- Fourier Transforms.
- Discrete Fourier Transform.

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- Allows one to break down functions/signals into their component parts and analyze, smooth or filter them.
- Also allows one to perform certain kinds of calculations and solve certain differential equations.

#### Fourier Series

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If the function is odd (i.e. antisymmetric) about the midpoint  $(x = \frac{L}{2})$  then one can write the sine series:

$$f(x) = \sum_{k=1}^{\infty} \beta_k \sin\left(\frac{2\pi kx}{L}\right)$$

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- If the function in not periodic, and we are only interested in a portion of this non periodic function over a finite interval, 0 to L, we can just take that portion and repeat it to create a periodic function!
- Then the Fourier coefficients will only give the correct information about the function in the interval 0 to L. Outside this interval, the function will be just repeated (and may not have anything to do with the original function.

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If  $k' \neq k$ :

$$\int_{0}^{L} \exp\left(i\frac{2\pi(k'-k)x}{L}\right) dx = \frac{L}{i2\pi(k'-k)} \left[\exp\left(i\frac{2\pi(k'-k)x}{L}\right)\right]_{0}^{L}$$
$$= \frac{L}{i2\pi(k'-k)} \left[e^{i2\pi(k'-k)} - 1\right]$$
$$= 0$$

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Thus, given a function f(x), we can find the Fourier coefficients  $\gamma_k$ , or given the coefficients, we can find the function f(x) – we can go back and forth freely between the function and the Fourier coefficients.

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- There are, however, many cases where this is not doable the integral is not doable because the function is too complicated or the function f(x) may not even be known in analytic form (for eg. if it is a signal measured in the laboratory experiment).
- In such cases, the integral can be evaluated numerically.

Applying the trapezoidal rule for integration (N slices of width h=L/N) to calculate  $\gamma_k$ :

$$\gamma_k = \frac{1}{L} \frac{L}{N} \left[ \frac{f(0)}{2} + \frac{f(L)}{2} + \sum_{n=1}^{N-1} f(x_n) \exp\left(-i\frac{2\pi k x_n}{L}\right) \right]$$

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This formula can be used to evaluate the coefficients on a computer. A simpler way to write this is as:

$$\gamma_k = \frac{1}{N} \left[ \sum_{n=0}^{N-1} y_n \exp\left(-i\frac{2\pi kn}{N}\right) \right]$$

where  $y_n = f(x_n)$ 

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The quantities  $\gamma_k$  and  $c_k$  only differ by the constant 1/N factor. For our purpose they are both equal, and we define the latter as the definition of DFT.

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$$\begin{split} \sum_{k=0}^{N-1} c_k \exp\left(i\frac{2\pi kn}{N}\right) &= \sum_{k=0}^{N-1} \sum_{n'=0}^{N-1} y_{n'} \exp\left(-i\frac{2\pi kn'}{N}\right) \exp\left(i\frac{2\pi kn}{N}\right) \\ &= \sum_{n'=0}^{N-1} y_{n'} \sum_{k=0}^{N-1} \exp\left(i\frac{2\pi k(n'-n)}{N}\right) \\ &= \sum_{n'=0}^{N-1} y_{n'} \frac{\exp\left(i2\pi(n'-n)\right) - 1}{\exp\left(i2\pi(n'-n)/N\right) - 1} = \sum_{n'=0}^{N-1} y_{n'} N\delta_{n,n'} \\ &= Ny_n \quad \text{assuming} \quad 0 \le n \le N \end{split}$$

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Equivalently,

$$y_n = \frac{1}{N} \sum_{i=1}^{N-1} c_k \exp\left(i\frac{2\pi kn}{N}\right)$$

This is called the "Inverse Discrete Fourier Transform."

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- It is important to appreciate that unlike the original Fourier series, the discrete version only gives us the sample values at  $y_n = f(x_n)$ . It tells us nothing about the value of the function f(x) in between the points.
- So, two different functions with same values at the sample points will have the same DFT – no matter what they do in between the points!

#### DFT for real functions

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$$c_{N-r} = \sum_{n=0}^{N-1} y_n \exp\left(-i\frac{2\pi(N-r)n}{N}\right)$$
$$= \sum_{n=0}^{N-1} y_n \exp(-i2\pi n) \exp\left(i\frac{2\pi rn}{N}\right)$$
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Thus:  $c_{N-1}=c_1^* \ c_{N-2}=c_2^*$  and so forth. That means that Fourier coefficients  $c_k$  of a real function only has to be calculated for  $0 \le k \le \frac{N}{2}$ .

### Discrete Fourier Transform

```
from numpy import zeros
from cmath import exp, pi
def dft(y):
    N = len(y)
    c = zeros(N//2 +1, complex)
    for k in range (N//2+1):
        for n in range(N):
            c[k] += y[n]*exp(-2j*pi*k*n/N)
return c
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One thing to notice about the DFT is that we can shift the sample points along the x-axis, and not much changes. Suppose:

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Then the DFT is:

$$c_k = \sum_{n=0}^{N-1} f(x_n + \Delta) \exp\left(-i\frac{2\pi k(x_n + \Delta)}{L}\right)$$
$$= \exp\left(-i\frac{2\pi k\Delta}{L}\right) \sum_{n=0}^{N-1} f(x'_n) \exp\left(-i\frac{2\pi kx_n}{L}\right)$$
$$= \exp\left(-i\frac{2\pi k\Delta}{L}\right) \sum_{n=0}^{N-1} y'_n \exp\left(-i\frac{2\pi kn}{N}\right)$$

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But this is the same as the original DFT except for a (k-dependent) phase factor. Thus the DFT is really independent of where we choose to place the samples – only the coefficients change by a phase factor.

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Suppose that we have an  $M \times N$  grid points of samples  $y_{mn}$ . We first transform on each of the M rows:

$$c'_{ml} = \sum_{n=0}^{N-1} y_{mn} \exp\left(-i\frac{2\pi ln}{N}\right)$$

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Then we take the  $l^{th}$  coefficient of each row and Fourier transform them :

$$c_{kl} = \sum_{m=0}^{M-1} c'_{ml} \exp\left(-i\frac{2\pi km}{M}\right)$$

Alternatively, we can write a single expression for the complete Fourier transform in two dimensions:

$$c_{kl} = \sum_{m=0}^{M-1} \sum_{n=0}^{N-1} y_{mn} \exp\left[-i2\pi \left(\frac{ln}{N} + \frac{km}{M}\right)\right]$$

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The corresponding inverse transform is:

$$y_{mn} = \frac{1}{MN} \sum_{k=0}^{M-1} \sum_{l=0}^{N-1} c_{kl} \exp \left[ i2\pi \left( \frac{ln}{N} + \frac{km}{M} \right) \right]$$