

**1a**

$$y'_o - y_o = 0 \quad \Rightarrow \quad y_o = A_o e^x.$$

Fulfilling both  $y_o(0) = 0$  and  $y_o(1) = 1$  is impossible  $\Rightarrow$  boundary layer.  
Layer at  $x = 0$  then

$$y_o = e^{x-1}$$

**1b**

Transformation  $\xi = x/\delta$

$$\begin{aligned} \frac{\epsilon}{\delta^2} Y'' + \xi Y' - \delta \xi Y &= 0 \\ (1) \quad (2) \quad (3) \end{aligned}$$

Clear (3)  $\ll$  (2)  $\Rightarrow$  (1) & (2) must be dominant balance and  $\delta = \sqrt{\epsilon}$

$$Y'' + \xi Y' = 0 \quad \Rightarrow \quad Y' = A e^{-\frac{1}{2}\xi^2}.$$

From the formula

$$Y = Y(0) + \int_0^\xi Y'(\hat{\xi}) d\hat{\xi} = A \int_0^\xi e^{-\frac{1}{2}\hat{\xi}^2} d\hat{\xi} = \sqrt{\frac{\pi}{2}} A \operatorname{erf}\left(\frac{\xi}{\sqrt{2}}\right)$$

Matching

$$\begin{aligned} y_{\text{match}} &= \lim_{\xi \rightarrow \infty} Y = \sqrt{\frac{\pi}{2}} A, \\ y_{\text{match}} &= \lim_{x \rightarrow 0} y_o = e^{-1}, \end{aligned}$$

Then  $A = \sqrt{2}/(e\sqrt{\pi})$  and

$$y_{\text{unif}} = y_o + Y - y_{\text{match}} = e^{-1} \left( e^x + \operatorname{erf}\left(\frac{x}{\sqrt{2\epsilon}}\right) - 1 \right).$$

**2a**

$$\frac{d^2 y}{dt^2} = \frac{\partial^2 y}{\partial t^2} + 2\epsilon \frac{\partial^2 y}{\partial \tau \partial t} + O(\epsilon^2)$$

$$y = y_0 + \epsilon y_1 \dots$$

$$\epsilon^0$$

$$\frac{\partial^2 y_0}{\partial t^2} + \omega^2 y_0 = 0, \quad y_0(0, 0) = 1, \quad \frac{\partial y_0(0, 0)}{\partial t} = 0.$$

Solution

$$y_0 = A(\tau) e^{i\omega(\tau)t} + c.c., \quad A(0) = \frac{1}{2}.$$

$\epsilon^1$

$$\frac{\partial^2 y_1}{\partial t^2} + \omega^2 y_1 = -2 \frac{\partial^2 y_0}{\partial \tau \partial t} = -(2i\omega A' - 2\omega A \omega' t) e^{i\omega(\tau)t} + c.c.$$

where  $A' = \frac{dA}{d\tau}$  etc.

Only way to avoid terms  $\sim t^2 e^{i\omega(\tau)t}$  in  $y_1$  is to require  $A \equiv 0$ , which yields nothing.

## 2b

Now  $y = y(T, \tau)$ .

$$\frac{d^2 y}{dt^2} = \sigma^2 \frac{\partial^2 y}{\partial T^2} + \epsilon \left( 2\sigma \frac{\partial^2 y}{\partial \tau \partial T} + \frac{d\sigma}{d\tau} \frac{\partial y}{\partial T} \right) + O(\epsilon^2)$$

$$y = y_0 + \epsilon y_1 \dots$$

$\epsilon^0$

$$\sigma^2 \frac{\partial^2 y_0}{\partial T^2} + \omega^2 y_0 = 0, \quad y_0(0, 0) = 1, \quad \frac{\partial y_0(0, 0)}{\partial T} = 0.$$

To avoid a slow coefficient in the exponent, which leads to secular terms in next order, we put  $\sigma = \omega$  and

$$y_0 = A(\tau) e^{iT} + c.c., \quad A(0) = \frac{1}{2}.$$

$\epsilon^1$

$$\omega^2 \frac{\partial^2 y_1}{\partial T^2} + \omega^2 y_1 = -2\omega \frac{\partial^2 y_0}{\partial \tau \partial T} - \omega' \frac{\partial y_0}{\partial \tau} = -i(2\omega A' + A\omega') e^{iT} + c.c.$$

Avoid secular terms:  $2\omega A' + A\omega' = 0 \Rightarrow A^2 \omega = \text{const.}$

Use of initial condition

$$A = \frac{1}{2} \sqrt{\frac{\omega(0)}{\omega(\tau)}}.$$

## 3a

$$\pi_1 = h\omega^2/g, \quad \pi_2 = kh.$$

$$\pi_1 = F(\pi_2) \Rightarrow$$

$$\omega^2 = \frac{g}{h} F(kh) = gkG(kh).$$

## 3b

Observe:  $F(kh) = kh \tanh kh$ .

$$kh \tanh kh = \frac{h\omega^2}{g} \rightarrow \infty$$

Since  $\tanh kh < 1$  must have  $kh \rightarrow \infty$ .

$\tanh kh \rightarrow 1$  as  $kh \rightarrow \infty \Rightarrow k_0 h = \frac{h\omega^2}{g}$  and  $k_0 = \frac{\omega^2}{g}$

### 3c

Work instead with  $\kappa = kh$  ( $\pi_2$ ) and  $\Omega = h\omega^2/g$  ( $\pi_1$ )

$$\Omega = \kappa \tanh \kappa.$$

Large  $\kappa \Rightarrow e^{-2\kappa} \ll 1$

$$\tanh \kappa = \frac{1 - e^{-2\kappa}}{1 + e^{-2\kappa}} = 1 - 2e^{-2\kappa} + O(e^{-4\kappa})$$

$\kappa = \kappa_0 + \kappa_1$  with  $\kappa_0 \gg \kappa_1$

$$\Omega = \kappa \tanh \kappa = \kappa \left( 1 - 2e^{-2\kappa} + O(e^{-4\kappa}) \right)$$

and possible dominant terms

$$\begin{array}{ccccccc} \Omega & = & \kappa_0 & + \kappa_1 & - 2\kappa_0 e^{-2(\kappa_0 + \kappa_1)} \\ (1) & & (2) & (3) & (4) \end{array}$$

Dominant balance: (1) & (2), secondary balance (3) & (4).

Moreover, guess that  $e^{-2(\kappa_0 + \kappa_1)} \approx e^{-2\kappa_0}$ . Must require  $\kappa_1 \ll 1$  ( $\kappa_1 \ll \kappa_0$  is not sufficient). Then

$$\kappa_1 = 2\kappa_0 e^{-2\kappa_0},$$

and indeed  $\kappa_1 \ll 1$ .

Some may attempt something like

$$\Omega = (\kappa_0 + \kappa_1) \tanh(\kappa_0 + \kappa_1) \approx (\kappa_0 + \kappa_1) \tanh \kappa_0 = \kappa_0 + \kappa_1 + (\tanh(\kappa_0) - 1)(\kappa_0 + \kappa_1) \approx \kappa_0 + \kappa_1 + (\tanh(\kappa_0) - 1)\kappa_0$$

Removing leading order we then obtain

$$\kappa_1 = -(\tanh(\kappa_0) - 1)\kappa_0.$$

This is by no means wrong, but less transparent. Still  $\kappa_1 \ll 1$  needs to be addressed.

### 4a

### 4b

Integration by parts  $\Rightarrow$

$$\delta J = \left( \frac{\partial L}{\partial y'} \delta y \right)_{x=b} + \int_a^b \left\{ \frac{\partial L}{\partial y} - \frac{d}{dx} \left( \frac{\partial L}{\partial y'} \right) \right\} \delta y dx.$$

$\delta J = 0$  for all  $\delta y$  that are zero at the ends yields the Euler equation. Then the integral vanishes and  $\delta y(b) \neq 0 \Rightarrow$

$$\left( \frac{\partial L}{\partial y'} \right)_{x=b} = 0$$

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