

# The behaviour of non-trivial zeroes in tapered zeroth order Riemann Siegel formula finite Dirichlet Series about the first quiescent region with lower symmetry dirichlet coefficients near high peaks of the 5-periodic Davenport-Heilbronn functions.

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## DRAFT Executive Summary

An investigation of the non-trivial zero behaviour about large 5-periodic Davenport-Heilbronn function peaks at  $t=\{4545032985.7, \dots\}$  for a simple perturbation of the dirichlet coefficients of tapered zeroth order Riemann Siegel formula for the 5-periodic Davenport-Heilbronn function dirichlet series is reported. The behaviour is compared to the behaviour of the tapered finite  $L(\chi_5(2, \cdot), s)$ ,  $L(\chi_5(3, \cdot), s)$  function Dirichlet Series of the zeroth order Riemann Siegel formula truncated at the first quiescent region which are the L function basis set of the 5-periodic Davenport-Heilbronn functions. An interesting correspondence between the location of the S value discontinuities associated with the non-trivial zeroes off the critical line of the  $f_{1,2}(s)$  function approximations and when the perturbed non-trivial zeroes intersect with the critical line is observed.

## Introduction

The tapered finite Dirichlet Series truncated about the second quiescent region in the final plateau of the oscillatory divergence of dirichlet series of L functions (or linear combinations thereof) provides a useful approximation of the mean value of the infinite series sum (i.e., averaging out the oscillatory divergence) [1-4]. For the 5 periodic Davenport-Heilbronn with non-trivial zeroes off the critical line and their underlying L function components, the second quiescent region is  $N_2 = \frac{t}{\pi} \cdot 5$  [4] where the underlying L function components have the conductor value 5 and are 1st degree L functions.

A technically weaker approximation but much faster calculation can be achieved using the tapered zeroth order Riemann Siegel formula based on the finite dirichlet series about the first quiescent region using resurgence. For the 5 periodic Davenport-Heilbronn and their underlying L function components, the first quiescent region is  $N_1 = \sqrt{\frac{t}{2\pi}} \cdot 5$  [4]. Given tapering using 2048 points (1024 integers  $\leq N_1$  and 1024 integers  $> N_1$ ), the tapered zeroth order Riemann Siegel formula for 5 periodic Davenport-Heilbronn has useful accuracy to discern zero spacings  $> 1e-5$  for  $t > 1.318e6$ .

In this paper, the non-trivial zero behaviour of a simple perturbation of the dirichlet coefficients of the 5-periodic Davenport-Heilbronn functions which exhibit non-trivial zeroes off the critical line is investigated via the tapered zeroth order Riemann Siegel function approximation [5]. The behaviour of the perturbed 5-periodic Davenport-Heilbronn functions is compared to their higher symmetry L function counterparts.

## The 5-periodic Davenport-Heilbronn functions and their underlying L function components

In L-function, Dirichlet series and Hurwitz Zeta function form the two 5-periodic Davenport-Heilbronn functions are of the form [6]

$$f_1(s) = \frac{1}{2\cos(\theta_1)} \left[ e^{i\theta_1} L(\chi_5(2, \cdot), s) + e^{-i\theta_1} L(\chi_5(3, \cdot), s) \right] \quad (1)$$

$$= 1 + \frac{\tan(\theta_1)}{2^s} - \frac{\tan(\theta_1)}{3^s} - \frac{1}{4^s} + \frac{0}{5^s} + \dots \quad (2)$$

$$= 5^{-s} \left( \zeta\left(s, \frac{1}{5}\right) + \tan(\theta_1) \cdot \zeta\left(s, \frac{2}{5}\right) - \tan(\theta_1) \cdot \zeta\left(s, \frac{3}{5}\right) - \zeta\left(s, \frac{4}{5}\right) \right) \quad (3)$$

where

$$\tan(\theta_1) = \frac{(\sqrt{10 - 2\sqrt{5}} - 2)}{(\sqrt{5} - 1)} \quad (4)$$

$$= 0.284079043840412296028291832393 \quad (5)$$

and

$$\theta_1 = 0.276787179448522625754266365045 \quad \text{radians} \quad (6)$$

The Davenport-Heilbronn  $f_1(s)$  function has the functional equation

$$f_1(s) = 5^{(\frac{1}{2}-s)} 2(2\pi)^{(s-1)} \cos\left(\frac{\pi s}{2}\right) \Gamma(1-s) f_1(1-s) = \chi(f_1(s)) \cdot f_1(1-s) \quad (7)$$

The second linear combination of L-functions 5-periodic Davenport Heilbronn function example  $f_2(s)$  [7,8] has the designation  $\tau_-(s)$  [8] arising from  $f_1(s)$  ( $\tau_+(s)$ ) &  $f_2(s)$  being the two coupled solutions of linear combinations of the  $\chi_5(2, \cdot)$  and  $\chi_5(3, \cdot)$  L-functions. The more recent work [8] as well as providing the functional equation, estimates the highest(lowest)  $\text{Re}(s)$  values for non-trivial zeroes of  $f_2(s)$  are approximately bounded by  $\text{Re}(s)=2.37$  (-1.37).

Expressed in L-function, Dirichlet series and Hurwitz Zeta function form the  $f_2(s)$  5-periodic function is

$$f_2(s) = \frac{1}{2\cos(\theta_2)} \left[ e^{i\theta_2} L(\chi_5(2, \cdot), s) + e^{-i\theta_2} L(\chi_5(3, \cdot), s) \right] \quad (8)$$

$$= 1 - \frac{\tan(\theta_2)}{2^s} + \frac{\tan(\theta_2)}{3^s} - \frac{1}{4^s} + \frac{0}{5^s} + \dots \quad (9)$$

$$= 5^{-s} \left( \zeta\left(s, \frac{1}{5}\right) - \tan(\theta_2) \cdot \zeta\left(s, \frac{2}{5}\right) + \tan(\theta_2) \cdot \zeta\left(s, \frac{3}{5}\right) - \zeta\left(s, \frac{4}{5}\right) \right) \quad (10)$$

where

$$\tan(\theta_2) = \frac{1}{0.284079043840412296028291832393} \quad (11)$$

and

$$\theta_2 = 1.2940091473463739934770553265951171821 \quad \text{radians} \quad (12)$$

The Davenport-Heilbronn  $f_2(s)$  function has the functional equation [8]

$$f_2(s) = -5^{(\frac{1}{2}-s)} 2(2\pi)^{(s-1)} \cos\left(\frac{\pi s}{2}\right) \Gamma(1-s) f_2(1-s) = -\chi(f_1(s)) \cdot f_2(1-s) \quad (13)$$

where the multiplicative factor on the RHS of equations (13) and (7) differ by a factor of -1. <sup>1</sup>

The underlying L-functions of  $f_1(s)$  and  $f_2(s)$  are a dual-pair of L functions [9] with the functional equations

$$L(\chi_5(2, \cdot), s) = \epsilon \cdot \chi(f_1(s)) L(\chi_5(3, \cdot), 1-s) \quad (14)$$

$$L(\chi_5(3, \cdot), s) = \bar{\epsilon} \cdot \chi(f_1(s)) L(\chi_5(2, \cdot), 1-s) \quad (15)$$

where

$$\begin{aligned} \epsilon(L(\chi_5(2, \cdot), s)) &= (0.85065080835203993218154049706301107225... \\ &\quad + i * 0.52573111211913360602566908484787660729...) \end{aligned} \quad (16)$$

is the sign of the functional equation for the L-function  $L(\chi_5(2, \cdot), s)$  [9].

### **Tapered zeroth order Riemann Siegel formula for $L(\chi_5(2, \cdot), s)$ , $L(\chi_5(3, \cdot), s)$ , $f_1(s)$ and $f_2(s)$**

To allow feasible calculations of the approximate behaviour of the 5-periodic Davenport-Heilbronn functions and their L-function components at higher values along the imaginary co-ordinate in the complex plane, an 2048 tapered finite zeroth order Riemann Siegel formula is used in this paper. Using 2048 point tapering, the tapered zeroth order Riemann Siegel formula for the 5 periodic Davenport-Heilbronn function and and their L-function components has useful accuracy to discern zero spacings  $> 1e-5$  for  $t > 1.318e6$ .

Firstly given the standard Riemann Siegel formula [10,11], the known functional equations for the 5-periodic Davenport-Heilbronn functions and their L-function components (equations (7), (13-15)) and their common first quiescent region location  $N = \sqrt{\frac{t}{2\pi} * 5}$  [4].

$$L(\chi_5(2, \cdot), s) \approx \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor)} \left( \frac{\chi_5(2, \cdot)}{k^s} \right) + \epsilon \cdot \chi(f_1(s)) \cdot \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor)} \left( \frac{\chi_5(3, \cdot)}{k^{(1-s)}} \right) \quad \text{to zeroth order} \quad (17)$$

Given the above formula, it is straightforward to derive a tapered finite zeroth order Riemann Siegel function approximation of the form

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<sup>1</sup>Earlier papers by Martin incorrectly omitted the -1 factor in the RHS definition of  $f_2(s)$ . In those papers only  $\text{abs}(\chi(f_2(s)))$  was used in the Z function calculations that were performed so the omission does not change the numerical results.

$$\begin{aligned}
L(\chi_5(2, \cdot), s) \approx & \left\{ \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor - p)} \left( \frac{\chi_5(2, \cdot)}{k^s} \right) + \sum_{i=(-p+1)}^p \frac{\frac{1}{2^{2p}} \left( 2^{2p} - \sum_{k=0}^{i+p-1} \binom{2p}{2p-k} \right) \cdot \chi_5(2, \cdot)}{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor + i)^s} \right\} \\
& + \epsilon \cdot \chi(f_1(s)) \cdot \left\{ \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor - p)} \left( \frac{\chi_5(3, \cdot)}{k^{(1-s)}} \right) + \sum_{i=(-p+1)}^p \frac{\frac{1}{2^{2p}} \left( 2^{2p} - \sum_{k=0}^{i+p-1} \binom{2p}{2p-k} \right) \cdot \chi_5(3, \cdot)}{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor + i)^{(1-s)}} \right\} \quad \text{as } t \rightarrow \infty
\end{aligned} \tag{18}$$

where  $2p=2048$  (for 2048 point tapering) is used in this paper. An important advantage of using the zeroth order expression for the Riemann Siegel function approximation is to simplify the calculation of first and second order derivatives for execution of the non-trivial zero quadrature search under the perturbation parameter  $(\alpha)$  to be described in the next section.

For the other three functions the following tapered zeroth order Riemann Siegel formula apply

$$L(\chi_5(3, \cdot), s) \approx \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor)} \left( \frac{\chi_5(3, \cdot)}{k^s} \right) + \bar{\epsilon} \cdot \chi(f_1(s)) \cdot \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor)} \left( \frac{\chi_5(2, \cdot)}{k^{(1-s)}} \right) \quad \text{to zeroth order} \tag{19}$$

Given the above formula, it is straightforward to derive a tapered finite zeroth order Riemann Siegel function approximation of the form

$$\begin{aligned}
L(\chi_5(3, \cdot), s) \approx & \left\{ \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor - p)} \left( \frac{\chi_5(3, \cdot)}{k^s} \right) + \sum_{i=(-p+1)}^p \frac{\frac{1}{2^{2p}} \left( 2^{2p} - \sum_{k=0}^{i+p-1} \binom{2p}{2p-k} \right) \cdot \chi_5(3, \cdot)}{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor + i)^s} \right\} \\
& + \bar{\epsilon} \cdot \chi(f_1(s)) \cdot \left\{ \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor - p)} \left( \frac{\chi_5(2, \cdot)}{k^{(1-s)}} \right) + \sum_{i=(-p+1)}^p \frac{\frac{1}{2^{2p}} \left( 2^{2p} - \sum_{k=0}^{i+p-1} \binom{2p}{2p-k} \right) \cdot \chi_5(2, \cdot)}{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor + i)^{(1-s)}} \right\} \quad \text{as } t \rightarrow \infty
\end{aligned} \tag{20}$$

$$\begin{aligned}
f_1(s) \approx & \left\{ \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor - p)} \left( \frac{\chi(f_1, k, \text{ mod } 5)}{k^s} \right) + \sum_{i=(-p+1)}^p \frac{\frac{1}{2^{2p}} \left( 2^{2p} - \sum_{k=0}^{i+p-1} \binom{2p}{2p-k} \right) \cdot \chi(f_1, k, \text{ mod } 5)}{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor + i)^s} \right\} \\
& + \chi(f_1(s)) \cdot \left\{ \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor - p)} \left( \frac{\chi(f_1, k, \text{ mod } 5)}{k^{(1-s)}} \right) + \sum_{i=(-p+1)}^p \frac{\frac{1}{2^{2p}} \left( 2^{2p} - \sum_{k=0}^{i+p-1} \binom{2p}{2p-k} \right) \cdot \chi(f_1, k, \text{ mod } 5)}{(\lfloor \sqrt{\frac{t}{2\pi} * 5} \rfloor + i)^{(1-s)}} \right\} \\
& \text{as } t \rightarrow \infty
\end{aligned} \tag{21}$$

where  $\chi(f_1, k, \text{ mod } 5) = \{1, \tan(\theta_1), -\tan(\theta_1), -1, 0\}$  following equation (2)

$$\begin{aligned}
f_2(s) \approx & \left\{ \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi}} * 5 \rfloor - p)} \left( \frac{\chi(f_2, k, \text{ mod } 5)}{k^s} \right) + \sum_{i=(-p+1)}^p \frac{\frac{1}{2^{2p}} \left( 2^{2p} - \sum_{k=0}^{i+p-1} \binom{2p}{2p-k} \right) \cdot \chi(f_2, k, \text{ mod } 5)}{(\lfloor \sqrt{\frac{t}{2\pi}} * 5 \rfloor + i)^s} \right\} \\
& - \chi(f_1(s)) \cdot \left\{ \sum_{k=1}^{(\lfloor \sqrt{\frac{t}{2\pi}} * 5 \rfloor - p)} \left( \frac{\chi(f_2, k, \text{ mod } 5)}{k^{(1-s)}} \right) + \sum_{i=(-p+1)}^p \frac{\frac{1}{2^{2p}} \left( 2^{2p} - \sum_{k=0}^{i+p-1} \binom{2p}{2p-k} \right) \cdot \chi(f_2, k, \text{ mod } 5)}{(\lfloor \sqrt{\frac{t}{2\pi}} * 5 \rfloor + i)^{(1-s)}} \right\} \\
& \text{as } t \rightarrow \infty
\end{aligned} \tag{22}$$

where  $\chi(f_2, k, \text{ mod } 5) = \{1, -\frac{1}{\tan(\theta_1)}, \frac{1}{\tan(\theta_1)}, -1, 0\}$  following equation (9)

Figure 1 shows the Riemann Siegel Z function behaviour based on tapered zeroth order Riemann Siegel formula calculations about the first quiescent region for the four functions  $f_2(s)$ ,  $f_1(s)$ ,  $L(\chi_5(2, \cdot), s)$  and  $L(\chi_5(3, \cdot), s)$ , along the critical line  $s=0.5+I^*t$  in the interval  $t=(4545032981.5, 4545032988.5)$ . In this interval,  $L(\chi_5(2, \cdot), s)$  has a large peak of height  $\sim 130$  at  $t=4,545,032,985.7$  which transforms into a large peak for  $f_2(s)$  of height  $\sim 240$  and  $f_1(s)$  of height  $\sim 65$  after accounting for the  $L(\chi_5(2, \cdot), s)$  weighted contribution to  $f_2(s)$  and  $f_1(s)$ .

## A simple perturbation of the 5-periodic Davenport-Heilbronn functions to produce lower symmetry behaviour

A simple perturbation of the 5-periodic Davenport-Heilbronn functions can be achieved by the modified function

$$f_{1,2}(s, \alpha)_{\text{pert}} = (1 - \alpha) + \alpha \cdot f_{1,2}(s) \tag{23}$$

and likewise for their underlying L-function components

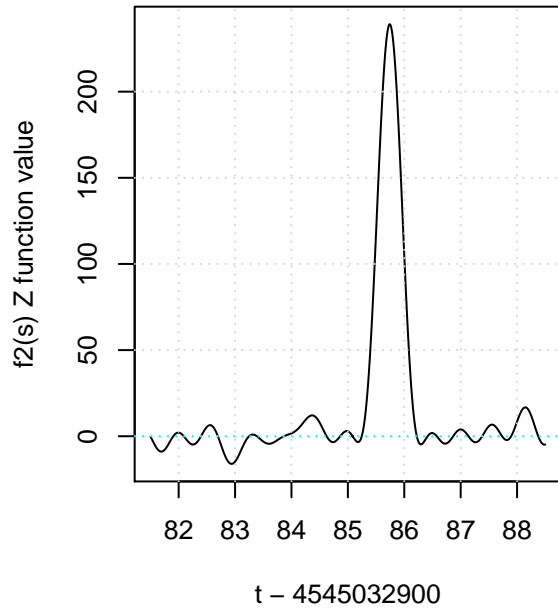
$$L(\chi_5(2, \cdot), s, \alpha)_{\text{pert}} = (1 - \alpha) + \alpha \cdot L(\chi_5(2, \cdot), s) \tag{24}$$

$$L(\chi_5(3, \cdot), s, \alpha)_{\text{pert}} = (1 - \alpha) + \alpha \cdot L(\chi_5(3, \cdot), s) \tag{25}$$

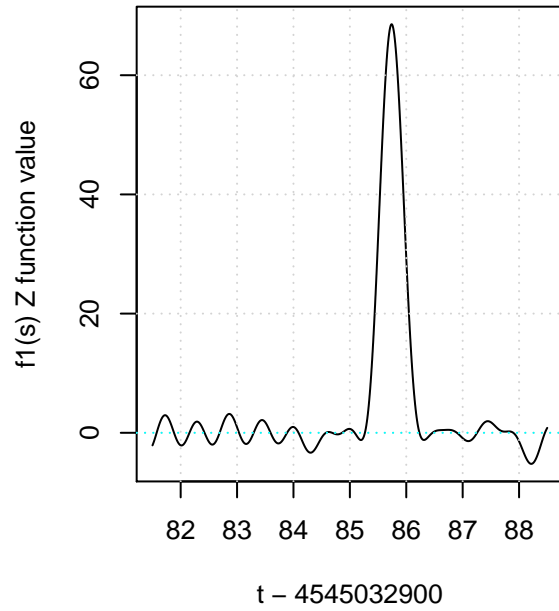
The impact of the perturbation is that the symmetry of the function is lowered compared to the 5-periodic Davenport-Heilbronn function or L function counterpart behaviour and the potential origin of the non-trivial zeroes with respect to Gram's law type behaviour can be ascertained.

An advantage of using the tapered zeroth order expression for the Riemann Siegel function approximation is that it simplifies the calculation of first and second order derivatives for execution of the non-trivial zero quadrature search under the perturbation parameter ( $\alpha$ ).

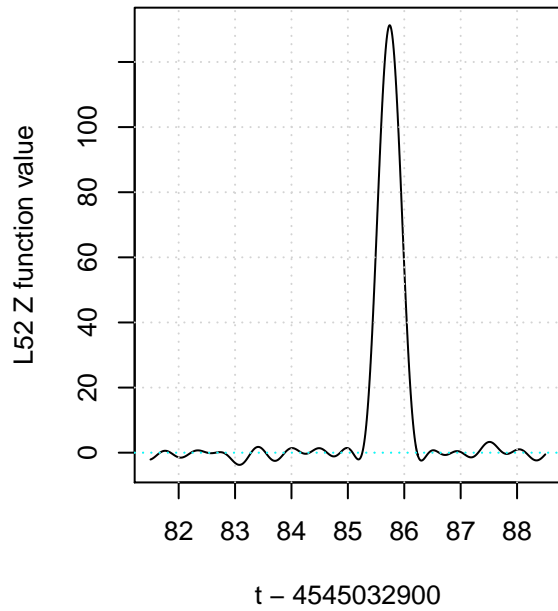
f2(s) Riemann Siegel Z function  
using 2048 pt tapered zeroth order Riemann Siegel formula  
along critical line  $s=0.5+it$  in the interval  $t=(4,545,032,981.5, 4,545,032,988.5)$



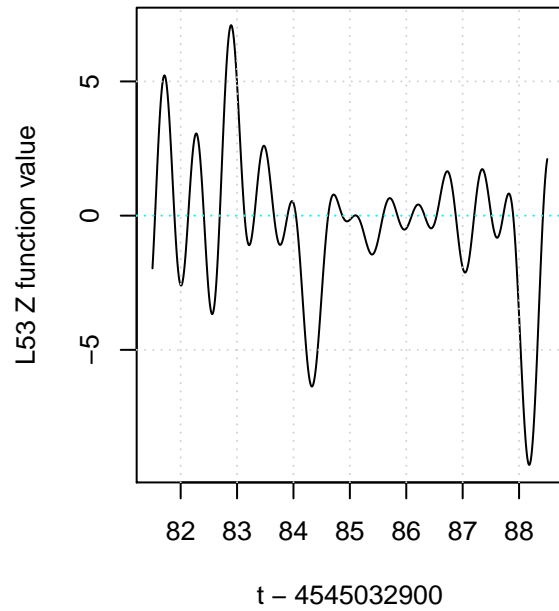
f1(s) Riemann Siegel Z function  
using 2048 pt tapered zeroth order Riemann Siegel formula  
along critical line  $s=0.5+it$  in the interval  $t=(4,545,032,981.5, 4,545,032,988.5)$



L52 Riemann Siegel Z function  
using 2048 pt tapered zeroth order Riemann Siegel formula  
along critical line  $s=0.5+it$  in the interval  $t=(4,545,032,981.5, 4,545,032,988.5)$



L53 Riemann Siegel Z function  
using 2048 pt tapered zeroth order Riemann Siegel formula  
along critical line  $s=0.5+it$  in the interval  $t=(4,545,032,981.5, 4,545,032,988.5)$



## Results - Examples of non-trivial zero behaviour under perturbation of tapered zeroth order Riemann Siegel formula, approximating 5-periodic Davenport-Heilbronn function and $L(\chi_5(2, \cdot), s)$ and $L(\chi_5(3, \cdot), s)$ function behaviour, away from the real axis.

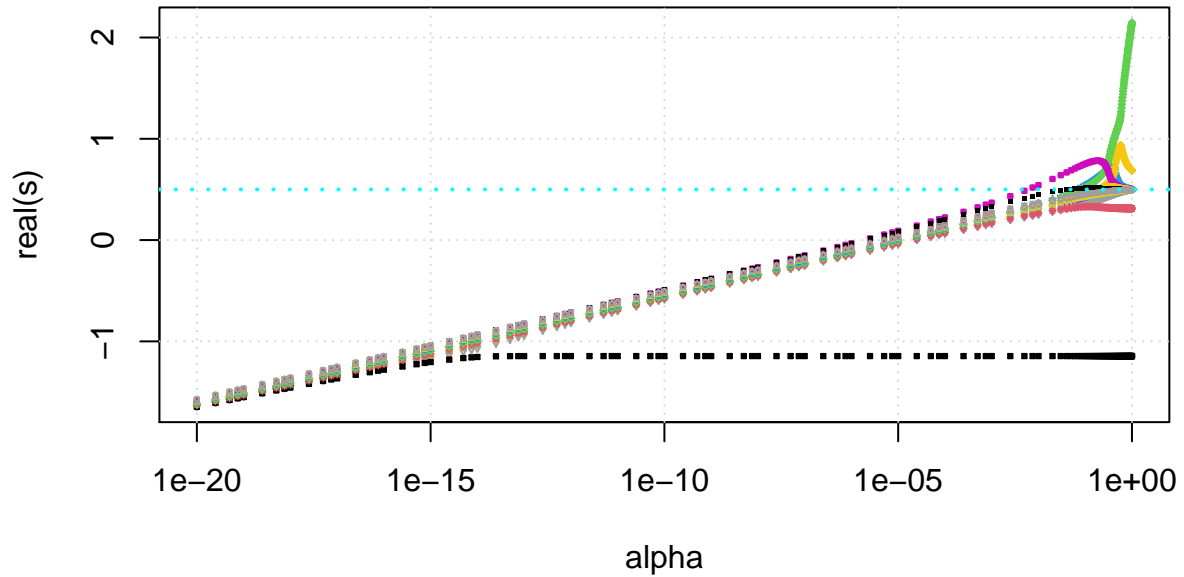
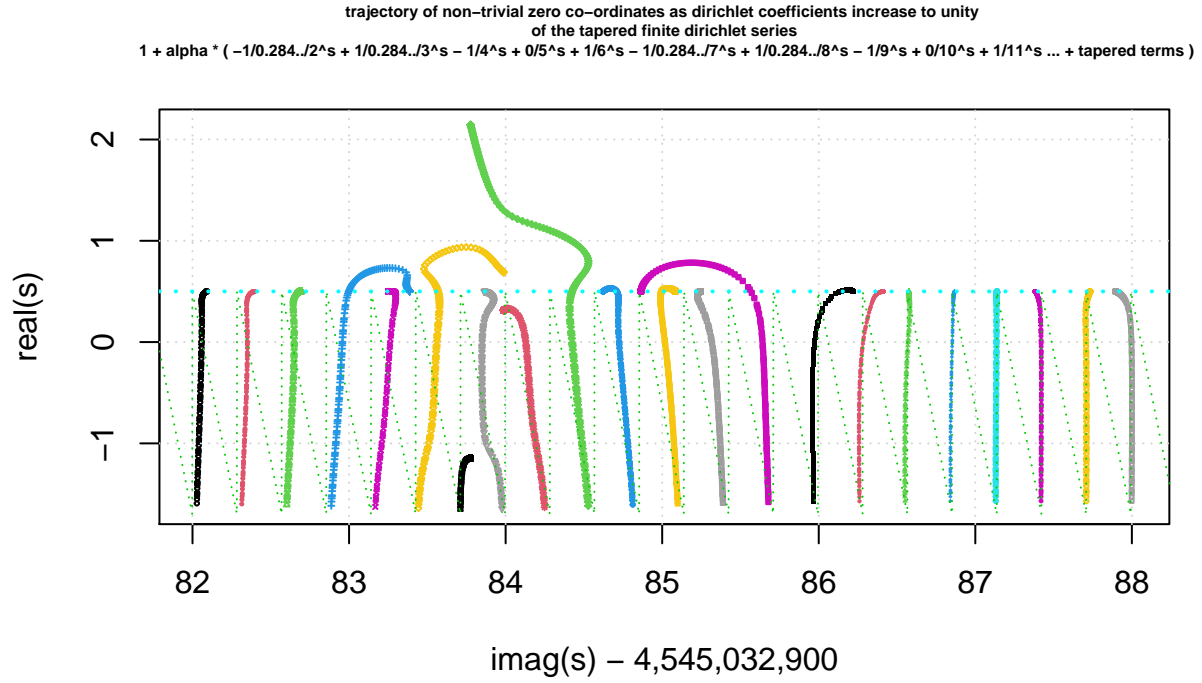
All the calculations of non-trivial zero locations for the four tapered finite zeroth order Riemann Siegel functions at the first quiescent region were performed using the pari-gp language [12] as a solution to second order Taylor series in  $\text{real}(s)$  and  $\text{imag}(s)$  that produces iterative fourth order polynomials for  $\text{imag}(t)$  and then  $\text{real}(s)$  respectively. The R language [13] and R-studio IDE [14] were used to piece the pari-gp based results together and produce graphs.

Figures 2, 4, 6, 8 display examples of the trajectory of the perturbed location of non-trivial zeroes for  $f_2(s)$ ,  $f_1(s)$ ,  $L(\chi_5(2, \cdot), s)$  and  $L(\chi_5(3, \cdot), s)$  respectively, near the  $L(\chi_5(2, \cdot), s)$  function peaks at  $t=\{45450322985.7\}$  which has peak height  $Z=\{130\}$ .

Figures 2, 4, 6, 8 have two panels,

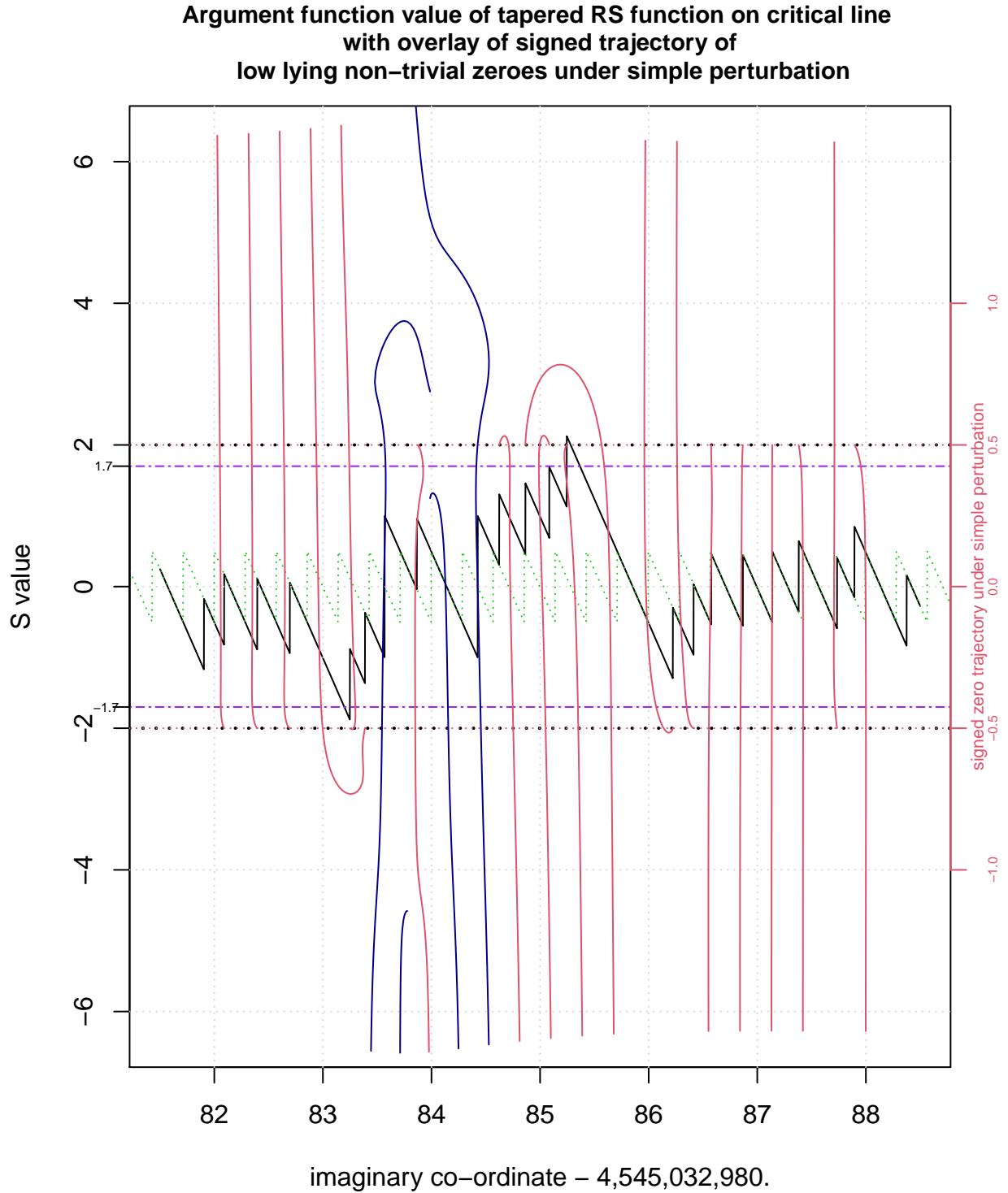
- The upper panel displays the  $\text{real}(s)$  versus  $\text{imag}(s)$  co-ordinate trajectory of nearby non-trivial zeroes as the perturbation varies from  $1e-20 < \alpha < 1$  for  $f_2(s)$  and  $1e-10 < \alpha < 1$  for  $f_1(s)$ ,  $L(\chi_5(2, \cdot), s)$  and  $L(\chi_5(3, \cdot), s)$  where  $\alpha = 1$  represent zero perturbation of the original function. As a guide on the upper panel are vertical lines indicating the expected imaginary co-ordinate of the zeroes if the Gram's law equivalent for these functions were perfectly obeyed. Under high perturbation  $\alpha \rightarrow 0$  when the Riemann Zeta function contribution is a heavily reduced the imaginary component of the non-trivial zeroes generally head towards these vertical line co-ordinates. For the  $f_2(s)$  function with a zero at  $s=-1.145206...+i*4545032983.776683...$   $\alpha < 1e-15$  before the perturbed non-trivial zero is asymptotically heading to its Gram's law equivalent.
- The lower panel displays the  $\text{real}(s)$  versus  $\alpha$  co-ordinate trajectory of nearby non-trivial zeroes as the perturbation varies from  $1e-20 < \alpha < 1$  for  $f_2(s)$  and  $1e-10 < \alpha < 1$  for  $f_1(s)$ ,  $L(\chi_5(2, \cdot), s)$ .

Figures 3, 5, 7, 9 display the  $S(1/2+it)$  values of the 2048 point tapered finite Riemann Siegel formula for  $f_2(s)$ ,  $f_1(s)$ ,  $L(\chi_5(2, \cdot), s)$  and  $L(\chi_5(3, \cdot), s)$  respectively, near the  $L(\chi_5(2, \cdot), s)$  function peaks at  $t=\{45450322985.7\}$ . Overlayed on these figures is the signed trajectory of the perturbed location of non-trivial zeroes (for  $1e-20 < \alpha < 1$  for  $f_2(s)$  and  $1e-10 < \alpha < 1$  for  $f_1(s)$ ,  $L(\chi_5(2, \cdot), s)$  and  $L(\chi_5(3, \cdot), s)$ ) in order to see if graphically there is any consistent behaviour between the  $S(1/2+it)$  values and the non-trivial zero perturbation trajectories. A signed trajectory just means that if a  $S(1/2+it)$  discontinuity has positive (negative) value then the associated non-trivial zero trajectory has a positive (negative) sign assigned (using a  $\pm 1$  multiplicative factor). This overlay helps visualise graphical evidence of non-trivial zero perturbation trajectory overshoot behaviour when  $|S(1/2+it)| > 1.7$ . This threshold of 1.7 which was observed in Riemann Zeta function examples (also based on tapered zeroth order Riemann Siegel formula calculations) did NOT occur for the  $L(\chi_5(2, \cdot), s)$  case in figure 7.

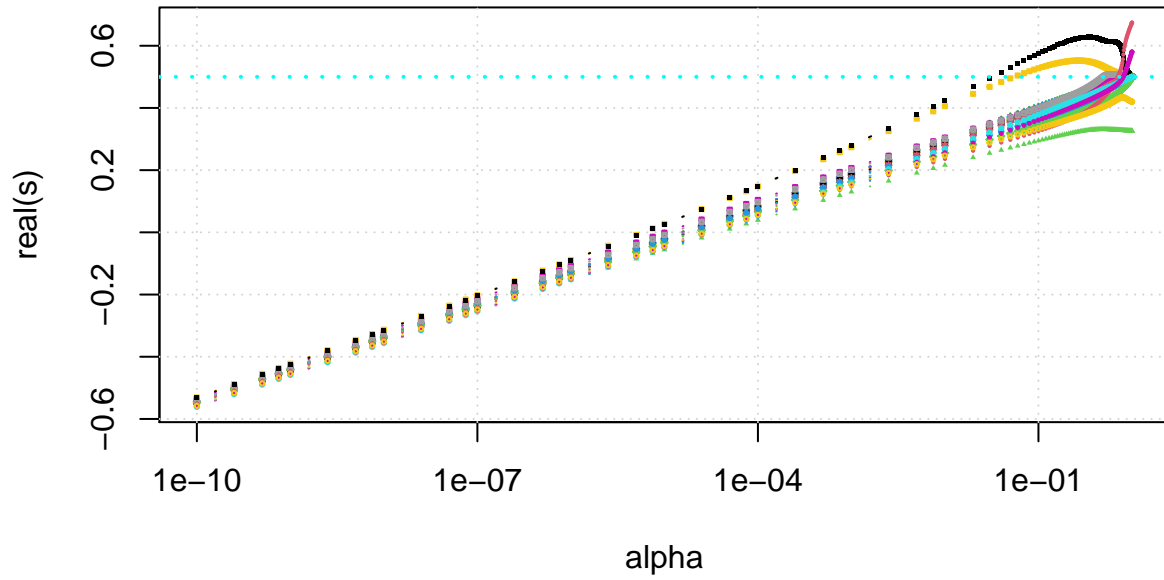
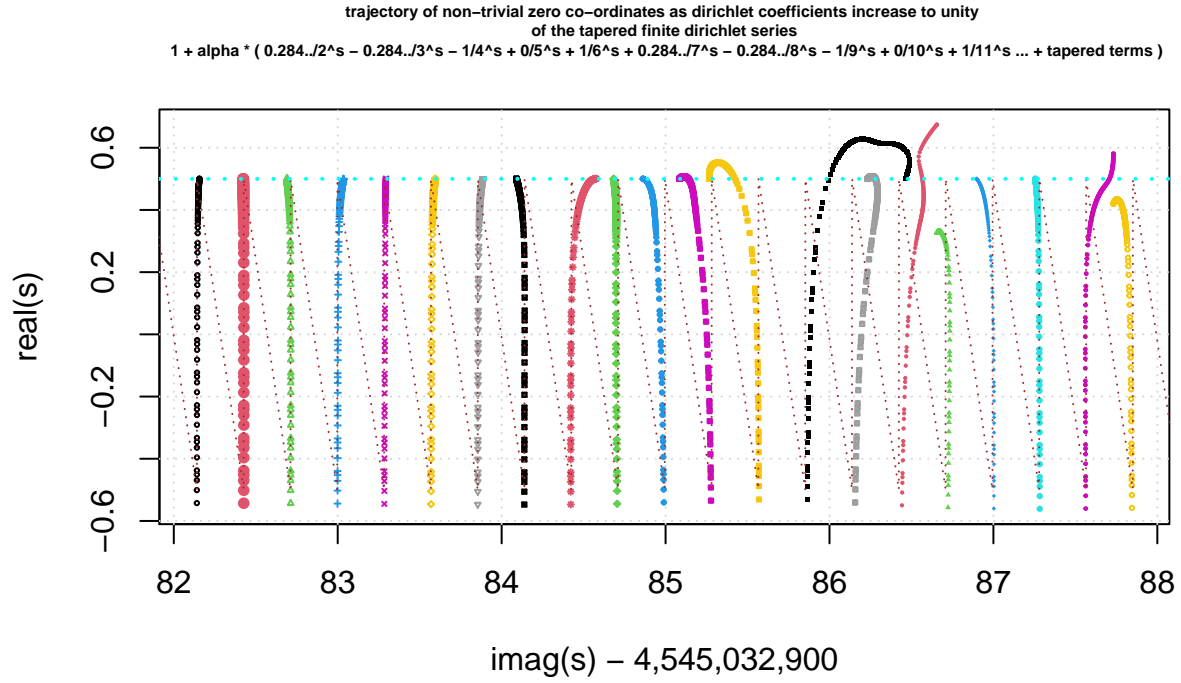


*Figure 2. The trajectory of (twenty two) non-trivial zero co-ordinates around the  $f_2(s)$  function critical line high peak ( $t=4545032985.7\dots$ ) as the magnitude ( $\alpha$ ) of the 2nd, 3rd, 4th, ... etc dirichlet coefficients of the tapered finite dirichlet series  $= 1 + \alpha * ( -1/0.284../2^s + 1/0.284../3^s - 1/4^s + 0/5^s + 1/6^s - 1/0.284../7^s + 1/0.284../8^s - 1/9^s + 0/10^s + 1/11^s \dots + \text{tapered terms} )$  increases to unity.*

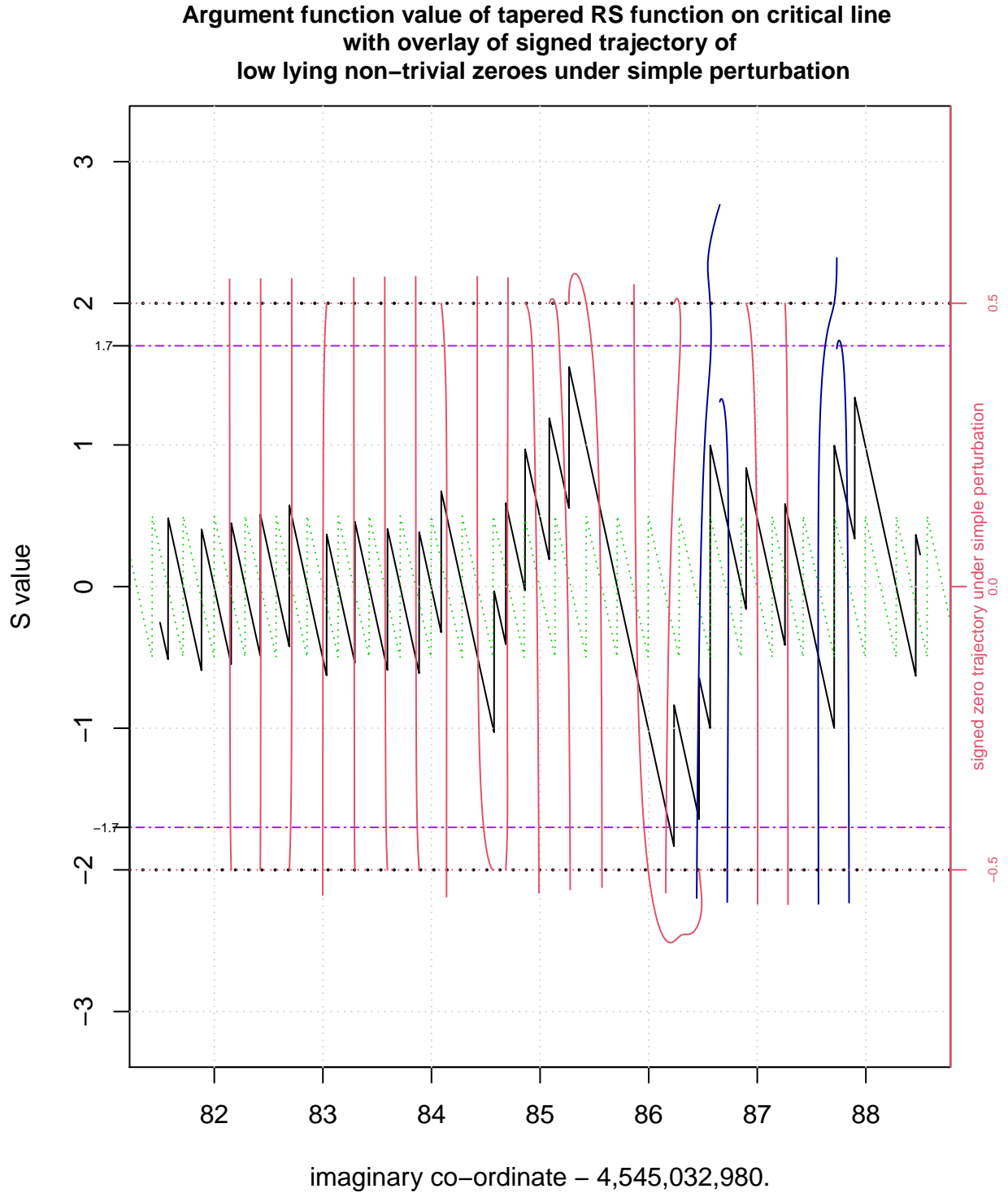




*Figure 3. Comparing signed trajectory of 2048 tapered zeroth order Riemann Siegel function for  $f_2(s)$  under simple perturbation (red lines for critical line non-trivial zeroes, blue line for non-trivial zeroes off the critical line), 2048 tapered zeroth order Riemann Siegel function argument function ( $S$  values) when  $s=0.5+it$  and  $f_2(s)$  Gram's law behaviour (green line). Overshooting signed trajectories occur when  $S \gtrsim 1.7$ .\**



*Figure 4. The trajectory of (twenty one) non-trivial zero co-ordinates around the  $f_1(s)$  function critical line high peak ( $t=4545032985.7\dots$ ) as the magnitude ( $\alpha$ ) of the 2nd, 3rd, 4th, ... etc dirichlet coefficients of the tapered finite dirichlet series  $= 1 + \alpha * ( 0.284../2^s - 0.284../3^s - 1/4^s + 0/5^s + 1/6^s + 0.284../7^s - 0.284../8^s - 1/9^s + 0/10^s + 1/11^s \dots + \text{tapered terms} )$  increases to unity.*



*Figure 5. Comparing signed trajectory of 2048 tapered zeroth order Riemann Siegel function for  $f_1(s)$  under simple perturbation (red lines for critical line non-trivial zeroes, blue line for non-trivial zeroes off the critical line), 2048 tapered zeroth order Riemann Siegel function argument function ( $S$  values) when  $s=0.5+it$  and  $f_1(s)$  Gram's law behaviour (green line). Overshooting signed trajectories occur when  $S \gtrsim 1.7$ . \**

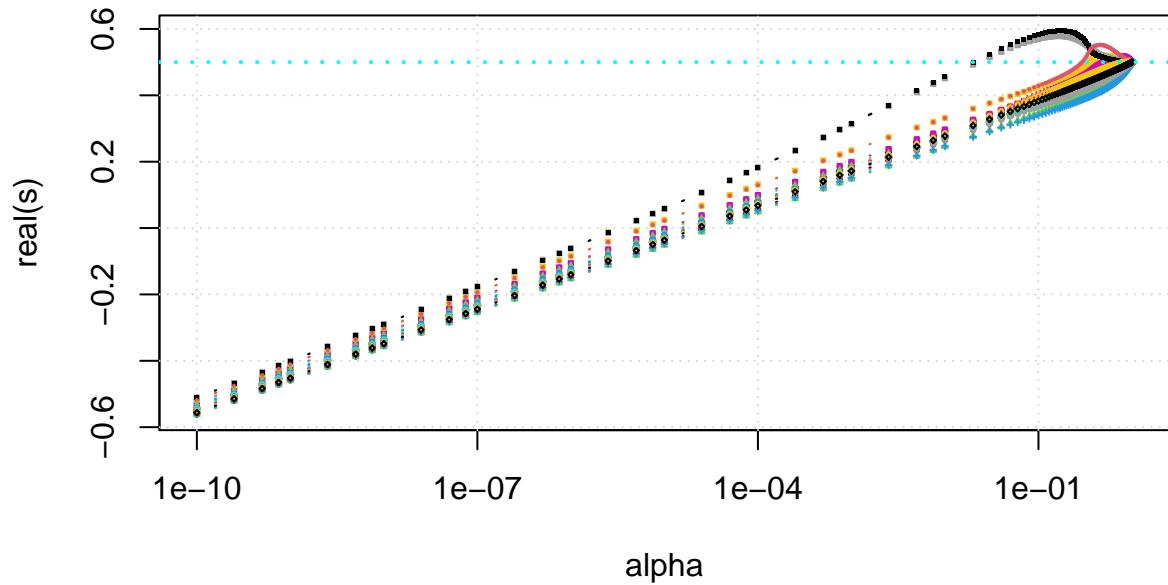
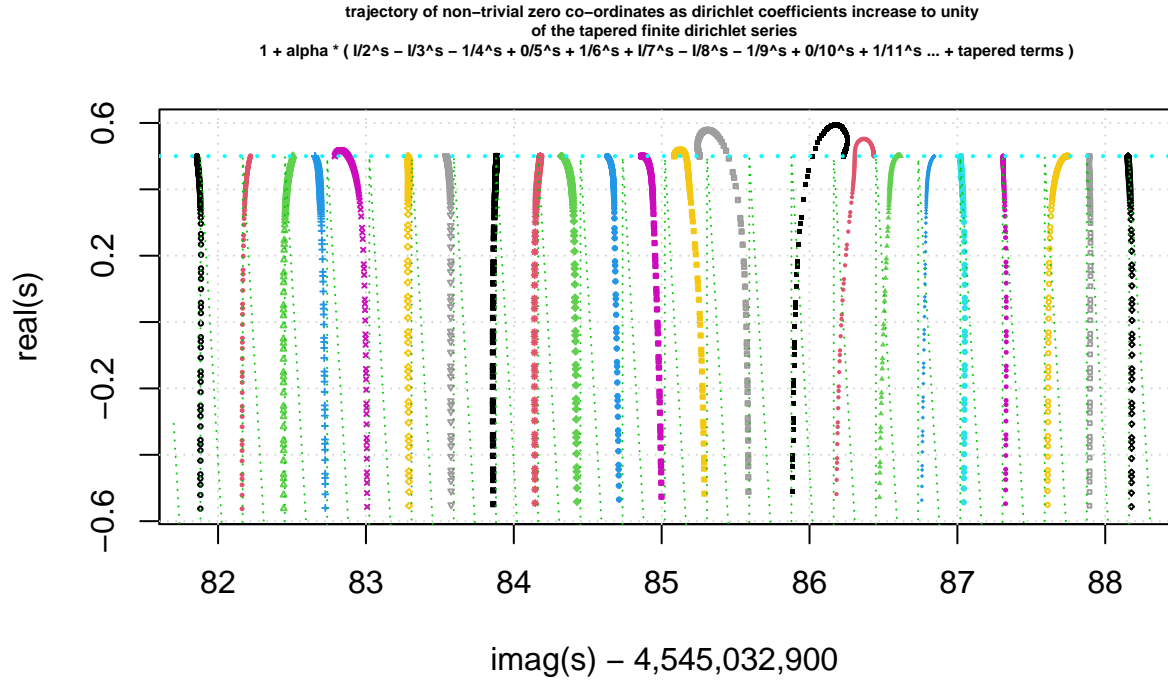
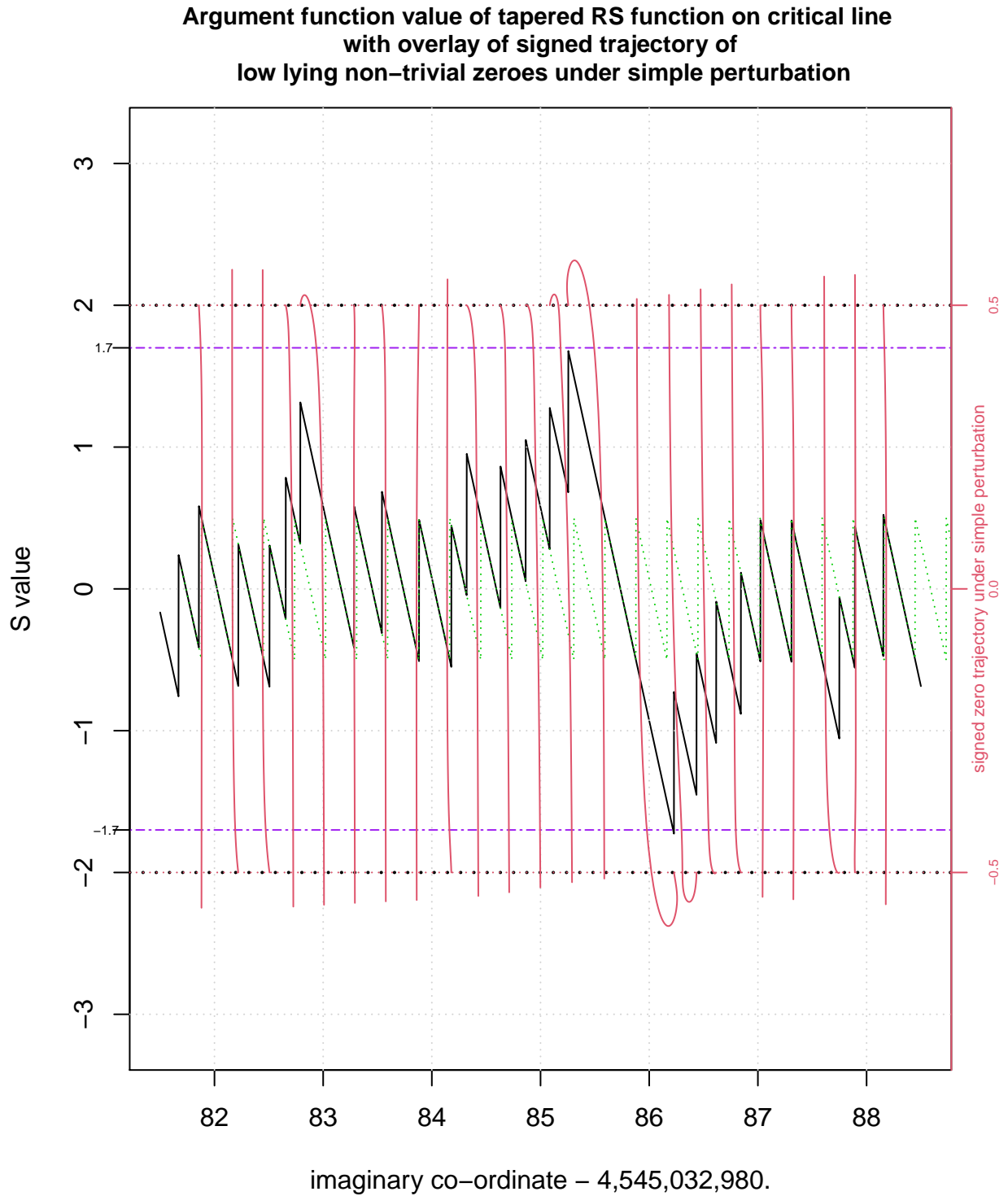
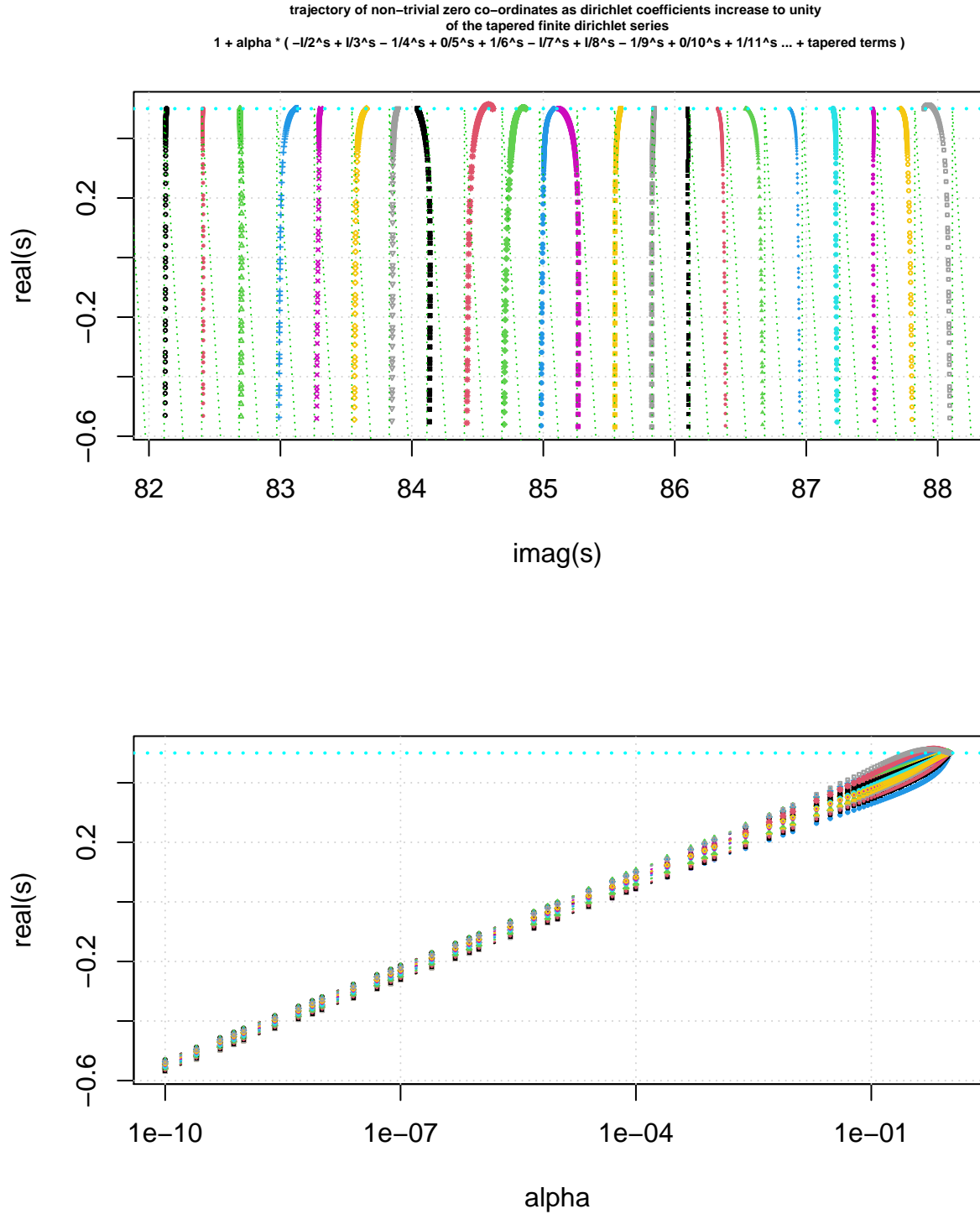


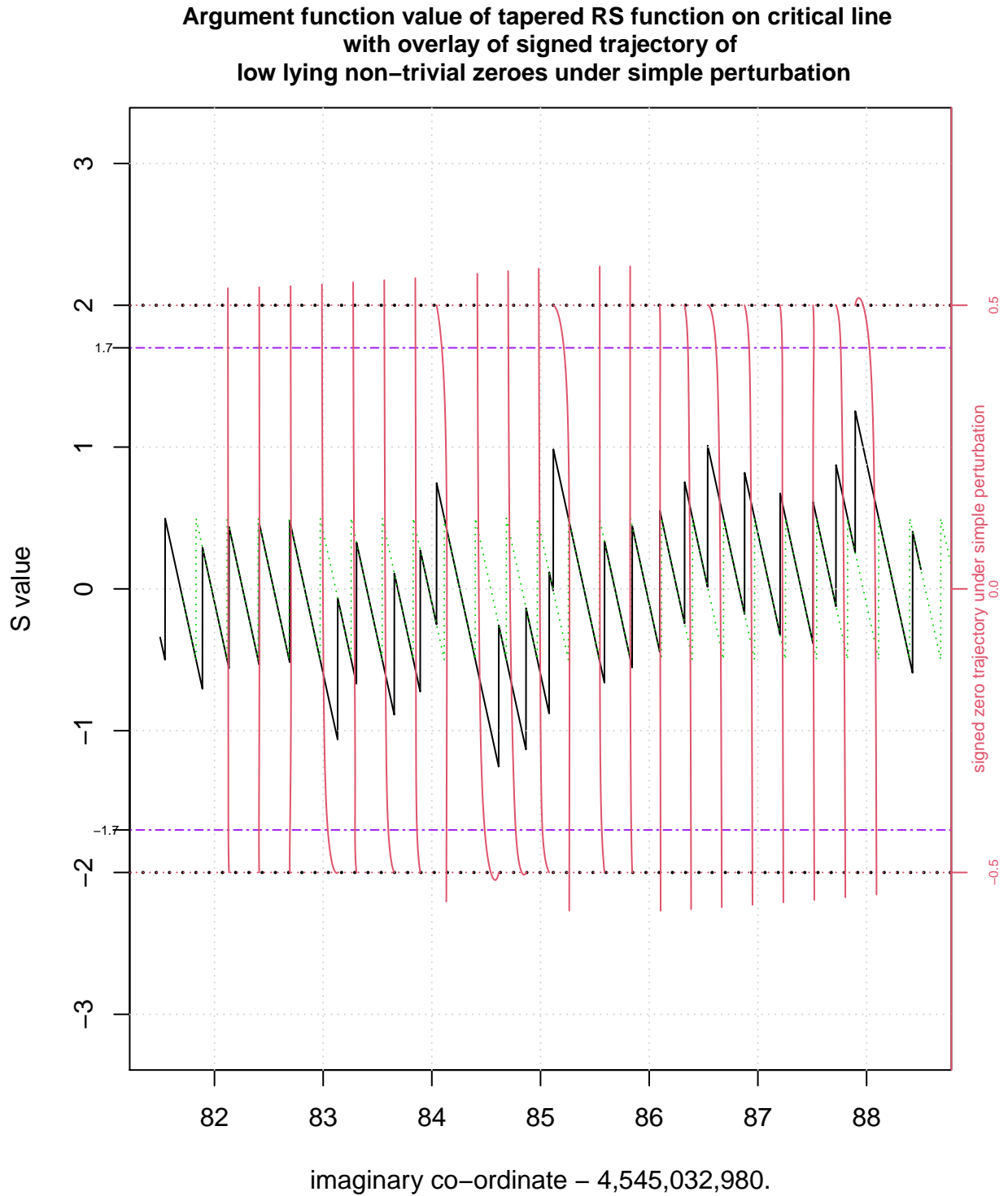
Figure 6. The trajectory of (twenty three) non-trivial zero co-ordinates around the  $f_2(s)$  function critical line high peak ( $t=4545032985.7\dots$ ) as the magnitude ( $\alpha$ ) of the 2nd, 3rd, 4th, ... etc dirichlet coefficients of the tapered finite dirichlet series  $= 1 + \alpha * ( 1/2^s - 1/3^s - 1/4^s + 0/5^s + 1/6^s + 1/7^s - 1/8^s - 1/9^s + 0/10^s + 1/11^s \dots + \text{tapered terms} )$  increases to unity.



*Figure 7. Comparing signed trajectory of 2048 tapered zeroth order Riemann Siegel function for  $f_2(s)$  under simple perturbation (red lines for critical line non-trivial zeroes, blue line for non-trivial zeroes off the critical line), 2048 tapered zeroth order Riemann Siegel function argument function ( $S$  values) when  $s=0.5+it$  and  $L(\chi_5(2, \cdot), s)$  Gram's law behaviour (green line). Overshooting signed trajectories did NOT occur when  $S \gtrsim 1.7$ .\**



*Figure 8. The trajectory of (twenty two) non-trivial zero co-ordinates around the  $f_2(s)$  function critical line high peak ( $t=4545032985.7\dots$ ) as the magnitude ( $\alpha$ ) of the 2nd, 3rd, 4th, ... etc dirichlet coefficients of the tapered finite dirichlet series  $= 1 + \alpha * (-1/2^s + 1/3^s - 1/4^s + 0/5^s + 1/6^s - 1/7^s + 1/8^s - 1/9^s + 0/10^s + 1/11^s \dots + \text{tapered terms})$  increases to unity.*



*Figure 9. Comparing signed trajectory of 2048 tapered zeroth order Riemann Siegel function for  $f_2(s)$  under simple perturbation (red lines for critical line non-trivial zeroes, blue line for non-trivial zeroes off the critical line), 2048 tapered zeroth order Riemann Siegel function argument function ( $S$  values) when  $s=0.5+it$  and  $L(\chi_5(3, \cdot), s)$  Gram's law behaviour (green line).\**

## Conclusions

Perturbing the dirichlet coefficients of the 2048 tapered zeroth order Riemann Siegel function about the first quiescent region of the 5-periodic Davenport-Heilbronn functions and its L function components provides useful insights into the origin and behaviour of the non-trivial zeroes. There discontinuities of magnitude 2 in the S value (along the critical line) appear to be associated with perturbed off critical line non-trivial zeroes when such perturbed non-trivial zeroes intersect the critical line.

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