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(24)

If the chemical composition is fixed and uniform, then $p(r)$, $H(r)$ and $\epsilon(r)$ are fixed functions of the density $\rho(r)$ and temperature $T(r)$.

unknown functions

The structure of the star is described by

- $M(r)$ mass contained within sphere of radius \underline{r} determines gravitational potential energy (and kinetic)
- $L(r)$ luminosity is radiant energy per second flowing outward through a spherical surface of radius \underline{r}
- $\rho(r)$ density at \underline{r}
- $T(r)$ temperature at \underline{r}

They are related

depend only on mass distribution



$$\frac{dp(r)}{dr} = - \frac{GM(r)\rho(r)}{r^2}$$

$$\frac{dL(r)}{dr} = 4\pi r^2 \epsilon(r) \rho(r)$$

$$\frac{dM(r)}{dr} = 4\pi r^2 \rho(r)$$

$$\frac{dT(r)}{dr} = - \frac{3H(r)\rho(r)L(r)}{4caT^3(r)4\pi r^2}$$

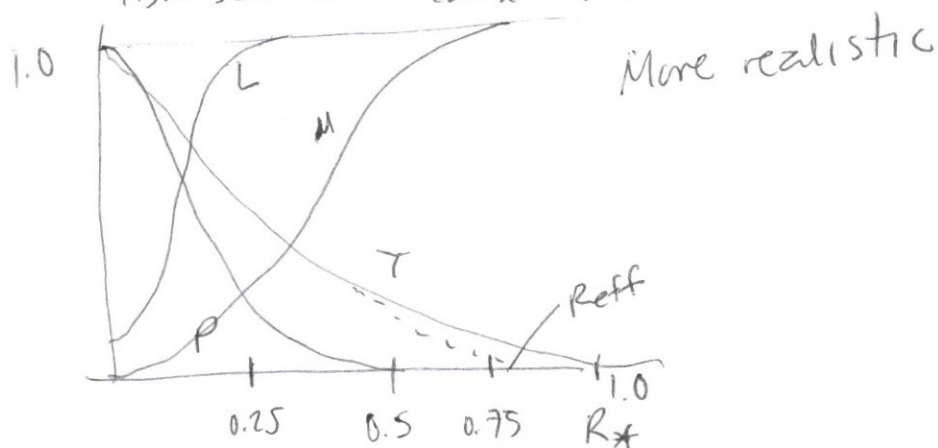
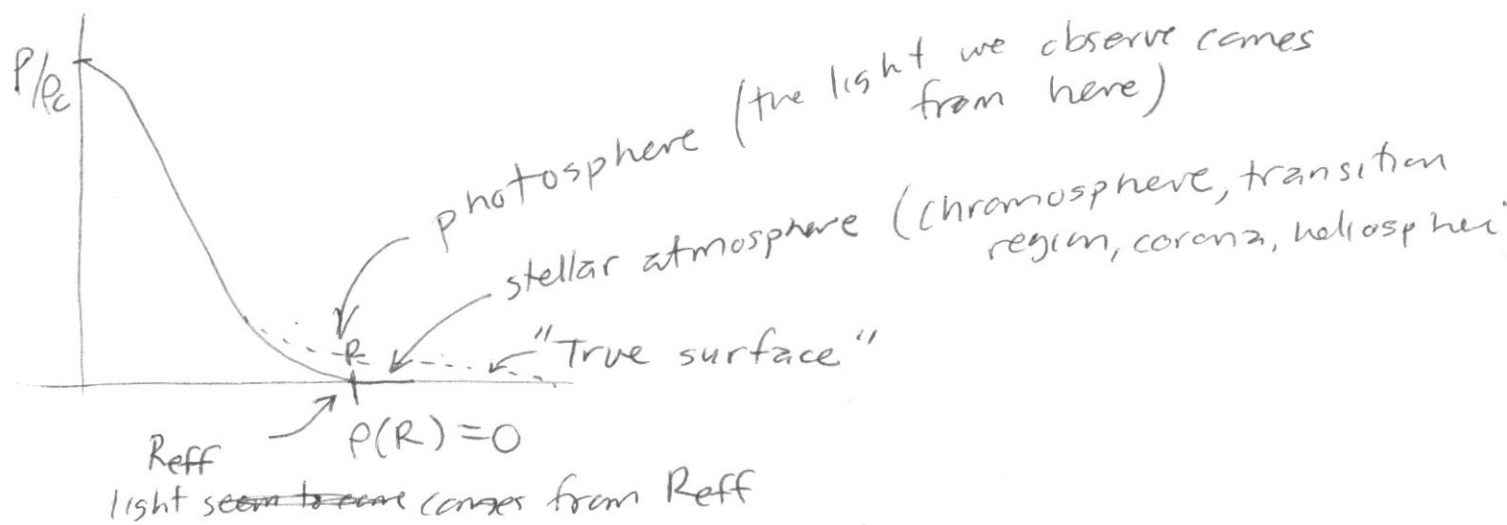
$$M(0) = L(0) = 0$$

$$\rho(R) = T(R) = 0$$

How can the temperature be zero at R ?

(25)

well it can't. Our condition that $\frac{1}{K\rho} \ll R-r$ is not valid close to the surface of the star



Assume all the ~~real~~ luminosity comes from sphere of radius R_{eff} , then

$$\sigma T^4 R_{\text{eff}} \times 4\pi R_{\text{eff}}^2 = L$$

$$\sigma = \frac{ac}{4} = \frac{2\pi^5 R_{\text{eff}}^4}{15 h^3 c^2 N_A^4}$$

The optical depth of the stellar atmosphere is $\int_{R_{\text{eff}}}^{R_{\text{true}}} K(r) \rho(r) dr = \tau_{\text{eff}}$

what this means is that

(26)

$$\rho(R) \ll \rho(0) \quad \text{and} \quad T(R) \ll T(0)$$

For the sun $\rho_c = 98 \pm 15 \frac{\text{g}}{\text{cm}^3}$ (water is 1 g/cm^3)

The T_{eff}
effective temperature

$$\rho_{T_{\text{eff}}=10} = 5 \times 10^{-7} \frac{\text{g}}{\text{cm}^3}$$

9 orders of
magnitude

is obtained by fitting

$$T_c = 13.6 \pm 1.2 \times 10^6 \text{ K}$$

the black body radiation
spectrum at T_{eff} to
observations

$$T_{T_{\text{eff}}=10} = 9700 \text{ K}$$

4 orders of
magnitude

Vogt-Russell theorem

4 first-order differential equations, 4 boundary conditions
of single parameter R

if composition is unique, there is a unique solution
to the system

BUT: The solution is not necessarily unique. there
might be other combinations of parameters
that produce the same solution!

Any given star is going to ~~have~~ experience pressure due to both radiation and matter (ideal gas)

$$P = P_{\text{gas}} + P_{\text{rad}}$$

We know that for black body radiation, $P_{\text{rad}} = \frac{a}{3} T^4$

and for the ideal gas, $P_{\text{gas}} = \frac{\rho k_B T}{m_i \mu}$ where μ is the molecular weight and m_i is the mass of the nucleon (mass of unit atomic weight)

Remember

$$\frac{dT}{dr} = - \frac{3 H \rho L}{4 c a T^3 4 \pi r^2}$$

$$\frac{dP_{\text{rad}}}{dT} = \frac{4}{3} a T^3 \Rightarrow dT = \frac{dP_{\text{rad}}}{\frac{4}{3} a T^3} = \frac{3}{4} \frac{dP_{\text{rad}}}{a T^3}$$

so

$$\frac{3}{4 a T^3} \frac{dP_{\text{rad}}}{dr} = - \frac{3 H \rho L}{4 a T^3 c 4 \pi r^2}$$

$$\frac{dP_{\text{rad}}}{dr} = - \frac{K(r) \rho(r) L(r)}{4 \pi c r^2}$$

Now, remember

$$\frac{dP(r)}{dr} = - \frac{G M(r) \rho(r)}{r^2}$$

$$P(r) - P_{\text{rad}}(r) = P_{\text{gas}}(r)$$

$$- \frac{\overset{P_{\text{rad}}}{K(r)P(r)L(r)}}{4\pi cr^2} + \frac{\overset{P(r)}{GM(r)P(r)}}{r^2} \geq 0 \quad \text{otherwise it would be pure radiation and star would explode}$$

$$\cancel{- \frac{dP(r)}{dr}} = - \frac{dP_{\text{gas}}(r)}{dr}$$

$$\Rightarrow \frac{GM(r)\cancel{P(r)}}{r^2} > \frac{K(r)\cancel{P(r)}L(r)}{4\pi cr^2}$$

$$\Rightarrow K(r)L(r) < 4\pi GcM(r)$$

$$\text{At } r = R_{\text{eff}}, \quad \boxed{K(R_{\text{eff}})L_* < 4\pi GcM_*}$$

Eddington Limit

If this limit is violated, the star disperses material until it is valid (reaches ^{hydrostatic} equilibrium). How far away it is from ~~equilibrium~~ being equal depends on the mass of the star (or galactic nucleus). Very massive stars have almost exclusively P_{rad} due to particles moving close to the speed of light. If you know the mass of the star and can estimate $K(R_{\text{eff}})$, you get L_* from which distance can be obtained.