Molecular Dynamics Simulation of Argon

The fundamental work on this problem was done by A. Rahman, Phys. Rev. **136**, A405 (1964). It was extended in many important ways by L. Verlet, Phys. Rev. **159**, 98 (1967), who introduced the Verlet algorithm and the use of a neighbor list to speed up the calculation.

Simple model of interacting Argon atoms

Consider N atoms of argon each with mass $m = 6.69 \times 10^{-26}$ kg. Argon is an *inert gas*: argons atoms behave approximately like hard spheres which attract one another with weak van der Waals forces. The forces between two argon atoms can be approximated quite well by a Lennard-Jones potential energy function:

$$V(r) = 4\varepsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^{6} \right] ,$$

where r is the distance between the centers of the two atoms, $\varepsilon = 1.65 \times 10^{-21}$ J is the strength of the potential energy, and $\sigma = 3.4 \times 10^{-10}$ m is the value of r at which the energy is zero. The $1/r^{12}$ term represents a repulsive *hard core* interaction between the argon atoms. The $1/r^{6}$ term represents an attractive dipole-dipole (van der Waals) interaction between the non-polar atoms. The potential has its minimum $V(2^{\frac{1}{6}}\sigma) = -\varepsilon$ at $r = 2^{\frac{1}{6}}\sigma$.

The shape of the potential and the strength of the Lennard-Jones force

$$F(r) = -\frac{dV(r)}{dr} = \frac{24\varepsilon}{\sigma} \left[2 \left(\frac{\sigma}{r} \right)^{13} - \left(\frac{\sigma}{r} \right)^{7} \right] ,$$

are shown in the following figure:

Lennard-Jones Potential
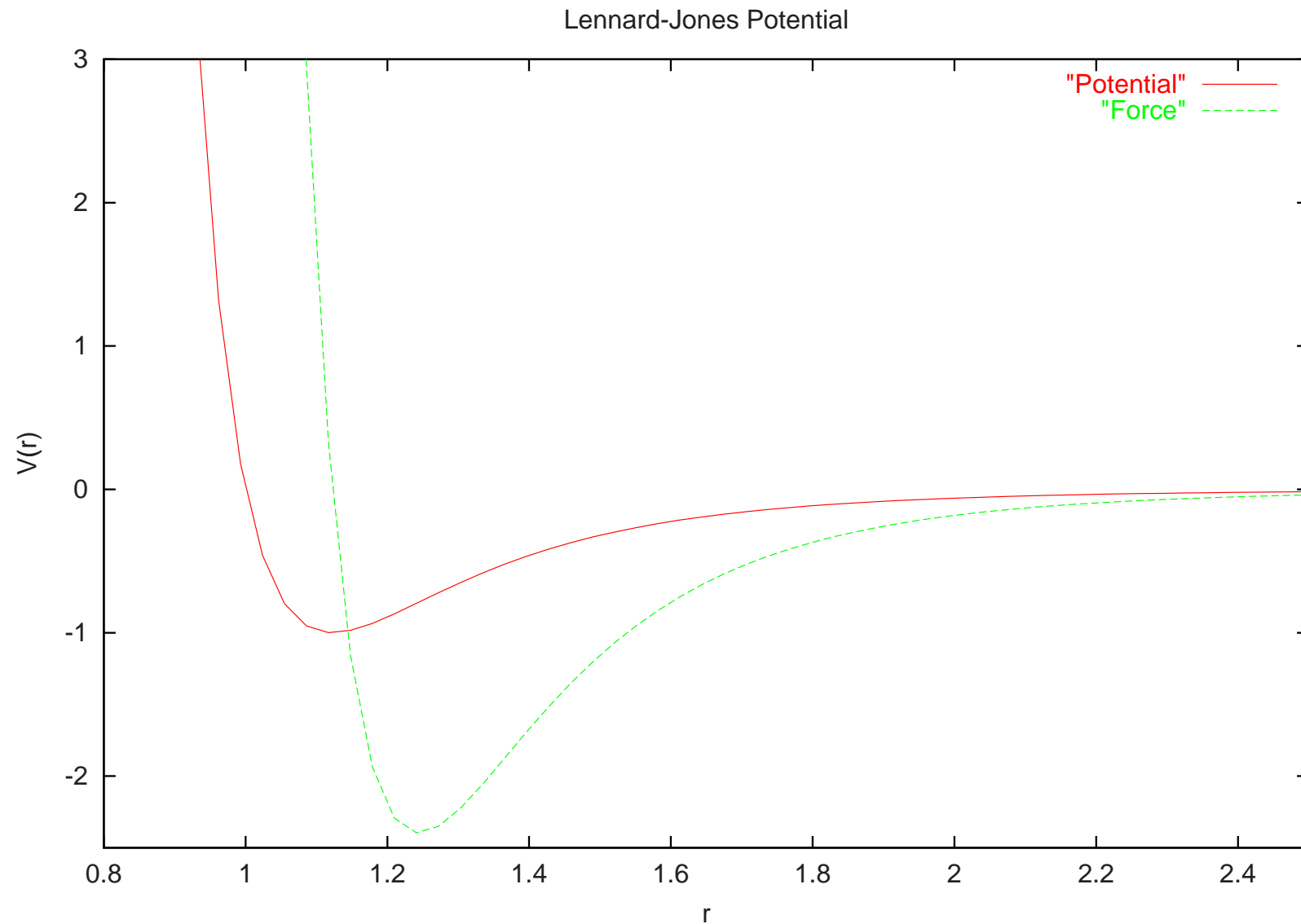
We will choose units of mass, length and energy so that $m = 1$, $\sigma = 1$, and $\varepsilon = 1$. The unit of time

in this system is given by

$$\tau = \sqrt{\frac{m\sigma^2}{\varepsilon}} = 2.17 \times 10^{-12}\,\text{s}\,,$$

which shows that the natural time scale for the dynamics of this system is a few picoseconds!

A simple MD program

First include some standard headers

<div align="right">md.cpp</div>

```cpp
#include <cmath>
#include <cstdlib>
#include <fstream>
#include <iostream>
#include <string>
using namespace std;
```

Define data structures to describe the kinematics of the system

```cpp
const int N = 64;           // number of particles
double r[N][3];             // positions
double v[N][3];             // velocities
double a[N][3];             // accelerations
```

We next need to set the initial positions and velocities of the particles. This is actually a complicated problem! Because the system can be simulated only for a few nanoseconds, the starting configuration must be very close to equilibrium to get good results. For a dense system, the atoms are usually placed at the vertices of a face-centered cubic lattice, which is tends to minimize the potential energy.

The atoms are also given random velocities to approximate the desired temerature.

In this preliminary program we will put the system in a cubical volume of side L and place the particles at the vertices of a simple cubic lattice:

```
double L = 10;              // linear size of cubical volume
double vMax = 0.1;          // maximum initial velocity component


void initialize() {

    // initialize positions
    int n = int(ceil(pow(N, 1.0/3)));   // number of atoms in each direction
    double a = L / n;                    // lattice spacing
    int p = 0;                           // particles placed so far
    for (int x = 0; x < n; x++)
        for (int y = 0; y < n; y++)
            for (int z = 0; z < n; z++) {
                if (p < N) {
                    r[p][0] = (x + 0.5) * a;
                    r[p][1] = (y + 0.5) * a;
                    r[p][2] = (z + 0.5) * a;
                }
                ++p;
            }
```

```
// initialize velocities
for (int p = 0; p < N; p++)
    for (int i = 0; i < 3; i++)
        v[p][i] = vMax * (2 * rand() / double(RAND_MAX) - 1);
}
```

Newton's Equations of Motion

The vector forces between atoms with positions $\mathbf{r}_i$ and $\mathbf{r}_j$ are given by

$$\mathbf{F}_{\text{on } i \text{ by } j} = -\mathbf{F}_{\text{on } j \text{ by } i} = 24(\mathbf{r}_i - \mathbf{r}_j)\left[2\left(\frac{1}{r}\right)^{-14} - \left(\frac{1}{r}\right)^{-8}\right],$$

where $r = |\mathbf{r}_i - \mathbf{r}_j|$.

The net force on atom i due to all of the other $N - 1$ atoms is given by

$$\mathbf{F}_i = \sum_{\substack{j=1 \\ j \neq i}}^{N} \mathbf{F}_{\text{on } i \text{ by } j} .$$

The equation of motion for atom i is

$$\mathbf{a}_i(t) \equiv \frac{d\mathbf{v}_i(t)}{dt} = \frac{d^2\mathbf{r}_i(t)}{dt^2} = \frac{\mathbf{F}_i}{m} ,$$

where $\mathbf{v}_i$ and $\mathbf{a}_i$ are the velocity and acceleration of atom i. These $3N$ second order ordinary differential equations (ODE's) have a unique solution as function of time t if *initial conditions*, that is, the values of positions $\mathbf{r}_i(t_0)$ and velocities $\mathbf{v}_i(t_0)$ of the particles are specified at some initial time t_0. The equations can be integrated numerically by choosing a small time step h and a discrete approximation to the equations to advance the solution by one step at a time.

The following function computes the accelerations of the particles from their current positions:

```
void computeAccelerations() {

    for (int i = 0; i < N; i++)              // set all accelerations to zero
        for (int k = 0; k < 3; k++)
            a[i][k] = 0;

    for (int i = 0; i < N-1; i++)            // loop over all distinct pairs i,j
        for (int j = i+1; j < N; j++) {
            double rij[3];                   // position of i relative to j
            double rSqd = 0;
            for (int k = 0; k < 3; k++) {
                rij[k] = r[i][k] - r[j][k];
                rSqd += rij[k] * rij[k];
            }
            double f = 24 * (2 * pow(rSqd, -7) - pow(rSqd, -4));
            for (int k = 0; k < 3; k++) {
```

```
            a[i][k] += rij[k] * f;
            a[j][k] -= rij[k] * f;
        }
    }
}
```

Velocity Verlet Integration Algorithm

There are many algorithms which can be used to solve ODE's. Verlet has developed several algorithms which are very widely used in MD simulations. One of them is the *velocity Verlet algorithm*

$$\mathbf{r}_i(t + dt) = \mathbf{r}_i(t) + \mathbf{v}_i(t)dt + \frac{1}{2}\mathbf{a}_i(t)dt^2$$

$$\mathbf{v}_i(t + dt) = \mathbf{v}_i(t) + \frac{1}{2}[\mathbf{a}_i(t + dt) + \mathbf{a}_i(t)]dt$$

It can be shown that the errors in this algorithm are of $\mathcal{O}(dt^4)$, and that it is very stable in MD applications and in particular conserves energy very well.

The following function advances the positions and velocities of the particles by one time step:

```
void velocityVerlet(double dt) {
    computeAccelerations();
    for (int i = 0; i < N; i++)
        for (int k = 0; k < 3; k++) {
```

```
        r[i][k] += v[i][k] * dt + 0.5 * a[i][k] * dt * dt;
        v[i][k] += 0.5 * a[i][k] * dt;
    }
    computeAccelerations();
    for (int i = 0; i < N; i++)
        for (int k = 0; k < 3; k++)
            v[i][k] += 0.5 * a[i][k] * dt;
}
```

The instantaneous temperature

This is a simulation in which the number of particles N and the volume L^3 of the system are fixed. Because the Lennard-Jones force is *conservative*, the total energy of the system is also constant.

If the system is in thermal equilibrium, then Boltzmann's *Equipartition Theorem* relates the absolute temperature T to the kinetic energy:

$$3(N-1) \times \frac{1}{2} k_{\mathrm{B}} T = \left\langle \frac{m}{2} \sum_{i=1}^{N} v_i^2 \right\rangle.$$

Here the angle brackets $\langle ... \rangle$ represent a thermal ensemble average. The factor $3(N-1)$ is the number of *internal* translational degrees of freedom which contribute to thermal motion: the motion of the center of mass of the system does not represent thermal energy!

```
    double instantaneousTemperature() {
        double sum = 0;
```

```
    for (int i = 0; i < N; i++)
        for (int k = 0; k < 3; k++)
            sum += v[i][k] * v[i][k];
    return sum / (3 * (N - 1));
}
```

Finally, here is the main function which steers the simulation:

```
int main() {
    initialize();
    double dt = 0.01;
    ofstream file("T.data");
    for (int i = 0; i < 1000; i++) {
        velocityVerlet(dt);
        file << instantaneousTemperature() << '\n';
    }
    file.close();
}
```

There are several ways in which this simple program needs to be improved:

- The volume is not really constant because the particles can move out of it! We need to impose suitable boundary conditions, for example periodic boundary conditions.

- The initial positions and velocities need to be chosen more carefully. We will place the particles on a face-centered cubic lattice, and use a Maxwell-Boltzmann distribution for the velocities.

- The system needs to be allowed to come to thermal equilibrium at the desired temperature.

- Thermal averages of various quantities need to be measured.

Output of simple program md.cpp

The program output file was plotted using Gnuplot. The instantaneous temperature is approximately constant for around one or two time units, and then it starts increasing with fluctuations.

Output of md.cpp
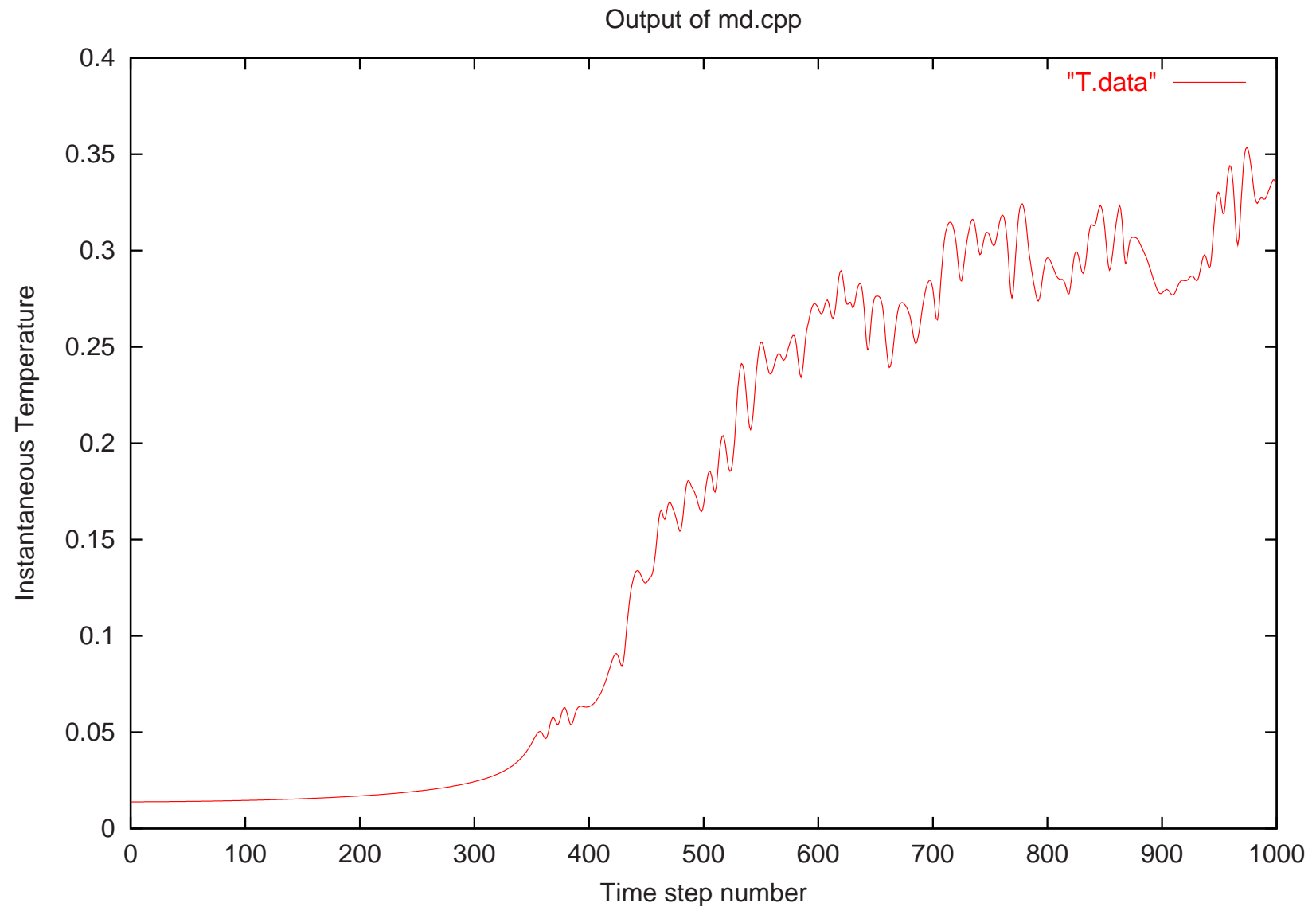
Improving the MD program

md2.cpp

```cpp
#include <cmath>
#include <cstdlib>
#include <fstream>
#include <iostream>
#include <string>
using namespace std;

// simulation parameters

int N = 64;                // number of particles
double rho = 1.0;          // density (number per unit volume)
double T = 1.0;            // temperature

// function declarations

void initialize();         // allocates memory, calls following 2 functions
void initPositions();      // places particles on an fcc lattice
void initVelocities();     // initial Maxwell-Boltzmann velocity distribution
void rescaleVelocities();  // adjust the instanteous temperature to T
double gasdev();           // Gaussian distributed random numbers
```

We will allocate particle arrays dynamically rather than statically

```
double **r;                    // positions
double **v;                    // velocities
double **a;                    // accelerations

void initialize() {
    r = new double* [N];
    v = new double* [N];
    a = new double* [N];
    for (int i = 0; i < N; i++) {
        r[i] = new double [3];
        v[i] = new double [3];
        a[i] = new double [3];
    }
    initPositions();
    initVelocities();
}
```

Position particles on a face-centered cubic lattice

The minimum energy configuration of this Lennard-Jones system is an fcc lattice. This has 4 lattice sites in each conventional cubic unit cell. If the number of atoms $N = 4M^3$, where $M = 1, 2, 3, ...$, then the atoms can fill a cubical volume. So MD simulations are usually done with $32 = 4 \times 2^3, 108 =$

$4 \times 3^3, 256, 500, 864, \ldots$ atoms.
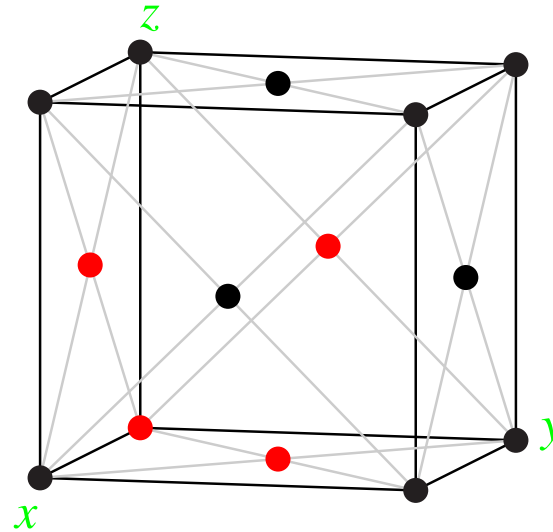
```cpp
double L;                            // linear size of cubical volume

void initPositions() {

    // compute side of cube from number of particles and number density
    L = pow(N / rho, 1.0/3);

    // find M large enough to fit N atoms on an fcc lattice
    int M = 1;
    while (4 * M * M * M < N)
        ++M;
    double a = L / M;               // lattice constant of conventional cell
```

The figure shows a conventional cubical unit cell:

The 4 atoms shown in red provide a basis for the conventional cell. Their positions in units of a are given by

$$(0,\ 0,\ 0) \qquad (0.5,\ 0.5,\ 0) \qquad (0.5,\ 0,\ 0.5) \qquad (0,\ 0.5,\ 0.5)\ .$$

In the following code we shift the basis by $(0.5,\ 0.5,\ 0.5)$ so the all atoms are inside the volume and none are on its boundaries.

```
// 4 atomic positions in fcc unit cell
double xCell[4] = {0.25, 0.75, 0.75, 0.25};
double yCell[4] = {0.25, 0.75, 0.25, 0.75};
double zCell[4] = {0.25, 0.25, 0.75, 0.75};
```

Next, the atoms are placed on the fcc lattice. If $N \neq 4M^3$ then some of the lattice sites are left unoccupied.

```
int n = 0;                        // atoms placed so far
for (int x = 0; x < M; x++)
    for (int y = 0; y < M; y++)
        for (int z = 0; z < M; z++)
            for (int k = 0; k < 4; k++)
                if (n < N) {
                    r[n][0] = (x + xCell[k]) * a;
                    r[n][1] = (y + yCell[k]) * a;
                    r[n][2] = (z + zCell[k]) * a;
                    ++n;
                }
}
```

Draw initial velocities from a Maxwell-Boltzmann distribution

A more realistic initial velocity distribution is that of an ideal gas at temperature T:

$$P(\mathbf{v}) = \left(\frac{m}{2\pi k_{\mathrm{B}} T}\right)^{3/2} e^{-m\left(v_x^2 + v_y^2 + v_z^2\right)/(2k_{\mathrm{B}} T)} .$$

Note that each velocity component is Gaussian distributed with mean zero and width $\sim \sqrt{T}$.

The function gasdev from *Numerical Recipes* returns random numbers with a Gaussian probability distribution

$$P(x) = \frac{e^{-(x-x_0)^2/(2\sigma^2)}}{\sqrt{2\pi\sigma^2}} \ ,$$

with center $x_0 = 0$ and unit variance $\sigma^2 = 1$. This function uses the Box-Müller algorithm.

```
double gasdev () {
    static bool available = false;
    static double gset;
    double fac, rsq, v1, v2;
    if (!available) {
        do {
            v1 = 2.0 * rand() / double(RAND_MAX) - 1.0;
            v2 = 2.0 * rand() / double(RAND_MAX) - 1.0;
            rsq = v1 * v1 + v2 * v2;
        } while (rsq >= 1.0 || rsq == 0.0);
        fac = sqrt(-2.0 * log(rsq) / rsq);
        gset = v1 * fac;
        available = true;
        return v2*fac;
    } else {
        available = false;
        return gset;
```

```
        }
    }

    void initVelocities() {

        // Gaussian with unit variance
        for (int n = 0; n < N; n++)
            for (int i = 0; i < 3; i++)
                v[n][i] = gasdev();
```

Since these velocities are randomly distributed around zero, the total momentum of the system will be close to zero but not exactly zero. To prevent the system from drifting in space, the center-of-mass velocity

$$\mathbf{v}_{\mathrm{CM}} = \frac{\sum_{i=1}^{N} m\mathbf{v}_i}{\sum_{i=1}^{N} m}$$

is computed and used to transform the atom velocities to the center-of-mass frame of reference.

```
        // Adjust velocities so center-of-mass velocity is zero
        double vCM[3] = {0, 0, 0};
        for (int n = 0; n < N; n++)
            for (int i = 0; i < 3; i++)
                vCM[i] += v[n][i];
        for (int i = 0; i < 3; i++)
```

```
        vCM[i] /= N;
    for (int n = 0; n < N; n++)
        for (int i = 0; i < 3; i++)
            v[n][i] -= vCM[i];


    // Rescale velocities to get the desired instantaneous temperature
    rescaleVelocities();
}
```

After setting the CM velocity to zero, the velocities are *scaled*

$$\mathbf{v}_i \longrightarrow \lambda \mathbf{v}_i$$

so that the instantaneous temperature has the desired value T

$$\lambda = \sqrt{\frac{3(N-1)k_{\mathrm{B}}T}{\sum_{i=1}^{N} mv_i^2}} \; .$$

```
    void rescaleVelocities() {
        double vSqdSum = 0;
        for (int n = 0; n < N; n++)
            for (int i = 0; i < 3; i++)
                vSqdSum += v[n][i] * v[n][i];
```

```
        double lambda = sqrt( 3 * (N-1) * T / vSqdSum );
        for (int n = 0; n < N; n++)
            for (int i = 0; i < 3; i++)
                v[n][i] *= lambda;
    }
```

Solving Newton's equations of motion

The same algorithms are used as in md.cpp with two improvements:

- periodic boundary conditions will be used to ensure that the number of particles in the simulation volume remains constant,

- the *minimum image convention* is used to compute the accelerations.

```
    void computeAccelerations() {

        for (int i = 0; i < N; i++)            // set all accelerations to zero
            for (int k = 0; k < 3; k++)
                a[i][k] = 0;

        for (int i = 0; i < N-1; i++)          // loop over all distinct pairs i,j
            for (int j = i+1; j < N; j++) {
                double rij[3];                 // position of i relative to j
                double rSqd = 0;
```

```
for (int k = 0; k < 3; k++) {
    rij[k] = r[i][k] - r[j][k];
```

Since we are using periodic boundary conditions, the system actually has an infinite number of copies of the N particles contained in the volume L^3. Thus there are an infinite number of pairs of particles, all of which interact with one another! The forces between a particular particle and its periodic copies actually cancel, but this is not true of pairs which are not images of one another. Since the Lennard Jones interaction is short ranged, we can safely neglect forces between particles in volumes that are not adjacent to one another. For adjacent volumes, we have to be more careful. It can happen that the separation $r_{ij} = |\mathbf{r}_i - \mathbf{r}_j|$ is *larger* than the separation $r_{ij'} = |\mathbf{r}_i - \mathbf{r}'_j|$ where j' is an image in an adjacent volume of particle j. The figure illustrates this in one dimension:
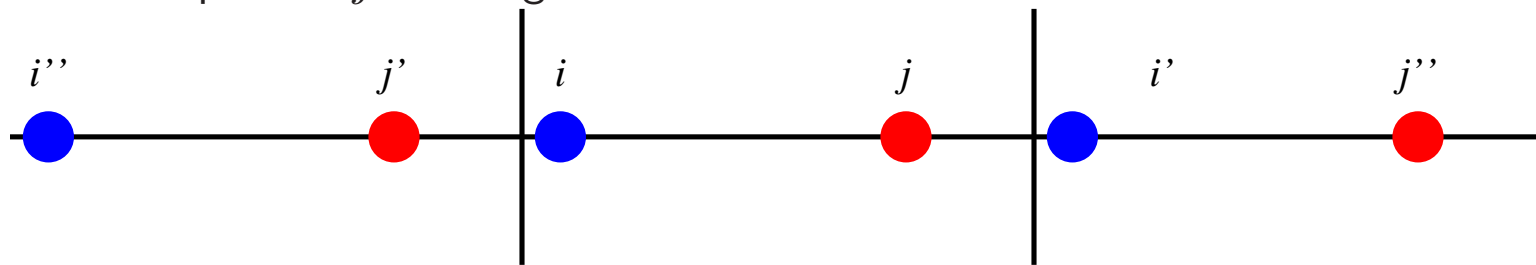
When this occurs we take into account the stronger force between i and the image j' and neglect the weaker force between i and j.

```
// closest image convention
if (abs(rij[k]) > 0.5 * L) {
    if (rij[k] > 0)
        rij[k] -= L;
```

```
            else
                rij[k] += L;
        }
        rSqd += rij[k] * rij[k];
    }
    double f = 24 * (2 * pow(rSqd, -7) - pow(rSqd, -4));
    for (int k = 0; k < 3; k++) {
        a[i][k] += rij[k] * f;
        a[j][k] -= rij[k] * f;
    }
  }
}
```

Velocity Verlet Integration Algorithm

The same algorithm

$$\mathbf{r}_i(t+dt) = \mathbf{r}_i(t) + \mathbf{v}_i(t)dt + \frac{1}{2}\mathbf{a}_i(t)dt^2$$

$$\mathbf{v}_i(t+dt) = \mathbf{v}_i(t) + \frac{1}{2}[\mathbf{a}_i(t+dt) + \mathbf{a}_i(t)]dt$$

is used as in the simple program `md.cpp`. Periodic boundary conditions will be imposed as the time step is implemented.

```
    void velocityVerlet(double dt) {
```

```
    computeAccelerations();
    for (int i = 0; i < N; i++)
        for (int k = 0; k < 3; k++) {
            r[i][k] += v[i][k] * dt + 0.5 * a[i][k] * dt * dt;
```

Once the atom is moved, periodic boundary conditions are imposed to move it back into the system volume if it has exited. This done for each component of the position as soon as it has been updated:

```
            // use periodic boundary conditions
            if (r[i][k] < 0)
                r[i][k] += L;
            if (r[i][k] >= L)
                r[i][k] -= L;
            v[i][k] += 0.5 * a[i][k] * dt;
        }
    computeAccelerations();
    for (int i = 0; i < N; i++)
        for (int k = 0; k < 3; k++)
            v[i][k] += 0.5 * a[i][k] * dt;
}
```

The instantaneous temperature is computed as in `md.cpp` from the equipartition formula

$$3(N-1) \times \frac{1}{2} k_{\mathrm{B}} T = \left\langle \frac{m}{2} \sum_{i=1}^{N} v_i^2 \right\rangle .$$

by the function

```
double instantaneousTemperature() {
    double sum = 0;
    for (int i = 0; i < N; i++)
        for (int k = 0; k < 3; k++)
            sum += v[i][k] * v[i][k];
    return sum / (3 * (N - 1));
}
```

Finally, here is the main function which steers the simulation:

```
int main() {
    initialize();
    double dt = 0.01;
    ofstream file("T2.data");
    for (int i = 0; i < 1000; i++) {
```

```
        velocityVerlet(dt);
        file << instantaneousTemperature() << '\n';
        if (i % 200 == 0)
            rescaleVelocities();
    }
    file.close();
}
```

The simulation is run for 1000 time steps. After every 200 steps, the velocities of the atoms are rescaled to drive the average temperature towards the desired value. The output shows that

- The temperature rises rapidly from the desired value $T = 1.0$ when the simulation is started. Why?

- It takes a few rescaling to push the temperature back to the desired value, and then the system appears to be in equilibrium.

Output of md2.cpp