## MEK4320, solution to extra problem 13.

Geir Pedersen

Department of Mathematics, UiO

May 15, 2020

4 D > 4 B > 4 E > 4 E > E 90

Geir Pedersen

Problem

## Introduction

When indices are used for differentiation the LSW equations read

$$\eta_t = -(hu)_{\times}, \quad u_t = -g\eta_{\times}. \tag{1}$$

In this problem we will study energy within the framework of linear shallow water theory. Hence, we shall use properties like hydrostatic pressure and vertically uniform  $\boldsymbol{u}$  whenever appropriate.

4 m > 4 m >

Geir Pedersen

Problems

# a Energy denity and flux.

## Question

Show, by direct calculation, that the energy density and flux become, respectively,

$$E = \frac{1}{2}\rho hu^2 + \frac{1}{2}\rho g\eta^2, \quad F = \rho gh\eta u.$$

## Solution

Density per volume of kinetic energy:  $\frac{1}{2}\rho u^2$ .

Here we have neglected the contribution from the vertical component in accordance with LSW theory.

Density per horizontal area (utilization of u = u(x, t)).

$$E_k = \int_{-h}^{\eta} \frac{1}{2} \rho u^2 dz = \frac{1}{2} \rho u^2 \int_{-h}^{\eta} dz = \frac{1}{2} \rho u^2 (h + \eta) \approx \frac{1}{2} \rho h u^2,$$

where the cubic term is deleted.

Geir Pedersen

← ← → ← 付き → ← き → ・ き → ぞ

Corresponding integral for potential energy

$$E_p = \int_{-h}^{\eta} \rho g z dz = \frac{1}{2} \rho g \eta^2 - \frac{1}{2} \rho g h^2,$$

where the last term is the equilibrium contribution. Omitting this we find

$$E = \frac{1}{2}\rho hu^2 + \frac{1}{2}\rho g\eta^2.$$

## The flux

The flux of kinetic energy is ignored since it is cubic in the field variables.

Integrating advection of potential energy and effect of pressure work, using  $p=\rho g(\eta-z)$  in accordance with LSW theory

$$F = \int_{-h}^{\eta} (\rho u + \rho g z u) dz = \int_{-h}^{\eta} \rho g \eta u dz = \rho g \eta u (h + \eta) \approx \rho g h \eta u.$$

Geir Pedersen

Problem

# b The wave mode.

## Problem text

Assume constant depth and find a wave mode solution to (1). To this end assume a form

$$\eta = A\cos(kx - \omega t), \quad u = U\cos(kx - \omega t).$$

## Solution

With constant depth (1) gives

$$\eta_t = -hu_x, \quad u_t = -g\eta_x.$$

Substitution of mode into this

 $\omega A \sin(kx - \omega t) = hkU \sin(kx - \omega t), \quad \omega U \sin(kx - \omega t) = gkA \sin(kx - \omega t).$ 

Deletion of common factors and re-organizing give

$$\frac{A}{U} = \frac{hk}{\omega}, \quad \frac{A}{U} = \frac{\omega}{gk}. \quad \Rightarrow \quad \omega^2 = ghk^2, \quad \frac{U}{A} = \frac{c}{h}.$$

Geir Pedersen

Problems

# c Averaged densities for mode.

## Problem text

Use the solution of the previous point to show that

$$\overline{E} = \frac{1}{2} \rho g A^2, \quad \overline{F} = \sqrt{gh} \, \overline{E}.$$

## Solution

By inserting the mode we find

$$E = \frac{1}{2}\rho hu^2 + \frac{1}{2}\rho g\eta^2 = \rho gA^2 \cos^2(kx - \omega t)$$
  
$$F = \rho gh\eta u = \rho gcA^2 \cos^2(kx - \omega t)$$

The average of  $\cos^2$  over a wavelength or period is  $\frac{1}{2}$ . Hence

$$\overline{E} = \frac{1}{2}\rho gA^2$$

$$\overline{F} = \frac{1}{2}\rho gcA^2 = c\overline{E}$$

where  $c = c_g = \sqrt{gh}$ 

ir Pedersen

Problem

# d Fulfillment of energy equation.

## Problem text

The energy equation reads

$$E_t + F_x = 0. (2)$$

Show that this is fulfilled by invoking (1).

Differentiation with careful use of product rule gives

$$E_t + F_x = \rho h u u_t + \rho g \eta \eta_t + \rho g \eta (h u)_x + \rho g h u \eta_x$$
(i) (ii) (iii) (iv)

Continuity eq.,  $\eta_t = -(hu)_x$ , causes (ii) and (iii) to cancel out. Momentum eq.,  $u_t = -g\eta_x$ , causes (i) and (iv) to cancel out. Hence  $E_t + F_x = 0$ .



From preceding slide

$$(\frac{1}{2}\rho hu^2)_t = -F_x - \rho g\eta\eta_t = -F_x - (\frac{1}{2}\rho\eta^2)_t.$$

The term within the last parantheses is the potential energy. Moving this to the other side we obtain

$$E_t = -F_x$$
.

# e Derivation of the energy equation from the PDE's.

### Problem text

Start with (1) and derive (2).

Multiply the momentum equation with ho hu (same as multiply by hoand integrate over flow depth)

$$\rho huu_t = -\rho ghu\eta_x.$$

Left hand side is time derivative of  $E_k$ . Squinting at the expression for F we rewrite the right hand side

$$(\frac{1}{2}\rho hu^2)_t = -(\rho ghu\eta)_x + \rho g\eta(hu)_x.$$

First term on rhs. is  $-F_x$ . The latter term is rewritten by the

$$(\frac{1}{2}\rho hu^2)_t = -F_x - \rho g\eta\eta_t = -F_x - (\frac{1}{2}\rho g\eta^2)_t.$$