

# Introduction to Multigrid Methods

## ***Chapter 8: Elements of Multigrid Methods***

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Textbooks: *A Multigrid Tutorial*, by William L Briggs. SIAM 1988

*A First Course in the Numerical Analysis of Differential Equations*, by Arieh Iserles. Cambridge 1996

*Matrix-based multigrid: Theory and Applications*, by Yair Shapira. Springer 2008

*Multi-Grid Methods and Applications*, by Wolfgang Hackbusch, 1985

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# 1. Successive over-relaxation (SOR)

An attempt to speed up convergence in Gauss–Seidel

$$Du = Du - \omega(Au - f)$$

$$Du = Du - \omega(D - L - U)u + \omega f$$

$$(D - \omega L)u = (\omega U + (1 - \omega)D)u + \omega f$$

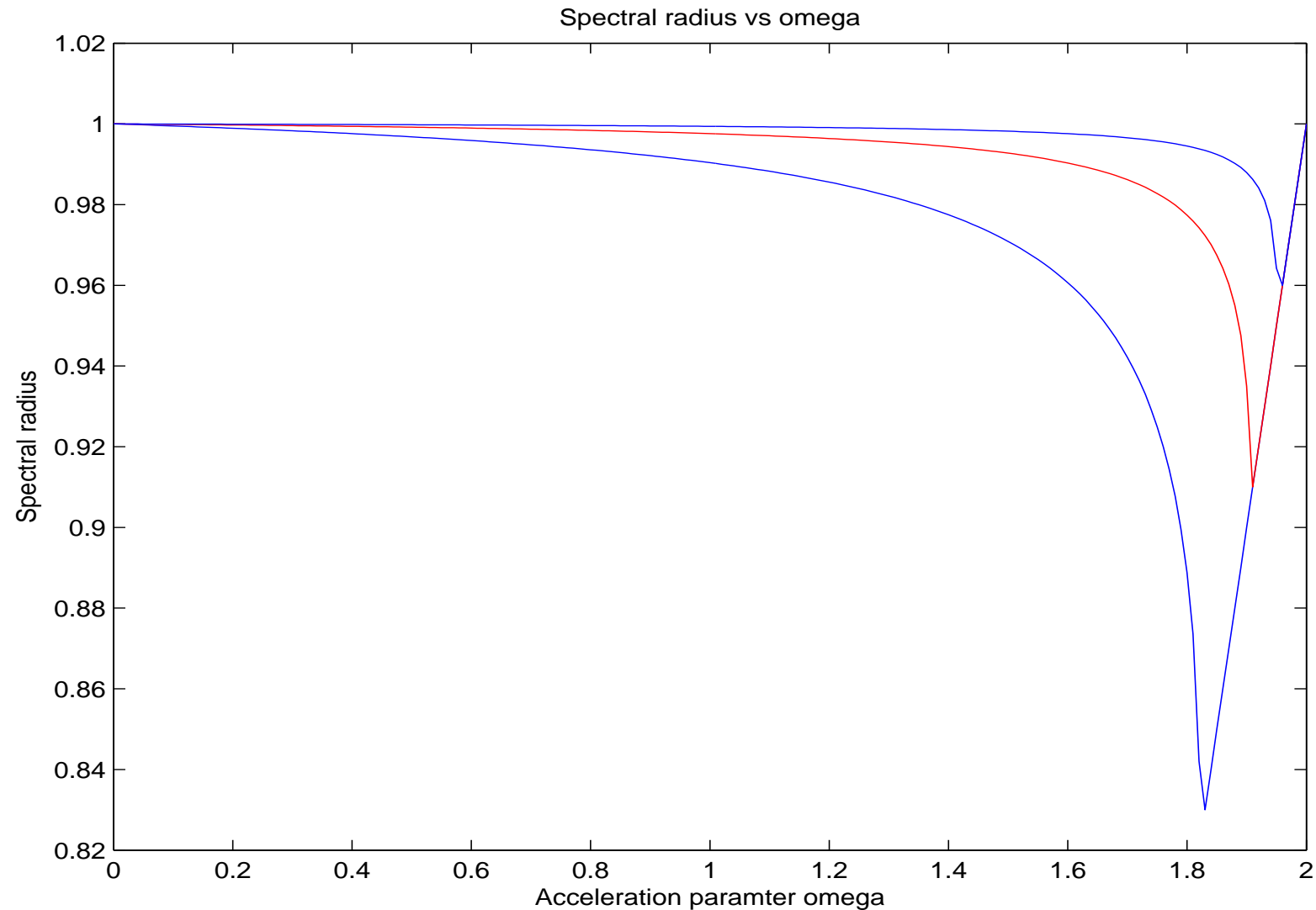
$$(D - \omega L)u^{m+1} = (\omega U + (1 - \omega)D)u^m + \omega f$$

**Iteration matrix**  $P_{SOR} = (D - \omega L)^{-1}((1 - \omega)D + \omega U)$

**Acceleration parameter**  $\omega = 1 \Rightarrow$  Gauss–Seidel method

# Spectral radius of SOR matrix for $T_{\Delta x}$

$\rho[P_{SOR}(\omega)]$  for  $N = 31, 63, 127$



# Faster convergence is possible!

For 1D Poisson, the spectral radius is minimized at

$$\hat{\omega} = \frac{2}{1 + \sqrt{1 - \rho^2[P_J]}}$$

As  $\rho[P_J] = \cos \pi \Delta x \approx 1 - \pi^2 \Delta x^2 / 2$  we have

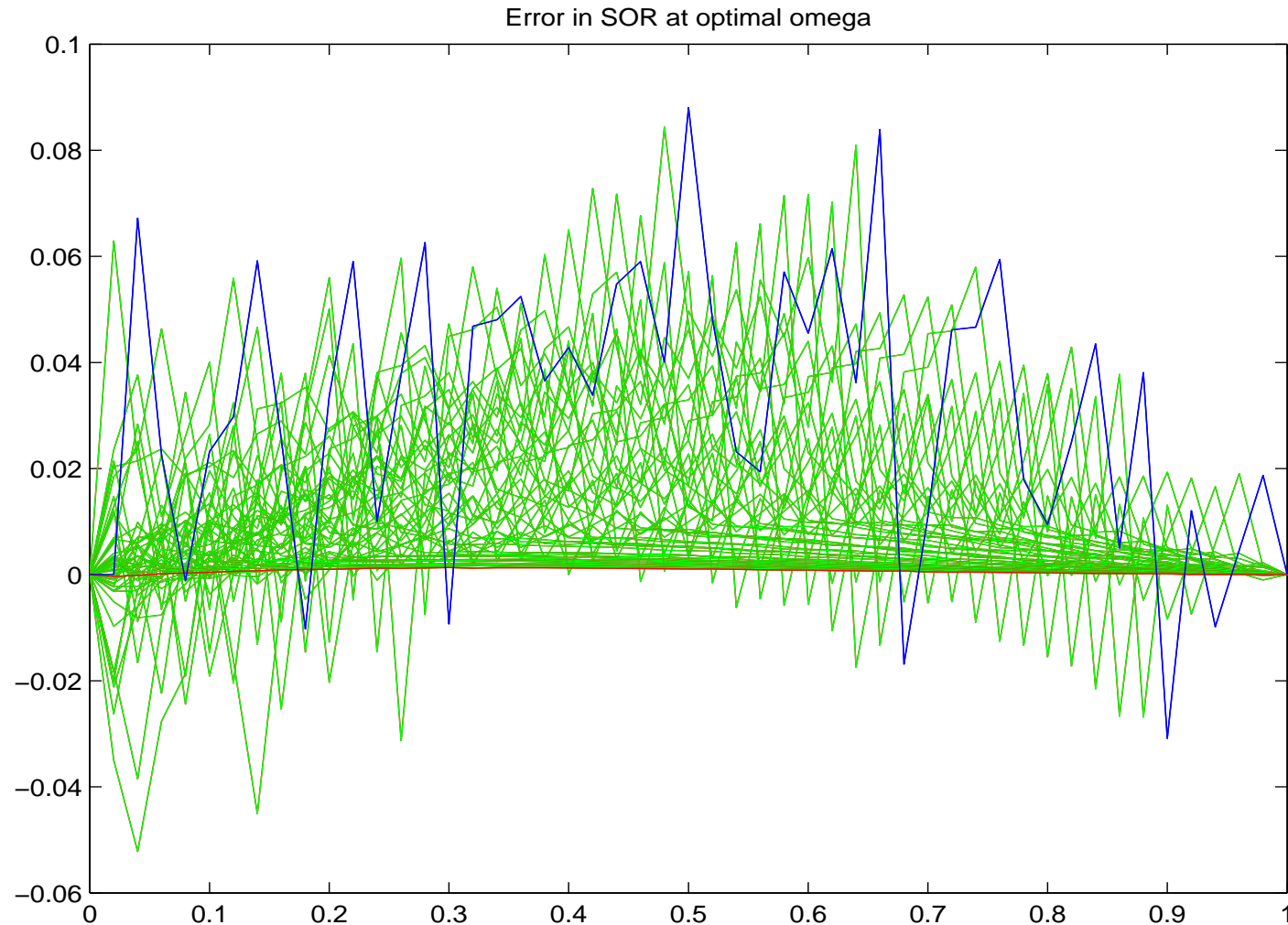
$$\hat{\omega} \approx 2 - 2\pi \Delta x$$

At this value the spectral radius of the SOR matrix is

$$\rho[P_{SOR}(\hat{\omega})] = \hat{\omega} - 1 \approx 1 - 2\pi \Delta x$$

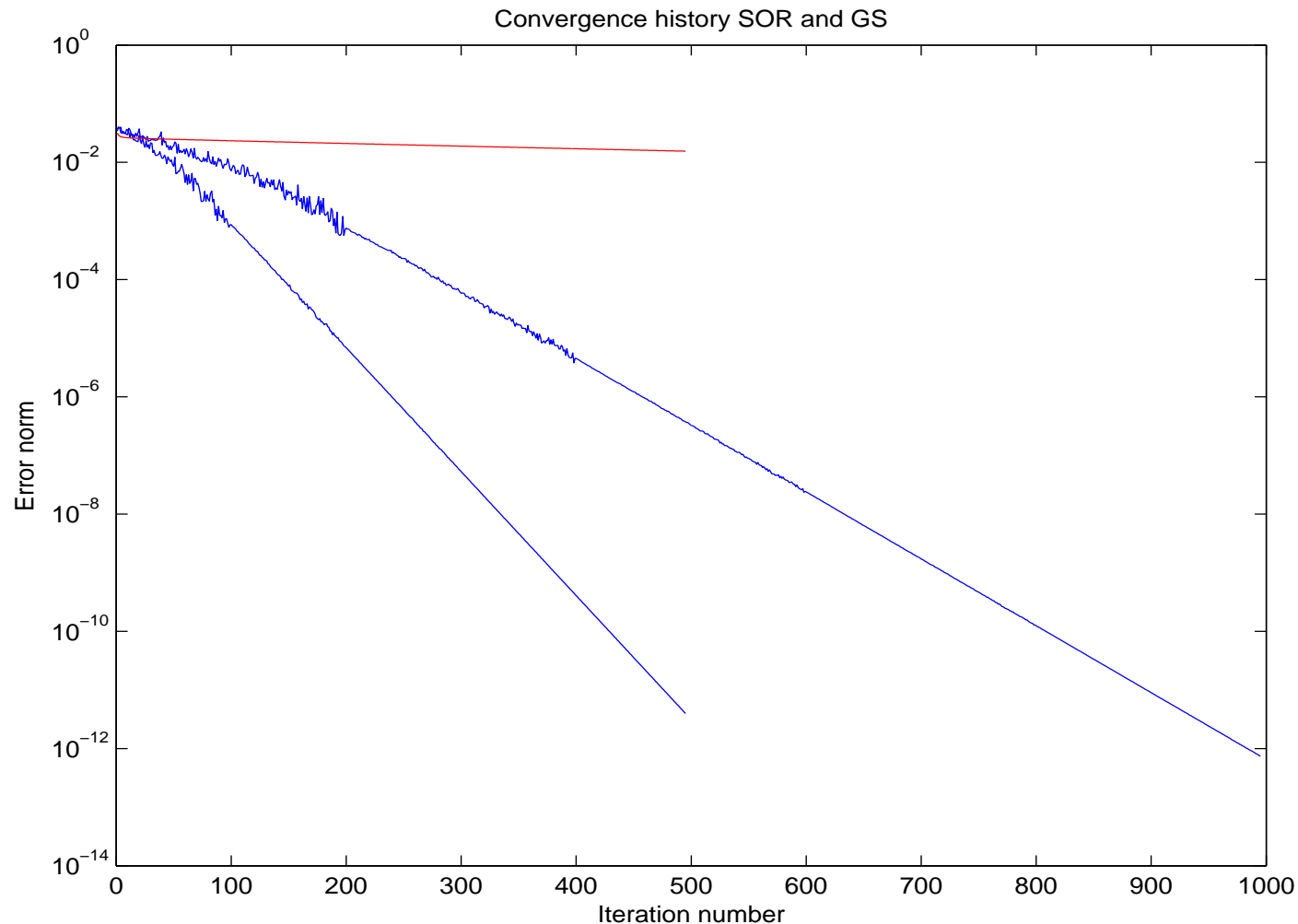
Much improved over Jacobi and Gauss–Seidel

# SOR error sequence simulation



**Note** Error “transportation” in live demo!

# SOR error sequence simulation, $N = 99,199$



**Note** Unit error reduction requires only  $O(N)$  iterations

## 2. Symmetric SOR (SSOR)

Neutralize “error transportation” to the left by alternating between forward and backward sweeps

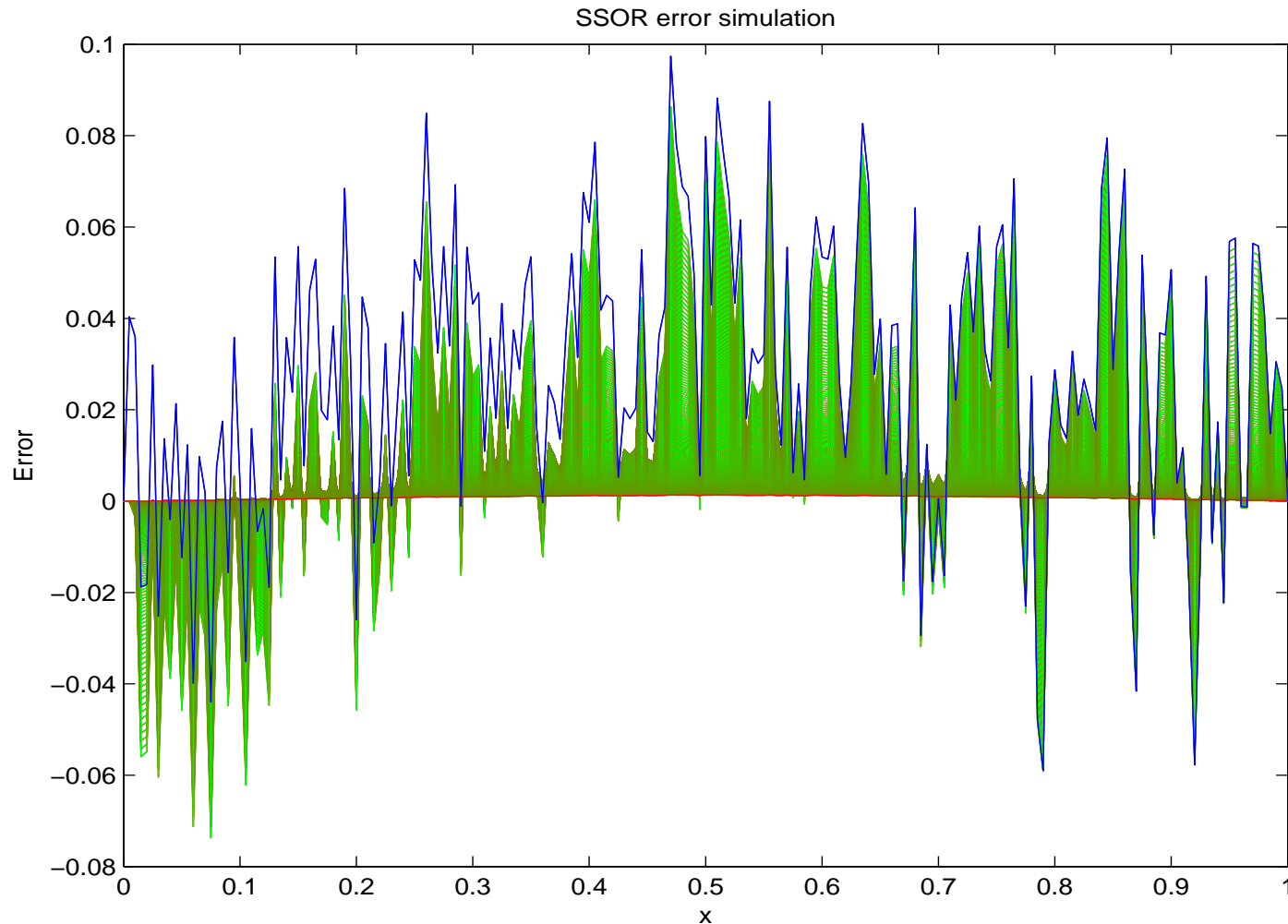
$$(D - \omega L)u^{m+1} = (\omega U + (1 - \omega)D)u^m + \omega f$$

$$(D - \omega U)u^{m+2} = (\omega L + (1 - \omega)D)u^{m+1} + \omega f$$

**Iteration matrix**  $P_{SORfwd} = (D - \omega L)^{-1}(\omega U + (1 - \omega)D)$  and  
 $P_{SORbwd} = (D - \omega U)^{-1}(\omega L + (1 - \omega)D)$

$$P_{SSOR} = P_{SORbwd}P_{SORfwd}$$

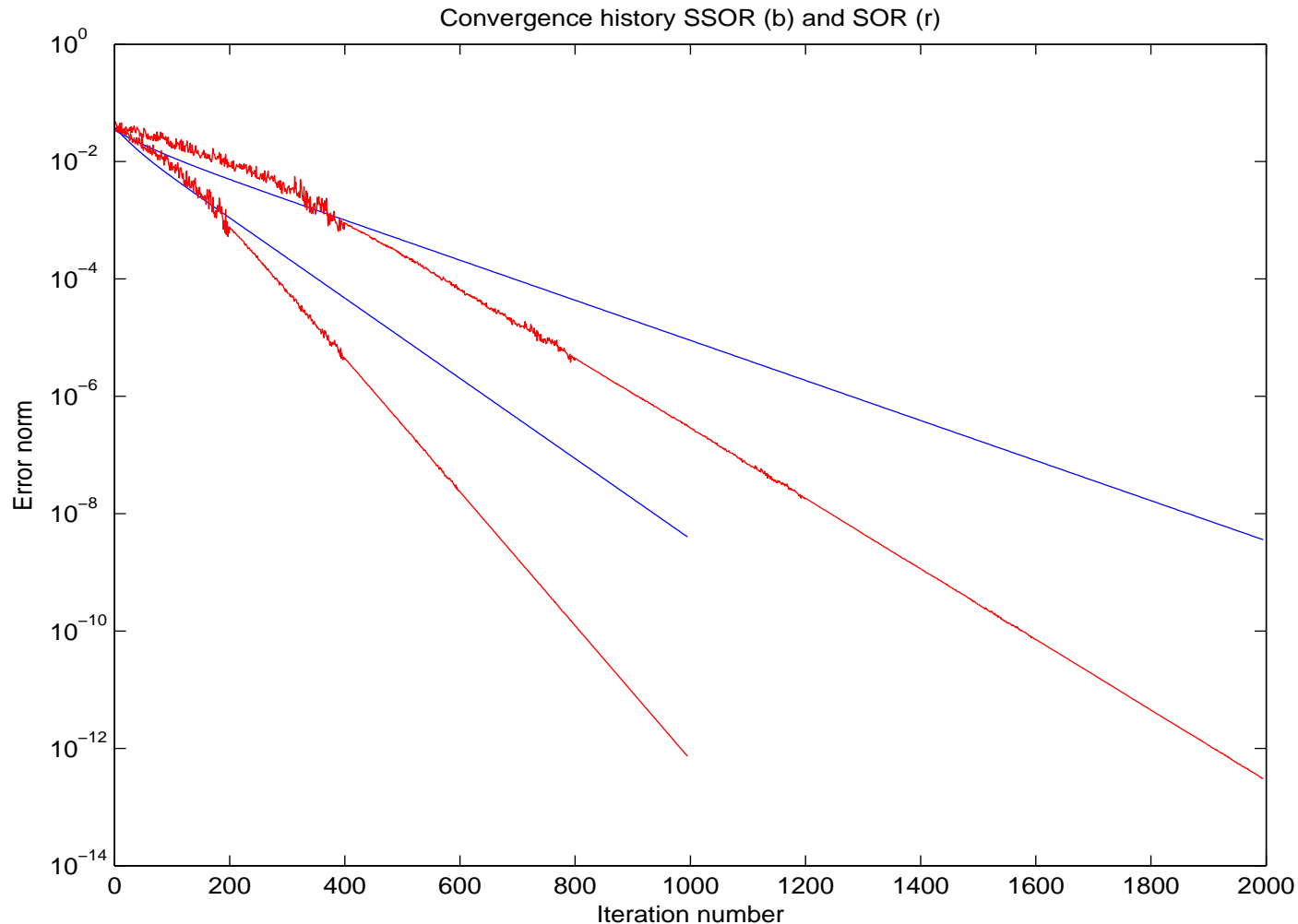
# SSOR error sequence simulation, $N = 199$



**Note** Error “transportation” eliminated

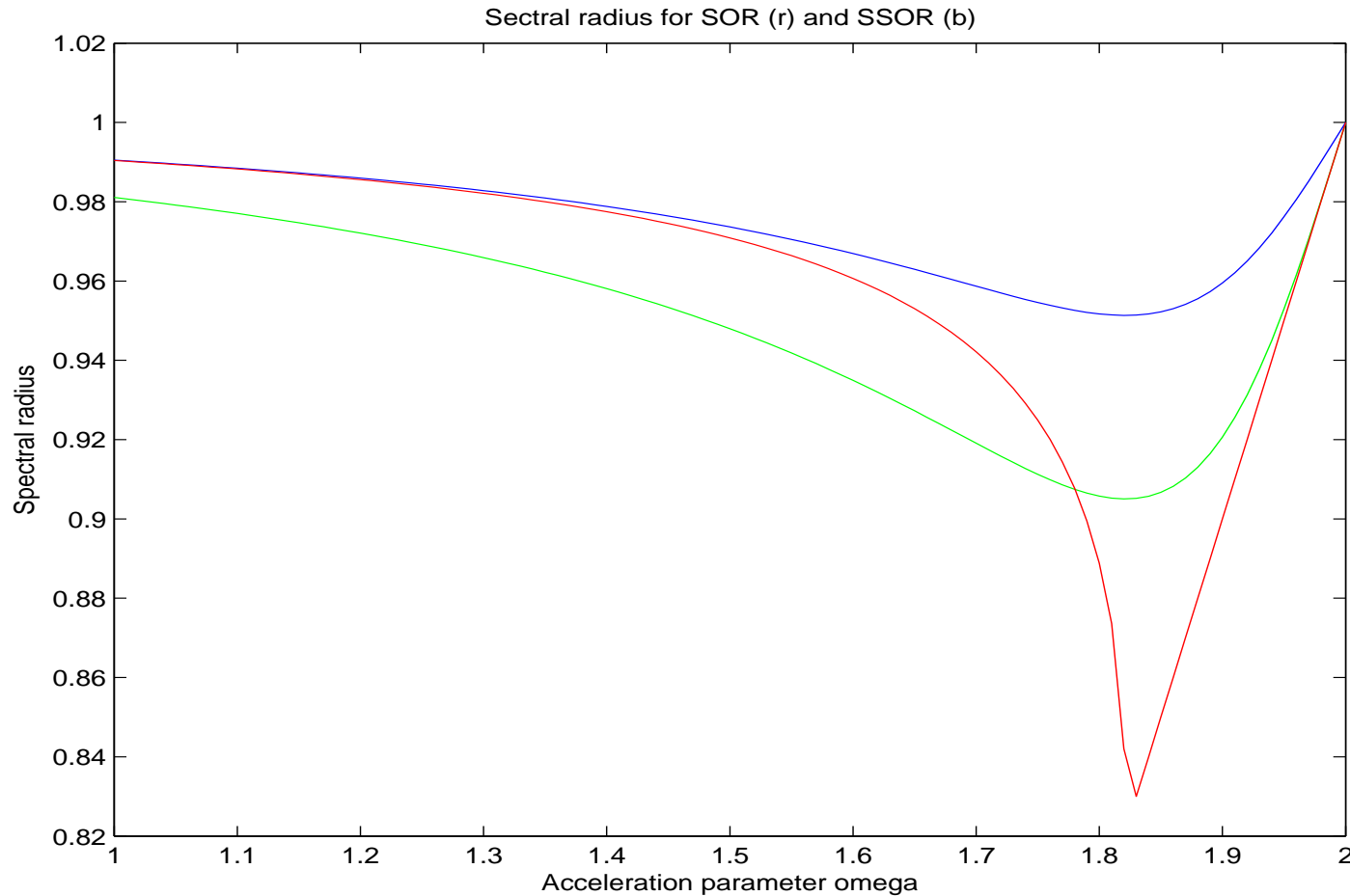


# SSOR error sequence simulation, $N = 199,399$



**Note** Unit error reduction in  $O(N)$  iterations, but **SOR** faster

# SOR and SSOR spectral radii vs $\omega$



$$\rho[P_{SOR}] \quad \sqrt{\rho[P_{SORbwd}P_{SORfwd}]} \quad \rho[P_{SORbwd}P_{SORfwd}]$$

**Note** Optimal acceleration parameter  $\hat{\omega}$  does not move

### 3. Error diffusion and the Jacobi method

For the damped Jacobi method the error recursion is

$$e^{m+1} = P_J(\omega)e^m$$

$$P_J(\omega) = (1 - \omega)I + \omega P_J$$

Consider the *diffusion equation*  $y_t = y_{xx}$  with homogeneous boundary conditions and 2nd order MOL semidiscretization

$$\dot{y} = T_{\Delta x}y$$

Use explicit Euler time–stepping to get

$$y^{m+1} = y^m + \Delta t \cdot T_{\Delta x}y^m$$

## Jacobi and diffusion

Introduce the *Courant number*  $\mu = \Delta t / \Delta x^2$  and the Toeplitz matrix

$$T = \begin{pmatrix} -2 & 1 & 0 \\ 1 & -2 & 1 \\ & & \ddots \\ 0 & 1 & -2 \end{pmatrix}$$

The recursion becomes

$$y^{m+1} = y^m + \mu \cdot T y^m$$

## Jacobi and diffusion...

Split the matrix  $T = D - L - U = -2I + S$  and note that

$$S = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ & & \ddots \\ 0 & 1 & 0 \end{pmatrix} = 2P_J$$

So the recursion can be written

$$y^{m+1} = y^m + \mu \cdot T y^m = [(1 - 2\mu)I + 2\mu P_J] y^m = P_J(2\mu) y^m$$

## Error diffusion in the Jacobi method

Solving diffusion equation  $y_t = y_{xx}$  using Explicit Euler time stepping with finite difference MOL produces

$$y^{m+1} = P_J(2\mu)y^m$$

Solving 2pBVP  $T_{\Delta x}u = f$  using damped Jacobi iterations produces the error recursion

$$e^{m+1} = P_J(\omega)e^m$$

The recursions are identical, with the interpretation twice Courant number  $2\mu =$  acceleration parameter  $\omega$

*Errors in damped Jacobi “diffuse” according to  $e_t = e_{xx}/2$*

# Discrete diffusion

Explicit time-stepping for diffusion is extremely slow

*CFL condition*  $0 < \mu \leq 1/2$

The same thing will happen to Jacobi iterations

*Convergence if*  $0 < \omega \leq 1$

After many steps/iterations *not on the CFL limit*, the solution will approach the low frequency mode  $\sin \pi x$ , while *high frequencies are efficiently damped*

**Recall** What is a high frequency is a grid property

## Side remark: The dynamic/static connection

If we want to solve the (dynamic) diffusion problem

$$y_t = y_{xx} - f(x)$$

its (static) stationary solution solves the 2pBVP

$$y'' = f$$

When we try to solve the latter problem, *every iterative method introduces (discrete) pseudo-dynamics*

$$y^{m+1} = y^m + \omega M r^m$$

Because  $r$  depends on  $y$ , this is the Explicit Euler method for  $\dot{y} = M r(y)$  with step size  $\omega$



# Summary of SOR methods

- ▶ Jacobi is *diffusive* in character, and Gauss–Seidel is *hyperbolic/diffusive*. Both require  $O(N^2)$  iterations for a unit error reduction
- ▶ SOR at optimal acceleration is largely *hyperbolic* in character and only requires  $O(N)$  iterations for a unit error reduction
- ▶ Compare  $\Delta t / \Delta x^2$  and  $\Delta t / \Delta x$  CFL conditions
- ▶ SOR “transports” errors in space “across boundary” with only moderate damping along the way
- ▶ SSOR eliminates error “transportation” but is slower

## 4. The Multigrid idea. Basic tools

To introduce multigrid, we return to damped Jacobi

$$\begin{aligned} r^m &\leftarrow T_{\Delta x} u^m - f \\ u^{m+1} &\leftarrow u^m - \omega D^{-1} r^m \end{aligned}$$

**Theorem** *Let the Jacobi matrix have eigenvalues and eigenvectors given by  $P_J v_k = \lambda_k v_k$ . Then the eigenvectors of  $P_J(\omega) = (1 - \omega)I + \omega P_J$  are  $v_k$  and the eigenvalues are*

$$\lambda_k(\omega) = 1 + \omega(\lambda_k - 1)$$

**Note** The eigenvectors coincide with those of  $T_{\Delta x}$

# Error expansion

For the error recursion (power iteration)

$$e^m = P_J^m(\omega)e^0$$

let  $e^0 = \sum_1^N \alpha_k v_k$ . Then

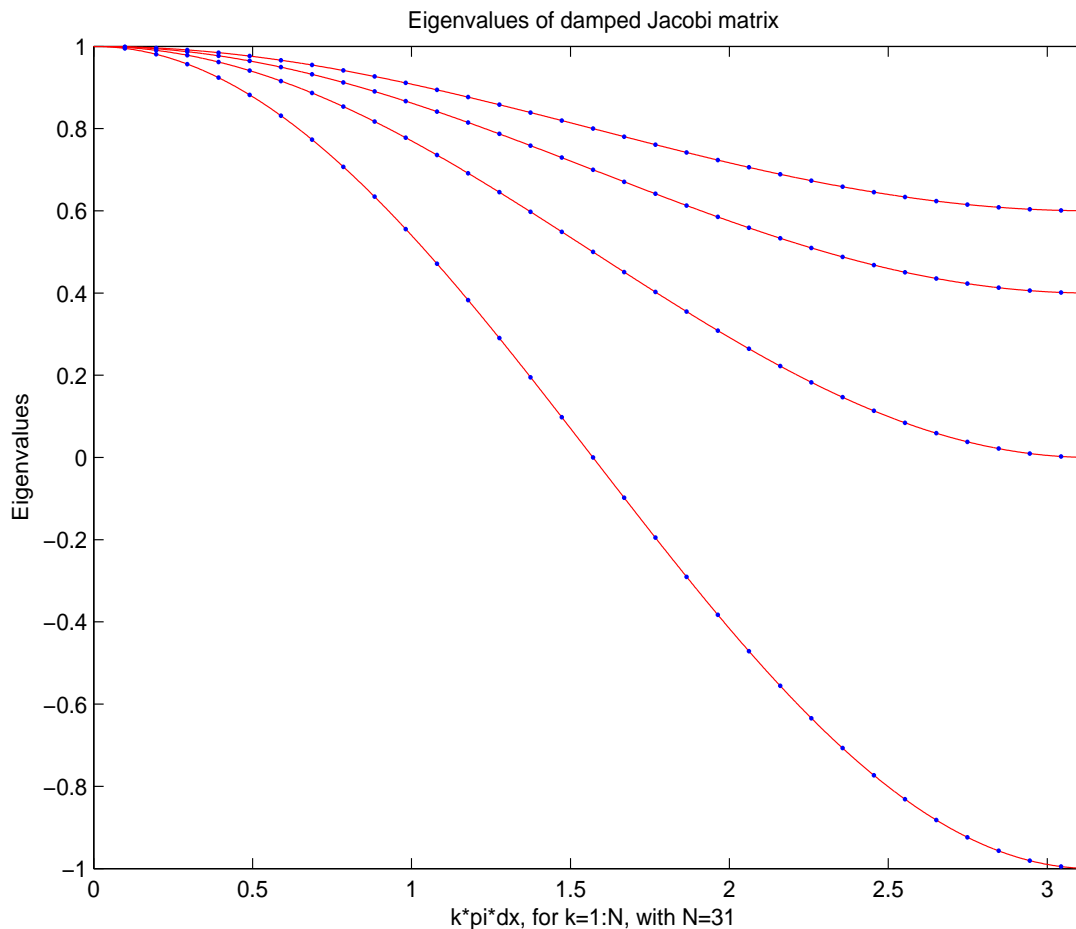
$$e^m = \sum_{k=1}^N \alpha_k \lambda_k^m(\omega) v_k$$

So the reduction of the  $k^{\text{th}}$  error mode is governed by the eigenvalue  $\lambda_k(\omega)$  and can in part be controlled by  $\omega$

We can *make the error smooth* by choosing  $\omega \in [2/3, 1/2]$

# Eigenvalue locations of $P_J(\omega)$

Eigenvalues for  $N = 31$  and  $\omega = 1, 0.5, 0.3, 0.2$



High frequency damping can be controlled by  $\omega$

# Solution, Error and Residual

Given a linear system  $T_{\Delta x}u = f$  construct sequence  $u^m \rightarrow u$

Now recall the following definitions

**The Good** The *solution* is defined by  $T_{\Delta x}u = f$

**The Bad** The *error* is defined by  $e^m = u^m - u$

**The Ugly** The *residual* is defined by  $r^m = T_{\Delta x}u^m - f$

$$T_{\Delta x}e^m = r^m$$

## Error (defect) correction

Note that  $u = u^m - e^m$ . If we can estimate the error,  $\tilde{e} \approx e^m$ , then

$$\tilde{u} \leftarrow u^m - \tilde{e}$$

is an improved solution, so we can put  $u^{m+1} \leftarrow u^m - \tilde{e}$

**Question** *Is it simpler to solve  $T_{\Delta x} e^m = r^m$  than  $T_{\Delta x} u = f$ ?*

**Answer** **YES!** *Because we only need moderate accuracy, solve the residual–error equation **on coarse grid***

***This will still recover the persistent low frequency modes!***

# Error and residual

Recall

$$e^m = \sum_{k=1}^N \alpha_k \lambda_k^m(\omega) v_k$$

with residual

$$r^m = T_{\Delta x} e^m = \sum_{k=1}^N \alpha_k \lambda_k^m(\omega) \kappa_k v_k$$

where  $\kappa_k$  is an eigenvalue of  $T_{\Delta x}$

**Note** As  $\kappa_k = -\frac{4}{\Delta x^2} \sin^2 \frac{k\pi \Delta x}{2}$  the residual is *huge*

## Error and residual...

Taking  $\omega \in [2/3, 1/2]$  means that the error

$$e^m = \sum_{k=1}^N \alpha_k \lambda_k^m(\omega) v_k$$

is *smooth*, as only low frequency modes ( $\lambda_k \approx 1$ ) present

The residual is usually much *less smooth* because *large*  $\kappa_k$  offset the benefit of small coefficients  $\alpha_k$

Solve  $T_{\Delta x} e^m = r^m$  on a *coarse grid*

This still recovers low modes and converges faster



## Example. Fine grid, coarse grid

Fine grid with an *odd number  $N$  of internal grid points*

$$\bar{\Omega}_h = \texttt{linspace}(0, 1, N+2)$$

$$h = 1/(N + 1)$$

Coarse grid with  *$(N - 1)/2$  internal grid points*

$$\bar{\Omega}_H = \texttt{linspace}(0, 1, (N+3)/2)$$

$$H = 2/(N + 1)$$

Mesh widths (step size)  $h$  and  $H$ , respectively

## Grid functions (vectors)

$v_h : \Omega_h \rightarrow \mathbb{R}$  is a *grid function* with  $N$  components

It is an  $N$ -vector (excluding boundary points and data)

Compare a continuous function  $u : [0, 1] \rightarrow \mathbb{R}$

**Note**  $u(\Omega_h)$  is the *restriction* of  $u$  to the grid  $\Omega_h$

Let  $f_h \equiv f(\Omega_h)$ . If  $u'' = f$  and  $T_h v_h = f_h$ , the *error* is

$$e_h = u(\Omega_h) - v_h$$

Both the solution and the error are grid functions

## Downsampling. From fine grid to coarse

**Example**  $N_h = 7$  internal points on fine grid  $\Omega_h$  map to  $N_H = 3$  internal points on coarse grid via *restriction map*

$$v_H = \begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 \end{pmatrix} v_h = I_h^H v_h$$

The restriction drops every second element. An alternative is a *smoothing map*, also known as a *weighted restriction*

$$v_H = \begin{pmatrix} 1/4 & 1/2 & 1/4 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1/4 & 1/2 & 1/4 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1/4 & 1/2 & 1/4 \end{pmatrix} v_h$$

Smoothing map = plain restriction  $\circ$  filter =  $I_h^H \cdot F_\pi$

$$\begin{pmatrix} 1/4 & 1/2 & 1/4 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1/4 & 1/2 & 1/4 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1/4 & 1/2 & 1/4 \end{pmatrix} =$$

$$\begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 \end{pmatrix} \cdot \frac{1}{4} \begin{pmatrix} 2 & 1 & 0 & 0 & 0 & 0 & 0 \\ 1 & 2 & 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 2 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 2 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 2 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 & 2 & 1 \\ 0 & 0 & 0 & 0 & 0 & 1 & 2 \end{pmatrix}$$

# Mapping a fine grid function to a coarse

**Example** Sometimes one *includes the boundary points*

$N_h = 3$  internal points on fine grid  $\bar{\Omega}_h$  map to  $N_H = 1$  internal point on coarse grid  $\bar{\Omega}_H$  via *restriction map*

$$\bar{v}_H = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix} \bar{v}_h = \bar{I}_h^H \bar{v}_h$$

A *smoothing* counterpart is (with unchanged boundaries)

$$v_H = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 1/4 & 1/2 & 1/4 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix} v_h$$

## What is a digital filter?

**Example** To filter out HF, *repeated averaging* can be used

$$u = \frac{1}{4} \begin{pmatrix} 2 & 1 & 0 & 0 \\ 1 & 2 & 1 & 0 \\ 0 & 1 & 2 & 1 \\ 0 & 0 & 1 & 2 \end{pmatrix} v = F_{\pi} \cdot v$$

Then  $u$  is *smoother* than  $v$ ,  $F_{\pi}$  is a 2<sup>nd</sup> order lowpass filter

**Note** The filter  $F_{\pi}$  is *Toeplitz* or a *discrete convolution*, and

$$F_{\pi} = \left(1 - \frac{1}{2}\right)I + \frac{1}{2}P_J = P_J(1/2)$$

Same as  $\omega = 1/2$  damped Jacobi iteration matrix for  $T_{\Delta x}$

# Filters. Analogue and Digital

An *analogue* linear filter is an integral operator  $F : v \mapsto u$

$$u = Fv \quad \Leftrightarrow \quad u(x) = \int_0^1 f(x, y)v(y) \, dy$$

Its discrete counterpart is a *digital* filter

$$u = Fv \quad \Leftrightarrow \quad u_i = \sum_{j=1}^N f_{i,j}v_j$$

The operator  $F$  (as defined by the kernel function  $f(x, y)$  and the matrix  $\{f_{i,j}\}$ ) determine the filter characteristics

# Convolutions and Toeplitz filters

A *convolution filter* is an integral operator  $F : v \mapsto u$

$$u = Fv \quad \Leftrightarrow \quad u(x) = \int_0^1 f(x - y) v(y) \, dy$$

A *discrete convolution* is of the form

$$u = Fv \quad \Leftrightarrow \quad u_i = \sum_{j=1}^N f_{i-j} v_j$$

**Note**  $F$  is a *Toeplitz matrix*, because if  $i - j = \text{const}$  then  $f_{i,j} = f_{i-j} = \text{const}$ . Elements are diagonalwise constant

**Theorem** *A Toeplitz filter is a discrete convolution*



# Oversampling. From coarse grid to fine

**Note** The maps  $I_h^H$  do not have inverses (information loss)

$I_H^h$  is an interpolation operator generating new elements

Left and right inverses

$$I_H^h \cdot I_h^H \neq I$$

$$I_h^H \cdot I_H^h = I$$

The second identity holds if  $I_h^H$  is a plain restriction, but not if it is a smoothing operator

# Prolongation through linear interpolation

**Example** The interpolation operator  $I_H^h$  can be defined by piecewise linear interpolation

$$v_{j+1/2} \leftarrow (v_j + v_{j+1})/2$$

This defines a *prolongation of the grid function* by inserting an interpolated value between previous coarse grid points

Other types of interpolation, such as splines, are possible but are often too expensive and of little advantage

Choose linear interpolation with linear FEM basis functions

# Matrix representation of the prolongation

**Example** Prolongation from 3 internal points to 7

$$v_h = \begin{pmatrix} 1/2 & 0 & 0 \\ 1 & 0 & 0 \\ 1/2 & 1/2 & 0 \\ 0 & 1 & 0 \\ 0 & 1/2 & 1/2 \\ 0 & 0 & 1 \\ 0 & 0 & 1/2 \end{pmatrix} v_H = I_H^h v_H$$

Compare using smoothing restriction:  $I_H^h = 2F_\pi(I_h^H)^T$

# Tools for restriction and prolongation

Two tools needed, the *Toeplitz filter*  $F_\pi = \frac{1}{4} \text{tridiag}(1 \quad 2 \quad 1)$  and the *plain restriction* (“injection”)

$$I_h^H = \begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 \end{pmatrix}$$

**General restriction map**  $R_h^H = I_h^H F_\pi$  (filtering optional)

**Linear interpolation prolongation map**  $P_H^h = 2F_\pi (I_h^H)^T$

## 5. A simple two-grid iteration

Solve  $T_h u_h = f_h$  iteratively using the following algorithm

1. Run one damped  $\omega = 2/3$  Jacobi iteration on fine grid

$$u_h^{m+1} \leftarrow u_h^m - \omega D_h^{-1} r_h^m \quad m = 0 : M - 1$$

2. Compute residual and restrict  $r_H \leftarrow I_h^H F_\pi(T_h u_h^M - f_h)$

3. Solve error equation  $T_H e_H = r_H$

4. Prolong  $e_h^M \leftarrow I_H^h e_H$

5. Correct  $u_h^{M+1} \leftarrow u_h^M - e_h^M$

6. Run one damped Jacobi iteration on fine grid

7. Repeat from 1 until convergence

## MATLAB code segment: the two-grid V step

```
rf = Tdx*v - f;           % Compute residual
v = v - omega*Ddx\rf;     % Pre-smoothing Jacobi

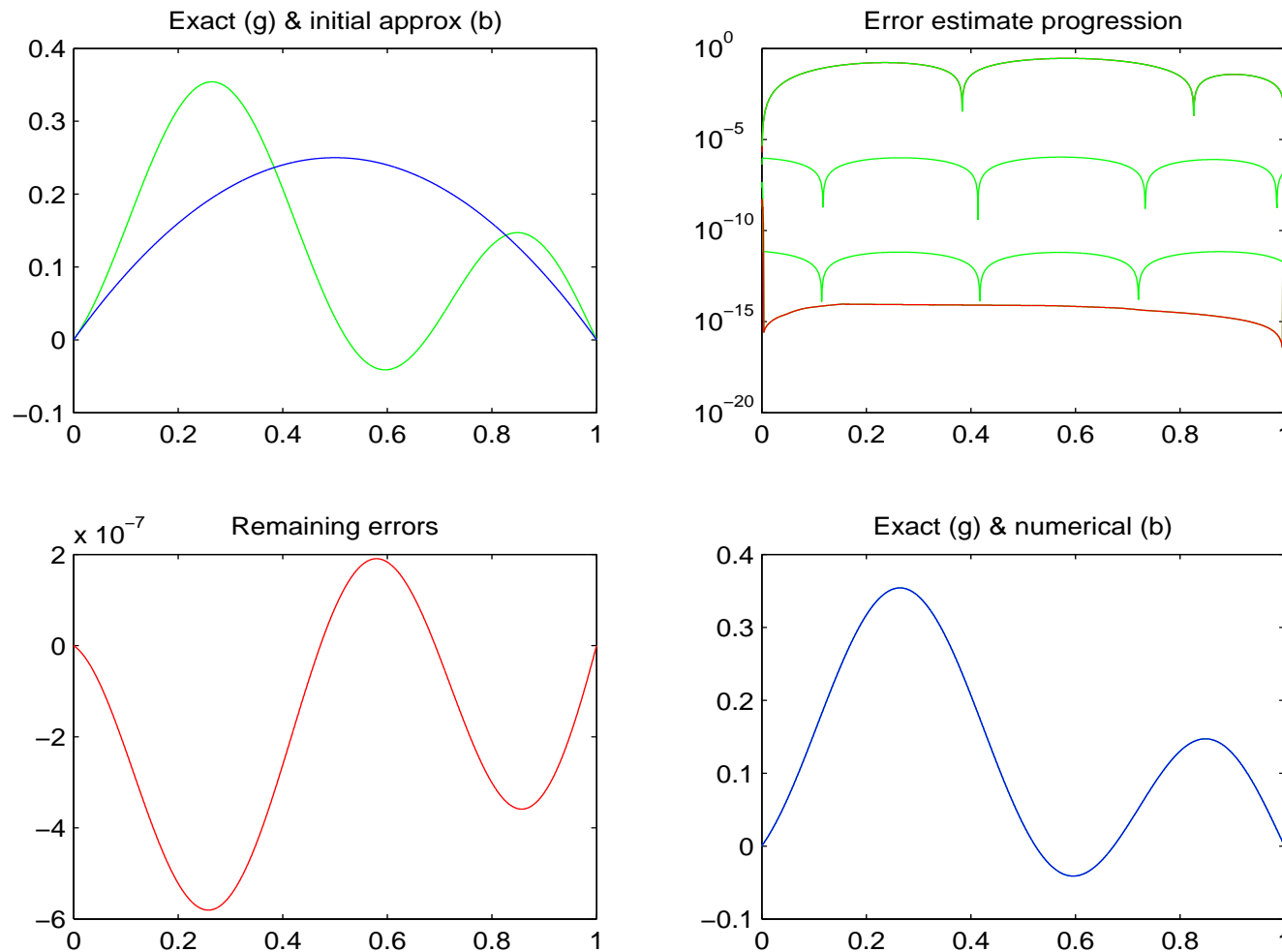
rf = Tdx*v - f;           % Compute residual
rf = lowpass(rf);         % Optional LP filter

rc = FMGrestrict(rf);     % Restrict to coarse grid
ec = Tdxc\rc;             % Solve error equation
ef = FMGprolong(ec);      % Prolong to fine grid

v = v - ef;               % Correct: remove error

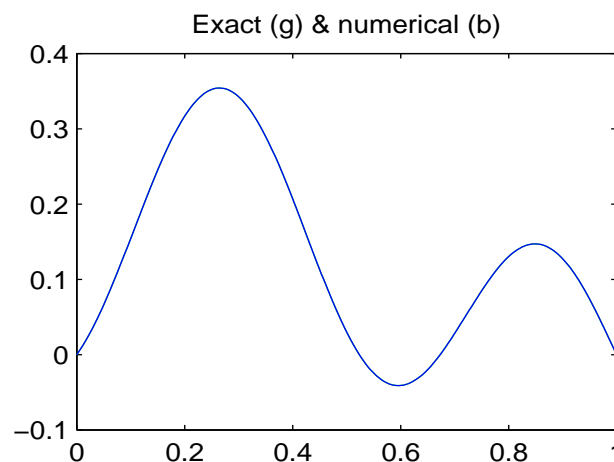
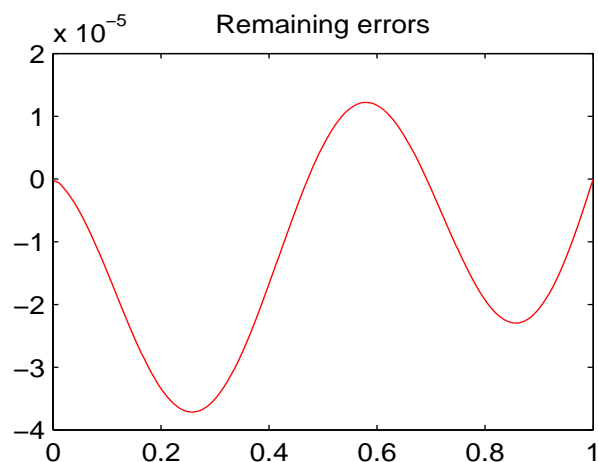
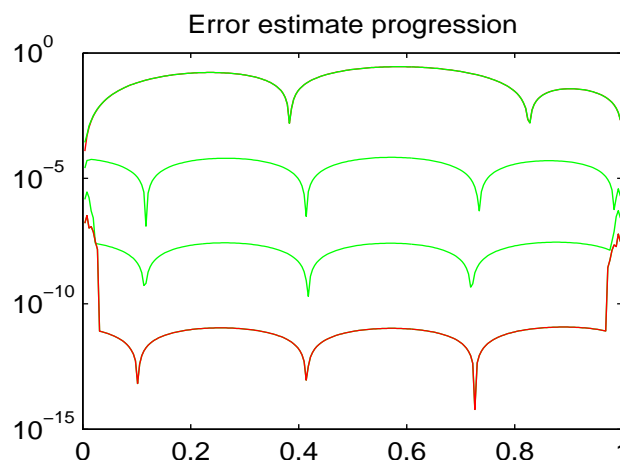
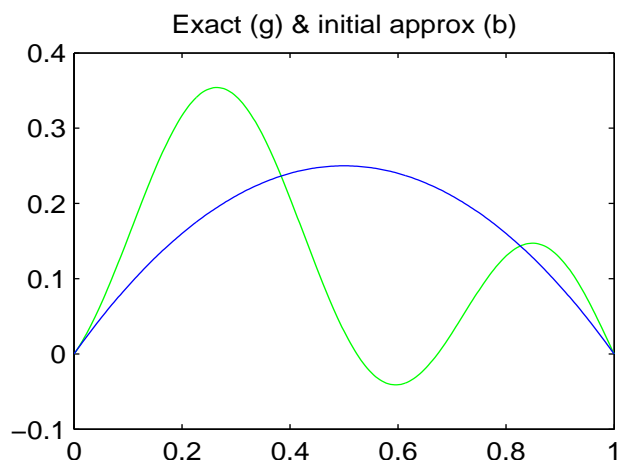
rf = Tdx*v - f;           % Compute residual
v = v - omega*Ddx\rf;     % Post-smoothing Jacobi
```

# Numerical test: Two-grid method, $N = 2048$



4 V-cycles on  $u'' = f$  on on fine grid with 2048 points, error estimation on 1024, single  $P_J(1/2)$  pre- and post-smoothing

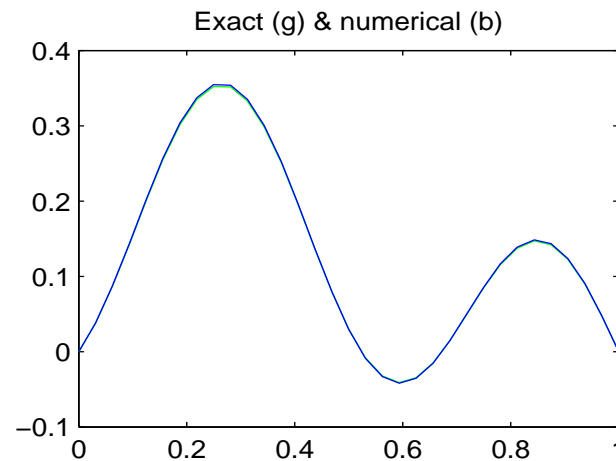
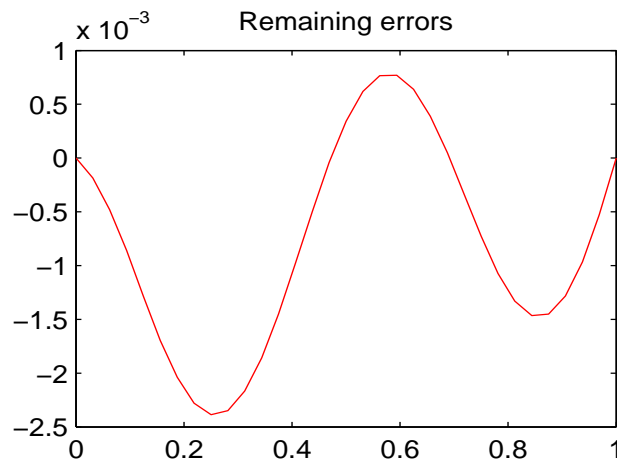
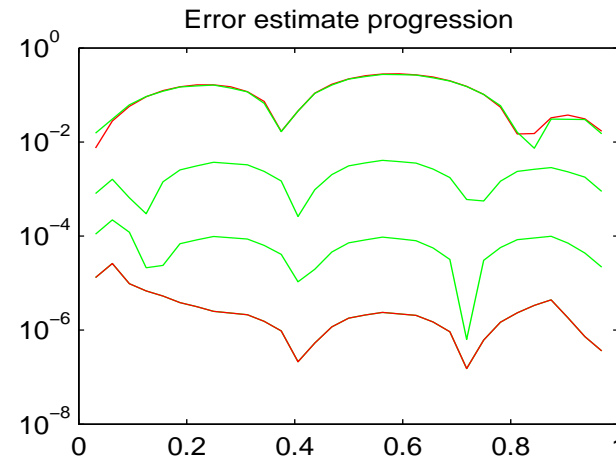
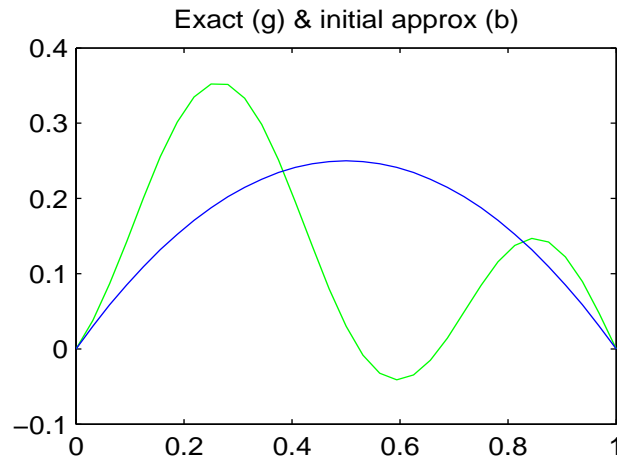
# Numerical test: Two-grid method, $N = 256$



4 V-cycles on  $u'' = f$  on on fine grid with 256 points, error estimation on 128, single  $P_J(1/2)$  pre- and post-smoothing

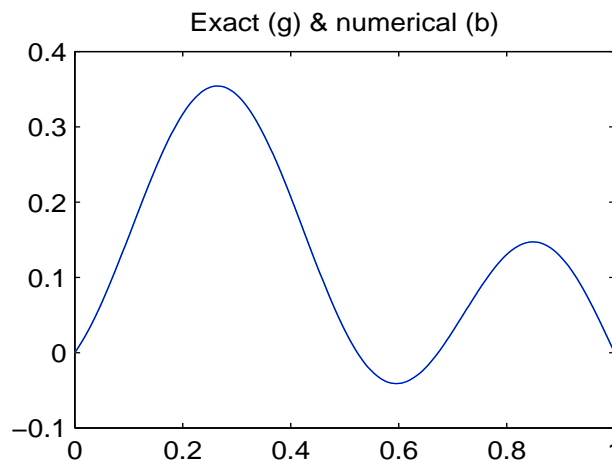
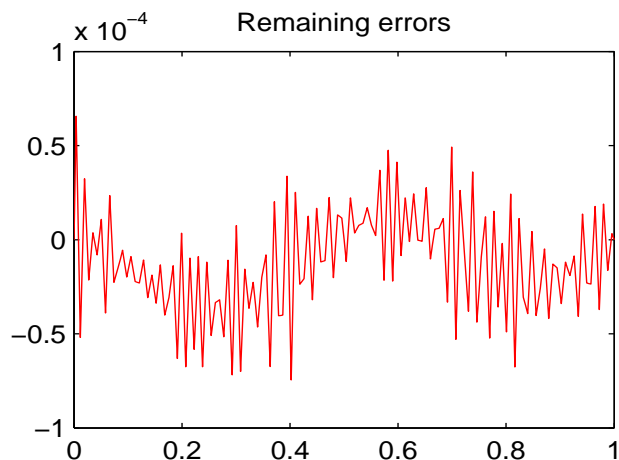
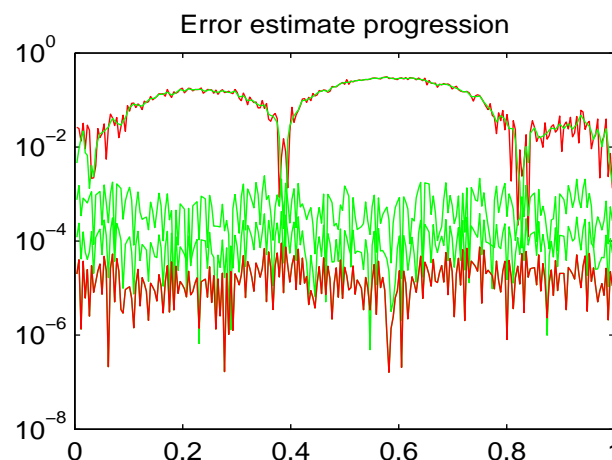
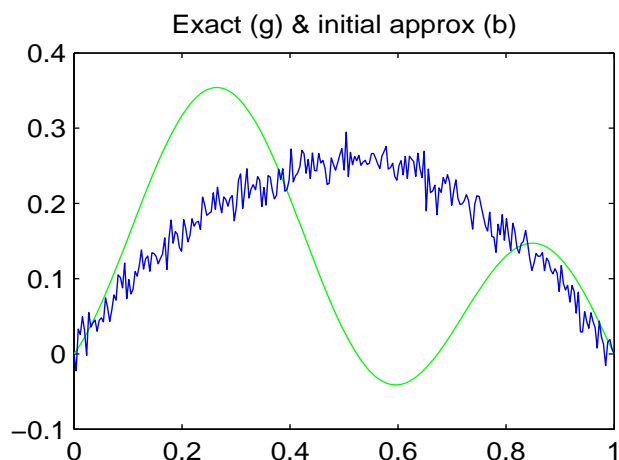


# Numerical test: Two-grid method, $N = 32$



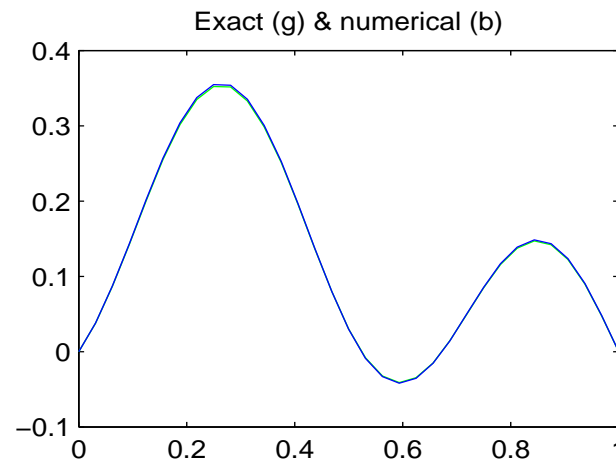
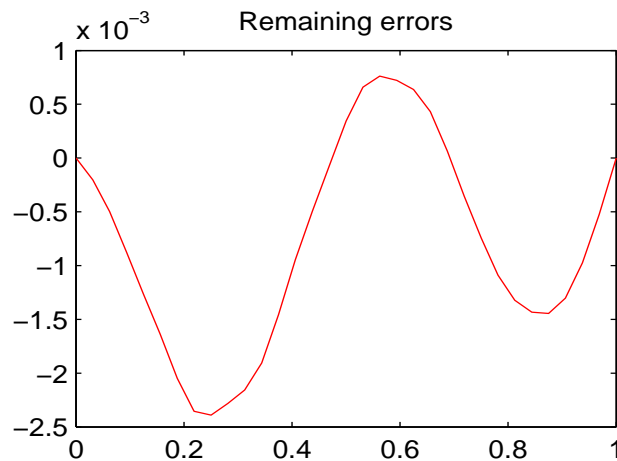
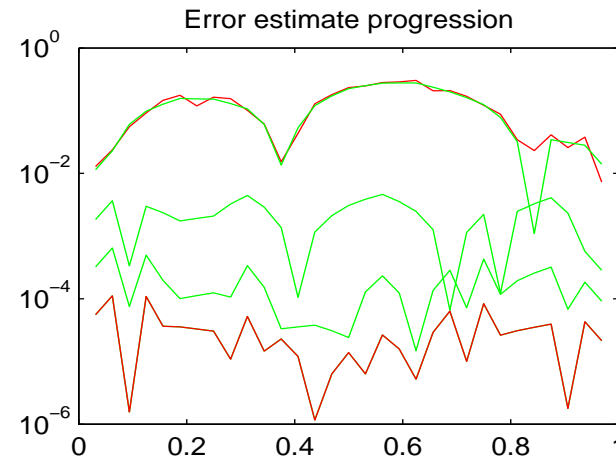
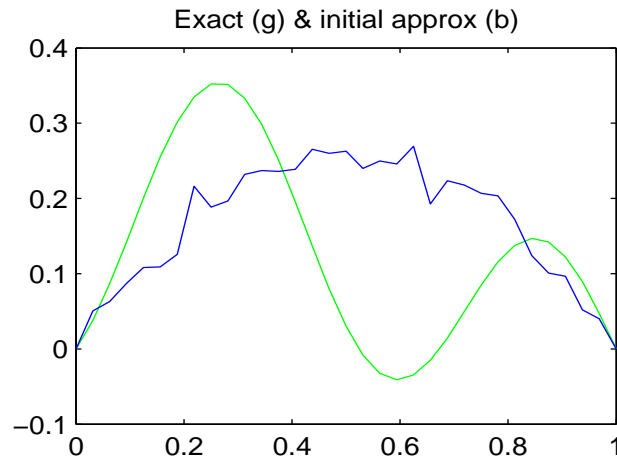
4 V-cycles on  $u'' = f$  on on fine grid with 32 points, error estimation on 16, single  $P_J(1/2)$  pre- and post-smoothing

# Numerical test: Two-grid method, $N = 256$



4 V-cycles on  $u'' = f$  on on fine grid with 256 points, error estimation on 128, with initial data contaminated by noise

# Numerical test: Two-grid method, $N = 32$



4 V-cycles on  $u'' = f$  on on fine grid with 32 points, error estimation on 16, with initial data contaminated by noise

## Observations in the numerical tests

- ▶ One can almost use the *same number of iterations*
- ▶ By and large, *convergence is unaffected by  $\Delta x \rightarrow 0$*
- ▶ One can solve the problem to high accuracy so that *only the global error remains*
- ▶ The basic iterative method is not crucial but must be *smoothing*
- ▶ Large initial errors are not a problem
- ▶ The method is robust even with respect to noise

# Why is the Multigrid Method faster?

All iterative methods work with the residual

$$y^{m+1} = y^m + \omega M r^m$$

The difference is that *multigrid estimates the error*

- ▶ *The residual is large and nonsmooth*
- ▶ *The error is moderate and smooth*
- ▶ *Basic residual-reducing iterative methods typically fail to eliminate low frequency error components*
- ▶ Therefore *only error estimating methods can be fast*

## 6. Multigrid on nested grids. The V cycle

We want to solve  $T_K u_K = f_K$  on a fine grid  $\Omega_K$  with mesh width  $h_K$

Construct an embedding of grids  $\Omega_0 \subset \Omega_1 \subset \dots \subset \Omega_K$  with mesh widths  $h_0 > h_1 > \dots > h_K$

On *regular grids* one typically takes  $h_k = 2h_{k+1}$  to make restrictions and prolongations simple

We want to formulate multigrid methods for solving

$$T_K u_K = f_K$$

on  $\Omega_K$  using all available grids

# Assumptions

- ▶ There are *restriction* operators  $I_k^{k-1} : \Omega_k \rightarrow \Omega_{k-1}$
- ▶ There are *prolongation* operators  $I_k^{k+1} : \Omega_k \rightarrow \Omega_{k+1}$
- ▶ The restrictions and prolongations must be very fast and run in  $\mathcal{O}(N)$  operations
- ▶ There is an iterative method with the *smoothing property* on every grid  $\Omega_k$ , i.e., it must attenuate high frequencies by a constant factor *independently of*  $h_k$
- ▶  $T_0 u_0 = f_0$  can be solved fast to provide a crude estimate  $u_0$

# The operators $\{T_k\}$

Different operators  $T_k$  are needed on the different grids  
Standard approach: if possible use same discretization

It may be “difficult” to construct the  $T_k$

*Use restriction and prolongation operators!*

**Definition** The operator triple  $(T_k, R_k^{k-1}, P_{k-1}^k)$  satisfies the *Galerkin condition* if

$$T_{k-1} = R_k^{k-1} T_k P_{k-1}^k$$

where  $R_k^{k-1}$  is the restriction and  $P_{k-1}^k$  is the prolongation



# The V-cycle. Sweep through the grids

Obtain a *starting approximation* by solving the problem on a very coarse grid, typically  $\Omega_0$ . Solving means by *direct method*,  $N = 1000$  could be typical. Prolong to finest grid

**A “V” sweep** Iterate from finest grid on successively coarser grids down to  $\Omega_0$ , then back to  $\Omega_K$  on successively finer grids

Repeat V-sweeps to form a **“V-cycle”** and iterate until convergence

Visit enough grids (many) to suppress the entire spectrum of frequencies, coarse grids for slow, fine grids for fast

# A full, recursive multigrid V-cycle

1. Procedure  $v := FMGV(v_k, T_k, f_k)$
2. If  $k = 0$  return  $v \leftarrow T_0 \backslash f_0$  and **STOP** else:
3. Pre-smoothing  $v_k \leftarrow v_k - \omega D_k^{-1}(T_k v_k - f_k)$
4. Restrict residual to coarse grid  $r_{k-1} \leftarrow I_k^{k-1}(T_k v_k - f_k)$
5. Recursive error estimation  $e_{k-1} \leftarrow FMGV(0, T_{k-1}, r_{k-1})$
6. Prolong error to fine grid  $e_k \leftarrow I_{k-1}^k e_{k-1}$
7. Correct by removing error  $v_k \leftarrow v_k - e_k$
8. Post-smoothing  $v_k \leftarrow v_k - \omega D_k^{-1}(T_k v_k - f_k)$
9. Output  $v \leftarrow v_k$

# Starting procedure

If direct solution  $v_0 = T_0 \backslash f_0$  is affordable, then obtain a first approximation on this grid by direct solver

Sweep down on successively finer grids:

1. Compute  $v_0 = T_0 \backslash f_0$
2. Prolong to next finer fine grid  $v_k \leftarrow I_{k-1}^k v_{k-1}$
3. Compute residual  $r_k = T_k v_k - f_k$
4. Smoothing iteration(s)  $v_k \leftarrow v_k - \omega D_k^{-1} r_k$
5. Prolong to next grid &c. until first  $v_K$  has been obtained

## MATLAB code segment. The recursive V cycle

```
rf = Tdx*v - f;           % Compute residual
v = v - omega*Ddx\rf;     % Pre-smoothing Jacobi

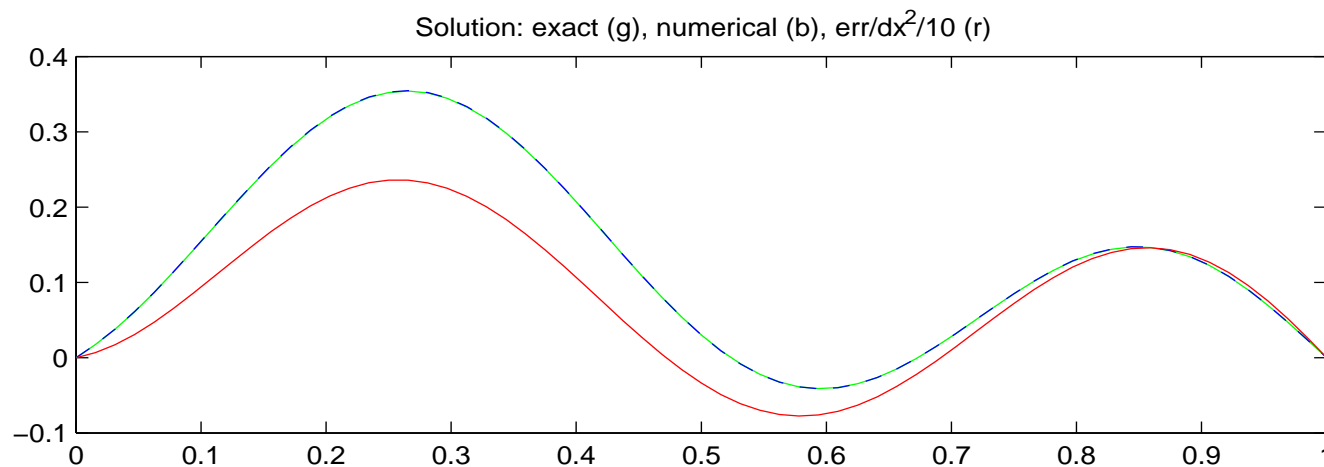
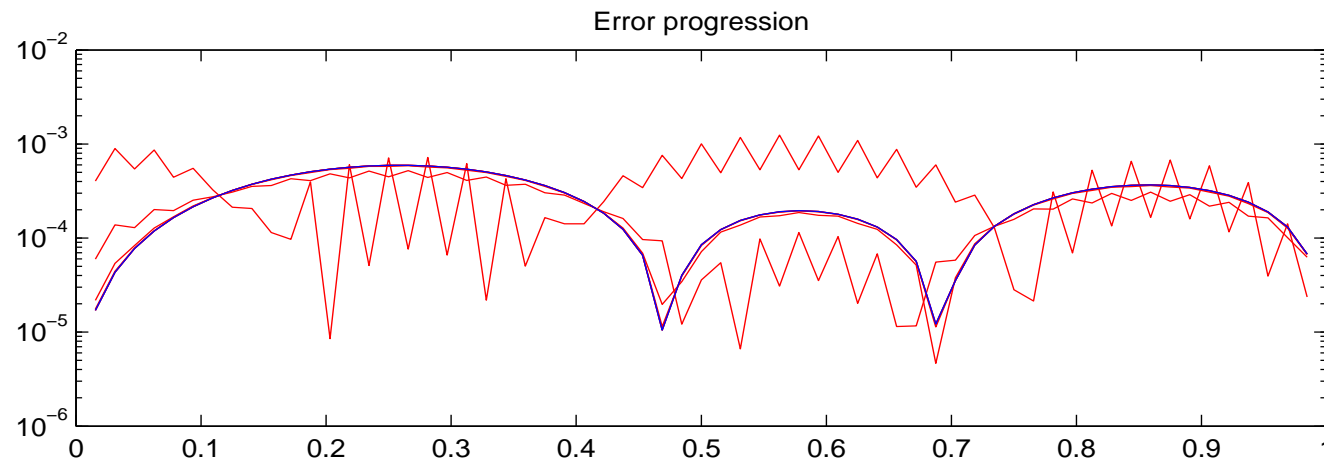
rf = Tdx*v - f;           % Compute residual
rf = lowpass(rf);         % Optional LP filter

rc = FMGrestrict(rf);     % Restrict to coarse
ec = FMGV(0,T2dx,rc);     % Solve error equation
ef = FMGprolong(ec);      % Prolong to fine grid

v = v - ef;               % Correct: remove error

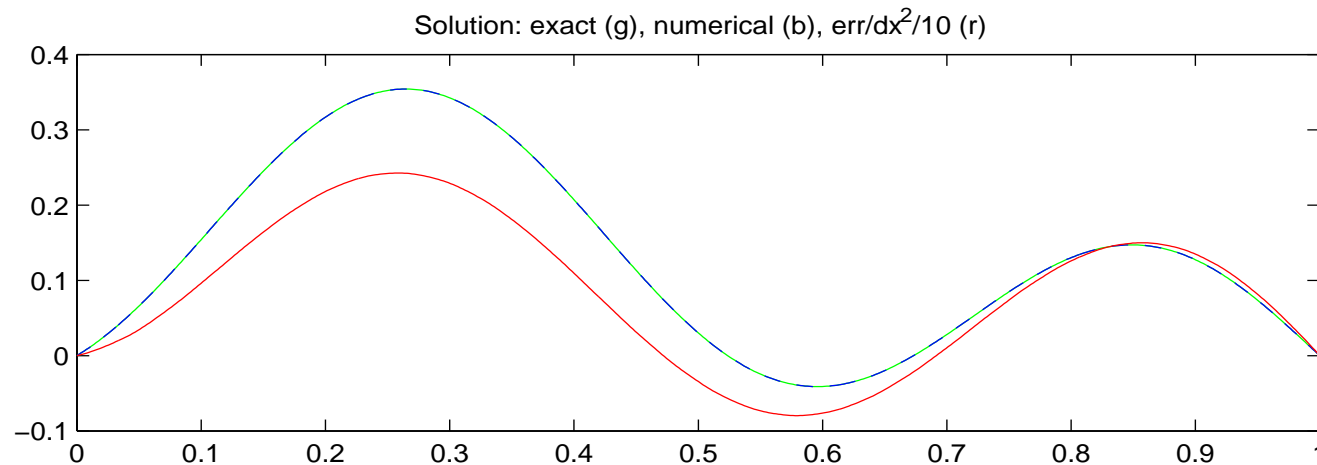
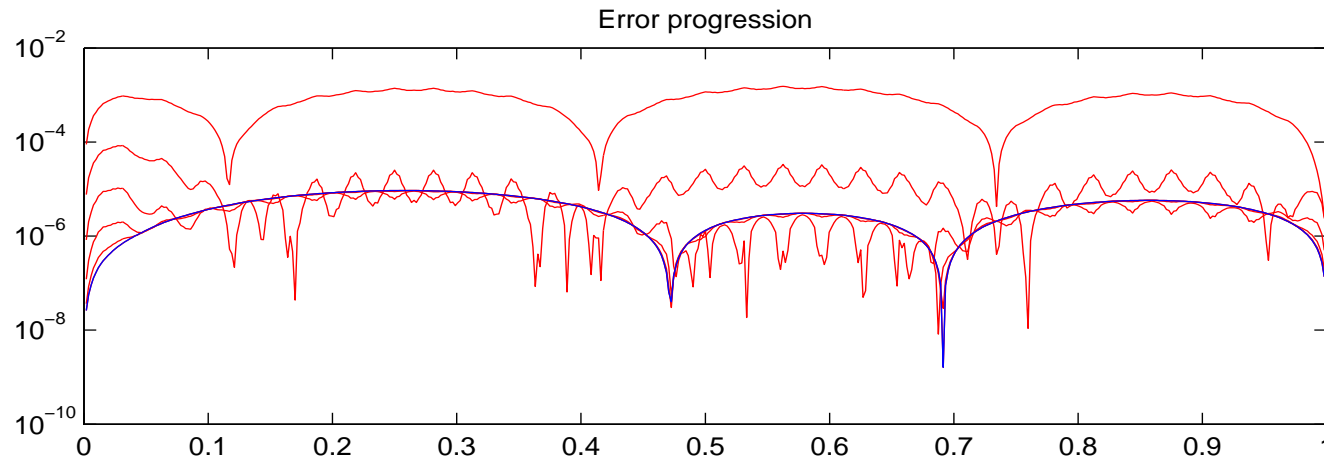
rf = Tdx*v - f;           % Compute residual
v = v - omega*Ddx\rf;     % Post-smoothing Jacobi
```

# Numerical test: Full MG V-cycles, $N = 63$



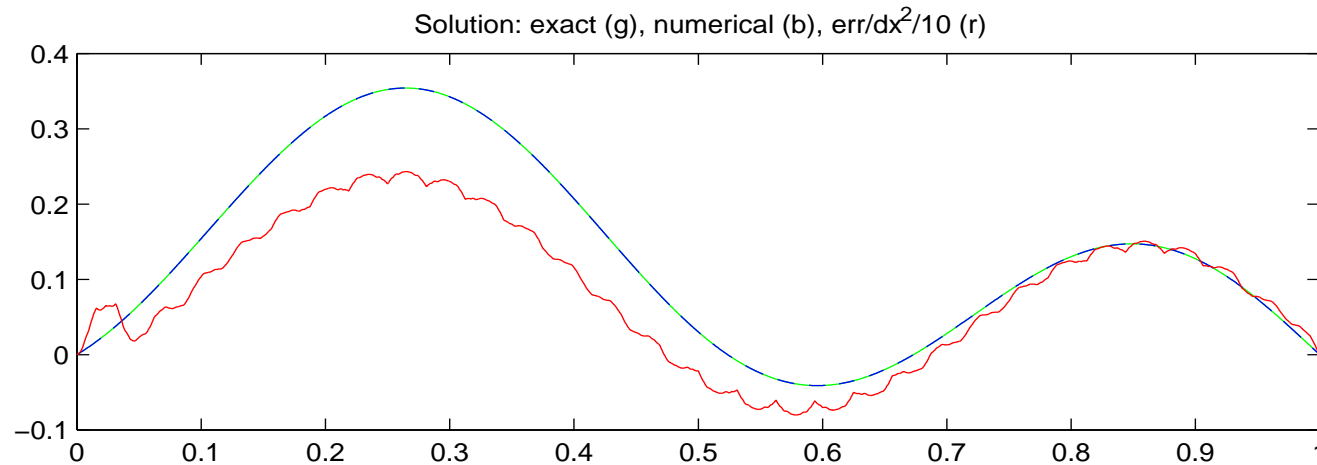
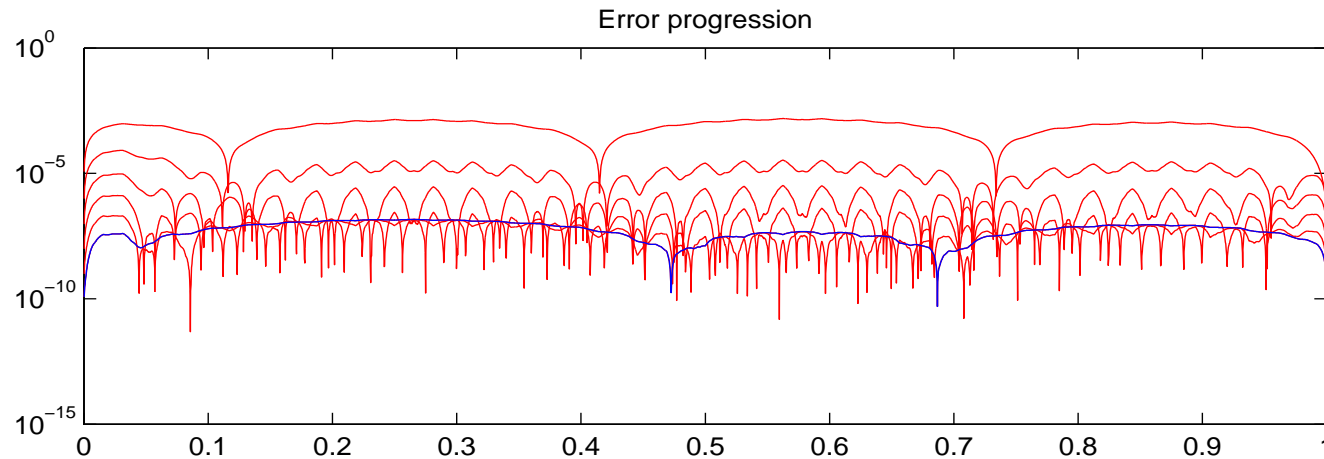
6 V-cycles on  $u'' = f$  on on fine grid with 63 points, coarsest grid with 31 points and Jacobi  $\omega = 2/3$

# Numerical test: Full MG V-cycles, $N = 511$



6 V-cycles on  $u'' = f$  on on fine grid with 511 points, coarsest grid with 31 points and Jacobi  $\omega = 2/3$

# Numerical test: Full MG V-cycles, $N = 4095$



6 V-cycles on  $u'' = f$  on on fine grid with 4095 points, coarsest grid with 31 points and Jacobi  $\omega = 2/3$

## Observations in the numerical tests

- ▶ *By and large, **multigrid iteration convergence is unaffected by  $\Delta x \rightarrow 0$***
- ▶ *One can solve the problem to high accuracy so that **only the global error remains***
- ▶ *Within 6 V-cycles the global error is reached*
- ▶ *Solution accuracy is  $O(\Delta x^2)$*
- ▶ *Large initial errors are not a problem*
- ▶ *The method is robust even with respect to noise, but more V-cycles may be required for large  $N$*



## 7. Efficiency. Memory and execution time

Let the  $K^{\text{th}}$  grid have  $N$  internal points

Then grid  $K - 1$  has  $(N + 1)/2 - 1 < N/2$  internal points

Total storage requirement is therefore

$$S = N \cdot (1 + 2^{-1} + \dots + 2^{-K}) < 2N$$

So, independently of the # of grids (recursion depths) *storage does not exceed that of using a single finer mesh*

# Memory requirements in $d$ dimensions

Let the  $K^{\text{th}}$  grid have  $N^d$  points (curse of dimensions)

Then grid  $K - 1$  has  $(N + 1)^d / 2 - 1 < (N/2)^d$  points

Total storage requirement in  $d$  dimensions is

$$S_d < N^d \cdot (1 + 2^{-d} + \cdots + 2^{-dK}) \approx N^d$$

$$S_1 < 2N \quad S_2 < \frac{4N^2}{3} \quad S_3 < \frac{8N^3}{7}$$

So, for once, the *curse of dimensions works our way!*

Multigrid is always cheaper than using a single finer grid

## Execution time “in theory”

Work unit  $U_K$  is the cost of one sweep on the fine  $K^{\text{th}}$  grid

$W_0 \sim 2^{-dK} N^{2d}$  is cost for direct solution on coarsest grid

Neglect restriction and prolongation costs

Cost  $W$  of V cycle with one pre- and one post-smoothing

$$W_d < 2U_K \cdot (1 + 2^{-d} + \cdots + 2^{-dK}) + 2W_0 \approx 2U_K + 2W_0$$

$$W_1 < 4U_K + 2W_0; \quad W_2 < \frac{8U_K}{3} + 2W_0; \quad W_3 < \frac{16U_K}{7} + 2W_0$$

So, *execution time does not exceed that of a single sweep on the next finer grid*

# So it's cheap, what accuracy do we get?

Solving  $T_{\Delta x} u_{\Delta x} = f$  we obtain  $v_{\Delta x}$  with

*Algebraic error*  $e_{\Delta x} = u_{\Delta x} - v_{\Delta x}$

*Global error*  $g_{\Delta x} = u_{\Delta x} - u(\Omega_{\Delta x})$

*Total error*  $\delta_{\Delta x} = v_{\Delta x} - u(\Omega_{\Delta x})$

Note that  $\delta_{\Delta x} = g_{\Delta x} - e_{\Delta x}$ , therefore

$$\begin{aligned} \|v_{\Delta x} - u(\Omega_{\Delta x})\|_{\Delta x} &\leq \|u_{\Delta x} - u(\Omega_{\Delta x})\|_{\Delta x} + \|v_{\Delta x} - u_{\Delta x}\|_{\Delta x} \\ &\leq C \cdot \Delta x^2 + \varepsilon_{\text{MG}} \end{aligned}$$

Try to make  $\varepsilon_{\text{MG}} \approx C \cdot \Delta x^2$  then global error will dominate

# Convergence analysis – a brief sketch

Assume linear convergence rate  $\gamma$  for each V-cycle

Run  $M$  V-cycles, reducing error to global error level if

$$\gamma^M \sim \Delta x^2 \sim N^{-2d}$$

for a 2<sup>nd</sup> order method, implying number of iterations

$$M = O(\log N)$$

Because the cost of each V-cycle is  $\approx U_K = O(N^d)$  total cost to reduce the error to the global error level is

$$O(N^d \log N) = C \cdot \text{cost(FFT)}$$