19.1 in MTW

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I. DERIVATION OF METRIC FAR OUTSIDE A WEAKLY GRAVITATING BODY

A. tt - component

The trace-reversed stress-energy tensor is defined as:

$$\bar{T}_{\mu\nu} \equiv T_{\mu\nu} - \frac{1}{2} \eta_{\mu\nu} \eta^{\sigma\rho} T_{\sigma\rho}, \tag{1}$$

where

$$\eta^{\sigma\rho}T_{\sigma\rho} = -T_{00} + T_{ii}. (2)$$

The tt component is:

$$\bar{T}_{00} = T_{00} - \frac{1}{2}T_{00} + \frac{1}{2}T_{ii} = \frac{1}{2}\left(T_{00} + T_{ii}\right) = \frac{1}{2}\left(T^{00} + T^{ii}\right). \tag{3}$$

$$h_{00} = \int d^3x' \frac{4\bar{T}_{00}(t - |x - x'|, x')}{|x - x'|} =$$
(4)

$$=4\int d^3x' \left(\bar{T}_{00}(t-r,x')+\bar{T}_{00,0}(t-r,x')\left(r-|x-x'|\right)+\cdots\right)\frac{1}{|x-x'|}$$
(5)

$$=4\int d^3x' \left(\bar{T}'_{00} + \bar{T}'_{00,0} \left(x^j \frac{x'^j}{r} + \frac{1}{2} \frac{x^j x^k}{r} \frac{x'^j x'^k - r'^2 \delta_{jk}}{r^2} + O\left(\frac{1}{r^2}\right)\right) + \cdots\right) \times$$

$$\left(\frac{1}{r} + \frac{x^l x'^l}{r^3} + O\left(\frac{1}{r^3}\right)\right). \tag{6}$$

The term with \bar{T}'_{00} (no time derivatives) is:

$$\frac{2}{r} \int d^3x' \left(T^{00} + T^{ii} \right) + \frac{2x^l}{r^3} \int d^3x' \left(T^{00} + T^{ii} \right) x'^l + O\left(\frac{1}{r^3}\right). \tag{7}$$

Using (19.4b), (19.6a), (19.7a), and (19.7b) from the textbook it becomes:

$$(7) = \frac{2M}{r} + \frac{1}{r} \int d^3x' \, T^{00}_{,00} r'^2 + \frac{2x^l}{r^3} \int d^3x' \, \left(T^{0i} x'^i x'^l - \frac{1}{2} T^{0l} r'^2 \right)_{,0} + O\left(\frac{1}{r^3}\right). \tag{8}$$

The term with $\bar{T}'_{00,0}$ (first-order time derivative) is:

$$\frac{2}{r^2} \int d^3x' \left(T^{00} + T^{ii} \right)_{,0} \left(x^j x'^j + \frac{x^k x^k x'^j x'^k}{r^2} + \frac{1}{2} x^j x^k \frac{x'^j x'^k - r'^2 \delta_{jk}}{r^2} \right) + O\left(\frac{1}{r^3}\right) \tag{9}$$

$$= 2 \int d^3x' \left(T^{00} + T^{ii} \right)_{,0} \left(\frac{x^j x'^j}{r^2} + \frac{\left(3x'^j x'^k - r'^2 \delta_{jk} \right) x^j x^k}{2r^4} \right) + O\left(\frac{1}{r^3} \right). \tag{10}$$

Every term in (8) and (10) except for $\frac{2M}{r}$ is precisely canceled by gauge transformation: $\tilde{h}_{00} = h_{00} - 2\partial_0 \xi_0$.

B. ti - components

Similarly, we can compute the n=0 term of the h_{0i} :

$$\frac{4}{r} \int d^3x' T_{0i} + \frac{4x^k}{r^3} \int d^3x' T_{0i} x'^k + O\left(\frac{1}{r^3}\right). \tag{11}$$

Applying zero momentum (19.4a) condition and using (19.7c) from the textbook:

$$(11) = -\frac{4x^k}{r^3} \int d^3x' T^{0i} x'^k + O\left(\frac{1}{r^3}\right)$$
 (12)

$$= -\frac{2x^k}{r^3} \int d^3x' \left(T^{0k} x'^l - T^{0k} x'^i \right) \tag{13}$$