"Shocking" Wave Phenomena in Periodic Materials

David I. Ketcheson

Nov. 4, 2010



D. Ketcheson () Nov. 4, 2010 1 / 34

Outline

- Nonlinear Waves in Periodic Media
- Quantification of the contraction of the contrac
- 3 Computational Investigation of Shock Formation

D. Ketcheson () Nov. 4, 2010 2 / 34

Elasticity in 1D

$$\epsilon_t - u_x = 0$$

$$\rho u_t - \sigma(\epsilon)_x = 0$$

D. Ketcheson () Nov. 4, 2010

Elasticity in 1D

$$\epsilon_t - u_x = 0$$

$$\rho u_t - \sigma(\epsilon)_x = 0$$

Strain: $\epsilon(x,t)$ Velocity: u(x,t)

Stress: $\sigma(\epsilon)$ Density: ρ

Elasticity in 1D

$$\epsilon_t - u_x = 0$$

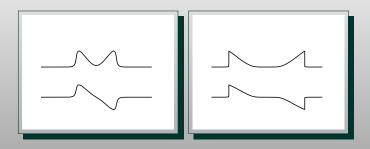
$$\rho u_t - \sigma(\epsilon)_x = 0$$

$$\begin{array}{ll} \text{Strain: } \epsilon(x,t) & \text{Velocity: } u(x,t) \\ \text{Stress: } \sigma(\epsilon) & \text{Density: } \rho \end{array}$$

$$\begin{array}{ccc} \sigma(\epsilon) = K\epsilon \implies \text{linear waves} \\ \sigma(\epsilon) = e^{K\epsilon} - 1 \implies \text{nonlinear waves} \end{array}$$

D. Ketcheson () Nov. 4, 2010 3 / 34

Time-Reversibility of Nonlinear Waves



- Waves are time-reversible as long as they remain smooth
- Time-reversibility is lost after shocks form

D. Ketcheson () Nov. 4, 2010

$$\eta(u,\epsilon) = \frac{1}{2}\rho u^2 + \int_0^\epsilon \sigma(s)ds.$$

Total entropy:

$$\int \eta dx$$
.

Conserved for smooth solutions

$$\eta(u,\epsilon) = \frac{1}{2}\rho u^2 + \int_0^\epsilon \sigma(s)ds.$$

Total entropy:

$$\int \eta dx$$
.

- Conserved for smooth solutions
- Decreases when shocks form

$$\eta(u,\epsilon) = \frac{1}{2}\rho u^2 + \int_0^\epsilon \sigma(s)ds.$$

Total entropy:

$$\int \eta dx$$
.

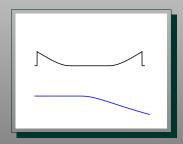
- Conserved for smooth solutions
- Decreases when shocks form

$$\eta(u,\epsilon) = \frac{1}{2}\rho u^2 + \int_0^\epsilon \sigma(s)ds.$$

Total entropy:

$$\int \eta dx$$
.

- Conserved for smooth solutions
- Decreases when shocks form



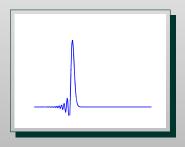
Dispersive Waves

$$\epsilon_t - u_x = \epsilon u_{xxx}$$
 $u_t - \sigma(\epsilon)_x = 0$

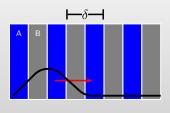
D. Ketcheson () Nov. 4, 2010 6 / 34

Dispersive Waves

$$\epsilon_t - u_x = \epsilon u_{xxx}$$
 $u_t - \sigma(\epsilon)_x = 0$



Linear Elasticity in Periodic Media

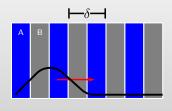


$$\epsilon_t - u_x = 0$$

$$\rho(x)u_t - K(x)\epsilon_x = 0$$

D. Ketcheson () Nov. 4, 2010 7 / 3

Linear Elasticity in Periodic Media



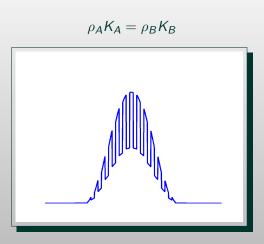
$$\epsilon_t - u_x = 0$$

$$\rho(x)u_t - K(x)\epsilon_x = 0$$

K(x) and $\rho(x)$ are piecewise-constant periodic functions Impedance: $Z(x) = \sqrt{\rho(x)K(x)}$

D. Ketcheson () Nov. 4, 2010 7 / 3

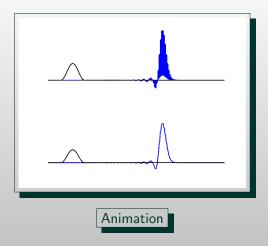
Impedance Matched Materials



In this case there are no reflections.

D. Ketcheson () Nov. 4, 2010

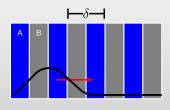
Impedance mismatch: Effective Dispersion



- Due to reflections, the pulse travels more slowly than it would in either component material
- The medium introduces effective dispersive behavior

D. Ketcheson () Nov. 4, 2010 9 / 34

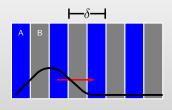
Nonlinear Elasticity in Periodic Media



$$\begin{aligned}
\epsilon_t - u_x &= 0 \\
\rho(x)u_t - \sigma(\epsilon, x)_x &= 0
\end{aligned}$$

D. Ketcheson () Nov. 4, 2010 10 / 34

Nonlinear Elasticity in Periodic Media



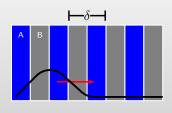
$$\epsilon_t - u_x = 0$$

$$\rho(x)u_t - \sigma(\epsilon, x)_x = 0$$

Nonlinear stress relation: $\sigma(\epsilon, x) = \exp(K(x)\epsilon) - 1$

D. Ketcheson () Nov. 4, 2010 10 / 34

Nonlinear Elasticity in Periodic Media



$$\epsilon_t - u_x = 0$$

$$\rho(x)u_t - \sigma(\epsilon, x)_x = 0$$

Nonlinear stress relation: $\sigma(\epsilon, x) = \exp(K(x)\epsilon) - 1$

Impedance: $Z(\epsilon, x) = \sqrt{\rho(x)\sigma_{\epsilon}(\epsilon, x)}$

D. Ketcheson () Nov. 4, 2010 10 / 34

Impedance:

$$Z(\epsilon, x) = \sqrt{\rho(x)\sigma_{\epsilon}(\epsilon, x)}$$

Impedance:

$$Z(\epsilon, x) = \sqrt{\rho(x)\sigma_{\epsilon}(\epsilon, x)}$$

Linearized impedance:

$$Z_0(x) = \sqrt{\rho(x)\sigma_{\epsilon}(0,x)}$$

D. Ketcheson () Nov. 4, 2010 11 /

Impedance:

$$Z(\epsilon, x) = \sqrt{\rho(x)\sigma_{\epsilon}(\epsilon, x)}$$

Linearized impedance:

$$Z_0(x) = \sqrt{\rho(x)\sigma_{\epsilon}(0,x)}$$

Linearized impedance contrast:

$$Z_{0,A}/Z_{0,B}$$

D. Ketcheson () Nov. 4, 2010

Impedance:

$$Z(\epsilon, x) = \sqrt{\rho(x)\sigma_{\epsilon}(\epsilon, x)}$$

Linearized impedance:

$$Z_0(x) = \sqrt{\rho(x)\sigma_{\epsilon}(0,x)}$$

Linearized impedance contrast:

$$Z_{0,A}/Z_{0,B}$$

$$Z_{0,A}/Z_{0,B}\approx 2$$

Animation

Impedance:

$$Z(\epsilon, x) = \sqrt{\rho(x)\sigma_{\epsilon}(\epsilon, x)}$$

Linearized impedance:

$$Z_0(x) = \sqrt{\rho(x)\sigma_{\epsilon}(0,x)}$$

Linearized impedance contrast:

$$Z_{0,A}/Z_{0,B}$$

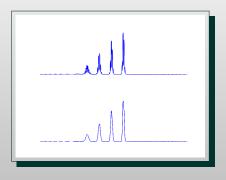
$$Z_{0,A}/Z_{0,B}\approx 2$$

Animation

Nonlinearity dominates → shocks form

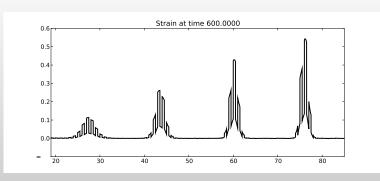
Higher Impedance Contrast





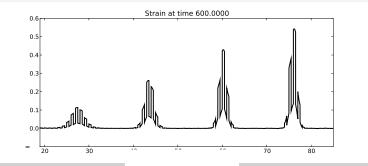
Close-up

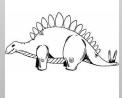
Stegotons



D. Ketcheson () Nov. 4, 2010 13 / 34

Stegotons





LeVeque & Yong, 2003

D. Ketcheson () Nov. 4, 2010 13 / 34

Soliton-like Properties

- All stegotons have same shape (in time-profile) under a rescaling
- Stegotons appear to interact only through a phase-shift



D. Ketcheson () Nov. 4, 2010 14 / 34

Outline

- Nonlinear Waves in Periodic Media
- Godunov-Type Methods and Limiters
- Computational Investigation of Shock Formation
- Generalizations

D. Ketcheson () Nov. 4, 2010 15 / 34

Conservation Laws

Defining

$$\mathbf{q} = \begin{bmatrix} \epsilon \\ \rho u \end{bmatrix}$$

and

$$f(\mathbf{q},x) = \begin{bmatrix} -u \\ -\sigma(\epsilon,x) \end{bmatrix}$$

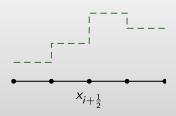
we can write the elasticity equations as a pair of conservation laws:

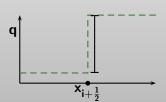
$$\mathbf{q}_t + f(\mathbf{q}, x)_x = 0.$$

Godunov's Method

- Solution represented by cell averages Q_i
- Each cell interface poses a Riemann problem
- Riemann problem at $x_{i-\frac{1}{2}}$ is solved to find flux $F_{i-\frac{1}{2}}$
- Flux-differencing:

$$Q_i^{n+1} = Q_i^n - \frac{\Delta t}{\Delta x} \left(F_{i+\frac{1}{2}} - F_{i-\frac{1}{2}} \right)$$

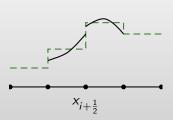


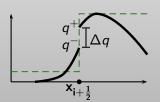


D. Ketcheson () Nov. 4, 2010 17 / 34

High order wave propagation

Piecewise-polynomial Reconstruction





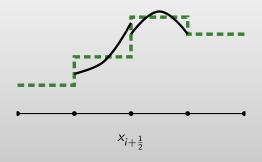
$$\frac{\partial Q_i}{\partial t} = -\frac{1}{\Delta x} \left(F_{i+\frac{1}{2}} - F_{i-\frac{1}{2}} \right)$$

- High order reconstruction in space
- High order ODE solver in time
- The fluxes are determined from the reconstructed states at each interface.
- Reconstruction must be done in a non-oscillatory manner

18 / 34

D. Ketcheson () Nov. 4, 2010

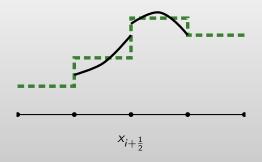
Non-oscillatory Reconstruction



- Limiters are employed to avoid overshoots
- Limiters try to maintain "true" discontinuities while avoiding "spurious" oscillations
- Difficult to achieve the right balance between compression and dissipation

D. Ketcheson () Nov. 4, 2010 19 / 34

Non-oscillatory Reconstruction



- Limiters are employed to avoid overshoots
- Limiters try to maintain "true" discontinuities while avoiding "spurious" oscillations
- Difficult to achieve the right balance between compression and dissipation

D. Ketcheson () Nov. 4, 2010 19 / 34

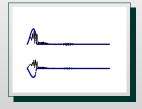
Outline

- Nonlinear Waves in Periodic Media
- Godunov-Type Methods and Limiters
- Computational Investigation of Shock Formation
- Generalizations

D. Ketcheson () Nov. 4, 2010 20 / 34

Time-reversibility of Stegotons

$$Z_{0,A}/Z_{0,B}=2$$

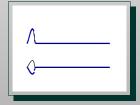


Time-reversibility of Stegotons

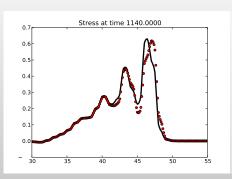
$$Z_{0,A}/Z_{0,B}=2$$

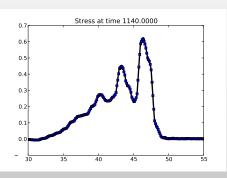


$$Z_{0,A}/Z_{0,B}=4$$



Time-reversibility





This is a useful test for high order numerical methods:

- Nonlinear, yet solution remains smooth for long times
- Known exact solution (apparently)
- Highlights ability of schemes to deal with spatially varying flux

D. Ketcheson () Nov. 4, 2010 22 / 34

Entropy Evolution

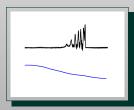
$$\eta(u,\epsilon) = \frac{1}{2}\rho u^2 + \int_0^\epsilon \sigma(s)ds.$$

Total entropy:

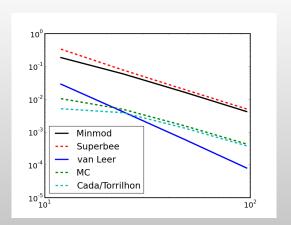
$$\int \eta dx$$
.

- Conserved for smooth solutions
- Decreases when shocks are present
- Solutions are time-reversible until shocks form

How does the entropy behave for nonlinear waves in periodic media?



Entropy and Limiters

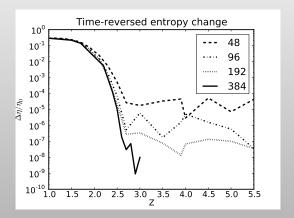


- Entropy change seems to be converging to zero with Δx
- Minmod, van Leer are dissipative (entropy decreases)
- Superbee, monotonized centered, and Cada/Torrilhon are compressive (entropy increases)

D. Ketcheson () Nov. 4, 2010 24 / 34

Probing the "phase transition"

- Entropy loss in high-order WENO simulations is even much smaller
- Numerical entropy errors are nearly time-reversible
- This allows us to probe shock formation very precisely



D. Ketcheson () Nov. 4, 2010 25 / 34

Outline

- Nonlinear Waves in Periodic Media
- Godunov-Type Methods and Limiters
- Computational Investigation of Shock Formation
- 4 Generalizations

D. Ketcheson () Nov. 4, 2010 26 / 34

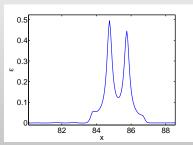
Generalizations

How general is this phenomenon? Is it generic for

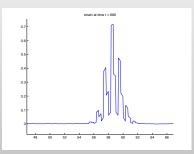
- Different nonlinearities?
- Different materials?
- More general hyperbolic systems?

D. Ketcheson () Nov. 4, 2010 27 / 34

Different Media



Sinusoidally varying medium



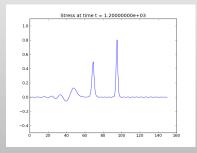
Three-layer medium

D. Ketcheson () Nov. 4, 2010 28 / 34

Different Nonlinearities

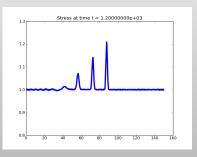
Cubic nonlinearity:

$$\sigma(\epsilon, x) = K_1(x)\epsilon + K_3(x)\epsilon^3$$



Isothermal gas:

$$\sigma(\epsilon, x) = -\kappa(x)/\epsilon^2$$
$$\rho(x) = 1$$

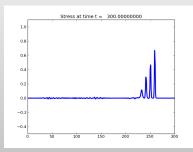


29 / 34

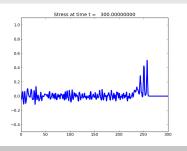
D. Ketcheson () Nov. 4, 2010

Randomly perturbed media

Near-periodic medium with randomly perturbed interface locations



Perturbed by 1%



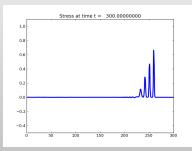
Perturbed by 10%

30 / 34

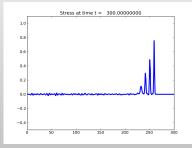
D. Ketcheson () Nov. 4, 2010

Randomly perturbed media

Near-periodic medium with randomly perturbed piecewise constant coefficients



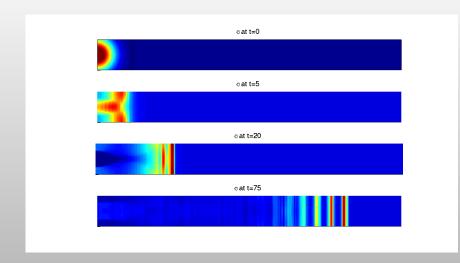
Perturbed by 1%



Perturbed by 25%

D. Ketcheson () Nov. 4, 2010

Stability in higher dimensions



D. Ketcheson () Nov. 4, 2010 32 / 34

Summary of Main Points

- Entropy evolution and time-reversibility allow us to measure the formation of shocks in nonlinear wave simulations
- Periodic materials with high impedance contrast appear to suppress the formation of shocks
- Large-amplitude signals in such materials can travel long distances without information loss
- Studying this phenomenon can also tell us new things about our numerical methods

D. Ketcheson () Nov. 4, 2010 33 / 34

Experimental Realization: Optical Stegotons?

- Medium consists of alternating layers of air and some material with high index of refraction and low loss
- High intensity laser pulse
- Currently we are designing simulations to test feasible design parameters

Joint work with Boon Ooi (KAUST).

D. Ketcheson () Nov. 4, 2010 34 / 34