

Advanced Engineering Mathematics Partial Differential Equations by Dennis G. Zill Notes

Chris Doble

November 2023

Contents

12 Orthogonal Functions and Fourier Series	1
12.1 Orthogonal Functions	1
12.2 Fourier Series	2
12.3 Fourier Cosine and Sine Series	3
12.4 Complex Fourier Series	4
12.5 Sturm-Liouville Problem	4
12.6 Bessel and Legendre Series	6
12.6.1 Fourier-Bessel Series	6
12.6.2 Fourier-Legendre Series	7

12 Orthogonal Functions and Fourier Series

12.1 Orthogonal Functions

- The **inner product** of two functions f_1 and f_2 on an interval $[a, b]$ is the number

$$(f_1, f_2) = \int_a^b f_1(x)f_2(x) dx.$$

- Two functions f_1 and f_2 are said to be orthogonal on an interval if $(f_1, f_2) = 0$.
- A set of real-valued functions $\{\phi_1(x), \phi_2(x), \dots, \phi_n(x)\}$ is said to be **orthogonal** on an interval if

$$(\phi_i, \phi_j) = 0 \text{ for } i \neq j.$$

- The **square norm** of a function is

$$||\phi_n(x)||^2 = (\phi_n, \phi_n)$$

and thus its **norm** is

$$||\phi_n(x)|| = \sqrt{(\phi_n, \phi_n)}.$$

- An **orthonormal set** of functions is an orthogonal set of functions that all have a norm of 1.
- An orthogonal set can be made into an orthonormal set by dividing each member by its norm.
- If $\{\phi_n(x)\}$ is an infinite orthogonal set of functions on an interval $[a, b]$ and $f(x)$ is an arbitrary function, then it's possible to determine a set of coefficients $c_n, n = 0, 1, 2, \dots$ such that

$$f(x) = \sum_{n=0}^{\infty} c_n \phi_n(x) = c_0 \phi_0(x) + c_1 \phi_1(x) + \dots + c_n \phi_n(x) + \dots$$

This is called an **orthogonal series expansion** of f or a **generalized Fourier series** where the coefficients are given by

$$c_n = \frac{(f, \phi_n)}{\|\phi_n\|^2}.$$

- A set of real-valued functions $\{\phi_n(x)\}$ is said to be **orthogonal with respect to a weight function** $w(x)$ on the interval $[a, b]$ if

$$\int_a^b w(x) \phi_m(x) \phi_n(x) dx = 0, \quad m \neq n.$$

12.2 Fourier Series

- The **Fourier series** of a function f defined on the interval $(-p, p)$ is given by

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left(a_n \cos \frac{n\pi}{p} x + b_n \sin \frac{n\pi}{p} x \right)$$

where

$$\begin{aligned} a_0 &= \frac{1}{p} \int_{-p}^p f(x) dx \\ a_n &= \frac{1}{p} \int_{-p}^p f(x) \cos \frac{n\pi}{p} x dx \\ b_n &= \frac{1}{p} \int_{-p}^p f(x) \sin \frac{n\pi}{p} x dx \end{aligned}$$

- At points of discontinuity in f , the Fourier series takes on the average of the values either side of it.
- The Fourier series of a function f gives a **periodic extension** of the function outside the interval $(-p, p)$.

12.3 Fourier Cosine and Sine Series

- A function f is said to be **even** if

$$f(-x) = f(x)$$

and **odd** if

$$f(-x) = -f(x).$$

- Even and odd functions have some interesting properties:
 - The product of two even functions is even.
 - The product of two odd functions is even.
 - The product of an even function and an odd function is odd.
 - If f is even, then $\int_{-a}^a f(x) dx = 2 \int_0^a f(x) dx$.
 - If f is odd, then $\int_{-a}^a f(x) dx = 0$.
- In light of this, if a function f is even its Fourier coefficients are

$$\begin{aligned} a_0 &= \frac{2}{p} \int_0^p f(x) dx \\ a_n &= \frac{2}{p} \int_0^p f(x) \cos \frac{n\pi}{p} x dx \\ b_n &= 0. \end{aligned}$$

The series consists of cosine terms and is called the **Fourier cosine series**.

- Similarly, if f is odd then

$$\begin{aligned} a_n &= 0, \quad n = 0, 1, 2, \dots \\ b_n &= \frac{2}{p} \int_0^p f(x) \sin \frac{n\pi}{p} x dx. \end{aligned}$$

The series consists of sine terms and is called the **Fourier sine series**.

- Sometimes a Fourier series “overshoots” the original value of the function near discontinuities. This is called the **Gibbs phenomenon**.
- Taking the Fourier cosine series of a function f over the interval $[0, L]$ effectively mirrors the function around the vertical axis.
- Taking the Fourier sine series of a function f over the interval $[0, L]$ effectively rotates it 180° around the origin.
- A particular solution for a nonhomogeneous differential equation with a periodic driving force can be found by taking the Fourier transform of the driving force then using the method of undetermined coefficients to determine the coefficients.

12.4 Complex Fourier Series

- The **complex Fourier series** of a function f defined on an interval $(-p, p)$ is given by

$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{in\pi x/p}$$

where

$$c_n = \frac{1}{2p} \int_{-p}^p f(x) e^{-in\pi x/p} dx, \quad n = 0, \pm 1, \pm 2, \dots$$

- The **fundamental period** of a Fourier series is $T = 2p$.
- The **fundamental angular frequency** of a Fourier series is $\omega = \frac{2\pi}{T}$.
- A **frequency spectrum** is a plot of the points $(n\omega, |c_n|)$ where ω is the fundamental angular frequency and c_n are the coefficients of the complex Fourier series. This can be useful to see how each harmonic contributes.

12.5 Sturm-Liouville Problem

- If a boundary value problem contains an arbitrary parameter λ , the values of λ for which the problem has nontrivial solutions are called the **eigenvalues** of the problem and the associated solutions are called the **eigenfunctions** of the problem.
- An orthogonal set of functions can be generated by solving a two-point boundary-value problem involving a linear second-order differential equation containing a parameter λ .
- A **regular Sturm-Liouville problem** is a boundary value problem

$$\frac{d}{dx}[r(x)y'] + [q(x) + \lambda p(x)]y = 0$$

subject to

$$\begin{aligned} A_1 y(a) + B_1 y'(a) &= 0 \\ A_2 y(b) + B_2 y'(b) &= 0 \end{aligned}$$

where p , q , r , and r' are real-valued functions continuous on an interval $[a, b]$, $r(x) > 0$ and $p(x) > 0$ for every x in that interval, the coefficients in the boundary conditions are real and independent of λ , A_1 and B_1 are not both zero, and A_2 and B_2 are not both zero.

- A boundary condition

$$A_1 y(a) + B_1 y'(a) = C$$

is said to be **homogeneous** if $C = 0$ and **nonhomogeneous** otherwise.

- A boundary-value problem consisting of a homogeneous differential equation and a homogeneous boundary condition is said to be homogeneous, otherwise it's nonhomogeneous.
- Multiple boundary conditions are said to be **separated** if each deals with values at a single point $x = a$ and **mixed** if each deals with values at multiple points $x = a, b, \dots$
- If a boundary-value problem can be identified as a Sturm-Liouville problem we know it has several properties:
 - There exists an infinite number of real eigenvalues that can be arranged in increasing order $\lambda_1 < \lambda_2 < \dots < \lambda_n < \dots$ such that $\lambda_n \rightarrow \infty$ as $n \rightarrow \infty$.
 - For each eigenvalue there is only one eigenfunction.
 - Eigenfunctions corresponding to different eigenvalues are linearly independent.
 - The set of eigenfunctions corresponding to the set of eigenvalues is orthogonal with respect to the weight function $p(x)$ on the interval $[a, b]$.
- If a Sturm-Liouville problem has $r(a) = 0$ and boundary conditions are specified at $x = b$, or $r(b) = 0$ and boundary conditions are specified at $x = a$, then it is called a **singular boundary-value problem**.
- If a Sturm-Liouville problem has $r(a) = r(b)$ and boundary conditions $y(a) = y(b)$, $y'(a) = y'(b)$, then it is called a **periodic boundary-value problem**.
- If the solutions to a singular or periodic boundary-value problem are bounded on the interval $[a, b]$ then the orthogonality relation holds.
- Any second-order linear differential equation

$$a(x)y'' + b(x)y' + [c(x) + \lambda d(x)]y = 0$$

can be transformed into a Sturm-Liouville problem providing the coefficients are continuous and $a(x) \neq 0$ on the interval of interest. This can be done by:

1. dividing by a ,
2. multiplying by the integrating factor $e^{\int (b/a) dx}$,
3. recognising that

$$e^{\int (b/a) dx} y'' + \frac{b}{a} e^{\int (b/a) dx} y' = \frac{d}{dx} \left[e^{\int (b/a) dx} y' \right],$$

4. and rewriting the equation as

$$\frac{d}{dx} \left[e^{\int (b/a) dx} y' \right] + \left(\frac{c}{a} e^{\int (b/a) dx} + \lambda \frac{d}{a} e^{\int (b/a) dx} \right) y = 0$$

which is the desired form and lets us recognise

$$\begin{aligned} r(x) &= e^{\int (b/a) dx} \\ q(x) &= \frac{c}{a} e^{\int (b/a) dx} \\ p(x) &= \frac{d}{a} e^{\int (b/a) dx}. \end{aligned}$$

12.6 Bessel and Legendre Series

12.6.1 Fourier-Bessel Series

Definition 12.6.1 Fourier-Bessel Series

The **Fourier-Bessel series** of a function f defined on the interval $(0, b)$ is given by

$$(i) \quad f(x) = \sum_{i=1}^{\infty} c_i J_n(\alpha_i x) \quad (15)$$

$$c_i = \frac{2}{b^2 J_{n+1}^2(\alpha_i b)} \int_0^b x J_n(\alpha_i x) f(x) dx, \quad (16)$$

where the α_i are defined by $J_n(\alpha b) = 0$.

$$(ii) \quad f(x) = \sum_{i=1}^{\infty} c_i J_n(\alpha_i x) \quad (17)$$

$$c_i = \frac{2\alpha_i^2}{(\alpha_i^2 b^2 - n^2 + h^2) J_n^2(\alpha_i b)} \int_0^b x J_n(\alpha_i x) f(x) dx, \quad (18)$$

where the α_i are defined by $h J_n(\alpha b) + \alpha b J_n'(\alpha b) = 0$.

$$(iii) \quad f(x) = c_1 + \sum_{i=2}^{\infty} c_i J_0(\alpha_i x) \quad (19)$$

$$c_1 = \frac{2}{b^2} \int_0^b x f(x) dx, \quad c_i = \frac{2}{b^2 J_0^2(\alpha_i b)} \int_0^b x J_0(\alpha_i x) f(x) dx, \quad (20)$$

where the α_i are defined by $J_0'(\alpha b) = 0$.

- The Fourier-Bessel series converges to f where it is continuous and

$$\frac{f(x+) + f(x-)}{2}$$

where it is discontinuous.

12.6.2 Fourier-Legendre Series

Definition 12.6.2 Fourier-Legendre Series

The **Fourier-Legendre series** of a function f defined on the interval $(-1, 1)$ is given by

$$f(x) = \sum_{n=0}^{\infty} c_n P_n(x), \quad (21)$$

where
$$c_n = \frac{2n+1}{2} \int_{-1}^1 f(x) P_n(x) dx. \quad (22)$$

- The Fourier-Legendre series converges to f where it is continuous and

$$\frac{f(x+) + f(x-)}{2}$$

where it is discontinuous.