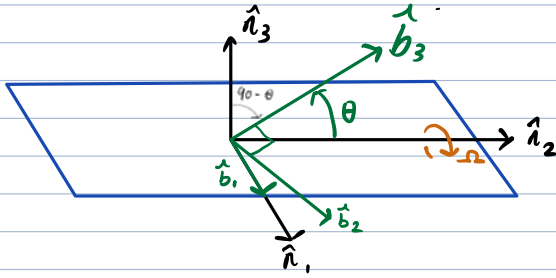


Fig. 1

A solid cylinder of mass m , radius a , and length l is pivoted about a transverse axis ($B-B'$) through its center of mass as shown in Fig. 1. The axis ($A-A'$) rotates with a constant angular velocity Ω . Assume $l > \sqrt{3}a$.

Q5) Find the EOM of the cylinder in terms of θ

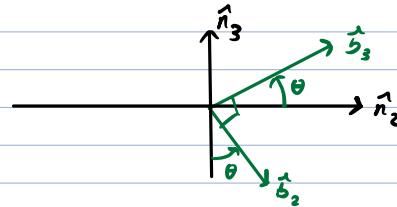
- The external frame is rotating abt $A-A'$ at Ω rad/s.
- The cylinder is free to pitch abt $B-B' \Rightarrow$ 1 DOF



$$\begin{aligned} \text{N-frame: } & \{0, \hat{n}_1, \hat{n}_2, \hat{n}_3\} \\ \text{B-frame: } & \{0, \hat{b}_1, \hat{b}_2, \hat{b}_3\} \end{aligned}$$

- Observe that the B-frame has to be rotated by $(90 - \theta)^\circ$ abt $\hat{b}_1 = \hat{n}_1$, for both frames to align.
- So,

$$\begin{aligned} \hat{b}_1 &= \hat{n}_1 \\ \hat{b}_2 &= \sin\theta \hat{n}_2 - \cos\theta \hat{n}_3 \\ \hat{b}_3 &= \cos\theta \hat{n}_2 + \sin\theta \hat{n}_3 \end{aligned}$$



$$\begin{bmatrix} \hat{b}_1 \\ \hat{b}_2 \\ \hat{b}_3 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \sin\theta & -\cos\theta \\ 0 & \cos\theta & \sin\theta \end{bmatrix} \begin{bmatrix} \hat{n}_1 \\ \hat{n}_2 \\ \hat{n}_3 \end{bmatrix}$$

- observe that the 3x3 matrix is a DCM mapping vectors in N-frame to B-frame.
- So,

$${}^B W = [BN] {}^N W$$

$${}^B \begin{bmatrix} w_1 \\ w_2 \\ w_3 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \sin\theta & -\cos\theta \\ 0 & \cos\theta & \sin\theta \end{bmatrix} \begin{bmatrix} \dot{\theta} \\ \Omega \\ 0 \end{bmatrix}$$

$${}^B \begin{bmatrix} w_1 \\ w_2 \\ w_3 \end{bmatrix} = \begin{bmatrix} \dot{\theta} \\ \Omega \sin\theta \\ \Omega \cos\theta \end{bmatrix} \quad \times$$

- In the above method, we assumed no prior knowledge of Euler Angles and have derived the rotation matrix mapping the inertial to body frame, vice versa.

↳ Can we use 1 of 12 Euler Angle sets to get the angular velocity vector in the body frame of the cylinder?

- If we are to use the kinematic Differential Eqn for this 2N, realize that the B- & N-frame are cascaded.

Ω acts on the N-frame.

$\dot{\theta}$ is the rate at which the cylinder rotates in B-frame.

So, mathematically,

$${}^B\{w\} = \underbrace{[B^{-1}]\{\dot{\theta}\}}_{\substack{\text{Euler Rate in body} \\ \text{mapped to angular} \\ \text{vel. in body}}} + \underbrace{[BN]{}^N\{w\}}_{\substack{\text{mapping angular vel.} \\ \text{of N-frame to} \\ \text{body}}}$$

If we considered the 3-2-1 Euler Angle Set,

$${}^B\begin{Bmatrix} w_1 \\ w_2 \\ w_3 \end{Bmatrix} = \begin{bmatrix} -s\theta_2 & 0 & 1 \\ c\theta_2 s\theta_3 & c\theta_3 & 0 \\ c\theta_2 c\theta_3 & -s\theta_3 & 0 \end{bmatrix} \begin{Bmatrix} \dot{\theta}_3 \\ \dot{\theta}_2 \\ \dot{\theta}_1 \end{Bmatrix} + \begin{bmatrix} 1 & 0 & 0 \\ 0 & s\theta & -c\theta \\ 0 & c\theta & s\theta \end{bmatrix} {}^N\begin{Bmatrix} 0 \\ \Omega \\ 0 \end{Bmatrix}$$

↳ Since the body, the cylinder, is constrained abt \hat{b}_2 & \hat{b}_3 ,

$$\theta_2 = 0 \quad \& \quad \theta_3 = 0 \quad \& \quad \dot{\theta}_1 = \dot{\theta}$$

$$\dot{\theta}_2 = 0 \quad \& \quad \dot{\theta}_3 = 0$$

↳ ${}^N\{w\}$ is strictly only considering angular vel. abt the outer frame.
There is only Ω abt \hat{n}_2 or A-A'.

Then,

$${}^B\begin{Bmatrix} w_1 \\ w_2 \\ w_3 \end{Bmatrix} = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{bmatrix} \begin{Bmatrix} 0 \\ 0 \\ \dot{\theta} \end{Bmatrix} + \begin{Bmatrix} 0 \\ s\theta \Omega \\ c\theta \Omega \end{Bmatrix}$$

$${}^B\begin{Bmatrix} w_1 \\ w_2 \\ w_3 \end{Bmatrix} = \begin{Bmatrix} \dot{\theta} \\ 0 \\ 0 \end{Bmatrix} + \begin{Bmatrix} 0 \\ \Omega s\theta \\ \Omega c\theta \end{Bmatrix}$$

$${}^B\begin{Bmatrix} w_1 \\ w_2 \\ w_3 \end{Bmatrix} = \begin{Bmatrix} \dot{\theta} \\ \Omega \sin(\theta) \\ \Omega \cos(\theta) \end{Bmatrix} \quad \text{X}$$

- Having found ${}^B\{w\}$, we can then find EOM of the cylinder in terms of θ .

Euler's Rotational EOM,

$$\begin{aligned} I_1 \dot{w}_1 + (I_3 - I_2) w_2 w_3 &= L_1 & -\textcircled{1} \\ I_2 \dot{w}_2 + (I_1 - I_3) w_3 w_1 &= L_2 & -\textcircled{2} \\ I_3 \dot{w}_3 + (I_2 - I_1) w_1 w_2 &= L_3 & -\textcircled{3} \end{aligned}$$

• Assuming no external torques ; $L_1 = L_2 = L_3 = 0$

• Cylinder is an axisymmetric body ; $\sim I_1 = I_2 = I_t = \frac{1}{12} m (3a^2 + l^2)$

$$\sim I_3 = I_a = \frac{1}{2} m a^2$$

- Eqn ③ goes to zero as $L_3 = 0$ & $(I_2 - I_1) = 0$.
 we can either use ② or ③.
 Since we want the EOM in terms of θ , we use ①,

$$I_1 \dot{\omega}_1 + (I_3 - I_2) \omega_2 \omega_3 = 0$$

$$I_t \ddot{\theta} + (I_a - I_t) (-\Omega \sin \theta) (\Omega \cos \theta) = 0$$

$$\ddot{\theta} + \frac{I_a - I_t}{I_t} \Omega^2 \sin \theta \cos \theta = 0$$

$$\ddot{\theta} - \frac{(I_t - I_a)}{I_t} \frac{\Omega^2}{2} \sin(2\theta) = 0$$

~~✗~~

Q6) Finding the natural freq., ω_n , of small oscillations abt $\theta_0 = \pi/2$.

$$\ddot{\theta} - \frac{(\frac{1}{12} m(3a^2 + l^2) - \frac{1}{2} m a^2)}{\frac{1}{12} m(3a^2 + l^2)} \frac{\Omega^2}{2} \sin(2\theta) = 0$$

$$\ddot{\theta} - \frac{\frac{1}{12} m(3a^2 + l^2 - 6a^2)}{\frac{1}{12} m(3a^2 + l^2)} \frac{\Omega^2}{2} \sin(2\theta) = 0$$

$$\ddot{\theta} - \frac{(-3a^2 + l^2)}{3a^2 + l^2} \frac{\Omega^2}{2} \sin(2\theta) = 0$$

$$\ddot{\theta} + \frac{(3a^2 - l^2)}{3a^2 + l^2} \frac{\Omega^2}{2} \sin(2\theta) = 0$$

For small perturbations abt $\theta_0 = \pi/2$,
 $\theta = \theta_0 + \delta\theta$
 $\dot{\theta} = \delta\dot{\theta}$
 $\ddot{\theta} = \delta\ddot{\theta}$

Then,

$$\delta\ddot{\theta} + \frac{(3a^2 - l^2)}{3a^2 + l^2} \frac{\Omega^2}{2} \sin(2(\frac{\pi}{2} + \delta\theta)) = 0$$

$$\delta\ddot{\theta} + \frac{3a^2 - l^2}{3a^2 + l^2} \frac{\Omega^2}{2} \sin(\pi + 2\delta\theta) = 0$$

$$\delta\ddot{\theta} + \frac{3a^2 - l^2}{3a^2 + l^2} \frac{\Omega^2}{2} (\cancel{\sin(\pi)} \cos(2\delta\theta) + \cancel{\cos(\pi)} \sin(2\delta\theta)) = 0$$

$$\delta\ddot{\theta} + \frac{3a^2 - l^2}{3a^2 + l^2} \frac{\Omega^2}{2} \sin(2\delta\theta) = 0, \quad \sin(2\delta\theta) \approx 2\delta\theta$$

$$\delta\ddot{\theta} + \frac{3a^2 - l^2}{3a^2 + l^2} \Omega^2 (\delta\theta) = 0$$

↳ This is in the form of; $\ddot{x} + \omega_n^2 x = 0$

- Note that $l > \sqrt{3} a$
 $l^2 > 3 a^2$
 $l^2 - 3a^2 > 0$

So if we take the square root of the coeff. of $\delta\theta$,
 ω_n wouldn't be a real value.
 So,

$$\delta\ddot{\theta} - \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 (\delta\theta) = 0$$

Then,

$$\omega_n = \sqrt{\frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2}$$

$$\omega_n = |\Omega| \sqrt{\frac{l^2 - 3a^2}{l^2 + 3a^2}}$$

~~✗~~

Q7) What is $\dot{\theta}$ when $\theta = \pi/2$, if the cylinder is released from $\theta=0$ with a small $+\theta_0$?

$$\ddot{\theta} - \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 \sin(2\theta) = 0$$

$$\ddot{\theta} = \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 \sin(2\theta)$$

$$\frac{d(\dot{\theta})}{dt} = \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 \sin 2\theta$$

$$d\dot{\theta} = \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 \sin 2\theta dt, \quad \dot{\theta} = \frac{d\theta}{dt}$$

$$\dot{\theta} d\dot{\theta} = \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 \sin 2\theta \frac{d\theta}{dt} dt$$

$$\dot{\theta} d\dot{\theta} = \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 \sin 2\theta d\theta$$

$$\int_{\dot{\theta}_0}^{\dot{\theta}} \dot{\theta} d\dot{\theta} = \int_{\theta_0}^{\theta} \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 \sin 2\theta d\theta$$

$$\left[\frac{\dot{\theta}^2}{2} \right]_{\dot{\theta}_0}^{\dot{\theta}} = \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 \left[\frac{-\cos 2\theta}{2} \right]_{\theta_0}^{\theta}$$

$$\frac{\dot{\theta}^2 - \dot{\theta}_0^2}{2} = \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 \left[\frac{-\cos 2\theta}{2} + \frac{\cos 2\theta_0}{2} \right]$$

$$\dot{\theta}^2 - 0 = \frac{l^2 - 3a^2}{l^2 + 3a^2} \Omega^2 (-\cos 2\theta + \cos 2\theta_0)$$

$$\dot{\theta} = \Omega \sqrt{\frac{l^2 - 3a^2}{l^2 + 3a^2} (-\cos 2\theta + \cos 2\theta_0)}$$