

From Parsons *et al* 2012:

$$\Delta_N^2(k) \propto (k)^{\frac{5}{2}} \left[\frac{1}{B} \right]^{\frac{1}{2}} \left[\frac{1}{\Delta \ln k} \right]^{\frac{1}{2}} \Omega T_{sys}^2 \left[\frac{1}{t_{per_day}} \right]^{\frac{1}{2}} \left[\frac{1}{t_{cam}} \right] \left[\frac{1}{N_a} \right] \left[\frac{f_o}{f} \right]^{\frac{1}{2}} \quad (1)$$

where k is the magnitude of the k -mode, B is the bandwidth, $\Delta \ln k$ is the log of the binsize, Ω is the field-of-view, T_{sys} is the system temperature, t_{per_day} is the number of hours observed per day, t_{cam} is the number of days observed, N_a is the number of antennas, and f_o/f is a configuration metric for a redundant array as therein defined.

Pulling out terms relating to diameter (D_a) and number, we can write Eq. 1 as

$$\Delta_N^2(k) \propto \frac{\Omega \sqrt{f_o/f}}{N_a \sqrt{t_{per_day}}} \propto \frac{(1/D_a^2)(1/\sqrt{N_a})}{N_a \sqrt{D_a}} = \frac{1}{D_a^{\frac{5}{2}} N_a^{\frac{3}{2}}} \quad (2)$$

where the dependencies on diameter and number have been substituted in, noting that the expressions for f_o/f and t_{per_day} were derived in Parsons et al where the baselines for the close-packed array are multiples of the diameter.

From Mellema *et al* 2013:

$$\Delta_N^2(k) \propto k^{\frac{3}{2}} \sqrt{D_c^2 \Delta D_c \Omega_{FOV}} \left[\frac{T_{sys}}{\sqrt{B t_{int}}} \right]^2 \left[\frac{A_{core} A_{eff}}{A_{coll}^2} \right] \quad (3)$$

Pulling out similar D_a and number terms:

$$\Delta_N^2(k) \propto \sqrt{\Omega} \left[\frac{1}{t_{int}} \right] \left[\frac{A_{core} A_{eff}}{A_{coll}^2} \right] \propto \left[\frac{1}{D_a} \right] \left[\frac{1}{t_{int}} \right] \left[\frac{(N D_a^2)(D_a^2)}{(N D_a^2)^2} \right] = \frac{1}{t_{int}} \frac{1}{N D_a^2} \quad (4)$$

where, for hex packing

$$A_{core} \propto \left[\frac{4N-1}{3} D_a^2 \right] \sim N D_a^2 \quad (5)$$

So, there is a factor of $\sqrt{N D_a}$ different. Allowing for incoherent drift-scanning, $t_{int} \rightarrow \sqrt{t_{int}} \sim 1/\sqrt{D_a}$, so we are left with a \sqrt{N} different, probably related to the specifics of f_o/f and/or $A_{core}/A_{coll}^2 \dots$

Running Jonnie's code for

- a bunch of diameters with 331 antennas (7-21 m)
- a bunch of elements for 14-m diameter (37-919)

and evaluate at $ks = 1$, yields:

$$\Delta_N^2(1) \approx \frac{1}{D^{1.5} N^{1.5}} \quad (6)$$

The exact coefficient is weakly dependent on the k -value in the middle k values. The power laws fit better for larger k 's (and are closer to 1.5).