

Introduction to Electrodynamics by David J.
Griffiths Problems

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December 2023

Contents

2	Electrostatics	3
2.1	3
2.2	3
2.3	3
2.4	3
2.5	4
2.6	4
2.7	4
2.8	4
2.9	5
2.10	5
2.11	5
2.12	6
2.13	6
2.14	6
2.15	6
2.16	7
2.17	7
2.18	8
2.20	8
2.21	9
2.22	9
2.23	10
2.24	10
2.25	10
2.26	11
2.28	11
2.31	12
2.32	12
2.33	13

2.34	13
2.36	15
2.37	17
2.38	17
2.39	18
2.40	18
2.41	18
2.42	19
2.43	20
2.44	20
2.46	21
2.47	21
2.49	22
2.50	22
2.51	23
2.52	23
2.53	23
2.55	24
2.56	24

3 Potentials	25
3.1	25
3.2	25
3.3	26
3.7	26
3.8	26
3.9	27
3.10	28
3.11	28
3.12	29
3.13	30
3.14	30
3.15	31
3.16	33
3.17	35
3.18	36
3.19	37
3.20	37
3.21	37
3.22	38
3.23	39
3.24	40
3.25	41

2 Electrostatics

2.1

- (a) $\mathbf{0}$
- (b) The same as if only the opposite charge were present — all others are cancelled out.

2.2

$$\begin{aligned}\mathbf{E} &= \frac{1}{4\pi\epsilon_0} 2 \frac{q}{z^2} \cos\theta \hat{\mathbf{x}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{dq}{[(d/2)^2 + z^2]^{3/2}} \hat{\mathbf{x}}\end{aligned}$$

2.3

$$\begin{aligned}\mathbf{r} &= z\hat{\mathbf{z}} \\ \mathbf{r}' &= x\hat{\mathbf{x}} \\ \boldsymbol{\rho} &= z\hat{\mathbf{z}} - x\hat{\mathbf{x}} \\ \rho &= \sqrt{x^2 + z^2} \\ \hat{\boldsymbol{\rho}} &= \frac{z\hat{\mathbf{z}} - x\hat{\mathbf{x}}}{\sqrt{x^2 + z^2}} \\ \mathbf{E} &= \frac{1}{4\pi\epsilon_0} \int_0^L \frac{\lambda}{x^2 + z^2} \frac{z\hat{\mathbf{z}} - x\hat{\mathbf{x}}}{\sqrt{x^2 + z^2}} dx \\ &= \frac{1}{4\pi\epsilon_0} \lambda \left(z\hat{\mathbf{z}} \int_0^L \frac{1}{(x^2 + z^2)^{3/2}} dx - \hat{\mathbf{x}} \int_0^L \frac{x}{(x^2 + z^2)} dx \right) \\ &= \frac{1}{4\pi\epsilon_0} \lambda \left[\frac{L}{z\sqrt{L^2 + z^2}} \hat{\mathbf{z}} - \left(\frac{1}{z} - \frac{1}{\sqrt{L^2 + z^2}} \right) \hat{\mathbf{x}} \right] \\ &= \frac{1}{4\pi\epsilon_0} \frac{\lambda}{z} \left[\left(-1 + \frac{z}{\sqrt{L^2 + z^2}} \right) \hat{\mathbf{x}} + \frac{L}{\sqrt{L^2 + z^2}} \hat{\mathbf{z}} \right]\end{aligned}$$

2.4

The electric field a distance z above the midpoint of a line segment of length $2L$ and uniform line charge λ is

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{2\lambda L}{z\sqrt{z^2 + L^2}} \hat{\mathbf{z}}.$$

Applying this to the four sides of the square, the horizontal components of opposite sides cancel leaving only the vertical component.

$$\begin{aligned}
\cos \theta &= \frac{z}{\hat{z}} \\
&= \frac{z}{\sqrt{(a/2)^2 + z^2}} \\
\mathbf{E} &= 4 \left(\frac{1}{4\pi\epsilon_0} \frac{\lambda a}{\sqrt{(a/2)^2 + z^2} \sqrt{(a/2)^2 + (a/2)^2 + z^2}} \hat{\mathbf{z}} \right) \cos \theta \\
&= \frac{1}{4\pi\epsilon_0} \frac{4a\lambda z}{[(a/2)^2 + z^2] \sqrt{(a/2)^2 + z^2}} \hat{\mathbf{z}}
\end{aligned}$$

2.5

$$\begin{aligned}
\mathbf{E} &= \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \frac{\lambda r}{r^2 + z^2} \cos \alpha \, d\theta \, \hat{\mathbf{z}} \\
&= \frac{1}{4\pi\epsilon_0} \frac{2\pi\lambda r z}{(r^2 + z^2)^{3/2}} \hat{\mathbf{z}}
\end{aligned}$$

2.6

$$\begin{aligned}
\mathbf{E} &= \frac{1}{4\pi\epsilon_0} \int \frac{dq}{\hat{z}^2} \cos \theta \hat{\mathbf{z}} \\
&= \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^R \frac{\sigma}{r^2 + z^2} \frac{z}{\sqrt{r^2 + z^2}} r \, dr \, d\theta \hat{\mathbf{z}} \\
&= \frac{1}{4\pi\epsilon_0} 2\pi\sigma z \int_0^R \frac{r}{(r^2 + z^2)^{3/2}} \, dr \hat{\mathbf{z}} \\
&= \frac{1}{4\pi\epsilon_0} 2\pi\sigma z \left(\frac{1}{z} - \frac{1}{\sqrt{R^2 + z^2}} \right) \hat{\mathbf{z}}
\end{aligned}$$

When $R \rightarrow \infty$

$$\mathbf{E} = \frac{\sigma}{2\epsilon_0} \hat{\mathbf{z}}.$$

2.7

$$\mathbf{E} = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{q}{z^2} \hat{\mathbf{z}} & z > R \\ \mathbf{0} & z < R \end{cases}$$

2.8

$$\mathbf{E} = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{q}{z^2} \hat{\mathbf{z}} & z > R \\ \frac{1}{4\pi\epsilon_0} \frac{qz}{R^3} \hat{\mathbf{z}} & z < R \end{cases}$$

2.9

(a)

$$\begin{aligned}\rho &= \epsilon_0 \nabla \cdot \mathbf{E} \\ &= \epsilon_0 \frac{1}{r^2} \frac{\partial}{\partial r} (kr^5) \\ &= 5\epsilon_0 kr^2\end{aligned}$$

(b)

$$\begin{aligned}Q_{\text{enc}} &= \epsilon_0 \oint \mathbf{E} \cdot d\mathbf{a} \\ &= \epsilon_0 \int_0^{2\pi} \int_0^\pi kR^3 R d\theta R \sin \theta d\phi \\ &= 2\pi\epsilon_0 kR^5 [-\cos \theta]_0^\pi \\ &= 4\pi\epsilon_0 kR^5 \\ Q_{\text{enc}} &= \int_V \rho d\tau \\ &= \int_0^{2\pi} \int_0^\pi \int_0^R 5\epsilon_0 kr^2 dr r d\theta r \sin \theta d\phi \\ &= 10\pi\epsilon_0 k \int_0^\pi \int_0^R r^4 \sin \theta dr d\theta \\ &= 2\pi\epsilon_0 kR^5 [-\cos \theta]_0^\pi \\ &= 4\pi\epsilon_0 kR^5\end{aligned}$$

2.10

If the charge was surrounded by 8 such cubes the total flux through all the cubes would be q/ϵ_0 . There are 24 outside faces to the larger cube, so the total flux through the shaded face is $q/(24\epsilon_0)$.

2.11

$$\begin{aligned}\int \mathbf{E}_{\text{inside}} \cdot d\mathbf{a} &= \frac{Q_{\text{enc}}}{\epsilon_0} \\ &= 0 \\ \mathbf{E}_{\text{inside}} &= \mathbf{0} \\ \int \mathbf{E}_{\text{outside}} \cdot d\mathbf{a} &= \frac{Q_{\text{enc}}}{\epsilon_0} \\ 4\pi r^2 E_{\text{outside}} &= \frac{4\pi R^2 \sigma}{\epsilon_0} \\ \mathbf{E}_{\text{outside}} &= \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}\end{aligned}$$

2.12

$$\begin{aligned}\int \mathbf{E} \cdot d\mathbf{a} &= \frac{Q_{\text{enc}}}{\epsilon_0} \\ 4\pi r^2 E &= \frac{\frac{4}{3}\pi r^3 \rho}{\epsilon_0} \\ \mathbf{E} &= \frac{r\rho}{3\epsilon_0} \hat{\mathbf{r}}\end{aligned}$$

2.13

$$\begin{aligned}\int \mathbf{E} \cdot d\mathbf{a} &= \frac{Q_{\text{enc}}}{\epsilon_0} \\ 2\pi s l E &= \frac{l\lambda}{\epsilon_0} \\ \mathbf{E} &= \frac{1}{2\pi\epsilon_0} \frac{\lambda}{s} \hat{\mathbf{s}}\end{aligned}$$

2.14

$$\begin{aligned}Q_{\text{enc}} &= \int_V \rho \, d\tau \\ &= \int_0^{2\pi} \int_0^\pi \int_0^r k r'^3 \sin \theta \, dr' \, d\theta \, d\phi \\ &= 2\pi k \int_0^\pi \left[\frac{1}{4} r'^4 \sin \theta \right]_0^r d\theta \\ &= \frac{1}{2} \pi k r^4 [-\cos \theta]_0^\pi \\ &= \pi k r^4 \\ \int \mathbf{E} \cdot d\mathbf{a} &= \frac{Q_{\text{enc}}}{\epsilon_0} \\ 4\pi r^2 E &= \frac{\pi k r^4}{\epsilon_0} \\ \mathbf{E} &= \frac{k r^2}{4\epsilon_0} \hat{\mathbf{r}}\end{aligned}$$

2.15

(a) $\mathbf{E} = \mathbf{0}$

(b)

$$\begin{aligned}
 Q_{\text{enc}} &= \int_0^{2\pi} \int_0^\pi \int_a^r k \sin \theta \, dr' \, d\theta \, d\phi \\
 &= 4\pi k(r-a) \\
 4\pi r^2 E &= \frac{4\pi k(r-a)}{\epsilon_0} \\
 \mathbf{E} &= \frac{k(r-a)}{\epsilon_0 r^2} \hat{\mathbf{r}}
 \end{aligned}$$

(c) $\mathbf{E} = \frac{k(b-a)}{\epsilon_0 r^2} \hat{\mathbf{r}}$

2.16

(a)

$$\begin{aligned}
 Q_{\text{enc}} &= \pi s^2 l \rho \\
 2\pi s l E &= \frac{\pi s^2 l \rho}{\epsilon_0} \\
 \mathbf{E} &= \frac{s\rho}{2\epsilon_0} \hat{\mathbf{s}}
 \end{aligned}$$

(b)

$$\mathbf{E} = \frac{a^2 \rho}{2\epsilon_0 s} \hat{\mathbf{s}}$$

(c)

$$\mathbf{E} = \mathbf{0}$$

2.17

$$\begin{aligned}
 2AE_{\text{inside}} &= \frac{2Ay\rho}{\epsilon_0} \\
 \mathbf{E}_{\text{inside}} &= \frac{y\rho}{\epsilon_0} \\
 \mathbf{E} &= \begin{cases} \frac{d\rho}{\epsilon_0} & d < y \\ \frac{y\rho}{\epsilon_0} & 0 < y < d \\ -\frac{y\rho}{\epsilon_0} & -d < y < 0 \\ -\frac{d\rho}{\epsilon_0} & y < -d \end{cases}
 \end{aligned}$$

2.18

The electric field inside a uniformly charged solid sphere is

$$\mathbf{E} = \frac{r\rho}{3\epsilon_0} \hat{\mathbf{r}}.$$

$$\begin{aligned}\mathbf{d} &= \mathbf{r}_1 - \mathbf{r}_2 \\ \mathbf{E} &= \frac{r_1\rho}{3\epsilon_0} \hat{\mathbf{r}}_1 - \frac{r_2\rho}{3\epsilon_0} \hat{\mathbf{r}}_2 \\ &= \frac{\rho}{3\epsilon_0} (\mathbf{r}_1 - \mathbf{r}_2) \\ &= \frac{\rho}{3\epsilon_0} \mathbf{d}\end{aligned}$$

2.20

a is impossible because its curl is nonzero.

$$\begin{aligned}V &= - \int_0^y 2kxy' dy' - \int_0^z 2kyz' dz \\ &= -2kx \left[\frac{1}{2}y'^2 \right]_0^y - 2ky \left[\frac{1}{2}z'^2 \right]_0^z \\ &= -k(xy^2 + yz^2) \\ -\nabla V &= k[y^2 \hat{\mathbf{x}} + (2xy + z^2) \hat{\mathbf{y}} + 2yz \hat{\mathbf{z}}] \\ &= \mathbf{E}\end{aligned}$$

2.21

$$\begin{aligned}
\mathbf{E} &= \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} & r > R \\ \frac{1}{4\pi\epsilon_0} \frac{qr}{R^3} & r < R \end{cases} \\
V_{\text{outside}}(r) &= - \int_{\infty}^r \frac{1}{4\pi\epsilon_0} \frac{q}{r'^2} dr' \\
&= - \frac{1}{4\pi\epsilon_0} q \left[-\frac{1}{r'} \right]_{\infty}^r \\
&= \frac{1}{4\pi\epsilon_0} \frac{q}{r} \\
-\nabla V_{\text{outside}} &= \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}} \\
&= \mathbf{E}_{\text{outside}} \\
V_{\text{inside}}(r) &= - \left(\int_{\infty}^R \frac{1}{4\pi\epsilon_0} \frac{q}{r'^2} dr' + \int_R^r \frac{1}{4\pi\epsilon_0} \frac{qr'}{R^3} dr' \right) \\
&= - \left(-\frac{1}{4\pi\epsilon_0} \frac{q}{R} + \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} \left[\frac{1}{2} r'^2 \right]_R^r \right) \\
&= \frac{1}{4\pi\epsilon_0} \frac{q}{2R} \left[3 - \left(\frac{r}{R} \right)^2 \right] \\
-\nabla V_{\text{inside}} &= \frac{1}{4\pi\epsilon_0} \frac{qr}{R^3} \hat{\mathbf{r}} \\
&= \mathbf{E}_{\text{inside}}
\end{aligned}$$

2.22

$$\begin{aligned}
\mathbf{E} &= \frac{1}{2\pi\epsilon_0} \frac{\lambda}{s} \hat{\mathbf{s}} \\
V &= - \int_O^s \frac{1}{2\pi\epsilon_0} \frac{\lambda}{s'} ds' \\
&= - \frac{1}{2\pi\epsilon_0} \lambda \ln \frac{s}{O} \\
-\nabla V &= \frac{1}{2\pi\epsilon_0} \frac{\lambda}{s} \hat{\mathbf{s}}
\end{aligned}$$

2.23

$$\begin{aligned}
\mathbf{E} &= \begin{cases} \mathbf{0} & r < a \\ \frac{k(r-a)}{\epsilon_0 r^2} \hat{\mathbf{r}} & a < r < b \\ \frac{k(b-a)}{\epsilon_0 r^2} \hat{\mathbf{r}} & b < r \end{cases} \\
V(0) &= - \int_{\infty}^0 E dr \\
&= - \left(\int_{\infty}^b \frac{k(b-a)}{\epsilon_0 r^2} dr + \int_b^a \frac{k(r-a)}{\epsilon_0 r^2} dr \right) \\
&= - \left(\frac{k(b-a)}{\epsilon_0} \left[-\frac{1}{r} \right]_{\infty}^b + \frac{k}{\epsilon_0} \left[\ln r + \frac{a}{r} \right]_b^a \right) \\
&= - \left[-\frac{k(b-a)}{\epsilon_0 b} + \frac{k}{\epsilon_0} \left(\ln a + 1 - \ln b - \frac{a}{b} \right) \right] \\
&= -\frac{k}{\epsilon_0} \left(-1 + \frac{a}{b} + \ln \frac{a}{b} + 1 - \frac{a}{b} \right) \\
&= \frac{k}{\epsilon_0} \ln \frac{b}{a}
\end{aligned}$$

2.24

$$\begin{aligned}
V(b) - V(0) &= - \int_0^b E dr \\
&= - \left(\int_0^a \frac{s\rho}{2\epsilon_0} ds + \int_a^b \frac{a^2\rho}{2\epsilon_0 s} ds \right) \\
&= - \left(\frac{\rho}{2\epsilon_0} \left[\frac{1}{2} s^2 \right]_0^a + \frac{a^2\rho}{2\epsilon_0} \ln \frac{b}{a} \right) \\
&= - \left(\frac{a^2\rho}{4\epsilon_0} + \frac{a^2\rho}{2\epsilon_0} \ln \frac{b}{a} \right) \\
&= -\frac{a^2\rho}{4\epsilon_0} \left(1 + 2 \ln \frac{a}{b} \right)
\end{aligned}$$

2.25

(a)

$$V = \frac{1}{4\pi\epsilon_0} \frac{2q}{\sqrt{(d/2)^2 + z^2}}$$

(b)

$$\begin{aligned} V &= \frac{1}{4\pi\epsilon_0} \int_{-L}^L \frac{\lambda}{\sqrt{x^2 + z^2}} dx \\ &= \frac{1}{4\pi\epsilon_0} \lambda \ln \left(1 + \frac{2L(L + \sqrt{L^2 + z^2})}{z^2} \right) \end{aligned}$$

(c)

$$\begin{aligned} V &= \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^R \frac{\sigma}{\sqrt{r^2 + z^2}} r dr d\theta \\ &= \frac{1}{4\pi\epsilon_0} 2\pi\sigma(\sqrt{R^2 + z^2} - z) \end{aligned}$$

2.26

$$\begin{aligned} V_{\text{bottom}} &= \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^h \frac{\sqrt{2}\sigma z}{\sqrt{2}z} d\phi dz \\ &= \frac{\sigma h}{2\epsilon_0} \\ V_{\text{top}} &= \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^h \frac{\sqrt{2}\sigma z}{\sqrt{z^2 + (h-z)^2}} d\phi dz \\ &= \frac{\sqrt{2}\sigma}{2\epsilon_0} \int_0^h \frac{z}{\sqrt{z^2 + (h-z)^2}} dz \\ &= \frac{\sigma h}{4\epsilon_0} \ln(3 + 2\sqrt{2}) \\ V_{\text{bottom}} - V_{\text{top}} &= \frac{\sigma h}{2\epsilon_0} \left[1 - \frac{1}{2} \ln(3 + 2\sqrt{2}) \right] \end{aligned}$$

2.28

$$\begin{aligned} V(r) &= \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^\pi \int_0^R \frac{\rho r'^2 \sin \theta}{\sqrt{r^2 + r'^2 - 2rr' \cos \theta}} dr' d\theta d\phi \\ &= \frac{\rho}{2\epsilon_0} \int_0^\pi \int_0^R \frac{r'^2 \sin \theta}{\sqrt{r^2 + r'^2 - 2rr' \cos \theta}} dr' d\theta \\ &= \frac{\rho}{2\epsilon_0} \left(R^2 - \frac{r^2}{3} \right) \\ &= \frac{q}{8\pi\epsilon_0 R} \left(3 - \frac{r^2}{R^2} \right) \end{aligned}$$

2.31

(a)

$$W = \frac{q^2}{4\pi\epsilon_0 a} \left(\frac{1}{\sqrt{2}} - 2 \right)$$

(b)

$$\begin{aligned} W &= \frac{1}{4\pi\epsilon_0} \left(-\frac{q^2}{a} + \frac{q^2}{\sqrt{2}a} - \frac{q^2}{a} - \frac{q^2}{a} + \frac{q^2}{\sqrt{2}a} - \frac{q^2}{a} \right) \\ &= \frac{q^2}{2\pi\epsilon_0 a} \left(\frac{1}{\sqrt{2}} - 2 \right) \end{aligned}$$

2.32

$$W = \frac{1}{4\pi\epsilon_0} \frac{q_A q_B}{a}$$

$$W = K_1 + K_2$$

$$\frac{1}{4\pi\epsilon_0} \frac{q_A q_B}{a} = \frac{1}{2} m_A v_A^2 + \frac{1}{2} m_B v_B^2$$

$$\frac{1}{2\pi\epsilon_0} \frac{q_A q_B}{a} = m_A v_A^2 + m_B v_B^2$$

$$0 = m_B v_B - m_A v_A$$

$$v_B = \frac{m_A}{m_B} v_A$$

$$\frac{1}{2\pi\epsilon_0} \frac{q_A q_B}{a} = m_A v_A^2 + m_B \left(\frac{m_A}{m_B} v_A \right)^2$$

$$= m_A v_A^2 + \frac{m_A^2}{m_B} v_A^2$$

$$= \frac{m_A(m_A + m_B)}{m_B} v_A^2$$

$$v_A = \sqrt{\frac{1}{2\pi\epsilon_0} \frac{q_A q_B}{(m_A + m_B)a} \frac{m_B}{m_A}}$$

$$v_B = \sqrt{\frac{1}{2\pi\epsilon_0} \frac{q_A q_B}{(m_A + m_B)a} \frac{m_A}{m_B}}$$

2.33

$$\begin{aligned}
 W &= \frac{1}{4\pi\epsilon_0} \left(-\frac{q^2}{a} + \frac{q^2}{2a} - \frac{q^2}{3a} + \dots \right) \\
 &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{a} \sum_{n=1}^{\infty} \frac{(-1)^n}{n} \\
 &= -\frac{1}{4\pi\epsilon_0} \frac{q^2}{a} \ln 2
 \end{aligned}$$

2.34

(a)

$$\begin{aligned}
 V &= \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{q}{2R} \left[3 - \left(\frac{r}{R} \right)^2 \right] & r < R \\ \frac{1}{4\pi\epsilon_0} \frac{q}{r} & r > R \end{cases} \\
 W &= \frac{1}{2} \int \rho V \, d\tau \\
 &= \frac{1}{2} \int_0^{2\pi} \int_0^\pi \int_0^R \rho \frac{1}{4\pi\epsilon_0} \frac{q}{2R} \left[3 - \left(\frac{r}{R} \right)^2 \right] r^2 \sin \theta \, dr \, d\theta \, d\phi \\
 &= \frac{q\rho}{8\epsilon_0 R} \int_0^\pi \int_0^R \left[3 - \left(\frac{r}{R} \right)^2 \right] r^2 \sin \theta \, dr \, d\theta \\
 &= \frac{q\rho R^2}{5\epsilon_0} \\
 &= \frac{qR^2}{5\epsilon_0} \frac{q}{\frac{4}{3}\pi R^3} \\
 &= \frac{1}{4\pi\epsilon_0} \frac{3q^2}{5R}
 \end{aligned}$$

(b)

$$\begin{aligned}
\mathbf{E} &= \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}} & r > R \\ \frac{1}{4\pi\epsilon_0} \frac{qr}{R^3} \hat{\mathbf{r}} & r < R \end{cases} \\
E^2 &= \begin{cases} \frac{1}{16\pi^2\epsilon_0^2} \frac{q^2}{r^4} & r > R \\ \frac{1}{16\pi^2\epsilon_0^2} \frac{q^2 r^2}{R^6} & r < R \end{cases} \\
W &= \frac{\epsilon_0}{2} \int E^2 d\tau \\
&= \frac{\epsilon_0}{2} \left(\int_0^{2\pi} \int_0^\pi \int_0^R \frac{1}{16\pi^2\epsilon_0^2} \frac{q^2 r^2}{R^6} r^2 \sin\theta \, dr \, d\theta \, d\phi \right. \\
&\quad \left. + \int_0^{2\pi} \int_0^\pi \int_R^\infty \frac{1}{16\pi^2\epsilon_0^2} \frac{q^2}{r^4} r^2 \sin\theta \, dr \, d\theta \, d\phi \right) \\
&= \frac{\epsilon_0}{2} \frac{1}{16\pi^2\epsilon_0^2} 2\pi q^2 \left(\int_0^\pi \int_0^R \frac{r^4}{R^6} \sin\theta \, dr \, d\theta + \int_0^\pi \int_R^\infty \frac{1}{r^2} \sin\theta \, dr \, d\theta \right) \\
&= \frac{1}{16\pi\epsilon_0} q^2 \left(\int_0^\pi \int_0^R \frac{r^4}{R^6} \sin\theta \, dr \, d\theta + \int_0^\pi \int_R^\infty \frac{1}{r^2} \sin\theta \, dr \, d\theta \right) \\
&= \frac{1}{16\pi\epsilon_0} q^2 \left(\frac{2}{5R} + \frac{2}{R} \right) \\
&= \frac{1}{4\pi\epsilon_0} \frac{3q^2}{5R}
\end{aligned}$$

(c)

$$\begin{aligned}
W &= \frac{\epsilon_0}{2} \left(\int_V E^2 d\tau + \oint_S V \mathbf{E} \cdot d\mathbf{a} \right) \\
&= \frac{\epsilon_0}{2} \left(\int_0^{2\pi} \int_0^\pi \int_0^R \frac{1}{(4\pi\epsilon_0)^2} \frac{q^2 r^2}{R^6} r^2 \sin \theta dr d\theta d\phi \right. \\
&\quad + \int_0^{2\pi} \int_0^\pi \int_R^a \frac{1}{(4\pi\epsilon_0)^2} \frac{q^2}{r^4} r^2 \sin \theta dr d\theta d\phi \\
&\quad \left. + \int_0^{2\pi} \int_0^\pi \frac{1}{4\pi\epsilon_0} \frac{q}{a} \frac{1}{4\pi\epsilon_0} \frac{q}{a^2} a^2 \sin \theta d\theta d\phi \right) \\
&= \frac{\epsilon_0}{2} \frac{1}{(4\pi\epsilon_0)^2} 2\pi q^2 \left(\int_0^\pi \int_0^R \frac{r^4}{R^6} \sin \theta dr d\theta \right. \\
&\quad \left. + \int_0^\pi \int_R^a \frac{1}{r^2} \sin \theta dr d\theta + \int_0^\pi \frac{1}{a} \sin \theta d\theta \right) \\
&= \frac{\epsilon_0}{2} \frac{1}{(4\pi\epsilon_0)^2} 2\pi q^2 \left[\frac{2}{5R} + 2 \left(\frac{1}{R} - \frac{1}{a} \right) + \frac{2}{a} \right] \\
&= \frac{1}{8\pi\epsilon_0} q^2 \left[\frac{1}{5R} + \frac{1}{R} \right] \\
&= \frac{1}{4\pi\epsilon_0} \frac{3q^2}{5R}
\end{aligned}$$

2.36

(a)

$$\begin{aligned}
\mathbf{E} &= \begin{cases} \mathbf{0} & r < a \\ \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}} & a < r < b \\ \mathbf{0} & b < r \end{cases} \\
E^2 &= \begin{cases} 0 & r < a \\ \frac{1}{(4\pi\epsilon_0)^2} \frac{q^2}{r^4} & a < r < b \\ 0 & b < r \end{cases} \\
W &= \frac{\epsilon_0}{2} \int E^2 d\tau \\
&= \frac{\epsilon_0}{2} \int_0^{2\pi} \int_0^\pi \int_a^b \frac{1}{(4\pi\epsilon_0)^2} \frac{q^2}{r^4} r^2 \sin \theta dr d\theta d\phi \\
&= \frac{\epsilon_0}{2} \frac{1}{(4\pi\epsilon_0)^2} 2\pi q^2 \int_0^\pi \int_a^b \frac{\sin \theta}{r^2} dr d\theta \\
&= \frac{q^2}{8\pi\epsilon_0} \left(\frac{1}{a} - \frac{1}{b} \right)
\end{aligned}$$

(b)

$$\begin{aligned}
W_{\text{shell}} &= \frac{1}{8\pi\epsilon_0} \frac{q^2}{R} \\
\mathbf{E} &= \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}} \\
\mathbf{E}_1 \cdot \mathbf{E}_2 &= -\frac{1}{(4\pi\epsilon_0)^2} \frac{q^2}{r^4} \\
W_{\text{total}} &= W_1 + W_2 + \epsilon_0 \int \mathbf{E}_1 \cdot \mathbf{E}_2 d\tau \\
&= \frac{q^2}{8\pi\epsilon_0} \left(\frac{1}{a} + \frac{1}{b} \right) - \epsilon_0 \int_0^{2\pi} \int_0^\pi \int_b^\infty \frac{1}{(4\pi\epsilon_0)^2} \frac{q^2}{r^4} r^2 \sin\theta dr d\theta d\phi \\
&= \frac{q^2}{8\pi\epsilon_0} \left(\frac{1}{a} + \frac{1}{b} \right) - \frac{1}{8\pi\epsilon_0} q^2 \int_0^\pi \int_b^\infty \frac{1}{r^2} \sin\theta dr d\theta \\
&= \frac{q^2}{8\pi\epsilon_0} \left(\frac{1}{a} + \frac{1}{b} \right) - \frac{1}{4\pi\epsilon_0} q^2 \int_b^\infty \frac{1}{r^2} dr \\
&= \frac{q^2}{8\pi\epsilon_0} \left(\frac{1}{a} + \frac{1}{b} \right) - \frac{1}{4\pi\epsilon_0} \frac{q^2}{b} \\
&= \frac{q^2}{8\pi\epsilon_0} \left(\frac{1}{a} + \frac{1}{b} - \frac{2}{b} \right) \\
&= \frac{q^2}{8\pi\epsilon_0} \left(\frac{1}{a} - \frac{1}{b} \right)
\end{aligned}$$

2.37

$$\begin{aligned}
r_1 &= r \\
E_1 &= \frac{1}{4\pi\epsilon_0} \frac{q_1}{r_1^2} \\
&= \frac{1}{4\pi\epsilon_0} \frac{q_1}{r^2} \\
r_2 &= \sqrt{a^2 + r^2 - 2ar \cos \theta} \\
E_2 &= \frac{1}{4\pi\epsilon_0} \frac{q_2}{r_2^2} \\
&= \frac{1}{4\pi\epsilon_0} \frac{q_2}{a^2 + r^2 - 2ar \cos \theta} \\
\cos \alpha &= \frac{r - a \cos \theta}{\sqrt{a^2 + r^2 - 2ar \cos \theta}} \\
\mathbf{E}_1 \cdot \mathbf{E}_2 &= E_1 E_2 \cos \alpha \\
&= \frac{1}{(4\pi\epsilon_0)^2} \frac{q_1 q_2}{r^2 (a^2 + r^2 - 2ar \cos \theta)} \frac{r - a \cos \theta}{\sqrt{a^2 + r^2 - 2ar \cos \theta}} \\
&= \frac{1}{(4\pi\epsilon_0)^2} \frac{q_1 q_2 (r - a \cos \theta)}{r^2 (a^2 + r^2 - 2ar \cos \theta)^{3/2}} \\
\epsilon_0 \int \mathbf{E}_1 \cdot \mathbf{E}_2 d\tau &= \epsilon_0 \int_0^{2\pi} \int_0^\pi \int_0^\infty \frac{1}{(4\pi\epsilon_0)^2} \frac{q_1 q_2 (r - a \cos \theta)}{r^2 (a^2 + r^2 - 2ar \cos \theta)^{3/2}} r^2 \sin \theta dr d\theta d\phi \\
&= \frac{q_1 q_2}{8\pi\epsilon_0} \int_0^\pi \int_0^\infty \frac{(r - a \cos \theta) \sin \theta}{(a^2 + r^2 - 2ar \cos \theta)^{3/2}} dr d\theta
\end{aligned}$$

2.38

(a)

$$\begin{aligned}
\sigma_R &= \frac{q}{4\pi R^2} \\
\sigma_a &= -\frac{q}{4\pi a^2} \\
\sigma_b &= \frac{q}{4\pi b^2}
\end{aligned}$$

(b)

$$\begin{aligned}
V &= -\int_\infty^b \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} dr - \int_a^R \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} dr \\
&= \frac{1}{4\pi\epsilon_0} q \left(\frac{1}{b} + \frac{1}{R} - \frac{1}{a} \right)
\end{aligned}$$

(c)

$$\sigma_b = 0$$

$$V = \frac{1}{4\pi\epsilon_0} q \left(\frac{1}{R} - \frac{1}{a} \right)$$

2.39

(a)

$$\sigma_a = -\frac{q_a}{4\pi a^2}$$

$$\sigma_b = -\frac{q_b}{4\pi b^2}$$

$$\sigma_R = \frac{q_a + q_b}{4\pi R^2}$$

(b)

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q_a + q_b}{r^2} \hat{\mathbf{r}}$$

(c)

$$\mathbf{E}_a = \frac{1}{4\pi\epsilon_0} \frac{q_a}{r^2} \hat{\mathbf{r}}$$

$$\mathbf{E}_b = \frac{1}{4\pi\epsilon_0} \frac{q_b}{r^2} \hat{\mathbf{r}}$$

(d)

$$\mathbf{0}$$

(e) a, b

2.40

(a) No. If it's close to the wall it will induce a surface charge and be attracted.

(b) No. If the conductor contains a cavity containing a like charge it will be repelled.

2.41

By Gauss's law, the electric field of each plate is

$$\oint \mathbf{E} \cdot d\mathbf{a} = \frac{Q_{\text{enc}}}{\epsilon_0}$$

$$2A'E = \frac{A' \frac{Q}{A}}{\epsilon_0}$$

$$\mathbf{E} = \frac{Q}{2A\epsilon_0} \hat{\mathbf{n}}$$

so the field between the plates is zero and the field outside is $Q/A\epsilon_0\hat{\mathbf{n}}$, resulting in a pressure of

$$\begin{aligned} P &= \frac{\epsilon_0}{2} E^2 \\ &= \frac{\epsilon_0}{2} \frac{Q^2}{A^2 \epsilon_0^2} \\ &= \frac{Q^2}{2A^2 \epsilon_0} \end{aligned}$$

2.42

$$\begin{aligned} \mathbf{E}_{\text{above}} &= \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \hat{\mathbf{r}} \\ \mathbf{f} &= \frac{1}{2} \sigma \mathbf{E}_{\text{above}} \\ &= \frac{1}{2} \frac{Q}{4\pi R^2} \frac{1}{4\pi\epsilon_0} \frac{Q}{R^2} \hat{\mathbf{r}} \\ &= \frac{Q^2}{32\pi^2 \epsilon_0 R^4} \hat{\mathbf{r}} \\ \mathbf{F} &= \int_0^{2\pi} \int_0^{\pi/2} \frac{Q^2}{32\pi^2 \epsilon_0 R^4} \cos \theta R^2 \sin \theta d\theta d\phi \hat{\mathbf{z}} \\ &= \frac{Q^2}{16\pi\epsilon_0 R^2} \int_0^{\pi/2} \cos \theta \sin \theta d\theta \hat{\mathbf{z}} \\ &= \frac{Q^2}{32\pi\epsilon_0 R^2} \hat{\mathbf{z}} \end{aligned}$$

2.43

$$\begin{aligned}
 \oint \mathbf{E} \cdot d\mathbf{a} &= \frac{Q}{\epsilon_0} \\
 2\pi s L E &= \frac{Q}{\epsilon_0} \\
 \mathbf{E} &= \frac{Q}{2\pi L \epsilon_0 s} \hat{\mathbf{s}} \\
 V &= - \int_b^a \frac{Q}{2\pi \epsilon_0 L s} \frac{1}{s} dr \\
 &= \frac{Q}{2\pi \epsilon_0 L} \ln \frac{b}{a} \\
 C &= \frac{Q}{V} \\
 &= \frac{2\pi \epsilon_0 L}{\ln b/a}
 \end{aligned}$$

So the capacitance per unit length is

$$C = \frac{2\pi \epsilon_0}{\ln b/a}.$$

2.44

(a)

$$\begin{aligned}
 P &= \frac{\epsilon_0}{2} E^2 \\
 W &= Fd \\
 &= PA\epsilon \\
 &= \frac{\epsilon_0}{2} E^2 A\epsilon
 \end{aligned}$$

(b)

$$\frac{\epsilon_0}{2} E^2 A\epsilon$$

2.46

$$\begin{aligned}
\nabla \cdot \mathbf{E} &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 3 \frac{k}{r} \right) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{k}{r} 2 \sin \theta \cos \theta \sin \phi \right) \\
&\quad + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \left(\frac{k}{r} \sin \theta \cos \phi \right) \\
&= \frac{3k}{r^2} + \frac{1}{r \sin \theta} \frac{2k}{r} \sin \phi (2 \sin \theta \cos^2 \theta - \sin^3 \theta) - \frac{1}{r \sin \theta} \frac{k}{r} \sin \theta \sin \phi \\
&= \frac{3k}{r^2} + \frac{2k \sin \phi}{r^2} (2 \cos^2 \theta - \sin^2 \theta) - \frac{k}{r^2} \sin \phi \\
&= \frac{k}{r^2} [3 + 2 \sin \phi (2 \cos^2 \theta - \sin^2 \theta) - \sin \phi] \\
&= \frac{k}{r^2} [3 + \sin \phi (4 \cos^2 \theta - 2 \sin^2 \theta - 1)] \\
&= \frac{k}{r^2} [3 + \sin \phi (6 \cos^2 \theta - 3)] \\
&= \frac{3k}{r^2} (1 + \cos 2\theta \sin \phi) \\
\rho &= \epsilon_0 \nabla \cdot \mathbf{E} \\
&= \frac{3k\epsilon_0}{r^2} (1 + \cos 2\theta \sin \phi)
\end{aligned}$$

2.47

$$\begin{aligned}
\mathbf{E} &= \frac{1}{4\pi\epsilon_0} \frac{Qr}{R^3} \hat{\mathbf{r}} \\
\rho &= \frac{Q}{\frac{4}{3}\pi R^3} \\
\rho \mathbf{E} &= \frac{3Q}{4\pi R^3} \frac{1}{4\pi\epsilon_0} \frac{Qr}{R^3} \hat{\mathbf{r}} \\
&= \frac{3r}{\epsilon_0} \left(\frac{Q}{4\pi R^3} \right)^2 \hat{\mathbf{r}} \\
F_z &= \int_0^{2\pi} \int_0^{\pi/2} \int_0^R \frac{3r}{\epsilon_0} \left(\frac{Q}{4\pi R^3} \right)^2 \cos \theta r^2 \sin \theta \, dr \, d\theta \, d\phi \\
&= \frac{3\pi}{\epsilon_0} \left(\frac{Q}{4\pi R^3} \right)^2 \int_0^{\pi/2} \int_0^R r^3 \sin 2\theta \, dr \, d\theta \\
&= \frac{3\pi}{\epsilon_0} \left(\frac{Q}{4\pi R^3} \right)^2 \frac{R^4}{4} \\
&= \frac{3Q^2}{64\pi\epsilon_0 R^2}
\end{aligned}$$

2.49

$$\begin{aligned}
Q_{\text{enc}} &= \int_0^{2\pi} \int_0^\pi \int_0^r k r'^3 \sin \theta \, dr' \, d\theta \, d\phi \\
&= 2\pi k \int_0^\pi \int_0^r r'^3 \sin \theta \, dr' \, d\theta \\
&= \pi k r^4 \\
\oint \mathbf{E} \cdot d\mathbf{a} &= \frac{Q_{\text{enc}}}{\epsilon_0} \\
4\pi r^2 E &= \frac{\pi k r^4}{\epsilon_0} \\
\mathbf{E} &= \begin{cases} \frac{k r^2}{4\epsilon_0} \hat{\mathbf{r}} & r < R \\ \frac{k R^4}{4\epsilon_0 r^2} \hat{\mathbf{r}} & r > R \end{cases} \\
W &= \frac{\epsilon_0}{2} \left(\int_0^{2\pi} \int_0^\pi \int_0^R \frac{k^2 r^4}{16\epsilon_0^2} \sin \theta \, dr \, d\theta \, d\phi \right. \\
&\quad \left. \int_0^{2\pi} \int_0^\pi \int_R^\infty \frac{k^2 R^8}{16\epsilon_0^2 r^4} \sin \theta \, dr \, d\theta \, d\phi \right) \\
&= \frac{\epsilon_0}{2} 2\pi \frac{k^2}{16\epsilon_0^2} \left(\int_0^\pi \int_0^R r^6 \sin \theta \, dr \, d\theta + \int_0^\pi \int_R^\infty \frac{R^8 \sin \theta}{r^2} \, dr \, d\theta \right) \\
&= \frac{\pi k^2}{16\epsilon_0} \left(\frac{2R^7}{7} + 2R^7 \right) \\
&= \frac{\pi k^2 R^7}{7\epsilon_0}
\end{aligned}$$

2.50

$$\begin{aligned}
V(\mathbf{r}) &= A \frac{e^{-\lambda r}}{r} \\
\mathbf{E} &= -\nabla V \\
&= A e^{-\lambda r} (1 + \lambda r) \frac{\hat{\mathbf{r}}}{r^2} \\
\rho &= \epsilon_0 \nabla \cdot \mathbf{E} \\
&= \epsilon_0 \left[A e^{-\lambda r} (1 + \lambda r) \nabla \cdot \frac{\hat{\mathbf{r}}}{r^2} + \frac{\hat{\mathbf{r}}}{r^2} \cdot \nabla (A e^{-\lambda r} (1 + \lambda r)) \right] \\
&= A \epsilon_0 \left[4\pi \delta(\mathbf{r}) + \frac{\hat{\mathbf{r}}}{r^2} \cdot (-\lambda^2 e^{-\lambda r} r \hat{\mathbf{r}}) \right] \\
&= A \epsilon_0 \left(4\pi \delta(\mathbf{r}) - \frac{\lambda^2 e^{-\lambda r}}{r} \right)
\end{aligned}$$

2.51

$$\begin{aligned}
V &= \int \frac{1}{4\pi\epsilon_0} \frac{\sigma}{z} dA \\
&= \frac{\sigma}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^R \frac{r}{\sqrt{r^2 + R^2 - 2rR\cos\theta}} dr d\theta \\
&= \frac{R\sigma}{4\pi\epsilon_0} \int_0^{2\pi} \left[\cos\theta \ln \left(1 + \csc \frac{\theta}{2} \right) + 2 \sin \frac{\theta}{2} - 1 \right] d\theta \\
&= \frac{R\sigma}{\pi\epsilon_0}
\end{aligned}$$

2.52

(a)

$$\begin{aligned}
V_- &= \frac{1}{2\pi\epsilon_0} \lambda \ln \frac{s_-}{a} \\
&= \frac{1}{2\pi\epsilon_0} \lambda \ln \frac{\sqrt{(y+a)^2 + z^2}}{a} \\
V_+ &= -\frac{1}{2\pi\epsilon_0} \lambda \ln \frac{s_+}{a} \\
&= -\frac{1}{2\pi\epsilon_0} \lambda \ln \frac{\sqrt{(y-a)^2 + z^2}}{a} \\
V &= V_- + V_+ \\
&= \frac{1}{4\pi\epsilon_0} \lambda \ln \frac{(y+a)^2 + z^2}{(y-a)^2 + z^2}
\end{aligned}$$

2.53

(a)

$$\begin{aligned}
\nabla^2 V &= -\frac{\rho}{\epsilon_0} \\
\nabla \cdot \nabla V &= -\frac{\rho}{\epsilon_0} \\
\nabla \cdot \frac{dV}{dx} \hat{\mathbf{x}} &= -\frac{\rho}{\epsilon_0} \\
\frac{d^2 V}{dx^2} &= -\frac{\rho}{\epsilon_0}
\end{aligned}$$

(b)

$$qV = \frac{1}{2}mv^2$$
$$v = \sqrt{\frac{2qV}{m}}$$

(c)

$$I = A\rho v$$

(d)

$$\begin{aligned}\frac{d^2V}{dx^2} &= -\frac{I}{Av\epsilon_0} \\ &= -\frac{I}{A\epsilon_0}\sqrt{\frac{m}{2qV}} \\ &= \beta V^{-1/2}\end{aligned}$$

2.55

$$\begin{aligned}\rho &= \epsilon_0 \nabla \cdot \mathbf{E} \\ &= a\epsilon_0\end{aligned}$$

2.56

$$\begin{aligned}E &= \frac{3GM^2}{5R} \\ E_{\text{sun}} &= 2.3 \times 10^{41} \text{ J} \\ t &= \frac{E_{\text{sun}}}{P} \\ &= 1.89 \times 10^7 \text{ years}\end{aligned}$$

3 Potentials

3.1

$$\begin{aligned} V_{\text{ave}} &= \frac{q}{4\pi\epsilon_0} \frac{1}{2zR} \left. \sqrt{z^2 + R^2 - 2zR \cos \theta} \right|_0^\pi \\ &= \frac{q}{4\pi\epsilon_0} \frac{1}{2zR} \left(\sqrt{z^2 + R^2 + 2zR} - \sqrt{z^2 + R^2 - 2zR} \right) \\ &= \frac{q}{4\pi\epsilon_0} \frac{1}{2zR} \left(\sqrt{(z+R)^2} - \sqrt{(R-z)^2} \right) \\ &= \frac{q}{4\pi\epsilon_0} \frac{1}{2zR} (z+R - R+z) \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{R} \end{aligned}$$

The average potential due to external charges is V_{center} and the average potential due to internal charges is

$$\frac{1}{4\pi\epsilon_0} \frac{Q_{\text{enc}}}{R}$$

so

$$V_{\text{ave}} = V_{\text{center}} + \frac{1}{4\pi\epsilon_0} \frac{Q_{\text{enc}}}{R}.$$

3.2

A stable equilibrium is a minimum of potential energy. Laplace's equation doesn't allow for minimums, so they must be saddle points and the charge can escape.

3.3

$$\begin{aligned}
0 &= \nabla^2 V \\
&= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) \\
&= \frac{1}{r^2} \left(2r \frac{\partial V}{\partial r} + r^2 \frac{\partial^2 V}{\partial r^2} \right) \\
&= \frac{2}{r} \frac{\partial V}{\partial r} + \frac{\partial^2 V}{\partial r^2} \\
V &= \frac{c_1}{r} + c_2 \\
0 &= \nabla^2 V \\
&= \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial V}{\partial s} \right) \\
&= \frac{1}{s} \left(\frac{\partial V}{\partial s} + s \frac{\partial^2 V}{\partial s^2} \right) \\
&= \frac{1}{s} \frac{\partial V}{\partial s} + \frac{\partial^2 V}{\partial s^2} \\
V &= c_1 + c_2 \ln s
\end{aligned}$$

3.7

$$\begin{aligned}
\mathbf{E} &= \frac{1}{4\pi\epsilon_0} q^2 \left(-\frac{2}{(2d)^2} + \frac{2}{(4d)^2} - \frac{1}{(6d)^2} \right) \hat{\mathbf{z}} \\
&= -\frac{1}{4\pi\epsilon_0} \frac{29q^2}{72d^2} \hat{\mathbf{z}}
\end{aligned}$$

3.8

(a)

$$\begin{aligned}
V(r, \theta) &= \frac{1}{4\pi\epsilon_0} \left[\frac{q}{\sqrt{r^2 + a^2 - 2ra \cos \theta}} + \frac{q'}{\sqrt{r^2 + b^2 - 2rb \cos \theta}} \right] \\
&= \frac{1}{4\pi\epsilon_0} \left[\frac{q}{\sqrt{r^2 + a^2 - 2ra \cos \theta}} - \frac{Rq/a}{\sqrt{r^2 + (R^2/a)^2 - 2r(R^2/a) \cos \theta}} \right] \\
&= \frac{1}{4\pi\epsilon_0} \left[\frac{q}{\sqrt{r^2 + a^2 - 2ra \cos \theta}} - \frac{q}{\sqrt{R^2 + (ra/R)^2 - 2ra \cos \theta}} \right]
\end{aligned}$$

(b)

$$\begin{aligned}
\sigma &= -\epsilon_0 \left. \frac{\partial V}{\partial r} \right|_{r=R} \\
&= \frac{q}{4\pi R} \frac{R^2 - a^2}{(a^2 + R^2 - 2aR \cos \theta)^{3/2}} \\
Q_{\text{induced}} &= \int_0^{2\pi} \int_0^\pi \sigma R^2 \sin \theta \, d\theta \, d\phi \\
&= \frac{qR(R^2 - a^2)}{2} \int_0^\pi \frac{\sin \theta}{(a^2 + R^2 - 2aR \cos \theta)^{3/2}} \, d\theta \\
&= \frac{qR(R^2 - a^2)}{a(a^2 - R^2)} \\
&= -\frac{R}{a} q \\
&= q'
\end{aligned}$$

(c)

$$\begin{aligned}
W &= \frac{1}{2} qV \\
&= \frac{1}{8\pi\epsilon_0} \frac{qq'}{a - b} \\
&= -\frac{1}{8\pi\epsilon_0} \frac{q^2 R/a}{a - R^2/a} \\
&= -\frac{1}{8\pi\epsilon_0} \frac{q^2 R}{a^2 - R^2}
\end{aligned}$$

3.9

Place the second image charge at the centre of the sphere with charge

$$q'' = 4\pi\epsilon_0 R V_0.$$

$$\begin{aligned}
F &= \frac{1}{4\pi\epsilon_0} q \left(\frac{q'}{(a-b)^2} + \frac{q''}{a^2} \right) \\
&= \frac{qq'}{4\pi\epsilon_0} \left(\frac{1}{(a-b)^2} - \frac{1}{a^2} \right) \\
&= \frac{qq'}{4\pi\epsilon_0} \frac{a^2 - (a-b)^2}{a^2(a-b)^2} \\
&= \frac{qq'}{4\pi\epsilon_0} \frac{b(2a-b)}{a^2(a-b)^2} \\
&= \frac{q(-Rq/a)}{4\pi\epsilon_0} \frac{(R^2/a)(2a - R^2/a)}{a^2(a - R^2/a)^2} \\
&= -\frac{q^2}{4\pi\epsilon_0} \left(\frac{R}{a} \right)^3 \frac{2a^2 - R^2}{(a^2 - R^2)^2}
\end{aligned}$$

3.10

(a)

$$V(x, y, z) = \frac{1}{4\pi\epsilon_0} \lambda \ln \frac{y^2 + (z+d)^2}{y^2 + (z-d)^2}$$

(b)

$$\begin{aligned}
\sigma &= -\epsilon_0 \left. \frac{\partial V}{\partial z} \right|_{z=0} \\
&= -\frac{d\lambda}{\pi(d^2 + y^2)}
\end{aligned}$$

3.11

You need three charges: $-q$ at $(-a, b)$, $-q$ at $(a, -b)$, and q at $(-b, -a)$. The potential is

$$\begin{aligned}
V &= \frac{1}{4\pi\epsilon_0} q \left(\frac{1}{\sqrt{(x-a)^2 + (y-b)^2}} - \frac{1}{\sqrt{(x+a)^2 + (y-b)^2}} \right. \\
&\quad \left. - \frac{1}{\sqrt{(x-a)^2 + (y+b)^2}} + \frac{1}{\sqrt{(x+a)^2 + (y+b)^2}} \right).
\end{aligned}$$

The force on q is

$$\mathbf{F} = \frac{q^2}{16\pi\epsilon_0} \left[\left(\frac{a}{(a^2 + b^2)^{3/2}} - \frac{1}{a^2} \right) \hat{\mathbf{x}} + \left(\frac{b}{(a^2 + b^2)^{3/2}} - \frac{1}{b^2} \right) \hat{\mathbf{y}} \right].$$

The work to bring q in from infinity is

$$W = \frac{q^2}{16\pi\epsilon_0} \left(\frac{1}{\sqrt{a^2 + b^2}} - \frac{1}{a} - \frac{1}{b} \right).$$

3.12

Two infinitely long wires running parallel to the x -axis a distance $2a$ apart with charge densities λ and $-\lambda$ have cylindrical equipotential surfaces with centres at

$$y_0 = \pm a \coth \frac{2\pi\epsilon_0 V_0}{\lambda}$$

radii

$$R = a \operatorname{csch} \frac{2\pi\epsilon_0 V_0}{\lambda}.$$

We know the equipotential surfaces (the pipes) and want to find the wires so we can find the potential, so

$$\begin{aligned} d &= a \coth \frac{2\pi\epsilon_0 V_0}{\lambda} \\ R &= a \operatorname{csch} \frac{2\pi\epsilon_0 V_0}{\lambda} \\ \frac{d}{R} &= \cosh \frac{2\pi\epsilon_0 V_0}{\lambda} \\ \operatorname{arcosh} \frac{d}{R} &= \frac{2\pi\epsilon_0 V_0}{\lambda} \\ \lambda &= \frac{2\pi\epsilon_0 V_0}{\operatorname{arcosh} d/R} \\ R &= a \operatorname{csch} \operatorname{arcosh} \frac{d}{R} \\ a &= \frac{R}{\operatorname{csch} \operatorname{arcosh} d/R} \\ &= (d + R) \sqrt{\frac{2d}{d + R} - 1} \\ &= \sqrt{d^2 - R^2} \end{aligned}$$

thus the potential is

$$V = \frac{V_0}{2 \operatorname{arcosh} d/R} \ln \frac{(y + d^2 - R^2)^2 + z^2}{(y - d^2 + R^2)^2 + z^2}.$$

3.13

$$\begin{aligned}
V_0(y) &= \begin{cases} V_0 & 0 \leq y \leq \frac{a}{2} \\ -V_0 & \frac{a}{2} \leq y \leq a \end{cases} \\
C_n &= \frac{2}{a} \left(\int_0^{a/2} V_0 \sin \frac{n\pi y}{a} dy - \int_{a/2}^a V_0 \sin \frac{n\pi y}{a} dy \right) \\
&= \frac{2V_0}{n\pi} \left(\cos \frac{n\pi y}{a} \Big|_{a/2}^a - \cos \frac{n\pi y}{a} \Big|_0^{a/2} \right) \\
&= \frac{2V_0}{n\pi} \left(\cos n\pi - \cos \frac{n\pi}{2} - \cos \frac{n\pi}{2} + 1 \right) \\
&= \frac{2V_0}{n\pi} \left(1 + (-1)^n - 2 \cos \frac{n\pi}{2} \right) \\
&= \frac{2V_0}{n\pi} \begin{cases} 0 & n \text{ is odd or divisible by 4} \\ 4 & \text{otherwise} \end{cases} \\
V &= \frac{8V_0}{\pi} \sum_{n=2,6,10,\dots}^{\infty} \frac{1}{n} e^{-n\pi x/a} \sin \frac{n\pi y}{a}
\end{aligned}$$

3.14

$$\begin{aligned}
\sigma &= -\epsilon_0 \frac{\partial V}{\partial x} \\
&= \frac{4\epsilon_0 V_0 \sin \frac{\pi y}{a}}{a \left(1 - \cos \frac{2\pi y}{a} \right)} \\
&= \frac{2\epsilon_0 V_0}{a} \frac{1}{\sin \pi y/a}
\end{aligned}$$

3.15

(a)

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} = 0$$

$$V(0, y) = 0$$

$$V(b, y) = V_0(y)$$

$$V(x, 0) = 0$$

$$V(x, a) = 0$$

$$V = X(x)Y(y)$$

$$X''Y + XY'' = 0$$

$$\frac{X''}{X} + \frac{Y''}{Y} = 0$$

$$\frac{Y''}{Y} = -\alpha^2$$

$$Y'' + \alpha^2 Y = 0$$

$$Y = c_1 \cos \alpha y + c_2 \sin \alpha y$$

$$Y = c_2 \sin \alpha y$$

$$Y = c_2 \sin \frac{n\pi y}{a}, n \in \mathbb{R}$$

$$\frac{X''}{X} = \alpha^2$$

$$X'' - \alpha^2 X = 0$$

$$X = c_3 \cosh \alpha x + c_4 \sinh \alpha x$$

$$X = c_4 \sinh \alpha x$$

$$= c_4 \sinh \frac{n\pi x}{a}, n \in \mathbb{R}$$

$$V = \sum_{n=1}^{\infty} C_n \sinh \frac{n\pi x}{a} \sin \frac{n\pi y}{a}$$

$$V_0(y) = V(b, y)$$

$$= \sum_{n=1}^{\infty} C_n \sinh \frac{n\pi b}{a} \sin \frac{n\pi y}{a}$$

$$C_n \sinh \frac{n\pi b}{a} = \frac{2}{a} \int_0^a V_0(y) \sin \frac{n\pi y}{a} dy$$

$$C_n = \frac{2}{a \sinh n\pi b/a} \int_0^a V_0(y) \sin \frac{n\pi y}{a} dy$$

(b)

$$\begin{aligned} C_n &= \frac{2V_0}{a \sinh n\pi b/a} \int_0^a \sin \frac{n\pi y}{a} dy \\ &= \frac{2V_0}{a \sinh n\pi b/a} \frac{a[1 - (-1)^n]}{n\pi} \\ &= \frac{2V_0[1 - (-1)^n]}{n\pi \sinh n\pi b/a} \\ V &= \frac{2V_0}{\pi} \sum_{n=1}^{\infty} \frac{1 - (-1)^n}{n \sinh n\pi b/a} \sinh \frac{n\pi x}{a} \sin \frac{n\pi y}{a} \end{aligned}$$

3.16

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

$$V(0, y, z) = 0$$

$$V(a, y, z) = 0$$

$$V(x, 0, z) = 0$$

$$V(x, a, z) = 0$$

$$V(x, y, 0) = 0$$

$$V(x, y, a) = V_0$$

$$V = X(x)Y(y)Z(z)$$

$$X''YZ + XY''Z + XYZ'' = 0$$

$$\frac{X''}{X} + \frac{Y''}{Y} + \frac{Z''}{Z} = 0$$

$$\frac{X''}{X} = -\alpha^2$$

$$\frac{Y''}{Y} = -\beta^2$$

$$\frac{Z''}{Z} = \alpha^2 + \beta^2$$

$$X'' + \alpha^2 X = 0$$

$$X = c_1 \cos \alpha x + c_2 \sin \alpha x$$

$$X = c_2 \sin \alpha x$$

$$X = c_2 \sin \frac{n\pi x}{a}, n \in \mathbb{R}$$

$$\frac{Y''}{Y} = -\beta^2$$

$$Y'' + \beta^2 Y = 0$$

$$Y = c_3 \cos \beta y + c_4 \sin \beta y$$

$$Y = c_4 \sin \beta y$$

$$Y = c_4 \sin \frac{m\pi y}{a}, m \in \mathbb{R}$$

$$\frac{Z''}{Z} = \alpha^2 + \beta^2$$

$$Z'' - (\alpha^2 + \beta^2)Z = 0$$

$$Z = c_5 \cosh \sqrt{\alpha^2 + \beta^2} z + c_6 \sinh \sqrt{\alpha^2 + \beta^2} z$$

$$= c_5 \cosh \pi \sqrt{(n/a)^2 + (m/a)^2} z$$

$$+ c_6 \sinh \pi \sqrt{(n/a)^2 + (m/a)^2} z$$

$$Z = c_6 \sinh \pi \sqrt{(n/a)^2 + (m/a)^2} z$$

$$\begin{aligned}
V &= \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} C_{n,m} \sinh \left(\pi \sqrt{(n/a)^2 + (m/a)^2} z \right) \sin \frac{n\pi x}{a} \sin \frac{m\pi y}{a} \\
V_0 &= V(x, y, a) \\
&= \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} C_{n,m} \sinh \left(\pi \sqrt{n^2 + m^2} \right) \sin \frac{n\pi x}{a} \sin \frac{m\pi y}{a} \\
C_{n,m} &= \frac{4V_0}{a^2 \sinh \left(\pi \sqrt{n^2 + m^2} \right)} \int_0^a \int_0^a \sin \frac{n\pi x}{a} \sin \frac{m\pi y}{a} dy dx \\
&= \frac{4V_0}{a^2 \sinh \left(\pi \sqrt{n^2 + m^2} \right)} \frac{a^2 [-1 + (-1)^m] [-1 + (-1)^n]}{nm\pi^2} \\
&= \frac{4V_0 [-1 + (-1)^n] [-1 + (-1)^m]}{nm\pi^2 \sinh \left(\pi \sqrt{n^2 + m^2} \right)} \\
&= \begin{cases} 0 & n \text{ even or } m \text{ even} \\ \frac{16V_0}{nm\pi^2 \sinh \left(\pi \sqrt{n^2 + m^2} \right)} & \text{otherwise} \end{cases} \\
V &= \frac{16V_0}{\pi^2} \sum_{n=1,3,5,\dots}^{\infty} \sum_{m=1,3,5,\dots}^{\infty} \frac{1}{nm} \frac{\sinh \left(\pi \sqrt{n^2 + m^2} z/a \right)}{\sinh \left(\pi \sqrt{n^2 + m^2} \right)} \sin \frac{n\pi x}{a} \sin \frac{m\pi y}{a}
\end{aligned}$$

3.17

$$\begin{aligned}
P_3(x) &= \frac{1}{2^3 3!} \left(\frac{d}{dx} \right)^3 (x^2 - 1)^3 \\
&= \frac{1}{48} \frac{d^3}{dx^3} [(x^2 - 1)^3] \\
&= \frac{1}{48} \frac{d^2}{dx^2} [6x(x^2 - 1)^2] \\
&= \frac{1}{48} \frac{d}{dx} [6(x^2 - 1)^2 + 24x^2(x^2 - 1)] \\
&= \frac{1}{48} [24x(x^2 - 1) + 48x(x^2 - 1) + 48x^3] \\
&= \frac{1}{48} (24x^3 - 24x + 48x^3 - 48x + 48x^3) \\
&= \frac{120}{48} x^3 - \frac{72}{48} x \\
&= \frac{5}{2} x^3 - \frac{3}{2} x \\
\frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) &= -12 \sin \theta \Theta \\
\Theta &= \frac{5}{2} \cos^3 \theta - \frac{3}{2} \cos \theta \\
\frac{d\Theta}{d\theta} &= -\frac{15}{2} \cos^2 \theta \sin \theta + \frac{3}{2} \sin \theta \\
\sin \theta \frac{d\Theta}{d\theta} &= -\frac{15}{2} \cos^2 \theta \sin^2 \theta + \frac{3}{2} \sin^2 \theta \\
\frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) &= \frac{3}{2} (1 - 5 \cos 2\theta) \sin 2\theta \\
\frac{3}{2} (1 - 5 \cos 2\theta) \sin 2\theta &= -12 \sin \theta \left(\frac{5}{2} \cos^3 \theta - \frac{3}{2} \cos \theta \right) \\
3(1 - 5 \cos 2\theta) \sin 2\theta &= 12 \sin \theta \cos \theta (3 - 5 \cos^2 \theta) \\
&= 6(3 - 5 \cos^2 \theta) \sin 2\theta \\
&= 6 \left(3 - 5 \frac{1 + \cos 2\theta}{2} \right) \sin 2\theta \\
&= 3(6 - 5 - 5 \cos 2\theta) \sin 2\theta \\
&= 3(1 - 5 \cos 2\theta) \sin 2\theta \\
\int_{-1}^1 P_1(x) P_3(x) dx &= \int_{-1}^1 x \left(\frac{5}{2} x^3 - \frac{3}{2} x \right) dx \\
&= \left[\frac{1}{2} x^5 - \frac{1}{2} x^3 \right]_{-1}^1 \\
&= \frac{1}{2} - \frac{1}{2} + \frac{1}{2} - \frac{1}{2} \\
&= 0
\end{aligned}$$

3.18

(a)

$$\begin{aligned}
 A_l &= \frac{V_0(2l+1)}{2R^l} \int_0^\pi P_l(\cos \theta) \sin \theta d\theta \\
 &= \begin{cases} V_0 & l = 0 \\ 0 & l \neq 0 \end{cases} \\
 V(r, \theta) &= \sum_{l=0}^{\infty} A_l r^l P_l(\cos \theta) \\
 &= V_0 \\
 B_l &= \frac{V_0(2l+1)}{2} R^{l+1} \int_0^\pi P_l(\cos \theta) \sin \theta d\theta \\
 &= \begin{cases} V_0 R & l = 0 \\ 0 & l \neq 0 \end{cases} \\
 V(r, \theta) &= \frac{V_0 R}{r}
 \end{aligned}$$

(b)

$$\begin{aligned}
 A_l &= \frac{\sigma_0}{2\epsilon_0 R^{l-1}} \int_0^\pi P_l(\cos \theta) \sin \theta d\theta \\
 &= \begin{cases} \frac{R\sigma_0}{\epsilon_0} & l = 0 \\ 0 & l \neq 0 \end{cases} \\
 V(r, \theta) &= \frac{R\sigma_0}{\epsilon_0} \\
 B_l &= A_l R^{2l+1} \\
 &= \begin{cases} \frac{R^2\sigma_0}{\epsilon_0} & l = 0 \\ 0 & l \neq 0 \end{cases} \\
 V(r, \theta) &= \frac{R^2\sigma_0}{\epsilon_0 r}
 \end{aligned}$$

3.19

$$\begin{aligned}
V_0 &= k \cos 3\theta \\
A_l &= \frac{k(2l+1)}{2R^l} \int_0^\pi \cos 3\theta P_l(\cos \theta) \sin \theta d\theta \\
&= \begin{cases} -\frac{3k}{5R} & l=1 \\ \frac{8k}{5R^3} & l=3 \\ 0 & \text{otherwise} \end{cases} \\
V(r, \theta) &= -\frac{3k}{5R} r P_1(\cos \theta) + \frac{8k}{5R^3} r^3 P_3(\cos \theta) \\
&= \frac{kr}{5R} \left[-3P_1(\cos \theta) + \frac{8}{R^2} r^2 P_3(\cos \theta) \right] \\
B_l &= \frac{k(2l+1)}{2} R^{l+1} \int_0^\pi \cos 3\theta P_l(\cos \theta) \sin \theta d\theta \\
&= \begin{cases} -\frac{3kR^2}{5} & l=1 \\ \frac{8kR^4}{5} & l=3 \\ 0 & \text{otherwise} \end{cases} \\
V(r, \theta) &= -\frac{3kR^2}{5r^2} P_1(\cos \theta) + \frac{8kR^4}{5r^4} P_3(\cos \theta) \\
&= \frac{kR^2}{5r^2} \left[\frac{8R^2}{r^2} P_3(\cos \theta) - 3P_1(\cos \theta) \right] \\
\sigma(\theta) &= -\epsilon_0 \left(\frac{\partial V_{\text{above}}}{\partial r} - \frac{\partial V_{\text{below}}}{\partial r} \right) \\
&= \frac{\epsilon_0 k (12 \cos \theta + 35 \cos 3\theta)}{5R}
\end{aligned}$$

3.20

$$\begin{aligned}
V(r, \theta) &= \begin{cases} \sum_{l=0}^{\infty} \frac{2l+1}{2} \frac{r^l}{R^l} \left(\int_0^\pi V_0(\theta) P_l(\cos \theta) \sin \theta d\theta \right) P_l(\cos \theta) & r \leq R \\ \sum_{l=0}^{\infty} \frac{2l+1}{2} \frac{R^{l+1}}{r^{l+1}} \left(\int_0^\pi V_0(\theta) P_l(\cos \theta) \sin \theta d\theta \right) P_l(\cos \theta) & r \geq R \end{cases} \\
\sigma_0 &= -\epsilon_0 \left(\frac{\partial V_{\text{out}}}{\partial r} - \frac{\partial V_{\text{in}}}{\partial r} \right) \Big|_{r=R} \\
&= \frac{\epsilon_0}{2R} \sum_{l=0}^{\infty} (2l+1)^2 C_l P_l(\cos \theta)
\end{aligned}$$

3.21

$$V(r, \theta) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r} - E_0 \left(r - \frac{R^3}{r^2} \right) \cos \theta$$

3.22

(a)

$$\begin{aligned}
\sum_{l=0}^{\infty} \frac{B_l}{r^{l+1}} &= \frac{\sigma}{2\epsilon_0} \left(\sqrt{r^2 + R^2} - r \right) \\
&= \frac{\sigma r}{2\epsilon_0} \left(\sqrt{1 + (R/r)^2} - 1 \right) \\
&= \frac{\sigma r}{2\epsilon_0} \left[\left(1 + \frac{(R/r)^2}{2} - \frac{(R/r)^4}{8} + \dots \right) - 1 \right] \\
&= \frac{\sigma}{2\epsilon_0} \left(\frac{R^2}{2r} - \frac{R^4}{8r^3} + \dots \right) \\
B_0 &= \frac{\sigma R^2}{4\epsilon_0} \\
B_1 &= 0 \\
B_2 &= -\frac{\sigma R^4}{16\epsilon_0}
\end{aligned}$$

(b)

$$\begin{aligned}
\sum_{l=0}^{\infty} A_l r^l &= \frac{\sigma}{2\epsilon_0} \left(\sqrt{r^2 + R^2} - r \right) \\
&= \frac{\sigma}{2\epsilon_0} \left(R \sqrt{1 + (r/R)^2} - r \right) \\
&= \frac{\sigma}{2\epsilon_0} \left[R \left(1 + \frac{(r/R)^2}{2} - \frac{(r/R)^4}{8} + \dots \right) - r \right] \\
&= \frac{\sigma}{2\epsilon_0} \left(R - r + \frac{r^2}{2R} - \frac{r^4}{8R^3} + \dots \right) \\
A_0 &= \frac{\sigma R}{2\epsilon_0} \\
A_1 &= -\frac{\sigma}{2\epsilon_0} \\
A_2 &= \frac{\sigma}{4\epsilon_0 R} \\
A'_0 &= A_0 \\
A'_1 &= -A_1 \\
A'_2 &= A_2
\end{aligned}$$

3.23

$$V(r, \theta) = \sum_{l=1}^{\infty} A_l r^l P_l(\cos \theta)$$

$$A_l = \frac{\sigma_0}{2\epsilon_0 R^{l-1}} \left(\int_0^{\pi/2} P_l(\cos \theta) \sin \theta d\theta - \int_{\pi/2}^{\pi} P_l(\cos \theta) \sin \theta d\theta \right)$$

$$A_0 = 0$$

$$A_1 = \frac{\sigma_0}{2\epsilon_0}$$

$$A_2 = 0$$

$$A_3 = -\frac{\sigma_0}{8\epsilon_0 R^2}$$

$$A_4 = 0$$

$$A_5 = \frac{\sigma_0}{16\epsilon_0 R^4}$$

$$A_6 = 0$$

$$V(r, \theta) = \sum_{l=0}^{\infty} \frac{B_l}{r^{l+1}} P_l(\cos \theta)$$

$$B_l = A_l R^{2l+1}$$

$$B_0 = 0$$

$$B_1 = \frac{\sigma_0 R^3}{2\epsilon_0}$$

$$B_2 = 0$$

$$B_3 = -\frac{\sigma_0 R^5}{8\epsilon_0}$$

$$B_4 = 0$$

$$B_5 = \frac{\sigma_0 R^7}{16\epsilon_0}$$

$$B_6 = 0$$

3.24

$$\begin{aligned} 0 &= \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial V}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2} \\ &= \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial V}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 V}{\partial \phi^2} \end{aligned}$$

$$V(s, \phi) = S(s)\Phi(\phi)$$

$$\begin{aligned} 0 &= \frac{1}{s} \frac{\partial}{\partial s} (sS')\Phi + \frac{1}{s^2} S\Phi'' \\ &= \frac{1}{s} (S' + sS'')\Phi + \frac{1}{s^2} S\Phi'' \\ &= \frac{S'}{sS} + \frac{S''}{S} + \frac{\Phi''}{s^2\Phi} \\ &= \frac{s^2S'' + sS'}{S} + \frac{\Phi''}{\Phi} \end{aligned}$$

$$\frac{\Phi''}{\Phi} = 0$$

$$\Phi'' = 0$$

$$\Phi = c_1 + c_2\phi$$

$$\frac{\Phi''}{\Phi} = -n^2$$

$$\Phi'' + \alpha^2\Phi = 0$$

$$\Phi = c_3 \cos \alpha\phi + c_4 \sin \alpha\phi$$

$$\Phi(0) = \Phi(2\pi)$$

$$c_1 = c_3 \cos 2\pi\alpha + c_4 \sin 2\pi\alpha$$

$$\alpha = n, n \in \mathbb{R}$$

$$\Phi = c_3 \cos n\phi + c_4 \sin n\pi$$

$$\frac{s^2S'' + sS'}{S} = 0$$

$$s^2S'' + sS' = 0$$

$$sS'' + S' = 0$$

$$S = c_5 + c_6 \ln s$$

$$\begin{aligned}
\frac{s^2 S'' + s S'}{S} &= n^2 \\
s^2 S'' + s S' - n^2 S &= 0 \\
S &= s^m \\
S' &= m s^{m-1} \\
S'' &= m(m-1) s^{m-2} \\
m(m-1) s^m + m s^m - n^2 s^m &= 0 \\
m^2 - m + m - n^2 &= 0 \\
m^2 - n^2 &= 0 \\
(m+n)(m-n) &= 0 \\
S &= c_7 s^n + c_8 s^{-n} \\
V &= S(s) \Phi(\phi) \\
&= (c_1 + c_2 \phi)(c_5 + c_6 \ln s) \\
&\quad + \sum_{n=1}^{\infty} (c_7 s^n + c_8 s^{-n})(c_3 \cos n\phi + c_4 \sin n\phi) \\
&= c_1 + c_2 \ln s \\
&\quad + \sum_{n=1}^{\infty} [s^n (A_n \cos n\phi + B_n \sin n\phi) + s^{-n} (C_n \cos n\phi + D_n \sin n\phi)]
\end{aligned}$$

3.25

$$\begin{aligned}
V &= 0 \text{ at } s = R \\
V &\rightarrow -E_0 s \cos \phi \text{ as } s \rightarrow \infty \\
V &= \left(A_1 s + \frac{C_1}{s} \right) \cos \phi \\
0 &= A_1 R + \frac{C_1}{R} \\
C_2 &= -A_1 R^2 \\
A_1 &= -E_0 \\
V &= E_0 s \left(\frac{R^2}{s^2} - 1 \right) \cos \phi \\
\sigma &= -\epsilon_0 \left(\frac{\partial V_{\text{out}}}{\partial s} - \frac{\partial V_{\text{in}}}{\partial s} \right) \Big|_{s=R} \\
&= -\epsilon_0 \frac{\partial V_{\text{out}}}{\partial s} \Big|_{s=R} \\
&= 2\epsilon_0 E_0 \cos \phi
\end{aligned}$$