

I. THE MODEL

We consider TASEP on a ring of L sites, labeled clockwise by the index $i = 1, 2, \dots, L$, where site 1 is the nearest-neighbor of site L in the clockwise direction. Each site of the lattice can be empty or occupied by just one particle. For convenience of notation we represent each configuration by a string $(n(1), n(2), \dots, n(L))$ of L occupation numbers $n(i) \in \{0, 1\}$ of the sites $i = 1, 2, \dots, L$, in increasing order from the left to the right.

The dynamics of the model corresponds to the discrete-time backward-ordered update with modified probabilities according to Ref. [1]. We fix that the update always starts from the bond $(L-1, L)$. A particle can hop to a vacant nearest-neighbor site in the clockwise direction, i.e. from site i to site $i+1$, or stay at its place. During each discrete moment of time t , an update of the configuration of the whole system takes place in 4 consecutive steps, passing through successive updates of all the pairs of nearest-neighbor sites in the counterclockwise order $(L-1, L), (L-2, L-1), \dots, (L, 1)$. The probability of a hop along a bond $(i, i+1)$ depends on whether a particle has jumped from site $i+1$ to site $i+2$ in the previous step, when the bond $(i+1, i+2)$ was updated within the same time moment, or not.

(1) In the case when site $i+1$ has not changed its occupation number, the probabilities are the standard ones: if site $i+1$ remains empty since the previous step, within the same moment of time, then the jump of a particle from an occupied site i to site $i+1$ takes place with probability p , and the particle stays immobile with probability $1-p$; if site $i+1$ remains occupied, no jump takes place and the configuration of the bond $(i, i+1)$ is conserved.

(2) If in the previous step, within the same moment of time, a particle has jumped from site $i+1$ to site $i+2$, thus leaving $i+1$ empty, then the jump of a particle from an occupied site i to site $i+1$ in the next step takes place with a different probability \tilde{p} , and the particle stays immobile with probability $1-\tilde{p}$.

Note that when $\tilde{p} = p$ one has the standard TASEP with backward-sequential update, and when $\tilde{p} = 0$ one has the TASEP with parallel update.

II. THE STATIONARY PROBABILITY DISTRIBUTION FOR THE MINIMAL MODEL

Here we consider the minimal nontrivial case of two particles, $N = 2$, on a ring of 4 sites, $L = 4$. Let us label all the 6 possible configurations as follows:

$$\begin{aligned} (1, 1, 0, 0) &= C_1, & (0, 1, 1, 0) &= C_2, & (0, 0, 1, 1) &= C_3, & (1, 0, 0, 1) &= C_4 \\ (1, 0, 1, 0) &= C_5, & (0, 1, 0, 1) &= C_6. \end{aligned} \quad (1)$$

Within every time moment, each of the above configurations generates a chain of possible configurations with probabilities following from the adopted update rules:

$$C_1 \rightarrow (1 - p)C_1 + p\tilde{p} C_2 + p(1 - \tilde{p})C_5, \quad (2)$$

$$C_2 \rightarrow p^2(1 - \tilde{p})C_1 + (1 - p)C_2 + p\tilde{p}(1 - p)C_3 + p^2\tilde{p} C_5 + p(1 - p)(1 - \tilde{p})C_6, \quad (3)$$

$$C_3 \rightarrow (1 - p)C_3 + p C_5, \quad (4)$$

$$C_4 \rightarrow p\tilde{p} C_1 + (1 - p)C_4 + p(1 - \tilde{p})C_6, \quad (5)$$

$$C_5 \rightarrow p^2\tilde{p} C_1 + p(1 - p)C_2 + p(1 - p)C_4 + (1 - p)^2C_5 + p^2(1 - \tilde{p})C_6, \quad (6)$$

$$C_6 \rightarrow p(1 - p)C_1 + p(1 - p)C_3 + p^2C_5 + (1 - p)^2C_6. \quad (7)$$

Let P_i , $i = 1, 2, \dots, 6$, be the probability of configuration C_i before the update, and let P'_i , $i = 1, 2, \dots, 6$, be the corresponding probability after the update. For the latter we have:

$$P'_1 = (1 - p)P_1 + p^2(1 - \tilde{p})P_2 + p\tilde{p} P_4 + p^2\tilde{p} P_5 + p(1 - p)P_6, \quad (8)$$

$$P'_2 = p\tilde{p} P_1 + (1 - p)P_2 + p(1 - p)P_5, \quad (9)$$

$$P'_3 = p\tilde{p}(1 - p)P_2 + (1 - p)P_3 + p(1 - p)P_6, \quad (10)$$

$$P'_4 = (1 - p)P_4 + p(1 - p)P_5, \quad (11)$$

$$P'_5 = p(1 - \tilde{p})P_1 + p^2\tilde{p} P_2 + p P_3 + (1 - p)^2P_5 + p^2P_6, \quad (12)$$

$$P'_6 = p(1 - p)(1 - \tilde{p})P_2 + p(1 - \tilde{p})P_4 + p^2(1 - \tilde{p})P_5 + (1 - p)^2P_6. \quad (13)$$

The stationarity conditions $P'_i = P_i$, $i = 1, 2, \dots, 6$, lead to a system of 6 linear equations, 5 of which are independent. All the probabilities depend linearly on P_2 , P_5 and P_6 . By assuming P_5 to be the independent variable, which is to be found from the probability normalization condition, we obtain the following equations for P_2 and P_5 :

$$\begin{aligned} [1 - p\tilde{p}(1 - \tilde{p})]P_2 - \tilde{p}(1 - p)P_6 &= (1 - p + \tilde{p}^2)P_5 \\ (1 - p)(1 - \tilde{p})P_2 - (2 - p)P_6 &= -(1 - \tilde{p})P_5. \end{aligned} \quad (14)$$

By solving the above equation and using the expressions for P_1 , P_3 and P_4 in terms of P_2 , P_5 and P_6 , we obtain

$$P_1 = \frac{2(1-p) + p\tilde{p}}{2-p-\tilde{p}(1-\tilde{p})} P_5, \quad (15)$$

$$P_2 = \frac{(1-p)(2-p) + \tilde{p}(1-p+\tilde{p})}{2-p-\tilde{p}(1-\tilde{p})} P_5, \quad (16)$$

$$P_3 = \frac{(1-p)[2-2p+p^2+2\tilde{p}(\tilde{p}-p)]}{2-p-\tilde{p}(1-\tilde{p})} P_5, \quad (17)$$

$$P_4 = (1-p)P_5, \quad (18)$$

$$P_6 = \frac{(1-\tilde{p})[2-2p+p^2+\tilde{p}(\tilde{p}-p)]}{2-p-\tilde{p}(1-\tilde{p})} P_5. \quad (19)$$

Finally, P_5 is obtained from the condition $\sum_{i=1}^6 P_i = 1$ which yields

$$P_5 = \frac{2-p-\tilde{p}(1-\tilde{p})}{(2-p)^3 + 4 - 3p - \tilde{p}[3(1-\tilde{p})^2 + 2\tilde{p}(p-\tilde{p}) - p^2]}. \quad (20)$$

For the current $J_{i,i+1}$ of particles through the different bonds $(i, i+1)$ we have the general expressions

$$J_{1,2} = p\tilde{p}P_1 + pP_4 + pP_5, \quad (21)$$

$$J_{2,3} = pP_1 + p\tilde{p}P_2 + pP_6, \quad (22)$$

$$J_{3,4} = pP_2 + pP_5, \quad (23)$$

$$J_{4,1} = p^2P_2 + pP_3 + p\tilde{p}P_4 + p^2\tilde{p}P_5 + pP_6. \quad (24)$$

As expected, in the stationary state all these currents are equal:

$$J_{i,i+1} = p \frac{(2-p)^2 - \tilde{p}(p-\tilde{p}) + \tilde{p}^2}{2-p-\tilde{p}(1-\tilde{p})} P_5. \quad (25)$$

There are two special cases in which the above results radically simplify. Notably, in these cases the generalized TASEP rules reduce to usual ones: $\tilde{p} = p$ corresponds to the ordinary backward-ordered update, and $\tilde{p} = 0$ — to the parallel update [2].

A. The case of ordinary backward-ordered update

At $\tilde{p} = p$ the stationary distribution (19) reduces to

$$P_1 = P_2 = P_5 = \frac{1}{3(2-p)}, \quad P_3 = P_4 = P_6 = \frac{1-p}{3(2-p)}, \quad (26)$$

and the current through each bond becomes

$$J_{i,i+1} = \frac{2p}{3(2-p)}. \quad (27)$$

B. The case of ordinary parallel update

At $\tilde{p} = 0$ the stationary distribution (19) reduces to

$$\begin{aligned} P_1 &= \frac{2(1-p)}{(2-p)^3 + 4 - 3p}, & P_2 = P_4 &= \frac{(1-p)(2-p)}{(2-p)^3 + 4 - 3p}, \\ P_3 &= \frac{(1-p)(2-2p+p^2)}{(2-p)^3 + 4 - 3p}, & P_5 &= \frac{2-p}{(2-p)^3 + 4 - 3p}, & P_6 &= \frac{2-2p+p^2}{(2-p)^3 + 4 - 3p}, \end{aligned} \quad (28)$$

and the current through each bond becomes

$$J_{i,i+1} = \frac{p(2-p)^3}{(2-p)^3 + 4 - 3p}. \quad (29)$$

It is known that the stationary state of TASEP with parallel update is characterized by the presence of particle-hole correlations. Indeed, from (28) it is evident that $P_5 > P_6 > P_1 > P_2 = P_4 > P_3$.

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 - [2] N. Rajewsky, L. Santen, A. Schadschneider, and M. Schreckenberg, J. Stat. Phys. **92**, 151, 1998.