ASYMPTOTISK BØLGEFRONT

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September 18, 2014

STASJONÆR FASE

Fourier integral

$$\eta(x,t) = \frac{1}{2\pi} \Re \int_{0}^{\infty} \tilde{\eta}_{0} e^{i\chi} dk$$
$$\chi \equiv kx - \omega(k)t$$

Stasjonært punkt

$$\frac{\mathrm{d}\chi(k_s)}{\mathrm{d}k}=0$$

som gir

$$c_g(k_s) = \frac{x}{t}$$

Dominant bidrag rundt k_s

ENDELIG DYP

Begrensning stasj. fase

$$c_g \le c_0 \equiv \sqrt{gH} \Rightarrow \text{Stasjon} \text{ær fase for } \frac{x}{t} \le c_0$$

 $\frac{x}{t} \sim c_0 \Rightarrow \text{bølgefront}$

Nær bølgefront

Bidrag fra $k \to 0$.

$$\chi(0) = \chi''(0) = 0 \Rightarrow$$

$$\chi \approx kx - c_0 \left(k - \frac{H^2}{6} k^3 \right) t$$

og

$$\eta(x,t) \sim \frac{1}{2\pi} \Re \int_{0}^{\infty} \tilde{\eta}_{0}(0) e^{i\left(kx - (c_{0}k - \frac{H^{2}}{6}c_{0}k^{3})t\right)} dk$$
(1)

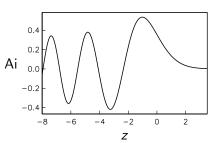
AIRY FUNKSJONEN

Integral i (1) uttrykkes ved:

$$I(z) = \frac{1}{\pi} \int_{0}^{\infty} \cos\left(\frac{1}{3}s^{3} + zs\right) ds$$

Integralformel for Airy funksjonen Ai som løser

$$\frac{\mathrm{d}^2 F}{\mathrm{d}z^2} - zF = 0 \tag{2}$$



BØLGEFRONT

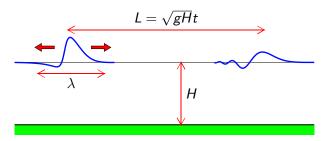
Nær $x = c_0 t$

$$\eta \sim \frac{\frac{1}{2}V}{(\frac{1}{2}c_0H^2t)^{\frac{1}{3}}} \text{Ai} \left(\frac{x - c_0t}{(\frac{1}{2}c_0H^2t)^{\frac{1}{3}}} \right)$$

 $\text{der } V = \tilde{\eta}(0) \text{ er volumet under } \eta(x,0)$

- Første topp nær $x = c_0 t$
- Form uavhengig av $\eta(x,0)$
- Bølgehøyde $\sim t^{-\frac{1}{3}}$
- Lengde øker $\sim t^{\frac{1}{3}}$

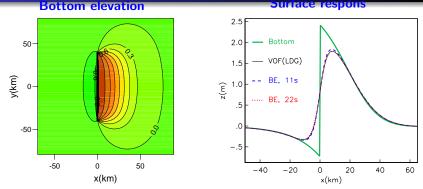
Wave front and dispersion in tsunami propagation



Dispersion often neglected for tsunamis due to large wavelength. Its significance depends on:

- Extent of source relative to depth.
- Propagation distance.

Exampel inspired by earthquake off Portugal (1969)

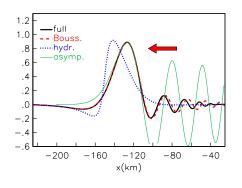


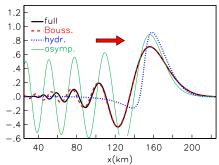
Magnitude: $M_s=7.9$, $H=5000\,\mathrm{m}$, inverse thrust fault, large dip angle $\approx 50^\circ$, fault length $\approx 70\,\mathrm{km} \Rightarrow$ rather confined bottom uplift Left panel: co-seismic bottom-uplift from Okada's formula Right panel: hydrodynamic response for centerline from 2D theories

2D response used as initial condition.



After $t = 11.3 \Rightarrow L = c_0 t = 150 \,\mathrm{k} m$





Curves:

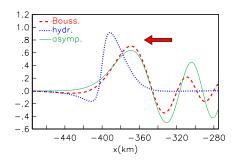
'full' and Boussinesq: numerical solutions with dispersion

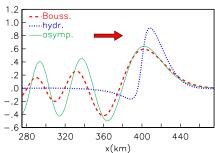
'hydr': is half the initial elevation

'asymp': Asymptotic solution for wave front



After $t = 30 \text{ min} \Rightarrow L = 400 \text{km}$





Now asymptote and numerical solution are becoming close.