

Classical Mechanics by John R. Taylor Problems

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1 Newton's Laws of Motion

1.1

$$\begin{aligned}\mathbf{b} + \mathbf{c} &= 2\hat{\mathbf{x}} + \hat{\mathbf{y}} + \hat{\mathbf{z}} \\ 5\mathbf{b} + 2\mathbf{c} &= 7\hat{\mathbf{x}} + 5\hat{\mathbf{y}} + 2\hat{\mathbf{z}} \\ \mathbf{b} \cdot \mathbf{c} &= 1 \\ \mathbf{b} \times \mathbf{c} &= \begin{vmatrix} \hat{\mathbf{x}} & \hat{\mathbf{y}} & \hat{\mathbf{z}} \\ 1 & 1 & 0 \\ 1 & 0 & 1 \end{vmatrix} \\ &= \hat{\mathbf{x}} - \hat{\mathbf{y}} - \hat{\mathbf{z}}\end{aligned}$$

1.5

$$\begin{aligned}\mathbf{v}_{\text{body}} &= \hat{\mathbf{x}} + \hat{\mathbf{y}} + \hat{\mathbf{z}} \\ \mathbf{v}_{\text{face}} &= \hat{\mathbf{x}} + \hat{\mathbf{z}} \\ \mathbf{v}_{\text{body}} \cdot \mathbf{v}_{\text{face}} &= v_{\text{body}} v_{\text{face}} \cos \theta \\ 2 &= \sqrt{6} \cos \theta \\ \cos \theta &= \frac{2}{\sqrt{6}} \\ \theta &= \arccos \frac{2}{\sqrt{6}} \\ &= 35.26^\circ\end{aligned}$$

1.11

The particle moves counterclockwise in an ellipse of width $2b$ and height $2c$. The angular speed is ω .

1.23

$$\begin{aligned}\mathbf{v} &= v \cos \theta \frac{\mathbf{b}}{b} - v \sin \theta \frac{\mathbf{b} \times \mathbf{c}}{bc} \\ &= \frac{\lambda \mathbf{b}}{b} - \frac{c \mathbf{b} \times \mathbf{c}}{bc} \\ &= \frac{\lambda \mathbf{b} - \mathbf{b} \times \mathbf{c}}{b^2}\end{aligned}$$

1.25

$$\begin{aligned}\frac{df}{dt} &= -3f \\ \frac{1}{f} \frac{df}{dt} &= -3 \\ \ln f &= -3t + c \\ f &= ce^{-3t}\end{aligned}$$

One constant.

1.35

$$\begin{aligned}
 F_x &= 0 \\
 ma_x &= 0 \\
 a_x &= 0 \\
 v_x &= c_1 \\
 &= v_o \cos \theta \\
 r_x &= v_o \cos(\theta)t + c_2 \\
 &= v_o \cos(\theta)t
 \end{aligned}$$

$$\begin{aligned}
 F_y &= 0 \\
 ma_y &= 0 \\
 a_y &= 0 \\
 v_y &= c_3 \\
 v_y &= 0 \\
 r_y &= c_4 \\
 r_y &= 0
 \end{aligned}$$

$$\begin{aligned}
 F_z &= -mg \\
 ma_z &= -mg \\
 a_z &= -g \\
 v_z &= -gt + c_5 \\
 &= v_o \sin \theta - gt \\
 r_z &= v_o \sin(\theta)t - \frac{1}{2}gt^2 + c_6 \\
 &= v_o \sin(\theta)t - \frac{1}{2}gt^2
 \end{aligned}$$

$$\begin{aligned}
 0 &= v_o \sin(\theta)t - \frac{1}{2}gt^2 \\
 t &= \frac{2 \sin(\theta)v_o}{g} \\
 r_x &= v_o \cos(\theta)t \\
 &= \frac{2 \cos(\theta) \sin(\theta)v_o^2}{g} \\
 &= \frac{\sin(2\theta)v_o^2}{g}
 \end{aligned}$$

1.37

(a)

$$\begin{aligned}F &= -mg \sin \theta \\ma &= -mg \sin \theta \\a &= -g \sin \theta \\v &= c_1 - gt \sin \theta \\&= v_o - gt \sin \theta \\x &= v_o t - \frac{1}{2}gt^2 \sin \theta\end{aligned}$$

(b)

$$t = \frac{2v_o}{g \sin \theta}$$

1.39

$$\begin{aligned}
 F_x &= -mg \sin \phi \\
 ma_x &= -mg \sin \phi \\
 a_x &= -g \sin \phi \\
 v_x &= c_1 - gt \sin \phi \\
 &= v_o \cos \theta - gt \sin \phi \\
 r_x &= v_o t \cos \theta - \frac{1}{2}gt^2 \sin \phi + c_2 \\
 &= v_o t \cos \theta - \frac{1}{2}gt^2 \sin \phi
 \end{aligned}$$

$$\begin{aligned}
 F_y &= -mg \cos \phi \\
 ma_y &= -mg \cos \phi \\
 a_y &= -g \cos \phi \\
 v_y &= c_3 - gt \cos \phi \\
 &= v_o \sin \theta - gt \cos \phi \\
 r_y &= v_o t \sin \theta - \frac{1}{2}gt^2 \cos \phi + c_4 \\
 &= v_o t \sin \theta - \frac{1}{2}gt^2 \cos \phi
 \end{aligned}$$

$$\begin{aligned}
 0 &= v_o t \sin \theta - \frac{1}{2}gt^2 \cos \phi \\
 t &= \frac{2v_o \sin \theta}{g \cos \phi}
 \end{aligned}$$

$$\begin{aligned}
 r_x &= \frac{2v_o^2 \cos \theta \sec \phi \sin \theta}{g} - \frac{2v_o^2 \sec \phi \sin^2 \theta \tan \phi}{g} \\
 &= \frac{2v_o^2 \sin \theta (\cos \theta \cos \phi - \sin \theta \sin \phi)}{g \cos^2 \phi} \\
 &= \frac{2v_o^2 \sin \theta \cos(\theta + \phi)}{g \cos^2 \phi}
 \end{aligned}$$

$$\begin{aligned}
\frac{dr_x}{d\theta} &= \frac{2v_o^2}{g \cos^2 \phi} [\cos \theta \cos(\theta + \phi) - \sin \theta \sin(\theta + \phi)] \\
&= \frac{2v_o^2 \cos(2\theta + \phi)}{g \cos^2 \phi} \\
0 &= \frac{2v_o^2 \cos(2\theta + \phi)}{g \cos^2 \phi} \\
&= \cos(2\theta + \phi) \\
2\theta + \phi &= \frac{\pi}{2} \\
\theta &= \frac{\pi}{4} - \frac{\phi}{2} \\
r_{x,\max} &= \frac{2v_o^2 \sin\left(\frac{\pi}{4} - \frac{\phi}{2}\right) \cos\left(\frac{\pi}{4} - \frac{\phi}{2} + \phi\right)}{g \cos^2 \phi} \\
&= \frac{v_o^2(1 - \sin \phi)}{g \cos^2 \phi} \\
&= \frac{v_o^2}{g(1 + \sin \phi)}
\end{aligned}$$

1.41

$$\begin{aligned}
F &= ma \\
T &= m \frac{v^2}{R} \\
&= m \frac{(\omega R)^2}{R} \\
&= m\omega^2 R
\end{aligned}$$

1.47

(a)

$$\begin{aligned}
\rho &= \sqrt{x^2 + y^2} \\
\phi &= \arctan \frac{y}{x} \\
z &= z
\end{aligned}$$

ρ is the distance of P from the z -axis.

The use of r may be unfortunate because it suggests it's the distance of P from the origin.

- (b) $\hat{\rho}$ points away from the z -axis, $\hat{\phi}$ points counter-clockwise around the z -axis, and \hat{z} points in the positive z direction.

$$\mathbf{r} = \rho\hat{\rho} + z\hat{z} + \sqrt{x^2 + y^2}\hat{\rho} + z\hat{z}$$

(c)

$$\begin{aligned}\mathbf{v} &= \frac{d\mathbf{r}}{dt} \\ &= \dot{\rho}\hat{\rho} + \rho\frac{d\hat{\rho}}{dt} + \dot{z}\hat{z} + z\frac{d\hat{z}}{dt} \\ &= \dot{\rho}\hat{\rho} + \rho\dot{\phi}\hat{\phi} + \dot{z}\hat{z} \\ \mathbf{a} &= \frac{d\mathbf{v}}{dt} \\ &= \ddot{\rho}\hat{\rho} + \dot{\rho}\frac{d\hat{\rho}}{dt} + \dot{\rho}\dot{\phi}\hat{\phi} + \rho\ddot{\phi}\hat{\phi} + \rho\dot{\phi}\frac{d\hat{\phi}}{dt} + \ddot{z}\hat{z} \\ &= \ddot{\rho}\hat{\rho} + \dot{\rho}\dot{\phi}\hat{\phi} + \dot{\rho}\dot{\phi}\hat{\phi} + \rho\ddot{\phi}\hat{\phi} - \rho\dot{\phi}^2\hat{\rho} + \ddot{z}\hat{z} \\ &= (\ddot{\rho} - \rho\dot{\phi}^2)\hat{\rho} + (2\dot{\rho}\dot{\phi} + \rho\ddot{\phi})\hat{\phi} + \ddot{z}\hat{z}\end{aligned}$$

2 Projectiles and Charged Particles

2.1

$$\begin{aligned}1 &= (1.6 \times 10^3)Dv \\ v &= \frac{1}{(1.6 \times 10^3)D} \\ &= 8.9 \text{ mm/s}\end{aligned}$$

When $v \gg 1 \text{ cm/s}$ the drag force can be treated as purely quadratic. For a beach ball this becomes $v \gg 1 \text{ mm/s}$.

2.3

(a)

$$\begin{aligned}\frac{f_{\text{quad}}}{f_{\text{lin}}} &= \frac{(1/4)\rho Av^2}{3\pi\eta Dv} \\ &= \frac{\rho\pi\left(\frac{D}{2}\right)^2 v}{12\pi\eta D} \\ &= \frac{\rho Dv}{48\eta} \\ &= \frac{R}{48}\end{aligned}$$

(b)

$$R = \frac{Dv\rho}{\eta} \approx 0.01$$

2.5

$$\begin{aligned} v_y(t) &= v_{\text{ter}} + (v_{y_0} - v_{\text{ter}})e^{-t/\tau} \\ &= v_{\text{ter}} + (2v_{\text{ter}} - v_{\text{ter}})e^{-t/\tau} \\ &= v_{\text{ter}}(1 + e^{-t/\tau}) \end{aligned}$$

The velocity starts at $2v_{\text{ter}}$ and asymptotically approaches v_{ter} .

2.7

$$\begin{aligned} F &= F(v) \\ m\dot{v} &= F(v) \\ m \frac{dv}{F(v)} &= dt \\ t &= \int_{v_0}^v m \frac{dv'}{F(v')} \end{aligned}$$

$$\begin{aligned} F &= F(v) \\ m\dot{v} &= F_0 \\ v &= \frac{F_0}{m}t + c \end{aligned}$$

2.11

(a)

$$m\dot{v} = -mg - bv$$

$$\dot{v} = -g - kv$$

$$\frac{1}{-g - kv}\dot{v} = 1$$

$$-\frac{1}{k}\ln(-g - kv) = t + c$$

$$\ln(-g - kv) = c - \frac{t}{\tau}$$

$$-g - kv = Ae^{-t/\tau}$$

$$v = \tau(-g - Ae^{-t/\tau})$$

$$= -v_{\text{ter}} - \tau Ae^{-t/\tau}$$

$$v_o = -v_{\text{ter}} - \tau A$$

$$A = -k(v_o + v_{\text{ter}})$$

$$v = -v_{\text{ter}} + (v_o + v_{\text{ter}})e^{-t/\tau}$$

$$y = -v_{\text{ter}}t - \tau(v_o + v_{\text{ter}})e^{-t/\tau} + c$$

$$0 = -\tau(v_o + v_{\text{ter}}) + c$$

$$c = \tau(v_o + v_{\text{ter}})$$

$$y = -v_{\text{ter}}t + \tau(v_o + v_{\text{ter}})(1 - e^{-t/\tau})$$

(b)

$$\begin{aligned}
0 &= -v_{\text{ter}} + (v_o + v_{\text{ter}})e^{-t/\tau} \\
e^{-t/\tau} &= \frac{v_{\text{ter}}}{v_o + v_{\text{ter}}} \\
-\frac{t}{\tau} &= \ln \frac{v_{\text{ter}}}{v_o + v_{\text{ter}}} \\
t &= -\tau \ln \frac{v_{\text{ter}}}{v_o + v_{\text{ter}}} \\
\\
y_{\text{max}} &= -v_{\text{ter}} \left(-\tau \ln \frac{v_{\text{ter}}}{v_o + v_{\text{ter}}} \right) + \tau(v_o + v_{\text{ter}}) \left(1 - \frac{v_{\text{ter}}}{v_o + v_{\text{ter}}} \right) \\
&= \tau v_{\text{ter}} \ln \frac{v_{\text{ter}}}{v_o + v_{\text{ter}}} + \tau(v_o + v_{\text{ter}} - v_{\text{ter}}) \\
&= \tau \left(v_o + v_{\text{ter}} \ln \frac{v_{\text{ter}}}{v_o + v_{\text{ter}}} \right) \\
&= \tau \left[v_o - v_{\text{ter}} \ln \left(1 + \frac{v_o}{v_{\text{ter}}} \right) \right]
\end{aligned}$$

(c)

$$\begin{aligned}
y_{\text{max}} &= \tau \left[v_o - v_{\text{ter}} \ln \left(1 + \frac{v_o}{v_{\text{ter}}} \right) \right] \\
&= \tau \left[v_o - g\tau \ln \left(1 + \frac{v_o}{g\tau} \right) \right] \\
&\approx \tau \left\{ v_o - g\tau \left[\frac{v_o}{g\tau} - \frac{1}{2} \left(\frac{v_o}{g\tau} \right)^2 \right] \right\} \\
&= \tau \left(v_o - v_o + \frac{1}{2} \frac{v_o^2}{g\tau} \right) \\
&= \frac{1}{2} \frac{v_o^2}{g}
\end{aligned}$$

2.13

$$\begin{aligned}
 v^2 &= \frac{2}{m} \int_{x_0}^x -kx' dx' \\
 &= -\frac{2k}{m} \left(\frac{1}{2}x^2 - \frac{1}{2}x_0^2 \right) \\
 &= -\frac{k}{m}(x^2 - x_0^2) \\
 v &= \sqrt{\frac{k}{m}(x_0^2 - x^2)} \\
 &= \omega \sqrt{x_0^2 - x^2}
 \end{aligned}$$

$$\begin{aligned}
 \int_{x_0}^x \frac{1}{\sqrt{x_0^2 - x'^2}} dx' &= \int_0^t \omega dt \\
 \arctan \frac{x}{\sqrt{x_0^2 - x^2}} - \arctan \frac{x_0}{\sqrt{x_0^2 - x_0^2}} &= \omega t \\
 \arctan \frac{x}{\sqrt{x_0^2 - x^2}} &= \omega t + \frac{\pi}{2} \\
 \frac{x}{\sqrt{x_0^2 - x^2}} &= \tan \left(\omega t + \frac{\pi}{2} \right) \\
 &= -\cot \omega t \\
 \frac{\sqrt{x_0^2 - x^2}}{x} &= -\tan \omega t \\
 \sqrt{x_0^2 - x^2} &= -x \tan \omega t \\
 x_0^2 - x^2 &= x^2 \tan^2 \omega t \\
 x^2 &= \frac{x_0^2}{1 + \tan^2 \omega t} \\
 &= \frac{x_0^2 \cos^2 \omega t}{\cos^2 \omega t + \sin^2 \omega t} \\
 &= x_0^2 \cos^2 \omega t \\
 x &= x_0 \cos \omega t
 \end{aligned}$$

2.15

$$\begin{aligned}
a_y &= -g \\
v_y &= v_{y0} - gt \\
y &= v_{y0}t - \frac{1}{2}gt^2 \\
0 &= v_{y0}t - \frac{1}{2}gt^2 \\
t &= \frac{2v_{y0}}{g} \\
x &= v_{x0}t \\
R &= \frac{2v_{x0}v_{y0}}{g}
\end{aligned}$$

2.19

(a)

$$\begin{aligned}
x &= v_{x0}t \\
y &= v_{y0}t - \frac{1}{2}gt^2 \\
&= \frac{v_{y0}}{v_{x0}}x - \frac{1}{2}g\left(\frac{x}{v_{x0}}\right)^2
\end{aligned}$$

(b)

$$\begin{aligned}
y &= \frac{v_{y0} + v_{\text{ter}}}{v_{x0}}x + v_{\text{ter}}\tau \ln\left(1 - \frac{x}{v_{x0}\tau}\right) \\
&\approx \frac{v_{y0}}{v_{x0}}x + \frac{g\tau}{v_{x0}}x - g\tau^2\left[\frac{x}{v_{x0}\tau} + \frac{1}{2}\left(\frac{x}{v_{x0}\tau}\right)^2\right] \\
&= \frac{v_{y0}}{v_{x0}}x - \frac{1}{2}g\left(\frac{x}{v_{x0}}\right)^2
\end{aligned}$$