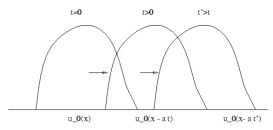
■ 1st order PDEs:
$$\frac{\partial r}{\partial t} = v \frac{\partial s}{\partial x}$$
, $\frac{\partial s}{\partial t} = v \frac{\partial r}{\partial x}$, $r \equiv v \frac{\partial u}{\partial x}$, $s \equiv \frac{\partial u}{\partial t}$

- 1st order PDEs: $\frac{\partial r}{\partial t} = v \frac{\partial s}{\partial x}$, $\frac{\partial s}{\partial t} = v \frac{\partial r}{\partial x}$, $r \equiv v \frac{\partial u}{\partial x}$, $s \equiv \frac{\partial u}{\partial t}$
- $\mathbf{F}(\mathbf{u}) = \begin{bmatrix} 0 & -v \\ -v & 0 \end{bmatrix} \cdot \mathbf{u}$; eigenvalues $\pm v$ waves

- 1st order PDEs: $\frac{\partial r}{\partial t} = v \frac{\partial s}{\partial x}$, $\frac{\partial s}{\partial t} = v \frac{\partial r}{\partial x}$, $r \equiv v \frac{\partial u}{\partial x}$, $s \equiv \frac{\partial u}{\partial t}$
- $\mathbf{F}(\mathbf{u}) = \begin{bmatrix} 0 & -v \\ -v & 0 \end{bmatrix} \cdot \mathbf{u}$; eigenvalues $\pm v$ waves
- \blacksquare advection eq., simplest wave eq. $\frac{\partial u}{\partial t} + v \frac{\partial u}{\partial x} = 0$

- 1st order PDEs: $\frac{\partial r}{\partial t} = v \frac{\partial s}{\partial x}$, $\frac{\partial s}{\partial t} = v \frac{\partial r}{\partial x}$, $r \equiv v \frac{\partial u}{\partial x}$, $s \equiv \frac{\partial u}{\partial t}$
- $\mathbf{F}(\mathbf{u}) = \begin{bmatrix} 0 & -v \\ -v & 0 \end{bmatrix} \cdot \mathbf{u}$; eigenvalues $\pm v$ waves
- advection eq., simplest wave eq. $\frac{\partial u}{\partial t} + v \frac{\partial u}{\partial x} = 0$
- initial wave retains form, just displaced; building block

- $\ \ \, \hbox{1st order PDEs:} \,\, \tfrac{\partial r}{\partial t} = v \tfrac{\partial s}{\partial x} , \,\, \tfrac{\partial s}{\partial t} = v \tfrac{\partial r}{\partial x} , \,\, r \equiv v \tfrac{\partial u}{\partial x} , \,\, s \equiv \tfrac{\partial u}{\partial t}$
- $\mathbf{F}(\mathbf{u}) = \begin{bmatrix} 0 & -v \\ -v & 0 \end{bmatrix} \cdot \mathbf{u}$; eigenvalues $\pm v$ waves
- advection eq., simplest wave eq. $\frac{\partial u}{\partial t} + v \frac{\partial u}{\partial x} = 0$
- initial wave retains form, just displaced; building block
- u(x,t) = u(x-vt,0)



Advection equation: FTCS, Lax

■ FTCS:
$$\frac{u_j^{n+1} - u_j^n}{\Delta t} = -v \left(\frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} \right)$$

Advection equation: FTCS, Lax

■ FTCS:
$$\frac{u_j^{n+1} - u_j^n}{\Delta t} = -v \left(\frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} \right)$$

- unfortunately numerically unstable
- $\qquad \text{VNSA: } \frac{\xi-1}{\Delta t} = -\frac{v}{2\Delta x}(e^{ik\Delta x} e^{-ik\Delta x}) \text{, or, } \xi = 1 i\frac{v\Delta t}{\Delta x}\sin k\Delta x$
- $\blacksquare |\xi| > 1 \text{ for all } \Delta t > 0!$

Advection equation: FTCS, Lax

$$\blacksquare \mathsf{FTCS} : \frac{u_j^{n+1} - u_j^n}{\Delta t} = -v \left(\frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} \right)$$

unfortunately numerically unstable

■ VNSA:
$$\frac{\xi-1}{\Delta t} = -\frac{v}{2\Delta x}(e^{ik\Delta x} - e^{-ik\Delta x})$$
, or, $\xi = 1 - i\frac{v\Delta t}{\Delta x}\sin k\Delta x$

- \bullet $|\xi| > 1$ for all $\Delta t > 0$!
- Lax: $u_j^{n+1} = \frac{(u_{j+1}^n + u_{j-1}^n)}{2} v\Delta t \left(\frac{u_{j+1}^n u_{j-1}^n}{2\Delta x}\right)$
- VNSA:

$$\xi = \frac{(e^{ik\Delta x} + e^{-ik\Delta x})}{2} - \frac{v\Delta t}{2\Delta x} (e^{ik\Delta x} - e^{-ik\Delta x}) = \cos k\Delta x - i\frac{v\Delta t}{\Delta x} \sin k\Delta x$$

- lacksquare $|\xi| < 1$ if $\Delta t \leq \Delta x/v$; Courant-Friedrichs-Lewy (CFL) condition
- Lax equivalent to $\frac{\partial u}{\partial t} + v \frac{\partial u}{\partial x} = \frac{\Delta x^2}{2\Delta t} \nabla^2 u$

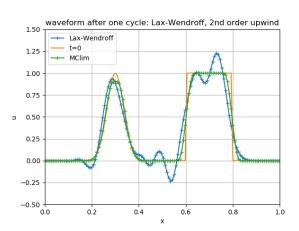


Advection equation: errors

- \blacksquare Amplitude, phase errors; amplification factor $\xi = A e^{i\phi}$
- Amplitude error: 1 A, phase error = $-iv\Delta t \phi$
- $|\xi| \leq 1$ is must; error for resolved modes $(k \ll \pi/\Delta x)$

Advection equation: errors

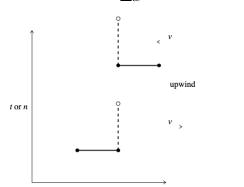
- Amplitude, phase errors; amplification factor $\xi = Ae^{i\phi}$
- Amplitude error: 1 A, phase error = $-iv\Delta t \phi$
- ullet $|\xi| \leq 1$ is must; error for resolved modes $(k \ll \pi/\Delta x)$
- shocks from nonlinearity! artificial/numerical viscosity
- identify amplitude, phase error in the following



- information propagates only in one direction
- v > 0: grid i affected only by i 1 and not i + 1!

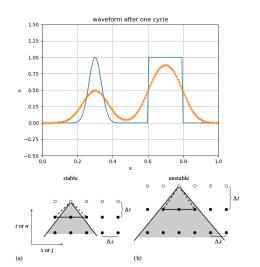
- information propagates only in one direction
- v > 0: grid i affected only by i 1 and not i + 1!
- upwind differencing (only 1st order accurate):

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} = -v_j^n \frac{u_j^n - u_{j-1}^n}{\Delta x}, \text{ for } v_j^n \ge 0$$
$$= -v_j^n \frac{u_{j+1}^n - u_j^n}{\Delta x}, \text{ for } v_j^n < 0$$



- VNSA: $\xi = 1 \left| \frac{v\Delta t}{\Delta x} \right| (1 \cos k\Delta x) i \frac{v\Delta t}{\Delta x} \sin k\Delta x$
- $|\xi|^2 = 1 2 \left| \frac{v\Delta t}{\Delta x} \right| \left(1 \left| \frac{v\Delta t}{\Delta x} \right| \right) \left(1 \cos k\Delta x \right)$

- VNSA: $\xi = 1 \left| \frac{v\Delta t}{\Delta x} \right| (1 \cos k\Delta x) i \frac{v\Delta t}{\Delta x} \sin k\Delta x$
- $|\xi|^2 = 1 2 \left| \frac{v\Delta t}{\Delta x} \right| (1 \left| \frac{v\Delta t}{\Delta x} \right|) (1 \cos k \Delta x)$
- for numerical stability, again, $\Delta t \leq \Delta x/v$ (CFL)



5/10

second order accuracy desirable

- second order accuracy desirable
- Two-step LW

$$u_{j+1/2}^{n+1/2} = \frac{1}{2}(u_{j+1}^n + u_j^n) - \frac{\Delta t}{2\Delta x}(F_{j+1}^n - F_j^n)$$

$$u_j^{n+1} = u_j^n - \frac{\Delta t}{\Delta x}\left(F_{j+1/2}^{n+1/2} - F_{j-1/2}^{n+1/2}\right)$$
two-step Lax Wendroff
halfstep points
halfstep points

x or i

- second order accuracy desirable
- Two-step LW

$$u_{j+1/2}^{n+1/2} = \frac{1}{2}(u_{j+1}^n + u_j^n) - \frac{\Delta t}{2\Delta x}(F_{j+1}^n - F_j^n)$$

$$u_j^{n+1} = u_j^n - \frac{\Delta t}{\Delta x}\left(F_{j+1/2}^{n+1/2} - F_{j-1/2}^{n+1/2}\right)$$
two-step Lax Wendroff halfstep points halfstep points

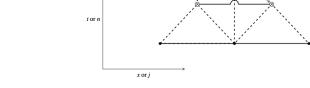
■ again $\Delta t \leq \Delta x/v$ for stability (CFL)

x or i

- second order accuracy desirable
- Two-step LW

$$u_{j+1/2}^{n+1/2} = \frac{1}{2}(u_{j+1}^n + u_j^n) - \frac{\Delta t}{2\Delta x}(F_{j+1}^n - F_j^n)$$

$$u_j^{n+1} = u_j^n - \frac{\Delta t}{\Delta x}\left(F_{j+1/2}^{n+1/2} - F_{j-1/2}^{n+1/2}\right)$$
two-step Lax Wendroff
halfstep points



- again $\Delta t \leq \Delta x/v$ for stability (CFL)
- problem: oscillations at discontinuities!

■ Euler equation develops shocks physically!

- Euler equation develops shocks physically!
- oscillatory schemes are unacceptable

- Euler equation develops shocks physically!
- oscillatory schemes are unacceptable
- can we get higher order, non-oscillatory schemes? Yes!

- Euler equation develops shocks physically!
- oscillatory schemes are unacceptable
- can we get higher order, non-oscillatory schemes? Yes!
- upwinding necessary: higher order Godunov methods

- Euler equation develops shocks physically!
- oscillatory schemes are unacceptable
- can we get higher order, non-oscillatory schemes? Yes!
- upwinding necessary: higher order Godunov methods
- MHD, Einstein equations for BH-BH merger

- Euler equation develops shocks physically!
- oscillatory schemes are unacceptable
- can we get higher order, non-oscillatory schemes? Yes!
- upwinding necessary: higher order Godunov methods
- MHD, Einstein equations for BH-BH merger
- Lesson: formal order of accuracy not everything!

- Euler equation develops shocks physically!
- oscillatory schemes are unacceptable
- can we get higher order, non-oscillatory schemes? Yes!
- upwinding necessary: higher order Godunov methods
- MHD, Einstein equations for BH-BH merger
- Lesson: formal order of accuracy not everything!
- Methods respecting physical constraints desirable $(\nabla \cdot \mathbf{B} = 0;$ mass, momentum, energy conservation)

- Euler equation develops shocks physically!
- oscillatory schemes are unacceptable
- can we get higher order, non-oscillatory schemes? Yes!
- upwinding necessary: higher order Godunov methods
- MHD, Einstein equations for BH-BH merger
- Lesson: formal order of accuracy not everything!
- Methods respecting physical constraints desirable $(\nabla \cdot \mathbf{B} = 0;$ mass, momentum, energy conservation)
- Finite Volume Methods to build in conservation laws

- Euler equation develops shocks physically!
- oscillatory schemes are unacceptable
- can we get higher order, non-oscillatory schemes? Yes!
- upwinding necessary: higher order Godunov methods
- MHD, Einstein equations for BH-BH merger
- Lesson: formal order of accuracy not everything!
- Methods respecting physical constraints desirable $(\nabla \cdot \mathbf{B} = 0;$ mass, momentum, energy conservation)
- Finite Volume Methods to build in conservation laws
- numerical instability won't go away with resolution!

■ Lax:
$$u_{j,l}^{n+1} = \frac{1}{4}(u_{j+1,l}^n + u_{j-1,l}^n + u_{j,l+1}^n + u_{j,l-1}^n) - \frac{\Delta t}{2\Delta x}(F_{j+1,l}^n - F_{j-1,l}^n + F_{j,l+1}^n - F_{j,l-1}^n)$$

- Lax: $u_{j,l}^{n+1} = \frac{1}{4}(u_{j+1,l}^n + u_{j-1,l}^n + u_{j,l+1}^n + u_{j,l-1}^n) \frac{\Delta t}{2\Delta x}(F_{j+1,l}^n F_{j-1,l}^n + F_{j,l+1}^n F_{j,l-1}^n)$
- \blacksquare 2-D advection $F_x=v_xu$, $F_y=v_yu$; VNSA with $u^n_{j,l}=\xi^ne^{ik_xj\Delta}e^{ik_yl\Delta}$
- stability condition: $\Delta t \leq \frac{\Delta}{\sqrt{N}|v|}$, N is dimensionality

- Lax: $u_{j,l}^{n+1} = \frac{1}{4}(u_{j+1,l}^n + u_{j-1,l}^n + u_{j,l+1}^n + u_{j,l-1}^n) \frac{\Delta t}{2\Delta x}(F_{j+1,l}^n F_{j-1,l}^n + F_{j,l+1}^n F_{j,l-1}^n)$
- \blacksquare 2-D advection $F_x=v_xu$, $F_y=v_yu$; VNSA with $u^n_{j,l}=\xi^ne^{ik_xj\Delta}e^{ik_yl\Delta}$
- stability condition: $\Delta t \leq \frac{\Delta}{\sqrt{N}|v|}$, N is dimensionality
- an unsplit scheme; one can split by directions

■ FTCS: but a very restrictive stability condition $\Delta t \leq \frac{\Delta^2}{2ND}$

- FTCS: but a very restrictive stability condition $\Delta t \leq \frac{\Delta^2}{2ND}$
- One can generalize Crank Nicolson or implicit but get tridiagonal with fringes; not ideal; iterative techniques

- FTCS: but a very restrictive stability condition $\Delta t \leq \frac{\Delta^2}{2ND}$
- One can generalize Crank Nicolson or implicit but get tridiagonal with fringes; not ideal; iterative techniques
- Alternating-direction implicit (ADI): split into two steps of size $\Delta t/2$ & update x/y implicitly alternately

$$u_{j,l}^{n+1/2} = u_{j,l}^n + \frac{D\Delta t}{2\Delta^2} \left(\delta_x^2 u_{j,l}^{n+1/2} + \delta_y^2 u_{j,l}^n \right)$$

$$u_{j,l}^n = u_{j,l}^{n+1/2} + \frac{D\Delta t}{2\Delta^2} \left(\delta_x^2 u_{j,l}^{n+1/2} + \delta_y^2 u_{j,l}^{n+1} \right)$$

$$\delta_x^2 u_{j,l} \equiv u_{j+1,l} - 2u_{j,l} + u_{j-1,l}$$

■ Two tridiagonal eqs. to solve every $\Delta t!$

Operator splitting in general

 most equations have several terms on RHS; e.g., advection+diffusion

Operator splitting in general

most equations have several terms on RHS; e.g., advection+diffusion

$$\frac{\partial u}{\partial t} = \mathcal{L}(u) = (\mathcal{L}_1 + \mathcal{L}_2 + \dots + \mathcal{L}_m)u$$

$$u^{n+(1/m)} = \mathcal{U}_1(u^n, \Delta t) \qquad u^{n+1/m} = \mathcal{U}_1(u^n, \Delta t/m)$$

$$u^{n+(2/m)} = \mathcal{U}_2(u^{n+(1/m)}, \Delta t) \qquad u^{n+2/m} = \mathcal{U}_2(u^{n+1/m}, \Delta t/m)$$

$$\dots \qquad \dots$$

$$u^{n+1} = \mathcal{U}_m(u^{n+(m-1)/m}, \Delta t) \qquad u^{n+1} = \mathcal{U}_m(u^{n+(m-1)/m}, \Delta t/m)$$

• update one operator at a time for full Δt OR full operator for smaller times (ADI)

Operator splitting in general

 most equations have several terms on RHS; e.g., advection+diffusion

$$\frac{\partial u}{\partial t} = \mathcal{L}(u) = (\mathcal{L}_1 + \mathcal{L}_2 + \dots + \mathcal{L}_m)u$$

$$u^{n+(1/m)} = \mathcal{U}_1(u^n, \Delta t) \qquad u^{n+1/m} = \mathcal{U}_1(u^n, \Delta t/m)$$

$$u^{n+(2/m)} = \mathcal{U}_2(u^{n+(1/m)}, \Delta t) \qquad u^{n+2/m} = \mathcal{U}_2(u^{n+1/m}, \Delta t/m)$$

$$\dots \qquad \dots$$

$$u^{n+1} = \mathcal{U}_m(u^{n+(m-1)/m}, \Delta t) \qquad u^{n+1} = \mathcal{U}_m(u^{n+(m-1)/m}, \Delta t/m)$$

- update one operator at a time for full Δt OR full operator for smaller times (ADI)
- first order accurate in time; can be made 2nd by Strang splitting $(\mathcal{A}/2)\mathcal{B}(\mathcal{A}/2)$

lacktriangle operators $\mathcal{A} \& \mathcal{B}$ do not commute in general!