# Quantum Computation and Quantum Information by Michael A. Nielsen and Isaac L. Chuang

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# Part I

# Fundamental concepts

# 1 Introduction and overview

#### 1.2 Quantum bits

- The special states  $|0\rangle$  and  $|1\rangle$  form an orthonormal basis and are known as **computational basis states**.
- $\bullet$  A quantum bit (qubit) is a linear combination of the computational basis states

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$$

where  $\alpha$  and  $\beta$  are complex numbers.

- When we measure a qubit we either get  $|0\rangle$  with probability  $|\alpha|^2$  or  $|1\rangle$  with probability  $|\beta|^2$ . Thus,  $|\alpha|^2 + |\beta|^2 = 1$  and a qubit can be thought of as a unit vector in a two-dimensional complex vector space.
- If a qubit is in the state

$$|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$

there's a 50/50 chance of measuring  $|0\rangle$  or  $|1\rangle$ .

• If we let

$$\alpha = e^{i\gamma} \cos \frac{\theta}{2}$$

and

$$\beta = e^{i\gamma} e^{i\varphi} \sin\frac{\theta}{2}$$

then

$$|\alpha|^2 + |\beta|^2 = \alpha^* \alpha + \beta^* \beta$$
$$= \cos^2 \frac{\theta}{2} + \sin^2 \frac{\theta}{2}$$
$$= 1$$

so the qubit is still normalised and it can be written

$$\left|\psi\right\rangle = e^{i\gamma} \left(\cos\frac{\theta}{2}\left|0\right\rangle + e^{i\varphi}\sin\frac{\theta}{2}\left|1\right\rangle\right).$$

It turns out that  $e^{i\gamma}$  has no observable effects and we can effectively write

$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1\rangle$$
.

This defines a point on a three-dimensional sphere known as the **Bloch** sphere where  $\theta$  and  $\varphi$  take on their usual roles in a spherical coordinate system.

• Before measurement a qubit is in a linear combination of  $|0\rangle$  and  $|1\rangle$  but when measured you get one or the other and the state of the system changes to match the measured result.

#### 1.2.1 Multiple Bits

• A two qubit system has four computational basis state  $|00\rangle$ ,  $|01\rangle$ ,  $|10\rangle$ , and  $|11\rangle$  so the general expression for the state of such a system is

$$|\psi\rangle = \alpha_{00} |00\rangle + \alpha_{01} |01\rangle + \alpha_{10} |10\rangle + \alpha_{11} |11\rangle.$$

• If you were to measure the first qubit, you would get  $|0\rangle$  with probability  $|\alpha_{00}|^2 + |\alpha_{01}|^2$  and the system would be left in the state

$$|\psi\rangle = \frac{\alpha_{00}|00\rangle + \alpha_{01}|01\rangle}{\sqrt{|\alpha_{00}|^2 + |\alpha_{01}|^2}},$$

i.e. it is renormalised such that the normalisation condition still holds.

## 1.3 Quantum Computation

#### 1.3.1 Single Qubit Gates

• The quantum NOT changes  $|0\rangle$  to  $|1\rangle$  and  $|1\rangle$  to  $|0\rangle$ . It acts linearly on superpositions of those states, i.e. it turns  $\alpha \, |0\rangle + \beta \, |1\rangle$  into  $\beta \, |0\rangle + \alpha \, |1\rangle$ . If a quantum state  $|\psi\rangle = \alpha \, |0\rangle + \beta \, |1\rangle$  is written in vector notation as

$$|\psi\rangle = \begin{bmatrix} \alpha \\ \beta \end{bmatrix}$$

then the quantum NOT gate can be expressed in matrix form as

$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}.$$

- In order to preserve the normalisation condition, matrix representations of quantum gates must be unitary, i.e.  $M^{\dagger}M = I$  where I is the identity matrix.
- An arbitrary unitary 2x2 matrix can be decomposed into a finite set of other 2x2 matrices. This means an arbitrary single qubit gate can be generated by a finite set of other gates.

#### 1.3.2 Multiply Qubit Gates

• The controlled-NOT or CNOT gate is a multi-qubit gate that has two input qubits known as the control qubit and the target qubit. If the control qubit is set to  $|0\rangle$  the target qubit is left alone, but if it's set to  $|1\rangle$  the target qubit is flipped. Another way of writing this is  $|A,B\rangle \to |A,A\oplus B\rangle$  where  $\oplus$  is modulo-two addition. Yet another way of writing this is in matrix form

$$U_{CN} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

where the first column describes what happens to the  $|00\rangle$  basis state, etc.

• Other classical gates like NAND or XOR can't be represented as quantum gates as they're irreversibile. For example, given the output  $A \oplus B$  from an XOR gate it's not possible to determine what the inputs A and B were.

 Any multiple qubit logic gate may be composed from CNOT and single qubit gates.

#### 1.3.3 Measurements in bases other than the computational basis

• Given any basis states  $|a\rangle$  and  $|b\rangle$  it is possible to express an arbitrary state as a linear combination  $\alpha\,|a\rangle\,+\,\beta\,|b\rangle$  of those states. For example, if  $|a\rangle\,=\,|+\rangle\,=\,(|0\rangle\,+\,|1\rangle)/\sqrt{2}$  and  $|b\rangle\,=\,|-\rangle\,=\,(|0\rangle\,-\,|1\rangle)/\sqrt{2}$  then

$$\left|\phi\right\rangle = \alpha\left|0\right\rangle + \beta\left|1\right\rangle = \alpha\frac{\left|+\right\rangle + \left|-\right\rangle}{\sqrt{2}} + \beta\frac{\left|0\right\rangle - \left|1\right\rangle}{\sqrt{2}} = \frac{\alpha + \beta}{\sqrt{2}}\left|+\right\rangle + \frac{\alpha - \beta}{\sqrt{2}}\left|-\right\rangle.$$

• If  $|a\rangle$  and  $|b\rangle$  are orthonormal, it's also possible to perform a measurement with respect to that basis. In the example above you would measure  $|+\rangle$  with probability  $|\alpha + \beta|^2/2$  and  $|-\rangle$  with probability  $|\alpha - \beta|^2/2$ .