

# Stabilisation of the Arrival Time of a High Energy Electron Beam at the 50 fs Level

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CLIC, a proposed future linear electron-positron collider, and other machines such as XFELs, place tight tolerances on the phase stabilities of their beams. CLIC proposes the use of a novel, high bandwidth and low latency, ‘phase feedforward’ system required to achieve a phase stability of  $0.2^\circ$  at 12 GHz, or about 50 fs. This work documents the results from operation of a prototype phase feedforward system at the CLIC test facility CTF3, with 30 MHz bandwidth and a total hardware latency of 100 ns. New phase monitors with 30 fs resolution, 20 kW amplifiers with 47 MHz bandwidth, and electromagnetic kickers have been designed and installed for the system. The system utilises a chicane in the beamline, for which new optics have been created and commissioned. The prototype has demonstrated CLIC-level phase stability, reducing an initial rms phase variation of  $0.92 \pm 0.04^\circ$  to  $0.20 \pm 0.01^\circ$  across a duration of 10 minutes.

## INTRODUCTION

CLIC [1] is a proposal for a future linear electron positron collider that uses a novel two beam acceleration concept to achieve a high accelerating gradient of 100 MV/m and a collision energy of up to 3 TeV. In this concept the 12 GHz RF power used to accelerate the high energy colliding beams is extracted from high intensity drive beams.

CLIC’s luminosity quickly drops if the RF phase jitters with respect to the main beam, causing energy errors and subsequent beam size growth at the interaction point. The RF phase stability must be  $0.2^\circ$  at 12 GHz (around 50 fs) or better to limit the luminosity loss to below 1% [1]. However, the phase stability of the drive beams cannot be guaranteed to be better than  $2^\circ$  at 12 GHz. CLIC therefore requires a “phase feedforward” (PFF) system that will reduce the drive beam phase jitter (rms) by an order of magnitude. Other machines, such as XFELs, have similar phase stability requirements.

The PFF system poses many challenges, particularly in terms of the hardware bandwidth ( $>17.5$  MHz [2]), power (500 kW amplifiers) and latency requirements. A prototype PFF system has therefore been designed, commissioned and operated at the CLIC test facility CTF3, at CERN, to prove its feasibility. The prototype system follows the same concept as the proposed CLIC scheme, and is the focus of this work.

CTF3 provides a 135 MeV electron beam bunched at 3 GHz with a pulse length of  $1.2 \mu\text{s}$  and a pulse repetition rate of 0.8 Hz. Further details on the CTF3 facility can be found in [REF]. All phases quoted in the paper are given in degrees at 12 GHz.

## SYSTEM DESIGN

A schematic of the PFF system is shown in Fig. 1. The system corrects the phase using two electromagnetic kickers installed before the first and last dipole in a four bend chicane (in the TL2 transfer line). The beam’s path length through the chicane depends on the magnitude and polarity of the voltage applied to the kickers. The phase is measured using a monitor upstream of the chicane (in the CT beam line), and then corrected by setting the kicker voltage to deflect bunches arriving early at the phase monitor on to longer trajectories in the chicane, and bunches arriving late on to shorter trajectories. Downstream of the chicane, in the TBL line, another phase monitor is placed to measure the effects of the correction.

The beam time of flight between the upstream phase monitor and the first kicker in the chicane is 380 ns. By bypassing the combiner ring (CR) and TL1 transfer line (see Fig. 1) the total cable length required to transport signals between the monitor and kickers is shorter, approximately 250 ns. The PFF correction in the chicane can therefore be applied to the same bunch initially measured at the phase monitor, providing the total system hardware latency is less than 130 ns.

## Hardware

The PFF system uses three phase monitors, two electromagnetic kickers, kicker amplifiers and a digitiser/feedforward controller.

The three phase monitors [3] are designed and built by INFN Frascati, with the associated electronics built by CERN. The monitors are 12 GHz resonating cavities with a dipole and monopole mode present. The output from opposing vertical pairs of feedthroughs are summed in hybrids to create a position independent signal. This

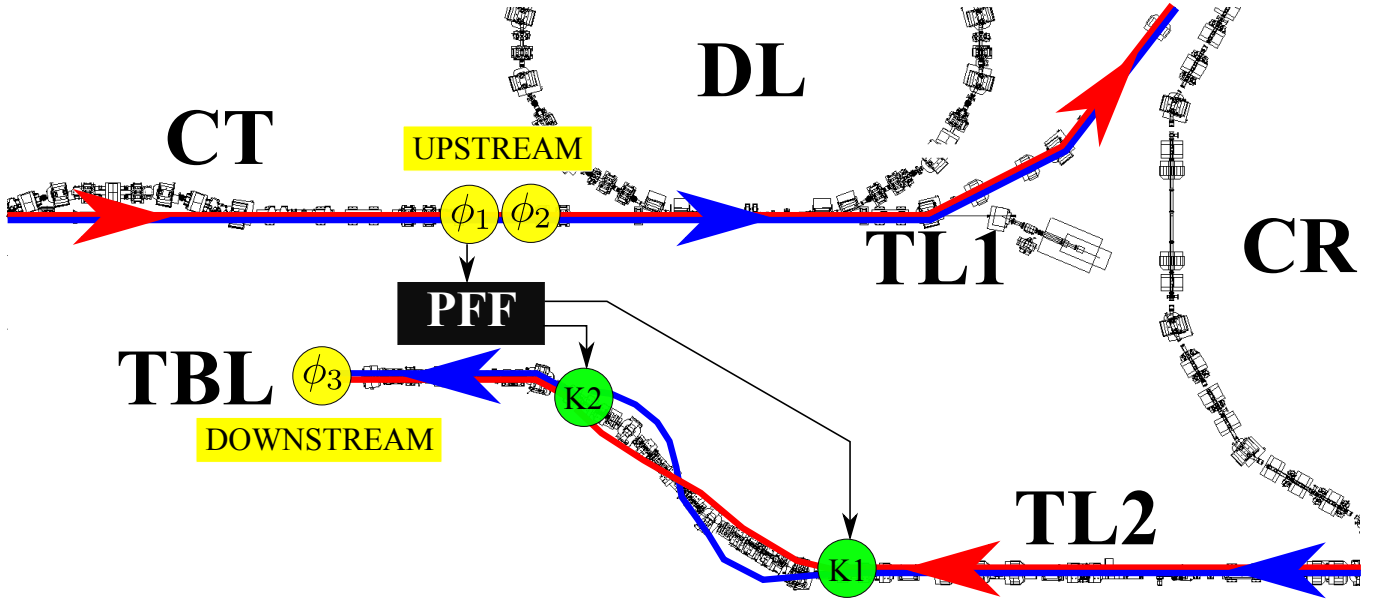


FIG. 1. Schematic of the PFF prototype at CTF3, showing the approximate location of the phase monitors ( $\phi_1$ ,  $\phi_2$  and  $\phi_3$ ) and the kickers (K1 and K2). The black box PFF represents the calculation and output of the correction, including the phase monitor electronics, feedforward controller and kicker amplifiers. A bunch arriving early at  $\phi_1$  is directed on to a longer path in the TL2 chicane using the kickers (blue trajectory), whereas a bunch arriving late will be directed on to a shorter path (red trajectory).

signal is split and mixed with a reference 12 GHz signal in eight separate mixers. The output from the eight mixers is combined, allowing a resolution of  $0.12^\circ$  to be achieved whilst maintaining linearity between  $\pm 70^\circ$ . The quoted resolution is determined by comparing the measurements of the two adjacent upstream monitors (installed in the CT line, see Fig. 1).

The two electromagnetic stripline kickers [4] were also designed and built by INFN Frascati, and are based on the DAFNE design. Each kicker is approximately 1 m in length, with a horizontal strip separation of 40 mm. A voltage of 1.26 kV applied to the downstream end of the kicker strips yields a horizontal deflection of 1 mrad for the 135 MeV CTF3 beam.

The kicker amplifiers [5] have been designed and built by the John Adams Institute/Oxford University. The 20 kW amplifiers consist of low voltage Si FETs driving high voltage SiC FETs, and for an input voltage of  $\pm 2$  V give an output of up to  $\pm 700$  V. The amplifier response is linear within 3% for input voltages between  $\pm 1.2$  V, then starts to saturate. The output has a bandwidth of 47 MHz for small signal variations up to 20% max output, and is slow rate limited for larger variations.

Finally, the Feedforward digitiser and controller (FONT5a board) [5] was also designed and built by John Adams Institute/Oxford University. This digitises the processed phase monitor signals and then calculates and outputs the appropriate voltage with which to drive the amplifier in order to correct the phase. The board consists of a Virtex-5 field programmable gate ar-

ray (FPGA), nine 14-bit analogue to digital converters (ADCs) clocked at 357 MHz, and four digital to analogue converters (DACs). The parameters of the correction, such as the system timing and gain, are controlled on the board via a LabVIEW data acquisition and control system (DAQ).

The combined hardware latency for the PFF system is approximately 100 ns. The output from the FONT5a board is delayed by 30 ns so that the drive voltage from the amplifiers and the beam arrive at the kickers synchronously.

### Chicane Optics

The PFF system places additional constraints on the optics of the correction chicane, and also on the beam lines between the upstream phase monitor and the chicane. These constraints are needed to ensure that the phase depends linearly on the kicker voltage, that the PFF system does not degrade the beam orbit stability downstream of the chicane, and that there is high correlation between the initial (uncorrected) upstream and downstream phase.

The correction range of the PFF system is defined by the kicker design, the maximum output voltage of the kicker amplifiers, and the optics transfer matrix coefficient  $R_{52}$  between the kickers in the chicane. The coefficient  $R_{52}$  describes the change in path length through the chicane per unit deflection at the first kicker. The

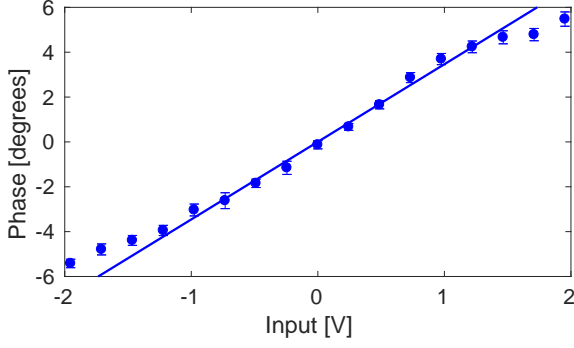


FIG. 2. Downstream phase vs. the kicker amplifier input voltage. Markers show the measured response, and the line a linear fit to the data restricted to inputs between  $\pm 1.2$  V

optics at CTF3 have  $R_{52} = 0.74$  m and at the maximum amplifier output of  $\pm 700$  V the kickers deflect the beam through  $\pm 0.56$  mrad. Together these define a correction range of approximately  $\pm 400$   $\mu\text{m}$ , or  $\pm 6^\circ$ , for the PFF prototype.

The measured phase shift in the chicane versus the amplifier input voltage is shown in Fig. 2, and agrees with the predicted range given the hardware and optics parameters. However, the output of the amplifier is non-linear for input voltages outside  $\pm 1.2$  V. The linear range of the PFF system is therefore closer to  $\pm 4^\circ$ .

The PFF system also should not degrade the transverse stability of beam after chicane. The purpose of the second kicker is to close the orbit bump created by the first kicker, so that the downstream beam orbit is independent of the kicker voltage. In terms of optics transfer matrix coefficients this can be achieved by requiring  $R_{11} = -1$  and  $R_{12} = 0$  between the kickers. Fig. 3 shows the horizontal beam orbit before, in and after the TL2 chicane for the maximum and minimum applied kick. The closure in the BPMs following the chicane is better than 0.1 mm, compared to a maximum offset of 1.5 mm inside the chicane.

### Phase Propagation

The PFF system acts to subtract the measured upstream phase ( $\phi_u$ ) from the initial downstream phase ( $\phi_d$ ) with a gain factor ( $g$ ):  $\phi_{\text{PFF}} = \phi_d - g\phi_u$ , where  $\phi_{\text{PFF}}$  is the corrected downstream phase. The optimal system gain is given by:  $g = \rho_{ud}\sigma_d/\sigma_u$ , where  $\sigma_u$  and  $\sigma_d$  are the initial upstream and downstream phase jitter respectively, and  $\rho_{ud}$  is the correlation between the upstream and downstream phase. The theoretical limit on the corrected downstream phase jitter ( $\sigma_{\text{PFF}}$ ) with this gain is given by:  $\sigma_{\text{PFF}} = \sigma_d\sqrt{1 - \rho_{ud}^2}$ .

One of the key challenges in operating the PFF prototype at CTF3 has been obtaining high correlation be-

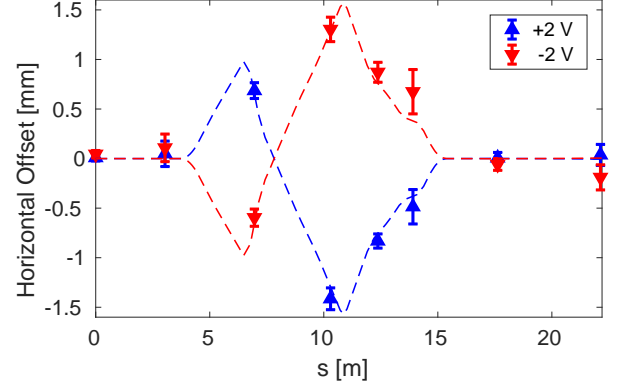


FIG. 3. Horizontal orbit in and around the TL2 chicane at maximum (blue) and minimum (red) input to the kicker amplifiers. Markers show the measured position in beam position monitors, and dashed lines the predicted orbit using the CTF3 MADX model and hardware parameters.

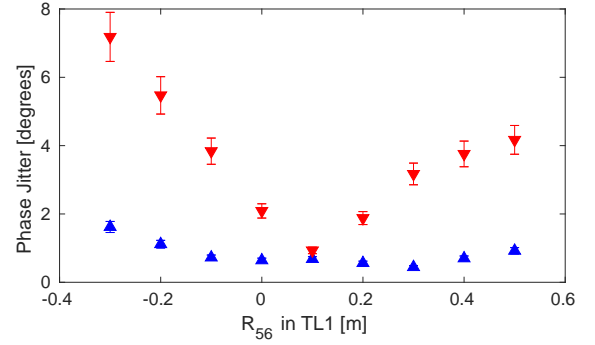


FIG. 4. Downstream (red points) and upstream (blue points) phase jitter vs. the  $R_{56}$  value in the set TL1 optics.

tween the initial, uncorrected, upstream and downstream phase. A correlation of 97% is required to reduce a typical initial phase jitter of  $0.8^\circ$  to the target of  $0.2^\circ$ . Early measurements showed below 40% correlation, and typically a factor 3 increase in the downstream phase jitter with respect to the upstream jitter.

The source of low correlation and jitter amplification was discovered to be energy dependent phase jitter introduced between the upstream and downstream phase monitors. This is described via the optics transfer matrix coefficient  $R_{56}$ :  $\phi_d = \phi_u + R_{56}(\Delta p/p)$ , where  $\Delta p/p$  is the relative beam energy offset.

To achieve high upstream-downstream phase correlation the PFF system requires  $R_{56}$  to be zero between the upstream and downstream phase monitors.

The phase propagation has therefore been optimised by creating new optics for the transfer line TL1 (see Fig. 1) with  $R_{56}$  values ranging between -30 cm and +60 cm in 0.5 cm increments. TL1 can then be used to compensate for the non-zero  $R_{56}$  values in other beam

lines, with the dominant contribution being from the TL2 chicane. In other words, TL1 is tuned such that  $R_{56}(\text{TL1}) = -R_{56}(\text{TL2})$ .

Fig. 4 shows the effect of varying the  $R_{56}$  value in TL1 on the downstream phase jitter. With an  $R_{56}$  of around 10 cm in TL1 the downstream phase jitter is reduced to the same level as the upstream jitter, as desired. The upstream-downstream phase correlation is also increased in these conditions, with correlations above 95% achieved.

In this way the necessary initial conditions for the PFF correction were obtained at CTF3. However, the optics at CTF3 also contain a large second order phase-energy dependence, which leads to a degradation in upstream-downstream phase correlation if there are drifts in beam energy. Energy drifts resulting from instabilities and trips of klystrons at CTF3 have therefore made it difficult to maintain high phase correlations for timescales longer than 10 minutes.

## RESULTS

### Gain Scan

With the optimal gain the PFF correction acts to remove all correlation between the upstream and downstream phase, reducing the downstream phase jitter. If the gain is too small some residual correlation will remain, and if it is too large the correlation will flip sign.

The optimal system gain can be derived empirically by observing the dependence of the downstream phase on the upstream phase with the correction on, as seen in Fig. 5. Optimal gain values for the system are typically in the range 1.0–1.5, being larger than unity when the downstream jitter is larger than upstream, as per the predicted theoretical values.

### Intra-Pulse Phase Variations

The PFF correction is shaped to remove phase variations along the  $1.2 \mu\text{s}$  CTF3 beam pulse. The predominant intra-pulse feature at CTF3 is a roughly parabolic “phase sag” of  $40^\circ$  peak-to-peak, resulting from the use of RF pulse compression. As this is much larger than the  $\pm 6^\circ$  range of the PFF system, only approximately a 400 ns portion of the pulse can be optimally corrected. The phase sag would not be present at CLIC, where in any case the drive beam pulse length is less than 400 ns.

Fig. 6 shows the effect of the PFF system on the intra-pulse phase variations. The convention at CTF3 is to operate the PFF system in interleaved mode, with the correction applied to alternating pulses only. This allows a measurement of the initial (‘PFF Off’) and corrected

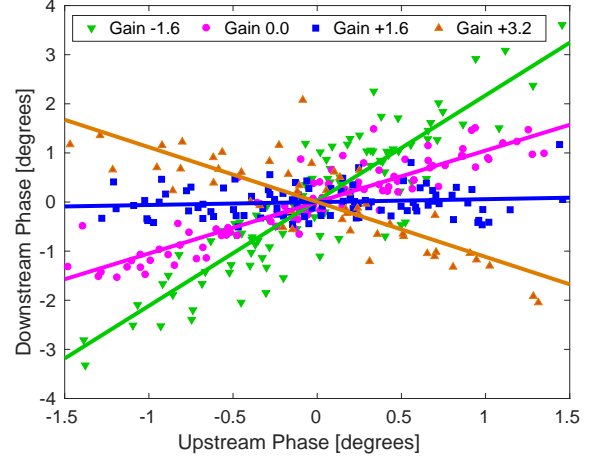


FIG. 5. Downstream phase vs. the upstream phase with the PFF system on at four different gains: -1.6 (green), 0.0 (magenta), +1.6 (blue) and +3.2 (orange). Data points (markers) and linear fits to the data (lines) are shown.

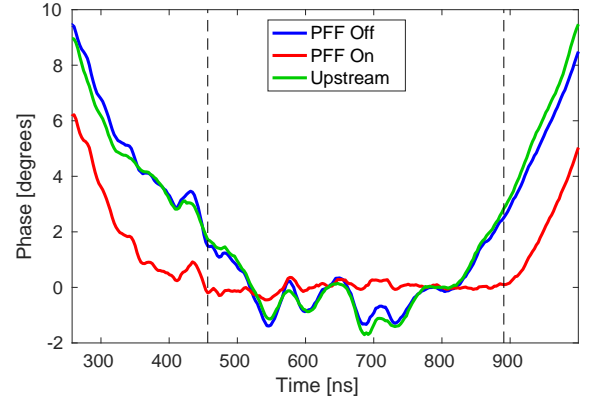


FIG. 6. Effect of the PFF system on intra-pulse phase variations. The pulse shape upstream (green), and downstream with the PFF system off (blue) and on (red) is shown.

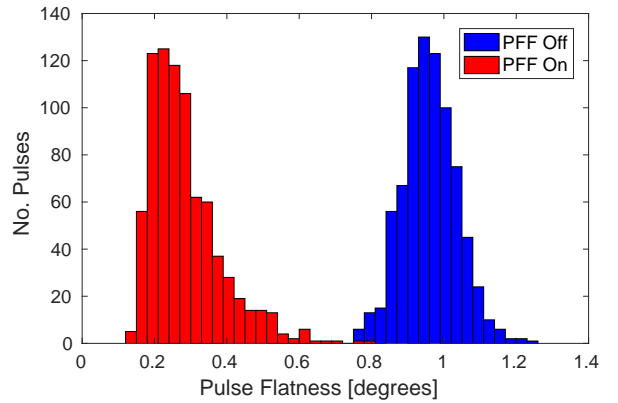


FIG. 7. Distribution of downstream rms phase values, here referred to as the “pulse flatness”, for each beam pulse with the PFF system off (blue) and on (red).

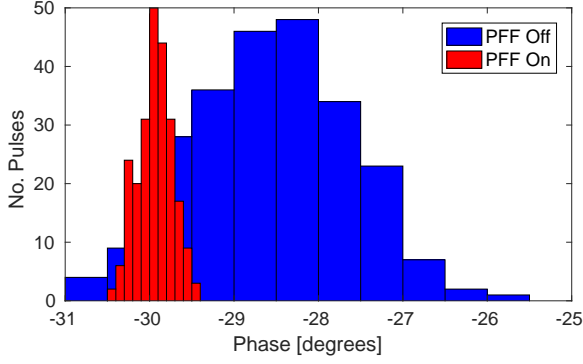


FIG. 8. Distribution of the mean downstream phase with the PFF system off (blue) and on (red).

(‘PFF On’) downstream phase to be performed concurrently. The upstream (PFF input) phase is also shown for comparison. Vertical dashed lines mark a 440 ns portion of the pulse where the correction is optimal, and this range is used to calculate statistics on the effect of the system.

In this range the PFF system flattens the phase, and almost all variations are removed. Residual offsets in the phase are still present where there are small uncorrelated differences between the shape of the initial upstream and downstream phase. Fig. 7 shows the rms phase variation within the 440 ns range for each beam pulse in the dataset, with the PFF system on and off. The PFF off pulses have an rms of  $0.960 \pm 0.003^\circ$  on average, and this is reduced to  $0.285 \pm 0.004^\circ$  by the PFF system.

The PFF system at CTF3 has been verified to reduce the amplitude of phase errors up to a frequency of 25 MHz, exceeding the CLIC requirements.

### Pulse-to-pulse Jitter

As well as removing intra-pulse phase variations the PFF system simultaneously corrects offsets in the overall mean phase, i.e. any pulse-to-pulse jitter. The mean phase of each beam pulse is calculated across the 440 ns range in the central portion of the pulse, as shown before in Fig. 6.

Fig. 8 shows the effect of the PFF system on the pulse-to-pulse stability across a dataset around ten minutes in length. An initial mean downstream phase jitter of  $0.92 \pm 0.04^\circ$  is reduced to  $0.20 \pm 0.01^\circ$  by the PFF correction. All correlation between the upstream and downstream jitter is removed, from  $96 \pm 2\%$  to  $0 \pm 7\%$ . The achieved stability is consistent with the theoretical prediction (considering the initial correlation and jitter) of  $0.26 \pm 0.06^\circ$  within error bars.

This represents the longest time period during which the target CLIC phase stability has been achieved with

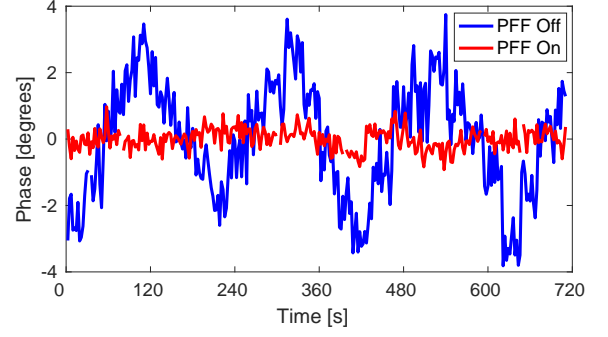


FIG. 9. Mean downstream phase with the PFF system off (blue) and on (red) vs. time, with additional phase variations added to the incoming phase.

the prototype.  $0.30^\circ$  phase jitter has been achieved in 20 minute datasets. The limiting factor for achieving this performance on longer timescales is the incoming beam conditions; in particular klystron instabilities at CTF3 lead to degraded upstream-downstream phase correlation and phase drifts outside the PFF correction range, as previously mentioned.

The PFF system has also been operated whilst intentionally varying the incoming mean phase, as shown in Fig. 9. The PFF system removes the additional phase variations and achieves more than a factor 5 reduction in downstream phase jitter, from  $1.71 \pm 0.07^\circ$  to  $0.32 \pm 0.01^\circ$  in this case.

## CONCLUSIONS

CLIC requires a PFF system to reduce the drive beam phase jitter by an order of magnitude, from  $2.0^\circ$  to  $0.2^\circ$  at 12 GHz, or better than 50 fs stability. A prototype of the system has been in operation at the CLIC test facility CTF3, and corrects the beam phase by varying the path length through a chicane using two electromagnetic kickers.

As well as the kickers, the system uses newly designed phase monitors with  $0.12^\circ$  resolution, high bandwidth 20 kW amplifiers and a low latency digitiser/feedforward controller. The system latency, including hardware and signal transit times, is less than the 380 ns beam time of flight between the input phase monitor and the correction chicane. Therefore, the feedforward correction can be directly applied to the same bunch initially measured at the monitor.

New optics for the correction chicane and other beam lines at CTF3 have been developed to yield the desired phase shifting behaviour and ensure high correlation between the initial upstream and downstream phase.

The prototype system has demonstrated  $0.20 \pm 0.01^\circ$  pulse-to-pulse phase jitter on a time scale of ten minutes.

It has also been shown to be able to flatten intra-pulse phase variations up to a frequency of 25 MHz. On longer timescales the performance of the system is limited by changes to the incoming beam conditions, in particular beam energy, which would be better controlled in any future application at CLIC.

### ACKNOWLEDGEMENTS

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- [1] M. Aicheler *et al.*, “A multi-TeV linear collider based on CLIC technology: CLIC conceptual design report,” CERN-2012-007 (2012).
- [2] A. Gerbershagen *et al.*, Phys. Rev. AB **18**, 041003 (2015).
- [3] F. Marcellini *et al.*, “RF phase monitor final report,” EuCARD-REP-2013-023 (2014).
- [4] A. Ghigo *et al.*, in *Proceedings of IPAC2011*, Vol. TUPC007 (San Sebastian, Spain, 2011) pp. 1000–1002.
- [5] J. Roberts, *Development of a Beam-based Phase Feedforward Demonstration at the CLIC Test Facility (CTF3)*, DPhil thesis, University of Oxford (2016).

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