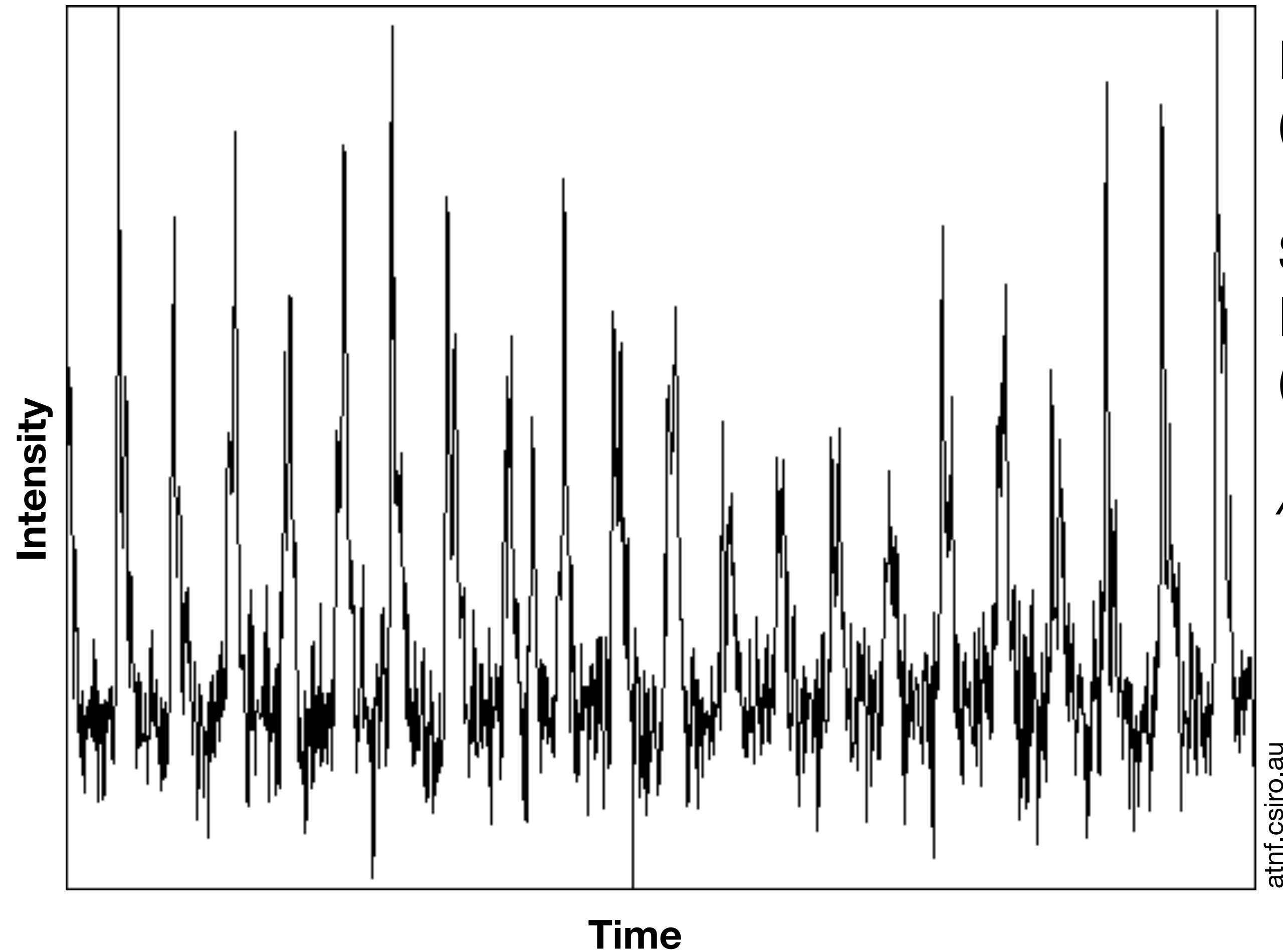


# Introduction to Pulsars

Dipankar Bhattacharya  
IUCAA, Pune

# PULSARS

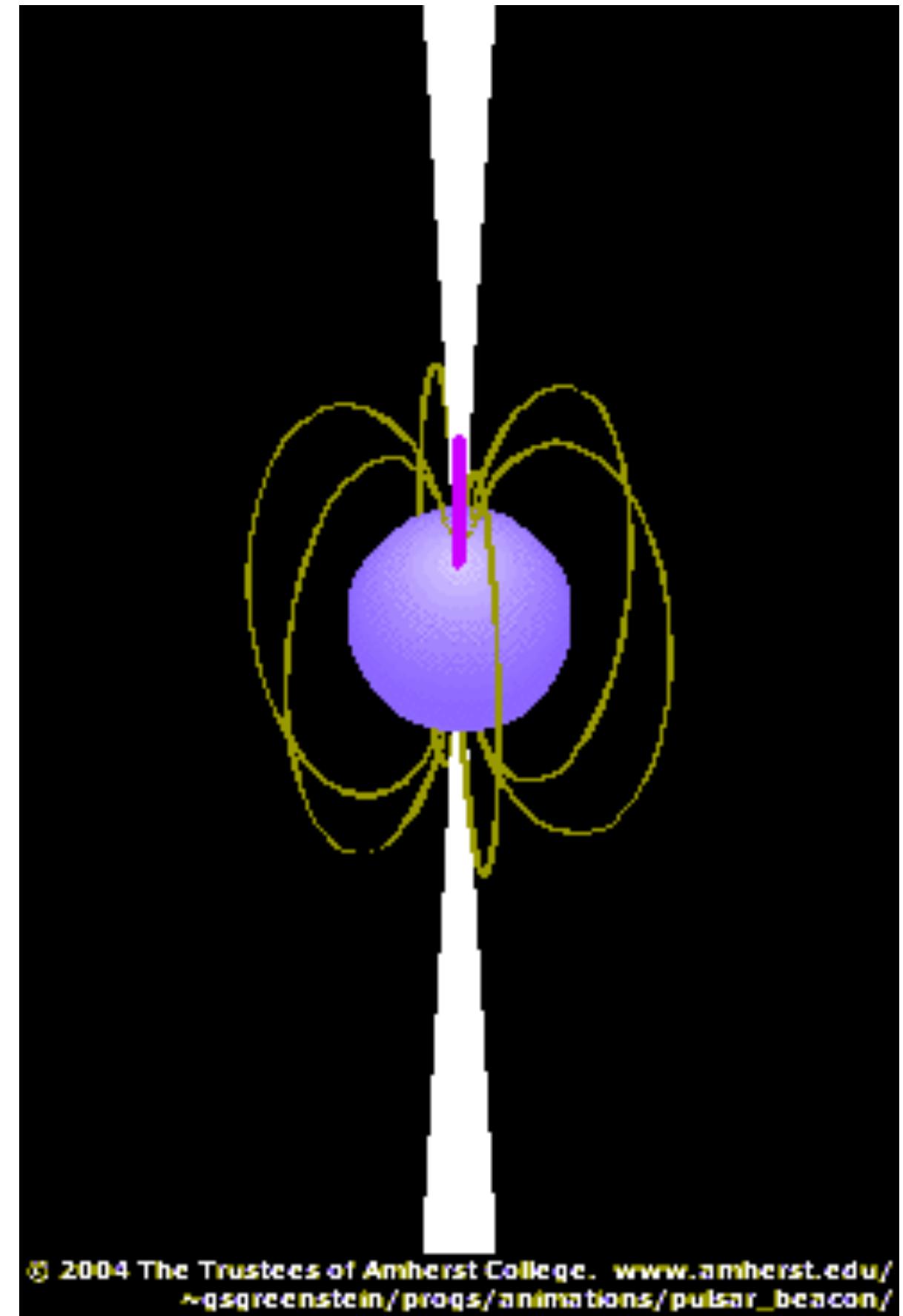
Emission received as a periodic sequence of pulses



Most seen in Radio band  
(2609)

Some at higher energy:  
IR-gamma rays  
(249)

ATNF catalogue June 2019



*Neutron stars extremely dense objects - average density  
 $10^{15}$  times that of water*

Spinning Neutron Star  
Mag. axis inclined to spin  
Radiation beam from  
magnetic poles

# Passage of matter to High Density

Air:  $\sim 10^{-3}$  g/cm<sup>3</sup>

Diffuse gas:  $\sim 10^{-24}$  g/cm<sup>3</sup>

Water:  $\sim 1$  g/cm<sup>3</sup>

Stars:  $\sim 10^{-2} - 10^2$  g/cm<sup>3</sup>

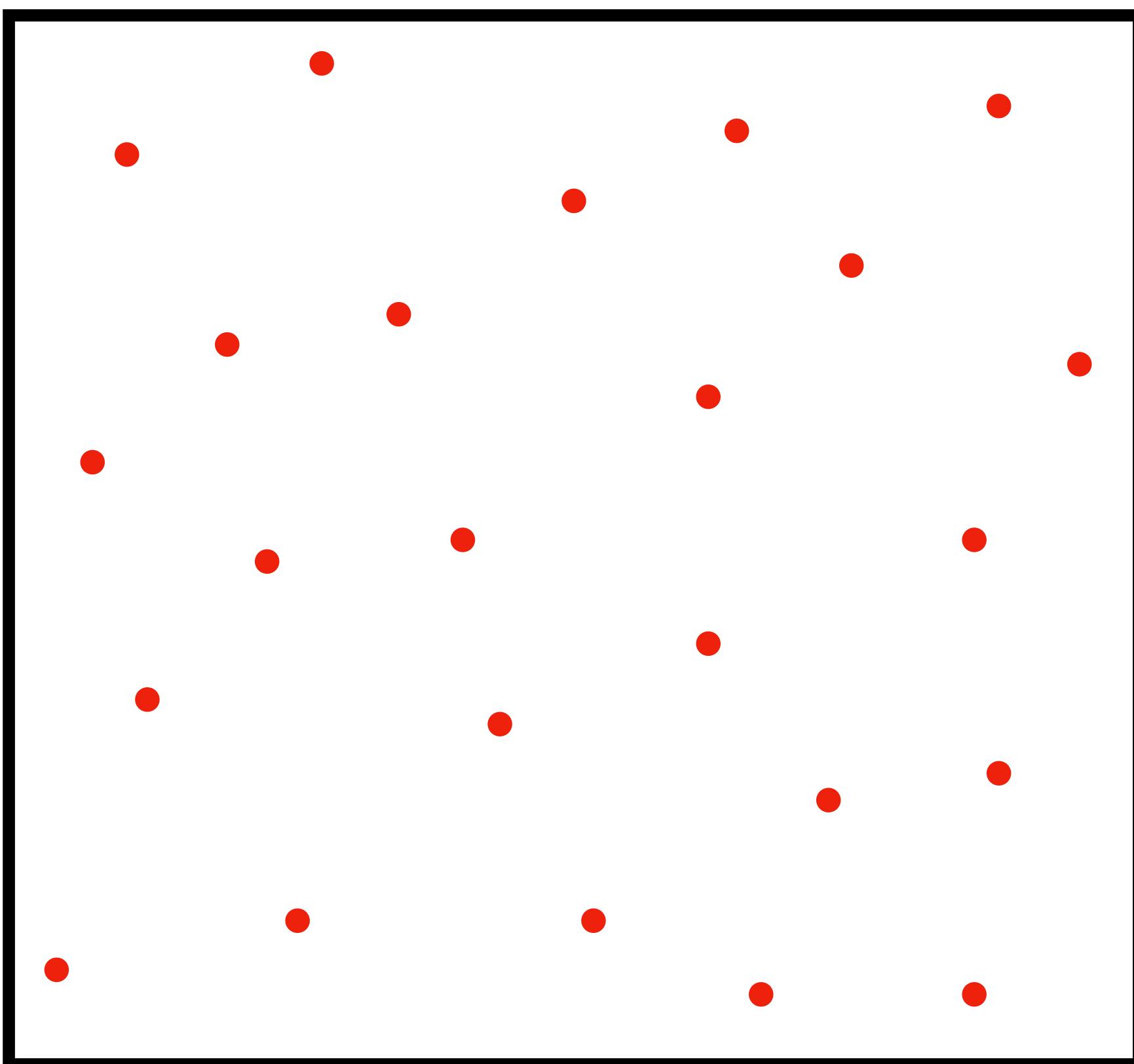
Heavy metals:  $\sim 20$  g/cm<sup>3</sup>

Planets:  $\sim 0.5-10$  g/cm<sup>3</sup>

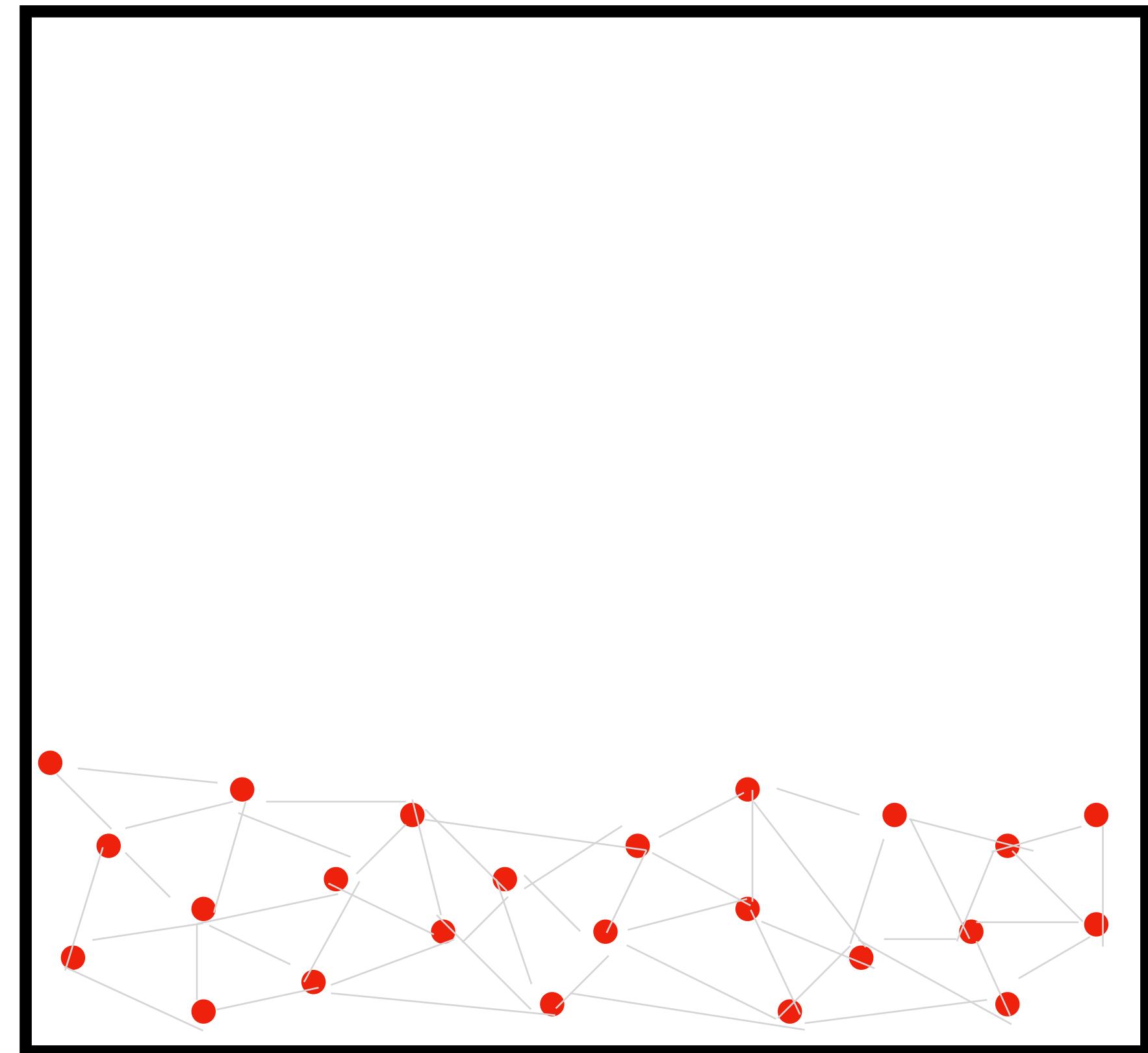
White Dwarf:  $\sim 10^6$  g/cm<sup>3</sup>

**Neutron Star:  $\sim 10^{15}$  g/cm<sup>3</sup>**

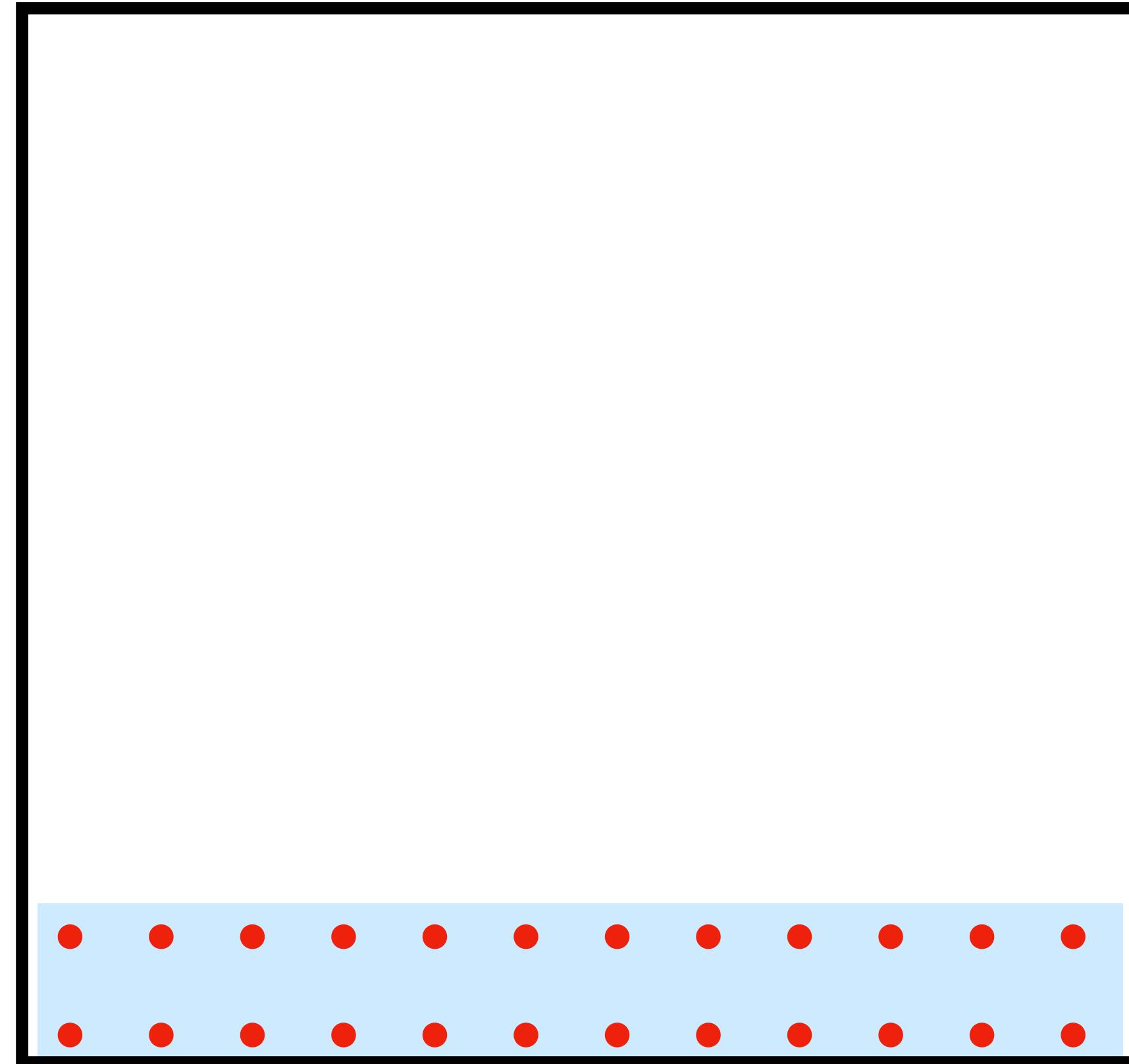
## Gas: Atoms/molecules in thermal motion



## Liquid: coulomb interaction



Solid: crystal order, free electrons,  
stronger coulomb interactions



**Terrestrial dense matter is held together by electrical forces**  
balance between attraction and repulsion

# Getting denser

Apply a squeeze

Strip electrons from atoms

Under the influence of external force, ions can get closer

Density rises until coulomb repulsion balances external force

Natural source of external squeezing force: **Gravity**

Coulomb force balancing gravity: *planets and moons*

If more and more matter is added, gravity will eventually overcome the coulomb force

Planets cannot get heavier than about Jupiter

# Getting denser

Stars too are held together by gravity

Much heavier than planets

Opposition to gravity comes from Thermal Pressure of gas

Gas is kept hot by nuclear fusion

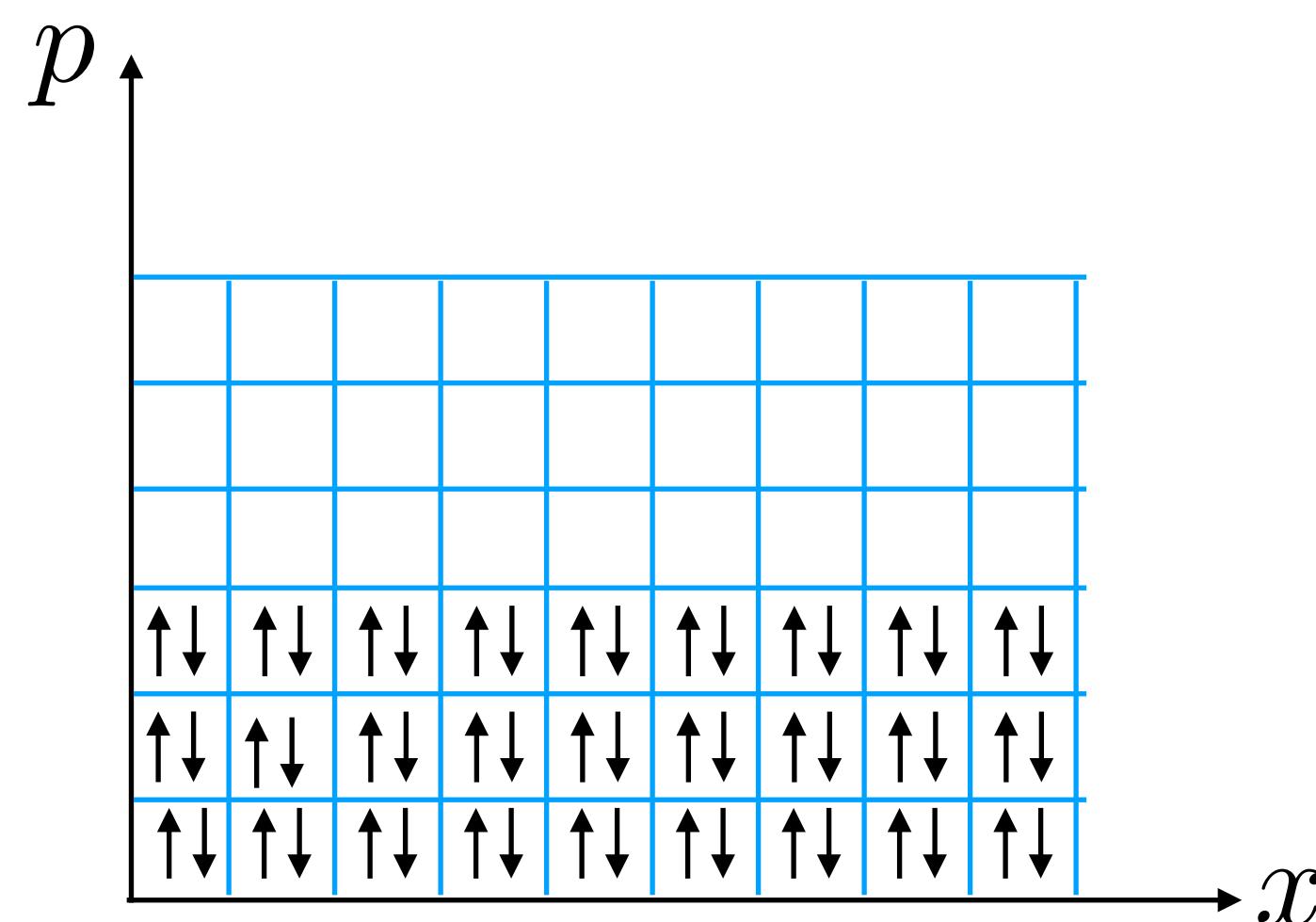
Eventually nuclear fuel will run out, gravity will win over thermal pressure

The stellar core will collapse. Coulomb force inadequate to provide support, matter will continue to get much denser

# A new type of pressure: Degeneracy

Electrons are elementary particles with half-integral spin:  $\hbar/2$

Such particles obey Fermi-Dirac statistics, leading to Pauli's exclusion principle:



Phase space (6-dim):  
position and momentum

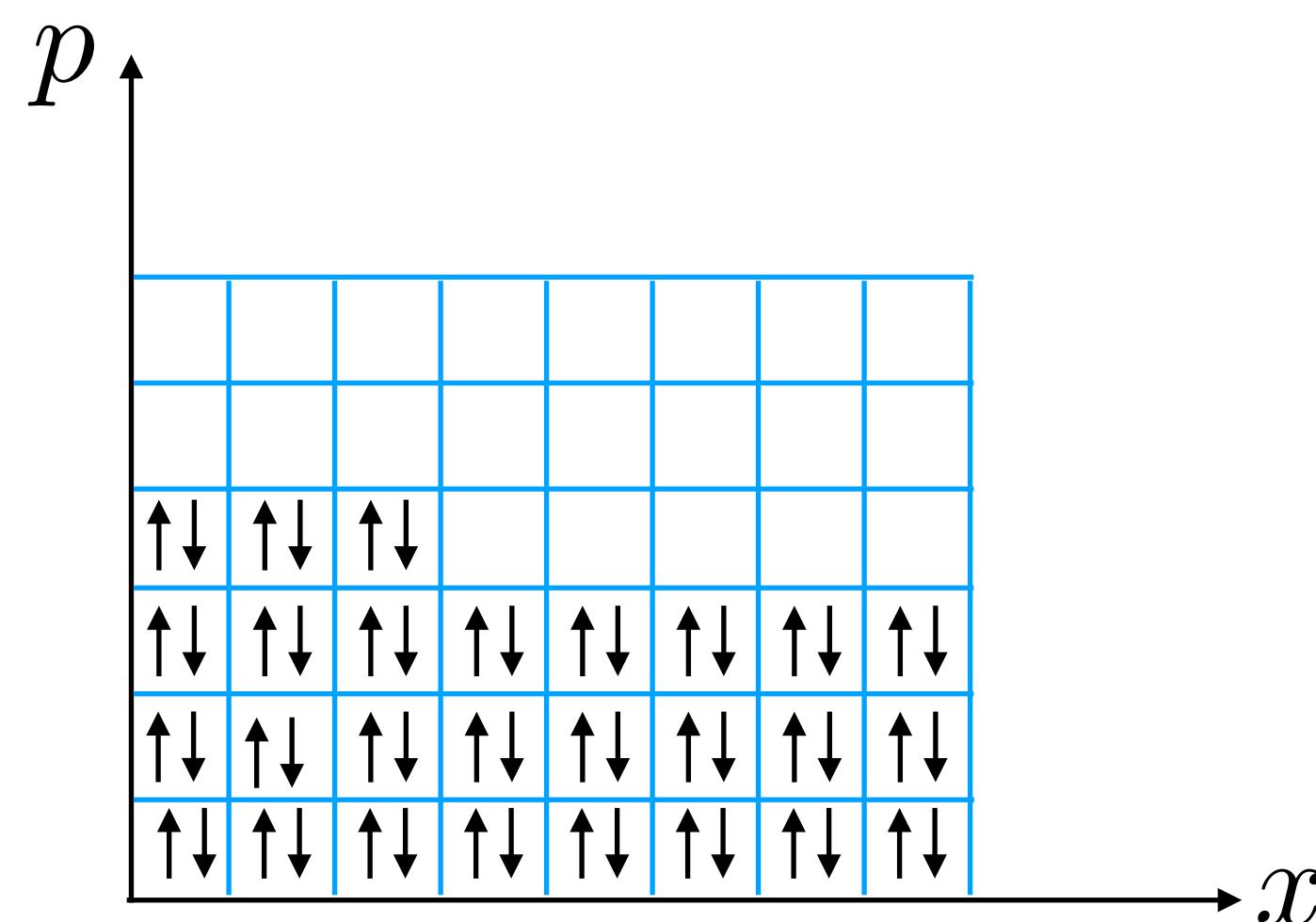
Divided into quantum cells of  
volume  $h^3$

Each cell may contain at  
most 2 electrons of opposite  
spin

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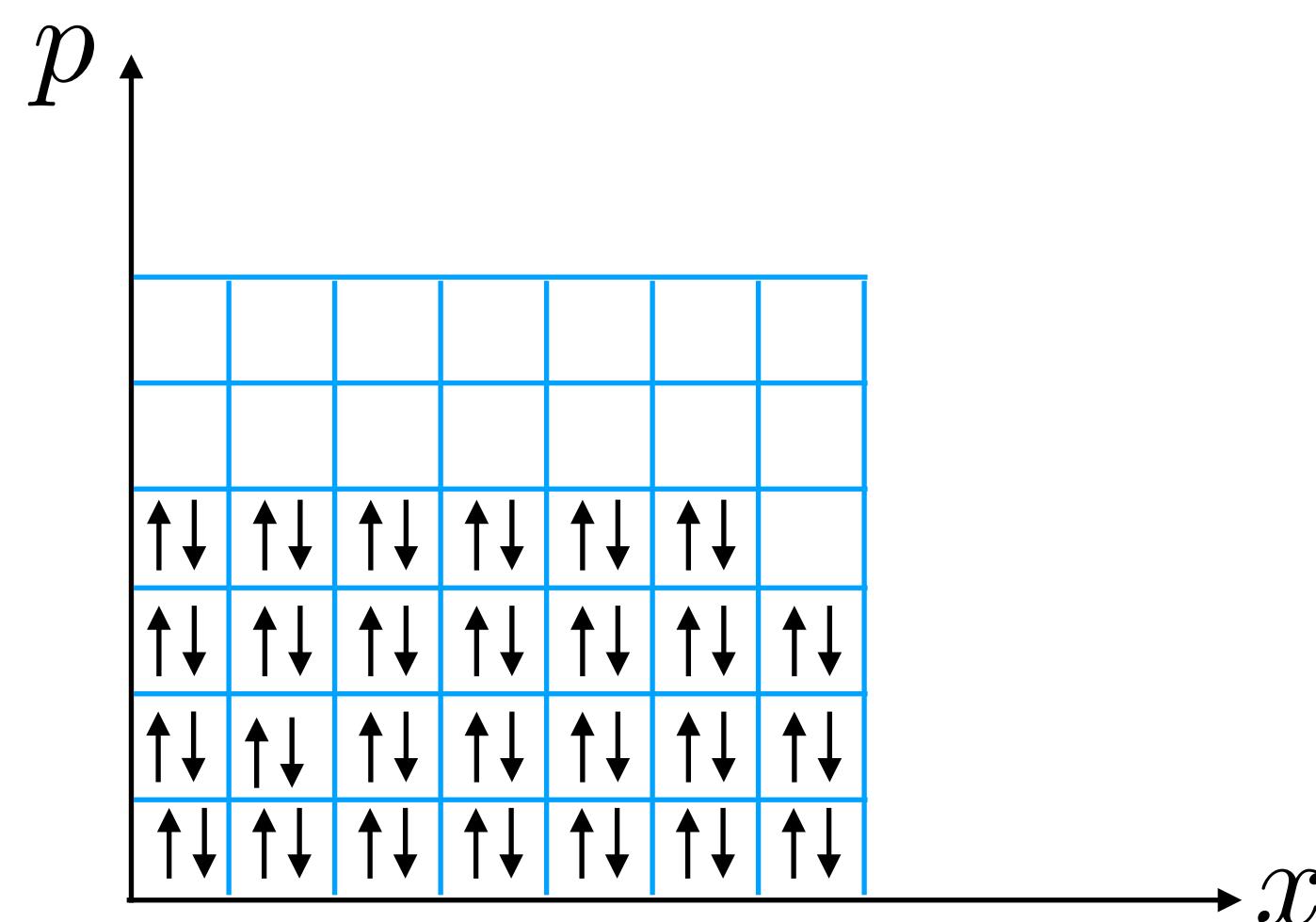
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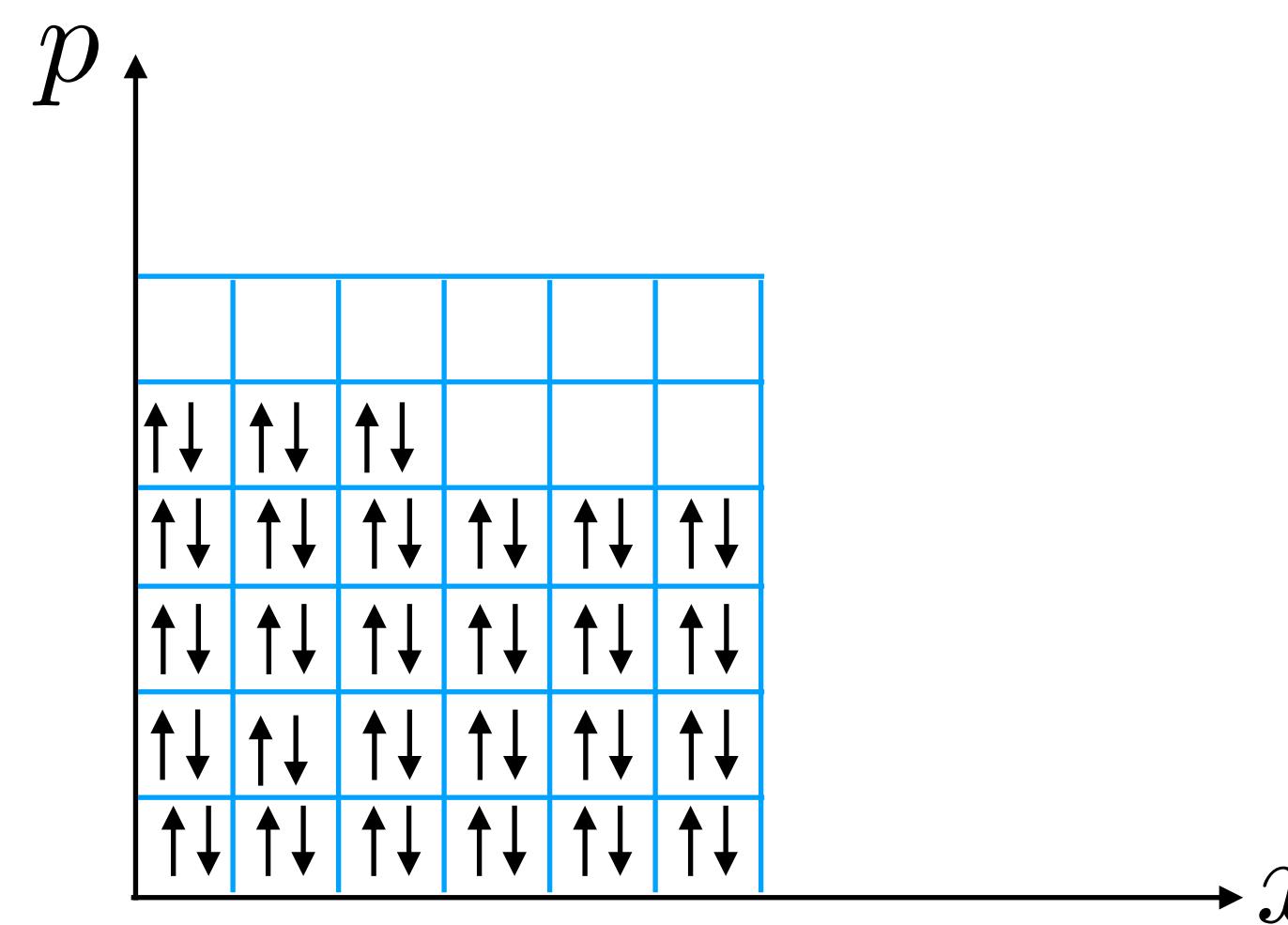
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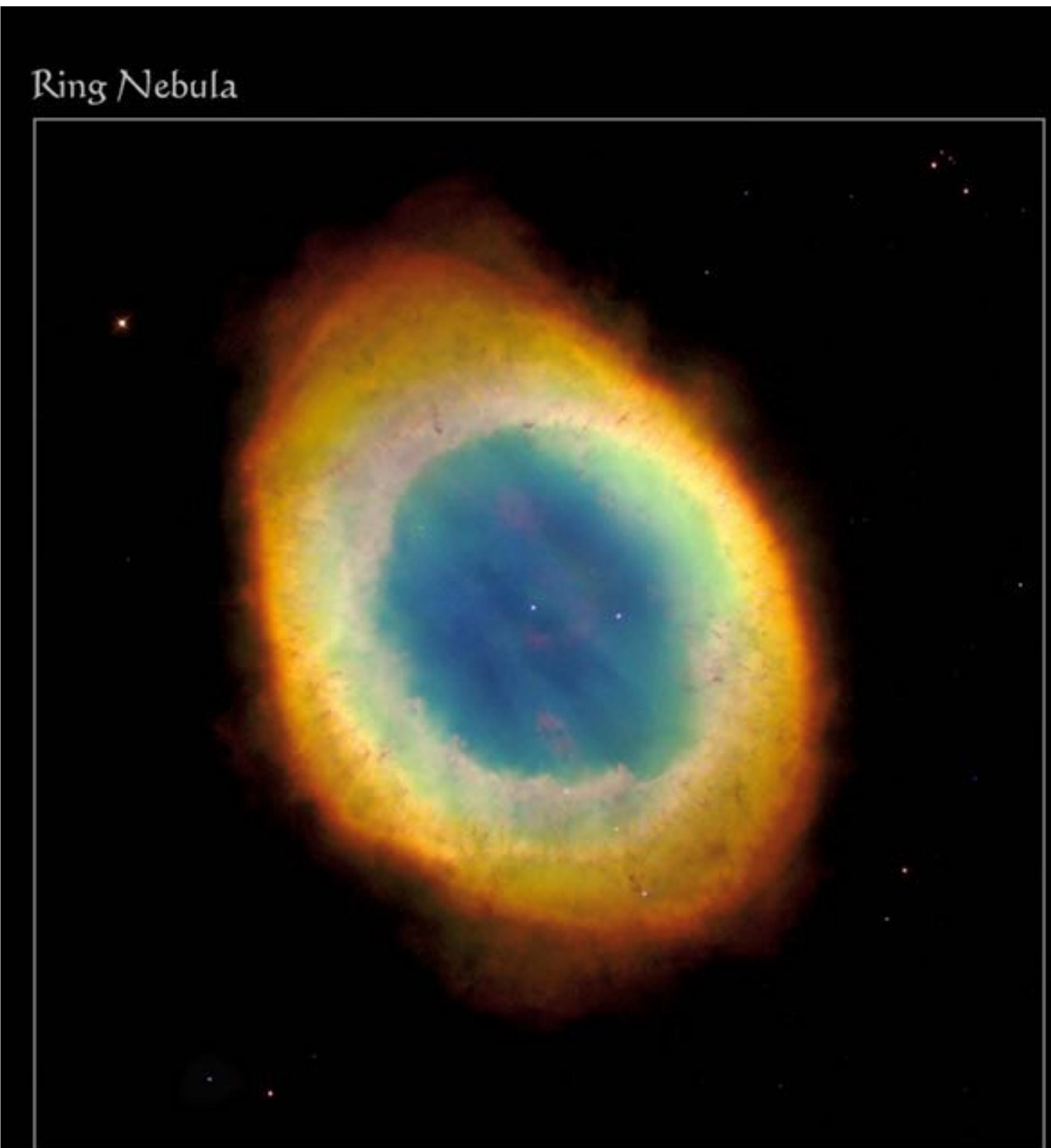
Each cell may contain at  
most 2 electrons of opposite  
spin

*The denser the material gets, the higher the momentum rises*  
This momentum is the source of Degeneracy Pressure

Stars held by electron degeneracy pressure:  
***White Dwarfs***

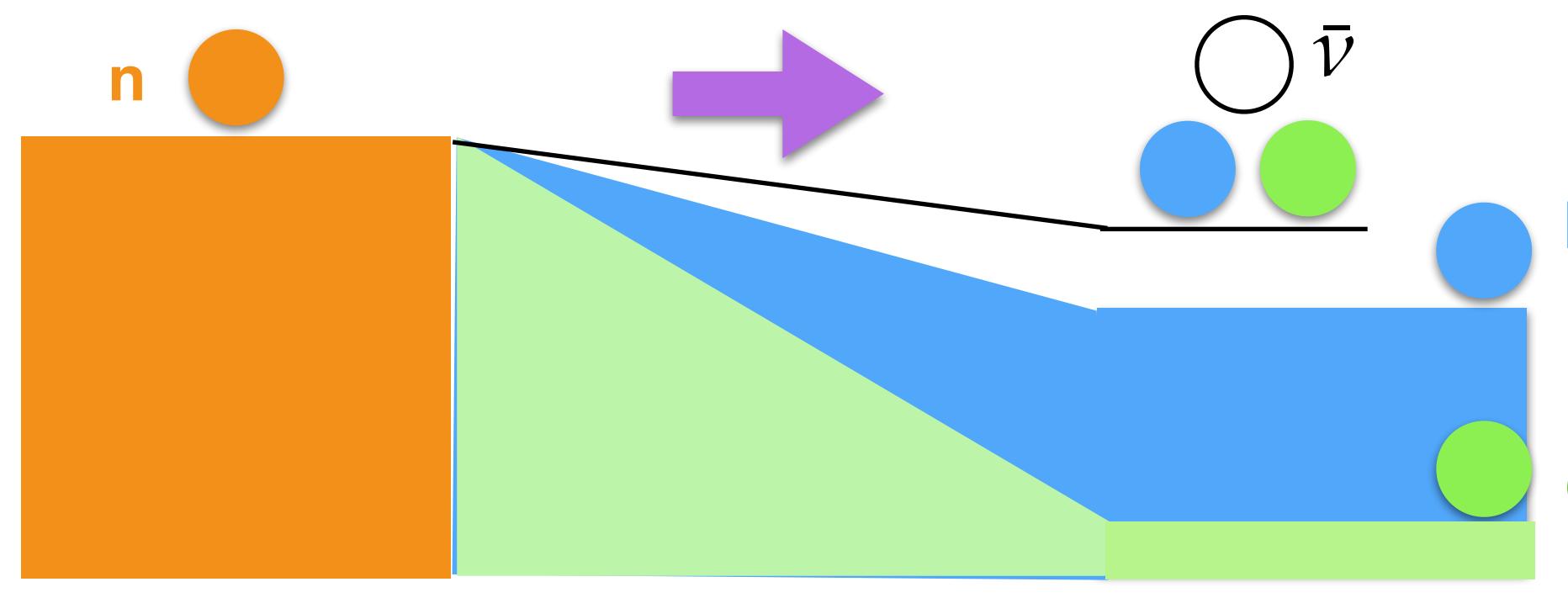
Can support up to 1.4 Solar Mass  
(Chandrasekhar Limit)

Heavier objects without nuclear burning  
will be squeezed further => neutronisation

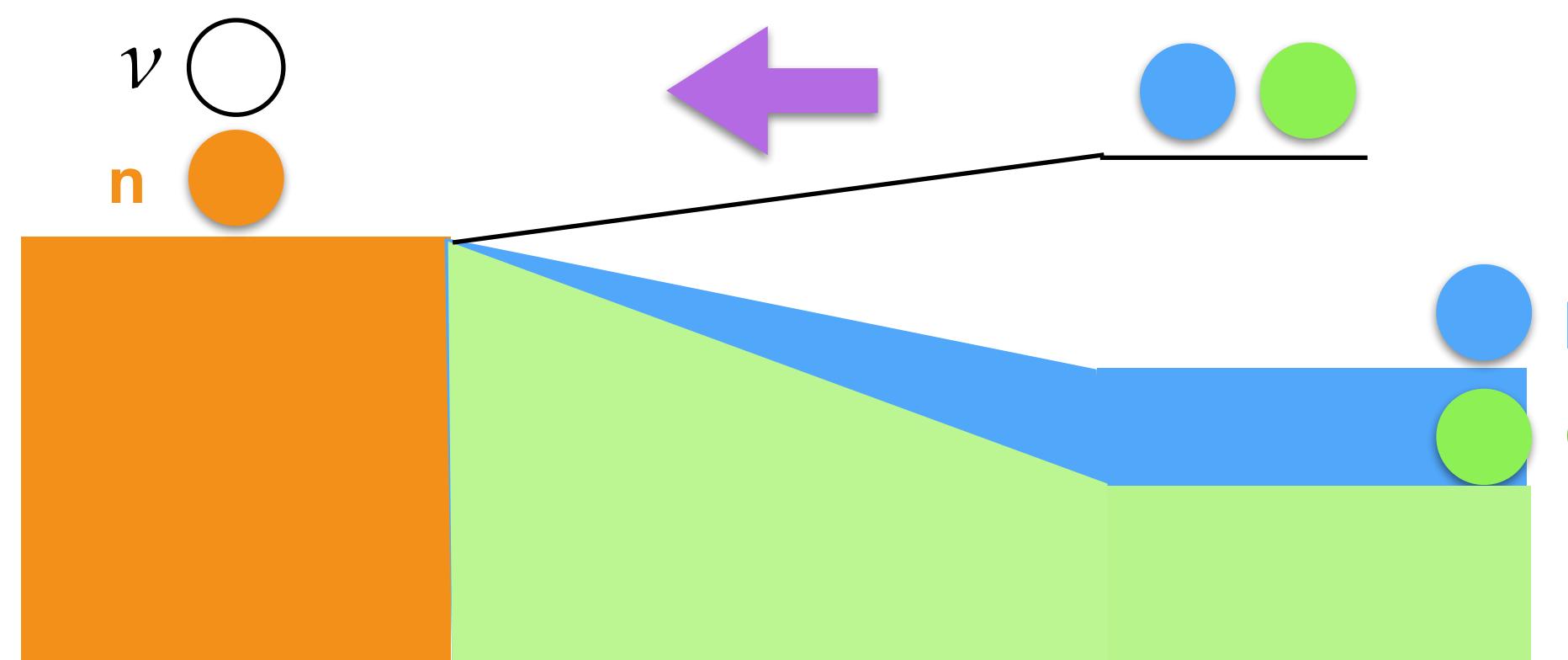


Hubble  
Heritage

# Neutronisation and $\beta$ -equilibrium



Low density



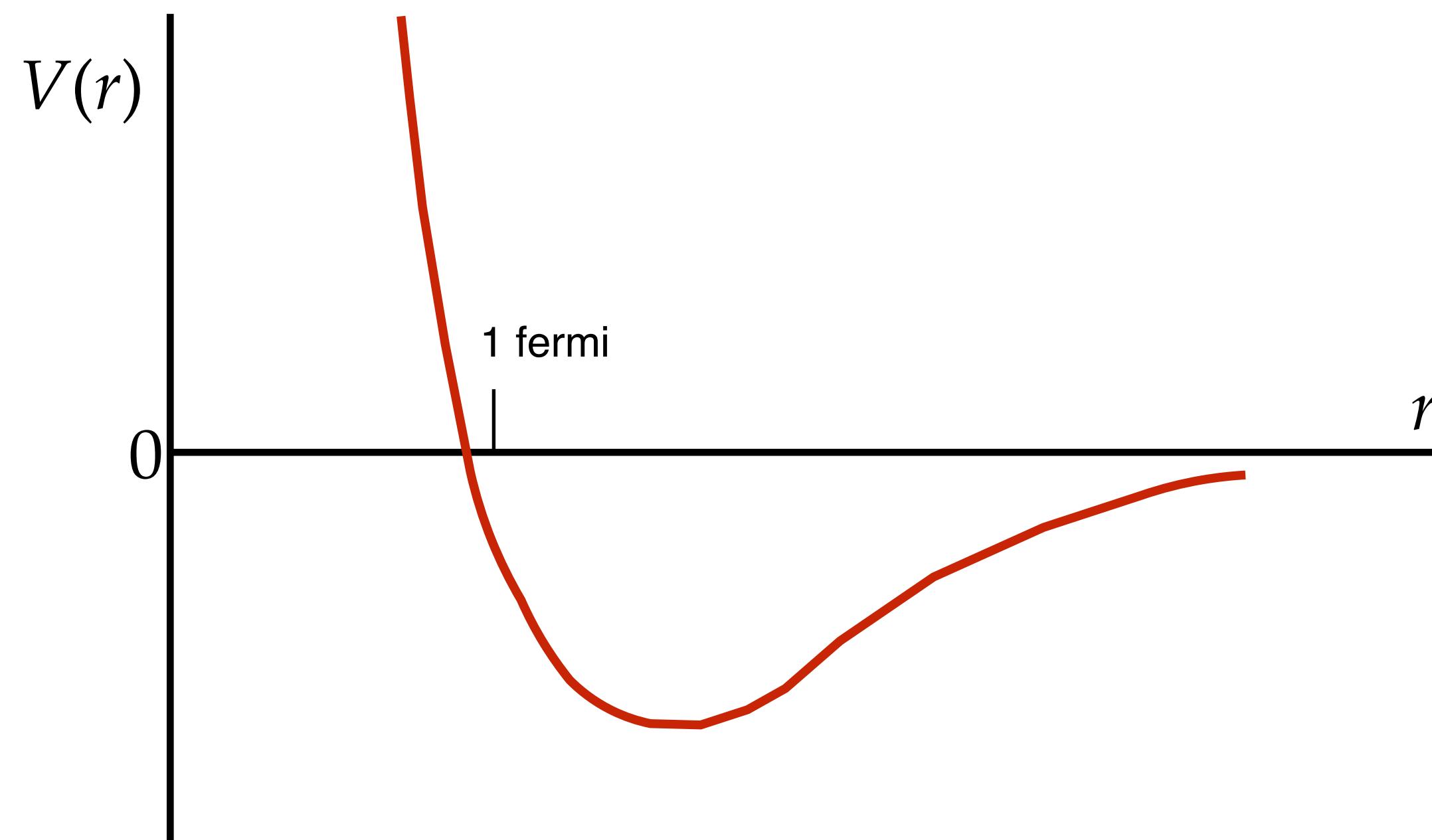
High density

$$\mu_n = \mu_p + \mu_e$$

$$\mu = \sqrt{p_F^2 c^2 + m^2 c^4}$$

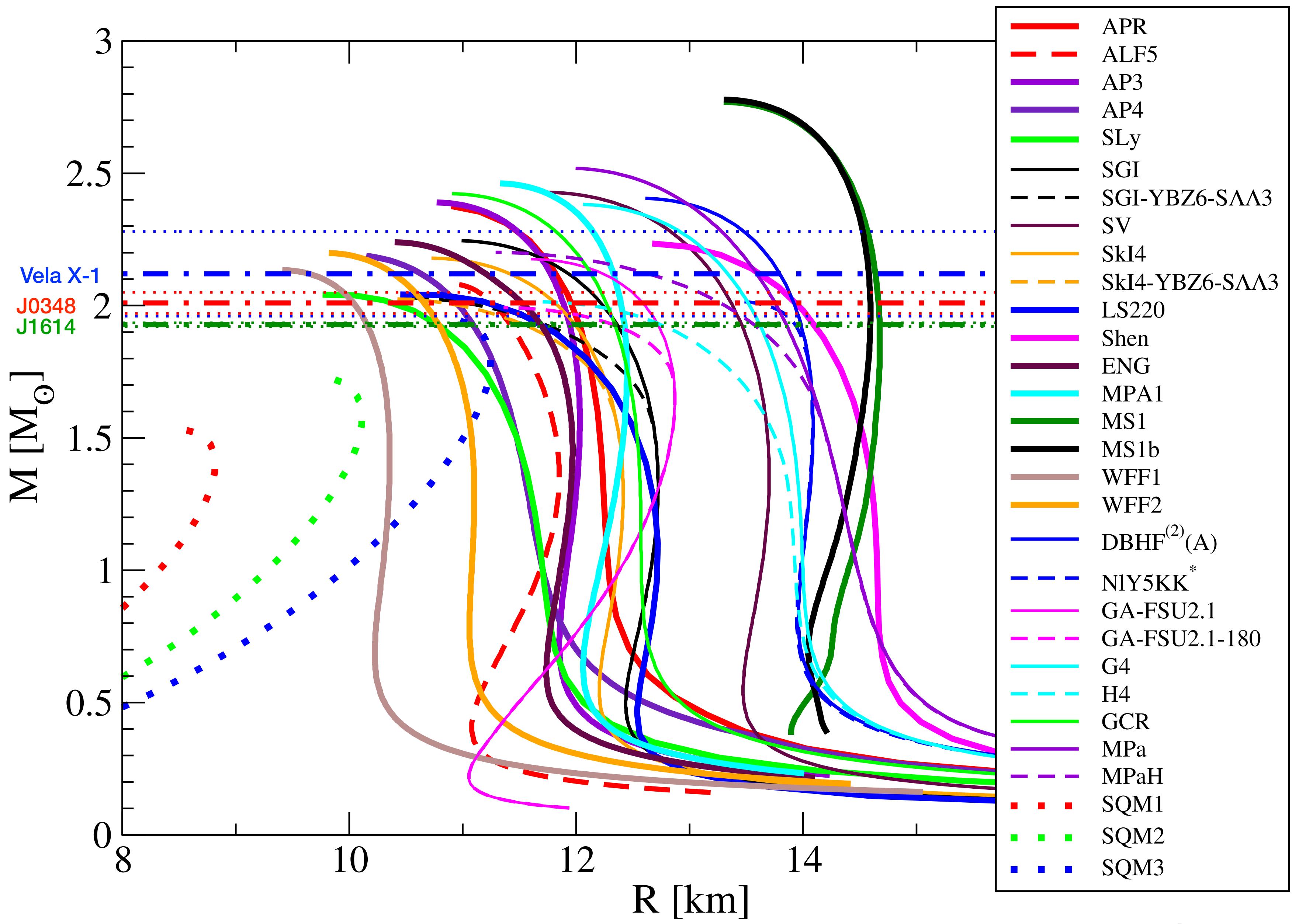
The neutron-rich matter produces pressure via

- Neutron degeneracy
- Repulsive strong interaction



At inter-nucleon distance below 1 fermi, strong interaction is repulsive, thus resists gravity.

*Nature of the repulsive strong interaction not yet known in detail. This leads to uncertainty in Pressure-Density relation (Equation of State)*



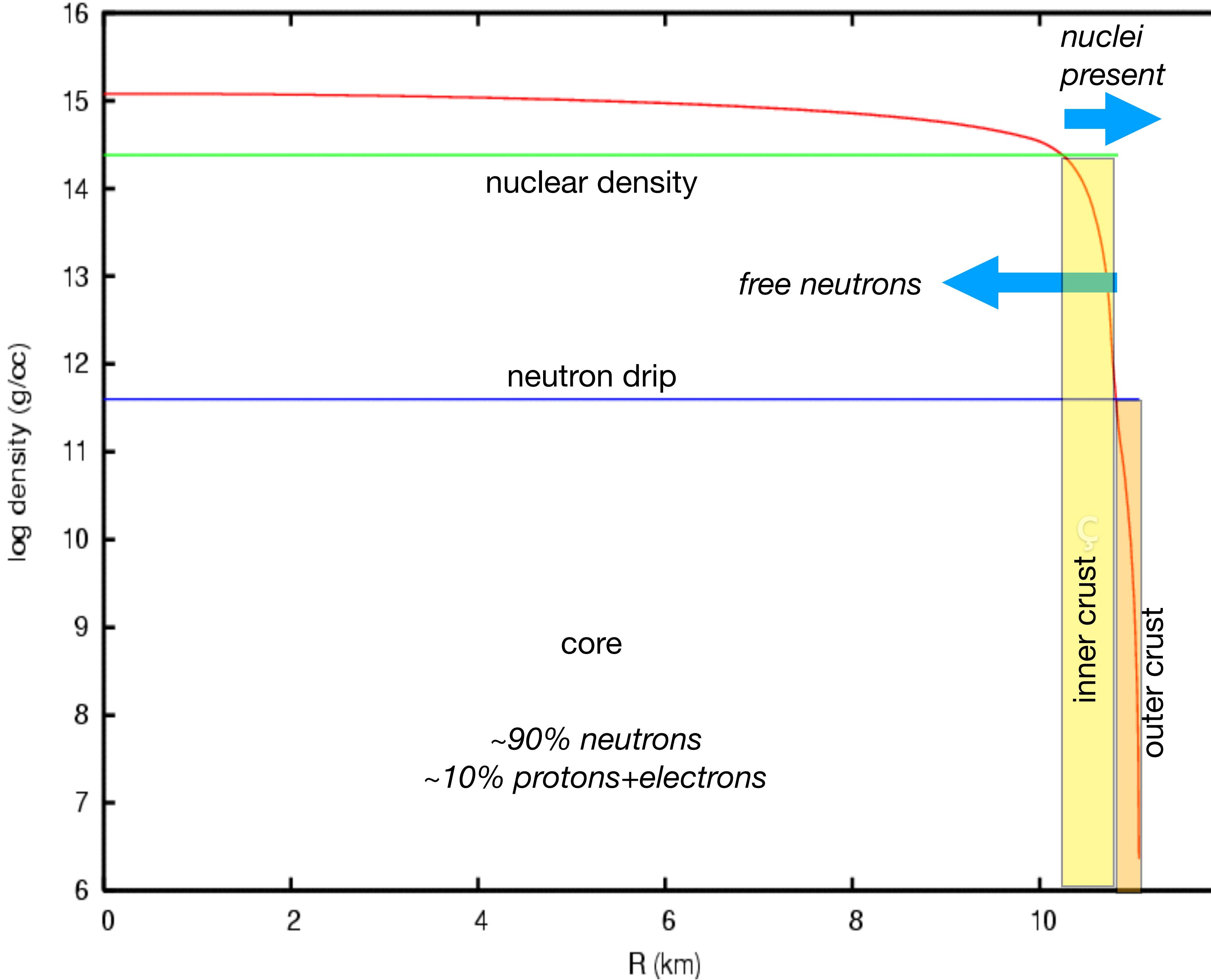
Current knowledge of nuclear forces is inadequate to accurately predict the pressure-density relation (Equation Of State) at very high density.

Leads to uncertainty in neutron star Mass-Radius relation

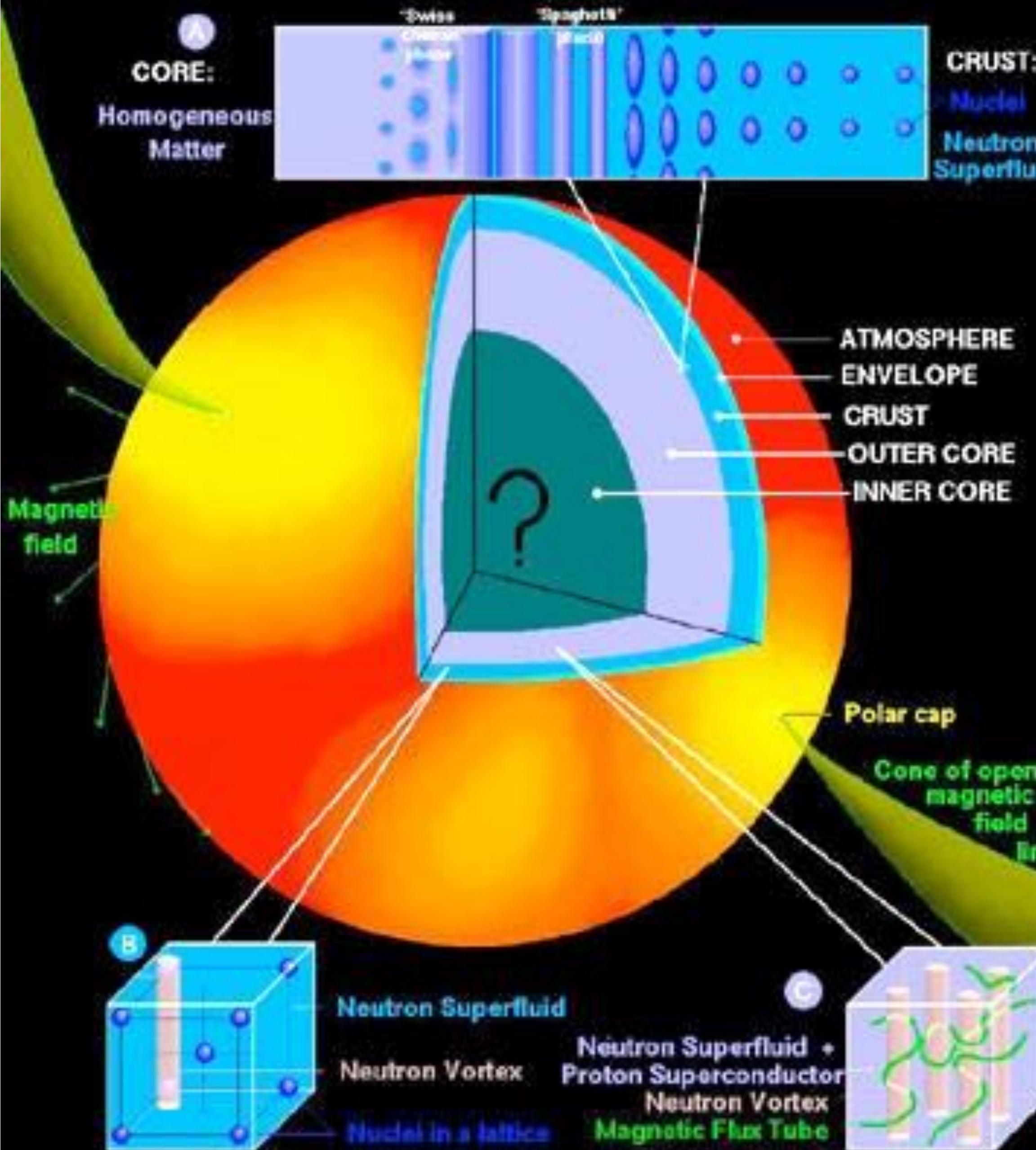
Measurement of NS mass, radius can thus unveil the fundamental physics involved

Pulsar timing a key technique for NS mass measurement

# Neutron Star Density Profile in a typical EoS



## A NEUTRON STAR: SURFACE and INTERIOR



In Neutron Star interior, free neutrons form a superfluid  
and the protons form a superconductor

These are the highest  $T_c$  superconductor/superfluid  
known in nature. Cooper pairs are formed due to strong  
interaction.  $T_c \sim 10^{8-10}$  K

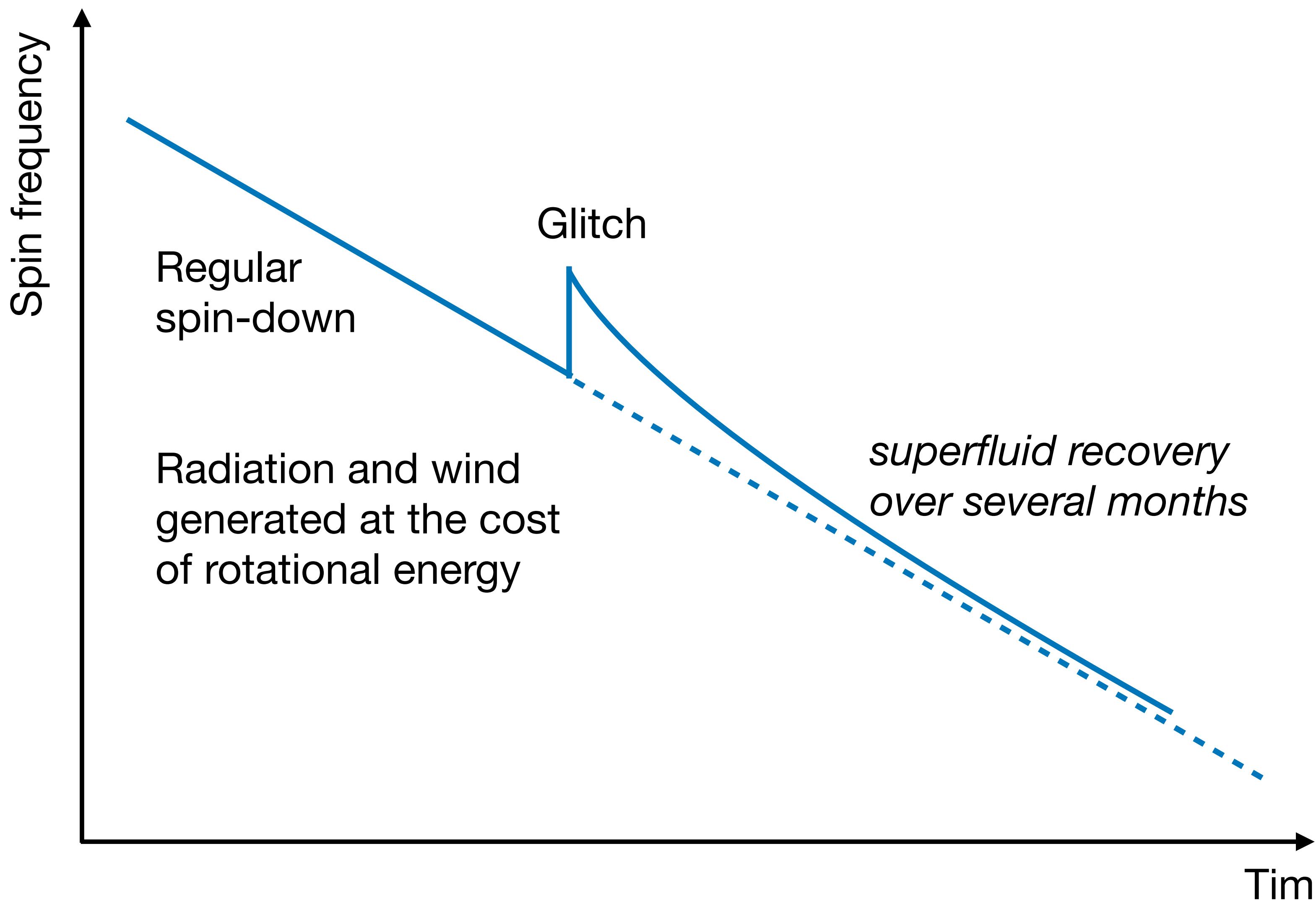
Neutron stars rotate fast ( $P \sim$  ms - min) and have  
strong magnetic fields (Teragauss)

Quantized Vortices carry rotation through the  
Superfluid, fluxoids carry magnetic field through the  
Superconductor

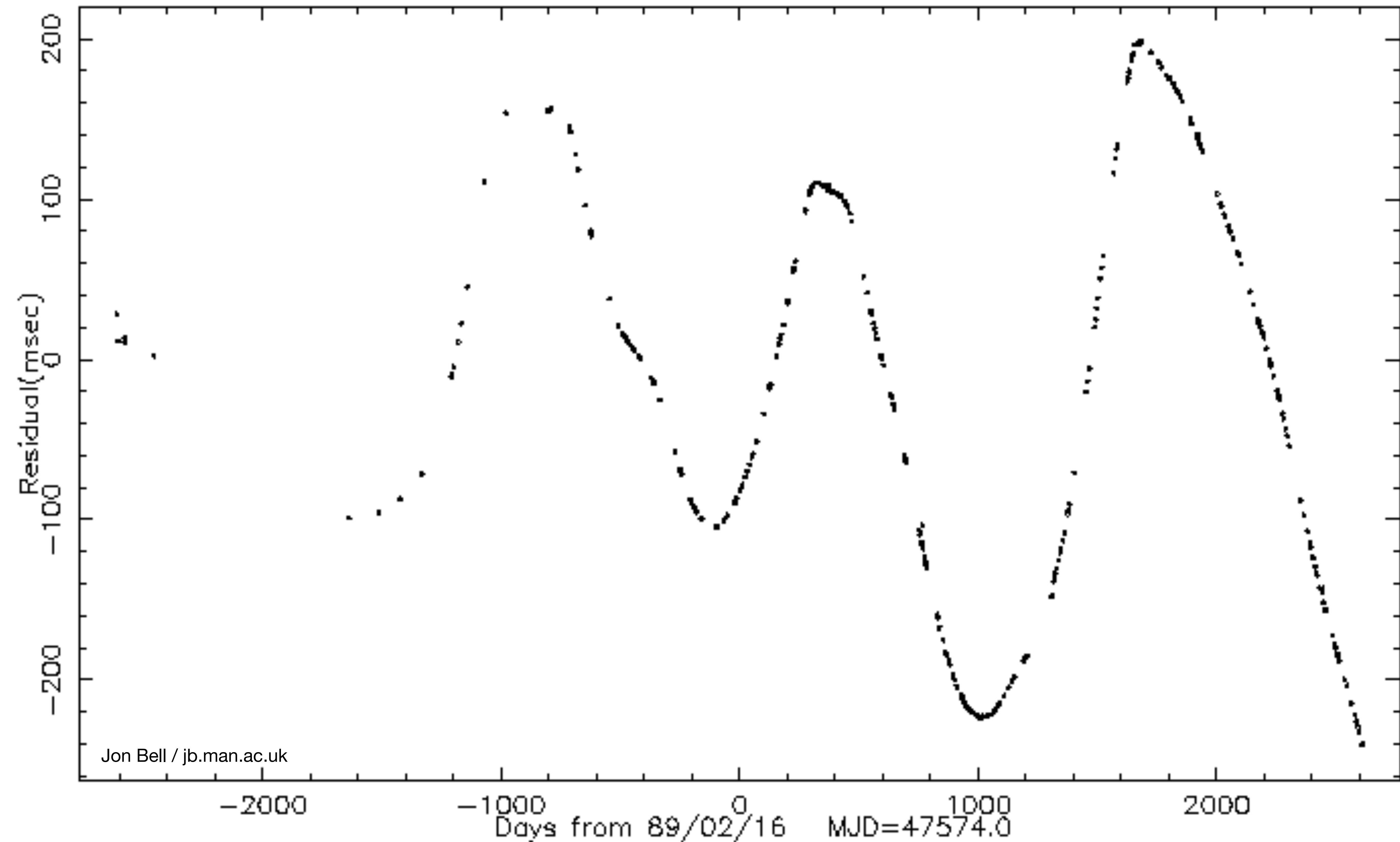
Dynamics of superfluid vortices show up in Pulsar  
timing: Glitches and Glitch recovery

# Pulsar Spin Evolution

Evidence of interior superfluid

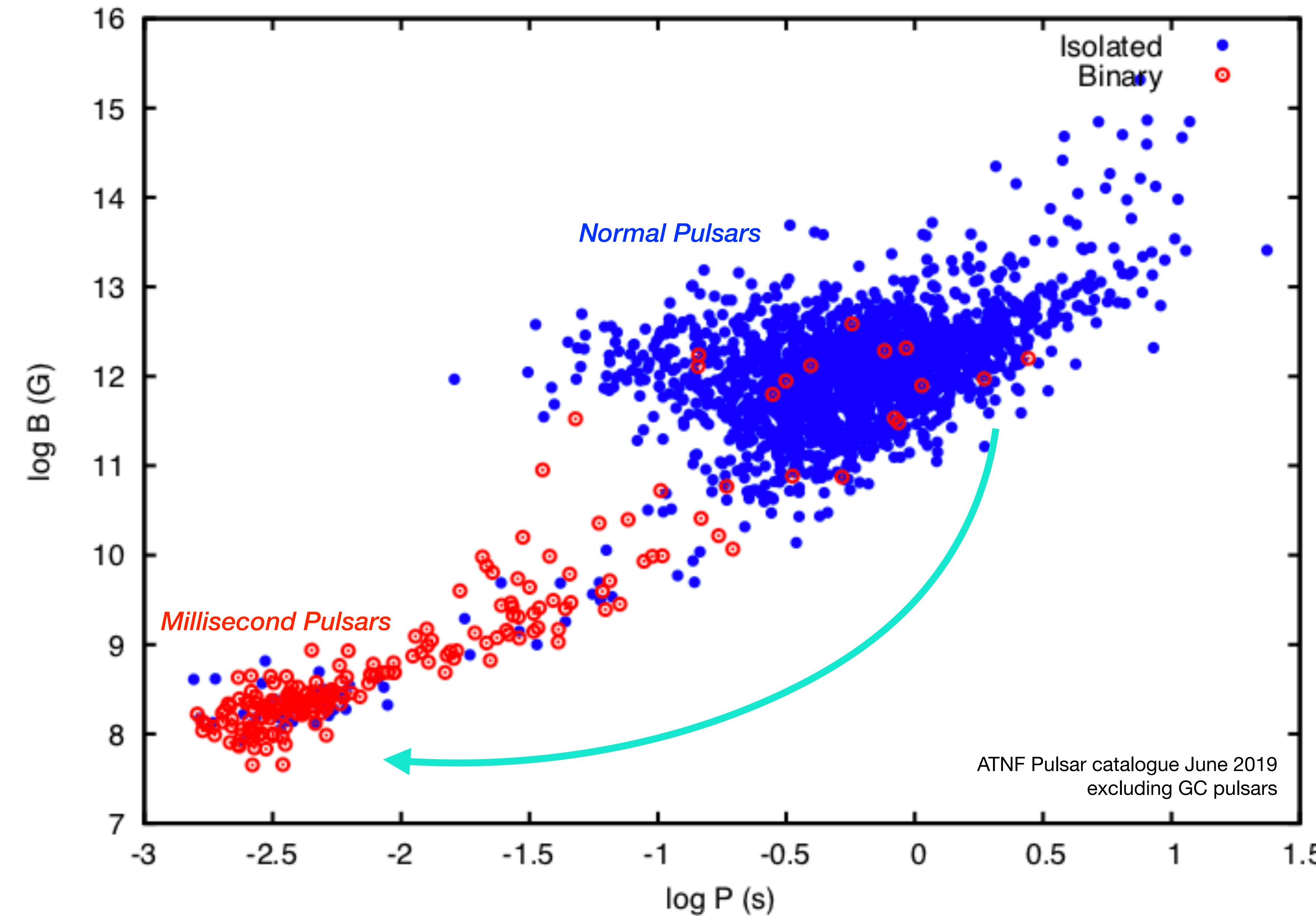


Residuals for 0611+22



Apart from glitches,  
random wanderings  
of pulse arrival time  
indicate jitter in  
spin period

This is called  
**Timing Noise**



Spin period  $P$  and spin-down rate  $\dot{P}$  measured by pulsar timing can be used to estimate the magnetic field strength assuming dipole torque

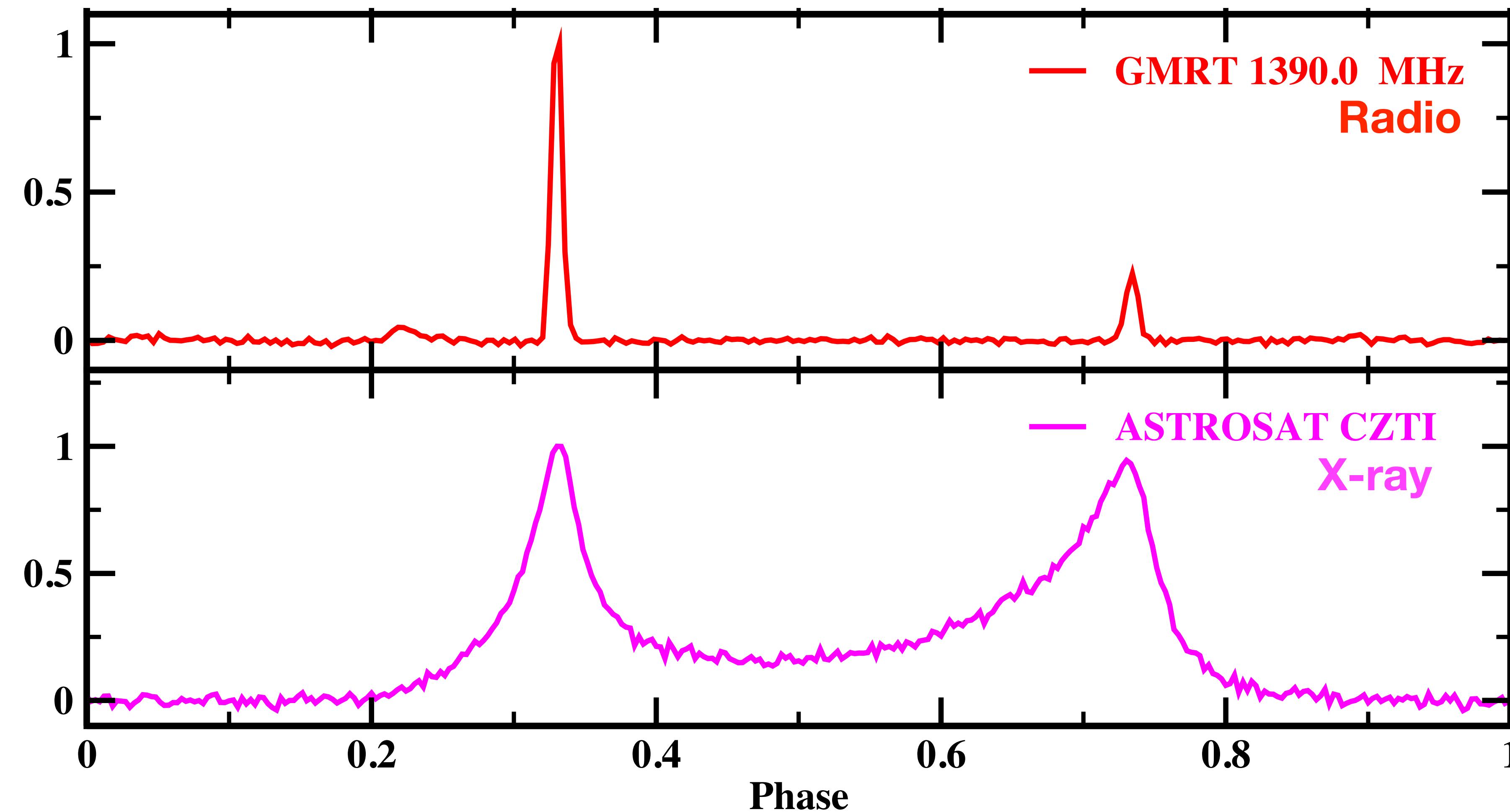
$$B_{12} \approx \sqrt{PP_{-15}}$$

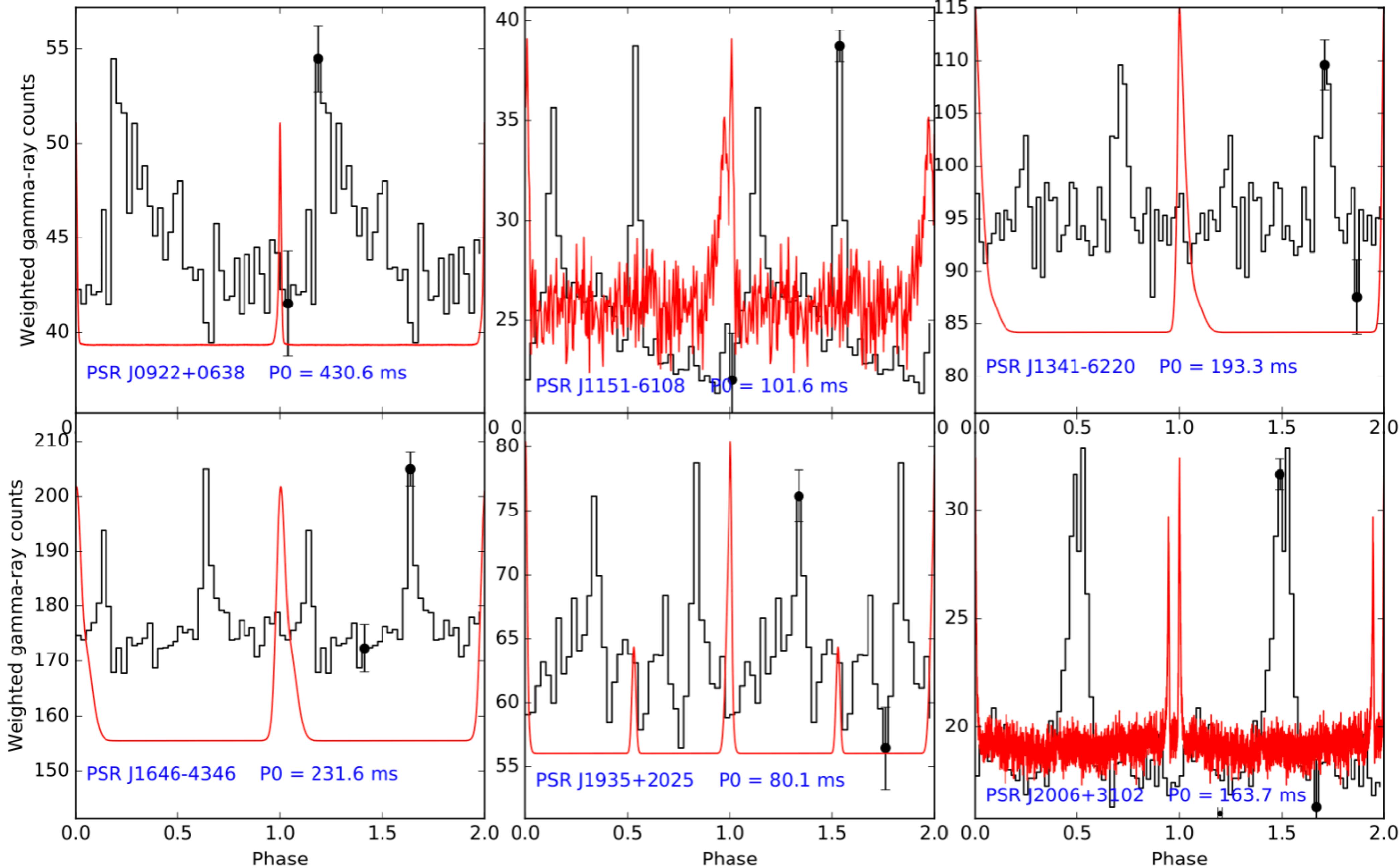
Magnetic field strength appears to reduce by processing in a binary system

Accretion in a binary also spins up the neutron star

# Pulsar Emission

The Crab Pulsar





Radio and High Energy emission typically originate in different regions

Coherent emission in Radio, Incoherent in High Energy

Most of the radiative luminosity emerges at High Energy. Radio carries a very small fraction

# A Radiation Primer

An accelerated charge emits radiation: Power  $P = \frac{2}{3c^3}q^2a_0^2$ ;  $a_0^2 = \gamma^4(|a_{\perp}|^2 + \gamma^2a_{\parallel}^2)$ ;  $\gamma = (1 - v^2/c^2)^{-1}$

For  $N$  charges:

Incoherent emission:  $P \rightarrow NP$

Coherent emission:  $q \rightarrow Nq$

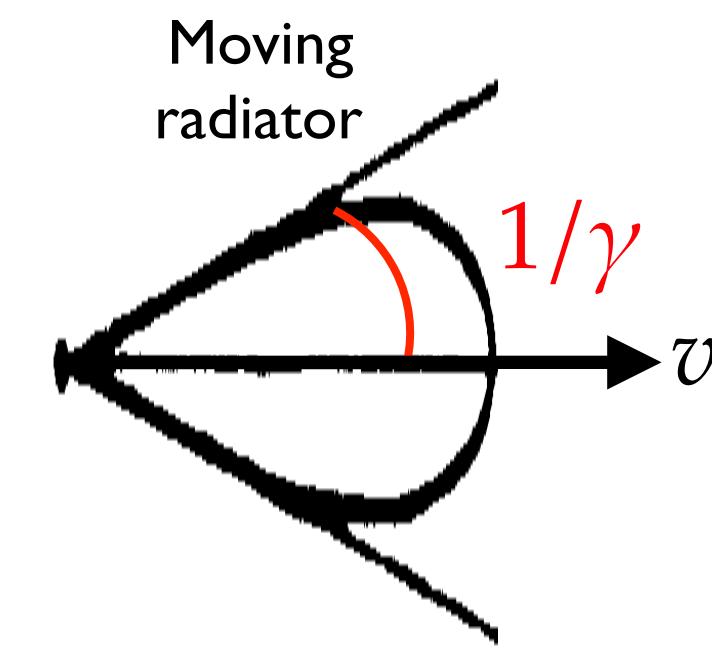
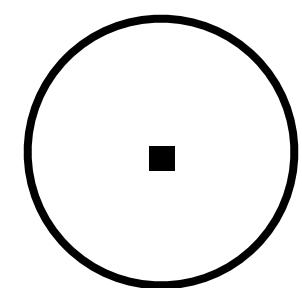
hence  $P \rightarrow N^2P$

requires bunching in a scale  $\ll \lambda$

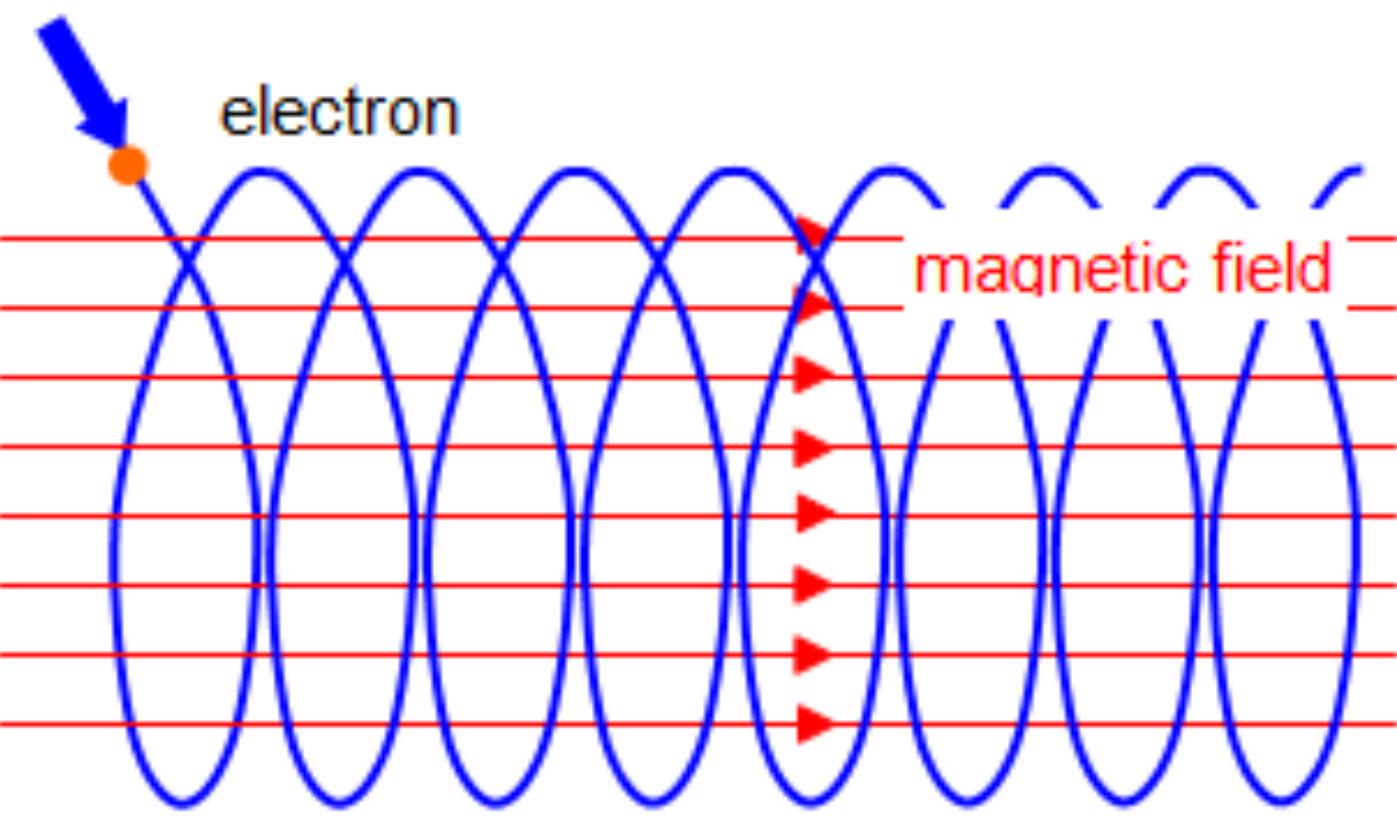
Polarisation  $\vec{E} \propto \hat{n} \times [(\hat{n} - \vec{\beta}) \times \dot{\vec{\beta}}]$

Beaming

Stationary radiator



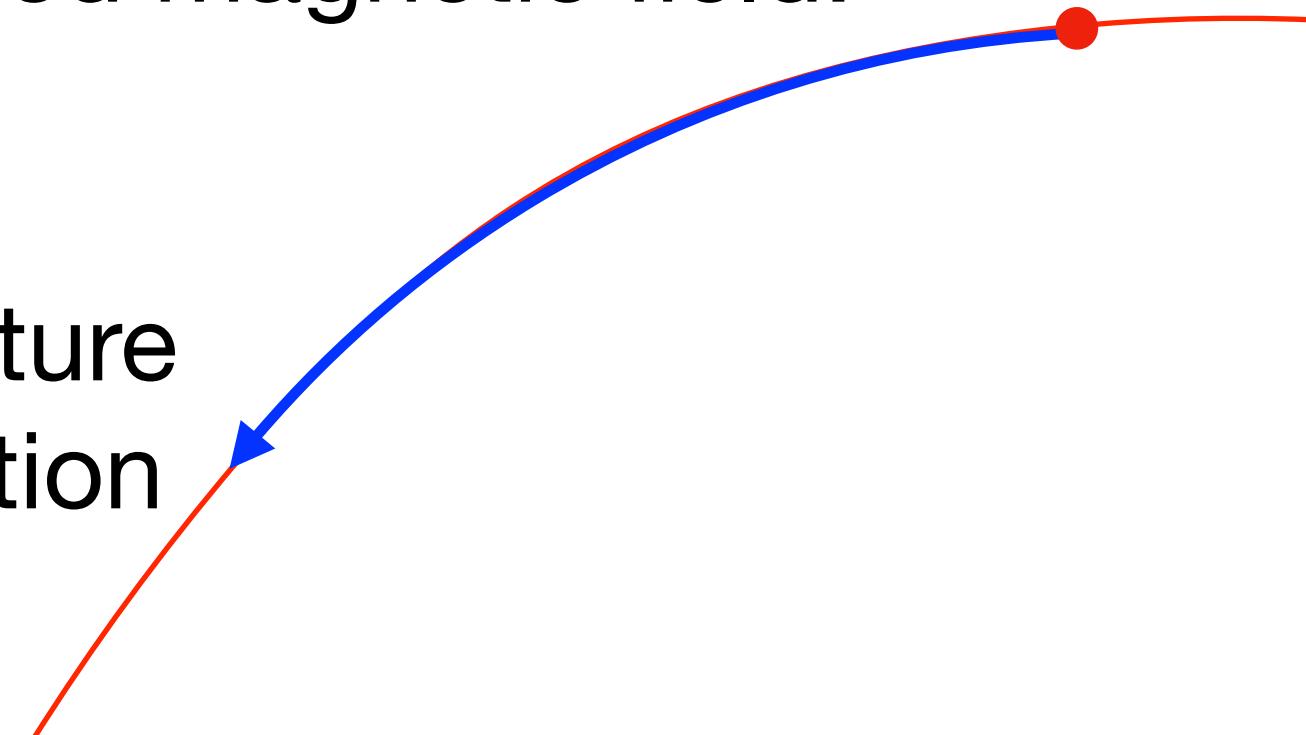
Relativistic electron gyrating in a magnetic field:



Synchrotron  
Radiation

Relativistic electron following  
a curved magnetic field:

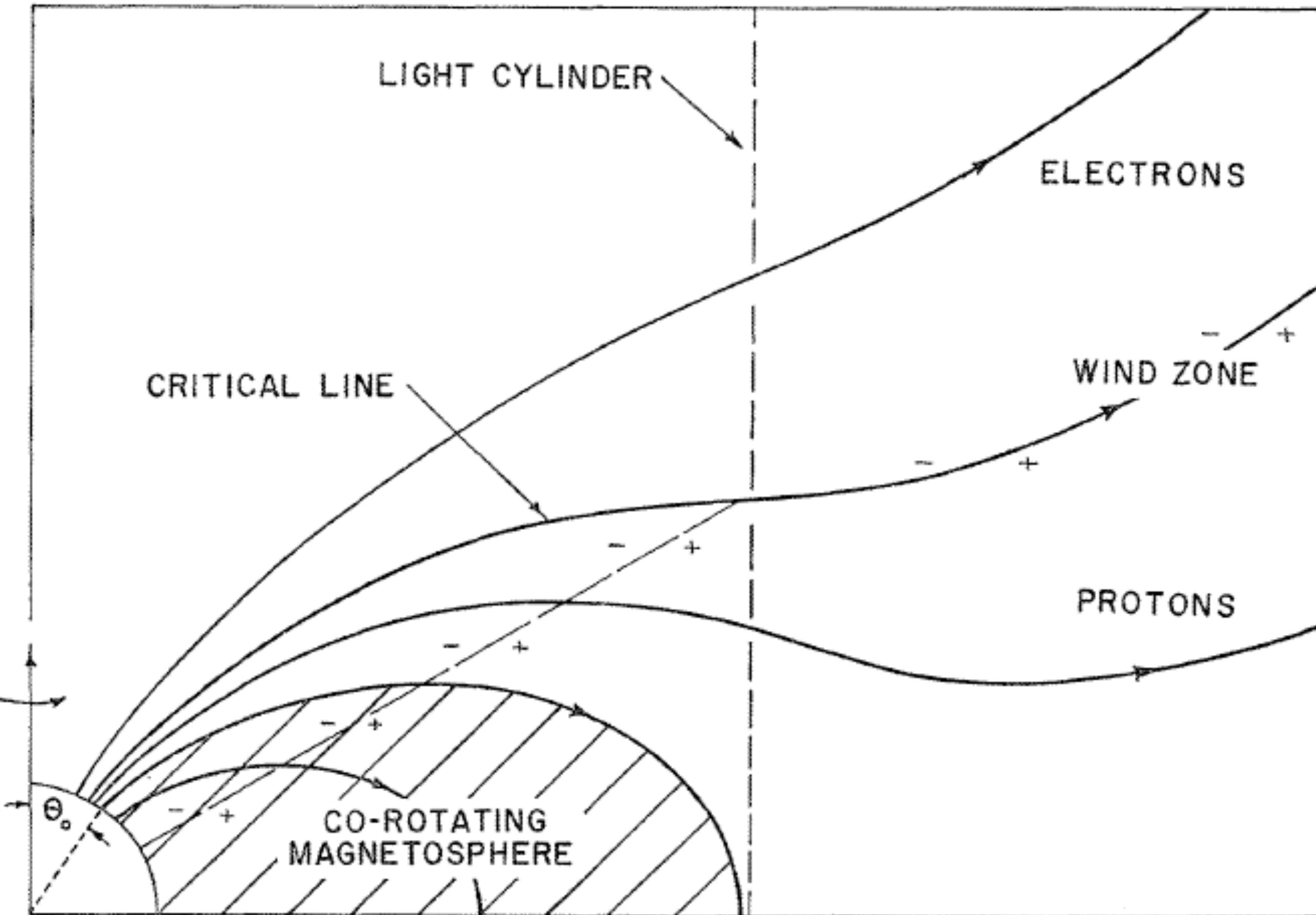
Curvature  
Radiation



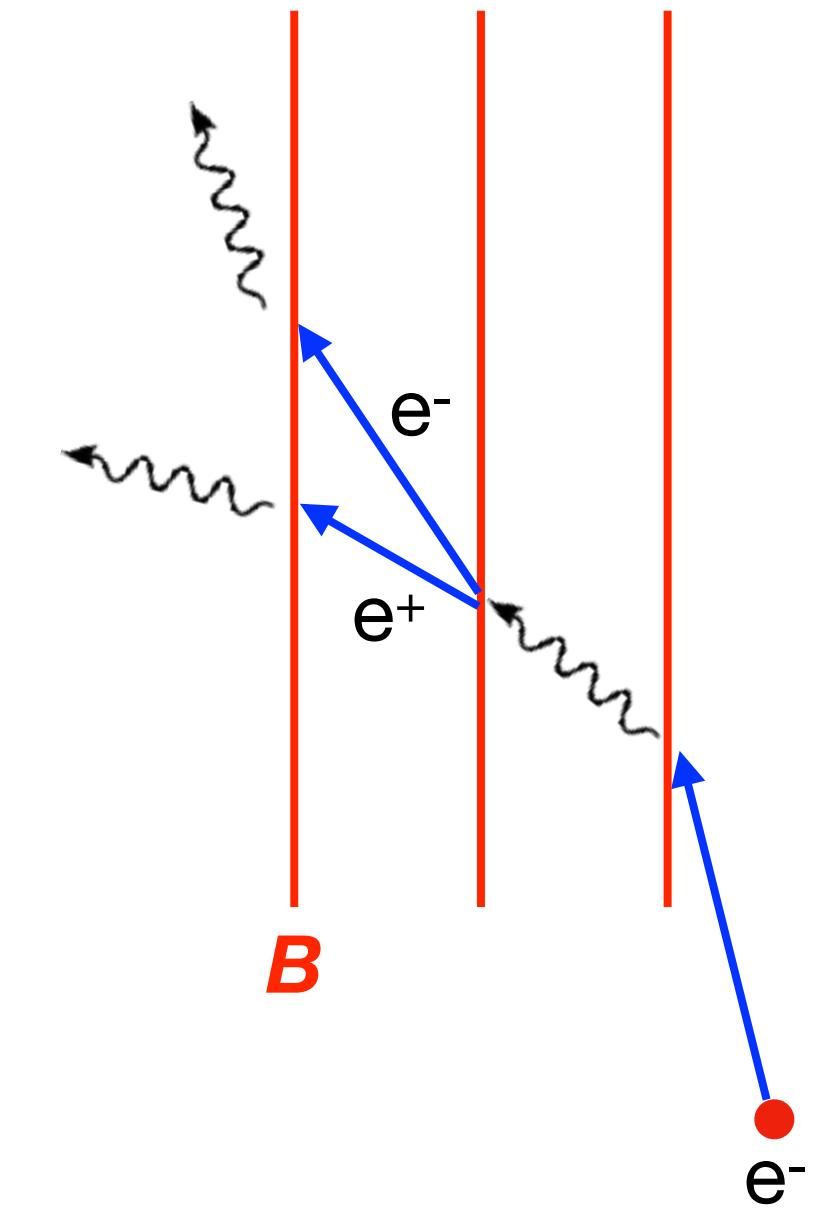
# The Pulsar Magnetosphere

Basic concepts: Goldreich & Julian 1969

Pair Cascade

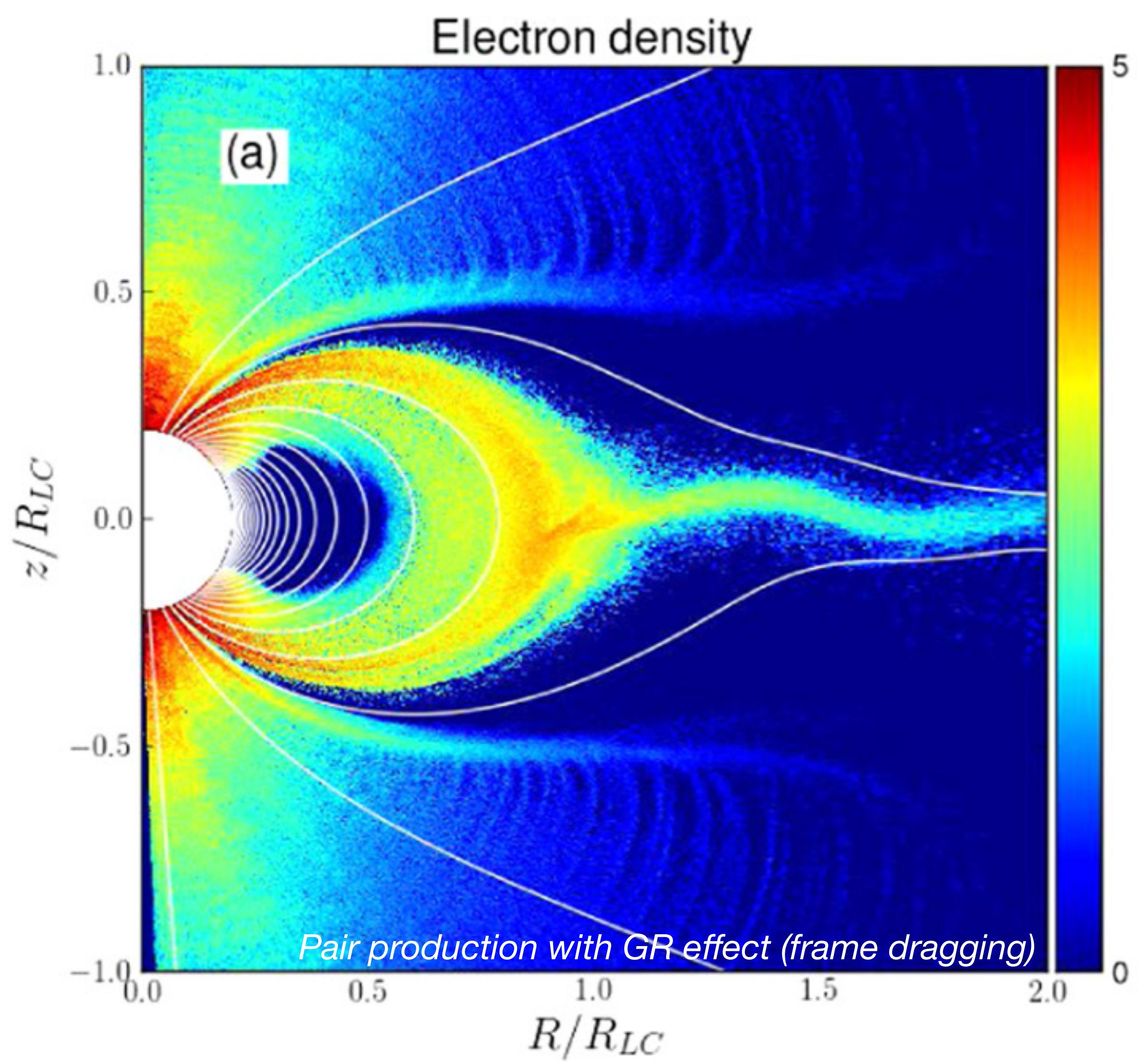
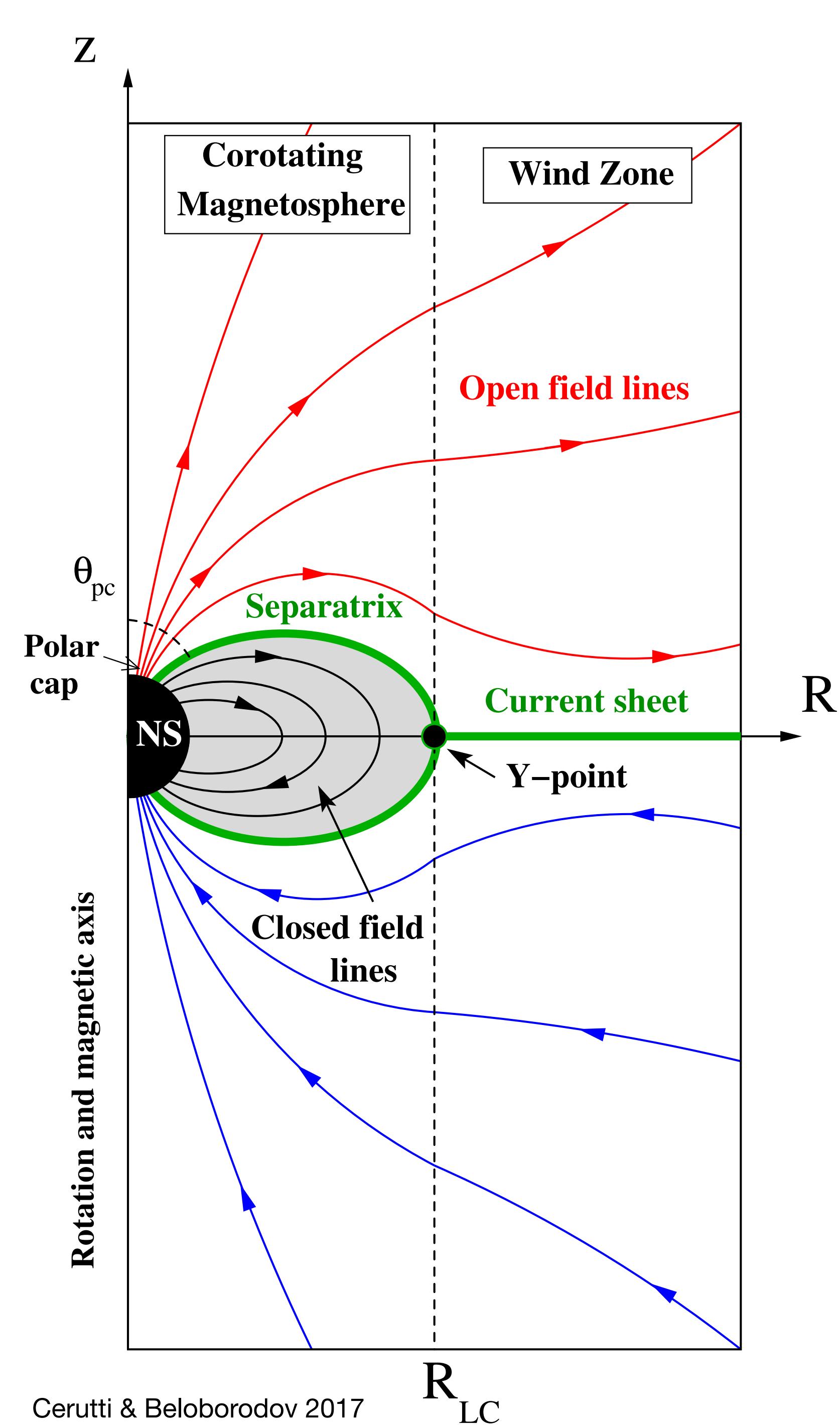


- A vacuum exterior would have potential drop exceeding  $10^{15} \text{ V}$
- Space charge must exist,  $\mathbf{E} \cdot \mathbf{B} = 0$
- Co-rotating magnetosphere can be maintained up to the light cylinder,  $\rho \approx -\Omega \cdot \mathbf{B} / 2\pi c$
- With pair production, no. density of charged particles may far exceed  $\rho$
- $\rho$  passes through 0 and changes sign in the magnetosphere

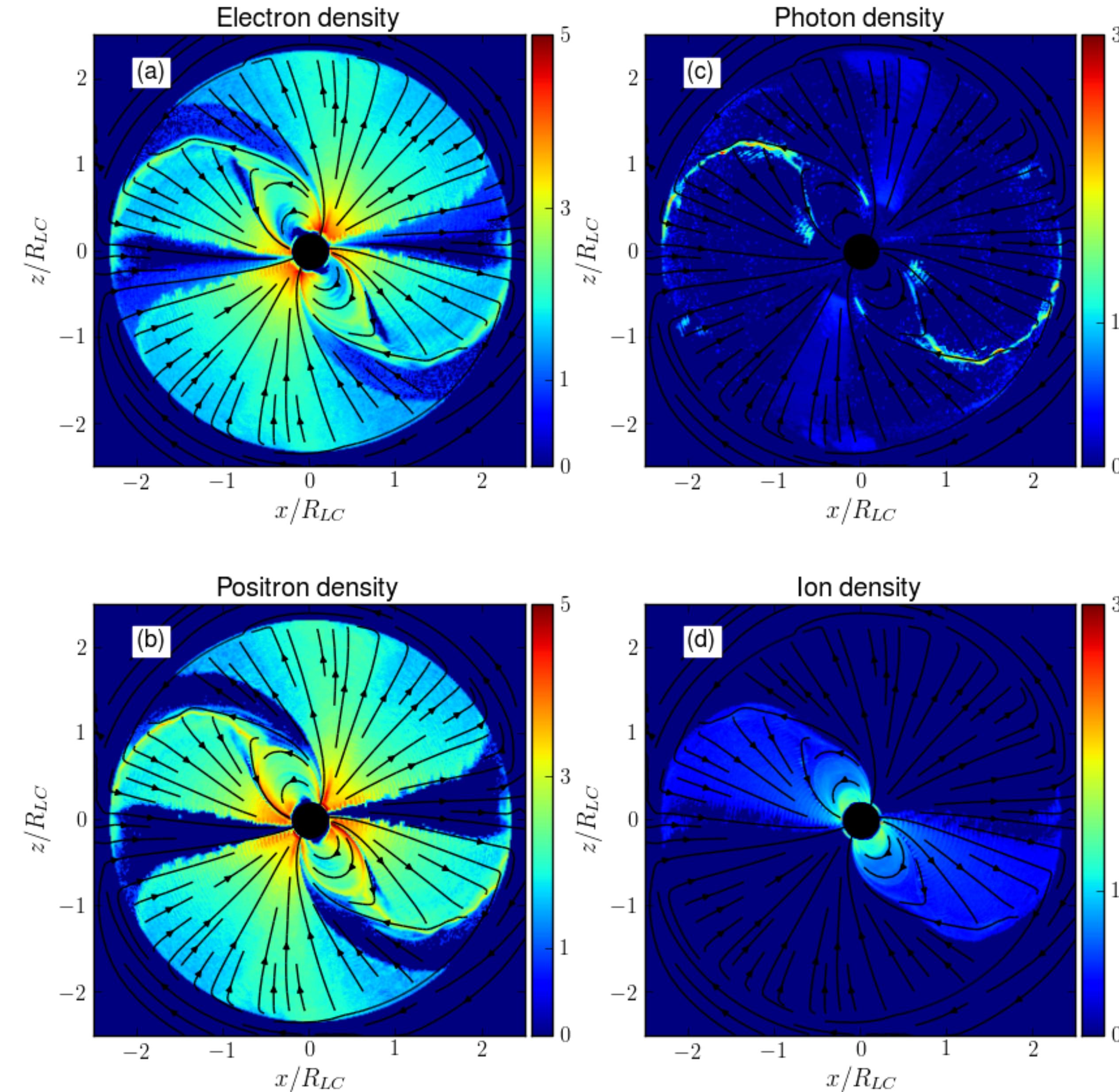


Charge bunching likely due to beam-plasma instability (primary particles moving through the secondary plasma)

Plasma on open field lines cannot co-rotate. Current flows out along these lines, creates toroidal mag. field which provides the dipole spin-down torque.



# GR-induced pair production



Inclination  
60°

# Coherent Emission at Radio wavelengths

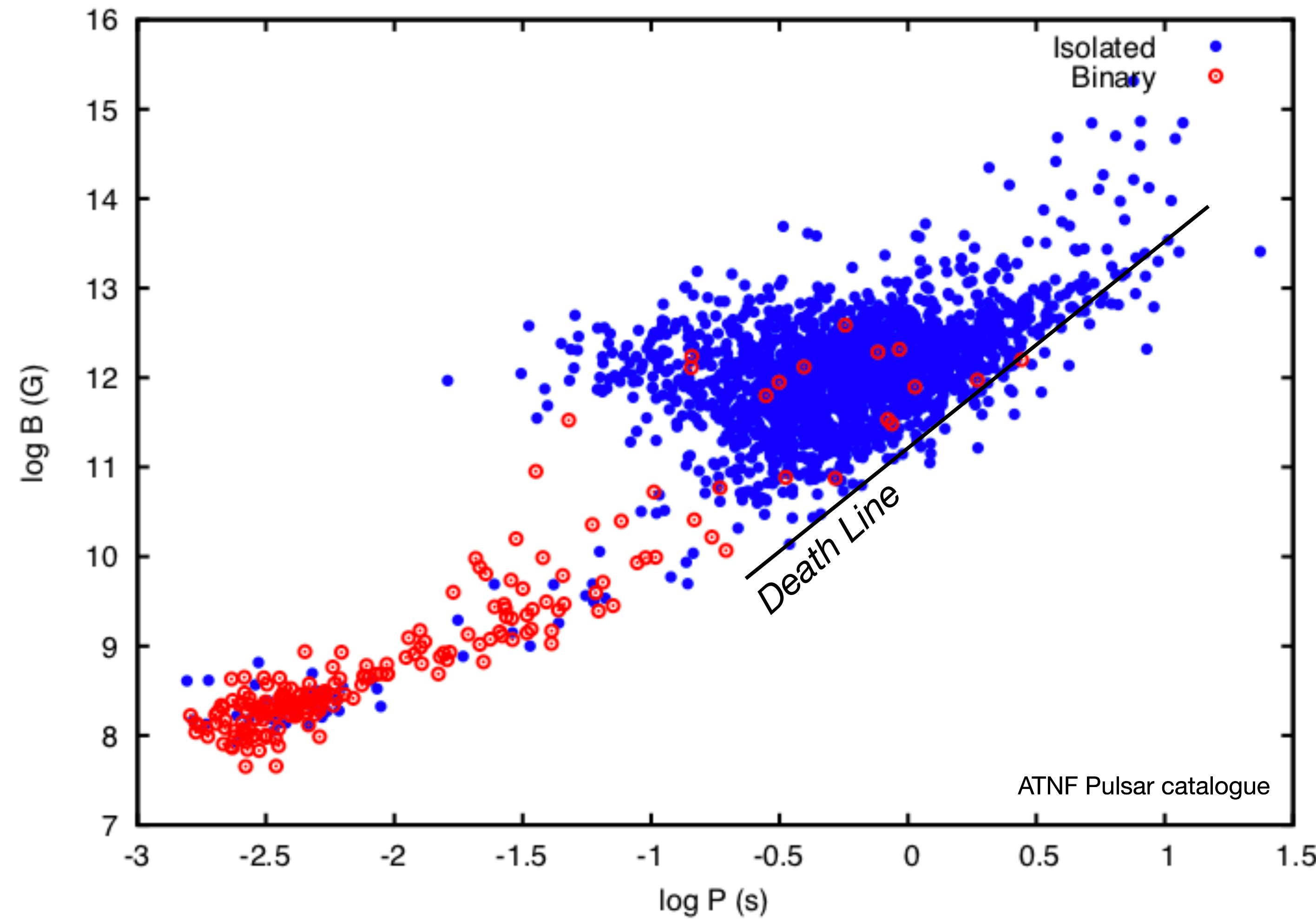
## Ruderman and Sutherland (1975) model

- Charge bunching at sub-wavelength scale
- Active pair production at Polar Cap
- Bunch dynamics in local E, B fields

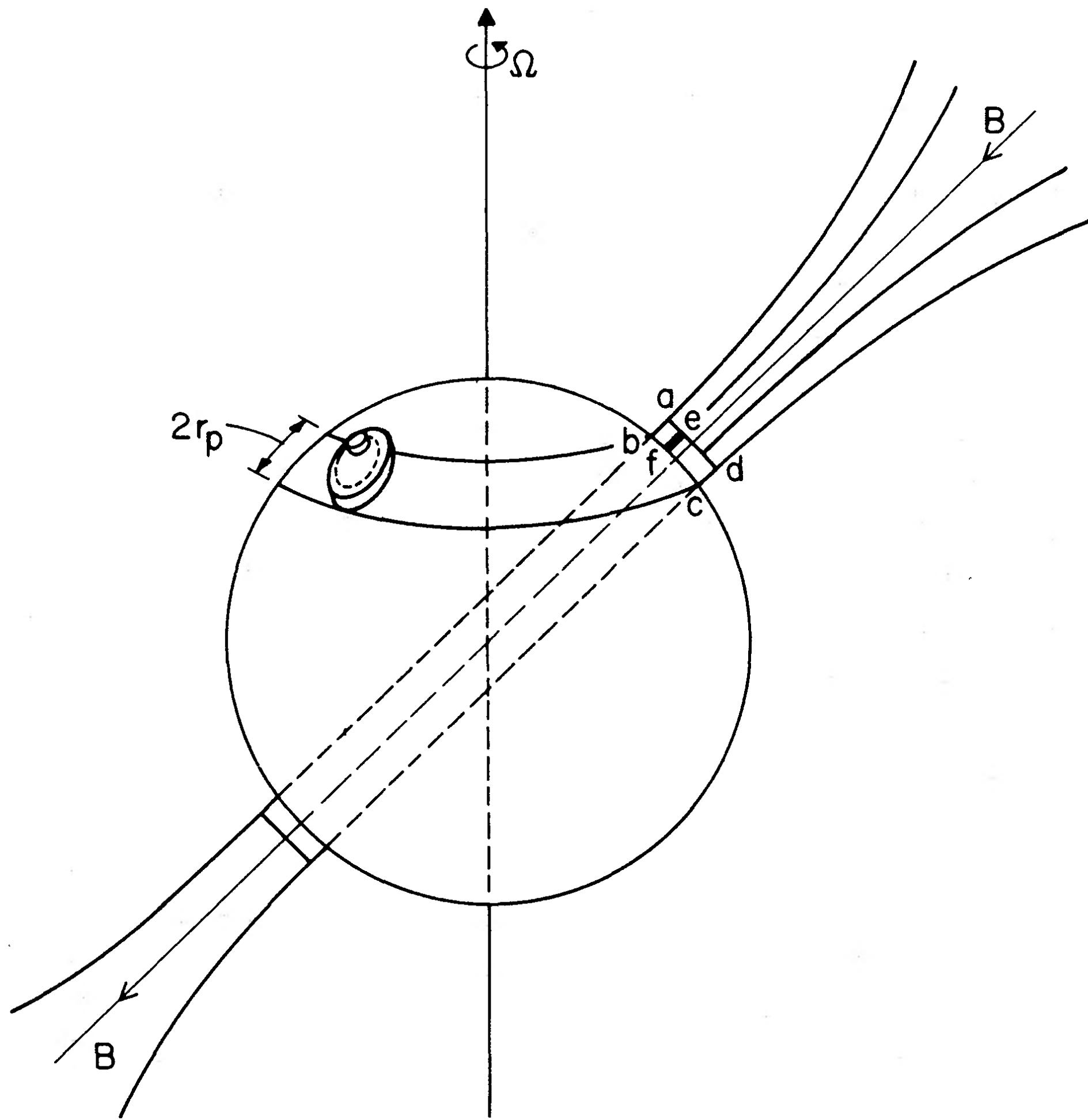
pair sustenance:

$$P_{\text{crit}} = 1.7 B_{12}^{8/13} R_6^{21/13} \rho_6^{-4/13} (15\chi)^{-2/3} \text{ s}$$

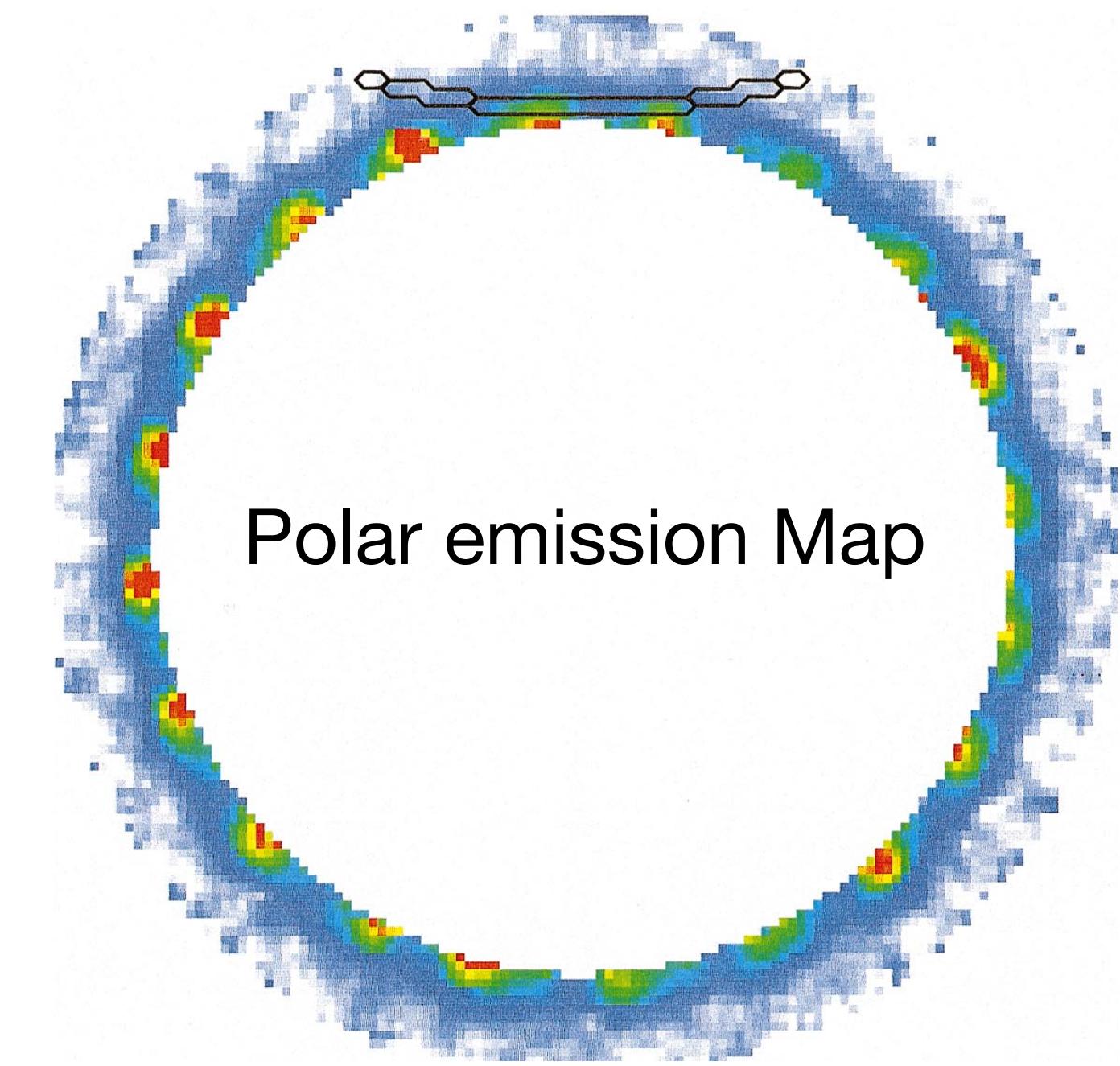
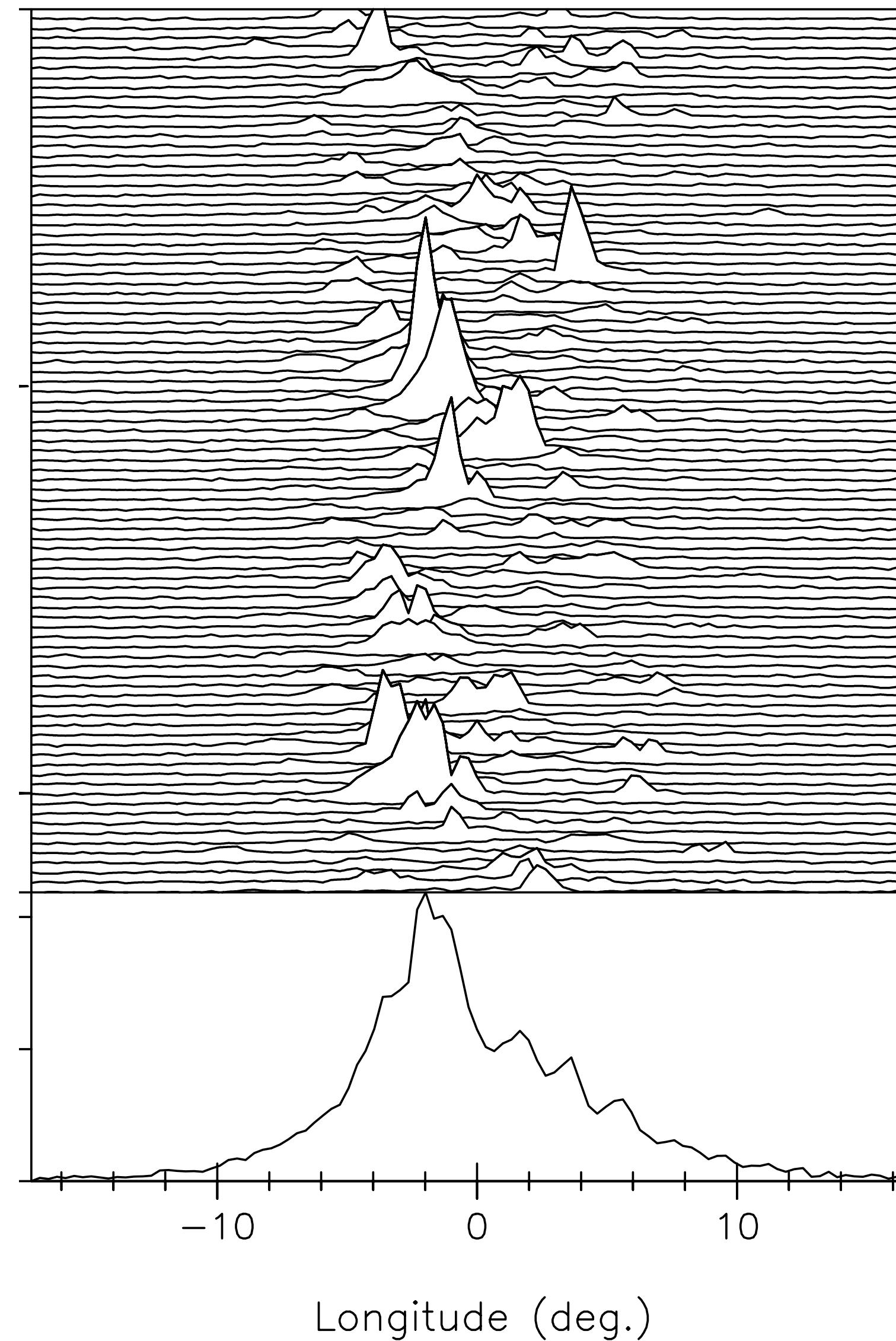
$$\chi \equiv \frac{\hbar\omega}{2mc^2} \frac{B_\perp e\hbar}{m^2 c^3} \gtrsim \frac{1}{15}$$



Charge bunches should rotate due to  
 $E \times B$  drift

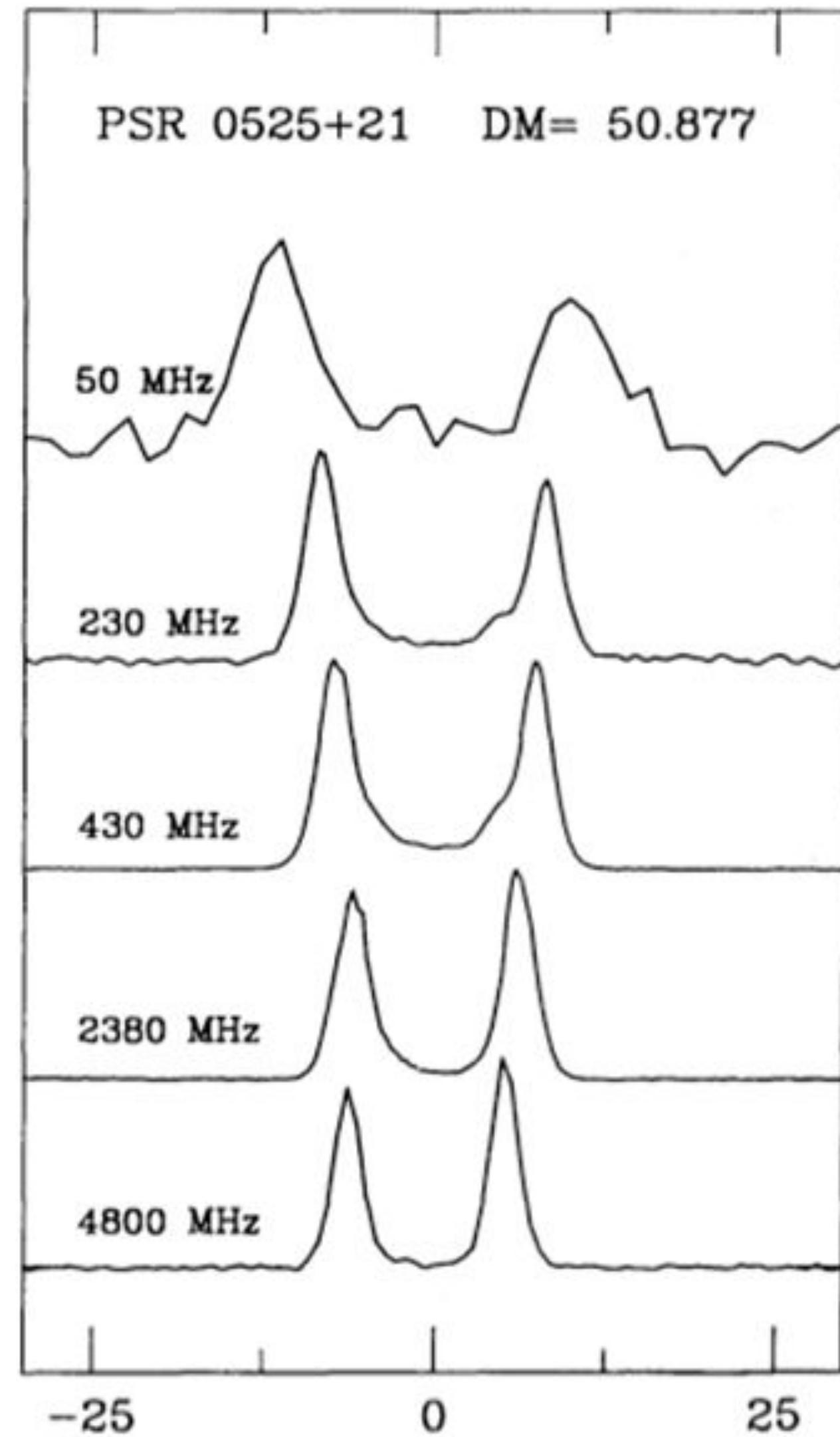


## PSR 0943+10 drifting surpluses

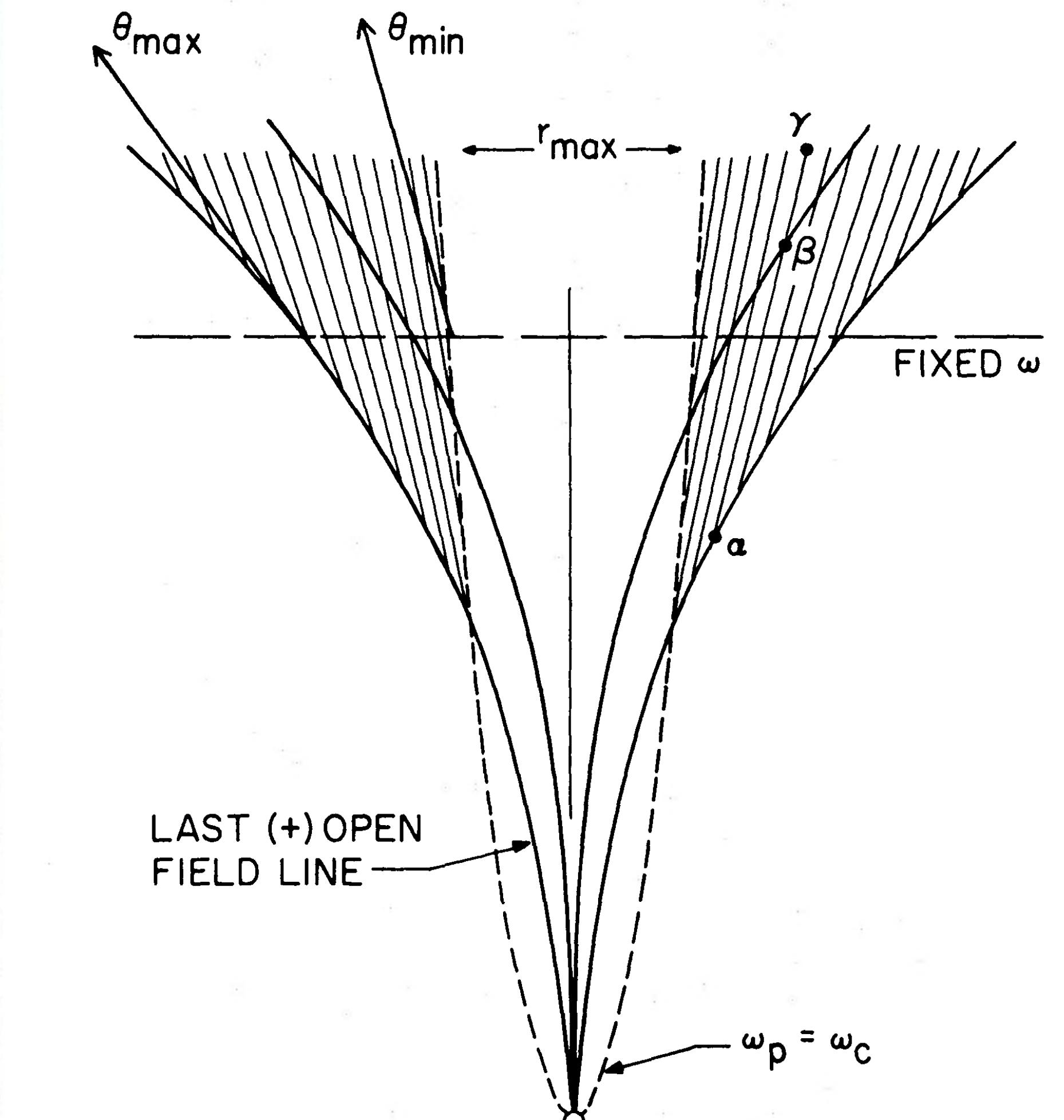
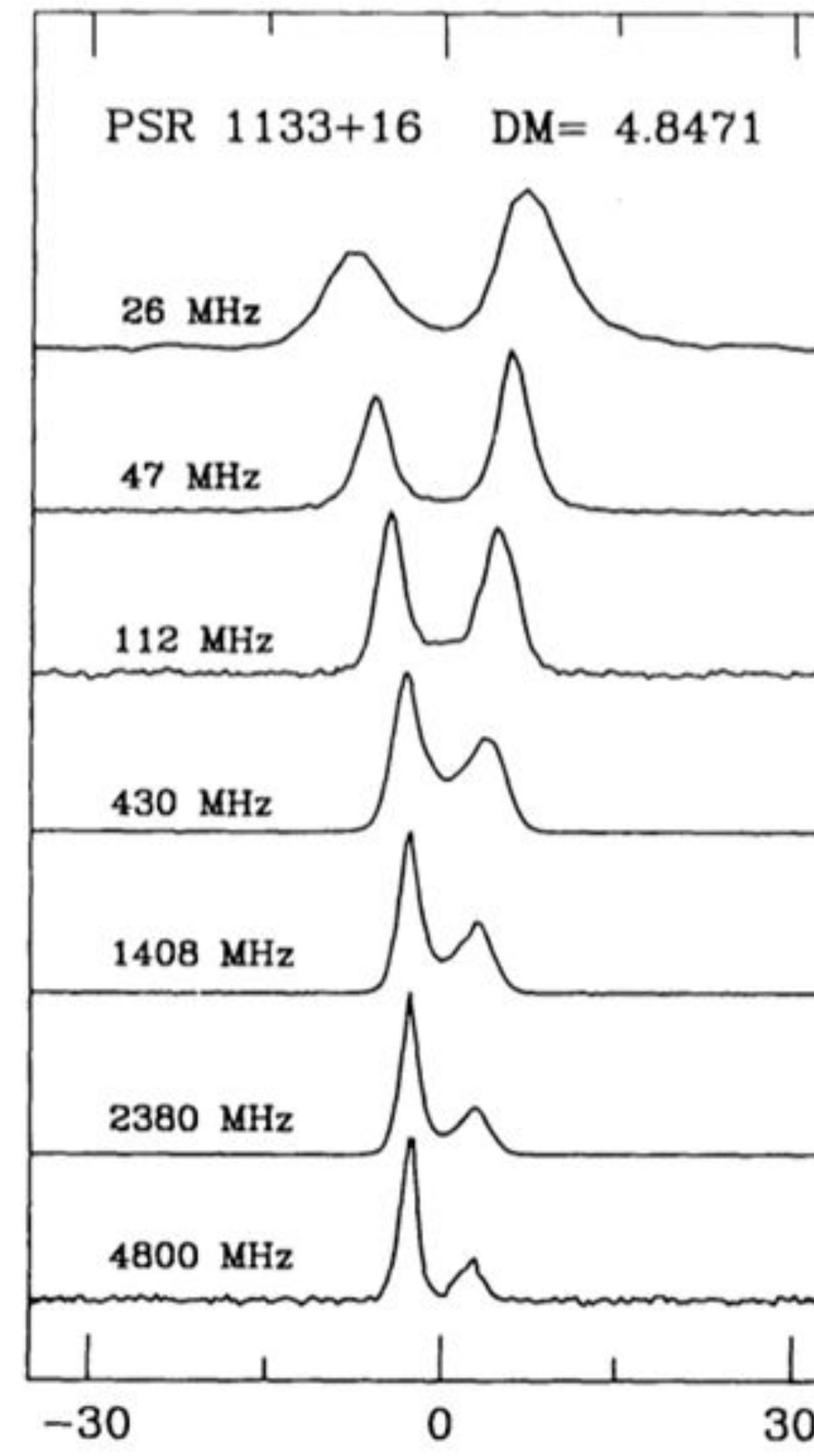


Deshpande & Rankin 1999

# Radio pulse shape: radius-to-frequency mapping



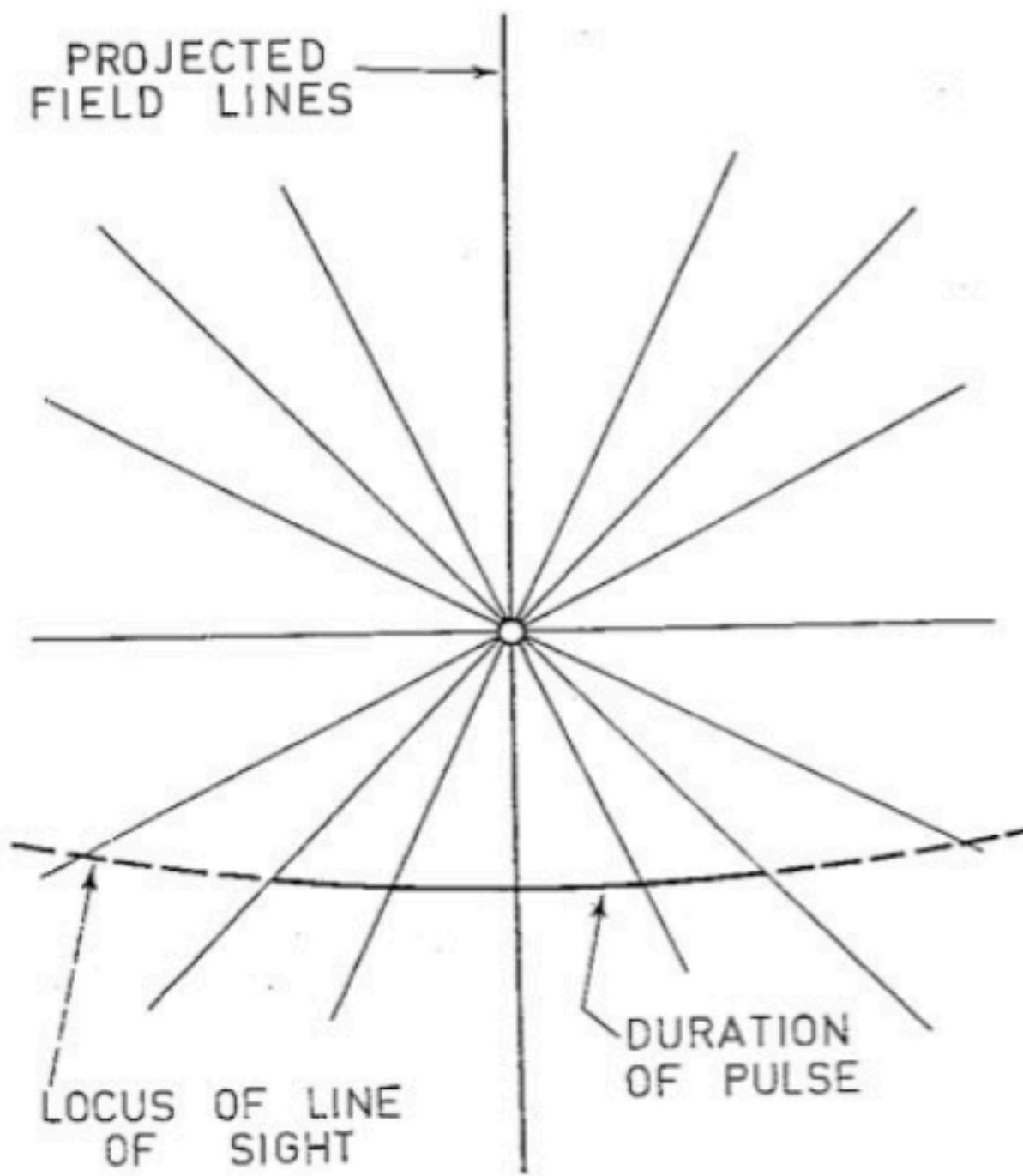
R.N. Manchester



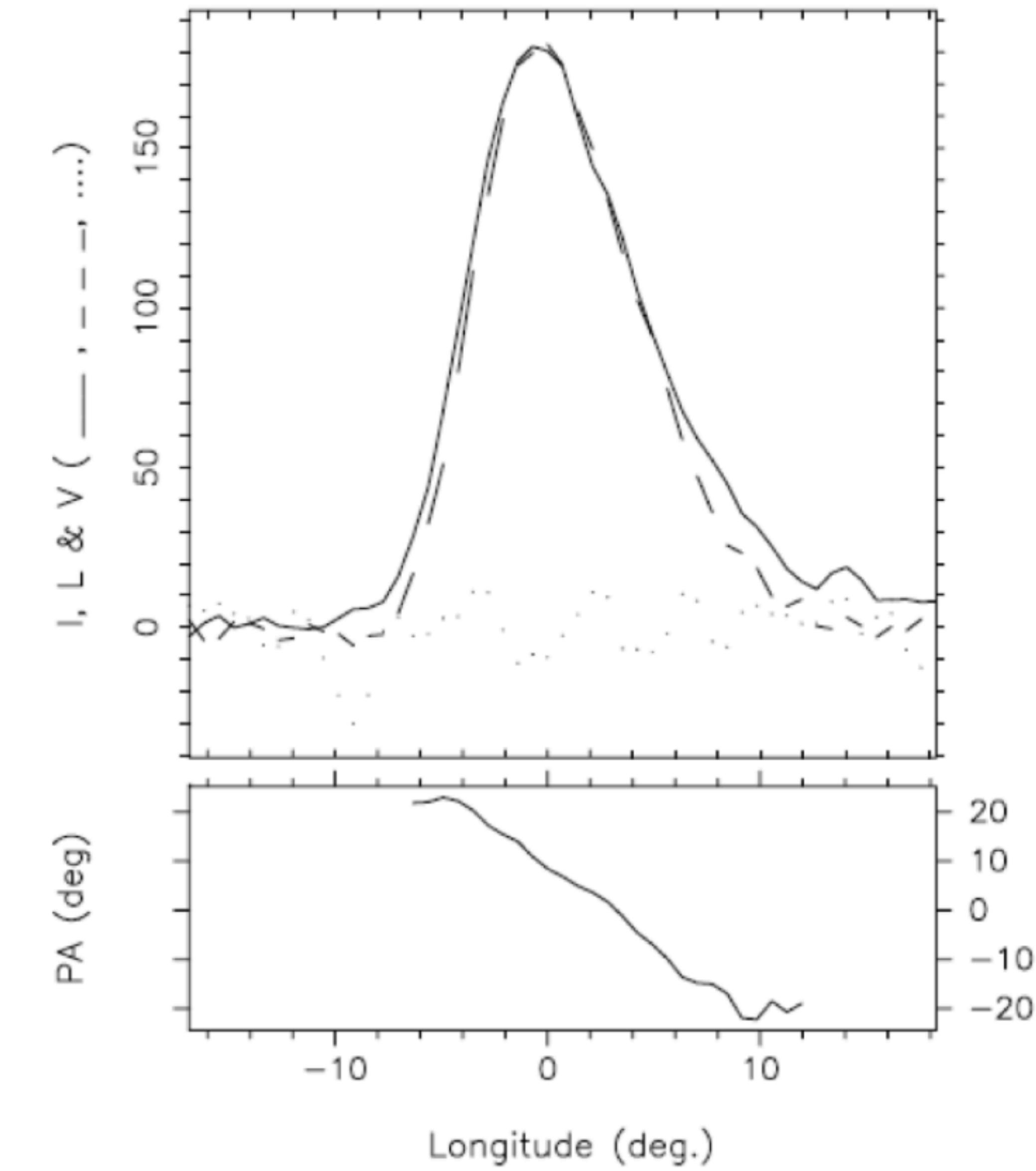
Ruderman & Sutherland 1975

# Polarization of pulsar radio emission

Strong linear polarization  
with S-pattern position-  
angle sweep:  
*curvature radiation,  
rotating vector model*



Radhakrishnan & Cooke 1969



Ramkumar & Deshpande 2001

# Propagation effects on pulsar signal

**Dispersion**

$$\frac{v_{\text{ph}}}{c} = \left(1 - \frac{\omega_p^2}{\omega^2}\right)^{-1/2}$$

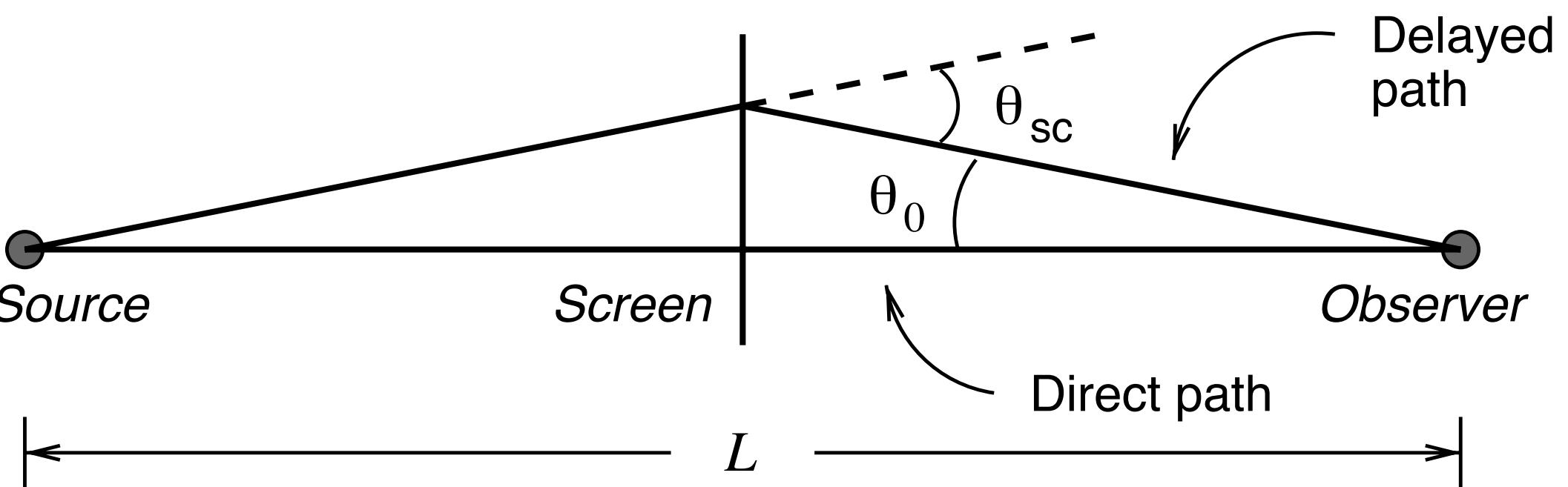
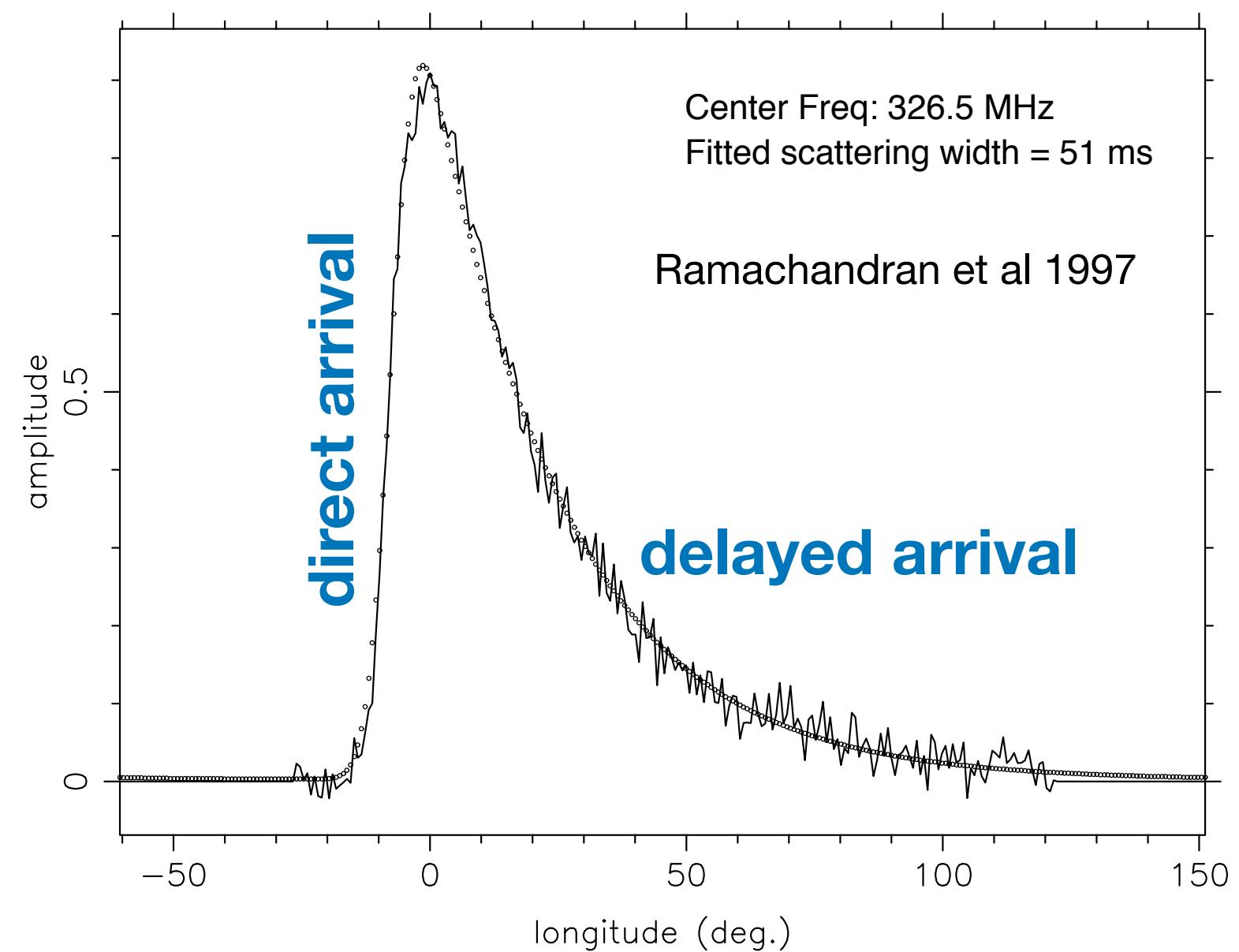
frequency dependent propagation time

$$DM = \int_0^L n_e dl$$

**Dispersion Measure**

**Scattering**

PSR J1833-0338 P=0.686676816 s DM=235.8 ORT 13/01/96



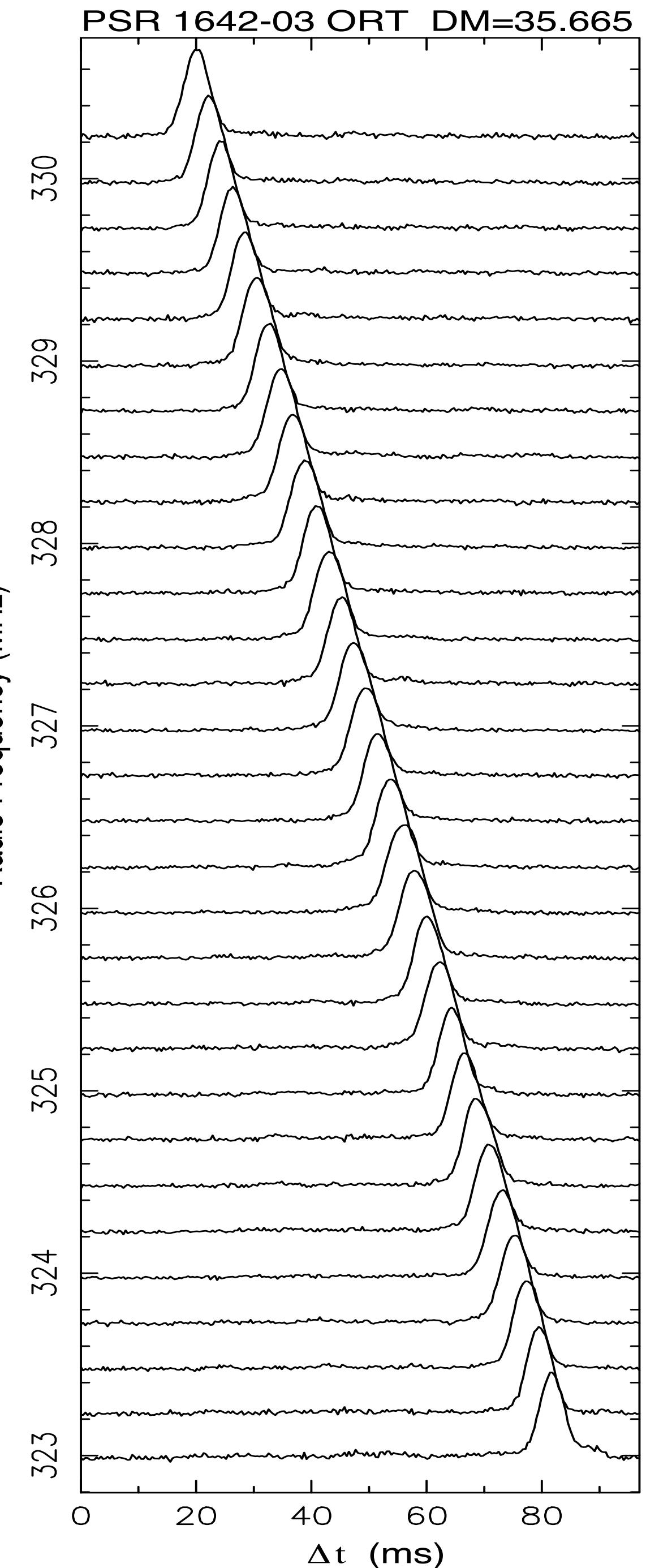
$$\tau_{\text{scat}} \propto (DM/\nu)^{4.4} \quad \text{approx.}$$

**Faraday Rotation**

$$\left(\frac{v_{\text{ph}}}{c}\right)_{R,L} = \left(1 - \frac{\omega_p^2}{\omega(\omega \pm \omega_B)}\right)^{-1/2}$$

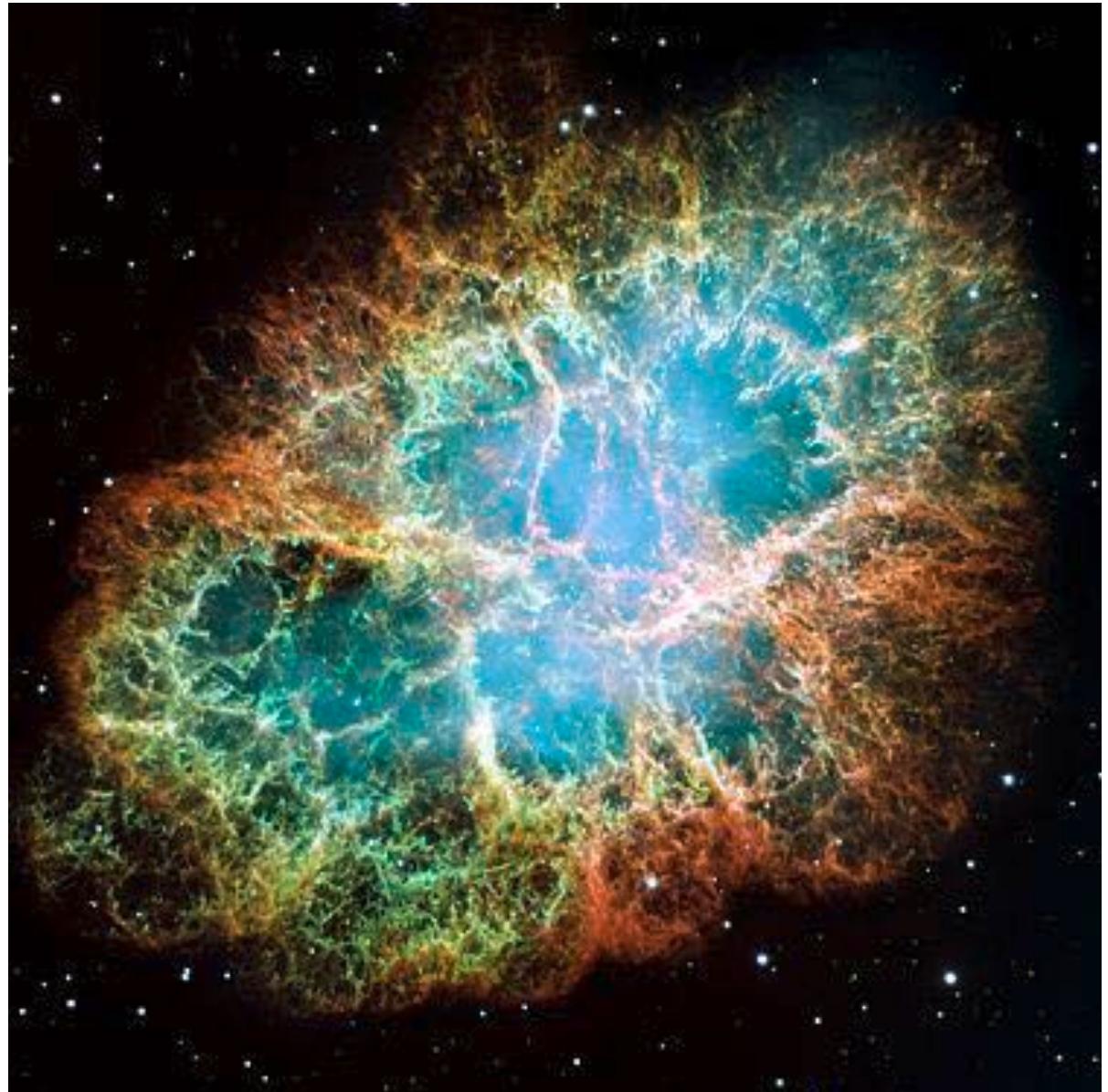
$$RM = \int_0^L n_e B_{\parallel} dl$$

**Rotation Measure**



# Matter around pulsars

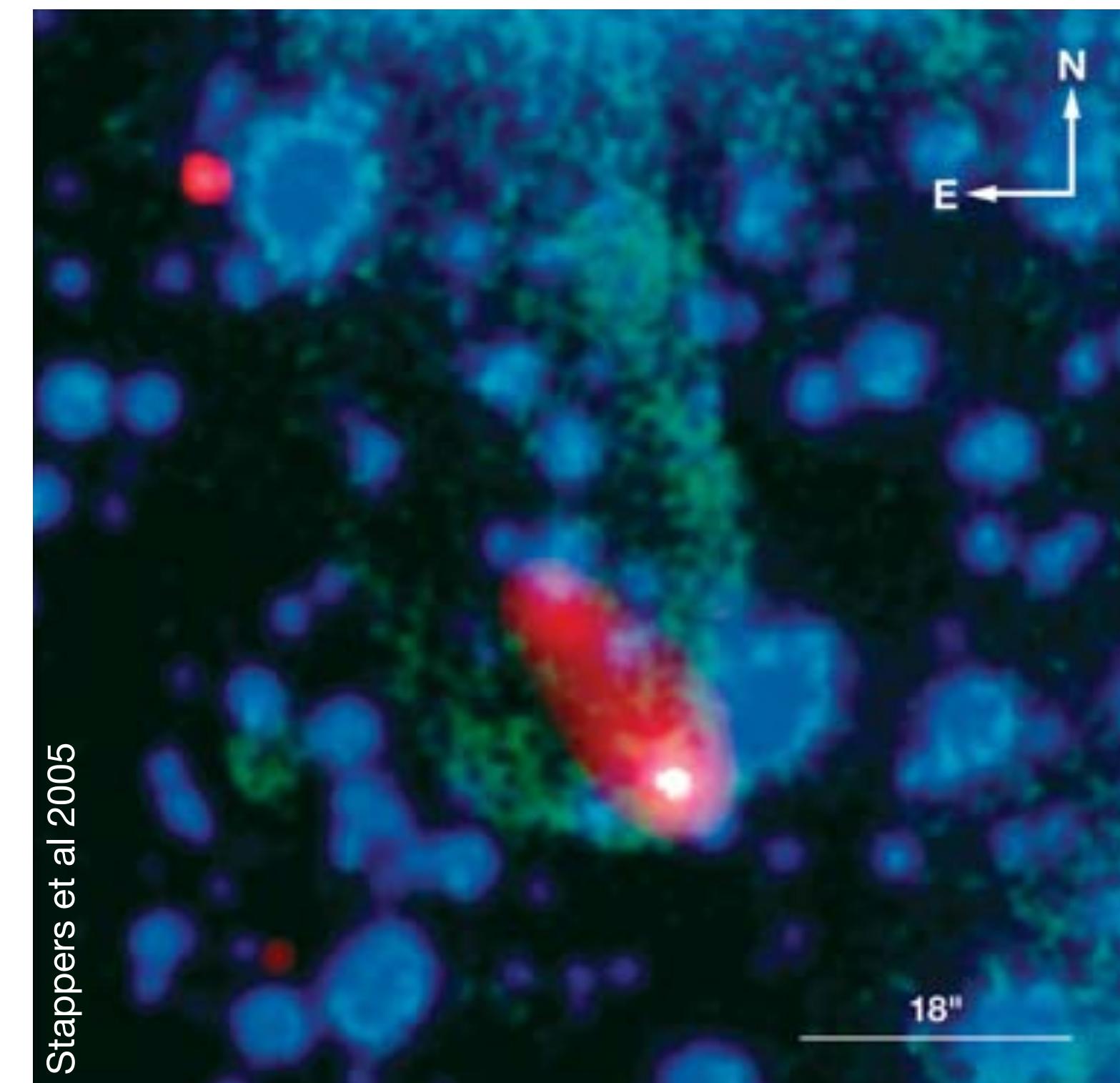
- Neutron Stars are born in supernova explosions
- Young pulsars are often found in supernova remnants
- Relativistic pair wind from the pulsar creates a nebula
- Fast moving pulsar through interstellar medium can also display a Pulsar Wind Nebula with bow shock
- Similar structures are seen in Black Widow pulsars



Crab Nebula HST



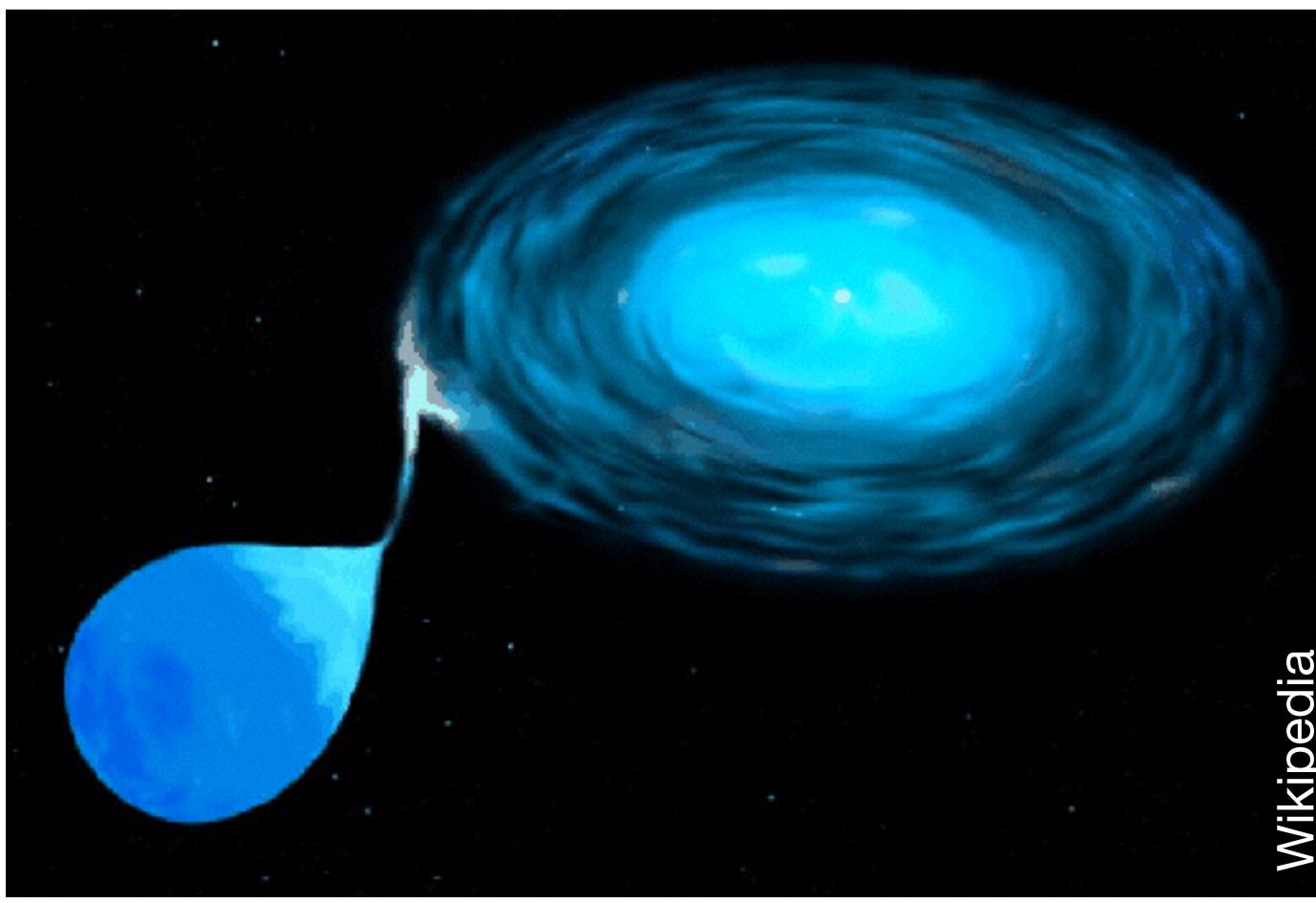
Vela PWN CXO



Stappers et al 2005

PSR B1957+20

# Evolution of a neutron star in a binary



Matter approaching from a companion has angular momentum, may affect the spin of the Neutron Star

Disk-magnetosphere interaction at Alfvén radius

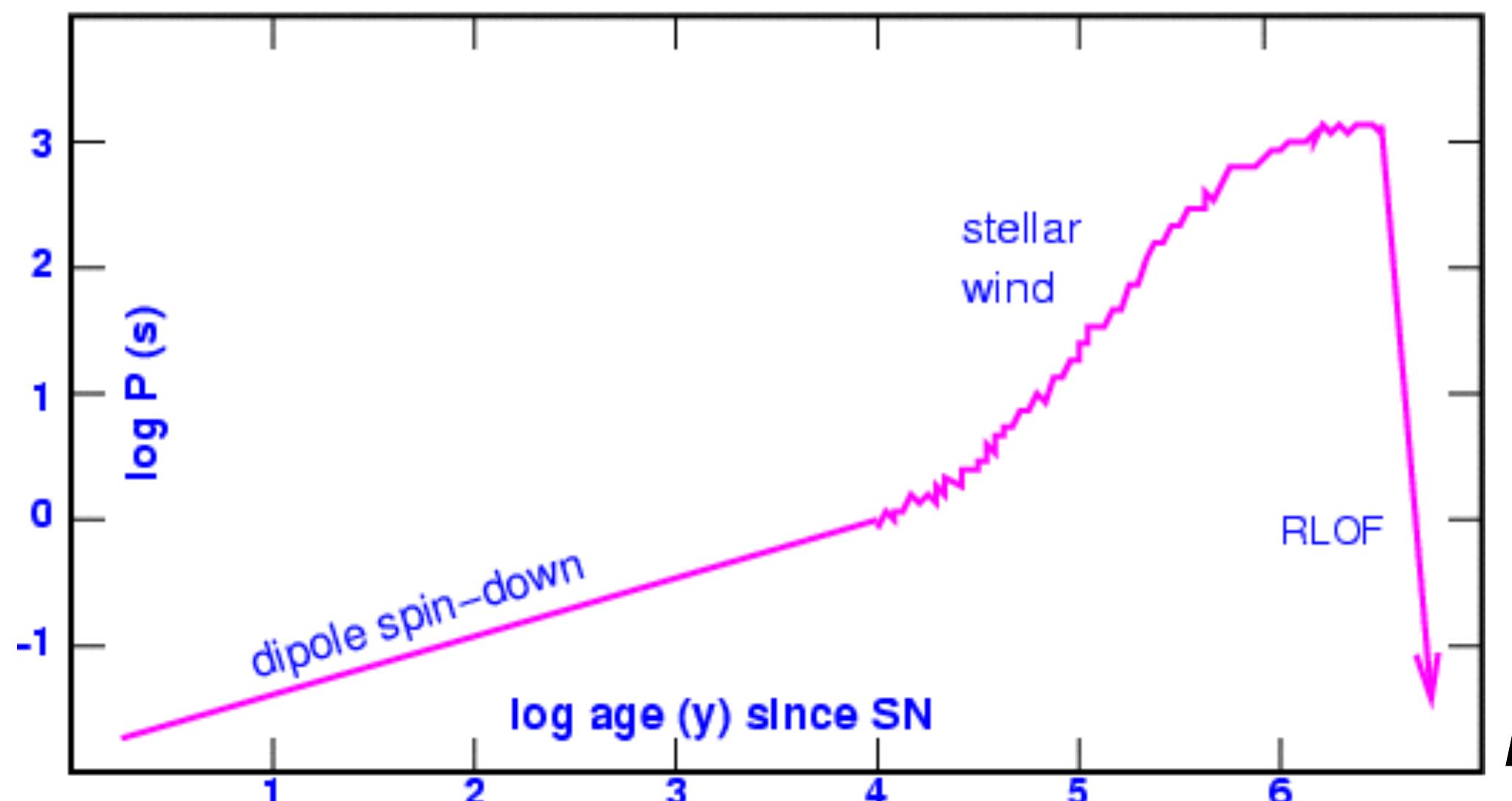
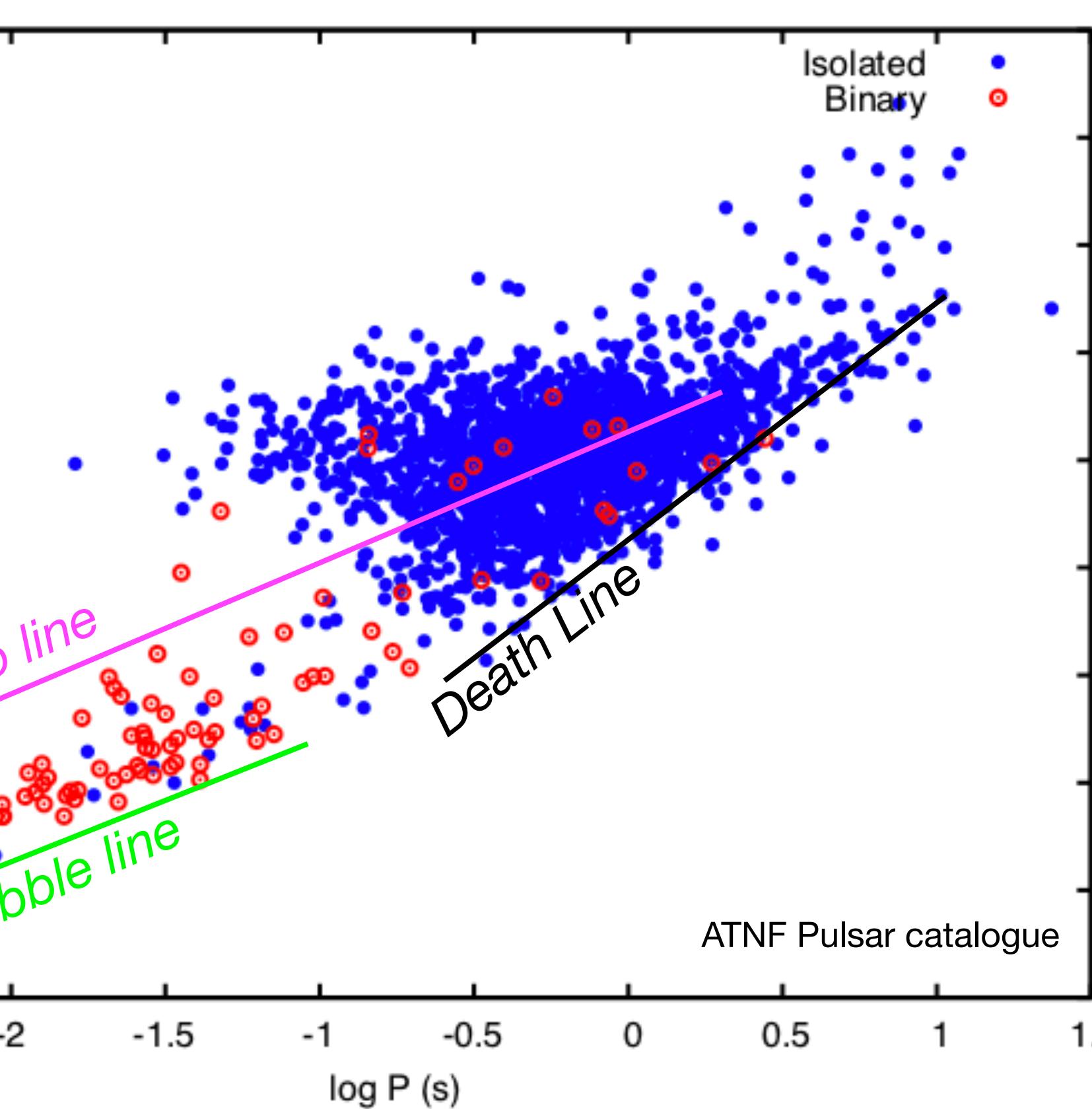
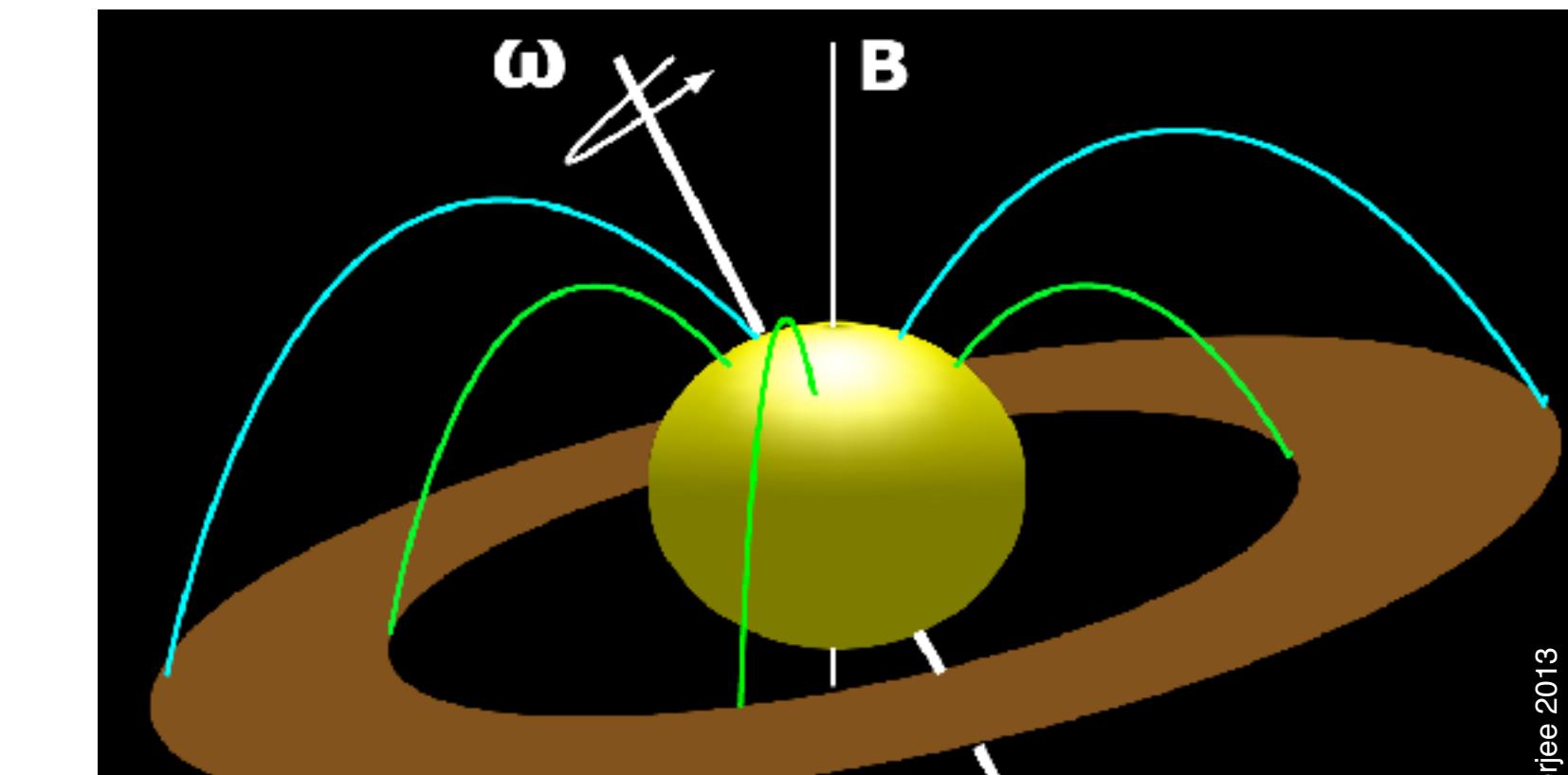
$$R_A \sim \left( \frac{B_*^2 R_*^6}{\dot{M} \sqrt{2GM}} \right)^{2/7}$$

$\omega_{\text{kep}}(R_A) > \omega_*$ : spin-up

$\omega_{\text{kep}}(R_A) < \omega_*$ : spin-down

$\omega_{\text{kep}}(R_A) = \omega_*$ : equilibrium

$$P_{\text{eq}} \propto B^{6/7} \dot{M}^{-3/7}$$



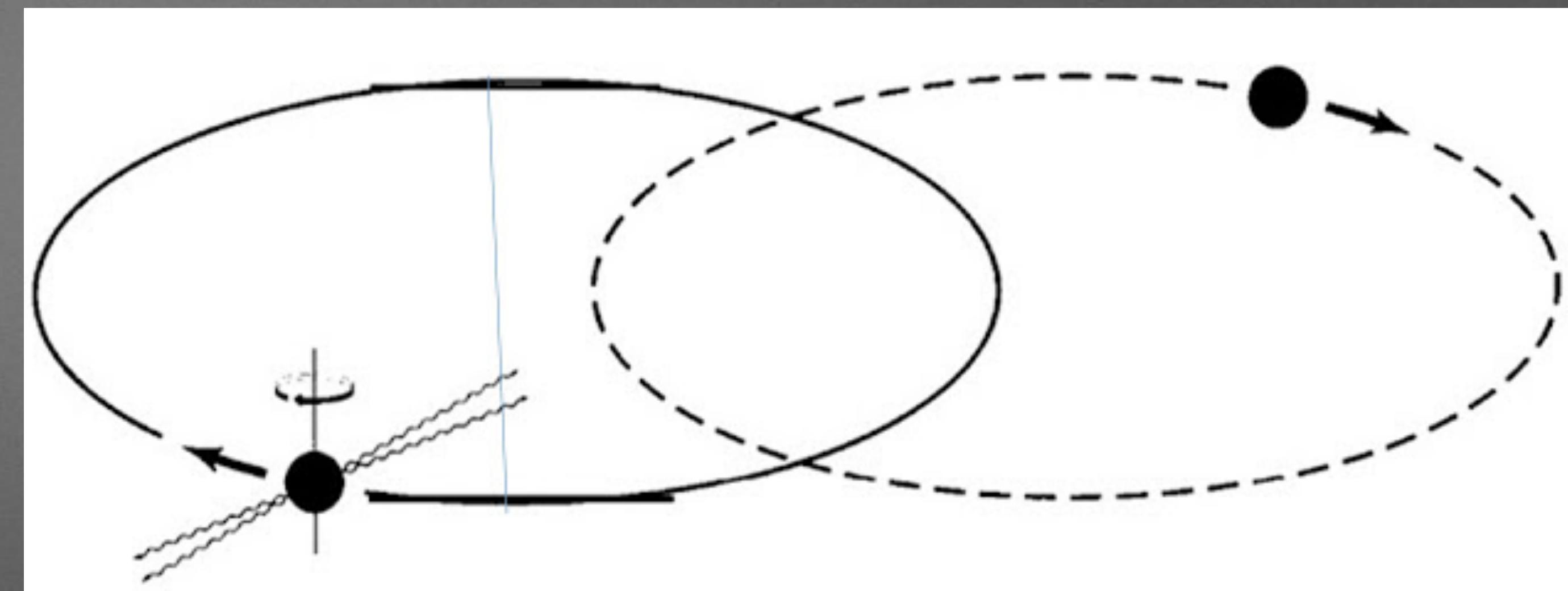
**Recycled  
pulsars**

# Timing relativistic binaries

Neutron Stars in Binaries:  
Keplerian and Post-Keplerian  
relativistic parameters

$P_b, e, \omega$

$dP_b/dt, d\omega/dt, \gamma, r, s$



R. Hulse, Nobel Lecture

$$\dot{\omega} = 3 \left( \frac{P_b}{2\pi} \right)^{-5/3} (T_\odot M)^{2/3} (1 - e^2)^{-1},$$

$$\gamma = e \left( \frac{P_b}{2\pi} \right)^{1/3} T_\odot^{2/3} M^{-4/3} m_2 (m_1 + 2m_2),$$

$$\dot{P}_b = -\frac{192\pi}{5} \left( \frac{P_b}{2\pi} \right)^{-5/3} \left( 1 + \frac{73}{24}e^2 + \frac{37}{96}e^4 \right) (1 - e^2)^{-7/2} T_\odot^{5/3} m_1 m_2 M^{-1/3},$$

$$r = T_\odot m_2,$$

$$s = x \left( \frac{P_b}{2\pi} \right)^{-2/3} T_\odot^{-1/3} M^{2/3} m_2^{-1}.$$

$$T_\odot \equiv GM_\odot/c^3$$

Damour & Deruelle 1986

- Measure Neutron Star Masses
- Test GR predictions
- First evidence of Gravitational Waves from orbital decay

# Concluding Remarks

- Combination of extreme density, gravity, magnetic field and spin make pulsars some of the most exotic objects in the universe
- Pulsars and neutron stars aid the study of fundamental physics - including nuclear forces and the theory of gravity
- Pulsar timing is the most accurate tool of astronomical measurements in existence today, and is continuously improving
- Pulsar timing is also finding application in the area of space navigation
- With new surveys, particularly with the SKA, promising to discover many more pulsars the future of pulsar science is rich and bright