

**Ex. 22 .**

a)

$$\begin{array}{c|c|c|c|c|c} x & m & t & C & k & x_0 \\ \hline L & M & T & \frac{M}{TL^2} & \frac{M}{T^2} & L \end{array}$$

$\Pi$  theorem, part 1: 6 parameters, 3 dimensions  $\Rightarrow$  3 dimensionless numbers. Simplest options:

Make the first, only, with  $x_*$ . The simplest option:  $\pi_1 = \frac{x_*}{x_0}$

Make the next, only, with  $t_*$ . Use also  $m$  and  $k$  (not  $x$ 'es):  $\pi_2 = \frac{kt_*^2}{m}$

The last is the only made with  $C$ . Avoid  $x_*$  and  $t_*$ . Then the theorem tells us that there is on  $\pi$  between  $C$ ,  $m$ ,  $x_0$  and  $k$ :  $\pi_3 = \frac{Cx_0^2}{\sqrt{mk}}$ .

b) We scale the problem according to

$$x_* = x_c x, \quad t_* = t_c t.$$

From the transformation of the initial condition from  $x_*(0) = x_0$  to  $x(0) = 1$  we find  $x_c = x_0$ . Then both initial conditions are fulfilled regardless of  $t_c$ . Using  $x_c = x_0$ , inserting the transformation in the ODE and reorganize the result such that the coefficient of the second derivative becomes unity we obtain

$$\frac{d^2x}{dt^2} + \frac{Cx_0^2 t_c}{m} x^2 \frac{dx}{dt} + \frac{t_c^2 k}{m} x = 0.$$

To reproduce the dimensionless ODE, given in the problem text, we must have

$$\frac{Cx_0^2 t_c}{m} = \epsilon, \quad \frac{t_c^2 k}{m} = 1,$$

which imply

$$t_c = \sqrt{\frac{m}{k}}, \quad \epsilon = \frac{Cx_0^2}{\sqrt{mk}}.$$

$t_c$  expresses the period of the undamped oscillator.  $\epsilon$  increases with  $C$  and  $x_0$  and decreases with  $m$  and  $k$ , as can be expected.

$x$ ,  $s$  and  $\epsilon$  are dimensionless numbers which may expressed by  $\pi_1$ ,  $\pi_2$  and  $\pi_3$ . In fact:

$$\epsilon = \pi_3, \quad x = \pi_1 \quad \text{and} \quad t = \sqrt{\pi_2}.$$

If we had obtained different  $\pi$ 's in sub-problem a, the above relations would also be different.

c) Introduction of  $\tau$ , such that  $x = x(t, \tau)$ , yields the set

$$\begin{aligned} \frac{\partial^2 x}{\partial t^2} + 2\epsilon \frac{\partial^2 x}{\partial t \partial \tau} + \epsilon x^2 \frac{dx}{dt} + x &= O(\epsilon^2), \\ x(0, 0) &= 1, \quad \frac{\partial x(0, 0)}{\partial t} + \epsilon \frac{\partial x(0, 0)}{\partial \tau} = 0. \end{aligned}$$

The expansion  $x = x_0(s, \tau) + \epsilon x_1(t, \tau)$  is then inserted.

Order  $\epsilon^0$

$$\begin{aligned}\frac{\partial^2 x_0}{\partial t^2} + x_0 &= O, \\ x_0(0, 0) &= 1, \quad \frac{\partial x_0(0, 0)}{\partial t} = 0.\end{aligned}$$

The solution is

$$x_0 = A(\tau)e^{it} + A^*(\tau)e^{-it} = A(\tau)e^{it} + c.c., \quad A(0) = \frac{1}{2}.$$

Order  $\epsilon^1$

$$\begin{aligned}\frac{\partial^2 x_1}{\partial t^2} + x_1 &= -2\frac{\partial^2 x_0}{\partial t \partial \tau} - x_0^2 \frac{dx_0}{dt}, \\ x_1(0, 0) &= 0, \quad \frac{\partial x_1(0, 0)}{\partial t} = -\frac{\partial x_0(0, 0)}{\partial \tau}.\end{aligned}$$

Inserting the expression for  $x_0$  on the right hand side of the ODE we arrive at

$$\frac{\partial^2 x_1}{\partial t^2} + x_1 = iA^3 e^{3it} - i\left(2\frac{dA}{d\tau} + A^2 A^*\right)e^{it} + c.c.$$

The solution must be damped. Hence, we do not allow  $x_1$  to grow linearly in  $t$ . Then, the  $e^{it}$  and  $e^{-it}$  parts of the right hand side must vanish, which implies

$$2\frac{dA}{d\tau} + A^2 A^* = 0.$$

Substituting the polar form  $A = ae^{i\psi}$ , where  $a$  and  $\psi$  are real, into this we obtain

$$\frac{da}{d\tau} + \frac{1}{2}a^3 = 0, \quad \frac{d\psi}{d\tau} = 0.$$

The first one is a separable equation

$$-2a^{-3}\frac{da}{d\tau} = 1 \quad \Rightarrow \quad a^{-2} = \tau + B \quad \Rightarrow \quad a = (B + \tau)^{-\frac{1}{2}}.$$

The second immediately yields  $\psi = D = \text{const.}$  The initial condition for  $A$  now implies  $B = 4$  and  $D = 0$ . Then

$$A = (4 + \tau)^{-\frac{1}{2}}.$$

Assembling the leading order solution we then find

$$A = \frac{1}{\sqrt{1 + \frac{\epsilon t}{4}}} \cos(t).$$