

Consider a cavity in a perfectly conductive material.

The cavity is a cube of edge L . There is an infinite number of electromagnetic modes within this cavity, although of course, not all will be occupied.

The charge in the cavity is zero, since it is a cavity.

From Gauss' Law, $\vec{\nabla} \cdot \vec{E} = \rho / \epsilon_0 = 0$

\uparrow charge \uparrow permittivity

In integral form (PHYS 2421) you saw it as $\Phi_E = \oint \vec{E} \cdot d\vec{A}$

electric flux surface integral
 \uparrow
 electric field \uparrow
 area element

In differential form (PHYS 4341) $\vec{\nabla} \cdot \vec{E} = \rho / \epsilon_0$

In case you didn't know, "Del" or "Nabla" is shorthand for

$$\vec{\nabla} = \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k}, \text{ so the dot product}$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0 \quad \text{KK Eq. 4.10}$$

is the divergence of the electric field.

Since the walls are conductive, charge is free to move.

If there is an electric field, charge will move until it neutralizes that ~~magnetic~~ ^{electric} field, and this will happen pretty much instantaneously, so $E_{\parallel} = 0$ (the electric field parallel to the wall is zero).

$$\vec{\nabla} \cdot \vec{E} = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z}$$

$$\frac{\partial E_x}{\partial x} = \frac{\partial}{\partial x} \left[E_0 \cos(k_x x) \sin(k_y y) \sin(k_z z) \right] =$$

$$E_0 \left[\cos(k_x x) \sin(k_y y) \frac{\partial}{\partial x} \sin(k_z z) \right]$$

product rule
for three terms

$$+ \sin(k_y y) \sin(k_z z) \frac{\partial}{\partial x} \cos(k_x x)$$

$$+ \sin(k_z z) \cos(k_x x) \frac{\partial}{\partial x} \sin(k_y y) \Big]$$

$$\frac{\partial E_x}{\partial x} = -E_0 k_x \sin(k_x x) \sin(k_y y) \sin(k_z z)$$

Similarly,

$$\frac{\partial E_y}{\partial y} = -E_0 k_y \sin(k_x x) \sin(k_y y) \sin(k_z z)$$

$$\frac{\partial E_z}{\partial z} = -E_0 k_z \sin(k_x x) \sin(k_y y) \sin(k_z z)$$

so

$$\vec{\nabla} \cdot \vec{E} = - \left[E_0 k_x + E_0 k_y + E_0 k_z \right] \sin(k_x x) \sin(k_y y) \sin(k_z z) = 0$$

$$\Rightarrow E_0 \left(\frac{n_x \pi}{L} + \frac{n_y \pi}{L} + \frac{n_z \pi}{L} \right) = 0$$

$$\Rightarrow E_0 n_x + E_0 n_y + E_0 n_z = \vec{E}_0 \cdot \vec{n} = 0 \quad \text{KK Eq. 4.11}$$

$$\text{with } \vec{E}_0 = E_0 \hat{i} + E_0 \hat{j} + E_0 \hat{k} \text{ and } \vec{n} = n_x \hat{i} + n_y \hat{j} + n_z \hat{k}$$

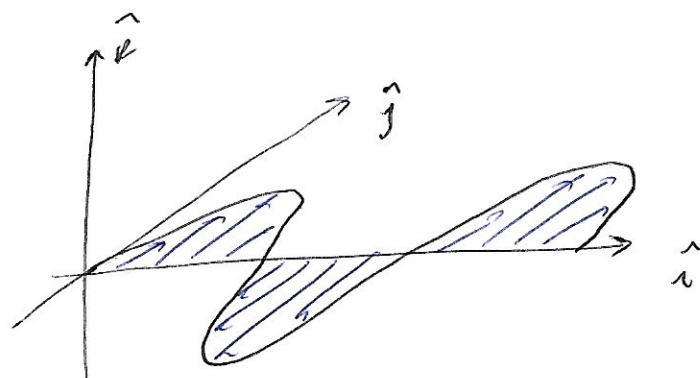
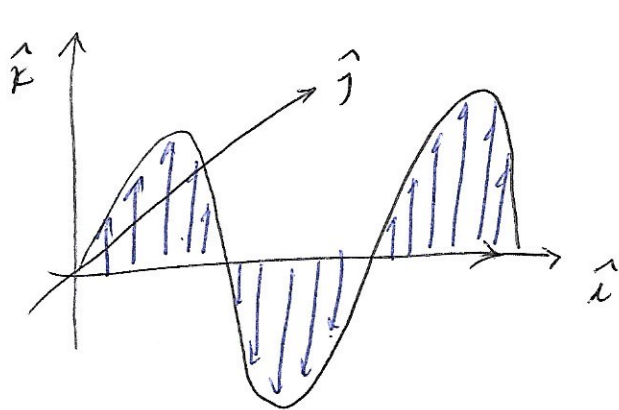
Eq. 4.11 is interesting. One case that satisfies

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$$\vec{E}_0 \cdot \vec{n} = 0 \quad \text{is} \quad n_x \hat{i} + 0 \hat{j} + 0 \hat{k} \quad \text{and} \quad E_{x0} = 0$$

$$\text{also} \quad 0 \hat{i} + n_y \hat{j} + 0 \hat{k} \quad \text{and} \quad E_{y0} = 0$$

$$0 \hat{i} + 0 \hat{j} + n_z \hat{k} \quad \text{and} \quad E_{z0} = 0$$



standing wave in the direction \hat{i} "rattles" the electric field in the transverse directions.

One way to write the wave equation is $\left(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) \vec{E} = 0$ Laplacian

~~For the case of E_z ,~~ In the case of E_z ,

$$c^2 \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) E_z = \frac{\partial^2 E_z}{\partial t^2} \quad \text{KK Eq. 4.12}$$

Looking at the spatial part,

$$c^2 \frac{\partial}{\partial z} \frac{\partial^2 E_z}{\partial z^2} = c^2 \frac{\partial}{\partial z} \left[-E_{z0} k_z \sin(k_x x) \sin(k_y y) \sin(k_z z) \right]$$

← got it before.

$$= -c^2 \left[E_{z0} k_z^2 \sin(k_x x) \sin(k_y y) \cos(k_z z) \right]$$

$$= -k_z^2 c^2 E_z$$

If we repeat for x, y , we get

$$-c^2 [E_x k_x^2 + E_y k_y^2 + E_z k_z^2] = -c^2 \vec{E} \cdot \vec{k}^2$$

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For the temporal part

$$\frac{\partial^2}{\partial t^2} E_z = \frac{\partial^2}{\partial t^2} = -\omega^2 \sin(\omega t) E_z \text{ (stuff)}$$

for the three parts it will be $-\omega^2 \vec{E} \sin(\omega t)$

with boundary conditions, eventually

$$-c^2 \vec{E} \cdot \vec{k}^2 = -\omega^2 \vec{E} \Rightarrow c^2 \vec{k}^2 = \omega^2$$

$$c^2 \frac{\pi^2}{L^2} (n_x^2 + n_y^2 + n_z^2) = \omega^2$$

$$\text{KK Eq. 4.13 } c^2 \pi^2 (n_x^2 + n_y^2 + n_z^2) = \omega^2 L^2$$

$$\text{with } n \equiv (n_x^2 + n_y^2 + n_z^2)^{1/2}, \quad \omega_n^2 = \frac{c^2 \pi^2 n^2}{L^2} \Rightarrow \omega_n = \frac{c n \pi}{L} \quad \text{KK Eq. 4.15}$$

KK Eq. 4.14

c is the speed of wave propagation, in this case the speed of light.

Before, we calculated the expectation value of the number of photons of a given angular frequency ω at temperature τ , this was the Planck distribution. By multiplying by $\hbar\omega$, we got the expectation value of the energy for state "s." The total energy, hence, will be $U = 2 \sum_n \langle \epsilon_n \rangle = 2 \sum_n \frac{\hbar \omega_n}{\exp(\hbar \omega_n / \tau) - 1}$

there are 2 polarization modes

KK Eq 4.16