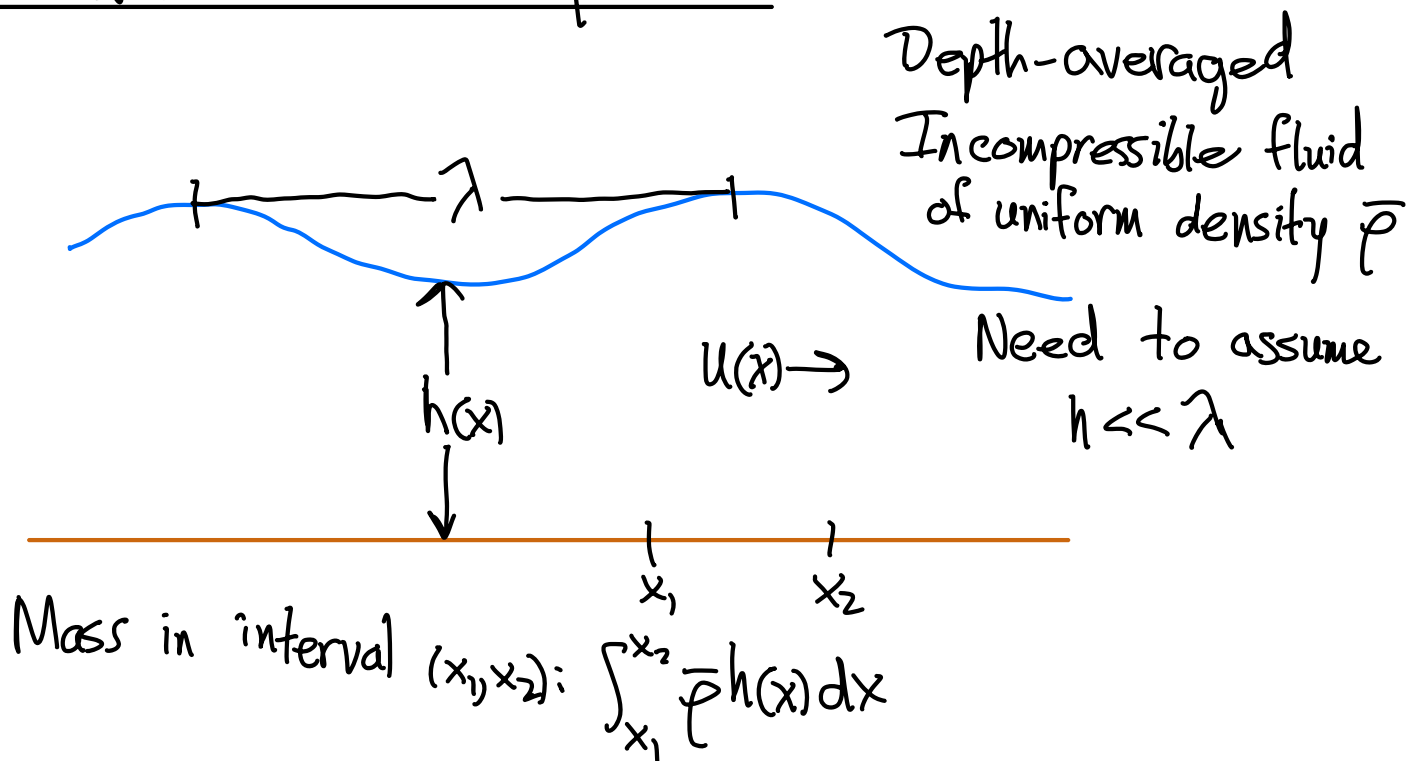


Shallow water equations



Flux: $\bar{\rho} h u$

Conservation of mass: $(\bar{\rho} h)_t + (\bar{\rho} h u)_x = 0 \Rightarrow h_t + (h u)_x = 0$

Conservation of momentum: $(\bar{\rho} h u)_t + (\bar{\rho} h u^2 + p)_x = 0$

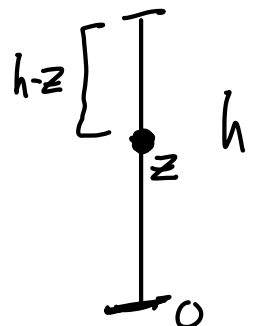
What is the pressure?

Hydrostatic pressure: exactly enough to support weight of water.

$$p(x, z) = g \bar{\rho} (h - z)$$

Total pressure in a column:

$$\begin{aligned} \int_0^h g \bar{\rho} (h - z) dz &= g \bar{\rho} \int_0^h (h - z) dz = g \bar{\rho} \left(h^2 - \frac{h^2}{2} \right) \\ &= \frac{1}{2} g \bar{\rho} h^2 \end{aligned}$$



$$h_t + (hu)_x = 0$$

1D Shallow water model

$$(hu)_t + \left(hu^2 + \frac{1}{2}gh^2\right)_x = 0$$

Characteristic structure

$$q_t + f(q)_x = 0 \Rightarrow \underline{q_t + f'(q)q_x = 0}$$

$$q = \begin{bmatrix} h \\ hu \end{bmatrix} \quad f(q) = \begin{bmatrix} hu \\ hu^2 + \frac{1}{2}gh^2 \end{bmatrix} = \begin{bmatrix} q^2 \\ \frac{(q^2)^2}{q^1} + \frac{1}{2}g(q^1)^2 \end{bmatrix}$$

$$f'(q) = \begin{bmatrix} 0 & 1 \\ gq^1 - \frac{(q^2)^2}{(q^1)^3} & 2\frac{q^2}{q^1} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ gh - u^2 & 2u \end{bmatrix}$$

Characteristic speeds are eigenvalues of $f'(q)$:

$$\det(\lambda I - f'(q)) = \det \begin{bmatrix} \lambda & -1 \\ u^2 - gh & \lambda - 2u \end{bmatrix} = \lambda(\lambda - 2u) + u^2 - gh$$

$$\lambda^2 - 2u\lambda + u^2 - gh = 0$$

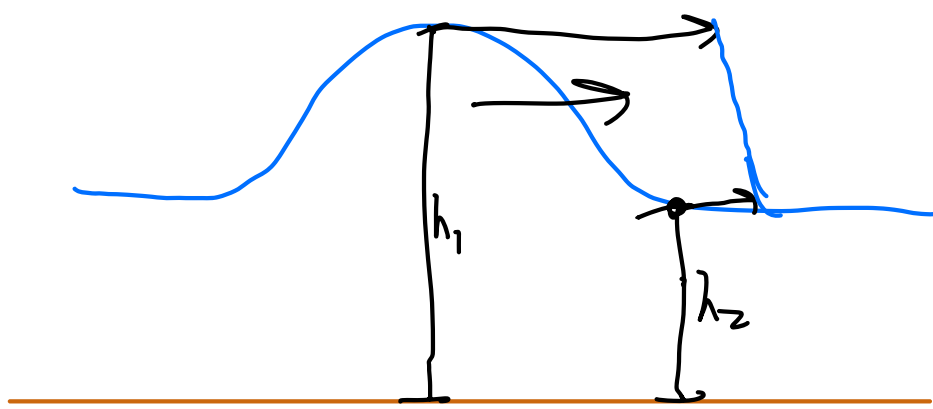
$$\begin{aligned} \lambda^1 &= u - \sqrt{gh} \\ \lambda^2 &= u + \sqrt{gh} \end{aligned}$$

$$r^1 = \begin{bmatrix} 1 \\ u - \sqrt{gh} \end{bmatrix}$$

$$r^2 = \begin{bmatrix} 1 \\ u + \sqrt{gh} \end{bmatrix}$$

Solution is not constant along characteristics, and they are not straight lines (in general)

Why do waves break in shallow water?



$$h_1 > h_2$$

$$u + \sqrt{gh_1} > u + \sqrt{gh_2}$$

In deep water

$$\frac{h_1 - h_2}{h_1} \text{ is small.}$$

Small perturbations move at the characteristic speeds.

$$q(x,t) = q_0 + \varepsilon \hat{q}(x,t) \quad \varepsilon \ll 1$$

$$q_t + f'(q)q_x = (q_0 + \varepsilon \hat{q})_t + f'(q_0 + \varepsilon \hat{q})(q_0 + \varepsilon \hat{q})_x = 0$$

$$\varepsilon \hat{q}_t + \varepsilon f'(q_0) \hat{q}_x = \mathcal{O}(\varepsilon^2)$$

$$\hat{q}_t + f'(q_0) \hat{q}_x = \mathcal{O}(\varepsilon)$$

Now suppose $\hat{q}_x = \alpha(x) r'(q_0)$. Then

$$\hat{q}_t + f'(q_0) \alpha(x) r'(q_0) = \hat{q}_t + \alpha(x) \lambda'(q_0) r'(q_0) = 0.$$

R.H. Jump Condition

Let (h_*, u_*) be fixed and suppose we have a shock with (h, u) being the other state.

$$s(q - q_*) = f(q) - f(q_*)$$

$$\alpha = h - h_*$$

$$\begin{aligned} \rightarrow s(h - h_*) &= hu - h_* u_* \\ s(hu - h_* u_*) &= hu^2 - h_* u_*^2 + \frac{1}{2}g(h^2 - h_*^2) \end{aligned} \left. \begin{array}{l} 2 \text{ Eqns.} \\ 3 \text{ unknowns: } s, h, u. \end{array} \right\}$$

$$\text{Result: } s = \frac{hu - h_* u_*}{\alpha} = u_* \pm \sqrt{gh \frac{h + h_*}{2h_*}} \leftarrow$$

$$hu = h_* u_* + \alpha \left[u_* \pm \sqrt{gh_* \left(1 + \frac{\alpha}{h_*}\right) \left(1 + \frac{\alpha}{2h_*}\right)} \right]$$

$$h = h_* + \alpha$$

Observe: $S^P \rightarrow \mathcal{X}^P$ as $\alpha \rightarrow 0$