= \int_{\mathbb{R}} \int_{\mathbb{R}} \alpha(\bar{x}_i, \bar{x}_j) |\beta(\bar{x}_i, \bar{x}_j)| f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j - \bar{\alpha} \overline{|\beta|} \\
&= \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{c}{\bar{r}} \sqrt{\frac{b}{2\pi}} \exp\left(-\frac{b}{2}(\bar{x}_i - \bar{x}_j)^2\right) cNb|\bar{x}_i - \bar{x}_j| \sqrt{\frac{b}{2\pi}} \exp\left(-\frac{b}{2}(\bar{x}_i - \bar{x}_j)^2\right) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j - \bar{\alpha} \overline{|\beta|} \\
&= \frac{c^2 Nb}{\bar{r}} \sqrt{\frac{b}{4\pi}} \int_{\mathbb{R}} \int_{\mathbb{R}} |\bar{x}_i - \bar{x}_j| \sqrt{\frac{b}{\pi}} \exp(-b(\bar{x}_i - \bar{x}_j)^2) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j - \bar{\alpha} \overline{|\beta|}. \quad (106)
\end{aligned}$$

220 Just as we used the folded normal to find  $\overline{|\beta|}$  and  $\text{Var}_{f_{\bar{X}}}(|\beta|)$ , we can calculate  $\text{Cov}_{f_{\bar{X}}}(\alpha, |\beta|)$  by considering

$$\int_{\mathbb{R}} \int_{\mathbb{R}} (\bar{x}_i - \bar{x}_j) \sqrt{\frac{b}{\pi}} \exp(-b(\bar{x}_i - \bar{x}_j)^2) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j = 0 \quad (107)$$

221 and

$$\begin{aligned}
&\int_{\mathbb{R}} \int_{\mathbb{R}} (\bar{x}_i - \bar{x}_j)^2 \frac{b}{\pi} \exp(-2b(\bar{x}_i - \bar{x}_j)^2) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\
&= \sqrt{\frac{2b}{\pi}} \int_{\mathbb{R}} \int_{\mathbb{R}} (\bar{x}_i - \bar{x}_j)^2 \frac{1}{\sqrt{2\pi \frac{1}{4b}}} \exp\left(-\frac{(\bar{x}_i - \bar{x}_j)^2}{2 \frac{1}{4b}}\right) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j \\
&= \sqrt{\frac{2b}{\pi}} \int_{\mathbb{R}} \left[ \left( \frac{\bar{x} + 4bV_{\bar{X}}\bar{x}_j}{1 + 4bV_{\bar{X}}} - \bar{x}_j \right)^2 + \frac{V_{\bar{X}}}{1 + 4bV_{\bar{X}}} \right] \frac{1}{\sqrt{2\pi(\frac{1}{4b} + V_{\bar{X}})}} \exp\left(-\frac{(\bar{x} - \bar{x}_j)^2}{2(\frac{1}{4b} + V_{\bar{X}})}\right) f_{\bar{X}}(\bar{x}_j) d\bar{x}_j \\
&= \sqrt{\frac{2b}{\pi}} \int_{\mathbb{R}} \left[ \left( \frac{\bar{x} - \bar{x}_j}{1 + 4bV_{\bar{X}}} \right)^2 + \frac{V_{\bar{X}}}{1 + 4bV_{\bar{X}}} \right] \frac{1}{\sqrt{2\pi(\frac{1}{4b} + V_{\bar{X}})}} \exp\left(-\frac{(\bar{x} - \bar{x}_j)^2}{2(\frac{1}{4b} + V_{\bar{X}})}\right) f_{\bar{X}}(\bar{x}_j) d\bar{x}_j \\
&= \sqrt{\frac{2b}{\pi}} \left[ \frac{(1 + 4bV_{\bar{X}})V_{\bar{X}}}{1 + 8bV_{\bar{X}}} \frac{1}{(1 + 4bV_{\bar{X}})^2} + \frac{V_{\bar{X}}}{1 + 4bV_{\bar{X}}} \right] \frac{1}{\sqrt{2\pi(\frac{1}{4b} + 2V_{\bar{X}})}} \\
&= \sqrt{\frac{2b}{\pi}} \frac{2V_{\bar{X}}}{1 + 8bV_{\bar{X}}} \sqrt{\frac{4b}{2\pi(1 + 8bV_{\bar{X}})}} = \frac{b}{\pi} \frac{4V_{\bar{X}}}{(1 + 8bV_{\bar{X}})^{3/2}}. \quad (108)
\end{aligned}$$

222 Hence

$$\int_{\mathbb{R}} \int_{\mathbb{R}} |\bar{x}_i - \bar{x}_j| \sqrt{\frac{b}{\pi}} \exp(-b(\bar{x}_i - \bar{x}_j)^2) f_{\bar{X}}(\bar{x}_i) f_{\bar{X}}(\bar{x}_j) d\bar{x}_i d\bar{x}_j = \sqrt{\frac{2}{\pi}} \sqrt{\frac{b}{\pi} \frac{4V_{\bar{X}}}{(1 + 8bV_{\bar{X}})^{3/2}}} = \frac{2}{\pi} \frac{\sqrt{2bV_{\bar{X}}}}{(1 + 8bV_{\bar{X}})^{3/4}} \quad (109)$$

223 and

$$\begin{aligned}
\text{Cov}_{f_{\bar{X}}}(\alpha, |\beta|) &= \frac{c^2 N b}{\bar{r}} \sqrt{\frac{b}{4\pi}} \frac{2}{\pi} \frac{\sqrt{2bV_{\bar{X}}}}{(1+8bV_{\bar{X}})^{3/4}} - \bar{\alpha} |\bar{\beta}| \\
&= \frac{2c^2 N b^2}{\pi \bar{r} (1+8bV_{\bar{X}})^{3/4}} \sqrt{\frac{V_{\bar{X}}}{2\pi}} - \frac{c}{\bar{r}} \sqrt{\frac{b}{2\pi(1+2bV_{\bar{X}})}} \frac{2}{\pi} \frac{c N b^{3/2}}{(1+4bV_{\bar{X}})^{3/4}} \sqrt{V_{\bar{X}}} \\
&= \frac{2c^2 N b^2}{\pi \bar{r} (1+8bV_{\bar{X}})^{3/4}} \sqrt{\frac{V_{\bar{X}}}{2\pi}} - \frac{2c^2 N b^2}{\pi \bar{r} (1+4bV_{\bar{X}})^{3/4}} \sqrt{\frac{V_{\bar{X}}}{2\pi(1+2bV_{\bar{X}})}} \\
&= \frac{2c^2 N b^2}{\pi \bar{r}} \sqrt{\frac{V_{\bar{X}}}{2\pi}} \left( \frac{1}{(1+8bV_{\bar{X}})^{3/4}} - \frac{1}{(1+4bV_{\bar{X}})^{3/4} (1+2bV_{\bar{X}})^{1/2}} \right). \quad (110)
\end{aligned}$$

### 224 10.1.3 Starting the calculation of $\text{Cov}_{f_{\bar{X}}}(\alpha, \mathfrak{C})$

225 We have

$$\mathfrak{C}(\bar{x}_i, \bar{x}_j) = c^2 N^2 b^2 \left( |\bar{x}_i - \bar{x}_j| + |1 - b(\bar{x}_i - \bar{x}_j)^2| \right)^2 \exp \left( -\frac{b}{2} (\bar{x}_i - \bar{x}_j)^2 \right). \quad (111)$$

226 Note that the random variable  $\delta = \bar{x}_i - \bar{x}_j$  is a mean zero Gaussian random variable with variance  $2V_{\bar{X}}$ . We  
227 write the probability density function of  $\delta$  as  $f_{\Delta}(\delta)$ . Substituting in  $\delta$ , we can write

$$\begin{aligned}
\mathfrak{C}(\delta, 0) &= c^2 N^2 b^2 \left( |\delta| + |1 - b\delta^2| \right)^2 \exp \left( -\frac{b}{2} \delta^2 \right) \\
&= c^2 N^2 b^2 \left( \delta^2 + 2||\delta| - b|\delta|^3| + (1 - b\delta^2)^2 \right) \exp \left( -\frac{b}{2} \delta^2 \right). \quad (112)
\end{aligned}$$

228 From this expression, we see properties of the folded normal distribution can be used to calculate several  
229 components of the integral  $\text{Cov}_{f_{\bar{X}}}(\alpha, \mathfrak{C})$ , but a major technical challenge lies in calculating

$$\int_{\mathbb{R}} ||\delta| - b|\delta|^3| \exp \left( -\frac{b}{2} \delta^2 \right) f_{\Delta}(\delta) d\delta. \quad (113)$$

230 Instead of overcoming this challenge to find an analytical form of  $\text{Cov}_{f_{\bar{X}}}(\alpha, \mathfrak{C})$  we turn to a numerical  
231 approach outlined in the following section.

## 232 10.2 Numerical estimates for heterogeneous $N$ and $G$

233 To verify our analytical estimates for the covariance of selection gradients and competition coefficients,  
234 equations (60) of the main text, and for approximating the correlation of competition coefficients and a  
235 metric of pairwise coevolution, Figure 5 of the main text, we developed a numerical approach by repeatedly  
236 simulating our system (52) of the main text across a set of randomly drawn model parameters. In particular,  
237 we ran each numerical integration with  $S = 100$  species for  $T_1 = 1000.0$  units of time. We then continued  
238 to integrate for  $T_2 = 1000.0$  units of time and calculated the covariances between competition coefficients  
239 and the quantities  $|\beta|$ ,  $\gamma$  and  $\mathfrak{C}$  for each of the last  $T_2$  time steps. We assume the temporal average of these  
240 covariances across the last  $T_2$  units of time approximates the expectation at equilibrium. We repeated this  
241 approach for randomly drawn  $a$  and  $c$  until our sample size reached 1000. We drew values for  $a$  and  $c$  from  
242 condition log-normal distributions. For  $a$ , we set the log-normal parameters equal to  $\mu = -10$  and  $\sigma = 6$   
243 and conditioned on  $a < 0.01$  to prevent extinction. For  $c$  we set the log-normal parameters equal to  $\mu = -4$

and  $\sigma = 4$  and conditioned on  $c < 1 \times 10^{-4}$  to prevent very large interspecific phenotypic variation and extinction. When extinction occurs, we aborted the simulation and restarted with newly drawn parameters. Results for  $\text{Corr}_{f_{\bar{X}}}(\alpha, |\beta|)$  and  $\text{Corr}_{f_{\bar{X}}}(\alpha, \gamma)$  are displayed in SM Figure 1.

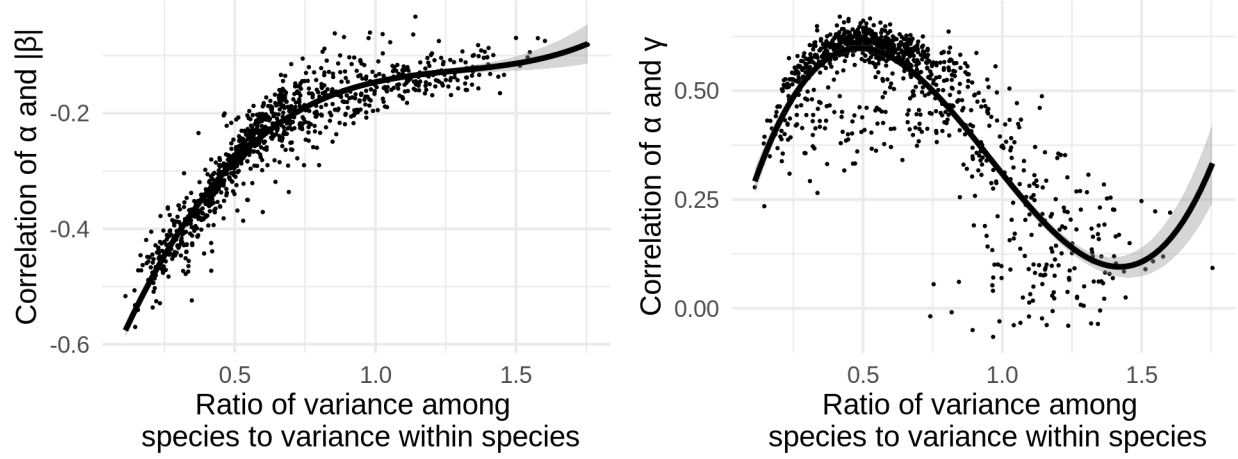


Figure 1: Numerical estimate for the correlations of selection gradients and competition coefficients.

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