# Chapter 1

**INTRODUCTION**

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Iris recognition is a mature surveillance technique that has many desirable properties of an ideal biometric. However, the inability accurately recognition irises for subjects positioned outside a small imaging volume is still an unsolved problem [1–3]. Assuming the iris acquisition system has sufficient magnification to resolve, this boundaries of the imaging volume is established by the Depth of Field (DOF) of the lens. Owing to the physics of light (explained shortly), lenses can focus perfectly only on a single surface—the plane of sharp focus—in the object space. In most imaging systems this surface is a *plane* due to the planar structure of the image sensor. The DOF is commonly defined as the region fore and aft the plane of sharp focus (in the object space) within which objects appear sharp in the image. This limitation renders current iris recognition systems to be used in highly controlled environments. It would be possible to use iris recognition systems in unconstrained environments such as public places, crowded hallways, public transit stations, etc. if its DOF can be extended.

The DOF limitation is not restricted only to iris acquisition systems. The inability of a camera to capture sufficiently high quality image outside of a small imaging volume is in fact a general problem of all imaging systems. This limitation is fundamental, in a sense that, it is imposed by the wave nature of light. The DOF problem is illustrated in [Figure 1.1](#Figure_1_1). The figure shows that the camera fails to acquire sharp images at all depths in a single image capture.

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| **Figure 1.1** Depth of field (DOF) problem. Image of the three human-figure cut-outs with sinusoidal patterns (2 *lp/mm*) and artificial irises and placed apart by 11 *cm* from each other. The camera, with lens of 80 mm focal length and F/5 aperture, was focused on the middle cut-out (3.6 *m* away from the camera). It is evident that the spatial resolution in the image falls off rapidly with increasing distance from the plane of sharp focus (middle cut-out) inhibiting the camera from resolving fine details uniformly across the imaging volume. |

This thesis aims to study the DOF problem pertinent to iris acquisition systems and propose a method to improve the DOF by combining Scheimpflug imaging techniques with computational imaging.

Scheimpflug imaging is a classical technique used by landscape photographs to produce photographs in which both near (foreground) objects and far (background) objects are in sharp focus.

Computational imaging is a set of modern techniques in which computation is an inherent part of the image formation unlike in traditional imaging where digital processing is applied to enhance an image after it has been captured. Computational imaging techniques such as integral imaging, light field capture, wavefront coding, focus stacking, etc. are generally used to overcome specific problems in imaging.

The problem for this thesis is stated in Section 1.1. A brief overview of iris recognition is presented in Section 1.2. Section 1.3 describes some of the desirable properties of iris recognition systems within the context of imaging. A short background on scheimpflug imaging and computational imaging are presented in Section 1.4.

### **1.1 Introduction**

Iris recognition, which authenticates a person’s identity by examining the structural pattern of the iris, is a key biometric (biological measurement) [4,5]. It is an ideal biometric because of its high accuracy, robustness and uniqueness of the iris pattern  [6]. As iris authentication does not require physical contact of the subject and does not entail human intervention in the process of the recognition tasks, it has the potential for use in overhead surveillance in crowded places. Unfortunately, iris recognition has so far only experienced limited success in such environments. Most systems, currently in operation, have stringent requirements for operation. For example, a subject is required to position his/her head in front of an imager placed at a pre-specified distance (standoff) for acquiring iris images. Such requirements also restrict its deployment in large scale environments, such as in crowded airports and busy railway stations, where there is a need for rapid authentication within a large volume. One of the crucial reasons for its use in restrictive and fully-cooperative environments is the inability of an imager to acquire sufficient quality image within a large range in a single exposure [1–3]. There is a strong desire for iris acquisition devices to have large depth-of-field [7–9]. Current solutions to improve the depth-of-field of iris imagers such as using multiple cameras, or wavefront coding are generally expensive economically and/or computationally [9], plagued by noise at higher spatial frequencies [1].

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| **Problem Statement** |
| The purpose of this study is to theoretically and experimentally explore the feasibility of improving the depth-of-field of iris acquisition systems using a combination of traditional scheimpflug imaging and modern computational imaging methods. |

Perfect imaging corresponds to the ability of an imager to produce a scaled replica of an object in the image space [10]. When only a small portion of the light wave emerging from an infinitesimally small point of light is collected through a finite opening of a camera’s aperture (Figure 1.2 (a)), the replica in the image space is not exact even in the absence of aberrations; instead, the image of the point spreads out in space due to diffraction at the aperture. This dispersed response in the three-dimensional image space is called *Point Spread Function* (PSF). The spreading of the PSF along the transverse direction (a 2D PSF) restricts an imager’s ability to resolve fine details (spatial frequency) on the image. For an extended object, which is made of several points, the 2D PSF smears the responses from neighboring points into each other causing blur. Similarly, the spread along the longitudinal direction limits the ability to discriminate points staggered closely in the direction of the optical axis causing a region of uncertainty; however, the extension of the 3D PSF along the optical axis enables multiple spatially-separated objects (or points) within a volume in the object space to form acceptably sharp images at once. Conversely, an object in the object space may be placed anywhere within this zone and still form a satisfactory

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| **Figure 1.2** Incoherent impulse response and DOF. (a) The image A’ of a point source A spreads out in space forming a zone of tolerance called Depth of Focus (DOF) in the image space; (b) The normalized intensity distribution of the 3D PSF of a 25 *mm*, F/5 lens imaging an axial point source at a distance of 100 *mm*. The expression for the 3D PSF was obtained for a circular aperture using scalar diffraction theory and paraxial assumption. The DOF, having prolate spheroidal shape, is defined as the region within which the intensity has above 80% of the intensity at the geometric focus point. The figure shows iso-surfaces representing 0.8, 0.2, 0.05 and 0.01 intensity levels. The ticks on the left vertical side indicate the locations of the first zeroes of the Airy pattern in the focal plane. The vertical axis has been exaggerated by 10 times in order to improve the display of the distribution. |

image. This zone of tolerance in the object space is called depth of field. The corresponding zone in the image space is called depth of focus [11]. In this thesis, the acronym “DOF” is used for both depth of field and depth of focus wherever its meaning shall be apparent from the context. In the image space, the DOF is defined as the region of the 3D PSF where the intensity is above 80% of the central maximum [12,13]. This zone is in the shape of a prolate spheroid. In the absence of aberrations, the maximum intensity occurs at the geometric focal point[[1]](#footnote-1), , where contributions from all parts of the pupil are all in phase. Figure 1.2 (b) shows the aberration-free intensity distribution, , as a function of defocus  about the geometric focal point for a light source placed at 100 *mm* from a lens of focal-length of 25 *mm* and aperture diameter of 5 *mm*. The expression for the distribution—normalized to make  equal to unity—is obtained using scalar diffraction theory and paraxial assumptions. The derivation of the expression is shown Appendix C.

The shape—length and breadth—of the 80% intensity region (**Figure 1.2(b)**) dictates the quality of the image acquired by an imager in terms of lateral spatial resolution and DOF. A first order optical simulation demonstrating the effect of the DOF in image acquisition at varying depths is shown in Figure 1.3. For this simulation a 100 mm focal length, f/5 lens that is focused at 1300 mm is used. In this setup, the imager has a DOF of 9.5 mm in the object space (calculated by applying the lens equation to the extremes of the DOF in image space). As it can be seen, the irises located outside the DOF region are severely blurred. It has been shown in [] that the performance of iris recognition deteriorates quickly with increasing amounts of defocus in the captured iris images.

Performance of iris recognition systems for acquiring iris images within a large volume in a single exposure is dependent on the extent of DOF. Since the very early days of landscape photography, scheimpflug imaging techniques have been used to adjust the orientation of the plane of sharp focus. The DOF, which is a prescribed zone around the plane of sharp focus, can thus be modified to bring a slanted object completely in focus using scheimpflug imaging. Recently, computational imaging techniques such as wavefront coding and pupil engineering have also been used to transform the distribution of light near the focal region for both extending the DOF and improving axial resolution (curtailing the DOF). It is posited in this work that an effective solution to extending the imaging volume of iris acquisition cameras may be realized by supplementing modern computational imaging techniques to well-known traditional methods of scheimpflug imaging. The following section presents a condensed overview of iris recognition.

### **1.2 Primer on iris recognition**

The human iris is the colored portion of the eye having a diameter which ranges between 10 mm and 13 mm []. The iris is perhaps the most complex tissue structure in the human body that is visible externally. The iris pattern has most of the desirable properties of an ideal biomarker, such as uniqueness, stability over time, and relatively easy accessibility. Being an internal organ, it is also protected from damage due to injuries or intentional fudging []. The presence or absence of specific features in the iris is largely determined by heredity (based on genetics); however, the spatial distribution of the cells that form a particular iris pattern during embryonic development is highly chaotic. This pseudo-random morphogenesis, which is determined by epigenetic factors, results in unique patterns of the irises in all individuals including that of identical twins []. The diverse microstructures in the iris that manifest at multiple spatial scales [] are shown in Figure 1.4. These textures, unique to each eye, provide distinctive biometric traits that are encoded by an iris recognition system into distinctive templates for the purpose of identity authentication. Iris color is not used as a biomarker since it is determined by genetics, which is not sufficiently discriminative.

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| **Figure 1.3** First order simulation of iris acquisition at multiple depths. The letters A, B &C (A’, B’&C’) denote both sources (images) and positions. The 1st row depicts three sources at three depths from a 100mm, f/5 lens. Point A (in focus) forms image A’. The 2nd row shows a 12mm iris in the object space (left) and the |3D PSF|2 of the source A (right). The |3D PSF|2 for objects B and C are very similar to that of A since their relative separation is trivial compared to their distances from the lens. The image plane, at geometric focus, senses the 3D PSF at A’, B’ & C’ respectively for points A, B & C. Positions B’ & C’, which are outside the DOF region in the image space, are 123.2 and 177.3 μm from A’ respectively. The corresponding incoherent 2D PSFs are shown in the 3rd row. The iris images in the 3rd row were obtained by convolving the incoherent 2D PSFs with the de-magnified iris image in 2nd row. |
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| **Figure 1.4** Complexity and uniqueness of human iris. Fine textures on the iris forms unique biometric patterns which are encoded by iris recognition systems. (Original image processed to emphasize features). |

# Chapter 2

**BACKGROUND**

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## **2.1 State-of-the-art large standoff iris acquisition**

One of the earliest systems for iris recognition at a large distance was reported by Sarnoff [] in 2005. It used two custom designed elliptical mirror telescopes with diameters of 101.6 *mm* and 203.2 *mm* to demonstrate the feasibility of iris recognition at large standoff distances of 5 *meters* and 10 *meters* respectively. The magnifications in the telescopes were sufficient to produce iris images with pixel resolution of 128 pixels for non-moving subjects placed at the said distances. For the Hitachi KP-M2RN infrared camera that has a pixel pitch of , and assuming 11 *mm* for the average iris diameter, back-of-the-envelope calculations (not published in the paper) suggest a focal length of around 585 *mm* and 1170 *mm* respectively for operating distances of 5 meters and 10 meters. Using , the geometrical DOF is calculated to be approximately 7.6 *mm*. In order to capture sufficient quality images at such low DOF and at such large distances, subjects’ heads were supported using a chin rest and the image acquisition setup was constructed on a moderately large optical table []. In order to improve the odds of capturing at least one good quality iris image, a sequence of video frames at 12 *fps* was used. The measured system MTF in the object space for

both distances indicated a spatial resolution of 2 *lp/mm* at 80% modulation strength. Since both the F-number and the magnification were kept constant for both 5 *m* and 10 *m* distances the optical performance was impacted by the doubling of the distance. Although, such a system is hardly suitable for deployment, the study demonstrated the feasibility of iris recognition at large standoff distances.

Sarnoff built another long distance system in 2007 using a Meade LX200-R F/10 8 *inch* reflecting telescope operating at 850 nm that generated useful images at 15+ *meters* []. An 18 m standoff distance iris image capture device utilizing adaptive optics was demonstrated by AOptixTM in 2007 [].

The HBOXTM from Global Rainmakers [] provided the ability to capture multiple subjects walking through a portal system at a distance of about 1.5 meters. The Iris On the Move (IOM) [] portal system developed by Sarnoff in 2008 is capable of acquiring iris images at distances of approximately 2-3 *meters* with the subject walking at normal pace. It uses an off-the-shelf 210 mm lens and a camera with pixel pitch of 7.4. The IOM system uses a focusing mechanism along with video capture to acquire iris images. The reported capture volume is at least 100x100x5 *mm*3 for a pixel resolution of 200 *pixels/cm* and optical resolution good enough for iris recognition. Also, it is to be noted that the DOF was not calculated based on a given definition of the circle-of-confusion; rather it was empirically found by identifying the region within which the recognition algorithm produced a Hamming distance of below 0.33 when matching templates and iris codes from the same eye. A general disadvantage of portal based systems is that the user movement is restricted to a narrow zone. Such restrictions may prevent these systems from being deployed in crowded locations.

A combined face and iris recognition system (CFAIRS) developed by Honeywell in 2008, which uses multiple FOV cameras, is capable of acquiring face and iris images of multiple subjects within a range of 1 to 5 meters []. Another multi-biometric system, the Retica’s Eagle-EyesTM, which too uses multiple cameras and video tracking, has a capture volume of 3x2x3 *m*3 and a standoff distance of 3 meters [].

Another iris recognition system at a standoff distance of about 3 meters was shown by Dong et al. in []. It uses a 300 mm lens set at F/15 mounted on a pan-tilt-unit, and it has a capture volume of 20x15x10cm3 for acquiring iris images with greater than 150 pixels across. Boehnen et al. [] demonstrated the Standoff Multimodal biometric system that is capable of acquiring iris images from a distance of 7 meters. It uses a multi-camera approach to increase the capture volume; the iris camera is a 150-500 mm f/5-6.3 zoom lens

More recently, a large standoff iris acquisition system with a pixel resolution of 200 pixels on the iris at a distance of 4-8 meters using an 800 mm COTS auto-focus camera has been demonstrated by Venugopalan et al. in []. Stoker et al. demonstrated an outdoors iris recognition system operating at a distance exceeding 25 meters using a 10 inch (254 mm) Ritchey-Chretien telescope, fitted with several custom optical elements for aberration correction and collimation []. Image restoration techniques are used to improve the quality of images degraded by atmospheric turbulence.

# Chapter 3

⬀⬀⬀**MODEL OF SCHEIMPFLUG IMAGING – I: PROPERTIES OF IMAGE**

*Essentially, all models are wrong, but some are useful.*

—George Box

For our investigation of the use of a Scheimpflug camera for extending the imaging volume of iris cameras, we require a model that includes the essential characteristics of Scheimpflug imaging. Although there are existing models, several of them use the thin lens approximation that is over-simplistic. For example, it is impossible to accurately describe the shift in the image field due to lens rotation using a thin lens model. On the other hand, models that do use thick-lens abstraction, while accurate, do not explicitly consider the effects of the pupils (defined later) on image formation. The lack of pupil parameters in these models makes it hard to predict the geometric properties of the image obtained using a Scheimpflug camera in which the lens and image planes are free to rotate about independent pivots. Therefore, we develop a new model that depicts the explicit dependence of the pupils on the properties of the image in Scheimpflug cameras. We have broken down the modeling process into two chapters—in Chapter 3 (this chapter) we develop the relationship between a world point and its image in a Scheimpflug camera, verify the model using optical ray tracing in Zemax, and finally study the consequence of rotating the sensor and lens on

the geometric properties of the image using our model. In Chapter 4 we derive object-image equations that relates the angles and directed distances between the object, lens and sensor planes for which the object plane is brought to focus on the sensor plane.

Conventional imaging systems consists of a lens and an *image plane* on which a sharp image of an *object plane* is formed. The object and image planes are called *conjugates*. Further, amongst the several possible planes that are perpendicular to the optical axis and pass through the lens, we designate the one as the *lens plane* that provides some advantage in the geometric model. For example, the plane through the lens center and the plane through the object-side principal point are designated as the lens planes in the thin-lens and thick-lens models respectively. In our model, the lens plane is the plane through the center of the entrance pupil (defined later).

The plane in the object space that is in sharp focus is called the Plane of Sharp Focus (PoSF). In the *fronto-parallel* configuration used in most camera designs, the lens and the sensor planes are parallel to each other and perpendicular to the optical axis. In such designs, the physics of optical imaging—described by the *Gaussian Formula*—dictates that the Plane of Sharp Focus must be parallel to the lens and image planes. In contrast, the lens and image planes in a Scheimpflug camera are free to swivel about their pivots (as shown in [Figure 3.1](#Figure_3_1))resulting in a corresponding swivel of the PoSF. We exploit this feature—the freedom to arbitrarily orient the PoSF—of Scheimpflug imaging to improve the depth of field of the iris acquisition devices. However, the degrees of freedom offered by the Scheimpflug camera comes at the cost of added complexity of operation. Therefore, a rich description of such cameras requires the development of a general model. In this chapter, we initiate the development of such a model of Scheimpflug imaging starting from the axioms of *geometric optics* (*ray optics*).

Assumptions are crucial and necessary for modeling that enable its expediency but limits its applicability. For the model described herein, we assume paraxial imaging, rotational symmetry and aberration-free optics in order to make the problem tractable. Additionally, we assume the refractive index of the lens elements and the interstitial medium to be isotropic (uniform along all directions) and homogeneous (uniform at all positions); this assumption imposes rectilinear propagation of light. Further, we assume the lens is surrounded by air (of refractive index equal to

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| **Figure 3.1** Scheimpflug camera movements. Insets (a), (b), (c), and (d) depict key camera movements—lens and sensor plane tilts about a horizontal axis—amongst several possible, in a Sinar P3 camera. Labels indicate the lens (1), the lens standard (2), bellows (3), sensor standard (4), and the sensor (5). The cyan lines on the two standards accentuates the orientations. Inset (e) is a superimposed sequence of images of the camera with the two standards in a multitude of orientations. The physical locations of the two pivots emerge at the intersection of the superimposed cyan lines. |

one). Consequently, the front and back focal lengths are equal, and the two nodal points coincide with the corresponding principal points.

In the following sections of this chapter we derive a general geometric imaging model that allow the both the sensor and a thick lens to be oriented arbitrarily about their own pivots, and directly incorporate important pupil related parameters in the model. We then verify the accuracy of the model using ray-tracing in Zemax. Following the verification, we study the consequences of the model and note some of the important predictions that will ultimately allow us to build an extended depth-of-field imaging system.

### **3.1 Introduction**

Optical imaging systems consist of several groups of elements; those elements endowed with optical power bends rays of light. The tiniest orifice in the system is called the *system aperture* or *stop*. Its interaction with the elements in the system gives rise to the pupils.

*Pupils* are the sine qua non of optical systems. They are indispensable in the design and specification of all optical systems, in both domains of *ray* and *wave optics*. The *entrance pupil* () is the image of the stop seen through the elements preceding it is. The *exit pupil* () is the image of the stop seen through the elements following it is. The region preceding the entrance pupil, which includes the objects and light sources, is called the object space; and the region following the exit pupil, which includes the image plane, is called the image space. The size and position of the stop (and hence the pupils) affect image resolution, aberration, brightness, and geometry.

Rotationally symmetric lenses have an axis of symmetry—the optical axis. A ray coincident with the optical axis traverses undeviated through the lens. Planes passing through the axis of such lenses are the meridional planes. Rays restricted to the meridional planes are *meridional rays*. Patterns formed by the meridional rays on either side of the optical axis are mirror-reversed, exhibiting bilateral symmetry. [Figure 3.2](#Figure_3_2) shows two types of meridional rays, traced in Zemax, that are fundamental to geometric analysis [18,19]. The *marginal ray* (MR) originates from the axial object position and skirts the edges of the aperture and pupils (virtually); the *chief ray* (CR) starts at an off-axis object point and pierces the centers of the aperture and pupils[[2]](#footnote-2) (virtually). This pair of rays determines the location and size of the pupils, the position of the image, and the magnification. Furthermore, the bundle of chief rays from the object space converge at the center of the entrance pupil—thus *homocentric*—forming the vertex of the object-space perspective cone; in the image space, the bundle of chief rays diverges from the center of the exit pupil producing the vertex of the image-space perspective cone.

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| **Figure 3.2** Fundamental rays (contained within the meridional plane) and pupils in a Double Gauss lens for an object at infinity. The chief rays—close to the optical axis (0°, ±5° in the object space at entrance pupil)—appear to converge at the center of entrance pupil and diverge from the center of exit pupil. The marginal ray, which is parallel to the optical axis since the object is at infinity, appear to skirt the edges of the two pupils. The red circles specify the vertices of the perspective cones (centers of the pupils). The rays were traced in Zemax. |

Imagine a film projector working backwards. Imagine the stream of light rays flowing from the illuminated portion of the scene towards a small circular hole in the projector. This pencil of rays forms a conical volume of light—the perspective cone—with its vertex at the hole and its base towards the scene. The “illuminated portion” is the angular extent of the scene visible in the image, confined by the circumferential chief rays. These extreme chief rays determine the opening angle of the cone. The “small hole” represents the entrance pupil of a camera or the pupil at the center of the iris in an eye. In the image space (behind the hole), the ray-pencil form another cone with the vertex at the center of the exit pupil. This image-space perspective cone projects the light from the scene onto the film surface or the retina in the eye. This process of image formation, known as the *central projection*, is fundamental to all imaging systems—inanimate and animate—including the camera and the eye. While the opening angle of the object-space perspective cone determines the field-of-view, its counterpart in the image space determines the angular dimension of the image. The ratio of the pupil sizes (pupil magnification) determines the relationship between the image and object-space opening angles of the two perspective cones [21,22].

### **3.2 Notations**

* *Coordinate system*: Right-handed, with the positive -axis oriented along the direction of light travel (left to right in the plane of drawing).
* *Scalars*: written as small letters (e.g. ). *Vectors*: small bold letters (e.g. ). *Matrices*: written as capital letters (e.g. ).
* *Object and image space*: Non-primed quantities represent object space (e.g. , ). Primed quantities represent image space (e.g.  , ).
* *Unit vectors*: represented using a hat (). The only exception is the direction cosine vectors which have unit norm, but are represented without the top hat (e.g. .
* *Reference frames*: Coordinate reference frames are denoted using curly brackets (e.g. represents the camera coordinate frame). Note that if the camera is not translated in three-dimensional space (as is the case in this thesis), a separate world frame is redundant. Therefore, in this thesis, we represent world points with respect to the camera frame .
  + A left superscript on a variable indicates the frame of reference. For example, indicates that the variable is expressed respect to the world coordinate frame. If no reference is explicitly stated, then variable is expressed with respect to the world coordinate frame (or the camera coordinate frame if the camera coordinate frame and the world coordinate frame are the same).
* A subscript is used to associate a variable with a specific optical parameter. For example, the subscript in variable is used to represent the location of the entrance pupil, and the subscript in is used to represent the rotation matrix applied to the lens plane in the camera frame, etc. If the camera coordinate frame is the same as the world coordinate frame, then the notation shall be used (the superscript is dropped for convenience).
* Indices into matrices and vectors start from one (1). For example, is the third element of the unit normal vector , and is the element in the first row and second column of the matrix . Further, a matrix is represented in terms of its columns as where are the columns of .

### **3.3 Relation between pupil magnification and chief ray angle**

The *pupil magnification* is defined as the ratio of the paraxial exit pupil diameter to the paraxial entrance pupil diameter [21,23,24]ref.

[Figure 3.3](#Figure_3_3) illustrates the meridional and sagittal planes associated with an arbitrarily located object of height above the optical axis and its image of height in a typical optical system. The figure also shows the chief ray from the object’s edge further from the optical axis, the marginal ray from the axial point in the object, and the two pupils contained in the meridional plane. The schematic, although simple, is quite general as a (meridional) plane always exist for a given object point irrespective of its position in the three-dimensional space, if the lens is rotationally symmetric.

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| **Figure 3.3** Schematic of chief and marginal rays. The ratio of the tangents of the chief ray angles in the object space to the image space yields the pupil magnification. |

Let the angles between the chief ray and the optical axis (called the *ray-angle*) in the object and image space be and respectively. Also, let the angles produced by the marginal ray with the optical axis in the object and image space be and respectively. Then, we can obtain the relation between the chief-ray ray-angles— and —and the pupil magnification as follows:

From the [Figure 3.3](#Figure_3_3) we obtain,

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Eliminating and after dividing by, we have

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A common observation in imaging is that we can increase the transverse magnification () by increasing the lens-to-image-plane distance while correspondingly decreasing the lens-to-object-plane distance in order to maintain focus on the object. However, increasing (decreasing) the image plane distance proportionally decreases (increases) the marginal ray angle (see [Figure 3.3](#Figure_3_3)). Consequently, the angular magnification () decreases with increase in lens-to-image-plane distance. Therefore, a large transverse magnification is associated with a correspondingly small angular magnification. This result follows from a more general theory called the *Lagrange invariant* property [21] of the two rays (the chief ray and the marginal ray) when applied between conjugate locations. As per the invariant property, the product of the transverse magnification and the angular magnification equals to one, i.e., or. Cancelling the corresponding terms in Eq. (3.2) yields the relationship between the pupil magnification and the object and image space chief ray angles as:

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where,

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|  | Pupil magnification. |
|  | Angle between the chief ray and the optical axis in the object (input) space. |
|  | Angle between the chief ray and the optical axis in the image (output) space. |

The above relation (Eq. (3.3)) has been previously derived in [24] using a different formulation.

For a given optical system with fixed focal length, the pupil magnification is constant. This constancy of the ratio of the tangents of the chief ray angles for varying object (and image) heights is a necessary and sufficient condition for distortion-free imaging known as the *Airy’s Tangent-Condition* [20,24]. Eq. (3.3) also suggests that when the perspective cones in the object- and image-space are symmetric. In the following section, we will use Eq. (3.3) to derive the relationship between the direction cosines of the object-space (input) chief rays and direction cosines of the image-space (output) chief rays.

### **3.4 Transfer of chief ray’s direction cosines between the pupils**

The direction cosines, a unit vector of cardinality three, specify the direction of a ray. Its elements are the cosines of the angles the ray makes with the three coordinate axes. In other words, the elements of the direction cosine vector are the projections of the unit vector in the direction of the ray on the -, -, and -axes. In the absence of aberrations, the chief ray starts from an object point, passing through the center of the entrance pupil (virtually), the aperture stop, exit pupil (virtually), and ends at the image point. Suppose we know the direction cosine of the chief ray in the object space (between the object point and the entrance pupil), what is the direction cosine of the chief ray in the image space (between the exit pupil and image point)? Furthermore, what is the relation between the input and output chief ray’s direction cosines if the lens is swiveled about a pivot point along the optical axis?

We begin by solving a specific problem of the *transfer* of the direction cosines between the pupils in which the optical axis coincides with the -axis of the camera frame. The configuration of this specific problem is show in [Figure 3.4](#Figure_3_4). Subsequently, we will deduce the general *transfer* expression in which the optical axis is free to swivel about the origin of. Let be the direction cosine of the chief ray from an object point to the center of the entrance pupil, and let be the corresponding direction cosine of the chief ray from the exit pupil to the image point. The parameters , ,and are specified with respect to frame .

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| **Figure 3.4** Specific problem—optical axis coincides with reference frame’s -axis. If and are the angles of the chief ray with the optical axis in the object and image space respectively, then and. |

If and are the zenith and azimuthal angles of the chief ray in the object space, and and the corresponding angles in the image space, then the direction cosines, in , are:

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Since the optical axis is aligned with the -axis, we have and . Substituting the expressions for from Eq. (3.4) into Eq. (3.3) we obtain:

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Further, since the input and output chief rays are confined to the same meridional plane [18,19], , yielding and in terms of and , the ratios of to , and :

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From (3.3) we have

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which after simplification yields in terms of the pupil magnification and input as

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Combining Eqs. (3.6) and (3.8), we obtain the expression for output direction cosine of the chief ray in terms of its input direction cosines and the pupil magnification as:

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which we can write compactly as:

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Our objective is to derive the expression for the transfer of the chief ray’s direction cosines from entrance pupil to exit pupil for arbitrary orientation of the optical axis as shown in [Figure 3.5](#Figure_3_5). Although a formal derivation is provided in [Appendix A.1](#_Appendix_A.1_Transfer), we can readily infer the general expression for the *transfer* from Eq. (3.10). Suppose we swivel the optical axis about the origin of the camera frame. This rotation can be described by the matrix. As before, we designate the ray from the object point to the (new position of the) center of the entrance pupil as the chief ray. Let us also suppose that we have another coordinate frame,, sharing its origin with and its -axis coincident with the optical axis. If be the direction cosine of the chief ray from the object point in the frame , then the direction cosine in the frame is and the third element of the direction cosine is , where is the third column of . Representing, the direction cosine of the chief ray emerging from the exit pupil is obtained by substituting for and for in Eq. (3.10):

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The above expression represents the output direction cosine in the coordinate frame In order to transform the output direction cosine from the coordinate frame to the camera frame we need to multiply the direction cosine vector by to obtain:

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| **Figure 3.5** Configuration of the general problem—optical axis pivots freely about the origin of camera frame. |

The positive or negative sign of the direction cosine determines the forward or backward direction of light-travel along a rectilinear path. Under the assumptions of isotropy and homogeneity, the only condition under which a ray of light emerges in an antipodal path from an interface is if it encounters a mirror surface *normally*. This condition does not arise within the context of our problem. Therefore, without any loss of generality, we can drop the negative sign in Eq. (3.11); accordingly, the output direction cosines assume the sign of the corresponding input direction cosines. Therefore, we obtain the general expression for the direction cosines of the chief ray in the image space as:

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where,

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|  | The third element of the input direction cosine vector following rotation. |
|  | Input or object space chief ray’s direction cosine vector from the object point to the entrance pupil. |
|  | Output or image space chief ray’s direction cosine vector emerging from the exit pupil. |
|  | Equal to , where is the pupil magnification. |
|  | Rotation matrix used to describe the orientation of the lens plane. |
|  | The third column of . |

Note that Eq. (3.12) only describes the output chief ray’s direction cosines—a free vector. The exact output chief ray is obtained from the knowledge of the direction cosine and the location of the exit pupil in the appropriate reference frame.

Although it is not obvious from the expression Eq. (3.12), we expect to have unit magnitude. We present a proof in [Appendix A.2](#_Appendix_A.2_The) that shows the (magnitude) of is indeed equal to one, and is the normalizing term.

Furthermore, we can draw the following inferences about from the Eq. (3.12):

1. If the pupil magnification, , then , which implies that the opening angles of the image and object space perspective cones are equal, irrespective of the orientation of the optical axis. Then, the lens is symmetric about a plane perpendicular to the optical axis (in addition to the rotational symmetry about the optical axis). It must not come as a surprise that symmetric lenses are can be reversed without affecting any optical system properties [23].
2. If we let, such that, then we can write , where is the scalar normalization term. Furthermore, as is a diagonal matrix, and is orthonormal, we can immediately recognize the form as the Eigen value decomposition of a symmetric matrix, with —the columns of —as the eigenvectors and the corresponding eigenvalues. i.e., as is a diagonal matrix,

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and

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In Eq. (3.13) the terms are the projections of the input direction cosine along the eigenvectors. Also,, the third column of the rotation matrix , is the direction of the optical axis.

### **3.5 Image formation for arbitrary orientation of the lens and image plane**

Geometric imaging is a mapping (*bijective* in projective space) between points in the three-dimensional world space to corresponding points on a mathematical surface that we call the *image*. Here we aim to study the nature of this mapping on a planar surface—the image plane—for arbitrary orientations of the lens and image planes. To that effect, we will use the knowledge of the transfer of direction cosines of the chief ray derived previously.

An extended object emanates a multitude of chief rays that reach the image space through the lens elements and the stop. The locus of points formed by the intersection of these rays with the image plane constitutes the *projection* of the object in the image plane [25,26]. Further, we generally identify the projection of the object point as an “image” when the pencil of rays from the object point geometrically converge at a single point in the image space.

For simplicity, we assume that the lens is unencumbered by radial distortions and optical aberrations. [Figure 3.6](#Figure_3_6) represents a schematic of the problem in which we have introduced an image plane whose orientation is described by its surface normal. Two local frames are also introduced: the frame is attached to the optical axis with its origin at entrance pupil, and the frame attached to the image plane with its origin at the image plane pivot. The image plane is free to tilt or swing about its local x-axis or y-axis respectively at the image plane pivot.

Let the exit pupil () be located units from the pivot point along the optical axis. Following the rotation of the optical axis, by applying the matrix, the position of the exit pupil in camera frame is given as.

We can represent the chief ray emerging from the exit pupil with direction cosine by the parametric equation:

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where represent points along the output chief ray in. The first term on the right hand side of Eq. (3.14) is the initial position of the ray (at the center of) and is a real number that determines the length of the ray.

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| **Figure 3.6** Schematic of geometric image formation. is the *central* *projection* of the object point on image plane. The optical axis and image plane are free to swivel about the origins of coordinate frames and respectively. |

We write the equation of the image plane in Hessian normal form as:

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where is the unit normal to the image plane, is the perpendicular distance of the plane from the origin of frame , and is a point on the image plane.

We obtain the expression for (in Eq. (3.14)) for which the ray intersects the image plane by equating to, multiplying Eq. (3.14) by, and rearranging the terms:

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Substituting Eq. (3.16) into Eq. (3.14) we get

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Now we proceed to find the expression for the perpendicular distance in terms of the known parameters. The origin of the image plane’s local reference frame,, is located at the intersection of the -axis of camera frame with the image plane (see [Figure 3.7](#Figure_3_7)). Given the location of the image plane’s pivot, , we can describe the orientation of the image plane using the surface normal. The image plane’s surface normal is obtained by applying the rotation matrix to the unit vector :

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The equation of the image plane as represented by Eq. (3.15) is

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Since is a point on the plane, we obtain as:

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| **Figure 3.7** Schematic of the image plane. The image plane having surface normal is located at units from the origin of camera frame along the z-axis that intersects the plane at . is the perpendicular distance from the origin to the plane. The local image coordinate frame with its origin at the intersection of the image plane and z-axis of the camera frame is represented by . |

Substituting from Eq. (3.19) into Eq. (3.17) yields the expression for the point of intersection of the chief ray with the image plane in terms of the input direction cosines as

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Let the entrance pupil be located at a distance from the pivot point along the optical axis in the camera frame. Then, similar to the description of the exit pupil, the location of the entrance pupil in is given as:

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Further, we express the direction cosine components in terms of the Cartesian coordinates of the object point and entrance pupil in the camera frame as:

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which we write compactly as:

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Substituting Eqs. (3.23) and (3.21) into Eq. (3.20), we obtain a general relationship between the object point and its corresponding image point as:

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Eq. (3.24) represents the image point in the camera frame. Once an image—a two dimensional representation of the scene—has been formed, we specify positions and dimensions in the image independent of the position and orientation of the sensor and lenses (e.g. in terms of pixels in a digital image). We can transform the image coordinates in the camera frame represented by Eq. (3.24) to the image frame by observing that the origin of is displaced from by the translation vector , and the standard basis vectors of frame are rotated by . Consequently, a point in relative to may be expressed as (see Eq. 2.53 in  [27]):

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Therefore, we can write the expression for the image point coordinates in the image plane’s reference frame as:

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Substituting from Eq. (3.24) into Eq. (3.26) we obtain the expression of the two-dimensional image point in the image frame as:

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Where,

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|  | 3D Cartesian coordinates (in physical units) of the world point in camera frame . In the adopted coordinate convention, the numerical value of -component of is negative. |
|  | 2D Cartesian coordinates (in physical units) of the image point in the image frame . Note that Eq. (3.27) produces a vector with the third element (i.e. the   component) identically equal to zero. |
|  | Equal to , where is the pupil magnification. |
|  | Location of the entrance pupil from the pivot (origin of ) along the optical axis. This scalar quantity physical units, usually in millimeters. |
|  | Location of the exit pupil from the pivot (origin of ) along the optical axis. |
|  | Location of the image plane from the pivot (origin of ) along the z-axis of . |
|  | Rotation matrix used to describe the orientation of the lens plane. |
|  | The third column of . |
|  | Rotation matrix used to describe the orientation of the image plane. |
|  | The image plane normal. denotes the -component of the normal. |
|  | Origin of the image frame with respect to camera frame; . |

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| **Figure 3.8** Ray tracing for verifying Eq. (3.27). Chief rays traced from a grid of points in the object plane through an ideal lens tilted about a point away from the entrance pupil along the optical axis to the tilted image plane. |

### **3.6 Verification of imaging equation in Zemax**

We derived Eq. (3.27), which relates a three-dimensional object point to its projection in the two-dimensional image plane of a Scheimpflug camera, analytically. Now we verify the accuracy of the relationship by comparing the numerically computed values of image points (intersection of chief ray with the image plane) using Eq. (3.27) with corresponding image points obtained by tracing chief rays from a grid of points belonging to the object plane. [Figure 3.8](#Figure_3_8) is a layout plot of the optical system modeled in Zemax showing (1) an object plane, (2) an ideal lens made from two paraxial surfaces and pivoted about a point away from the entrance pupil (), and (3) an image plane pivoted about the image plane pivot along the -axis. We can arbitrarily assign any rotation angle to both the lens and image planes (pivot at their local frame’s origin) with respect to both - and -axis. The orientation of both planes is represented using *intrinsic* rotations matrices (composed of elemental rotations first about the -axis followed by rotation about the new -axis). Symbols and represent the angles of rotation of the lens plane about the - and -axes while and represent the angles of rotation of the image plane about the - and -axes. The results of the simulation are tabulated in [Table 3.1](#Table_3_1), which shows the set of object points, the numerically computed image points, the ray traced image points, and the absolute difference between the numerically computed and ray traced image points. We observe that the numerically computed and ray traced values of the image points are very close; the small difference in their values can be attributed to the error associated with floating point operations. This comparison demonstrates that the analytically derived expression (Eq. (3.27)) representing geometric relationship between a three-dimensional object point and its image point in the absence of optical aberrations is accurate.

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| **Table 3.1** Comparison of numerically computed image points with ray traced (in Zemax) image points for the optical system shown in [Figure 3.8](#Figure_3_8). | | | |
| World point | Computed image point | Ray-traced image point | Absolute difference |
| (0.0, 0.0, -509.0) | (-0.3108, -0.6291, 0.0) | (-0.3108, -0.6291, 0.0) | (1.8e-09, 3.1e-09, 7.5e-15) |
| (10.0, -10.0, -509.0) | (-0.8003, -0.0863, 0.0) | (-0.8003, -0.0863, 0.0) | (2.1e-09, 2.7e-09, 3.0e-15) |
| (-50.0, 50.0, -509.0) | (2.1291, -3.3352, 0.0) | (2.1291, -3.3352, 0.0) | (1.2e-09, 3.2e-09, 2.9e-15) |
| (70.71, 70.71, -509.0) | (-4.2013, -5.0221, 0.0) | (-4.2013, -5.0221, 0.0) | (2.6e-09, 5.1e-09, 4.7e-15) |
| (100.0, 0.0, -509.0) | (-5.5251, -1.0101, 0.0) | (-5.5251, -1.0101, 0.0) | (1.3e-09, 8.4e-09, 3.1e-15) |
| (0.0, 100.0, -509.0) | (-0.6031, -6.4387, 0.0) | (-0.6031, -6.4387, 0.0) | (2.2e-09, 4.0e-09, 2.2e-16) |
| (100.0, 100.0, -509.0) | (-5.8238, -6.8542, 0.0) | (-5.8238, -6.8542, 0.0) | (5.6e-10, 2.5e-10, 2.2e-15) |

### **3.7 Geometric properties of images under lens and image plane rotation**

Following the verification of Eq. (3.27), we use the expression to qualitatively study the effects of lens and sensor rotations on the geometric properties of the image. We also investigate the effects of pupil magnification, , and location of the lens pivot on the nature of the geometric distortions. [Figures 3.10 – 3.15](#Figure_3_10) show the type of distortions in images—of two planes in the object space—for several lens and sensor orientations. In all these figures, the basic setup is similar to that shown in [Figure 3.8](#Figure_3_8) except that the object space consists of two planes—a near plane and a far plane. The near plane is a square of 88.15 mm on each side, and the far plane is a square of 178.3 mm on each side placed at twice the distance of the near plane from the entrance pupil. The exact distance of the near plane (and consequently the far plane) from the lens vary depending upon the pupil magnification, such that the images of the two planes are 4.5 mm on each side on the sensor. Also, since the -axis of the camera frame passes through the center of both object planes, the two images are coincident in the frontoparallel configuration (i.e. when the object planes are parallel to lens and image planes).

The object points consist of square grids on each of the object planes. The corresponding “image points” are the points of intersection of the chief-rays (emanating from the object points) with the image plane ([Figure 3.9 (a)](#Figure_3_9)). The orange “Y” markers represent the group of image points from the near object plane. The blue “inverted Y” markers represent the image points from the far object plane. Note that in frontoparallel configuration the two images of the two object planes coincide; however, for the sake of visual clarity, we separate the two set of image points horizontally by 5 *mm* on either side from the center ([Figure 3.9 (b)](#Figure_3_9)). In [Figures 3.10 – 3.15](#Figure_3_10) while lighter shaded (orange and blue) markers are used to represent the points in frontoparallel configuration, darker shaded markers of either color represent the image points following the rotation of the sensor or lens. Rotation of either the lens or the sensor induces a geometric distortion of the image field in which the points across the image field translates by different amounts and

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| **Figure 3.9** “Image points” corresponding to two object planes—a far plane twice the size of the near plane. (a) The image points are coincident. (b) The coincident image points are separated laterally for the purpose of our investigation. |

directions. These translations are shown by the gray-to-white arrows between the original and shifted positions (drawn if the magnitude of the shift is greater than a certain threshold). The white level of the arrows specifies the normalized magnitude of translation—brighter indicates relatively larger translation. The figures also display information about the standard deviation (SD) of the arrow lengths. This statistic gives a sense of the non-uniform translation of the image points across the image field. If all image points shift by the same amount, then the standard deviation will be zero. A larger value of the standard deviation indicates greater diversity in shifts, and hence greater distortion. In addition to the standard deviation, we also measure how much the centroid of the set of points from the two images shifts. The translation of the centroid gives a sense of how the total image field “appear” to shift. However, we must note that we consider an image field to have *translated* only when *all* image points have at least a fixed minimum amount of shift in the direction of translation. Note that in all cases shown here, the *image points* were not determined using a “best focus” criterion, but rather by the point of intersection of the chief rays with the image plane. However, this definition of the *image* adopted for the current discussion does not limit the study of geometric properties, such as the kind of transformations induced by the rotations of the sensor and lens planes.

#### **3.7.1 Properties of image field induced by sensor rotation ()**

[Figure 3.10](#Figure_3_10) shows the images of the two object planes obtained under sensor rotation for three different values of pupil magnification. Studying the figures, we observe that:

1. The image plane is no longer frontoparallel with the object plane. As a result, the transverse magnification varies across the image field. Therefore, the image points undergo a field dependent and asymmetrical geometric distortion.
2. Since the location of both the entrance pupil and the exit pupil remain fixed, the on-axis image point continues to remain on-axis subsequent to the rotations. Therefore, we can conclude that the there is no translation of the image field following sensor rotation.
3. The amount of perspective distortion is directly proportional to the pupil magnification.
4. However, as shown in [Figure 3.11](#Figure_3_11), the distortion is independent of the object distance. Therefore, rotating the sensor plane does not introduce parallax between images that are obtained under varying orientations of the sensor. Consequently, if we capture multiple images under several rotations of the sensor, the inter-image homography—mapping between corresponding points of two images—is simply a perspective mapping of the following form:

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#### **3.7.2 Properties of image field induced by lens rotation away from center of the entrance pupil (; )**

The influence of object distance and pupil magnification on the image shape obtained under lens rotation away from the entrance pupil is depicted in [Figure 3.12](#Figure_3_12). The emergence of parallax in the images is highlighted in [Figure 3.13](#Figure_3_13), in which we have plotted the top, middle and bottom rows of the image points for the three cases of pupil magnifications. We observe that:

1. The dominant effect of rotating the lens about a point along the optical axis is, in general, a *non-uniform shift* of the image field that depends on both the pupil magnification and object distance.
2. Since points in the image field undergo non-uniform translation, the standard deviation of the translation vector length is non-zero. Also, because the amount of shift of the image field is dependent on the object distance, the value of the standard deviation of the translation vector lengths is different for the images of the two object planes.
3. Since the amount of translation depends on the object distance, images obtained while varying rotation angle of the lens exhibit parallax as shown in [Figure 3.13](#Figure_3_13). We observe that in each case (of different pupil magnifications), the initially overlapping rows of the two images (from the two object planes) diverge with progressive rotation of the lens.
4. From the above observations, we infer that unless the pupil magnification is equal to one, the inter-image homography is a depth (object distance) dependent perspective mapping. In other words, for every object plane, the inter-image homography is of the form:

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1. However, if the pupil magnification is equal to one, then the inter-image homography reduces to a depth dependent scaling transformation of the form:

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#### **3.7.3 Properties of image field induced by lens rotation about the center of the entrance pupil (; )**

Finally, the properties of the image obtained when the lens is rotated about the center of the entrance pupil are depicted in [Figure 3.14](#Figure_3_14) and [Figure 3.15](#Figure_3_15). We observe that:

1. Rotation of the lens about the entrance pupil induces a shift of the image field, as must be expected.
2. However, unlike the previous case, the shift of the image field is *independent* of the object distance. We see that the standard deviation of the translation vector lengths is equal for both the group of points (from the two object planes).
3. Since the shift of the image field is independent of the object distance, there is no parallax between images obtained while rotating the lens about the entrance pupil. This is a very important property that can be used for several computational imaging techniques that rely on multiple image capture, including omnifocus imaging, digital super-resolution, panoramic imaging, etc.
4. If the pupil magnification is not equal to one, the inter-image homography is a *depth independent* perspective transformation of the form:

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1. If the pupil magnification is equal to one, then the inter-image homography is a depth independent perspective transformation of the form:

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| **Figure 3.10** Geometric image under image plane (sensor) rotation for varying pupil magnifications. (a) , (b) , (c) . |

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| **Figure 3.11** Comparison of geometric distortion induced by sensor rotation for varying object plane distances. The figure shows only the top, middle and bottom rows of image-points for pupil magnification ([Figure 3.10 (a)](#Figure_3_10)). Unlike in the previous figure, the two sets of image-points from the two object planes are left unseparated. We observe that the image-points corresponding to the two object planes that are at different distances from the lens experience the same type and amount of distortion. If that was not the case, then the image-point rows corresponding to the two planes would have diverged. |

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| **Figure 3.12** Geometric image under lens rotation away from the entrance pupil for varying pupil magnifications. (a) , (b) , (c) . |

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| **Figure 3.13** Variation of geometric distortion of image field induced by lens rotation away from the entrance pupil as a function of object distance and pupil magnification. |
|  |
| **Figure 3.14** Geometric image under lens rotation away from the entrance pupil for varying pupil magnifications. (a) , (b) , (c) . |

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| **Figure 3.15** Variation of geometric distortion of images induced by lens rotation *about* the entrance pupil as a function of object distance and pupil magnification. |

### **3.8 Summary**

The geometric relation between a three-dimensional object point () in camera coordinates and the corresponding image point () in the two-dimensional image plane for an aberration-free Scheimpflug camera is given by Eq. (3.27):

This imaging model has two important characteristics that differentiate it from other existing ones:

1. It is a most general geometric imaging relationship between object and image points for rectilinear imaging that accommodates the rotation of the image and the lens planes about their individual pivots. Common imaging models, such as for frontoparallel imaging, directly falls out of the above model.
2. The model incorporates optical parameters such as the pupil distances, pupil locations and pupil magnification, which enable us to accurately analyze the nature of the geometric image obtained the various possible orientation of the lens and image planes.

Our imaging model allowed us to study the effects of rotation of the lens and image planes on the geometric properties of the image. We discovered that if the pupil magnification equal to one and the lens is rotated about the entrance pupil the inter-image homography—the transformation that relates the image of a scene obtained under varying rotations of the lens—is a simple composition of a scaling and translation transformations.

# Chapter 4

**⬀ MODEL OF SCHEIMPFLUG IMAGING – II: FOCUSING**

*The scientist explains the world by successive approximations.*

—Erwin Hubble

In the last chapter, we derived a relationship (Eq. (3.27)), between object and image points that allowed us to study and discover important geometric properties of the image in a Scheimpflug camera. However, in deriving Eq. (3.27) we did not impose any constraint on the orientations of the sensor and lens planes that would produce a sharp, focused image of a (tilted) object plane on the sensor plane. Hence, while Eq. (3.27) is useful in analyzing and predicting the geometric properties of the images for a given sensor and lens plane orientation, it is suitable for determining the required orientation of the sensor and lens planes for focusing on a given tilted object plane. In this chapter, we derive a general relationship between the lens, image and object planes that ensures geometric focus on the tilted object plane. Like Eq. (3.27), we aim to explicitly include the pupil parameters in our focusing equation so that we can model the effects of the pupils on focusing. Further, we show (as examples) that this general focusing expression yields object-image relationships that are specific to common types of Scheimpflug configurations.

### **4.1 Introduction**

To keep the problem tractable, we impose the constraint that three pivots (that hinges the object, lens, and sensor planes) lie along the -axis of the camera frame, and the origin of is co-located with optical axis’ pivot (the lens plane is perpendicular to the optical axis). We also restrict the rotation angles of the object, lens, and image planes between and about both - and -axes (in-plane rotations or rotations about the -axis is irrelevant for our purpose). Provided we make no distinction between the faces (front or back) of the planes, this restriction on the angles of rotations is not limiting in any way since we can uniquely describe all possible plane orientations in three dimensions. On the other hand, this constraint warrants non-negative values for the -component of the plane normals and permits us to estimate the plane normal unambiguously.

### **4.2 Relationship between the object, lens, and image planes for focusing**

We begin by deriving an expression for the chief ray joining an arbitrary point in the object plane to a point in the image plane. For to be the geometric image of , the chief ray between the conjugate points must satisfy the Gaussian imaging equation. This constraint allows us to uniquely determine the position (of the image plane along the z-axis of) and orientation of the three planes in Scheimpflug configuration. The setup is shown in [Figure 4.1](#Figure_4_1).

The object plane is located at a distance of (the numerical value of is negative in our convention) from the origin of camera frame, along the -axis. The object plane, pivoted about the point in the camera frame, is completely described by the pivot point and the normal,. We describe the object plane normal itself as the product of the rotation matrix and ():

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The rotation matrix is typically composed of elementary rotation matrices that represent rotations about the - and -axis.

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| **Figure 4.1** Schematic of Scheimpflug imaging. The figure shows the object plane, optical axis, and image plane pivoted about points along the -axis, , of the camera frame. The local object plane and image plane coordinates frames ( and) are centered on the object- and image- plane pivots. |

If be the perpendicular distance of the object plane from the origin, we can determine using the Hessian normal form as:

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since is a point on the object plane. Further, any general point on the object plane, satisfies the plane equation:

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Similar to the object plane normal, we describe the image plane normal as:

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where is a rotation matrix applied to the image plane at the pivot point .

Repeating the steps used to derive the object plane equation, we obtain the equation for the image plane as:

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where is any point on the image plane.

Suppose the entrance () and exit () pupils are located at distances and from the lens’ pivot (origin of) respectively, along the optical axis. Also, we describe the rotation of the optical axis by applying the matrix. Then, the positions of the pupils in following the application of the rotation matrix are:

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Consider the chief ray from the object point to the corresponding image point , passing through the center of the entrance and exit pupils. Let the direction cosines of the ray in the object- and image- space be and respectively. Since is the (unit-length) direction vector of the ray from to , we can write:

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where, is the position of the entrance pupil center, and is the length of the ray. Substituting Eq. (4.6) into Eq. (4.8), we obtain:

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Further, since is a point on the object plane, it satisfies the object plane Eq. (4.3). Substituting into Eq. (4.3) and rearranging terms we obtain:

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The chief ray in the image space emerges from the exit pupil with a direction cosine vector. Therefore, we can write the equation of the chief ray in parametric form (using Eq. (4.7)) as:

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where is any point on the ray, is the position of the exit pupil, and is the length of the ray.

If the length of the chief ray in the image space be, then at the point of intersection of the chief ray with the image plane and in Eq. (4.11). Therefore, we obtain

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Substituting in Eq. (4.5) and rearranging terms, we obtain

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Eqs. (4.10) and (4.13) gives the length of the chief rays between and in the object space and between and in the image space respectively. For to be the geometrically focused image of , the lengths in the object and image space must satisfy the Gaussian imaging equation.

The well-known Gaussian imaging equation () relates the focal length and the conjugate plane *directed* distances measured from the principal planes. Instead, if the directed distances are specified with respect to the pupil planes (from entrance pupil to object plane; and from exit pupil to image plane), then a variant of the Gaussian imaging equation is used, which incorporates the pupil magnification into the formula:

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where and are directed distances along the optical axis measured from the entrance pupil to the object plane, and from the exit pupil to the image plane respectively; is the pupil magnification; and is the focal length. We have provided a derivation and a brief exposition of Eq. (4.14) in [Appendix B.1](#_Appendix_B.1_Derivation).

The most common application of Eq. (4.14) is for frontoparallel imaging in which the conjugate planes are parallel to each other and perpendicular to the optical axis. Moreover, pairs of object-image conjugate points satisfy this relation even if the ensemble of object- and image- points belong to planes on object and image sides respectively that are *not parallel* to each other.

The ray vector of length and direction in the object space is. The projection of this ray vector on the optical axis () is. Similarly, the ray projection of the image space ray vector on the optical axis is. In order to substitute and in to Eq. (4.14) we need to ensure that they are the directed distances. Following our sign convention, the directed distance from the entrance pupil to the object point is . Substituting and into Eq. (4.14) we obtain:

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Further, substituting the expressions for and (Eqs. (4.10) and (4.13)) into the above equation, we obtain:

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The direction cosine of the chief ray in the image space,, is related to the direction cosine of the chief ray in the object space as (Eq. (3.12)):

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Substituting, into Eq. (4.16) we obtain:

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To simplify the above expression, let us consider as

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Therefore, we can write in Eq. (4.17) as. Similarly, we can also write as. Then, Eq. (4.17) can be written as:

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We can further simplify the above equation by noting that:

1. ,
2. ,
3. , and
4. .

Using the above results Eq. (4.19) reduces to:

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Further, multiplying the above equation by the scalar , and using the commutative and distributive properties of dot product, we obtain:

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Note that in the above expressions, the third column of the rotation matrix, is the unit vector along the optical axis.

The direction cosine vector has -Norm equal to one. Therefore, the above equation is satisfied only if the second vector,, is either perpendicular to or identically equal to zero. Now, is it possible for the second vector to be perpendicular to (the chief ray’s direction vector)? As we did not make any specific assumptions about in the derivation of Eq. (4.21), *all* chief rays—an infinitude of vectors within the object- and image-space perspective cones—must satisfy Eq. (4.21). Are all possible vectors perpendicular to the second vector? Since the second vector is a linear combination of (the object plane normal), (the transformed image plane normal), and (unit vector along the optical axis), we can deduce that , *in general*, is not perpendicular to the second vector. Therefore, the second vector must be equal to zero, which yields:

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Further, we can simplify Eq. (4.22) if we let and. Then, after factoring and out of the denominator terms, we can write Eq. (4.22) as

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Where,

|  |  |
| --- | --- |
|  | Directed distance of the object plane’s pivot from the origin of the camera frame, along the z-axis of the camera frame. (In most of the problems that we are interested, the value of is negative.) |
|  | Directed distance of the image plane’s pivot from the camera frame, along the z-axis of the camera frame. |
|  | Focal length of the lens. |
|  | Equal to , where is the pupil magnification. |
|  | Location of the entrance pupil from the pivot (origin of ) along the optical axis. This scalar quantity physical units, usually in millimeters. |
|  | Location of the exit pupil from the pivot (origin of ) along the optical axis. |
|  | Rotation matrix used to describe the orientation of the lens plane. |
|  | The third column of . |
|  | Equals , where is the object plane normal. |
|  | Equals , where is the image plane normal. |

The expedient simplification from Eq. (4.22) to Eq. (4.23) relies on our ability to describe the *unit* normal vectors and using only the components along - and -axes. If we know the - and - components of the normal vector, we can determine the - component uniquely because we have restricted the angles of plane rotations between and about both - and -axes (one of the starting assumptions). For example, if the object- and lens-plane orientations and distances are known, and we estimate the image plane distance and orientation vector () of the image plane using Eq. (4.23), then we can determine the image plane normal as:

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where , , and we have dropped the negative sign from the expression for since is guaranteed to be positive as discussed under the assumptions at the beginning of this section.

Eq. (4.23) is most general in the sense that it readily yields the various expressions for specific cases of object, lens and image plane orientations as we now demonstrate in the following set of examples.

### **4.3 Examples of typical Scheimpflug imaging configurations**

#### **4.3.1 Example: Focusing in frontoparallel configuration**

Suppose the object plane is frontoparallel with the lens plane associated with a thick lens and units from the origin of camera frame (along the -axis of the ), what should be the orientation and position of the image plane to focus on the object plane?

Based on the given data, we have the following:

1. Since the lens is not tilted, , .
2. Since the object plane is not tilted, and .
3. Let and. Then, , and , and

Substituting the above parameters in Eq. (4.23) we have

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which yields the following relations from each row of Eq. (4.25):

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Writing and where, is the directed distance from the entrance pupil to the object plane and is the directed distance from the exit pupil the image plane, we can rewrite Eq. (4.28) as

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As a concrete example, if , , ,  , and , then we can calculate the image plane distance from the camera frame’s origin using Eq. (4.28) to be .

Further, if, then the sign of is positive, which is a condition for real and inverted images. For example, if, and, then is positive for. If the sign of is negative, then a virtual and upright image is formed in front of the lens. In a survey of 120 imaging lenses (see [Appendix B.2](#_Appendix_B.2_A)) from a database (Zemax Zebase) of well-designed lenses, we found over 90% of all lenses to have pupil magnification greater than 0.5 and no lens having pupil magnification less than 0.2. Thus, in the common imaging scenarios, the sign of is positive.

From Eqs. (4.26) and (4.27) we see that the image plane normal is equal to. This implies that the image plane is parallel to the lens and object plane and perpendicular to the optical axis if the object plane is parallel to the lens plane.

Additionally, if a thin lens model (,), Eq. (4.28) reduces to the Gaussian lens equation for thin lenses:

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Furthermore, if the distances to the object and image plane are specified from the object- and image- space principal planes instead of the pupils, Eq. (4.28) reduces to ()

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in which, is the distance of the object plane from the object-space principal point, and and represent the locations of the object- and image-space principal points in the camera frame . While is numerically negative, the signs of and depend on the position of the principal points with respect to the origin of.

#### **4.3.2 Example: Focusing on tilted object plane by tilting the image plane**

Suppose the object plane, pivoted at, is tilted by an angle about the -axis ([Figure 4.2](#Figure_4_2)), and the lens plane perpendicular to the -axis of the camera frame , what is the conjugate orientation and position of the image plane for achieving a geometrically focused image assuming a thick lens model?

1. Since the lens is not tilted, , .
2. Let and. Then, and .
3. Let and. Then, , and .

Substituting the above parameters in Eq. (4.23) we obtain:

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| **Figure 4.2** Object and image plane tilt. In the above cross-sectional (y-z plane) view, the object plane is tilted by an angle of about the x-axis at. We would like to find the position and orientation of the image plane in order to focus on the object. |

As in the [Example 4.3.1](#_4.3.1_Example:_Focusing), we obtain the distance of the image plane pivot in the camera frame, from the third row of Eq. (4.32) as:

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Also, as in [Example 4.3.1](#_4.3.1_Example:_Focusing), writing and where, is the directed distance from the entrance pupil to the origin of the object plane and is the directed distance from the exit pupil to the origin of the image plane, we obtain:

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in which is numerically negative and is positive for macroscopic imaging.

From the first and second rows of Eq. (4.32), we obtain:

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and

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Since the object plane is rotated by about the -axis, the rotation matrix is

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Note that if the direction of rotation of the object plane about the -axis is from +z-axis to +y-axis (as depicted in [Figure 4.2](#Figure_4_2)), then is numerically negative.

The object plane normal is (Eq. (4.1)) is:

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from which, we get (by dividing by ) as:

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Substituting into Eqs. (4.35) and (4.36) we obtain:

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Eq. (4.40) in conjunction with Eq. (4.24) suggest that, which implies that the image plane normal is confined to the - plane. That is, if the object plane is rotated only about the -axis, then the image plane must also be rotated only about the -axis to achieve geometric focus on the tilted object plane. The exact expression for the components of the image plane normal is obtained by substituting and into Eq. (4.24). The angle of the image plane normal with respect to the three axes is further obtained as the cosine inverses of the corresponding components of the normal vector.

Furthermore, there remains a desideratum to have a more direct relationship between the rotation angles of the object and image planes that can be readily used instead of computing the plane normals. To that end, let us suppose that the required image plane tilt (rotation about the x-axis) is. If we represent the rotation matrix of the image plane as:

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we obtain the image plane normal as:

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which produces as:

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Substituting (from Eq. (4.43)) into Eq. (4.40) we obtain:

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where, for common macroscopic imaging, is numerically negative, and is positive. Therefore, the sign of is opposite to the sign of . This result implies that *the image plane must be rotated in the direction opposite to the direction of rotation of the object plane*.

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| **Figure 4.3** Object and image plane tilt (distances measured from principal planes). In the above cross-sectional (- plane) view, the object plane is tilted by an angle of about the -axis. The distances to the object and image plane pivots are specified from the respective principal planes. |

We can further modify the relation Eq. (4.44) if the distances are specified with respect to Principal planes ([Figure 4.3](#Figure_4_3)). Then, represents the magnification between the image and object side principal planes and therefore, is equal to one; the ratio represents the ratio of the distances of the principal-plane-to-object plane in the object side and principal-plane-to-image plane in the image side along the z-axis, and therefore is equal to the magnification along the -axis. Letting,, and, we can rewrite Eq. (4.44) for the case if the image and object plane distances are measured from the Principal planes as:

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#### **4.3.3 Example: Focusing on a tilted object plane by tilting a lens using thin lens model**

Determine the orientation (angle ) of a thin lens required to focus on a tilted object if the lens plane is not tilted.

1. Since we have a thin lens, , , .
2. Let and. Then, .
3. Since the image plane is not rotated, .

The unknowns in this problem are the image plane distance , and the lens plane angle . In general, a rotation matrix has three degrees of freedom, however, since we have restricted the rotation of the lens plane to only about the - and - axes, the rotation matrix in our problem has only two degrees of freedom. Further, we can describe the orientation of the lens plane using just the third column of (the normal to the lens plane). Therefore, in total, we have three unknowns—the image plane distance and the two rotation angles of the lens plane. Another way to think about the number of knowns and unknowns is that we are required to determine the normal vector to the lens plane and the distance of the image plane. Since only two components of the normal vector are essential to determine the orientation of the plane (when the angles of rotation of the plane is restricted between ), we have three unknowns.

Substituting the known parameters in Eq. (4.23) we obtain:

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The three rows of the above equation yields:

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and

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In order to solve for the image plane distance in Eq. (4.49) we need . Can we uniquely determine from and ? A property of any rotation matrix is that each column (or row) has unit length. Therefore, . Furthermore, if the rotation matrix is composed of elementary rotations only about the - and - axes, then , which is a product of the cosine of the angles of rotations about the two axes, is guaranteed to be positive. Hence,

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For a concrete example, suppose the object plane is tilted only about the -axis by an angle . Then, (as in the previous example), and . Therefore, , , and . The image plane distance is given by .

The dependence of on (and implicitly on the amount of lens tilt as shown below in Eq. (4.52)) implies that we need to also translate the image plane location along the -axis of in order to focus on a tilted plane. Instead, if we chose to focus on the tilted plane employing just sensor rotation, the distance of the sensor’s pivot from the camera center remains fixed.

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| **Figure 4.4** Object and lens (thin lens model) plane tilt. In the above cross-sectional (- plane) view, the object plane is tilted by an angle of about the -axis at. We would like to find the position of the image plane and orientation of the lens plane in order to focus on the tilted object surface. |

Furthermore, implies that the lens is rotated only about the -axis. If we let (the third column of ), then the relationship between the lens plane rotation angle () and object plane rotation angle () is obtained as (and shown in [Figure 4.4](#Figure_4_4))

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In the above equation is numerically negative, therefore the sign of is same as the sign of , which implies that the direction of rotation of the lens and object planes are congruent.

Furthermore, we obtain the image plane distance along the -axis of as:

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The sign of is positive, forming real and inverted image behind the lens, if .

#### **4.3.4 Example: Focusing on a tilted object plane by tilting a lens using thick lens model**

If the image plane is not rotated, what is the required angle of rotation of the lens in order to focus on an object surface that is titled about the -axis by an angle using a thick-lens model?

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| **Figure 4.5** Object and lens (thick lens model) plane tilt. In the above cross-sectional (y- plane) view, the object plane is tilted by an angle of about the x-axis at. We would like to find the position of the image plane   and orientation of the lens plane in order to focus on the tilted object surface. |

The schematic of the problem is shown in [Figure 4.5](#Figure_4_5). We can represent the orientation of the object plane that is tilted about the x-axis using the rotation matrix

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such that

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In the previous problem (using thin lens model) we observed that the direction of rotation of the object- and lens-planes are congruent. Although we expect the exact angle of the lens plane to differ from the solution of the previous problem because of the thick lens model, there is no reason to suspect the direction of rotation of the lens to be different from the thin lens case. Therefore, the structure of the rotation matrix representing the lens plane’s orientation is similar to that of the object plane:

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such that,

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As the image plane is not tilted, we represent the normal of the image plane as:

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Therefore, , , and

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Substituting the above parameters into Eq. (4.23) we obtain:

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The third row, following simple algebraic steps, yields the formula for the image plane distance in terms of the angles of the object and lens planes as:

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From the second row we obtain:

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Substituting from Eq. (4.59) into Eq. (4.60) and writing , , and , we obtain:

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Multiplying by we obtain:

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Following few algebraic steps, we obtain the finite conjugate imaging relationship between the object plane tilt angle and lens plane tilt angle when the lens is rotated about a pivot (away from the entrance pupil) as:

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In the discussion that follows we will use the notation to represent the right-hand-side of the above equation. Eq. (4.63) is an implicit relationship between the angles and . Comparing Eq. (4.63) with Eq. (4.51) immediately shows that for a given object plane tilt angle the lens tilt angles obtained by the thick-lens (more accurate) model deviates from that obtained using a thin lens model. Further, the object plane angle (in focus) obtained via Eq. (4.63) for a given lens tilt angle depends on the location of the lens pivot point along the optical axis. The variation of with respect to the pivot position (offset from entrance pupil, ) for an object plane pivoted at from the entrance pupil of a lens with pupil magnification is show in [Figure 4.6](#Figure_4_6). The thin-lens model doesn’t account for such deviations.

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| **Figure 4.6** Variation of (-axis) with respect to lens pivot position for (a) , (b) , and (c) . The x-axis is the offset of the pivot position from the entrance pupil. The plots show that the range of is large for larger values of . In each graph the two other plots are also plotted in lighter values for comparison of the slopes. |

As the object plane distance increases, the effective object-to-entrance-pupil distance () in Eq. (4.63) tends to a constant value for relatively small changes in the entrance-pupil-to-pivot-point distance . Therefore, for relatively large object plane distances the variation of is expected to be negligible for small changes in .

##### ***Verification of formulae for focusing on a tilted object plane by tilting the lens***

Before we begin to examine the consequences of the focusing equations Eq. (4.59) and Eq. (4.63) that relates the lens plane orientation and image plane distance from the origin of camera frame if the lens is tilted about a point away from the entrance pupil, it is imperative to ensure that the equations are verified. [Table 4.1](#Table_4_1) enumerates the results of our test. In order to test the equations, we implemented a thick lens model in Zemax using two paraxial surface—to ensure aberration-free geometric imaging model—with pupil magnification . The lens surfaces were grouped within two coordinate break surfaces that allowed the lens to be tilted about a point away from the entrance pupil. The object plane surface was placed at from . Then, for every object plane orientation (col. 1), the appropriate lens tilt angle (col. 2) and image plane distance (col. 3) were obtained using Zemax’s optimization function, to minimize spot radius across the field. Following optimization for every , the value of obtained from Zemax (along with the values of , , ) was used to numerically compute (col. 4) and (col. 5) using the derived equations Eq. (4.59) and Eq. (4.63). We can observe that the values of and obtained numerically using the derived equations are very closely matched. It must be emphasized that the values ( and ) obtained in Zemax were using ray-tracing and optimization, rather than numerical evaluation of analytic expressions.

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| **Table 4.1** Verification of imaging equations Eq. (4.59) and Eq. (4.63) for focusing on a tilted object plane by tilting a lens about a point away from the entrance pupil. | | | | |
| (Zemax)1 | (Zemax)2 | (Zemax)3 | (numerical)4 | (numerical)5 |
| 0.0° | 0.0° | 24.17073 *mm* | 2.07042e-13° | 24.17073 *mm* |
| -10.0° | -0.46989° | 24.17163 *mm* | -9.99973° | 24.17163 *mm* |
| 25.0° | 1.24260° | 24.17701 *mm* | 24.995702° | 24.17701 *mm* |
| -40.0° | -2.23573° | 24.19107 *mm* | -39.98214° | 24.19107 *mm* |
| 65.0° | 5.70827° | 24.30378 *mm* | 64.91024° | 24.30377 *mm* |
| -80.0° | -14.99585° | 25.11194 *mm* | -79.74010° | 25.11146 *mm* |
| 1. *Object plane tilt set in Zemax.* 2. *Lens plane tilt obtained using optimization using ray-tracing in Zemax.* 3. *Image plane distance obtained using optimization using ray-tracing in Zemax.* 4. *Object plane tilt computed numerically using the value of in column 2.* 5. *Image plane distance computed numerically using the value of in column 2.* | | | | |

##### ***Consequences and analysis of the focusing equation***

Eq. (4.63) suggests that for a given lens tilt angle , we can determine the orientation of the plane-of-sharp focus . But what if we need to compute given ? Quite often we need the inverse relationship; that is, we need to determine the lens tilt angle required to focus on a tilted object plane. We will return to this question a few times in this section—first qualitatively in the imminent discussion followed by a detailed analysis near the end of this section. Another related question that we may ask at this point is whether the relationship between and , as depicted using Eq. (4.63), always one-to-one? In other words, is the function , which represents the right-hand-side of Eq. (4.63), *monotonic* and *invertible* within the interval ? If is monotonic, then it follows that will be monotonic within and we will have a one-to-one relationship between and . A test to determine the monotonicity of within an interval is to examining if the first derivative, , changes sign within the interval. Later in this section, when we examine the case of rotating the lens about the entrance pupil, we will present a more detailed analysis of the first derivative of . But first we carry out a qualitative analysis of Eq. (4.63) which relates the orientations of the object and lens planes if lens is pivoted about a point away from the entrance pupil. Although we will seldom rotate a lens about a point away from its entrance pupil because of the unwieldy distortions induced to the image field (see [section 3.7.2](#_3.7.2_Properties_of)), the methods and insights developed in this study, nevertheless, will carry over to our examination of the latter case.

In [Figure 4.7](#Figure_4_7) we have plotted values of — or —and the object tilt angle versus the lens tilt angle for four different values of while keeping the parameters , and fixed. We can observe that while the function is monotonic for and , it becomes non-monotonic for and as evidenced by the presence of stationary points (the exact locations are not important) in the corresponding plots. Consequently, for example, when we obtain the same value of ( 72.3°) for two different values of (17.65° and 45.0°). These plots suggest that is not always monotonic. In our simulations we have found that tends to become non-monotonic as the value of becomes very small.

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| **Figure 4.7** Object plane angle and versus lens tilt angle if a lens is rotated about a point away from the entrance pupil. In these plots , and . The plots show that the function is not always monotonic, especially for small values of . For e.g., in the plots corresponding to the function yields the same value of for and . |

Apart from the monotonicity of , we also need to consider the sign of the image plane distance from the exit pupil (), equal to , obtained through Eq. (4.59). If in Eq. (4.59), then is negative which implies that a *virtual image* is formed in front of the lens. Therefore, in order to form a *real image* on a sensor, the condition must be satisfied.

Once the two independent conditions—the monotonicity of and the formation of real image—are satisfied, we can return to the question regarding how to determine the lens tilt for a known object plane tilt angle . Obtaining an expression for as a function of from Eq. (4.63) is not straightforward. However, we can develop some insights into the problem if we substitute and into Eq. (4.63), which yields the implicit equation:

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We can recognize Eq. (4.64) as a *quartic plane curve* of the form . The object plane tilt angle along with the parameters , , and form the coefficients of this fourth degree plane curve. Further, since and , we also obtain a second curve—a unit circle with equation . The lens tilt angle (more precisely and ) satisfies both these equations, therefore it must be at the point of intersection of the two curves. In [Figure 4.8](#Figure_4_8) we have plotted several quartic curves corresponding to different pupil magnifications and two choices of lens rotations and . Note that the corresponding value of for each curve in both groups will be different. The other parameters— and —are same for all curves. In order to keep the discussion simple, these parameters along with value the

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| **Figure 4.8** Determination of lens tilt angle for known object tilt angle using point of intersection of *quartic plane curve* with the unit circle. If a lens is rotated about a point away from the entrance pupil, the equation for as a function of , , , and can be expressed as a quartic plane curve in Cartesian coordinates with , . Furthermore, the intersection of the quartic curve with the unit circle yields the unknown lens tilt angle . The figure plots two groups of quartic curves—the blue curves correspond to equations for different values of (, , and ) and but the same while the green curves corresponds to . For all curves , and . |

for each curve used in the figure satisfy the conditions of monotonicity of and real image. The curves in blue correspond to the that satisfy Eq. (4.64), and the green

curves correspond to . Since the constant in Eq. (4.64)—corresponding to in the general quartic equation—is zero, all curves pass through the origin. Additionally, since the coefficient of the term is zero while the coefficient of the term is non-zero, the curves are either above the x-axis (for positive values of ) or below the y-axis (for negative values of ). As shown in the figure, all such quartic curves (for which the parameters satisfy the monotonicity condition) intersect the unit circle at two points—in the first and second quadrants if , or in the third and fourth quadrants if . However, since in Eq. (4.64) the abscissa of the point of intersection must always be positive for . Consequently, we can a determine a unique point of intersection of the quartic curve (given , , , and ) with the unit circle that correspond to the unknown lens tilt angle .

Based on the above discussion we see that it is possible to implement an algorithm to find the point of intersection of a quartic curve and the unit circle—for example, using Newton’s method—to determine given . Alternatively, if a good initial estimate of as a starting point is known, an iterative algorithm that converges towards the true point of intersection along the unit circle may be used to find the . We demonstrate such an algorithm at the end of this section.

Rotating the lens about a point away from the entrance pupil distorts the image field in complex ways and induces parallax between corresponding scene points in sequence of images obtained under lens rotations. Therefore, in several imaging applications that require multiple image captures under lens rotation it is imperative to rotate the lens about the entrance pupil. In the rest of this section we will undertake a closer examination of this scenario.

We can obtain the equations for image plane distance and plane of sharp focus orientation for the case when the lens is pivoted at the entrance pupil by substituting and (where, is the distance of the exit pupil from the entrance pupil) in Eqs. (4.59) and (4.63) respectively yielding:

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and

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| **Table 4.2** Verification of imaging equations Eq. (4.65) and Eq. (4.66) for focusing on a tilted object plane by tilting a lens about the entrance pupil. | | | | |
| (Zemax)1 | (Zemax)2 | (Zemax)3 | (numerical)4 | (numerical)5 |
| 0.0° | 0.0° | 29.17073 *mm* | -2.24301e-15° | 29.17073 *mm* |
| -10.0° | -0.46989° | 29.17145 *mm* | -10.0° | 29.17145 *mm* |
| 25.0° | 1.24249° | 29.17572 *mm* | 25.0° | 29.17572 *mm* |
| -40.0° | -2.23504° | 29.18687 *mm* | -40.0° | 29.18687 *mm* |
| 65.0° | 5.69682° | 29.27607 *mm* | 65.0° | 29.27607 *mm* |
| -80.0° | -14.79587° | 29.90304 *mm* | -80.0° | 29.90304 *mm* |
| 1. *Object plane tilt set in Zemax.* 2. *Lens plane tilt obtained using optimization using ray-tracing in Zemax.* 3. *Image plane distance obtained using optimization using ray-tracing in Zemax.* 4. *Object plane tilt computed numerically using the value of in column 2.* 5. *Image plane distance computed numerically using the value of in column 2.* | | | | |

[Table 4.2](#Table_4_2) enumerates the results of verifying the above equations against Zemax. The optical system used in the test is similar to that used to verify the Scheimpflug imaging equations for lens rotation about a point away from the entrance pupil, except that since the pivot was shifted to the entrance pupil, the object distance changed from to . The table shows that the numerically computed value of object plane tilt angle using the analytic expressions matches exactly to that used in Zemax to tilt the object plane. Also, the value of obtained in Zemax using ray tracing and optimization is in exact match with that obtained using the analytic expression we derived, proving the accuracy of the above expressions.

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| **Figure 4.9** Object plane angle and versus lens tilt angle if a lens is rotated about the entrance pupil. The plots show that the function is not always monotonic, especially for small values of . For e.g., in the plots corresponding to the function yields the same value of for and and for all plots. |

Comparing Eqs. (4.65) and (4.66) with Eqs. (4.59) and (4.63) respectively, we can immediately observe that the equations are far less complex when the lens is rotated about the entrance pupil. However, just like before, the function that represents the right-hand-side of Eq. (4.66) is not always monotonic. As shown in [Figure 4.9](#Figure_4_9), is monotonic for and but non-monotonic for and .

The following two subsections are devoted to a deeper examination of the condition of monotonicity of and the development of an iterative algorithm for determining the lens tilt angle .

##### ***Condition for monotonicity of***

The first derivative of with respect to is:

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| where, . |  |

We observe that the first derivative is a cubic equation in . A cubic equation has at least one real root which implies that the graph of must cross the real x-axis at least once. However, in our problem, the lens rotation angle is restricted to within . Consequently, we are only concerned with those roots of that are in the open interval , because . More concretely, if has a real positive root in the interval implying that it changes sign, then is non-monotonic within . In such a case we cannot find a unique for a given value of . In [Figure 4.10](#Figure_4_10) we have plotted the first derivative for varying values of . We can see that the plots of for and have at least two roots as they cross the real x-axis twice with the interval . Of course, this result was expected as we have already seen in [Figure 4.10](#Figure_4_10) that is non-monotonic for and .

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| **Figure 4.10** Plots of the first derivative of . The plots first derivative for and cross the real x-axis twice within the interval implying real roots in the interval . Consequently, is non-monotonic in . |

Additionally, we can easily examine the first derivative when , which is given as

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For , the first derivative is a linear function of and crosses the -axis at the origin as can be seen in [Figure 4.10](#Figure_4_10). Therefore, it has no real roots in the open interval , implying that is monotonic. Furthermore, since is numerically negative (directed distance), is a monotonically increasing function implying that , and consequently , increases with .

Heretofore we have used visualizations to analyze the conditions under which it is possible to invert the function . We now proceed to find an analytic expression that can be used to test the monotonicity of . We can, of course, use any numerical computation tool to find the cubic roots of the first derivative and verify if the first derivative has real roots in the open interval . In fact, our derivation is based on the algebraic method for solving cubic roots published by Gerolamo Cardano in his treatise *Ars Magna* in 1545 [28].

The roots of a cubic polynomial is given as:

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where,

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and

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We can determine the nature of the roots from —if , then is real resulting in one real root () and two roots that are complex conjugates (, ); if , then the imaginary terms of and vanishes resulting in all three roots being real and at least two equal; and if , then and becomes complex conjugates resulting in all three roots being real.

Since in first derivative represented by Eq. (4.67) is positive in our problem, the roots of the first derivative are same as the roots of the scaled cubic equation

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with coefficients , , , and and .

Eq. (4.72) is not suitable if ; however, we have already seen that when , the first derivative represents a line passing through the origin. For , the coefficient of the linear term—— is larger than both the constant and quadratic terms. In fact, the magnitude of tends to two, starting from a very large magnitude, as increases from one. At the same time the magnitude of tends towards zero. Therefore, for the curve represented by Eq. (4.72) never crosses the real x-axis in the interval ; where is a very small number whose exact value depends on and corresponds to an angle that is very close to . Therefore, we only need to test for the monotonicity of if .

Based on the discussion of nature of the cubic roots, we would expect that if all three roots are real, then there is a high chance that one or more of these real valued roots would lie within the interval . Indeed, based on tens of thousands of randomly generated combinations of parameters , and we have found that the only instances in which we obtain real roots of Eq. (4.72) in the interval is when . Finally, substituting the expressions for the coefficients , , , and from Eq. (4.72) into Eq. (4.71) and Eq. (4.70) we get the *sufficient condition* for to be monotonic and invertible if as:

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##### ***Algorithm for finding for known***

The Eq. (4.66) can be used to easily determine the tilted object plane orientation that is brought to focus if the lens is tiled by an angle about the center of the entrance pupil. However, it is often required to find the lens tilt angle as a function of a known object plane tilt angle . Deriving a closed-form inverse relation of the Eq. (4.66) is not easy. Therefore, we will develop an iterative algorithm for determining . The central idea behind the algorithm was introduced earlier when we saw that the general relationship between and , if the lens is pivoted about a point away from the entrance pupil, lead to a quartic plane curve equation parameterized by , , and . The point of intersection of the quartic curve with the unit-circle (and having a positive value of abscissa) yielded . In a similar vein, by substituting and in Eq. (4.66), we obtain the following implicit equation if the lens is rotated about the entrance pupil:

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Similar to our earlier observation on the reduction of complexity of the equation relating and , the implicit equation too reduces from a fourth degree *quartic plane curve* (Eq. (4.64)) to a second degree *quadratic plane curve*. We have plotted several such curves in two groups in [Figure 4.11](#Figure_4_11). The curves in each group, distinguished by the green and blue lines, belong to and respectively. In each group the various curves correspond to different values of (and ). For all curves in the figure the focal length , and the object plane distance . The shapes of these curves are almost always elliptic, although for small values of we have also observed parabolic and hyperbolic shapes. At the curve, irrespective

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| **Figure 4.11** Determination of lens tilt for known object plane tilt using point of intersection of *quadratic plane curve* with the unit circle. If a lens is rotated about the entrance pupil, the equation relating as a function of , , , and can be expressed as a quadratic plane curve with , . The intersection of this curve with the unit circle yields the unknown . The figure plots two groups of curves—the blue curves correspond to equations for different values of (, , , and ) and but the same while the green curves corresponds to . The red cross indicates a second point of intersection of one of the curves with the unit circle implying that there exist two values of ( and ) that evaluates to the same value of using Eq. (4.66). Therefore, the relationship between and is not unique for for choice of object distance () and focal length (). |

of other parameters, is a circle. The point of intersection of each quadratic curve with the unit circle in the first or third quadrants yields the lens tilt angle for positive or negative sign of object tilt angle respectively. The red cross in the figure depicts the second possible solution for (equal to ) when , implying that the system does not meet the *sufficient condition* for determining .

As discussed previously, one possible method to find the lens tilt angle for known , , , and would be to numerically compute the points of intersection of the corresponding quadratic curve with the unit circle, for example using Newton’s method, followed by selecting the point of intersection from the appropriate quadrant depending on the sign of . In our problem, since the desired lens tilt angle is constrained to reside on the unit circle, instead of fining the roots of a cubic equation, we can use an iterative algorithm to find , provided we can find a good starting point that is relatively close to the desired . In fact, we do have a good starting point—the value of for . If and the lens is rotated about the entrance pupil, the equation relating and is equivalent to rotating a thin lens represented by Eq.(4.51).

The iterative algorithm used for determining the lens tilt angle for a known value of is shown in [Table 4.3](#Table_4_3). In addition, [Figure 4.12](#Figure_4_12) and [Figure 4.13](#Figure_4_13) visually explain how the algorithm functions. [Figure 4.12](#Figure_4_12) describes how , , and evolve over iteration number as the algorithm converges to the value of starting from an initial estimate given by thin-lens approximation. The figure also shows how the step size is modified when the error changes sign. While [Figure 4.12](#Figure_4_12) illustrates how the algorithm converges to the right values of , [Figure 4.13](#Figure_4_13) shows how evolves along the unit circle for two different choices of . Both figures demonstrate that for typical optical systems the algorithm converges rapidly to the true value of the . Furthermore, more number of iterations improves the precision of the estimate. It is important to note that the system with (green) is only used to demonstrate the algorithm. Although the parameters meet the sufficient condition test, the image plane of this hypothetical system is located several meters behind the exit pupil. On the other hand, an optical system with is quite common. For both curves, and .

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| **Figure 4.12** Example determination of lens tilt angle . The plot shows the iteration variables (red), (cyan), and (blue) against the iteration number. While the range of values of and are shown on the left vertical axis, the right vertical axis represents . The true value of is , which corresponds to with , and . The algorithm starts with an initial estimate of (obtained using thin lens approx.) and converges to in the first 20 iterations. The final value of , after 48 iterations, is . The plot also shows that the step size remains constant as the error in each iteration decreases until the value of the error changes sign; at which point the step’s value and sign are modified. |

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| **Table 4.3** Algorithm for finding lens tilt required to focus on an object plane tilted by . |
| Objective  Given object plane tilt angle , determine , the lens tilt angle about the entrance pupil, required to focus on object plane that is at a distance from the lens pivot. The lens has focal length and pupil magnification . The iteration quits if either the specified accuracy for the result, , or the specified maximum steps, , is reached.  Algorithm   1. Define function: 2. Compute initial estimate of to provide a starting point: 3. Compute initial error: 4. Compute initial value of step to increment/ decrement : 5. Set iteration variables for error: , 6. Set iteration variable for angle : 7. Set step counter: 8. repeat:    1. Increment step counter:    2. Update estimate of lens tilt angle :    3. Store current error:    4. Compute new error:    5. Modify step direction and size if error changes sign: :       * Set new step size: |

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| **Figure 4.13** Inner workings of the iterative algorithm for determining given . The ellipse (blue) corresponds to , . The hyperbola (green) corresponds to , . The true values of for the blue and green curves are and respectively, shown by the blue and green spheres on the unit circle (magenta). The initial estimates for are shown by the small red spheres. The opacity, position, and numeric tags of white spheres depict the direction of convergence, the value of in the iteration and the iteration number. For , the algorithm converges at a faster rate than for as indicated by the number of steps and precision of the estimated in the two cases. It is important to note that the system with (green) is only used to demonstrate the algorithm. |

### **4.4 Summary**

We derived the formula to create a geometrically sharp image of a tilted object plane (described by the normal ) on a tilted image plane (described by the normal ) through a tilted lens (described by the rotation matrix ) as (Eq. (4.23)):

We saw that the above equation is most general and special configurations of both frontoparallel and Scheimpflug imaging readily fall out of it.

Using the above formula, we derived a set of formula for focusing on a tilted object plane, if lens’ pivot is away from the entrance pupil as (Eq. (4.59)):

and (Eq. (4.63))

If, however the lens is pivoted at the center of the entrance pupil, these formulae simplifies to (Eq. (4.65)):

and (Eq. (4.66))

Furthermore, since it is rather difficult to invert Eq. (4.63) and Eq. (4.66) to obtain a formula for the lens tilt angle as a function of object tilt angle , we use a computational technique to estimate starting from an initial estimate.

# ⬀Chapter 5

**CHAPTER 5 TITLE**

*Quote.*

—Author

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# ⬀Chapter 6

**CHAPTER 6 TITLE**

*Quote.*

—Author

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# APPENDIX⬀ A

### **Appendix A.1 Transfer of chief ray’s direction cosine for arbitrary orientation of the optical axis**

In section 3.4 we derived the expression for the transfer of the chief ray’s direction cosine from the entrance pupil to the exit pupil for a specific problem in which the optical axis was coincident with the -axis of . Furthermore, we inferred the general expression for the transfer relation—in which the optical axis is free to swivel about the origin of—from the expression obtained for the specific problem. Here we apply the method of induction to formally derive the general expression.

Eq. (3.10) accurately represents the *transfer* for the specific problem; however, we will cast the expression in a slightly different form whose raison d'être is to enable generalization—through direct application of the result. Specifically, we express the output chief ray as a linear combination of the input chief ray and the optical axis; this is possible because the two rays and the optical axis span the same (meridional) plane. Let, the standard basis vector along -axis of, represent the optical axis since the optical axis is coincident with the -axis. Then,

where and are weights, and.

Rewriting the above equation as

we can the readily obtain the weight by comparing equations Eq. (3.9) and Eq. as:

Substituting into and comparing with Eq. (3.8) yields as:

We are now ready to apply the result of the specific problem to the general problem. [Figure 3.5](#Figure_3_5) shows the schematic of the general problem—the optical axis pivots about the origin of. Let us describe the general orientation of the optical axis by the action of the rotation matrix on. The matrix may be a composition of two or more matrices that denotes a sequence of rotations about the -axis and -axis. Then,, the unit vector representing the new orientation of the optical axis, is obtained as the transformation of by the rotation matrix: or .

Since the output direction cosine, the input direction cosine, and the optical axis lie on the same plane we can write as the linear combination of and :

Note that the input direction cosine in Eq. (following the rotation of the optical axis) is, in general, different from the corresponding in Eq. even for the same object-point . This difference is due to the displacement of entrance pupil () following the rotation of the optical axis; in fact, the designation of a ray as the chief ray (from to) alters as we displace the entrance pupil. Multiplying Eq. by, we obtain:

Letting and, yields

Comparing Eqs. and , we obtain the expressions for the weights and as:

Where represents the projection of the direction cosine vector, **,** on the rotated optical axis. If we write the matrix where are the columns of , then:

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and

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Therefore, since is the third element of.

Rewriting Eq. as:

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which can be compactly written as:

Finally, substituting and yields the general expression for the direction cosines of the chief ray in the image space in terms of the pupil magnification and direction cosines in the object space as:

where.

### **Appendix A.2 The direction cosine, originating from exit pupil, has unit -Norm**

Claim: The direction cosine in the image space, obtained by the linear transformation of the direction cosinein the object space, has unit , and is the normalization term.

*Proof*.

The expression for the direction cosine in the image space is

where , is the column of the rotation matrix applied to the optical axis, , and is the pupil magnification.

Our objective is to prove. For the convenience of notation within the proof, let

where, and , the diagonal matrix with non-negative real values.

Also, we represent the columns of as Then, , and

Since is a rotation matrix, it is orthonormal (the column of, having unit length, are orthogonal to each other). Therefore, . Then,

where

and and .

Since is a diagonal matrix, we can rewrite as

Now,

Also,

Substituting in Eq. we obtain

Further, substituting into Eq. we obtain

It follows that the scalar quantity is the normalization term. ☐

# APPENDIX⬀ B

### **Appendix B.1 Derivation of Gaussian imaging equation with pupil magnification**

The familiar Gaussian imaging equation, , relates the object and image plane distances with the focal length In this formula, is the *directed distance* (numerically negative as per our sign convention) between the object plane (perpendicular to the optical axis) and the principal plane () in the object space, is the directed distance (numerically positive for *real* images) between the in-focus image plane and the principal plane () in the image space. The distances and are measured along the optical axis.

If the distances of the object and image planes are specified with respect to the entrance () and exit () pupil centers instead of the Principal planes, then the Gaussian lens formula needs to be modified. Here we derive the modified formula starting from the Gaussian lens formula. The same result was derived in [24] using a slightly different approach.

[Figure B1.1](#Figure_B1_1) shows a schematic of the entrance and exit pupils, the object and image space Principal planes, and the object and image points. In the figure, and are the distances from the Principal planes to the object and image planes, and are distances from the Principal planes to the entrance- and exit-pupils, and and are the distances from the entrance- and exit-pupils to the object and image planes. Since the entrance- and exit-pupil planes are conjugates, like the object and image planes, the Gaussian lens formula holds as follows:

and

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| **Figure B1.1** Schematic of imaging through a lens. The figure shows the object () and its image (), the object space principal plane () and the image side principal plane (), the entrance () and exit () pupils, and the associated distances along the optical axis. |

The transverse magnification between the object and image planes is given as:

For images that are *real* and inverted, the transverse magnification is numerically negative since the directed distance is numerically negative, and is numerically positive.

The pupil magnification is defined as the ratio of the exit pupil diameter to the entrance pupil diameter. It is also the ratio between the exit pupil and entrance pupil distances (measured from the principal planes), just like the transverse magnification between any conjugate planes:

Equating Eqs. and , we obtain

Further, substituting and in the above equation, and using Eqs. and , we obtain a relationship between pupil magnification, transverse magnification, the object and image plane distances (specified with respect to the pupils) as:

Substituting and in Eq. and equating with Eq. yields

which after cross-multiplication and cancellations of common terms produces

Dividing throughout by, and substituting by the pupil magnification, and by we obtain:

Where,

|  |  |
| --- | --- |
|  | Pupil magnification. |
|  | Directed distance from the entrance pupil to the object plane. |
|  | Directed distance from the exit pupil to the image plane. |
|  | Focal length. |

Note that Eq. is valid even if the and denote distances from the principal planes, provided we let. This outcome is indeed consistent with geometric optics theory, according to which the magnification between the principal planes is unity. In fact, Eq. (4.14) is more general than the Gaussian Lens formula in that it relates a pair of conjugate planes with any other pair of conjugate planes for which the transverse magnification (between the planes) is known. When one of the pairs happen to be the principal planes ( and) between which the magnification is one, we obtain the Gaussian Lens formula.

Finally, we also obtain the equations for computing the entrance- and exit-pupil distances from the respective principal planes by substituting Eq. into Eq. as

### **Appendix B.2 A brief account on the significance of pupil magnification**

Although a pupil magnification close to one is a desirable property from the point of view of distortion in the presence of orientation misalignment, it seems to be hardly a critical design choice for most practical lenses except for those used for Scheimpflug photography as evident in plot of pupil magnifications in [Figure B2.1](#Figure_B2_1). In addition, the figure shows that telephoto lenses have pupil magnification less one, and retrofocus wide-angle lenses have pupil magnification greater than one. The telephoto lenses employ a negative focal length group near the sensor plane to accommodate a long focal length lens into a compact body. Consequently, the exit pupil height (the image of the limiting aperture at the image side) is smaller compared to the entrance pupil height (the image of the limiting aperture at the object side). Therefore, in telephoto lenses the pupil magnification less than one. On the other hand, a negative focal length group is placed at the front in short focal length, retrofocus lenses to create space between the lens and sensor which results in a larger exit pupil height compared to the entrance pupil height. Thus, the pupil magnification of retrofocus lenses are greater than one.

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| **Figure B2.1** Pupil magnification in a wide variety of lenses that form *real* images. The figure demonstrates the absence of any correlation between pupil magnification and focal length. In addition, only 20 in the sample 120 (or one in six) lenses have pupil magnification in the range. Over 90% of all lenses have pupil magnification greater than 0.5. We obtained the samples from the Zemax Zebase library, which is a comprehensive catalogue of well-designed professional lenses. |

The F-number (or F/#) is the ratio of the effective focal length (distance between the image side Principal plane to the image plane with the object at infinity) to the paraxial entrance-pupil diameter [ref]

This infinite conjugate F/# is commonly specified by lens manufacturers on lens bodies.

The pupil magnification is the ratio of the exit-pupil diameter to the entrance-pupil diameter :

Substituting from Eq. into Eq. , we obtain

When the object is at infinity, the distance between the image plane and the exit pupil is obtained from Eq. as

Substituting in place of in Eq. yields the alternative and equivalent definition for F-number—*as the ratio of the exit-pupil-to-image-plane distance to the exit-pupil diameter*:

where is the distance from the exit pupil to the image plane, and is the diameter of the exit pupil.

The F-number along with the wavelength determines the diffraction limited spatial resolution of optical imaging systems at the image plane as given by the equation [ref]

For finite conjugate imaging, the object-plane-to-entrance-pupil distance decreases concomitant with an increase in exit-pupil-to-image-plane distance. This increase in the image plane distance effectively increases the F-number. The expression for the effective F-number in terms of the pupil magnification is derived next.

Substituting (where is the transverse magnification) from Eq. into Eq. , followed by simple algebraic steps yields:

We have established that the F-number (for infinite conjugate) is the ratio of the exit-pupil-to-image-plane distance to the exit pupil diameter. To obtain the effective F-number at finite conjugates, we substitute Eq. , the expression for the image-plane distance for finite conjugate imaging, into Eq. :

Further, substituting (Eq. ) and replacing with we obtain:

where is the transverse magnification (numerically negative for *real* images).

Now, we can obtain a more accurate equation for diffraction limited spatial resolution that is equally valid for both finite and infinite conjugate imaging by substituting Eq. into Eq.

where is the wavelength, is the standard F-number defined for infinite conjugate imaging, is the pupil magnification, and is the transverse magnification (numerically negative for *real* images). When the object is at infinity, and Eq. reduces to the optical resolution expression for infinite conjugate imaging.

To Investigate.

### **Appendix B.3 Estimation of pupil locations, and pupil magnification**

TO DO:

Describe why you would like to estimate these parameters

The entrance-pupil is the

Eliminating by substituting into Eq. (7.6) following simple algebraic manipulation we get

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where is the transverse magnification, is the distance of the object plane from the entrance-pupil, is the focal length. and are numerically negative for *real* images.

# REFERENCES

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1. While this statement is true for imaging systems with Fresnel number () much greater than unity, systems with smaller Fresnel number, such as long focal-length laser systems, exhibit a focal shift in the location of the maximum intensity from the geometrical focus point towards the aperture. [↑](#footnote-ref-1)
2. In the presence of spherical aberrations, the chief ray goes through the center of the aperture but may not exactly go through the center of the pupils [19,20]. [↑](#footnote-ref-2)