

$$\vec{x}_{2} \qquad \vec{x}_{1}$$

$$\vec{\chi}_1 = \vec{r} + \vec{\chi}_2$$

$$\vec{\chi}_1 = \vec{r} + \left(-\frac{m_1}{m_2} \vec{\chi}_1 \right)$$

$$(1 + m_1/m_2) \overrightarrow{x_1} = \overrightarrow{r}$$

$$\left(\frac{M_1+M_2}{M_2}\right)\overrightarrow{X}_1 = \overrightarrow{r}$$

$$\frac{M}{m_2}\vec{x} = \vec{r}$$

$$\overrightarrow{X}_1 = \left(\frac{m_2}{M}\right) \overrightarrow{r}$$

M, M2 masses of stars 1 & 2 M=M,+M2 total mass of the system

$$m_{1}\overrightarrow{X}_{1} + m_{2}\overrightarrow{X}_{2} = 0$$

$$0 \text{ rigin at the center of mass}$$

$$-\overrightarrow{X}_{2} = \overrightarrow{T} - \overrightarrow{X}_{1}$$

$$\overrightarrow{X}_{2} = \overrightarrow{X}_{1} - \overrightarrow{T}$$

$$\overrightarrow{X}_{2} = -m_{2} \overrightarrow{X}_{2} - \overrightarrow{T}$$

$$(1 + m_{2}/m_{1})\overrightarrow{X}_{2} = -\overrightarrow{T}$$

$$(\frac{m_{1} + m_{2}}{m_{1}})\overrightarrow{X}_{2} = -\overrightarrow{T}$$

$$\overrightarrow{M}_{1}\overrightarrow{X}_{2} = -\overrightarrow{T}$$

$$\overrightarrow{M}_{1}\overrightarrow{X}_{2} = -\overrightarrow{T}$$

$$\overrightarrow{M}_{2}\overrightarrow{X}_{2} = -\overrightarrow{T}$$

$$\overrightarrow{M}_{1}\overrightarrow{X}_{2} = -\overrightarrow{T}$$

$$\overrightarrow{M}_{2}\overrightarrow{X}_{2} = -\overrightarrow{T}$$

$$\overrightarrow{M}_{2}\overrightarrow{X}_{2} = -\overrightarrow{T}$$

$$\overrightarrow{M}_{1}\overrightarrow{X}_{2} = -\overrightarrow{T}$$

$$\overrightarrow{M}_{2}\overrightarrow{M}_{1}\overrightarrow{M}_{2}$$

$$\overrightarrow{N}_{2} = -\overrightarrow{T}$$

$$\overrightarrow{M}_{1}\overrightarrow{M}_{2}\overrightarrow{M}_{2} = -\overrightarrow{T}$$

For stars moving much slower than the speed of light, $F = m d^2x$

$$F = m \frac{d^2x}{dt^2}, 50$$

$$|\mathbf{M}_{i}| \frac{d^{2}\vec{\gamma_{1}}}{dt^{2}}|_{\mathbf{M}_{i}} = \left|\frac{G \mathbf{M}_{i} \mathbf{m}_{2}}{r^{2}}\right| \left|\frac{\mathbf{M}_{2}}{dt^{2}}\right| = \left|\frac{G \mathbf{M}_{i} \mathbf{M}_{2}}{r^{2}}\right|$$
Since \vec{r} is parallel to \hat{i} ,
$$\frac{d^{2}\vec{\gamma_{1}}}{dt^{2}} = -G \mathbf{M}_{2}\vec{r}/r^{3} \qquad \frac{d^{2}\vec{\gamma_{2}}}{dt^{2}} = +G \mathbf{M}_{i}\vec{r}/r^{3}$$

$$\frac{d^{2}x_{1}}{dt^{2}} = -Gm_{2}\vec{r}/r^{3} \qquad \frac{d^{2}x_{2}}{dt^{2}} = +Gm_{1}\vec{r}/r^{3}$$

$$\frac{d^{2}\vec{r}}{dt^{2}} = \frac{d^{2}}{dt^{2}}(\vec{x_{1}} - \vec{y_{2}}) = \iint_{M_{2}} \left[-Gm_{2}\vec{r}/r^{3} - Gm_{1}\vec{r}/r^{3} \right]$$

$$\frac{d^2r^2}{dt^2} = -G(m_1+m_2)r^2/r^3 = -GMr^2/r^3 \frac{W2.1.2}{dt^2}$$

So we reduce a 2-body problem into a 1-body problem.

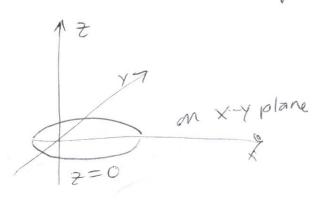
This innocent-looking differential equation is actually pretty difficult to solve analytically as you need elliptic integrals. It is called an ORBIT EQUATION.

Luckily, we know the solution thanks to Euler and others.

$$\vec{r} = r \left(\cos \varphi, \sin \varphi, \sigma \right)$$

where
$$r = \frac{L}{1 - e \cos \varphi} \frac{d\varphi}{dt} = \sqrt{\frac{GM}{L^3}} \left(1 - e \cos \varphi\right)^2$$

e is the eccentricity and L the semi-latus rectum For an ellipse with major axis Za and minor $axis = \frac{2b}{e}$, $e = 1 - \frac{b^2}{a^2}$ $L = a(1 - e^2)$



Semi-tatus coordinate system majoraxis along

(oh) pastran apastron @ \p=0 Foci So

For low eccentricity (e × 0), orbit

looks circular $\sqrt{1-b^2/a^2} \approx 0 \Rightarrow \frac{b^2}{a^2} \approx 1 \Rightarrow a \approx b$

\(\alpha\) \(\a

rmax & a , rmin ~ a

$$a(1-e^2) \approx a$$

r = a (1+e) aphelion in the case of treeastn per, astron &

(furthest point)

 $Q \varphi = \pi$ closest point rmin = 2 (1-e) perihelium Far

earth sun

For high eccentricity (tosates1)

orbit is elongated

$$\sqrt{1-\frac{b^2}{a^2}} \stackrel{<}{\sim} 1 \Rightarrow \frac{b^2}{a^2} \stackrel{<}{\sim} 0 \Rightarrow b \stackrel{<}{\sim} 0$$

=> 7 2 00

conic section

Definition of eltipse - Locus of set of points P whose distance from pto to

directrix focus is a constant e of the

distance from p to directrix

Notice that $\frac{d\theta}{dt} = \sqrt{\frac{GM}{13}} \left(1 - e\cos\varphi\right)^2$

For periastron (closest point) $e=\pi \Rightarrow \cos \phi = 0$

$$\frac{dd}{dt}\Big|_{U=T} = \sqrt{\frac{GM}{L^3}}$$

a pastron (furthest point) Q=0=> cosq=1 For

$$\frac{dq}{dt}\Big|_{\varphi=0} = \sqrt{\frac{GM}{L^3}} (1-e)^2 < \frac{d\varphi}{dt}\Big|_{\varphi=\pi}$$

GRAVITATION

A DOES IT MAKES SENSE FROM

Total energy
$$E = \frac{m_1}{N_1} + \frac{m_2}{N_2} + \frac{m_2}{N_2} = \frac{Gm_1m_2}{r^2}$$

$$Kinetic po tentral$$

$$|J| = \frac{m_1 m_2}{m_1 + m_2} \sqrt{GML}$$

$$|J| = \frac{m_1 m_2}{m_1 + m_2} \int GML$$

$$E = -\frac{m_1 m_2}{m_1 + m_2} \frac{MG}{2a}$$

$$\lim_{m \to \infty} \frac{MG}{2a}$$

Binaries Boto

- · Visual binary stars can be resolved independently Far enough apart from each other and close enough to earth
- · Eclipsing binary Orbital planes oriented along earth line of sight and periodically pass in front of each other. Change in luminosity.

· Spectrum binary - Two spectrums supermposed (84)

If some component is in the line of

Sight, then the spectra will be shifted (Dopler shift, Vobs = Vrest $\sqrt{\frac{1-v/c}{1+v/c}}$)

· Spectroscopic binary - If I is small, then the period will be short (s

$$\frac{\Delta U_1}{U_1} = \frac{m_2}{M} \sqrt{\frac{M6}{Lc^2}} \sin i \left[\frac{M}{Lc^2} \right]$$

i angle between line of sight and orbit