0.1 Ideal Gas Model

From (McQuarrie, Statistical mechanics, 2000) the ideal gas method (IGM) for calculating an absolute free energy is outlined below.

$$G = H - TS \tag{1}$$

$$H = U + RT \tag{2}$$

$$U = E_{\text{pot}} + E_{\text{ZPE}} + E_{\text{trns}} + E_{\text{rot}} + E_{\text{vib}}$$
 (3)

$$S = S_{\text{trns}} + S_{\text{rot}} + S_{\text{vib}} + S_{\text{elec}} \tag{4}$$

where T is temperature, R the ideal gas constant and $S_{\rm elec}$ is taken to be zero for all molecules. The internal energy components are then

$$E_{\rm ZPE} = \frac{N_a}{2} \sum_{i} h \nu_i \tag{5}$$

$$E_{\rm trns} = \frac{3}{2}RT\tag{6}$$

$$E_{\rm rot} = \begin{cases} 0 & \text{if } N = 1\\ RT & \text{if linear}\\ \frac{3}{2}RT & \text{otherwise} \end{cases}$$
 (7)

$$E_{\rm vib} = R \sum_{i} \frac{\theta_i}{e^{\theta_i/T} - 1} \quad , \quad \theta_i = h\nu_i/k_B$$
 (8)

where N_a is Avogadro's's constant, N is the number of atoms in the molecule, k_B Boltzmann's constant, ν_i the *i*-th harmonic frequency and h is Planks constant. The entropic components are

$$S_{\rm trns} = R \ln(q_{\rm trns}) + \frac{5}{2}R \tag{9}$$

$$S_{\text{rot}} = \begin{cases} 0 & \text{if } N = 1\\ R \ln(q_{\text{rot}}) + R & \text{if linear}\\ R \ln(q_{\text{rot}}) + \frac{3}{2}R & \text{otherwise} \end{cases}$$
(10)

$$S_{\text{vib}}^{\text{HO}} = R \sum_{i} \frac{\theta_{i}}{T(e^{\theta_{i}/T} - 1)} - \ln(1 - e^{-\theta_{i}/T})$$
 (11)

$$q_{\rm trans} = \left(\frac{2\pi m k_B T}{h^2}\right)^{3/2} V_{\rm eff} \quad , \quad V_{\rm eff} = \begin{cases} k_B T/p^{\circ} & \quad \text{if 1 atm standard state} \\ c^{\circ}/N_a & \quad \text{if 1 M standard state} \end{cases}$$

 $q_{\rm rot} = \frac{T^{3/2}}{\sigma_r} \sqrt{\frac{\pi}{\omega_r}} \quad , \quad \omega_r = \prod_k \frac{h^2}{8\pi^2 k_B I_k}$ (12)

where q are molecular partition functions, p° is the standard pressure (1 atm) and c° the standard concentration (1 mol dm⁻³), σ_r is the rotational symmetry number for the molecule and I_k a diagonal element of the moment of inertia matrix.

Due to the vibrational entropy contribution being overestimated for low frequency modes Thrular proposed a correction, which instead of summing over frequencies in $S_{\rm vib}^{\rm HO}$ does so over $\max(\nu_{\rm thresh}, \nu_i)$ to shift all low frequencies to a threshold value (*J. Phys. Chem. B* 2011, **115**, 14556). An alternative method from Grimme (*Chem. Eur. J.*, 2012, **18**, 9955) uses an interpolation between a harmonic oscillator and rigid rotor to scale down the contribution from the low frequency modes as

$$S_{\text{vib}}^{\text{Grimme}} = \sum_{i} w_i S_{\text{vib}}^{\text{HO}}(i) + (1 - w_i) \left(R \ln \left(\sqrt{\frac{8\pi^3 \mu_i' k_B T}{h^2}} \right) + \frac{R}{2} \right)$$
(14)

$$\mu_i' = \frac{\mu_i \bar{B}}{\mu_i + \bar{B}} \quad , \quad \mu_i = \frac{h}{8\pi^2 \nu_i} \quad , \quad \bar{B} = \text{Tr}[I]/3$$
 (15)

$$w_i = \frac{1}{1 + (\omega_0/\nu_i)^{\alpha}} \tag{16}$$

where ω_0 and α are adjustable parameters.