

# Introduction to Quantum Mechanics by David J. Griffiths Problems

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## Contents

|          |  |           |
|----------|--|-----------|
| <b>I</b> | <b>Theory</b>                                | <b>5</b>  |
| <b>1</b> | <b>The Wave Function</b>                     | <b>5</b>  |
| 1.1      | .....  | 5         |
| 1.2      | .....  | 6         |
| 1.3      | .....  | 7         |
| 1.4      | .....  | 8         |
| 1.5      | .....  | 9         |
| 1.6      | .....  | 10        |
| 1.8      | .....  | 10        |
| 1.9      | .....  | 11        |
| 1.10     | .....  | 13        |
| 1.14     | .....  | 14        |
| 1.15     | .....  | 15        |
| 1.16     | .....  | 15        |
| 1.18     | .....  | 17        |
| <b>2</b> | <b>Time-Independent Schrödinger Equation</b> | <b>17</b> |
| 2.1      | .....  | 17        |
| 2.2      | .....  | 18        |
| 2.3      | .....  | 18        |
| 2.4      | .....  | 19        |
| 2.5      | .....  | 20        |
| 2.6      | .....  | 21        |
| 2.7      | .....  | 22        |
| 2.8      | .....  | 24        |
| 2.9      | .....  | 24        |
| 2.10     | .....  | 25        |
| 2.11     | .....  | 26        |
| 2.12     | .....  | 29        |

|                    |           |
|--------------------|-----------|
| 2.13               | 30        |
| 2.14               | 32        |
| 2.15               | 32        |
| 2.17               | 33        |
| 2.18               | 33        |
| 2.20               | 33        |
| 2.21               | 34        |
| 2.22               | 35        |
| 2.26               | 36        |
| 2.29               | 36        |
| 2.30               | 37        |
| 2.31               | 38        |
| 2.34               | 39        |
| 2.35               | 42        |
| 2.36               | 43        |
| 2.37               | 44        |
| 2.38               | 45        |
| 2.39               | 46        |
| 2.40               | 47        |
| 2.41               | 49        |
| 2.43               | 50        |
| 2.44               | 51        |
| 2.45               | 51        |
| 2.53               | 51        |
| <b>3 Formalism</b> | <b>52</b> |
| 3.1                | 52        |
| 3.2                | 53        |
| 3.4                | 53        |
| 3.5                | 54        |
| 3.6                | 55        |
| 3.7                | 56        |
| 3.8                | 57        |
| 3.9                | 57        |
| 3.11               | 58        |
| 3.12               | 58        |
| 3.14               | 59        |
| 3.15               | 61        |
| 3.17               | 61        |
| 3.18               | 62        |
| 3.19               | 63        |
| 3.20               | 63        |
| 3.23               | 65        |
| 3.24               | 65        |
| 3.25               | 66        |
| 3.26               | 66        |

|  |           |
|--|-----------|
| 3.27   | 68        |
| 3.28   | 69        |
| 3.30   | 69        |
| 3.31   | 70        |
| 3.32   | 72        |
| 3.33   | 73        |
| 3.34   | 74        |
| 3.35   | 75        |
| 3.36   | 75        |
| 3.37   | 77        |
| 3.38   | 79        |
| 3.41   | 80        |
| 3.43   | 81        |
| 3.44   | 81        |
| 3.45   | 82        |
| 3.46   | 83        |
| <b>4 Quantum Mechanics in Three Dimensions</b> | <b>85</b> |
| 4.1  | 85        |
| 4.2  | 87        |
| 4.3  | 90        |
| 4.4  | 91        |
| 4.5  | 92        |
| 4.6  | 92        |
| 4.7  | 93        |
| 4.9  | 94        |
| 4.10   | 95        |
| 4.12   | 96        |
| 4.13   | 97        |
| 4.14   | 98        |
| 4.15   | 99        |
| 4.16   | 99        |
| 4.18   | 100       |
| 4.19   | 101       |
| 4.20   | 101       |
| 4.21   | 102       |
| 4.22   | 103       |
| 4.24   | 104       |
| 4.25   | 105       |
| 4.26   | 107       |
| 4.28   | 107       |
| 4.29   | 107       |
| 4.30   | 108       |
| 4.31   | 110       |
| 4.32   | 111       |
| 4.33   | 112       |

|                              |            |
|------------------------------|------------|
| 4.35                         | 112        |
| 4.36                         | 113        |
| 4.37                         | 115        |
| 4.38                         | 116        |
| 4.40                         | 116        |
| 4.46                         | 117        |
| 4.48                         | 118        |
| 4.52                         | 118        |
| 4.53                         | 119        |
| 4.54                         | 120        |
| 4.55                         | 121        |
| 4.57                         | 123        |
| 4.58                         | 124        |
| 4.61                         | 125        |
| 4.64                         | 125        |
| 4.70                         | 126        |
| 4.71                         | 128        |
| <b>5 Identical Particles</b> | <b>129</b> |
| 5.1                          | 129        |
| 5.2                          | 131        |
| 5.4                          | 132        |
| 5.5                          | 133        |
| 5.6                          | 134        |
| 5.7                          | 134        |
| 5.8                          | 135        |
| 5.13                         | 135        |
| 5.14                         | 136        |
| 5.15                         | 137        |
| 5.16                         | 137        |
| 5.17                         | 138        |
| 5.18                         | 138        |
| 5.19                         | 139        |
| 5.20                         | 139        |
| 5.21                         | 139        |
| 5.22                         | 140        |
| 5.23                         | 140        |
| 5.24                         | 140        |
| 5.25                         | 141        |
| 5.29                         | 141        |
| 5.30                         | 142        |
| 5.32                         | 142        |

|          |   |            |
|----------|---|------------|
| <b>6</b> | <b>Symmetries &amp; Conservation Laws</b> | <b>143</b> |
| 6.1      | .....                                     | 143        |
| 6.2      | .....                                     | 143        |
| 6.3      | .....                                     | 143        |
| 6.4      | .....                                     | 144        |
| 6.5      | .....                                     | 144        |
| 6.7      | .....                                     | 144        |
| 6.8      | .....                                     | 145        |
| 6.9      | .....                                     | 145        |
| 6.10     | .....                                     | 146        |
| 6.11     | .....                                     | 146        |
| 6.12     | .....                                     | 147        |
| 6.13     | .....                                     | 147        |
| 6.15     | .....                                     | 148        |
| 6.16     | .....                                     | 149        |
| 6.18     | .....                                     | 149        |
| 6.19     | .....                                     | 151        |
| 6.20     | .....                                     | 151        |
| 6.21     | .....                                     | 152        |
| 6.22     | .....                                     | 153        |

## Part I

# Theory

## 1 The Wave Function

### 1.1

(a)

$$\begin{aligned}
 \langle j^2 \rangle &= \sum j^2 P(j) \\
 &= 14^2 \frac{1}{14} + 15^2 \frac{1}{14} + 16^2 \frac{3}{14} + 22^2 \frac{2}{14} + 24^2 \frac{2}{14} + 25^2 \frac{5}{14} \\
 &= \frac{3217}{7} \\
 &\approx 459.571 \\
 \langle j \rangle^2 &= \left( \sum j P(j) \right)^2 \\
 &= 441
 \end{aligned}$$

(b)

$$\begin{aligned}
 \Delta j_{14} &= -7 \\
 \Delta j_{15} &= -6 \\
 \Delta j_{16} &= -5 \\
 \Delta j_{22} &= 1 \\
 \Delta j_{24} &= 3 \\
 \Delta j_{25} &= 4 \\
 \sigma^2 &= \sum (\Delta j)^2 P(j) \\
 &= \frac{130}{7} \\
 &\approx 18.571
 \end{aligned}$$

(c)

$$\sigma^2 = \sqrt{\langle j^2 \rangle - \langle j \rangle^2} = 18.571$$

## 1.2

(a)

$$\begin{aligned}
 \langle x^2 \rangle &= \int_0^h x^2 \rho(x) dx \\
 &= \int_0^h \frac{x^{3/2}}{2\sqrt{h}} dx \\
 &= \frac{1}{2\sqrt{h}} \left[ \frac{2}{5} x^{5/2} \right]_0^h \\
 &= \frac{h^2}{5} \\
 \langle x \rangle^2 &= \frac{h^2}{9} \\
 \sigma &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\
 &= \sqrt{\frac{h^2}{5} - \frac{h^2}{9}} \\
 &= h \sqrt{\frac{4}{45}} \\
 &= \frac{2}{3\sqrt{5}} h
 \end{aligned}$$

(b)

$$\begin{aligned}
1 - \int_{\langle x \rangle - \sigma}^{\langle x \rangle + \sigma} \rho(x) dx &= 1 - \frac{1}{2\sqrt{h}} [2\sqrt{x}]_{\langle x \rangle - \sigma}^{\langle x \rangle + \sigma} \\
&= 1 - \frac{1}{\sqrt{h}} \left( \sqrt{\frac{1}{3}h + \frac{2}{3\sqrt{5}}h} - \sqrt{\frac{1}{3}h - \frac{2}{3\sqrt{5}}h} \right) \\
&= 1 - \left( \sqrt{\frac{1}{3} + \frac{2}{3\sqrt{5}}} - \sqrt{\frac{1}{3} - \frac{2}{3\sqrt{5}}} \right) \\
&\approx 0.393
\end{aligned}$$

### 1.3

(a)

$$\begin{aligned}
\rho(x) &= Ae^{-\lambda(x-a)^2} \\
1 &= \int_{-\infty}^{\infty} \rho(x) dx \\
&= A \int_{-\infty}^{\infty} e^{-\lambda(x-a)^2} dx \\
&= A \sqrt{\frac{\pi}{\lambda}} \\
A &= \sqrt{\frac{\lambda}{\pi}}
\end{aligned}$$

(b)

$$\begin{aligned}
\langle x \rangle &= \sqrt{\frac{\lambda}{\pi}} \int_{-\infty}^{\infty} x e^{-\lambda(x-a)^2} dx \\
&= a \\
\langle x^2 \rangle &= \sqrt{\frac{\lambda}{\pi}} \int_{-\infty}^{\infty} x^2 e^{-\lambda(x-a)^2} dx \\
&= a^2 + \frac{1}{2\lambda} \\
\sigma &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\
&= \sqrt{a^2 + \frac{1}{2\lambda} - a^2} \\
&= \frac{1}{\sqrt{2\lambda}}
\end{aligned}$$

## 1.4

(a)

$$\begin{aligned}
 1 &= \int_{-\infty}^{\infty} |\Psi(x, 0)|^2 dx \\
 &= \left(\frac{A}{a}\right)^2 \int_0^a x^2 dx + \left(\frac{A}{b-a}\right)^2 \int_a^b (b-x)^2 dx \\
 &= \frac{1}{3}A^2a + \left(\frac{A}{b-a}\right)^2 \left[-\frac{1}{3}(b-x)^3\right]_a^b \\
 &= \frac{1}{3}A^2a + \frac{1}{3}A^2(b-a) \\
 &= \frac{1}{3}A^2b \\
 A &= \sqrt{\frac{3}{b}}
 \end{aligned}$$

(c)  $x = a$

(d)

$$\begin{aligned}
 \int_0^a |\Psi(x, 0)|^2 dx &= \frac{3}{a^2b} \left[\frac{1}{3}x^3\right]_0^a \\
 &= \frac{a}{b}
 \end{aligned}$$

(e)

$$\begin{aligned}
 \langle x \rangle &= \int_{-\infty}^{\infty} x |\Psi(x, 0)|^2 dx \\
 &= \frac{3}{a^2b} \left[\frac{1}{4}x^4\right]_0^a + \frac{3}{b(b-a)^2} \int_a^b x(b-x)^2 dx \\
 &= \frac{3a^2}{4b} + \frac{3}{b(b-a)^2} \int_a^b (b^2x - 2bx^2 + x^3) dx \\
 &= \frac{3a^2}{4b} + \frac{3}{b(b-a)^2} \left[\frac{1}{2}b^2x^2 - \frac{2}{3}bx^3 + \frac{1}{4}x^4\right]_a^b \\
 &= \frac{3a^2}{4b} + \frac{3}{b(b-a)^2} \left(\frac{1}{2}b^4 - \frac{2}{3}b^4 + \frac{1}{4}b^4 - \frac{1}{2}a^2b^2 + \frac{2}{3}a^3b - \frac{1}{4}a^4\right) \\
 &= \frac{3a^2}{4b} + \frac{3}{b(b-a)^2} \frac{1}{12}(b-a)^3(3a+b) \\
 &= \frac{3a^2}{4b} + \frac{1}{4b}(3ab + b^2 - 3a^2 - ab) \\
 &= \frac{1}{2}a + \frac{1}{4}b
 \end{aligned}$$



## 1.5

(a)

$$\Psi(x, t) = Ae^{-\lambda|x|}e^{-i\omega t}$$

$$\Psi(x, 0) = Ae^{-\lambda|x|}$$

$$\begin{aligned} 1 &= A^2 \int_{-\infty}^{\infty} e^{-2\lambda|x|} dx \\ &= 2A^2 \int_0^{\infty} e^{-2\lambda x} dx \\ &= 2A^2 \left[ -\frac{1}{2\lambda} e^{-2\lambda x} \right]_0^{\infty} \\ &= \frac{A^2}{\lambda} \\ A &= \sqrt{\lambda} \end{aligned}$$

(b)

$$\begin{aligned} \langle x \rangle &= \int_{-\infty}^{\infty} x \lambda e^{-2\lambda|x|} dx \\ &= \lambda \int_{-\infty}^{\infty} x e^{-2\lambda|x|} dx \\ &= 0 \\ \langle x^2 \rangle &= \int_{-\infty}^{\infty} x^2 \lambda e^{-2\lambda|x|} dx \\ &= 2\lambda \int_0^{\infty} x^2 e^{-2\lambda x} dx \\ &= \frac{1}{2\lambda^2} \end{aligned}$$

(c)

$$\begin{aligned} \sigma &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\ &= \frac{1}{\sqrt{2}\lambda} \\ 1 - \int_{-\sigma}^{\sigma} \lambda e^{-2\lambda|x|} dx &= 1 - 2\lambda \int_0^{\sigma} e^{-2\lambda x} dx \\ &= 1 - 2\lambda \left[ -\frac{1}{2\lambda} e^{-2\lambda x} \right]_0^{\sigma} \\ &= e^{-2\lambda\sigma} \\ &= e^{-\sqrt{2}} \\ &\approx 0.243 \end{aligned}$$

## 1.6

The chain rule requires that you apply it to both  $x$  and  $|\Psi|^2$  which gives the same result

$$\begin{aligned}
 \frac{d\langle x \rangle}{dt} &= \frac{d}{dt} \int x |\Psi|^2 dx \\
 &= \int \frac{d}{dt} (x |\Psi|^2) dx \\
 &= \int \left( 0 \cdot |\Psi|^2 + x \frac{\partial |\Psi|^2}{\partial t} \right) dx \\
 &= \int x \frac{\partial |\Psi|^2}{\partial t} dx
 \end{aligned}$$

## 1.8

$$\begin{aligned}
 i\hbar \frac{\partial}{\partial t} \left( e^{-iV_0 t/\hbar} \Psi \right) &= -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \left( e^{-iV_0 t/\hbar} \Psi \right) + (V + V_0) \left( e^{-iV_0 t/\hbar} \Psi \right) \\
 i\hbar \left( -\frac{iV_0}{\hbar} e^{-iV_0 t/\hbar} \Psi + e^{-iV_0 t/\hbar} \frac{\partial \Psi}{\partial t} \right) &= -\frac{\hbar^2}{2m} e^{-iV_0 t/\hbar} \frac{\partial^2 \Psi}{\partial x^2} + V e^{-iV_0 t/\hbar} \Psi + V_0 e^{-iV_0 t/\hbar} \Psi \\
 V_0 \Psi + i\hbar \frac{\partial \Psi}{\partial t} &= -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V \Psi + V_0 \Psi \\
 i\hbar \frac{\partial \Psi}{\partial t} &= -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V \Psi
 \end{aligned}$$

$$\begin{aligned}
 \langle Q(x, p) \rangle &= \int \left( e^{-iV_0 t/\hbar} \Psi \right)^* [Q(x, -i\hbar \partial/\partial x)] e^{-iV_0 t/\hbar} \Psi dx \\
 &= \int e^{iV_0 t/\hbar} \Psi^* [Q(x, -i\hbar \partial/\partial x)] e^{-iV_0 t/\hbar} \Psi dx \\
 &= \int \Psi^* [Q(x, -i\hbar \partial/\partial x)] \Psi dx
 \end{aligned}$$

No effect on the expectation value.

## 1.9

(a)

$$\begin{aligned}
 \Psi(x, t) &= Ae^{-a[(mx^2/\hbar)+it]} \\
 1 &= A^2 \int_{-\infty}^{\infty} e^{-2a(mx^2/\hbar)} dx \\
 &= A^2 \int_{-\infty}^{\infty} e^{-2a(mx^2/\hbar)} dx \\
 &= A^2 \sqrt{\frac{\pi\hbar}{2am}} \\
 A^2 &= \sqrt{\frac{2am}{\pi\hbar}} \\
 A &= \left(\frac{2am}{\pi\hbar}\right)^{1/4}
 \end{aligned}$$

(b)

$$\begin{aligned}
 \Psi &= Ae^{-a[(mx^2/\hbar)+it]} \\
 \frac{\partial \Psi}{\partial t} &= -ia\Psi \\
 \frac{\partial \Psi}{\partial x} &= -\frac{2amx}{\hbar}\Psi \\
 \frac{\partial^2 \Psi}{\partial x^2} &= -\frac{2am}{\hbar} \left( \Psi + x \frac{\partial \Psi}{\partial x} \right) \\
 &= -\frac{2am}{\hbar} \left( 1 - \frac{2amx^2}{\hbar} \right) \Psi \\
 V\Psi &= i\hbar \frac{\partial \Psi}{\partial t} + \frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} \\
 &= a\hbar\Psi - a\hbar \left( 1 - \frac{2amx^2}{\hbar} \right) \Psi \\
 V &= a\hbar - a\hbar + 2a^2mx^2 \\
 &= 2a^2mx^2
 \end{aligned}$$

(c)

$$\begin{aligned}
\langle x \rangle &= A^2 \int_{-\infty}^{\infty} e^{-2a(mx^2/\hbar)} x \, dx \\
&= 0 \\
\langle x^2 \rangle &= A^2 \int_{-\infty}^{\infty} e^{-2a(mx^2/\hbar)} x^2 \, dx \\
&= 2A^2 \int_0^{\infty} e^{-2a(mx^2/\hbar)} x^2 \, dx \\
&= \frac{\hbar}{4am} \\
\langle p \rangle &= \int_{-\infty}^{\infty} \Psi^* \left[ -i\hbar \frac{\partial}{\partial x} \right] \Psi \, dx \\
&= -i\hbar \int_{-\infty}^{\infty} A e^{-a[(mx^2/\hbar)-it]} \left( -\frac{2amx}{\hbar} A e^{-a[(mx^2/\hbar)+it]} \right) dx \\
&= 2iA^2 am \int_{-\infty}^{\infty} x e^{-2amx^2/\hbar} \, dx \\
&= 0 \\
\langle p^2 \rangle &= \int_{-\infty}^{\infty} \Psi^* \left[ -\hbar^2 \frac{\partial^2}{\partial x^2} \right] \Psi \, dx \\
&= -\hbar^2 \int_{-\infty}^{\infty} A e^{-a[(mx^2/\hbar)-it]} \left[ -\frac{2am}{\hbar} \left( 1 - \frac{2amx^2}{\hbar} \right) A e^{-a[(mx^2/\hbar)+it]} \right] dx \\
&= 2A^2 am\hbar \int_{-\infty}^{\infty} e^{-2amx^2/\hbar} \left( 1 - \frac{2amx^2}{\hbar} \right) dx \\
&= am\hbar
\end{aligned}$$

(d)

$$\begin{aligned}
\sigma_x &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\
&= \sqrt{\frac{\hbar}{4am}} \\
\sigma_p &= \sqrt{am\hbar} \\
\sigma_x \sigma_p &= \sqrt{\frac{1}{4} \hbar^2} \\
&= \frac{1}{2} \hbar \\
&\geq \frac{1}{2} \hbar
\end{aligned}$$

Yes, this is consistent with Heisenberg's uncertainty principle.

### 1.10

(a)

$$P(0) = 0$$

$$\begin{aligned} P(1) &= \frac{2}{25} \\ &= 0.08 \end{aligned}$$

$$\begin{aligned} P(2) &= \frac{3}{25} \\ &= 0.12 \end{aligned}$$

$$\begin{aligned} P(3) &= \frac{1}{5} \\ &= 0.2 \end{aligned}$$

$$\begin{aligned} P(4) &= \frac{3}{25} \\ &= 0.12 \end{aligned}$$

$$\begin{aligned} P(5) &= \frac{3}{25} \\ &= 0.2 \end{aligned}$$

$$\begin{aligned} P(6) &= \frac{3}{25} \\ &= 0.2 \end{aligned}$$

$$\begin{aligned} P(7) &= \frac{1}{25} \\ &= 0.04 \end{aligned}$$

$$\begin{aligned} P(8) &= \frac{2}{25} \\ &= 0.08 \end{aligned}$$

$$\begin{aligned} P(9) &= \frac{3}{25} \\ &= 0.12 \end{aligned}$$

(b) The most probable digit is 3, the median digit is 4, and the average value is  $\frac{118}{25} = 4.72$ .

(c)  $\sigma = 2.474$

### 1.14

(a)

$$\begin{aligned}
 P_{ab}(t) &= \int_a^b |\Psi(x, t)|^2 dx \\
 \frac{dP_{ab}}{dt} &= \frac{d}{dt} \int_a^b |\Psi(x, t)|^2 dx \\
 &= \int_a^b \frac{d}{dt} (|\Psi(x, t)|^2) dx \\
 &= \int_a^b \frac{\partial}{\partial x} \left[ \frac{i\hbar}{2m} \left( \Psi^* \frac{\partial \Psi}{\partial x} - \frac{\partial \Psi^*}{\partial x} \Psi \right) \right] dx \\
 &= J(a, t) - J(b, t)
 \end{aligned}$$

The units are  $s^{-1}$ .

(b)

$$\begin{aligned}
 \Psi(x, t) &= Ae^{-a[(mx^2/\hbar)+it]} \\
 \frac{\partial \Psi}{\partial x} &= -\frac{2amx}{\hbar} \Psi \\
 \Psi^*(x, t) &= Ae^{-a[(mx^2/\hbar)-it]} \\
 \frac{\partial \Psi^*}{\partial x} &= -\frac{2amx}{\hbar} \Psi^* \\
 J(x, t) &= \frac{i\hbar}{2m} \left( \Psi \frac{\partial \Psi^*}{\partial x} - \Psi^* \frac{\partial \Psi}{\partial x} \right) \\
 &= \frac{i\hbar}{2m} \left[ \Psi \left( -\frac{2amx}{\hbar} \Psi^* \right) - \Psi^* \left( -\frac{2amx}{\hbar} \Psi \right) \right] \\
 &= 0
 \end{aligned}$$

1.15

$$\begin{aligned}
\frac{d}{dt} \int_{-\infty}^{\infty} \Psi_1^* \Psi_2 dx &= \int_{-\infty}^{\infty} \left( \frac{\partial \Psi_1^*}{\partial t} \Psi_2 + \Psi_1^* \frac{\partial \Psi_2}{\partial t} \right) dx \\
&= \int_{-\infty}^{\infty} \left[ \left( -i \frac{\hbar}{2m} \frac{\partial^2 \Psi_1^*}{\partial x^2} + i \frac{V}{\hbar} \Psi_1^* \right) \Psi_2 \right. \\
&\quad \left. + \Psi_1^* \left( i \frac{\hbar}{2m} \frac{\partial^2 \Psi_2}{\partial x^2} - i \frac{V}{\hbar} \Psi_2 \right) \right] dx \\
&= i \frac{\hbar}{2m} \int_{-\infty}^{\infty} \left( \Psi_1^* \frac{\partial^2 \Psi_2}{\partial x^2} - \frac{\partial^2 \Psi_1^*}{\partial x^2} \Psi_2 \right) dx \\
&= i \frac{\hbar}{2m} \left[ \Psi_1^* \frac{\partial \Psi_2}{\partial x} \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{\partial}{\partial x} (\Psi_1^* \Psi_2) dx \right. \\
&\quad \left. - \frac{\partial \Psi_1^*}{\partial x} \Psi_2 \Big|_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \frac{\partial}{\partial x} (\Psi_1^* \Psi_2) dx \right] \\
&= 0
\end{aligned}$$

1.16

(a)

$$\begin{aligned}
1 &= \int_{-a}^a A^2 (a^2 - x^2)^2 dx \\
&= A^2 \int_0^a (a^2 - x^2)^2 dx \\
&= \frac{16}{15} A^2 a^5 \\
A &= \sqrt{\frac{15}{16a^5}}
\end{aligned}$$

(b)

$$\begin{aligned}
\langle x \rangle &= \int_{-a}^a x A (a^2 - x^2) dx \\
&= 0
\end{aligned}$$

(c)

$$\begin{aligned}
\langle p \rangle &= \int_{-a}^a \Psi^* \left( -i \hbar \frac{\partial}{\partial x} \right) \Psi dx \\
&= 2i A^2 \hbar \int_{-a}^a x (a^2 - x^2) dx \\
&= 0
\end{aligned}$$

(d)

$$\begin{aligned}\langle x^2 \rangle &= \int_{-a}^a \Psi^* x^2 \Psi dx \\ &= A^2 \int_{-a}^a x^2 (a^2 - x^2)^2 dx \\ &= A^2 \frac{16}{105} a^7 \\ &= \frac{a^2}{7}\end{aligned}$$

(e)

$$\begin{aligned}\langle p^2 \rangle &= \int_{-a}^a \Psi^* \left( -\hbar^2 \frac{\partial^2}{\partial x^2} \right) \Psi dx \\ &= -\hbar^2 \int_{-a}^a A(a^2 - x^2)(-2A) dx \\ &= 4A^2 \hbar^2 \int_0^a (a^2 - x^2) dx \\ &= 4A^2 \hbar^2 \left[ a^2 x - \frac{1}{3} x^3 \right]_0^a \\ &= 4A^2 \hbar^2 \left( a^3 - \frac{1}{3} a^3 \right) \\ &= \frac{8}{3} A^2 a^3 \hbar^2 \\ &= \frac{8}{3} \frac{15}{16a^5} a^3 \hbar^2 \\ &= \frac{5}{2} \frac{\hbar^2}{a^2}\end{aligned}$$

(f)

$$\begin{aligned}\sigma_x &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\ &= \sqrt{\frac{a^2}{7}} \\ &= \frac{a}{\sqrt{7}}\end{aligned}$$

(g)

$$\begin{aligned}\sigma_p &= \sqrt{\langle p^2 \rangle - \langle p \rangle^2} \\ &= \sqrt{\frac{5}{2}} \frac{\hbar}{a}\end{aligned}$$



(h)

$$\begin{aligned}\sigma_x \sigma_p &= \sqrt{\frac{5}{14}} \hbar \\ &\geq \frac{1}{2} \hbar\end{aligned}$$

**1.18**

(a)

$$\begin{aligned}\frac{\hbar}{\sqrt{3mk_B T}} &> d \\ \frac{\sqrt{3mk_B T}}{\hbar} &< \frac{1}{d} \\ T_{\text{electron}} &< \frac{\hbar^2}{3d^2 m k_B} \\ &< 1.3 \times 10^5 \text{ K} \\ T_{\text{nuclei}} &< 2.5 \text{ K}\end{aligned}$$

(b)

$$\begin{aligned}PV &= Nk_B T \\ \frac{V}{N} &= \frac{k_B T}{P} \\ d &= \left( \frac{k_B T}{P} \right)^{1/3} \\ \frac{\hbar}{\sqrt{3mk_B T}} &> \left( \frac{k_B T}{P} \right)^{1/3} \\ T &< \frac{1}{k_B} \left( \frac{\hbar^2}{3m} \right)^{3/5} P^{2/5}\end{aligned}$$

## 2 Time-Independent Schrödinger Equation

### 2.1

(a)

$$\begin{aligned}\int_{-\infty}^{\infty} |\Psi|^2 dx &= \int_{-\infty}^{\infty} \Psi^* \Psi dx \\ &= \int_{-\infty}^{\infty} \psi^* e^{i(E_0 - i\Gamma)t/\hbar} \psi e^{-i(E_0 + i\Gamma)t/\hbar} dx \\ &= e^{2\Gamma t/\hbar} \int_{-\infty}^{\infty} |\psi|^2 dx\end{aligned}$$

In order for this to equal 1 for all  $t$ ,  $\Gamma$  must be 0.

- (b) If  $\psi(x)$  is a complex solution to the time-independent Schrödinger equation then so is  $\psi^*(x)$  and  $\psi(x) + \psi^*(x)$  which is real.

## 2.2

If  $\psi$  and its second derivative always have the same sign,  $\psi$  will increase or decrease without bound forever. This means there is no non-zero choice of constant  $A$  such that

$$\int_{-\infty}^{\infty} |A\psi|^2 dx = 1$$

and thus the equation can't be normalised.

The classical analog of this is statements is that the potential energy of a system can't exceed its total energy.

## 2.3

The time-independent Schrödinger equation in an infinite square well is

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi.$$

If  $E = 0$  then  $\psi = Ax + B$  which isn't normalisable.

If  $E < 0$  then  $\psi = Ae^{kt} + Be^{-kt}$  where  $k \in \mathbb{R}$  which also isn't normalisable.

## 2.4

$$\begin{aligned}
\Psi_n(x, t) &= \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right) e^{-i(n^2\pi^2\hbar/2ma^2)t} \\
\langle x \rangle &= \int_0^a \Psi_n^* x \Psi_n dx \\
&= \frac{2}{a} \int_0^a x \sin^2\left(\frac{n\pi}{a}x\right) dx \\
&= \frac{a}{2} \\
\langle x^2 \rangle &= \int_0^a \Psi_n^* x^2 \Psi_n dx \\
&= \frac{2}{a} \int_0^a x^2 \sin^2\left(\frac{n\pi}{a}x\right) dx \\
&= a^2 \left( \frac{1}{3} - \frac{1}{2n^2\pi^2} \right) \\
\langle p \rangle &= \int_0^a \Psi_n^* \left( -i\hbar \frac{\partial}{\partial x} \right) \Psi_n dx \\
&= -i \frac{2\hbar n\pi}{a^2} \int_0^a \sin\left(\frac{n\pi}{a}x\right) \cos\left(\frac{n\pi}{a}x\right) dx \\
&= 0 \\
\langle p^2 \rangle &= \int_0^a \Psi_n^* \left( -\hbar^2 \frac{\partial^2}{\partial x^2} \right) \Psi_n dx \\
&= \frac{2\hbar^2 n^2 \pi^2}{a^3} \int_0^a \sin^2\left(\frac{n\pi}{a}x\right) dx \\
&= \left( \frac{n\pi\hbar}{a} \right)^2 \\
\sigma_x &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\
&= \frac{a}{2} \sqrt{\frac{1}{3} - \frac{2}{n^2\pi^2}} \\
\sigma_p &= \sqrt{\langle p^2 \rangle - \langle p \rangle^2} \\
&= \frac{n\pi\hbar}{a}
\end{aligned}$$

## 2.5

(a)

$$\begin{aligned}
 1 &= \int_0^a A^2 (\psi_1 + \psi_2)^2 dx \\
 &= A^2 \int_0^a (\psi_1^2 + 2\psi_1\psi_2 + \psi_2^2) dx \\
 &= \frac{2A^2}{a} \left[ \int_0^a \sin^2\left(\frac{\pi}{a}x\right) dx + \int_0^a \sin^2\left(\frac{2\pi}{a}x\right) dx \right] \\
 &= 2A^2 \\
 A &= \frac{1}{\sqrt{2}}
 \end{aligned}$$

(b)

$$\begin{aligned}
 \Psi(x, t) &= \frac{1}{\sqrt{2}} \left[ \sqrt{\frac{2}{a}} \sin\left(\frac{\pi}{a}x\right) e^{-i\omega t} + \sqrt{\frac{2}{a}} \sin\left(\frac{2\pi}{a}x\right) e^{-4i\omega t} \right] \\
 |\Psi(x, t)|^2 &= \Psi^* \Psi \\
 &= \frac{1}{a} \left[ \sin\left(\frac{\pi}{a}x\right) e^{i\omega t} + \sin\left(\frac{2\pi}{a}x\right) e^{4i\omega t} \right] \\
 &\quad \left[ \sin\left(\frac{\pi}{a}x\right) e^{-i\omega t} + \sin\left(\frac{2\pi}{a}x\right) e^{-4i\omega t} \right] \\
 &= \frac{1}{a} \left[ \sin^2\left(\frac{\pi}{a}x\right) + \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) e^{-3i\omega t} \right. \\
 &\quad \left. + \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) e^{3i\omega t} + \sin^2\left(\frac{2\pi}{a}x\right) \right] \\
 &= \frac{1}{a} \left[ \sin^2\left(\frac{\pi}{a}x\right) + \sin^2\left(\frac{2\pi}{a}x\right) \right. \\
 &\quad \left. + 2 \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) \cos(3\omega t) \right]
 \end{aligned}$$

(c)

$$\begin{aligned}
 \langle x \rangle &= \int_0^a \Psi^* x \Psi dx \\
 &= \int_0^a x |\Psi|^2 dx \\
 &= \frac{a}{2} \left[ 1 - \frac{32}{9\pi^2} \cos(3\omega t) \right]
 \end{aligned}$$

(d)

$$\begin{aligned}
 \langle p \rangle &= m \frac{d \langle x \rangle}{dt} \\
 &= \frac{16am\omega}{3\pi^2} \sin(3\omega t) \\
 &= \frac{8\hbar}{3a} \sin(3\omega t)
 \end{aligned}$$

(e) You can get  $E_1$  or  $E_2$  and the probability of getting each is  $1/2$ .

$H = \frac{1}{2}(E_1 + E_2)$  is the mean of the two possible energy values.

## 2.6

$$\begin{aligned}
 \Psi(x, 0) &= A[\psi_1 + e^{i\phi}\psi_2] \\
 1 &= \int_0^a |\Psi|^2 dx \\
 &= \int_0^a \Psi^* \Psi dx \\
 &= A^2 \int_0^a (\psi_1 + e^{-i\phi}\psi_2)(\psi_1 + e^{i\phi}\psi_2) dx \\
 &= A^2 \int_0^a (\psi_1^2 + e^{i\phi}\psi_1\psi_2 + e^{-i\phi}\psi_1\psi_2 + \psi_2^2) dx \\
 &= \frac{2A^2}{a} \int_0^a \left[ \sin^2\left(\frac{\pi}{a}x\right) + e^{i\phi} \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) \right. \\
 &\quad \left. e^{-i\phi} \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) + \sin^2\left(\frac{2\pi}{a}x\right) \right] dx \\
 &= \frac{2A^2}{a} \int_0^a \left[ \sin^2\left(\frac{\pi}{a}x\right) + \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) \cos\phi \right. \\
 &\quad \left. + \sin^2\left(\frac{2\pi}{a}x\right) \right] dx \\
 &= 2A^2 \\
 A &= \frac{1}{\sqrt{2}} \\
 \Psi(x, t) &= \frac{1}{\sqrt{a}} \left[ \sin\left(\frac{\pi}{a}x\right) e^{-i\omega t} + \sin\left(\frac{2\pi}{a}x\right) e^{i(\phi-4\omega t)} \right]
 \end{aligned}$$

$$\begin{aligned}
|\Psi|^2 &= \Psi^* \Psi \\
&= \frac{1}{a} \left[ \sin\left(\frac{\pi}{a}x\right) e^{i\omega t} + \sin\left(\frac{2\pi}{a}x\right) e^{-i(\phi-4\omega t)} \right] \\
&\quad \left[ \sin\left(\frac{\pi}{a}x\right) e^{-i\omega t} + \sin\left(\frac{2\pi}{a}x\right) e^{i(\phi-4\omega t)} \right] \\
&= \frac{1}{a} \left[ \sin^2\left(\frac{\pi}{a}x\right) + \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) e^{i(\phi-3\omega t)} \right. \\
&\quad \left. \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) e^{-i(\phi-3\omega t)} + \sin^2\left(\frac{2\pi}{a}x\right) \right] \\
&= \frac{1}{a} \left[ \sin^2\left(\frac{\pi}{a}x\right) + \sin^2\left(\frac{2\pi}{a}x\right) \right. \\
&\quad \left. + 2 \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) \cos(\phi - 3\omega t) \right] \\
\langle x \rangle &= \int_0^a \Psi^* x \Psi dx \\
&= \int_0^a x |\Psi|^2 dx \\
&= \frac{a}{2} \left[ 1 - \frac{32}{9\pi^2} \cos(3\omega t - \phi) \right]
\end{aligned}$$

## 2.7

(a)

$$\begin{aligned}
1 &= \int_0^a |\Psi|^2 dx \\
&= A^2 \left[ \int_0^{a/2} x^2 dx + \int_{a/2}^a (a-x)^2 dx \right] \\
&= A^2 \left\{ \frac{1}{3} \left[ \frac{a}{2} \right]^3 + \left[ -\frac{1}{3}(a-x)^3 \right]_{a/2}^a \right\} \\
&= A^2 \left( \frac{a^3}{24} + \frac{a^3}{24} \right) \\
&= \frac{A^2 a^3}{12} \\
A &= \frac{2\sqrt{3}}{\sqrt{a^3}}
\end{aligned}$$

(b)

$$\begin{aligned}
c_n &= \sqrt{\frac{2}{a}} \int_0^a \sin\left(\frac{n\pi}{a}x\right) \Psi(x, 0) dx \\
&= \sqrt{\frac{2}{a}} \left[ \int_0^{a/2} \sin\left(\frac{n\pi}{a}x\right) Ax dx + \int_{a/2}^a \sin\left(\frac{n\pi}{a}x\right) A(a-x) dx \right] \\
&= \frac{2\sqrt{6}}{a^2} \left[ \int_0^{a/2} x \sin\left(\frac{n\pi}{a}x\right) dx + \int_{a/2}^a (a-x) \sin\left(\frac{n\pi}{a}x\right) dx \right] \\
&= \frac{8\sqrt{6}}{n^2\pi^2} \sin^2\left(\frac{n\pi}{4}\right) \sin\left(\frac{n\pi}{2}\right) \\
&= \begin{cases} 0 & n \text{ even} \\ (-1)^{(n-1)/2} \frac{4\sqrt{6}}{n^2\pi^2} & n \text{ odd} \end{cases} \\
\Psi(x, t) &= \frac{4\sqrt{6}}{\pi^2} \sqrt{\frac{2}{a}} \sum_{n=1,3,5,\dots}^{\infty} (-1)^{(n-1)/2} \frac{1}{n^2} \sin\left(\frac{n\pi}{a}x\right) e^{-i(n^2\pi^2\hbar/2ma^2)t}
\end{aligned}$$

(c)

$$\begin{aligned}
|c_1|^2 &= \left(\frac{4\sqrt{6}}{\pi^2}\right)^2 \\
&\approx 0.985
\end{aligned}$$

(d)

$$\begin{aligned}
E_n &= \frac{n^2\pi^2\hbar^2}{2ma^2} \\
\langle H \rangle &= \sum_{n=0}^{\infty} |c_{2n+1}|^2 E_{2n+1} \\
&= \sum_{n=0}^{\infty} \left(\frac{4\sqrt{6}}{(2n+1)^2\pi^2}\right)^2 \frac{(2n+1)^2\pi^2\hbar^2}{2ma^2} \\
&= \sum_{n=0}^{\infty} \frac{48\hbar^2}{(2n+1)^2ma^2\pi^2} \\
&= \frac{48\hbar^2}{ma^2\pi^2} \sum_{n=0}^{\infty} \frac{1}{(2n+1)^2} \\
&= \frac{6\hbar^2}{ma^2}
\end{aligned}$$

## 2.8

$$\begin{aligned}
 1 &= \int_0^{a/2} |\Psi|^2 dx \\
 &= A^2 \int_0^{a/2} dx \\
 &= \frac{aA^2}{2} \\
 A &= \sqrt{\frac{2}{a}} \\
 c_n &= \frac{2}{a} \int_0^{a/2} \sin\left(\frac{n\pi}{a}x\right) dx \\
 |c_1|^2 &= \left(\frac{2}{\pi}\right)^2 \\
 &\approx 0.405
 \end{aligned}$$

## 2.9

$$\begin{aligned}
 \Psi(x, 0) &= Ax(a - x) \\
 \langle H \rangle &= \int_0^a \Psi(x, 0)^* \hat{H} \Psi(x, 0) dx \\
 &= \int_0^a \Psi(x, 0)^* \left( -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \right) \Psi(x, 0) dx \\
 &= \frac{A^2 \hbar^2}{m} \int_0^a x(a - x) dx \\
 &= \frac{30 \hbar^2}{ma^5} \frac{a^3}{6} \\
 &= \frac{5 \hbar^2}{ma^2}
 \end{aligned}$$



## 2.10

(a)

$$\begin{aligned}
 \psi_2(x) &= \frac{1}{\sqrt{2!}}(\hat{a}_+)\psi_1 \\
 &= \frac{1}{\sqrt{2}} \frac{1}{\sqrt{2\hbar m\omega}} \left(-\hbar \frac{d}{dx} + m\omega x\right) \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \sqrt{\frac{2m\omega}{\hbar}} x e^{-\frac{m\omega}{2\hbar}x^2} \\
 &= \frac{1}{\sqrt{2\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \left(-\hbar \frac{d}{dx} + m\omega x\right) x e^{-\frac{m\omega}{2\hbar}x^2} \\
 &= \frac{1}{\sqrt{2\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \left[-\hbar \left(e^{-\frac{m\omega}{2\hbar}x^2} - \frac{m\omega}{\hbar} x^2 e^{-\frac{m\omega}{2\hbar}x^2}\right) + m\omega x^2 e^{-\frac{m\omega}{2\hbar}x^2}\right] \\
 &= \frac{1}{\sqrt{2}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \left(\frac{2m\omega}{\hbar} x^2 - 1\right) e^{-\frac{m\omega}{2\hbar}x^2}
 \end{aligned}$$

## 2.11

(a)

$$\begin{aligned}
\langle x \rangle &= \int_{-\infty}^{\infty} \psi_0^* x \psi_0 dx \\
&= \alpha^2 \int_{-\infty}^{\infty} x e^{-\frac{m\omega}{\hbar} x^2} dx \\
&= 0 \\
\langle p \rangle &= m \frac{d\langle x \rangle}{dt} \\
&= 0 \\
\langle x^2 \rangle &= \int_{-\infty}^{\infty} \psi_0^* x^2 \psi_0 dx \\
&= \alpha^2 \int_{-\infty}^{\infty} x^2 e^{-\frac{m\omega}{\hbar} x^2} dx \\
&= \frac{\hbar}{2m\omega} \\
\langle p^2 \rangle &= \int_{-\infty}^{\infty} \psi_0^* \left( -\hbar^2 \frac{d^2}{dx^2} \right) \psi_0 dx \\
&= -\hbar^2 \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \int_{-\infty}^{\infty} e^{-\frac{m\omega}{2\hbar} x^2} \frac{d}{dx} \left( -\frac{m\omega}{\hbar} x e^{-\frac{m\omega}{2\hbar} x^2} \right) dx \\
&= \hbar^2 \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \frac{m\omega}{\hbar} \int_{-\infty}^{\infty} e^{-\frac{m\omega}{2\hbar} x^2} \left( e^{-\frac{m\omega}{2\hbar} x^2} - \frac{m\omega}{\hbar} x^2 e^{-\frac{m\omega}{2\hbar} x^2} \right) dx \\
&= \hbar^2 \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \frac{m\omega}{\hbar} \int_{-\infty}^{\infty} \left( 1 - \frac{m\omega}{\hbar} x^2 \right) e^{-\frac{m\omega}{2\hbar} x^2} dx \\
&= \hbar^2 \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \frac{m\omega}{\hbar} \frac{\hbar\sqrt{\pi}}{2\sqrt{\hbar m\omega}} \\
&= \frac{1}{2} m\hbar\omega
\end{aligned}$$

$$\begin{aligned}
\psi_1(x) &= \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \sqrt{\frac{2m\omega}{\hbar}} x e^{-\frac{m\omega}{2\hbar} x^2} \\
\langle x \rangle &= 0 \\
\langle p \rangle &= m \frac{d\langle x \rangle}{dt} \\
&= 0 \\
\langle x^2 \rangle &= \int_{-\infty}^{\infty} \psi_1^* x^2 \psi_1 dx \\
&= \left(\frac{m\omega}{\pi\hbar}\right)^{1/2} \frac{2m\omega}{\hbar} \int_{-\infty}^{\infty} x^4 e^{-\frac{m\omega}{\hbar} x^2} dx \\
&= \left(\frac{m\omega}{\pi\hbar}\right)^{1/2} \frac{2m\omega}{\hbar} \frac{3}{4} \sqrt{\pi} \left(\frac{\hbar}{m\omega}\right)^{5/2} \\
&= \frac{3}{2} \frac{\hbar}{m\omega} \\
\langle p^2 \rangle &= \int_{-\infty}^{\infty} \psi_1^* \left(-\hbar^2 \frac{d^2}{dx^2}\right) \psi_1 dx \\
&= -\hbar^2 \left(\frac{m\omega}{\pi\hbar}\right)^{1/2} \frac{2m\omega}{\hbar} \int_{-\infty}^{\infty} x e^{-\frac{m\omega}{2\hbar} x^2} \frac{d}{dx} \left(e^{-\frac{m\omega}{2\hbar} x^2} - \frac{m\omega}{\hbar} x^2 e^{-\frac{m\omega}{2\hbar} x^2}\right) dx \\
&= -\hbar^2 \left(\frac{m\omega}{\pi\hbar}\right)^{1/2} \frac{2m\omega}{\hbar} \int_{-\infty}^{\infty} x e^{-\frac{m\omega}{2\hbar} x^2} \left[-\frac{m\omega}{\hbar} x e^{-\frac{m\omega}{2\hbar} x^2} - \frac{2m\omega}{\hbar} x e^{-\frac{m\omega}{2\hbar} x^2} + \left(\frac{m\omega}{\hbar}\right)^2 x^3 e^{-\frac{m\omega}{2\hbar} x^2}\right] dx \\
&= 2\hbar^2 \left(\frac{m\omega}{\pi\hbar}\right)^{1/2} \left(\frac{m\omega}{\hbar}\right)^2 \int_{-\infty}^{\infty} x^2 e^{-\frac{m\omega}{\hbar} x^2} \left(3 - \frac{m\omega}{\hbar} x^2\right) dx \\
&= 2\hbar^2 \left(\frac{m\omega}{\pi\hbar}\right)^{1/2} \left(\frac{m\omega}{\hbar}\right)^2 \frac{3}{4} \sqrt{\pi} \left(\frac{\hbar}{m\omega}\right)^{3/2} \\
&= \frac{3}{2} m\hbar\omega
\end{aligned}$$

(b)

$$\begin{aligned}\sigma_x &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\ &= \sqrt{\frac{\hbar}{2m\omega}} \\ \sigma_p &= \sqrt{\langle p^2 \rangle - \langle p \rangle^2} \\ &= \sqrt{\frac{m\hbar\omega}{2}} \\ \sigma_x \sigma_p &= \frac{\hbar}{2} \\ \sigma_x &= \sqrt{\frac{3\hbar}{2m\omega}} \\ \sigma_p &= \sqrt{\frac{3m\hbar\omega}{2}} \\ \sigma_x \sigma_p &= \frac{3}{2}\hbar\end{aligned}$$

(c)

$$\begin{aligned}\langle T \rangle &= \frac{\langle p^2 \rangle}{2m} \\ &= \frac{\hbar\omega}{4} \\ \langle V \rangle &= \frac{1}{2}m\omega^2 \langle x^2 \rangle \\ &= \frac{1}{4}\hbar\omega \\ \langle T \rangle &= \frac{\langle p^2 \rangle}{2m} \\ &= \frac{3}{4}\hbar\omega \\ \langle V \rangle &= \frac{1}{2}m\omega^2 \langle x^2 \rangle \\ &= \frac{3}{4}\hbar\omega\end{aligned}$$

## 2.12

$$\begin{aligned}
\langle x \rangle &= \int_{-\infty}^{\infty} \psi_n^* x \psi_n dx \\
&= \sqrt{\frac{\hbar}{2m\omega}} \int_{-\infty}^{\infty} \psi_n^* (\hat{a}_+ + \hat{a}_-) \psi_n dx \\
&= \sqrt{\frac{\hbar}{2m\omega}} \int_{-\infty}^{\infty} \psi_n^* (\sqrt{n+1} \psi_{n+1} + \sqrt{n} \psi_{n-1}) dx \\
&= 0
\end{aligned}$$

$$\begin{aligned}
\langle p \rangle &= \int_{-\infty}^{\infty} \psi_n^* p \psi_n dx \\
&= i \sqrt{\frac{\hbar m \omega}{2}} \int_{-\infty}^{\infty} \psi_n^* (\hat{a}_+ - \hat{a}_-) \psi_n dx \\
&= i \sqrt{\frac{\hbar m \omega}{2}} \int_{-\infty}^{\infty} \psi_n^* (\sqrt{n+1} \psi_{n+1} - \sqrt{n} \psi_{n-1}) dx \\
&= 0
\end{aligned}$$

$$\begin{aligned}
\langle x^2 \rangle &= \int_{-\infty}^{\infty} \psi_n^* x^2 \psi_n dx \\
&= \frac{\hbar}{2m\omega} \int_{-\infty}^{\infty} \psi_n^* (\hat{a}_+^2 + \hat{a}_+ \hat{a}_- + \hat{a}_- \hat{a}_+ + \hat{a}_-^2) \psi_n dx \\
&= \frac{\hbar}{2m\omega} (2n+1) \int_{-\infty}^{\infty} |\psi_n|^2 dx \\
&= \frac{\hbar}{m\omega} \left( n + \frac{1}{2} \right)
\end{aligned}$$

$$\begin{aligned}
\langle p^2 \rangle &= \int_{-\infty}^{\infty} \psi_n^* p^2 \psi_n dx \\
&= -\frac{\hbar m \omega}{2} \int_{-\infty}^{\infty} \psi_n^* (\hat{a}_+^2 - \hat{a}_+ \hat{a}_- - \hat{a}_- \hat{a}_+ + \hat{a}_-^2) \psi_n dx \\
&= \frac{\hbar m \omega}{2} (2n+1) \int_{-\infty}^{\infty} |\psi_n|^2 dx \\
&= \hbar m \omega \left( n + \frac{1}{2} \right)
\end{aligned}$$

$$\begin{aligned}
\langle T \rangle &= \left\langle \frac{p^2}{2m} \right\rangle \\
&= \frac{1}{2} \hbar \omega \left( n + \frac{1}{2} \right)
\end{aligned}$$

$$\begin{aligned}
\sigma_x &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\
&= \sqrt{\frac{\hbar}{m\omega} \left( n + \frac{1}{2} \right)} \\
\sigma_p &= \sqrt{\hbar m\omega \left( n + \frac{1}{2} \right)} \\
\sigma_x \sigma_p &= (2n + 1) \frac{\hbar}{2} \\
&\geq \frac{\hbar}{2}
\end{aligned}$$

## 2.13

(a)

$$\begin{aligned}
\Psi(x, 0) &= A[3\psi_0(x) + 4\psi_1(x)] \\
1 &= \int_{-\infty}^{\infty} |\Psi(x, 0)|^2 dx \\
&= A^2 \int_{-\infty}^{\infty} [9\psi_0(x)^2 + 24\psi_0(x)\psi_1(x) + 16\psi_1(x)^2] dx \\
&= 25A^2 \\
A &= \frac{1}{5}
\end{aligned}$$

(b)

$$\begin{aligned}
\Psi(x, t) &= \frac{1}{5}[3\psi_0(x)e^{-i\omega t/2} + 4\psi_1(x)e^{-3i\omega t/2}] \\
|\Psi(x, t)|^2 &= \Psi(x, t)^* \Psi(x, t) \\
&= \frac{1}{25}[3\psi_0(x)e^{i\omega t/2} + 4\psi_1(x)e^{3i\omega t/2}][3\psi_0(x)e^{-i\omega t/2} + 4\psi_1(x)e^{-3i\omega t/2}] \\
&= \frac{1}{25}[9\psi_0(x)^2 + 12\psi_0(x)\psi_1(x)e^{-i\omega t} + 12\psi_0(x)\psi_1(x)e^{i\omega t} + 16\psi_1(x)^2] \\
&= \frac{1}{25}[9\psi_0(x)^2 + 16\psi_1(x)^2 + 24\psi_0(x)\psi_1(x)\cos\omega t]
\end{aligned}$$

(c)

$$\begin{aligned}
\langle x \rangle &= \int_{-\infty}^{\infty} \Psi^* x \Psi dx \\
&= \frac{1}{25} \int_{-\infty}^{\infty} x(9\psi_0^2 + 16\psi_1^2 + 24\psi_0\psi_1 \cos \omega t) dx \\
&= \frac{24}{25} \int_{-\infty}^{\infty} x\psi_0\psi_1 \cos(\omega t) dx \\
&= \frac{24}{25} \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \sqrt{\frac{2m\omega}{\hbar}} \cos(\omega t) \int_{-\infty}^{\infty} x^2 e^{-\frac{m\omega}{\hbar} x^2} dx \\
&= \frac{24}{25} \left( \frac{m\omega}{\pi\hbar} \right)^{1/2} \sqrt{\frac{2m\omega}{\hbar}} \cos(\omega t) \frac{1}{2} \sqrt{\pi} \left( \frac{\hbar}{m\omega} \right)^{3/2} \\
&= \frac{24}{25} \sqrt{\frac{\hbar}{2m\omega}} \cos(\omega t) \\
\langle p \rangle &= m \frac{d\langle x \rangle}{dt} \\
&= -\frac{24}{25} \sqrt{\frac{\hbar m\omega}{2}} \sin(\omega t) \\
\frac{d\langle p \rangle}{dt} &= -\frac{24}{25} \sqrt{\frac{\hbar m\omega}{2}} \omega \cos(\omega t) \\
V &= \frac{1}{2} m\omega^2 x^2 \\
\frac{\partial V}{\partial \theta} &= m\omega^2 x \\
\left\langle -\frac{\partial V}{\partial x} \right\rangle &= -m\omega^2 \langle x \rangle \\
&= -\frac{24}{25} \sqrt{\frac{\hbar m\omega}{2}} \omega \cos(\omega t) \\
&= \frac{d\langle p \rangle}{dt}
\end{aligned}$$

(d)

$$\begin{aligned}
E_0 &= \frac{\hbar\omega}{2} \\
P(E_0) &= \frac{9}{25} \\
E_1 &= \frac{3\hbar\omega}{2} \\
P(E_1) &= \frac{16}{25}
\end{aligned}$$

2.14

$$1 - \left(\frac{m\omega}{\pi\hbar}\right)^{1/2} \int_{-\sqrt{\hbar/m\omega}}^{\sqrt{\hbar/m\omega}} e^{-m\omega x^2/\hbar} dx = 1 - \left(\frac{m\omega}{\pi\hbar}\right)^{1/2} \sqrt{\frac{\pi\hbar}{m\omega}} \operatorname{erf} 1$$

$$= 0.157$$

2.15

$$a_{j+2} = \frac{-2(n-j)}{(j+1)(j+2)} a_j$$

$$a_3 = -\frac{4}{3} a_1$$

$$a_5 = \frac{4}{15} a_1$$

$$H_5(\xi) = a_1 \left( \xi - \frac{4}{3} \xi^3 + \frac{4}{15} \xi^5 \right)$$

$$= \frac{1}{120} a_1 (120\xi - 160\xi^3 + 32\xi^5)$$

$$= 32\xi^5 - 160\xi^3 + 120\xi$$

$$a_2 = -6a_0$$

$$a_4 = \frac{-8}{12} a_2$$

$$= 4a_0$$

$$a_6 = \frac{-4}{30} a_4$$

$$= -\frac{8}{15} a_0$$

$$H_6(\xi) = a_0 \left( 1 - 6\xi^2 + 4\xi^4 - \frac{8}{15} \xi^6 \right)$$

$$= \frac{1}{120} a_0 (120 - 720\xi^2 + 480\xi^4 - 64\xi^6)$$

$$= 64\xi^6 - 480\xi^4 + 720\xi^2 - 120$$



## 2.17

$$\begin{aligned}
 Ae^{ikx} + Be^{-ikx} &= A[\cos(kx) + i\sin(kx)] + B[\cos(kx) - i\sin(kx)] \\
 &= (A + B)\cos(kx) + i(A - B)\sin(kx) \\
 C &= A + B \\
 D &= i(A - B) \\
 -iD &= A - B \\
 A &= \frac{C - iD}{2} \\
 B &= \frac{C + iD}{2}
 \end{aligned}$$

## 2.18

$$\begin{aligned}
 \Psi_k(x, t) &= Ae^{i\left(kx - \frac{\hbar k^2}{2m}t\right)} \\
 J(x, t) &= \frac{i\hbar}{2m} \left( \Psi \frac{\partial \Psi^*}{\partial x} - \Psi^* \frac{\partial \Psi}{\partial x} \right) \\
 &= \frac{\hbar k |A|^2}{m}
 \end{aligned}$$

The probability travels in the same direction as the wave.

## 2.20

(a)

$$\begin{aligned}
 \Psi(x, 0) &= Ae^{-a|x|} \\
 1 &= \int_{-\infty}^{\infty} \Psi^* \Psi \, dx \\
 &= |A|^2 \int_{-\infty}^{\infty} e^{-2a|x|} \, dx \\
 &= 2|A|^2 \int_0^{\infty} e^{-2ax} \, dx \\
 &= \frac{|A|^2}{a} \\
 A &= \sqrt{a}
 \end{aligned}$$

(b)

$$\begin{aligned}
\phi(k) &= \sqrt{\frac{a}{2\pi}} \int_{-\infty}^{\infty} e^{-a|x|-ikx} dx \\
&= \sqrt{\frac{a}{2\pi}} \int_{-\infty}^{\infty} e^{-a|x|} [\cos(kx) - i \sin(kx)] dx \\
&= \sqrt{\frac{a}{2\pi}} \int_{-\infty}^{\infty} e^{-a|x|} \cos(kx) dx \\
&= \sqrt{\frac{a}{2\pi}} 2 \int_0^{\infty} e^{-ax} \cos(kx) dx \\
&= \sqrt{\frac{a}{2\pi}} \frac{2a}{a^2 + k^2}
\end{aligned}$$

(c)

$$\Psi(x, t) = \frac{a^{3/2}}{\pi} \int_{-\infty}^{\infty} \frac{1}{a^2 + k^2} e^{i\left(kx - \frac{\hbar k^2}{2m} t\right)} dk$$

## 2.21

(a)

$$\begin{aligned}
\Psi(x, 0) &= A e^{-ax^2} \\
1 &= A^2 \int_{-\infty}^{\infty} e^{-2ax^2} dx \\
&= \sqrt{\frac{\pi}{2a}} A^2 \\
A &= \left(\frac{2a}{\pi}\right)^{1/4}
\end{aligned}$$

(b)

$$\begin{aligned}
\phi(k) &= \frac{1}{\sqrt{2\pi}} \left(\frac{2a}{\pi}\right)^{1/4} \int_{-\infty}^{\infty} e^{-(ax^2 + ikx)} dx \\
&= \frac{1}{(2\pi a)^{1/4}} e^{-k^2/4a} \\
\Psi(x, t) &= \frac{1}{\sqrt{2\pi}} \frac{1}{(2\pi a)^{1/4}} \int_{-\infty}^{\infty} e^{-\frac{k^2}{4a} + i\left(kx - \frac{\hbar k^2}{2m} t\right)} dk \\
\Psi(x, t) &= \left(\frac{2a}{\pi}\right)^{1/4} \frac{1}{\gamma} e^{-ax^2/\gamma^2}
\end{aligned}$$

(c)

$$\begin{aligned}
|\Psi(x, t)|^2 &= \Psi^* \Psi \\
&= \left(\frac{2a}{\pi}\right)^{1/2} \frac{1}{\gamma^*} e^{-ax^2/(\gamma^*)^2} \frac{1}{\gamma} e^{-ax^2/\gamma^2} \\
&= \left(\frac{2a}{\pi}\right)^{1/2} \frac{1}{\sqrt{1-2i\hbar a t/m}} e^{-ax^2/(1-2i\hbar a t/m)} \\
&\quad \frac{1}{\sqrt{1+2i\hbar a t/m}} e^{-ax^2/(1+2i\hbar a t/m)} \\
&= \left(\frac{2a}{\pi}\right)^{1/2} \frac{1}{\sqrt{1+(2\hbar a t/m)^2}} e^{-2ax^2/[1+(2\hbar a t/m)^2]} \\
&= \sqrt{\frac{2}{\pi}} w e^{-2w^2 x^2}
\end{aligned}$$

As  $t$  increases  $|\Psi|^2$  flattens out and broadens.

(d)

$$\begin{aligned}
\langle x \rangle &= \int_{-\infty}^{\infty} \Psi^* x \Psi dx \\
&= 0 \\
\langle p \rangle &= m \frac{d\langle x \rangle}{dt} \\
&= 0 \\
\langle x^2 \rangle &= \int_{-\infty}^{\infty} \Psi^* x^2 \Psi dx \\
&= \sqrt{\frac{2}{\pi}} w \int_{-\infty}^{\infty} x^2 e^{-2w^2 x^2} dx \\
&= \frac{1}{4w^2}
\end{aligned}$$

## 2.22

(a)  $-25$

(b)  $1$

(c)  $0$

2.26

$$\begin{aligned}
 F(k) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \delta(x) e^{-ikx} dx \\
 &= \frac{1}{\sqrt{2\pi}} \\
 f(x) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ikx} dk
 \end{aligned}$$

2.29

$$\begin{aligned}
 \psi(x) &= \begin{cases} Fe^{-\kappa x} & x > a \\ C \sin(lx) & 0 < x < a \\ -\psi(-x) & x < 0 \end{cases} \\
 Fe^{-\kappa a} &= C \sin(la) \\
 -\kappa Fe^{-\kappa a} &= lC \cos(la) \\
 -\frac{1}{\kappa} &= \frac{1}{l} \tan(la) \\
 \tan z &= -\frac{l}{\kappa} \\
 &= -\frac{la}{\kappa a} \\
 &= -\frac{z}{\sqrt{z_0^2 - z^2}}
 \end{aligned}$$

For large  $z_0$  the intersections occur just below  $z_n = n\pi$  so

$$\begin{aligned}
 z &= la \\
 n\pi &\approx \frac{\sqrt{2m(E + V_0)}}{\hbar} a \\
 E + V_0 &\approx \frac{n^2 \pi^2 \hbar^2}{2ma^2}.
 \end{aligned}$$

As  $z_0$  decreases there are fewer and fewer bound states. When  $z_0 < \pi/2$  there are no odd bound states.

## 2.30

$$\begin{aligned}
\psi(x) &= \begin{cases} Fe^{-\kappa x} & x > a \\ D \cos(lx) & 0 < x < a \\ \psi(-x) & x < 0 \end{cases} \\
1 &= \int_{-\infty}^{\infty} |\psi|^2 dx \\
&= 2 \left( |D|^2 \int_0^a \cos^2(lx) dx + |F|^2 \int_a^{\infty} e^{-2\kappa x} dx \right) \\
&= 2 \left[ |D|^2 \frac{2al + \sin(2al)}{4l} + \frac{|F|^2}{2\kappa} e^{-2\kappa a} \right] \\
&= |D|^2 \left[ a + \frac{\sin(2al)}{2l} + \frac{\cos^2(al)}{\kappa} \right] \\
&= |D|^2 \left[ a + \frac{2 \sin(al) \cos(al)}{2l} + \frac{\cos^3(al)}{l \sin(al)} \right] \\
&= |D|^2 \left\{ a + \frac{\cos(al)}{l \sin(al)} [\sin^2(al) + \cos^2(al)] \right\} \\
&= |D|^2 \left[ a + \frac{\cos(al)}{l \sin(al)} \right] \\
&= |D|^2 \left[ a + \frac{1}{l \tan(al)} \right] \\
&= |D|^2 \left[ a + \frac{1}{\kappa} \right] \\
D &= \frac{1}{\sqrt{a + 1/\kappa}}
\end{aligned}$$

$$\begin{aligned}
1 &= \left\{ \frac{1}{a + 1/\kappa} \left[ a + \frac{\sin(2al)}{2l} \right] + |F|^2 \frac{e^{-2\kappa a}}{\kappa} \right\} \\
\frac{(Fe^{-\kappa a})^2}{\kappa} &= 1 - \frac{1}{a + 1/\kappa} \left[ a + \frac{\sin(2al)}{2l} \right] \\
(Fe^{-\kappa a})^2 &= \kappa - \frac{\kappa}{a + 1/\kappa} \left[ a + \frac{\sin(2al)}{2l} \right] \\
&= \frac{\kappa a + 1 - \kappa a - \kappa \sin(al) \cos(al)/l}{a + 1/\kappa} \\
&= \frac{1 - \sin^2(al)}{a + 1/\kappa} \\
F &= \frac{e^{\kappa a} \cos(al)}{\sqrt{a + 1/\kappa}}
\end{aligned}$$

**2.31**

$$\begin{aligned}1 &= 2aV_0 \\ V_0 &= \frac{1}{2a} \\ z_0 &= \frac{a}{\hbar} \sqrt{2mV_0} \\ &= \frac{a}{\hbar} \sqrt{\frac{m}{a}} \\ &= \frac{\sqrt{am}}{\hbar} \\ \lim_{a \rightarrow 0} z_0 &= 0\end{aligned}$$

## 2.34

(a)

$$\begin{aligned}
 V(x) &= \begin{cases} 0 & x \leq 0 \\ V_0 & x > 0 \end{cases} \\
 -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} &= E\psi \\
 \frac{d^2\psi}{dx^2} &= -\frac{2mE}{\hbar^2}\psi \\
 &= -k^2\psi \\
 k &= \frac{\sqrt{2mE}}{\hbar} \\
 \psi &= Ae^{ikx} + Be^{-ikx} \\
 -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V_0\psi &= E\psi \\
 \frac{d^2\psi}{dx^2} &= -\frac{2m(E - V_0)}{\hbar^2}\psi \\
 &= \kappa^2\psi \\
 \kappa &= \frac{\sqrt{2m(V_0 - E)}}{\hbar} \\
 \psi &= Fe^{-\kappa x} \\
 A + B &= F \\
 ik(A - B) &= -\kappa F \\
 F &= -i\frac{k}{\kappa}(A - B) \\
 A + B &= -i\frac{k}{\kappa}(A - B) \\
 \left(1 - i\frac{k}{\kappa}\right)B &= -\left(1 + i\frac{k}{\kappa}\right)A \\
 B &= -\frac{1 + ik/\kappa}{1 - ik/\kappa}A \\
 R &= \frac{|B|^2}{|A|^2} \\
 &= \left(-\frac{1 + ik/\kappa}{1 - ik/\kappa}\right)\left(-\frac{1 - ik/\kappa}{1 + ik/\kappa}\right) \\
 &= 1
 \end{aligned}$$

(b)

$$\psi(x) = \begin{cases} Ae^{ikx} + Be^{-ikx} & x < 0 \\ Fe^{ilx} & x > 0 \end{cases}$$

$$k = \frac{\sqrt{2mE}}{\hbar}$$

$$l = \frac{\sqrt{2m(E - V_0)}}{\hbar}$$

$$A + B = F$$

$$ik(A - B) = ilF$$

$$F = \frac{k}{l}(A - B)$$

$$A + B = \frac{k}{l}(A - B)$$

$$\left(\frac{k}{l} + 1\right)B = \left(\frac{k}{l} - 1\right)A$$

$$B = \frac{k/l - 1}{k/l + 1}A$$

$$R = \frac{|B|^2}{|A|^2}$$

$$= \left(\frac{k/l - 1}{k/l + 1}\right)^2$$

$$= \left(\frac{k - l}{k + l}\right)^2$$

$$= \frac{(k - l)^4}{(k^2 - l^2)^2}$$

$$k^2 - l^2 = \frac{2mE}{\hbar^2} - \frac{2m(E - V_0)}{\hbar^2}$$

$$= \frac{2m}{\hbar^2}V_0$$

$$k - l = \frac{\sqrt{2m}}{\hbar}(\sqrt{E} - \sqrt{E - V_0})$$

$$R = \frac{(\sqrt{E} - \sqrt{E - V_0})^4}{V_0^2}$$



(d)

$$\begin{aligned} B &= F - A \\ F &= \frac{k}{l}(A - F + A) \\ \left(1 + \frac{k}{l}\right) F &= \frac{2k}{l} A \\ F &= \frac{2k}{k+l} A \\ \frac{l}{k} &= \sqrt{\frac{E - V_0}{E}} \\ T &= \left| \frac{F}{A} \right|^2 \frac{l}{k} \\ &= \left( \frac{2k}{k+l} \right)^2 \frac{l}{k} \\ &= \frac{4kl}{(k+l)^2} \\ &= \frac{4kl(k-l)^2}{(k^2 - l^2)^2} \\ &= \frac{4\sqrt{E}\sqrt{E - V_0}(\sqrt{E} - \sqrt{E - V_0})^2}{V_0^2} \\ T + R &= \frac{4kl}{(k+l)^2} + \frac{(k-l)^2}{(k+l)^2} \\ &= \frac{4kl + k^2 - 2kl + l^2}{(k+l)^2} \\ &= \frac{k^2 + 2kl + l^2}{(k+l)^2} \\ &= \frac{(k+l)^2}{(k+l)^2} \\ &= 1 \end{aligned}$$

## 2.35

(a)

$$\begin{aligned}
 V(x) &= \begin{cases} 0 & x < 0 \\ -V_0 & x > 0 \end{cases} \\
 -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} &= E\psi \\
 \frac{d^2\psi}{dx^2} &= -\frac{2mE}{\hbar^2}\psi \\
 &= -k^2\psi \\
 k &= \frac{\sqrt{2mE}}{\hbar} \\
 \psi &= Ae^{ikx} + Be^{-ikx} \\
 -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} - V_0\psi &= E\psi \\
 \frac{d^2\psi}{dx^2} &= -\frac{2m}{\hbar^2}(E + V_0)\psi \\
 &= -l^2\psi \\
 l &= \frac{\sqrt{2m(E + V_0)}}{\hbar} \\
 \psi &= Fe^{ilx} \\
 A + B &= F \\
 ik(A - B) &= ilF \\
 k(A - B) &= l(A + B) \\
 (k + l)B &= (k - l)A \\
 B &= \frac{k - l}{k + l}A \\
 R &= \left| \frac{B}{A} \right|^2 \\
 &= \left( \frac{k - l}{k + l} \right)^2 \\
 &= \left( \frac{\sqrt{E} - \sqrt{E + V_0}}{\sqrt{E} + \sqrt{E + V_0}} \right)^2 \\
 &= \frac{1}{9}
 \end{aligned}$$

(c)

$$T = 1 - R = \frac{8}{9}$$

## 2.36

$$\begin{aligned}
 V(x) &= \begin{cases} 0 & |x| < a \\ \infty & |x| > a \end{cases} \\
 -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} &= E\psi \\
 \frac{d^2\psi}{dx^2} &= -\frac{2mE}{\hbar^2}\psi \\
 &= -k^2\psi \\
 k &= \frac{\sqrt{2mE}}{\hbar} \\
 \psi &= A \sin kx + B \cos kx \\
 0 &= -A \sin ka + B \cos ka \\
 0 &= A \sin ka + B \cos ka \\
 B \cos ka &= 0 \\
 k &= \frac{n\pi}{2a}, \quad n = 1, 3, 5, \dots \\
 E &= \frac{n^2\pi^2\hbar^2}{2m(2a)^2} \\
 \psi &= B \cos\left(\frac{n\pi}{2a}x\right), \quad n = 1, 3, 5, \dots \\
 1 &= |B|^2 \int_{-a}^a \cos^2\left(\frac{n\pi}{2a}x\right) dx \\
 B &= \frac{1}{\sqrt{a}} \\
 \psi &= \frac{1}{\sqrt{a}} \cos\left(\frac{n\pi}{2a}x\right), \quad n = 1, 3, 5, \dots \\
 A \sin ka &= 0 \\
 k &= \frac{n\pi}{2a}, \quad n = 2, 4, 6, \dots \\
 E &= \frac{n^2\pi^2\hbar^2}{2m(2a)^2} \\
 \psi &= A \sin\left(\frac{n\pi}{2a}x\right), \quad n = 2, 4, 6, \dots \\
 1 &= |A|^2 \int_{-a}^a \sin^2\left(\frac{n\pi}{2a}x\right) dx \\
 A &= \frac{1}{\sqrt{a}} \\
 \psi &= \frac{1}{\sqrt{a}} \sin\left(\frac{n\pi}{2a}x\right), \quad n = 2, 4, 6, \dots
 \end{aligned}$$

## 2.37

$$\begin{aligned}
\Psi(x, 0) &= A \sin^3 \left( \frac{\pi}{a} x \right) \\
&= A \left[ \frac{3}{4} \sin \left( \frac{\pi}{a} x \right) - \frac{1}{4} \sin \left( \frac{3\pi}{a} x \right) \right] \\
&= A \sqrt{\frac{a}{2}} \left[ \frac{3}{4} \psi_1(x) - \frac{1}{4} \psi_3(x) \right] \\
1 &= |A|^2 \frac{a}{2} \int_0^a \left[ \frac{3}{4} \psi_1(x) - \frac{1}{4} \psi_3(x) \right]^2 dx \\
&= |A|^2 \frac{a}{2} \int_0^a \left[ \frac{9}{16} \psi_1(x)^2 - \frac{3}{8} \psi_1(x) \psi_3(x) + \frac{1}{16} \psi_3(x)^2 \right] dx \\
&= \frac{5}{16} a |A|^2 \\
A &= \frac{4}{\sqrt{5a}} \\
\Psi(x, 0) &= \frac{1}{\sqrt{10}} [3\psi_1(x) - \psi_3(x)] \\
\Psi(x, t) &= \frac{1}{\sqrt{10}} [3\psi_1(x) e^{-iE_1 t/\hbar} - \psi_3(x) e^{-iE_3 t/\hbar}] \\
\langle x \rangle &= \int_0^a \Psi^* x \Psi dx \\
&= \frac{1}{10} \int_0^a x \left( 9\psi_1^2 + \psi_3^2 - 3\psi_1\psi_3 e^{-i(E_3-E_1)t/\hbar} - 3\psi_1\psi_3 e^{-i(E_1-E_3)t/\hbar} \right) dx \\
&= \frac{1}{10} \int_0^a x \left[ 9\psi_1^2 + \psi_3^2 - 6\psi_1\psi_3 \cos \left( \frac{E_3-E_1}{\hbar} t \right) \right] dx \\
&= \frac{1}{10} \left[ 9 \langle x \rangle_1 + \langle x \rangle_3 - 6 \cos \left( \frac{E_3-E_1}{\hbar} t \right) \int_0^a x \psi_1 \psi_3 dx \right] \\
&= \frac{1}{10} \left[ \frac{9}{2} a + \frac{1}{2} a \right] \\
&= \frac{a}{2} \\
P(E_1) &= \frac{9}{10} \\
P(E_3) &= \frac{1}{10} \\
\langle E \rangle &= E_1 P(E_1) + E_3 P(E_3) \\
&= \frac{9\pi^2 \hbar^2}{20ma^2} + \frac{9\pi^2 \hbar^2}{20ma^2} \\
&= \frac{9\pi^2 \hbar^2}{10ma^2}
\end{aligned}$$

2.38

(a)

$$\begin{aligned}
 \Psi(x, t) &= \sum_{n=1}^{\infty} c_n \psi_n(x) e^{-iE_n t/\hbar} \\
 \Psi(x, 0) &= \sum_{n=1}^{\infty} c_n \psi_n(x) \\
 E_n T &= \frac{n^2 \pi^2 \hbar^2}{2ma^2} \frac{4ma^2}{\pi \hbar} \\
 &= 2\pi n^2 \hbar \\
 \Psi(x, T) &= \sum_{n=1}^{\infty} c_n \psi_n(x) e^{-iE_n T/\hbar} \\
 &= \sum_{n=1}^{\infty} c_n \psi_n(x) e^{-2\pi i n^2} \\
 &= \sum_{n=1}^{\infty} c_n \psi_n(x) \\
 &= \Psi(x, 0)
 \end{aligned}$$

(b)

$$\begin{aligned}
 E &= \frac{1}{2} m v^2 \\
 v &= \sqrt{\frac{2E}{m}} \\
 T &= \frac{2a}{v} \\
 &= a \sqrt{\frac{2m}{E}}
 \end{aligned}$$

(c)

$$\begin{aligned}
 \frac{4ma^2}{\pi \hbar} &= a \sqrt{\frac{2m}{E}} \\
 \frac{16m^2 a^2}{\pi^2 \hbar^2} &= \frac{2m}{E} \\
 E &= \frac{\pi^2 \hbar^2}{8ma^2} \\
 &= \frac{E_1}{4}
 \end{aligned}$$

### 2.39

(a)

$$\Psi(x, 0) = \begin{cases} \frac{2\sqrt{3}}{a\sqrt{a}}x & 0 \leq x \leq a/2 \\ \frac{2\sqrt{3}}{a\sqrt{a}}(a-x) & a/2 \leq x \leq a \end{cases}$$

$$\frac{d}{dx}\Psi(x, 0) = \frac{2\sqrt{3}}{a\sqrt{a}} \left[ 1 - 2\theta\left(x - \frac{a}{2}\right) \right]$$

(b)

$$\frac{d^2}{dx^2}\Psi(x, 0) = -\frac{4\sqrt{3}}{a\sqrt{a}}\delta\left(x - \frac{a}{2}\right)$$

(c)

$$\begin{aligned} \hat{H}\Psi(x, 0) &= \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right] \Psi(x, 0) \\ &= \frac{\hbar^2}{2m} \frac{4\sqrt{3}}{a\sqrt{a}} \delta\left(x - \frac{a}{2}\right) + V(x)\Psi(x, 0) \\ \langle H \rangle &= \int \Psi(x, 0)^* \hat{H}\Psi(x, 0) dx \\ &= \int_0^a \Psi(x, 0)^* \left[ \frac{\hbar^2}{2m} \frac{4\sqrt{3}}{a\sqrt{a}} \delta\left(x - \frac{a}{2}\right) + V(x)\Psi(x, 0) \right] dx \\ &= \Psi\left(\frac{a}{2}, 0\right)^* \frac{\hbar^2}{2m} \frac{4\sqrt{3}}{a\sqrt{a}} + \int_0^a \Psi(x, 0)^* V(x)\Psi(x, 0) dx \\ &= \frac{6\hbar^2}{ma^2} \end{aligned}$$

## 2.40

(a)

$$\begin{aligned}
V(x) &= \frac{1}{2}m\omega^2 x^2 \\
\xi &= \sqrt{\frac{m\omega}{\hbar}}x \\
\psi_n(x) &= \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(\xi) e^{-\xi^2/2} \\
\Psi(x, 0) &= A \left(1 - 2\sqrt{\frac{m\omega}{\hbar}}x\right)^2 e^{-\frac{m\omega}{2\hbar}x^2} \\
&= A \left(1 - 4\sqrt{\frac{m\omega}{\hbar}}x + \frac{4m\omega}{\hbar}x^2\right) e^{-\frac{m\omega}{2\hbar}x^2} \\
&= A \left(\frac{\pi\hbar}{m\omega}\right)^{1/4} \left[3\psi_0(x) - 2\sqrt{2}\psi_1(x) + 2\sqrt{2}\psi_2(x)\right] \\
1 &= A^2 \sqrt{\frac{\pi\hbar}{m\omega}} \int_{-\infty}^{\infty} (3\psi_0 - 2\sqrt{2}\psi_1 + 2\sqrt{2}\psi_2)^2 dx \\
&= A^2 \sqrt{\frac{\pi\hbar}{m\omega}} \int_{-\infty}^{\infty} (9\psi_0^2 - 12\sqrt{2}\psi_0\psi_1 + 12\sqrt{2}\psi_0\psi_2 + 8\psi_1^2 - 16\psi_1\psi_2 + 8\psi_2^2) dx \\
&= 25A^2 \sqrt{\frac{\pi\hbar}{m\omega}} \\
A &= \frac{1}{5} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \\
\Psi(x, 0) &= \frac{3}{5}\psi_0(x) - \frac{2\sqrt{2}}{5}\psi_1(x) + \frac{2\sqrt{2}}{5}\psi_2(x)
\end{aligned}$$

(b)

$$\begin{aligned}E_0 &= \frac{\hbar\omega}{2} \\P(E_0) &= \frac{9}{25} \\E_1 &= \frac{3\hbar\omega}{2} \\P(E_1) &= \frac{8}{25} \\E_2 &= \frac{5\hbar\omega}{2} \\P(E_2) &= \frac{8}{25} \\\langle E \rangle &= \frac{\hbar\omega}{2} \frac{9}{25} + \frac{3\hbar\omega}{2} \frac{8}{25} + \frac{5\hbar\omega}{2} \frac{8}{25} \\&= \frac{73}{50} \hbar\omega\end{aligned}$$



(c)

$$\begin{aligned}
\xi &= \sqrt{\frac{m\omega}{\hbar}} x \\
\psi_n(x) &= \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(\xi) e^{-\xi^2/2} \\
\Psi(x, T) &= B \left(1 + 2\sqrt{\frac{m\omega}{\hbar}} x\right)^2 e^{-\frac{m\omega}{2\hbar} x^2} \\
&= B \left(1 + 4\sqrt{\frac{m\omega}{\hbar}} x + 4\frac{m\omega}{\hbar} x^2\right) e^{-\frac{m\omega}{2\hbar} x^2} \\
&= B \left(\frac{\pi\hbar}{m\omega}\right)^{1/4} \left[3\psi_0(x) + 2\sqrt{2}\psi_1(x) + 2\sqrt{2}\psi_2(x)\right] \\
1 &= |B|^2 \sqrt{\frac{\pi\hbar}{m\omega}} \int_{-\infty}^{\infty} [3\psi_0 + 2\sqrt{2}\psi_1 + 2\sqrt{2}\psi_2]^2 dx \\
&= |B|^2 \sqrt{\frac{\pi\hbar}{m\omega}} \int_{-\infty}^{\infty} [9\psi_0^2 + 12\sqrt{2}\psi_0\psi_1 + 12\sqrt{2}\psi_0\psi_2 + 8\psi_1^2 + 16\psi_1\psi_2 + 8\psi_2^2] dx \\
&= 25|B|^2 \sqrt{\frac{\pi\hbar}{m\omega}} \\
B &= \frac{1}{5} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \\
\Psi(x, T) &= \frac{3}{5}\psi_0(x) + \frac{2\sqrt{2}}{5}\psi_1(x) + \frac{2\sqrt{2}}{5}\psi_2(x) \\
\Psi(x, t) &= \frac{3}{5}\psi_0(x)e^{-i\omega t/2} - \frac{2\sqrt{2}}{5}\psi_1(x)e^{-3i\omega t/2} + \frac{2\sqrt{2}}{5}\psi_2(x)e^{-5i\omega t/2} \\
&= e^{-i\omega t/2} \left[ \frac{3}{5}\psi_0(x) - \frac{2\sqrt{2}}{5}\psi_1(x)e^{-i\omega t} + \frac{2\sqrt{2}}{5}\psi_2(x)e^{-2i\omega t} \right] \\
e^{-i\omega T} &= -1 \\
e^{-2i\omega T} &= 1 \\
T &= \frac{\pi}{\omega}
\end{aligned}$$

## 2.41

The argument for calculating the allowed energies and wavefunctions is the same, except there is a boundary condition  $\psi(0) = 0$ . This leaves only  $\psi_n(x)$  for odd  $n$ .

2.43

$$\begin{aligned}
k &= \frac{\sqrt{2mE}}{\hbar} \\
\psi(x) &= \begin{cases} -A \sin kx + B \cos kx & -a < x \leq 0 \\ A \sin kx + B \cos kx & 0 \leq x < a \end{cases} \\
\Delta \left( \frac{d\psi}{dx} \right) &= 2Ak \\
\Delta \left( \frac{d\psi}{dx} \right) &= \frac{2m\alpha}{\hbar^2} \psi(0) \\
2Ak &= \frac{2m\alpha}{\hbar^2} B \\
B &= \frac{\hbar^2 k}{m\alpha} A \\
\psi(x) &= A \left( \sin kx + \frac{\hbar^2 k}{m\alpha} \cos kx \right) \\
0 &= A \sin ka + \frac{\hbar^2 k}{m\alpha} A \cos ka \\
\tan ka &= -\frac{\hbar^2 k}{m\alpha} \\
ka &\approx \frac{n\pi}{2}, \quad n = 1, 3, 5, \dots \\
E &\approx \frac{n^2 \pi^2 \hbar^2}{2m(2a)^2}, \quad n = 1, 3, 5, \dots \\
\psi(x) &= \begin{cases} A \sin kx - B \cos kx & -a < x \leq 0 \\ A \sin kx + B \cos kx & 0 \leq x < a \end{cases} \\
-B &= B \\
B &= 0 \\
\psi(x) &= A \sin kx \\
0 &= A \sin ka \\
ka &= \frac{n\pi}{2}, \quad n = 2, 4, 6, \dots \\
\psi(x) &= A \sin \left( \frac{n\pi}{2a} x \right) \\
E &= \frac{n^2 \pi^2 \hbar^2}{2m(2a)^2}, \quad n = 2, 4, 6, \dots
\end{aligned}$$

2.44

$$\begin{aligned}
-\frac{\hbar^2}{2m} \frac{d^2 \psi_1}{dx^2} \psi_2 + V \psi_1 \psi_2 &= E \psi_1 \psi_2 \\
-\frac{\hbar^2}{2m} \frac{d^2 \psi_2}{dx^2} \psi_1 + V \psi_1 \psi_2 &= E \psi_1 \psi_2 \\
\psi_2 \frac{d^2 \psi_1}{dx^2} - \psi_1 \frac{d^2 \psi_2}{dx^2} &= 0 \\
\frac{d}{dx} \left( \psi_2 \frac{d\psi_1}{dx} - \psi_1 \frac{d\psi_2}{dx} \right) &= 0 \\
\psi_2 \frac{d\psi_1}{dx} - \psi_1 \frac{d\psi_2}{dx} &= c \\
c &= 0 \\
\psi_2 \frac{d\psi_1}{dx} &= \psi_1 \frac{d\psi_2}{dx} \\
\frac{1}{\psi_1} \frac{d\psi_1}{dx} &= \frac{1}{\psi_2} \frac{d\psi_2}{dx} \\
\ln \psi_1 &= \ln \psi_2 + c \\
\psi_1 &= A \psi_2
\end{aligned}$$

2.45

(a)

$$\begin{aligned}
-\frac{\hbar^2}{2m} \frac{d^2 \psi_n}{dx^2} \psi_m + V(x) \psi_n \psi_m &= E_n \psi_n \psi_m \\
-\frac{\hbar^2}{2m} \frac{d^2 \psi_m}{dx^2} \psi_n + V(x) \psi_n \psi_m &= E_m \psi_n \psi_m \\
\frac{d^2 \psi_m}{dx^2} \psi_n - \frac{d^2 \psi_n}{dx^2} \psi_m &= \frac{2m}{\hbar^2} (E_n - E_m) \psi_n \psi_m \\
\frac{d}{dx} \left( \frac{d\psi_m}{dx} \psi_n - \frac{d\psi_n}{dx} \psi_m \right) &= \frac{2m}{\hbar^2} (E_n - E_m) \psi_n \psi_m
\end{aligned}$$

(b)

$$\begin{aligned}
\int_{x_1}^{x_2} \frac{d}{dx} (\psi'_m \psi_n - \psi'_n \psi_m) dx &= \frac{2m}{\hbar^2} (E_n - E_m) \int_{x_1}^{x_2} \psi_n \psi_m dx \\
\psi'_m(x_2) \psi_n(x_2) - \psi'_m(x_1) \psi_n(x_1) &= \frac{2m}{\hbar^2} (E_n - E_m) \int_{x_1}^{x_2} \psi_n \psi_m dx
\end{aligned}$$

2.53

(a)

$$\frac{1}{1 - i\beta} \begin{pmatrix} i\beta & 1 \\ 1 & i\beta \end{pmatrix}$$

(b)

$$\frac{e^{-2ika}}{\cos(2la) - i \frac{(k^2 + l^2)}{2kl} \sin(2la)} \begin{pmatrix} i \frac{\sin(2la)}{2kl} (l^2 - k^2) & 1 \\ 1 & i \frac{\sin(2la)}{2kl} (l^2 - k^2) \end{pmatrix}$$

### 3 Formalism

#### 3.1

(a)

$$\begin{aligned} \left| \int_a^b (f^* + g^*)(f + g) dx \right| &= \left| \int_a^b (f^* f + f^* g + g^* f + g^* g) dx \right| \\ &\leq \int_a^b |f|^2 dx + \left| \int_a^b f^* g dx \right| + \left| \int_a^b g^* f dx \right| + \int_a^b |g|^2 dx \\ &\leq \int_a^b |f|^2 dx + 2 \sqrt{\int_a^b |f|^2 dx \int_a^b |g|^2 dx} + \int_a^b |g|^2 dx \end{aligned}$$

The set of all normalised functions isn't a vector space because e.g. multiplying a function by a constant also multiplies its integral by that constant meaning it's no longer a member of the vector space.

(b)

$$\begin{aligned} \langle \beta | \alpha \rangle &= \int_a^b \beta^* \alpha dx \\ &= \left( \int_a^b \alpha^* \beta dx \right)^* \\ &= \langle \alpha | \beta \rangle^* \\ \langle a | a \rangle &= \int_a^b |a|^2 dx \\ &\geq 0 \end{aligned}$$

If  $\langle \alpha | \alpha \rangle = 0$  that implies  $|\alpha|^2 = 0$  everywhere in the interval and thus  $|\alpha\rangle = |0\rangle$ .

$$\begin{aligned} \langle \alpha | (b|\beta\rangle + c|\gamma\rangle) &= \int_{x_1}^{x_2} \alpha^* (b\beta) dx + \int_{x_1}^{x_2} \alpha^* (c\gamma) dx \\ &= b \int_{x_1}^{x_2} \alpha^* \beta dx + c \int_{x_1}^{x_2} \alpha^* \gamma dx \\ &= b \langle \alpha | \beta \rangle + c \langle \alpha | \gamma \rangle \end{aligned}$$

### 3.2

(a)

$$\int_0^1 x^{2\nu} dx = \frac{1}{2\nu+1} [x^{2\nu+1}]_0^1$$

The integral is defined for  $\nu > -1/2$ . For the case  $\nu = -1/2$

$$\int_0^1 x^{-1} dx = [\ln x]_0^1 = \ln 1 - \ln 0 = 0 - \infty.$$

So  $f(x) = x^\nu$  is in Hilbert space for  $\nu > -1/2$ .

(b)

$$\begin{aligned} \int_0^1 x dx &= \frac{1}{2} \\ \int_0^1 x^3 dx &= \frac{1}{4} \\ \int_0^1 x^{-1} dx &= [\ln x]_0^1 \\ &= 0 - \infty \end{aligned}$$

$f(x)$  and  $xf(x)$  are in Hilbert space, but not  $(d/dx)f(x)$ .

### 3.4

(a)

$$\begin{aligned} \langle f | (\hat{Q} + \hat{R}) f \rangle &= \langle f | \hat{Q} f \rangle + \langle f | \hat{R} f \rangle \\ &= \langle \hat{Q} f | f \rangle + \langle \hat{R} f | f \rangle \\ &= \langle (\hat{Q} + \hat{R}) f | f \rangle \end{aligned}$$

(b)

$$\begin{aligned} \langle f | \alpha \hat{Q} g \rangle &= \alpha \langle f | \hat{Q} g \rangle \\ &= \alpha \langle \hat{Q} f | g \rangle \\ \langle \alpha \hat{Q} f | g \rangle &= \alpha^* \langle \hat{Q} f | g \rangle \\ \alpha &= \alpha^* \end{aligned}$$

$\alpha$  is real.

(c)

$$\begin{aligned}\langle f|\hat{Q}\hat{R}g\rangle &= \langle \hat{Q}f|\hat{R}g\rangle \\ &= \langle \hat{R}\hat{Q}f|g\rangle\end{aligned}$$

The product of the operators is hermitian when  $\hat{Q}\hat{R} = \hat{R}\hat{Q}$  i.e.  $[\hat{Q}, \hat{R}] = 0$ .

(d)

$$\begin{aligned}\langle \Psi|\hat{x}\Psi\rangle &= \int \Psi^* \hat{x} \Psi dx \\ &= \int \Psi^* \hat{x}^* \Psi dx \\ &= \int (\hat{x}\Psi)^* \Psi dx \\ &= \langle \hat{x}\Psi|\Psi\rangle \\ \langle \Psi|\hat{H}\Psi\rangle &= \int \Psi^* \hat{H} \Psi dx \\ &= \int \Psi^* \left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right] \Psi dx \\ &= -\frac{\hbar^2}{2m} \int \Psi^* \frac{d^2 \Psi}{dx^2} dx + \int \Psi^* V(x) \Psi dx \\ &= -\frac{\hbar^2}{2m} \left[ \Psi^* \frac{d\Psi}{dx} \Big|_{-\infty}^{\infty} - \int \frac{d\Psi^*}{dx} \frac{d\Psi}{dx} dx \right] + \langle V(x)\Psi|\Psi\rangle \\ &= \frac{\hbar^2}{2m} \left[ \frac{d\Psi^*}{dx} \Psi \Big|_{-\infty}^{\infty} - \int \frac{d^2 \Psi^*}{dx^2} \Psi dx \right] + \langle V(x)\Psi|\Psi\rangle \\ &= -\frac{\hbar^2}{2m} \int \frac{d^2}{dx^2} \Psi^* \Psi dx + \langle V(x)\Psi|\Psi\rangle \\ &= \left\langle -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \Psi \middle| \Psi \right\rangle + \langle V(x)\Psi|\Psi\rangle \\ &= \langle \hat{H}\Psi|\Psi\rangle\end{aligned}$$

### 3.5

(a)

$$\begin{aligned}x^\dagger &= x \\ i^\dagger &= -i \\ \left( \frac{d}{dx} \right)^\dagger &= -\frac{d}{dx}\end{aligned}$$

(b)

$$\begin{aligned}
\langle f|\hat{Q}\hat{R}g\rangle &= \int f^\dagger \hat{Q}\hat{R}g \, dx \\
&= \int (\hat{Q}^\dagger f)^\dagger \hat{R}g \, dx \\
&= \int (\hat{R}^\dagger \hat{Q}^\dagger f)^\dagger g \, dx \\
&= \langle \hat{R}^\dagger \hat{Q}^\dagger f|g\rangle
\end{aligned}$$

(c)

$$\begin{aligned}
\hat{a}_+ &= \frac{1}{\sqrt{2\hbar m\omega}}(-i\hat{p} + m\omega x) \\
\langle f|\hat{a}g\rangle &= \left\langle f \left| \frac{1}{\sqrt{2\hbar m\omega}}(-i\hat{p} + m\omega x)g \right. \right\rangle \\
&= \frac{1}{\sqrt{2\hbar m\omega}} \langle f|(-i\hat{p} + m\omega x)g\rangle \\
&= \frac{1}{\sqrt{2\hbar m\omega}} (\langle f|-i\hat{p}g\rangle + \langle f|m\omega xg\rangle) \\
&= \frac{1}{\sqrt{2\hbar m\omega}} (\langle f|-i\hat{p}g\rangle + \langle m\omega x f|g\rangle) \\
&= \frac{1}{\sqrt{2\hbar m\omega}} (\langle i\hat{p}f|g\rangle + \langle m\omega x f|g\rangle) \\
&= \left\langle \frac{1}{\sqrt{2\hbar m\omega}}(i\hat{p} + m\omega x)f \left| g \right. \right\rangle \\
&= \langle \hat{a}_- f|g\rangle
\end{aligned}$$

### 3.6

$$\begin{aligned}
\langle f|\hat{Q}g\rangle &= \int_0^{2\pi} f^* \frac{d^2 g}{d\phi^2} d\phi \\
&= f^* \frac{dg}{d\phi} \Big|_0^{2\pi} - \int_0^{2\pi} \left( \frac{df}{d\phi} \right)^* \frac{dg}{d\phi} d\phi \\
&= - \left( \frac{df}{d\phi} \right)^* g \Big|_0^{2\pi} + \int_0^{2\pi} \left( \frac{d^2 f}{d\phi^2} \right)^* g d\phi \\
&= \langle \hat{Q}f|g\rangle
\end{aligned}$$

Yes, the operator is hermitian.

$$\begin{aligned}
\hat{Q}f &= qf \\
\frac{d^2 f}{d\phi^2} &= qf \\
\frac{d^2 f}{d\phi^2} - qf &= 0 \\
f &= Ae^{\sqrt{q}\phi} + Be^{-\sqrt{q}\phi} \\
f(\phi + 2\pi) &= Ae^{\sqrt{q}(\phi+2\pi)} + Be^{\sqrt{q}(\phi+2\pi)} \\
&= Ae^{\sqrt{q}\phi} e^{2\pi\sqrt{q}} + Be^{\sqrt{q}\phi} e^{2\pi\sqrt{q}} \\
2\pi\sqrt{q} &= 1 \\
q &= -n^2, \quad n = 0, 1, 2, \dots
\end{aligned}$$

The eigenfunctions are  $f = Ae^{\pm\sqrt{q}\phi}$  and the eigenvalues are  $q = 0, 1, 2, \dots$ . The spectrum is degenerate as there are two eigenfunctions associated with each eigenvalue  $q > 0$ .

### 3.7

(a)

$$\begin{aligned}
h &= af + bg \\
\hat{Q}h &= \hat{Q}(af + bg) \\
&= \hat{Q}(af) + \hat{Q}(bg) \\
&= a\hat{Q}f + b\hat{Q}g \\
&= aqf + bqg \\
&= q(af + bg) \\
&= qh
\end{aligned}$$

(b)

$$\begin{aligned}
\frac{d^2}{dx^2} e^x &= e^x \\
\frac{d^2}{dx^2} e^{-x} &= e^{-x} \\
f &= e^x + e^{-x} \\
g &= e^x - e^{-x}
\end{aligned}$$



### 3.8

(a)

$$\begin{aligned}
 \hat{Q} &= i \frac{d}{d\phi} \\
 \hat{Q}f &= qf \\
 i \frac{df}{d\phi} &= qf \\
 \frac{df}{d\phi} + iqf &= 0 \\
 f &= Ae^{-iq\phi} \\
 e^{-2\pi iq} &= 1 \\
 q &= 0, \pm 1, \pm 2, \dots \\
 \int_0^{2\pi} Ae^{-iq\phi} Be^{-iq'\phi} d\phi &= AB \int_0^{2\pi} e^{-i(q+q')\phi} d\phi \\
 &= 0
 \end{aligned}$$

(b)

$$\begin{aligned}
 \hat{Q} &= \frac{d^2}{d\phi^2} \\
 \hat{Q}f &= qf \\
 \frac{d^2 f}{d\phi^2} - qf &= 0 \\
 f &= Ae^{\pm\sqrt{q}\phi} \\
 q &= -n^2, n = 0, 1, 2, \dots \\
 \int_0^{2\pi} Ae^{\pm\sqrt{q}\phi} Be^{\pm\sqrt{q'}\phi} d\phi &= AB \int_0^{2\pi} e^{\pm in\phi} e^{\pm in'\phi} d\phi \\
 &= AB \int_0^{2\pi} e^{i(\pm n \pm n')\phi} d\phi \\
 &= AB \left[ \frac{1}{i(\pm n \pm n')} e^{i(\pm n \pm n')\phi} \right]_0^{2\pi} \\
 &= AB \frac{1}{i(\pm n \pm n')} [e^{i(\pm n \pm n')2\pi} - 1] \\
 &= 0
 \end{aligned}$$

### 3.9

(a) Infinite square well

(b) Delta function barrier

(c) Delta function well

### 3.11

$$\begin{aligned}
\Psi_0(x, t) &= \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-m\omega x^2/2\hbar} e^{-i\omega t/2} \\
\Phi_0(p, t) &= \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{-ipx/\hbar} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-m\omega x^2/2\hbar} e^{-i\omega t/2} dx \\
&= \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-i\omega t/2} \int_{-\infty}^{\infty} e^{-ipx/\hbar} e^{-m\omega x^2/2\hbar} dx \\
&= \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-i\omega t/2} \sqrt{\frac{2\pi\hbar}{m\omega}} e^{-p^2/2\hbar m\omega} \\
&= \frac{1}{(\pi\hbar m\omega)^{1/4}} e^{-p^2/2\hbar m\omega} e^{-i\omega t/2} \\
\frac{p^2}{2m} &= \frac{\hbar\omega}{2} \\
p &= \pm\sqrt{\hbar m\omega} \\
1 - \int_{-\sqrt{\hbar m\omega}}^{\sqrt{\hbar m\omega}} |\Phi_0|^2 dp &= 1 - \frac{1}{(\pi\hbar m\omega)^{1/4}} \int_{-\sqrt{\hbar m\omega}}^{\sqrt{\hbar m\omega}} e^{-p^2/\hbar m\omega} dp \\
&= 1 - \frac{1}{(\pi\hbar m\omega)^{1/2}} \sqrt{\pi\hbar m\omega} \operatorname{erf} 1 \\
&= 0.16
\end{aligned}$$

### 3.12

$$\begin{aligned}
\Psi(x, t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{i(kx - \frac{\hbar k^2}{2m}t)} dk \\
\Phi(p, t) &= \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{-ipx/\hbar} \left( \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{i(kx - \frac{\hbar k^2}{2m}t)} dk \right) dx \\
&= \frac{1}{2\pi\sqrt{\hbar}} \int_{-\infty}^{\infty} \phi(k) e^{-i\frac{\hbar k^2}{2m}t} \left( \int_{-\infty}^{\infty} e^{i(k - p/\hbar)x} dx \right) dk \\
&= \frac{1}{2\pi\sqrt{\hbar}} \int_{-\infty}^{\infty} \phi(k) e^{-i\frac{\hbar k^2}{2m}t} 2\pi\delta(k - p/\hbar) dk \\
&= \frac{1}{\sqrt{\hbar}} \int_{-\infty}^{\infty} \delta(k - p/\hbar) \phi(k) e^{-i\frac{\hbar k^2}{2m}t} dk \\
&= \frac{1}{\sqrt{\hbar}} \phi(p/\hbar) e^{-i\frac{p^2}{2\hbar m}t} \\
|\Phi(p, t)|^2 &= \frac{1}{\hbar} |\phi(p/\hbar)|^2
\end{aligned}$$

### 3.14

(a)

$$\begin{aligned}
 [\hat{A} + \hat{B}, \hat{C}] &= (\hat{A} + \hat{B})\hat{C} - \hat{C}(\hat{A} + \hat{B}) \\
 &= \hat{A}\hat{C} + \hat{B}\hat{C} - \hat{C}\hat{A} - \hat{C}\hat{B} \\
 &= \hat{A}\hat{C} - \hat{C}\hat{A} + \hat{B}\hat{C} - \hat{C}\hat{B} \\
 &= [\hat{A}, \hat{C}] + [\hat{B}, \hat{C}] \\
 \hat{A}[\hat{B}, \hat{C}] + [\hat{A}, \hat{C}]\hat{B} &= \hat{A}\hat{B}\hat{C} - \hat{A}\hat{C}\hat{B} + \hat{A}\hat{C}\hat{B} - \hat{C}\hat{A}\hat{B} \\
 &= \hat{A}\hat{B}\hat{C} - \hat{C}\hat{A}\hat{B} \\
 &= [\hat{A}\hat{B}, \hat{C}]
 \end{aligned}$$

(b)

$$\begin{aligned}
 [x^n, \hat{p}] &= \left[ x^n, -i\hbar \frac{d}{dx} \right] \\
 &= x^n \left( -i\hbar \frac{d}{dx} \right) - \left( -i\hbar \frac{d}{dx} \right) x^n \\
 &= x^n \left( -i\hbar \frac{d}{dx} \right) (1) + i\hbar n x^{n-1} \\
 &= i\hbar n x^{n-1}
 \end{aligned}$$

(c)

$$\begin{aligned}
 [f(x), \hat{p}]g(x) &= f(x) \left( -i\hbar \frac{d}{dx} \right) g(x) - \left( -i\hbar \frac{d}{dx} \right) [f(x)g(x)] \\
 &= -i\hbar f(x) \frac{dg}{dx} + i\hbar \left[ \frac{df}{dx} g(x) + f(x) \frac{dg}{dx} \right] \\
 &= i\hbar \frac{df}{dx} g(x) \\
 [f(x), \hat{p}] &= i\hbar \frac{df}{dx}
 \end{aligned}$$

(d)

$$\begin{aligned}
[\hat{H}, \hat{a}_-]g &= \hat{H}\hat{a}_-g - \hat{a}_-\hat{H}g \\
&= \hbar\omega \left( \hat{a}_-\hat{a}_+ - \frac{1}{2} \right) \hat{a}_-g - \hat{a}_-\hbar\omega \left( \hat{a}_-\hat{a}_+ - \frac{1}{2} \right) g \\
&= \hbar\omega \hat{a}_-\hat{a}_+\hat{a}_-g - \frac{1}{2}\hbar\omega \hat{a}_-g - \hbar\omega \hat{a}_-^2\hat{a}_+g + \frac{1}{2}\hbar\omega \hat{a}_-g \\
&= \hbar\omega \hat{a}_-\hat{a}_+\hat{a}_-g - \hbar\omega \hat{a}_-^2\hat{a}_+g \\
&= \hbar\omega \hat{a}_-(\hat{a}_+\hat{a}_- - \hat{a}_-\hat{a}_+)g \\
&= -\hbar\omega \hat{a}_-g \\
[\hat{H}, \hat{a}_-] &= -\hbar\omega \hat{a}_- \\
[\hat{H}, \hat{a}_+]g &= \hat{H}\hat{a}_+g - \hat{a}_+\hat{H}g \\
&= \hbar\omega \left( \hat{a}_-\hat{a}_+ - \frac{1}{2} \right) \hat{a}_+g - \hat{a}_+\hbar\omega \left( \hat{a}_-\hat{a}_+ - \frac{1}{2} \right) g \\
&= \hbar\omega \hat{a}_-\hat{a}_+^2g - \frac{1}{2}\hbar\omega \hat{a}_+g - \hbar\omega \hat{a}_+\hat{a}_-\hat{a}_+g + \frac{1}{2}\hbar\omega \hat{a}_+g \\
&= \hbar\omega \hat{a}_-\hat{a}_+^2g - \hbar\omega \hat{a}_+\hat{a}_-\hat{a}_+g \\
&= \hbar\omega (\hat{a}_-\hat{a}_+ - \hat{a}_+\hat{a}_-)\hat{a}_+g \\
&= \hbar\omega \hat{a}_+g \\
[\hat{H}, \hat{a}_+] &= \hbar\omega \hat{a}_+
\end{aligned}$$

### 3.15

$$\begin{aligned}
\left[ x, \frac{p^2}{2m} + V \right] g &= x \left( \frac{p^2}{2m} + V \right) g - \left( \frac{p^2}{2m} + V \right) xg \\
&= x \frac{p^2}{2m} g + xVg - \frac{p^2}{2m} xg - Vxg \\
&= \frac{1}{2m} (xp^2g - p^2xg) \\
&= \frac{1}{2m} \left[ -\hbar^2 x \frac{d^2g}{dx^2} + \hbar^2 \frac{d}{dx} \left( g + x \frac{dg}{dx} \right) \right] \\
&= \frac{1}{2m} \left[ -\hbar^2 x \frac{d^2g}{dx^2} + \hbar^2 \left( \frac{dg}{dx} + \frac{dg}{dx} + x \frac{d^2g}{dx^2} \right) \right] \\
&= \frac{\hbar^2}{m} \frac{dg}{dx} \\
\left[ x, \frac{p^2}{2m} + V \right] &= \frac{\hbar^2}{m} \frac{d}{dx} \\
&= -\frac{\hbar}{im} \langle p \rangle \\
\sigma_x^2 \sigma_H^2 &\geq \left( \frac{1}{2i} \left\langle \left[ x, \frac{p^2}{2m} + V \right] \right\rangle \right)^2 \\
&= \frac{\hbar^2}{4m^2} |\langle p \rangle|^2 \\
\sigma_x \sigma_H &\geq \frac{\hbar}{2m} |\langle p \rangle|
\end{aligned}$$

This doesn't tell us much because for stationary states  $\sigma_H = 0$  and  $\langle p \rangle = 0$  so this says  $0 \geq 0$ .

### 3.17

$$\begin{aligned}
\left( -i\hbar \frac{d}{dx} - \langle p \rangle \right) \Psi &= ia(x - \langle x \rangle) \Psi \\
-i\hbar \frac{d\Psi}{dx} - \langle p \rangle \Psi &= ia(x - \langle x \rangle) \Psi \\
\frac{d\Psi}{dx} + \frac{\langle p \rangle + ia(x - \langle x \rangle)}{i\hbar} \Psi &= 0 \\
\frac{d\Psi}{dx} + \left[ \frac{a}{\hbar} (x - \langle x \rangle) - i \frac{\langle p \rangle}{\hbar} \right] \Psi &= 0 \\
\frac{d\Psi}{dx} + \frac{a}{\hbar} x \Psi - \frac{a}{\hbar} \langle x \rangle \Psi - i \frac{\langle p \rangle}{\hbar} \Psi &= 0 \\
\Psi &= A e^{-ax^2/2\hbar} e^{a\langle x \rangle x/\hbar} e^{i\langle p \rangle x/\hbar} \\
&= B e^{-a(x - \langle x \rangle)^2/2\hbar} e^{i\langle p \rangle x/\hbar}
\end{aligned}$$

### 3.18

(a)

$$\begin{aligned}Q &= 1 \\ \hat{Q} &= 1 \\ [\hat{H}, \hat{Q}] &= 0 \\ \frac{d}{dt} \langle Q \rangle &= 0\end{aligned}$$

(b)

$$\begin{aligned}Q &= H \\ \hat{Q} &= \hat{H} \\ [\hat{H}, \hat{Q}] &= [\hat{H}, \hat{H}] \\ &= 0 \\ \frac{d}{dt} \langle H \rangle &= 0\end{aligned}$$

(c)

$$\begin{aligned}Q &= x \\ \hat{Q} &= x \\ [\hat{H}, \hat{Q}] &= [\hat{H}, x] \\ &= -i \frac{\hbar}{m} \hat{p} \\ \frac{d}{dt} \langle x \rangle &= \frac{i}{\hbar} \left\langle -i \frac{\hbar}{m} p \right\rangle \\ &= \frac{\langle p \rangle}{m}\end{aligned}$$

(d)

$$\begin{aligned}Q &= p \\ \hat{Q} &= \hat{p} \\ [\hat{H}, \hat{Q}] &= [\hat{H}, \hat{p}] \\ &= i\hbar \frac{\partial V}{\partial x} \\ \frac{d}{dt} \langle p \rangle &= \frac{i}{\hbar} \left\langle i\hbar \frac{\partial V}{\partial x} \right\rangle \\ &= - \left\langle \frac{\partial V}{\partial x} \right\rangle\end{aligned}$$

### 3.19

(a)

$$\begin{aligned}\frac{d^2}{dt^2} \langle x \rangle &= \frac{d}{dx} \left( \frac{\langle p \rangle}{m} \right) \\ &= -\frac{1}{m} \left\langle \frac{\partial V}{\partial x} \right\rangle \\ &= 0\end{aligned}$$

(b)

$$\begin{aligned}\frac{d^2}{dt^2} \langle x \rangle &= \frac{d}{dx} \left( \frac{\langle p \rangle}{m} \right) \\ &= -\frac{1}{m} \left\langle \frac{\partial V}{\partial x} \right\rangle \\ &= -\omega^2 \langle x \rangle \\ \frac{d^2}{dt^2} \langle x \rangle + \omega^2 \langle x \rangle &= 0 \\ \langle x \rangle &= A \sin \omega t + B \cos \omega t\end{aligned}$$

### 3.20

$$\begin{aligned}\Psi &= \frac{1}{\sqrt{2}}(\psi_1 e^{-iE_1 t/\hbar} + \psi_2 e^{-iE_2 t/\hbar}) \\ \hat{H}\Psi &= \frac{1}{\sqrt{2}}(E_1 \psi_1 e^{-iE_1 t/\hbar} + E_2 \psi_2 e^{-iE_2 t/\hbar}) \\ \hat{H}^2\Psi &= \frac{1}{\sqrt{2}}(E_1^2 \psi_1 e^{-iE_1 t/\hbar} + E_2^2 \psi_2 e^{-iE_2 t/\hbar}) \\ \langle H^2 \rangle &= \langle \Psi | \hat{H}^2 | \Psi \rangle \\ &= \frac{1}{2} \int_0^a (\psi_1^* e^{iE_1 t/\hbar} + \psi_2^* e^{iE_2 t/\hbar})(E_1^2 \psi_1 e^{-iE_1 t/\hbar} + E_2^2 \psi_2 e^{-iE_2 t/\hbar}) dx \\ &= \frac{1}{2} \int_0^a (E_1^2 |\psi_1|^2 + E_2^2 \psi_1^* \psi_2^* e^{i(E_1 - E_2)t/\hbar} + E_1^2 \psi_2^* \psi_1 e^{i(E_2 - E_1)t/\hbar} \\ &\quad + E_2^2 |\psi_2|^2) dx \\ &= \frac{1}{2}(E_1^2 + E_2^2) \\ \langle H \rangle &= \frac{1}{2}(E_1 + E_2) \\ \sigma_H^2 &= \langle H^2 \rangle - \langle H \rangle^2 \\ &= \frac{1}{2}(E_1^2 + E_2^2) - \frac{1}{4}(E_1 + E_2)^2 \\ &= \frac{1}{4}(E_2 - E_1)^2\end{aligned}$$

$$\begin{aligned}
\omega &= \frac{\pi^2 \hbar}{2ma^2} \\
\langle x \rangle &= \frac{a}{2} \left[ 1 - \frac{32}{9\pi^2} \cos(3\omega t) \right] \\
\langle x^2 \rangle &= \frac{1}{2} \int_0^a (\psi_1^* e^{iE_1 t/\hbar} + \psi_2^* e^{iE_2 t/\hbar}) x^2 (\psi_1 e^{-iE_1 t/\hbar} + \psi_2 e^{-iE_2 t/\hbar}) dx \\
&= \frac{1}{2} \int_0^a x^2 (|\psi_1|^2 + \psi_1^* \psi_2 e^{i(E_1 - E_2)t/\hbar} + \psi_2^* \psi_1 e^{i(E_2 - E_1)t/\hbar} + |\psi_2|^2) dx \\
&= \frac{1}{a} \int_0^a x^2 \left[ \sin^2\left(\frac{\pi}{a}x\right) + \sin^2\left(\frac{2\pi}{a}x\right) \right. \\
&\quad \left. + 2 \sin\left(\frac{\pi}{a}x\right) \sin\left(\frac{2\pi}{a}x\right) \cos(3\omega t) \right] dx \\
&= \frac{a^2}{144\pi^2} [-45 + 48\pi^2 - 256 \cos(3\omega t)] \\
\sigma_x^2 &= \langle x^2 \rangle - \langle x \rangle^2 \\
&= \frac{a^2}{4} \left[ \frac{1}{3} - \frac{5}{4\pi^2} - \left( \frac{32}{9\pi^2} \right)^2 \cos^2(3\omega t) \right] \\
\frac{d\langle x \rangle}{dt} &= \frac{16a\omega}{3\pi^2} \sin(3\omega t) \\
&= \frac{8\hbar}{3ma} \sin(3\omega t)
\end{aligned}$$

$$\begin{aligned}
\sigma_H^2 \sigma_x^2 &\geq \left( \frac{\hbar}{2} \right)^2 \left| \frac{d\langle x \rangle}{dt} \right|^2 \\
\frac{1}{4} \left( \frac{3\pi^2 \hbar^2}{2ma^2} \right)^2 \frac{a^2}{4} \left[ \frac{1}{3} - \frac{5}{4\pi^2} - \left( \frac{32}{9\pi^2} \right)^2 \cos^2(3\omega t) \right] &\geq \left( \frac{\hbar}{2} \right)^2 \left[ \frac{8\hbar}{3ma} \sin(3\omega t) \right]^2 \\
\left( \frac{3}{4} \right)^2 \left[ \frac{1}{3} - \frac{5}{4\pi^2} - \left( \frac{32}{9\pi^2} \right)^2 \cos^2(3\omega t) \right] &\geq \left( \frac{8}{3\pi^2} \right)^2 \sin^2(3\omega t) \\
\frac{1}{3} - \frac{5}{4\pi^2} &\geq \left( \frac{32}{9\pi^2} \right)^2
\end{aligned}$$



### 3.23

$$\begin{aligned}
\hat{P}^2 |\beta\rangle &= \hat{P}(\hat{P} |\beta\rangle) \\
&= \hat{P}(\langle\alpha|\beta\rangle |\alpha\rangle) \\
&= \langle\alpha|\beta\rangle \hat{P} |\alpha\rangle \\
&= \langle\alpha|\beta\rangle (\langle\alpha|\alpha\rangle |\alpha\rangle) \\
&= \langle\alpha|\beta\rangle |\alpha\rangle \\
&= \hat{P} |\beta\rangle
\end{aligned}$$

So,  $\hat{P}^2 = \hat{P}$ .

$$\begin{aligned}
\hat{P} |\beta\rangle &= \lambda |\beta\rangle \\
\langle\alpha|\beta\rangle |\alpha\rangle &= \lambda |\beta\rangle
\end{aligned}$$

If  $|\beta\rangle$  is a constant multiple of  $|\alpha\rangle$ , then

$$\begin{aligned}
\langle\alpha|c\alpha\rangle |\alpha\rangle &= \lambda c |\alpha\rangle \\
c \langle\alpha|\alpha\rangle |\alpha\rangle &= \lambda c |\alpha\rangle \\
c |\alpha\rangle &= \lambda c |\alpha\rangle
\end{aligned}$$

thus  $\lambda = 1$ .

If  $|\beta\rangle$  is orthogonal to  $|\alpha\rangle$ , then

$$\begin{aligned}
\langle\alpha|\beta\rangle |\alpha\rangle &= \lambda |\beta\rangle \\
0 &= \lambda |\beta\rangle
\end{aligned}$$

thus  $\lambda = 0$ .

### 3.24

$$\begin{aligned}
\hat{Q} &= \hat{Q}^\dagger \\
Q_{mn} &= \langle e_m | \hat{Q} | e_n \rangle \\
Q_{nm} &= \langle e_n | \hat{Q} | e_m \rangle \\
Q_{nm}^* &= \langle e_n | \hat{Q} | e_m \rangle^* \\
&= \langle e_m | \hat{Q}^\dagger | e_n \rangle \\
&= \langle e_m | \hat{Q} | e_n \rangle \\
&= Q_{mn}
\end{aligned}$$

### 3.25

$$\hat{H}|\alpha\rangle = \lambda|\alpha\rangle$$

$$\epsilon[(\langle 1|\alpha\rangle + \langle 2|\alpha\rangle)|1\rangle + (\langle 1|\alpha\rangle - \langle 2|\alpha\rangle)|2\rangle] = \lambda(\langle 1|\alpha\rangle|1\rangle + \langle 2|\alpha\rangle|2\rangle)$$

From this we get two equations

$$\epsilon(\langle 1|\alpha\rangle + \langle 2|\alpha\rangle) = \lambda\langle 1|\alpha\rangle$$

$$\epsilon(\langle 1|\alpha\rangle - \langle 2|\alpha\rangle) = \lambda\langle 2|\alpha\rangle$$

If we let  $|\alpha\rangle = \begin{pmatrix} a \\ b \end{pmatrix}$  then this becomes

$$\epsilon(a + b) = \lambda a$$

$$\epsilon(a - b) = \lambda b$$

or in matrix form

$$\epsilon \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} = \lambda \begin{pmatrix} a \\ b \end{pmatrix}.$$

The eigenvalues of this matrix are  $\lambda = \pm\sqrt{2}\epsilon$  and the eigenvectors are  $|1\rangle + (\sqrt{2} \pm 1)|2\rangle$ .

### 3.26

(a)

$$\langle\alpha| = -i\langle 1| - 2\langle 2| + i\langle 3|$$

$$\langle\beta| = -i\langle 1| + 2\langle 3|$$

(b)

$$\begin{aligned} \langle\alpha|\beta\rangle &= (-i\langle 1| - 2\langle 2| + i\langle 3|)(i|1\rangle + 2|3\rangle) \\ &= \langle 1|1\rangle - 2i\langle 1|3\rangle - 2i\langle 2|1\rangle - 4\langle 2|3\rangle - \langle 3|1\rangle + 2i\langle 3|3\rangle \\ &= 1 + 2i \\ \langle\beta|\alpha\rangle &= (-i\langle 1| + 2\langle 3|)(i|1\rangle - 2|2\rangle - i|3\rangle) \\ &= \langle 1|1\rangle + 2i\langle 1|2\rangle - \langle 1|3\rangle + 2i\langle 3|1\rangle - 4\langle 3|2\rangle - 2i\langle 3|3\rangle \\ &= 1 - 2i \\ &= \langle\alpha|\beta\rangle^* \end{aligned}$$

(c)

$$\begin{aligned}
\hat{A} &= |\alpha\rangle \langle \beta| \\
\hat{A} |1\rangle &= |\alpha\rangle \langle \beta|1\rangle \\
&= |\alpha\rangle \langle 1|\beta\rangle^* \\
&= -i |\alpha\rangle \\
&= |1\rangle + 2i |2\rangle - |3\rangle \\
A_{11} &= \langle 1|\hat{A}|1\rangle \\
&= 1 \\
A_{21} &= 2i \\
A_{31} &= -1 \\
\hat{A} |2\rangle &= |\alpha\rangle \langle \beta|2\rangle \\
&= |\alpha\rangle \langle 2|\beta\rangle^* \\
&= 0 \\
A_{12} &= 0 \\
A_{22} &= 0 \\
A_{32} &= 0 \\
\hat{A} |3\rangle &= |\alpha\rangle \langle \beta|3\rangle \\
&= |\alpha\rangle \langle 3|\beta\rangle^* \\
&= 2 |\alpha\rangle \\
&= 2i |1\rangle - 4 |2\rangle - 2i |3\rangle \\
A_{13} &= 2i \\
A_{23} &= -4 \\
A_{33} &= -2i \\
A &= \begin{pmatrix} 1 & 0 & 2i \\ 2i & 0 & -4 \\ -1 & 0 & -2i \end{pmatrix}
\end{aligned}$$

It's not hermitian.

### 3.27

(a)

$$\begin{aligned}
 \hat{Q}|\alpha\rangle &= \hat{Q} \sum_{n=1}^{\infty} \langle e_n|\alpha\rangle |e_n\rangle \\
 &= \hat{Q} \left( \sum_{n=1}^{\infty} |e_n\rangle \langle e_n| \right) |\alpha\rangle \\
 &= \left( \sum_{n=1}^{\infty} \hat{Q} |e_n\rangle \langle e_n| \right) |\alpha\rangle \\
 &= \left( \sum_{n=1}^{\infty} q_n |e_n\rangle \langle e_n| \right) |\alpha\rangle \\
 \hat{Q} &= \sum_{n=1}^{\infty} q_n |e_n\rangle \langle e_n|
 \end{aligned}$$

(b)

$$\begin{aligned}
 \hat{Q} &= \sum_{n=1}^{\infty} q_n |e_n\rangle \langle e_n| \\
 \hat{Q}^2 &= \left( \sum_{n=1}^{\infty} q_n |e_n\rangle \langle e_n| \right) \left( \sum_{l=1}^{\infty} q_l |e_l\rangle \langle e_l| \right) \\
 &= \sum_{n=1}^{\infty} \sum_{l=0}^{\infty} q_n q_l |e_n\rangle \langle e_n| e_l\rangle \langle e_l| \\
 &= \sum_{n=1}^{\infty} q_n^2 |e_n\rangle \langle e_n| \\
 e^{\hat{Q}} &= \sum_{n=1}^{\infty} e^{q_n} |e_n\rangle \langle e_n| \\
 &= \sum_{n=1}^{\infty} \left( \sum_{k=0}^{\infty} \frac{q_n^k}{k!} \right) |e_n\rangle \langle e_n| \\
 &= \sum_{k=0}^{\infty} \frac{1}{k!} \left( \sum_{n=1}^{\infty} q_n^k |e_n\rangle \langle e_n| \right) \\
 &= \sum_{k=0}^{\infty} \frac{1}{k!} \hat{Q}^k \\
 &= 1 + \hat{Q} + \frac{1}{2} \hat{Q}^2 + \frac{1}{3!} \hat{Q}^3 + \dots
 \end{aligned}$$

### 3.28

(a)

$$\begin{aligned}\sin \hat{D} &= \hat{D} - \frac{D^3}{3!} + \frac{\hat{D}^5}{5!} - \frac{\hat{D}^7}{7!} + \cdots \\ (\sin \hat{D})x^5 &= 5x^4 - 10x^2 + 1\end{aligned}$$

(b)

$$\begin{aligned}\frac{1}{1 - \hat{D}/2} &= 1 + \frac{\hat{D}}{2} + \left(\frac{\hat{D}}{2}\right)^2 + \left(\frac{\hat{D}}{2}\right)^3 + \cdots \\ &= 1 + \frac{\hat{D}}{2} + \frac{\hat{D}^2}{4} + \frac{\hat{D}^3}{8} + \cdots \\ \frac{1}{1 - \hat{D}/2} \cos x &= \cos x - \frac{1}{2} \sin x - \frac{1}{4} \cos x + \frac{1}{8} \sin x + \cdots \\ &= \left(-\frac{1}{2} + \frac{1}{8} - \frac{1}{32} + \cdots\right) \sin x + \left(1 - \frac{1}{4} + \frac{1}{16} - \cdots\right) \cos x \\ &= -\frac{2}{5} \sin x + \frac{4}{5} \cos x\end{aligned}$$

### 3.30

$$\begin{aligned}c_n(t) &= \langle n | S(t) \rangle \\ &= \left\langle n \left| \int dx |x\rangle \langle x| \right| S(t) \right\rangle \\ &= \int \langle n | x \rangle \langle x | S(t) \rangle dx \\ &= \int \langle x | n \rangle^* \Psi(x, t) dx \\ &= \int \psi_n(x)^* \Psi(x, t) dx\end{aligned}$$

### 3.31

$$\begin{aligned}
|e_1\rangle &= 1 \\
\langle e_1|e_1\rangle &= \int_{-1}^1 dx \\
&= 2 \\
|e'_1\rangle &= \frac{1}{\sqrt{2}} \\
|e_2\rangle &= x \\
\langle e'_1|e_2\rangle &= \int_{-1}^1 \frac{1}{\sqrt{2}} x dx \\
&= \frac{1}{\sqrt{2}} \left[ \frac{1}{2} x^2 \right]_{-1}^1 \\
&= 0 \\
\langle e_2|e_2\rangle &= \int_{-1}^1 x^2 dx \\
&= \left[ \frac{1}{3} x^3 \right]_{-1}^1 \\
&= \frac{2}{3} \\
|e'_2\rangle &= \sqrt{\frac{3}{2}} x \\
|e_3\rangle &= x^2 \\
\langle e'_1|e_3\rangle &= \int_{-1}^1 \frac{1}{\sqrt{2}} x^2 dx \\
&= \frac{1}{\sqrt{2}} \left[ \frac{1}{3} x^3 \right]_{-1}^1 \\
&= \frac{\sqrt{2}}{3} \\
\langle e'_2|e_3\rangle &= \int_{-1}^1 \sqrt{\frac{3}{2}} x^3 dx \\
&= 0 \\
|e'_3\rangle &= x^2 - \frac{\sqrt{2}}{3} |e'_1\rangle \\
&= x^2 - \frac{1}{3}
\end{aligned}$$

$$\begin{aligned}
\langle e'_3 | e'_3 \rangle &= \int_{-1}^1 \left( x^2 - \frac{1}{3} \right)^2 dx \\
&= \int_{-1}^1 \left( x^4 - \frac{2}{3}x^2 + \frac{1}{9} \right) dx \\
&= \left[ \frac{1}{5}x^5 - \frac{2}{9}x^3 + \frac{1}{9}x \right]_{-1}^1 \\
&= \frac{2}{5} - \frac{4}{9} + \frac{2}{9} \\
&= \frac{18}{45} - \frac{20}{45} + \frac{10}{45} \\
&= \frac{8}{45}
\end{aligned}$$

$$\begin{aligned}
|e''_3\rangle &= \sqrt{\frac{45}{8}} \left( x^2 - \frac{1}{3} \right) \\
&= \sqrt{\frac{5}{2}} \left( \frac{3}{2}x^2 - \frac{1}{2} \right)
\end{aligned}$$

$$|e_4\rangle = x^3$$

$$\begin{aligned}
\langle e'_1 | e_4 \rangle &= \int_{-1}^1 \frac{1}{\sqrt{2}} x^3 dx \\
&= 0
\end{aligned}$$

$$\begin{aligned}
\langle e'_2 | e_4 \rangle &= \int_{-1}^1 \sqrt{\frac{3}{2}} x^4 dx \\
&= \sqrt{\frac{3}{2}} \left[ \frac{1}{5} x^5 \right] \\
&= \frac{\sqrt{6}}{5}
\end{aligned}$$

$$\begin{aligned}
\langle e''_3 | e_4 \rangle &= \int_{-1}^1 \sqrt{\frac{5}{2}} \left( \frac{3}{2}x^2 - \frac{1}{2} \right) x^3 dx \\
&= \sqrt{\frac{5}{2}} \int_{-1}^1 \left( \frac{3}{2}x^5 - \frac{1}{2}x^3 \right) dx \\
&= 0
\end{aligned}$$

$$\begin{aligned}
|e'_4\rangle &= |e_4\rangle - \frac{\sqrt{6}}{5} |e'_2\rangle \\
&= x^3 - \frac{3}{5}x
\end{aligned}$$

$$\begin{aligned}
\langle e'_4 | e'_4 \rangle &= \int_{-1}^1 \left( x^3 - \frac{3}{5}x \right)^2 dx \\
&= \int_{-1}^1 \left( x^6 - \frac{6}{5}x^4 + \frac{9}{25}x^2 \right) dx \\
&= \left[ \frac{1}{7}x^7 - \frac{6}{25}x^5 + \frac{3}{25}x^3 \right]_{-1}^1 \\
&= \frac{2}{7} - \frac{12}{25} + \frac{6}{25} \\
&= \frac{50}{175} - \frac{84}{175} + \frac{42}{175} \\
&= \frac{8}{175} \\
|e''_4\rangle &= \sqrt{\frac{175}{8}} \left( x^3 - \frac{3}{5}x \right) \\
&= \sqrt{\frac{7}{2}} \left( \frac{5}{2}x^3 - \frac{3}{2}x \right)
\end{aligned}$$

### 3.32

(a)

$$\begin{aligned}
\langle \hat{Q} \rangle &= \langle \Psi | \hat{Q} | \Psi \rangle \\
&= \langle \Psi | \hat{Q}^\dagger | \Psi \rangle^* \\
&= - \langle \Psi | \hat{Q} | \Psi \rangle^* \\
&= - \langle \hat{Q} \rangle^*
\end{aligned}$$

(b)

$$\begin{aligned}
\hat{Q} |\psi\rangle &= \lambda |\psi\rangle \\
\langle \psi | \hat{Q} | \psi \rangle &= \langle \psi | \lambda | \psi \rangle \\
&= \lambda \langle \psi | \psi \rangle \\
\langle \psi | \hat{Q} | \psi \rangle^* &= \lambda^* \langle \psi | \psi \rangle^* \\
\langle \psi | \hat{Q}^\dagger | \psi \rangle &= \lambda^* \langle \psi | \psi \rangle \\
- \langle \psi | \hat{Q} | \psi \rangle &= \lambda^* \langle \psi | \psi \rangle \\
\langle \psi | \hat{Q} | \psi \rangle &= -\lambda^* \langle \psi | \psi \rangle \\
\hat{Q} |\psi\rangle &= -\lambda^* |\psi\rangle
\end{aligned}$$



(c)

$$\begin{aligned}
\hat{Q}|f\rangle &= \lambda_1|f\rangle \\
\hat{Q}|g\rangle &= \lambda_2|g\rangle \\
\langle g|\hat{Q}|f\rangle &= \langle g|\lambda_1|f\rangle \\
&= \lambda_1\langle g|f\rangle \\
\langle f|\hat{Q}^\dagger|g\rangle &= \lambda_1^*\langle f|g\rangle \\
-\langle f|\hat{Q}|g\rangle &= \lambda_1^*\langle f|g\rangle \\
\langle f|\lambda_2|g\rangle &= -\lambda_1^*\langle f|g\rangle \\
\lambda_2\langle f|g\rangle &= -\lambda_1^*\langle f|g\rangle \\
\lambda_2 &= -\lambda_1^* \\
\lambda_2 &= \lambda_1 \\
\langle f|g\rangle &= 0
\end{aligned}$$

(d)

$$\begin{aligned}
[\hat{A}, \hat{B}] &= \hat{A}\hat{B} - \hat{B}\hat{A} \\
[\hat{A}, \hat{B}]^\dagger &= (\hat{A}\hat{B} - \hat{B}\hat{A})^\dagger \\
&= \hat{A}^\dagger\hat{B}^\dagger - \hat{B}^\dagger\hat{A}^\dagger \\
&= -(\hat{A}\hat{B} - \hat{B}\hat{A}) \\
&= -[\hat{A}, \hat{B}] \\
[\hat{A}, \hat{B}]^\dagger &= \hat{A}^\dagger\hat{B}^\dagger - \hat{B}^\dagger\hat{A}^\dagger \\
&= -(\hat{A}\hat{B} - \hat{B}\hat{A}) \\
&= -[\hat{A}, \hat{B}]
\end{aligned}$$

(e)

$$\begin{aligned}
\hat{Q}|q_n\rangle &= \lambda_n|q_n\rangle \\
&= (x + iy)|q_n\rangle \\
&= x|q_n\rangle + iy|q_n\rangle \\
&= \hat{X}|q_n\rangle + \hat{Y}|q_n\rangle \\
&= (\hat{X} + \hat{Y})|q_n\rangle
\end{aligned}$$

### 3.33

(a)  $\psi_1$

(b)  $b_1$  and  $b_2$  with  $P(b_1) = \frac{9}{25}$  and  $P(b_2) = \frac{16}{25}$ .

(c)

$$\begin{aligned}
\phi_1 &= \frac{3}{5}\psi_1 + \frac{4}{5}\psi_2 \\
\phi_2 &= \frac{4}{5}\psi_1 - \frac{3}{5}\psi_2 \\
P(a_1) &= P(b_1) \left(\frac{3}{5}\right)^2 + P(b_2) \left(\frac{4}{5}\right)^2 \\
&= \left(\frac{9}{25}\right)^2 + \left(\frac{16}{25}\right)^2 \\
&= \frac{337}{625} \\
&\approx 53.9\%
\end{aligned}$$

### 3.34

(a)

$$\begin{aligned}
\Phi_n(p, t) &= \frac{1}{\sqrt{2\pi\hbar}} \int_0^a e^{-ipx/\hbar} \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right) e^{-iE_n t/\hbar} dx \\
&= \frac{1}{\sqrt{\pi\hbar a}} e^{-iE_n t/\hbar} \int_0^a e^{-ipx/\hbar} \sin\left(\frac{n\pi}{a}x\right) dx \\
&= \sqrt{\frac{a\pi}{\hbar}} \frac{ne^{-iE_n t/\hbar}}{(n\pi)^2 - (ap/\hbar)^2} [1 - (-1)^n e^{-ipa/\hbar}]
\end{aligned}$$

(b)

$$|\Phi_n(p, t)|^2 = \frac{a\pi}{\hbar} \frac{4n^2}{[(n\pi)^2 - (ap/\hbar)^2]^2} \begin{cases} \cos^2\left(\frac{a}{2\hbar}p\right) & n \text{ odd} \\ \sin^2\left(\frac{a}{2\hbar}p\right) & n \text{ even} \end{cases}$$

### 3.35

$$\begin{aligned}
\Psi(x, 0) &= \begin{cases} \frac{1}{\sqrt{2n\lambda}} e^{i2\pi x/\lambda} & -n\lambda < x < n\lambda \\ 0 & \text{otherwise} \end{cases} \\
\Phi(p, 0) &= \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi\hbar}} e^{-ipx/\hbar} \Psi(x, 0) dx \\
&= \frac{1}{2\sqrt{\pi\hbar n\lambda}} \int_{-n\lambda}^{n\lambda} e^{i(2\pi/\lambda - p/\hbar)x} dx \\
&= \frac{1}{2\sqrt{\pi\hbar n\lambda}} \frac{1}{i(2\pi/\lambda - p/\hbar)} \left[ e^{i(2\pi/\lambda - p/\hbar)x} \right]_{-n\lambda}^{n\lambda} \\
&= \frac{1}{2\sqrt{\pi\hbar n\lambda}} \frac{1}{i(2\pi/\lambda - p/\hbar)} e^{i(2\pi/\lambda - p/\hbar)n\lambda} - e^{-i(2\pi/\lambda - p/\hbar)n\lambda} \\
&= \frac{1}{\sqrt{\pi\hbar n\lambda}} \frac{1}{2\pi/\lambda - p/\hbar} \sin \left[ \left( \frac{2\pi}{\lambda} - \frac{p}{\hbar} \right) n\lambda \right] \\
&= \sqrt{\frac{\lambda\hbar}{n\pi}} \frac{1}{\lambda p - 2\pi\hbar} \sin \left( \frac{n\lambda}{\hbar} p \right) \\
w_x &= 2n\lambda \\
w_p &= \frac{2\pi\hbar}{n\lambda}
\end{aligned}$$

As  $n \rightarrow \infty$ ,  $w_x \rightarrow \infty$  and  $w_p \rightarrow 0$ .

$$\begin{aligned}
w_x w_p &= 2n\lambda \frac{2\pi\hbar}{n\lambda} \\
&= 4\pi\hbar \\
&\geq \frac{\hbar}{2}
\end{aligned}$$

### 3.36

(a)

$$\begin{aligned}
1 &= \int_{-\infty}^{\infty} \left( \frac{A}{x^2 + a^2} \right)^2 dx \\
&= |A|^2 \int_{-\infty}^{\infty} \frac{1}{(x^2 + a^2)^2} dx \\
&= \frac{\pi |A|^2}{2a^3} \\
A &= a \sqrt{\frac{2a}{\pi}}
\end{aligned}$$

(b)

$$\begin{aligned}
\langle x \rangle &= \langle \Psi | x | \Psi \rangle \\
&= \int_{-\infty}^{\infty} \Psi^* x \Psi \, dx \\
&= \frac{a^3}{\pi} \int_{-\infty}^{\infty} \frac{2x}{(x^2 + a^2)^2} \, dx \\
u &= x^2 + a^2 \\
du &= 2x \, dx \\
\langle x \rangle &= \frac{a^3}{\pi} \int_{-\infty}^{\infty} \frac{1}{u^2} \\
&= \frac{a^3}{\pi} \left[ -\frac{1}{u} \right]_{-\infty}^{\infty} \\
&= 0 \\
\langle x^2 \rangle &= \langle \Psi | x^2 | \Psi \rangle \\
&= \int_{-\infty}^{\infty} \Psi^* x^2 \Psi \, dx \\
&= \frac{2a^3}{\pi} \int_{-\infty}^{\infty} \frac{x^2}{(x^2 + a^2)^2} \, dx \\
&= a^2 \\
\sigma_x &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\
&= a
\end{aligned}$$

(c)

$$\begin{aligned}
\Phi(x, 0) &= \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi\hbar}} e^{-ipx/\hbar} a \sqrt{\frac{2a}{\pi}} \frac{1}{x^2 + a^2} \, dx \\
&= \frac{a\sqrt{a}}{\pi\sqrt{\hbar}} \int_{-\infty}^{\infty} e^{-ipx/\hbar} \frac{1}{x^2 + a^2} \, dx \\
&= \sqrt{\frac{a}{\hbar}} e^{-|p|a/\hbar} \\
\frac{a}{\hbar} \int_{-\infty}^{\infty} e^{-2|p|a/\hbar} \, dp &= 1
\end{aligned}$$

(d)

$$\begin{aligned}
\langle p \rangle &= \langle \Phi | p | \Phi \rangle \\
&= \int_{-\infty}^{\infty} \Phi^* p \Psi dp \\
&= \frac{a}{\hbar} \int_{-\infty}^{\infty} p e^{-2|p|a/\hbar} dp \\
&= 0 \\
\langle p^2 \rangle &= \frac{a}{\hbar} \int_{-\infty}^{\infty} p^2 e^{-2|p|a/\hbar} dp \\
&= \frac{\hbar^2}{2a^2} \\
\sigma_p &= \sqrt{\langle p^2 \rangle - \langle p \rangle^2} \\
&= \frac{\hbar}{\sqrt{2}a}
\end{aligned}$$

(e)

$$\begin{aligned}
\sigma_x \sigma_p &= \frac{\hbar}{\sqrt{2}} \\
&\geq \frac{\hbar}{2}
\end{aligned}$$

### 3.37

$$\begin{aligned}
[\hat{H}, xp] &= x[\hat{H}, p] + [\hat{H}, x]p \\
&= i\hbar x \frac{dV}{dx} - \frac{i\hbar p^2}{m} \\
\frac{d}{dt} \langle xp \rangle &= \frac{i}{\hbar} \langle [\hat{H}, xp] \rangle \\
&= \frac{i}{\hbar} \left\langle i\hbar x \frac{dV}{dx} - \frac{i\hbar p^2}{m} \right\rangle \\
&= \left\langle \frac{p^2}{m} \right\rangle - \left\langle x \frac{dV}{dx} \right\rangle \\
&= 2 \langle T \rangle - \left\langle x \frac{dV}{dx} \right\rangle
\end{aligned}$$

The left hand side is 0 because expectation values are constant in stationary states.

$$\begin{aligned}
V &= \frac{1}{2}m\omega^2x^2 \\
\frac{dV}{dx} &= m\omega^2x \\
x\frac{dV}{dx} &= m\omega^2x^2 \\
&= 2V \\
2\langle T \rangle &= \left\langle x\frac{dV}{dx} \right\rangle \\
&= \langle 2V \rangle \\
&= 2\langle V \rangle \\
\langle T \rangle &= \langle V \rangle
\end{aligned}$$

### 3.38

$$\begin{aligned}
\Psi(x, 0) &= \frac{1}{\sqrt{2}}(\psi_1 + \psi_2) \\
\Psi(x, t) &= \frac{1}{\sqrt{2}}(\psi_1 e^{-iE_1 t/\hbar} + \psi_2 e^{-iE_2 t/\hbar}) \\
0 &= \int_{-\infty}^{\infty} \Psi(x, 0) \Psi(x, t) dx \\
&= \frac{1}{2} \int_{-\infty}^{\infty} (\psi_1 + \psi_2)(\psi_1 e^{-iE_1 t/\hbar} + \psi_2 e^{-iE_2 t/\hbar}) dx \\
&= \frac{1}{2} \int_{-\infty}^{\infty} (\psi_1^2 e^{-iE_1 t/\hbar} + \psi_1 \psi_2 e^{-iE_2 t/\hbar} + \psi_1 \psi_2 e^{-iE_1 t/\hbar} + \psi_2^2 e^{-iE_2 t/\hbar}) dx \\
&= \frac{1}{2}(e^{-iE_1 t/\hbar} + e^{-iE_2 t/\hbar}) \\
&= e^{-iE_1 t/\hbar} + e^{-iE_2 t/\hbar} \\
e^{-iE_1 t/\hbar} &= -e^{-iE_2 t/\hbar} \\
&= e^{\pi i} e^{-iE_2 t/\hbar} \\
-\frac{iE_1 t}{\hbar} &= i \left( \pi - \frac{E_2 t}{\hbar} \right) \\
-\frac{E_1 t}{\hbar} &= \pi - \frac{E_2 t}{\hbar} \\
t &= \frac{\pi \hbar}{E_2 - E_1} \\
\Delta t &= \frac{t}{\pi} \\
&= \frac{\hbar}{E_2 - E_1} \\
\Delta E &= \frac{1}{2}(E_2 - E_1) \\
\Delta E \Delta t &= \frac{\hbar}{2}
\end{aligned}$$

### 3.41

$$\begin{aligned}
\Psi &= \frac{e^{i\theta_0}}{\sqrt{2}} \psi_0 e^{-i\omega t/2} + \frac{e^{i\theta_1}}{\sqrt{2}} \psi_1 e^{-3i\omega t/2} \\
\langle p \rangle &= \langle \Psi | p | \Psi \rangle \\
&= \frac{1}{2} \langle \psi_0 | p | \psi_0 \rangle + \frac{1}{2} e^{i(\theta_1 - \theta_0)} e^{-i\omega t} \langle \psi_0 | p | \psi_1 \rangle \\
&\quad + \frac{1}{2} e^{i(\theta_0 - \theta_1)} e^{i\omega t} \langle \psi_1 | p | \psi_0 \rangle + \frac{1}{2} \langle \psi_1 | p | \psi_1 \rangle \\
\langle \psi_0 | p | \psi_0 \rangle &= \frac{d}{dt} \langle \psi_0 | x | \psi_0 \rangle \\
&= 0 \\
\langle \psi_1 | p | \psi_1 \rangle &= \frac{d}{dt} \langle \psi_1 | x | \psi_1 \rangle \\
&= 0 \\
\langle p \rangle &= \frac{1}{2} e^{i(\theta_1 - \theta_0 - \omega t)} \langle \psi_0 | p | \psi_1 \rangle + \frac{1}{2} e^{-i(\theta_1 - \theta_0 - \omega t)} \langle \psi_1 | p | \psi_0 \rangle \\
\langle \psi_0 | p | \psi_1 \rangle &= -i \sqrt{\frac{\hbar m \omega}{2}} \\
\langle \psi_1 | p | \psi_0 \rangle &= i \sqrt{\frac{\hbar m \omega}{2}} \\
\langle p \rangle &= -\frac{1}{2} i e^{i(\theta_1 - \theta_0 - \omega t)} \sqrt{\frac{\hbar m \omega}{2}} + \frac{1}{2} i e^{-i(\theta_1 - \theta_0 - \omega t)} \sqrt{\frac{\hbar m \omega}{2}} \\
&= -\frac{1}{2} i \sqrt{\frac{\hbar m \omega}{2}} (e^{i(\theta_1 - \theta_0 - \omega t)} - e^{-i(\theta_1 - \theta_0 - \omega t)}) \\
&= \sqrt{\frac{\hbar m \omega}{2}} \sin(\theta_1 - \theta_0 - \omega t)
\end{aligned}$$

The largest possible value of  $\langle p \rangle$  is  $\sqrt{\hbar m \omega / 2}$ . If it takes on this value at  $t = 0$  then

$$\Psi(x, t) = \frac{1}{\sqrt{2}} e^{-i\omega t/2} (\psi_0 + i\psi_1 e^{-i\omega t}).$$



### 3.43

(a)

$$\begin{aligned}
|z|^2 &= \text{Re}(z)^2 + \text{Im}(z)^2 \\
&= \left(\frac{z+z^*}{2}\right)^2 + \left(\frac{z-z^*}{2i}\right)^2 \\
\sigma_A^2 \sigma_B^2 &\geq \left(\frac{\langle f|g\rangle + \langle g|f\rangle}{2}\right)^2 + \left(\frac{\langle f|g\rangle - \langle g|f\rangle}{2i}\right)^2 \\
&= \left(\frac{\langle \hat{A}\hat{B}\rangle + \langle \hat{B}\hat{A}\rangle - 2\langle A\rangle\langle B\rangle}{2}\right)^2 + \left(\frac{\langle [\hat{A}, \hat{B}]\rangle}{2i}\right)^2 \\
&= \frac{1}{4}[\langle -i[\hat{A}, \hat{B}]\rangle^2 + (\langle \hat{A}\hat{B}\rangle + \langle \hat{B}\hat{A}\rangle - 2\langle A\rangle\langle B\rangle)^2] \\
&= \frac{1}{4}(\langle C\rangle^2 + \langle D\rangle^2)
\end{aligned}$$

(b)

$$\begin{aligned}
B &= A \\
\hat{C} &= -i[\hat{A}, \hat{B}] \\
&= 0 \\
\hat{D} &= 2(\hat{A}^2 - \langle A\rangle^2) \\
\langle D\rangle &= \langle \Psi|D|\Psi\rangle \\
&= \langle \Psi|2(\hat{A}^2 - \langle A\rangle^2)|\Psi\rangle \\
&= 2(\langle \Psi|\hat{A}^2|\Psi\rangle - \langle A\rangle^2) \\
&= 2(\langle A^2\rangle - \langle A\rangle^2) \\
&= 2\sigma_A^2 \\
\sigma_A^2 \sigma_B^2 &= \frac{1}{4}(\langle C\rangle^2 + \langle D\rangle^2) \\
\sigma_A^4 &\geq \sigma_A^4
\end{aligned}$$

### 3.44

(a) The eigenvalues of  $\mathbf{H}$  are  $a-b$ ,  $a+b$ , and  $c$  and the associated eigenvectors

are  $\begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix}$ ,  $\begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix}$ , and  $\begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$ , respectively.

$$|S(t)\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} e^{-ict/\hbar}$$

(b)

$$\begin{aligned}
\begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} &= \frac{1}{2} \left[ \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} - \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix} \right] \\
|S(t)\rangle &= \frac{1}{2} \left[ \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} e^{-i(a+b)t/\hbar} - \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix} e^{-i(a-b)t/\hbar} \right] \\
&= \frac{1}{2} e^{-iat/\hbar} \begin{pmatrix} e^{-ibt/\hbar} + e^{ibt/\hbar} \\ 0 \\ e^{-ibt/\hbar} - e^{ibt/\hbar} \end{pmatrix} \\
&= e^{-iat/\hbar} \begin{pmatrix} \cos(bt/\hbar) \\ 0 \\ -i \sin(bt/\hbar) \end{pmatrix}
\end{aligned}$$

**3.45**

$$\begin{aligned}
\langle n|\hat{x}|S(t)\rangle &= \left\langle n \left| \hat{x} \sum_{n'=0}^{\infty} |n'\rangle \langle n'| \right| S(t) \right\rangle \\
&= \sum_{n'=0}^{\infty} \langle n|\hat{x}|n'\rangle \langle n'|S(t)\rangle \\
&= \sqrt{\frac{\hbar}{2m\omega}} \sum_{n'=0}^{\infty} (\sqrt{n'}\delta_{n,n'-1} + \sqrt{n}\delta_{n',n-1})c_{n'}(t) \\
&= \sqrt{\frac{\hbar}{2m\omega}} [\sqrt{n+1}c_{n+1}(t) + \sqrt{n}c_{n-1}(t)]
\end{aligned}$$

### 3.46

(a)

$$h_1 = \hbar\omega$$

$$h_2 = 2\hbar\omega$$

$$h_3 = 2\hbar\omega$$

$$|h_1\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

$$|h_2\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

$$|h_3\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$a_1 = 2\lambda$$

$$a_2 = \lambda$$

$$a_3 = -\lambda$$

$$|a_1\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$|a_2\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix}$$

$$|a_3\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} -1 \\ 1 \\ 0 \end{pmatrix}$$

$$b_1 = 2\mu$$

$$b_2 = \mu$$

$$b_3 = -\mu$$

$$|b_1\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

$$|b_2\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix}$$

$$|b_3\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ -1 \\ 1 \end{pmatrix}$$

(b)

$$\begin{aligned}
\langle H \rangle &= \langle S(0) | H | S(0) \rangle \\
&= \begin{pmatrix} c_1^* & c_2^* & c_3^* \end{pmatrix} \hbar\omega \begin{pmatrix} 1 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 2 \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \\ c_3 \end{pmatrix} \\
&= \hbar\omega \begin{pmatrix} c_1^* & c_2^* & c_3^* \end{pmatrix} \begin{pmatrix} c_1 \\ 2c_2 \\ 2c_3 \end{pmatrix} \\
&= \hbar\omega (|c_1|^2 + 2|c_2|^2 + 2|c_3|^2) \\
\langle A \rangle &= \lambda(c_1^*c_2 + c_2^*c_1 + 2|c_3|^2) \\
\langle B \rangle &= \mu(2|c_1|^2 + c_2^*c_3 + c_3^*c_2)
\end{aligned}$$

(c)

$$|S(t)\rangle = c_1 |h_1\rangle e^{-i\omega t} + c_2 |h_2\rangle e^{-2i\omega t} + c_3 |h_3\rangle e^{-2i\omega t}$$

You could measure  $H$  as  $\hbar\omega$  with probability  $|c_1|^2$  or  $2\hbar\omega$  with probability  $|c_2|^2 + |c_3|^2$ .

## 4 Quantum Mechanics in Three Dimensions

### 4.1

(a)

$$\begin{aligned}[r_i, r_j] &= 0 \\ [p_i, p_j] &= 0 \\ [x, p_x]f &= x \left( -i\hbar \frac{\partial}{\partial x} \right) f - \left( -i\hbar \frac{\partial}{\partial x} \right) (xf) \\ &= -i\hbar x \frac{\partial f}{\partial x} + i\hbar \left( f + x \frac{\partial f}{\partial x} \right) \\ &= i\hbar f \\ [p_x, x] &= \left( -i\hbar \frac{\partial}{\partial x} \right) (xf) - x \left( -i\hbar \frac{\partial}{\partial x} \right) f \\ &= -i\hbar \left( f + x \frac{\partial f}{\partial x} \right) + i\hbar x \frac{\partial f}{\partial x} \\ &= -i\hbar f \\ [x, p_y] &= x \left( -i\hbar \frac{\partial}{\partial y} \right) f - \left( -i\hbar \frac{\partial}{\partial y} \right) (xf) \\ &= -i\hbar x \frac{\partial f}{\partial y} + i\hbar x \frac{\partial f}{\partial y} \\ &= 0 \\ [r_i, p_j] &= i\hbar \delta_{ij} \\ [p_i, r_j] &= -i\hbar \delta_{ij}\end{aligned}$$

(b)

$$\begin{aligned}
[\hat{H}, x]f &= \left(-\frac{\hbar^2}{2m}\nabla^2 + V\right)(xf) - x\left(-\frac{\hbar^2}{2m}\nabla^2 + V\right)f \\
&= -\frac{\hbar^2}{2m}\left[\frac{\partial}{\partial x}\left(f + x\frac{\partial f}{\partial x}\right) + \frac{\partial}{\partial y}\left(x\frac{\partial f}{\partial y}\right) + \frac{\partial}{\partial z}\left(x\frac{\partial f}{\partial z}\right)\right] \\
&\quad + Vxf + \frac{\hbar^2}{2m}x\nabla^2 f - Vxf \\
&= -\frac{\hbar^2}{2m}\left(2\frac{\partial f}{\partial x} + x\nabla^2 f\right) + \frac{\hbar^2}{2m}x\nabla^2 f \\
&= -\frac{\hbar^2}{m}\frac{\partial f}{\partial x} \\
[\hat{H}, x] &= -\frac{\hbar^2}{m}\frac{\partial}{\partial x} \\
[\hat{H}, \mathbf{r}] &= -\frac{\hbar^2}{m}\nabla \\
\frac{d}{dt}\langle \mathbf{r} \rangle &= \frac{i}{\hbar}\langle [\hat{H}, \mathbf{r}] \rangle \\
&= \frac{1}{m}\langle -i\hbar\nabla \rangle \\
&= \frac{1}{m}\langle \mathbf{p} \rangle \\
[\hat{H}, \hat{p}_x]f &= \left(-\frac{\hbar^2}{2m}\nabla^2 + V\right)\left(-i\hbar\frac{\partial}{\partial x}\right)f - \left(-i\hbar\frac{\partial}{\partial x}\right)\left(-\frac{\hbar^2}{2m}\nabla^2 + V\right)f \\
&= -i\hbar\left(-\frac{\hbar^2}{2m}\nabla^2 + V\right)\frac{\partial f}{\partial x} + i\hbar\left(\frac{\partial}{\partial x}\right)\left(-\frac{\hbar^2}{2m}\nabla^2 f + Vf\right) \\
&= i\frac{\hbar^3}{2m}\nabla^2\frac{\partial f}{\partial x} - i\hbar V\frac{\partial f}{\partial x} - i\frac{\hbar^3}{2m}\nabla^2\frac{\partial f}{\partial x} + i\hbar\left(\frac{\partial V}{\partial x}f + V\frac{\partial f}{\partial x}\right) \\
&= i\hbar\frac{\partial V}{\partial x}f \\
[\hat{H}, \hat{p}_x] &= i\hbar\frac{\partial V}{\partial x} \\
[\hat{H}, \mathbf{p}] &= i\hbar\nabla V \\
\frac{d}{dt}\langle \mathbf{p} \rangle &= \frac{i}{\hbar}\langle [\hat{H}, \mathbf{p}] \rangle \\
&= \langle -\nabla V \rangle
\end{aligned}$$

(c)

$$\begin{aligned}\sigma_{r_i}\sigma_{p_j} &\geq \frac{1}{2i} \langle [r_i, p_j] \rangle \\ &= \frac{1}{2i} i\hbar\delta_{ij} \\ &= \frac{\hbar}{2}\delta_{ij}\end{aligned}$$

## 4.2

(a)

$$\begin{aligned}\psi(\mathbf{r}) &= X(x)Y(y)Z(z) \\ -\frac{\hbar^2}{2m}\nabla^2\psi &= E\psi \\ -\frac{\hbar^2}{2m}\left(\frac{\partial^2 X}{\partial x^2}YZ + X\frac{\partial^2 Y}{\partial y^2}Z + XY\frac{\partial^2 Z}{\partial z^2}\right) &= EXYZ \\ \frac{X''}{X} + \frac{Y''}{Y} + \frac{Z''}{Z} &= -\frac{2m}{\hbar^2}E\end{aligned}$$

The terms on the left hand side are each functions of a different variable, so they must be constant. Starting with the  $X$  term:

$$\begin{aligned}\frac{X''}{X} &= -\alpha \\ X'' + \alpha X &= 0\end{aligned}$$

If  $\alpha < 0$

$$X = A_x e^{\sqrt{-\alpha}x} + B_x e^{-\sqrt{-\alpha}x}$$

If  $\alpha = 0$

$$X = A_x x + B_x$$

If  $\alpha > 0$

$$X = A_x \sin(\sqrt{\alpha}x) + B_x \cos(\sqrt{\alpha}x)$$

Boundary conditions require that  $\alpha > 0$ .  $X(0) = 0$  so  $B_x = 0$ .  $X(a) = 0$  so  $\sqrt{\alpha} = n_x \pi / a \Rightarrow \alpha = n_x^2 \pi^2 / a^2$ ,

$$X = A_x \sin\left(\frac{n_x \pi}{a}x\right).$$

Repeating the above for  $Y$  and  $Z$  finds

$$Y = A_y \sin\left(\frac{n_y \pi}{a} y\right)$$

$$Z = A_z \sin\left(\frac{n_z \pi}{a} z\right)$$

so

$$\psi(\mathbf{r}) = A_x A_y A_z \sin\left(\frac{n_x \pi}{a} x\right) \sin\left(\frac{n_y \pi}{a} y\right) \sin\left(\frac{n_z \pi}{a} z\right).$$

Assuming  $A_x = A_y = A_z = A$  and normalising finds

$$1 = \int_0^a \int_0^a \int_0^a A^6 \sin^2\left(\frac{n_x \pi}{a} x\right) \sin^2\left(\frac{n_y \pi}{a} y\right) \sin^2\left(\frac{n_z \pi}{a} z\right) d^3 \mathbf{r}$$

$$= A^6 \frac{a^3}{8}$$

$$A^6 = \frac{8}{a^3}$$

$$= \left(\frac{2}{a}\right)^3$$

$$A = \sqrt{\frac{2}{a}}$$

so

$$\psi(\mathbf{r}) = \left(\frac{2}{a}\right)^{3/2} \sin\left(\frac{n_x \pi}{a} x\right) \sin\left(\frac{n_y \pi}{a} y\right) \sin\left(\frac{n_z \pi}{a} z\right).$$

Finally

$$-\frac{2m}{\hbar^2} E = -\alpha - \beta - \gamma$$

$$= -\frac{\pi^2 n_x^2}{a^2} - \frac{\pi^2 n_y^2}{a^2} - \frac{\pi^2 n_z^2}{a^2}$$

$$E = \frac{\pi^2 \hbar^2}{2ma^2} (n_x^2 + n_y^2 + n_z^2), \quad n_x, n_y, n_z = 1, 2, 3, \dots$$



(b)

$$E_1 = 3 \frac{\pi^2 \hbar^2}{2ma^2}, n = 1$$

$$E_2 = 6 \frac{\pi^2 \hbar^2}{2ma^2}, n = 3$$

$$E_3 = 9 \frac{\pi^2 \hbar^2}{2ma^2}, n = 3$$

$$E_4 = 11 \frac{\pi^2 \hbar^2}{2ma^2}, n = 3$$

$$E_5 = 12 \frac{\pi^2 \hbar^2}{2ma^2}, n = 1$$

$$E_6 = 14 \frac{\pi^2 \hbar^2}{2ma^2}, n = 6$$

(c)

$$E_{14} = 27 \frac{\pi^2 \hbar^2}{2ma^2}$$

$$d = 4$$

### 4.3

(a)

$$\begin{aligned}
\phi(r, \theta, \phi) &= Ae^{-r/a} \\
\frac{\partial \psi}{\partial r} &= -\frac{A}{a}e^{-r/a} \\
\frac{\partial}{\partial r} \left( r^2 \frac{\partial \psi}{\partial r} \right) &= \frac{\partial}{\partial r} \left( -\frac{A}{a}r^2e^{-r/a} \right) \\
&= -\frac{A}{a} \left( 2re^{-r/a} - \frac{1}{a}r^2e^{-r/a} \right) \\
&= \frac{A}{a^2}r^2 \left( 1 - \frac{2a}{r} \right) e^{-r/a} \\
\frac{\partial \psi}{\partial \theta} &= 0 \\
\frac{\partial^2 \psi}{\partial \phi^2} &= 0 \\
-\frac{\hbar^2}{2m} \frac{A}{a^2} \left( 1 - \frac{2a}{r} \right) e^{-r/a} + V(r)Ae^{-r/a} &= EAe^{-r/a} \\
-\frac{\hbar^2}{2ma^2} \left( 1 - \frac{2a}{r} \right) + V(r) &= E \\
-\frac{\hbar^2}{2ma^2} &= E \\
V(r) &= -\frac{\hbar^2}{mar}
\end{aligned}$$

(b)

$$\begin{aligned}
\phi(r, \theta, \phi) &= Ae^{-r^2/a^2} \\
\frac{\partial \psi}{\partial r} &= -\frac{2}{a^2}rAe^{-r^2/a^2} \\
\frac{\partial}{\partial r} \left( r^2 \frac{\partial \psi}{\partial r} \right) &= \frac{\partial}{\partial r} \left( -\frac{2}{a^2}r^3Ae^{-r^2/a^2} \right) \\
&= -\frac{2A}{a^2} \left( 3r^2e^{-r^2/a^2} - \frac{2}{a^2}r^4e^{-r^2/a^2} \right) \\
&= \frac{2A}{a^2}r^2 \left( \frac{2}{a^2}r^2 - 3 \right) e^{-r^2/a^2}
\end{aligned}$$

$$\begin{aligned}
-\frac{\hbar^2}{2m} \frac{2A}{a^2} \left( \frac{2}{a^2} r^2 - 3 \right) e^{-r^2/a^2} + V(r) A e^{-r^2/a^2} &= E A e^{-r^2/a^2} \\
-\frac{\hbar^2}{2m} \frac{2}{a^2} \left( \frac{2}{a^2} r^2 - 3 \right) + V(r) &= E \\
\frac{\hbar^2}{2m} \frac{2}{a^2} 3 &= E \\
\frac{3\hbar^2}{ma^2} &= E \\
\frac{3\hbar^2}{ma^2} + \frac{\hbar^2}{2m} \frac{2}{a^2} \left( \frac{2}{a^2} r^2 - 3 \right) &= V(r) \\
\frac{2\hbar^2}{ma^4} r^2 &= V(r)
\end{aligned}$$

#### 4.4

$$\begin{aligned}
Y_0^0 &= \frac{1}{2\sqrt{\pi}} \\
Y_2^1 &= \frac{1}{2} \sqrt{\frac{5}{6\pi}} e^{i\phi} P_2^1(\cos \theta) \\
&= \frac{1}{2} \sqrt{\frac{5}{6\pi}} e^{i\phi} (-1)^1 (1 - \cos^2 \theta)^{1/2} \left( \frac{d}{dx} \right) P_2(\cos \theta) \\
&= -\frac{1}{2} \sqrt{\frac{5}{6\pi}} e^{i\phi} \sin \theta \left( \frac{d\theta}{dx} \frac{d}{d\theta} \right) \frac{1}{2} (3 \cos^2 \theta - 1) \\
&= -\frac{1}{4} \sqrt{\frac{5}{6\pi}} e^{i\phi} \sin \theta \left( -\frac{1}{\sin \theta} \right) (-6 \cos \theta \sin \theta) \\
&= -\frac{1}{2} \sqrt{\frac{15}{2\pi}} e^{i\phi} \sin \theta \cos \theta \\
\int |Y_0^0|^2 d\Omega &= \frac{1}{4\pi} \int_0^\pi \int_0^{2\pi} \sin \theta d\theta d\phi \\
&= 1 \\
\int |Y_2^1|^2 d\Omega &= \frac{15}{8\pi} \int_0^\pi \int_0^{2\pi} \sin^3 \theta \cos^2 \theta d\theta d\phi \\
&= 1 \\
\int (Y_0^0)^* Y_2^1 d\Omega &= -\frac{1}{2\sqrt{\pi}} \frac{1}{2} \sqrt{\frac{15}{2\pi}} \int_0^\pi \int_0^{2\pi} e^{i\phi} \sin^2 \theta \cos \theta d\theta d\phi \\
&= 0
\end{aligned}$$

#### 4.5

$$\begin{aligned}
\Theta(\theta) &= A \ln \left( \tan \frac{\theta}{2} \right) \\
0 &= \sin \theta \frac{d}{d\theta} \left( \sin \theta \frac{d\Theta}{d\theta} \right) \\
&= \frac{1}{2} A \sin \theta \frac{d}{d\theta} \left( \sin \theta \csc \frac{\theta}{2} \sec \frac{\theta}{2} \right) \\
&= \frac{1}{2} A \sin \theta \left( \cos \theta \csc \frac{\theta}{2} \sec \frac{\theta}{2} - \frac{1}{2} \csc^2 \frac{\theta}{2} \sin \theta + \frac{1}{2} \sec^2 \frac{\theta}{2} \sin \theta \right) \\
&= 0
\end{aligned}$$

It's not a valid physical solution because it blows up at  $\theta = 0$  and  $\theta = \pi$ .

#### 4.6

$$\begin{aligned}
Y_\ell^{-m} &= \sqrt{\frac{(2\ell+1)(\ell+m)!}{4\pi(\ell-m)!}} e^{-im\phi} P_\ell^{-m}(\cos \theta) \\
&= \sqrt{\frac{(2\ell+1)(\ell+m)!}{4\pi(\ell-m)!}} e^{-im\phi} (-1)^m \frac{(\ell-m)!}{(\ell+m)!} P_\ell^m(\cos \theta) \\
&= (-1)^m \sqrt{\frac{(2\ell+1)(\ell-m)!}{4\pi(\ell+m)!}} e^{-im\phi} P_\ell^m(\cos \theta) \\
&= (-1)^m (Y_\ell^m)^*
\end{aligned}$$

#### 4.7

$$\begin{aligned}
Y_3^2(\theta, \phi) &= \sqrt{\frac{7}{480\pi}} e^{2i\phi} P_3^2(\cos \theta) \\
&= \frac{1}{4} \sqrt{\frac{105}{2\pi}} e^{2i\phi} \sin^2 \theta \cos \theta \\
Y_\ell^\ell(\theta, \phi) &= \sqrt{\frac{2\ell+1}{4\pi(2\ell)!}} e^{i\ell\phi} P_\ell^\ell(\cos \theta) \\
P_\ell^\ell(x) &= (-1)^\ell (1-x^2)^{\ell/2} \left( \frac{d}{dx} \right)^\ell P_\ell(x) \\
&= (-1)^\ell (1-x^2)^{\ell/2} \left( \frac{d}{dx} \right)^\ell \left[ \frac{1}{2^\ell \ell!} \left( \frac{d}{dx} \right)^\ell (x^2-1)^\ell \right] \\
&= (-1)^\ell \frac{1}{2^\ell \ell!} (1-x^2)^{\ell/2} \left( \frac{d}{dx} \right)^{2\ell} (x^2-1)^\ell \\
&= (-1)^\ell \frac{(2\ell)!}{2^\ell \ell!} (1-x^2)^{\ell/2} \\
Y_\ell^\ell(\theta, \phi) &= \sqrt{\frac{2\ell+1}{4\pi(2\ell)!}} e^{i\ell\phi} (-1)^\ell \frac{(2\ell)!}{2^\ell \ell!} (1-\cos^2 \theta)^{\ell/2} \\
&= \frac{1}{\ell!} \sqrt{\frac{(2\ell+1)!}{4\pi}} \left( -\frac{1}{2} e^{i\phi} \sin \theta \right)^\ell
\end{aligned}$$

## 4.9

(a)

$$\begin{aligned}
 n_1(x) &= -(-x) \left( \frac{1}{x} \frac{d}{dx} \right) \frac{\cos x}{x} \\
 &= -\frac{\sin x}{x} - \frac{\cos x}{x^2} \\
 n_2(x) &= -(-x)^2 \left( \frac{1}{x} \frac{d}{dx} \right)^2 \frac{\cos x}{x} \\
 &= -x^2 \left( \frac{1}{x} \frac{d}{dx} \right) \left[ \frac{1}{x} \frac{d}{dx} \left( \frac{\cos x}{x} \right) \right] \\
 &= -x^2 \left( \frac{1}{x} \frac{d}{dx} \right) \left[ \frac{1}{x} \left( -\frac{\sin x}{x} - \frac{\cos x}{x^2} \right) \right] \\
 &= x \frac{d}{dx} \left( \frac{\sin x}{x^2} + \frac{\cos x}{x^3} \right) \\
 &= x \left( \frac{\cos x}{x^2} - \frac{2 \sin x}{x^3} - \frac{\sin x}{x^3} - \frac{3 \cos x}{x^4} \right) \\
 &= \left( \frac{1}{x} - \frac{3}{x^3} \right) \cos x - \frac{3 \sin x}{x^2}
 \end{aligned}$$

(b)

$$\begin{aligned}
 n_1(x) &\approx -\frac{x}{x} - \frac{1}{x^2} \\
 &= -1 - \frac{1}{x^2} \\
 n_2(x) &\approx \left( \frac{1}{x} - \frac{3}{x^3} \right) - \frac{3}{x^2} x \\
 &= -\frac{2}{x} - \frac{3}{x^3}
 \end{aligned}$$

#### 4.10

(a)

$$\begin{aligned}
-\frac{\hbar^2}{2m} \frac{d^2 u}{dr^2} + \frac{\hbar^2}{2m} \frac{2}{r^2} u &= E u \\
\frac{d^2 u}{dr^2} - \frac{2}{r^2} u &= -\frac{2mE}{\hbar^2} u \\
\frac{d^2 u}{dr^2} &= \left( \frac{2}{r^2} - k^2 \right) u \\
j_1(kr) &= \frac{\sin(kr)}{(kr)^2} - \frac{\cos(kr)}{kr} \\
u &= A r j_1(kr) \\
&= A r \left[ \frac{\sin(kr)}{(kr)^2} - \frac{\cos(kr)}{kr} \right]
\end{aligned}$$

$$\begin{aligned}
\frac{du}{dr} &= A \left[ \frac{\sin(kr)}{(kr)^2} - \frac{\cos(kr)}{kr} \right] \\
&\quad + A r \left[ \frac{k \cos(kr)}{(kr)^2} - \frac{2 \sin(kr)}{k^2 r^3} + \frac{k \sin(kr)}{kr} + \frac{\cos(kr)}{kr^2} \right] \\
&= \frac{A}{kr} \left[ \frac{\sin(kr)}{kr} - \cos(kr) \right] \\
&\quad + \frac{A}{k} \left[ \frac{2 \cos(kr)}{r} - \frac{2 \sin(kr)}{kr^2} + k \sin(kr) \right] \\
\frac{d^2 u}{dr^2} &= \frac{A(-2 + k^2 r^2)(kr \cos(kr) - \sin(kr))}{k^2 r^3} \\
&= \frac{2 - k^2 r^2}{r^2} A r \left[ \frac{\sin(kr)}{(kr)^2} - \frac{\cos(kr)}{kr} \right] \\
&= \left( \frac{2}{r^2} - k^2 \right) u
\end{aligned}$$

(b)

$$\begin{aligned}
0 &= \frac{\sin x}{x^2} - \frac{\cos x}{x} \\
&= \frac{\sin x}{x} - \cos x \\
\cos x &= \frac{\sin x}{x} \\
x &= \tan x
\end{aligned}$$

$$x \approx \left(N + \frac{1}{2}\right) \pi, n \in \mathbb{Z}$$

$$ka = x$$

$$\frac{\sqrt{2mE_{N1}}}{\hbar}a \approx \left(N + \frac{1}{2}\right) \pi$$

$$E_{N1} = \frac{\hbar^2 \pi^2}{2ma^2} \left(N + \frac{1}{2}\right)^2, N \in \mathbb{Z}$$

#### 4.12

$$R_{n\ell}(r) = \frac{1}{r} \left(\frac{r}{an}\right)^{\ell+1} e^{-r/an} v(r/an)$$

$$R_{30} = e^{-r/3a} \frac{c_0}{3a} \left[ 1 - 2 \left(\frac{r}{3a}\right) - \frac{1}{3} \left(\frac{r}{3a}\right)^2 \right]$$

$$R_{31} = \frac{1}{r} \left(\frac{r}{3a}\right)^2 c_0 \left[ 1 - \frac{1}{2} \left(\frac{r}{3a}\right) \right]$$

$$R_{32} = \frac{1}{r} \left(\frac{r}{3a}\right)^3 c_0$$



### 4.13

(a)

$$\begin{aligned}
 R_{20}(r) &= \frac{c_0}{2a} \left(1 - \frac{r}{2a}\right) e^{-r/2a} \\
 1 &= \int_0^\infty |R_{20}|^2 r^2 dr \\
 &= \frac{c_0^2}{4a^2} \int_0^\infty \left(1 - \frac{r}{2a}\right)^2 e^{-r/a} r^2 dr \\
 &= \frac{ac_0^2}{2} \\
 c_0 &= \sqrt{\frac{2}{a}} \\
 R_{20}(r) &= \frac{1}{\sqrt{2}} a^{-3/2} \left(1 - \frac{r}{2a}\right) e^{-r/2a} \\
 \psi_{200} &= R_{20} Y_0^0 \\
 &= \frac{1}{\sqrt{2\pi a}} \frac{1}{2a} \left(1 - \frac{r}{2a}\right) e^{-r/2a}
 \end{aligned}$$

(b)

$$\begin{aligned}
 R_{21}(r) &= \frac{c_0}{4a^2} r e^{-r/2a} \\
 1 &= \frac{c_0^2}{16a^4} \int_0^\infty r^4 e^{-r/a} dr \\
 &= \frac{24}{16} a c_0^2 \\
 c_0 &= \sqrt{\frac{2}{3a}} \\
 R_{21}(r) &= \sqrt{\frac{2}{3a}} \frac{1}{4a^2} r e^{-r/2a} \\
 \psi_{21\pm 1} &= R_{21} Y_1^1 \\
 &= \mp \frac{1}{\sqrt{\pi a}} \frac{1}{8a^2} r e^{-r/2a} \sin \theta e^{\pm i\phi} \\
 \psi_{210} &= \frac{1}{\sqrt{2\pi a}} \frac{1}{4a^2} r e^{-r/2a} \cos \theta
 \end{aligned}$$

#### 4.14

(a)

$$L_q = \frac{e^x}{q!} \left( \frac{d}{dx} \right)^q (e^{-x} x^q)$$

$$L_0 = 1$$

$$\begin{aligned} L_1 &= e^x \frac{d}{dx} (x e^{-x}) \\ &= e^x (e^{-x} - x e^{-x}) \\ &= 1 - x \end{aligned}$$

$$\begin{aligned} L_2 &= \frac{1}{2} e^x \left( \frac{d}{dx} \right)^2 (x^2 e^{-x}) \\ &= \frac{1}{2} e^x \frac{d}{dx} (2x e^{-x} - x^2 e^{-x}) \\ &= \frac{1}{2} e^x (2e^{-x} - 2x e^{-x} - 2x e^{-x} + x^2 e^{-x}) \\ &= 1 - 2x + \frac{1}{2} x^2 \end{aligned}$$

$$\begin{aligned} L_3 &= \frac{1}{6} e^x \left( \frac{d}{dx} \right)^3 (x^3 e^{-x}) \\ &= \frac{1}{6} e^x \left( \frac{d}{dx} \right)^2 (3x^2 e^{-x} - x^3 e^{-x}) \\ &= \frac{1}{6} e^x \frac{d}{dx} (6x e^{-x} - 6x^2 e^{-x} + x^3 e^{-x}) \\ &= \frac{1}{6} e^x (6e^{-x} - 18x e^{-x} + 9x^2 e^{-x} - x^3 e^{-x}) \\ &= 1 - 3x + \frac{3}{2} x^2 - \frac{1}{6} x^3 \end{aligned}$$

#### 4.15

(a)

$$\begin{aligned}
 \psi_{100} &= \frac{1}{\sqrt{\pi a^3}} e^{-r/a} \\
 \langle r \rangle &= \int \psi_{100}^* r \psi_{100} \\
 &= \frac{1}{\pi a^3} \int_0^\infty \int_0^\pi \int_0^{2\pi} r^3 e^{-2r/a} \sin \theta \, dr \, d\theta \, d\phi \\
 &= \frac{4}{a^3} \int_0^\infty r^3 e^{-2r/a} \, dr \\
 &= \frac{3}{2} a \\
 \langle r^2 \rangle &= \frac{4}{a^3} \int_0^\infty r^4 e^{-2r/a} \, dr \\
 &= 3a^2
 \end{aligned}$$

(b)

$$\begin{aligned}
 \langle x \rangle &= 0 \\
 \langle x^2 \rangle &= a^2
 \end{aligned}$$

(c)

$$\begin{aligned}
 \psi_{211} &= R_{21} Y_1^1 \\
 &= -\frac{1}{\sqrt{\pi a}} \frac{1}{8a^2} r e^{-r/2a} \sin \theta e^{i\phi} \\
 \langle x^2 \rangle &= \int \psi_{211}^* x^2 \psi_{211} \\
 &= \frac{1}{64\pi a^5} \int_0^\infty \int_0^\pi \int_0^{2\pi} r^6 e^{-r/a} \sin^5 \theta \cos^2 \phi \, dr \, d\theta \, d\phi \\
 &= 12a^2
 \end{aligned}$$

#### 4.16

$$\begin{aligned}
 P(r) &= \int |\psi_{100}|^2 r^2 \sin \theta \, d\theta \, d\phi \\
 &= \frac{4}{a^3} r^2 e^{-2r/a} \\
 P'(r) &= \frac{4}{a^3} \left( 2r e^{-2r/a} - \frac{2}{a} r^2 e^{-2r/a} \right) \\
 &= \frac{8}{a^3} r \left( 1 - \frac{r}{a} \right) e^{-2r/a} \\
 r_{\max} &= a
 \end{aligned}$$

#### 4.18

(a)

$$\begin{aligned}
\Psi(\mathbf{r}, t) &= \frac{1}{\sqrt{2}}(\psi_{211} + \psi_{21-1})e^{-iE_2t/\hbar} \\
&= \frac{1}{\sqrt{2}}R_{21}(Y_1^1 + Y_1^{-1})e^{-iE_1t/4\hbar} \\
&= \frac{1}{4a^2} \frac{1}{\sqrt{3a}} re^{-r/2a} \left( -\sqrt{\frac{3}{8\pi}} \sin\theta e^{i\phi} + \sqrt{\frac{3}{8\pi}} \sin\theta e^{-i\phi} \right) e^{-iE_1t/4\hbar} \\
&= -i \frac{1}{4a^2} \frac{1}{\sqrt{2\pi a}} re^{-r/2a} \sin\theta \sin\phi e^{-iE_1t/4\hbar}
\end{aligned}$$

(b)

$$\begin{aligned}
\langle V \rangle &= \int \Psi^* \left( -\frac{e^2}{4\pi\epsilon_0} \frac{1}{r} \right) \Psi \\
&= -\frac{e^2}{4\pi\epsilon_0} \frac{1}{16a^4} \frac{1}{2\pi a} \int_0^\infty \int_0^\pi \int_0^{2\pi} r^2 e^{-r/a} \sin^2\theta \sin^2\phi \frac{1}{r} r^2 \sin\theta \, dr \, d\theta \, d\phi \\
&= -\frac{e^2}{128\pi^2\epsilon_0 a^5} \int_0^\infty \int_0^\pi \int_0^{2\pi} r^3 e^{-r/a} \sin^3\theta \sin^2\phi \, dr \, d\theta \, d\phi \\
&= -\frac{e^2}{4\pi\epsilon_0} \frac{1}{4a} \\
&= -6.8 \text{ eV}
\end{aligned}$$

#### 4.19

$$\begin{aligned}
 E_1(Z) &= - \left[ \frac{m_e}{2\hbar^2} \left( \frac{Ze^2}{4\pi\epsilon_0} \right)^2 \right] \\
 &= Z^2 E_1(1) \\
 &= (-13.6 \text{ eV}) Z^2 \\
 E_n(Z) &= \frac{E_1(Z)}{n^2} \\
 &= \frac{(-13.6 \text{ eV}) Z^2}{n^2} \\
 &= Z^2 E_n(1) \\
 a(Z) &= \frac{a}{Z} \\
 \mathcal{R}(Z) &= Z^2 R(1) \\
 \frac{1}{\lambda} &= Z^2 R \\
 \lambda &= \frac{1}{Z^2 R} \\
 \lambda_{Z=2} &= \frac{1}{4R} \\
 &\approx 2.28 \times 10^{-8} \text{ m} \\
 &= 22.8 \text{ nm}
 \end{aligned}$$

#### 4.20

(a)

$$V = -\frac{GMm}{r}$$

(b)

$$a_g = \frac{\hbar^2}{GMm^2}$$

(c)

$$\begin{aligned}
E_c &= \frac{1}{2}mv^2 - \frac{GMm}{r_0} \\
\frac{GMm}{r_0^2} &= \frac{mv^2}{r_0} \\
\frac{1}{2}mv^2 &= \frac{GMm}{2r_0} \\
E_c &= -\frac{GMm}{2r_0} \\
E_n &= -\left[\frac{m}{2\hbar^2}(GMm)^2\right] \frac{1}{n^2} \\
&= -\frac{G^2M^2m^3}{2\hbar^2} \frac{1}{n^2} \\
-\frac{GMm}{2r_0} &= -\frac{G^2M^2m^3}{2\hbar^2} \frac{1}{n^2} \\
\frac{\hbar^2}{GMm^2r_0} &= \frac{1}{n^2} \\
\frac{a_g}{r_0} &= \frac{1}{n^2} \\
n &= \sqrt{\frac{r_0}{a_g}}
\end{aligned}$$

#### 4.21

$$\begin{aligned}
\langle f_\ell^m | L_+ L_- f_\ell^m \rangle &= \langle f_\ell^m | (L^2 - L_z^2 + \hbar L_z) f_\ell^m \rangle \\
&= \langle f_\ell^m | L^2 f_\ell^m \rangle - \langle f_\ell^m | L_z^2 f_\ell^m \rangle + \langle f_\ell^m | \hbar L_z f_\ell^m \rangle \\
&= \langle f_\ell^m | \hbar^2 \ell(\ell+1) f_\ell^m \rangle - \langle f_\ell^m | \hbar^2 m^2 f_\ell^m \rangle + \langle f_\ell^m | \hbar^2 m f_\ell^m \rangle \\
&= \hbar^2 \ell(\ell+1) \langle f_\ell^m | f_\ell^m \rangle - \hbar^2 m^2 \langle f_\ell^m | f_\ell^m \rangle + \hbar^2 m \langle f_\ell^m | f_\ell^m \rangle \\
&= \hbar^2 \ell(\ell+1) - \hbar^2 m^2 + \hbar^2 m \\
&= \hbar^2 [\ell(\ell+1) - m(m-1)] \\
\langle f_\ell^m | L_+ L_- f_\ell^m \rangle &= \langle L_- f_\ell^m | L_- f_\ell^m \rangle \\
&= \langle B_\ell^m f_\ell^{m-1} | B_\ell^m f_\ell^{m-1} \rangle \\
&= |B_\ell^m|^2 \\
|B_\ell^m|^2 &= \hbar^2 [\ell(\ell+1) - m(m-1)] \\
B_\ell^m &= \hbar \sqrt{\ell(\ell+1) - m(m-1)}
\end{aligned}$$

The same argument for  $\langle f_\ell^m | L_- L_+ f_\ell^m \rangle$  finds  $A_\ell^m$ .

## 4.22

(a)

$$\begin{aligned}
[r_i, p_j] &= i\hbar\delta_{ij} \\
[r_i, r_j] &= 0 \\
[p_i, p_j] &= 0 \\
[L_z, x] &= [xp_y - yp_x, x] \\
&= [xp_y, x] - [yp_x, x] \\
&= x[p_y, x] + [x, x]p_y - y[p_x, x] - [y, x]p_x \\
&= i\hbar y \\
[L_z, y] &= [xp_y - yp_x, y] \\
&= [xp_y, y] - [yp_x, y] \\
&= x[p_y, y] + [x, y]p_y - y[p_x, y] - [y, y]p_x \\
&= -i\hbar x \\
[L_z, z] &= [xp_y - yp_x, z] \\
&= [xp_y, z] - [yp_x, z] \\
&= x[p_y, z] + [x, z]p_y - y[p_x, z] - [p_x, z]y \\
&= 0 \\
[L_z, p_x] &= [xp_y - yp_x, p_x] \\
&= [xp_y, p_x] - [yp_x, p_x] \\
&= x[p_y, p_x] + [x, p_x]p_y - y[p_x, p_x] - [y, p_x]p_x \\
&= i\hbar p_y \\
[L_z, p_y] &= [xp_y - yp_x, p_y] \\
&= [xp_y, p_y] - [yp_x, p_y] \\
&= x[p_y, p_y] + [x, p_y]p_y - y[p_x, p_y] - [y, p_y]p_x \\
&= -i\hbar p_x \\
[L_z, p_z] &= [xp_y - yp_x, p_z] \\
&= [xp_y, p_z] - [yp_x, p_z] \\
&= x[p_y, p_z] + [x, p_z]p_y - y[p_x, p_z] - [y, p_z]p_x \\
&= 0
\end{aligned}$$

(b)

$$\begin{aligned}
[L_z, L_x] &= [L_z, yp_z - zp_y] \\
&= [L_z, yp_z] - [L_z, zp_y] \\
&= y[L_z, p_z] + [L_z, y]p_z - z[L_z, p_y] - [L_z, z]p_y \\
&= -i\hbar xp_z + i\hbar zp_x \\
&= i\hbar L_y
\end{aligned}$$

(c)

$$\begin{aligned}
[L_z, r^2] &= [L_z, x^2 + y^2 + z^2] \\
&= [L_z, x^2] + [L_z, y^2] + [L_z, z^2] \\
&= x[L_z, x] + [L_z, x]x + y[L_z, y] + [L_z, y]y + z[L_z, z] + [L_z, z]z \\
&= 2i\hbar xy - 2i\hbar xy \\
&= 0 \\
[L_z, p^2] &= [L_z, p_x^2 + p_y^2 + p_z^2] \\
&= [L_z, p_x^2] + [L_z, p_y^2] + [L_z, p_z^2] \\
&= p_x[L_z, p_x] + [L_z, p_x]p_x + p_y[L_z, p_y] + [L_z, p_y]p_y + p_z[L_z, p_z] + [L_z, p_z]p_z \\
&= 2i\hbar p_x p_y - 2i\hbar p_x p_y \\
&= 0
\end{aligned}$$

## 4.24

(a)

$$\begin{aligned}
L_+ L_- f &= \hbar e^{i\phi} \left( \frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right) \left[ -\hbar e^{-i\phi} \left( \frac{\partial}{\partial \theta} - i \cot \theta \frac{\partial}{\partial \phi} \right) \right] f \\
&= -\hbar^2 e^{i\phi} \left( \frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right) \left[ e^{-i\phi} \left( \frac{\partial f}{\partial \theta} - i \cot \theta \frac{\partial f}{\partial \phi} \right) \right] \\
&= -\hbar^2 e^{i\phi} \left[ e^{-i\phi} \left( \frac{\partial^2 f}{\partial \theta^2} + i \csc^2 \theta \frac{\partial f}{\partial \phi} - i \cot \theta \frac{\partial^2 f}{\partial \phi \partial \theta} \right) \right. \\
&\quad \left. + \cot \theta e^{-i\phi} \left( \frac{\partial f}{\partial \theta} - i \cot \theta \frac{\partial f}{\partial \phi} \right) + i \cot \theta e^{-i\phi} \left( \frac{\partial^2 f}{\partial \phi \partial \theta} - i \cot \theta \frac{\partial^2 f}{\partial \phi^2} \right) \right] \\
&= -\hbar^2 \left( \frac{\partial^2 f}{\partial \theta^2} + i \csc^2 \theta \frac{\partial f}{\partial \phi} + \cot \theta \frac{\partial f}{\partial \theta} - i \cot^2 \theta \frac{\partial f}{\partial \phi} + \cot^2 \theta \frac{\partial^2 f}{\partial \phi^2} \right) \\
&= -\hbar^2 \left( \frac{\partial^2 f}{\partial \theta^2} + i \frac{\partial f}{\partial \phi} + \cot \theta \frac{\partial f}{\partial \theta} + \cot^2 \theta \frac{\partial^2 f}{\partial \phi^2} \right) \\
L_+ L_- &= -\hbar^2 \left( \frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \cot^2 \theta \frac{\partial^2}{\partial \phi^2} + i \frac{\partial}{\partial \phi} \right)
\end{aligned}$$



(b)

$$\begin{aligned}
L^2 &= L_+ L_- + L_z^2 - \hbar L_z \\
&= -\hbar^2 \left( \frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \cot^2 \theta \frac{\partial^2}{\partial \phi^2} + i \frac{\partial}{\partial \phi} \right) - \hbar^2 \frac{\partial^2}{\partial \phi^2} + i \hbar^2 \frac{\partial}{\partial \phi} \\
&= -\hbar^2 \left( \frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \cot^2 \theta \frac{\partial^2}{\partial \phi^2} + i \frac{\partial}{\partial \phi} + \frac{\partial^2}{\partial \phi^2} - i \frac{\partial}{\partial \phi} \right) \\
&= -\hbar^2 \left[ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \cot^2 \theta \frac{\partial^2}{\partial \phi^2} + \frac{\partial^2}{\partial \phi^2} \right] \\
&= -\hbar^2 \left[ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]
\end{aligned}$$

#### 4.25

(a)

$$L_+ Y_\ell^\ell = 0$$

(b)

$$\begin{aligned}
L_z Y_\ell^\ell &= \hbar \ell Y_\ell^\ell \\
-i\hbar \frac{\partial Y_\ell^\ell}{\partial \phi} &= \hbar \ell Y_\ell^\ell \\
\frac{\partial Y_\ell^\ell}{\partial \phi} &= i\ell Y_\ell^\ell \\
Y_\ell^\ell &= e^{i\ell\phi} \Theta(\theta) \\
0 &= L_+ Y_\ell^\ell \\
&= \hbar e^{i\phi} \left( \frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right) e^{i\ell\phi} \Theta(\theta) \\
&= \hbar e^{i\phi} \left[ e^{i\ell\phi} \frac{\partial \Theta}{\partial \theta} - \ell \cot \theta e^{i\ell\phi} \Theta(\theta) \right] \\
&= \frac{\partial \Theta}{\partial \theta} - \ell \cot \theta \Theta(\theta) \\
\frac{1}{\Theta(\theta)} \frac{\partial \Theta}{\partial \theta} &= \ell \cot \theta \\
\int \frac{1}{\Theta(\theta)} d\Theta &= \ell \int \frac{\cos \theta}{\sin \theta} d\theta \\
\ln \Theta &= \ell \ln \sin \theta + c_1 \\
&= \ln(\sin^\ell \theta) + c_1 \\
\ln \frac{\Theta}{\sin^\ell \theta} &= c_1 \\
\frac{\Theta}{\sin^\ell \theta} &= c_2 \\
\Theta &= c_2 \sin^\ell \theta \\
Y_\ell^\ell &= c e^{i\ell\phi} \sin^\ell \theta
\end{aligned}$$

4.26

$$\begin{aligned}
Y_2^1(\theta, \phi) &= -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi} \\
&= -\frac{1}{2}\sqrt{\frac{15}{8\pi}} \sin(2\theta) e^{i\phi} \\
L_+ Y_2^1 &= \hbar e^{i\phi} \left( \frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right) \left[ -\frac{1}{2}\sqrt{\frac{15}{8\pi}} \sin(2\theta) e^{i\phi} \right] \\
&= -\frac{1}{2}\sqrt{\frac{15}{8\pi}} \hbar e^{i\phi} [2 \cos(2\theta) e^{i\phi} - \cot \theta \sin(2\theta) e^{i\phi}] \\
&= \sqrt{\frac{15}{8\pi}} \hbar (e^{i\phi} \sin \theta)^2 \\
&= \hbar \sqrt{(2-1)(2+1+1)} Y_2^1 \\
&= 2\hbar Y_2^1 \\
Y_2^1 &= \frac{1}{4}\sqrt{\frac{15}{2\pi}} (e^{i\phi} \sin \theta)^2
\end{aligned}$$

4.28

$$\begin{aligned}
r &= \frac{e^2}{4\pi\epsilon_0 mc^2} \\
&= 0.0796 \text{ m} \\
\frac{1}{2}\hbar &= I\omega \\
&= \frac{2}{5}mr^2 \frac{v}{r} \\
v &= \frac{5\hbar}{4mr} \\
&= 5.14 \times 10^{10} \text{ m/s}
\end{aligned}$$

This is 100 times the speed of light, so no it doesn't make sense.

4.29

(a)

$$\begin{aligned}
[S_x, S_y] &= S_x S_y - S_y S_x \\
&= \begin{pmatrix} i\hbar^2/2 & 0 \\ 0 & -i\hbar^2/2 \end{pmatrix} \\
&= i\hbar S_z \\
[S_y, S_z] &= i\hbar S_x \\
[S_z, S_x] &= i\hbar S_y
\end{aligned}$$

(b)

$$\begin{aligned}
\sigma_x \sigma_x &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\
&= \delta_{xx} \\
\delta_x \delta_y &= \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} \\
&= i\epsilon_{xyz}\sigma_z \\
\delta_x \delta_z &= \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \\
&= i\epsilon_{xzy}\sigma_y \\
\delta_y \delta_x &= \begin{pmatrix} -i & 0 \\ 0 & 1 \end{pmatrix} \\
&= i\epsilon_{yxz}\sigma_z \\
\delta_y \delta_y &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\
&= \delta_{yy} \\
\delta_y \delta_z &= \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \\
&= i\epsilon_{yzx}\sigma_x \\
\delta_z \delta_x &= \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \\
&= i\epsilon_{zxy}\sigma_y \\
\delta_z \delta_y &= \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix} \\
&= i\epsilon_{zyx}\sigma_x \\
\delta_z \delta_z &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\
&= \delta_{zz}
\end{aligned}$$

#### 4.30

(a)

$$\begin{aligned}
1 &= \chi^\dagger \chi \\
&= |A|^2 \begin{pmatrix} -3i & 4 \end{pmatrix} \begin{pmatrix} 3i \\ 4 \end{pmatrix} \\
&= 25|A|^2 \\
A &= \pm \frac{1}{5}
\end{aligned}$$

(b)

$$\begin{aligned}\chi &= \frac{1}{5} \begin{pmatrix} 3i \\ 4 \end{pmatrix} \\ P\left(\frac{\hbar}{2}\right) &= \frac{1}{2} \left| \frac{1}{5}(3i + 4) \right|^2 \\ &= \frac{1}{50} (-3i + 4)(3i + 4) \\ &= \frac{1}{2} \\ \langle S_x \rangle &= 0 \\ \langle S_y \rangle &= \chi^\dagger S_y \chi \\ &= \frac{\hbar}{50} \begin{pmatrix} -3i & 4 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 3i \\ 4 \end{pmatrix} \\ &= \frac{\hbar}{50} \begin{pmatrix} -3i & 4 \end{pmatrix} \begin{pmatrix} 0 & -4i \\ -3 & 0 \end{pmatrix} \\ &= -\frac{12}{25} \hbar \\ \langle S_z \rangle &= \frac{\hbar}{50} \begin{pmatrix} -3i & 4 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 3i \\ 4 \end{pmatrix} \\ &= \frac{\hbar}{50} \begin{pmatrix} -3i & 4 \end{pmatrix} \begin{pmatrix} 3i \\ -4 \end{pmatrix} \\ &= -\frac{7}{50} \hbar\end{aligned}$$

(c)

$$\begin{aligned}\langle S_x^2 \rangle &= \frac{\hbar^2}{100} \begin{pmatrix} -3i & 4 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 3i \\ 4 \end{pmatrix} \\ &= \frac{\hbar^2}{4} \\ \sigma_{S_x} &= \sqrt{\langle S_x^2 \rangle - \langle S_x \rangle^2} \\ &= \frac{\hbar}{2} \\ \langle S_y^2 \rangle &= \frac{\hbar^2}{4} \\ \sigma_{S_y} &= \sqrt{\frac{1}{4} \hbar^2 - \frac{144}{625} \hbar^2} \\ &= \frac{7}{50} \hbar \\ \sigma_{S_z} &= \frac{12}{25} \hbar\end{aligned}$$

(d)

$$\begin{aligned}
\sigma_{S_x} \sigma_{S_y} &= \frac{7}{100} \hbar^2 \\
&= \frac{\hbar}{2} |\langle S_z \rangle| \\
&\geq \frac{\hbar}{2} |\langle S_z \rangle| \\
\sigma_{S_y} \sigma_{S_z} &= \frac{42}{625} \hbar^2 \\
&\geq \frac{\hbar}{2} |\langle S_x \rangle| \\
&= 0 \\
\sigma_{S_z} \sigma_{S_x} &= \frac{6}{25} \hbar^2 \\
&= \frac{\hbar}{2} |\langle S_y \rangle| \\
&\geq \frac{\hbar}{2} |\langle S_z \rangle|
\end{aligned}$$

**4.31**

$$\begin{aligned}
\chi &= \begin{pmatrix} a \\ b \end{pmatrix} \\
\langle S_x \rangle &= \frac{\hbar}{2} (a^* b + a b^*) \\
\langle S_y \rangle &= i \frac{\hbar}{2} (a b^* - a^* b) \\
\langle S_z \rangle &= \frac{\hbar}{2} (|a|^2 - |b|^2) \\
\langle S_x^2 \rangle &= \frac{\hbar^2}{4} (|a|^2 + |b|^2) \\
\langle S_y^2 \rangle &= \frac{\hbar^2}{4} (|a|^2 + |b|^2) \\
\langle S_z^2 \rangle &= \frac{\hbar^2}{4} (|a|^2 + |b|^2) \\
\langle S^2 \rangle &= 3 \frac{\hbar^2}{4} (|a|^2 + |b|^2) \\
\langle S_x^2 \rangle + \langle S_y^2 \rangle + \langle S_z^2 \rangle &= 3 \frac{\hbar^2}{4} (|a|^2 + |b|^2) \\
&= \langle S^2 \rangle
\end{aligned}$$

### 4.32

(a)

$$\begin{aligned}
 0 &= \begin{vmatrix} -\lambda & -i\hbar/2 \\ i\hbar/2 & -\lambda \end{vmatrix} \\
 &= \lambda^2 - \frac{\hbar^2}{4} \\
 \lambda &= \pm \frac{\hbar}{2} \\
 \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} &= \pm \frac{\hbar}{2} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \\
 \begin{pmatrix} -i\beta \\ i\alpha \end{pmatrix} &= \pm \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \\
 \beta &= \pm i\alpha \\
 \chi_+^{(y)} &= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix} \\
 \chi_-^{(y)} &= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}
 \end{aligned}$$

(b)

$$\begin{aligned}
 \chi &= \begin{pmatrix} a \\ b \end{pmatrix} \\
 &= \left( \frac{a - ib}{\sqrt{2}} \right) \chi_+^{(x)} + \left( \frac{a + ib}{\sqrt{2}} \right) \chi_-^{(x)}
 \end{aligned}$$

You could get  $\frac{\hbar}{2}$  with probability  $\frac{|a-ib|^2}{2}$  or  $-\frac{\hbar}{2}$  with probability  $\frac{|a+ib|^2}{2}$ .

$$\begin{aligned}
 \frac{1}{2}|a - ib|^2 + \frac{1}{2}|a + ib|^2 &= \frac{1}{2}(a + ib)(a - ib) + \frac{1}{2}(a - ib)(a + ib) \\
 &= \frac{1}{2}(a^2 + b^2) + \frac{1}{2}(a^2 + b^2) \\
 &= 1
 \end{aligned}$$

(c)  $\hbar^2/2$  with probability 1.

### 4.33

$$\begin{aligned}
S_r &= S \cdot \hat{\mathbf{r}} \\
&= S_x \sin \theta \cos \phi + S_y \sin \theta \sin \phi + S_z \cos \theta \\
&= \frac{\hbar}{2} \left[ \begin{pmatrix} 0 & \sin \theta \cos \phi \\ \sin \theta \cos \phi & 0 \end{pmatrix} + \begin{pmatrix} 0 & -i \sin \theta \sin \phi \\ i \sin \theta \sin \phi & 0 \end{pmatrix} \right. \\
&\quad \left. + \begin{pmatrix} \cos \theta & 0 \\ 0 & -\cos \theta \end{pmatrix} \right] \\
&= \frac{\hbar}{2} \begin{pmatrix} \cos \theta & \sin \theta (\cos \phi - i \sin \phi) \\ \sin \theta (\cos \phi + i \sin \phi) & -\cos \theta \end{pmatrix} \\
&= \frac{\hbar}{2} \begin{pmatrix} \cos \theta & e^{-i\phi} \sin \theta \\ e^{i\phi} \sin \theta & -\cos \theta \end{pmatrix}
\end{aligned}$$

$$\begin{aligned}
\left| \begin{pmatrix} \frac{\hbar}{2} \cos \theta - \lambda & \frac{\hbar}{2} e^{-i\phi} \sin \theta \\ \frac{\hbar}{2} e^{i\phi} \sin \theta & -\frac{\hbar}{2} \cos \theta - \lambda \end{pmatrix} \right| &= \left( \frac{\hbar}{2} \cos \theta - \lambda \right) \left( -\frac{\hbar}{2} \cos \theta - \lambda \right) \\
&\quad - \left( \frac{\hbar}{2} e^{-i\phi} \sin \theta \right) \left( \frac{\hbar}{2} e^{i\phi} \sin \theta \right) \\
&= -\frac{\hbar^2}{4} \cos^2 \theta - \lambda \frac{\hbar}{2} \cos \theta + \lambda \frac{\hbar}{2} \cos \theta + \lambda^2 \\
&\quad - \frac{\hbar^2}{4} \sin^2 \theta \\
&= \lambda^2 - \frac{\hbar^2}{4} \\
\lambda &= \pm \frac{\hbar}{2}
\end{aligned}$$

### 4.35

(a)

$$\begin{aligned}
\chi(t) &= \begin{pmatrix} \cos(\alpha/2) e^{i\gamma B_0 t/2} \\ \sin(\alpha/2) e^{-i\gamma B_0 t/2} \end{pmatrix} \\
\frac{1}{2} |a + b|^2 &= \frac{1}{2} [\cos(\alpha/2) e^{-i\gamma B_0 t/2} + \sin(\alpha/2) e^{i\gamma B_0 t/2}] \\
&\quad [\cos(\alpha/2) e^{i\gamma B_0 t/2} + \sin(\alpha/2) e^{-i\gamma B_0 t/2}] \\
&= \frac{1}{2} [\cos^2(\alpha/2) + \cos(\alpha/2) \sin(\alpha/2) e^{-i\gamma B_0 t} \\
&\quad + \cos(\alpha/2) \sin(\alpha/2) e^{i\gamma B_0 t} + \sin^2(\alpha/2)] \\
&= \frac{1}{2} [1 + \sin \alpha \cos(\gamma B_0 t)]
\end{aligned}$$



(b)

$$\begin{aligned}\frac{1}{2}|a - ib|^2 &= \frac{1}{2}[\cos(\alpha/2)e^{-i\gamma B_0 t/2} + i\sin(\alpha/2)e^{i\gamma B_0 t/2}] \\ &\quad [\cos(\alpha/2)e^{i\gamma B_0 t/2} - i\sin(\alpha/2)e^{-i\gamma B_0 t/2}] \\ &= \frac{1}{2}[\cos^2(\alpha/2) - i\cos(\alpha/2)\sin(\alpha/2)e^{-i\gamma B_0 t} \\ &\quad + i\cos(\alpha/2)\sin(\alpha/2)e^{i\gamma B_0 t} + \sin^2(\alpha/2)] \\ &= \frac{1}{2}[1 - \sin\alpha\sin(\gamma B_0 t)]\end{aligned}$$

(c)

$$|a|^2 = \cos^2 \frac{\alpha}{2}$$

#### 4.36

(a)

$$\begin{aligned}\mathbf{B} &= B_0 \cos(\omega t) \hat{k} \\ H &= -\gamma \mathbf{B} \cdot \mathbf{S} \\ &= -\gamma B_0 \cos(\omega t) S_z \\ &= -\frac{\gamma B_0 \hbar}{2} \cos(\omega t) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}\end{aligned}$$

(b)

$$\begin{aligned}
i\hbar \frac{\partial \chi}{\partial t} &= H\chi \\
i\hbar \begin{pmatrix} a'(t) \\ b'(t) \end{pmatrix} &= -\frac{\gamma B_0 \hbar}{2} \cos(\omega t) \begin{pmatrix} a(t) \\ -b(t) \end{pmatrix} \\
i\hbar a'(t) &= -\frac{\gamma B_0 \hbar}{2} \cos(\omega t) a(t) \\
\frac{1}{a(t)} a'(t) &= i \frac{\gamma B_0}{2} \cos(\omega t) \\
\ln[a(t)] &= i \frac{\gamma B_0}{2\omega} \sin(\omega t) + c_1 \\
a(t) &= A e^{i\gamma B_0 \sin(\omega t)/2\omega} \\
i\hbar b'(t) &= \frac{\gamma B_0 \hbar}{2} \cos(\omega t) b(t) \\
\frac{1}{b(t)} b'(t) &= -i \frac{\gamma B_0}{2} \cos(\omega t) \\
\ln[b(t)] &= -i \frac{\gamma B_0}{2\omega} \sin(\omega t) + c_2 \\
b(t) &= B e^{-i\gamma B_0 \sin(\omega t)/2\omega} \\
\chi &= \begin{pmatrix} A e^{i\gamma B_0 \sin(\omega t)/2\omega} \\ B e^{-i\gamma B_0 \sin(\omega t)/2\omega} \end{pmatrix} \\
1 &= \frac{1}{2} |A + B|^2 \\
&= \frac{1}{2} (A^* + B^*)(A + B) \\
&= \frac{1}{2} (|A|^2 + A^* B + B^* A + |B|^2) \\
0 &= \frac{1}{2} |A - B|^2 \\
&= \frac{1}{2} (A^* - B^*)(A - B) \\
&= \frac{1}{2} (|A|^2 - A^* B - B^* A + |B|^2) \\
\chi &= \frac{1}{\sqrt{2}} \begin{pmatrix} e^{i\gamma B_0 \sin(\omega t)/2\omega} \\ e^{-i\gamma B_0 \sin(\omega t)/2\omega} \end{pmatrix}
\end{aligned}$$

(c)

$$\begin{aligned}
\frac{A-B}{\sqrt{2}} &= \frac{1}{2} [e^{i\gamma B_0 \sin(\omega t)/2\omega} - e^{-i\gamma B_0 \sin(\omega t)/2\omega}] \\
&= i \sin \left[ \frac{\gamma B_0}{2\omega} \sin(\omega t) \right] \\
\frac{1}{2} |A-B|^2 &= \sin^2 \left[ \frac{\gamma B_0}{2\omega} \sin(\omega t) \right]
\end{aligned}$$

(d)

$$|B_0| = \frac{\omega\pi}{\gamma}$$

**4.37**

(a)

$$\begin{aligned}
S_- |10\rangle &= S_- \left[ \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) \right] \\
&= \frac{1}{\sqrt{2}} [(S_-^{(1)} |\uparrow\rangle) |\downarrow\rangle + |\uparrow\rangle (S_-^{(2)} |\downarrow\rangle) + (S_-^{(1)} |\downarrow\rangle) |\uparrow\rangle + |\downarrow\rangle (S_-^{(2)} |\uparrow\rangle)] \\
&= \frac{1}{\sqrt{2}} [\hbar |\downarrow\downarrow\rangle + \hbar |\downarrow\downarrow\rangle] \\
&= \sqrt{2}\hbar |1-1\rangle
\end{aligned}$$

(b)

$$\begin{aligned}
S_+ |00\rangle &= S_+ \left[ \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right] \\
&= \frac{1}{\sqrt{2}} [(S_+^{(1)} |\uparrow\rangle) |\downarrow\rangle + |\uparrow\rangle (S_+^{(2)} |\downarrow\rangle) - (S_+^{(1)} |\downarrow\rangle) |\uparrow\rangle - |\downarrow\rangle (S_+^{(2)} |\uparrow\rangle)] \\
&= \frac{1}{\sqrt{2}} (\hbar |\uparrow\uparrow\rangle - \hbar |\uparrow\uparrow\rangle) \\
&= 0 \\
S_- |00\rangle &= S_- \left[ \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right] \\
&= \frac{1}{\sqrt{2}} [(S_-^{(1)} |\uparrow\rangle) |\downarrow\rangle + |\uparrow\rangle (S_-^{(2)} |\downarrow\rangle) - (S_-^{(1)} |\downarrow\rangle) |\uparrow\rangle - |\downarrow\rangle (S_-^{(2)} |\uparrow\rangle)] \\
&= \frac{1}{\sqrt{2}} (\hbar |\downarrow\downarrow\rangle - \hbar |\downarrow\downarrow\rangle) \\
&= 0
\end{aligned}$$

(c)

$$\begin{aligned}
S^{(1)} \cdot S^{(2)} |\uparrow\uparrow\rangle &= (S_x^{(1)} |\uparrow\rangle)(S_x^{(2)} |\uparrow\rangle) + (S_y^{(1)} |\uparrow\rangle)(S_y^{(2)} |\uparrow\rangle) \\
&\quad + (S_z^{(1)} |\uparrow\rangle)(S_z^{(2)} |\uparrow\rangle) \\
&= \left(\frac{\hbar}{2} |\downarrow\rangle\right) \left(\frac{\hbar}{2} |\downarrow\rangle\right) + \left(\frac{i\hbar}{2} |\downarrow\rangle\right) \left(\frac{i\hbar}{2} |\downarrow\rangle\right) \\
&\quad + \left(\frac{\hbar}{2} |\uparrow\rangle\right) \left(\frac{\hbar}{2} |\uparrow\rangle\right) \\
&= \frac{\hbar^2}{4} |\downarrow\downarrow\rangle - \frac{\hbar^2}{4} |\downarrow\downarrow\rangle + \frac{\hbar^2}{4} |\uparrow\uparrow\rangle \\
S^2 |\uparrow\uparrow\rangle &= \left(\frac{3\hbar^2}{4} + \frac{3\hbar^2}{4} + 2\frac{\hbar^2}{4}\right) |\uparrow\uparrow\rangle \\
&= 2\hbar^2 |\uparrow\uparrow\rangle \\
S^{(1)} \cdot S^{(2)} |\downarrow\downarrow\rangle &= (S_x^{(1)} |\downarrow\rangle)(S_x^{(2)} |\downarrow\rangle) + (S_y^{(1)} |\downarrow\rangle)(S_y^{(2)} |\downarrow\rangle) \\
&\quad + (S_z^{(1)} |\downarrow\rangle)(S_z^{(2)} |\downarrow\rangle) \\
&= \left(\frac{\hbar}{2} |\uparrow\rangle\right) \left(\frac{\hbar}{2} |\uparrow\rangle\right) + \left(-\frac{i\hbar}{2} |\uparrow\rangle\right) \left(-\frac{i\hbar}{2} |\uparrow\rangle\right) \\
&\quad + \left(-\frac{\hbar}{2} |\downarrow\rangle\right) \left(-\frac{\hbar}{2} |\downarrow\rangle\right) \\
&= \frac{\hbar^2}{4} |\uparrow\uparrow\rangle - \frac{\hbar^2}{4} |\uparrow\uparrow\rangle + \frac{\hbar^2}{4} |\downarrow\downarrow\rangle \\
S^2 |\downarrow\downarrow\rangle &= \left(\frac{3\hbar^2}{4} + \frac{3\hbar^2}{4} + 2\frac{\hbar^2}{4}\right) |\downarrow\downarrow\rangle \\
&= 2\hbar^2 |\downarrow\downarrow\rangle
\end{aligned}$$

#### 4.38

(a)  $3/2, 1/2$

(b)  $0, 1$

#### 4.40

(a)

$$|31\rangle = \frac{1}{\sqrt{15}} |22\rangle |1-1\rangle + \sqrt{\frac{8}{15}} |21\rangle |10\rangle + \sqrt{\frac{6}{15}} |20\rangle |11\rangle$$

$2\hbar$  with probability  $1/15$ ,  $\hbar$  with probability  $8/15$ , and  $0$  with probability  $6/15$ .

#### 4.46

(a)

$$\begin{aligned}
V(r) &= \frac{1}{2}m\omega^2 r^2 \\
&= \frac{1}{2}m\omega^2(x^2 + y^2 + z^2) \\
\psi &= XYZ \\
-\frac{\hbar^2}{2m}\nabla^2\psi + V\psi &= E\psi \\
-\frac{\hbar^2}{2m}\left(\frac{X''}{X} + \frac{Y''}{Y} + \frac{Z''}{Z}\right) + \frac{1}{2}m\omega^2(x^2 + y^2 + z^2) &= E \\
-\frac{\hbar^2}{2m}\frac{X''}{X} + \frac{1}{2}m\omega^2 x^2 &= c_1 \\
-\frac{\hbar^2}{2m}X'' + \frac{1}{2}m\omega^2 x^2 X &= c_1 X \\
c_1 &= E_{x_n} \\
&= \left(n_x + \frac{1}{2}\right)\hbar\omega \\
\frac{\hbar^2}{2m}\left(\frac{Y''}{Y} + \frac{Z''}{Z}\right) - \frac{1}{2}m\omega^2(y^2 + z^2) + E &= c_1 \\
\frac{\hbar^2}{2m}\frac{Y''}{Y} - \frac{1}{2}m\omega^2 y^2 + E &= c_2 \\
-\frac{\hbar^2}{2m}Y'' + \frac{1}{2}m\omega^2 y^2 Y &= (E - c_2)Y \\
E - c_2 &= \left(n_y + \frac{1}{2}\right)\hbar\omega \\
c_1 - \frac{\hbar^2}{2m}\frac{Z''}{Z} + \frac{1}{2}m\omega^2 z^2 &= c_2 \\
-\frac{\hbar^2}{2m}Z'' + \frac{1}{2}m\omega^2 z^2 Z &= (c_2 - c_1)Z \\
c_2 - c_1 &= \left(n_z + \frac{1}{2}\right)\hbar\omega \\
c_2 &= (n_x + n_z + 1)\hbar\omega \\
E &= \left(n_x + n_y + n_z + \frac{3}{2}\right)\hbar\omega \\
&= \left(n + \frac{3}{2}\right)\hbar\omega
\end{aligned}$$

4.48

(b)

$$\begin{aligned}
 V &= -\frac{1}{4\pi\epsilon_0} \frac{e^2}{r} \\
 \nabla V &= \frac{\partial V}{\partial r} \hat{\mathbf{r}} \\
 &= \frac{1}{4\pi\epsilon_0} \frac{e^2}{r^2} \hat{\mathbf{r}} \\
 2 \langle T \rangle &= \langle \mathbf{r} \cdot \nabla V \rangle \\
 &= \left\langle \frac{1}{4\pi\epsilon_0} \frac{e^2}{r} \right\rangle \\
 &= -\langle V \rangle \\
 E_n &= \langle T \rangle + \langle V \rangle \\
 &= \langle T \rangle - 2 \langle T \rangle \\
 &= -\langle T \rangle \\
 \langle T \rangle &= -E_n \\
 \langle V \rangle &= 2E_n
 \end{aligned}$$

(c)

$$\begin{aligned}
 V &= \frac{1}{2} m \omega^2 r^2 \\
 \nabla V &= m \omega^2 r \hat{\mathbf{r}} \\
 2 \langle T \rangle &= \langle m \omega^2 r^2 \rangle \\
 &= 2 \langle V \rangle \\
 \langle T \rangle &= \langle V \rangle \\
 E_n &= \langle T \rangle + \langle V \rangle \\
 &= 2 \langle T \rangle \\
 \langle T \rangle &= \frac{E_n}{2} \\
 \langle V \rangle &= \frac{E_n}{2}
 \end{aligned}$$

4.52

(a)

$$\begin{aligned}
 \psi_{321} &= \frac{1}{9a} \frac{1}{\sqrt{30a}} e^{-r/3a} \frac{4r^2}{9a^2} L_0^5 \left( \frac{2r}{3a} \right) Y_2^1(\theta, \phi) \\
 &= -\frac{1}{\sqrt{\pi a}} \frac{r^2}{81a^3} e^{-r/3a} \cos \theta \sin \theta e^{i\phi}
 \end{aligned}$$

(b)

$$\begin{aligned}
\int |\psi_{321}|^2 &= \frac{1}{81^2 \pi a^7} \int_0^\infty \int_0^\pi \int_0^{2\pi} r^6 e^{-2r/3a} \cos^2 \theta \sin^3 \theta \, dr \, d\theta \, d\phi \\
&= \frac{2}{81^2 a^7} \int_0^\infty \int_0^\pi r^6 e^{-2r/3a} \cos^2 \theta \sin^3 \theta \, dr \, d\theta \\
&= \frac{8}{15 \cdot 81^2 a^7} \int_0^\infty r^6 e^{-2r/3a} \, dr \\
&= 1
\end{aligned}$$

(c)

$$\begin{aligned}
\langle r^s \rangle &= \int \psi_{321}^* r^s \psi_{321} \\
&= \frac{8}{98415 a^7} \int_0^\infty r^s r^6 e^{-2r/3a} \, dr \\
&= \frac{8}{98415 a^7} (s+6)! \left( \frac{3a}{2} \right)^{s+7} \\
&= \frac{1}{720} (s+6)! \left( \frac{3a}{2} \right)^s \\
&= \frac{(s+6)!}{6!} \left( \frac{3a}{2} \right)^s
\end{aligned}$$

Finite for  $s > -7$

## 4.53

(a)

$$\psi_{433} = -\frac{1}{\sqrt{\pi a}} \frac{1}{6144 a^4} e^{-r/4a} r^3 \sin^3 \theta e^{3i\phi}$$

(b)

$$\begin{aligned}
\langle r \rangle &= \frac{1}{6144^2 \pi a^9} \int_0^\infty \int_0^\pi \int_0^{2\pi} r^9 e^{-r/2a} \sin^7 \theta \, dr \, d\theta \, d\phi \\
&= \frac{2}{6144^2 a^9} \int_0^\infty \int_0^\pi r^9 e^{-r/2a} \sin^7 \theta \, dr \, d\theta \\
&= \frac{64}{35 \cdot 6144^2 a^9} \int_0^\infty r^9 e^{-r/2a} \, dr \\
&= \frac{64}{35 \cdot 6144^2 a^9} 371589120 a^10 \\
&= 18a
\end{aligned}$$

(c)

$$\begin{aligned}(L_x^2 + L_y^2)\psi_{433} &= (L^2 - L_z^2)\psi_{433} \\ &= [\hbar^2 3(3+1) - \hbar^2 3^2]\psi_{433} \\ &= 3\hbar^2\psi_{433}\end{aligned}$$

Probability 1.

#### 4.54

(a)

$$\begin{aligned}\psi_{100} &= \frac{1}{\sqrt{\pi a^3}} e^{-r/a} \\ P_{\text{inside}} &= \frac{1}{\pi a^3} \int_0^b \int_0^\pi \int_0^{2\pi} e^{-2r/a} r^2 \sin \theta \, dr \, d\theta \, d\phi \\ &= \frac{2}{a^3} \int_0^b \int_0^\pi e^{-2r/a} r^2 \sin \theta \, dr \, d\theta \\ &= \frac{4}{a^3} \int_0^b e^{-2r/a} r^2 \, dr \\ &= \frac{4}{a^3} \frac{1}{4} a \left[ a^2 - (a^2 + 2ab + 2b^2) e^{-2b/a} \right] \\ &= 1 - \left[ 1 + 2\frac{b}{a} + 2\left(\frac{b}{a}\right)^2 \right] e^{-2b/a}\end{aligned}$$

(b)

$$\begin{aligned}P_{\text{inside}} &= 1 - \left( 1 + \epsilon + \frac{1}{2}\epsilon^2 \right) e^{-\epsilon} \\ &= 1 - \left( 1 + \epsilon + \frac{1}{2}\epsilon^2 \right) \left( 1 - \epsilon + \frac{\epsilon^2}{2} - \frac{\epsilon^3}{6} + \cdots \right) \\ &= 1 - \left( 1 - \epsilon + \frac{\epsilon^2}{2} - \frac{\epsilon^3}{6} + \cdots \right) - \left( \epsilon - \epsilon^2 + \frac{\epsilon^3}{2} - \cdots \right) \\ &\quad - \left( \frac{1}{2}\epsilon^2 - \frac{1}{2}\epsilon^3 + \cdots \right) \\ &= \frac{\epsilon^3}{6} + \cdots \\ &= \frac{4}{3} \left( \frac{b}{a} \right)^3 + \cdots\end{aligned}$$



(c)

$$\begin{aligned}
 P_{\text{inside}} &\approx \frac{4\pi b^3}{3} |\psi(0)|^2 \\
 &= \frac{4\pi b^3}{3} \frac{1}{\pi a^3} \\
 &= \frac{4}{3} \left(\frac{b}{a}\right)^3
 \end{aligned}$$

(d)

$$P_{\text{inside}} \approx 1.07 \times 10^{-14}$$

#### 4.55

(a)

$$\begin{aligned}
 R_{n(n-1)}(r) &= c_0 \frac{1}{r} \rho^{\ell+1} e^{-\rho} \\
 &= c_0 \frac{1}{r} \left(\frac{r}{an}\right)^n e^{-r/an} \\
 &= N_n r^{n-1} e^{-r/an} \\
 1 &= N_n^2 \int_0^\infty r^{2n} e^{-2r/an} dr \\
 u &= \frac{2r}{an} \\
 r &= \frac{anu}{2} \\
 dr &= \frac{an}{2} du \\
 1 &= N_n^2 \int_0^\infty \left(\frac{anu}{2}\right)^{2n} e^{-u} \frac{an}{2} du \\
 &= N_n^2 \left(\frac{an}{2}\right)^{2n+1} \int_0^\infty u^{2n} e^{-u} du \\
 &= N_n^2 \left(\frac{an}{2}\right)^{2n+1} \Gamma(2n+1) \\
 &= N_n^2 \left(\frac{an}{2}\right)^{2n+1} (2n)! \\
 N_n &= \left(\frac{2}{an}\right)^n \sqrt{\frac{2}{an(2n)!}}
 \end{aligned}$$

(b)

$$\begin{aligned}\langle r \rangle &= \int R^* r R \\ &= N_n^2 \int_0^\infty r^{2n+1} e^{-2r/an} dr \\ u &= \frac{2r}{an} \\ r &= \frac{anu}{2} \\ dr &= \frac{an}{2} du \\ \langle r \rangle &= N_n^2 \int_0^\infty \left( \frac{anu}{2} \right)^{2n+1} e^{-u} \frac{an}{2} du \\ &= \left( n + \frac{1}{2} \right) na \\ \langle r^2 \rangle &= \int R^* r^2 R \\ &= \left( \frac{2}{an} \right)^{2n} \frac{2}{an(2n)!} \int_0^\infty r^{2n+2} e^{-2r/an} dr \\ &= \left( n + \frac{1}{2} \right) (n+1)(na)^2\end{aligned}$$

(c)

$$\begin{aligned}\sigma_r &= \sqrt{\langle r^2 \rangle - \langle r \rangle^2} \\ &= \sqrt{\left( n + \frac{1}{2} \right) (n+1)(na)^2 - \left( n + \frac{1}{2} \right)^2 (na)^2} \\ &= \left( n + \frac{1}{2} \right) na \sqrt{\frac{n+1}{n+\frac{1}{2}} - 1} \\ &= \langle r \rangle \sqrt{\frac{n+1-n-\frac{1}{2}}{n+\frac{1}{2}}} \\ &= \frac{\langle r \rangle}{\sqrt{2n+1}}\end{aligned}$$

4.57

(a)

$$\begin{aligned}
A &= x^2 \\
B &= L_z \\
[A, B] &= [x^2, L_z] \\
&= x[x, L_z] + [x, L_z]x \\
&= x[x, xp_y - yp_x] + [x, xp_y - yp_x]x \\
&= x[x, xp_y] - x[x, yp_x] + [x, xp_y]x - [x, yp_x]x \\
&= x^2[x, p_y] + x[x, x]p_y - xy[x, p_x] - x[x, y]p_x \\
&\quad + x[x, p_y]x + [x, x]p_yx - y[x, p_x]x - [x, y]p_xx \\
&= -i\hbar xy - i\hbar xy \\
&= -2i\hbar xy \\
\sigma_A^2 \sigma_B^2 &\geq \left( \frac{1}{2i} \langle [A, B] \rangle \right)^2 \\
&= \left( \frac{1}{2i} \langle -2i\hbar xy \rangle \right)^2 \\
&= \hbar^2 \langle xy \rangle^2 \\
\sigma_A \sigma_B &= \hbar |\langle xy \rangle|
\end{aligned}$$

(b)

$$\begin{aligned}
L_z \psi_{n\ell m} &= m\hbar \psi_{n\ell m} \\
L_z^2 \psi_{n\ell m} &= m^2 \hbar^2 \psi_{n\ell m} \\
\sigma_B &= \sqrt{\langle L_z^2 \rangle - \langle L_z \rangle^2} \\
&= \sqrt{m^2 \hbar^2 - m^2 \hbar^2} \\
&= 0
\end{aligned}$$

(c)

$$\langle xy \rangle = 0$$

# 4.58

(a)

$$\begin{aligned}
 1 &= \chi^\dagger \chi \\
 &= |A|^2 \begin{pmatrix} 1+2i & 2 \end{pmatrix} \begin{pmatrix} 1-2i \\ 2 \end{pmatrix} \\
 &= |A|^2 [(1+2i)(1-2i) + 4] \\
 &= 9|A|^2 \\
 A &= \frac{1}{3}
 \end{aligned}$$

(b)

$$\chi = \frac{1}{3} \begin{pmatrix} 1-2i \\ 2 \end{pmatrix}$$

You could get  $\frac{\hbar}{2}$  with probability  $\frac{1}{9}(1+2i)(1-2i) = \frac{5}{9}$  or  $-\frac{\hbar}{2}$  with probability  $\frac{4}{9}$ .

$$\langle S_z \rangle = \left( \frac{5}{9} - \frac{4}{9} \right) \frac{\hbar}{2} = \frac{\hbar}{18}$$

(c) You could get  $\frac{\hbar}{2}$  with probability

$$\frac{1}{9} \frac{(3+2i)(3-2i)}{2} = \frac{13}{18}$$

or  $-\frac{\hbar}{2}$  with probability

$$\frac{1}{9} \frac{(-1+2i)(-1-2i)}{2} = \frac{5}{18}.$$

$$\langle S_x \rangle = \left( \frac{13}{18} - \frac{5}{18} \right) \frac{\hbar}{2} = \frac{2}{9} \hbar$$

(d) You could get  $\frac{\hbar}{2}$  with probability

$$\frac{1}{9} \frac{(1+4i)(1-4i)}{2} = \frac{17}{18}$$

or  $-\frac{\hbar}{2}$  with probability

$$\frac{1}{18}.$$

$$\langle S_y \rangle = \left( \frac{17}{18} - \frac{1}{18} \right) \frac{\hbar}{2} = \frac{4}{9} \hbar$$

#### 4.61

The basis states are  $|\frac{3}{2}\rangle$ ,  $|\frac{1}{2}\rangle$ ,  $|\frac{1}{2}\rangle$ , and  $|\frac{3}{2}\rangle$ . The matrix corresponding to the raising operator is

$$S_+ = \hbar \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & \sqrt{3} \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

and the matrix corresponding to the lowering operator is

$$S_- = \hbar \begin{pmatrix} 0 & 0 & 0 & 0 \\ \sqrt{3} & 0 & 0 & 0 \\ 0 & 2 & 0 & 0 \\ 0 & 0 & \sqrt{3} & 0 \end{pmatrix}.$$

$S_{\pm} = S_x \pm S_y$  so

$$\begin{aligned} S_x &= \frac{S_+ + S_-}{2} \\ &= \frac{\hbar}{2} \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ \sqrt{3} & 0 & 2 & 0 \\ 0 & 2 & 0 & \sqrt{3} \\ 0 & 0 & \sqrt{3} & 0 \end{pmatrix}. \end{aligned}$$

The eigenvalues of  $S_x$  are  $\frac{3}{2}\hbar$ ,  $\frac{1}{2}\hbar$ ,  $-\frac{1}{2}\hbar$ , and  $-\frac{3}{2}\hbar$ , as expected.

#### 4.64

- (a) You would get  $2\hbar^2$  with probability 1.
- (b) You could get 0 with probability  $\frac{1}{3}$  or  $3\hbar$  with probability  $\frac{2}{3}$ .
- (c) You would get  $\frac{3}{4}\hbar^2$  with probability 1.
- (d) You could get  $\frac{\hbar}{2}$  with probability  $\frac{1}{3}$  or  $-\frac{\hbar}{2}$  with probability  $\frac{2}{3}$ .
- (e)

$$\begin{aligned} \sqrt{\frac{1}{3}} \left| \begin{smallmatrix} 1 & 1 \\ 2 & 2 \end{smallmatrix} \right\rangle |10\rangle + \sqrt{\frac{2}{3}} \left| \begin{smallmatrix} 1 & 1 \\ 2 & -2 \end{smallmatrix} \right\rangle |11\rangle &= \frac{1}{\sqrt{3}} \left( \sqrt{\frac{2}{3}} \left| \begin{smallmatrix} 3 & 1 \\ 2 & 2 \end{smallmatrix} \right\rangle - \frac{1}{\sqrt{3}} \left| \begin{smallmatrix} 1 & 1 \\ 2 & 2 \end{smallmatrix} \right\rangle \right) \\ &\quad + \sqrt{\frac{2}{3}} \left( \frac{1}{\sqrt{3}} \left| \begin{smallmatrix} 3 & 1 \\ 2 & 2 \end{smallmatrix} \right\rangle + \sqrt{\frac{2}{3}} \left| \begin{smallmatrix} 1 & 1 \\ 2 & 2 \end{smallmatrix} \right\rangle \right) \\ &= \sqrt{\frac{8}{9}} \left| \begin{smallmatrix} 3 & 1 \\ 2 & 2 \end{smallmatrix} \right\rangle + \sqrt{\frac{1}{9}} \left| \begin{smallmatrix} 1 & 1 \\ 2 & 2 \end{smallmatrix} \right\rangle \end{aligned}$$

You could measure  $\frac{15}{4}\hbar^2$  with probability  $\frac{8}{9}$  or  $\frac{3}{4}\hbar^2$  with probability  $\frac{1}{9}$ .

(f) You would measure  $\frac{\hbar}{2}$  with probability 1.

(g)

$$\begin{aligned}
& R_{21}^* \left( \sqrt{\frac{1}{3}} (Y_1^0)^* \chi_+^\dagger + \sqrt{\frac{2}{3}} (Y_1^1)^* \chi_-^\dagger \right) R_{21} \left( \sqrt{\frac{1}{3}} Y_1^0 \chi_+ + \sqrt{\frac{2}{3}} Y_1^1 \chi_- \right) \\
&= \frac{1}{3} |R_{21}|^2 (|Y_1^0|^2 + 2|Y_1^1|^2) \\
&= \frac{1}{3} \frac{1}{24a^5} r^2 e^{-r/a} \left( \frac{3}{4\pi} \cos^2 \theta + \frac{3}{4\pi} \sin^2 \theta \right) \\
&= \frac{1}{96\pi a^5} r^2 e^{-r/a}
\end{aligned}$$

(h)

$$\begin{aligned}
& \frac{1}{3} |R_{21}|^2 \int_0^\pi \int_0^{2\pi} |Y_1^0|^2 \sin \theta \, d\theta \, d\phi \\
&= \frac{1}{3} \frac{1}{24a^5} r^2 e^{-r/a} \int_0^\pi \int_0^{2\pi} \frac{3}{4\pi} \cos^2 \theta \sin \theta \, d\theta \, d\phi \\
&= \frac{1}{48a^5} r^2 e^{-r/a} \int_0^\pi \cos^2 \theta \sin \theta \, d\theta \\
&= \frac{1}{72a^5} r^2 e^{-r/a}
\end{aligned}$$

#### 4.70

- (a) At  $t = 0$  all particles are in the spin-up state with respect to the  $z$  axis  $|\uparrow\rangle$ . At  $t = t_1$  the spin of a fraction  $f$  of the particles is measured with respect to the  $z$  axis (the  $z$  group) and the remaining fraction  $1 - f$  is measured with respect to the  $x$  axis (the  $x$  group). There is no change in the  $z$  group so the fraction of particles in the state  $|\uparrow\rangle$  is  $f$ . The  $x$  group is evenly split, so the fraction of particles in either the  $|\leftarrow\rangle$  or  $|\rightarrow\rangle$  states is  $\frac{1}{2}(1 - f)$ . At  $t = t_2$  the spin of all particles is measured with respect to the  $x$  axis. The  $z$  group is evenly split and the  $x$  group is unchanged so the fraction of particles in either the  $|\leftarrow\rangle$  or  $|\rightarrow\rangle$  states is  $\frac{1}{2}f + \frac{1}{2}(1 - f) = \frac{1}{2}$ . The subensemble of particles in the  $|\rightarrow\rangle$  state is comprised of particles from both the  $z$  and  $x$  groups. Of those from the  $z$  group, all of them were in the spin-up state at  $t = t_1$ , and of those from the  $x$  group, they also were all in the spin-up state at  $t = t_1$ . In conclusion: all of the particles in the subensemble had spin-up in the first measurement.

(b)

$$\begin{aligned}P_+ &= \cos^2(\theta_{ab}/2) \\P_- &= \sin^2(\theta_{ab}/2) \\P_{++} &= \cos^2(\theta_{ab}/2) \cos^2(\theta_{bc}/2) \\P_{+-} &= \cos^2(\theta_{ab}/2) \sin^2(\theta_{bc}/2) \\P_{-+} &= \sin^2(\theta_{ab}/2) \sin^2(\theta_{bc}/2) \\P_{--} &= \sin^2(\theta_{ab}/2) \cos^2(\theta_{bc}/2) \\\frac{P_{++}}{P_{++} + P_{-+}} &= \frac{1}{1 + P_{-+}/P_{++}} \\&= [1 + \tan^2(\theta_{ab}/2) \tan^2(\theta_{bc}/2)]^{-1}\end{aligned}$$

#### 4.71

(a)

$$\begin{aligned}
\psi_{21m} &= \frac{1}{\sqrt{24a^5}} r e^{-r/2a} Y_1^m(\theta, \phi) \\
\psi_{21-1} &= \frac{1}{\sqrt{24a^5}} r e^{-r/2a} \sqrt{\frac{3}{8\pi}} \sin \theta e^{-i\phi} \\
&= \frac{1}{\sqrt{64\pi a^5}} r e^{-r/2a} \sin \theta e^{-i\phi} \\
\psi_{210} &= \frac{1}{\sqrt{24a^5}} r e^{-r/2a} \sqrt{\frac{3}{4\pi}} \cos \theta \\
&= \frac{1}{\sqrt{32\pi a^5}} r e^{-r/2a} \cos \theta \\
\psi_{211} &= -\frac{1}{\sqrt{24a^5}} r e^{-r/2a} \sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi} \\
&= -\frac{1}{\sqrt{64\pi a^5}} r e^{-r/2a} \sin \theta e^{i\phi} \\
\psi_{2p_x} &= \frac{1}{\sqrt{32\pi a^3}} \frac{x}{a} e^{-r/2a} \\
&= \frac{1}{\sqrt{32\pi a^5}} r e^{-r/2a} \sin \theta \cos \phi \\
&= \frac{1}{\sqrt{2}} (\psi_{21-1} - \psi_{211}) \\
\psi_{2p_y} &= \frac{1}{\sqrt{32\pi a^3}} \frac{y}{a} e^{-r/2a} \\
&= \frac{1}{\sqrt{32\pi a^5}} r e^{-r/2a} \sin \theta \sin \phi \\
&= \frac{i}{\sqrt{2}} (\psi_{21-1} + \psi_{211}) \\
\psi_{2p_z} &= \frac{1}{\sqrt{32\pi a^3}} \frac{z}{a} e^{-r/2a} \\
&= \frac{1}{\sqrt{32\pi a^5}} r e^{-r/2a} \cos \theta \\
&= \psi_{210}
\end{aligned}$$



## 5 Identical Particles

### 5.1

(a)

$$\begin{aligned}
 \mathbf{R} + \frac{\mu}{m_1} \mathbf{r} &= \frac{m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2}{m_1 + m_2} + \frac{m_2}{m_1 + m_2} (\mathbf{r}_1 - \mathbf{r}_2) \\
 &= \frac{m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2 + m_2 \mathbf{r}_1 - m_2 \mathbf{r}_2}{m_1 + m_2} \\
 &= \mathbf{r}_1 \\
 \mathbf{R} - \frac{\mu}{m_2} \mathbf{r} &= \frac{m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2}{m_1 + m_2} - \frac{m_1}{m_1 + m_2} (\mathbf{r}_1 - \mathbf{r}_2) \\
 &= \frac{m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2 - m_1 \mathbf{r}_1 + m_1 \mathbf{r}_2}{m_1 + m_2} \\
 &= \mathbf{r}_2
 \end{aligned}$$

Let  $\mathbf{r} = (x, y, z)$  and  $\mathbf{R} = (X, Y, Z)$ , then

$$\begin{aligned}
 (\nabla_1)_x &= \frac{\partial}{\partial x_1} \\
 &= \frac{\partial X}{\partial x_1} \frac{\partial}{\partial X} + \frac{\partial x}{\partial x_1} \frac{\partial}{\partial x} \\
 &= \frac{\mu}{m_2} (\nabla_R)_x + (\nabla_r)_x \\
 \nabla_1 &= \frac{\mu}{m_2} \nabla_R + \nabla_r \\
 (\nabla_2)_x &= \frac{\partial}{\partial x_2} \\
 &= \frac{\partial X}{\partial x_2} \frac{\partial}{\partial X} + \frac{\partial x}{\partial x_2} \frac{\partial}{\partial x} \\
 &= \frac{\mu}{m_1} (\nabla_R)_x - (\nabla_r)_x \\
 \nabla_2 &= \frac{\mu}{m_1} \nabla_R - \nabla_r
 \end{aligned}$$

(b)

$$\begin{aligned}
\nabla_1^2 \psi &= \nabla_1 \cdot \nabla_1 \psi \\
&= \nabla_1 \cdot \left( \frac{\mu}{m_2} \nabla_R \psi + \nabla_r \psi \right) \\
&= \frac{\mu}{m_2} \nabla_R \cdot \left( \frac{\mu}{m_2} \nabla_R \psi + \nabla_r \psi \right) + \nabla_r \cdot \left( \frac{\mu}{m_2} \nabla_R \psi + \nabla_r \psi \right) \\
&= \left( \frac{\mu}{m_2} \right)^2 \nabla_R^2 \psi + 2 \frac{\mu}{m_2} (\nabla_r \cdot \nabla_R) \psi + \nabla_r^2 \psi \\
\nabla_2^2 \psi &= \nabla_2 \cdot \nabla_2 \psi \\
&= \nabla_2 \cdot \left( \frac{\mu}{m_1} \nabla_R \psi - \nabla_r \psi \right) \\
&= \frac{\mu}{m_1} \nabla_R \cdot \left( \frac{\mu}{m_1} \nabla_R \psi - \nabla_r \psi \right) - \nabla_r \cdot \left( \frac{\mu}{m_1} \nabla_R \psi - \nabla_r \psi \right) \\
&= \left( \frac{\mu}{m_1} \right)^2 \nabla_R^2 \psi - 2 \frac{\mu}{m_1} (\nabla_r \cdot \nabla_R) \psi + \nabla_r^2 \psi
\end{aligned}$$

$$\begin{aligned}
& -\frac{\hbar^2}{2m_1} \nabla_1^2 \psi - \frac{\hbar^2}{2m_2} \nabla_2^2 \psi + V\psi = E\psi \\
& -\frac{\hbar^2}{2m_1} \left[ \left( \frac{\mu}{m_2} \right)^2 \nabla_R^2 \psi + 2 \frac{\mu}{m_2} (\nabla_r \cdot \nabla_R) \psi + \nabla_r^2 \psi \right] \\
& -\frac{\hbar^2}{2m_2} \left[ \left( \frac{\mu}{m_1} \right)^2 \nabla_R^2 \psi - 2 \frac{\mu}{m_1} (\nabla_r \cdot \nabla_R) \psi + \nabla_r^2 \psi \right] + V\psi = E\psi \\
& -\frac{\hbar^2 \mu^2}{2} \left( \frac{1}{m_1 m_2^2} + \frac{1}{m_1^2 m_2} \right) \nabla_R^2 \psi - \frac{\hbar^2}{2} \left( \frac{1}{m_1} + \frac{1}{m_2} \right) \nabla_r^2 \psi + V\psi = E\psi \\
& -\frac{\hbar^2}{2(m_1 + m_2)} \nabla_R^2 \psi - \frac{\hbar^2}{2\mu} \nabla_r^2 \psi + V\psi = E\psi
\end{aligned}$$

(c)

$$\begin{aligned}
\psi(\mathbf{R}, \mathbf{r}) &= \psi_R(\mathbf{R}) \psi_r(\mathbf{r}) \\
& -\frac{\hbar^2}{2(m_1 + m_2)} \nabla_R^2 \psi_R \psi_r - \frac{\hbar^2}{2\mu} \nabla_r^2 \psi_r \psi_R + V(\mathbf{r}) \psi_R \psi_r = E \psi_R \psi_r \\
& -\frac{\hbar^2}{2(m_1 + m_2)} \frac{\nabla_R^2 \psi_R}{\psi_R} - \frac{\hbar^2}{2\mu} \frac{\nabla_r^2 \psi_r}{\psi_r} + V(\mathbf{r}) = E \\
& \left( -\frac{\hbar^2}{2(m_1 + m_2)} \frac{\nabla_R^2 \psi_R}{\psi_R} - E_R \right) + \left( -\frac{\hbar^2}{2\mu} \frac{\nabla_r^2 \psi_r}{\psi_r} + V(\mathbf{r}) - E_r \right) = 0
\end{aligned}$$

## 5.2

(a)

$$\begin{aligned}\mu &= \frac{m_e m_p}{m_e + m_p} \\ E_1 &= -\frac{\mu}{2\hbar^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \\ &\approx -13.5982 \text{ eV}\end{aligned}$$

(b)

$$\begin{aligned}\Delta E_H &= E_{3_H} - E_{2_H} \\ &= \frac{E_{1_H}}{9} - \frac{E_{1_H}}{4} \\ &= \left( \frac{1}{9} - \frac{1}{4} \right) E_{1_H} \\ &\approx 1.888 \text{ eV} \\ \lambda_H &= 6.569 \times 10^{-7} \text{ m} \\ E_{1_D} &= \frac{m_e(m_n + m_p)}{m_e + m_n + m_p} \frac{1}{2\hbar^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \\ \Delta E_D &= E_{3_D} - E_{2_D} \\ &= \left( \frac{1}{9} - \frac{1}{4} \right) E_{1_D} \\ &\approx 1.889 \text{ eV} \\ \lambda_D &= 6.567 \times 10^{-7} \text{ m} \\ \Delta\lambda &= 1.79 \times 10^{-10} \text{ m}\end{aligned}$$

(c)

$$\begin{aligned}\mu &= \frac{m_e^2}{m_e + m_2} \\ &= \frac{1}{2} m_e \\ E'_1 &= \frac{1}{2} E_1 \\ &= -6.8 \text{ eV}\end{aligned}$$

(d)

$$\begin{aligned}
\mu_p &= \frac{m_e m_p}{m_e + m_p} \\
\mu_\mu &= \frac{m_\mu m_p}{m_\mu + m_p} \\
\frac{\mu_\mu}{\mu_p} &= \frac{m_\mu m_p}{m_\mu + m_p} \cdot \frac{m_e + m_p}{m_e m_p} \\
&= \frac{m_\mu (m_e + m_p)}{m_e (m_\mu + m_p)} \\
&= 206.77 \frac{m_e + m_p}{206.77 m_e + m_p} \\
&\approx 185.943 \\
\frac{1}{\lambda} &= R' \left( 1 - \frac{1}{4} \right) \\
&= \frac{3}{4} R' \\
\lambda &= \frac{4}{3 R'} \\
&\approx 6.54 \times 10^{-10} \text{ m}
\end{aligned}$$

## 5.4

(a)

$$\begin{aligned}
1 &= \int |\psi|^2 \\
&= |A|^2 \int [\psi_a(\mathbf{r}_1) \psi_b(\mathbf{r}_2) \pm \psi_b(\mathbf{r}_1) \psi_a(\mathbf{r}_2)]^* [\psi_a(\mathbf{r}_1) \psi_b(\mathbf{r}_2) \pm \psi_b(\mathbf{r}_1) \psi_a(\mathbf{r}_2)] \\
&= |A|^2 \int |\psi_a(\mathbf{r}_1)|^2 |\psi_b(\mathbf{r}_2)|^2 \pm \psi_a(\mathbf{r}_1)^* \psi_b(\mathbf{r}_1) \psi_a(\mathbf{r}_2) \psi_b(\mathbf{r}_2)^* \\
&\quad \pm \psi_a(\mathbf{r}_1) \psi_b(\mathbf{r}_1)^* \psi_a(\mathbf{r}_2)^* \psi_b(\mathbf{r}_2) + |\psi_a(\mathbf{r}_2)|^2 |\psi_b(\mathbf{r}_1)|^2 \\
&= 2|A|^2 \\
A &= \frac{1}{\sqrt{2}}
\end{aligned}$$

(b)

$$\begin{aligned}
1 &= \int |\psi|^2 \\
&= |A|^2 \int [2\psi_a(\mathbf{r}_1)\psi_a(\mathbf{r}_2)]^* [2\psi_a(\mathbf{r}_1)\psi_a(\mathbf{r}_2)] \\
&= 4|A|^2 \int \psi_a^*(\mathbf{r}_1)\psi_a(\mathbf{r}_1)\psi_a^*(\mathbf{r}_2)\psi_a(\mathbf{r}_2) \\
A &= \frac{1}{2}
\end{aligned}$$

## 5.5

(a)

$$\begin{aligned}
\hat{H} &= -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_1^2} - \frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_2^2} \\
\hat{H}\psi_{12} &= -\frac{\hbar^2}{2m} \frac{\sqrt{2}}{a} \left[ -\left(\frac{\pi}{a}\right)^2 \sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{2\pi x_2}{a}\right) \right. \\
&\quad + \left(\frac{2\pi}{a}\right)^2 \sin\left(\frac{2\pi x_1}{a}\right) \sin\left(\frac{\pi x_2}{a}\right) \\
&\quad - \left(\frac{2\pi}{a}\right)^2 \sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{2\pi x_2}{a}\right) \\
&\quad \left. + \left(\frac{\pi}{a}\right)^2 \sin\left(\frac{2\pi x_1}{a}\right) \sin\left(\frac{\pi x_2}{a}\right) \right] \\
&= -\frac{\hbar^2}{2m} \frac{\sqrt{2}}{a} \left[ -5\left(\frac{\pi}{a}\right)^2 \sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{2\pi x_2}{a}\right) \right. \\
&\quad \left. + 5\left(\frac{\pi}{a}\right)^2 \sin\left(\frac{2\pi x_1}{a}\right) \sin\left(\frac{\pi x_2}{a}\right) \right] \\
&= 5\left(\frac{\pi}{a}\right)^2 \frac{\hbar^2}{2m} \psi \\
E &= \frac{5\pi^2 \hbar^2}{2ma^2} \\
&= 5K
\end{aligned}$$

## 5.6

(a)

$$\begin{aligned}
\langle x \rangle_l &= \frac{2}{a} \int_0^a x \sin^2 \left( \frac{l\pi}{a} x \right) dx \\
&= \frac{a}{2} \\
\langle x^2 \rangle_l &= \frac{2}{a} \int_0^a x^2 \sin^2 \left( \frac{l\pi}{a} x \right) dx \\
&= a^2 \left( \frac{1}{3} - \frac{1}{2(l\pi)^2} \right) \\
\langle (x_1 - x_2)^2 \rangle &= \langle x^2 \rangle_l + \langle x^2 \rangle_n - 2 \langle x \rangle_l \langle x \rangle_n \\
&= a^2 \left( \frac{1}{3} - \frac{1}{2(l\pi)^2} \right) + a^2 \left( \frac{1}{3} - \frac{1}{2(n\pi)^2} \right) - \frac{a^2}{2} \\
&= a^2 \left[ \frac{1}{6} - \frac{1}{2\pi^2} \left( \frac{1}{l^2} + \frac{1}{n^2} \right) \right]
\end{aligned}$$

(b)

$$\begin{aligned}
\langle x \rangle_{ln} &= \frac{2}{a} \int x \sin \left( \frac{l\pi}{a} x \right) \sin \left( \frac{n\pi}{a} x \right) dx \\
&= \frac{4aln[-1 + (-1)^{l+n}]}{\pi^2(l^2 - n^2)^2} \\
&= \begin{cases} \frac{a(-8ln)}{\pi^2(l^2 - n^2)^2} & \text{if } l \text{ and } n \text{ have opposite parity} \\ 0 & \text{if } l \text{ and } n \text{ have same parity} \end{cases} \\
\langle (x_1 - x_2)^2 \rangle &= a^2 \left[ \frac{1}{6} - \frac{1}{2\pi^2} \left( \frac{1}{l^2} + \frac{1}{n^2} \right) \right] - \frac{128a^2l^2n^2}{\pi^4(l^2 - n^2)^4}
\end{aligned}$$

(c)

$$\langle (x_1 - x_2)^2 \rangle = a^2 \left[ \frac{1}{6} - \frac{1}{2\pi^2} \left( \frac{1}{l^2} + \frac{1}{n^2} \right) \right] + \frac{128a^2l^2n^2}{\pi^4(l^2 - n^2)^4}$$

## 5.7

(a) (i)

$$\psi(x_1, x_2) = \psi_0(x_1)\psi_1(x_2)$$

(ii)

$$\psi(x_1, x_2) = \frac{1}{2}[\psi_0(x_1)\psi_1(x_2) + \psi_1(x_1)\psi_0(x_2)]$$

(iii)

$$\psi(x_1, x_2) = \frac{1}{2}[\psi_0(x_1)\psi_1(x_2) - \psi_1(x_1)\psi_0(x_2)]$$

(b) (i)

$$\begin{aligned}\langle (x_1 - x_2)^2 \rangle &= \langle x^2 \rangle_0 + \langle x^2 \rangle_1 - 2 \langle x \rangle_0 \langle x \rangle_1 \\ &= \frac{4\hbar}{2m\omega}\end{aligned}$$

## 5.8

(a)

$$\psi(x_1, x_2, x_3) = \psi_a(x_1)\psi_b(x_2)\psi_c(x_3)$$

(b)

$$\begin{aligned}\psi(x_1, x_2, x_3) &= \frac{1}{\sqrt{6}}[\psi_a(x_1)\psi_b(x_2)\psi_c(x_3) + \psi_a(x_1)\psi_c(x_2)\psi_b(x_3) \\ &\quad + \psi_b(x_1)\psi_a(x_2)\psi_c(x_3) + \psi_b(x_1)\psi_c(x_2)\psi_a(x_3) \\ &\quad + \psi_c(x_1)\psi_a(x_2)\psi_b(x_3) + \psi_c(x_1)\psi_b(x_2)\psi_a(x_3)]\end{aligned}$$

(c)

$$\begin{aligned}\begin{vmatrix} \psi_a(x_1) & \psi_b(x_1) & \psi_c(x_1) \\ \psi_a(x_2) & \psi_b(x_2) & \psi_c(x_2) \\ \psi_a(x_3) & \psi_b(x_3) & \psi_c(x_3) \end{vmatrix} &= \psi_a(x_1)[\psi_b(x_2)\psi_c(x_3) - \psi_c(x_2)\psi_b(x_3)] \\ &\quad - \psi_b(x_1)[\psi_a(x_2)\psi_c(x_3) - \psi_c(x_2)\psi_a(x_3)] \\ &\quad + \psi_c(x_1)[\psi_a(x_2)\psi_b(x_3) - \psi_b(x_2)\psi_a(x_3)] \\ \psi(x_1, x_2, x_3) &= \frac{1}{\sqrt{6}}[\psi_a(x_1)\psi_b(x_2)\psi_c(x_3) - \psi_a(x_1)\psi_c(x_2)\psi_b(x_3) \\ &\quad - \psi_b(x_1)\psi_a(x_2)\psi_c(x_3) + \psi_b(x_1)\psi_c(x_2)\psi_a(x_3) \\ &\quad + \psi_c(x_1)\psi_a(x_2)\psi_b(x_3) - \psi_c(x_1)\psi_b(x_2)\psi_a(x_3)]\end{aligned}$$

## 5.13

(a)

$$\begin{aligned}E_{\text{before}} &= 4 \left( \frac{E_1}{4} + \frac{E_1}{4} \right) \\ &= 2E_1 \\ E_{\text{after}} &= 4 \left( \frac{E_1}{4} + E_1 \right) \\ &= 5E_1 \\ E_{\text{electron}} &= 4 \frac{E_1}{4} + \Delta E \\ &= -2E_1 \\ &= 27.2 \text{ eV}\end{aligned}$$

(b)

$$\begin{aligned}\mathcal{R}' &= \frac{m_e}{4\pi c\hbar^3} \left( \frac{2e^2}{4\pi\epsilon_0} \right)^2 \\ &= 4\mathcal{R} \\ \frac{1}{\lambda} &= 4\mathcal{R} \left( \frac{1}{n_f^2} - \frac{1}{n_i^2} \right)\end{aligned}$$

### 5.14

- (a) If the electrons were identical bosons the total wavefunction would be symmetric — a symmetric spatial wavefunction would be paired with a triplet spin state and an antisymmetric spatial wavefunction would be paired with a singlet spin state.
- (b) If the electrons were distinguishable particles there would be no (anti)symmetry requirement on the total wavefunction, so every energy level could be in the singlet or triplet state.



### 5.15

(a)

$$\begin{aligned}
\left\langle \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \right\rangle &= \int \psi_0^* \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \psi_0 \\
&= \int \frac{1}{|r_1 - r_2|} |\psi_0|^2 d^3\mathbf{r}_2 d^3\mathbf{r}_1 \\
&= \frac{64}{\pi^2 a^6} \int_0^\infty \int_0^\pi \int_0^{2\pi} \frac{1}{\sqrt{r_1^2 + r_2^2 - 2r_1 r_2 \cos \theta_2}} \\
&\quad e^{-4(r_1+r_2)/a} r_2^2 \sin \theta_2 dr_2 d\theta_2 d\phi_2 d^3\mathbf{r}_1 \\
&= \frac{128}{\pi a^6} \int_0^\infty \frac{r_1 + r_2 - |r_1 - r_2|}{r_1 r_2} e^{-4(r_1+r_2)/a} r_2^2 dr_2 d^3\mathbf{r}_1 \\
&= \frac{256}{\pi a^6} \left( \int_0^{r_1} \frac{r_2^2}{r_1} e^{-4(r_1+r_2)/a} dr_2 d^3\mathbf{r}_1 \right. \\
&\quad \left. + \int_{r_1}^\infty r_2 e^{-4(r_1+r_2)/a} dr_2 d^3\mathbf{r}_1 \right) \\
&= \frac{256}{\pi a^6} \int_0^\infty \int_0^\pi \int_0^{2\pi} \frac{a^2 e^{-8r_1/a} [a(e^{4r_1/a} - 1) - 2r_1]}{32r_1} \\
&\quad r_1^2 \sin \theta_1 dr_1 d\theta_1 d\phi_1 \\
&= \frac{8}{\pi a^4} \int_0^\infty \int_0^\pi \int_0^{2\pi} r_1 e^{-8r_1/a} [a(e^{4r_1/a} - 1) - 2r_1] \\
&\quad \sin \theta_1 dr_1 d\theta_1 d\phi_1 \\
&= \frac{32}{a^4} \int_0^\infty r_1 e^{-8r_1/a} [a(e^{4r_1/a} - 1) - 2r_1] dr_1 \\
&= \frac{5}{4a}
\end{aligned}$$

(b)

$$\begin{aligned}
E_{\text{interaction}} &= \frac{e^2}{4\pi\epsilon_0} \left\langle \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \right\rangle \\
&\approx 34 \text{ eV} \\
E'_0 &= E_0 + E_{\text{interaction}} \\
&\approx -75 \text{ eV}
\end{aligned}$$

### 5.16

$$\begin{aligned}
\psi_0(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) &= \psi_{100}(\mathbf{r}_1) \psi_{100}(\mathbf{r}_2) \psi_{200}(\mathbf{r}_3) \\
E'_1 &= 9(E_1 + E_1 + E_2) \\
&\approx -275.4 \text{ eV}
\end{aligned}$$

## 5.17

- (a)
- Hydrogen  $(1s)^1$
  - Helium  $(1s)^2$
  - Lithium  $(1s)^2(2s)^1$
  - Beryllium  $(1s)^2(2s)^2$
  - Boron  $(1s)^2(2s)^2(2p)^1$
  - Carbon  $(1s)^2(2s)^2(2p)^2$
  - Nitrogen  $(1s)^2(2s)^2(2p)^3$
  - Oxygen  $(1s)^2(2s)^2(2p)^4$
  - Fluorine  $(1s)^2(2s)^2(2p)^5$
  - Neon  $(1s)^2(2s)^2(2p)^6$
- (b)
- Hydrogen  $^2S_{1/2}$
  - Helium  $^1S_0$
  - Lithium  $^2S_{1/2}$
  - Beryllium  $^1S_0$
  - Boron  $^2P_{1/2}, ^2P_{3/2}$
  - Carbon  $^1S_0, ^3S_1, ^1P_1, ^3P_0, ^3P_1, ^3P_2, ^1D_2, ^3D_1, ^3D_2, ^3D_3$
  - Nitrogen  $^2S_{1/2}, ^4S_{3/2}, ^2P_{1/2}, ^2P_{3/2}, ^4P_{1/2}, ^4P_{3/2}, ^4P_{5/2}, ^2D_{3/2}, ^2D_{5/2}, ^4D_{1/2}, ^4D_{3/2}, ^4D_{5/2}, ^4D_{7/2}, ^2F_{5/2}, ^2F_{7/2}, ^4F_{3/2}, ^4F_{5/2}, ^4F_{7/2}, ^4F_{9/2}$

## 5.18

- (a) Orthohelium should have lower energy than parahelium.
- (b) Hund's first rule says that the ground state of carbon has  $S = 1$ . This is a triplet state and is thus symmetric. That means the spatial wavefunction needs to be antisymmetric. Hund's second rule suggests we should choose  $L = 2$ . If we go to the "top of the ladder" for  $L = 2$  we find  $|2, 2\rangle = |1, 1\rangle_1 |1, 1\rangle_2$  which is symmetric and thus combining it with the triplet state would violate the Pauli exclusion principle. We must choose  $L = 1$ .
- (c) Boron can be  $^2P_{1/2}$  or  $^2P_{3/2}$ . The  $2p$  subshell can contain up to 6 electrons so for boron it is less than half filled and thus the lowest energy level has  $J = |L - S| = |1 - \frac{1}{2}| = \frac{1}{2}$  or  $^2P_{1/2}$ .
- (d) For carbon the highest total spin is  $S = 1$  which is a triplet state and is thus symmetric. Top of the ladder for  $L = 0$  is  $|0, 0\rangle = |0, 0\rangle_1 |0, 0\rangle_2$  which is symmetric, for  $L = 1$  is  $|1, 1\rangle = |1, 1\rangle |0, 0\rangle$  which is antisymmetric, and  $L = 2$  is  $|1, 1\rangle_1 |1, 1\rangle_2$  which is symmetric. Thus  $L = 1$ . The  $2p$  subshell is less than half filled and thus the lowest energy level has  $J = |L - S| = |1 - 1| = 0$  or  $^3P_0$ .

### 5.19

$S = 2$ ,  $L = 6$ ,  $J = 8$ , the  $n = 1, 2$ , and 3 shells are filled, the  $n = 4$  and 5 shells are partially filled.

### 5.20

$$\begin{aligned} E_{\text{tot}} &= \frac{\hbar^2 (3\pi^2 Nd)^{5/3}}{10\pi^2 m} V^{-2/3} \\ &= \frac{3Nd}{5} E_F \\ \frac{E_{\text{tot}}}{Nd} &= \frac{3}{5} E_F \end{aligned}$$

### 5.21

(a)

$$\begin{aligned} \rho_{\text{electron}} &= \frac{N_A \rho_{\text{mass}}}{m_a V} \\ &= 8.50 \times 10^{28} \text{ m}^{-3} \\ E_F &= \frac{\hbar^2}{2m} (3\rho\pi^2)^{2/3} \\ &= 7.05 \text{ eV} \end{aligned}$$

(b)

$$\begin{aligned} E_F &= \frac{1}{2} m v^2 \\ v &= \sqrt{\frac{2E_F}{m}} \\ &= 1.57 \times 10^6 \text{ m/s} \end{aligned}$$

It is safe to assume the electrons in copper are nonrelativistic.

(c)

$$\begin{aligned} k_B T &= E_F \\ T &= \frac{E_F}{k_B} \\ &= 8.18 \times 10^4 \text{ K} \end{aligned}$$

(d)

$$\begin{aligned} P &= \frac{(3\pi^2)^{2/3} \hbar^2}{5m} \rho^{5/3} \\ &= 3.84 \times 10^{10} \text{ Pa} \end{aligned}$$

## 5.22

$$\begin{aligned}
 \rho_{\text{electron}} &= \frac{\rho_{\text{mass}}}{m_a} \\
 &= 1.64 \times 10^{28} \text{ m}^{-3} \\
 E_F &= \frac{\hbar^2}{2m} (3\rho_{\text{electron}}\pi^2)^{2/3} \\
 &= 4.28 \times 10^{-4} \text{ eV} \\
 k_B T_F &= E_F \\
 T_F &= \frac{E_F}{k_B} \\
 &= 4.96 \text{ K}
 \end{aligned}$$

## 5.23

$$\begin{aligned}
 P &= \frac{(3\pi^2)^{2/3} \hbar^2}{5m} \rho^{5/3} \\
 &= \frac{(3\pi^2)^{2/3} \hbar^2}{5m} \left( \frac{Nd}{V} \right)^{5/3} \\
 \frac{dP}{dV} &= -\frac{5}{3} \frac{(3\pi^2)^{2/3} \hbar^2}{5m} \frac{(Nd)^{5/3}}{V^{8/3}} \\
 &= -\frac{5}{3} \frac{P}{V} \\
 B &= -V \frac{dP}{dV} \\
 &= \frac{5}{3} P \\
 B_{\text{copper}} &= 6.4 \times 10^{10} \text{ Pa}
 \end{aligned}$$

## 5.24

(a)

$$\begin{aligned}
 \psi(x) &= A \sin(kx) + B \cos(kx) \\
 A \sin(ka) &= [e^{iqa} - \cos(ka)] B \\
 \psi(x) &= \frac{[e^{iqa} - \cos(ka)] B \sin(kx)}{\sin(ka)} + B \cos(kx) \\
 &= \frac{B e^{iqa}}{\sin(ka)} \{ \sin(kx) + e^{-iqa} \sin[k(a-x)] \}
 \end{aligned}$$

## 5.25

$$\begin{aligned}
 10 &= \beta \\
 &= \frac{m\alpha a}{\hbar^2} \\
 a^2 &= \frac{10\hbar^2}{m} \frac{a}{\alpha} \\
 a &\approx 8.73 \times 10^{-10} \text{ m} \\
 -1 &= \cos z + \beta \frac{\sin z}{z} \\
 z &= 2.62768 \\
 ka &= 2.62768 \\
 \frac{\sqrt{2mE}}{\hbar} &= \frac{2.62768}{a} \\
 E &= \frac{2.62768^2 \hbar^2}{2ma^2} \\
 &\approx 0.345 \text{ eV}
 \end{aligned}$$

## 5.29

(a) Each particle can be in one of three states, so  $3^3 = 27$ .

(b)

*aaa*

*aab*

*abb*

*aac*

*acc*

*bbb*

*bbc*

*bcc*

*ccc*

*abc*

10

(c) Fermions can't be in the same state, so 1.

5.30

$$\begin{aligned}
E_{n_x, n_y} &= \frac{\hbar^2 \pi^2}{2m} \left( \frac{n_x^2}{l_x^2} + \frac{n_y^2}{l_y^2} \right) \\
&= \frac{\hbar^2 k^2}{2m} \\
\frac{\pi^2}{l_x l_y} &= \frac{\pi^2}{A} \\
\frac{1}{4} (\pi k_F^2) &= \frac{Nd}{2} \left( \frac{\pi^2}{A} \right) \\
k_F &= \sqrt{2\sigma\pi} \\
E_F &= \frac{\hbar^2}{2m} k_F^2 \\
&= \frac{\pi \hbar^2 \sigma}{m}
\end{aligned}$$

5.32

(a)

$$\begin{aligned}
\frac{1}{8} \left( \frac{4}{3} \pi k_{\pm}^3 \right) &= N_{\pm} \left( \frac{\pi^3}{V} \right) \\
k_{\pm} &= \left( \frac{6\pi^2 N_{\pm}}{V} \right)^{1/3} \\
dE &= \frac{\hbar^2 k^2}{2m} \frac{V}{2\pi^2} k^2 dk \\
&= \frac{\hbar^2 V}{4\pi^2 m} k^4 dk \\
E_{\text{total}} &= \frac{\hbar^2 V}{4\pi^2 m} \left( \int_0^{k_+} k^4 dk + \int_0^{k_-} k^4 dk \right) \\
&= \frac{\hbar^2 V}{20\pi^2 m} \left[ \left( \frac{6\pi^2 N_+}{V} \right)^{5/3} + \left( \frac{6\pi^2 N_-}{V} \right)^{5/3} \right] \\
&= \frac{\hbar^2}{20\pi^2 m} [(6\pi^2 N_+)^{5/3} + (6\pi^2 N_-)^{5/3}] V^{-2/3}
\end{aligned}$$

If  $N_+ = N_- = N$  then

$$E_{\text{total}} = \frac{\hbar^2 (6\pi^2 N)^{5/3}}{10\pi^2 m} V^{-2/3}$$

## 6 Symmetries & Conservation Laws

### 6.1

- (a)  $\psi(x, y, z)$  mirrored across the  $xy$  plane becomes  $\psi(x, y, -z)$ , then rotated  $180^\circ$  around the  $z$  axis becomes  $\psi(-x, -y, -z)$ .
- (b) In spherical coordinates, mirroring across the  $xy$  plan turns  $\theta$  into  $\pi - \theta$  and rotating  $180^\circ$  around the  $z$  axis  $\phi$  becomes  $\phi + \pi$  so  $\psi(r, \theta, \phi)$  becomes  $\psi(r, \pi - \theta, \phi + \pi)$ .
- (c)

$$\begin{aligned}
 \hat{\Pi}\psi_{n\ell m}(r, \theta, \phi) &= \hat{\Pi}[R_{n\ell}(r)Y_\ell^m(\theta, \phi)] \\
 &= R_{n\ell}(r)Y_\ell^m(\pi - \theta, \phi + \pi) \\
 &= (-1)^\ell R_{n\ell}(r)Y_\ell^m(\theta, \phi) \\
 &= (-1)^\ell \psi_{n\ell m}(r, \theta, \phi)
 \end{aligned}$$

### 6.2

$$\begin{aligned}
 \hat{U}^\dagger &= \exp(i\hat{Q})^\dagger \\
 &= \left(1 + i\hat{Q} - \frac{1}{2!}\hat{Q}^2 - \frac{1}{3!}i\hat{Q}^3 + \frac{1}{4!}\hat{Q}^4 + \dots\right)^\dagger \\
 &= 1 - i\hat{Q}^\dagger - \frac{1}{2!}(\hat{Q}^\dagger)^2 + \frac{1}{3!}i(\hat{Q}^\dagger)^3 + \frac{1}{4!}(\hat{Q}^\dagger)^4 + \dots \\
 &= 1 - i\hat{Q} - \frac{1}{2!}\hat{Q}^2 + \frac{1}{3!}i\hat{Q}^3 + \frac{1}{4!}\hat{Q}^4 + \dots \\
 &= \exp(-i\hat{Q}) \\
 \hat{U}^\dagger \hat{U} &= \exp(-i\hat{Q}) \exp(i\hat{Q}) \\
 &= 1
 \end{aligned}$$

### 6.3

$$\begin{aligned}
 \hat{p}'f(x) &= \hat{T}(a)^\dagger \hat{p} \hat{T}(a)f(x) \\
 &= \hat{T}(-a) \left(-i\hbar \frac{d}{dx}\right) f(x-a) \\
 &= \hat{T}(-a)[-i\hbar f'(x-a)] \\
 &= -i\hbar f'(x) \\
 &= \hat{p}
 \end{aligned}$$

## 6.4

$$\begin{aligned}
\hat{Q}'(\hat{x}, \hat{p})f(x) &= \hat{T}(a)^\dagger \hat{Q}(\hat{x}, \hat{p}) \hat{T}(a)f(x) \\
&= \hat{T}(-a) \left( \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} a_{mn} \hat{x}^m \hat{p}^n \right) f(x-a) \\
&= \left( \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} a_{mn} (\hat{x}+a)^m \hat{p}^n \right) f(x) \\
&= \hat{Q}(\hat{x}+a, \hat{p})f(x)
\end{aligned}$$

## 6.5

$$\begin{aligned}
\psi(x) &= e^{iqx} u(x) \\
\psi(x-a) &= e^{-iqa} \psi(x) \\
e^{iq(x-a)} u(x-a) &= e^{-iqa} e^{iqx} u(x) \\
&= e^{iq(x-a)} u(x) \\
u(x-a) &= u(x)
\end{aligned}$$

## 6.7

(a)

$$\begin{aligned}
\hat{T}(a)\psi(x_1, x_2) &= \psi(x_1 - a, x_2 - a) \\
&= \sum_{n=0}^{\infty} \frac{1}{n!} (-a)^n \frac{d^n}{dx_1^n} \psi(x_1, x_2 - a) \\
&= \exp\left(-\frac{ia}{\hbar} \hat{p}_1\right) \psi(x_1, x_2 - a) \\
&= \exp\left(-\frac{ia}{\hbar} \hat{p}_1\right) \sum_{n=0}^{\infty} \frac{1}{n!} (-a)^n \frac{d^n}{dx_2^n} \psi(x_1, x_2) \\
&= \exp\left[-\frac{ia}{\hbar} (\hat{p}_1 + \hat{p}_2)\right] \psi(x_1, x_2) \\
&= \exp\left(-\frac{ia}{\hbar} \hat{P}\right) \psi(x_1, x_2)
\end{aligned}$$



## 6.8

(a)

$$\begin{aligned}
 \langle \phi | \hat{\Pi} \psi \rangle &= \int_{-\infty}^{\infty} \phi^*(x) \hat{\Pi} \psi(x) dx \\
 &= \int_{-\infty}^{\infty} \phi^*(x) \psi(-x) dx \\
 &\quad u = -x \\
 &\quad du = -dx \\
 \langle \phi | \hat{\Pi} \psi \rangle &= \int_{\infty}^{-\infty} \phi^*(-u) \psi(u) (-du) \\
 &= \int_{-\infty}^{\infty} \phi^*(-u) \psi(u) du \\
 &= \int_{-\infty}^{\infty} \hat{\Pi} \phi^*(u) \psi(u) du \\
 &= \langle \hat{\Pi} \phi | \psi \rangle \\
 \hat{\Pi} &= \hat{\Pi}^\dagger
 \end{aligned}$$

(b)

$$\begin{aligned}
 \hat{\Pi} f(x) &= \lambda f(x) \\
 f(-x) &= \lambda f(x) \\
 \hat{\Pi} f(-x) &= \lambda f(-x) \\
 f(x) &= \lambda f(-x) \\
 f(x) &= \lambda(\lambda f(x)) \\
 &= \lambda^2 f(x) \\
 \lambda &= \pm 1
 \end{aligned}$$

## 6.9

(a)

$$\begin{aligned}
 \hat{\Pi}^\dagger \hat{f} \hat{\Pi} &= \hat{f} \\
 \hat{f} \hat{\Pi} &= \hat{\Pi} \hat{f} \\
 [\hat{\Pi}, \hat{f}] &= \hat{\Pi} \hat{f} - \hat{f} \hat{\Pi} \\
 &= 0 \\
 \hat{\Pi}^\dagger \hat{g} \hat{\Pi} &= -\hat{g} \\
 \hat{g} \hat{\Pi} &= -\hat{\Pi} \hat{g} \\
 \{\hat{\Pi}, \hat{g}\} &= \hat{\Pi} \hat{g} + \hat{g} \hat{\Pi} \\
 &= 0
 \end{aligned}$$

(b)

$$\begin{aligned}
\hat{\Pi}^\dagger \hat{\mathbf{V}} \hat{\Pi} &= -\hat{\mathbf{V}} \\
\hat{\mathbf{V}} \hat{\Pi} &= -\hat{\Pi} \hat{\mathbf{V}} \\
\{\hat{\Pi}, \hat{\mathbf{V}}\} &= \hat{\Pi} \hat{\mathbf{V}} + \hat{\mathbf{V}} \hat{\Pi} \\
&= 0 \\
\hat{\Pi}^\dagger \hat{\mathbf{W}} \hat{\Pi} &= \hat{\mathbf{W}} \\
\hat{\mathbf{W}} \hat{\Pi} &= \hat{\Pi} \hat{\mathbf{W}} \\
[\hat{\Pi}, \hat{\mathbf{W}}] &= \hat{\Pi} \hat{\mathbf{W}} - \hat{\mathbf{W}} \hat{\Pi} \\
&= 0
\end{aligned}$$

### 6.10

$$\begin{aligned}
\hat{x}' |x\rangle &= \hat{\Pi}^\dagger \hat{x} \hat{\Pi} |x\rangle \\
&= \hat{\Pi}^\dagger \hat{x} |-x\rangle \\
&= -x \hat{\Pi}^\dagger |-x\rangle \\
&= -x \hat{\Pi} |-x\rangle \\
&= -x |x\rangle \\
&= -\hat{x} |x\rangle \\
\hat{x}' &= -\hat{x} \\
\hat{p}' |p\rangle &= \hat{\Pi}^\dagger \hat{p} \hat{\Pi} |p\rangle \\
&= \hat{\Pi}^\dagger \hat{p} |-p\rangle \\
&= -p \hat{\Pi}^\dagger |-p\rangle \\
&= -p \hat{\Pi} |-p\rangle \\
&= -p |p\rangle \\
&= -\hat{p} |p\rangle \\
\hat{p}' &= -\hat{p}
\end{aligned}$$

### 6.11

$$\begin{aligned}
\langle n' \ell' m' | \hat{\mathbf{L}} | n \ell m \rangle &= \langle n' \ell' m' | \hat{\Pi}^\dagger \hat{\mathbf{L}} \hat{\Pi} | n \ell m \rangle \\
&= \langle n' \ell' m' | (-1)^{\ell'} \hat{\mathbf{L}} (-1)^\ell | n \ell m \rangle \\
&= (-1)^{\ell+\ell'} \langle n' \ell' m' | \hat{\mathbf{L}} | n \ell m \rangle
\end{aligned}$$

So  $\langle n' \ell' m | \hat{\mathbf{L}} | n \ell m \rangle = 0$  if  $\ell + \ell'$  is odd.

### 6.12

$$\begin{aligned}
[\hat{\Pi}, \hat{S}_x] &= \begin{pmatrix} a & b \\ c & d \end{pmatrix} \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} - \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} a & b \\ c & d \end{pmatrix} \\
&= \frac{\hbar}{2} \left[ \begin{pmatrix} b & a \\ d & c \end{pmatrix} - \begin{pmatrix} c & d \\ a & b \end{pmatrix} \right] \\
&= \frac{\hbar}{2} \begin{pmatrix} (b-c) & (a-d) \\ (d-a) & (c-b) \end{pmatrix} \\
&= \mathbf{0} \\
a &= d \\
b &= c \\
[\hat{\Pi}, \hat{S}_y] &= \begin{pmatrix} a & b \\ c & d \end{pmatrix} \frac{i\hbar}{2} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} - \frac{i\hbar}{2} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} a & b \\ c & d \end{pmatrix} \\
&= \frac{i\hbar}{2} \left[ \begin{pmatrix} b & -a \\ d & -c \end{pmatrix} - \begin{pmatrix} -c & -d \\ a & b \end{pmatrix} \right] \\
&= \frac{i\hbar}{2} \begin{pmatrix} (b+c) & (d-a) \\ (d-a) & -(b+c) \end{pmatrix} \\
&= \mathbf{0} \\
a &= d \\
b+c &= 0 \\
b &= 0 \\
c &= 0 \\
\hat{\Pi} &= \pi \mathbf{I}
\end{aligned}$$

### 6.13

(a)

$$\langle \mathbf{p}_e \rangle = \langle \psi_{100} | \hat{\mathbf{p}}_e | \psi_{100} \rangle$$

$\ell + \ell' = 0 + 0 = 0$  is even so by Laporte's rule  $\langle \mathbf{p}_e \rangle = \mathbf{0}$ .

(b)

$$\begin{aligned}
\psi &= \frac{1}{\sqrt{2}}(\psi_{200} + \psi_{210}) \\
\langle \hat{\mathbf{p}}_e \rangle &= \langle \psi | \hat{\mathbf{p}}_e | \psi \rangle \\
&= \left\langle \frac{1}{\sqrt{2}}(\psi_{200} + \psi_{210}) \left| \hat{\mathbf{p}}_e \right| \frac{1}{\sqrt{2}}(\psi_{200} + \psi_{210}) \right\rangle \\
&= \frac{1}{2} (\langle \psi_{200} | \hat{\mathbf{p}}_e | \psi_{200} \rangle + \langle \psi_{200} | \hat{\mathbf{p}}_e | \psi_{210} \rangle + \langle \psi_{210} | \hat{\mathbf{p}}_e | \psi_{200} \rangle + \langle \psi_{210} | \hat{\mathbf{p}}_e | \psi_{210} \rangle) \\
&= \frac{1}{2} (\langle \psi_{200} | \hat{\mathbf{p}}_e | \psi_{210} \rangle + \langle \psi_{210} | \hat{\mathbf{p}}_e | \psi_{200} \rangle) \\
&= -\frac{e}{2} (\langle \psi_{200} | \hat{\mathbf{r}} | \psi_{210} \rangle + \langle \psi_{210} | \hat{\mathbf{r}} | \psi_{200} \rangle)
\end{aligned}$$

6.15

$$\begin{aligned}
\hat{f}' &= \hat{R}^\dagger \hat{f} \hat{R} \\
&= \exp \left( \frac{i\varphi}{\hbar} \hat{\mathbf{n}} \cdot \hat{\mathbf{L}} \right) \hat{f} \exp \left( -\frac{i\varphi}{\hbar} \hat{\mathbf{n}} \cdot \hat{\mathbf{L}} \right) \\
&\approx \left( 1 + \frac{i\delta}{\hbar} \hat{\mathbf{n}} \cdot \hat{\mathbf{L}} \right) \hat{f} \left( 1 - \frac{i\delta}{\hbar} \hat{\mathbf{n}} \cdot \hat{\mathbf{L}} \right) \\
&= \hat{f} + \frac{i\delta}{\hbar} \hat{\mathbf{n}} \cdot [\hat{\mathbf{L}}, \hat{f}] \\
&= \hat{f}
\end{aligned}$$

### 6.16

$$\begin{aligned}
[\hat{L}_i, \hat{V}_j] &= i\hbar\epsilon_{ijk}\hat{V}_k \\
\epsilon_{ijk} &= \begin{cases} 1 & ijk = 123, 231, 312 \\ -1 & ijk = 132, 213, 321 \\ 0 & \text{otherwise} \end{cases} \\
\hat{R}_{\hat{\mathbf{y}}}(\delta) &= \exp\left(-\frac{i\delta}{\hbar}\hat{\mathbf{y}} \cdot \hat{\mathbf{L}}\right) \\
&= 1 - \frac{i\delta}{\hbar}\hat{L}_y \\
\hat{\mathbf{V}}' &= \hat{R}_{\hat{\mathbf{y}}}(\delta)^\dagger \hat{\mathbf{V}} \hat{R}_{\hat{\mathbf{y}}}(\delta) \\
&= \left(1 + \frac{i\delta}{\hbar}\hat{L}_y\right) \hat{\mathbf{V}} \left(1 - \frac{i\delta}{\hbar}\hat{L}_y\right) \\
&\approx \hat{\mathbf{V}} + \frac{i\delta}{\hbar}[\hat{L}_y, \hat{\mathbf{V}}] \\
&= \begin{pmatrix} \hat{V}_x + \frac{i\delta}{\hbar}[\hat{L}_y, \hat{V}_x] \\ \hat{V}_y + \frac{i\delta}{\hbar}[\hat{L}_y, \hat{V}_y] \\ \hat{V}_z + \frac{i\delta}{\hbar}[\hat{L}_y, \hat{V}_z] \end{pmatrix} \\
&= \begin{pmatrix} \hat{V}_x + \delta\hat{V}_z \\ \hat{V}_y \\ \hat{V}_z - \delta\hat{V}_x \end{pmatrix} \\
&= \begin{pmatrix} 1 & 0 & \delta \\ 0 & 1 & 0 \\ -\delta & 0 & 1 \end{pmatrix} \hat{\mathbf{V}} \\
D &= \begin{pmatrix} \cos\varphi & 0 & \sin\varphi \\ 0 & 1 & 0 \\ -\sin\varphi & 0 & \cos\varphi \end{pmatrix}
\end{aligned}$$

### 6.18

(a)

$$\begin{aligned}
[\hat{\Pi}, \hat{T}(a)]f(x) &= \hat{\Pi}\hat{T}(a)f(x) - \hat{T}(a)\hat{\Pi}f(x) \\
&= \hat{\Pi}f(x-a) - \hat{T}(a)f(-x) \\
&= f(-x+a) - f(-x-a)
\end{aligned}$$

(b)

$$\begin{aligned}f_p(x) &= \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar} \\ \hat{\Pi} f_p(x) &= \frac{1}{\sqrt{2\pi\hbar}} e^{ip(-x)/\hbar} \\ &= f_{-p}(x)\end{aligned}$$

(c)

$$\begin{aligned}\hat{T}(a) \left[ \frac{1}{\sqrt{\pi\hbar}} \cos\left(\frac{px}{\hbar}\right) \right] &= \frac{1}{\sqrt{\pi\hbar}} \cos\left(\frac{p(x-a)}{\hbar}\right) \\ &= \frac{1}{\sqrt{\pi\hbar}} \cos\left(\frac{px}{\hbar} - \frac{pa}{\hbar}\right) \\ -\frac{pa}{\hbar} &= \arctan\left(-\frac{b}{a}\right) \\ \tan\left(-\frac{pa}{\hbar}\right) &= -\frac{b}{a} \\ b &= -a \tan\left(-\frac{pa}{\hbar}\right) \\ \frac{1}{\sqrt{\pi\hbar}} &= \sqrt{a^2 + b^2} \\ \frac{1}{\pi\hbar} &= a^2 + b^2 \\ &= a^2 + a^2 \tan^2\left(-\frac{pa}{\hbar}\right) \\ &= a^2 \left[ 1 + \tan^2\left(-\frac{pa}{\hbar}\right) \right] \\ &= a^2 \sec^2\left(-\frac{pa}{\hbar}\right) \\ a &= \frac{1}{\sqrt{\pi\hbar}} \cos\left(-\frac{pa}{\hbar}\right) \\ \frac{1}{\pi\hbar} &= \frac{1}{\pi\hbar} \cos^2\left(-\frac{pa}{\hbar}\right) + b^2 \\ b &= \frac{1}{\sqrt{\pi\hbar}} \sin\left(-\frac{pa}{\hbar}\right) \\ \hat{T}(a) \left[ \frac{1}{\sqrt{\pi\hbar}} \cos\left(\frac{px}{\hbar}\right) \right] &= \frac{1}{\sqrt{\pi\hbar}} \cos\left(-\frac{pa}{\hbar}\right) \cos\left(\frac{px}{\hbar}\right) \\ &\quad + \frac{1}{\sqrt{\pi\hbar}} \sin\left(-\frac{pa}{\hbar}\right) \sin\left(\frac{px}{\hbar}\right)\end{aligned}$$

## 6.19

(a)

$$\begin{aligned}
[\hat{L}_i, \hat{V}_j] &= i\hbar\epsilon_{ijk}\hat{V}_k \\
\epsilon_{ijk} &= \begin{cases} 1 & ijk = 123, 231, 312 \\ -1 & ijk = 132, 213, 321 \\ 0 & \text{otherwise} \end{cases} \\
\hat{V}_{\pm} &= \hat{V}_x \pm i\hat{V}_y \\
[\hat{L}_z, \hat{V}_{\pm}] &= [\hat{L}_z, \hat{V}_x \pm i\hat{V}_y] \\
&= [\hat{L}_z, \hat{V}_x] \pm i[\hat{L}_z, \hat{V}_y] \\
&= i\hbar\hat{V}_y \pm i(-i\hbar\hat{V}_x) \\
&= \hbar(\pm\hat{V}_x + i\hat{V}_y) \\
[\hat{L}_z, \hat{V}_+] &= \hbar\hat{V}_+ \\
[\hat{L}_z, \hat{V}_-] &= -\hbar\hat{V}_-
\end{aligned}$$

## 6.20

$$\begin{aligned}
0 &= \langle n'\ell'm' | [\hat{L}_-, \hat{f}] | n\ell m \rangle \\
&= \langle n'\ell'm' | \hat{L}_- \hat{f} | n\ell m \rangle - \langle n'\ell'm' | \hat{f} \hat{L}_- | n\ell m \rangle \\
&= A_{\ell'}^{m'} \langle n'\ell'(m'+1) | \hat{f} | n\ell m \rangle - B_{\ell}^m \langle n'\ell'm' | \hat{f} | n\ell(m-1) \rangle \\
\ell' &= \ell \\
m' &= m-1 \\
A_{\ell'}^{m'} &= \hbar\sqrt{\ell'(\ell'+1) - m'(m'+1)} \\
&= \hbar\sqrt{\ell(\ell+1) - m(m-1)} \\
B_{\ell}^m &= \hbar\sqrt{\ell(\ell+1) - m(m-1)} \\
\langle n'\ell m | \hat{f} | n\ell m \rangle &= \langle n'\ell(m-1) | \hat{f} | n\ell(m-1) \rangle \\
\langle n'\ell m | \hat{f} | n\ell m \rangle &= \langle n'\ell(m+1) | \hat{f} | n\ell(m+1) \rangle
\end{aligned}$$

## 6.21

$$\begin{aligned}\psi &= \frac{1}{\sqrt{2}}(\psi_{211} + \psi_{21-1}) \\ \langle \psi | r | \psi \rangle &= \frac{1}{2}(\langle \psi_{211} | + \langle \psi_{21-1} |) r (| \psi_{211} \rangle + | \psi_{21-1} \rangle) \\ &= \frac{1}{2}(\langle \psi_{211} | r | \psi_{211} \rangle + \langle \psi_{211} | r | \psi_{21-1} \rangle + \langle \psi_{21-1} | r | \psi_{211} \rangle + \langle \psi_{21-1} | r | \psi_{21-1} \rangle) \\ &= \frac{1}{2}(\langle \psi_{211} | r | \psi_{211} \rangle + \langle \psi_{21-1} | r | \psi_{21-1} \rangle) \\ &= \langle 21 || r || 21 \rangle \\ &= 30a^2\end{aligned}$$



## 6.22

(a)

$$[\hat{L}_i, \hat{V}_j] = i\hbar\epsilon_{ijk}\hat{V}_k$$

$$\epsilon_{ijk} = \begin{cases} 1 & ijk = 123, 231, 312 \\ -1 & ijk = 132, 213, 321 \\ 0 & \text{otherwise} \end{cases}$$

$$[\hat{L}_z, \hat{V}_z] = i\hbar\epsilon_{11k}\hat{V}_k \\ = 0$$

$$\begin{aligned} [\hat{L}_z, \hat{V}_\pm] &= [\hat{L}_z, \hat{V}_x \pm i\hat{V}_y] \\ &= [\hat{L}_z, \hat{V}_x] \pm i[\hat{L}_z, \hat{V}_y] \\ &= i\hbar\hat{V}_y \pm i(-i\hbar\hat{V}_x) \\ &= \hbar(\pm\hat{V}_x + i\hat{V}_y) \\ &= \pm\hbar\hat{V}_\pm \end{aligned}$$

$$\begin{aligned} [\hat{L}_\pm, \hat{V}_\pm] &= [\hat{L}_x \pm i\hat{L}_y, \hat{V}_x \pm i\hat{V}_y] \\ &= [\hat{L}_x \pm i\hat{L}_y, \hat{V}_x] \pm i[\hat{L}_x \pm i\hat{L}_y, \hat{V}_y] \\ &= [\hat{L}_x, \hat{V}_x] \pm i[\hat{L}_y, \hat{V}_x] \pm i[\hat{L}_x, \hat{V}_y] - [\hat{L}_y, \hat{V}_y] \\ &= \pm i(-i\hbar\hat{V}_z + i\hbar\hat{V}_z) \\ &= 0 \end{aligned}$$

$$\begin{aligned} [\hat{L}_\pm, \hat{V}_z] &= [\hat{L}_x \pm i\hat{L}_y, \hat{V}_z] \\ &= [\hat{L}_x, \hat{V}_z] \pm i[\hat{L}_y, \hat{V}_z] \\ &= -i\hbar\hat{V}_y \pm i(i\hbar\hat{V}_x) \\ &= \hbar(\mp\hat{V}_x - i\hat{V}_y) \\ &= \mp\hbar\hat{V}_\pm \end{aligned}$$

$$\begin{aligned} [\hat{L}_\pm, \hat{V}_\mp] &= [\hat{L}_x \pm i\hat{L}_y, \hat{V}_x \mp i\hat{V}_y] \\ &= [\hat{L}_x \pm i\hat{L}_y, \hat{V}_x] \mp i[\hat{L}_x \pm i\hat{L}_y, \hat{V}_y] \\ &= [\hat{L}_x, \hat{V}_x] \pm i[\hat{L}_y, \hat{V}_x] \mp i[\hat{L}_x, \hat{V}_y] + [\hat{L}_y, \hat{V}_y] \\ &= \pm i([\hat{L}_y, \hat{V}_x] - [\hat{L}_x, \hat{V}_y]) \\ &= \pm i(-i\hbar\hat{V}_z - i\hbar\hat{V}_z) \\ &= \pm 2\hbar\hat{V}_z \end{aligned}$$

(b)

$$\begin{aligned} 0 &= \langle n' \ell' m' | [\hat{L}_z, \hat{V}_z] | n \ell m \rangle \\ &= \langle n' \ell' m' | \hat{L}_z \hat{V}_z | n \ell m \rangle - \langle n' \ell' m' | \hat{V}_z \hat{L}_z | n \ell m \rangle \\ &= \hbar m' \langle n' \ell' m' | \hat{V}_z | n \ell m \rangle - \hbar m \langle n' \ell' m' | \hat{V}_z | n \ell m \rangle \\ &= (m' - m) \langle n' \ell' m' | \hat{V}_z | n \ell m \rangle \end{aligned}$$

so  $\langle n' \ell' m' | \hat{V}_z | n \ell m \rangle = 0$  unless  $m' = m$ .