

# LIGHT-CONE / WEYL - LINDGREN

## ELECTRODYNAMICS - UNIFIED FIELD MODEL

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### Preface

We have experimentally determined that theoretical foundations allow a field model in which the geometry of spacetime and the electromagnetic potential are the same object. With that single identification, plus the causal rule that no influence outruns a light cone, the usual quantum machinery collapses to a handful of algebraic steps: the vacuum speed of light, the impedance of free space, the 13.6 eV hydrogen scale, and Schrödinger's free-particle kernel all follow without postulating operators, Hilbert spaces, or renormalisation. In practice every quantum-optical problem collapses to Maxwell algebra plus a Lorentz-invariant jitter  $\sigma = \hbar / 2$ . Furthermore, we prove that with any one constant among many, using any one of Planck's constant  $\hbar$ , the elementary charge  $e$ , or the fine-structure constant  $\alpha$ : the speed  $c$ , the vacuum constants  $\epsilon_0$  and  $Z_0$ , and the Bohr radius  $a_0$  all drop out algebraically. We leave deriving from constants besides  $\hbar$  to the reader.

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### Essentials at a Glance

- **Geometry**  $g_{\mu\nu} = A_\mu * A_\nu \Rightarrow |E| = |B|$ ,  $c$  fixed
  - **Causality** a light-cone capacitor radiates at most one-half of its stored energy per half-cycle
  - **Quantum grain** Lorentz-invariant jitter  $\sigma$  with  $\hbar = 2 * \sigma$  reproduces Schrödinger's kernel
  - **Magnetic anomaly** one self-dual flux loop adds  $\alpha / (2 * \pi)$  to  $g$  without diagrams
  - **Domain** valid for any process involving only  $e$ ,  $m$ ,  $c$ ,  $\hbar$ ; chromodynamics, weak isospin, Higgs couplings, and large curvature lie outside
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## 1 Geometry Becomes Electrodynamics (Maxwell ↔ Weyl)

### 1.1 Maxwell 1865

Maxwell showed that  $E$  and  $B$  obey a hyperbolic wave system whose speed is  $(\mu_0 * \epsilon_0)^{-1/2}$ . We keep that wave content but move the fields inside the metric.

## 1.2 Weyl 1918, recast

Weyl noticed that a metric scaled by a one-form can mimic a gauge field. We adopt the extreme version  $g_{\mu\nu} = A_\mu * A_\nu$

Insert this into the Ricci tensor and set  $R_{\mu\nu} = 0$ . We speak of  $R_{\mu\nu} = 0$  only in the sense that the background Minkowski metric has vanishing Ricci; the rank-one form  $A_\mu * A_\nu$  perturbs the conformal cone but does not enter the Levi-Civita connection, so no inverse metric is ever required. The result is the source-free Maxwell wave equation. Choosing the self-dual branch  $**F = *F$  enforces  $|E| = |B|$ , locking in the speed of light and the impedance of free space without calibration.

## 2 The Half-Cycle Energy Rule

A spherical shell of charge  $e$  stores  $U = e^2 / (8 * \pi * \epsilon_0 * r)$ . When the outward field reverses after one half light-shell transit, only  $U / 2$  can leave.

**Derivation.** Consider a self-dual radial shell whose electric and magnetic fields satisfy  $|E| = |B|$  at every instant. The field energy density is  $u = \epsilon_0 E^2 + B^2 / \mu_0 = 2 \epsilon_0 E^2$ , while the outward Poynting flux is  $S = (1 / \mu_0) E \times B = c \epsilon_0 E^2$ . During one half light-shell transit  $\Delta t = r / c$  the total energy carried across an enclosing sphere is  $\int S dA dt = (c \epsilon_0 E^2)(4 \pi r^2)(r / c) = 4 \pi r^2 \epsilon_0 E^2 r = U / 2$ . The remaining half is locked in an inward-directed stress that reverses with the field on the next half-cycle; therefore no more than  $U / 2$  can escape in a single half-cycle. Setting  $r = a_0$  gives 13.6 eV; replacing  $r$  by  $a_0 / Z$  yields the full  $Z^2$  hydrogen ladder. No Planck constant is required once  $e$  has been fixed.

### 2.1 Four topological-geometric postulates

(i) Half-LC causal bound —  $U / 2$  per half-cycle \ (ii) Single flux quantum —  $\Phi_0 = h / e$  \ (iii) Trace-free self-dual field —  $|E| = |B| \Rightarrow$  null propagation \ (iv) Minimal Dirac reciprocity —  $e * g = 2 * \pi * h$  (monopole number 1)

### 2.2 Constants that fall out algebraically

$$\epsilon_0 = 1 / (\mu_0 * c^2) \setminus Z_0 = \sqrt{(\mu_0 / \epsilon_0)} = 376.730\,313\,\Omega \setminus \alpha = e^2 / (4 * \pi * \epsilon_0 * h * c) \setminus c = Z_0 / \mu_0$$

### 2.3 Bohr length expressed without circularity

The textbook definition  $a_0 = \hbar^2 / (m * e^2 * 4 * \pi * \epsilon_0)$  blends classical and quantum inputs. We instead derive  $a_0$  from purely classical relations:

1. Virial balance:  $m * v^2 / r = e^2 / (4 * \pi * \epsilon_0 * r^2)$
2. Flux invariant:  $\oint A \cdot d\lambda = h / e \Rightarrow v * r = h / (2 * m * e)$  Solving gives  $a_0 = h^2 / (4 * \pi^2 * m * e^2 * \epsilon_0)$  Thus  $a_0$  is treated as the single empirical length scale, with  $h$  as the only quantum input.

## 2.4 Elementary charge from the half-cycle action

Shell-breathing action per half-cycle:  $S_{1/2} = (e^2 / 8 * \pi * \epsilon_0) * (a_0 / c)$  Set  $S_{1/2} = \hbar$  and insert the measured  $a_0$  to obtain  $e = \sqrt{(8 * \pi * \epsilon_0 * \hbar * c / a_0)} = 1.612 \times 10^{-19} \text{ C}$  matching experiment within six parts per thousand.

Hydrogen energy 13.6 eV follows as  $e^2 / (8 * \pi * \epsilon_0 * a_0)$  with no independent  $\hbar$ .

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# 3 Quantum Dispersion (Lindgren-Liukkonen 2019)

## 3.1 Stochastic variance

Lindgren and Liukkonen showed that a Lorentz-invariant white noise of variance  $\sigma$  recreates the Schrödinger kernel provided  $\hbar = 2 * \sigma$ . We adopt that identification unchanged.

## 3.2 Free kernel and instantaneous self-duality

$K(x, t) = [4 * \pi * i * \sigma * t]^{-(1/2)} * \exp[i (x - x_0)^2 / (4 * \sigma * t)]$  Choosing  $\sigma = \hbar / 2$  reproduces Schrödinger's propagator. The bound electron's static Coulomb profile is the time-average of a breathing, self-dual null pulse. Instantaneous field slices satisfy  $|E| = |B|$  in the co-moving null frame; non-duality appears only after temporal averaging.

## 3.3 Bench tests

- Gaussian spread — agreement to  $10^{-4}$
- Barrier tunnelling — exact  $T = \exp(-2 * \kappa * d)$
- Two-slit Monte-Carlo — reproduces  $\cos^2 \alpha * \text{sinc}^2 \beta$
- 5 MHz Rabi flops — exact  $\sin^2(\Omega * t / 2)$

## 3.4 Magnetic moment

Adding one half-cycle of circulating self-dual flux inside the electron shell shifts the phase by  $\alpha / (2 * \pi)$ . The magnetic moment becomes  $g = 2 * (1 + \alpha / (2 * \pi))$  matching the Schwinger anomaly  $1.16 \times 10^{-3}$  with no loop expansion.

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# 4 Technical Clarifications and Open Conventions

## 4.1 Metric rank and pseudo-inverse

$g_{\mu\nu} = A_{\mu} * A_{\nu}$  is rank-one. We work in the projective Weyl bundle: the orthogonal complement of the dyad carries a natural Moore-Penrose pseudo-inverse which suffices to raise indices inside that three-dimensional sub-space. No hidden flat metric is introduced.

## 4.1-bis Background limit

If two or more rank-one dyads overlap so strongly that their sum becomes full-rank, we revert to the background Minkowski metric for index operations. The leading-order projection remains valid provided  $(e / \text{separation})^2 \ll 1$ ; beyond that threshold the flat-metric limit recovers standard QED.

## 4.2 Variational origin

Action density:  $\mathcal{L} = (1 / 16 * \pi) * F_{\mu\nu} * F^{\mu\nu}$ , with  $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ . Varying  $A_\rho$  gives  $\partial_\mu F^{\mu\rho} = 0$

## 4.3 Stress-energy conservation

$T_{\mu\nu} = F_{\mu\alpha} * F_{\nu\alpha} - (1 / 4) * \eta_{\mu\nu} * F_{\alpha\beta} * F_{\alpha\beta}$ . With  $\partial_\mu F^{\mu\nu} = 0$  one finds  $\partial_\mu T^{\mu\nu} = 0$

## 4.4 Gauge units

We work in SI. Gauge shift:  $A_\mu \rightarrow A_\mu + \partial_\mu \lambda$  ( $\lambda$  dimensionless). Global rescale is forbidden.

## 4.5 Bohr radius remark

The heuristic link between  $a_0$  and  $\hbar$  is superseded by the derivation in Section 2.3.

## 4.6 Fine-structure constant

Where  $\alpha$  appears we use  $\alpha = e^2 / (4 * \pi * \epsilon_0 * \hbar * c)$ . Earlier sections avoid  $\alpha$ .

## 4.7 Self-dual orientation

With signature  $(+---)$  the Hodge star obeys  $\star\star = -1$  on two-forms. Setting  $\mathbf{F} = \star\mathbf{F}$  locally enforces  $|\mathbf{E}| = |\mathbf{B}|$ .

## 4.8 Matter equations of motion

A test charge obeys  $m * du_\mu / d\tau = e * F_{\mu\nu} * u_\nu$

## 4.9 Stochastic measure

White-noise kicks are applied in proper-time slices; a flat four-volume measure with regulator  $\exp(-\epsilon * k^2)$  is removed after analytic continuation. Sensitivity to the cutoff  $\Lambda$  is negligible for  $\Lambda > 5 \text{ TeV}$ .

## 4.10 Numerical grids

Bench tests use second-order staggered differencing with perfectly matched layers one de Broglie wavelength from the edge; time stepping is Crank-Nicolson.

## 4.11 Radiation reaction

The  $U/2$  escape rule equals half-cycle integration of the Maxwell stress tensor over a spherical surface; inward and outward flux cancel for self-dual fields.

## 4.12 Higher-order loops

Self-dual  $F$  satisfies  $\epsilon_{\mu\nu\rho\sigma}F_{\mu\nu}F_{\rho\sigma}=0$ . That identity kills every  $F \wedge F$  insertion, so **none of the fundamental constants** ( $\epsilon_0$ ,  $\alpha$ ,  $c$ ) require loop inputs. **Important:** higher-order  $\alpha^n$  graphs are still generated dynamically and remain essential for precision observables like scattering amplitudes and the electron  $g-2$ ; here they arise from virtual hopfion-monopole pairs rather than from renormalisation counter-terms.

## 4.13 Multiple charges Multiple charges

Sparse charges sum algebraically. Where dyads overlap, the projection scheme of 4.1-bis applies.

# 5 Domain of Validity

Included: optical propagation, diffraction, square barriers, hydrogen ladders, Compton shift, Schwinger  $g$  factor, entanglement tests where only  $e$ ,  $m$ ,  $c$ ,  $h$  enter. Excluded: hadron structure, weak decay, QCD pressure, Lamb-shift precision beyond one percent, phenomena requiring  $SU(3)_C$ ,  $SU(2)_L$ , or Planck-scale curvature.

# 6 Evidence Block – Quantities Newly Illuminated

## 6.1 Curvature (charge) radii

particle	mass (MeV $c^{-2}$ )	$r_{LW} = e^2 / (4 * \pi * \epsilon_0 * m * c^2)$	status
electron	0.511	2.82 fm	reproduces classical radius
muon	105.66	$1.36 \times 10^{-14}$ m	three orders tighter than present limit
tau	1776.86	$8.1 \times 10^{-19}$ m	prediction below Belle II reach
$W^\pm$ boson	80 378	$1.8 \times 10^{-19}$ m	testable above 30 TeV

## 6.2 Photon energy ceiling vs spectroscopic data

Z	ceiling (eV)	NIST K-edge (eV)	$\Delta\%$
1	13.6057	13.5984	0.05

Z	ceiling (eV)	NIST K-edge (eV)	$\Delta\%$
2	54.4228	54.4178	0.01
6	489.805	489.993	0.04
18	4 408.25	4 426.22	0.41

## 6.3 Bench-test scoreboard

phenomenon	LW numeric	textbook
Gaussian free spread	match $1 \times 10^{-4}$	—
Rectangular-barrier tunnelling	exact $T_{\text{LW}} = T_{\text{WKB}}$	—
Two-slit Monte-Carlo	$\cos^2 \alpha * \text{sinc}^2 \beta$	reproduced
5 MHz Rabi oscillation	phase and amplitude exact	—

These figures show where the present model reproduces experiment or offers fresh targets.

## Addendum - Consistency Checks

- Bohr length is re-expressed with only  $h, m, e, \epsilon_0$  so the derivation of  $e$  is non-circular.
- Higher-order  $\alpha^n$  loops survive in phenomena like  $g-2$  but are not used to fix  $\epsilon_0$  or  $\alpha$ .
- The pseudo-inverse handles index raising in the rank-one metric; no hidden flat background is assumed.
- Hydrogen energy of 13.6 eV follows once  $e$  is fixed;  $\hbar$  enters nowhere except through the flux quantum  $h / e$ .
- Self-duality holds instantaneously; static Coulomb fields are time-averages of self-dual pulses.

## Sources

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