dispersion relation $\hbar\omega = E = p^2/2m = \hbar^2k^2/2m$. We then get the wave equation

$$i\hbar\partial_t\Psi = -\frac{\hbar^2}{2m}(\boldsymbol{\sigma}\boldsymbol{\nabla})^2\Psi\ . \tag{132}$$

This is *Pauli's equation* for the evolution of a free quantum particle with spin 1/2.

As final step, we include the electric and the magnetic potentials, as we did in the case of the Schrödinger equation. We again use *minimal coupling*, substituting $i\hbar\partial_t$ by $i\hbar\partial_t - qV$ and $-i\hbar\boldsymbol{\nabla}$ by $-i\hbar\boldsymbol{\nabla} - q\boldsymbol{A}$, thus introducing electric charge q and the potentials

Challenge 139 s

V and $\boldsymbol{A}$. A bit of algebra involving the spin operator then leads to the famous complete form of the Pauli equation

$$(i\hbar\partial_t - qV)\Psi = \frac{1}{2m}(-i\hbar\boldsymbol{\nabla} - q\boldsymbol{A})^2\Psi - \frac{q\hbar}{2m}\boldsymbol{\sigma}\boldsymbol{B}\Psi\ , \tag{133}$$

where now the magnetic field $\boldsymbol{B} = \boldsymbol{\nabla}\times\boldsymbol{A}$ appears explicitly. The equation is famous for describing, among others, the motion of silver atoms, which have spin 1/2, in the Stern–

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Gerlach experiment. This is due to the new, last term on the right-hand side, which does not appear in the Schrödinger equation. The new term is a pure spin effect and predicts a g-factor of 2. Depending on the spin orientation, the sign of the last term is either positive or negative; the term thus acts as a spin-dependent potential. The two options for the spin orientation then produce the upper and the lower beams of silver atoms that are observed in the Stern–Gerlach experiment.

In summary, a non-relativistic tangle that rotates continuously reproduces the Pauli equation. In particular, such a tangle predicts that the g-factor of an elementary charged fermion is 2.

Rotating arrows and path integrals

Another simple way to visualize the equivalence between the strand model and the Pauli equation uses the formulation of quantum theory with path integrals. We recall that tangle tails are not observable, and that the tangle core defines the position and phase of the quantum particle. If the core is approximated to be of vanishing size, thus 'point-like', then the motion of the core describes the 'path' of the particle. (This equivalence was already mentioned188 above.)

Feynman described the motion of quantum particles in his famous popular book on

Ref. 165

QED as advancing rotating arrows. The continuous rotation of the tangle core visualizes Feynman's rotating little arrow. The different possible motions of the 'point-like' tangle core corresponds to different paths. Quantum theory appears when the effects of all possible paths are superposed. In particular, the phase and amplitude for each path must added like small vectors. In the strand model the effects of all possible paths are added automatically, through the fluctuations of the tangles motion. And through the definitions given above, the addition occurs in exactly the way that Feynman described.

In this way, the tangle model reproduces the path integral formulation of quantum mechanics.

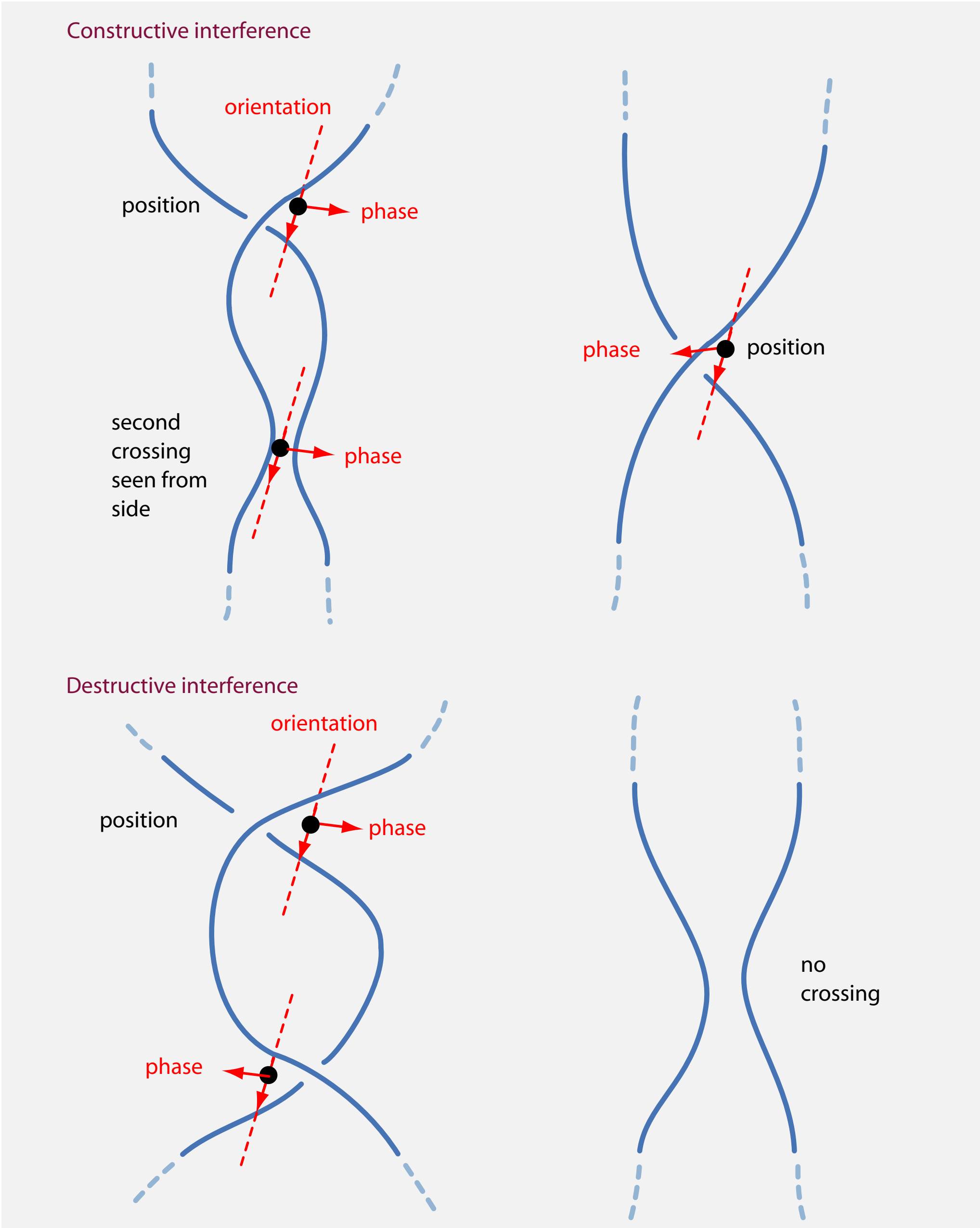

FIGURE 37 Interference in the strand model: two connected crossings – in this case with the same distance, i.e., with the same 'amplitude' – superpose constructively (top) and destructively (bottom).

INTERFERENCE AND DOUBLE SLITS

Often, interference is seen as the essence of quantum theory, or as the biggest difference between quantum theory and classical physics. Also interference can be visualized with strands.

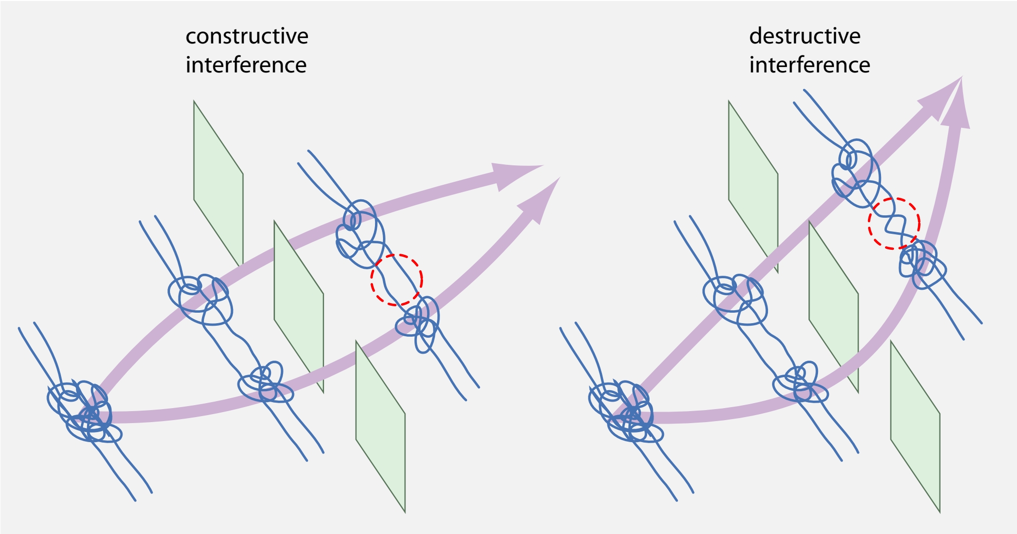

FIGURE 38 A fermion tangle passing a double slit: constructive interference (left) and destructive interference (right).

In nature, interference – whether constructive or destructive or partial – results from the superposition of states with different phase values at a point in space. For strands, superposition occurs for strands, when strand segments are connected. The phase value is defined by the crossing geometry, as explained in Figure 31. In the strand model, 'at a point in space' becomes 'at two points within Planck scale distance'. Together, this yields the fundamental superposition mechanism shown in Figure 37. The phases of two crossings on two connected strands can add up or cancel. The general case, with different amplitudes, is easily deduced.

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The strand explanation of interference allows to describe the double slit experiment. Because of its tails, a fermion tangle obeys spinor statistics and spinor rotation behaviour. This leads to the observed interference behaviour for spin 1/2 particles, as visualized in Figure 38. The corresponding visualization for photon interference is given in Figure 39.

Interference is essential in all effects of quantum physics. We explore a few additional ones.

Measurements and wave function collapse

In nature, a measurement of a quantum system in a superposition is observed to yield one of the possible eigenvalues and to prepare the system in the corresponding eigenstate. In nature, the probability of each measurement outcome depends on the coefficient of that eigenstate in the superposition.

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To put the issue into context, here is a short reminder from quantum mechanics. Every measurement apparatus shows measurement results. Thus, every measurement apparatus is a device with memory. (In short, it is *classical*.) All devices with memory contain one or several baths. Thus, every measurement apparatus couples at least one bath to the system it measures. The coupling depends on and defines the observable to be measured by the apparatus. Every coupling of a bath to a *quantum* systems leads to decoherence.

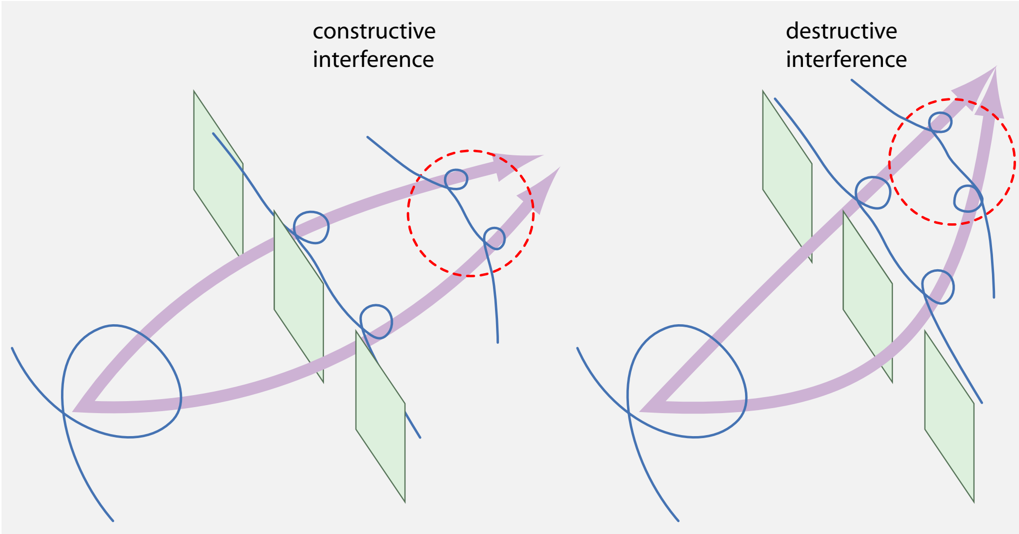

FIGURE 39 The double-slit experiment with photons: constructive interference (left) and destructive interference (right).

Decoherence leads to probabilities and wave function collapse. In short, collapse and measurement probabilities are necessary and automatic in quantum theory.

The strand model describes the measurement process in precisely the same way as standard quantum theory; in addition, it *visualizes* the process.

▷ A ***measurement*** is modelled as a strand deformation induced by the measurement apparatus that 'pulls' a tangle towards the resulting eigenstate.

▷ This pulling of strands models and visualizes the ***collapse*** of the wave function.

An example of measurement is illustrated in Figure 40. When a measurement is performed on a superposition, *the untangled 'addition region' can be imagined to shrink into disappearance.* For this to happen, one of the underlying eigenstates has to 'eat up' the other: that is the collapse of the wave function. In the example of the figure, the addition region can disappear either towards the outside or towards the inside. The choice is due to the bath that is coupled to the system during measurement; the bath thus determines the outcome of the measurement. We also deduce that the probability of measuring a particular eigenstate will depend on the (weighed) volume that the eigenstate took up in the superposition.

This visualization of the wave function collapse also makes clear that the collapse is not limited by any speed limit, as no energy and no information is transported. Indeed, the collapse happens by displacing strands and at most crossings, but does not produce any crossing changes.

In summary, the strand model describes measurements in precisely the same way as usual quantum theory. In addition, *strands visualize the collapse of the wave function as a shape deformation from a superposed tangle to an eigenstate tangle.*

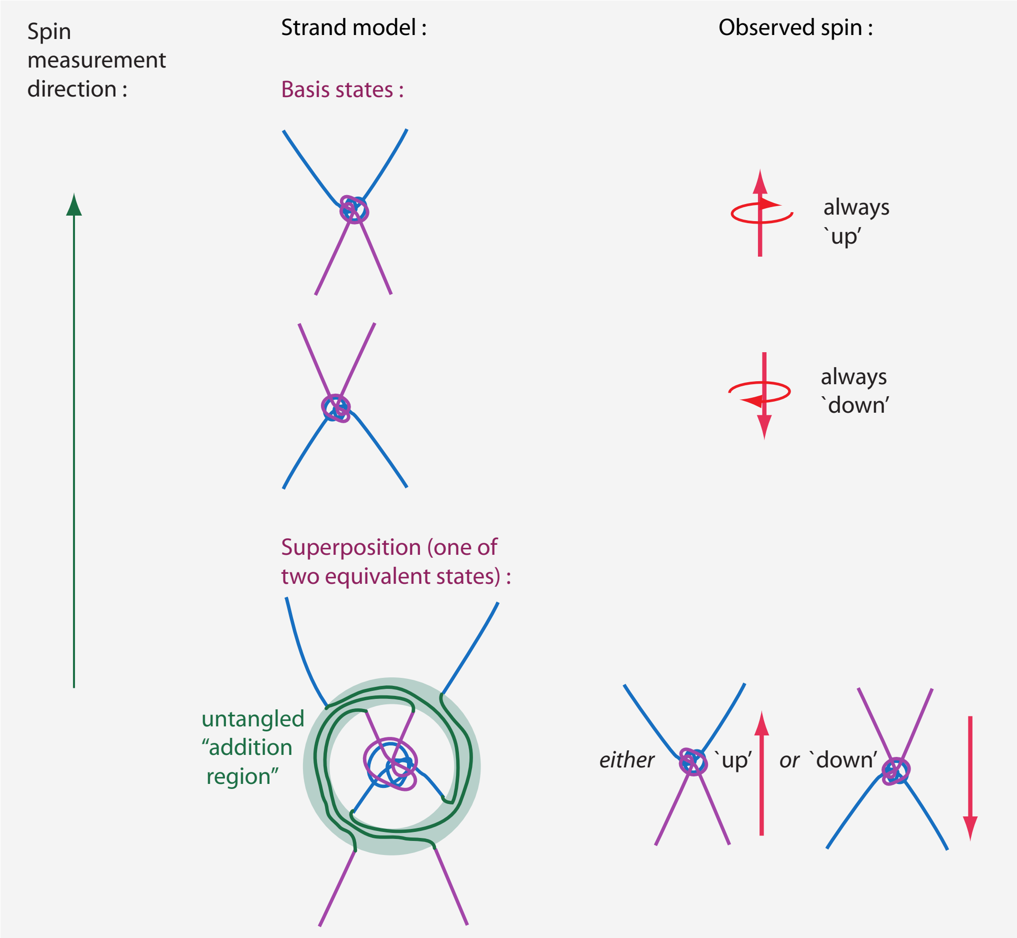

FIGURE 40 Measurement of a spin superposition: the addition region disappears either outwards or inwards.

Hidden variables and the Kochen–Specker theorem

At first sight, the strand model seems to fall into the trap of introducing hidden variables into quantum theory. One could indeed argue that the shapes (and fluctuations) of the strands play the role of hidden variables. On the other hand, it is well known that non-contextual hidden variables are impossible in quantum theory, as shown by the Kochen–Specker theorem (for sufficiently high Hilbert-space dimensions). Is the strand model flawed? No. Ref. 166

We recall that strands are not observable. In particular, strand shapes are not physical observables and thus not physical (hidden) variables either. Even if we tried promoting strand shapes to physical variables, the evolution of the strand shapes would only be observable through the ensuing crossing switches. And crossing switches evolve due to the influence of the environment, which consists of all other strands in nature, including those of space-time itself. Thus

▷ The evolution of strand shapes and crossing switches is *contextual*.

Therefore, the strand model does not contradict the Kochen–Specker theorem.

In simple language, in quantum theory, hidden variables are not a problem if they are properties of the environment, and not of the quantum system itself. This is precisely the case for the strand model. For a quantum system, the strand model provides no hidden variables. In fact, for a quantum system, the strand model provides no variables beyond the usual ones from quantum theory. And as expected and required from any model that reproduces decoherence, the strand model leads to a contextual, probabilistic description of nature.

In summary, despite using fluctuating tangles as underlying structure, the strand model is equivalent to usual quantum theory. The strand model contains nothing more and nothing less than usual quantum theory.

Many-particle states and entanglement

In nature, the quantum states of two or more particles can be *entangled*. Entangled states are many-particle states that are not separable. Entangled states are one of the most fascinating quantum phenomena; especially in the case of macroscopic entanglement, they are still being explored in many experiments. We will discover that the strand model visualizes them simply and clearly.

To describe entanglement, we first need to clarify the notion of many-particle state. In the strand model,

▷ A ***many-particle state*** is composed of several tangles.

In this way, an N-particle wave function defines N values at every point in space, one value for each particle. This is possible, because in the strand model, the strands of one particle tangle are *separate* from the strands of other particles.

Usually, a N-particle wave function is described by a single-valued function in $3N$ dimensions. It is less known that a single-valued N-particle wave function in $3N$ dimensions is mathematically equivalent to an N-valued wave function in three dimensions. Usually, N-valued functions are not discussed; we feel uneasy with the concept. But the strand model naturally defines N wave function values at each point in space: each particle has its own tangle, and each tangle yields, via short-term averaging, one complex value, with magnitude and phase, at each point in space. In this way, the strand model is able to describe N particles in just 3 dimensions.

In other words, the strand model does not describe N particles with 1 function in $3N$ dimensions; it describes many-particle states with N functions in 3 dimensions. In this way, the strand model remains as close to everyday life as possible. Many incorrect statements on this issue are found in the research literature; many authors incorrectly claim the impossibility of many-particle quantum theory in 3 dimensions. Some authors even claim, in contrast to experiment, that it is impossible to visualize many-particle states in 3 dimensions. These arguments all fail to consider the possibility to define completely separate wave functions for each particle in three dimensions. (It must be said that this unusual possibility is hard to imagine if wave functions are described as

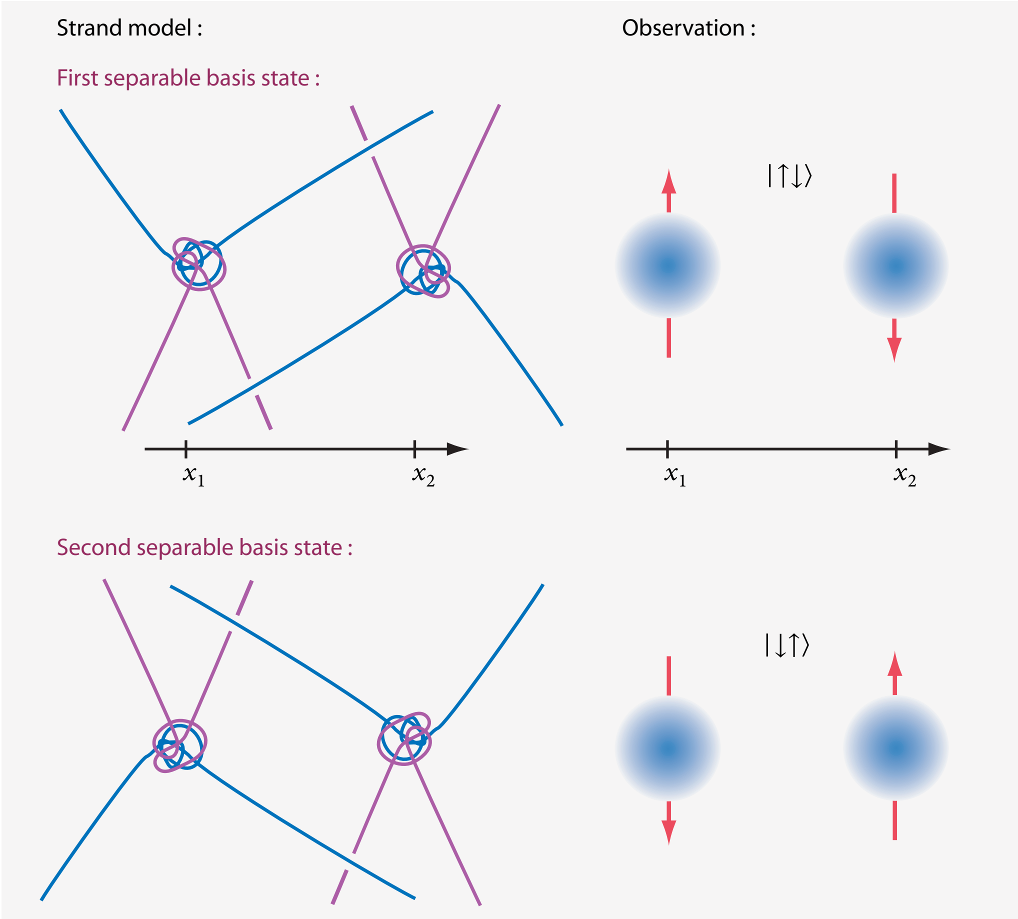

FIGURE 41 Two examples of two distant particles with spin in separable states: observation and strand model.

Ref. 165 continuous functions.) However, clear thinkers like Richard Feynman always pictured many-particle wave functions in 3 dimensions. Also in this domain, the strand model provides an underlying picture to Feynman's approach. This is another situation where the strand model eliminates incorrect thinking habits and supports the 'naive' view of quantum theory.

Now that we have defined many-particle states, we can also define entangled states.

▷ An ***entangled state*** is a non-separable superposition of separable many-particle states. State are separable when their tangles can be pulled away without their tails being entangled.

We will now show that the above definitions of superpositions and of measurements using strands are sufficient to describe entanglement.

As first example, we explore entangled states of the spin of two distant massive fermions. This is the famous thought experiment proposed by David Bohm. In the strand

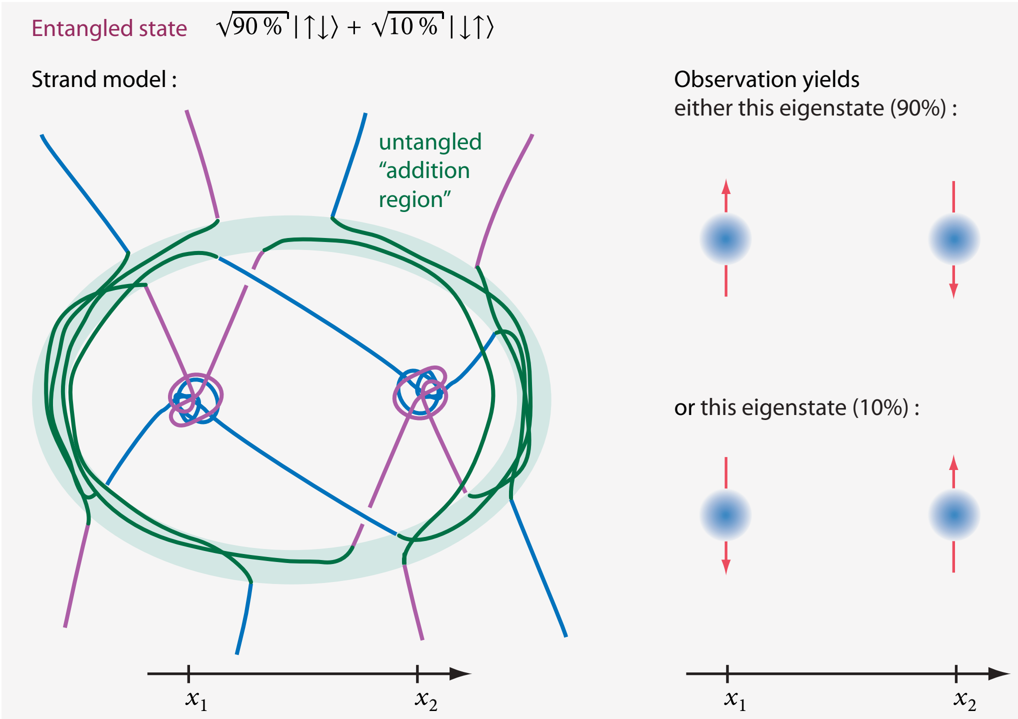

FIGURE 42 An entangled spin state of two distant particles.

model, two distant particles with spin 1/2 in a *separable* state are modelled as two distant, separate tangles of identical topology. Figure 41 shows two separable basis states, namely the two states with total spin 0 given by $|\uparrow\downarrow\rangle$ and by $|\downarrow\uparrow\rangle$. Such states can also be produced in experiments. We note that to ensure total spin 0, the tails must be imagined to cross somewhere, as shown in the figure.

We can now draw a superposition $\sqrt{90\,\%}\,|\uparrow\downarrow\rangle + \sqrt{10\,\%}\,|\downarrow\uparrow\rangle$ of the two spin-0 basis states. We simply use the definition of addition and find the state shown in Figure 42. We can now use the definition of measurement to check that the state is indeed entangled. If we measure the spin orientation of one of the particles, the untangled addition region disappears. The result of the measurement will be either the state on the inside of the addition region or the state on the outside. And since the tails of the two particles are linked, after the measurement, independently of the outcome, the spin of the two particles will always point in opposite directions. This happens for every particle distance. Despite this extremely rapid and apparently superluminal collapse, no energy travels faster than light. The strand model thus reproduces exactly the observed behaviour of entangled spin 1/2 states.

A second example is the entanglement of two photons, the well-known Aspect experiment. Also in this case, entangled spin 0 states, i.e., entangled states of photons of opposite helicity (spin), are most interesting. Again, the strand model helps to visualize the situation. Here we use the strand model for the photon that we will deduce only later Page 230 on. Figure 43 shows the strand model of the two separable basis states and the strand model of the entangled state. Again, the measurement of the helicity of one photon in

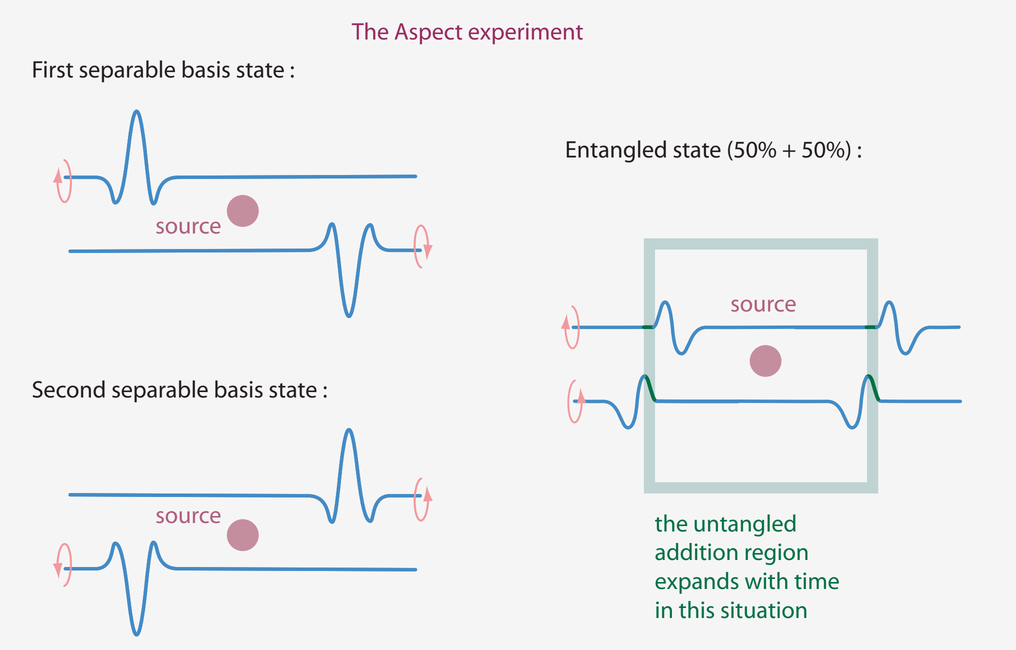

FIGURE 43 The basis states and an entangled state of two distant photons travelling in opposite directions, with total spin 0.

the entangled state will lead to one of the two basis states. And as soon as the helicity of one photon is measured, the helicity of its companion collapses to the opposite value, whatever the distance! Experimentally, the effect has been observed for distances of many kilometres. Again, despite the extremely rapid collapse, no energy travels faster than light. And again, the strand model completely reproduces the observations.

Ref. 167

Mixed states

Mixed states are statistical ensembles of pure states. In the strand model,

▷ A ***mixed state*** is a (weighted) temporal alternation of pure states.

Mixed states are important in discussions of thermodynamic quantities. We mention them to complete the equivalence of the states that appear in quantum theory with those provided by the strand model. We do not pursue this topic any further.

The dimensionality of space-time

'Nature consists of particles moving in empty space.' Democritus stated this 2500 years ago. Today, we know that is a simplified description of one half of physics: it is a simplified description of quantum theory. In fact, Democritus' statement, together with strands, allows us to argue that physical space must have three dimensions, as we will see now.

Deducing the dimensionality of physical space from first principles is an old and difficult problem. The difficulty is also due to the lack of alternative descriptions of nature. Page 169 Our exploration of the foundations of the strand model has shown that humans, animals and machines always use three spatial dimensions to describe their environment. They cannot do otherwise. Humans, animals and machines cannot talk and think without three dimensions as background space.

But how can we show that *physical space* – not the *background space* we need for thinking – is three-dimensional and must be so? We need to show that (1) all experiments reproduce the result and that (2) no other number of dimensions yields a consistent description of nature.

In nature, and also in the strand model, as long as particles can be defined, they can be rotated around each other and they can be exchanged. No experiment has ever been performed or has ever been proposed that changes this observation. The observed properties of rotations, of spin 1/2, of particle exchange and all other observations confirm that space has three dimensions. Fermions only exist in three dimensions. In the strand model, the position and the orientation of a particle is intrinsically a three-dimensional quantity; physical space is thus three-dimensional, in all situations where it can be defined. (The only situations where this definition is impossible are horizons and the Planck scales.) In short, both nature and the strand model are found to be three-dimensional at all experimentally accessible energy scales. Conversely, detecting an additional spatial dimension would directly invalidate the strand model.

Nature has three dimensions. The only way to predict this result is to show that no other number is possible. The number of dimensions of nature can only result from a self-consistency argument. And interestingly, the strand model produces such an argument.

In the strand model, knots and tangles are impossible to construct in physical spaces with dimensions *other* than three. Indeed, mathematicians can show that in four spatial dimensions, every knot and every tangle can be undone. (In this argument, time is not and does not count as a fourth spatial dimension, and strands are assumed to remain one-dimensional entities.) Worse, in the strand model, spin does not exist in spaces that have more or fewer than three dimensions. Also the vacuum and its quantum fluctuations do not exist in more than three dimensions. Moreover, in other dimensions it is impossible to formulate the fundamental principle. In short, the strand model of matter and of observers, be they animals, people or machines, is possible in three spatial dimensions only. No description of nature with a background or physical space of more or less than three dimensions is possible with strands. Conversely, constructing such a description would invalidate the strand model.

The same type of arguments can be collected for the one-dimensionality of physical Challenge 141 e time. It can be fun exploring them – for a short while. In summary, the strand model *only* works in $3+1$ space-time dimensions; it does not allow any other number of dimensions. Page 164 We have thus ticked off another of the millennium issues. We can thus continue with our adventure.

Operators and the Heisenberg picture

In quantum theory, Hermitean operators play an important role. In the strand model, *Hermitean* or self-adjoint operators are operators that leave the tangle topology invariant.

Also unitary operators play an important role in quantum theory. In the strand model, *unitary* operators are operators that deform tangles in a way that the corresponding wave function retains its norm, i.e., such that tangles retain their topology and their core shape.

Physicists know two ways to describe quantum theory. One is to describe evolution with time-dependent quantum states – the *Schrödinger picture* we are using here – and the other is to describe evolution with time-dependent operators. In this so-called *Heisenberg picture*, the temporal evolution is described by the operators.

The two pictures of quantum theory are equivalent. In the Heisenberg picture, the fundamental principle, the equivalence of a crossing switch with $\hbar$, becomes a statement on the behaviour of operators. Already in 1987, Louis Kauffman had argued that the commutation relation for the momentum and position operators

Ref. 168

$$px - xp = \hbar i \qquad (134)$$

is related to a crossing switch. The present section confirms that speculation.

In quantum mechanics, the commutation relation follows from the definition of the momentum operator as $p = \hbar k$, $k = -i\partial_x$ being the wave vector operator. The factor $\hbar$ defines the unit of momentum. The wave vector counts the number of wave crests of a wave. Now, in the strand model, a rotation of a state by an angle π is described by a multiplication by i. Counting wave crests of a propagating state is only possible by using the factor i, as this factor is the only property that distinguishes a crest from a trough. In short, the commutation relation follows from the fundamental principle of the strand model.

Lagrangians and the principle of least action

Before we derive the Dirac equation, we show that the strand model naturally leads to describe motion with Lagrangians.

In nature, physical action is an observable measured in multiples of the natural unit, the quantum of action $\hbar$. Action is the fundamental observable about nature, because *action measures the total change occurring in a process.*

In the strand model,

▷ The physical ***action*** W of a physical process is the observed number of crossing switches of strands. Action values are multiples of $\hbar$.

We note that these multiples, if averaged, do not need to be integer multiples. We further note that through this definition, *action is observer-invariant.* This important property is thus automatic in the strand model.

In nature, energy is action per time. Thus, in the strand model we have:

▷ ***Energy*** is the number of crossing switches per time in a system.

In nature, when free quantum particles move, their phase changes linearly with time. In other words, the 'little arrow' representing the free particle phase rotates with constant angular frequency. We saw that in the strand model, the 'little arrow' is taken as (half) the orientation angle of the tangle core, and the arrow rotation is (half) the rotation of the tangle core.

▷ The ***kinetic energy*** T of a particle is the number of crossing switches per time induced by shape fluctuations of the continuously rotating tangle core.

We call $\mathcal{T}$ the corresponding volume density: $\mathcal{T} = T/V$. In nature, the Lagrangian is a practical quantity to describe motion. For a *free* particle, the Lagrangian density $\mathcal{L} = \mathcal{T}$ is simply the kinetic energy density, and the action $W = \int \mathcal{L}\, \mathrm{d}V \mathrm{d}t = Tt$ is the product of kinetic energy and time. In the strand model, a free particle is a constantly rotating and advancing tangle. We see directly that this constant evolution minimizes the action W for a particle, given the states at the start and at the end.

This aspect is more interesting for particles that interact. Interactions can be described by a potential energy U, which is, more properly speaking, the energy of the field that produces the interaction. In the strand model,

▷ ***Potential energy*** U is the number of crossing switches per time induced by an interaction field.

We call $\mathcal{U}$ the corresponding volume density: $\mathcal{U} = U/V$. In short, in the strand model, an interaction changes the rotation rate and the linear motion of a particle tangle.

In the strand model, the *difference* between kinetic and potential energy is thus a quantity that describes how much a system consisting of a tangle and a field *changes* at a given time. The total change is the integral over time of all instantaneous changes. In other words, in the strand model we have:

▷ The ***Lagrangian density*** $\mathcal{L} = \mathcal{T} - \mathcal{U}$ is the number of crossing switches per volume and time, averaged over many Planck scales.

▷ The physical ***action*** $W = \int L\, \mathrm{d}t = \int\int \mathcal{L}\, \mathrm{d}V \mathrm{d}t$ of a physical process is the observed number of crossing switches of strands. The action value W_{if} between an initial state ψ_{i} and a final state ψ_{f} is given by

$$W_{\mathrm{if}} = \langle \psi_{\mathrm{i}} \mid \int \mathcal{L}\, \mathrm{d}t \mid \psi_{\mathrm{f}} \rangle = \langle \psi_{\mathrm{i}} \mid \int (\mathcal{T} - \mathcal{U})\, \mathrm{d}t \mid \psi_{\mathrm{f}} \rangle \;. \qquad (135)$$

Since energy is related to crossing switches, it is natural that strand fluctuations that do *not* induce crossing switches are *favoured*. In short, the strand model states

▷ Evolution of strands *minimizes crossing switch number.* As a result, strands minimize the action W.

In the strand model, *the least action principle appears naturally*. In the strand model, an evolution has least action when it occurs with the smallest number of crossing changes. With this connection, one can also show that the strand model implies Schwinger's quantum action principle.

Challenge 142 e

To calculate quantum motion with the principle of least action, we need to define the kinetic and the potential energy in terms of strands. There are various possibilities for Lagrangian densities for a given evolution equation; however, all are equivalent. In case of the free Schrödinger equation, one possibility is:

$$\mathcal{L} = \frac{i\hbar}{2}(\overline{\psi}\,\partial_t\psi - \partial_t\overline{\psi}\,\psi) - \frac{\hbar^2}{2m}\nabla\overline{\psi}\nabla\psi \ . \tag{136}$$

In this way, the principle of least action can be used to describe the evolution of the Schrödinger equation. The same is possible for situations with potentials, for the Pauli equation, and for all other evolution equations of quantum particles.

We thus retain that the strand model explains the least action principle. It explains it because strand evolution minimizes the number of crossing switches.

Special relativity: the vacuum

In nature, there is an invariant limit energy speed c, namely the speed of light and of all other massless radiation. Special relativity is the description of the consequences from this observation, in the case of a flat space-time.

We remark that special relativity also implies and requires that the flat vacuum looks exactly the same for all inertial observers. In the strand model, the idea of flat vacuum as a set of fluctuating featureless strands that are *unknotted* and *unlinked* automatically implies that for any inertial observer the flat vacuum has no matter content, has no energy content, is isotropic and is homogeneous. The strand model thus realizes this basic requirement of special relativity. In the strand model, vacuum is *Lorentz-invariant*.

Many models of the vacuum, even fluctuating ones, have difficulties reproducing Lorentz invariance. The strand model differs, because the strands are not the observable entities; only their crossing switches are. This topological definition, together with the averaging of the fluctuations, makes the vacuum Lorentz-invariant.

We note that in the strand model, the vacuum is unique, and the vacuum energy of flat infinite vacuum is exactly zero. In the strand model, there is no divergence of the vacuum energy, and there is thus *no* contribution to the cosmological constant from quantum field theory. In particular, in the strand model there is no need for supersymmetry to explain the small energy density of the vacuum.

Special relativity: the invariant limit speed

In the strand model, massless particles are unknotted and untangled. Even though we will deduce the strand model for photons only later on, we use it here already, to speed up the discussion. In the strand model, the *photon* is described by a single, helically deformed unknotted strand, as shown in Figure 51. Therefore, we can define:

Page 230

▷ The ***Planck speed*** c is the observed average speed of crossing switches due to photons.

Because the definition uses crossing switches and a massless particle, the speed of light c is an *energy speed*. Also speed of light c is an average for long times. Indeed, as is well-

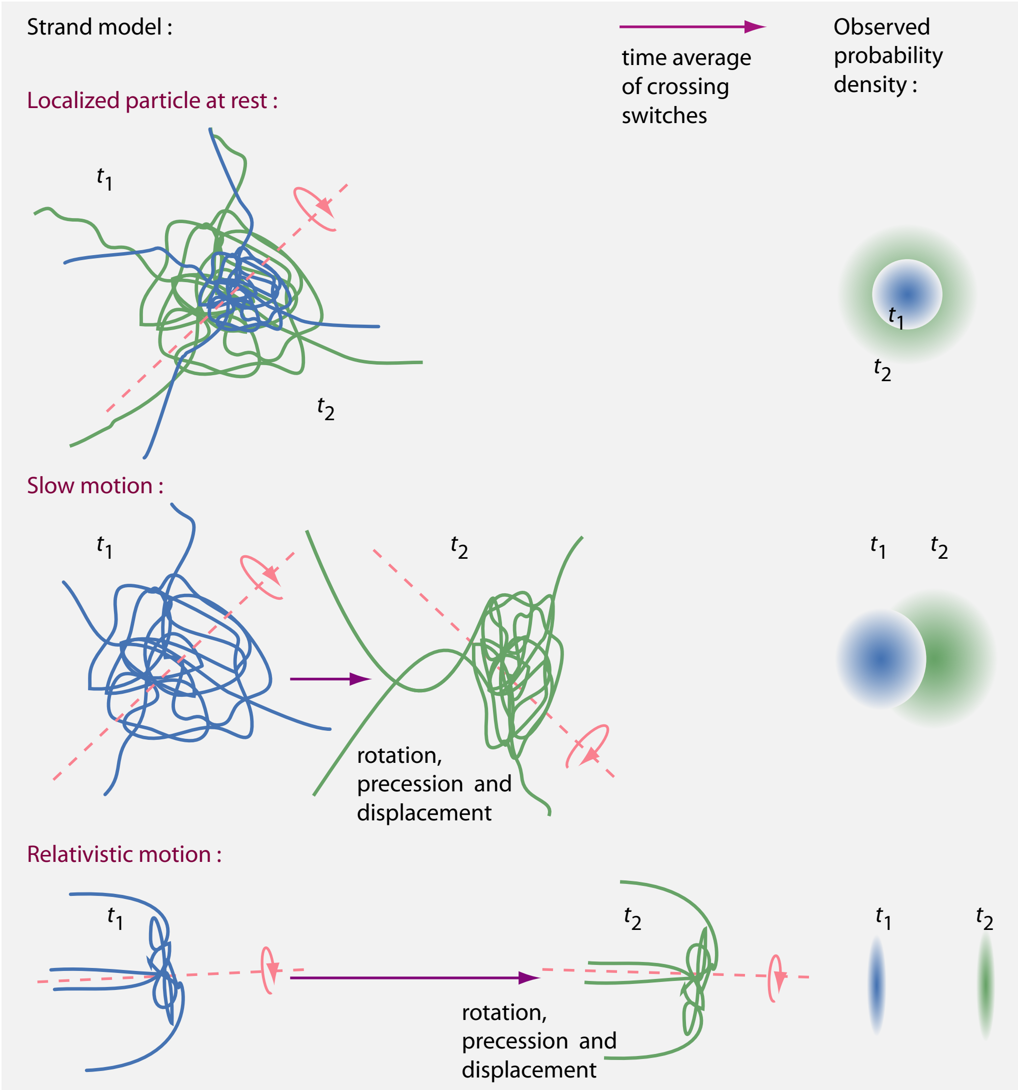

FIGURE 44 Tangles at rest, at low speed and at relativistic speed.

Ref. 169 known in quantum field theory, due to the indeterminacy relation, single photons can travel faster or slower than light, but the probability for large deviations is extremely low.

The linear motion of a helically deformed photon strand through the vacuum strands is similar to the motion of a bottle opener through cork. It differs from the linear motion of a matter tangle through vacuum, which makes use of the belt trick. The belt trick slows Page 351 fermions down, though the details are not simple, as we will discover below. In short, we find that matter tangles always move *more slowly than light*. The speed c is a *limit speed*.

In fact, we see that ultrarelativistic tangles move, as shown in Figure 44, almost like light. We thus find that matter can *almost* reach the speed of light. The speed c is indeed a *limit speed* for matter.

However, one problem remains open: how exactly do tangles move through the web

Page 351 that describes the vacuum? We will clarify this issue later on. In a few words, the motion of a photon requires that the strands of the surrounding space make room for it. This requires favourable fluctuations, thus a finite time. The motion process of photons thus makes it clear that the speed of light is *finite*.

The speed of light c is defined as an average, because, as well-known in quantum field theory, there are small probabilities that light moves faster or slower that c. But the average result c will be the same for every observer. The value of the speed c is thus *invariant*.

In 1905, Einstein showed that the mentioned properties of the speed of light – energy speed, limit speed, finite speed and invariant speed – imply the Lorentz transformations. In particular, the three properties of the speed of light c imply that the energy E of a particle of mass m is related to its momentum p as

$$E^2 = m^2c^4 + c^2p^2 \quad \text{or} \quad \hbar^2\omega^2 = m^2c^4 + c^2\hbar^2k^2 \; . \tag{137}$$

This dispersion relation is thus also valid for massive particles made of tangled strands – Page 358 even though we cannot yet calculate tangle masses. (We will do this later on.)

Should we be surprised at this result? No. In the fundamental principle, the definition Page 152 of the crossing switch, we inserted the speed of light as the ratio between the Planck length and the Planck time. Therefore, by defining the crossing switch in the way we did, we have implicitly stated the invariance of the speed of light.

Fluctuating strands imply that flat vacuum has no matter or energy content, for *every* inertial observer. Due to the strand fluctuations, flat vacuum is also homogeneous and isotropic for every inertial observer. Therefore, together with the $3+1$-dimensionality of Page 209 space-time deduced above, we have now definitely shown that flat vacuum has Poincaré Page 164 symmetry. This settles another issue from the millennium list.

Page 196 The relativistic dispersion relation differs from the non-relativistic case in two ways. First, the energy scale is shifted, and now includes the rest energy $E_0 = c^2m$. Secondly, the spin precession is not independent of the particle speed any more; for relativistic particles, the spin lies close to the direction of motion. Both effects follow from the existence of a limit speed.

If we neglect spin, we can use the relativistic dispersion relation to deduce directly the well-known Klein–Gordon equation for the evolution of a wave function:

$$-\hbar^2\partial_{tt}\psi = m^2c^4 - c^2\hbar^2\nabla^2\psi \; . \tag{138}$$

In other words, the strand model implies that relativistic tangles follow the Klein–Gordon equation. We now build on this result to deduce Dirac's equation for relativistic quantum motion.

Dirac's equation deduced from tangles

The relativistic Klein–Gordon equation assumes that spin effects are negligible. This approximation fails to describe most experiments. A precise description of relativistic elementary particles must include spin.

So far, we deduced the Schrödinger equation using the relation between phase and the quantum of action, using the non-relativistic energy–momentum relation, and neglecting spin. In the next step we deduced the Pauli equation by including the properties of spin 1/2. The following step was to deduce the Klein–Gordon equation using again the relation between phase and the quantum of action, this time the relativistic energy–momentum relation, but assuming zero spin. The final and correct description of elementary fermions, the Dirac equation, results from combining all three ingredients: (1) the relation between the quantum of action and the phase of the wave function, (2) the relativistic mass–energy relation, and (3) the effects of spin 1/2. Now we can reproduce this derivation because all three ingredients are reproduced by the strand model.

We first recall the derivation of the Dirac equation found in textbooks. The main observation about spin in the relativistic context is the existence of states of right-handed and of left-handed chirality: spin can precess in two opposite senses around the direction of momentum. In addition, for massive particles, the two chiral states mix. The existence of two chiralities requires a description of spinning particles with a wave function that has *four* complex components, thus *twice* the number of components that appear in the Pauli equation. Indeed, the Pauli equation implicitly assumes only one, given sign for the chirality, even though it does not specify it. This simple description is possible because in non-relativistic situations, states of different chirality do not mix.

Consistency requires that each of the four components of the wave function of a relativistic spinning particle must follow the relativistic energy–momentum relation, and
Ref. 170 thus the Klein–Gordon equation. This requirement is known to be sufficient to deduce
Ref. 171 the Dirac equation. One of the simplest derivations is due to Lerner; we summarize it here.

When a spinning object moves relativistically, we must take both chiralities into account. We call u the negative chiral state and v the positive chiral state. Each state is described by two complex numbers that depend on space and time. The 4-vector for probability and current becomes

$$J_\mu = u^\dagger \sigma_\mu u + v^\dagger \sigma_\mu v \,. \qquad (139)$$

We now introduce the four-component spinor φ and the 4×4 spin matrices a_μ

$$\varphi = \begin{pmatrix} u \\ v \end{pmatrix} \quad \text{and} \quad \alpha_\mu = \begin{pmatrix} \sigma_\mu & 0 \\ 0 & \overline{\sigma}_\mu \end{pmatrix} \,, \qquad (140)$$

where $\sigma_\mu = (I, \boldsymbol{\sigma})$ and $\overline{\sigma}_\mu = (I, -\boldsymbol{\sigma})$ and I is the 2×2 identity matrix. The 4-current can then be written as

$$J_\mu = \varphi^\dagger \alpha_\mu \varphi \,. \qquad (141)$$

The three requirements of current conservation, Lorentz invariance and linearity then
Ref. 171 yield the evolution equation

$$\mathrm{i}\hbar \partial^\mu (\alpha_\mu \varphi) + mc\gamma_5 \varphi = 0 \,. \qquad (142)$$

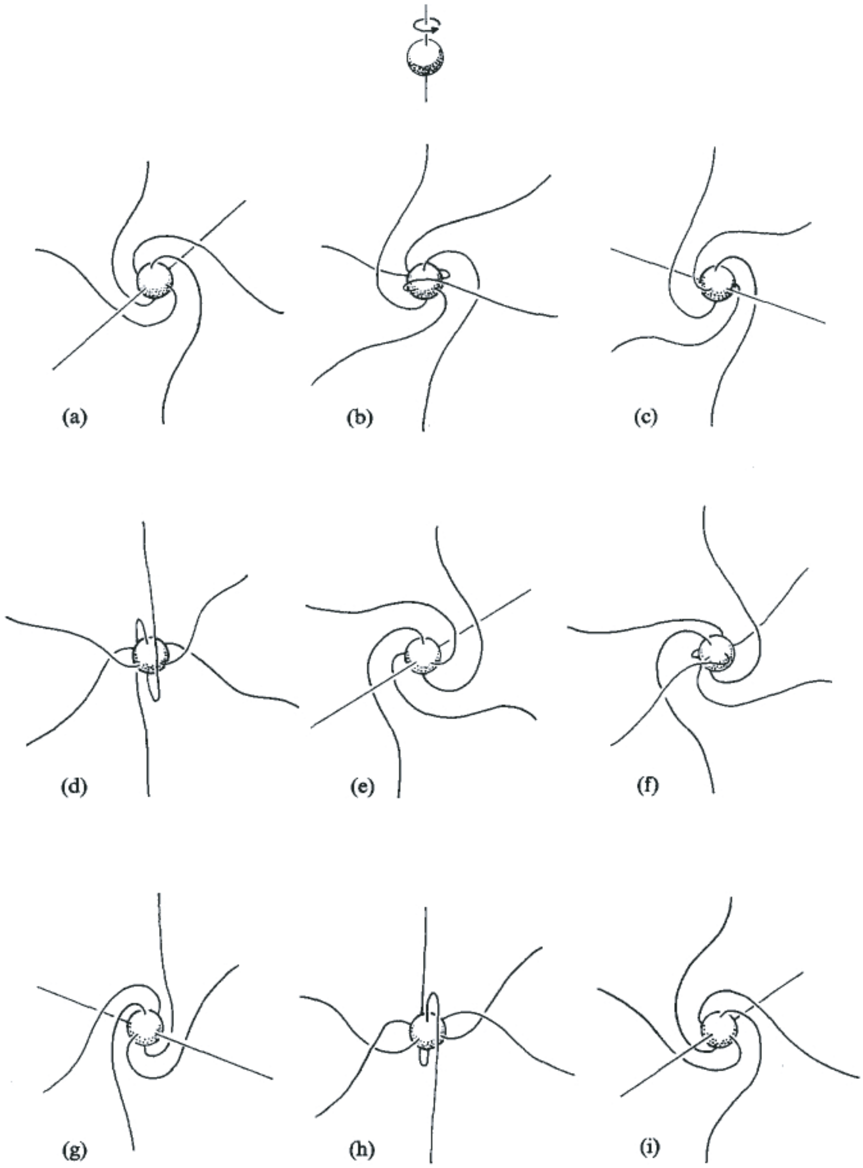

FIGURE 45 The belt trick for a rotating body with many tails, as used by Battey-Pratt and Racey to deduce the Dirac equation (© Springer Verlag, from Ref. 172).

This is the Dirac equation in the (less usual) spinorial representation.* The last term shows that mass mixes right and left chiralities. The equation can be expanded to include potentials using minimal coupling, in the same way as done above for the Schrödinger

* The matrix γ_5 is defined here as

$$\gamma_5 = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} , \tag{143}$$

where I is the 2×2 identity matrix.

and Pauli equations.

The above textbook derivation of the Dirac equation from usual quantum theory can be repeated and visualized also with the help of strands. There is no difference in arguments or results. The derivation with the help of strands was performed for the first time by Battey-Pratt and Racey, in 1980. They explored a central object connected by unobservable strands (or 'tails') to the border of space, as shown in Figure 45. In their approach, the central object plus the tails correspond to a quantum particle. The central object is assumed to be continuously rotating, thus reproducing spin 1/2. They also assumed that only the central object is observable. (In the strand model, the central object becomes the tangle core.) Battey-Pratt and Racey then explored a relativistically moving object of either chirality. They showed that a description of such an object requires four complex fields. Studying the evolution of the phases and axes for the chiral objects yields the Dirac equation. The derivation by Battey-Pratt and Racey is mathematically equivalent to the textbook derivation just given.

Ref. 172

We can thus say that the Dirac equation follows from the belt trick. We will visualize this connection in more detail in the next section. When the present author found this connection in 2008, Lou Kauffman pointed out the much earlier paper by Battey-Pratt and Racey. In fact, Paul Dirac was still alive when they found this connection, but unfortunately he did not answer their letter asking for comment.

In summary, tangles completely reproduce both the rotation and the linear motion of elementary fermions. Therefore, the strand model provides a simple view on the evolution equations of quantum theory. In the terms of the strand model, when spin is neglected, the Schrödinger equation describes the evolution of crossing density. For relativistic fermions, when the belt trick is included, the Dirac equation describes the evolution of crossing density. In fact, strands visualize these evolution equations in the most concrete way known so far.

Visualizing spinors and Dirac's equation using tangles

Despite its apparent complexity, the Dirac equation makes only a few statements: spin 1/2 particles are fermions, obey the relativistic energy–momentum relation, keep the quantum of action invariant, and thus behave like a wave. Each statement is visualized by the tangle model of fermions: tangles behave as spinors, the relativistic energy–momentum relation is built-in, the fundamental principle holds, and rotating tangle cores reproduce the evolution of the phase. Let us look at the details.

Given a particle tangle, the short-time fluctuations lead, after averaging of the crossings, to the wave function. The tangle model of fermions also provides a *visualization* of the *spinor* wave function. Indeed, at each point in space, the wave function has the following parameters:

— There is an average density $\rho(x, t)$; physically, this is the probability density. In the strand model, this is the local crossing density.
— There is a set of three Euler angles α, β and γ; physically, they describe the average local orientation and phase of the spin axis. In the strand model, this is the average local orientation and phase of the tangle core.
— There is a second set of three parameters $\boldsymbol{v} = (v_x, v_y, v_z)$; physically, they describe, at one's preference, either the average local Lorentz boost or a second set of three Euler

angles. In the strand model, these parameters describe the average local deformation of the core that is due to the Lorentz boost. It can also be seen as the axis around which the belt trick is performed.
— There is a phase δ; physically, this represents the relative importance of particle and antiparticle density. In the strand model, this phase describes with what probability the average local belt trick is performed right-handedly or left-handedly.

In total, these are eight real parameters; they correspond to one positive real number and seven phases. They lead to the description of a spinor wave function as
Ref. 173

$$\varphi = \sqrt{\rho}\, e^{i\delta}\, L(\boldsymbol{v})\, R(\alpha/2, \beta/2, \gamma/2) \,, \qquad (144)$$

where the product LR is an abbreviation for the boosted and rotated unit spinor and all parameters depend on space and time. This expression is equivalent to the description with four complex parameters used in most textbooks. In fact, this description of a spinor wave function and the related physical visualization of its density and its first six phases dates already from the 1960s. The visualisation can be deduced from the study of relativistic spinning tops or of relativistic fluids. Rotating tangles are more realistic, however. In contrast to all previous visualizations, the rotating tangle model explains also the last, seventh phase. This is the phase that describes matter and anti-matter, that explains the appearance of the quantum of action $\hbar$, and that explains the fermion behaviour.
Ref. 173

In short, only rotating tangles together with the fundamental principle provide a simple, complete and precise visualisation of spinor wave functions and their evolution. The tangle model for spinning relativistic quantum particles remains a simple extension of Feynman's idea to describe a quantum particle as a rotating little arrow. The arrow can be imagined as being attached to the rotating tangle core. The tails are needed to reproduce fermion behaviour. The specific type of tangle core determines the type of particle. The blurring of the crossings defines the wave function. Rotating arrows describe non-relativistic quantum physics; rotating tangles describe relativistic quantum physics.
Ref. 165

Visualizing spinor wave functions with tangles of strands helps the understanding of the Dirac equation in several ways.

1. Tangles support the view that elementary particles are little rotating entities, also in the relativistic case. This fact has been pointed out by many scholars over the years. The strand model provides a consistent visualization for these discussions.
Ref. 173
2. The belt trick can be seen as the mechanism underlying the famous Zitterbewegung that is part of the Dirac equation. The limitations in the observing the belt trick translate directly into the difficulties of observing the Zitterbewegung.
Ref. 174
3. The belt trick also visualizes why the velocity operator for a relativistic particle has eigenvalues $\pm c$.
4. The Compton length is often seen as the typical length at which quantum field effects take place. In the tangle model, it would correspond to the average size needed for the belt trick. The strand model thus suggests that the mass of a particle is related to the average size needed for the belt trick.
5. Tangles support the – at first sight bizarre – picture of elementary particles as little charges rotating around a centre of mass. Indeed, in the tangle model, particle rota-
Ref. 175

Page 177 tion requires a regular application of the belt trick of Figure 19, and the belt trick can be interpreted as inducing the rotation of a charge, defined by the tangle core, around a centre of mass, defined by the average of the core position. It can thus be helpful to use the strand model to visualize this description.

6. The tangle model can be seen as a vindication of the stochastic quantization research programme; quantum motion is the result of underlying fluctuations. For example, the similarity of the Schrödinger equation and the diffusion equation is modelled and explained by the strand model: since crossings can be rotated, diffusion of crossings leads to the imaginary unit that appears in the Schrödinger equation. Ref. 176

In short, rotating tangles are a correct underlying model for the propagation of fermions. And so far, tangles are also the only known correct model. *Tangles model propagators.* This modelling is possible because the Dirac equation results from only three ingredients:

— the relation between the quantum of action and the phase of the wave function (the wave behaviour),
— the relation between the quantum of action and spinor behaviour (the exchange behaviour),
— and the mass–energy relation of special relativity (the particle behaviour), itself due to the fundamental principle.

And all three ingredients are reproduced by the strand model. We see that the apparent complexity of the Dirac equation hides its fundamental simplicity. The strand model reproduces the ingredients of the Dirac equation, reproduces the equation itself, and makes the simplicity manifest. In fact, we can say:

▷ The Dirac equation describes the relativistic infinitesimal belt trick or string trick.

The belt trick is fundamental for understanding the Dirac equation. In the strand model, core rotations vary along two dimensions – the rotation is described by two angles – and so does the belt trick. The resulting four combinations form the four components of the Dirac spinor and of the Dirac equation. Page 180

In summary, tangles can be used as a precise visualization and explanation of quantum physics. Wave functions, also those of fermions, are *blurred tangles* – with the detail that not the strands, but their crossings are blurred.

Quantum mechanics vs. quantum field theory

Quantum mechanics is the approximation to quantum physics in which fields are continuous and particles are immutable. In the strand model, quantum mechanics is thus the approximation in which a particle is described by a tangle with a shape that is *fixed* in time. This approximation allows us to derive the Dirac equation, the Klein–Gordon equation, the Proca equation, the Pauli equation and the Schrödinger equation. In this approximation, the strand model for the electron in a hydrogen atom is illustrated in Figure 46. This approximation already will allow us to deduce the existence of the three gauge interactions, as we will see in the next chapter.

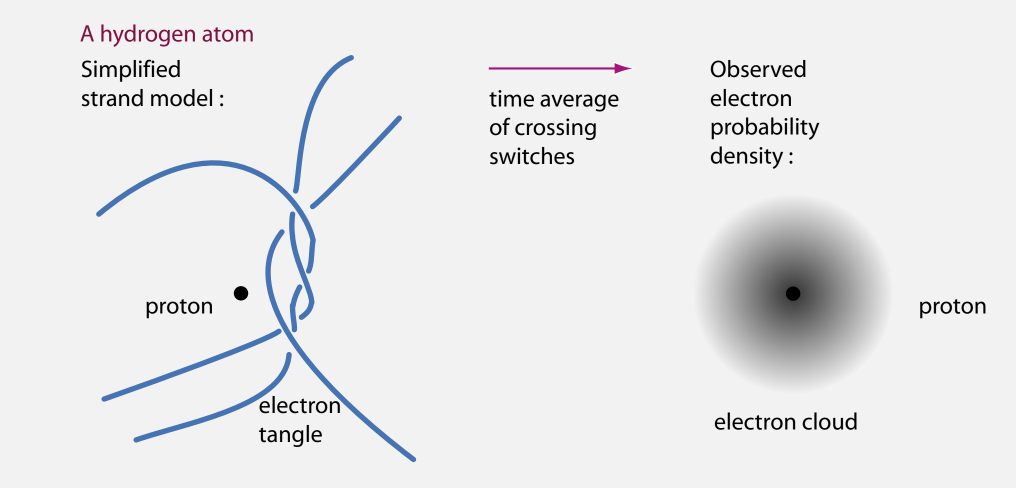

FIGURE 46 A simple, quantum-mechanical view of a hydrogen atom.

In contrast, *quantum field theory* is the description in which fields are themselves described by bosons, and particles types can transform into each other. The strand model allows us to deduce the existence of all known gauge bosons, as shown in the next chapter. In the strand description of quantum field theory, particles are not tangles with a fixed shape of their core, but for each particle, the shape *varies*. This variation leads to gauge boson emission and absorption.

A flashback: settling three paradoxes of Galilean physics

In all descriptions of physics, space and time are measured, explained and defined using matter. This occurs, for example, with the help of metre bars and clocks. On the other hand, matter is measured, explained and defined using space and time. This occurs, for example, by following a localized body over space and time. The circularity of the two definitions is at the basis of modern physics.

Page 110

As already mentioned above, the circularity is a natural consequence of the strand model. Both matter and space-time turn out to be approximations of the same basic building blocks; this common origin explains the apparent circular reasoning of Galilean physics. Most of all, the strand model changes it from a paradox to a logical necessity.

The strand model defines vacuum, and thus physical space, as a result of averaging strand crossings. Space is thus a *relative* concept. Newton's bucket experiment is sometimes seen as a counter-argument to this conclusion and as an argument for absolute space. However, the strand model shows that any turning object is connected to the rest of the universe through its tails. This connection makes every rotation an example of relative motion. Rotation is thus always performed relatively to the horizon of the universe. On the other hand, the detection of tangles among the tails allows a *local* determination of the rotation state, as is observed. Strands thus confirm that rotation and space are relative concepts. Strands thus also explain why we can turn ourselves on ice by rotating an arm over our head, without outside help. Strands lie to rest all issues around the rotating bucket.

A long time ago, Zeno of Elea based one of his paradoxes – the flying arrow that cannot reach the target – on an assumption that is usually taken as granted: he stated the impossibility to distinguish a short-time image (or state) of a *moving* body from the image (or state) of a *resting* body. The flattening of the tangles involved shows that the assumption is incorrect; motion and rest are *distinguishable*, even in (imagined) photographs taken with extremely short shutter times. The argument of Zeno is thus not possible, and the paradox disappears.

Fun challenges about quantum theory

> Urlaub ist die Fortsetzung des Familienlebens unter erschwerten Bedingungen.*
> Dieter Hildebrandt

Challenge 143 s Are the definitions for the addition and multiplication of Schrödinger wave functions that were given above also valid for spinor tangle functions?

* *

Challenge 144 e The definition of tangle functions, or wave functions, did not take into account the crossings of the vacuum strands, but only those of the particle tangle. Why is this allowed?

* *

Challenge 145 e Modelling the measurement of action at the quantum level as the counting of full turns of a wheel is a well-known idea that is used by good teachers to take the mystery out of quantum physics. The strand model visualizes this idea by assigning the quantum of action $\hbar$ to a full turn of one strand segment around another.

* *

Challenge 146 s Is any axiomatic system of quantum theory in contrast with the strand model?

* *

Challenge 147 s In the strand model, tangle energy is related to tangle core rotation. What is the difference between the angular frequency for tangles in the non-relativistic and in the relativistic case?

* *

Ref. 177 Challenge 148 e If you do not like the deduction of quantum mechanics given here, there is an alternative: you can deduce quantum mechanics in the way Schwinger did in his course, using the quantum action principle.

* *

Ref. 178 Challenge 149 r Modern teaching of the Dirac equation replaces the spinor picture with the vector picture. Hrvoje Nikolić showed that the vector picture significantly simplifies the understanding of Lorentz covariance of the Dirac equation. How does the vector picture clarify the relation between the belt trick and the Dirac equation?

* 'Vacation is the continuation of family life under aggravated conditions.' Dieter Hildebrandt (b. 1927 Bunzlau, d. 2013 Munich) was a cabaret artist, actor and author.

* *

In the strand description of quantum mechanics, strands are impenetrable: they cannot pass through each other (at finite distances). Can quantum mechanics also be derived if the model is changed and this process is allowed? Is entanglement still found?
Challenge 150 s

* *

A puzzle: Is the belt trick possible in a continuous and deformable medium – such as a sheet or a mattress – in which a coloured sphere is suspended? Is the belt trick possible with an *uncountably* infinite number of tails?
Challenge 151 e

* *

Page 181
At first sight, the apheresis machine diagram of Figure 24 suggests that, using the belt trick, animals could grow and use wheels instead of legs, because rotating wheels could be supplied with blood and connected to nerves. Why did wheels not evolve nevertheless?
Challenge 152 s

SUMMARY ON QUANTUM THEORY OF MATTER: EXPERIMENTAL PREDICTIONS

In this chapter, we used the fundamental principle – crossing switches define the quantum of action $\hbar$ and the other Planck units – to deduce that particles are tangles of strands and that wave functions are time-averaged rotating tangles. In simple words,

▷ Both non-relativistic and relativistic wave functions are *blurred rotating tangles*.

More precisely, a wave function appears from the blurred crossings of a tangle. The components and phases of the wave function at a point in space are due to the orientation and phase of crossings at that point. We also deduced that blurred tangles obey the least action principle and the Dirac equation.

In other words, visualizing the quantum of action as a crossing switch implies quantum theory. The strand model confirms Bohr's statement: quantum theory is indeed a consequence of the quantum of action. Specifically, the strand model thus shows that all quantum effects are *consequences of extension* and *consequences of the three dimensions of space*. More precisely, all quantum effects are *due to tails*, the tails of the tangles that represent a quantum system. In particular, the strand model confirms that

▷ The Dirac equation is essentially the infinitesimal version of the belt trick (or string trick).

In other words, strands also reproduce also the *propagator* of quantum particles.

As a result, we have shown that strands reproduce the relativistic Lagrangian density

$\mathcal{L}$ of charged, elementary, relativistic fermions in an external electromagnetic field $\boldsymbol{A}$

$$\mathcal{L} = \overline{\varphi}\,(i\hbar c \not{D} - c^2 m)\,\varphi\,, \tag{145}$$

where

$$\not{D} = \gamma^\sigma \mathrm{D}_\sigma = \gamma^\sigma(\partial_\sigma - iqA_\sigma)\,. \tag{146}$$

We thus conclude that *strands reproduce the quantum theory of matter*.

The strand model predicts deviations from the relativistic matter Lagrangian, and thus from the Dirac equation, *only* in three cases: first, when quantum aspects of electrodynamic field play a role, second, when nuclear interactions play a role, and third, when space curvature, i.e., strong gravity, plays a role. All this agrees with observation.

We will deduce the description of quantum electrodynamics and of the nuclear interactions in the next chapter. In the case of gravity, the strand model predicts that deviations from quantum theory occur exclusively when the energy–momentum of an elementary particle approaches the Planck value, i.e., for really strong gravity. Such deviations are not accessible to experiment at present. We will explore this situation in the subsequent chapter.

In addition, the strand model predicts that in nature, the Planck values for momentum and energy are limit values that cannot be exceeded by a quantum particle. All experiments agree with this prediction.

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The deduction of quantum theory from strands given here is, at present, the *only* known microscopic explanation for quantum physics. So far, no other microscopic model, no different explanation nor any other Planck-scale deduction of quantum theory has been found. In particular, the extension of fundamental entities – together with observability limited to crossing switches – is the key to understanding quantum physics.

Page 164

Let us evaluate the situation. In our quest to explain the open issues of the millennium list, we have explained the origin of Planck units, the origin of wave functions, the origin of the least action principle, the origin of space-time dimensions, the Lorentz and Poincaré symmetries, the origin of particle identity, and the simplest part of the Lagrangian of quantum field theory, namely, the Lagrangian of free fermions, such as the electron, and that of fermions in continuous external fields. Therefore, for the next leg, we turn to the most important parts of the standard model Lagrangian that are missing: those due to gauge interactions.

Chapter 9

GAUGE INTERACTIONS DEDUCED FROM STRANDS

What are interactions? At the start of this volume, when we summarized what relates the Planck units to relativity and to quantum theory, we pointed out that the nature of interactions at Planck scales was still in the Page 18 dark. In the year 2000, it was known for several decades that the essential properties of the electromagnetic, the weak and the strong nuclear interaction are their respective gauge symmetries: all three interactions are *gauge interactions*. But the underlying Ref. 179 reason for this property was still unknown.

In this chapter we discover that fluctuating strands in three spatial dimensions explain the existence of precisely three gauge interactions, each with precisely the gauge symmetry group that is observed. This is the first time ever that such an explanation Ref. 180 is possible. In other terms, we will deduce quantum field theory from strands. Indeed, strands provide a natural mechanism for interactions that explains and implies Feynman diagrams. The term 'mechanism' has to be taken with a grain of salt, because there is nothing mechanical involved; nevertheless, the term is not wrong, because we shall discover a surprisingly simple result: *Gauge interactions and gauge symmetries are due to specific strand deformations.*

In this chapter, we work in *flat* space-time, as is always done in quantum field theory. We leave the quantum aspects of *curved* space-time and of gravitation for the next chapter. We thus start by exploring the non-gravitational interactions in the quantum domain.

Interactions and phase change

Experiments in the quantum domain show that interactions *change the phase* of wave functions. But how precisely does this happen? The strand model will give us a simple answer: the emission and the absorption of gauge bosons is only possible *together* with a phase change. To explain this connection, we need to study the phase of tangle *cores* in more detail.

When we explored spin and its connection to the belt trick, we pictured the rotation Page 176 of the tangle core in the same way as the rotation of a belt buckle: we assumed that the core of the tangle rotates like a *rigid* object. The rotation is achieved through the shape fluctuations of the tails only. Why did we assume this?

In Feynman's description of quantum theory, *free particles are advancing rotating arrows.* Ref. 165 In the strand model, *free* particle motion is modelled as the change of position of the tangle core and *spin* as the rotation of the core. We boldly assumed that the core

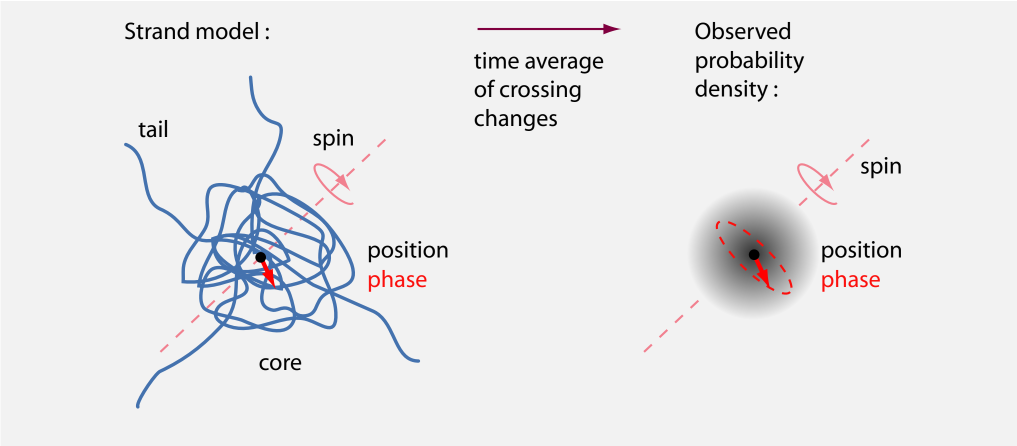

FIGURE 47 In the chapter on quantum theory, the phase was defined assuming a *rigidly rotating core*; this approximation was also used in the description of particle translation.

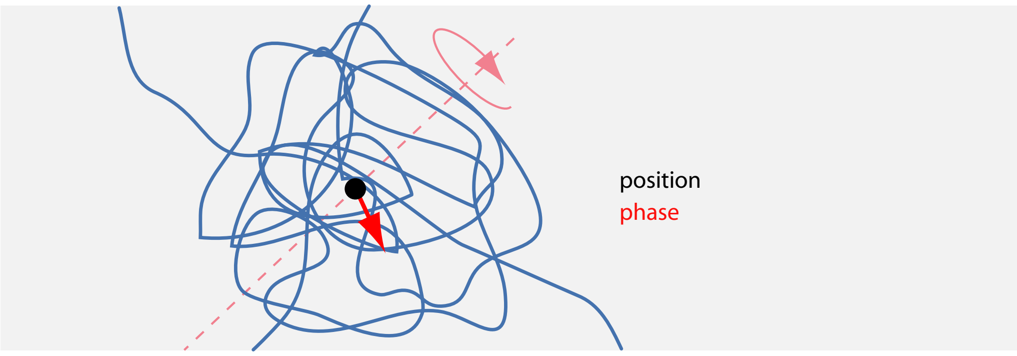

FIGURE 48 A magnified tangle core shows that the phase can also change due to *core deformations*; such core deformations lead to *gauge interactions*.

remained rigid, attached the phase arrow to it, and described spin as the rotation of the core with its attached arrow, as shown again in Figure 47. This bold simplification led us to the Dirac equation. In short, the assumption of a rigid core works. But what happens if the core is *not* rigid?

We know from observation and from quantum theory that

▷ An ***interaction*** is a process that changes the phase of a wave function, but differs from a rotation.

In the strand model, shape deformations of tangle cores also lead to phase changes – and such deformations differ from a rotation. In fact, we will discover that core deformations automatically lead to precisely those three gauge interactions that we observe in nature.

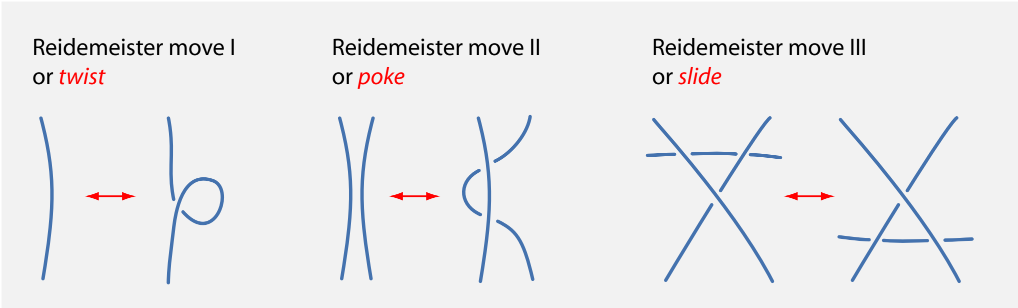

FIGURE 49 The Reidemeister moves: the three types of deformations that induce crossing switches – if the moves are properly defined in three dimensions.

Tail deformations versus core deformations

We can summarize the previous chapter, on the free motion of matter tangles, as the chapter that focused on shape fluctuations of *tails*. Indeed, the belt trick completed the proof that

▷ ***Space-time symmetries*** are due to *tail* deformations.

All space-time symmetries – translation, rotation, boost, spin and particle exchange – are due to tail deformations; in such tail deformations, the tangle core is assumed to remain unchanged and rigid (in its own rest frame).

In contrast, the present chapter focuses on shape fluctuations in *tangle cores*.We will discover that

▷ ***Gauge symmetries*** are due to *core* deformations.

Let us explore the tangle core in more detail. Figure 48 shows a magnified view of the core and its phase arrow. The phase of the core results from the phases of all its crossings. The figure illustrates that the phase arrow will be sensitive to the shape fluctuations and deformations of the strand segments that make up the core.

In nature, any phase change of the wave function that is not due to a space-time symmetry is due to an interaction. For the strand model, this connection implies:

▷ When the phase of a core changes through *rigid orientation change*, we speak of ***core rotation***.

▷ When the phase of a core changes through *core shape deformation*, we speak of ***interaction***.

We thus need to understand two things: First, what kinds of core deformation exist? Secondly, how precisely is the phase – i.e., each arrow definition – influenced by core deformations? In particular, we have to check the answers and deductions with experiment.

The first question, on the classification of the core deformations, is less hard than

Ref. 180 it might appear. The fundamental principle – events are crossing switches of strands – implies that deformations are observable only if they induce crossing switches. Other deformations do not have any physical effect. (Of course, certain deformations will have crossing switches for one observer and none for another. We will take this fact into consideration.) Already in 1926, the mathematician Kurt Reidemeister classified all those
Ref. 182 tangle deformations that lead to crossing switches. The classification yields exactly three classes of deformations, today called the three *Reidemeister moves*. They are shown in Figure 49.

- ▷ The ***first Reidemeister move***, or ***type I move***, or ***twist***, is the addition or removal of a twist in a strand.
- ▷ The ***second Reidemeister move***, or ***type II move***, or ***poke***, is the addition or removal of a bend of one strand under (or over) a second strand.
- ▷ The ***third Reidemeister move***, or ***type III move***, or ***slide***, is the displacement of one strand segment under (or over) the crossing of two other strands.

The type number of each Reidemeister move is also the number of involved strands. We will discover that despite appearances, each Reidemeister move induces a crossing switch. To find this connection, we have to generalize the original Reidemeister moves, which were defined in a two-dimensional projection plane, to the three-dimensional situation of tangle cores.

The three Reidemeister moves turn out to be related to the three gauge interactions:

- ▷ The first Reidemeister move corresponds to *electromagnetism*. The second Reidemeister move corresponds to the *weak nuclear interaction*. The third Reidemeister move corresponds to the *strong nuclear interaction*.

We will prove this correspondence in the following.

For each Reidemeister move we will explore two types of core deformation processes: One deformation type are *core fluctuations*, which correspond, as we will see, to the emission and absorption of *virtual* interaction bosons. The other deformations are *externally induced core disturbances*, which correspond to the emission and absorption of *real* interaction bosons. As the first step, we show that both for fluctuations and for disturbances, the first Reidemeister move, the twist, is related to the electromagnetic interaction.

Electrodynamics and the first Reidemeister move

Experiments show that electromagnetism is described by potentials. Experiments also show that potentials change the phase, the rotation frequency and the wave number of wave functions. In particular, for electromagnetism, the potentials are due to the flow of real and virtual, massless, uncharged spin-1 photons. Photons are emitted from or absorbed by charged elementary particles; neutral elementary particles do not emit or absorb photons. There are two types of electric charge, positive and negative. The attrac-

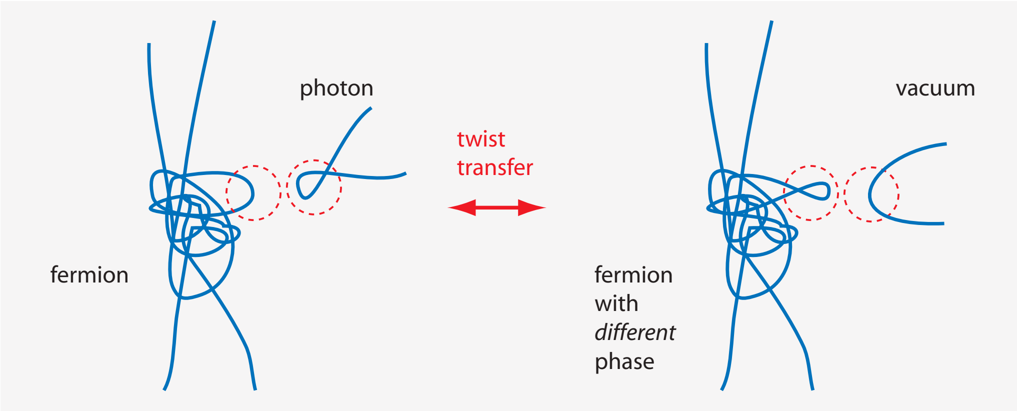

FIGURE 50 A single strand changes the rotation of a tangle: *twist transfer* is the basis of electromagnetism in the strand model. No strand is cut or reglued; the transfer occurs statistically, through the excluded volume due to the impenetrability of strands. Twist transfer generates a U(1) gauge group, as explained in the text.

tion and repulsion of static charges diminishes with the inverse square of the distance. Charge is conserved. All charged particles are massive and move slower than light. The Lagrangian of matter coupled to the electromagnetic field has a U(1) gauge symmetry – it is described by minimal coupling. Electromagnetism has a single fundamental Feynman diagram. The electromagnetic coupling constant at low energy, the so-called *fine structure constant*, is measured to be $\alpha = 1/137.035\,999\,139(31)$; its energy dependence is described by renormalization.

Ref. 5

The previous paragraph contains everything known about the electromagnetic interaction. For example, Maxwell's field equations follow from Coulomb's inverse square relation, its relativistic generalization, and the conservation of charge . More precisely, all experimental observations about electricity and magnetism follow from the Lagrangian of quantum electrodynamics, or QED. In short, we now need to show that the Lagrangian of QED follows from the strand model.

Ref. 183

STRANDS AND THE TWIST, THE FIRST REIDEMEISTER MOVE

In the strand model of electromagnetism, massless spin 1 bosons such as the photon are made of a single strand. How can a single strand change the phase of a tangle? The answer is given in Figure 50: a *twisted loop* in a single strand will influence the rotation of a tangle because it changes the possible shape fluctuations of the tangle core. Due to the impenetrability of strands, an approaching twisted loop will sometimes transfer its twist to the tangle: this process will deform the tangle core and thereby change its phase. The observed effect of an electromagnetic field on the phase of a charged fermion is the *time average* of all such twist transfers.

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Single strands represent bosons, as we saw above. Rotating the core of the twist around the tethers by 2π gives back the original configuration: twists have spin 1. Twisted loops are single strands and can have *two* twist senses, or two polarizations. Single, twisted and *unknotted* strands have no mass; in other words, twisted loops effectively move with the speed of light. And twisted loops, being curved, carry energy.

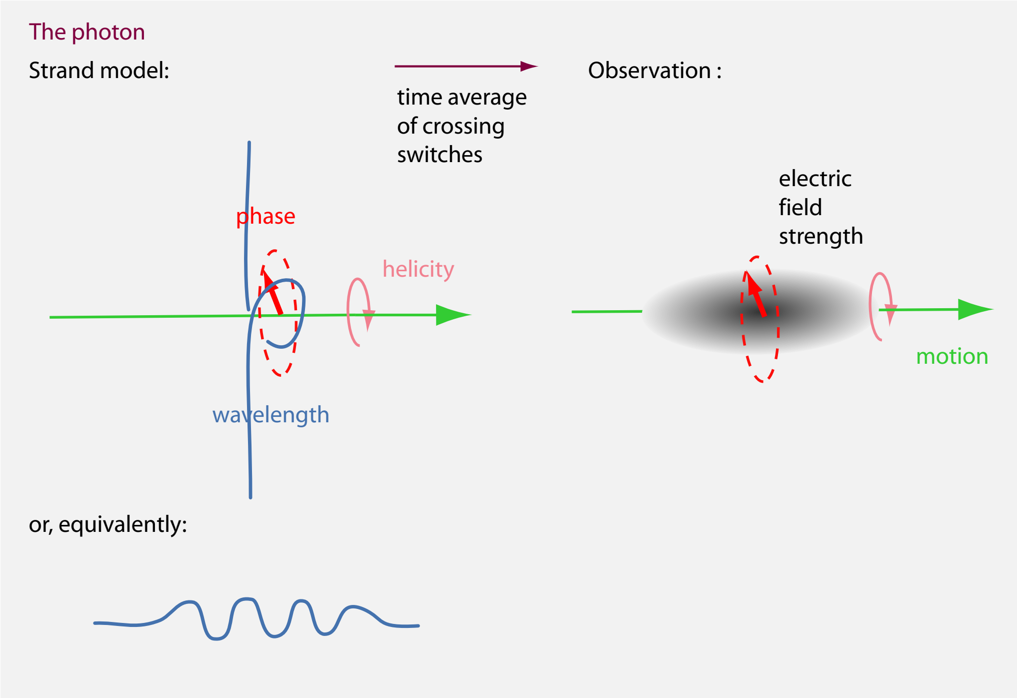

FIGURE 51 The photon in the strand model.

Approaching twisted loops will change the phase, i.e., the orientation of a matter tangle. Twisted loops correspond to a local rotation of a strand segment by π. But twists can be generalized to arbitrary angles. These generalized twists can be concatenated. Because they are described by a single angle, and because a double twist is equivalent to no twist at all, twists form a U(1) group. We show this in detail shortly. Page 235

In summary, twists behave like *photons* in all their properties. Therefore, the strand model suggests:

▷ A ***photon*** is a twisted strand. An illustration is given in Figure 51.

▷ The ***electromagnetic interaction*** is the transfer of twists, i.e., the transfer of first Reidemeister moves, between two particles, as shown in Figure 50.

The transfer of a twist from a single strand to a tangle core thus models the absorption of a photon. We stress again that this transfer results from the way that strands hinder each other's motion, because of their impenetrability. No strand is ever cut or reglued.

Can photons decay, disappear or break up?

The strand model of the photon, as shown in Figure 51, might be seen to suggest that photons can disappear. For example, if a photon strand is straightened out by pulling the ends of the helical deformation, the helix might disappear. A helix might also disappear by a shape fluctuation or transform into several helices. However, this is a fallacy.

A lone twist cannot disappear by pulling; ''pulling'' requires an apparatus that performs it. That is impossible. A lone twist cannot disappear by fluctuations either, because a photon also includes the vacuum strands around it. In the strand model, the energy of the photon is localized in the configuration formed by the photon strand and the surrounding vacuum strands. In the strand model, energy is localized in regions of strand curvature. If the helical strands disappears, the surrounding vacuum strands are curved instead, or more strongly, and the twist energy is taken up by these surrounding strands. The net result is that the helix is transferred, permanently or for a short time, to another strand. In other terms, in the strand model, photons can also move by hopping from one strand to the next.

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Also, a single photon strand cannot break up into *several* photon strands of smaller helical diameters or of different rotation frequencies. Such a process is prevented by the fundamental principle, when the vacuum is taken into account.

Challenge 153 e

The only way in which a photon can disappear completely is by transferring its crossing, i.e., its energy to a tangle. Such a process is called the *absorption* of a photon by a charged particle.

In short, due to energy and to topological restrictions, the strand model prevents the decay, disappearance or splitting of photons, as long as no electric charge is involved. Linear and angular momentum conservation also lead to the same conclusion. Photons are *stable* particles in the strand model.

Electric charge

Surrounded by a bath of photon strands, not all fermion tangles will change their phase. A tangle subject to randomly approaching virtual photons will feel a net effect over time only if it lacks some symmetry. In other words, only tangles that lack a certain symmetry will be electrically charged. Which symmetry will this be?

In a bath of photon strands, thus in a bath that induces random Reidemeister I moves, only *chiral* fermion tangles are expected to be influenced. In other terms:

▷ ***Electric charge*** is due to lack of mirror symmetry, i.e., to tangle chirality.

Conversely, we have:

▷ ***Electrically charged particles*** randomly emit twisted strands. Due to the tangle chirality, a random emission will lead to a slight asymmetry, so that right-handed twists will be in the majority for particles of one charge, and left-handed twists will be in the majority for particles of the opposite charge.

Equating electric charge with tangle chirality allows modelling several important observations. First, because chirality can be right-handed or left-handed, there are positive and negative charges. Second, because strands are never cut or reglued in the strand model, chirality, and thus electric charge, is a *conserved quantity*. Third, chirality is only possible for tangles that are localized, and thus massive. Therefore, chiral tangles – charged particles – always move slower than light. Fourth, a chiral tangle at rest induces a twisted strand density around it that changes as $1/r^2$, as is illustrated in Figure 52. Finally,

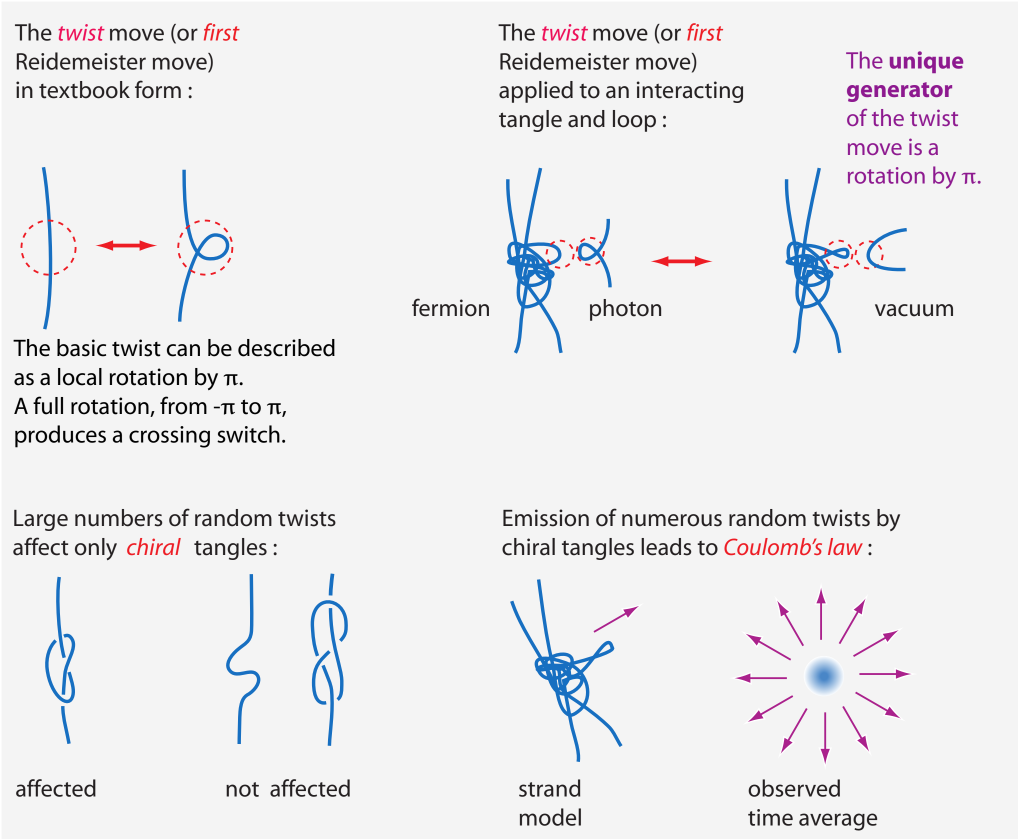

FIGURE 52 Electromagnetism in the strand model: the electromagnetic interaction, electric charge and Coulomb's inverse square relation. (This image needs to be updated: no knotted tangles occur in nature.)

photons are uncharged; thus they are not influenced by other photons (to first order).

In short, all properties of electric charge found in nature are reproduced by the tangle model. We now check this in more detail.

Challenge: What topological invariant is electric charge?

Chirality explains the sign of electric charge, but not its magnitude in units of the elementary charge e. A full definition of electric charge must include this aspect.

Mathematicians defined various topological invariants for knot and tangles. *Topological invariants* are properties that are independent of the shape of the knot or tangle, but allow to distinguish knots or tangles that differ in the ways they are knotted or tangled up. Several invariants are candidates as building blocks for electric charge: *chirality* c, which can be $+1$ or -1, *minimal crossing number* n, or *topological writhe* w, i.e., the signed minimal crossing number.

Page 437

A definition of electric charge q, proposed by Claus Ernst, is $q = c\,(n \bmod 2)$. Another option for the definition of charge is $q = w/3$. Equivalent definitions use the linking number. At this point of our exploration, the issue is open. We will come back to the

Page 390 detailed connection between charge, chirality and tangle topology later on.

Electric and magnetic fields and potentials

The definition of photons with twisted strands leads to the following definition.

▷ The ***electric field*** is the volume density of (oriented) crossings of twisted loops.

▷ The ***magnetic field*** is the flow density of (oriented) crossings of twisted loops.

▷ The ***electric potential*** is the density of twisted loops.

▷ The ***magnetic potential*** is the flow density of twisted loops.

The simplest way to check these definitions is to note that the random emission of twisted loops by electric charges yields Coulomb's inverse square relation: the force between two static spherical charges changes with inverse square of the distance. The strand model implies that in this case, *the crossing density is proportional to the square of the loop density*; in other words, the potential falls of as the inverse distance, and the electric field as the square distance.

The definition of the magnetic field simply follows from that of the electric field by changing to moving frame of reference. The two field definitions are illustrated in Figure 53.

We note that the electric field is defined almost in the same way as the wave function: Page 186 both are oriented crossing densities. However, the electric field is defined with the crossing density of *twisted loops*, whereas the wave function is defined with the crossing density of *tangles*. The definitions differ only by the topology of the underlying strand structures.

In the strand model, energy, or action per time, is the number of crossing switches *per time*. The electromagnetic field energy per volume is thus given by the density of crossing switches *per time* that are due to twisted loops. Now, the strand model implies that *the crossing switch density per time is given by half the square of the crossing density* Challenge 154 e *plus half the square of the crossing density flow*. For twisted loops, we thus get that the energy density is half the square of the electric plus half the square of the magnetic field. Inserting the proportionality factors that lead from Planck units to SI units we get the well-known expression

$$\frac{E}{V} = \frac{\varepsilon_0}{2} E^2 + \frac{1}{2\mu_0} B^2 . \tag{147}$$

The strand model thus reproduces electromagnetic energy.

We note that in the strand model, the definition of the fields implies that there is no *magnetic charge* in nature. This agrees with observation.

The strand model predicts limit values to all observables. They always appear when strands are as closely packed as possible. This implies a maximum electric field value $E_{max} = c^4/4Ge \approx 1.9 \cdot 10^{62}$ V/m and a maximum magnetic field value $B_{max} = c^3/4Ge \approx 6.3 \cdot 10^{53}$ T. All physical systems – including all astrophysical objects, such as gamma-ray

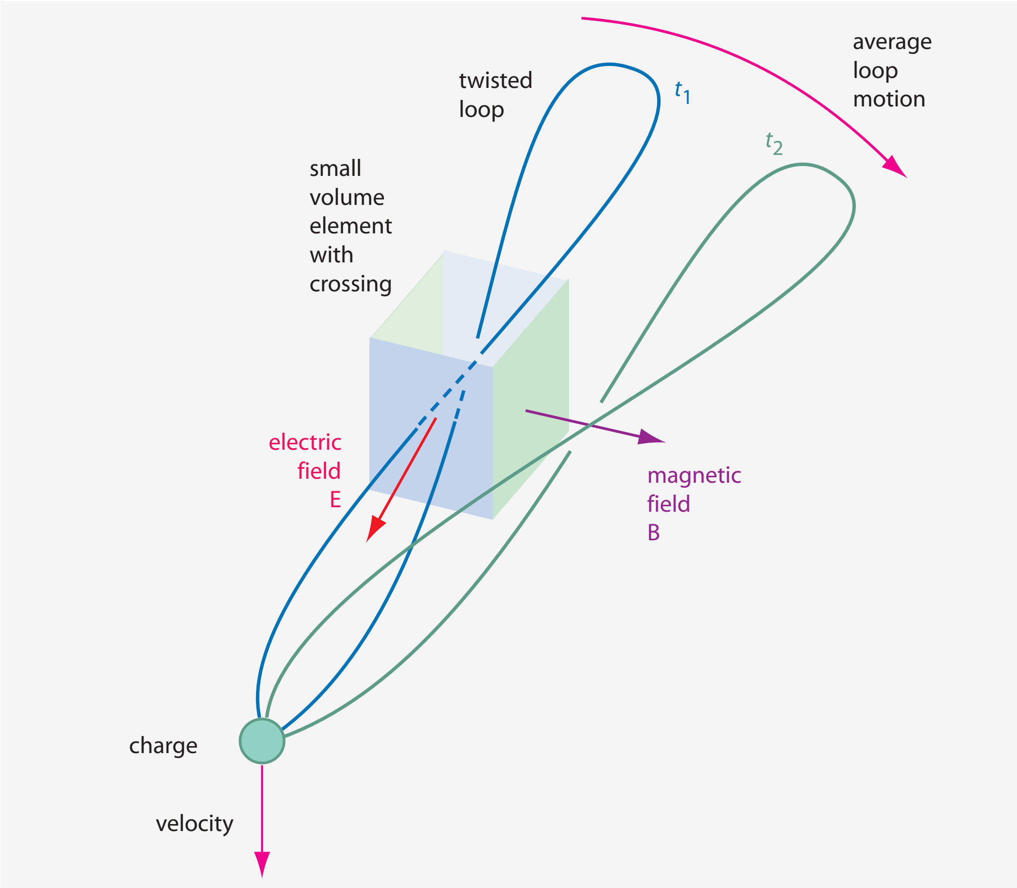

FIGURE 53 Moving twists allow us to define electric fields – as the density of twisted loop crossings – and magnetic fields – as the corresponding flow.

bursters or quasars – are predicted to conform to this limit. This strand model prediction indeed agrees with observations so far.

The Lagrangian of the electromagnetic field

In classical electrodynamics, the energy density of the electromagnetic field is used to deduce its Lagrangian density. The Lagrangian density describes the intrinsic, observer-independent change that occurs in a system. In addition, the Lagrangian density must be quadratic in the fields and be a Lorentz-scalar.

A precise version of these arguments leads to the Lagrangian density of the electromagnetic field F

$$\mathcal{L}_{\mathrm{EM}} = \frac{\varepsilon_0}{2} E^2 - \frac{1}{2\mu_0} B^2 = -\frac{1}{4\mu_0} F_{\mu\nu} F^{\mu\nu} \qquad (148)$$

where the electromagnetic field F is defined with the electromagnetic potential A as

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \; . \qquad (149)$$

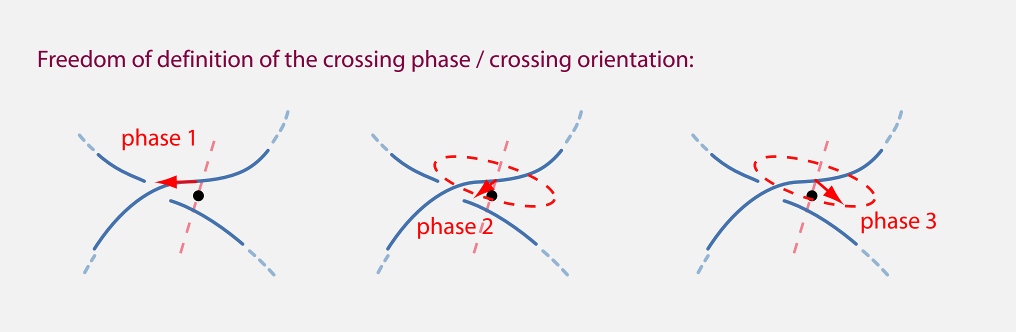

FIGURE 54 The definition of the phase or orientation of a single crossing is not unique: there is a freedom of choice.

Since the strand model reproduces the electromagnetic energy, it also reproduces the Lagrangian of classical electrodynamics. In particular, Maxwell's equations for the electromagnetic field follow from this Lagrangian density. Maxwell's field equations are thus a consequence of the strand model. Obviously, this is no news, because any model that reproduces Coulomb's inverse square distance relation and leaves the speed of light invariant automatically contains Maxwell's field equations. Ref. 183

U(1) GAUGE INVARIANCE INDUCED BY TWISTS

In nature, the electromagnetic potential A_μ is not uniquely defined: one says that there is a freedom in the choice of gauge. The change from one gauge to another is a *gauge transformation*. Gauge transformations are thus transformations of the electromagnetic potential that have no effect on observations. In particular, gauge transformations leave unchanged all field intensities and field energies on the one hand and particle probabilities and particle energies on the other hand.

All these observations can be reproduced with strands. In the strand model, the following definitions are natural:

- ▷ A ***gauge choice*** for *radiation* and for *matter* is the choice of definition of the respective phase arrow.
- ▷ A ***gauge transformation*** is a change of definition of the phase arrow.

In the case of electrodynamics, the gauge freedom is a result of allowing phase choices that lie in a plane around the crossing orientation. (The other interactions follow from the other possible phase choices.) The phase choice can be different at every point in space. Changing the (local) phase definition is a (local) gauge transformation. Changing the phase definition for a single crossing implies changing the phase of wave functions and of the electromagnetic potentials. A schematic illustration of the choice of gauge is given in Figure 54 and Figure 55.

We note that gauge transformations have no effect on the density or flow of crossings or crossing switches. In other words, gauge transformations leave electromagnetic field intensities and electromagnetic field energy invariant, as observed. Similarly, gauge transformations have no effect on the number of crossing switches of rotating tangles.

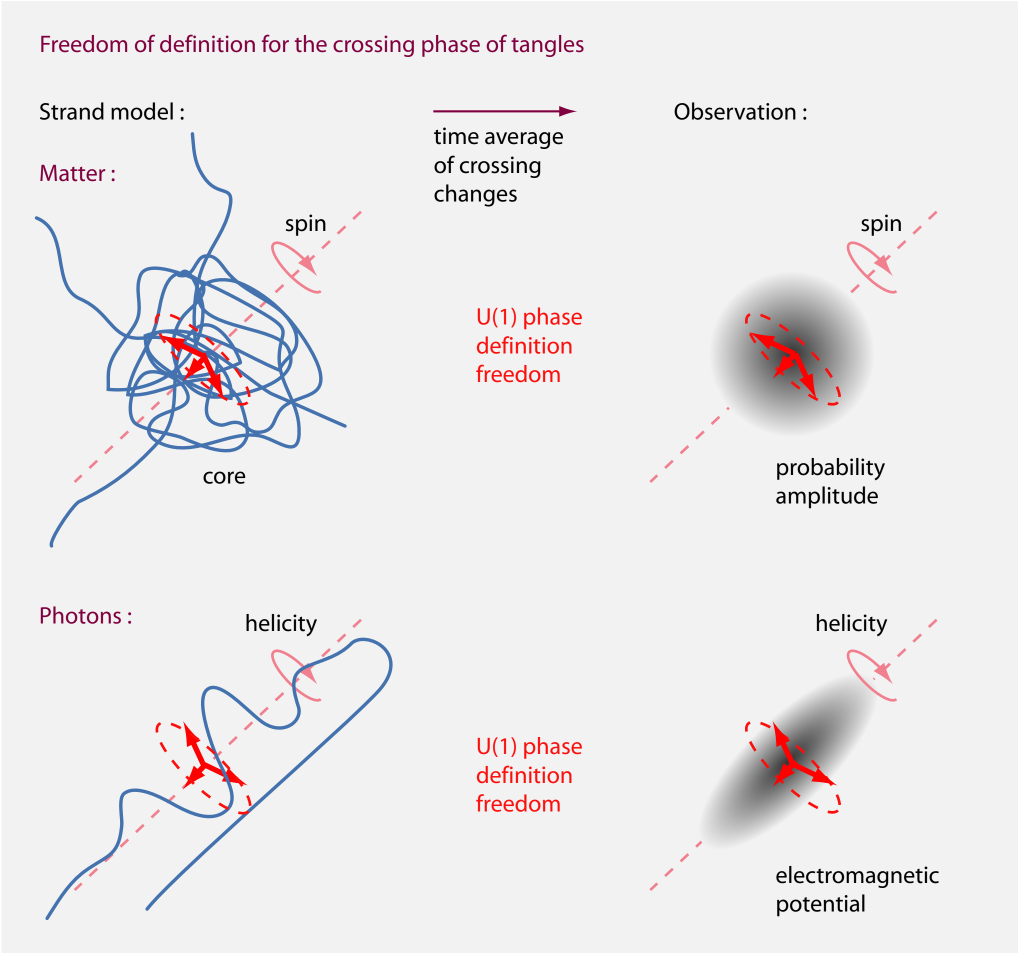

FIGURE 55 The freedom in definition of the phase of crossings leads to the gauge invariance of electrodynamics. Three exemplary choices of phase are shown.

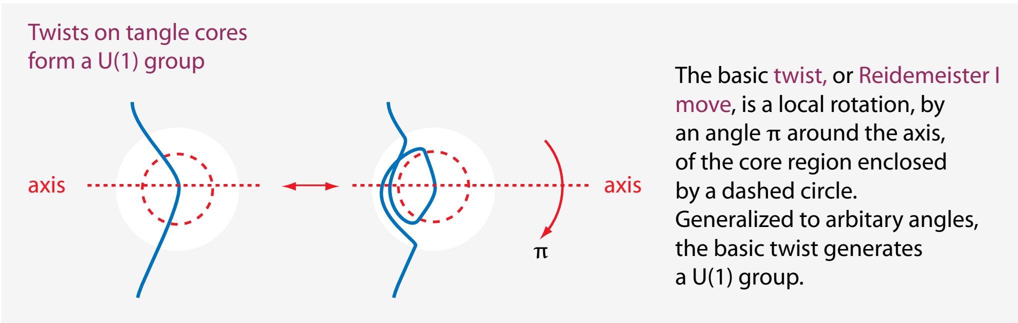

FIGURE 56 How the set of generalized twists – the set of all local rotations of a single strand segment around an axis – forms a U(1) gauge group.

A rotation by 4π does not change the phase, independently of which definition of arrow is chosen. Therefore, gauge transformations leave probability densities – and even observable phase differences – unchanged. This agrees with experiment.

A gauge transformation on a wave functions also implies a gauge transformation on the electrodynamic potential. The strand model thus implies that the two transformations are connected, as is observed. This connection is called *minimal coupling*. In short, minimal coupling is a consequence of the strand model.

U(1) GAUGE INTERACTIONS INDUCED BY TWISTS

There is only a small step from a gauge *choice* to a gauge *interaction*. We recall:

▷ A ***gauge interaction*** is a change of phase resulting from a strand deformation of the particle core.

In particular, electromagnetism results from the transfer of *twists*; twists are one of the three types of core deformations that lead to a crossing switch.

The basic twist, or first Reidemeister move, corresponds to a local rotation of some strand segment in the core by an angle π, as illustrated by Figure 56. Twists can be generalized to arbitrary angles: we simply define a *generalized twist* as a local rotation of a strand segment by an arbitrary angle. The rotation axis is chosen as shown in Figure 56. Generalized twists can be concatenated, and the identity twist – no local rotation at all – also exists. Generalized twists thus form a group. Furthermore, a generalized twist by 2π is equivalent to no twist at all, as is easily checked with a piece of rope: keeping the centre region is it disappears by pulling the ends, in contrast to a twist by π.

Challenge 155 e

Generalized twists thus behave like $e^{i\vartheta}$. Their concatenation produces a multiplication table

$$\begin{array}{c|c} \cdot & e^{i\pi} \\ \hline e^{i\pi} & 1 \end{array} \qquad (150)$$

that generate and define a U(1) group. In other words, Figure 56 shows that generalized twists define the group U(1), which has the topology of a circle.

In summary, the addition of a twist to a fermion tangle or to a photon strand changes their phase, and thus represents a gauge interaction. We have shown that core fluctuations induced by twists produce a U(1) gauge symmetry. Electromagnetic field energy and particle energy are U(1) invariant. In short, the strand model implies that *the gauge group of quantum electrodynamics is U(1).* With this result, we are now able to deduce the full Lagrangian of QED.

THE LAGRANGIAN OF QED

Given the U(1) gauge invariance of observables, the Lagrangian of quantum electrodynamics, or QED, follows directly, because U(1) gauge invariance is equivalent to minimal coupling. We start from the Lagrangian density $\mathcal{L}$ of a *neutral*, *free*, and relativistic fermion in an electromagnetic field. It is given by

$$\mathcal{L} = \overline{\Psi}(i\hbar c\partial\!\!\!/ - c^2 m)\Psi - \frac{1}{4\mu_0}F_{\mu\nu}F^{\mu\nu} \,. \qquad (151)$$

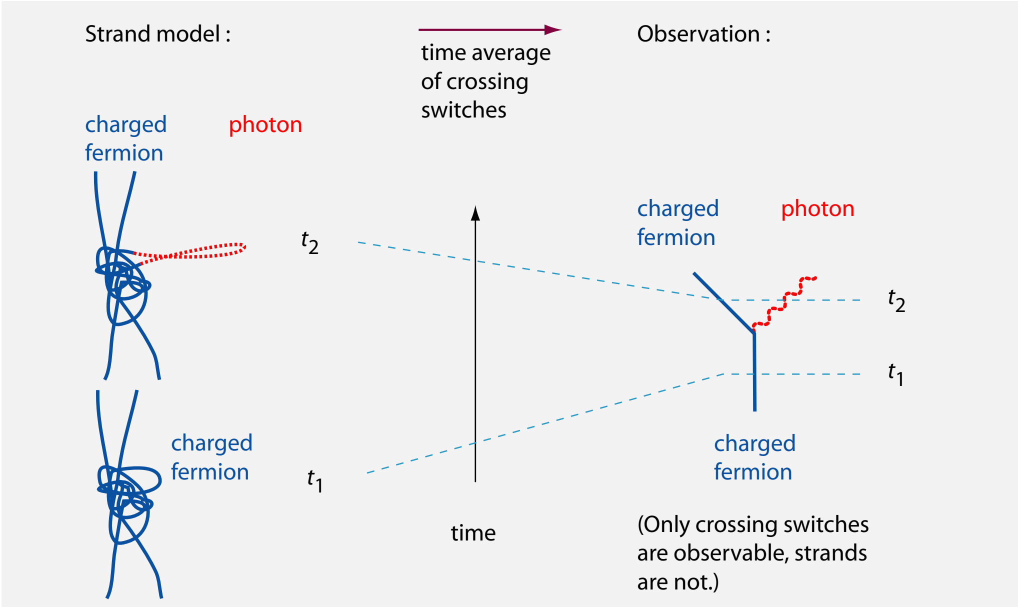

FIGURE 57 The fundamental Feynman diagram of QED and its tangle version.

Page 223 We deduced the fermion term in the chapter of quantum theory, and we deduced the electromagnetic term just now, from the properties of twisted loops.

As we have seen, the strand model implies minimal coupling. This changes the Lagrangian density for a *charged*, i.e., *interacting*, relativistic fermion in the electromagnetic field, into the Lagrangian density of QED:

$$\mathcal{L}_{\text{QED}} = \overline{\Psi}(i\hbar c \not{D} - c^2 m)\Psi - \frac{1}{4\mu_0} F_{\mu\nu} F^{\mu\nu} . \tag{152}$$

Here, $\not{D} = \gamma^\sigma D_\sigma$ is the *gauge covariant derivative* that is defined through minimal coupling to the charge q:

$$D_\sigma = \partial_\sigma - iqA_\sigma . \tag{153}$$

Minimal coupling implies that the Lagrangian density of QED is invariant under U(1) gauge transformations. We will discuss the details of the charge q later on. Page 382

We have thus recovered the Lagrangian density of quantum electrodynamics from strands. Strands thus reproduce the most precisely tested theory of physics.

Feynman diagrams and renormalization

Feynman diagrams are abbreviations of formulas to calculate effects of quantum electrodynamics in perturbation expansion. Feynman diagrams follow from the Lagrangian of QED. All Feynman diagrams of QED can be constructed from one fundamental diagram, shown on the right-hand side of Figure 57. Important Feynman diagrams are shown on the left-hand sides of Figure 58 and of Figure 59.

In the strand model, the fundamental Feynman diagram can be visualized directly

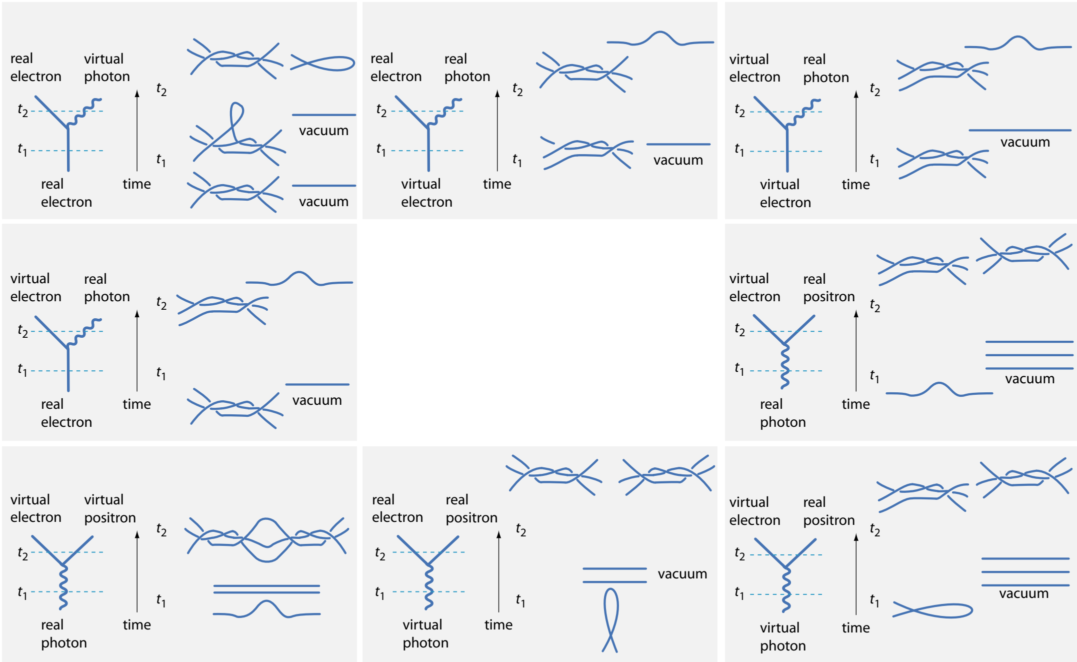

FIGURE 58 The different variations of the fundamental Feynman diagram of QED and their tangle versions.

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Page 314

in terms of strands, as shown on the left-hand side of Figure 57. This is the same diagram that we have explored right at the start of the section on electrodynamics, when we defined electrodynamics as twist exchange. (The precise tangles for the charged fermions will be deduced later on.) Since all possible Feynman diagrams are constructed from the fundamental diagram, the strand model allows us to interpret all possible Feynman diagrams as strand diagrams. For example, the strand model implies that the vacuum is full of virtual particle-antiparticle pairs, as shown in Figure 59.

In quantum field theory, Lagrangians must not only be Lorentz and gauge invariant, but must also be renormalizable. The strand model makes several statements on this issue. At this point, we focus on QED only; the other gauge interactions will be treated below. The strand model reproduces the QED Lagrangian, which is renormalizable. Renormalizability is a natural consequence of the strand model in the limit that strand diameters are negligible. The reason for renormalizability that the strand model reproduces the single, fundamental Feynman diagram of QED, without allowing other types of diagrams.

The twist deformations underlying the strand model for QED also suggest new ways to calculate higher order Feynman diagrams. Such ways are useful in calculations of g-factors of charged particles, as shown in the next section. In particular, the strand model for QED, as shown in Figure 57, implies that higher order QED diagrams are simple *strand deformations* of lower order diagrams. Taking statistical averages of strand deformations up to a given number of crossings thus allows us to calculate QED effects up to a given order in the coupling. The strand model thus suggests that non-perturbative calculations are possible in QED.

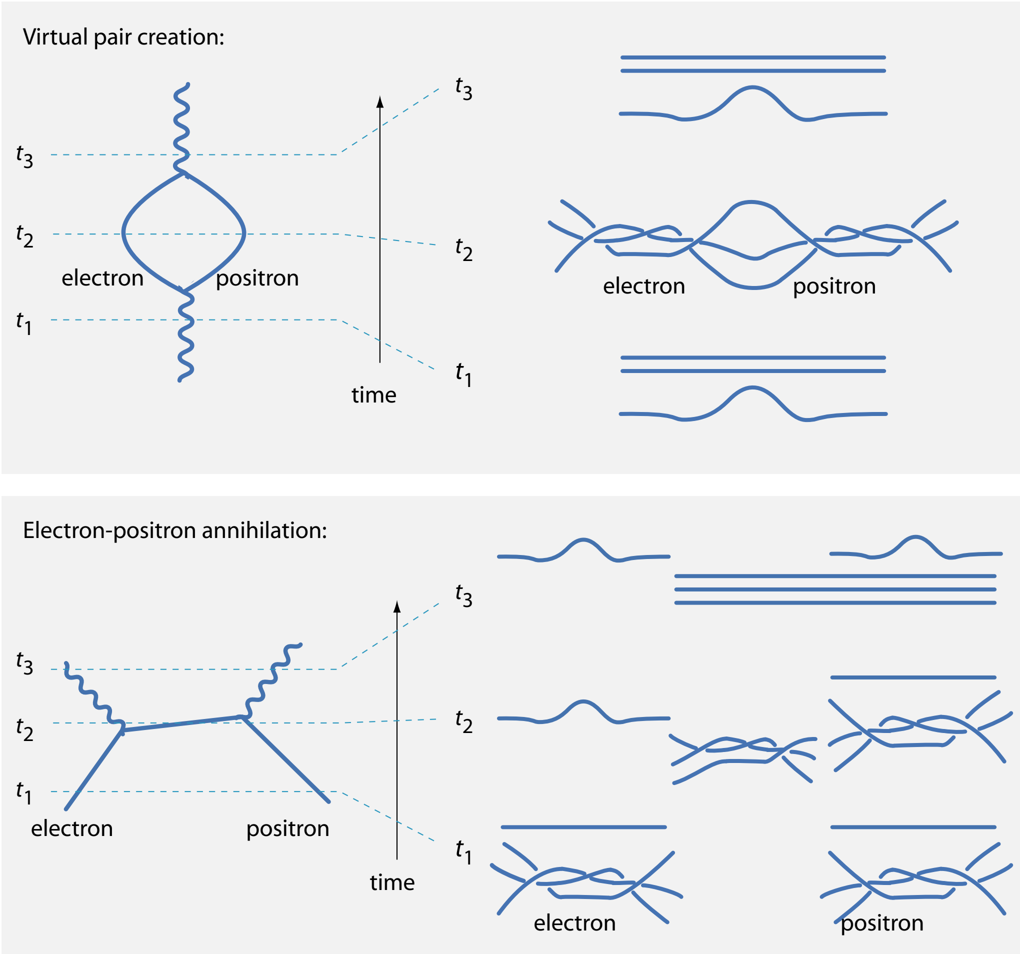

FIGURE 59 Some Feynman diagrams of QED with their tangle versions.

For precise non-perturbative calculations, the effective diameter of the strands must be taken into account. The diameter eliminates the Landau pole and all ultraviolet divergences of QED. In the strand model, the vacuum energy of the electromagnetic field is automatically zero. In other words, the strand model eliminates all problems of QED; in fact, QED appears as an approximation of the strand model for negligible strand diameter. In passing, we thus predict that perturbation theory for QED is valid and *converges* if the strand model, and in particular the finite strand diameter, is taken into account. (The diameter is the only gravitational influence predicted to affect QED.) However, we do not pursue these topics in the present text.

The strand model also suggests that the difference between renormalized and unrenormalized mass and charge is related to the difference between minimal and non-minimal crossing switch number, or equivalently, between tangle deformations with few and with many crossings, where strands are deformed on smaller distance scales. In other terms, unrenormalized quantities – the so-called *bare* quantities at Planck energy – can be imagined as those deduced when the tangles are pulled tight, i.e., pulled to Planck distances, whereas renormalized mass and charge values are those deduced for particles surrounded by many large-size fluctuations.

The strand model also suggests a visualization for the cut-off used in QED. The cut-off is a characteristic energy or length used in intermediate calculations. In the strand model, the cut-off corresponds to the size of the image.

In summary, the strand model provides a new underlying picture or mechanism for Feynman diagrams. The strand model does not change any physical result at any experimentally accessible energy scale. In particular, the measured change or 'running' with Page 386 energy of the fine structure constant and of the masses of charged particles are reproduced by the strand model, because Feynman diagrams of all orders are reproduced up to energies just below the Planck scale. Deviations between QED and the strand model are only expected near the Planck energy, when tangles of Planck diameter are pulled tight.

The anomalous magnetic moment

The anomalous magnetic moment g of the electron and of the muon is given by the well-known expression

$$\frac{g}{2} = 1 + \frac{\alpha}{2\pi} - O(\alpha^2) , \tag{154}$$

where $g/2$ is half the so-called g-factor, with a measured value of 1.00116(1), and α is the fine structure constant, with a measured value of 1/137.036(1). Julian Schwinger discovered this expression in 1948; the involved calculations that led Schwinger to this and similar results in quantum field theory earned him the 1965 Nobel Prize in Physics. The result is also inscribed on the memorial marker near his grave in Mount Auburn Cemetery. The strand model proposes an intuitive explanation for this result.

Generally speaking, the factor $g/2$ describes the ratio between the 'mechanical' or 'geometric' rotation frequency – the rotation of the particle *mass* that leads to spin – and 'magnetic' rotation frequency – the rotation of the particle *charge* that leads to the magnetic moment. More precisely, the definition of the g-factor of a particle with charge e and mass m is

$$\frac{g}{2} = \frac{\mu/e}{S/m} . \tag{155}$$

Here, μ is the magnetic moment and S is the intrinsic angular momentum, or spin.

The *mechanical* or *geometric* rotation frequency is related to the ratio of the intrinsic angular momentum L and the mass m. Using the definitions from classical physics, we have $\boldsymbol{S}/m = \boldsymbol{r} \times \boldsymbol{v}$. The *magnetic* rotation frequency is related to the ratio of the magnetic moment $\boldsymbol{\mu}$ and the electric charge e. Classically, this ratio is $\boldsymbol{\mu}/e = \boldsymbol{r} \times \boldsymbol{v}$. Therefore, in classical physics – and also in the first order of the Pauli–Dirac description of the electron – the two rotation frequencies coincide, and the factor $g/2$ is thus equal to 1. However, as mentioned, both experiment and QED show a slight deviation of $g/2$ from unity, called the *anomalous* magnetic moment.

In the strand model, the geometric or mechanical rotation of a charged elementary particle is due to the rotation of the tangle core as a rigid whole, whereas the magnetic Page 227 rotation also includes phase changes due to the *deformations of the tangle core*. In par-

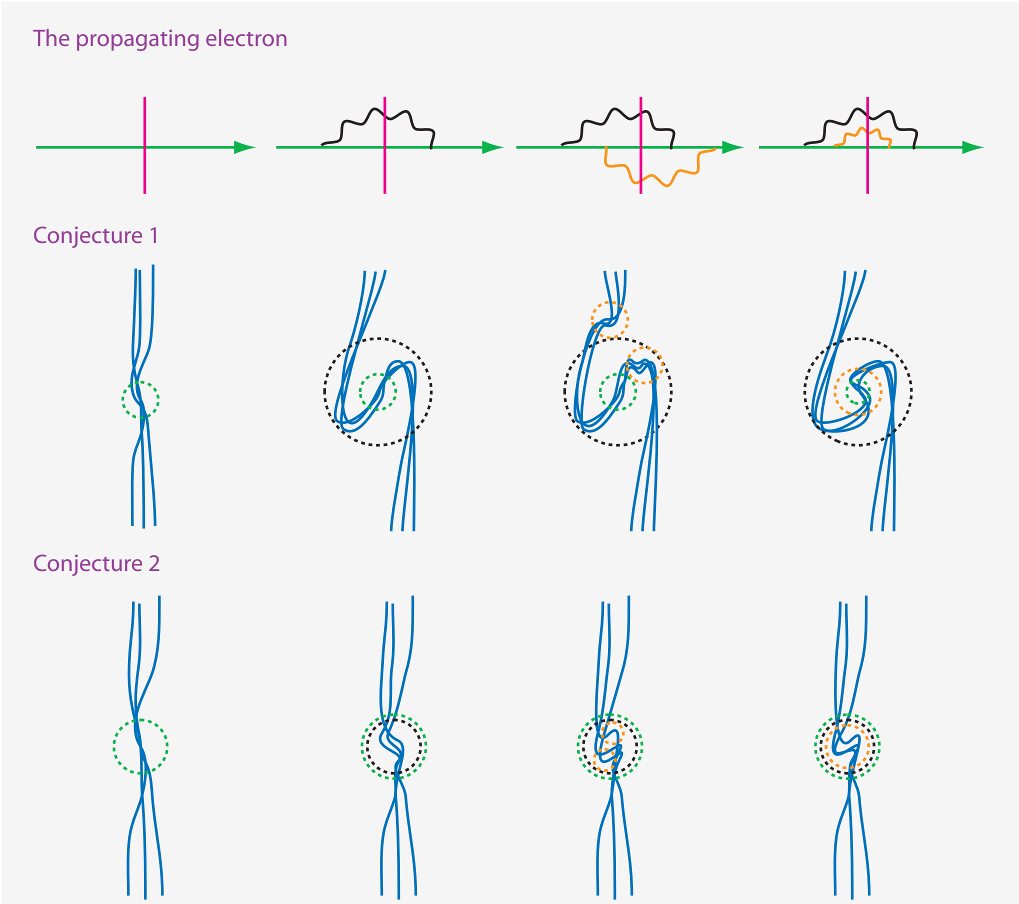

FIGURE 60 Two conjectured correspondences between the Feynman diagrams of quantum electrodynamics and the strand model for a propagating free electron. The lower strand model configurations are shown for a single instant – marked in magenta – of the electron propagator drawn above them. (For simplicity, the external field is not drawn.) In the first conjecture, the loops of the belt trick are conjectured to correspond to the virtual photons in the propagator and to be responsible for the anomalous magnetic moment. In the second conjecture, the deformations of the core correspond to the virtual photons.

ticular, the magnetic rotation of a charged elementary particle includes phase changes due to emission and reabsorption of virtual photons, i.e., of twisted loops.

In nature, the probability of the emission and reabsorption of a photon is determined by the fine structure constant α. The emission and reabsorption process leads to an additional angle that makes the 'magnetic' rotation angle differ from the 'mechanical' rotation angle. Since the fine structure constant describes the rotation of the phase due to virtual photon exchange, the emission and reabsorption of a virtual photon leads to an angle difference, and this angle difference is given by the fine structure constant itself. The ratio between the purely mechanical or geometric and the full magnetic rotation frequency is therefore not one, but increased by the ratio between the additional angle α and 2π. This is Schwinger's formula.

In short, the strand model reproduces Schwinger's celebrated formula for the anomalous magnetic moment almost from thin air. The strand model also implies that Schwinger's formula is valid for *all* charged elementary particles, independently of their mass; this is indeed observed. Higher order corrections also appear naturally in the strand model. Finally, the strand model implies that the complete expression, with all orders included, *converges*, because the full result is due to the shape and dynamics of the tangle core. The discussions about the existence of the perturbation limit in QED are thus laid to rest.

If we look into the details, it might be that the belt trick itself is at the origin of the anomalous magnetic moment. A conjecture for this connection is proposed and illustrated in Figure 60: if the two loops formed by the belt trick are seen as virtual photons, the factor $2\alpha/4\pi$ arises naturally. So do the higher-order terms. This explanation would relate the belt trick directly to the additional magnetic rotation angle. However, it might also be that this correspondence of the strand images in the figure to the upper diagrams is not fully correct. The topic is subject of research.

Challenge 156 e

A second conjecture is also given in Figure 60. The virtual photons could correspond to deformations of the tangle core. This conjecture is more in line with the distinction between gravity and gauge interactions given above, where it was stated that gravity is due to tail deformations and gauge interactions are due to core deformations. This conjecture is more in line with the distinction between a geometric and a magnetic rotation: the geometric rotation would be due to the rigid rotation of the tangle core, and the magnetic rotation would be due to an additional effect due to core deformation.

Page 227

Both conjectures on the origin of the g-factor imply that $1 < g/2 < 2$; in fact, we can even argue, using $\alpha < 1$, that the strand model implies

$$1 < g/2 < 1 + \frac{1}{2\pi} . \qquad (156)$$

This is not a new result; it is already implied by ordinary quantum field theory. However, the strand description of particle rotation suggests a way to calculate the g-factor and the fine structure constant. We will explore this below.

Page 393

Maxwell's equations

Ref. 183

The strand model of charge and photons reproduces Maxwell's equations . But strands also allow us to visualize and check Maxwell's field equations of classical electrodynamics directly. The equations are:

$$\begin{aligned}
\nabla\, \mathbf{E} &= \frac{\rho}{\varepsilon_0} , \\
\nabla\, \mathbf{B} &= 0 , \\
\nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} , \\
\nabla \times \mathbf{B} &= \frac{1}{c^2}\frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J} .
\end{aligned} \qquad (157)$$

▪ The first of these equations is satisfied whatever the precise mechanism at the basis of twisted loop emission by electric charges may be. Indeed, any mechanism in which a charge randomly sends out or swallows a twisted handle yields a $1/r^2$ dependence for the electrostatic field and the required connection between charge and the divergence of the electric field. This is not a deep result: any spherically-symmetric system that randomly emits or swallows some entity produces the equation, including the underlying inverse-square dependence. The result can also be confirmed in another, well-known way. In any exchange interaction between two charges, the exchange time is proportional to their distance apart r; in addition, quantum theory states that the exchanged momentum is inversely proportional to the distance r. Therefore, the force, or momentum per unit time, varies as $1/r^2$. This relation is valid independently of the underlying motion of the twisted loops, because space has three dimensions: all localized sources automatically fulfil the inverse square dependence.

The constant on the right-hand side of the first equation results from the definition of the units; in the language of the strand model, the constant fixes the twisted loop emission rate for an elementary charge.

▪ The second of the field equations (157) expresses the lack of magnetic charges. This equation is automatically fulfilled by the strand model, as the definition of the magnetic field with strands does not admit any magnetic sources. In fact, strands suggest that no localized entity can have a magnetic charge. Also this equation is valid independently of the details of the motion of the strands. Again, this is a topological effect.

▪ The third field equation relates the temporal change of the magnetic field to the curl of the electric field. In the strand model, this is satisfied naturally, because a curl in the electric field implies, by construction, a change of the magnetic field, as shown by Page 234 Figure 53. Again, this relation is valid independently of the details of the motion of the strands, as long as the averaging scale is taken to be large enough to allow the definition of electric and the magnetic fields.

▪ The most interesting equation is the last of the four Maxwell equations (157): in particular, the second term on the right-hand side, the dependence on the charge current. In the description of electrodynamics, the charge current **J** appears with a positive sign and with no numerical factor. (This is in contrast to linearized gravity, where the current has a numerical factor and a negative sign.) The positive sign means that a larger current produces a larger magnetic field. The strand model reproduces this factor: strands lead to an effect that is proportional both to charge (because more elementary charges produce more crossing flows) and to speed of movement of charge (large charge speed lead to larger flows). Because of this result, the classical photon spin, which is defined as L/ω, and which determines the numerical factor, namely 1, that appears before the charge current **J**, is recovered. Also this connection is obviously independent of the precise motion of the underlying strands.

The first term on the right-hand side of the fourth equation, representing the connection between a changing electric field and the curl of the magnetic field, is automatically Page 234 in agreement with the model. This can again be checked from Figure 53 – and again, this is a topological effect, valid for any underlying strand fluctuation. As an example, when a capacitor is charged, a compass needle between the plates is deflected. In the strand model, the accumulating charges on the plates lead to a magnetic field. The last of Maxwell's equations is thus also confirmed by the strand model.

In summary, the strand model reproduces Maxwell's equations. However, this is not a great feat. Maxwell-like equations appear in many places in field theory, for example in solid-state physics and hydrodynamics. Mathematical physicists are so used to the appearance of Maxwell-like equations in other domains that they seldom pay it much attention. The real tests for any model of electrodynamics, quantum or classical, are the deviations that the model predicts from electrodynamics, especially at high energies.

Curiosities and fun challenges about QED

Can you show that the calculation of the vacuum energy density of an infinite flat vacuum, when using strands, yields exactly zero, as expected? Challenge 157 e

* *

Can you confirm that the strand model of quantum electrodynamics does not violate charge conjugation C nor parity P at any energy? Challenge 158 e

* *

Can you confirm that the strand model of quantum electrodynamics conserves colour and weak charge at all energies, using the results of the next sections? Challenge 159 e

* *

Can you determine whether the U(1) gauge group deduced here is that of electrodynamics or that of weak hypercharge? Challenge 160 e

* *

Challenge 161 d Can you find a measurable deviation of the strand model from QED?

Summary on QED and experimental predictions

In the strand model, photons are single, helically twisted strands, randomly exchanged between charges; charges are chiral tangles, and therefore they effectively emit and absorb real and virtual photons. This is the complete description of QED using strands.

In particular, we have shown that Reidemeister I moves – or twists – of tangle cores lead to U(1) gauge invariance, Coulomb's inverse square relation, Maxwell's equations of electrodynamics and to Feynman diagrams. In short, we have deduced all experimental properties of quantum electrodynamics, except one: the strength of the coupling. Despite this open point, we have settled one line of the millennium list of open issues: we know the origin of the electromagnetic interaction and of its properties. Page 164

Is there a difference between the strand model and quantum electrodynamics? The precise answer is: there are *no measurable* differences between the strand model and QED. For example, the g-factor of the electron or the muon predicted by QED is not changed by the strand model. The U(1) gauge symmetry and the whole of QED remain valid at all energies. There are no magnetic charges. There are no other gauge groups. QED remains exact in all cases – as long as gravity plays no role.

The strand model prediction of a lack of larger gauge symmetries is disconcerting. There is thus *no* grand unification in nature; there is no general gauge group in nature, be it SU(5), SO(10), E6, E7, E8, SO(32) or any other. This result indirectly also rules out

Ref. 184 supersymmetry and supergravity. This unpopular result contrasts with many cherished habits of thought.

In the strand model, the equivalence of Feynman diagrams and strand diagrams implies that deviations of the strand model from QED are expected *only* when gravity starts to play a role. The strand model predicts that this will only happen just near the Planck energy $\sqrt{\hbar c^5/4G}$. At lower energies, QED is predicted to remain valid.

The strand model also confirms that the combination of gravity and quantum theory turns all Planck units into *limit* values, because there is a maximum density of strand crossings in nature, due to the fundamental principle. In particular, the strand model confirms the maximum electric field value $E_{\rm max} = c^4/4Ge \approx 1.9 \cdot 10^{62}\,\mathrm{V/m}$ and a maximum magnetic field value $B_{\rm max} = c^3/4Ge \approx 6.3 \cdot 10^{53}\,\mathrm{T}$. So far, these predictions are not in contrast with observations.

Thus the strand model predicts that approaching the electric or magnetic field limit values – given by quantum gravity – is the only option to observe deviations from QED. But measurements are not possible in those domains. Therefore we can state that there are no measurable differences between the strand model and QED.

Our exploration of QED has left open only two points: the calculation of the electromagnetic coupling constant and the determination of the spectrum of possible tangles for the elementary particles. Before we clarify these points, we look at the next Reidemeister move.

The weak nuclear interaction and the second Reidemeister move

In nature, the weak interaction is the result of the absorption and the emission of massive spin-1 bosons that form a broken weak triplet. The W and the Z bosons are emitted or absorbed by particles with weak charge; these are the left-handed fermions and right-handed antifermions. In other words, the weak interaction breaks parity P maximally. The W boson has unit electric charge, the Z boson has vanishing electric charge. The emission or absorption of W bosons changes the particle type of the involved fermion. The weak bosons also interact among themselves. All weakly charged particles are massive and move slower than light. The Lagrangian of matter coupled to the weak field has a broken SU(2) gauge symmetry. There are fundamental Feynman diagrams with triple and with quartic vertices. The weak coupling constant is determined by the electromagnetic coupling constant and the weak boson masses; its energy dependence is fixed by renormalization. The Higgs boson ensures full consistency of the quantum field theory of the weak interaction. Page 255

The previous paragraph summarizes the main observations about the weak interaction. More precisely, all observations related to the weak interaction are described by its Lagrangian. Therefore, we need to check whether the weak interaction Lagrangian follows from the strand model.

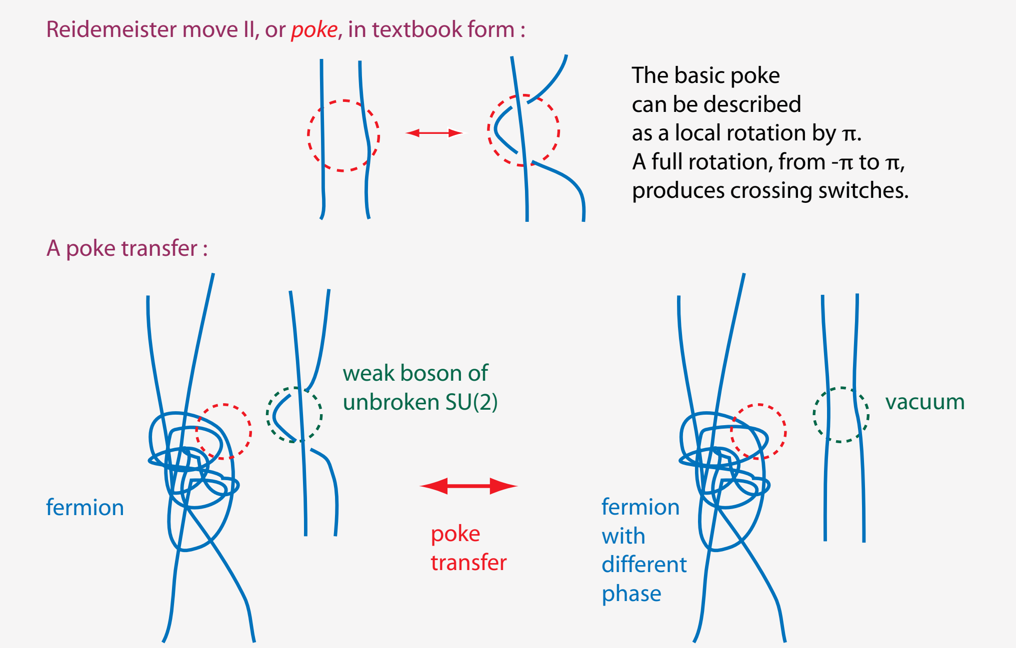

FIGURE 61 *Poke transfer* is the basis of the weak interaction in the strand model. No strand is cut or reglued; the transfer occurs only through the excluded volume due to the impenetrability of strands.

Strands, pokes and SU(2)

Page 225 As explained above, any gauge interaction involving a fermion is a deformation of the tangle core that changes the phase and rotation of the fermion tangle. We start directly with the main definition.

▷ The ***weak interaction*** is the transfer of a poke, i.e., the transfer of a Reidemeister II move, between two particles. An illustration is given in Figure 61. Strands are not cut in this process; they simply transfer the deformation as a result of their impenetrability.

Strands describe the weak interaction as exchange of pokes. In tangle cores, the *basic* pokes induce local rotations by an angle π, as shown in Figure 62: each basic poke rotates the region enclosed by the dotted circle. A full poke produces two crossings. There are *three*, linearly independent, basic pokes, in three mutually orthogonal directions. The three basic pokes τ_x, τ_y and τ_z act on the local region in the same way as the three possible mutually orthogonal rotations act on a belt buckle. For completeness, we note that the following arguments do not depend on whether the two strands involved in a poke are parallel, orthogonal, or at a general angle. The following arguments also do not depend Challenge 162 e on whether the pokes are represented by deforming *two* strands or only *one* strand. Both cases lead to crossing switches, for each possible poke type.

Figure 62 illustrates that the product of two different basic pokes gives the third basic Challenge 163 e poke, together with a sign – which depends on whether the sequence is cyclic or not –

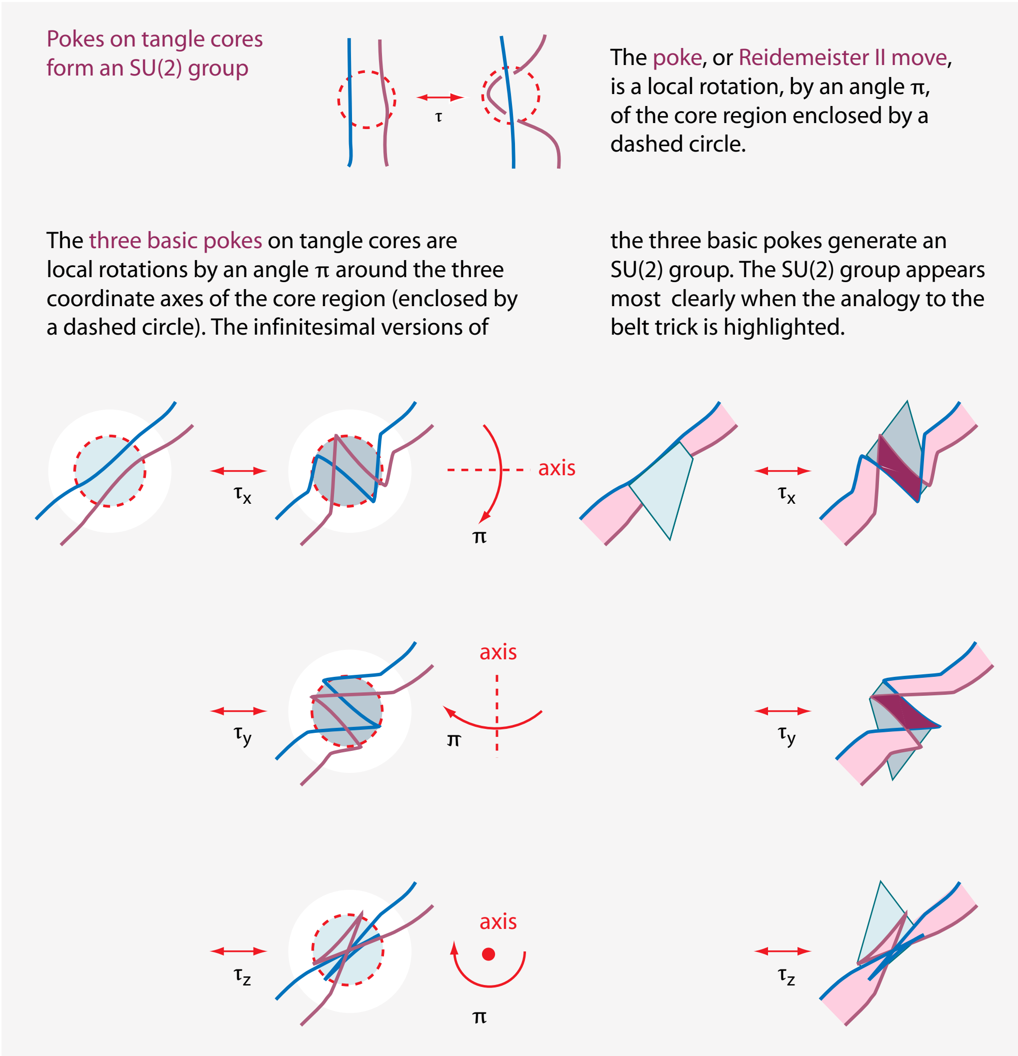

FIGURE 62 How the set of all pokes – the set of all deformations induced on tangle cores by the weak interaction – forms an SU(2) gauge group: the three pokes lead to the belt trick, illustrated here with a pointed buckle and two belts. The illustrated deformations of two strands represent the three unbroken weak vector bosons.

and a factor of i. Using the definition of -1 as a local rotation of the buckle region by Page 191 2π, we also find that the square of each basic poke is -1. In detail, we can read off the following multiplication table for the three basic pokes:

$$
\begin{array}{c|ccc}
\cdot & \tau_x & \tau_y & \tau_z \\
\hline
\tau_x & -1 & i\tau_z & -i\tau_y \\
\tau_y & -i\tau_z & -1 & i\tau_x \\
\tau_z & i\tau_y & -i\tau_x & -1
\end{array}
\qquad (158)
$$

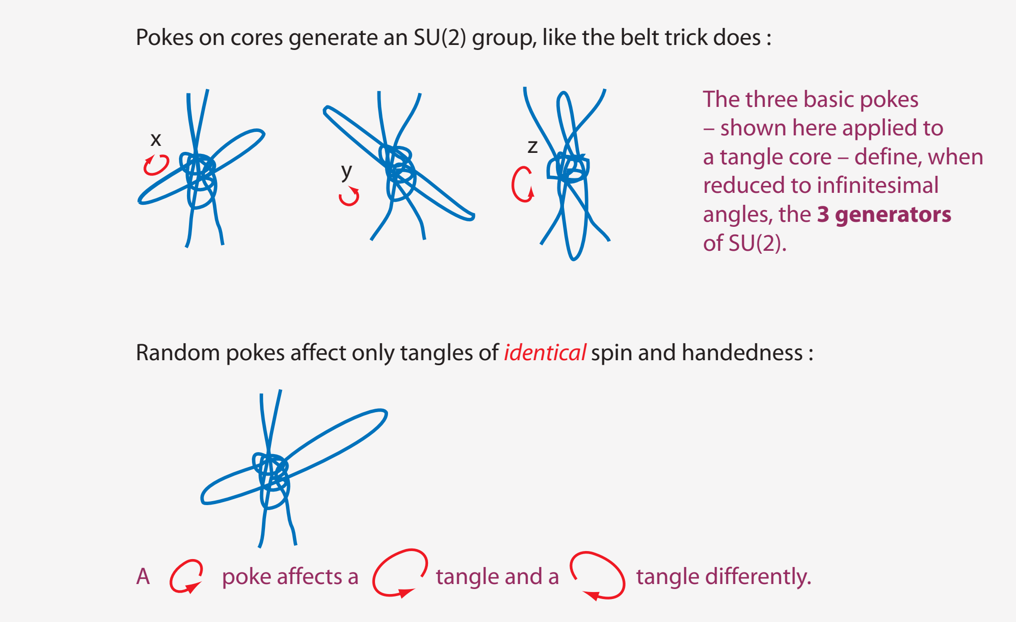

FIGURE 63 The three basic pokes and weak charge in the strand model.

In other terms, the three basic pokes – and in particular also their infinitesimal versions – behave like the generators of an SU(2) group. Because pokes can be seen as local rotations of a buckle region, they can be generalized to arbitrary angles. Such arbitrary pokes can be concatenated. We thus find that arbitrary pokes form a full SU(2) group. This is the reason for their equivalence with the belt trick.

The different gauge choices for a particle are not illustrated in Figure 62. The gauge choices arise from the different ways in which the basic pokes τ_x, τ_y and τ_z can be assigned to the set of deformations that describe the belt trick.

In summary, we can state that in any definition of the phase of a tangled fermion core, there is an SU(2) gauge freedom; in addition, there exists an interaction with SU(2) gauge symmetry. In other words, the strand model implies, through the second Reidemeister move, *the existence of the unbroken weak interaction with a gauge group SU(2).*

Weak charge and parity violation

A particle has weak charge if, when subject to many random pokes, a non-zero average phase change occurs. Surrounded by a bath of strands that continuously induce random pokes, not all tangles will change their phase on a long-time average: only tangles that lack symmetry will. One symmetry that must be lacking is spherical symmetry. Therefore, only tangles whose cores lack *spherical symmetry* have the chance to be influenced by random pokes. Since all tangles, independently of their core details, lack spherical symmetry, all such tangles, i.e., all massive particles, are candidates to be influenced, and thus are candidates for weakly charged particles. We therefore explore them in detail now.

If a tangle is made of *two or more* linked strands, it represents a massive spin-1/2

Page 298 particle (except for a simple twist, which represents the graviton). All such fermion cores lack spherical and cylindrical symmetry. When a fermion spins, two things happen: the core rotates and the belt trick occurs, which untangles the tails. Compared to the Page 220 direction of motion, the rotation and the untangling can be either left-handed or right-handed.

Every poke is a shape transformation of the core with a preferred handedness. The chirality is of importance in the following.

A particle has weak charge if random pokes lead to a long-time phase change. In order to feel any average effect when large numbers of random pokes are applied, a core must undergo different effects for a poke and its reverse. As already mentioned, this requires a lack of core symmetry. Whenever the core has no symmetry, non-compensating phase effects will occur: if the core rotation with its tail untangling and the poke are of the same handedness, the phase will increase, whereas for opposite handedness, the phase will decrease a bit less.

▷ Non-vanishing ***weak charge*** for fermions appears only for tangle cores whose handedness leads to average poke effects.

In other words, the strand model predicts that random pokes will only affect a core if the core handedness and the randomly applied belt trick are of the *same* handedness. In physical terms, random pokes will only affect left-handed particles or right-handed antiparticles. Thus, the strand model predicts that *the weak interaction violates parity maximally*, This is exactly as observed. In other terms, weak charge and the parity violation of the weak interaction are consequences of the belt trick. This relation is summarized in Figure 63.

If an elementary particle is described by a *two tangled* strands, we expect it to be influenced by average pokes. Such tangle cores are spin-1 bosons; their cores lack spherical and cylindrical symmetry. The core rotation will induce a left-right asymmetry that will lead to a higher effect of a poke than of its reverse. Two-stranded particles are thus predicted to carry weak charge. We therefore expect that quarks – to be explored below – and the weak bosons themselves interact weakly.

Because the weak bosons interact weakly, the strand model implies that the weak interaction is a *non-Abelian* gauge theory, as is observed.*

If a tangle is made of a *single unknotted* strand, it is not affected by random pokes. The strand model thus predicts that the photon has no weak charge, as is observed. The same also holds for gluons.

The strand definition of weak charge leads to two conclusions that can be checked by experiment. First, all electrically charged particles – having cores that are chiral and thus lack cylindrical symmetry – are predicted to be weakly charged. Secondly, in the strand model, only massive particles interact weakly; in fact, *all* massive particles interact weakly, because their cores lack cylindrical symmetry. In other words, all weakly

* Non-Abelian gauge theory was introduced by Wolfgang Pauli. In the 1950s, he explained the theory in series of talks. Two physicists, Yang Chen Ning and Robert Mills, then wrote down his ideas. Yang later received the Nobel Prize in Physics with Lee Tsung Dao for a different topic, namely for the violation of parity of the weak interaction.

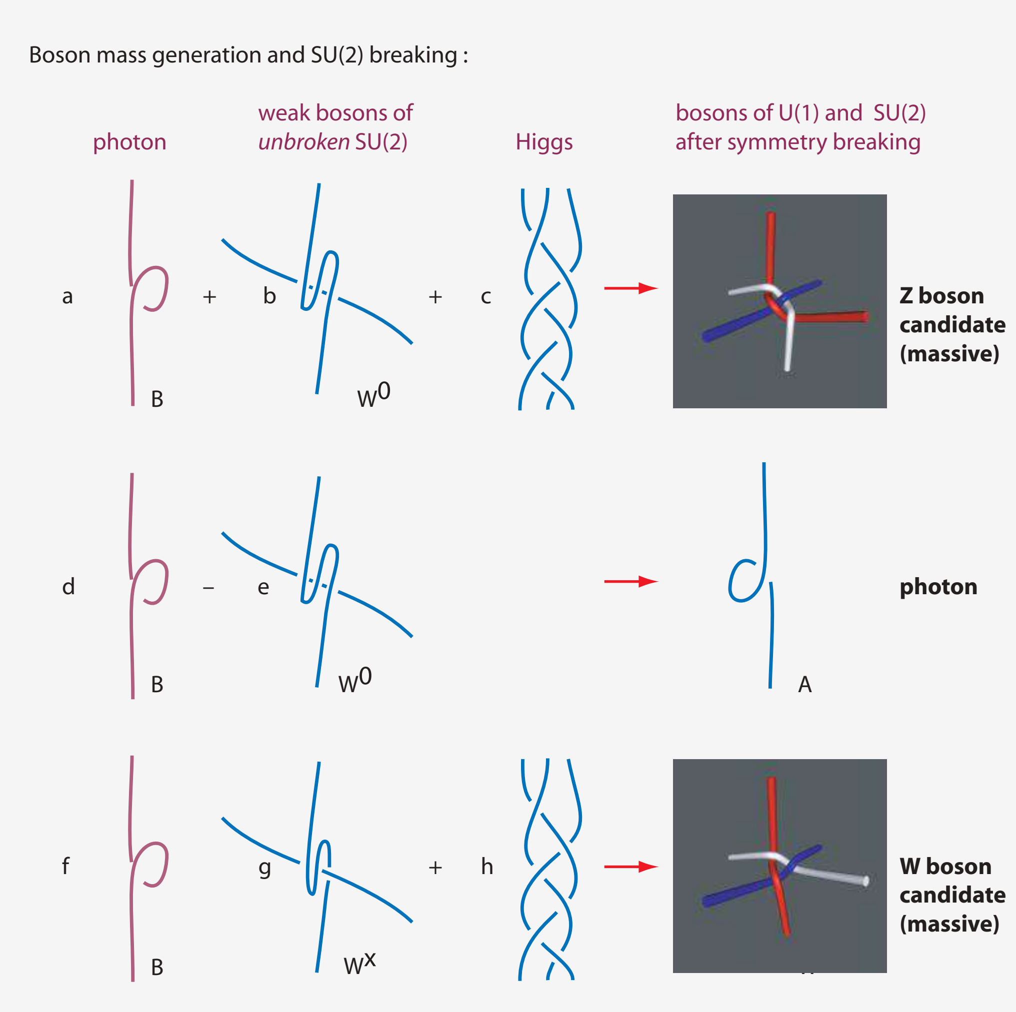

FIGURE 64 Poke-inducing strand motions (left) become massive weak vector tangles (right) through symmetry breaking and tail braiding. Tail braiding is related to the Higgs boson, whose tangle model will be clarified later on.

charged particles move more slowly than light and vice versa. Both conclusions agree with observation.

In summary, all properties of weak charge found in nature are reproduced by the tangle model.

Weak bosons

Gauge bosons are those particles that are exchanged between interacting fermions: gauge bosons induce phase changes of fermions. This implies that the (unbroken) weak bosons are the particles* that induce the three poke moves:

* This reworked strand model of the W and Z bosons arose in 2015.

▷ ***Weak, unbroken, intermediate bosons*** are described by double strands. Illustrations are given in and Figure 62 and Figure 64.

Single strands that induce phase changes in fermions interacting weakly are shown on the left side of Figure 64. They correspond to the three basic pokes τ_x, τ_y and τ_z.

We note two additional points. First of all, the (unbroken) bosons could also be described by the motion of a single strand in a strand group. This makes them *spin 1* particles.

Furthermore, unknotted tangles are *massless*. In the strand model, tangles that induce pokes *differ* from the massive weak intermediate bosons, shown on the right of Figure 64. This difference is due to the *breaking* of the SU(2) gauge symmetry, as we will find out soon.

The Lagrangian of the unbroken SU(2) gauge interaction

The energy of the weak field is given by the density of weak gauge boson strands. As long as the SU(2) symmetry is not broken, the energy of the weak field and the energy of fermions are both SU(2) invariant. As a consequence, we are now able to deduce a large part of the Lagrangian of the weak interaction, namely the Lagrangian for the case that the SU(2) symmetry is unbroken.

As long as SU(2) is unbroken, the vector bosons are described as unknotted tangles that induce pokes, as shown on the left of Figure 64. There are three such bosons. Since they can be described by a single strand that moves, they have spin 1; since they are unknotted, they have zero mass and electric charge.

Energy is the number of crossing switches per time. As long as SU(2) is unbroken and the weak bosons are massless, the energy of the weak boson field and thus their Page 234 Lagrangian density is given by the same expression as the energy of the photon field. In particular, the strand model implies that energy density is quadratic in the field intensities. We only have to add the energies of all three bosons together to get:

$$\mathcal{L} = -\frac{1}{4}\sum_{a=1}^{3} W^a{}_{\mu\nu} W_a{}^{\mu\nu} , \qquad (159)$$

This expression is SU(2) gauge invariant. Indeed, SU(2) gauge transformations have no effect on the number of crossing switches due to weak bosons or to the motion of pokes. Thus, gauge transformations leave weak field intensities and thus also the energy of the weak fields invariant, as observed.

We can now write down the Lagrangian for weakly charged fermions interacting with the weak vector bosons. Starting from the idea that tangle core deformations lead to phase redefinitions, we have found that pokes imply that the *unbroken* weak Lagrangian density for matter and radiation fields is SU(2) gauge invariant. In parallel to electrodynamics we thus get the Lagrangian

$$\mathcal{L}_{\text{unbroken weak}} = \sum_f \overline{\Psi}_f (i\hbar c \not{D} - m_f c^2)\Psi_f - \frac{1}{4}\sum_{a=1}^{3} W^a{}_{\mu\nu} W_a{}^{\mu\nu} , \qquad (160)$$

where $\not{D}$ is now the SU(2) gauge covariant derivative and the first sum is taken over all fermions. In this Lagrangian, only the left-handed fermions and the right-handed antifermions carry weak charge. This Lagrangian, however, does *not* describe nature: the observed SU(2) breaking is missing.

SU(2) BREAKING

In nature, the weak interaction does *not* have an SU(2) gauge symmetry. The symmetry is only approximate; is said to be *broken*. The main effect of SU(2) symmetry breaking are the non-vanishing – and different – masses for the W and Z bosons, and thus the weakness and the short range of the weak interaction. In addition, the symmetry breaking implies a *mixing* of the weak and the electromagnetic interaction: it yields the so-called Page 257 *electroweak* interaction. This mixing is often called electroweak 'unification'.

The strand model suggests the following description:

▷ ***Mass generation*** for bosons and the related SU(2) ***symmetry breaking*** are due to *tail braiding* at the border of space. Figure 64 illustrates the idea.

In this description, tail braiding* is assumed to occur at a distance outside the domain of observation; in that region – which can be also the border of physical space – tail braiding is *not* forbidden and *can* occur. The probability of tail braiding is low, because the crossings have first to fluctuate to that distance and then fluctuate back. Nevertheless, the process of tail braiding can take place.

Tail braiding appears *only* in the weak interaction. It does not appear in the other two gauge interactions, as the other Reidemeister moves are not affected by processes at the border of space. In the strand model, this is the reason that only SU(2) is broken in nature. In short, SU(2) breaking is a natural consequence of the second Reidemeister move.

Tail braiding transforms the unbraided, and thus massless, poke strands into the braided, and thus massive W and Z strands. Tail braiding leads to particle cores: therefore is a mass-generating process. The precise mass values that it generates will be de- Page 358 termined below. The strand model thus confirms that mass generation is related to the breaking of the weak interaction.

Tail braiding mixes the W^0 with the 'original' photon. This is shown in Figure 64. The mixing is due to the topological similarities of the strand models of the two particles. The resulting Z boson is achiral, and thus electrically neutral, as observed. We note that the existence of a neutral, massive Z boson implies that elastic neutrino scattering in matter occurs in nature, as was observed for the first time in 1974. Since any electrically charged Page 250 particle also has weak charge, the existence of a Z boson implies that any two electrically charged particles can interact both by exchange of photons and by exchange of Z bosons. In other words, SU(2) breaking implies electroweak mixing, or, as is it usually called, Page 257 electroweak 'unification'.

Tail braiding takes place in several weak interaction processes, as shown in Figure 67. Page 256 Tail braiding thus can change particle topology, and thus particle type. The strand model

* In the original strand model of the weak bosons, from the year 2008, the role of tail braiding was taken by strand overcrossing.

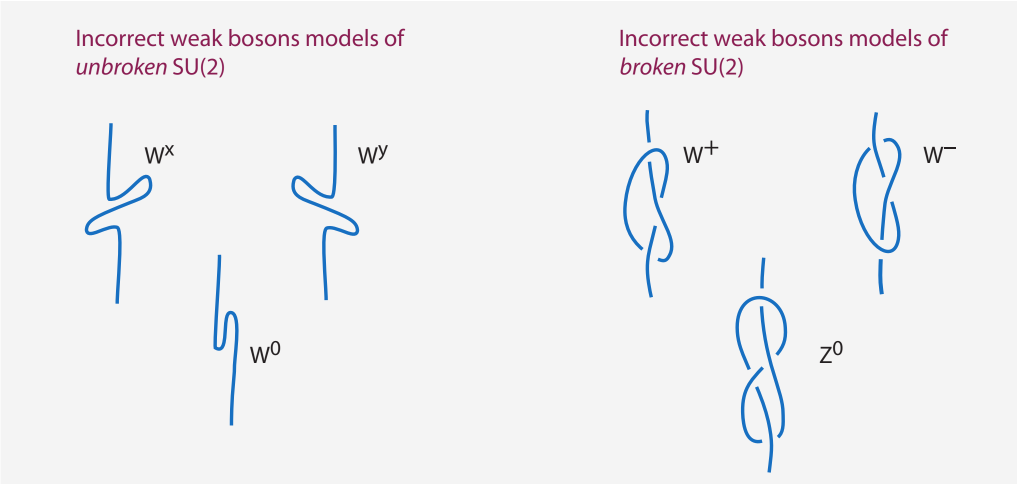

FIGURE 65 The supposed models for the massive weak gauge bosons after symmetry breaking, from 2008 (on the right side), now seen to be incorrect.

thus predicts that the weak interaction *changes* particle flavours (types), as is observed. In fact, the strand model also predicts that *only* the weak interaction has this property. This is also observed.

On the other hand, strands are never cut or glued back together in the strand model, not even in the weak interaction. As a result, the strand model predicts that the weak interaction conserves electric charge, spin and, as we will see below, colour charge, baryon number and lepton number. All this is observed. Page 318

Tail braiding also implies that the tangles for the Z boson and for the W boson shown above are only the *simplest* tangles associated with each boson; more complicated tangles are higher order propagating states of the same basic open knots. This will be of great importance later on, for the proof that all gauge bosons of nature are already known today. Page 358

In summary, the second Reidemeister move leads to *tail braiding*; tail braiding leads to the observed properties of SU(2) symmetry breaking. (Equivalently, the strand model implies that the simplest tangles of the weak interaction bosons show SU(2) symmetry, whereas the more complicated, massive tangles break this symmetry.) The value of the mixing angle and the particle masses have still to be determined. This will be done below. Page 360

Open issue: are the W and Z tangles correct?

In 2014, Sergei Fadeev raised an issue: A *tangle* version of the W and Z that does *not* contain any knot and does not require an actual strand overcrossing process at spatial infinity, the strand model would gain in simplicity and elegance. Thinking about the issue, it became clear that such a tangle could occur when vacuum strands were included, as shown above.

In contrast, in 2008, in the first version of the strand model, the W boson after symmetry breaking was thought to be an open overhand knot, and the Z boson an open figure-eight knot.

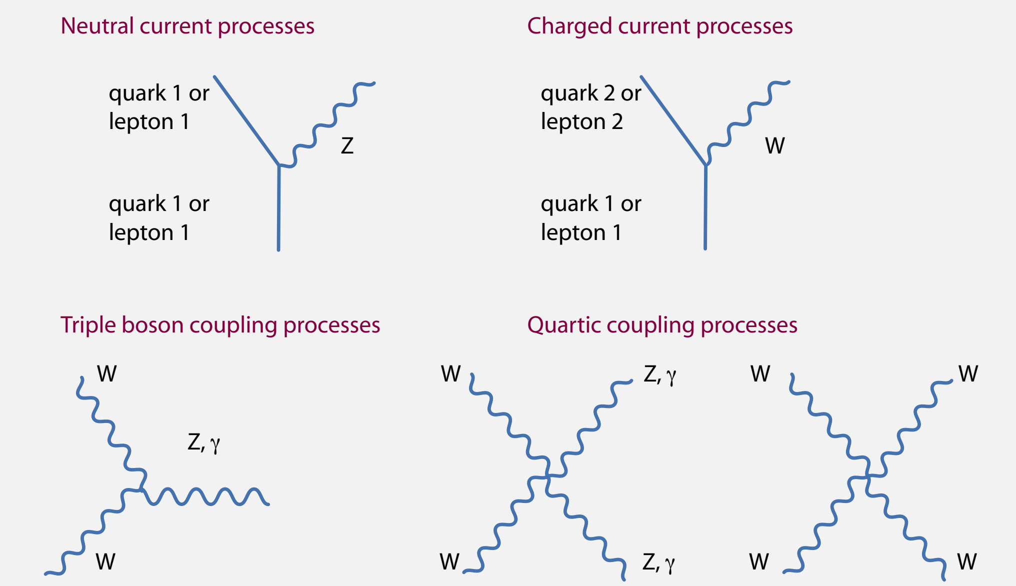

FIGURE 66 The fundamental Feynman diagrams of the weak interaction that do not involve the Higgs boson.

It might well be that the new, 2015/2016 strand models for the two intermediate vector Page 251 bosons, shown in Figure 64 are still not correct. The possibility remains intriguing and Challenge 164 ny a definitive issue still needs to be found.

The electroweak Lagrangian

We can now use the results on SU(2) symmetry breaking to deduce the *electroweak* Lagrangian density. We have seen that symmetry breaking leaves the photon massless but introduces masses to the weak vector bosons, as shown in Figure 64. The non-vanishing boson masses M_W and M_Z add kinetic terms for the corresponding fields in the Lagrangian.

Due to the symmetry breaking induced by tail braiding, the Z boson results from the mixing with the (unbroken) photon. The strand model predicts that the mixing can be described by an angle, the so-called weak mixing angle θ_w. In particular, the strand model implies that $\cos\theta_w = M_W/M_Z$.

As soon as symmetry breaking is described by a mixing angle due to tail braiding, we get the known electroweak Lagrangian, though at first without the terms due to the Page 328 Higgs boson. (We will come back to the Higgs boson later on.) We do not write down the Lagrangian of the weak interaction predicted by the strand model, but the terms are the same as those found in the standard model of elementary particles. There is one important difference: the Lagrangian so derived does not yet contain quark and lepton mixing. Indeed, experiments show that the weak fermion eigenstates are not the same as the strong or electromagnetic eigenstates: quarks mix, and so do neutrinos. The reason for this observation, and the effect that mixing has on the weak Lagrangian, will become Page 376 clear once we have determined the tangles for each fermion.

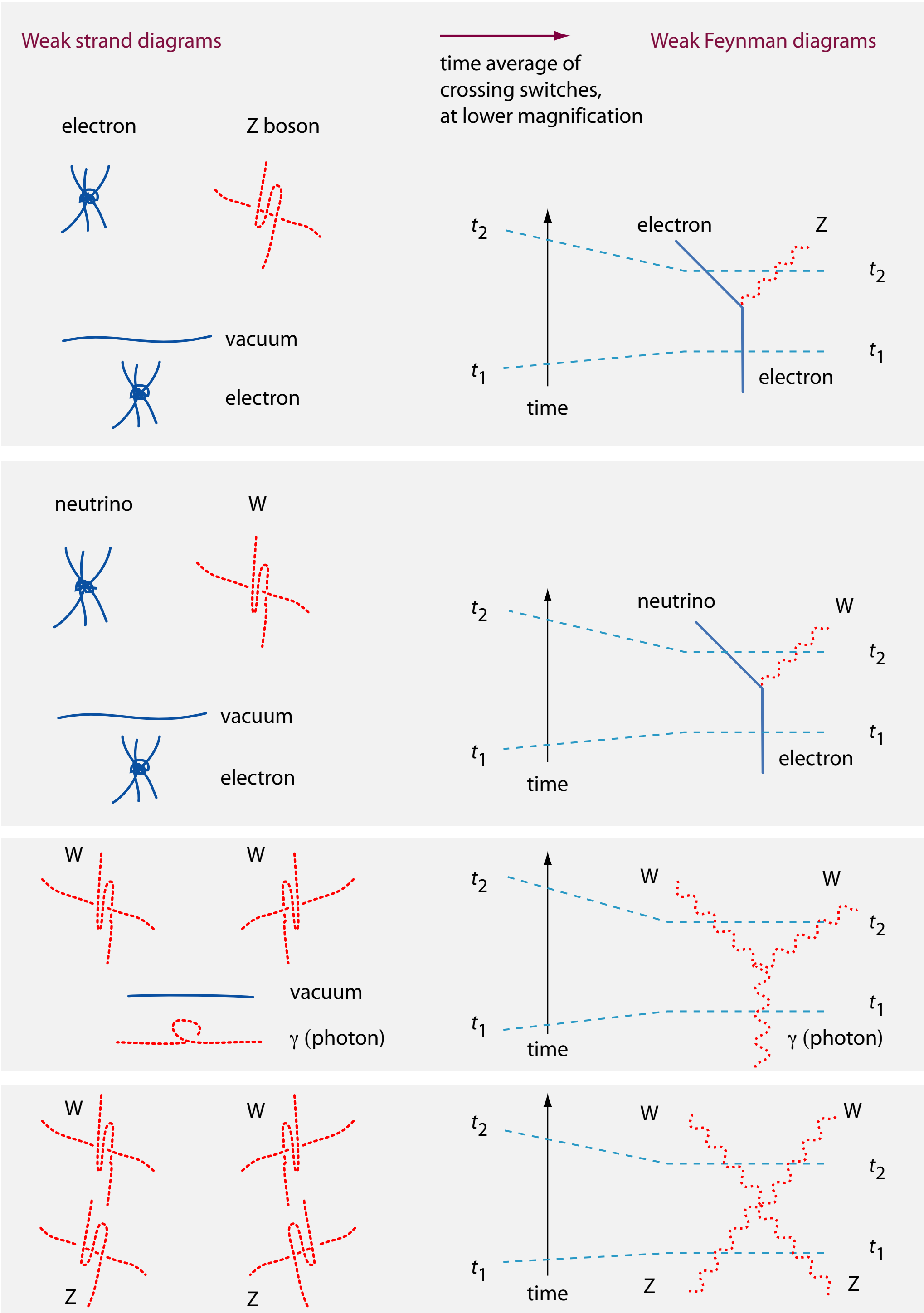

FIGURE 67 The strand model for the fundamental Feynman diagrams of the weak interaction. The tangles for the fermions are introduced later on.

In summary, the strand model implies the largest part of the Lagrangian of the weak interaction. The issue of the Higgs boson is still open, and the electroweak Lagrangian contains a number of constants that are not yet clarified. These unexplained constants are the number of the involved elementary particles, their masses, couplings, mixing angles and CP violation phases, as well as the value of the weak mixing angle.

The weak Feynman diagrams

In nature, the weak interaction is described by a small number of fundamental Feynman diagrams. Those not containing the Higgs boson are shown in Figure 66. These Feynman diagrams encode the corresponding Lagrangian of the weak interaction.

In the strand model, pokes lead naturally to strand versions of the fundamental Feynman diagrams. This happens as shown in Figure 67. We see again that the strand model reproduces the weak interaction: each Feynman diagram is due to a strand diagram for which only crossing switches are considered, and for which Planck size is approximated as zero size. In particular, the strand model does not allow any *other* fundamental diagrams for the weak interaction.

Challenge 165 e

The finite and small number of possible strand diagrams and thus of Feynman diagrams implies that the weak interaction is *renormalizable*. For example, the change or 'running' of the weak coupling with energy is reproduced by the strand model, because the running can be determined through the appropriate Feynman diagrams.

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Fun challenges and curiosities about the weak interaction

The W boson and its antiparticle are observed to annihilate through the electromagnetic interaction, yielding two or more photons. The tangle model of the weak bosons has a lot of advantages compared to the knot model: The annihilation is much easier to understand.

Challenge 166 e

* *

The strand model, like the standard model of particle physics, predicts that everything about the weak interaction is already known. Nevertheless, the most important weak process, the *decay of the neutron*, is being explored by many precision experiments. The strand model predicts that none of these experiments will yield any surprise.

* *

The strand model makes clear that the weak interaction and the electromagnetic interaction *mix*, but do not unify. There is only electroweak mixing, and *no* electroweak unification, despite claims to the contrary by the Nobel Prize committee and many other physicists. In fact, Sheldon Glashow, who received the Nobel Prize in Physics for this alleged 'unification', agrees with this assessment. So do Richard Feynman and, above all, Martin Veltman, who was also involved in the result; he even makes this very point in his Nobel Prize lecture. The incorrect habit to call electroweak mixing a 'unification' was one of the main reason for the failure of past unification attempts: it directed the attention of researchers in the wrong direction.

Ref. 185

In the strand model, the mixing of the electromagnetic and the weak interaction can be seen as a consequence of knot geometry: the poke generators of the weak interaction also contain twists, i.e., also contain generators of the electromagnetic interaction. In contrast, generators of other Reidemeister moves do not mix among them or with pokes; and indeed, no other type of interaction mixing is observed in nature.

Summary on the weak interaction and experimental predictions

We have deduced the main properties of the weak Lagrangian from the strand model. We have shown that Reidemeister II moves – or pokes – in tangle cores lead to a broken SU(2) gauge group and to massive weak bosons. We found that the deviation from tangle core sphericity plus chirality is weak charge, and that the weak interaction is non-Abelian. We have also shown that the weak interaction naturally breaks parity maximally and mixes with the electromagnetic interaction. In short, we have deduced the main experimental properties of the weak interaction.

Is there a difference between the strand model and the electroweak Lagrangian of the standard model of particle physics? Before we can fully answer the question on deviations between the strand model and the standard model, we must settle the issue of the Page 328 Higgs boson. This is done later on.

In any case, the strand model predicts that the broken SU(2) gauge symmetry remains valid at all energies. No other gauge groups appear in nature. The strand model thus predicts again that there is no grand unification, and thus no larger gauge group, be it SU(5), SO(10), E6, E7, E8, SO(32) or any other group. Also this result indirectly rules out supersymmetry and supergravity.

The strand model also predicts that the combination of gravity and quantum theory turns all Planck units into *limit* values, because there is a maximum density of strand crossings in nature, due to the fundamental principle. Therefore, the strand model predicts a *maximum weak field* value given by the Planck force divided by the smallest weak charge. All physical systems – including all astrophysical objects, such as neutron stars, quark starts, gamma-ray bursters or quasars – are predicted to conform to this limit. So far, no observed field value is near this limit, so that the prediction does not contradict observation.

So far, our exploration of the weak interaction has left us with a few open issues: we need to calculate the weak coupling constant and determine the tangle for each particle of the standard model, including the Higgs boson. But we also need to explain weak fermion mixing, CP violation and the masses of all particles. Despite these open points, we Page 164 have settled another line of the millennium list: we know the origin of the weak interaction and of its main properties. Before we clarify the open issues, we explore the third Reidemeister move.

The strong nuclear interaction and the third Reidemeister move

In nature, the strong interaction is the result of the absorption and the emission of massless, electrically uncharged, spin-1 gauge bosons that are called *gluons*. Gluons interact with quarks, the only fermions with *colour* charge. Fermions can have three different colour charges, antifermions three different anticolours. Gluons form an octet, are them-

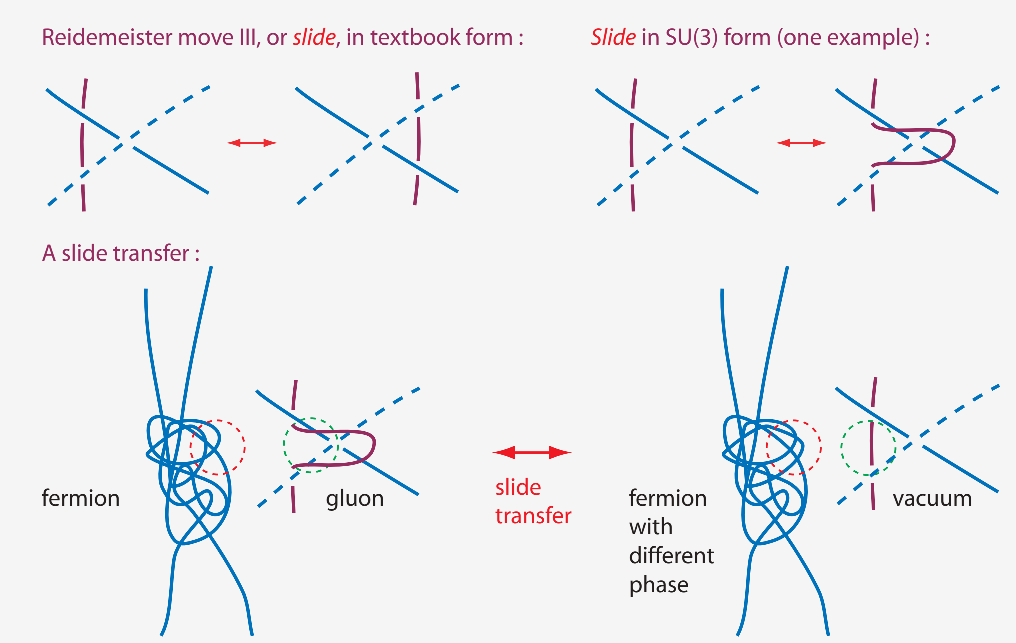

FIGURE 68 A gluon changes the phase of a tangle: *slide transfer* is the basis of the strong interaction in the strand model. During the interaction, no strand is cut or reglued; the transfer occurs purely through the excluded volume that results from the impenetrability of strands. Slide transfer generates a SU(3) gauge group, as explained in the text.

selves colour charged and therefore also interact among themselves. The Lagrangian of quarks coupled to the gluon field has an unbroken SU(3) gauge symmetry. There are three fundamental Feynman diagrams: one for quark-gluon interaction and two for gluon-gluon interactions: a triple and a quartic gluon vertex. The strong coupling constant is about 0.5 at low energy; its energy dependence is determined by renormalization. Its value decreases with increasing energy.

The previous paragraph summarizes the main observations about the strong interaction. All known observations related to the strong interaction, without any known exception, are contained in its Lagrangian. Therefore, we need to show that the strong interaction Lagrangian follows from the strand model.

Strands and the slide, the third Reidemeister move

Page 225

As explained above, interactions of fermions are deformations of the tangle core that change its phase. We start directly by presenting the strand model for the strong interaction.

▷ The ***strong interaction*** is the transfer of *slides*, i.e., the transfer of third Reidemeister moves, between a gluon and a particle. As shown in Figure 68, strands are not cut in this process; gluons simply transfer slide deformations to tangle cores as a result of their impenetrability.

Such a slide transfer will influence the phase of the affected particle tangle. Therefore, slide transfers are indeed a type of interaction.

An introduction to SU(3)

Before we show that slides are responsible for the strong nuclear interaction, we summarize the mathematical properties of the Lie group SU(3). This Lie group is the structure generated by the unitary 3×3 matrices with determinant $+1$. It is a *group*, because matrices can be properly multiplied, because the identity matrix is included, and inverse matrices exist. SU(3) is also a *manifold*; a quick check shows that it has eight dimensions. In short, SU(3) is a *Lie group*: its elements behave like points on a manifold that can be multiplied. The Lie bracket is the commutator. A general element E of SU(3) can be written as an exponential in the well-known way

Challenge 167 e

$$E = \mathrm{e}^{\sum_{n=1}^{8} \alpha_n i \lambda_n / 2} \tag{161}$$

where the eight real parameters α_n can be thought of as the eight coordinates of the group elements on the group manifold. Since SU(3) is compact and simple, these coordinates are best visualized as angles. Of course, i is the imaginary unit. The generators λ_n are complex, traceless and hermitian 3×3 matrices; they are used to define a basis for the group elements. The eight generators are *not* group elements themselves. They describe the structure of the group manifold near the identity matrix; for a Lie group, this local structure defines the full group manifold. Like for any basis, also set of eight generators λ_n is not unique. Of the many possible choices for the generators, the *Gell-Mann matrices* λ_1 to λ_8 are the most commonly used in physics.

The Gell-Mann matrices λ_n, the corresponding group elements D_n for general angles, and the group elements E_n for the finite angle π are given by:

$$\lambda_1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \; D_1(\alpha) = \mathrm{e}^{\alpha \mathrm{i} \lambda_1/2} = \begin{pmatrix} \cos\alpha/2 & i\sin\alpha/2 & 0 \\ i\sin\alpha/2 & \cos\alpha/2 & 0 \\ 0 & 0 & 1 \end{pmatrix},$$

$$E_1 = \mathrm{e}^{\pi \mathrm{i} \lambda_1/2} = \begin{pmatrix} 0 & i & 0 \\ i & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\lambda_2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \; D_2(\alpha) = \mathrm{e}^{\alpha \mathrm{i} \lambda_2/2} = \begin{pmatrix} \cos\alpha/2 & \sin\alpha/2 & 0 \\ -\sin\alpha/2 & \cos\alpha/2 & 0 \\ 0 & 0 & 1 \end{pmatrix},$$

$$E_2 = \mathrm{e}^{\pi \mathrm{i} \lambda_2/2} = \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\lambda_3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix},\ D_3(\alpha) = \mathrm{e}^{\alpha \mathrm{i} \lambda_3/2} = \begin{pmatrix} \cos\alpha/2 + i\sin\alpha/2 & 0 & 0 \\ 0 & \cos\alpha/2 - i\sin\alpha/2 & 0 \\ 0 & 0 & 1 \end{pmatrix},$$

$$E_3 = \mathrm{e}^{\pi \mathrm{i} \lambda_3/2} = \begin{pmatrix} i & 0 & 0 \\ 0 & -i & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\lambda_4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix},\ D_4(\alpha) = \mathrm{e}^{\alpha \mathrm{i} \lambda_4/2} = \begin{pmatrix} \cos\alpha/2 & 0 & i\sin\alpha/2 \\ 0 & 1 & 0 \\ i\sin\alpha/2 & 0 & \cos\alpha/2 \end{pmatrix},$$

$$E_4 = \mathrm{e}^{\pi \mathrm{i} \lambda_4/2} = \begin{pmatrix} 0 & 0 & i \\ 0 & 1 & 0 \\ i & 0 & 0 \end{pmatrix}$$

$$\lambda_5 = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix},\ D_5(\alpha) = \mathrm{e}^{\alpha \mathrm{i} \lambda_5/2} = \begin{pmatrix} \cos\alpha/2 & 0 & \sin\alpha/2 \\ 0 & 1 & 0 \\ -\sin\alpha/2 & 0 & \cos\alpha/2 \end{pmatrix},$$

$$E_5 = \mathrm{e}^{\pi \mathrm{i} \lambda_5/2} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ -1 & 0 & 0 \end{pmatrix}$$

$$\lambda_6 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix},\ D_6(\alpha) = \mathrm{e}^{\alpha \mathrm{i} \lambda_6/2} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\alpha/2 & i\sin\alpha/2 \\ 0 & i\sin\alpha/2 & \cos\alpha/2 \end{pmatrix},$$

$$E_6 = \mathrm{e}^{\pi \mathrm{i} \lambda_6/2} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & i \\ 0 & i & 0 \end{pmatrix}$$

$$\lambda_7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix},\ D_7(\alpha) = \mathrm{e}^{\alpha \mathrm{i} \lambda_7/2} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\alpha/2 & \sin\alpha/2 \\ 0 & -\sin\alpha/2 & \cos\alpha/2 \end{pmatrix},$$

$$E_7 = \mathrm{e}^{\pi \mathrm{i} \lambda_7/2} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix}$$

$$\lambda_8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix},$$

$$D_8(\alpha) = \mathrm{e}^{\sqrt{3}\alpha \mathrm{i} \lambda_8/2} = \begin{pmatrix} \cos\alpha/2 + i\sin\alpha/2 & 0 & 0 \\ 0 & \cos\alpha/2 + i\sin\alpha/2 & 0 \\ 0 & 0 & \cos\alpha - i\sin\alpha \end{pmatrix},$$

$$E_8 = D_8(\pi) = \begin{pmatrix} i & 0 & 0 \\ 0 & i & 0 \\ 0 & 0 & -1 \end{pmatrix}. \tag{162}$$

The eight Gell-Mann matrices λ_n are hermitean, traceless and trace-orthogonal. The corresponding group elements D_n and E_n can be thought as the unnormed and normed basis vectors of the group manifold. We note that the definition of E_8 differs from that of the other group elements E_n: it contains an extra factor $\sqrt{3}$. The fourfold concatenation of each matrix $i\lambda_n$ is the identity matrix – except for the case $i\lambda_8$. Instead, the generator λ_8 commutes with λ_1, λ_2 and λ_3 – though not with the other generators.

There is *no* ninth or tenth Gell-Mann matrix. Such a matrix would not be linearly independent from the first eight ones. Indeed, the two matrices deduced from λ_3 using symmetry considerations, namely

$$\lambda_9 = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \; D_9(\alpha) = \mathrm{e}^{\alpha \mathrm{i} \lambda_9/2} = \begin{pmatrix} \cos\alpha/2 - i\sin\alpha/2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \cos\alpha/2 + i\sin\alpha/2 \end{pmatrix},$$
$$E_9 = D_9(\pi) = \begin{pmatrix} -i & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & i \end{pmatrix}$$
$$\lambda_{10} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}, \; D_{10}(\alpha) = \mathrm{e}^{\alpha \mathrm{i} \lambda_{10}/2} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\alpha/2 + i\sin\alpha/2 & 0 \\ 0 & 0 & \cos\alpha/2 - i\sin\alpha/2 \end{pmatrix},$$
$$E_{10} = D_{10}(\pi) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & i & 0 \\ 0 & 0 & -i \end{pmatrix} \qquad (163)$$

are linear combinations of λ_3 and λ_8; in particular, we have $\lambda_3 + \lambda_9 + \lambda_{10} = 0$ and $\sqrt{3}\lambda_8 + \lambda_9 = \lambda_{10}$. Therefore, λ_9 and λ_{10} are *not* Gell-Mann matrices. (Also two further matrices corresponding to λ_8 in the other two triplets can be defined. The sum of these three matrices is 0 as well.)

The multiplication properties of the Gell-Mann generators λ_1 to λ_8 are listed in Table 10. To make the threefold symmetry more evident, the table also lists the products containing the linearly dependent matrices λ_9 and λ_{10}. Writing the table with the commutators would directly show that the generators form a Lie algebra.

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The *centre* of SU(3) – the subgroup that commutes with all other elements of the group – is Z_3; its threefold symmetry is useful in understanding the behaviour of the group elements and of the generators in more detail.

The group elements E_1 to E_8 listed above share the property that their fourth powers $(E_n)^4$ are the identity matrix. The first matrix triplet E_1, E_2, E_3, the second triplet E_4, E_5, E_9 and the third triplet E_6, E_7, E_{10} each form a SU(2) subgroup. Reflecting the threefold symmetry of its centre, SU(3) contains three linearly independent SU(2) subgroups. The group element E_8 commutes with the first triplet E_1, E_2, E_3; therefore, these four elements generate a U(2) subgroup of SU(3). This U(2) subgroup, often sloppily labeled as SU(2)xU(1), is given by those 3 by 3 matrices that contain a unitary 2 by 2 matrix in the upper left, contain zeroes in the remaining four off-diagonal elements, and contain the inverse value of the determinant of the 2 by 2 matrix in the remaining, lower right diagonal element. In short, SU(3) contains three linearly independent U(2) subgroups.

SU(3) is characterized by the way that the SU(2) triplets are connected. In particular, the product $E_3E_9E_{10}$ is the identity, reflecting the linear dependence of the three corresponding generators λ_n. We also have $E_8E_9 = E_{10}$. Also the product of E_8 with its

TABLE 10 The multiplication table for the generators λ_1 to λ_8 of SU(3), and for the additional, *linearly dependent* matrices $\lambda_9 = -\lambda_3/2 - \lambda_8\sqrt{3}/2$ and $\lambda_{10} = -\lambda_3/2 + \lambda_8\sqrt{3}/2$ that are *not* generators. Note that, despite the appearance, $\lambda_4^2 = \lambda_5^2 = \lambda_9^2$ and $\lambda_6^2 = \lambda_7^2 = \lambda_{10}^2$.

	λ_1	λ_2	λ_3	λ_4	λ_5	λ_9	λ_6	λ_7	λ_{10}	λ_8
λ_1	$2/3 + \lambda_8/\sqrt{3}$	$i\lambda_3$	$-i\lambda_2$	$\lambda_6/2 + i\lambda_7/2$	$-i\lambda_6/2 + \lambda_7/2$	$-\lambda_1/2 + i\lambda_2/2$	$\lambda_4/2 + i\lambda_5/2$	$-i\lambda_4/2 + \lambda_5/2$	$\lambda_1/2 + i\lambda_2/2$	$\lambda_1/\sqrt{3}$
λ_2	$-i\lambda_3$	$2/3 + \lambda_8/\sqrt{3}$	$i\lambda_1$	$i\lambda_6/2 - \lambda_7/2$	$\lambda_6/2 + i\lambda_7/2$	$-i\lambda_1/2 - \lambda_2/2$	$-i\lambda_4/2 + \lambda_5/2$	$-\lambda_4/2 - i\lambda_5/2$	$-i\lambda_1/2 + \lambda_2/2$	$\lambda_2/\sqrt{3}$
λ_3	$i\lambda_2$	$-i\lambda_1$	$2/3 + \lambda_8/\sqrt{3}$	$\lambda_4/2 + i\lambda_5/2$	$-i\lambda_4/2 + \lambda_5/2$	$-1/3 - \lambda_3/3 + \lambda_9/3$	$-\lambda_6/2 - i\lambda_7/2$	$i\lambda_6/2 - \lambda_7/2$	$-1/3 + \lambda_3/3 + \lambda_{10}/3$	$\lambda_3/\sqrt{3}$
λ_4	$\lambda_6/2 - i\lambda_7/2$	$-i\lambda_6/2 - \lambda_7/2$	$\lambda_4/2 - i\lambda_5/2$	$2/3 + \lambda_3/2 - \lambda_8/2\sqrt{3}$	$-i\lambda_9$	$i\lambda_5$	$\lambda_1/2 + i\lambda_2/2$	$i\lambda_1/2 - \lambda_2/2$	$-\lambda_4/2 - i\lambda_5/2$	$-\lambda_4/2\sqrt{3} - i\sqrt{3}\lambda_5/2$
λ_5	$i\lambda_6/2 + \lambda_7/2$	$\lambda_6/2 - i\lambda_7/2$	$i\lambda_4/2 + \lambda_5/2$	$i\lambda_9$	$2/3 + \lambda_3/2 - \lambda_8/2\sqrt{3}$	$-i\lambda_4$	$-i\lambda_1/2 + \lambda_2/2$	$\lambda_1/2 + i\lambda_2/2$	$i\lambda_4/2 - \lambda_5/2$	$i\sqrt{3}\lambda_4/2 - \lambda_5/2\sqrt{3}$
λ_9	$-\lambda_1/2 - i\lambda_2/2$	$i\lambda_1/2 - \lambda_2/2$	$-1/3 - \lambda_3/3 + \lambda_9/3$	$-i\lambda_5$	$i\lambda_4$	$2/3 + 2\lambda_3/3 + \lambda_9/3$	$\lambda_6/2 - i\lambda_7/2$	$i\lambda_6/2 + \lambda_7/2$	$-1/3 - \lambda_9/3 + \lambda_{10}/3$	$-1 + \lambda_{10}$
λ_6	$+\lambda_4/2 - i\lambda_5/2$	$i\lambda_4/2 + \lambda_5/2$	$-\lambda_6/2 + i\lambda_7/2$	$\lambda_1/2 - i\lambda_2/2$	$i\lambda_1/2 + \lambda_2/2$	$\lambda_6/2 + i\lambda_7/2$	$2/3 - \lambda_3/2 - \lambda_8/2\sqrt{3}$	$i\lambda_{10}$	$-i\lambda_7$	$-\lambda_6/2\sqrt{3} - i\sqrt{3}\lambda_7/2$
λ_7	$i\lambda_4/2 + \lambda_5/2$	$-\lambda_4/2 + i\lambda_5/2$	$-i\lambda_6/2 - \lambda_7/2$	$-i\lambda_1/2 - \lambda_2/2$	$\lambda_1/2 - i\lambda_2/2$	$-i\lambda_6/2 + \lambda_7/2$	$-i\lambda_{10}$	$2/3 - \lambda_3/2 - \lambda_8/2\sqrt{3}$	$i\lambda_6$	$i\sqrt{3}\lambda_6/2 - \lambda_7/2\sqrt{3}$
λ_{10}	$-\lambda_1/2 + i\lambda_2/2$	$-i\lambda_1/2 - \lambda_2/2$	$-1/3 + \lambda_3/3 - \lambda_{10}/3$	$-\lambda_4/2 + i\lambda_5/2$	$-i\lambda_4/2 - \lambda_5/2$	$-1/3 - \lambda_9/3 + \lambda_{10}/3$	$i\lambda_7$	$-i\lambda_6$	$2/3 - \lambda_3/3 + \lambda_9/3$	$1 + \lambda_9$
λ_8	$\lambda_1/\sqrt{3}$	$\lambda_2/\sqrt{3}$	$\lambda_3/\sqrt{3}$	$-\lambda_4/2\sqrt{3} + i\sqrt{3}\lambda_5/2$	$-i\sqrt{3}\lambda_4/2 - \lambda_5/2\sqrt{3}$	$-1 + \lambda_{10}$	$-\lambda_6/2\sqrt{3} + i\sqrt{3}\lambda_7/2$	$-i\sqrt{3}\lambda_6/2 - \lambda_7/2\sqrt{3}$	$1 + \lambda_9$	$2/3 - \lambda_8/\sqrt{3}$

companions from the other two triplets is the identity.

Finally, the product $(E_k E_l)^3$ for any k taken from the set $(1, 2, 4, 5, 6, 7)$ and any l from the same set, but from a *different* triplet, is also the identity matrix. This property of the third powers – taken together with the threefold symmetry of its centre – can be seen as the essential property that distinguishes SU(3) from other Lie groups. We now return to the strand model and show that slides indeed define an SU(3) group.

From slides to SU(3)

The *slide*, or *third Reidemeister move*, involves *three* pieces of strands. The textbook version of the third Reidemeister move – which is called E_0 here and is illustrated in Figure 69 – moves or 'slides' one strand, taken to be the horizontal blue one in the figure,

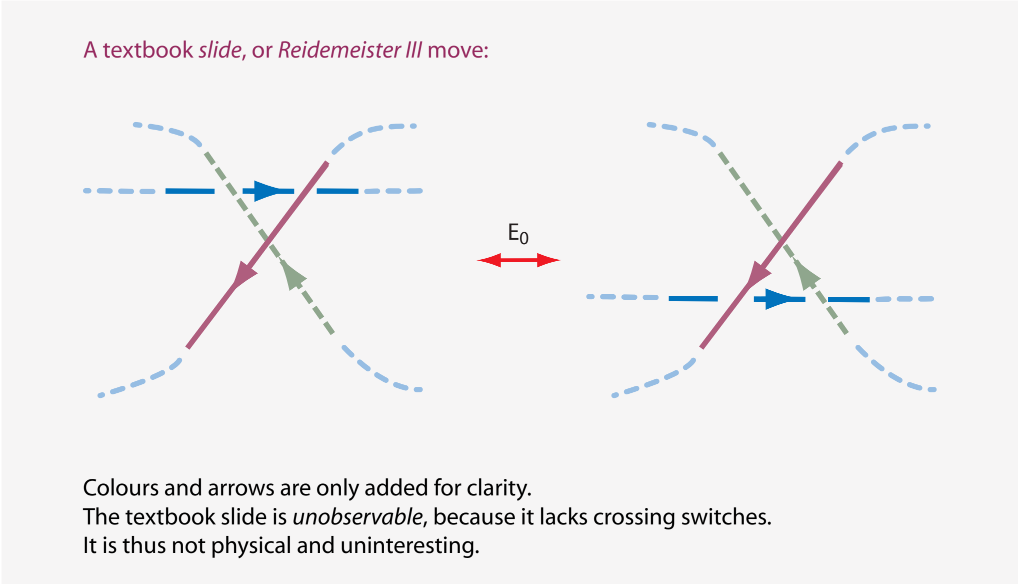


FIGURE 69 The textbook version E_0 of the slide move, or third Reidemeister move, is unobservable, because it does not involve crossing switches.

against a crossing of the other two. Equivalently, we can say that a slide pushes two strands against the blue strand that is kept in place. This textbook slide – we also call it a *pure slide* here – does not contain any crossing switch; following the fundamental principle of the strand model, it is therefore unobservable, or, simply said, of no physical relevance. However, related strand moves that do involve crossing switches do exist.

We introduce eight *generalized slides*, or slide-rotations, for a three-strand configuration; they are shown in Figure 70. We directly call these generalized slides E_1 to E_8, because they will turn out to correspond to the SU(3) group elements with the same name that were introduced above. In other words, we will show that the *generalized slides* E_n are elements of a Lie group SU(3); in particular, they obey all the properties expected from the correspondence with the SU(3) generators λ_n in Gell-Mann's choice:

$$E_n = e^{\pi i \lambda_n / 2} . \tag{164}$$

In the strand model, the generators λ_n describe the difference between an infinitesimal generalized slide – thus a slide-rotation with a rotation by an infinitesimal angle – and the identity. For slides, concatenation is equivalent to group multiplication, as expected. Slides form a group. We will now show that the slide generators obey the multiplication table already given in Table 10.

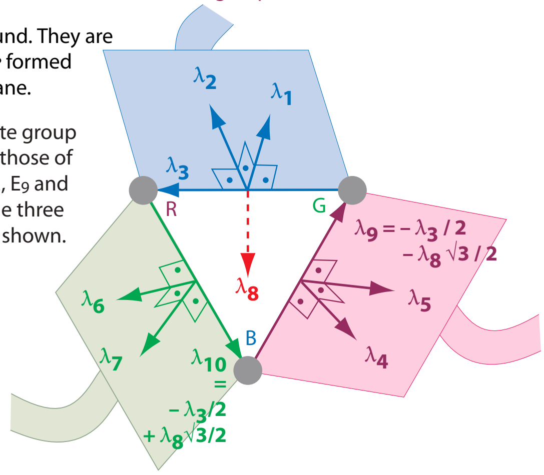
To see how the SU(3) multiplication table follows from Figure 70, we first note that the starting strand configuration of the Reidemeister III move contains, if all spatial configurations are considered, the same threefold symmetry as the centre of SU(3). In particular, like the generators and the basis vectors of SU(3), also the slides of the figure can be grouped into three triplets.

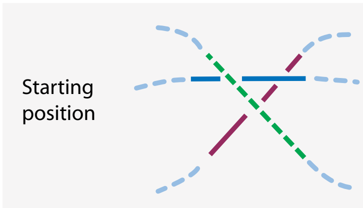
We now focus on the first triplet, the one formed by the three slides E_1, E_2 and E_3.

The *generalized slides*, or *Reidemeister III moves*, acting on three strands, form an SU(3) group.

The 8 generalized slides are shown below, with grey background. They are local slides and rotations by an angle π of an *imaginary buckle* formed by (usually) two strands. The strands lie (mostly) in a paper plane.

For each SU(2) triplet inside SU(3), the rotation axes of the finite group elements E_n are arranged at right angles to each other, as are those of the generators λ_n shown on the right. The rotation axes for E_3, E_9 and E_{10} are parallel; they are perpendicular to the paper plane. The three imaginary belt buckles for the three SU(2) subgroups are also shown.





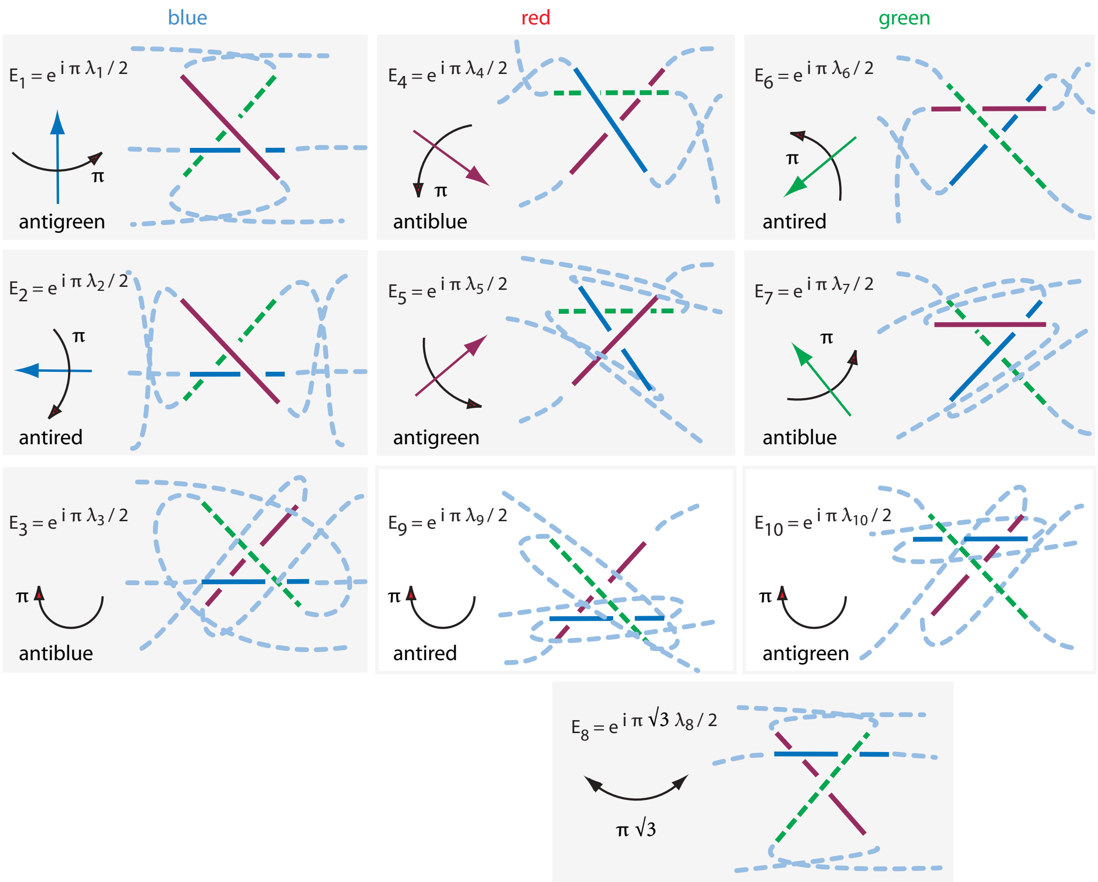


FIGURE 70 The strand deformations for the generalized slide moves E_n. The corresponding generators λ_n lead to an SU(3) structure, as shown in the text. Note that the rotation vectors for the generators λ_n and for the generalized slide moves E_n differ from each other. For clarity, the figure shows, instead of the deformation of the strand under discussion, the complementary deformations of the other two strands.

To makes things clear, these moves can be pictured as combined deformations and slides of the red and green strands against the horizontal blue strand. We can imagine these moves like those of the belt trick, but acting on an *imaginary buckle* formed only by the red and green strands. These generalized slides do contain crossing changes; therefore they are observable and are of physical relevance.

We note that 'slide' is not a perfect term for the generalized deformations E_1 to E_8; in fact, we might prefer to call them *slide-rotations*, because they are slide-rotations by an angle π that are applied to an imaginary belt buckle. Despite the involved construction, these generalized, observable moves remain modelled on the textbook slide E_0; in particular, they require *three* strand segments. The generalized, observable moves just defined *differ* from the twists and pokes discussed above, in the sections on the electromagnetic and weak interactions; thus they differ from Reidemeister I and II moves. As a result, we will usually continue to call the generalized, observable moves simply *slides*.

For simplicity, we assume – similarly to what we did in the discussion about the weak interaction – that the three strand segments are (roughly) in a plane. This is an idealized situation; in fact, the arguments given in the following apply also to all other three-dimensional configurations of three strands. In particular, the same results appear if all three strands segments are assumed perpendicular to each other, instead of lying in a plane.

We note that the rotation axes of the generalized slides E_1 and E_2 are neither aligned nor orthogonal to the paper plane. More precisely, the rotation axes of E_1, E_4 and E_6 are perpendicular to the sides of a cube. E_2, E_5 and E_7 are perpendicular to them. For the first triplet, the rotation axes E_1, E_2 and E_3 form an orthonormal basis; the same is valid for the other two triplets. We now show that the slides of the first triplet define an SU(2) group.

The observable, generalized slides in the triplet E_1, E_2 and E_3 can be concatenated. We distinguish two cases. The first case is the concatenation of any such slide with itself. The result corresponds to a rotation by 2π of the chosen strand pair and its imaginary belt buckle, and thus induces a corresponding amount of tail twisting. In fact, when any slide of the triplet is concatenated *four* times with itself, the result is the identity operation. Comparing a twofold and a fourfold concatenation, we see that they differ only by an entangling, or algebraically, by a minus sign for the imaginary buckle. This already realizes half of the belt trick that visualizes SU(2).

The other case to be checked is the concatenation of two different slides of the triplet. The result is always the third slide of the triplet (up to a sign that depends on whether the combination is cyclical or not). This behaviour realizes the other half of the belt trick. In short, we have shown that the triplet containing the first three generalized slides defines an SU(2) group. More precisely, the infinitesimal slide-rotations λ_1, λ_2 and λ_3 corresponding to the finite SU(3) elements E_1, E_2 and E_3 generate the SU(2) Lie algebra of an SU(2) Lie group. The SU(2) subgroup just found is just one of the three linearly independent SU(2) subgroups of SU(3). The generators of the first slide triplet thus reproduce Page 263 the nine results in the upper left of Table 10. We thus retain that we can indeed visualize the first three generalized slides with the help of the three orthogonal rotations by π of an imaginary belt buckle formed by the red and green strands.

For the visualization of SU(3) it is essential to recall that the direction in three-dimensional space of the vectors visualizing λ_n and those visualizing E_n differ from each other. This already the case for U(1).

The remaining generalized slides that are possible in the three-strand configuration are easily constructed using the threefold symmetry of the strand configuration; they are illustrated in Figure 70. For each of the three strand segments there is a triplet of observable slides; this yields a total of *nine* possible generalized slides for the observer defined by the paper plane. In the second triplet, the slides corresponding to E_1 and E_2 are called E_4 and E_5, and in the third triplet they are called E_6 and E_7. For the three slides corresponding to E_3 – we call the other two E_9 and E_{10} – *only two* generators are linearly independent. Indeed, the figure shows that $E_3 E_9 E_{10}$ – whose axes are all three parallel – is the identity matrix; this expected from an SU(3) structure. The three operations E_3, E_9 and E_{10} also commute with all other operations; thus they form the centre of the group defined by all E. The second linearly independent, generalized slide of common use, E_8, is also shown in the figure; it is a linear combination of E_9 and E_{10}. We note that the strand model also visualizes the factor $\sqrt{3}$ in the definition of E_8. In total, we get *eight* linearly independent generalized slides. All slides, except for E_8, act on an imaginary belt buckle that is formed by two strands.

We saw that the generators corresponding to the slides E_1, E_2 and E_3 generate an SU(2) subgroup. The same holds for the corresponding triplet E_4, E_5 and the linear combination $E_9 = -E_3/2 - E_8\sqrt{3}/2$ (corresponding to E_3), and for the triplet E_6, E_7 and $E_{10} = -E_3/2 + E_8\sqrt{3}/2$. For each of these slides, a fourfold concatenation yields the identity; and inside each triplet, the concatenation of two different slides yields a multiple of the third slide. In short, for each triplet, the corresponding infinitesimal slides generate an SU(2) group. These three SU(2) groups are linearly independent. We have thus reproduced an important part of the structure of SU(3). In addition, we have found a visualization of SU(3); since each SU(2) group can be represented by a separate imaginary buckle, the group SU(3) can be visualized – in many, but not all in aspects – with the help of three imaginary buckles. The top right of Figure 70 illustrates this visualization.

The correspondence of the slides and the multiplication table increases further if we change slightly the definition of the first triplet. In this first triplet we can take as imaginary buckle the set of *all three* central segments. Moving all three strands together simplifies the visualization, because for the first triplert, the blue strand is trapped between the other two strands. In this way, generalized slide still consists of a rotation followed by a slide. And we still have a SU(2) subgroup for the first triplet.

The slide E_8 differs from the other slides, as expected from SU(3). It describes a motion that rotates the red and green strands in opposite directions; this is illustrated in Figure 70. E_8 is thus *not* well described with an imaginary belt buckle. It is straightforward to check that the slide E_8 commutes with E_1, E_2, E_3 and obviously with itself, but not with the other generalized slides. Together, E_8 and the first triplet thus form a U(2) Lie group, as expected. In addition, we find that E_8 commutes with E_9 and E_{10}, and that $E_8 E_9 = E_{10}$, as expected from SU(3). The strand model also implies that the product of E_8 with its two counterparts from the other triplets is the identity matrix, as expected from SU(3).

The last step to show the equivalence of slides and SU(3) requires us to confirm the multiplication properties – between slides E_n or between generators λ_n – from *different* triplets. In fact, because of the three-fold symmetry of the centre, we only need to check two multiplication results between slides from different triplets: one that either involves

λ_3 or λ_8, and one that does not.

We begin with products involving λ_3 and one of the first two elements of another triplet. Such products yield a weighted sum of generators of the triplet. It is easier to check these product properties by using the exemplary relation between finite group elements $E_5E_3E_4 = E_3$. Note that only this specific permutation of 5, 3 and 4 yields this result. Playing with the strand model confirms the relation. Similar comments apply to $E_6E_3E_7 = E_3$ – and to the corresponding products involving E_9, such as $E_1E_9E_2 = E_9$, or E_{10}, such as $E_1E_{10}E_2 = E_{10}$ – as well as $E_4E_8E_5 = E_8$ and $E_6E_8E_7 = E_8$. The strand model allows anybody to check that these relations are satisfied.

We continue with the exemplary product $\lambda_5\lambda_7$, respectively E_5E_7. We note a basic difference between a product like $\lambda_5\lambda_7$ and any product of two generators from the same triplet. The product $\lambda_5\lambda_7$ – like the other concatenations of generators from different triplets – does not yield a single generator, but yields a combination, i.e., a *sum* of generators. The combination is not easy to visualize with strands; an easier way is to check the SU(3) algebra using the properties of the product E_5E_7.

As mentioned above, in SU(3), for products involving the first two members from different triplets, *the threefold concatenation* $(E_iE_j)^3$ *is the identity.* And indeed, Figure 70 confirms that $(E_2E_4)^3$ or $(E_5E_7)^3$ is the identity. Similarly, also the other products can be tested with the help of three strands.

Using the visualization with three strands, we have thus confirmed all products of generators from two different triplets that appear in Table 10. We note that Figure 70 also illustrates that the three slides E_2, E_5 and E_7 generate an SO(3) group, the rotation group in three dimensions. In order to see this, we observe that the infinitesimal versions of the three slides generate all possible rotations in three dimensions of the central triangle. An SO(3) group also appears for the slides 1, 4 and 7, for the slides 1, 5 and 6, and for the slides 2, 4 and 6. These are the four basic SO(3) subgroups of SU(3). The remaining combinations of three operations from three different triplets – such as 1, 4 and 6, or the combination 1, 5 and 7, or the combination 2, 4 and 7, or the combination 2, 5 and 6 – do *not* generate any subgroup. This can be confirmed by exploring the corresponding strand moves.

We can conclude: in a region with three strands crossing each other, the eight linearly independent, generalized slides that can be applied to that region define the group SU(3). In other words, *the group SU(3) follows from the third Reidemeister move.*

In the same way as for the other gauge groups, we find that particles whose strand models contain configurations with three strand segments can be subject to an SU(3) gauge interaction. In experiments, this interaction is called the *strong nuclear interaction*. The strong interaction is due to the Reidemeister III move. Like for the other interactions, a particle will only interact strongly if its tangle is not too symmetric, because in the symmetric case, averaged over time, there will be no net interaction. We will clarify the details below, when we discuss the specific tangles and colour charges of the different elementary matter particles.

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The strand model for gluons

Physically, the eight slides corresponding to the Gell-Mann matrices represent the effects of the eight *gluons*, the intermediate vector bosons of the strong interaction, that can act on a particle.

▷ Given that the eight slides E_1 to E_8 represent the effects of the eight gluons, they also represent the gluons themselves.

Interactions are transfers of a tangle process to another tangle. Therefore

▷ The absorption of a gluon is a slide that is transferred to another particle.
▷ The emission of a gluon is a slide that is transferred to three vacuum strands.

To visualize the concept of gluon even further, we can say that every gluon can be described as a strand structure that continuously performs an SU(3) operation, i.e., a generalized slide continuously repeating itself. We found a similar correspondence for the other gauge interactions. In case of the electromagnetic interaction, the intermediate vector boson, the photon, can be described as a strand that continuously performs a U(1) operation, i.e., a rotation. In case of the weak interaction, a weak intermediate vector boson can be described as a strand that continuously performs an SU(2) operation, i.e., an operation from the belt trick. This is most evident in the unbroken form of the weak bosons.

Every gluon can also be seen as the deformation of a single strand that drags its surrounding with it. This single strand description of gluons implies that gluons have vanishing mass and vanishing charge. This single strand description of gluons also implies that they have spin 1, as is observed. The strand model of the gluon also implies that free gluons would have a huge energy.

Challenge 168 e

The SU(3) multiplication table confirms that the eight gluons transform according to the adjoint (and faithful) representation of SU(3). Therefore, each row or column in a Gell-Mann matrix thus corresponds to one of the three *colours* of the strong interaction. The exploration of slide concatenation also showed that two general slides do not commute and do not anticommute. The group SU(3) is *non-Abelian*. This implies that gluons interact among themselves. Both the multiplication table and the strand model for gluons imply that two interacting gluons can yield either one or two new gluons, but not more. This is illustrated in Figure 71. The strand model, through its generation of SU(3), thus implies that gluons interact among themselves, but only in triple and quartic gluon vertices.

Slides – i.e., gluon emission or absorption – never change the topology of tangles, and in particular, of matter tangles. Therefore, the strand model predicts that the strong interactions conserve electric charge, baryon number, weak isospin, flavour, spin and all parities. This is indeed observed. In particular, there is a natural lack of C, P and CP violation by slides. This is precisely what is observed for the strong interaction.

Because gluons do not change the topology of the particle tangles they act upon, but only change their shape, gluons are predicted to be massless in the strand model, despite interacting among themselves. And because gluons interact among themselves, free gluons are predicted not to appear in nature. And of course, all these conclusions agree with experiments.

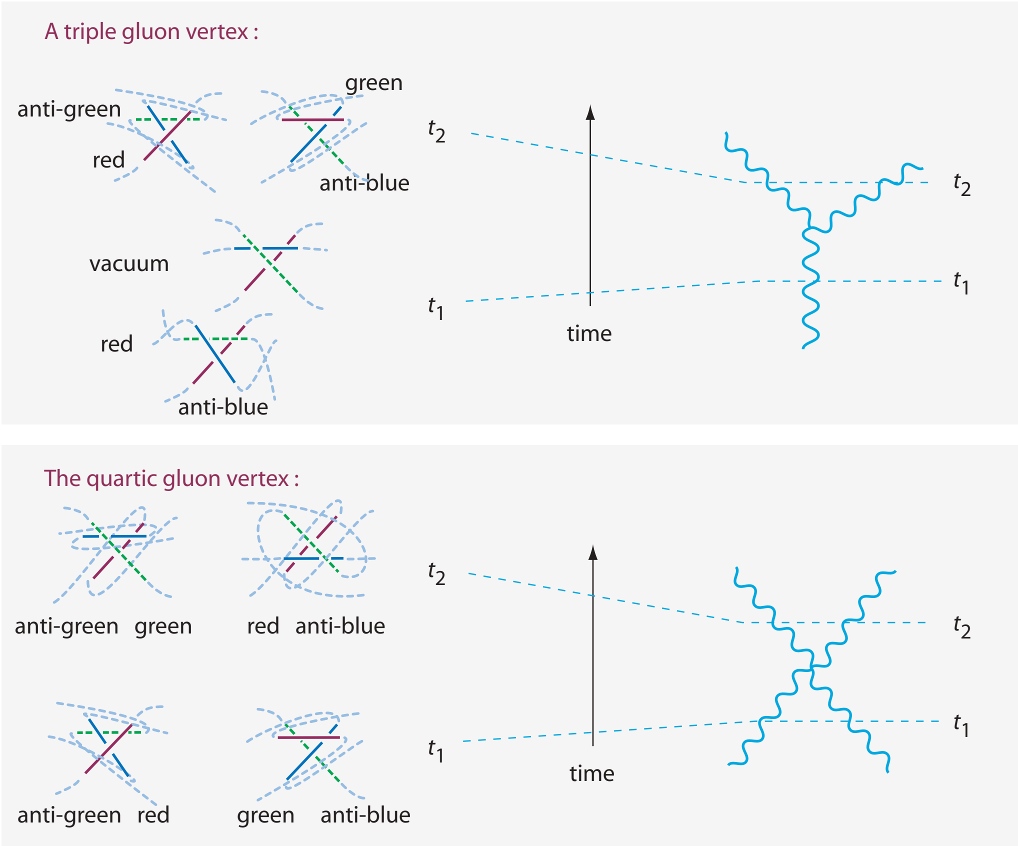

FIGURE 71 The two types of self-interaction of gluons in the strand model.

In summary, we have shown that in the strand model, the strong nuclear interaction and all its properties appear automatically form slides, i.e., from Reidemeister III moves. In particular, the strand model implies that the Lagrangian of strongly interacting fermions has a SU(3) gauge invariance that is due to generalized slide deformations.

The gluon Lagrangian

Gluons are massless particles with spin 1. As a result, the field intensities and the Lagrangian are determined in the same way as for photons: energy density is the square of crossing density, i.e., the 'square' of field intensity. Since there are 8 gluons, the Lagrangian density becomes

$$\mathcal{L}_{\text{gluons}} = -\frac{1}{4}\sum_{a=1}^{8} G^a_{\mu\nu} G^{\mu\nu}_a \tag{165}$$

where the gluon field intensities, with *two* greek indices, are given naturally as

$$G^a_{\mu\nu} = \partial_\mu G^a_\nu - \partial_\nu G^a_\mu - g f^{abc} G^b_\mu G^c_\nu \,, \tag{166}$$

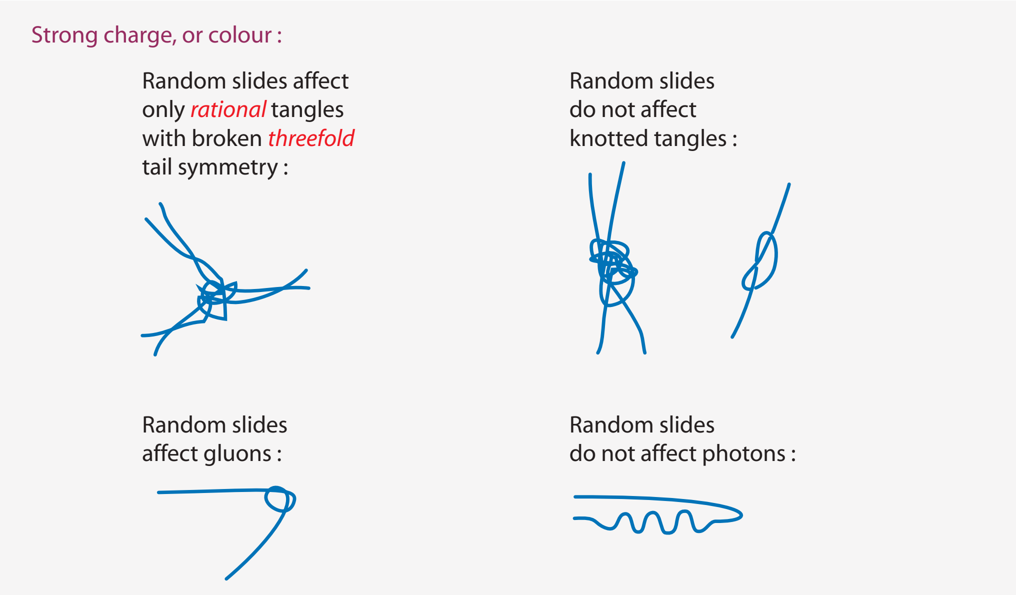

FIGURE 72 Tangles with and without colour charge. (This image needs to be updated: no knotted tangles occur in nature.)

and f_{abc} are the structure constants of SU(3) that can be deduced from the multiplication table given above. The quantities G^a_μ, with *one* greek index, are the gluon vector potentials. The last term in the definition of the field intensities corresponds to the triple and quartic vertices in the Feynman diagrams of gluon interactions. They are shown in Figure 71. The Lagrangian is simply the natural generalization from the U(1) case of photons to the SU(3) case of gluons. In short, we obtain the usual free gluon Lagrangian from the strand model.

Page 234

Colour charge

Surrounded by a bath of gluons that randomly induce slides of all kinds, not all fermion cores will change their rotation state. Generally speaking, particles have colour if a bath of random gluons changes their phase. Only tangles which lack some symmetry will therefore possess colour charge. Tangle that are symmetric will be neutral, or 'white'. Which symmetry is important here?

We see directly that the *photon* tangle is not sensitive to a gluon bath. The same is valid for W and Z bosons. These tangles are too simple. The strand model predicts that these particles are colour-neutral, i.e., that they are 'white', as is observed.

On the other hand, the multiplication properties given above shows that *gluons* interact among themselves and thus that they have colour charge. In fact, group theory shows that their properties are best described by saying that gluons have a colour and an anticolour; this is the simplest way to describe the representation to which they belong. In short, the strand model of gluons automatically implies that they carry both a colour and an anti-colour.

Fermions behave differently. In the strand model, a fermion has colour charge if the

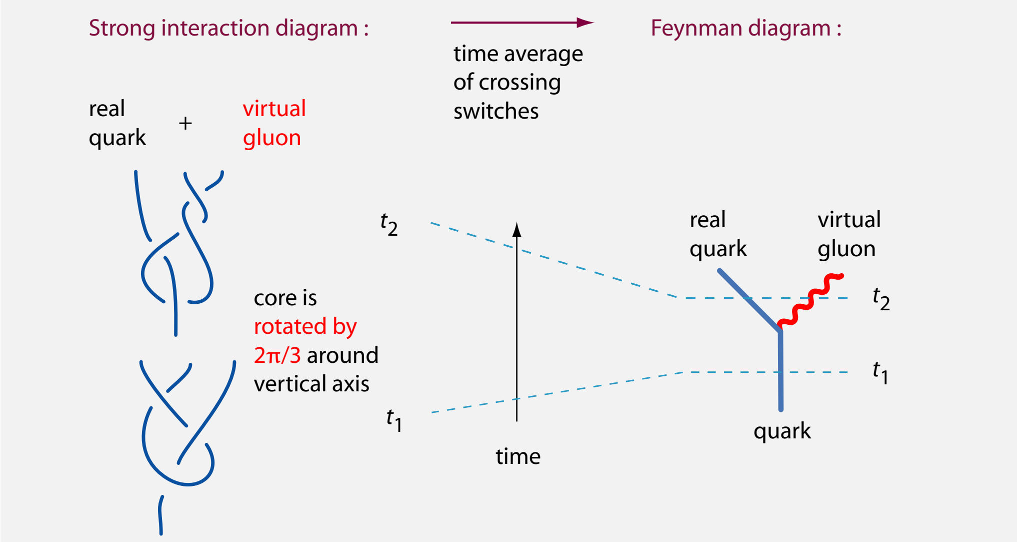

FIGURE 73 The Feynman diagram of the strong interaction for a quark. The upper triplet of tails correspond to the three belts.

corresponding triple belt model is affected by large numbers of random gluons. The first tangles that come to mind are tangles made of three strands, such as the simple tangles shown in Figure 70. But a short investigation shows that such tangles are colour-neutral, or 'white'. We will see below that this implies that *leptons* are colour-neutral, or 'white'. In contrast, a rational fermion tangle does not suffer this fate. (We recall that a so-called *rational tangle* is by definition made of exactly *two* strands; a two-stranded tangle is rational if the two strands can be untangled just by moving the tails around.) In a bath of gluon strands that induce slides, i.e., third Reidemeister moves, a general *rational tangle* made of two strands is expected to be influenced, and thus to be colour-charged.

Rational tangles made of two strands are the simplest possible tangles with colour. A tangle is called *rational* if it can be untangled just by moving the tails around. An example of a rational tangle is shown in Figure 73. Such tangles break the three-fold symmetry of the three-belt structure, and are thus colour-charged. We will show below how these tangles are related to *quarks*. We can thus say:

▷ A fermion tangle has ***colour charge*** if its three-belt model is not symmetric for rotations by $\pm 2\pi/3$.

Coloured rational tangles automatically have *three* possible colours:

▷ The ***three colour charges*** are the three possibilities to map a tangle to the three belt model.* *Each colour is thus a particular orientation in ordinary space.*

If we want to explore more complicated types of tangles of two strands, such as *prime*

Challenge 169 ny

* Can you define a geometric or even a topological knot invariant that reproduces colour charge?

Page 318 tangles or *locally knotted* tangles, we recall that such tangles are not part of the strand model. The strand model thus predicts that rational tangles made of two strands are the basic colour states. And indeed, in nature, quarks are the only fermions with colour charge.

We can summarize that colour charge is related to orientation in space. The three possible colours and anticolours are consequences of the possible orientations along the three dimensions of space.

Properties of the strong interaction

In the strand model, all interactions are *deformations* of the tangle core. Specifically, the strong interaction is due to exchange of *slides*. Particles have strong charge, or colour, if their tangles lack the three-belt symmetry just specified. In the case of coloured fermions, *colour change* is a change of the mapping to the three-belt model, i.e., a change of orientation of the tangle in space.

If we use the strand definition of the strong interaction, visual inspection shows us that slide exchanges, and thus gluon exchanges, are deformations that conserve topology; therefore gluon exchange *conserves colour*. Since the strong interaction conserves the topology of all involved tangles and knots, the strong interaction also *conserves electric charge*, *parity*, and, as we shall see below, *all other quantum numbers* – except colour itself, of course. All these results correspond to observation.

The Lagrangian of QCD

We started from the idea that tangle core deformations lead to phase redefinitions. We then found that slides imply that the strong interaction Lagrangian for matter and for radiation fields is SU(3) gauge invariant. If we include these two gauge invariances into the fermion Lagrangian density from the Dirac equation, we get

$$\mathcal{L}_{\mathrm{QCD}} = \sum_q \overline{\Psi}_q (i\hbar c \not{D} - m_q c^2 \delta_{qq'}) \Psi_{q'} - \frac{1}{4} \sum_{a=1}^{8} G^a{}_{\mu\nu} G_a{}^{\mu\nu} , \qquad (167)$$

where the index q counts the coloured fermion, i.e., the quark. In this Lagrangian density, $\not{D}$ is now the SU(3) gauge covariant derivative

$$\not{D} = \not{\partial} - g\, \gamma^\mu G^a_\mu \lambda_a , \qquad (168)$$

where g is the gauge coupling, λ_a are the generators of SU(3), i.e., the Gell-Mann matrices given above, and the G^a_μ are, as before, the gluon vector potentials. The last term in the covariant derivative corresponds to the Feynman diagram and the strand diagram of Figure 73. This is the Lagrangian density of QCD.

In summary: the strand model reproduces QCD. However, we have not yet deduced the number and masses m_q of the quarks, nor the strong gauge coupling g.

Renormalization of the strong interaction

The slide move description of the strong interaction implies that only three Feynman diagrams are possible: one QCD Feynman diagram is possible for quarks, and only the triple and the quartic vertices are possible among gluons. This limited range of options allowed us to deduce the QCD Lagrangian. The limited range of options is also essential for the *renormalization* of QCD. The strand model thus automatically ensures that the strong interaction is renormalizable.

In short, the strand model provides a new underlying picture for the Feynman diagrams of the strong interaction, but does not change the physical results at any energy scale accessible in the laboratory. In particular, the measured running of the strong coupling constant is reproduced. Page 386 Indeed, in the strand model, a flux-tube–like bond between the quarks appears automatically, as we will see when exploring hadrons. Page 332, page 340 At high kinetic energies, the bond has little effect, so that quarks behave more like free particles. In short, we find that the strand model reproduces *asymptotic freedom* and also provides an argument for quark confinement. We will return to the issue in more detail below. Page 337

Curiosities and fun challenges about SU(3)

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Deducing the Lie group SU(3) from a three-dimensional model is a new result. In particular, deducing the gauge group SU(3) as a *deformation gauge group* is new. Frank Wilczek, Alfred Shapere, Alden Mead, Jerry Marsden and several others have confirmed that before this discovery, only the geometric Lie group SO(3) and its subgroups had been found in deformations. Ref. 186 The fundamental principle of the strand model shows its power by overcoming this limitation. (Apparently, nobody had even realized that the belt trick already implies the possibility of an SU(2) gauge group for deformations.)

* *

We have discussed the *shape deformations* that lead to the SU(3) group. But what are the precise *phase choices* for a crossing that lead to SU(3) invariance? Challenge 170 ny

* *

Do the two linear independent gluons with lined-up tails have the same properties as the other six gluons? Challenge 171 ny

* *

Three strands can cross each other also in another way, such that the three strands are interlocked. Why can we disregard the situation in this section? Challenge 172 s

* *

Deducing the Lie groups U(1), SU(2) and SU(3) directly from a basic principle contradicts another old dream. Many scholars hoped that the three gauge groups have something to do with the sequence complex numbers, quaternions and octonions. The strand model quashes this hope – or at least changes it in an almost unrecognizable way.

* *

Challenge 173 e The tangles for the W and Z bosons have no colour charge. Can you confirm this?

* *

Challenge 174 ny The Lie group SU(3) is also the symmetry group of the three-dimensional harmonic oscillator. What is the geometric relation to the Lie group SU(3) induced by slides?

* *

Challenge 175 e Confirm that the strand model does not contradict the Coleman–Mandula theorem on the possible conserved quantities in quantum field theory.

* *

Challenge 176 e Confirm that the strand model does not contradict the Weinberg–Witten theorem on the possible massless particles in quantum field theory.

* *

Challenge 177 d Are the *Wightman axioms* of quantum field theory fulfilled by the strand model with interactions? The *Haag–Kastler axioms*? Is Haag's theorem circumvented?

* *

Ref. 187 Challenge 178 ny Show that the BCFW recursion relation for tree level gluon scattering follows from the strand model.

Summary on the strong interaction and experimental predictions

We have deduced the Lagrangian density of QCD from the strand model with the help of slides. Is there a difference between the strand model and QCD? No, not as long as gravity plays no role. The strand model predicts that gravitation only comes into play near the Planck energy $\sqrt{\hbar c^5/4G}$. And indeed, accelerator experiments have not yet found any effect that contradicts QCD, and therefore no effect that contradicts the strand model of the strong interaction.

The strand model also predicts that the strong interaction is naturally CP-invariant. This means that axions – particles invented to explain the invariance – are unnecessary: Page 353 as shown below, the strand model even predicts that they do not to exist. Both predictions agree with experiment.

The strand model of the strong interaction implies that the SU(3) gauge symmetry is valid at all energies. No other gauge group plays a role in the strong interaction. The strand model thus predicts again that there is no grand unification in nature, and thus no larger gauge group. Often discussed groups such as SU(5), SO(10), E6, E7, E8 or SO(32) are predicted not to apply to nature. Also this prediction is not contradicted by experiment.

The strand model further predicts that the combination of gravity and quantum theory turns all Planck units into *limit* values. The strand model thus predicts a maximum strong field value given by the Planck force divided by the strong charge of the quark. All physical systems – including all astrophysical objects, such as neutron stars, quark stars, gamma-ray bursters or quasars – are predicted to conform to this field limit. So far, this prediction is validated by experiment.

In summary, we have shown that Reidemeister III moves – or slides – in tangle cores lead to an SU(3) gauge invariance and a Lagrangian that reproduces the strong interaction. Colour charge is related to the topology of certain rational tangles. In this way, we have deduced the origin and most observed properties of the strong interaction. We have Page 164 thus settled another issue of the millennium list. However, we still need to deduce the tangles and the number of quarks, their masses and the strength of the strong coupling.

Summary and predictions about gauge interactions

At this point of our adventure, we have deduced gauge theory and the three known gauge Page 152 interactions from strands. Using only the fundamental principle, we explained the dimensions of space-time, the Planck units, the principle of least action, the appearance of the gauge groups U(1), broken SU(2) and SU(3), of renormalization, of Lorentz symmetry and of permutation symmetry. Thus we have deduced all the concepts and all the mathematical structures that are necessary to *formulate* the standard model of elementary particles.

In particular, the strand model provides a description and explanation of the three gauge interactions at Planck scales that is based on *deformations* of strands. The description of QED is illustrated in Figure 74: it shows the way that strands model the emission of a photon by an electron. The deduction of the three gauge interactions given in this text, with the help of the Reidemeister moves, is the first and, at present, the *only* explanation of the three gauge forces. No other explanation or deduction has ever been given.

We have shown that quantum field theory is an *approximation* of the strand model. The approximation appears when the strand diameter is neglected; quantum field theory is thus valid for all energies below the Planck scale. In other words, in contrast to many other attempts at unification, the strand model is *not a generalization* of quantum field theory. The strand model for the three gauge interactions is also unmodifiable. These Page 149 properties are in agreement with our list of requirements for a final theory.

We have not yet deduced the complete standard model: we still need to show which types of particles exist, which properties they have and what couplings they produce. However, we have found that the strand model explains all the mathematical structures Page 164 from the millennium list that occur in quantum field theory and in the standard model of particle physics. In fact, the strand explanation for the origin of the gauge interactions allows us to make several definite predictions.

Predicting the number of interactions in nature

Already in 1926, Kurt Reidemeister proved an important theorem about possible deformations of knots or tangles that lead to changes of crossings. When tangles are described with two-dimensional diagrams, all possible deformations can be reduced to *exactly three* Ref. 182 moves, nowadays called after him. In the strand model, the two-dimensional tangle diagram describes what an observer *sees* about a physical system. Together with the equivalence of interactions as crossing-changing deformations, Reidemeister's theorem thus proves that there are *only three gauge interactions* in nature. In particular, there is no fifth force. Searches for additional gauge interactions are predicted to fail. And indeed, they have all failed up to now.

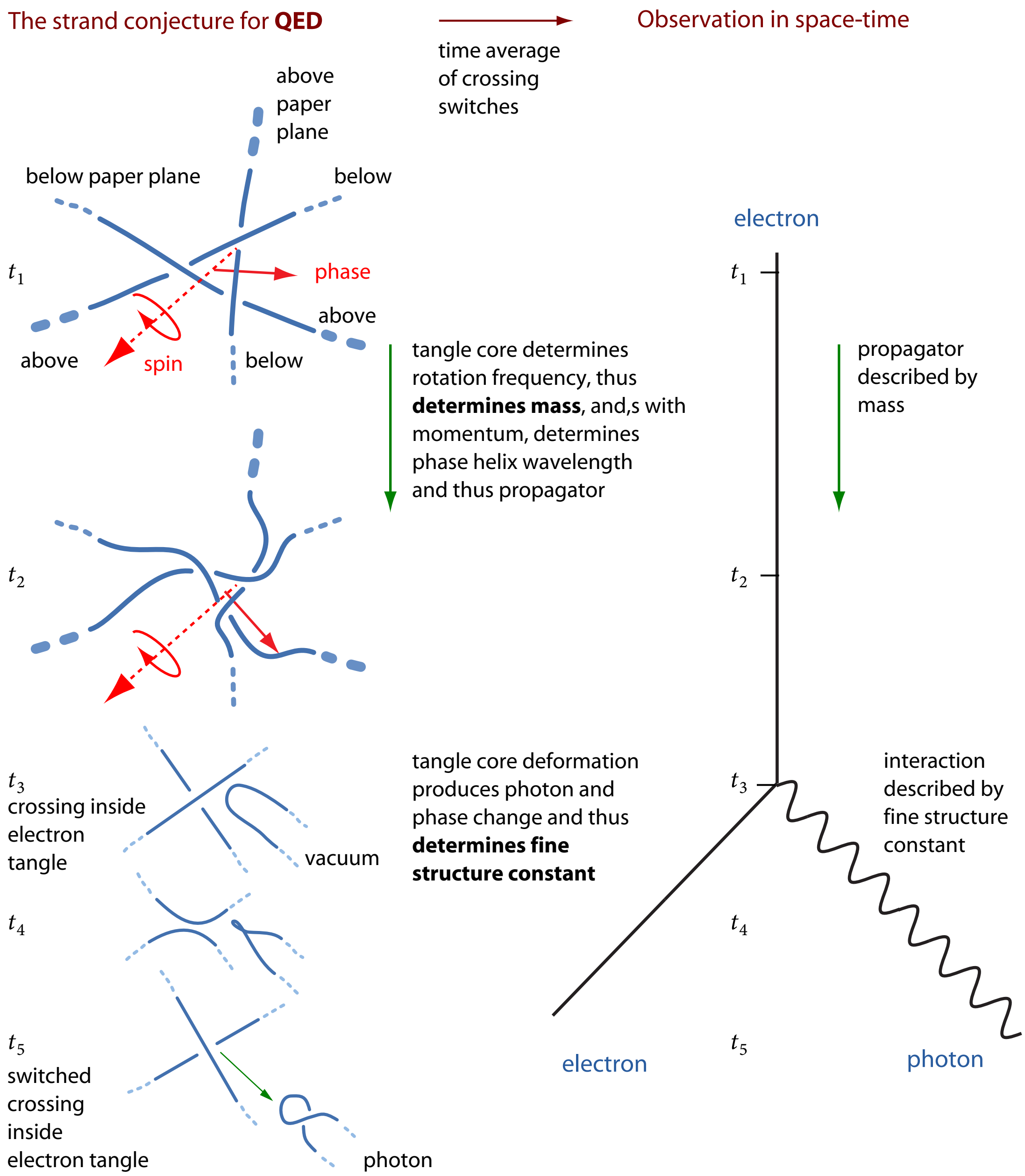

FIGURE 74 QED in one picture: In the strand model, the electron mass and the fine structure constant are determined by the tangle – only the simplest family member is shown here – and its shape change under fluctuations.

Unification of interactions

Ref. 142 We can also state that there is only *one* Reidemeister move. This becomes especially clear if we explore the three-dimensional shape of knots instead of their two-dimensional diagrams: all three Reidemeister moves can be deduced from the *same* deformation of a single strand. Only the projection on a two-dimensional diagram creates the distinction between the three moves. In the terms of the strand model, this means that all gauge interactions are in fact aspects of only one basic process, a fluctuation of strand shape, and that the three gauge interactions are only distinguished by their projections. In this way, the three gauge interactions are thus *unified* by the strand model.

The plane of projection used in a strand diagram defines a mapping from strand fluctuations to Reidemeister moves. The projection plane is defined by the observer, i.e., by the frame of reference. Depending on the projection plane, a general deformation is mapped into different Reidemeister moves. At first sight, the nature of an interaction – whether electromagnetic, strong or weak – seems to depend on the observer. In nature, however, this is not the case. But this contradiction is only apparent. In the strand model, the nature of interaction of a particle results from the type of asymmetry of its tangle core. Certain strand deformations do not lead to interactions, because their effects are suppressed by the averaging of short-time fluctuations underlying every observation. In other words, the averaging process at the basis of observations also ensures that interactions are effectively observer-independent at low energy.

Page 385
In short, the strand model provides a natural *unification* of the interactions. And this unification of the interactions differs completely from any past proposal. The final test, of course, can only be provided by experiment.

No divergences

The strand model implies that there are *no* divergences in the quantum description of nature. This lack of divergence occurs because all measurement values appear after strand effects have been averaged out. As mentioned above, strand effects on space-time disappear through 'shivering' and strand effects on particles disappear through wavefunctions.

Page 153

In summary, in the strand model, no interaction implies or contains divergences: neither gravity nor the gauge interactions. There are neither ultraviolet nor infrared divergences. The strand model avoids divergences, infinities and singularities of any kind from its very start.

Grand unification, supersymmetry and other dimensions

The three gauge interactions are due to the three Reidemeister moves. Therefore, the strand model asserts that there is *no* single gauge group for all interactions. In short, the strand model asserts that there is *no* so-called *grand unification*. The absence of grand unification implies the absence of large proton decay rates, the absence of additional, still undiscovered gauge bosons, the absence of neutron–antineutron oscillations, and the absence of sizeable electric dipole moments in elementary particles. All these searches are ongoing at present; the strand model predicts that they yield *null results*.

Page 396

Supersymmetry and approaches based on it assume gauge group unification. However, as just explained, the strand model predicts that there is no supersymmetry and therefore no supergravity. The strand model also predicts the absence of all conjectured 'superparticles'. In 2016 and again in 2017, the numerous experiments at CERN confirmed the prediction: there is no sign of supersymmetry in nature.

Reidemeister moves are confined to three spatial dimensions. Indeed, the strand model is based on exactly three spatial dimensions. It predicts that there are no other, undetected dimensions of space. The strand model also predicts the absence of non-Page 348 commutative space-time, even though, with some imagination, strands can be seen as remotely related to that approach. Finally, the strand model predicts the lack of different vacua: the vacuum is unique.

In short, the strand model differs both experimentally and theoretically from the unification proposals made in the twentieth century. In particular, the strand model predicts the *absence* of additional symmetries, of additional energy scales, and of additional space-time properties at high energy. The strand model predicts that unification is not achieved by searching for higher symmetries, nor for higher dimensions, nor for concepts that contain both. This lack of complex mathematical or symmetry concepts in nature is disappointing; the hopes and search activities in the last fifty years are predicted to have been misguided. In other words, the predictions of the strand model are unpop-Page 149 ular. However, these predictions agree with our list of requirements for a final theory; Page 412 and so far, all these predictions agree with experiment.

No new observable gravity effects in particle physics

Page 8 In the 'cube' structure of physics shown in Figure 1, the transition from the final, unified description to quantum field theory occurs by neglecting gravity, i.e., by assuming flat space-time. The same transition occurs in the strand model, where neglecting gravity in addition requires neglecting the strand diameter. In this way, the gravitational constant G disappears completely from the description of nature.

We can summarize our findings on quantum field theory also in the following way:

▷ The strand model predicts that particle masses are the only observable effect of gravity in quantum physics and in particle physics.

Page 311 This result will be complemented below by a second, equally restrictive result that limits the observable quantum effects in the study of gravity. In short, the strand model keeps particle physics and general relativity almost completely separated from each other. This is a consequence of the different effects produced by tail deformations and by core deformations. And again, the prediction of a lack of additional gravitational effects in particle physics agrees with all experiments so far.

The status of our quest

In this chapter, we have deduced that strands predict exactly three interactions. Interactions are deformations of tangle cores and just three classes of such core deformations exist. The three classes of deformations are given by the three Reidemeister moves. Because of the properties of the Reidemeister moves, the three interactions are described by a U(1), a broken SU(2) and a SU(3) gauge symmetry, respectively.

Strands also show that the three interactions are renormalizable, relativistically invariant, and that they follow the least action principle. Strands thus imply the three interaction Lagrangians of the standard model of particle physics. In addition, strands predict the absence of other interactions, symmetries and space-time structures.

Page 164 If we look at the millennium list of open issues in fundamental physics, we have now solved all issues concerning the mathematical structures that appear in quantum field theory and in the standard model of particle physics.

▷ All mathematical structures found in quantum physics result from the fundamental principle of the strand model.

Equivalently, extension contains all quantum effects. This is an intriguing result that induces us to continue our exploration. Only two groups of issues are still unexplained: the theory of general relativity and the spectrum of elementary particles. We proceed in this order.

CHAPTER 10

GENERAL RELATIVITY DEDUCED FROM STRANDS

General relativity describes the deformations of the vacuum. In everyday life, gravitation is the only such effect that we observe. But on astronomical scale, gravity shows more phenomena: vacuum can deflect light, producing gravitational lenses, can wobble, giving gravitational waves, and can accelerate, yielding the darkness of the sky and the fascinating black holes. All these observations require general relativity for their description. Therefore, general relativity must be part of any unified description of nature.

In the following, we explain the existence of gravity as a consequence of strands. Then we deduce the field equations of general relativity, the entropy of black holes and relativistic cosmology from the strand model. We also predict the outcome of many quantum gravity experiments. Finally, we deduce the consequences of strands for cosmology. We include several experimental predictions. Of all Planck-scale models of space or space-time, strands seem to be the simplest one that provides these deductions.

FLAT SPACE, SPECIAL RELATIVITY AND ITS LIMITATIONS

Page 209
We have seen above that any observer automatically introduces a 3+1-dimensional *background* space-time. We have also seen that in the case of quantum theory, *physical* space-time, the space-time that is formed by the fluctuations of the vacuum strands, is naturally 3+1-dimensional and flat. In the absence of gravity, physical space and background space coincide.

Page 213
Using strands, we have deduced:

▷ c is the invariant limit for all energy speeds.

This limit is achieved only by free massless particles, such as photons. Strands also showed us that massive particles move more slowly than light. In short, strands reproduce special relativity.

The strand model thus predicts that *pure* special relativity is correct for all situations and all energies in which gravity and quantum theory play no role. The strand model also predicts that when gravity or quantum effects do play a role, general relativity or quantum theory *must* be taken into account. This means that there is no domain of nature in which intermediate descriptions are valid.

It is sometimes suggested that the invariant Planck energy limit for elementary particles might lead to a 'doubly special relativity' that deviates from special relativity

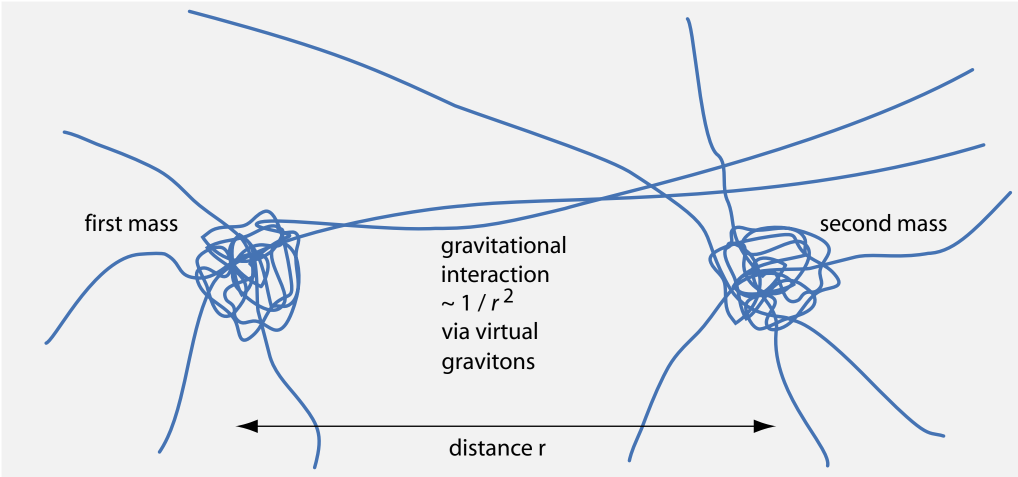

FIGURE 75 Gravitational attraction as result of twisted tail pairs – or twisted tether pairs.

Ref. 85 at high particle energy. However, this suggestion is based on two assumptions: that at Planck energy *point masses* are a viable approximation to particles, and that at Planck energy *vacuum and matter differ*. In the strand model, and in nature, both assumptions are incorrect. Nature, as general relativity shows, does not allow the existence of point masses: the densest objects in nature are black holes, and these are not point-like for any mass value.

In addition, quantum theory implies the fuzziness of matter and space. As a result, Page 65 near Planck energy, matter and vacuum *cannot* be distinguished. Put simply, no system near Planck energy can be described without general relativity and without quantum Ref. 188 gravity. In short, the strand model predicts that the approach of 'doubly special relativ- Page 8 ity' cannot be correct. Also Figure 1 makes this point: there is no description of nature besides the usual ones.

To sum up, the strand model reproduces special relativity when masses are approximated as point-like in flat space. But at the same time, the strand model states that a negligibly small, light and localizable mass cannot exist – neither in flat nor in curved space. This matches observations.

Classical gravitation

In nature, at low speeds and in the flat space limit, gravitation is observed to lead to an acceleration a of test masses that changes as the inverse square distance from the gravitating mass;

$$a = G\frac{M}{R^2} . \tag{169}$$

This acceleration is called *universal gravitation* or *classical gravitation*. It is an excellent approximation for the solar system and for many star systems throughout the universe.

In the strand model, every space-time effect, including gravitation, is due to the behaviour of tangle tails. In the strand model, every mass, i.e., every system of tangles, is connected to the border of space by tails. The nearer a mass is to a second mass, the more frequently the twisted tails of one mass affect the other mass. Figure 75 illustrates the situation. The strand model states:

▷ ***Gravitation*** is due to the fluctuations of tail crossings: gravity is due to twisted tail pairs.

Around a mass, the tail twists fluctuate; averaged of time, the fluctuations lead to a crossing switch density around every mass. The resulting potential energy – where energy is action per time and thus given by the number of crossing switches per time – changes like
Challenge 179 e the inverse distance from the central mass. This is the reason for the $1/r$-dependence of the gravitational potential and the $1/r^2$-dependence of gravitational acceleration. (This applies to all those cases where spatial curvature is negligible.) In simple words, in the strand model, the inverse square dependence of gravitational acceleration is due to the three-dimensionality of space combined with the one-dimensionality of strands.

The strand model also shows that masses and energies are always positive: every tangle contains curved strands. The model also shows qualitatively that larger masses produce stronger attraction: larger masses contain more particles and thus produce more
Page 358 crossing switches. We will show below that the number density of crossing switches for each particle is indeed determined by the mass.

In the strand model, twisted tail pairs are (virtual) gravitons. Strands thus reproduce the idea that gravity is due to the exchange of (virtual) gravitons. Indeed, the strand model of the graviton, illustrated below in Figure 79, provides a consistent model that fulfils all requirements: it has the correct spin and quantum numbers, it fits with the idea of curvature defect, and it couples to masses producing universal gravity.

In the strand model, crossing switches are not only related to action and energy; they are also related to entropy. A slightly different – but equivalent – view on gravitation therefore appears when we put the stress on the entropic aspect.

Deducing universal gravitation from black hole properties

Black holes have entropy; this implies universal gravitation. There are at least two ways to explain this connection.

Ref. 189 An especially concise explanation was recently given by Erik Verlinde. In this view, *gravity appears because any mass M generates an effective vacuum temperature around it.* A gravitating mass M attracts test masses because during the *fall* of a test mass, the total entropy *decreases*. It is not hard to describe these ideas quantitatively.

Given a spherical surface A enclosing a gravitating mass M at its centre, the acceleration a of a test mass located somewhere on the surface is given by the local vacuum temperature T:

$$a = T\,\frac{2\pi\, kc}{\hbar}\ , \qquad (170)$$

where k is the Boltzmann constant. This relation is called the *Fulling–Davies–Unruh effect* and relates vacuum temperature and local acceleration. Thus, an inertial or a freely falling mass (or observer) measures a vanishing vacuum temperature.

In the strand model, the vacuum temperature at the surface of the enclosing sphere is given by the crossing switches induced by the tails starting at the mass. We can determine the vacuum temperature by dividing the energy E contained inside the sphere by *twice* the *maximum* possible entropy S for that sphere. This maximum value is the entropy that the sphere would have if it were a black hole horizon; it can be calculated by the strand Page 289 model, as we will see shortly.

The temperature T is thus given by the expression

$$T = \frac{E}{2S} = \frac{M}{A}\frac{2G\hbar}{kc} . \tag{171}$$

The factor 2 needs explanation; it derives from the dependence of entropy on the square of the radius. We do not discuss the details here.

Neglecting spatial curvature, we can set $A = 4\pi R^2$; this gives a temperature at the enclosing sphere given by

$$T = \frac{M}{R^2}\frac{G\hbar}{2\pi kc} . \tag{172}$$

Page 283 Inserting this expression into the expression (170) for the Fulling–Davies–Unruh acceleration a, we get

$$a = G\frac{M}{R^2} . \tag{173}$$

This is universal gravitation, as discovered by Robert Hooke and popularized by Isaac Newton. Since spatial curvature was neglected, and the central mass was assumed at rest, this expression is only valid for large distances and small speeds. We have thus deduced universal gravity from the effects of gravitating masses on vacuum temperature. Below, we show that in the relativistic case this sequence of arguments – which was given by Page 295 Jacobson fifteen years before Verlinde – leads to the field equations of general relativity.

An alternative deduction of universal gravitation from black hole entropy is the following. The gravitational force F on a test mass m is given by the vacuum temperature T created by the central mass M and by the change of entropy S per length that is induced by the motion of the test mass:

$$F = T\frac{\mathrm{d}S}{\mathrm{d}x} . \tag{174}$$

The change of entropy $\mathrm{d}S/\mathrm{d}x$ when a test mass m moves by a distance x can be determined from the strand model in a simple manner. When the test mass m moves by a (reduced) Compton wavelength, in the strand model, the mass has rotated by a full turn: the entropy change is thus $2\pi k$ per (reduced) Compton wavelength. Thus we have

$$\frac{\mathrm{d}S}{\mathrm{d}x} = m\frac{2\pi kc}{\hbar} . \tag{175}$$

Using the temperature T found in expression (172), we get an expression for the gravitational force given by

$$F = G\frac{Mm}{R^2} . \tag{176}$$

This is universal gravitation again.

We have thus deduced universal gravitation from the entropy and the vacuum temperature generated by gravitating masses. We note that the temperature and entropy of black holes are limit values. We can thus state that universal gravitation is a consequence of nature's limit values.

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Summary on universal gravitation from strands

Universal gravitation is due to the temperature and entropy of the (curved) vacuum around masses. The limit case is the temperature and entropy of black holes. In the strand model, these temperature and entropy values are a consequence of the underlying strand crossing switches; we will show this shortly.

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More precisely, gravitation is due to twisted tail pairs. In the strand model, universal gravitation thus appears as an effect of the crossing switches induced by masses. In fact, we have several explanations of universal $1/r^2$ gravitation using strands. We have deduced universal gravitation from the energy of strands, from the temperature of strands and from the entropy of strands around a mass. We have also have deduced universal gravitation from the maximum force, which strands fulfil as well. In short, strands explain the origin of universal gravitation in several consistent ways.

Incidentally, modelling mass as a source for strand crossing switches – twisted strand pairs – is remotely reminiscent of Georges-Louis Lesage's eighteenth-century model of gravitation. Lesage proposed that gravity appears because many tiny, usually unnoticed corpuscules push masses together. In fact, as we will see shortly, there is a certain similarity between these assumed tiny corpuscules and virtual gravitons. And interestingly, all criticisms of Lesage's model then cease to hold. First, there is no deceleration of free masses in inertial motion, thanks to the built-in special-relativistic invariance. Secondly, there is no heating of masses, because the entangled tails represent virtual gravitons that scatter elastically. Thirdly, and most of all, by replacing the *corpuscules ultra-mondains* of Lesage by virtual gravitons – and finally by strands – we can predict an additional effect of gravity that is not described by the inverse square dependence: space-time curvature.

Ref. 190

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Curved space

In nature, observation shows that physical space is not flat around masses, i.e., in the presence of gravity. Near mass and energy, physical space is *curved*. Observations also show that curved space-time remains 3+1-dimensional. The observation of this type of curvature was predicted long before it was measured, because curvature follows unambiguously when the observer-invariance of the speed of light c and the observer-invariance of the gravitational constant G are combined.

We continue directly with the strand model of spatial curvature and show that all observations are reproduced.

▷ ***Curvature*** (of physical space-time) is due to simple, unknotted and weakly

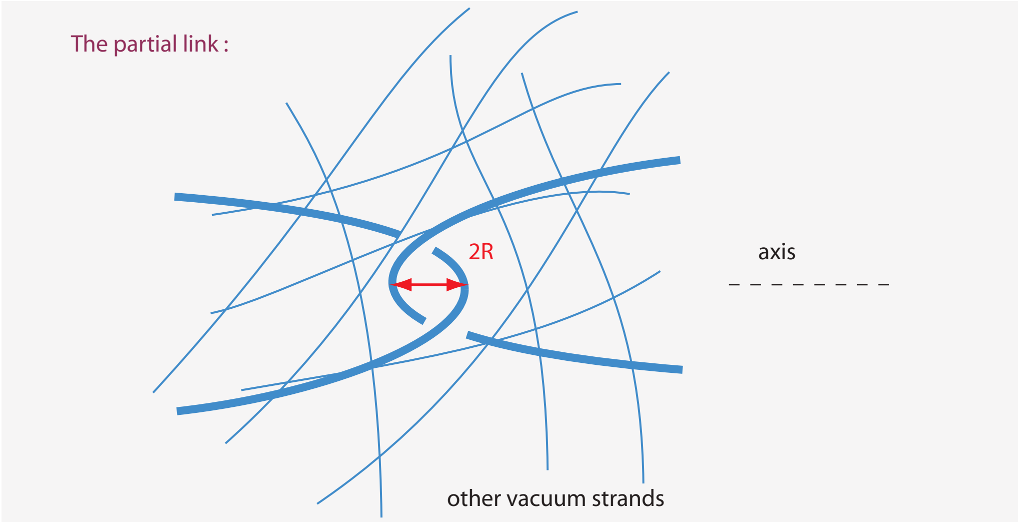

FIGURE 76 A schematic model of the fundamental defect, and thus the fundamental type of curvature: the *partial link*.

localized defects in the tangle of strands that make up the vacuum. An example is shown in Figure 76.

- ▷ In the case of curvature, *physical* space-time, which is due to averaged strand crossing switches, *differs* from flat *background* space-time, which usually corresponds to the tangent or to the asymptotic space-time. In Figure 76, the grey background colour can be taken as visualization of the background space.
- ▷ ***Mass*** is a localized defect in space and is due to tangled strands. Thus mass curves space around it.
- ▷ ***Energy*** in a volume is the number of crossing switches per unit time. As a result, mass is equivalent to energy. As a second result, energy also curves space.
- ▷ ***Gravitation*** is the space-time curvature originating from compact regions with mass or energy.

These natural definitions show that curvature is due to strand configurations. In particular, curvature is built of unknotted – i.e., massless – *defects*. The massless defects leading to curvature are usually dynamic: they evolve and change. Such curvature defects – virtual gravitons – originate at regions containing matter or energy. In fact, the curvature of
Page 358
space around masses is a natural result of fluctuations of the strands that make up matter tangles.

We note that curved space, being a time average, is *continuous* and *unique*. Vacuum or curved space, more precisely, curved physical space, thus differs from background space, which is flat (and drawn in grey in the figures).

Incidentally, the distinction between physical and background space also avoids Einstein's hole argument; in fact, the distinction allows discussing it clearly, as only physical

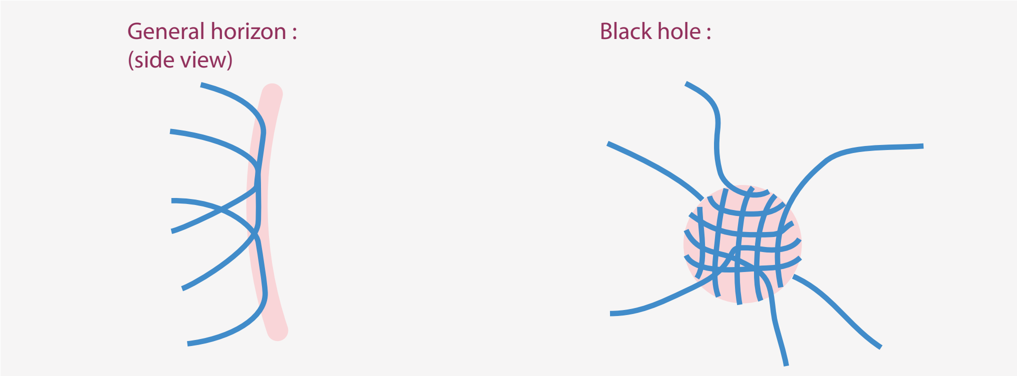

FIGURE 77 A schematic model of a general and a spherical horizon as tight weaves, as pictured by a distant observer. In the strand model there is *nothing*, no strands and thus no space, behind a horizon.

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space describes nature.

The structure of horizons and black holes

In general relativity, another concept plays a fundamental role. In the strand model we have:

▷ A ***horizon*** is a tight, one-sided weave of strands.

Therefore, there are no strands behind the horizon. This implies that behind a horizon, there is no matter, no light, no space and no time – just *nothing*. Indeed, this is the experience of any observer about a horizon. A horizon is thus a structure that limits physical space. It does *not* limit background space.

One particular type of horizon is well-known:

▷ A ***black hole*** is a tight, one-sided and *closed* weave of strands.

In principle, closed horizons can have any shape. The simplest case is the spherical, non-rotating horizon, which defines the *Schwarzschild black hole*. It is illustrated on the right-hand side of Figure 77.

If an observer is located outside a spherical horizon, the strand model states that there is nothing *inside* the horizon: no matter, no light and no vacuum. The strand model thus provides a simple and drastic view of black hole horizons. Figure 77 also illustrates that the concept of radius (or size) of a black hole has to be approached with the (well-known) care. In general, the size of a structure made of strands is the number of crossings encountered when travelling through it. However, an observer cannot travel *through* a black hole: there are no strands inside, thus there is no vacuum there! The size of a black hole must therefore be defined indirectly. The simplest way is to take the square root of the area, divided by 4π, as the radius. Thus the strand model, like general relativity, requires that the size of a compact horizon be defined by travelling *around* it.

We note that the strand model also provides an intuitive explanation for the differences between a rotating and a non-rotating black hole.

Is there something behind a horizon?

A drawing of a horizon weave, such as the one of Figure 77, clearly points out the difference between the background space and the physical space. The *background space* is the space we need for thinking, and is the space in which the drawing is set. The *physical space* is the one that appears as a consequence of the averaging of the strand crossings. Physical, curved space exists only on the observer side – usually outside – of the horizon. The physical space around a black hole is curved; it agrees with the background space only at infinite distance from the horizon. Inside the horizon, there is background space, but no physical space. In short, the strand model implies that – for an observer at spatial infinity – there is *nothing*, not even a singularity, inside a black hole horizon.

▷ There is no physical space, no matter and no singularity inside a horizon.

Horizons are observer-dependent. Both the existence and the shape of a horizon depends on the observer. As we will see, this happens in precisely the same way as in usual general relativity. In the strand model, there is no contradiction between the one observer at spatial infinity who says that there is *nothing* behind a horizon, not even physical space, and another, falling observer, who does not observe a horizon and thus states that there is *something* there. In the strand model, the two statements naturally transform into each other under change of viewpoint. Indeed, the transformation between the two viewpoints contains a deformation of the involved strands.

We note that the equivalence of viewpoints and the statement that there is nothing behind a horizon is based on the combination of general relativity and quantum theory. If we would continue thinking that space and time is a manifold of points – thus disregarding quantum theory – these statements would *not* follow.

In summary, one-sided tight weaves are a *natural* definition of horizons.

Energy of black hole horizons

The strand model allows us to calculate the energy content of a closed horizon. Energy is action per unit time. In the strand model, the energy of a non-rotating spherical horizon is thus given by the number N_{cs} of crossing switches per time unit. In a tight weave, crossing switches cannot happen in parallel, but have to happen sequentially. As a result, a crossing switch 'propagates' to the neighbouring Planck area on the surface. Since the horizon weave is tight and the propagation speed is one crossing per crossing switch time, this happens at the speed of light. In the time T that light takes to circumnavigate the spherical horizon, all crossings switch. We thus have:

Challenge 180 e

$$E = \frac{N_{cs}}{T} = \frac{4\pi R^2}{2\pi R}\,\frac{c^4}{4G} = R\,\frac{c^4}{2G}\,. \tag{177}$$

Strands thus imply the well-known relation between energy (or mass) and radius of Schwarzschild black holes.

How do the crossing switches occur at a horizon of a black hole? This interesting

Challenge 181 e puzzle is left to the reader.

The tight-weave model of horizons also illustrates and confirms both the *hoop conjecture* and the *Penrose conjecture*. For a given mass, because of the minimum size of crossings, a spherical horizon has the smallest possible diameter, compared to other possible shapes. The strand model naturally implies that, for a given mass, spherical black holes indeed are the densest objects in nature.

The nature of black holes

The strand model naturally implies the *no-hair theorem*. Since all strands are the same, independently of the type of matter that formed or fell into the horizon, a black hole has no characteristics other than mass, angular momentum and charge. Here we used a result from the next chapter, when it will become clear that all elementary particles are indeed made of the same featureless strands. Taking that result as given, we deduce that flavour quantum numbers and particle number do not make sense for black holes. We also deduce that weak and strong charge are not defined for black holes. Strands explain naturally why neutral black holes made of antimatter and neutral black holes made of matter do not differ, if their masses and angular momenta are the same. In short, the strand model of nature implies the no-hair theorem: *strands, not hairs*.

Horizons and black holes are borderline systems between space and matter. This borderline property must be fulfilled by every final theory. The strand model fulfils this requirement: in the strand model, black holes can either be described as curved space or as tightly packed particles in permanent free fall.

Entropy of vacuum and matter

Both vacuum and matter are made of fluctuating strands. We note directly:

▷ The flat and infinite vacuum has *vanishing* entropy, because the number of crossing switches is zero on average.

At the same time,

▷ Curved space and horizons have *non-vanishing* entropy.

The entropy of vacuum and of horizons differs from that of matter. In the *absence of gravity*, the number of microstates of matter is determined – as in usual thermodynamics (thermostatics) – by the behaviour of tangle *cores*.

In *strong gravity*, when the distinction between matter and vacuum is not so clear-cut, the number of microstates is determined by the possible crossing switches of the strands. In strong gravity, only *tails* play a role. This becomes clear when we calculate the entropy of black holes.

Entropy of black holes deduced from the strand model

Despite the tight weaving, the strands making up a horizon are fluctuating and moving: the weave shape fluctuates and crossing switch all the time. This fluctuating motion is

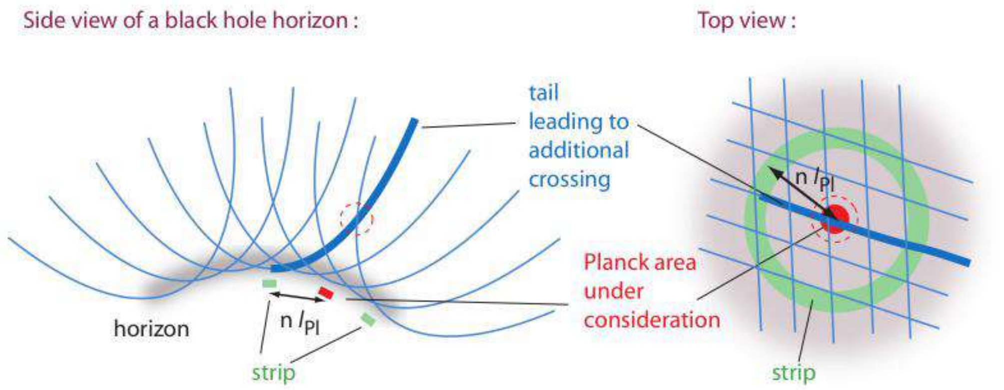

FIGURE 78 The entropy of black holes results from the number of possible crossing states above a Planck area.

the reason why horizons – in particular those of black holes – have entropy.

The weave model of a horizon, illustrated in detail in Figure 78, allows us to calculate the corresponding entropy. Since the horizon is a tight weave, there is a crossing on each Planck area. To a first approximation, on each (corrected) Planck area of the horizon, the strands can cross in *two* different ways. The fundamental principle of the strand model thus yields two microstates per Planck area. The number N of Planck areas is given by $N^2 = Ac^3/4G\hbar$. The resulting number of black hole microstates is 2^{N^2}. The entropy is given by the natural logarithm of the number of the possible microstates times k. This approximation gives an entropy of a horizon of

$$S = A\,\frac{kc^3}{4G\hbar}\ln 2\,. \qquad (178)$$

This result is the well-known first approximation of black hole entropy: one bit per corrected Planck area. In the strand model, the proportionality of entropy and area is thus a direct consequence of the *extension* of the strands. This proportionality is also well known from studies of quantum gravity and of strings. In those approaches however, the relation between the area proportionality and extension is less obvious.

Ref. 191

For Schwarzschild black holes, the entropy value of expression (178) is *not* correct. In the strand model, this incorrect value is explained as a consequence of neglecting the effects of the strand *tails*. Indeed, additional contributions to the entropy appear at a *finite distance* from the horizon, due to the crossing of the tails on their way to the border of space, as shown in Figure 78. The actual entropy will thus be larger than the first approximation, but still be proportional to the area A.

The correct proportionality factor between the area and the entropy of a black hole results when the strand tails are taken into account. (The correction factor is called the *Barbero–Immirzi parameter* in the research literature on quantum gravity.) The calculation is simplest for Schwarzschild black holes. By construction, a black hole with macroscopic radius R, being a tight weave, has R/l_{Pl} tails. For each given Planck area, there are, apart from the basic, or lowest crossing, additional crossings 'above it', along the radial direction, as shown in Figure 78. These additional crossings are due to the tails from neighbouring and distant Planck areas.

Taking into effect all strand tails allows us to calculate the average number of crossings *above* a given Planck area. The main point is to perform this calculation for all those tails that start in a circular strip of Planck width centred around the Planck area under consideration. We then add the probabilities for all possible circular strips. One such circular strip is drawn in Figure 78.

The definition of horizons as tight weaves implies that a horizon with N^2 Planck areas is made of N strands. This means that for each circular strip of radius nl_{Pl}, there is only *one* strand that starts there and reaches spatial infinity as a tail. For this tail, the average probability p that it crosses above the central Planck area under consideration is

Challenge 182 e

$$p = \frac{1}{n!} \,. \tag{179}$$

Summing over all strips, i.e., over all values n, we get a total of $\sum_{n=0}^{\infty} 1/n! = \mathrm{e} = 2.71828...$ microstates *on* and *above* the central Planck area under consideration. Thus the number e replaces the number 2 of the first approximation: the number of horizon microstates of a Schwarzschild black hole is not 2^{N^2}, but e^{N^2}. As a consequence, the entropy of a macroscopic Schwarzschild horizon becomes

$$S = A \,\frac{kc^3}{4G\hbar} \,. \tag{180}$$

This is the Bekenstein–Hawking expression for the entropy of Schwarzschild black holes. The strand model thus reproduces this well-known result. With this explanation of the difference between 2 and e = 2.71828..., the strand model confirms an old idea:

▷ The entropy of a black hole is located *at and near* the horizon.

The above calculation, however, counts some states more than once. Topologically identical spherical horizons can differ in the direction of their north pole and in their state of rotation around the north–south axis. If a spherical horizon is made of N strands, it has N^2 possible physical orientations for the north pole and N possible angular orientations around the north–south axis. The actual number of microstates is thus e^{N^2}/N^3. Using the relation between N^2 and the surface area A, namely $A = N^2 4G\hbar/c^3$, we get the final result

$$S = A \,\frac{kc^3}{4G\hbar} - \frac{3k}{2} \ln \frac{A\,c^3}{4G\hbar} \,. \tag{181}$$

The strand model thus makes a specific prediction for the logarithmic correction of the entropy of a Schwarzschild black hole. This final prediction of the strand model agrees with many (but not all) calculations using superstrings or other quantum gravity approaches. Ref. 192

In summary, the entropy value (180), respectively (181), of black holes is due to the *extension* of the fundamental entities in the strand model and to the *three dimensions* of space. If either of these properties were not fulfilled, the entropy of black holes would not result. This is not a surprise; also our deduction of quantum theory was based on the same two properties. In short: like every quantum effect, also the entropy of black holes is a result of extension and three-dimensionality. Only a three-dimensional description of nature agrees with observation.

Temperature, radiation and evaporation of black holes

The strands that make up a horizon fluctuate in shape. Since every horizon contains energy, the shape fluctuations imply energy fluctuations. In other words, horizons are predicted to have a *temperature*. The value of the temperature can be deduced from the strand model by noting that the characteristic size of the fluctuations for a spherical horizon is the radius R of the horizon. Therefore we have

$$kT = \frac{\hbar c}{2\pi R} \ . \qquad (182)$$

Using the definition of *surface gravity* as $a = c^2/R$, we get

$$T = \frac{\hbar a}{2\pi k c} \ . \qquad (183)$$

The strand model predicts that horizons have a temperature proportional to their surface gravity. This result has been known since 1973. Ref. 57, Ref. 58

All hot bodies radiate. The strand model thus predicts that Schwarzschild black holes *radiate* thermal radiation of the horizon temperature, with power and wavelength

$$P = 2\pi\hbar c^2/R^2 \quad , \quad \lambda \approx R \ . \qquad (184)$$

This confirms a well-known consequence of the temperature of black holes.

Like all thermal systems, horizons follow thermodynamics. In the strand model, black hole radiation and evaporation occur by reduction of the number of strands that make up the horizon. The strand model thus predicts that black holes *evaporate completely*, until only elementary particles are left over. In particular, the strand model implies that in black hole radiation, there is *no* information loss.

In short, strands reproduce all aspects of black hole evaporation. The strand model also shows that there is no information loss in this process.

Black hole limits

In many ways, black holes are *extreme* physical systems. Not only are black holes the limit systems of general relativity; black holes also realize various other limits. As such, black holes resemble light, which realizes the speed limit. We now explore some of these limits.

For a general physical system, not necessarily bound by a horizon, the definitions of energy and entropy with strands allow some interesting conclusions. The entropy of a system is the result of the number of crossing possibilities. The energy of a system is the number of crossing changes per unit time. A large entropy is thus only possible if a system shows many crossing changes per time. Since the typical system time is given by the circumference of the system, the entropy of a physical system is therefore limited:

$$S \leqslant ER\, 2\pi k/\hbar c\, . \tag{185}$$

This relation is known as *Bekenstein's entropy bound*; the precise definitions of the quantities in the bound need some care, as Don Page explains. The bound thus also follows from the strand model. Strands imply that the equality is realized only for black holes.

Ref. 193

In the strand model, horizons are tight, one-sided weaves. For example, this implies that any tangle that encounters a horizon is essentially flat. Because of tangle flatness and the extension of the tails, at most one Planck mass can cross a horizon during a Planck time. This yields the mass rate limit

$$\mathrm{d}m/\mathrm{d}t \leqslant c^3/4G \tag{186}$$

that is valid in general relativity and in nature.

Black holes can rotate. The strand model states that there is a highest angular frequency possible; it appears when the equator of the black hole rotates with the speed of light. As a result, the angular momentum J of a black hole is limited by

$$J < 2GM^2/c\, . \tag{187}$$

Ref. 194

This limit is well known from general relativity.

The electric charge of a black hole is also limited. The force limit in nature implies that the electrical forces between two charged black holes must be lower than their gravitational interaction. This means that

$$\frac{Q^2}{4\pi\varepsilon_0 r^2} \leqslant \frac{GM^2}{r^2}\, , \tag{188}$$

or

$$Q^2 \leqslant 4\pi\varepsilon_0 GM^2\, . \tag{189}$$

This is the well-known charge limit for (static) black holes given by the Reissner–Nordström metric. The maximum charge of a black hole is proportional to its radius. It follows directly from the maximum force principle.

To explain the charge limit, we deduce that the *extremal* charge surface density Q/A of a black hole is proportional to $1/R$. The higher the horizon curvature, the more charge per Planck area is possible. In the strand model, a horizon is a tight weave of strands. We are thus led to conjecture that at Planck scale, electric charge is related to and limited by strand curvature. We will explore this connection in more detail below.

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The strand model limits energy density to the Planck energy per Planck volume, or to the value $c^7/(16G^2\hbar)$. This limit implies a lower size limit for black holes, particles and any localized system. Therefore, the strand model does not allow singularities, be they dressed or naked. And indeed, no singularity has ever been observed.

In summary, the strand model reproduces the known limit properties of horizons. And all these results are independent of the precise fluctuation details of the strands.

Curvature around black holes

The tails of a black hole extend up to the border of space; the density of tails is highest at the horizon. A black hole is therefore surrounded by partial links at any *finite* distance from the horizon. In other words, the space around a black hole is *curved*. The value of the space-time curvature increases as one approaches the horizon, because of the way in which the partial links hinder each other in their motion. The nearer they are to the horizon, the more they hinder each other. The curvature that appears is proportional to the density of partial links and to their average strand curvature.

At the horizon, the curvature radius is the horizon radius R. By construction, the number of tails departing from a non-rotating black hole is proportional to R. The spatial curvature is given by the average crossing density gradient. Hence at a radial distance r from a static black hole, the spatial curvature K is

$$K \sim \frac{R}{r^3} . \tag{190}$$

So at the horizon itself, the curvature K is (of the order of) the inverse square of the horizon radius; further away, it decreases rapidly, with the third power of the distance. This result is a well-known property of the Schwarzschild solution and is due to the extension of the strands. The rapid decay with radius is the reason why in everyday situations there is no noticeable curvature of space-time. In short, strands allow us to deduce the correct curvature of space-time around black holes and spherical masses.

The shape of non-rotating black holes

The strand model also explains and visualizes the importance of spherical horizons in nature. First of all, strands illustrate the non-existence of (uncharged) one-dimensional or toroidal horizons in 3 + 1 space-time dimensions. Such configurations are unstable, in particular against transverse shear and rearrangement of the strands.

The strand model also implies that non-rotating, closed horizons are spherical. Obviously, spheres are the bodies with the smallest surface for a given volume. The minimum horizon surface appears because the strands, through their fluctuations, effectively 'pull' on each Planck area of the horizon. As a result, all non-rotating macroscopic horizons will evolve to the spherical situation in a few Planck times. (Deviations from the spherical shape will mainly occur near Planck scales.) With the definition of gravity waves given below, it also becomes clear that strongly deformed, macroscopic and non-spherical horizons are unstable against emission of gravity waves or of other particles. In short,

▷ All non-rotating horizons of non-spherical shape are unstable.

The strand model thus confirms that spherical horizons are favoured and that the most compact bodies with a given mass. The reasoning can be extended to rotating horizons, yielding the well-known shapes.

In summary, strands reproduce all known qualitative and quantitative properties of horizons and of black holes, and thus of general systems with strong gravitational fields. All predictions from strands agree with observations and with other approaches to quantum gravity. These hints already suggest that strands imply the field equations.

The field equations of general relativity

The field equations can be deduced from the fundamental principle in two different, but related ways. Essentially, both derivations repeat the reasoning for universal gravitation Page 282 given above, but for the relativistic case. The first deduction of the field equations is based Ref. 22 on an old argument on the thermodynamics of space-time. Strands show that horizons have three thermodynamic properties:

— an area–entropy relation of $S = A\, kc^3/4G\hbar$,

— a curvature–temperature relation of $T = a\,\hbar/2\pi kc$,

— a relation between heat and entropy of $\delta Q = T\delta S$.

Using these three properties, and using the relation

$$\delta Q = \delta E\ , \tag{191}$$

that is valid *only* in case of horizons, we get the first principle of horizon mechanics

$$\delta E = \frac{c^2}{8\pi G} a\, \delta A\ . \tag{192}$$

From this relation, using the Raychaudhuri equation, we obtain the field equations of Page 33 general relativity. This deduction was given above.*

In other words, the field equations result from *the thermodynamics of strands*. It is worth noting that the result is independent of the details of the fluctuations or of the microscopic model of space, as long as the three thermodynamic properties just given

* Here is the argument in a few lines. The first principle of horizon mechanics can be rewritten, using the energy–momentum tensor T_{ab}, as

$$\int T_{ab} k^a \mathrm{d}\Sigma^b = \frac{c^2}{8\pi G} a\, \delta A$$

where $\mathrm{d}\Sigma^b$ is the general surface element and k is the Killing vector that generates the horizon. The Raychaudhuri equation allows us to rewrite the right-hand side as

$$\int T_{ab} k^a \mathrm{d}\Sigma^b = \frac{c^4}{8\pi G} \int R_{ab} k^a \mathrm{d}\Sigma^b$$

where R_{ab} is the Ricci tensor describing space-time curvature. This equality implies that

$$T_{ab} = \frac{c^4}{8\pi G} (R_{ab} - (R/2 + \Lambda) g_{ab})$$

where Λ is an undetermined constant of integration. These are Einstein's field equations of general relativity. The field equations are valid everywhere and for all times, because a suitable coordinate transformation can put a horizon at any point and at any time. To achieve this, just change to a suitable accelerating frame, as Vol. II, page 98 explained in the volume on relativity.

are valid. In fact, these properties must be fulfilled by any model of space-time; and indeed, several competing models of space claim to fulfil them.

We can use the relation between fluctuations and strands to settle an issue mentioned above, in the section on quantum theory. Strand fluctuations *must* obey the thermodynamic properties to allow us to define space-time. If they obey these properties, then space-time exists and curves according to general relativity.

Page 188

A second derivation of the field equations of general relativity follows the spirit of the strand model most closely. It is even shorter. Strands imply that all physical quantities are limited by the corresponding Planck limit. These limits are due to the limit to the fundamental principle, in other words, they are due to the packing limit of strands. In particular, the fundamental principle limits force by $F \leqslant c^4/4G$ and power by $P \leqslant c^5/4G$. We have already shown above that this limit implies the field equation.

Ref. 19
Page 33

In other words,

▷ Given that black holes and thus horizons are thermodynamic systems, so is curved space.

The reason: both can be transformed into each other. Therefore:

▷ Since black holes have thermodynamic aspects, so has gravity.

And since black holes are built from microscopic degrees of freedom, so is curved space. Or, in simple words:

▷ Space is made of many small entities.

And finally we can state:

▷ Space is made of strands, because strands are the simplest entities that yield black hole entropy.

Strands are the simplest way to incorporate quantum effects into gravitation. If we take into consideration that strands are the only way known so far to incorporate gauge interactions, we can even conclude that strands are the only way known so far to incorporate all quantum effects into gravitation.

In summary, the strand model asserts that the field equations appear as consequences of fluctuations of impenetrable, featureless strands. In particular, the strand model implies and confirms that a horizon and a particle gas at Planck energy do not differ. However, the value of the cosmological constant is *not* predicted from strand thermodynamics.

Equations from no equation

The strand model asserts that the field equations of general relativity are not the result of another, more basic evolution equation, but result directly from the fundamental principle. To say it bluntly, the field equations are deduced from a drawing – the funda-

Page 149 mental principle shown in Figure 10. This strong, almost unbelievable statement is due to a specific property of the field equations and to two properties of the strand model.

First of all, the field equations are, above all, consequences of the thermodynamics of space-time. In the strand model, the thermodynamic properties are deduced as a consequence of the strand fluctuations. This deduction does not require underlying evolution equations; the field equations follow from the statistical behaviour of strands.

The second, essential property of the strand model is its independence from the underlying motion of the strands. In the strand model we obtain the evolution equations of the vacuum – the field equations of general relativity – without deducing them from another equation. We do not need an evolution equation for the strand shape; the deduction of the field equations works for *any* underlying behaviour of strand shapes, as long as the thermodynamic properties of the strand fluctuations are reproduced.

The third and last essential property that allows us to deduce the field equations directly from a graph, and not from another equation, is the relation between the graph and natural physical units. The relation with natural units, in particular with the quantum of action $\hbar$ and the Boltzmann constant k, is fundamental for the success of the strand model.

In summary, the fundamental principle of the strand model contains all the essential properties necessary for deducing the field equations of general relativity. In fact, the discussion so far makes another important point: unique, underlying, more basic evolution equations for the tangle shape *cannot* exist. There are two reasons. First, an underlying equation would itself require a deduction, thus would not be a satisfying solution to unification. Secondly, and more importantly, evolution equations are differential equations; they assume well-behaved, smooth space-time. At Planck scales, this is impossible.

▷ Any principle that allows deducing the field equations cannot itself be an evolution equation.

The Hilbert action of general relativity

We have just shown that the strand model implies the field equations of general relativity. Page 211 We have also shown above that, in the strand model, the least action principle is a natural property of all motion of strands. Combining these two results, we find that a natural way to describe the motion of space-time is the (extended) *Hilbert action* given by

$$W = \frac{c^4}{16\pi G} \int (R - 2\Lambda)\, \mathrm{d}V \ , \qquad (193)$$

where R is the Ricci scalar, $\mathrm{d}V = \sqrt{\det g}\, \mathrm{d}^4x$ is the invariant 4-volume element of the metric g, and Λ is the cosmological constant, whose value we have not determined yet. As is well known, the description of evolution with the help of an action does not add anything to the field equations; both descriptions are equivalent.

For a curved three-dimensional space, the Ricci scalar R is the average amount, at a given point in space, by which the curvature deviates from the zero value of flat space. Page 286 In the strand model, this leads to a simple statement, already implied by Figure 76:

FIGURE 79 The graviton in the strand model: a twist in a pair of tethers.

▷ The *Ricci scalar* R is the ratio of additional or missing crossings per spatial volume, compared to flat space.

As usual, the averaging is performed over all spatial orientations. A similar statement can be made for the cosmological constant Λ. In short, we can say: the Hilbert action follows directly from the fundamental principle of the strand model.

Space-time foam

Quantum physics implies that at scales near the Planck length and the Planck time, space-time fluctuates heavily. John Wheeler called the situation *space-time foam*; the term *quantum foam* is also used. In a sense, *quantum gravity* can be defined, if at all, as the description of space-time foam. This reduced view arises because no separate theory of quantum gravity is possible in nature.

Historically, there have been many speculations on the details of space-time foam. Apart from its fluctuations, researchers speculated about the appearance of topology changes – such as microscopic wormholes – about the appearance of additional dimensions of space – between six and twenty-two – or about the appearance of other unusual properties – such as microscopic regions of negative energy, networks or loop structures.

The strand model makes a simple prediction that contradicts most previous speculations:

▷ Space-time foam is made of fluctuating strands.

At everyday scales, the foam is not noticed, because background space and physical space are indistinguishable. At Planck scales, space-time is not fundamentally different from everyday space-time. No unusual topology, no additional dimensions, and no new or unusual properties appear at Planck scales. Above all, the strand model predicts that there are *no* observable effects of space-time foam; for example, 'space-time noise' or 'particle diffusion' do not exist. The strand model of space-time foam is both simple and unspectacular.

Ref. 195

Gravitons, gravitational waves and their detection

In the strand model, gravitons can be seen as a special kind of partial links: it appears to be a twisted pair of tethers. An example is shown in Figure 79. As a twisted pair of parallel strands, the graviton returns to itself after rotation by π; it thus behaves like a spin-2 boson, as required.

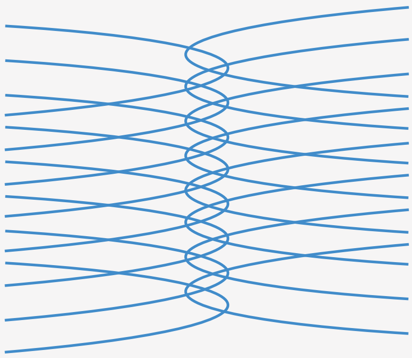

FIGURE 80 A speculative, highly schematic model for a cosmic string, a one-dimensional defect in space-time

Can single gravitons be observed? The strand model implies that the absorption of a single graviton by an elementary particle changes its spin or position. However, such a change *cannot* be distinguished from a quantum fluctuation, because the graviton is predicted to be massless. Furthermore, the strand model predicts that gravitons do not interact with photons, because they have no electric charge. In summary, the strand model predicts:

▷ Single gravitons *cannot* be detected.

The situation changes for gravitational waves. Such waves are coherent superpositions of large numbers of gravitons and are observable classically. In such a case, the argument against the detection of single gravitons does not apply. In short, the strand model predicts that gravitational waves *can* be observed. (This prediction, made by many since 1915 and repeated in this text on the basis of the strand model in 2008, came true in February 2016. The observations also produced the extremely low mass limit of at most $1.2 \times 10^{-22}\,\mathrm{eV}/c^2$ for any possible mass of the photon.)

Challenge 183 e

Ref. 196

Open challenge: Improve the argument for the graviton tangle

The argument that leads to the graviton tangle is rather hand-waving. Can you make the argument more compelling? Could the four tails form a cross and thus span a plane instead of a ribbon?

Challenge 184 ny

Other defects in vacuum

The strand model provides a quantum description of gravitation. The strand model does so by explaining physical space as the average of the crossing switches induced by strand fluctuations among untangled strands. Matter, radiation and horizons are defects in the 'sea' of untangled strands.

So far, we have been concerned with *particles*, i.e., localized, zero-dimensional defects, and with *horizons*, i.e., two-dimensional defects. Now, modelling of the vacuum as a set of untangled strands also suggests the possible existence of *one-dimensional* – equivalent to dislocations and disclinations in solids – of *additional* two-dimensional defects, or of *three-dimensional* defects. Such defects could model cosmic strings, domain walls, wormholes, toroidal black holes, time-like loops and regions of negative energy.

An example of such a possible new defect is illustrated in Figure 80. The illustration can be seen as the image of a one-dimensional defect or as the cross section of a two-dimensional defect. Are such defects stable against fluctuations? The strand model suggests that they are not. These defects are expected to decay into a mixture of gravitons, black holes, matter and radiation particles. However, this issue is still a topic of research, and will not be covered here.

Exploring the stability of wormholes, time-like loops and toroidal black holes leads to similar results. It seems that the strand model should not allow time-like loops of macroscopic size, since any configuration that cannot be embedded locally into three flat spatial dimensions is either a particle or a black hole. Alternatively, macroscopic time-like loops would collapse or decay because of the fluctuations of the strands. In the same way, wormholes or black holes with non-trivial topology should be unstable against more usual strand structures, such as particles or black holes.

We also note the strand model does not allow volume defects (black holes being surface-like defects). The most discussed types of volume defects are macroscopic regions of negative energy. Energy being action per unit time, and action being connected to crossing changes, the model does not allow the construction of negative-energy regions. However, the strand model does allow the construction of regions with lower energy than their environment, as in the Casimir effect, by placing restrictions on the wavelengths of photons.

The strand model thus predicts the absence of additional defects and tangle types. The final and general connection between tangle types and defects is shown (again) in Table 11. The next chapter will give details of the tangles corresponding to each particle.

In summary, the strand model reproduces the results of modern quantum gravity and predicts that the more spectacular defects conjectured in the past – linear defects such as cosmic strings, surface defects such as wormholes, volume defects such as negative-energy regions – *do not appear* in nature.

The gravity of superpositions

What is the gravitational field of a quantum system in a macroscopic superposition? The issue has been raised by many scholars as an important step towards the understanding of how to combine gravitation and quantum theory.

The strand model deflates the importance of the issue. The model shows – or predicts, if one prefers – that the gravitational field of a superposition is the temporal and spatial
Vol. IV, page 143
average of the evolving quantum system, possibly under inclusion of decoherence.

What is the gravitational field of a single quantum particle in a double-slit experi-
Page 204
ment? As Figure 39 shows, the gravitational field almost always appears in both slits, and only very rarely in just one slit.

In summary, in the strand model, the combination of gravitation and quantum theory is much simpler than was expected by most researchers. For many decades it was suggested that the combination was an almost unattainable goal. In fact, in the strand model we can almost say that the two descriptions combine naturally.

TABLE 11 Correspondences between physical systems and mathematical tangles.

Physical system	Strands	Tangle type
Vacuum	many infinite unknotted strands	unlinked
Dark energy	many fluctuating infinite strands	unlinked
Elementary vector boson	one infinite strand	a curve
Quark	two infinite strands	rational tangle
Lepton	three infinite strands	braided tangle
Meson, baryon	three or more infinite strands	composed of rational tangles
Higher-order propagating fermion	two or more infinite strands	general rational tangle
Virtual particles	open or unlinked strands	trivial tangles
Composed systems	many strands	separable tangles
Graviton	two infinite twisted strands	specific rational tangle
Gravity wave	many infinite twisted strands	many graviton tangles
Horizon	many tightly woven infinite strands	web-like rational tangle
Young universe	closed strand(s)	knot (link)

Torsion, curiosities and challenges about quantum gravity

Ref. 197 On the one hand, the strand model denies the existence of any specific effects of *torsion* on gravitation. On the other hand, the strand model of matter describes spin with the belt trick. The belt trick is thus the strand phenomenon that is closest to the idea of torsion. Therefore, exaggerating a bit in the other direction, it could also be argued that in the strand model, torsion effects are quantum field theory effects.

* *

The strand model describes three-dimensional space as made of tangled strands. Several similar models have been proposed in the past.

The model of space as a *nematic world crystal* stands out as the most similar. This model was proposed by Hagen Kleinert in the 1980s. Ref. 198 He took his inspiration from the famous analogy by Ekkehart Kröner between the equations of solid-state elasticity around line defects and the equations of general relativity. Ref. 199

Ref. 157 Also in the 1980s, the mentioned posets have been proposed as the fundamental structure of space. Ref. 200 Various models of quantum gravity from the 1990s, inspired by spin networks, spin foams and by similar systems, describe empty space as made of extended constituents. These extended constituents tangle, or bifurcate, or are connected, or sometimes all of this at the same time. Depending on the model, the constituents are lines, circles or ribbons. In some models their shapes fluctuate, in others they don't.

Ref. 158 Around the year 2000, another type of Planck-scale crystal model of the vacuum has been proposed by David Finkelstein. In 2008, a specific model of space, a crystal-like network of connected bifurcating lines, has been proposed by Gerard 't Hooft. Ref. 201

All these models describe space as made of some kind of extended constituents in a three-dimensional background. All these models derive general relativity from these constituents by some averaging procedure. The lesson is clear: it is *not* difficult to derive general relativity from a Planck-scale model of space. It is *not* difficult to unify gravity and quantum theory. As Luca Bombelli said already in the early 1990s, the challenge for a Planck-scale model of nature is not to derive gravity or general relativity; the challenge is to derive the other interactions. So far, the strand model seems to be the only model that has provided such a derivation.

* *

The Planck force is the force value necessary to produce a change $\hbar$ in a Planck time over a Planck length. The Planck force thus appears almost exclusively at horizons. Challenge 185 e

* *

Already in the 1990s, Leonard Susskind speculated that black holes could be formed by a single wound-up string. Ref. 202 Strands differ from strings; they differ in the number of dimensions, in their intrinsic properties, in their symmetry properties, in the fields they carry and in the ways they generate entropy. Nevertheless, the similarity with the strand model of black holes is intriguing.

* *

In September 2010, two years after the strand model appeared, independent research confirmed its description of physical space, as already mentioned above. Page 163 In an extended article exploring the small scale structure of space from several different research perspectives in general relativity, Steven Carlip comes to the conclusion that all these perspectives suggest the common idea that 'space at a fixed time is thus threaded by rapidly fluctuating lines'. Ref. 155

Ref. 203 In 2011, also independently, Marcelo Botta Cantcheff modelled space as a statistic ensemble of one-dimensional 'strings'. He explained the main properties of space, including the thermodynamic properties of black holes.

* *

The first version of the strand model assumed that space is not defined at the cosmic horizon, and that therefore, strand impenetrability does not hold there. The same was thought to occur at black hole horizons. The newest version of the strand model does not seem to need this exception to impenetrability. Can you explain black hole entropy without it? Challenge 186 e

* *

Page 36 The strand model also allows us to answer the question whether quantum particles are black holes: no, they are not. Quantum particles are tangles, like black holes are, but particles do *not* have horizons. As a side result, the mass of all particles is lower than a Planck mass, or more precisely, lower than a Planck mass black hole.

Strands imply that gravity is weaker than the three gauge interactions. This consequence, like the low particle mass just mentioned, is due to the different origins of gravity and gauge interactions. Gravity is due to the strand tails, whereas gauge interactions are due to the tangle cores. Thus gravity is the weakest interaction in everyday

life. The observation of the weakness of gravity at everyday and other energy scales is
Ref. 204 sometimes called the *weak gravity conjecture*. It is naturally valid in the strand model.
Page 8 The conjecture is also part of the Bronshtein cube shown in Figure 1.

* *

For an observer at spatial infinity, a black hole horizon is an averaged-out tight web of strands. What does a falling observer experience? The question will still capture the imagination in many years. Such an observer will also see strands; above all, a falling observer will never hit any singularity. The details of the fall are so involved that they are not discussed here, because the fall affects both the black hole appearance and the observer.

* *

Can black hole radiation be seen as the result of trying to tear vacuum apart? Yes and no. The answer is no, because physical vacuum cannot be torn apart, due to the maximum force principle. But the answer is also yes in a certain sense, because the maximum force is the closest attempt to this idea that can be realized or imagined.

* *

The strand model makes the point that *entanglement* and the vacuum – and thus quantum gravity – have the same nature: both are due to crossing strands. This idea
Ref. 205 has been explored independently by Mark van Raamsdonk.

* *

As we have seen, the strand model predicts no *observable* violation of Lorentz-invariance – even though it predicts its violation at Planck scale. Strands predict the lack of dispersion, birefringence and opacity of the vacuum. Strands predict that the vacuum has three dimensions whenever it is observed and that it is unique, without phase transitions. We already mentioned the impossibility of detecting single gravitons.

All these negative predictions are examples of the *'no avail'* conjecture:

▷ Quantum gravity effects cannot be distinguished from ordinary quantum fluctuations.

Despite many attempts to disprove it, all experiments so far confirm the conjecture. Because both quantum gravity effects and quantum effects are due to tail fluctuations, the strand model seems to imply the conjecture.

* *

Ref. 206 The strand model of black holes also confirms a result by Zurek and Thorne from the 1980s: the entropy of a black hole is the logarithm of the number of ways in which it could have been made.

* *

Argue that because of the strand model, no black hole can have a mass below the (cor-
Challenge 187 s rected) Planck mass, about 11 μg, and thus that *microscopic black holes* do not exist. Can you find a higher lower limit for the mass?

* *

Do atoms or the elementary fermions moving inside matter emit gravitational radiation, and why? The question was already raised by Albert Einstein in 1916. The strand model answers the issue in the same way as textbook physics. Elementary particles in atoms – in the ground state – do not emit gravitational waves for the same reason that they do not emit electromagnetic waves: for atoms in the ground state, there is no lower state into which they could decay. Excited atomic states do not emit gravitational waves because of the extremely low emission probability; it is due to the extremely low mass quadrupole values.

* *

Ref. 207 In 2009 Mikhail Shaposhnikov and Christof Wetterich argued that if gravitation is 'asymptotically safe', there is no physics beyond the standard model and the Higgs mass must be around 126 GeV – exactly the value that was found experimentally a few years afterwards. A quantum field theory is called *asymptotically safe* if it has a fixed point at extremely high energies. Does the strand model imply that gravity is – maybe only effectively – asymptotically safe? Challenge 188 ny

* *

Page 59 It is often stated that general relativity does not allow the description of fermions if the topology of space is kept fixed. This is wrong: the strand model shows that fermions can be included in the case that space is seen as an average of extended fundamental entities.

* *

Following the fundamental principle of the strand model, G is the fundamental constant that describes gravitation. The strand model predicts that gravity is the same for all energy scales; in other words, the constant G is *not* expected to change with energy. This agrees with recent results from quantum gravity and distinguishes the behaviour of G from that of the coupling constants in the gauge interactions of particle physics. Ref. 208

Predictions of the strand model about gravity

As just presented, the strand model makes several verifiable predictions about general relativity and quantum gravity.

— The maximum energy speed in nature is c, at all energy scales, in all directions, at all times, at all positions, for every physical observer. This agrees with observations.

— No deviations from special relativity appear for any measurable energy scale, as long as gravity plays no role. No 'double' or 'deformed special relativity' holds in nature, even though a maximum energy-momentum for elementary particles does exist in nature. Whenever special relativity is not valid, general relativity, or quantum field theory, or both together need to be used. This agrees with observations.

— There is a maximum power or luminosity $c^5/4G$, a maximum force or momentum flow $c^4/4G$, and a maximum mass change rate $c^3/4G$ in nature. The limits hold for all energy scales, in all directions, at all times, at all positions, for every physical observer. These predictions agree with observations, though only few experimental observations are close to these limit values.

— There is a minimum distance and a minimum time interval in nature. There is a maximum curvature and a maximum mass density in nature. There are no singularities in nature. All this agrees with observations, including the newly discovered black hole mergers.
— The usual black hole entropy expression given by Bekenstein and Hawking holds. The value has never been measured, but is consistently found in all calculations performed so far. In fact, black hole entropy is related to the Fulling–Davies–Unruh effect, which itself is related to the Sokolov–Ternov effect. This latter effect has already been observed in several accelerators, for the first time in 1971. However, it now seems that this observation does not actually prove black hole entropy. Vol. V, page 147
— There are no deviations from general relativity, as described by the Hilbert action, for any measurable scale. The only deviations appear in situations with a few strands, i.e., in situations where quantum theory is necessary. This agrees with observations, including those of black hole mergers, but experimental data are far from sufficient; undetected deviations could still exist.
— There is no modified Newtonian dynamics, or MOND, with evolution equations that differ from general relativity. The rotation curves of stars in galaxies are due to dark matter, to other conventional explanations, or both. Ref. 209
— There is no effect of torsion that modifies general relativity. This agrees with observations.
— There is no effect of higher derivatives of the metric on the motion of bodies. This agrees with observations, but experimental data are far from sufficient.
— Observations are independent of the precise strand fluctuations. Mathematical consistency checks of this prediction are possible.
— No wormholes, no negative energy regions and no time-like loops exist. This agrees with observations, but experimental data are far from covering every possible loophole.
— The Penrose conjecture and the hoop conjecture hold. Here, a mathematical consistency check is possible.
— There are no cosmic strings and no domain walls. This agrees with observations, but experimental data are far from exhaustive.
— Gravitons have spin 2; they return to their original state after a rotation by π and are bosons. This agrees with expectations.
— Gravitons cannot be detected, due to the indistinguishability with ordinary quantum fluctuations of the detector. This agrees with data so far.
— Atoms emit neither gravitational waves nor gravitons.
— Gravitational waves exist and can be detected. This agrees with various experiments; the final, direct confirmation occurred in late 2015.
— The gravitational constant G does not run with energy – as long as the strand diameter can be neglected. In this domain, G is not renormalized. This prediction agrees with expectations and with data, though the available data is sparse.

All listed predictions are unspectacular; they are made also by other approaches that contain general relativity as limiting cases. In particular, the strand model, like many other approaches, predicts:

Page 309

▷ With the exception of the cosmological constant and of particle masses (and possibly the Sokolov–Ternov effect), *no quantum gravity effects will be observed.*

Gravity will not yield new measurable quantum effects. So far, this prediction agrees with experiment – and with almost all proposed models of quantum foam in the research literature. In other words, we have found *no unexpected* experimental predictions from the strand model in the domain of quantum gravity. This is the so-called *'no avail'* conjecture; and it is not a surprise. Ref. 97

Page 8 In fact, the Bronshtein cube of Figure 1 also implies:

▷ There is *no* separate theory of quantum gravity that includes relativity but does not include the other interactions.

There is no room for a theory of relativistic quantum gravity in nature.

In short, strands lead us to expect deviations from general relativity only in two domains: in cosmology (such as changes of the cosmological constant) and in particle physics. The rest of this chapter deals with cosmology. The subsequent chapters focus on particle physics.

Cosmology

Cosmology is an active field of research, and new data are collected all the time. We start with a short summary.

The sky is dark at night. This and other observations about the red shift show that the universe is surrounded by a horizon and is of finite size and age. Modern measurements show that cosmic age is about 13 800 million years. The universe expands; the expansion is described by the field equations of general relativity. The universe's expansion accelerates; the acceleration appears to be described by the *cosmological constant* Λ – also called *dark energy* – that has a small positive value. The universe is observed to be flat, and, averaged over large scales, homogeneous and isotropic. At present, the observed average matter density in the universe is about 18 times smaller than the energy density due to the cosmological constant. In addition, there appears to be a large amount of matter around galaxies that does not radiate; the nature of this *dark matter* is unclear. Galaxy formation started from early density fluctuations; the typical size and amplitude of the fluctuations are known. The topology of space is observed to be simple.

The strand model, like any unified description of nature, must reproduce and explain these observations. Otherwise, the strand model is wrong.

The finiteness of the universe

In the strand model, cosmology is based on the following idea:

▷ The ***universe*** is made of *one* fluctuating strand that criss-crosses from and to the horizon. Fluctuations increase the complexity of the strand tangledness

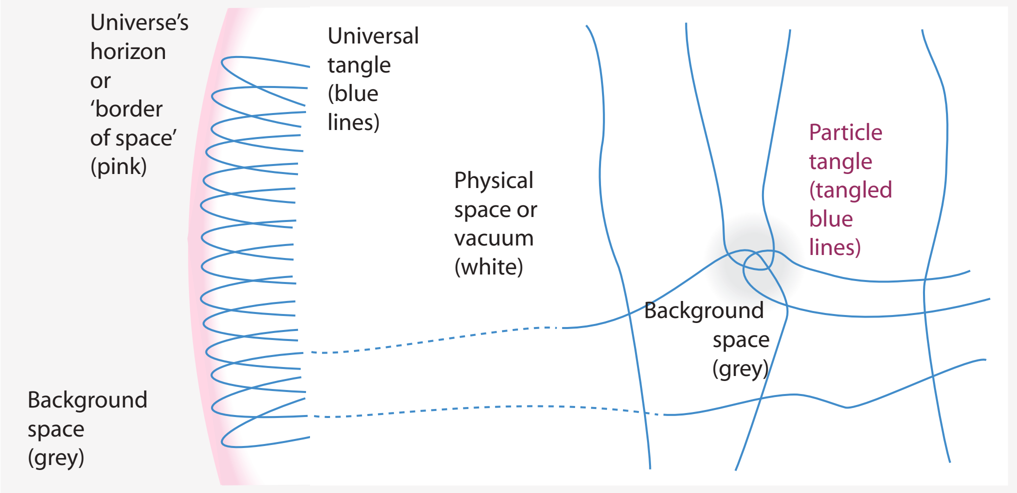

FIGURE 81 In the strand model, the universe is limited by a horizon, as schematically illustrated here. Physical space (white) matches background space (grey) only inside the horizon. Physical space thus only exists inside the cosmic horizon.

over time.

In other words, the strands of all particles are woven into the sky. The existence of finite size and of finite age then follows automatically:

▷ The ***universe's horizon*** appears at the age or distance at which the strand crossings cannot be embedded any more into a common three-dimensional background space. The horizon expands over time.

The strand model thus has a simple explanation for the finiteness of the universe and the horizon that bounds it: The universe's horizon is a weave that joins all strand tails. A schematic illustration of the cosmic horizon is given in Figure 81.

The strand model predicts that the horizon of the universe is an *cosmological particle* Ref. 210 *horizon*, an event horizon similar to that of a black hole. Until 1998, this possibility seemed questionable; but in 1998, it was discovered that the expansion of the universe is Ref. 211 accelerating. This discovery implies that the cosmic horizon is indeed an event horizon, as required by the strand model.

The strand model predicts that the universe is a kind of *inverted back hole*. Like for any situation that involves a horizon, the strand model thus does not allow us to make statements about properties 'before' the big bang or 'outside' the horizon. As explained Page 288 above, strands predict that there is nothing behind a horizon.

In particular, the strand model implies that the matter that appears at the cosmic horizon during the evolution of the universe appears through Bekenstein–Hawking radiation. This contrasts with the 'classical' explanation form general relativity that new matter appears simply because it existed behind the horizon beforehand and then crosses the horizon into the 'visible part' of the universe.

We note that modelling the universe as a single strand implies that it contains tangles.

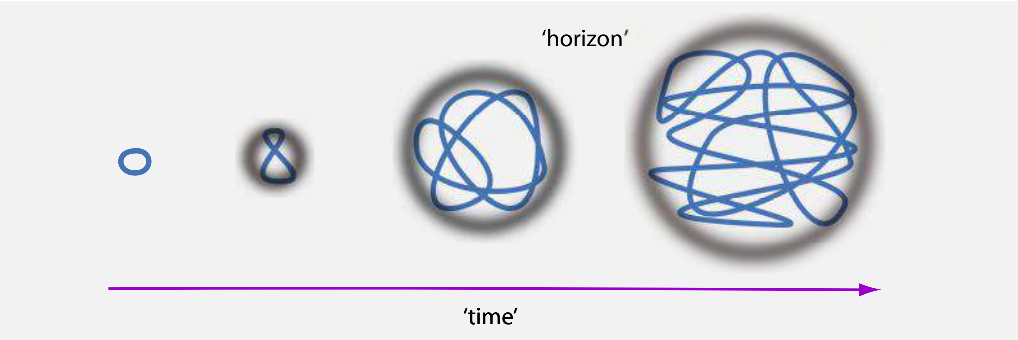

FIGURE 82 An extremely simplified view of how the universe evolved near the big bang. In this evolution, physical time, space and the surrounding horizon are in the process of getting defined.

In other words, the strand model makes the prediction that the universe cannot be empty, but that it must contain particles. Strand cosmology also confirms that the question of initial conditions for the universe does not really make sense: particles appear at the horizon.

Page 102

We also note that describing the universe as made of a single strand is a natural, but somewhat unusual way to incorporate what particle physicists and cosmologists like to call *holography*. Holography is the idea that all observables of a physical system are defined on a boundary enclosing the system. In other words, if we would know, at Planck scale, everything that happens on the walls of a room, we could know everything that is and goes on inside the room. Instead of holography, we could also call it the *NSA dream*. Holography is a consequence of the extension of the fundamental constituents of nature and is a natural consequence of the strand model. As a consequence, strand cosmology naturally reproduces holographic cosmology – though not fully, as is easy to check.

Challenge 189 e

> Or cette liaison ou cet accommodement de toutes les choses créées à chacune, et de chacune à toutes les autres, fait que chaque substance simple a des rapports qui expriment toutes les autres, et qu'elle est par conséquent un miroir vivant perpétuel de l'univers.*
>
> Gottfried Wilhelm Leibniz, *Monadologie*, 56.

The big bang – without inflation

Any expanding, homogeneous and isotropic matter distribution had earlier stages of smaller size and higher density. Also the universe has been hotter and denser in the past. But the strand model also states that singularities do not appear in nature, because there is a highest possible energy density. As a result, the big bang might be imagined as illustrated in Figure 82. Obviously, physical space and time are not well defined near that situation, so that the figure has to be taken with a grain of salt. Nevertheless, it shows how

Ref. 212
Page 293

* 'Now this connexion or adaptation of all created things to each and of each to all, means that each simple substance has relations which express all the others, and, consequently, that it is a perpetual living mirror of the universe.'

the evolution of the universe can be seen as resulting from the increase in tangledness of the strand that makes up nature.

The strand model leads to the conjecture that the evolution of the universal strand just after the big bang automatically yields both a homogeneous and isotropic matter distribution and a flat space. Also the scale invariance of early density fluctuations seems natural in the strand model. In short, the strand model looks like a promising alternative to *inflation*: the hypothesis of inflation becomes unnecessary in the strand model, because strand cosmology directly makes the predictions that seem so puzzling in classical cosmology. This issue is still subject of research.

The cosmological constant

At present (2019), the issue of the cosmological constant, thus of dark energy, is under investigation. Watch this space for updates.

The value of the matter density

The strand model predicts that horizons emit particles. As a consequence, the strand model predicts an upper limit for the number N_b of baryons that could have been emitted by the cosmic horizon during its expansion. For a horizon shining throughout the age of the universe t_0 while emitting the maximum power $c^5/4G$, we get

$$N_{b0} \leqslant \frac{t_0\, c^5/4G}{m_b c^2} = 2.6 \cdot 10^{79} \ . \qquad (194)$$

Equality would hold only if the contributions of photons, electrons, neutrinos and dark matter could be neglected. In short, using the age $t_0 = 13.8$ Ga, the strand model predicts that at most $2.6 \cdot 10^{79}$ baryons exist in the universe at present. Modern measurements indeed give values around this limit. Ref. 215

More about matter and energy densities will be added here soon.

Open challenge: What are the effects of dark matter?

Conventionally, it is argued that *cold* dark matter exists for three reasons: First, it is necessary to grow the density fluctuations of the cosmic microwave background rapidly enough to achieve the present-day high values. Secondly, it is needed to yield the observed amplitudes for the acoustic peaks in the cosmic background oscillations. Third, it explains the rotation curves observed in hundreds of galaxies and galaxy clusters. Can the strand model change these arguments? Challenge 190 ny

Later on, it will be argued that following the strand model, dark matter can only be a mixture of conventional matter and black holes. How does this dark matter prediction explain the galaxy rotation curves? This leads to a really speculative question: Could tangle effects at the scale of a full galaxy be related to dark matter? Research is ongoing, and will be added here in the coming months. Challenge 191 ny

The topology of the universe

In the strand model, physical space-time, whenever it is defined, *cannot* be multiply connected. Also all quantum gravity approaches make this prediction, and the strand model confirms it: because physical space-time is a result of averaging strand crossing switches, non-trivial topologies (except black holes) do not occur as solutions. For example, the strand model predicts that wormholes do not exist. In regions where space-time is undefined – at and beyond horizons – it does not make sense to speak of space-time topology. In these regions, the fluctuations of the universal strand determine observations. In short, the strand model predicts that all searches for non-trivial *macroscopic* (and microscopic) topologies of the universe, at both high and low energies, will yield negative results. So far, this prediction agrees with all observations.

Predictions of the strand model about cosmology

In the domain of cosmology, the strand model makes the following testable predictions.

— The universe is not empty. (Agrees with observation.)
— Its integrated luminosity does not exceed the power limit $c^5/4G$. (Agrees with observation.)
— The universe's energy density does not exceed the entropy bound. (Agrees with observation.)
— The are no singularities in nature. (Agrees with observation.)
— The matter density of the universe decreases with age, roughly as $\rho t \sim 1/t^2$. (Checks are under way. This prediction differs from the usual cosmological models.)
— There is nothing behind the cosmic horizon. Matter, energy and space appear at the horizon. (Agrees with observations and requirements of logic.)
— Early density fluctuations are scale-invariant. (Agrees with observation.)
— The universe is flat and homogeneous. (Agrees with observation.)
— Apart (maybe) from the cosmological constant Λ, all other fundamental constants of nature are constant over time and space. (Agrees with observation, despite occasional claims of the contrary.)
— Inflation is unnecessary.
— The universe's topology is trivial. There are no wormholes, no time-like loops, no cosmic strings, no toroidal black holes, no domain walls and no regions of negative energy. (Agrees with observation.)
— The above statements are independent of the precise fluctuation details. (Can be tested with mathematical investigations.)

All these predictions can and will be tested in the coming years, either by observation or by computer calculations.

SUMMARY ON MILLENNIUM ISSUES ABOUT RELATIVITY AND COSMOLOGY

We have deduced special relativity, general relativity and cosmology from the strand model. The fundamental principle of the strand model implies the invariant Planck units, the Lagrangian and action of general relativity, the finiteness of the universe and, above all, it explains in simple terms the entropy of black holes.

Space-time foam is replaced by the strand model of the vacuum: empty space is the time-average of untangled strands. More precisely, space is the thermodynamic average of crossing switches that are due to shape fluctuations of untangled strands.

The strand model – and in particular, the strand model of the vacuum – explains the number of space-time dimensions, the vacuum energy density, the matter density and the finiteness of the universe. The cosmological constant is a consequence of the finite size of the universe. The issue of the initial conditions of the universe has been defused. The macroscopic and microscopic topology of the universe is simple. And dark matter is predicted to be, as shown in the next chapter, a combination of conventional matter Page 355 and black holes.

The most important predictions of the strand model are the decrease of the cosmological constant with time and the absence of inflation. Various experiments will test these predictions with increased precision in the coming years. So far, measurements do not contradict these predictions.

The strand model confirms that the speed of light c and the corrected Planck force $c^4/4G$ are *limit* values. The strand model also predicts that no variation in space and time of c, G, $\hbar$ and k can be detected, because they define all measurement units.

The strand model predicts that the cosmological constant and the masses of the elementary particles are the *only* quantum effects that will be observed in the study of Page 279 gravitation. Strands strongly suggest that additional effects of quantum gravity cannot be measured. In particular, no effects of space-time foam will be observed.

The strand model is, at present, the simplest – but not the only – known model of quantum gravity that allows deducing all these results. In particular, the strands' explanation of black hole entropy is by far the simplest one known.

General relativity is an approximation of the strand model. The approximation appears when the quantum of action and, in particular, the strand diameter are neglected. General relativity and cosmology thus appear by approximating $\hbar$ as 0 in the strand Page 8 model – as required by the Bronshtein cube of physics that is shown in Figure 1. Strands imply that general relativity is valid for all energies below the Planck energy. In other words, the strand model is not a generalization of general relativity. This conforms to Page 149 the list of requirements for the final theory.

Page 164 If we look at the millennium list of open issues in physics, we see that – except for the issue of dark matter – all issues about general relativity and cosmology have been settled. The strand model explains the mathematical description of curved space-time and of general relativity. The strand model also provides a simple model of quantum gravity – Page 276 maybe the simplest known one. Above, we had already shown that the strand model explains all mathematical structures that appear in quantum theory and in particle physics. Together with the results from this chapter we can now say: *the strand model explains* all *concepts, i.e., all mathematical structures that appear in physical theories.* In particular, strands explain the metric, curvature, wave functions, field intensities – and the probabilistic behaviour of all of them. They all result from averaging crossing switches.

In summary, starting from the fundamental principle of the strand model, we have understood that strands are the origin of gravitation, general relativity, quantum gravity and cosmology. We have also understood the mathematical description of gravitation – and, before, that of quantum physics – found in all textbooks. These results encourage us to continue our quest. Indeed, we are not done yet: we still need to deduce the possible elementary particles and to explain their properties.

CHAPTER 11

THE PARTICLE SPECTRUM DEDUCED FROM STRANDS

> No problem can withstand the assault of sustained thinking.
>
> Voltaire**

Ref. 218

Strands describe quantum theory, gauge interactions and general relativity. But do strands also settle all issues left open by twentieth-century physics? Do they settle the origin of all the elementary particles, their quantum numbers, their masses and their mixing angles? How does the infinite number of possible tangles lead to a finite number of elementary particles? And finally, do strands explain the coupling constants? In the millennium list of open issues in fundamental physics, these are the issues that remain. The strand model is correct only if these issues are resolved.

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In this chapter, we show that the strand model indeed explains the known spectrum of elementary particles, including the three generations of quarks and leptons. The strand model is the first approach of modern physics that can provide such an explanation.

It should be stressed that from this point onwards, the ideas are particularly speculative. In the chapters so far, the agreement of the strand model with quantum field theory and general relativity has been remarkable. The following chapters assign specific tangles to specific particles. Such assignments are, by nature, not completely certain. The speculative nature of the ideas now becomes particularly apparent.

PARTICLES AND QUANTUM NUMBERS FROM TANGLES

In nature, we observe three entities: vacuum, horizons and particles. Of these, (quantum) *particles* are *localized* entities with specific *intrinsic* properties, i.e., properties that do not depend on their motion.

In nature, all the intrinsic properties of every particle, every object and every image are completely described by three types of *basic* properties: (1) the elementary particles they contain, (2) their behaviour under space-time transformations, (3) their interactions. The full list of these basic intrinsic properties of particles is given in Table 12.

Given the basic intrinsic properties for each elementary particle, physicists can deduce *all* those intrinsic particle properties that are *not* listed; examples are the half life, decay modes, branching ratios, electric dipole moment, T-parity, gyromagnetic ratio or electric polarizability. Of course, the basic intrinsic properties also allow physicists to deduce *every* property of every object and image, such as size, shape, colour, brightness, density,

** Voltaire (b. 1694 Paris, d. 1778 Paris) was an influential philosopher, politician and often satirical writer.

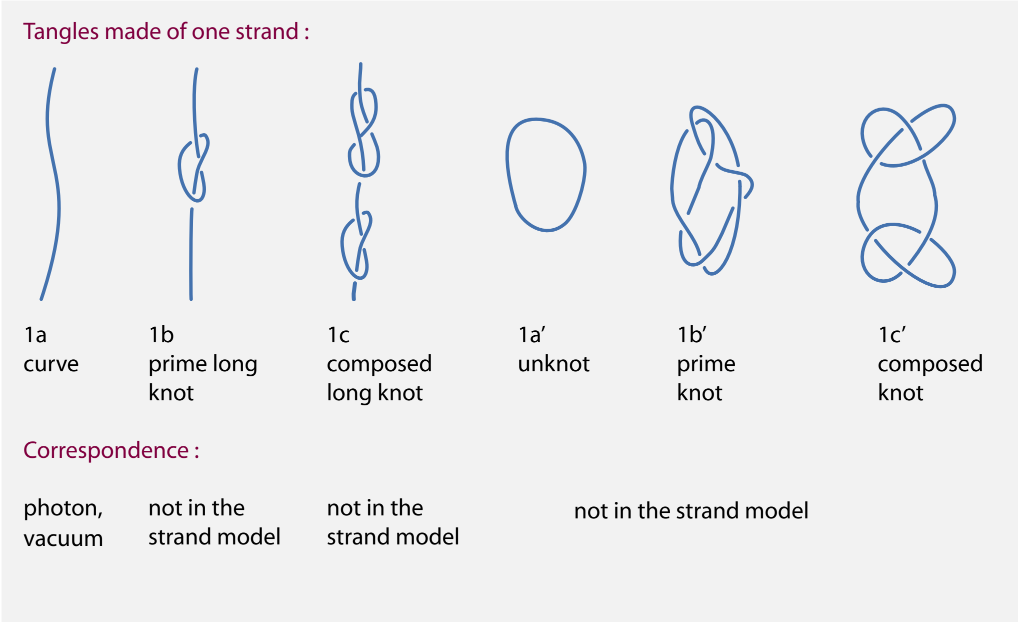

FIGURE 83 Examples for each class of tangles made of one strand.

elasticity, brittleness, magnetism or conductance.

In short, understanding *all* properties of matter and images thus only requires understanding the *basic* properties of quantum particles; and understanding the *basic* properties of quantum particles only requires understanding the *basic* properties of the *elementary* particles.

The strand model states that all elementary (and all composed) particles are tangles of strands. This leads us to ask: Which tangle is associated to each elementary particle? What kinds of elementary particles are possible? Do these tangles reproduce, for each elementary particle, the observed values of the basic properties listed in Table 12?

Page 175

It turns out that the strand model only allows a *limited number* of *elementary* particles. In addition, the tangles of these elementary particle have intrinsic properties that *match* the observed properties. To prove these strong statements, we first recall that all massive elementary particles are represented by an *infinite sequence* of tangles. We now explore tangles according to the number of strands they are made of.

PARTICLES MADE OF ONE STRAND

In the strand model, all particles made of *one* strand have spin 1, are elementary, and are bosons. Conversely, all massless elementary spin-1 bosons can only have two tails, and thus must be made of a single strand. Such one-stranded tangles return to the original strand after a core rotation by 2π. Massive elementary spin-1 bosons can have one or more strands. Tangles of more than one strand can only have spin 1 if they represent massive elementary or composed particles. In short, classifying one-stranded tangles allows classifying all elementary gauge bosons.

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TABLE 12 The full list of *basic* intrinsic properties of quantum particles, from which all other observed intrinsic properties of particles, objects and images can be deduced.

Property	Possible Value	Determines
Quantum numbers due to space-time symmetries:		
Spin S or J	integer or half-integer multiple of $\hbar$	statistics, rotation behaviour, conservation
P parity	even (+1) or odd (−1)	behaviour under reflection, conservation
C parity	even (+1) or odd (−1)	behaviour under charge conjugation, conservation
Interaction properties:		
Mass M	between 0 and the Planck mass	gravitation, inertia
Electric charge Q	integer multiples of one third of electron or proton charge	Lorentz force, coupling to photons, conservation
Weak charge	rational multiple of weak coupling constant	weak scattering and decays, coupling to W and Z, partial conservation
Mixing angles	between 0 and $\pi/2$	mixing of quarks and neutrinos, flavour change
CP-violating phases	between 0 and $\pi/2$	degree of CP violation in quarks and neutrinos
Strong charge, i.e., colour	rational multiple of strong coupling constant	confinement, coupling to gluons, conservation
Flavour quantum numbers, describing elementary particle content:		
Lepton number(s) L'	integer(s)	conservation in strong and e.m. interactions
Baryon number B	integer times 1/3	conservation in all three gauge interactions
Isospin I_z or I_3	+1/2 or −1/2	up and down quark content, conservation in strong and e.m. interactions
Strangeness S'	integer	strange quark content, conservation in strong and e.m. interactions
Charmness C'	integer	charm quark content, conservation in strong and e.m. interactions
Bottomness B'	integer	bottom quark content, conservation in strong and e.m. interactions
Topness T'	integer	top quark content, conservation in strong and e.m. interactions

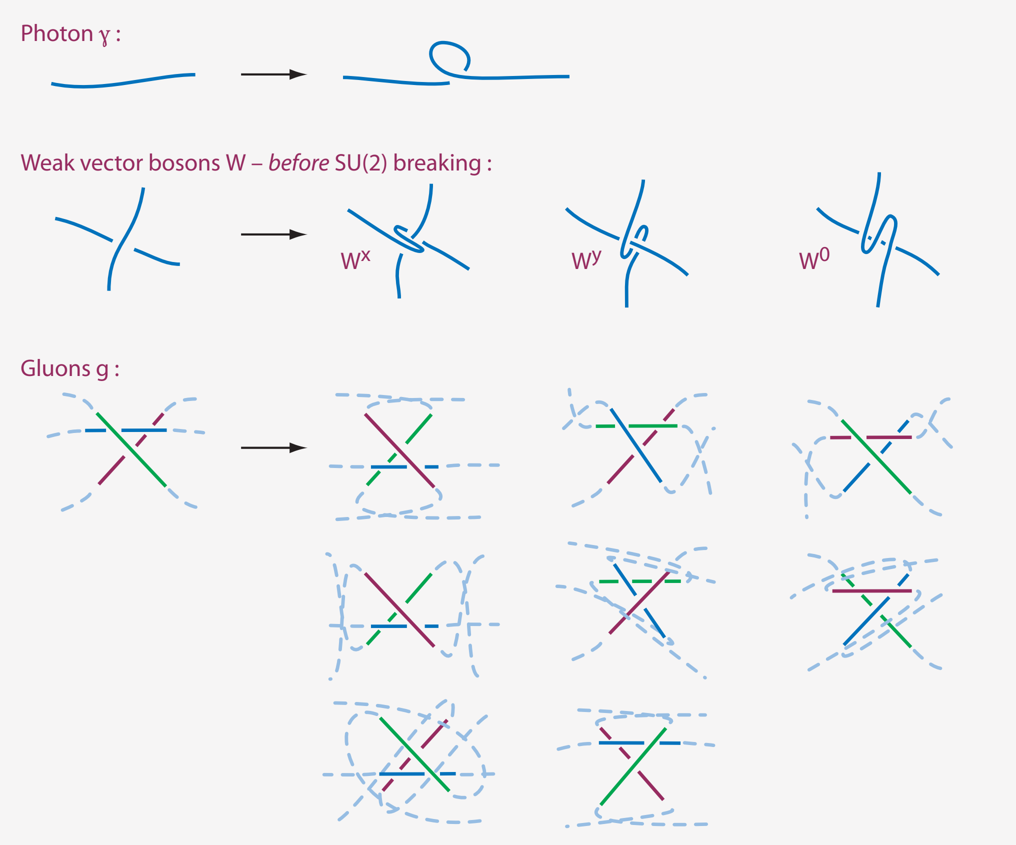

FIGURE 84 The gauge bosons in the strand model. All differ from vacuum by one curved strand – though, for clarity, the gluons are shown here using their complementary two-strand moves.

Mathematicians have already classified one-stranded tangles; they are usually called *open knots* or *long knots*. To get an overview, we list an example for each class of one-stranded tangles on the left-hand side of Figure 83. For completeness, closed curves are shown on the right-hand side of the figure. We now explore each of these classes.

Unknotted curves

The simplest type of tangle made of one strand is an *unknotted curve*, shown as example 1a in Figure 83. The study of gauge interactions has shown that unknotted strands are, depending on their precise average shape, either vacuum strands or gauge bosons.

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The time-average of a vacuum strand is straight. A single strand represents a particle if the time-averaged strand shape is not a straight line.

In the strand model, vacuum strands in flat space are, on average, *straight*. In this property, vacuum strands differ from gauge bosons, which, on average, have *curved* strands, and thus carry energy.

Gauge bosons – and Reidemeister moves

Gauge bosons are the carrier particles of the interactions. In the strand model, the gauge interactions are due to the three Reidemeister moves. The electromagnetic, the weak and the strong interaction correspond to respectively the first, second and third Reidemeister move. As we have seen above, when the three Reidemeister moves deform fermion tangle cores they generate U(1), SU(2) and SU(3) gauge symmetries. The detailed exploration Page 225 of the correspondence between tangle deformation and gauge theory led us to the gauge boson tangles shown in Figure 84.

▷ All gauge bosons – before symmetry breaking when applicable – are due to a single moving and curved strand.

A single strand represents a particle if the time-averaged strand shape is not a straight line. The lack of straightness implies non-vanishing energy. A single-strand particle can thus be either a strand with a bulge or a strand whose tails are not aligned along a straight line. The size of the bulge is related to the wavelength.

Page 228 As explained above, the *first* Reidemeister move, the twist, leads to the modelling of photons as helical strands. Therefore, photons have vanishing mass and two possible polarizations. Photons do not have tangled, localized family members; photons are massless. Their specific unknotted and twisted strand shapes also imply that photons generate an Abelian gauge theory and that photons do not interact among themselves. Automatically, photons have no weak and no strong charge. The strand model further implies that photons have negative P-parity and C-parity, as is observed.

Page 246 The study of the *second* Reidemeister move, the *poke*, showed that deformations induced by pokes can also involve braiding of tangle tails; this leads to the symmetry breaking of the weak interaction. As a result, the observed W and the Z boson strands become massive. The tangle of the W is chiral, and thus it is electrically charged; the tangle of the Z is achiral and thus electrically neutral. Being tangled, the W and the Z also carry weak charge and thus interact among themselves, generating a *non-Abelian* gauge theory. The strand model also implies that the W and the Z have no P-parity, no C-parity and no colour charge, as is observed.

Page 258 The study of the *third* Reidemeister move, the slide, led us to the existence of eight gluons. The eight gluons are unknotted, thus they carry no mass, no electric charge and no weak charge. Each gluon tangle has two possible polarizations. The strand model of gluons also implies that they have negative P-parity and no C-parity, as is observed. Gluons tangles carry colour and interact among themselves, thus they generate a non-Abelian gauge theory. In contrast to the other two interactions, free, single gluons are short-lived, because their structure induces rapid hadronization: when gluons act on the vacuum, quark–antiquark pairs are produced. Gluons do not have tangled family members; they are massless in the high energy limit, when their tails are aligned.

For completeness we mention that by assignment, all gauge bosons differ from vacuum by a single curved strand, have vanishing lepton and baryon numbers, and thus also lack all flavour quantum numbers. All this is as observed.

The strand model explains the lack of *classical SU(2) field waves* as a consequence of the breaking of the SU(2) symmetry and the consequent mass of the weak bosons.

Strands explain the lack of *classical SU(3) waves*, also called *gluonic waves*, as a consequence of the topological impossibility to produce such waves, which is related to the infinite mass of single free gluons.

In somewhat sloppy language we can say that the shape and the effects of photons are one-dimensional, those of the unbroken weak bosons are two-dimensional, and those of the gluons are three-dimensional. This is the essential reason that they reproduce the U(1), SU(2) and SU(3) groups, and that no higher gauge groups exist in nature.

In summary, Reidemeister's theorem implies that the list of known gauge bosons with spin 1 is complete. But the list of possible tangles made of a single strand is much more extensive; we are not done yet.

Open or long knots

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Single strands could also contain knotted regions. We have explained earlier on that all such possibilities – mathematically speaking, all so-called *open knots* or *long knots* – have no relation to particles. In the strand model, they cannot appear and thus play no role. The original strand model from 2008 did include such configurations as particles (for example as W and Z bosons), but it now – i.e., after 2014 – seems that this inclusion is an unnecessary complication.

Closed tangles: knots

Figure 83 shows, on the right-hand side, examples for all classes of *closed* tangles of one strand, i.e., of tangles *without tails*. They are usually just called *knots* in mathematics. In the strand model knots do not appear. They do not seem to have physical relevance and we do not explore them here.

Summary on tangles made of one strand

In summary, a single strand represents a particle if the strand shape is, on average, not a straight line. This distinguishes a vacuum strand from a particle strand. A particle strand can thus be a strand with a bulge or a strand whose tails are not aligned along a straight line. All tangles made of *one open strand* represent *elementary* particles of spin 1, thus elementary vector bosons.

Massless elementary spin-1 particles are made of one open strand also because other tangles cannot reproduce both zero mass and the spin-1 behaviour under rotations: only one-stranded tangles return to the original strand after a core rotation by 2π and allow vanishing mass at the same time.

In the strand model, the tangle made of one curved strand is assigned to the *photon*. The strand model correctly reproduces and thus explains the properties of the photon.

Particles made of two strands

In the strand model, particle tangles can also be made of *two* strands. Examples for all the classes of two-stranded tangles are given in Figure 85. Each class has a physical particle assignment.

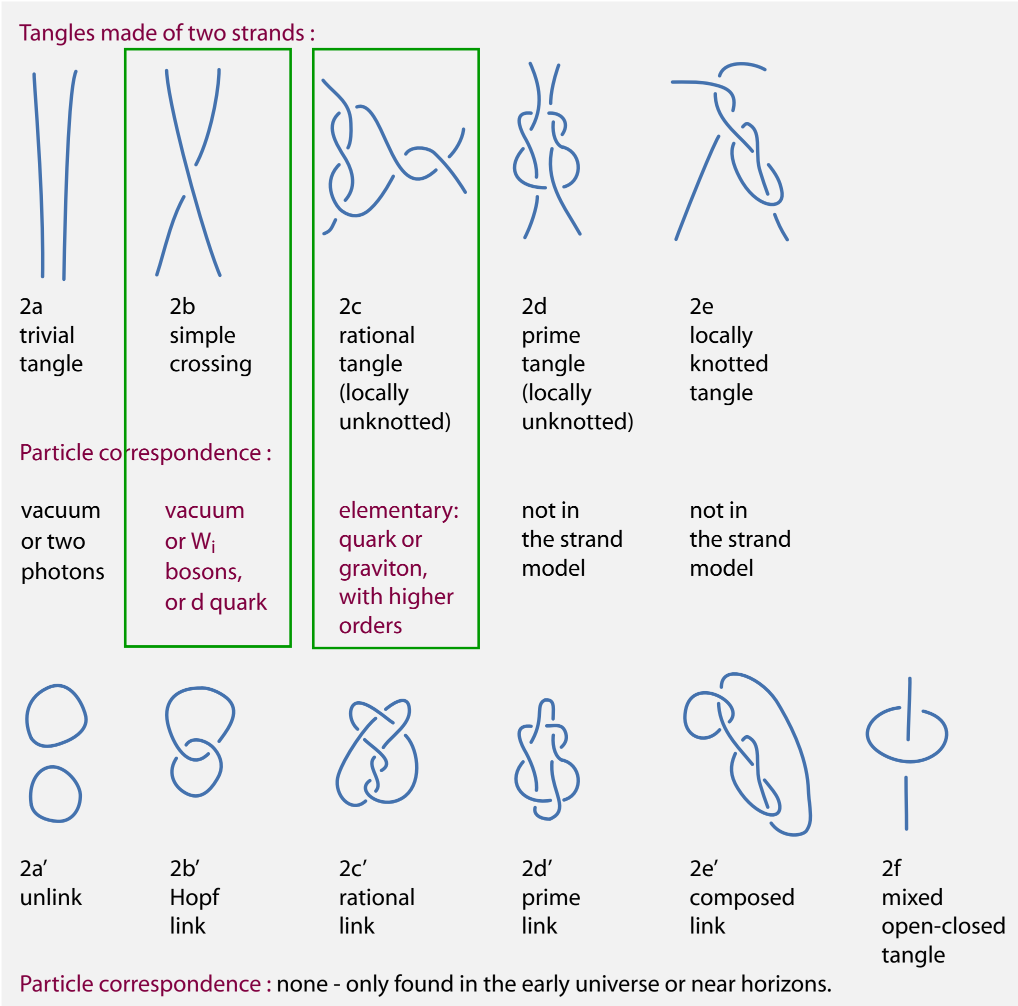

FIGURE 85 Possible tangles made of two strands.

— The simplest tangle made of two strands is the *trivial tangle*, shown as example 2a in Figure 85. In the strand model, the trivial tangle, like all *separable* tangles, is a *composite* system. Each of the two strands can represent either the vacuum or a photon. Simply stated, the trivial tangle of two strands is not an elementary particle.

— The simplest non-trivial tangle made of two strands is the *crossing*, shown as 2b in Figure 85. In the strand model, the crossing appears as part of the vacuum, or as unbroken W bosons; in addition, for certain tail configurations, it can represent a down quark, as we will see below.

— A new class of tangles are the *rational tangles*, represented by example 2c in the figure. A rational tangle is a tangle that can be untangled by moving its tails *around*. (Also example 2b is a rational tangle.) Rational tangles are distinct from prime and from locally knotted tangles, shown as examples 2d and 2e, which require pulling the tail *through* the tangle to untangle it. Rational tangles are thus *weakly* tangled. As we will

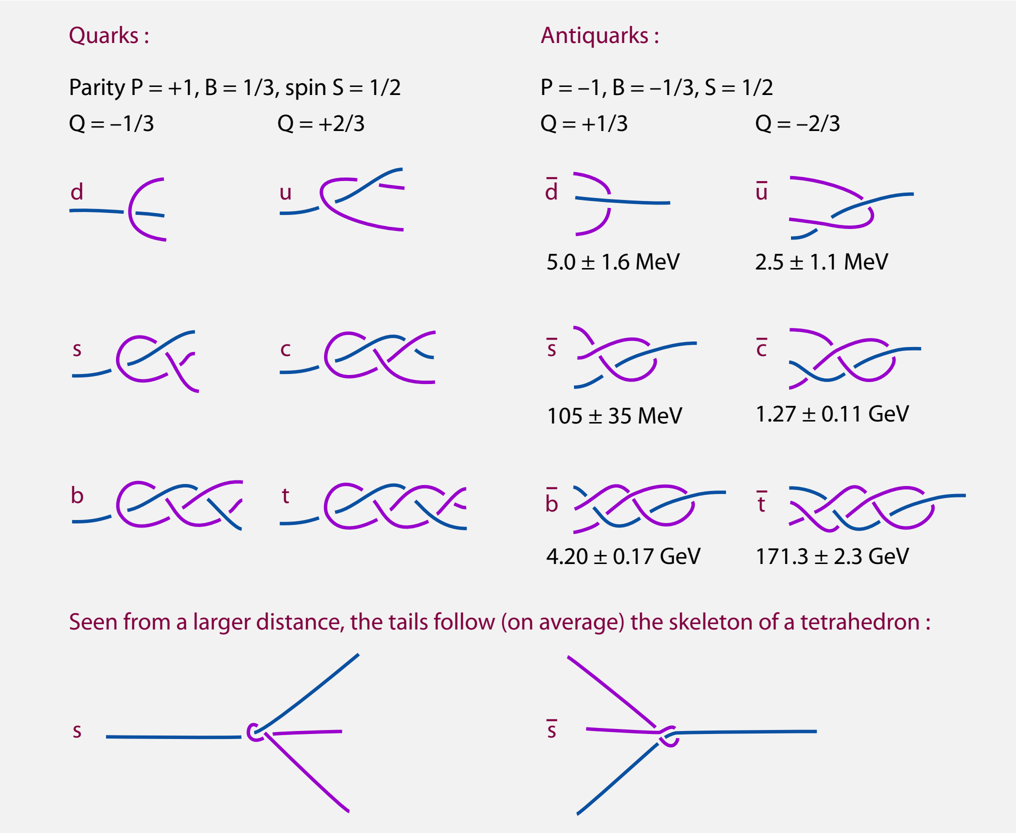

FIGURE 86 The simplest tangles assigned to the quarks and antiquarks. For reference, the experimental mass values are also given.

see,

▷ Rational tangles of two strands represent the *graviton* and the *quarks*.

We will discuss them in detail in the next two sections. More complicated rational tangles are higher-order propagating states of the simpler ones.

— Another class of tangles are *prime tangles*, for which the tangle 2d is an example. Like knotted one-stranded tangles, we conclude that prime tangles are not part of the strand model.
— Still another class of tangles are *locally knotted tangles*, shown as example 2e. Also this class is not part of the strand model.
— Finally, *closed tangles*, *links* and *mixed tangles*, shown in the lower row of Figure 85, have no role in the strand model.

In short, the only two-stranded tangles of interest in the strand model are the rational tangles. We now explore them in more detail.

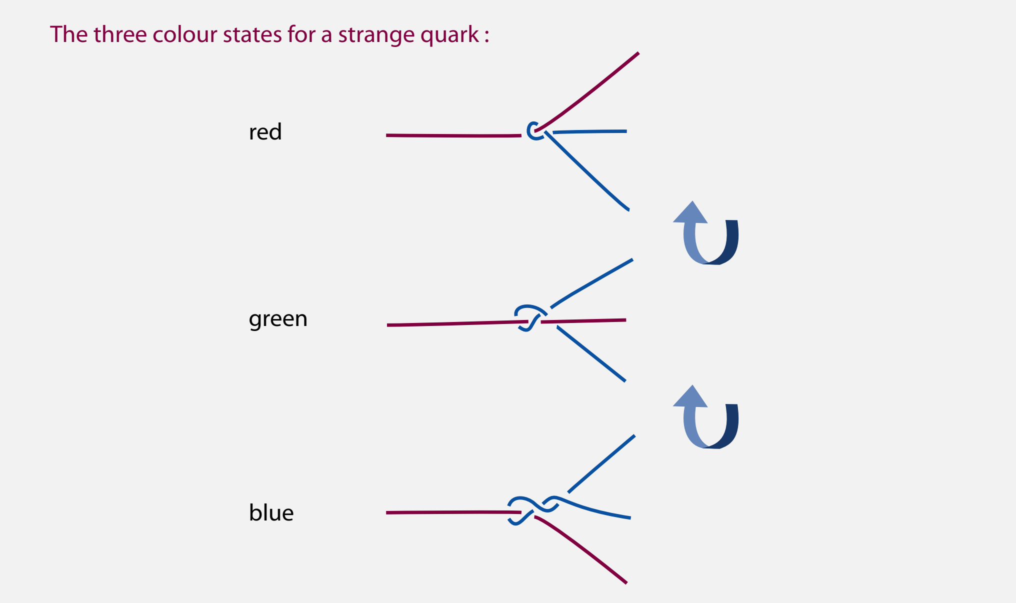

FIGURE 87 The three colour charges correspond to the three possible spatial orientations; the centre tail on the right is always above the paper plane, the other two tails on the right are below the paper plane.

Quarks

Page 271 The exploration of the strand model and of the strong interaction showed: the tangle of a coloured fermion, thus of a quark, must be rational, must reproduce the three possible colour options, and must break the three-belt symmetry.

The simplest tangles that realize these requirements are shown in Figure 86: quark tangles are *rational tangles* made of *two* strands. Higher quark generations have larger crossing numbers. The four tails form the skeleton of a tetrahedron. A particle with two strands tangled in this way automatically has spin 1/2. The electric charges of the quarks are 1/3 and −2/3, an assignment that is especially obvious for up and down quarks and that will become clearer later on, in the study of hadrons. Parity is naturally assigned as done in Figure 86. Baryon number and the other flavour quantum numbers – isospin, strangeness, charm, bottomness, topness – are naturally assigned as usual. The flavour quantum numbers simply 'count' the number of corresponding quark tangles. Like all localized tangles, quarks have weak charge. We will explore weak charge in more detail Page 382 below. Antiquarks are mirror tangles and have opposite quantum numbers. We will see Page 332, page 340 below that these assignments reproduce the observed quantum numbers of all mesons and baryons, as well as all their other properties.

We note that the simplest version of the down quark is a simple crossing; nevertheless, it differs from its antiparticle, because the simple crossing mixes with the braid with seven crossings, 13 crossings, etc.; this mixing is due to the leather trick, as shown below. Page 323 And for every quark type, these more complicated braids differ from those of their antiparticles.

For each quark, the four tails form the skeleton of a tetrahedron. In Figure 86 and

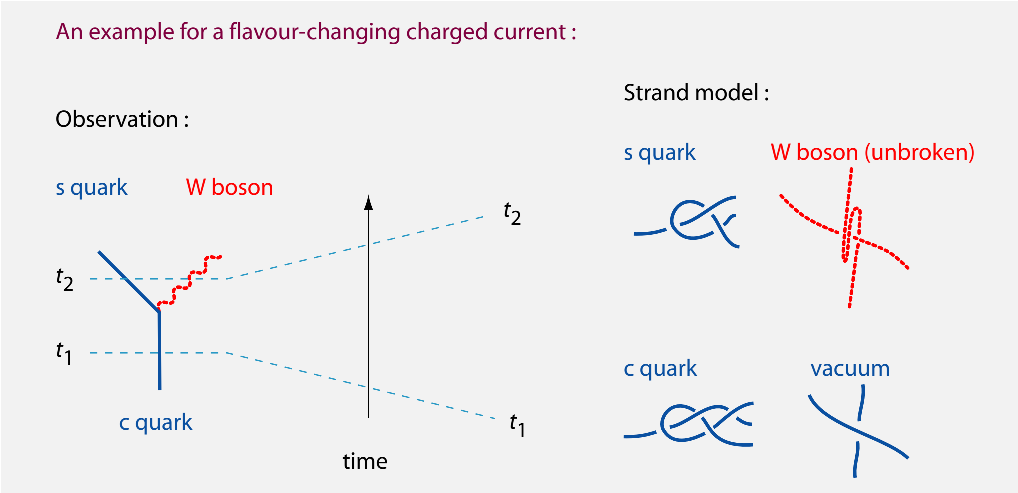

FIGURE 88 Absorption or emission of a W boson changes quark flavour.

Figure 87, the tetrahedral skeletons are drawn with one tail in the paper plane; of the other three tails, the middle one is assumed to be *above* the paper plane, and the outer two tails to be *below* the paper plane. This is important for the drawing of quark compounds later on. The three tails allow us to reproduce the strong interaction and the colour charge of the quarks: each colour is one of three possible orientations in space; more precisely, the three colours result from the three possible ways to map a quark tangle to the three belt structure. Each colour corresponds to a different choice for the tail that lies above the paper plane, as shown in Figure 87. The colour interaction of quarks will be clarified in the section on mesons.

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Page 332

In the strand model, the quark tangles thus carry *colour*. In nature, no free coloured particle has been observed. The strand model reproduces this observation in several ways. First of all, all leptons and baryons are colour-neutral, as we will see shortly. Secondly, only free quark tangles, as shown in Figure 86, have a definite colour state, because they have a fixed orientation in space. Thirdly, free quark states, thus quark states in the tetrahedral configuration of Figure 86, do not fit into vacuum even at large distances from the core; thus free quarks carry infinitely high energy. In practice, this means that free quark states do not occur in nature. Indeed, a free, coloured quark tangle can reduce its energy by interacting with one or several other quarks. The result is a strong colour attraction between quarks that leads to colourless composites.

In short, also in the strand model, only colourless composites of quarks exist as stable free particles. We will explore quark composites and the issue of confinement of quarks in more detail shortly.

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Page 337

In nature, quarks are weakly charged and interact with W bosons. In the strand model, the absorption or the emission of a W boson is the operation that takes a quark tangle and adds or subtracts a braiding step. This process is illustrated in Figure 88, which shows that a braiding (unbraiding) operation corresponds to the emission (absorption) of an W boson before symmetry breaking. It is straightforward to check that this operation fulfils all conservation laws and properties that are observed for these so-called

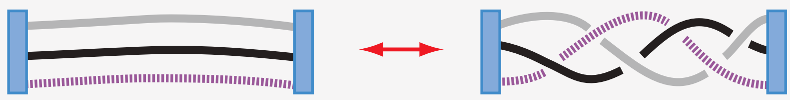

FIGURE 89 The leather trick is the deformation process that changes these two structures into each other. The leather trick limits structures made of three-stranded braids to six basic types.

flavour-changing charged currents. The absorption or emission of an (unbroken) Z boson has no braiding effect. The strand model thus reproduces the result that only the charged weak bosons can change quark flavours, as is observed.

For completeness, we mention that quarks, being tangles of *two* strands, have vanishing lepton number. Indeed, as we will see below, lepton tangles are made of *three* strands. Page 325

In summary, all quantum numbers of quarks are reproduced by the strand model, as long as quarks are modelled as braids of two strands with ends directed along the corners of a tetrahedron.

Quark generations

We stress that the quark tangles shown Figure 86 represent only the *simplest* tangle for each quark. First of all, longer braids are mapped to each of the six quarks. This might seem related to the *leather trick* shown in Figure 89. This trick is well-known to all people in the leather trade: if a braid of three strands has $n \geqslant 6$ crossings, it can be deformed into a braid with $n-6$ crossings. We might conjecture that, due to the leather trick, there is no way to introduce more than 6 quarks in the strand model.

In fact, the leather trick argument assumes that the braid end – and thus the ends of the strands – can be moved *through* the braids. In the strand model, this can only happen at the horizon, the only region where space (and time) are not well-defined, and where such manipulations become possible. The low probability of such a process will be important in the determination of quark masses. Page 365

Instead of resting on the leather trick, it is simpler to assume that braids with large numbers of crossings are mapped modulo 6 to the braids with the smallest number of crossings. This is consistent, because in the strand model, a braid with six additional crossings is mapped to a particle together with a virtual Higgs boson. The modulo 6 rule thus represents the Yukawa mass generation mechanism in the strand model. Page 330

In summary, in the strand model, each quark is not only represented by the tangles shown in Figure 86, but also by tangles with 6 additional crossings, with 12 additional crossings, etc.

As a mathematical check, we can also ask whether *all* other rational tangles are mapped to quarks. Rational tangles of higher complexity arise by repeatedly twisting any pair of tails of a quark tangle. This process produces an infinite number of complex two-stranded tangles. In the strand model, these tangles are quarks surrounded by virtual particles. Equivalently, we can say that all the more complex rational tangles that do not appear in Figure 86 are higher-order propagators of quarks. Challenge 192 e

FIGURE 90 The graviton in the strand model.

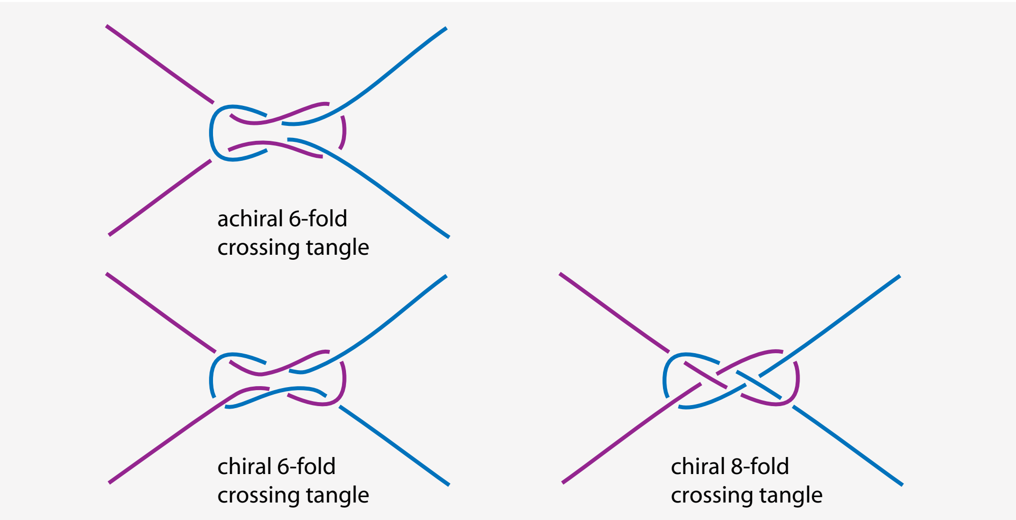

FIGURE 91 Which particle states are described by these tangles?

The graviton

Page 298 One rational tangle made of two strands is special. This special tangle is shown (again) in Figure 90. It differs from a quark tangle in one property: the tails are parallel (and near) to each other, and thus lie (almost completely) in a plane. Its tangle core returns to its original state after rotation by π, and therefore models a spin-2 particle. The tangle is not localized along its propagation direction; thus it has no mass, no electric and no weak charge. It also has no colour charge. *The tangle represents the graviton.* Similar tangles with higher winding numbers represent higher orders in the perturbation theory of gravitation.

Page 281 The chapter on gravitation has already shown how gravitons lead to curvature, horizons and the field equations of general relativity.

A puzzle

The topic of two-stranded tangles also requires to solve the puzzle of Figure 91. To which Challenge 193 s physical states do the three pictured tangles correspond?

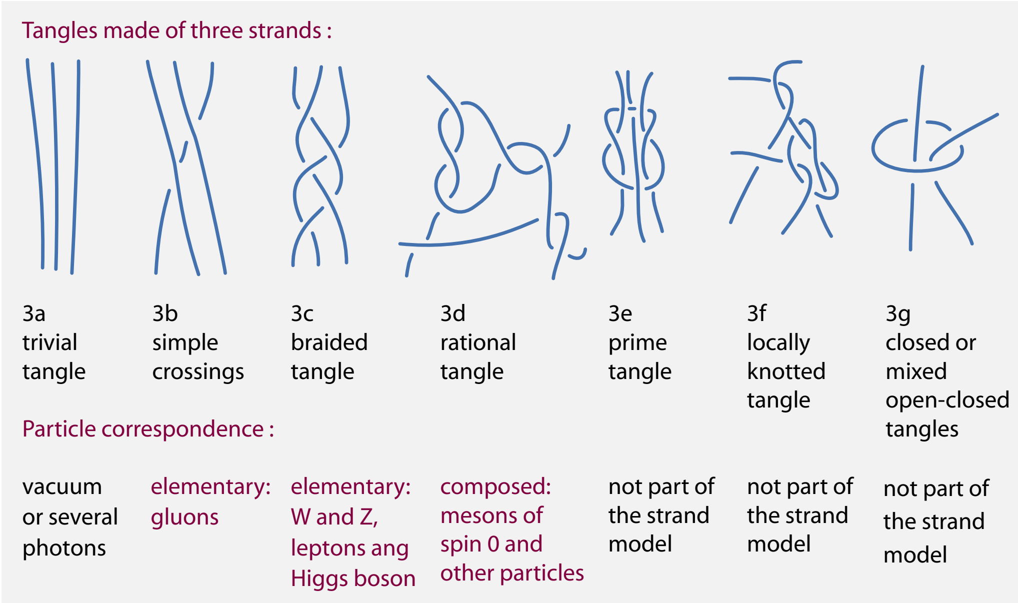

FIGURE 92 Examples for all the classes of tangles made of three strands.

Summary on two-stranded tangles

In summary, the strand model predicts that apart from the six quarks, the graviton, and the unbroken weak bosons W_1, W_2 and W_3, no other two-stranded elementary particle exists in nature.

Quarks and the graviton, the elementary particles made of two strands, are *rational* tangles. Their strand models are thus not tangled in a complicated way, but tangled in the *least complicated* way possible. This connection will be of importance in our search for elementary particles that are still undiscovered.

Particles made of three strands

In the strand model, the next group are particles made of *three* strands. Examples for all classes of three-stranded tangles are given in Figure 92. Several classes of three-stranded tangles turn out to be composites of two-stranded particles. However, a number of tangles are new and represent elementary particles.

Leptons

The candidate tangles from 2008 for the leptons shown in Figure 93 are the simplest possible non-trivial tangles with three strands. These lepton tangles are simple braids with tails reaching the border of space. The six tails probably point along the coordinate axes. These braided tangles have the following properties.

— Each lepton is localized. Each lepton has mass: its three tails con be braided, thus have non-vanishing Yukawa coupling, thus generate mass. And each lepton has spin

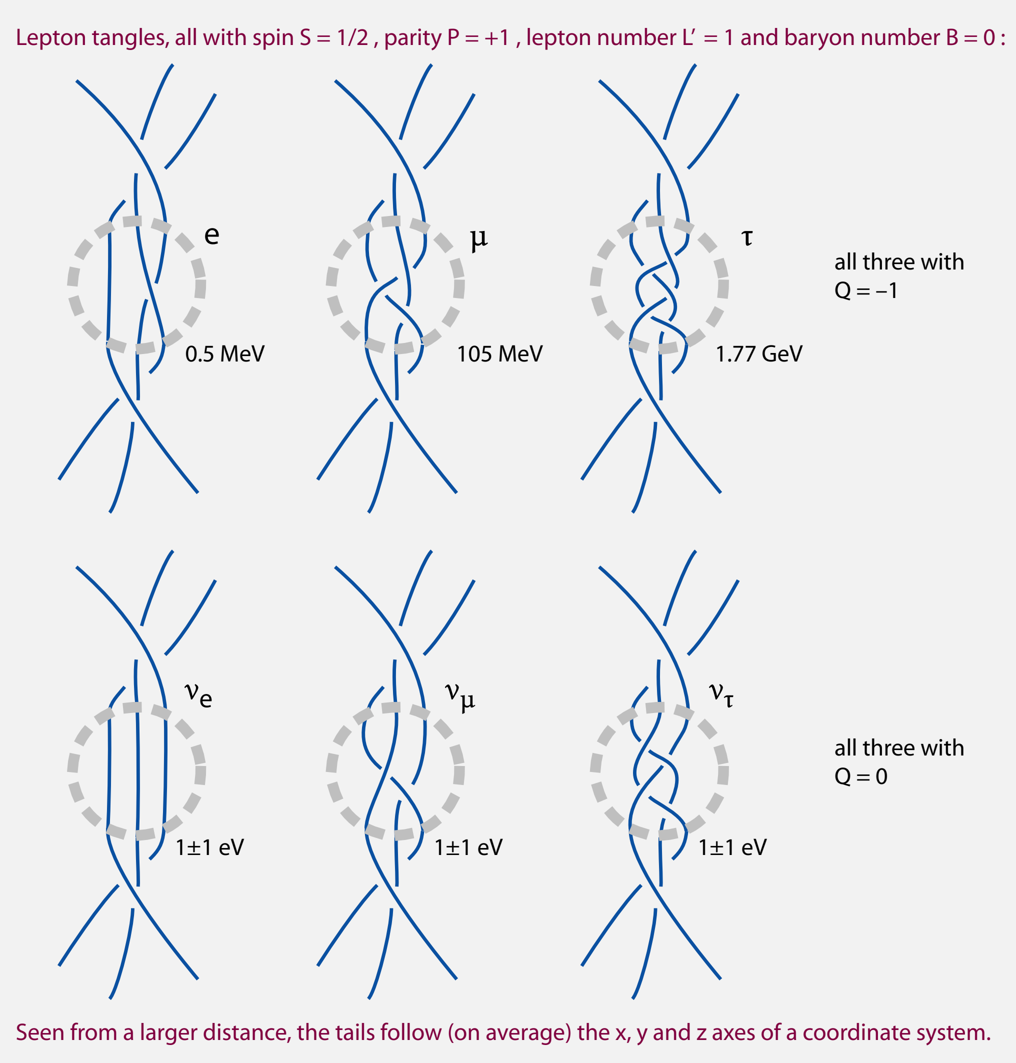

FIGURE 93 The simplest tangles of the leptons, with the experimental mass values. Antileptons are mirror tangles.

1/2. Each lepton thus follows the Dirac equation.
— Each lepton has weak charge.
— Charged leptons and antileptons differ. Each has two possible chiralities.
— Three of the tangles are topologically chiral, thus electrically charged, and three other tangles are topologically achiral, thus uncharged.
— The spatial parity *P* of the charged lepton tangles is opposite to that of their antiparticles.
— Being made of three strands, lepton tangles have vanishing colour charge and vanishing baryon number.
— In contrast to quarks, lepton tangles can be inserted in the vacuum using a localized, i.e., finite amount of energy and are thus predicted to exist as free particles.

— The three types of lepton (flavour) numbers can be assigned as usual; the lepton numbers are conserved in reactions, apart for neutrino mixing effects, as we will see below.
— The strand model predicts that the electron, the charged tangle with the lowest mass, is stable, as there is no way for it to decay and conserve charge and spin. The other two generations are predicted to be unstable, due to weak decays that simplify their topology.
— The three generations are reproduced by the strand model, as every more complicated braid can be seen as equivalent to one of the first six braids, with the same braiding argument that limits the number of quarks.
— There is a natural mapping between the six quarks and the six leptons. It appears when the final bend of the 'longer' quark strand is extended to the border of space, thus transforming a two-stranded quark braid into a three-stranded lepton braid. Thus we get three common generations for quarks and leptons.
— The neutrino strands differ by tail braiding; the strand model thus predicts that the weak interaction mixes neutrinos.
— All lepton tangles differ from each other. Thus the mass values are different for each lepton.
— Due to the small amount of tangling, the strand model predicts that the masses of the leptons are much smaller than those of the W and Z boson. This is indeed observed. (This also suggests a relation between the mass and the total curvature of a tight tangle.)
— The simplest tangle for the electron neutrino also suggests that the mass values for the electron neutrino is naturally small, as its tangle is almost not tangled.
— The strand model predicts that lepton masses increase with the generation number. Since the neutrino masses are not precisely known, this prediction cannot yet be checked.
— Neutrinos and antineutrinos are both massive and differ from each other. If the tangle of the electron neutrino is correct, the electron neutrino of opposite chirality is expected to be seen only rarely – as is observed.

In summary, tangles of three strands have precisely the quantum numbers and most properties of leptons. In particular, the strand model predicts exactly three generations of leptons, and predicts that all leptons have mass.

Ref. 224 This implies that searches for the neutrino-less double beta decay should yield negative results, that the magnetic moments of the neutrinos should have the exceedingly Ref. 225 small values predicted by the standard model of particle physics, and that rare muon and other decays should occur at the small rates predicted by the standard model.

Open issue: are the lepton tangles correct?

The argument that leads to the lepton tangles is vague. The tangle assignments might need corrections. There are two issues.

First, there is an aesthetic issue: in most particle tangles, the electric charge unit is Page 390 given by *three* crossings of the same sign. It seems odd that leptons should form an exception.

Secondly, the candidate tangles suggest that the muon neutrino is more massive than

FIGURE 94 A candidate tangle for the Higgs boson in the strand model: the open version (left) and the corresponding closed version (right). For the left version, the tails approach the six coordinate axes at infinity.

the electron. Most probably therefore, the tangles need amends. Can you improve the situation, either by finding better tangles or by finding better arguments? Challenge 194 ny

The Higgs boson – the mistaken section from 2009

The existence of the Higgs boson is predicted from the standard model of elementary particle physics using two arguments. First of all, the Higgs boson prevents unitarity violation in longitudinal W–W and Z–Z boson scattering. Secondly, the Higgs boson confirms the symmetry breaking mechanism of SU(2) and the related mass generation mechanism of fermions. Quantum field theory predicts that the Higgs boson has spin 0, has no electric or strong charge, and has positive C and P parity. In other words, the Higgs boson is predicted to have, apart from its weak charge, the same quantum numbers as the vacuum.

In the strand model, there seems to be only one possible candidate tangle for the Higgs boson, shown on the left of Figure 94. The tangle has positive C and P parity, and has vanishing electric and strong charge. The tangle also corresponds to the tangle added by the leather trick; it thus could be seen to visualize how the Higgs boson gives mass to the quarks and leptons. However, there are two issues with this candidate. First, the tangle is a deformed, higher-order version of the electron neutrino tangle. Secondly, the spin value is not 0. In fact, there is no way at all to construct a spin-0 tangle in the strand model. These issues lead us to reconsider the arguments for the existence of the Higgs boson altogether.

Page 253 We have seen that the strand model proposes a clear mechanism for mass generation:

▷ *Mass* is due to strand braiding.

This mechanism, due to the weak interaction, explains the W and Z boson mass ratio, Page 360 as we will see below. The leather trick that explains fermion masses can be seen as the addition of a sixfold tail braiding. In particular, the rarity of the braiding process explains why particle masses are so much smaller than the Planck mass. In short, the strand model explains mass *without* a Higgs boson.

If the Higgs boson does not exist, how is the unitarity of longitudinal W and Z boson scattering maintained? The strand model states that interactions of tangles in particle collisions are described by deformations of tangles. Tangle deformations in turn are de- Page 210 scribed by unitary operators. Therefore, the strand model predicts that unitarity is never

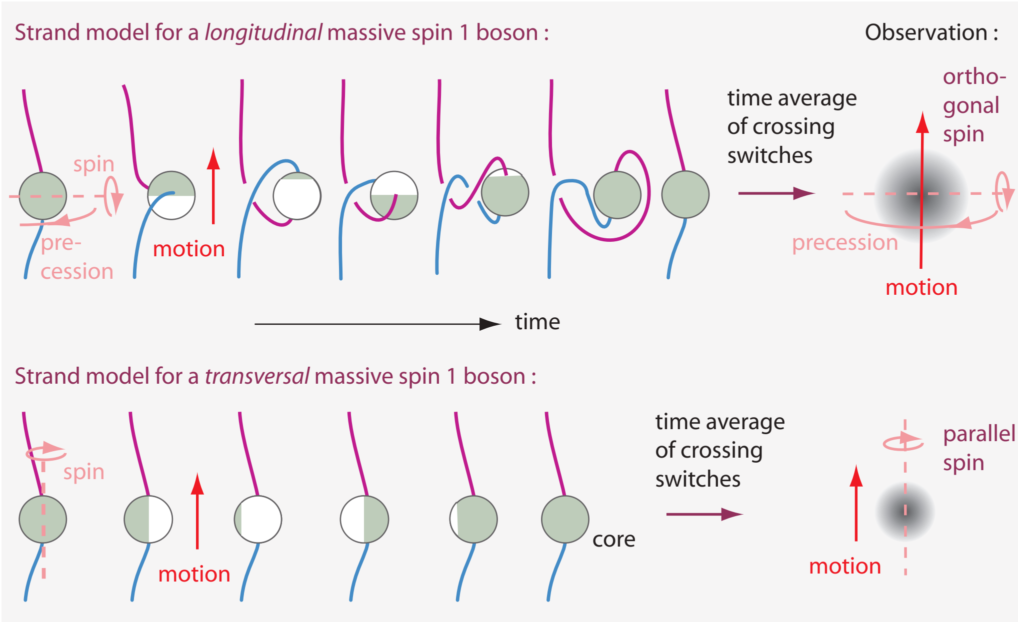

FIGURE 95 In the strand model, transverse and longitudinal W and Z bosons differ. (Note added in 2012: this statement is mistaken.)

violated in nature. In particular, the strand model automatically predicts that the scattering of longitudinal W or Z bosons does *not* violate unitarity.

In other terms, the strand model predicts that the conventional argument about unitarity violation, which requires a Higgs boson, must be wrong. How can this be? There are Ref. 226 at least two loopholes available in the research literature, and the strand model realizes them both.

The first known loophole is the appearance of non-perturbative effects. It is known for a long time that non-perturbative effects can mimic the existence of a Higgs boson in usual, perturbative approximations. In this case, the standard model could remain valid Ref. 227 at high energy without the Higgs sector. This type of electroweak symmetry breaking would lead to longitudinal W and Z scattering that does not violate unitarity.

The other loophole in the unitarity argument appears when we explore the details of the longitudinal scattering process. In the strand model, longitudinal and transverse W or Z bosons are modelled as shown in Figure 95. For longitudinal bosons, spin and its precession leads to a different situation than transversal bosons: longitudinal bosons are *more delocalized* than transversal bosons. This is not the case for fermions, where the belt trick leads to the *same* delocalization for longitudinal and transverse polarization. Interestingly, it is also known for a long time that different delocalization for longitudinal and transversal bosons *maintains* scattering unitarity, and that in the case of delocaliz- Ref. 228 ation the conventional argument for the necessity of the Higgs boson is wrong. These are well-known consequences of the so-called *non-local regularization* in quantum field theory. The strand model thus provides a specific model for this non-locality, and at the same time explains why it *only* appears for longitudinal W and Z bosons.

The issue of different scattering behaviour for longitudinal and transverse weak bosons also raises the question whether the mass of the longitudinal and the transversal bosons are precisely equal. The possibility, triggered by Figure 95, might seem appealing at first sight in order to solve the unitarity problem. However, the strand model forbids such a mass difference. In the strand model, mass is due to tangle fluctuations, but does not depend on spin direction.

In other words, the strand model predicts that the scattering of longitudinal W and Z bosons is the first system that will show effects specific to the strand model. Such precision scattering experiments might be possible at the Large Hadron Collider in Geneva. These experiments will allow checking the *non-perturbative effects* and the *regularization effects* predicted by the strand model. For example, the strand model predicts that the wave function of a longitudinal and a transversally polarized W or Z boson of the same energy differ in cross section.

In summary, the strand model predicts well-behaved scattering amplitudes for longitudinal W and Z boson scattering in the TeV region, together with the absence of the Higgs boson.* The strand model explains mass generation and lack of unitarity violations in longitudinal W or Z boson scattering as consequences of tail braiding, i.e., as non-perturbative and non-local effects, and not as consequences of an elementary spin-0 Higgs boson. The forthcoming experiments at the Large Hadron Collider in Geneva will test this prediction.

The Higgs boson – the corrected section of 2012

In July 2012, CERN researchers from two different experiments announced the observation of a new neutral boson with a mass of 125 GeV. Additional data analysis showed that the boson has spin 0 and positive parity. All experimental checks confirm that the boson behaves like the Higgs boson predicted in 1963 by Peter Higgs and a number of other researchers.

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The results lead to question several statements made in 2009 in the previous section.

- Is the tangle on the left-hand side of Figure 94 really a higher order version of the electron neutrino? It turns out that this statement is wrong: in contrast to the tangle of the neutrino, the tangle of Figure 94 is not twisted.
- Does the tangle of Figure 96 have spin 1/2 or spin 0? As mentioned already in 2009, an effective spin 0 might be possible, in a similar way that it is possible for spin-0 mesons. Spin 0 behaviour might appear because the tangle can be oriented in different directions or because of the Borromean property: no two strands have more crossings than two vacuum strands; the time average of these situations has the same symmetry as the vacuum, and thus implies spin 0.
- Does the tangle of Figure 96 have the correct, positive, C and P values expected for a Higgs boson? It seems so.
- Is the mentioned non-locality effect for W and Z bosons real? If the effect were real, it should also appear for other spin-1 particles. In the strand model, mass values should

* If the arguments against the Higgs boson turn out to be wrong, then the strand model might be saved with a dirty trick: we could argue that the tangle on the left-hand side of Figure 94 might effectively have spin 0. In this case, the ropelength of the Borromean rings, 29.03, together with the ropelengths of the weak bosons, lead to a Higgs mass prediction, to first order, in the range from $(29.03/10.1)^{1/3} \cdot 80.4\,\text{GeV} = 114\,\text{GeV}$ to $(29.03/13.7)^{1/3} \cdot 91.2\,\text{GeV} = 117\,\text{GeV}$, plus or minus a few per cent.

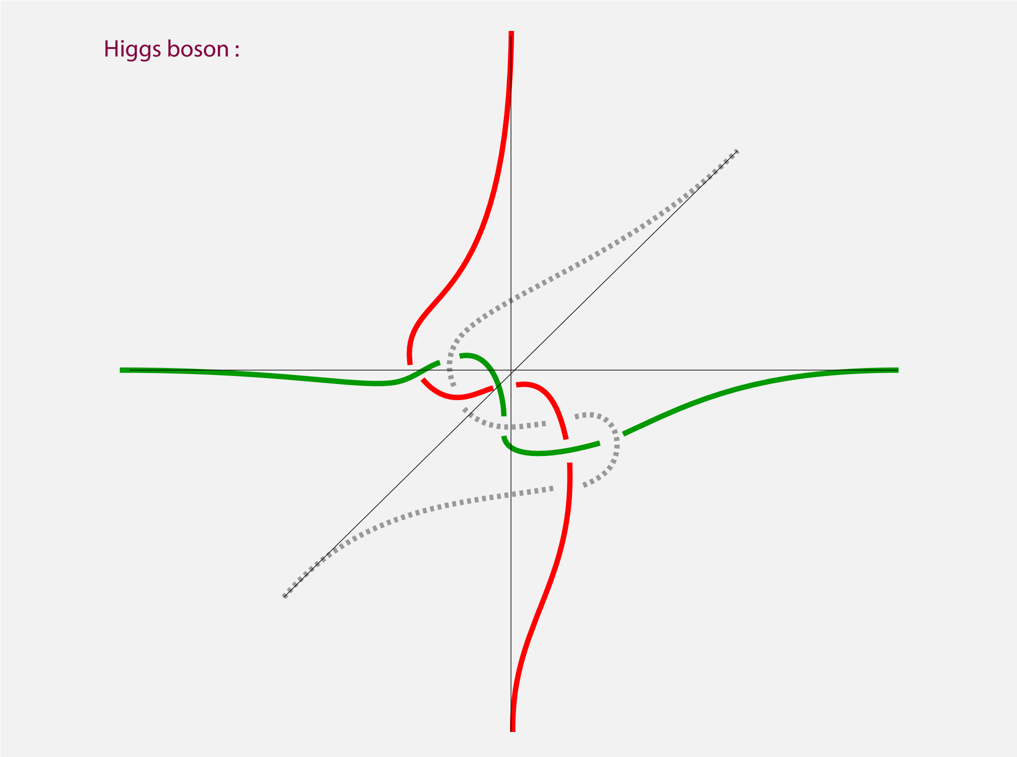

FIGURE 96 The tangle of the Higgs boson in the strand model. Spin 0 appears because the braid can be oriented in different directions, so that the time average has spherical symmetry. The tangle has 9 crossings: 3 crossings appear already in the vacuum configuration of three strands, and the additional 6 crossings (see Figure 94) are due to the Higgs boson.

not depend on spin orientation, but only on tangle core topology. The statements made in 2009 on delocalization and longitudinal scattering seem wrong in retrospect.

— Would the Higgs boson tangle assignment of Figure 96 be testable? Yes; any tangle assignment must yield the observed mass value and the observed branching ratios and decay rates. This is a subject of research. But already at the qualitative level, the proposed tangle structure of the Higgs boson suggests decays into leptons that are similar to those observed at CERN.
— Is the tangle of Figure 96 elementary? Yes.
— Are there other possible Higgs boson tangles? This issue is open. The braid structure seems the most appealing structure, as it embodies the effect of tail braiding, an effect that is important for the appearance of mass. Page 370
— Are knots and links, i.e., closed tangles, really forbidden? The discussion about the Higgs boson concerns the open tangle shown in Figure 96, not the Borromean link shown on the right-hand side of Figure 94. So far, there is no evidence for closed tangles in the strand model. Such evidence would mean a departure from the idea that nature is a single strand.
— Does the Higgs boson issue put into question the strand model as a whole? First of all, SU(2) breaking is unaffected. Secondly, a mistaken tangle–particle assignment can be

accommodated in the strand model; new forces or symmetries cannot. Therefore the strand model is not put into question.

— Could several, possibly charged, Higgs bosons exist? No such tangles seem possible – as long as a tangle with *two* Figure 96 Higgs cores in sequence is not a separate particle.
— Has some other strand model effect been overlooked? Could other elementary or composed particles exist? For example, the structure of the Higgs boson might be seen to suggest that lepton families reappear (roughly) every 125 GeV. Is that the case? The issue is not completely settled. It seems more probable that those higher tangles simply yield corrections to the Higgs mass.

In short, the existence of the standard model Higgs boson seems compatible with the strand model. The 2009 mistake about the Higgs also shows that the exploration of the strand model is not yet complete. In any case, the strand model has not been falsified by the discovery of the Higgs boson.

Page 331 Assuming that the Higgs tangle shown in Figure 96 is correct, we have an intuitive proposal for the mechanism that produces mass, namely *tail braiding*. The proposed Higgs tangle also allows a number of experimental predictions.

2012 predictions about the Higgs

— The Higgs tangle implies a Higgs boson with vanishing charge, positive parity, being elementary – as is observed.
— The Higgs tangle allows us to estimate the Higgs/Z mass ratio. Using the new, unknotted, tangle model for the W and Z bosons, the estimates are in the region of the observed values. Improving the estimates is still subject of research.
— The Higgs tangle and the strand model imply that the standard model is correct up to Planck energy, and that the Higgs mass value should reflect this. The observed Higgs mass of 125 GeV complies also with this expectation.
— Therefore, the strand model suggests that no deviations between the standard model and data should ever be observed in any experiment.
— The strand model again and consistently predicts the lack of supersymmetry.
— In the case that several Higgs bosons exist or that the braided Higgs tangle does not apply, the strand model is in trouble.
— In the case that effects, particles or interactions beyond the standard model are observed, the strand model is in trouble.

Quark-antiquark mesons

In the strand model, all three-stranded tangles apart from the leptons, as well as all four-stranded tangles represent *composite* particles. The first example are *mesons*.

In the strand model, rational tangles of three strands are quark-antiquark mesons with spin 0. The quark tangles yield a simple model of these *pseudoscalar* mesons, shown on the left-hand sides of Figure 97, Figure 99 and Figure 100. The right-hand sides of the figures show *vector* mesons, thus with spin 1, that consist of *four* strands. All tangles are rational. Inside mesons, quarks and antiquarks 'bond' at three spots that form a triangle oriented perpendicularly to the bond direction and to the paper plane. To increase clar-

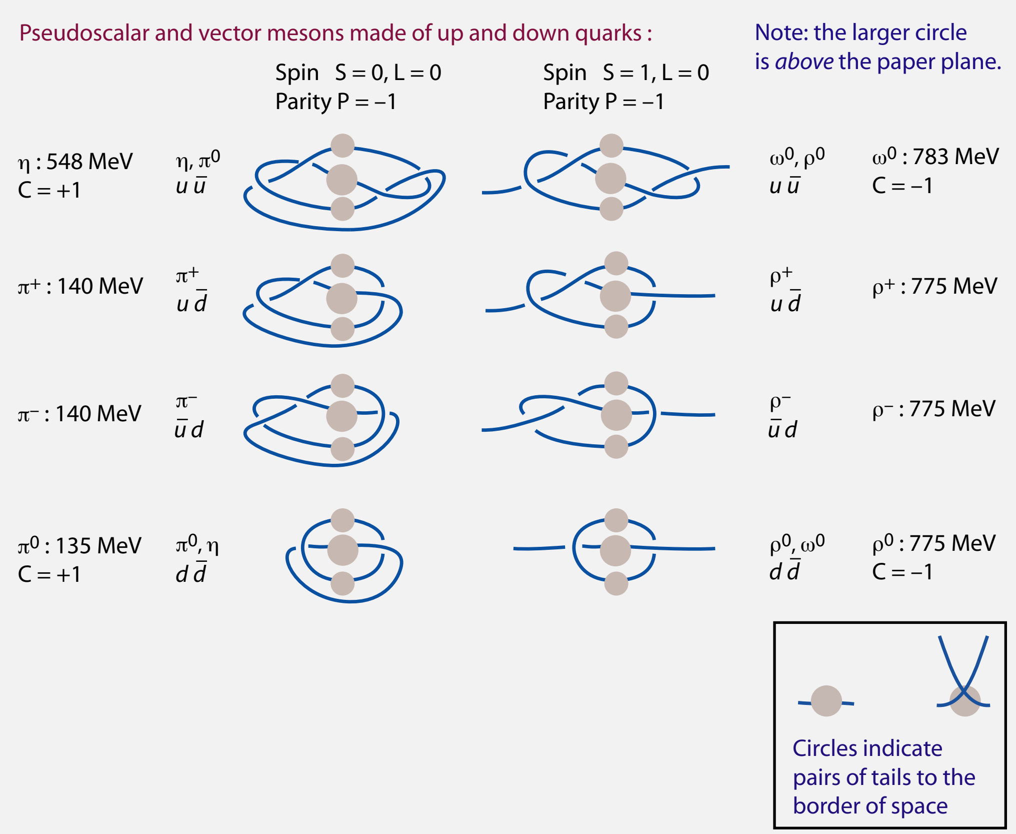

FIGURE 97 The simplest strand models for the light pseudoscalar and vector mesons (circles indicate crossed tail pairs to the border of space), with the observed mass values.

ity, the 'bonds' are drawn as circles in the figures; however, they consist of two crossed (linked) tails of the involved strands that reach the border of space, as shown in Figure 98.

With this construction, mesons made of two quarks are only possible for the type $\overline{q}\,q$. Ref. 229 Other combinations, such as $q\,q$ or $\overline{q}\,\overline{q}$, turn out to be unlinked. We note directly that Ref. 230 this model of mesons resembles the original string model of hadrons from 1973, but also Ref. 231 the Lund string model and the recent QCD string model.

To compare the meson structures with experimental data, we explore the resulting quantum numbers. As in quantum field theory, also in the strand model the parity of a particle is the product of the intrinsic parities and of wave function parity. The states with orbital angular momentum $L = 0$ are the lowest states. Experimentally, the lightest mesons have quantum numbers $J^{PC} = 0^{-+}$, and thus are pseudoscalars, or have $J^{PC} = 1^{--}$, and thus are vector mesons. The strand model reproduces these observed quantum numbers. (We note that the spin of any composite particle, such as a meson, is low-energy quantity; to determine it from the composite tangle, the tails producing the bonds – drawn as circles in the figures – must be neglected. As a result, the low-energy spin of mesons and of baryons is correctly reproduced by the strand model.)

In the strand model, the meson states are colour-neutral, or 'white', by construction, because the quark and the antiquark, in all orientations, always have opposite colours

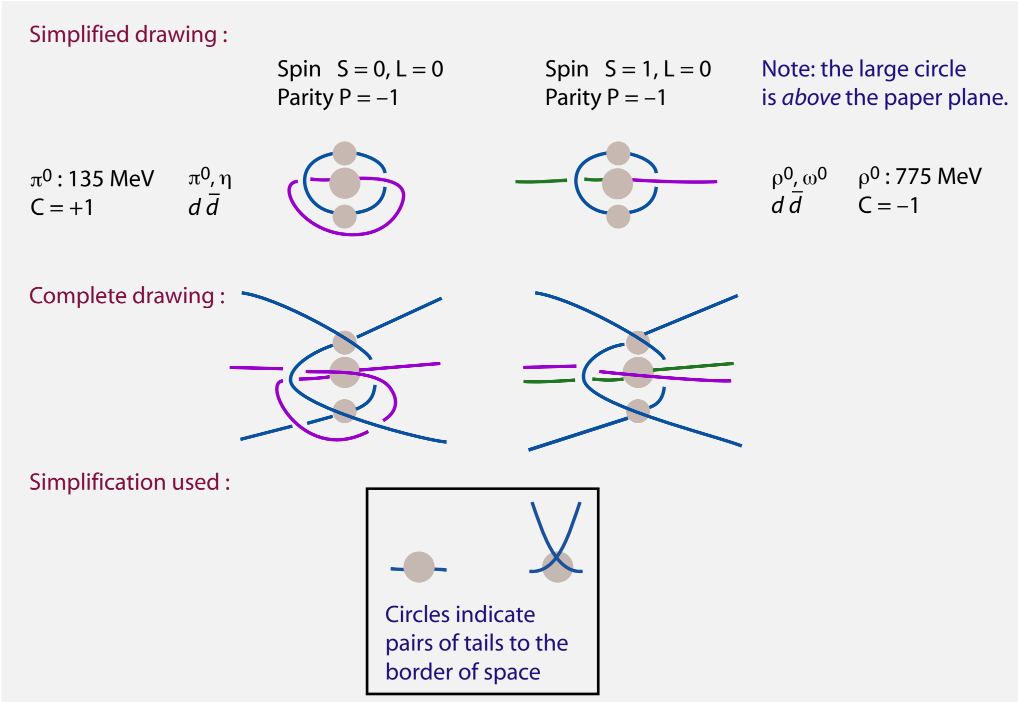

FIGURE 98 The meaning of the circles used in the tangle graphs of mesons and baryons.

that add up to white.

In the strand model, the electric charge is an integer for all mesons. Chiral tangles are charged, achiral tangles uncharged. The charge values deduced from the strand model thus reproduce the observed ones.

In experiments, no mesons with quantum numbers 0^{--}, 0^{+-}, or 1^{-+} are observed. Also this observation is reproduced by the quark tangles, as is easily checked by direct inspection. The strand model thus reproduces the very argument that once was central to the acceptance of the quark model itself.

It is important to realize that in the strand model, each meson is represented by a *tangle family* consisting of *several* tangle structures. This has three reasons. First, the 'circles' can be combined in different ways. For example, both the $u\,\overline{u}$ and the $d\,\overline{d}$ have as alternate structure a line plus a ring. This common structure is seen as the underlying reason that these two quark structures *mix*, as is indeed observed. (The same structure is also possible for $s\,\overline{s}$, and indeed, a full description of these mesons must include mixing with this state as well.) The second reason that mesons have several structures are the mentioned, more complicated braid structures possible for each quark, namely with 6, 12, etc. additional braid crossings. The third reason for additional tangle structures is the occurrence of higher-order Feynman diagrams of the weak interaction, which add yet another group of more complicated topologies that also belong to each meson.

In short, the mesons structures of Figure 97, Figure 99 and Figure 100 are only the *simplest* tangles for each meson. Nevertheless, all tangles, both the simplest and the more complicated meson tangles, reproduce spin values, parities, and all the other quantum

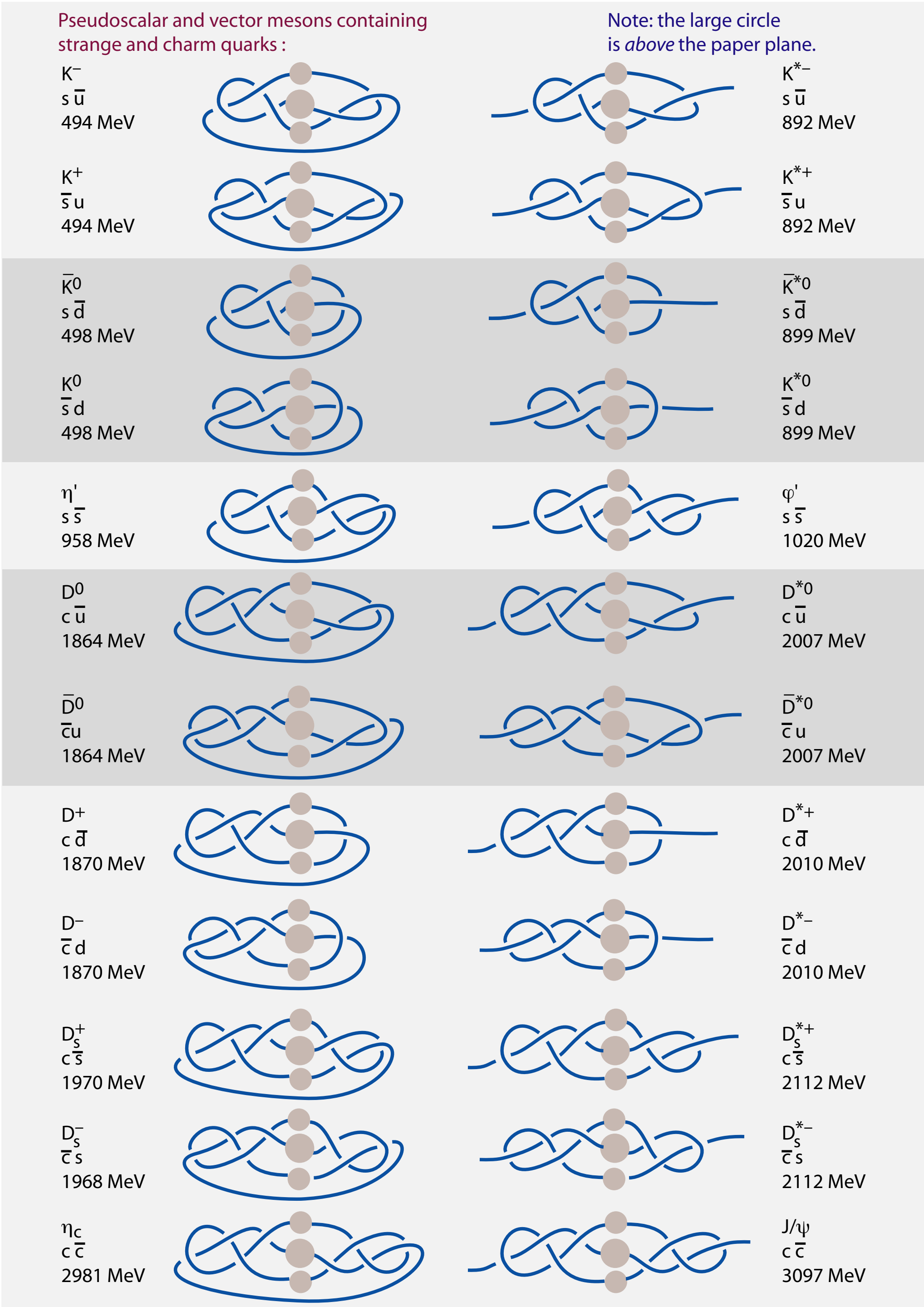

FIGURE 99 The simplest strand models for strange and charmed mesons with vanishing orbital angular momentum. Mesons on the left side have spin 0 and negative parity; mesons on the right side have spin 1 and also negative parity. Circles indicate crossed tail pairs to the border of space; grey boxes indicate tangles that mix with their antiparticles and which are thus predicted to show CP violation.

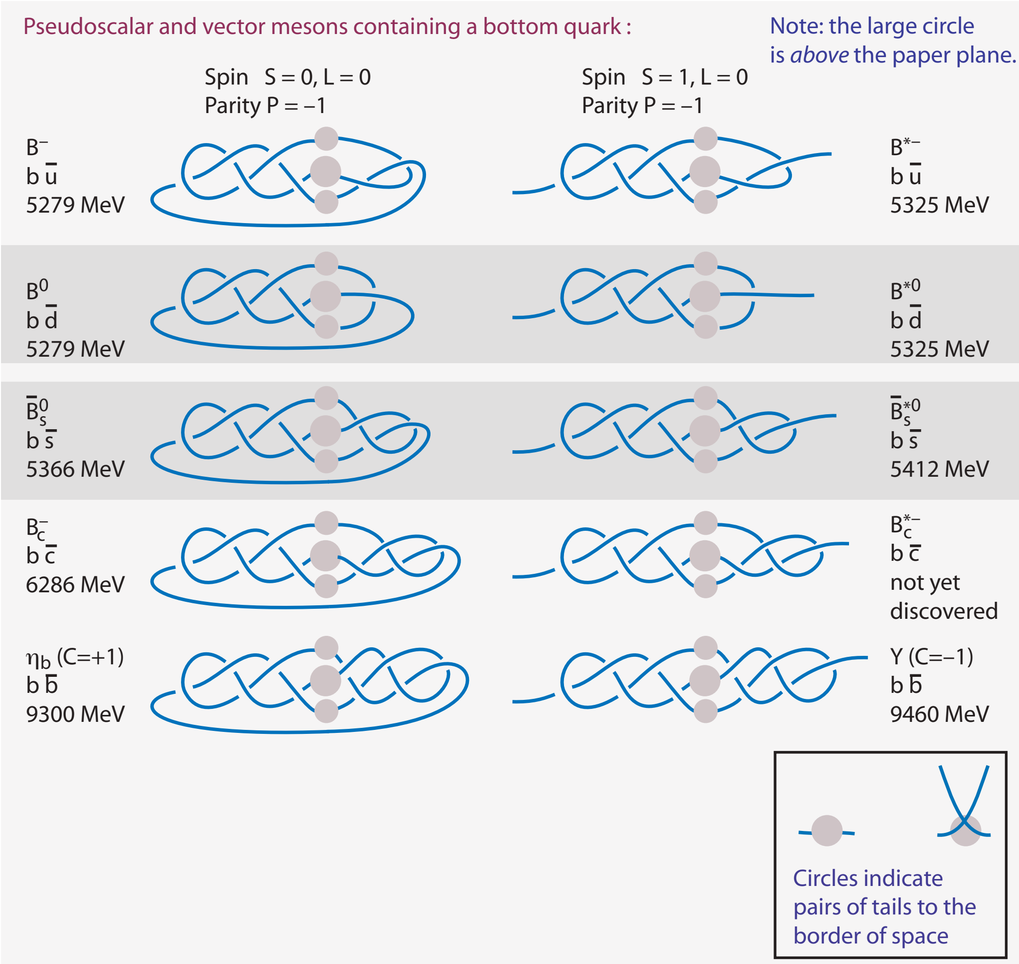

FIGURE 100 The simplest strand models for some heavy pseudoscalar and vector mesons, together with their experimental mass values. Antiparticles are not drawn; their tangles are mirrors of the particle tangles. Circles indicate crossed tail pairs to the border of space; grey boxes indicate tangles that mix with their antiparticles and which are thus predicted to show CP violation.

numbers of mesons. Indeed, in the strand model, the more complicated tangles automatically share the quantum numbers of the simplest one.

Meson form factors

The strand model also predicts directly that all mesons from Figure 97, Figure 99 and Figure 100, in fact all mesons with vanishing orbital momentum, are *prolate*. This (unsurprising) result is agreement with observations. Mesons with non-vanishing orbital momentum are also predicted to be prolate. This latter prediction about meson shapes is made also by all other meson models, but has not yet been checked by experiment. Ref. 232

There is another way to put what we have found so far. The strand model makes the following prediction: When the meson tangles are averaged over time, the crossing densities reproduce the measured spatial, quark flavour, spin and colour part of the meson wave functions. This prediction can be checked against measured form factors and against lattice QCD calculations.

MESON MASSES, EXCITED MESONS AND QUARK CONFINEMENT

The strand model also allows us to understand meson masses. We recall that a *topologically complicated* tangle implies a *large* mass. With this relation, Figure 97 predicts that the π^0, η and $\pi^{+/-}$ have different masses and follow the observed meson mass sequence $m(\pi^0) < m(\pi^{+/-}) < m(\eta)$. The other mass sequences can be checked with the help of Figure 97, Figure 99 and Figure 100; there are no contradictions with observations. However, there is one limit case: the strand model predicts different masses for the ρ^0, ω, and $\rho^{+/-}$. So far, observations only partly confirm the prediction. Recent precision experiments seem to suggest that ρ^0 and $\rho^{+/-}$ have different mass; this result has
Ref. 233 not been confirmed yet.

More precise mass determinations will be possible with numerical calculations. This
Page 358 will be explored in more detail later on. In any case, the strand model for mesons suggests that the quark masses are not so important for the determination of meson masses, whereas the details of the quark-antiquark bond are. Indeed, the light meson and baryon masses are much higher than the masses of the constituent quarks.

The relative unimportance of quark masses for many meson masses is also confirmed for the case of *excited* mesons, i.e., for mesons with orbital angular momentum L. It is well known that mesons of non-vanishing orbital angular momentum can be grouped into sets which have the same quark content, but different total angular momentum $J = L + S$. These families are observed to follow a well-known relation between total angular momentum J and mass m, called *Regge trajectories*:

$$J = \alpha_0 + \alpha_1 m^2 \qquad (195)$$

Ref. 234 with an (almost) constant factor α_1 for all mesons, about 0.9 GeV/fm. These relations, the famous *Regge trajectories*, are explained in quantum chromodynamics as deriving from the linear increase with distance of the effective potential between quarks, thus from the properties of the relativistic harmonic oscillator. The linear potential itself is usually seen as a consequence of a fluxtube-like bond between quarks.

In the strand model, the fluxtube-like bond between the quarks is built-in automatically, as shown in Figure 101. All mesons have three connecting 'bonds' and these three bonds can be seen as forming one common string tube. In the simplified drawings, the bond or string tube is the region containing the circles. In orbitally excited mesons, the three bonds are expected to lengthen and thus to produce additional crossing changes, thus additional effective mass. The strand model also suggests a *linear* relation. Since the mechanism is expected to be similar for all mesons, which all have three bonding circles, the strand model predicts the *same* slope for all meson (and baryon) Regge trajectories. This is indeed observed.

In summary, the strand model reproduces meson mass sequences and quark confinement in its general properties.

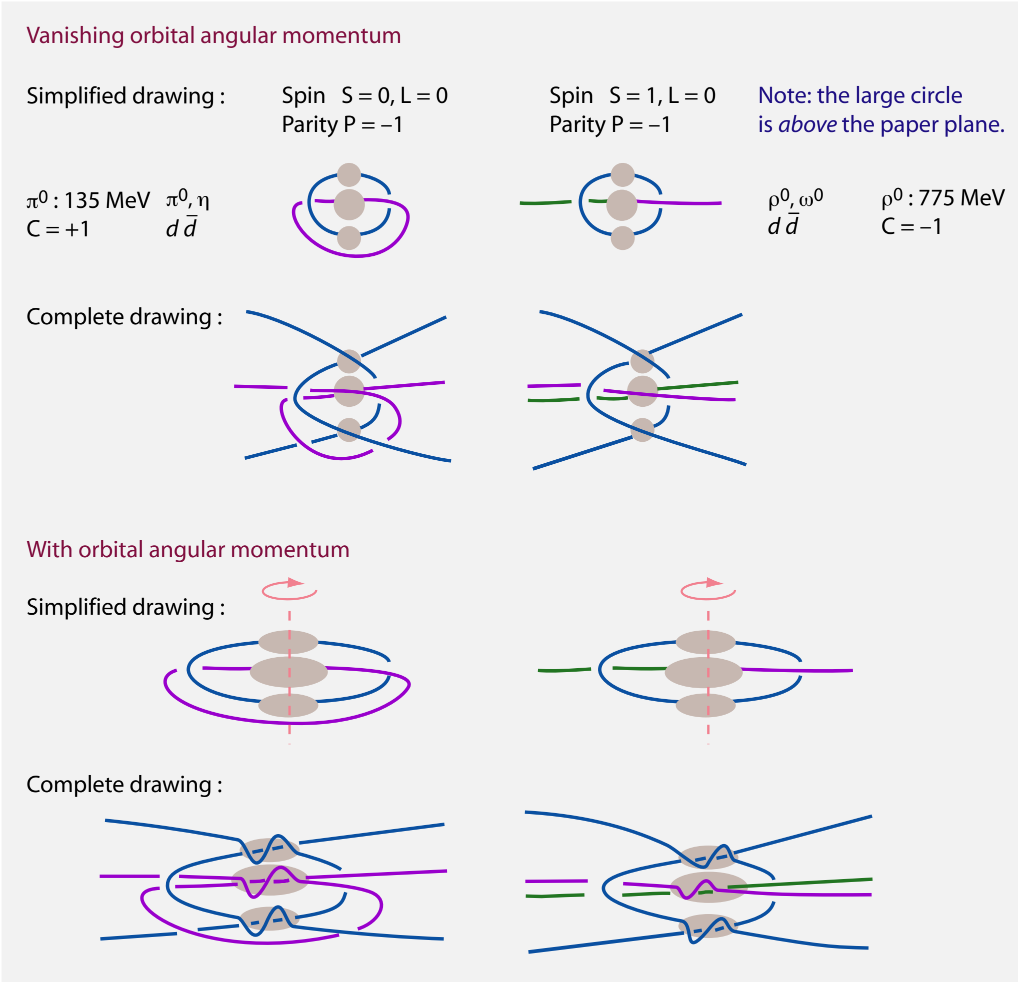

FIGURE 101 The strand model for mesons without (top) and with (bottom) orbital angular momentum.

CP violation in mesons

Ref. 235 In the weak interaction, the product CP of C and P parity is usually conserved. However, rare exceptions are observed for the decay of the K^0 meson and in various processes that involve the B^0 and B^0_s mesons. In each of these exceptions, the meson is found to Ref. 233 mix with its own antiparticle. CP violation is essential to explain the matter–antimatter asymmetry of the universe.

The strand model allows us to deduce whether the mixing of a meson with its own antiparticle is possible or not. As expected, only neutral mesons are candidates for such mixing, because of charge conservation. In the strand model, particle–antiparticle mixing is possible whenever the transition from a neutral meson to its antiparticle is possible in *two* ways: by taking the mirror of the meson tangle or by shifting the position of the binding strands. All mesons for which this is possible are shown in grey boxes in Figure 97, Figure 99 and Figure 100. The strand model also makes it clear that such mixing requires shifting of the bonds; this is a low-probability process that is due to the weak interaction. The strand model thus predicts that the weak interaction violates CP invariance in mesons that mix with their antiparticles.

Since the spin 1 mesons decay strongly and thus do not live long enough, the small effect of CP violation is de facto only observed in pseudoscalar, spin-0 mesons. The strand model thus predicts observable mixings and CP violation for the mesons pairs $K^0 - \overline{K}^0$, $D^0 - \overline{D}^0$, $B^0 - \overline{B}^0$, $B^0_s - \overline{B}^0_s$. The prediction by the strand model corresponds precisely to those systems for which CP violation is actually observed. (CP violation in D mesons was finally discovered at CERN in 2011, after it was predicted both by the standard model and the strand model, in earlier editions of this volume.)

Ref. 233

In the strand model, meson–antimeson mixing is possible because the various quarks are braided strands. Because of this braid structure, the existence of meson–antimeson mixing is a consequence of the existence of three quark generations. The meson structures also make it clear that such mixings would not be possible if there were no third quark generation. The strand model thus reproduces the usual explanation of CP violation as the result of three quark generations.

For the strong and the electromagnetic interaction, the strand model predicts that there is no mixing and no CP violation, because gluons and photons do not change particle topology. Therefore, the strand model suggests the absence of axions. The lack of a suitable tangle for axions, shown later on, then turns this suggestions into a prediction.

Page 353

In summary, the existence of CP violation in the weak interactions and the lack of CP violation in the strong interaction are natural consequences of the strand model.

Other three-stranded tangles

In the strand model, the omitted complicated tangles made of three strands are either higher-order propagating versions of the tangles just presented or composites of one-stranded or two-stranded particles.

The three-strand analog of the graviton – three parallel, but twisted strands – is not an elementary particle, but a composed structure.

Challenge 195 s

Spin and three-stranded particles

Why do three strands sometimes form a spin 0 particle, such as the elementary Higgs boson, sometimes a spin 1/2 particle, such as the elementary electron, and sometimes a spin 1 particle, such as a composed meson? The answer depends on how the strands are free to move against each other.

The Higgs tangle appears through tangling of vacuum strands, and inherits the zero spin of vacuum. The W and Z tangles have a special property: two strands can rotate around the third; this makes them bosons as well, but of spin 1. Fermion tangles have neither property; their core can only rotate through the belt trick; thus they are fermions.

Summary on three-stranded tangles

A number of elementary particles are made of trhee strands: The massive W and Z, the gluons, the leptons and the Higgs boson. Their tangles reproduce all their observed quantum numbers. The tangles also imply that neutrinos and anti-neutrinos differ, are massive, and are Dirac particles.

Page 330 The strand model (corrected in 2012) also predicts that, apart from the mentioned particles, no other elementary particle made of three strands exist in nature.

In the case of *composite* particles made of three strands, the strand model proposes tangles for all pseudoscalar mesons; the resulting quantum numbers and mass sequences match the observed values.

TANGLES OF FOUR AND MORE STRANDS

If we add one or more strand to a three-strand tangle, no additional class of tangles appears. The tangle classes remain the same as in the three-strand case. In other words, *no additional elementary particles* arise in the strand model. To show this, we start our exploration with the *rational* tangles.

We saw above that the rational tangles made of four strands represent the vector mesons. We have already explored them together with the scalar mesons. But certain more complicated rational tangles are also important in nature, as we consist of them.

Baryons

In the strand model, rational tangles made of five or six strands are baryons. The quark tangles of the strand model yield the tangles for baryons in a natural way, as Figure 102 shows. Again, not all quark combinations are possible. First of all, quark tangles do not allow mixed $q\,q\,\overline{q}$ or $q\,\overline{q}\,\overline{q}$ structures, but only $q\,q\,q$ or $\overline{q}\,\overline{q}\,\overline{q}$ structures. In addition, the tangles do not allow (fully symmetric) spin 1/2 states for $u\,u\,u$ or $d\,d\,d$, but only spin 3/2 states. The model also naturally predicts that there are only two spin 1/2 baryons made of u and d quarks. All this corresponds to observation. The tangles for the simplest baryons are shown in Figure 102.

The electric charges of the baryons are reproduced. In particular, the tangle topologies imply that the proton has the same charge as the positron. Neutral baryons have topologically achiral structures; nevertheless, the neutron differs from its antiparticle, as can be deduced from Figure 102, through its three-dimensional shape. The Δ baryons
Page 390 have different electric charges, depending on their writhe.

Baryons are naturally colour-neutral, as observed. The model also shows that the baryon wave function usually cannot be factorized into a spin and quark part: the nucleons need *two* graphs to describe them, and tangle shapes play a role. Baryon parities are reproduced; the neutron and the antineutron differ. All this corresponds to known baryon behaviour. Also the observed baryon shapes (in other words, the baryon quadrupole
Ref. 232 moments) are reproduced by the tangle model.

The particle masses of proton and neutron differ, because their topologies differ. However, the topological difference is 'small', as seen in Figure 102, so the mass difference is small. The topological difference between the various Δ baryons is even smaller, and indeed, their mass difference is barely discernible in experiments.

The strand model naturally yields the baryon octet and decuplet, as shown in Figure 103 and Figure 104. In general, complicated baryon tangles have higher mass than simpler ones, as shown in the figures; this is also the case for the baryons, not illustrated here, that include other quarks. And like for mesons, baryon Regge trajectories are due

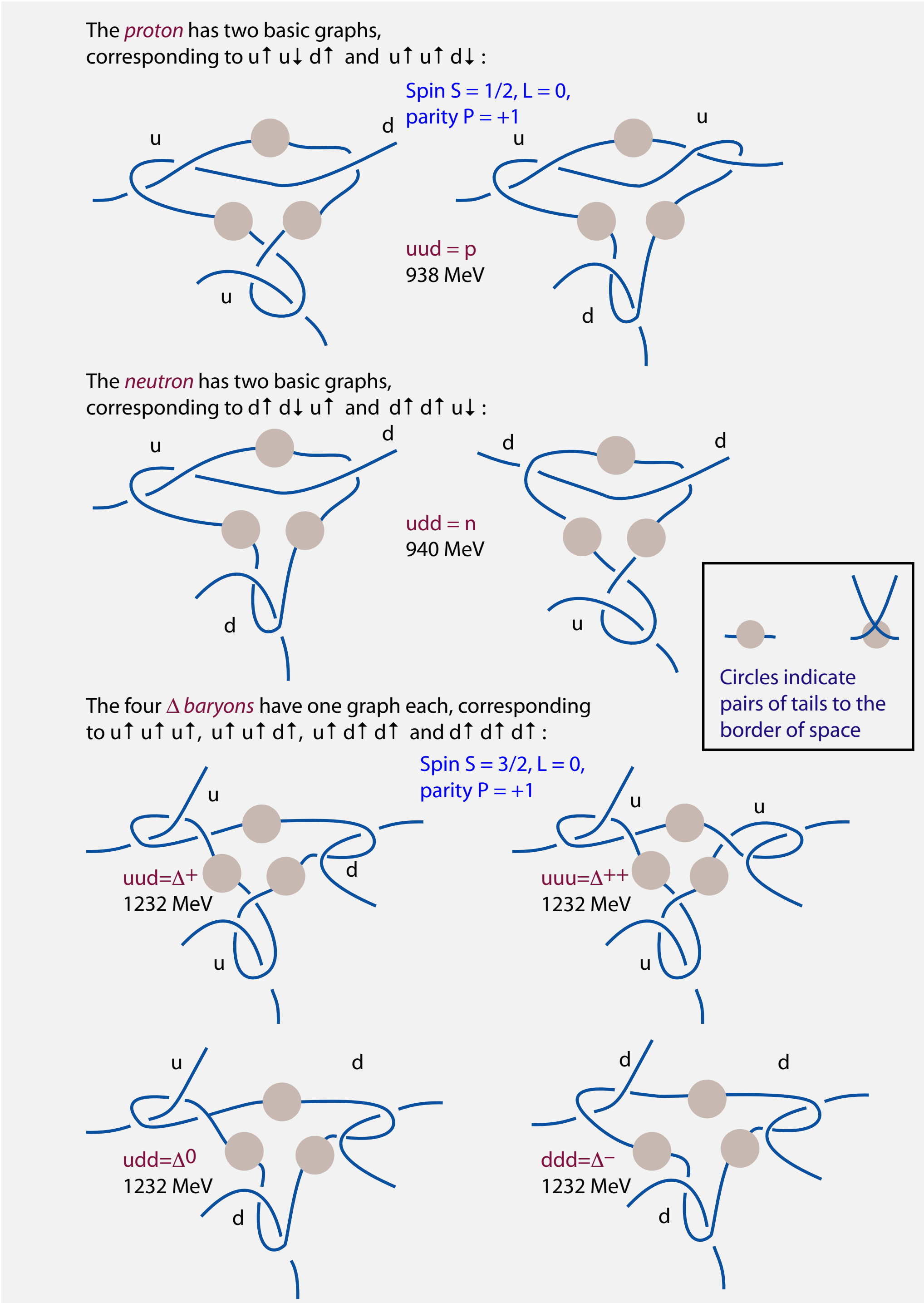

FIGURE 102 The simplest strand models for the lightest baryons made of up and down quarks (circles indicate linked tail pairs to the border of space), together with the measured mass values.

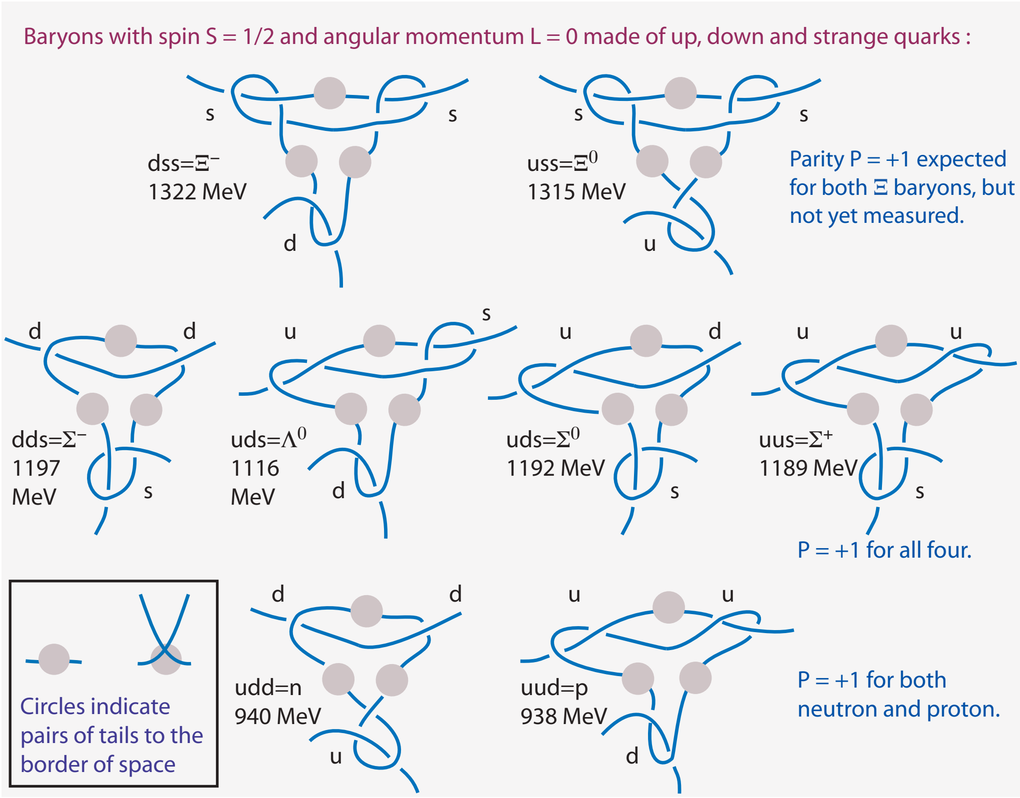

FIGURE 103 One tangle (only) for each baryon in the lowest J=L+S=1/2 baryon octet (circles indicate linked tail pairs to the border of space), together with the measured mass values.

to 'stretching' and tangling of the binding strands. Since the bonds to each quark are again (at most) three, the model qualitatively reproduces the observation that the Regge slope for all baryons is the same and is equal to that for mesons. We note that this also implies that the quark masses play only a minor role in the generation of hadron masses; this old result from QCD is thus reproduced by the strand model.

The arguments presented so far only reproduce mass sequences, not mass values. Actual hadron mass calculations are possible with the strand model: it is necessary to compute the number of crossing changes each tangle produces. There is a chance, but no certainty, that such calculations might be simpler to implement than those of lattice QCD.

Tetraquarks and exotic mesons

Ref. 220 Among the exotic mesons, tetraquarks are the most explored cases. It is now widely be-
Ref. 236 lieved that the low-mass scalar mesons are tetraquarks. In the strand model, tetraquarks are possible; an example is given in Figure 105. This is a six-stranded rational tangle. Spin, parities and mass sequences from the strand model seem to agree with observations. If the arrangement of Figure 105 would turn out to be typical, the tetraquark looks more like a bound pair of two mesons and not like a state in which all four quarks are bound in equal way to each other. On the other hand, a tetrahedral arrangement of quarks might

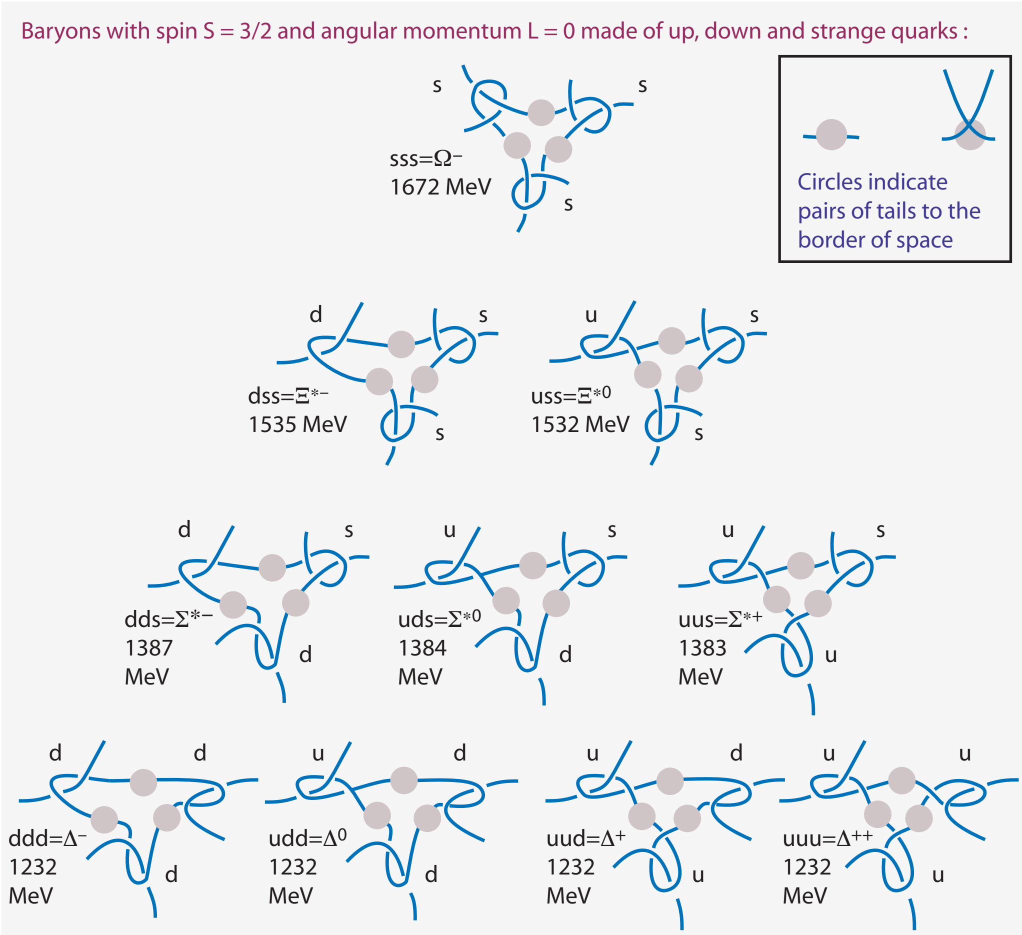

FIGURE 104 One tangle for each baryon in the lowest J=3/2 baryon decuplet (circles indicate linked tail pairs to the border of space), together with the measured mass values.

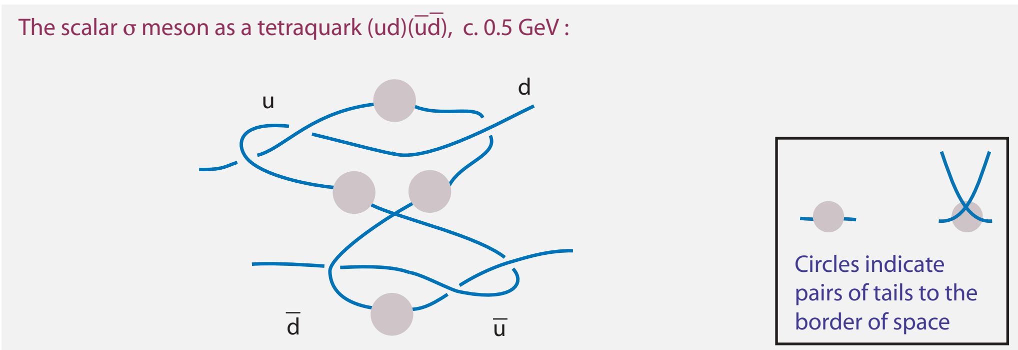

FIGURE 105 The strand model for a specific tetraquark (circles indicate linked tail pairs to the border of space).

also be possible. The details of this topic are left for future exploration.

The strand model makes an additional statement: knotted (hadronic) strings in quark–antiquark states are impossible. Such states have been proposed by Niemi. In the strand model, such states would not be separate mesons, but usual mesons with one or several added virtual weak vector bosons. This type of exotic mesons is therefore predicted not to exist.

Ref. 222

The situation for glueballs, which are another type of exotic mesons, has already been discussed above.

Page 344, page 339

Other tangles made of four or more strands

We do not need to explore other *prime* tangles or *locally knotted* tangles made of four or more strands. They are either not allowed or are higher-order versions of rational tangles, as explained already in the case of two and three strands. We also do not need to explore *separable* tangles. Separable tangles are composite of tangles with fewer strands.

One class of tangles remains to be discussed: *braided* tangles of four or more strands. Now, a higher-order perturbation of the weak interaction can always lead to the topological entanglement of some vacuum strand with a tangle of fewer strands. Braided tangles of four or more strands are thus higher-order propagating states of three-stranded leptons or hadrons.

We can also state this in another way. There are no tangles of four or more strands that are more tangled than the trivial tangle but less tangled than the lepton tangles. Therefore, no additional elementary particles are possible. In short, *the tangle model does not allow elementary particles with four or more strands.*

Glueballs

Ref. 219, Ref. 220

There is no observational evidence for glueballs yet, even though simulations of QCD on the lattice predict the existence of several such states in the 1.5 GeV/c^2 mass range. The lack of experimental confirmation is usually explained by the strong background noise in the reaction that produces glueballs, and by the expected strong mixing with mesons of similar quantum numbers. The experimental search for glueballs is still ongoing.

The often conjectured glueball could be made of two or three gluons. (The lowest-mass glueball is usually expected to be made of two gluons.) In the strand model, such structures would be tangles made of six or nine strands.

However, the masslessness of gluons and their spin do not seem to allow such a tangle. The argument is not watertight, however, and the issue is still subject of research.

Ref. 221

Ref. 222

Whatever the situation for glueballs might be, the strand model of gluons seems in contrast with the models of glueballs as knots that were proposed by Buniy and Kephart or by Niemi. These models are based on *closed* knots, not on tangles with tails. The strand model does not seem to allow real particles *of zero spin* that are composed of gluons. On the other hand, if closed knots were somehow possible in the strand model, they would imply the existence of glueballs.

In summary, the issue of glueballs is not settled; a definitive solution might even lead to additional checks of the strand model.

The mass gap problem and the Clay Mathematics Institute

Ref. 223 The Clay Mathematics Institute offers a large prize to anybody who proves the following statement: *For any compact simple non-Abelian gauge group, quantum gauge theory exists in continuous, four-dimensional space-time and produces a mass gap.* This is one of their so-called *millennium problems*.

The strand model does not allow arbitrary gauge groups in quantum field theory. According to the strand model, the only compact simple non-Abelian gauge group of interest is SU(3), the gauge group of the strong nuclear interaction. And since the strand model does not seem to allow for glueballs, for SU(3) an effective mass gap of the order of the Planck mass is predicted. (If glueballs would exist in the strand model, the mass gap would still exist but be smaller.) Indeed, the strand model explains the short range of the strong interaction as a consequence of the details of Reidemeister III moves and the quark tangle topology.

The strand model further states that space-time and gauge groups are low-energy approximations that arise through time-average of strands and their crossings. Neither points nor fields exist at a fundamental level; points and fields are approximations to strands. According to the strand model, the *quantum* properties of nature result from the extension of strands. As a consequence, the strand model denies the existence of *any quantum* gauge theory as a separate, exact theory on *continuous* space-time.

The strand model does predict a mass gap for SU(3); but the strand model also denies the existence of quantum gauge theory for any other compact simple non-Abelian gauge group. And even in the case of SU(3) it denies – like for any other gauge groups – the Page 276 existence of a quantum gauge theory on continuous space-time. As deduced above, the strand model allows only the three known gauge groups, and allows their existence only in the non-continuous strand model of space-time. In short, it is *impossible* to realize the wish of the Clay Mathematics Institute.

Summary on tangles made of four or more strands

By exploring all possible tangle classes in detail, we have shown that *every* localized structure made of strands has an interpretation in the strand model. In particular, the strand model makes a simple statement on any tangle made of four or more strands: such a tangle is *composite* of the elementary tangles made of one, two or three strands. In other terms, there are *no* elementary particles made of four or more strands in nature.

The strand model states that each possible tangle represents a physical particle system: an overview is given in Table 13. The mapping between tangles and particles is only possible because (infinitely) many tangles are assigned to each massive elementary particle.

The result of this exploration is that the strand model limits the number of elementary particles to those contained in the standard model of particle physics.

TABLE 13 The match between tangles and particles in the strand model.

Strands	Tangle	Particle	Type
1	unknotted	elementary	vacuum, (unbroken) gauge boson
1	knotted	–	not in the strand model
2	unknotted	composed	composed of simpler tangles
2	rational	elementary	quark or graviton
2	prime, knotted	–	not in the strand model
3	unknotted	composed	composed of simpler tangles
3	braided	elementary	lepton
3	rational	elementary or composed	leptons
3	prime, knotted	–	not in the strand model
4 & more	like for 3 strands	all composed	composed of simpler tangles

FUN CHALLENGES AND CURIOSITIES ABOUT PARTICLE TANGLES

Challenge 196 s

In the strand model, mass appears due to tail braiding. But mass is also due to tangle rotation and fluctuation. How do the two definitions come together?

* *

The following statement seems absurd, but is correct:

▷ The tangle model implies that all elementary particles are point-like, *without* internal structure.

Indeed, if at all, the strand model implies deviations from point-like behaviour only at Planck scale; particles are point-like for all practical purposes.

* *

Challenge 197 e

In the strand model, only crossing switches are observable. How then can the specific tangle structure of a particle have any observable effects? In particular, how can quantum numbers be related to tangle structure, if the only observables are due to crossing changes?

* *

No neutral weak currents that change strangeness or other flavours are observed. In the strand model this observation is a consequence of the tangle shape of the Z boson.

* *

Ref. 237
Challenge 198 r

In 2014, Marek Karliner predicted the existence of six-quark states. Can the strand model reproduce them? Can it settle whether they are molecules of three mesons or genuine six-quark states?

* *

Challenge 199 e Can you use the strand model to show that pentaquarks do not exist?

* *

Ref. 238 What is the relation of the model shown here to the ideas of Viro and Viro on skew lines?

* *

The most prominent proponent of the idea that particles might be knots was, in 1868, William Thomson–Kelvin. He proposed the idea that different atoms might be differently 'knotted vortices' in the 'ether'. The proposal was ignored – and rightly so – because it did not explain anything: neither the properties nor the interactions of atoms were explained. The proposal simply had no relation to reality. In retrospect, the main reason for this failure was that elementary particles and quantum theory were unknown at the time. (Ref. 239)

* *

Purely topological models for elementary particles have been proposed and explored by various scholars in the past. But only a few researchers ever proposed specific topological structures for each elementary particle. Such proposals are easily criticized, so that it is easy to make a fool of oneself; any such proposal thus needs a certain amount of courage.

- Ref. 240 Herbert Jehle modelled elementary particles as closed knots already in the 1970s. However, his model did not reproduce quantum theory, nor does it reproduce all particles known today.
- Ref. 143 Ng Sze Kui has modelled mesons as knots. There is however, no model for quarks, leptons or bosons, nor a description for the gauge interactions.
- Tom Mongan has modelled elementary particles as made of three strands that each carry electric charge. However, there is no connection with quantum field theory or general relativity. (Ref. 241)
- Jack Avrin has modelled hadrons and leptons as Moebius bands, and interactions as cut-and-glue processes. The model however, does not explain the masses of the particles or the coupling constants. (Ref. 139)
- Ref. 141 Robert Finkelstein has modelled fermions as knots. This approach, however, does not explain the gauge properties of the interactions, nor most properties of elementary particles.
- Sundance Bilson-Thompson, later together with his coworkers, modelled elementary fermions and bosons as structures of triple ribbons. The leather trick is used, like in the strand model, to explain the three generations of quarks and leptons. This is by far the most complete model from this list. However, the origin of particle mass, of particle mixing and, most of all, of the gauge interactions is not explained. (Ref. 140)

* *

Strands are *not* superstrings. In contrast to superstrings, strands have a fundamental principle. (This is the biggest conceptual difference.) The fundamental principle for strands is not fulfilled by superstrings. In contrast to superstrings, strands have no tension, no supersymmetry and no own Lagrangian. (This is the biggest physical difference.) Because strands have no tension, they cannot oscillate. Because strands have no supersymmetry, general relativity follows directly. Because strands have no own Lagrangian, particles are tangles, not oscillating superstrings, and quantum theory follows directly. In fact, the definitions of particles, wave functions, fields, vacuum, mass and horizons differ completely in the two approaches.

In contrast to superstrings, strands describe the number of gauge interactions and of particle generations. In contrast to superstrings, strands describe quarks, hadrons, confinement, Regge behaviour, asymptotic freedom, particle masses, particle mixing and coupling constants. In the strand model, in contrast to 'open superstrings', no important configuration has ends. In contrast to open or closed superstrings, strands move in three spatial dimensions, not in nine or ten; strands resolve the anomaly issue without higher dimensions or supersymmetry, because unitarity is automatically maintained, by construction; strands are not related to membranes or supermembranes. In the strand model, no strand is 'bosonic' or 'heterotic', there is no E(8) or SO(32) gauge group, there are no general 'pants diagrams' for all gauge interactions, there is no AdS/CFT duality, there is no 'landscape' with numerous vacuum states, and there is no 'multiverse'. In contrast to superstrings, strands are based on Planck units. And in contrast to superstrings, strands yield the standard model of elementary particles without any alternative.
Ref. 144 In fact, not a single statement about superstrings is applicable to strands.

* *

Strands do not require higher dimensions. On the other hand, it can be argued that strands do produce an additional non-commutative structure at each point in space. In a sense, when strands are averaged over time, a non-commutative inner space is created at each point in space. As a result, when we focus at a specific spatial position over somewhat longer times scales than the Planck time, we can argue that, at that point of space, nature is described by a product of three-dimensional space with an internal, non-commutative space. Since many years, Alain Connes and his colleagues have explored
Ref. 242 such product spaces in detail. They have discovered that with an appropriately chosen non-commutative inner space, it is possible to reproduce many, but not all, aspects of the standard model of particle physics. Among others, choosing a suitable non-commutative space, they can reproduce the three gauge interactions; on the other hand, they cannot reproduce the three particle generations.

Connes' approach and the strand model do not agree completely. One way to describe the differences is to focus on the relation of the inner spaces at different points of space. Connes' approach assumes that each point has its own inner space, and that these spaces are not related. The strand model, instead, implies that the inner spaces of neighbouring points are related; they are related by the specific topology and entanglement of the involved strands. For this very reason the strand model does allow to understand the origin of the three particle generations and the details of the particle spectrum.

There are further differences between the two approaches. Connes' approach assumes that quantum theory and general relativity, in particular, the Hilbert space and the spatial manifold, are given from the outset. The strand model, instead, deduces these structures from the fundamental principle. And, as just mentioned, Connes' approach is not

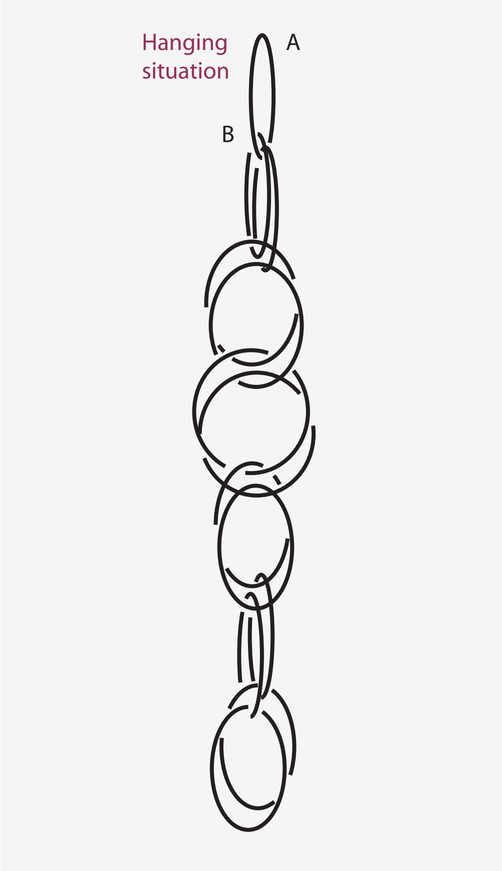

FIGURE 106 A ring chain gives an impression of motion along the chain, when holding ring B while dropping ring A.

unique or complete, whereas the strand model seems to be. Of the two, only the strand model seems to be unmodifiable, or 'hard to vary'.

* *

The strand model implies that there is *nothing new* at small distances. At small distances, or high energies, nature consists only of strands. Thus there are no new phenomena there. Quantum theory states that at small scales, nothing new appears: at small scales, there are *no* new degrees of freedom. For example, quantum theory states that there is no kingdom *Lilliput* in nature. The strand model thus confirms the essence of quantum theory. And indeed, the strand model predicts that between the energy scale of the heaviest elementary particle, the top quark, 173 GeV, and the Planck energy, 10^{19} GeV, nothing is to be found. There is a so-called *energy desert* – empty of interesting features, particles or phenomena – in nature.

* *

Most ropes used in sailing, climbing or other domains of everyday life are produced by braiding. Searching for 'braiding machine' on the internet yields a large amount of videos. Searching for 'LEGO braiding machine' shows the most simple and beautiful

FIGURE 107 The ring chain trick produces an illusion of motion (mp4 film © Franz Aichinger). Can more rings be added in horizontal directions?

Challenge 200 e

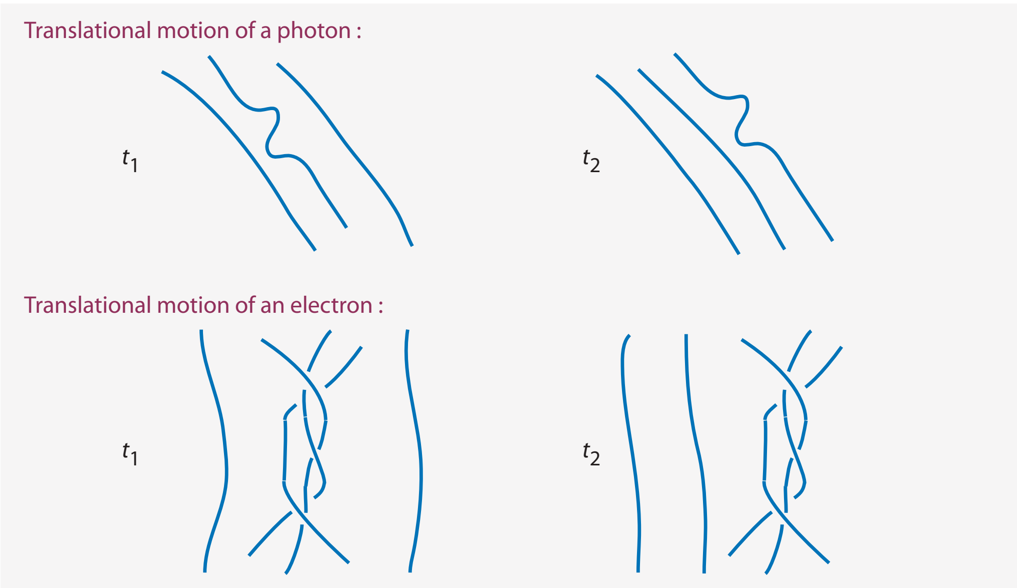

FIGURE 108 Motion of photons and electrons through strand hopping.

examples and allows you to see how they work.

* *

Not all tangle assignments are self-evident at first sight. Figure 109 shows a tangle whose status in the strand model is not clear. Can you explain what the tangle represents?

Challenge 201 e

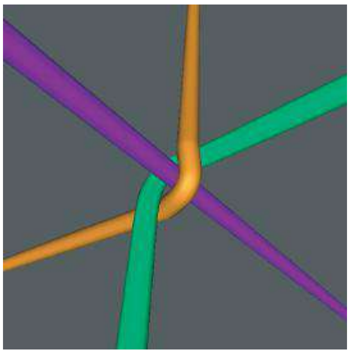

FIGURE 109 A discarded candidate tangle for the W boson.

* *

Page 153
Challenge 202 ny
What is the effect of shivering on braiding, and thus on weak particle mixing, on particle tangle families and on the number of generations?

* *

Are all bosons made of strands whose ends are exactly opposite to each other at spatial infinity? Photon, graviton, gluon, W, Z and the virtual Higgs comply. The unbroken ones are axial, the broken ones are flat. Is there a reason or a sense for this issue?
Challenge 203 e

CPT invariance

CPT invariance is a fundamental property of quantum field theory. In the strand model, charge conjugation C is modelled as a mirror transformation of the tangle; parity P is modelled as the change of sign of the belt trick of the tangle core; and motion inversion T is modelled as the inverse motion of the core of a particle tangle.

In other words, CPT invariance is natural in the strand model. Therefore, the strand model predicts that particles and antiparticles have the same g-factor, the same dipole moment, the same mass, the same spin, exactly opposite charge value, etc. All this is also predicted by quantum field theory, and is confirmed by experiment.

Motion through the vacuum – and the speed of light

Up to now, one problem was left open: How can a particle, being a tangle of infinite extension, move through the web of strands that makes up the vacuum? An old trick, known already in France in the nineteenth century, can help preparing for the idea of particle motion in space. Figure 106 shows a special chain that is most easily made with a few dozen key rings. If the ring B is grabbed and the ring A released, this latter ring seems to fall down along the whole chain in a helical path, as shown in the film of Figure 200. If you have never seen the trick, try it yourself; the effect is astonishing. In reality, this is an optical illusion. No ring is actually falling, but the sequence of rings moves in a way that creates the impression of ring motion. And this old trick helps us to solve a number of issues about particle motion that we swept under the carpet so far.
Ref. 243

The main idea on particle motion in the strand model is the following:

▷ ***Translational particle motion*** is also due to strand substitution, or 'strand hopping'.

A schematic illustration of translational motion is given in Figure 108. In the strand model, contrary to the impression given so far, a tangle does not always need to move as a whole along the strand. This is seen most easily in the case of a photon. It is easy to picture that the tangle structure corresponding to a photon can also hop from strand to strand. At any stage, the structure is a photon; but the involved strand is never the same.

The idea of motion through strand hopping also works for massive particles. The motion of a massive particle, such as an electron, is shown schematically in Figure 108. The figure shows that through a tail unbraiding, the structure that describes an electron can get rid of one strand and grab a new one. This process has a low probability, of course. In the strand model, this is one reason that massive particles move more slowly than light, even if the first approximation yields a zero mass value.

We note that this explanation of motion is important also for the mapping from strand diagrams to Feynman diagrams. For many such diagrams, for example for the annihilation of particles and antiparticles in QED, strand hopping and tail unbraiding play a role. Without them, the mapping from strands to quantum field theory would not be possible.

In summary, tangles of massive particles *can* move through the vacuum using hopping – via tail unbraiding – and this naturally happens more slowly than the motion of photons, which do not need any process at the border of space to hop. The speed of photons is thus a limit speed for massive particles; special relativity is thus recovered.

> “The ground of science was littered with the corpses of dead unified theories.”
>
> Freeman Dyson

SUMMARY ON MILLENNIUM ISSUES ABOUT PARTICLES AND THE VACUUM

We have discovered that the strand model makes a strong statement: elementary particles can only be made of one, two or three tangled strands. Each massive elementary particle is represented by an infinite *family* of rational tangles of fixed strand number. The family members differ by the added number of braids made of three strands with 6 crossings, i.e., by the different numbers of virtual Higgs bosons.

For *one-stranded* particles, the strand model shows that the photon is the only possibility. For *two-stranded* particles, the strand model shows that there are precisely three generations of two massive quarks. For *three-stranded* elementary particles, the strand model shows that there is a Higgs boson, 8 gluons, the W, the Z, and three generations of leptons. Neutrinos and antineutrinos differ and are massive Dirac particles. The strand model thus predicts that the neutrino-less double-beta decay will *not* be observed. Glueballs most probably do not exist.

The strand model uses the tangle assignments of Figure 110 and Figure 111 to explain the origin of all quantum numbers of the observed elementary particles. The strand model reproduces the quark model, including all the allowed and all the forbidden hadron states. For mesons and baryons, the strand model predicts the correct mass sequences and quantum numbers. Therefore, we have also completed the argument that

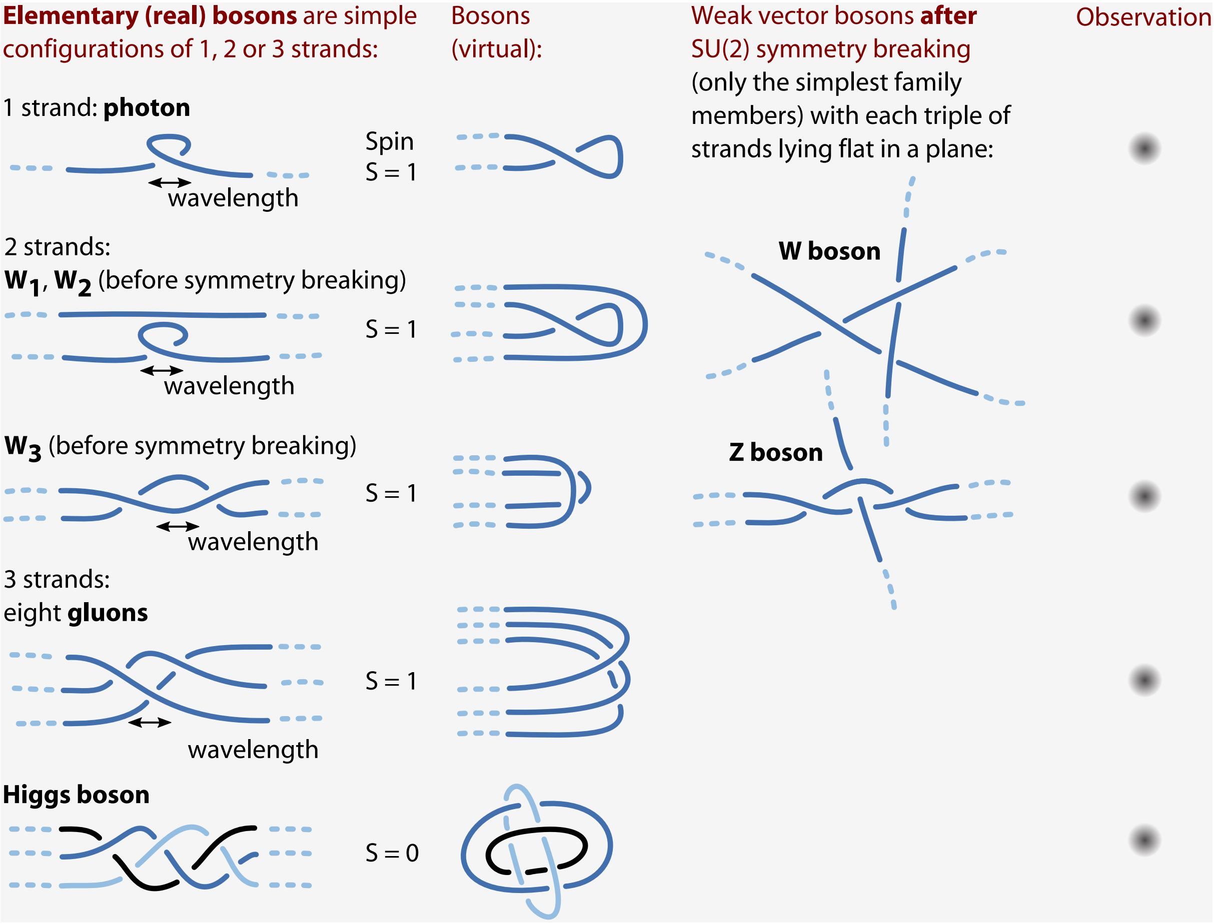

FIGURE 110 The tangle models for the elementary bosons. These tangles determine the spin values, the corresponding propagators, and ensure that the massless photons and gluons move with the speed of light. No additional elementary bosons appear to be possible.

Page 155 all observables in nature are due to crossing switches. Tetraquarks are predicted to exist. A way to calculate hadron form factors is proposed.

In the strand model, all tangles are mapped to known particles. The strand model predicts that *no* elementary particles outside the standard model exist, because no tangles are left over. For example, there are no axions, no leptoquarks and no supersymmetric particles in nature. The strand model also predicts the lack of other gauge bosons and other interactions. In particular, the strand model – corrected in 2012 – reproduces the Page 330 existence of exactly one Higgs boson. In fact, any new elementary particle found in the future would contradict and invalidate the strand model.

In simple words, the strand model explains why the known elementary particles exist and why others do not. We have thus settled two further items from the millennium list Page 164 of open issues. In fact, the deduction of the elementary particle spectrum given here is, the first and, at present, also the *only* such deduction in the research literature.

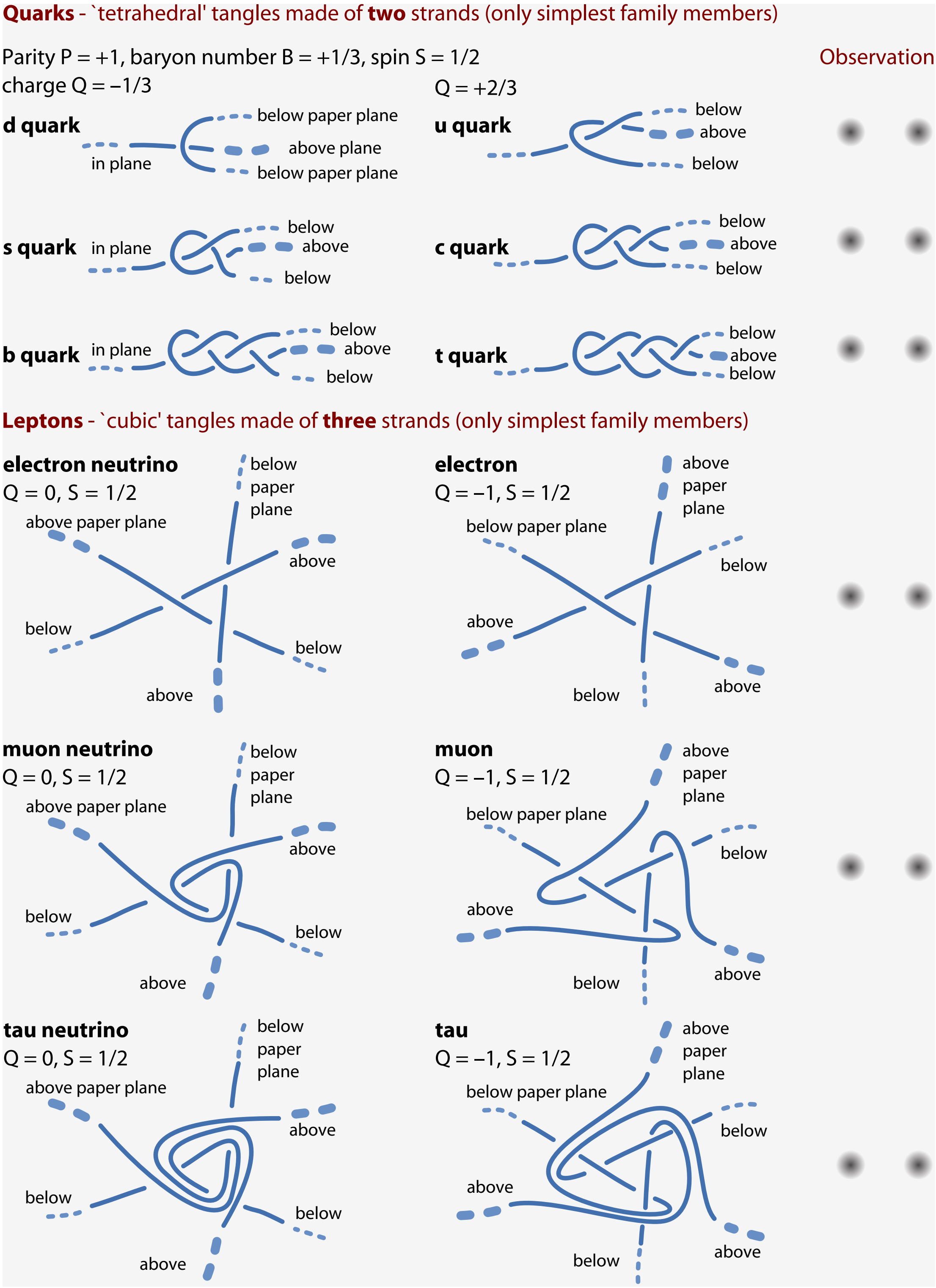

FIGURE 111 Elementary fermions are described by rational, i.e., unknotted tangles. Their structures lead to coupling to the Higgs, as illustrated in Figure 130, produce positive mass values, and limit the number of generations to 3. The tangles determine the specific fermion propagators. The tethers of the quark tangles follow the axes of a tetrahedron. The neutrino cores are simpler when seen in three dimensions: they are simply twisted triples of strands. The tethers of all lepton tangles approach the three coordinate axes at large distances from the core. No additional elementary fermions appear to be possible.

The omnipresent number 3

The strand model shows that the number 3 that appears so regularly in the standard model of particle physics – 3 generations, 3 interactions, charge values $e/3$ and $2e/3$ of quarks (as shown below), 3 colours and SU(3) – is, in each case, a consequence of the three-dimensionality of space. In fact, the strand model adds a further, but related number 3 to this list, namely the maximum number of strands that make up elementary particles.

Page 209
The three-dimensionality of space is, as we saw already above, a result of the existence of strand tangles: only three dimensions allow tangles of strands. In short, all numbers 3 that appear in fundamental physics are explained by strands.

Predictions about dark matter and searches for new physics

Following the vast majority of scholars, astrophysical observations imply that galaxies and galaxy clusters are surrounded by large amounts of matter that does not radiate. This unknown type of matter is called *dark matter*.

In the strand model, the known elementary particles are the only possible ones. Therefore,

▷ The strand model predicts that dark matter, if it exists, is a mixture of particles of the standard model and black holes.

Page 164
This statement settles a further item from the millennium list of open issues.

The prediction from 2008 of a lack of new elementary particles in dark matter is at odds with the most favoured present measurement interpretations, but cannot yet be ruled out. The detection of black hole mergers in 2015 can even be seen as a partial confirmation. However, the issue is obviously not yet settled. In fact, the prediction provides another hard test of the model: if dark matter is found to be made of yet unknown particles, the strand model is in trouble.

We can condense all the results on particle physics found so far in the following statement:

▷ There is nothing to be discovered about nature outside general relativity and the standard model of particle physics.

Strands predict that there is no hidden aspect of nature left. In particular, the strand model predicts a so-called high-energy *desert*: it predicts the lack of any additional elementary particle. Equivalently, the strand model predicts that apart from the Planck scale, there is no further energy scale in particle physics. Some researchers call this beautiful result the *nightmare scenario*.

In other words, there is no room for discoveries beyond the Higgs boson at the Large Hadron Collider in Geneva, nor at the various dark matter searches across the world. If any new elementary particle is discovered, the strand model is wrong. More precisely, if any new elementary particle *that contradicts the strand model* is discovered, the strand model is wrong. That some unknown elementary particle has been missed in the present exploration of tangle classes is still a logical possibility, however.

Because the strand model confirms the standard model and general relativity, a further prediction can be made: *the vacuum is unique and stable.* There is no room for other options. For example, there are no domains walls between different vacuum states and the universe will not decay or change in any drastic manner.

In summary, the strand model predicts a lack of any kind of science fiction in modern physics.

Chapter 12
PARTICLE PROPERTIES DEDUCED FROM STRANDS

> Tutto quel che vedete, lo devo agli spaghetti.**
>
> Sophia Loren

The Planck units, via strands and the fundamental principle, explain almost all that is known about motion: strands explain *what* moves and *how* it moves. But the strand model is only correct if it also explains every measured property of every elementary particle. So far, we only deduced the spectrum and the quantum numbers of the elementary particles. Three kinds of particle properties from the millennium list remain open: the *masses*, the *mixing angles* and the *couplings*. These measured particle properties are important, because they determine the amount of change – or physical action – induced by the motion of each elementary particle.

So far, the strand model has answered all open questions on motion that we explored. In particular, the strand model has explained why quantum field theory, the interactions, the particle spectrum, general relativity and cosmology are what they are. But as long as we do not understand the measured properties of the elementary particles, we do not understand motion completely.

In short, the next step is to find a way to *calculate* these particle properties – and obviously, to show that the calculations agree with the measurements. The step is particularly interesting; so far, no other unified model in the research literature has ever achieved such calculations – not even calculations that disagree with measurements.

Because the strand model makes no experimental predictions that go beyond general relativity and the standard model of particle physics, explaining the properties of elementary particles is the *only way* to confirm the strand model. Many ways to test or to refute the strand model are possible; but only a calculation of the measured particle properties can confirm it.

The ideas in this chapter are more speculative than those of the past chapters, because the reasoning depends on the way that specific tangles are assigned to specific particles. Such assignments are never completely certain. We continue keeping this in mind.

** 'Everything you see, I owe it to spaghetti.' Sofia Villani Scicolone is an Italian actress and Hollywood star.

TABLE 14 The measured elementary particle masses, as given by the Particle Data Group in 2016.

Elementary particle	Mass value
Electron neutrino	$< 2\,\mathrm{eV}/c^2$
Muon neutrino quark	$< 2\,\mathrm{eV}/c^2$
Tau neutrino	$< 2\,\mathrm{eV}/c^2$
Electron	$0.510\,998\,9461(31)\,\mathrm{MeV}/c^2$
Muon	$105.658\,3745(24)\,\mathrm{MeV}/c^2$
Tau	$1776.86(12)\,\mathrm{MeV}/c^2$
Up quark ($q = 2/3\,e$)	$2.2(6)\,\mathrm{MeV}/c^2$
Down quark ($q = -1/3\,e$)	$4.7(5)\,\mathrm{MeV}/c^2$
Strange quark ($q = -1/3\,e$)	$96(8)\,\mathrm{MeV}/c^2$
Charm quark ($q = 2/3\,e$)	$1.27(3)\,\mathrm{GeV}/c^2$
Bottom quark ($q = -1/3\,e$)	$4.18(4)\,\mathrm{GeV}/c^2$
Top quark ($q = 2/3\,e$)	$173.21(1.22)\,\mathrm{GeV}/c^2$
W boson	$80.385(15)\,\mathrm{GeV}/c^2$
Z boson	$91.1876(21)\,\mathrm{GeV}/c^2$
Higgs boson	$125.09(24)\,\mathrm{GeV}/c^2$
Photon	not detectable
Gluons	not detectable
Graviton	not detectable
For comparison: the corrected Planck mass $\sqrt{\hbar c/4G}$	$0.611 \cdot 10^{19}\,\mathrm{GeV}/c^2$

THE MASSES OF THE ELEMENTARY PARTICLES

The mass describes the inertial and gravitational effects of a body. The strand model must reproduce all mass values observed in nature; if it doesn't, it is wrong.

To reproduce the masses of *all* bodies, it is sufficient that the strand model reproduces the measured masses, the mixing angles and the coupling strengths of the *elementary* particles. We start with their masses. All measured mass values are given in Table 14. All these values – more precisely: their ratios to the Planck mass – are unexplained and are part of the millennium list of open issues in fundamental physics.

In nature, the *gravitational mass* of a particle is determined by the space curvature that it induces around it. In the strand model, this curvature is due to the modified fluctuations that result from the presence of the tangle core; in particular, the curvature is due to the modified fluctuations of the particle tails – twisted tail pairs – and to the modified vacuum strand fluctuations just around the particle position. The modified strand shape fluctuations produce a crossing switch distribution around the tangle core; the crossing switch distribution leads to spatial curvature; at sufficiently large distances, this curvature distribution is detected as a gravitational mass.

In contrast, *inertial mass* appears in the Dirac equation. In the strand model, inertial mass is determined by the frequency and the wavelength of the helix drawn by the rotating phase vector. These quantities in turn are influenced by the type of tangle, by the fluctuations induced by the particle charges, by the topology changes induced by the weak interaction, and, in the case of fermions, by the average frequency and size of the belt (and possibly leather) trick. All these processes are due to strand shape fluctuations.

In short, both gravitational and inertial particle mass are due to strand fluctuations. More specifically, the mass seems mainly due to the fluctuations of the *tails* of the particle tangle: gravitational and inertial mass are due to the belt trick. *The strand model thus suggests that gravitational and inertial mass are automatically equal.* In particular, the strand model suggests that every mass is surrounded by fluctuating tails with crossing switches whose density decreases with distance and is proportional to the mass itself. As Page 282 discussed above, this idea leads to universal gravity.

General properties of particle mass values

So far, our adventure allows us to deduce several results on the mass values of elementary particles:

— The strand model implies that the masses of elementary particles are *not free parameters*, but that they are determined by the specific topology, or tangledness, of the underlying tangles and their tangle families. Particle masses are thus *fixed* and *discrete* in the strand model – as is observed. Of course, we have to take into account all the members in each tangle family.

— The strand model implies that masses are always *positive* numbers.

— The strand model implies that the *more complex* a tangle is, the *higher* its mass value is. This follows from the behaviour of tangle tail fluctuations around the tangle core.

— Because particle masses are due to strand fluctuations, the strand model also implies that all elementary particle masses are *much smaller than the Planck mass*, as is observed. Also this result follows from the behaviour of tangle tail fluctuations around the tangle core.

— Because particle masses are due to strand fluctuations, particle and antiparticle masses – their tangles are mirrors of each other – are always *equal*, as is observed.

— Because particle masses are due to strand fluctuations, particle masses do *not* depend on the age of the universe, nor on their position in the universe, nor on any other state variable: The strand model predicts that particle masses are constant and invariant, Page 373 as is observed.

— Because particle masses are due to strand fluctuations, and the fluctuations differ somewhat for tight and loose tangles of the same shape and topology, the strand model predicts that particle masses change – or *run* – with energy, as is observed.

The general properties of particle masses are thus reproduced by the strand model. Therefore, continuing our exploration makes sense. We start by looking for ways to determine the mass values from the tangle structures. We discuss each particle class separately, first looking at mass ratios, then at absolute mass values.

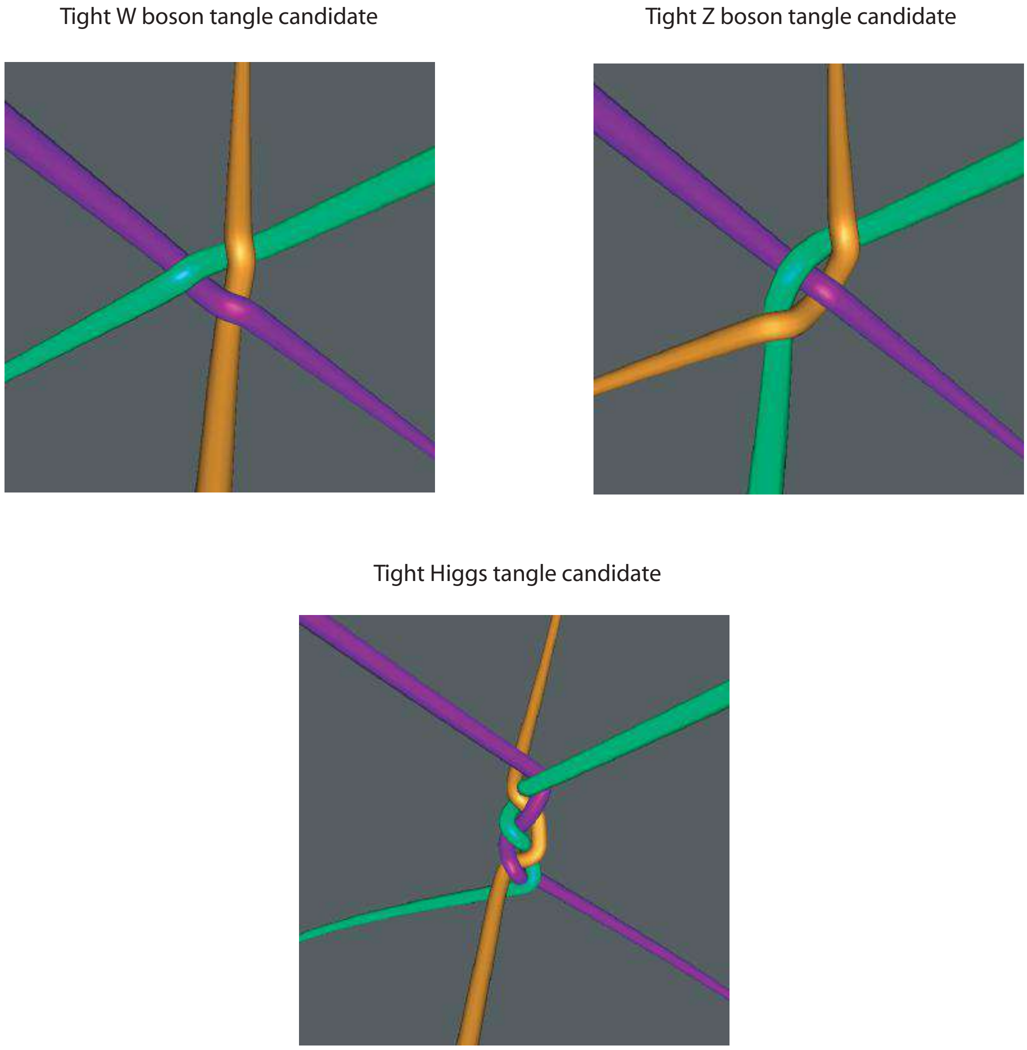

FIGURE 112 Tight tangle candidates (of 2015/2016) for the simplest tangles of the W, the Z and the Higgs bosons. In contrast to the pictures, the W and Z tails lie in a plane.

Boson masses

Three elementary particles of integer spin have non-vanishing mass: the W boson, the Z boson and the Higgs boson. Mass calculations are especially simple for bosons, because in the strand model, they are *clean* systems: each boson is described by a relatively simple tangle family; furthermore, bosons do not need the belt trick to rotate continuously.

We expect that the induced curvature, and thus the gravitational mass, of an elementary boson is due to the disturbance it introduces into the vacuum. *At Planck energy*, this disturbance will be, to a large extent, a function of the *ropelength* introduced by the corresponding *tight* tangle. Let us clarify these concepts.

Tight or *ideal* tangles or knots are those tangles or knots that appear if we imagine strands as being made of a rope of *constant* diameter that is *infinitely flexible*, *infinitely* *slippery* and pulled as tight as possible. Examples of tight tangles are shown in Figure 112. With physical ropes from everyday life, tight knots and tangles can only be approximated, because they are not infinitely flexible and slippery; tight tangles are mathematical idealizations. But tight tangles of strands are of special interest: if we recall that each strand has an effective diameter of one Planck length, tight tangles realize the Planck limit of the strand model.

— The *ropelength* of a tight *closed knot* is the length of a perfectly flexible and slippery rope of constant diameter required to tie the tight knot. In other words, the ropelength is the smallest amount of idealized rope needed to tie a knot.
— The *ropelength* of a tight *open knot* is the length by which a very long rope tied into a tight knot is shortened.
— With a bit of care, the concept of ropelength can be also be defined for tangles of several strands.

In the following, the ropelength is assumed to be measured in units of the rope *diameter*. Measuring ropelength in units of the rope radius is less common.

In the strand model, the ropelength measures, to a large extent, the amount by which a tight knot or tangle disturbs the vacuum around it. The ropelength fulfils all the properties of particle mass mentioned above: the ropelength is discrete, positive, increases with tangle complexity, is equal for particles and antiparticles, and is a constant and invariant quantity. The ropelength will thus play an important role in any estimate of a particle mass.

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It is known from quantum field theory that the masses of W and Z bosons do not change much between Planck energy and everyday energy, whatever renormalization scheme is used. This allows us, with a good approximation, to approximate the weak boson masses at low, everyday energy with their mass values at Planck energy. Thus we can use tight tangles to estimate boson masses.

Ref. 244

In the strand model, the *gravitational mass* of a spin 1 boson is proportional to the radius of the disturbance that it induces in the vacuum. For a boson, this radius, and thus the mass, scales as the third root of the ropelength of the corresponding tight tangle.

W/Z boson mass ratio and mixing angle (in the 2016 tangle model)

Page 360

Candidates for the simplest tangles of the W boson and of the Z boson families are shown in Figure 112. The corresponding ropelength values for tight tangles, determined numerically, are $L_W = 4.28$ and $L_Z = 7.25$ rope diameters. The strand model estimates the W/Z mass ratio by the cube root of the ropelength ratio:

Ref. 246

$$\frac{m_W}{m_Z} \approx \left(\frac{L_W}{L_Z}\right)^{1/3} = 0.84 \,. \tag{196}$$

Ref. 233

This value has to be compared with the experimental ratio of 80.4 GeV/91.2 GeV=0.88. The agreement between experiment and strand model is not good. But the result is acceptable. First of all, the strand model reproduces the higher value of the neutral Z boson's mass: a tangle with spatial symmetry is more complex than one without. Finally, it is also clear why the calculated mass ratio does not match the experimental result.

First, the simple tangles represent and approximate W and Z bosons only to the first order. As mentioned above, in the strand model, every massive particle is represented by an infinite family of tangles. The strand model thus also predicts that the match between the calculated and the measured ratio m_W/m_Z should improve when higher-order Feynman diagrams, and thus more complicated tangle topologies, are taken into account. Improving the calculation is still a subject of research. Secondly, approximating the tight knot effects with an effective radius, thus just using the ropelength to determine the mass, implies neglecting the actual shape, and effectively approximating their shape by a sphere. Thirdly, as already mentioned, this calculation assumes that the low energy mass ratio and the mass ratio at Planck energy are equal.

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Despite the used approximations, the tight tangle estimate for the W/Z mass ratio gives an acceptable agreement with experiment. The main reason is that we expect the strand fluctuations from the various family members to be similar for particles with the *same* number of strands. For these mass ratios, the tail braiding processes cancel out. Also the other two approximations are expected to be roughly similar for the two weak bosons. This similarity explains why determining the W/Z boson mass *ratio* is possible with acceptable accuracy.

The W/Z mass ratio also determines the weak mixing angle θ_w of the weak interaction Lagrangian, through the relation $\cos\theta_w = m_W/m_Z$. The strand model thus predicts the value of the weak mixing angle to the same accuracy as it predicts the W/Z mass ratio.

This argument leads to a puzzle: Can you deduce from the strand model how the W/Z mass ratio changes with energy?

Challenge 204 ny

Also the *inertial masses* of the W and Z bosons can be compared. In quantum theory, the inertial mass relates the wavelength and the frequency of the wave function. In the strand model, a quantum particle that moves through vacuum is a tangle core that rotates while advancing. The frequency and the wavelength of the helix thus generated determine the inertial mass. The process is analogous to the motion of a body moving at constant speed in a viscous fluid at small Reynolds numbers. Despite the appearance of friction, the analogy is possible. If a small body of general shape is pulled through a viscous fluid by a constant force, such as gravity, it follows a *helical* path. This analogy implies that, for spin 1 particles, the frequency and the wavelength are above all determined by the effective radius of the small body. This radius is the main influence on the frequency of the belt trick. The strand model thus suggests that the inertial mass – inversely proportional to the path frequency and the path wavelength squared – of the W or the Z boson is *approximately* proportional to its tight knot radius. This radius is given by the cube root of the ropelength. We thus get the same result as for the gravitational mass.

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Ref. 163

Also the inertial mass is not exactly proportional to the average tight knot radius; the precise shape of the tight knot and the other tangle family members play a role. The strand model thus predicts that a more accurate mass calculation has to take into account these effects.

In summary, the strand model predicts, from the belt trick, a W/Z mass ratio and thus a weak mixing angle close to the observed ratio, and explains the deviation of the approximation from the measured value – provided that the tangle assignments are correct.

The g-factor of the W boson

Ref. 233 Experiments show that the W boson has a g-factor with the value $g_W = 2.2(2)$. The limited accuracy does not yet allow to detect any anomalous magnetic moment, which would be especially interesting. Nevertheless, the results can be compared to the prediction of the strand model.

The strand model makes a simple prediction for charged elementary particles: because mass rotation and charge rotation are both due to the rotation of the tangle core, the g-factor of all such particles is 2 – in the approximation that neglects Feynman diagrams of higher order, i.e., that neglect anomalous effects. In particular, the g-factor of the W boson is predicted to be 2 in this approximation. Also this prediction thus agrees both with experiment and with the standard model of particle physics. Ref. 245

The Higgs/Z boson mass ratio

The observed mass value of the Higgs boson is 125(1) GeV. The observed mass value for the Z boson is 91.2(1) GeV. Like for the other bosons, the strand model suggests using the ropelength to estimate the mass of the Higgs boson tangle. The candidate tangle for the Higgs boson was illustrated above in Figure 96, and its tight version is shown in Figure 112.

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Ref. 246 The ropelength of the tight Higgs tangle turns out to be 17.1 diameters, determined by Eric Rawdon with a computer approximation. This value yields a naive mass estimate for the Higgs boson of $(17.1/7.25)^{1/3} \cdot 91.2$ GeV, i.e.,

$$m_{\text{Higgs}} \approx 121 \text{ GeV} . \qquad (197)$$

Starting with the W boson yields an estimate for the Higgs mass of 128 GeV. Both estimates are not good but acceptable, given that the non-sphericity of the W, Z and Higgs boson tangles have not been taken into account. (The strand model suggests that for a strongly non-spherical shape – such as the shape of the W, Z and Higgs tangle – the effective mass is higher than the value deduced from ropelength alone.) Deducing better mass ratio estimates for the W, Z and Higgs tangles is still a subject of research.

In summary, the strand model predicts Higgs/Z, Higgs/W and W/Z mass *ratios* close to the observed values; and the model suggests explanations for the deviations of the approximation from the observed value – provided that the tangle assignments for the three bosons are correct.

A first approximation for absolute boson mass values

The tangles for the W, Z and Higgs bosons also provide a first approximation for their *absolute* mass values. The tangles are rational; in particular, each tangle is made of strands that can be pulled straight. This implies, for each strand separately, that a configuration with no extra strand length and no net core rotation is possible. As a result, in the first approximation, the gravitational mass and the inertial mass of the elementary bosons both *vanish*.

A better approximation for mass values requires to determine, for each boson, the probability of crossing switches in and around its tangle core. This probability depends

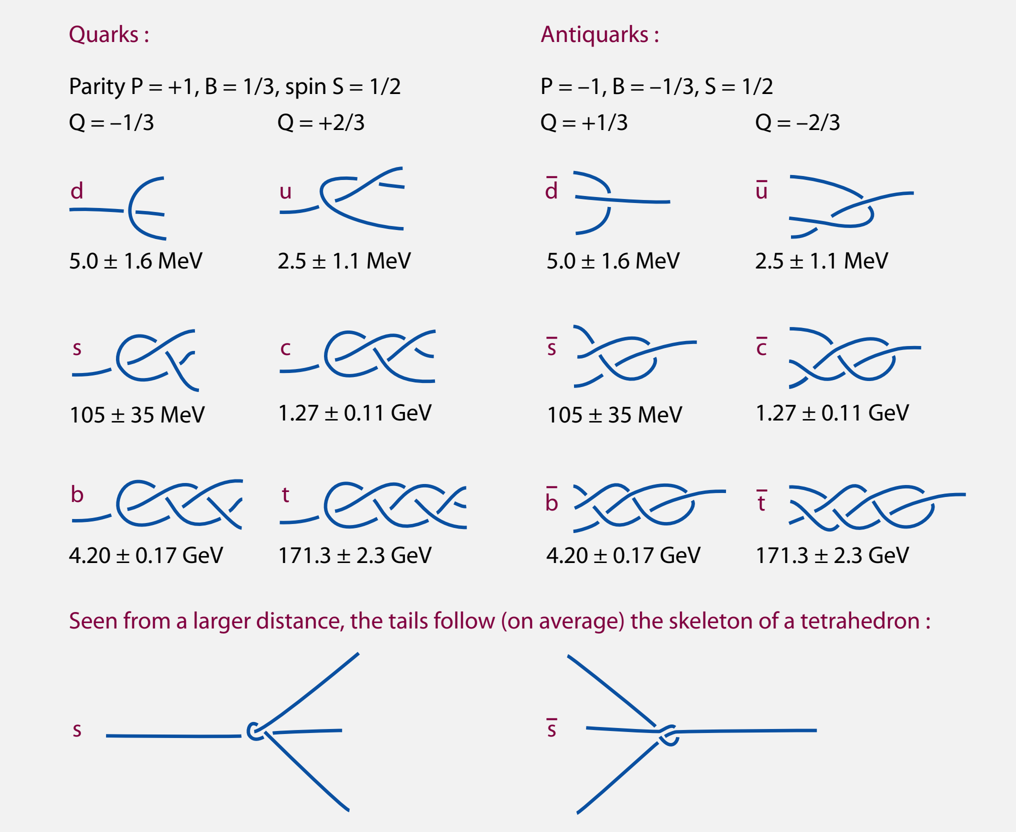

FIGURE 113 The simplest tangles assigned to the quarks and antiquarks. The experimental mass values are also given. Calculated mass values are still a topic of research.

on the probabilities for tail braiding and for core rotation. These probabilities are low, because, sloppily speaking, the corresponding strand fluctuations are rare. The rarity is a due to the specific tangle type: tangles whose strands can be pulled straight have low crossing switch probabilities at their core or at their tails when they propagate.

The strand model thus predicts that elementary boson masses, like all other elementary particle masses, are much smaller than the Planck mass, though not exactly zero. This prediction agrees with observation: experimentally, the mass values for the W, Z and Higgs are of the order of 10^{-17} Planck masses. We will search for more precise mass estimates below.

Quark mass ratios

Quarks are fermions. In the strand model, mass estimates for fermions are more difficult than for bosons, because their shapes and their tails are less symmetric. Still, using Figure 113, the strand model allows two predictions about the relations between quark masses.

— The quark masses are predicted to be the same for every possible colour charge. This is observed.

Page 321 — The progression in ropelength of the tight basic tangles for the six quarks suggests a progression in their masses. This is observed, though with the exception of the up quark mass. For this exceptional case, effects due to tail braiding, to the Higgs boson, Page 376 and to quark mixing are expected to play a role, as argued below.

Let us try to extract numerical values for the quark mass ratios. We start by exploring the tight quark tangles, thus Planck-scale mass values. For each quark number q, the quark mass will be the weighted average over the mass of its family tangles with q, $q + 6$, $q + 12$, ... crossings, where the period 6 is due to the Higgs boson. Each tight tangle has a certain ropelength. The mass of each tangle will be determined by the frequency of crossing changes at the core, including those due to the belt trick. The quark mass then is the average over all family tangles; it will be determined by the frequency of tail braiding and of all other fluctuations that generate crossing switches.

For determining mass ratios, the frequency of the crossing switches at the core are the most important. Given that the particles are fermions, not bosons, this frequency is expected to be an exponential of the ropelength L. Among quarks, we thus expect a general mass dependence of the type

$$m \sim \mathrm{e}^{aL} \tag{198}$$

where a is an unknown number of order 1. We note directly that such a relation promises general agreement with the observed quark mass ratios. Nevertheless, the exponential dependence must not be seen as more than an educated guess.

Ref. 247 Actual ropelength calculations by Eric Rawdon and Maria Fisher show that the ropelength of quark tangles increases roughly linearly with q, as expected from general knot theoretic arguments. Their results are given in Table 15. Comparing these calculated ropelength differences with the known low-energy quark masses confirms that the number a has an effective value in the range between 0.4 and 0.9, and thus indeed is of order one.

The results of Table 15 suggest that the top quark should be particularly heavy – as is observed. The results of Table 15 also suggest that something special is going on for the u-d quark pair, which is out of sequence with the other quarks. Indeed, the strand model predicts a very small mass, – at the Planck scale – for the down quark. However, in nature, the down mass is observed to be *larger* than the up mass. (We note that despite this issue, meson mass sequences are predicted correctly.)

It could well be that the large symmetry of the simplest down quark tangle is the reason for the exceptional mass ordering. The large symmetry of the down tangle should lead to easier, i.e., more frequent interactions with the Higgs boson. In other words, the braiding, i.e., the mixing with the more massive family members, appears to be *higher* for the down quark than that for the up quark, and this would explain the higher mass of the down quark. Maybe this explanation can be checked with data.

The experimental values for the quark masses are given in Table 16; the table also Ref. 248 includes the values extrapolated to Planck energy for the pure standard model. The calculation of the strand model *does not agree* with the data. The only encouraging aspect is that the ropelength approximation provides an approximation for older speculations

TABLE 15 Calculated ropelengths, in units of the rope *diameter*, of tight quark tangles of Figure 86 (Page 320) with tails oriented along the skeleton of a tetrahedron.

Tangle	Length	Ropelength	Difference
skeleton (vacuum)	138.564065	base value	
simplest d	139.919533	1.355468	1.355468
simplest u	142.627837	4.063773	2.708305
simplest s	146.175507	7.611443	3.547670
simplest c	149.695643	11.131578	3.520136
simplest b	153.250364	14.686299	3.554721
simplest t	157.163826	18.599761	3.913462

TABLE 16 For comparison: the quark masses at Planck energy,, calculated from the measured quark masses using the standard model of particle physics – assuming that it is correct up to Planck energy.

Quark	Low energy mass	Planck energy mass
u ($q = 2/3e$)	2.5(1.1) MeV	0.45(0.16) MeV
d ($q = -1/3e$)	5.0(1.6) MeV	0.97(0.10) MeV
s ($q = -1/3e$)	105(35) MeV	19.4(1.2) MeV
c ($q = 2/3e$)	1270(110) MeV	213(8) MeV
b ($q = -1/3e$)	4200(170) MeV	883(10) MeV
t ($q = 2/3e$)	171300(2300) MeV	66993(880) MeV

on approximately *fixed mass ratios* between the up-type quarks u, c, t and *fixed mass ratios* between down-type quarks d, s, b. The attempted strand model estimate shows that ropelength alone is *not sufficient* to understand quark mass ratios. Research has yet to determine which tangle shape aspect has to be included to improve the correspondence with experiment.
Ref. 249

In fact, the strand model predicts that everyday quark masses result from a combination of three effects: the effect of ropelength and of tangle core shape on rotation and the belt trick, the effect of sixfold tail braiding, and the effect of the energy dependence of mass between Planck energy and everyday energy, due to core loosening.

In short, an analytic calculation for quark masses seems difficult, due to their elongated shape. Better analytical approximations should be possible, though. And with sufficient computer power, it will also be possible to determine the frequency of core shape deformations and core rotations, including the belt trick. It will also be possible to determine the energy dependence of the quark masses, and the probability for tail braiding. More research is needed on all these points; in the end, it will allow to calculate quark masses.

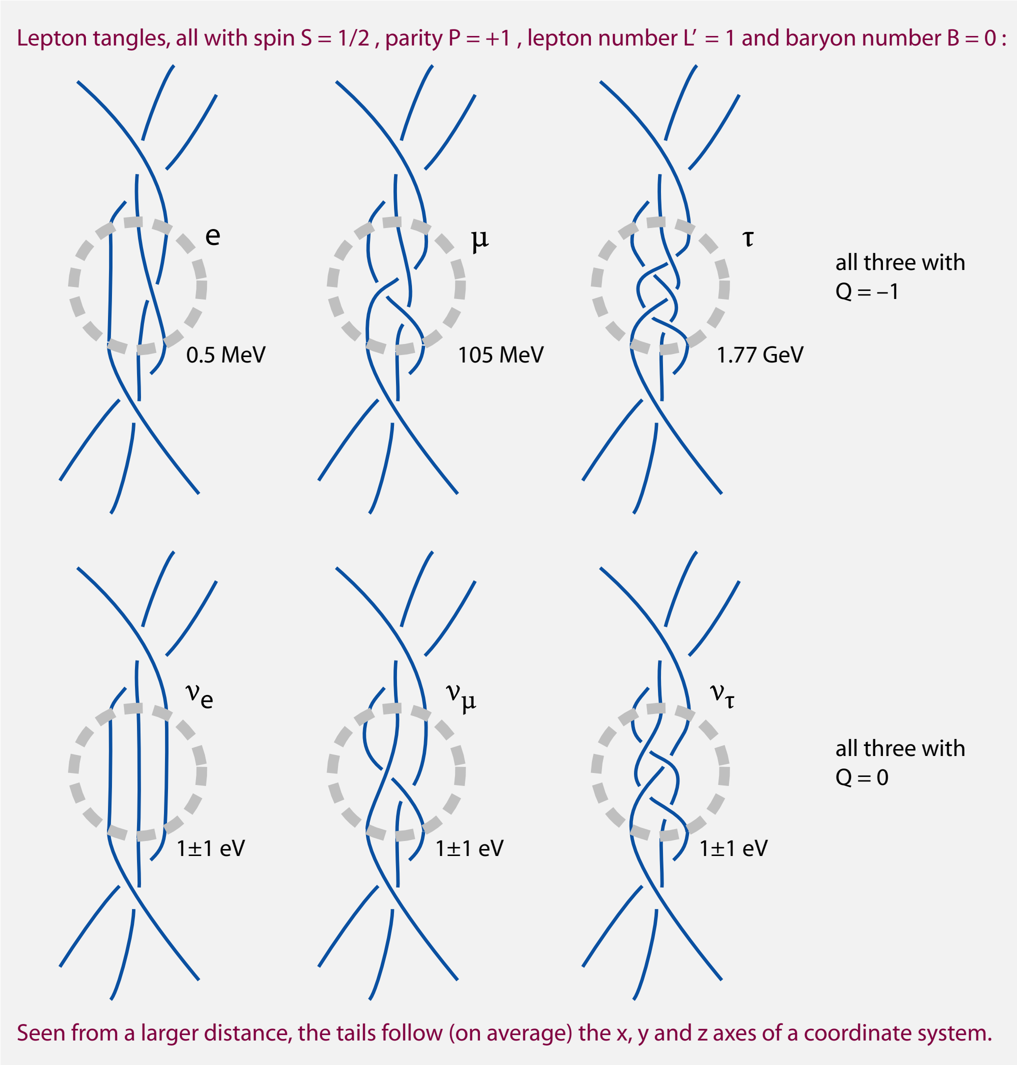

FIGURE 114 Simple – wrong – candidate tangles for the leptons. Antileptons are mirror tangles. The experimental mass values are also given. These wrong tangles yield the wrong mass sequence.

Lepton mass ratios

Mass calculations for leptons are as involved as for quarks. Each lepton, being a fermion, has a large family of associated tangles: there is a simplest tangle and there are the tangles that appear through repeated application of tail braiding. Despite this large tangle family, some results can be deduced from the simplest lepton tangles alone, disregarding the higher-order family members.

Mass calculations require first to find the correct lepton tangles. This is not straightforward. A consistent proposal appears when the set of lepton tangles reproduces all Feynman diagrams. This proposal is shown in Figure 115.

Both for neutrinos and for charged leptons, the progression in ropelength of the tight

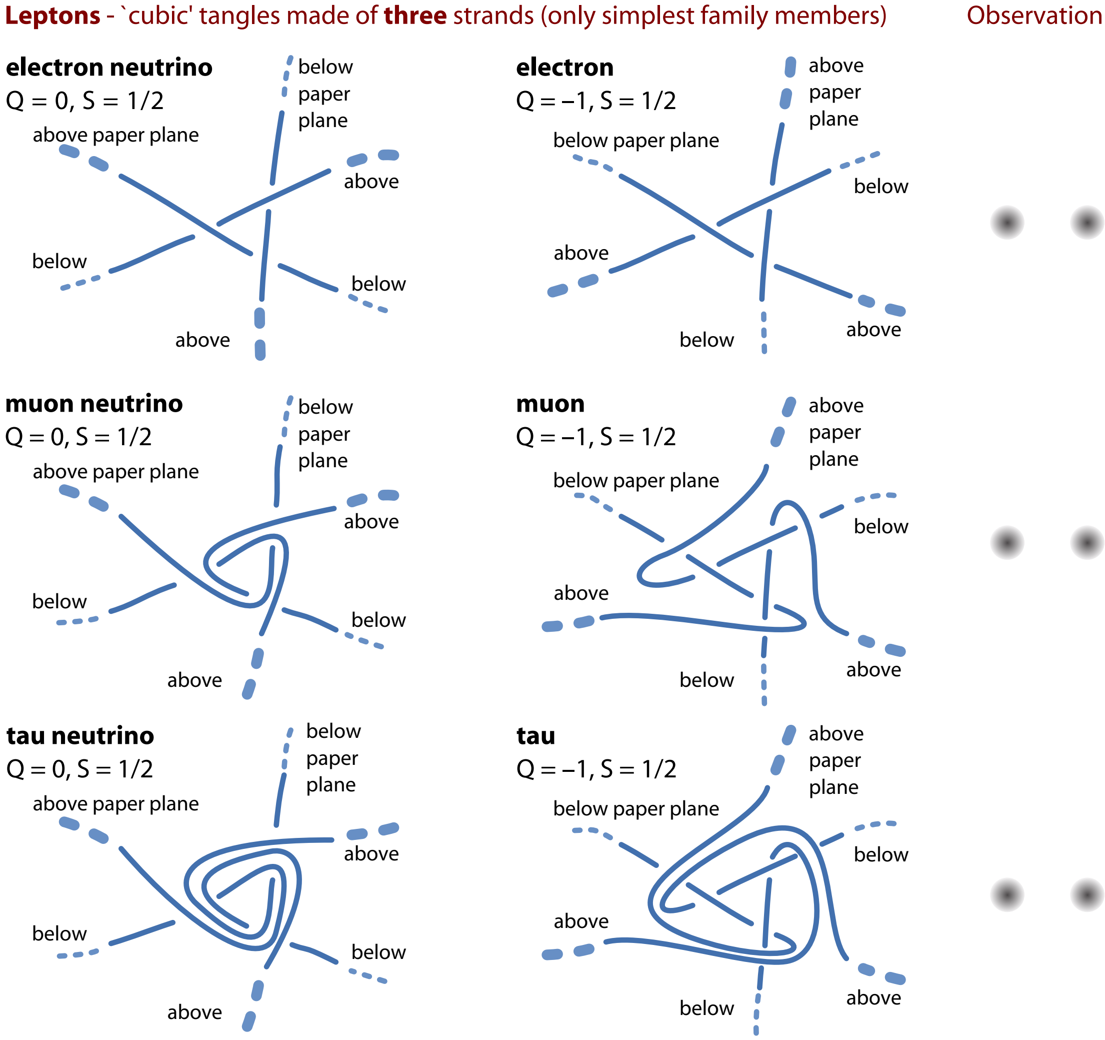

FIGURE 115 Presently proposed tangles for the leptons. Antileptons are mirror tangles. Feynman diagrams are reproduced correctly. The handedness of neutrinos is apparent. Mass sequences are as observed. The experimental mass values are also listed. At present, estimated mass values are within a few orders of magnitude.

versions of the basic tangles predicts a progression in their masses. This is indeed observed. In addition, the tangles of Figure 115 imply that the neutrinos are much less massive than the charged leptons. Spin 1/2 arises. Also the electric charge quantum number arises, as long as charge is related to topological chirality. And the handedness of neutrinos is evident. These tangle assignments are thus promising.

For each lepton tangle with l crossings, knot theory predicts a ropelength L that increases roughly proportionally to the crossing number: $L \sim l$. Each lepton mass value will again be given by the frequency of crossing switches due to rotations, including the belt trick, and of tail braiding. We thus expect a general relation of the type

$$m_l \sim e^{bL_l} \tag{199}$$

where b is a number of order 1 that takes into account the shape of the tangle core. Such a relation is in general agreement with the observed ratios between lepton masses.

We note that the lepton mass generation mechanism of the strand model differs from Page 330 other proposals in the research literature. It agrees with the Higgs mechanism but goes beyond it. For neutrinos, the mechanism contradicts the see-saw mechanism but confirms the Yukawa mechanism directly. From a distance, the mass mechanism of the Ref. 250 strand model also somewhat resembles conformal symmetry breaking.

In sort, research on lepton masses is ongoing; calculations of ropelengths and other geometric properties of the lepton tangles will allow a more detailed analysis. The main challenge remains to estimate the neutrino masses before experiments – such as KATRIN – determine them.

On the absolute values of particle masses

In nature, the masses of elementary particles are observed to be much lower than the Planck mass: the observed values lie between about 10^{-30} for neutrinos and 10^{-17} for the top quark. Particle masses are constant over space and time. Antiparticles have the same mass as particles. Gravitational and inertial masses are the same. Following the standard model, particle masses are due to the Higgs mechanism. Finally, elementary particles masses run with energy.

All qualitative observations about particle mass are reproduced by the strand model. However, the explanation of the numerical values is still lacking.

In the strand model, the gravitational mass of elementary particles is due to disturbance of the vacuum, in particular to the disturbance of the vacuum fluctuations. Larger masses are due to more complex tangles. Since rest mass is localized energy, rest mass is due to crossing switches per time. Larger masses have more crossing switches per time than lower masses.

In the strand model, the inertial mass of elementary particles is their reluctance to rotate. Inertial mass describes the relation between rotation frequency and wavelength; in other terms, inertial mass described the steepness of the helix drawn by the rotating phase arrow of a propagating particle. Larger masses have low steepness, smaller masses have higher steepness. Larger masses are due to more complex tangles.

As we just saw, the strand model predicts mass *sequences* and *ratios* of elementary particle masses that corroborate or at least do not contradict observations too much. The next step is to determine *absolute* mass values from the strand model. So far we only found that elementary particle masses are much smaller than a Planck mass. But to validate the strand model, we need more precise statements.

To determine gravitational mass values, we need to count those crossing switches that occur at rest; to determine inertial mass values, we need to look for crossing switches in the case of a moving particle – or, if we prefer, to understand the origin of the steepness of the helix drawn by the phase arrow. All these methods should first lead to mass value estimates and then to mass value calculations.

In general, the strand model reduces mass determination to the calculation of the details of a process: How often do the fluctuations of strands lead to crossing switches? There are various candidates for the crossing switches that lead to particle mass.

The first candidate for mass-producing crossings is tail switching. In general however, tail switching leads to different particle types. Only for the Higgs process, i.e., the addition of a full Higgs braid to a particle, is this process expected to be relevant. We can also say that tail braid addition are the strand model's version of the Yukawa coupling terms.

The next candidate for mass-producing crossings is the belt trick. The belt trick emits twisted tether pairs, thus virtual gravitons, and is a good candidate to understand gravitational mass. The belt trick also leads to core rotation and core displacement, which is the essence of inertial mass. Determining probabilities for the belt trick thus promises to allow mass calculations.

The third candidate for mass-producing crossings could appear when particles shed one strand and grab a new one. The influence of this process is not clear yet.

A fourth candidate for mass-producing crossings is the leather trick. However, the leather trick cannot be realized for strands that reach spatial infinity; therefore it is expected that it plays no role. In fact, the supposed effects of the leather tricks in an early phase of the strand model have now become effects due to the Higgs absorption and emission.

A fifth candidate for mass-producing crossings are those crossings that occur *above* or *around* the core, similar to the crossing that occur above the horizon of a black hole. This candidate group includes the belt trick, includes the Higgs mechanism, and thus is equivalent to the set of previous candidates.

It might be that some other mass-producing switching processes are being overlooked, but this seems unlikely. Therefore, in the following we explore the fifth candidate group in more detail, namely the crossings *around* a given tangle core.

Before looking for estimates, we note that in the past, various researchers have reached the conclusion that all elementary particle masses should be due to a common process or energy scale. Among theoretical physicists, the breaking of conformal symmetry has always been a candidate for such a process. Among experimental physicists, the Higgs mechanism – now confirmed by experiment – is the favourite explanation of all elementary particle masses. In the strand model, crossing switches around tangles are related to the Higgs boson. At the same time we can also argue that tangles break the conformal symmetry of vacuum. With a bit of distance, we can thus say that the strand model agrees with both research expectations.

Ref. 250, Ref. 251
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Let us continue with the quest for absolute mass estimates. In the strand model, *absolute* mass values are *not* purely geometric quantities that can be deduced directly from the shapes of tangle knots. Particle masses are due to *dynamical* processes. Absolute mass values are due to strand fluctuations; and these fluctuations are influenced by the core topology, the core shape, the core ropelength and core tightness.

To determine absolute particle mass values, we need to determine the ratio between the particle mass and the *Planck mass*. This means to determine the ratio between the crossing switch probability for a given particle and the crossing switch probability for a Planck mass, namely one switch per Planck time.*

* What is a Planck mass? In the strand model, a Planck mass corresponds to a structure that produces one

Energy is action per time. Mass is localized energy. In other words, the *absolute* mass of a particle is given by the average number of crossing switches it induces per time:

▷ Mass is crossing switch rate.

More precisely, the crossing switch rate of a particle at rest is its gravitational mass, and the crossing switch rate induced by propagation is its inertial mass. Let us explore the relations.

Given that mass is determined by the crossing switch rate, we deduce that particle mass values are determined by tangle topology, are fixed, are discrete, are positive, increase with tangle core complexity, are identical for particle and antiparticles, are constant over time, and are much smaller than the Planck mass. Because all these properties match observations, the local crossing switch rate indeed realizes all qualitative requirements for absolute particle mass values. We can thus proceed with the hope to learn more. In order to calculate absolute particle masses, we just need to determine the number of crossing switches per time that every particle tangle induces.

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One general way to perform a particle mass calculation is to use a *computer*, insert a strand model of the fluctuating vacuum plus the strand model of the particle under investigation, and count the number of crossing switches per time. The basis for one such approach, using the analogy of the evolution of a polymer in liquid solution, is shown in Figure 116. In contrast to polymers, also the change of strand length has to be taken into account. By determining, for a given core topology, the average frequency with which crossing switches appear for a tethered core, we can estimate the masses of the leptons, quarks and bosons. In such a mass calculation, the mass scale is set indirectly, through the time scale of the fluctuation spectrum. This is tricky but feasible. One would first need to find the parameter space and the fluctuation spectrum for which the polymer tangle follows the Schrödinger equation. Calculations with different tangles should then yield the different mass values. Such a simulation would also of interest for exploring the strand model of quantum mechanics.

A more precise computer simulation would also model the vacuum itself with strands. This approach would even allow to explore gravitational and inertial mass separately. In such a simulation, the particle mass appears when the helical motion of a tangle moving through a strand vacuum is observed. The required effort can be reduced by using the most appropriate computer libraries.

A further general way to determine particle masses is to search for *analytical approximations*. This is a fascinating conceptual and mathematical challenge. The main issue is to clarify which crossing switches contribute most to particle mass.

crossing switch for every Planck time, constantly, without interruption. But the strand model predicts that such structures do *not* appear as localized particles, because every localized particle – i.e., every tangle – has, by construction, a much smaller number of induced crossing switches per time. Following the strand model, elementary particles with Planck mass *do not exist*. This conclusion agrees with observation. But the strand model also implies that black holes with a Planck mass *do not exist*. Indeed, such Planck-scale black holes, apart from being extremely short-lived, have no simple strand structure. We can state that a Planck mass is never localized. Given these results, we cannot use a model of a *localized* Planck mass as a unit or a benchmark to determine particle masses.

The impossibility of using Planck mass as a unit is also encountered in everyday life: no mass measurement in any laboratory is performed by using this unit as a standard.

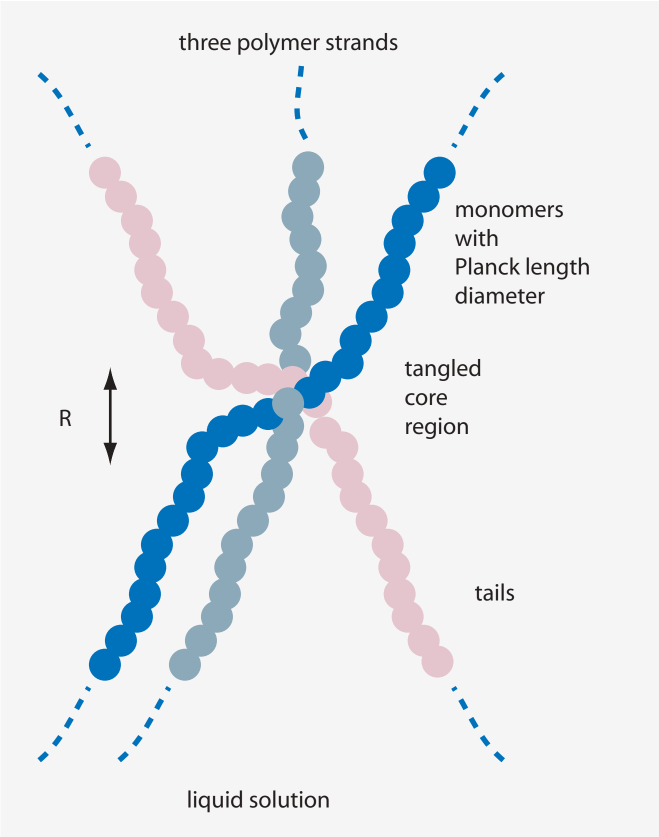

FIGURE 116 Determining lepton mass values with the help of a polymer analogy of strands. After rescaling, the probability of crossing switches around the tangle core yields an estimate for the mass of the elementary particle with that tangle.

Analytical estimates for particle masses

A first analytical attempt is the following. We assume that the inertial mass for a moving fermion is proportional to the fluctuation-induced appearance of the belt trick. If the tight core has a diameter of, say, three Planck lengths – and thus a circumference of around 9 Planck lengths – then the probability p of the belt trick for a particle with six tails will be in the range

$$p \approx (\mathrm{e}^{-9})^6 \approx 10^{-24} . \qquad (200)$$

This value would be the order of magnitude for the mass estimate, in Planck units. Such an estimate is only very rough, and the exponent can be quite different. Nevertheless, we do get an explanation for the large difference between the Planck mass and the typical fermion mass. A more precise analytical approximation for the belt trick probability – not an impossible feat – will therefore solve the so-called *mass hierarchy problem*. We thus want to know:

Challenge 205 r

▷ What is the numerical probability of the belt trick for a tethered core of given topology with fluctuating tails?

So far, several experts on polymer evolution have failed to provide even the crudest estimate for the probability of the belt trick in a polymer-tethered ball. Can you provide one?

A second analytical approach starts from the following question:

Challenge 206 r ▷ How often does a tail cross above the tangle core?

This question is loosely related to the previous one; in addition, this approach illustrates why complex cores have larger mass. The probability of such crossings, when squared, would be an estimate for the crossing switch rate, and thus for the particle mass. (There are additional details to the calculation.) We note directly that the number of tails will have a smaller impact on mass than the complexity of the tangle. So far, a reliable estimate for the crossing number, as a function of the tangle core properties, is still missing – even a crude one. Can you find one? May be the roughly linear relation observed between ropelength and (average) crossing number can be of help. Ref. 252

Open issues about mass calculations

Calculating absolute particle masses from tangle fluctuations, either numerically or with an analytical approximation, will allow the final check of the statements in this section. Challenge 207 e The strand model predicts that the resulting values will match experiments. For these calculations, it is essential that the tangle assignment for each elementary particle is correct. Finally, in 2019, the tangles for the W tangle and for the lepton tangles seem settled.

* *

Because the strand model predicts a lack of new physics beyond the standard model of particle physics, the calculation of neutrino masses, and thus their mass sequence, is one of the few possible *predictions* – in contrast to retrodictions – that are left over in the strand model.

* *

Challenge 208 s Is the mass of a tangle related to the vacuum density of strands?

* *

Challenge 209 s Do particle masses depend on the cosmological constant?

* *

The mass of an elementary particle does not depend on the spin direction. In particular, the W and Z bosons have equal longitudinal and transversal mass. The strand model does not allow an influence of spin orientation on mass.

* *

Challenge 210 s Can the concept of *total curvature* of a tangle help to calculate particle masses?

* *

Does the effect of tail braiding confirm the conjecture that every experiment is described by a small energy scale, determining the resolution or precision, and a large energy scale, less obvious, that determines the accuracy? Challenge 211 d

* *

If tail braiding is due to the weak interaction, and if the Higgs is a tail-braided vacuum, can we deduce that the Higgs interaction is a higher order effect of the weak interaction? Can we deduce a concrete experimental prediction from this relation?

On fine-tuning and naturalness

It has become fashionable, since about a decade, to state that the standard model of elementary particle physics is 'fine-tuned'. The term expresses several ideas. First of all, the extremely low value of the vacuum energy is not obvious when all the zero-point field contributions from the various elementary particles of the standard model are included. A low vacuum energy seems only possible if the masses and the particle types of the standard model are somehow interrelated. In other words, the term 'fine tuning' expresses, above all, the *lack of understanding* of the origin of the masses, mixings and coupling constants of elementary particles.

The term 'fine tuning' is also used to state that the universe would be *very different* if the fundamental constants would be different. But this statement lacks deep truth. In this usage, the term 'fine tuning' states that particle masses are *not* parameters that can be varied at will. In common usage, 'parameters' are variable constants; but the low value of the vacuum energy – as well as many other observations – shows that the masses of elementary particles *cannot* be varied without destroying the validity of the standard model of particle physics.

Some people suggest that 'fine-tuning' implies that the standard model of particle physics is 'unnatural', whatever this might mean in detail. Some even suggest that the parameters of the standard model lack any explanation. The strand model – but also common sense – show that this suggestion is false.

The strand model *naturally* has a low vacuum energy, because the unknotted strands of flat space naturally have a zero energy density, and the particle masses, mixings and coupling constants are not variable or random, but *naturally* unique and fixed in value. Any *correct* description of nature must be 'fine-tuned'. If the standard model would not be 'fine-tuned', it would not describe nature.

In short, the fashionable term 'fine-tuned' is equivalent to the terms 'unmodifiable' and 'hard to vary' that were discussed above. Page 166 All these terms highlight the lack of alternatives to the world as we observe it, the existence of explanations for the processes around us, and our ability to discover and grasp them. This is part of the wonders of nature. And the strand model makes those wonders apparent at the Planck scale.

Summary on elementary particle masses and millennium issues

The strand model implies that masses are dynamic quantities fixed by processes due to the geometric and topological properties of specific tangle families. As a result, strands explain why the masses of elementary particles are not free parameters, but fixed and unique constants, and why they are much smaller than the Planck mass by many orders of magnitude. Strands also reproduce all known qualitative properties of particle masses. Page 371

Strands provide estimates for a number of elementary particle *mass ratios*, such as m_W/m_Z and m_{Higgs}/m_W. Most quark and lepton mass sequences and first rough estimates of mass ratios appear to agree with experimental data. All hadron mass sequences are predicted correctly.

The strand model also promises to calculate *absolute mass values*, including their change or 'running' with energy. Particle masses are in the range of 10^{-20} times the Planck mass, near the observed mass values. In the future, more precise calculations will allow either improving the match with observations or refuting the strand model.

The results are encouraging for two reasons. First of all, no other unified model that agrees with experiment explains the qualitative properties of mass and mass sequences. Secondly, no research on statistical rational tangles exists; an understanding of the parameters of nature might be lacking because results in this research field are lacking.

In the millennium list of open issues we have thus caught a glimpse of how to settle Page 164 the origin of particle masses – though we have not calculated them yet. Because further interesting challenges are awaiting us, we continue nevertheless. In the next leg, we investigate how elementary particle states mix.

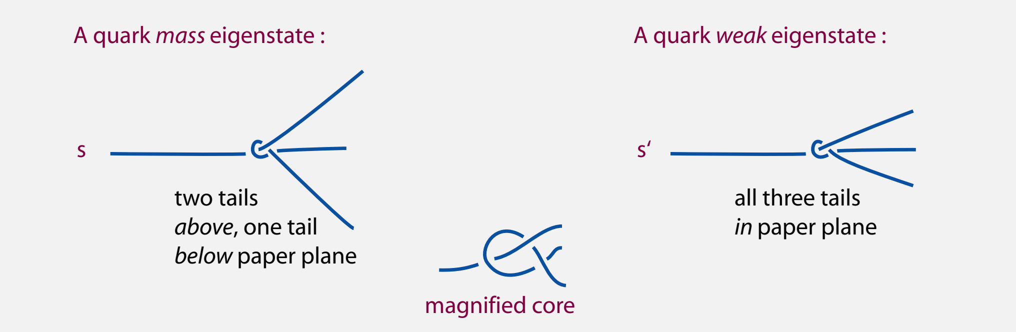

FIGURE 117 Tail shifting leads to quark mixing: mass eigenstates and weak eigenstates differ.

Mixing angles

In nature, the *mass* eigenstates for fermions differ from their *weak* eigenstates: quarks mix among themselves, and so do neutrinos. Quarks also show CP violation; for neutrinos, the issue is still open. These effects are described by two so-called *mixing matrices*. The two mixing matrices contain fundamental constants of nature. For the strand model to be correct, it must allow calculating the measured values of all components of the two mixing matrices.

Quark mixing – the experimental data

In nature, the quark mass eigenstates and their weak eigenstates differ. This difference was discovered in 1963 by Nicola Cabibbo and is called *quark mixing*. The values of the
Ref. 233 elements of the quark mixing matrix have been measured in many experiments, and more experiments aiming to increase the measurement precision are under way.

Vol. V, page 251 The quark *mixing matrix*, also called CKM mixing matrix, is defined by

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = (V_{ij}) \begin{pmatrix} d \\ s \\ b \end{pmatrix} . \tag{201}$$

where, by convention, the states of the +2/3 quarks u, c and t are unmixed. Unprimed quarks names represent strong (and electromagnetic) eigenstates, primed quark names
Ref. 233 represent weak eigenstates. In its standard parametrization, the mixing matrix reads

$$V = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \tag{202}$$

where $c_{ij} = \cos\theta_{ij}$, $s_{ij} = \sin\theta_{ij}$ and i and j label the generation $(1 \leqslant i, j \leqslant 3)$. The mixing matrix thus contains three mixing angles, θ_{12}, θ_{23} and θ_{13}, and one phase, δ. In the limit $\theta_{23} = \theta_{13} = 0$, i.e., when only *two* generations mix, the only remaining parameter is the angle θ_{12}, called the *Cabibbo angle*; this angle is Cabibbo's original discovery. The last parameter, the so-called *CP-violating phase* δ, by definition between 0 and 2π, is measured to be different from zero; it expresses the observation that CP invariance is violated in the case of the weak interactions. The CP-violating phase only appears in the third column of the matrix; therefore CP violation requires the existence of (at least) three generations.

Challenge 212 e

The present 90 % confidence values for the measured *magnitude* of the complex quark mixing matrix elements are

Ref. 233

$$|V| = \begin{pmatrix} 0.97427(14) & 0.22536(61) & 0.00355(15) \\ 0.22522(61) & 0.97343(15) & 0.0414(12) \\ 0.00886(33) & 0.0405(12) & 0.99914(5) \end{pmatrix} . \quad (203)$$

All these numbers are unexplained constants of nature, like the particle masses. Within experimental errors, the matrix V is unitary.

A huge amount of experimental work lies behind this short summary. The data have been collected over many years, in numerous scattering and decay experiments, by thousands of researchers. Nevertheless, this short summary represents all the data that any unified description has to reproduce about quark mixing.

Quark mixing – explanations

In the standard model of particle physics, the quark mixing matrix is usually seen as due to the coupling between the vacuum expectation value of the Higgs field and the left-handed quark doublets or the right handed quark singlets. However, this description does not lead to a numerical prediction.

A slightly different description of quark mixing is given in the strand model. In the strand model, the Higgs field and its role as mass generator and unitarity maintainer is a special case of the process of tail braiding. And braiding is related to the weak interaction. Because the various quarks are differently tangled rational tangles, tail braiding can reduce or increase the crossings in a quark tangle, and thus change quark flavours. We thus deduce from the strand model that quark mixing is an automatic result of the strand model and related to the weak interaction. We also deduce that quark mixing is due to the *same* process that generates quark masses, as expected. But we can say more.

Page 330
Page 253

In the strand model, the *mass eigenstate* – and colour eigenstate – is the tangle shape in which colour symmetry is manifest and in which particle position is defined. The mass eigenstates of quarks correspond to tangles whose three colour-tails point in three directions that are equally distributed in space. The shape in which the tails point in three, equally spaced directions is the shape that makes the SU(3) representation under core slides manifest.

In contrast, the *weak eigenstates* are those shapes that makes the SU(2) behaviour of core pokes manifest. For a quark, the weak eigenstate appears to be that shape of a tangle for which all tails lie in a plane; for such plane configuration, the tails and the core mimic a belt and its buckle, the structure that generates SU(2) behaviour. The two types of eigenstates are illustrated in Figure 117.

In the strand model, masses are dynamical effects related to tangle shape. In the case of quarks, the two configurations just mentioned will thus behave differently. We call the transformation from a mass eigenstate to a weak eigenstate or back *tail shifting*. Tail shifting is a deformation: the tails as a whole are rotated and shifted. On the other hand, tail shifting can also lead to untangling of a quark tangle; in other words, tail shifting can lead to tail braiding and thus can transform quark flavours. The process of tail shifting can thus explain quark mixing. (Tail shifting also explains the existence of neutrino mixing, and the lack of mixing for the weak bosons.)

Tail shifting can thus be seen as a *partial* tail braiding; as such, it is due to the weak interaction. This connection yields the following predictions:

Page 210
Ref. 233

— Tail shifting, both with or without tail braiding at the border of space, is a generalized deformation. Therefore, it is described by a unitary operator. The first result from the strand model is thus that the quark mixing matrix is unitary. This is indeed observed.
— For quarks, tail braiding is a process with small probability. As a consequence, the quark mixing matrix will have its highest elements on the diagonal. This is indeed observed.
— Tail shifting also naturally predicts that quark mixing will be higher between neighbouring generations, such as 1 and 2, than between distant generations, such as 1 and 3. This is also observed.
— The connection between mixing and mass also implies that the 1–2 mixing is stronger than the 2–3 mixing, as is observed.
— Finally, tail shifting predicts that the numerical values in the quark mixing matrix can be deduced from the difference between the shapes of the two kinds of tangles shown in Figure 117. In particular, tail shifting also predicts that the quark mixing angles change, or run, with energy. In addition, the effect is predicted to be small. On the other hand, so far there is no reliable experimental data on the effect.

Performing a precise calculation of mixing angles and their running with energy is still a subject of research.

A challenge

Ref. 253

Can you deduce the approximate expression

$$\tan\theta_{\mathrm{u\,mix}} = \sqrt{\frac{m_{\mathrm{u}}}{m_{\mathrm{c}}}} \qquad (204)$$

Challenge 213 r

for the mixing of the up quark from the strand model?

CP violation in quarks

The CP violating phase δ for quarks is usually expressed with the *Jarlskog invariant*, defined as $J = \sin\theta_{12}\sin\theta_{13}\sin\theta_{23}^2\cos\theta_{12}\cos\theta_{13}\cos\theta_{23}\sin\delta$. This involved expression is independent of the definition of the phase angles and was discovered by Cecilia Jarlskog, an important Swedish particle physicist. Its measured value is $J = 3.06(21)\cdot 10^{-5}$.

Ref. 233

Because the strand model predicts three quark generations, the quark model implies the possibility of CP violation. In the section on mesons we have seen that the strand model actually predicts the existence CP violation. In particular, Figure 99 shows that

Page 338
Page 335

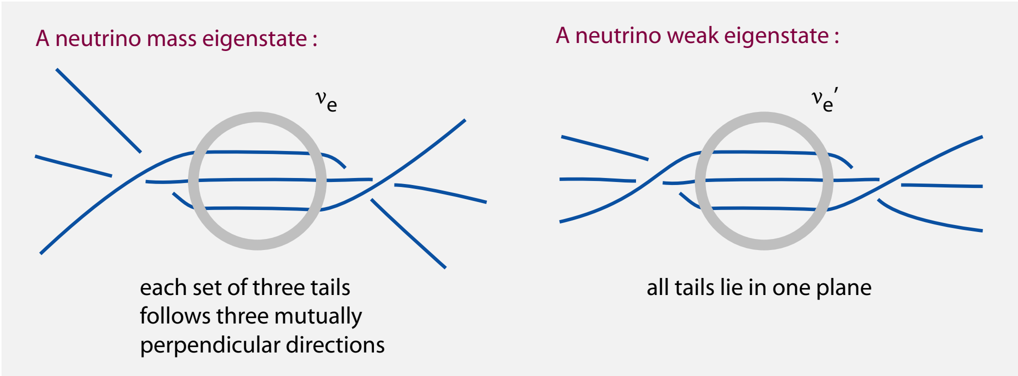

FIGURE 118 Tail shifting leads to neutrino mixing: mass eigenstates and weak eigenstates differ.

with the help of tail shifting, K^0 and $\overline{K}^0$ mesons mix, and that the same happens with certain other neutral mesons. Figure 100 shows a further example. As just mentioned, the possibility of tail shifting implies that CP violation is small, but non-negligible – as is observed.

The strand model thus predicts that the quark mixing matrix has a non-vanishing CP-violating phase. The value of this phase is predicted to follow from the geometry of the quark tangles, as soon as their shape fluctuations are properly accounted for. This topic is still a subject of research.

Neutrino mixing

The observation, in 1998, of neutrino mixing is comparably recent in the history of particle physics, even though the important physicist Bruno Pontecorvo predicted the effect already in 1957. Again, the observation of neutrino mixing implies that also for neutrinos the mass eigenstates and the weak eigenstates differ. The values of the mixing matrix elements are only known with limited accuracy so far, because the extremely small neutrino mass makes experiments very difficult. Experimental progress across the world is summarized on the website www.nu-fit.org. The absolute value of the so-called PMNS mixing matrix U is

Ref. 254
Ref. 233

$$|U| = \begin{pmatrix} 0.82(1) & 0.54(2) & -0.15(3) \\ -0.35(6) & 0.70(6) & 0.62(6) \\ 0.44(6) & -0.45(6) & 0.77(6) \end{pmatrix} . \qquad (205)$$

Again, these numbers are unexplained fundamental constants of nature. Within experimental errors, the matrix U is unitary. The mixing among the three neutrino states is strong, in contrast to the situation for quarks. Neutrino masses are known to be positive; however, present measurements are not precise and only yield values of the order of 1 ± 1 eV.

In the strand model, the lepton mass eigenstates correspond to tangles whose tails point along the three coordinate axes. In contrast, the weak eigenstates again correspond to tangles whose tails lie in a plane. The two kinds of eigenstates are illustrated in Figure 118. Again, the transition between the two eigenstates is due to tail shifting, a special kind of strand deformation.

We thus deduce that neutrino mixing, like quark mixing, is an automatic result of the strand model and is related to the weak interaction. Given that the neutrino masses are small and similar, and that neutrinos do not form composites, the strand model predicts that the mixing values are large. This is a direct consequence of tail shifting, which in the case of similar masses, mixes neutrino tangles leads to large mixings between *all* generations, and not only between neighbouring generations. In the strand model, the large degree of neutrino mixing is thus seen as a consequence of their low and similar masses, of their tangle structure, and of their existence as free particles.

Like for quarks, the strand model predicts a *unitary* mixing matrix for neutrinos. The strand model also predicts that the geometry of the neutrino tangles and their fluctuations will allow us to calculate the mixing angles. More precise predictions are still subject of research.

CP violation in neutrinos

The strand model predicts that the three neutrinos are massive Dirac particles, not Majorana particles. This has not yet been confirmed by experiment. The strand model thus predicts that the neutrino mixing matrix has *only one* CP-violating phase. (It would have three such phases if neutrinos were Majorana particles.) The value of this phase is predicted to follow from the neutrino tangles and a proper accounting of their fluctuations. Also this calculation is still a subject of research.

On the one hand, the strand model suggests the appearance of CP violation in neutrinos. On the other hand, it is unclear when the value of the CP-violating phase will ever be measured with sufficient precision. This is one of the hardest open challenge of experimental particle physics.

The mechanism of CP violation has important consequences in cosmology, in particular for the matter–antimatter asymmetry. Since the strand model predicts the absence of the see-saw mechanism, the strand model rules out leptogenesis, an idea invented to explain the lack of antimatter in the universe. The strand model is more on the line with electroweak baryogenesis.

Ref. 255
Ref. 256

Open challenge: calculate mixing angles and phases ab initio

Calculating the mixing angles and phases ab initio, using the statistical distribution of strand fluctuations, is possible in various ways. In particular, it is interesting to find the relation between the probability for a tail shift and for a tail braiding. This will allow checking the statements of this section.

Challenge 214 ny

Because the strand model predicts a lack of new physics beyond the standard model of particle physics, the calculation of neutrino mixing angles is one of the few possible *predictions* that are left over in fundamental physics. Since the lepton tangles are still tentative, a careful investigation is necessary.

One possibility is that only the electron neutrino tangle given above is correct, and that the other two neutrinos are similar to it, just with more built-in torsion. Figure 119 illustrates this possibility. If this assignment were correct, two of the mixing angles should be large (and maybe have the zero-order value of 120°/3=40°). In addition, very low mass

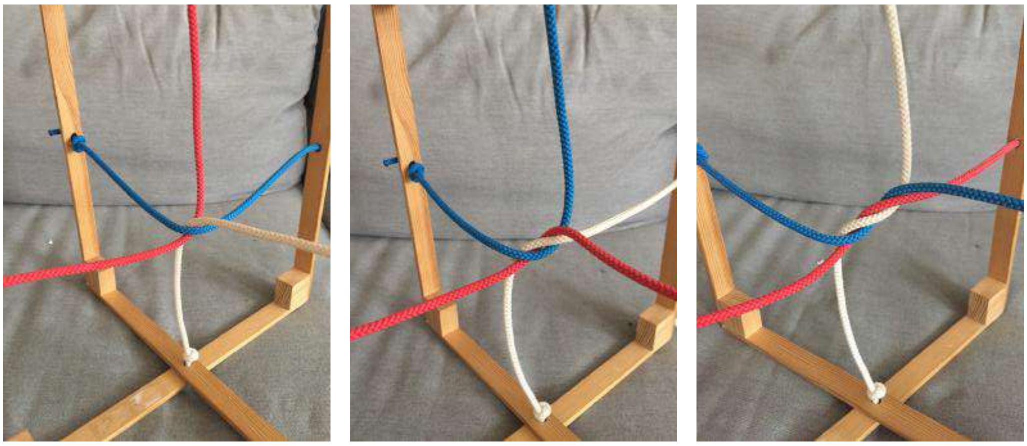

FIGURE 119 An alternative candidate assignment for the three neutrino tangles that generates large mixing between neighbouring generations and strong preference for one handedness.

values would arise naturally, in normal ordering. This tangle assignment would also explain the difficulty of observing neutrinos of opposite handedness: neutrinos would have an extremely high preference for one handedness (belt-trick); the mirror images of the tangles in Figure 119 would correspond to antineutrinos. Despite these appealing aspects, this tentative assignment has one unclear issue: explaining the lack of a fourth neutrino generation is not straightforward.

Summary on mixing angles and the millennium list

The strand model implies that mixing angles for quarks and neutrinos are properties of their tangle families. The existence of mixing is due to the shape of tangles and their fluctuations. As a result, strands explain why mixing angles are not free parameters, but discrete and unique constants of nature. The strand model also predicts that mixing angles are constant during the evolution of the universe.

We have shown that tangles of strands predict non-zero mixing angles for quarks and neutrinos, as well as CP-violation in both cases. The strand model also predicts that the mixing angles of quarks and neutrinos can be calculated from strand fluctuations. Strands predict that mixing matrices are unitary and that they run with energy. Strands also predict a specific sequence of magnitudes among matrix elements; the few predictions so far agree with the experimental data. Finally, the strand model rules out leptogenesis.

We have thus partly settled four further items from the millennium list of open issues. Page 164 All qualitative aspects and some sequences are reproduced correctly, but no hard quantities were deduced yet. The result is somewhat disappointing, but it is also encouraging. At present, no other explanation for quark and neutrino mixing is known. Future calculations will allow either improving the checks or refuting the strand model. We leave this topic unfinished and proceed to the most interesting topic that is left: understanding the coupling constants.

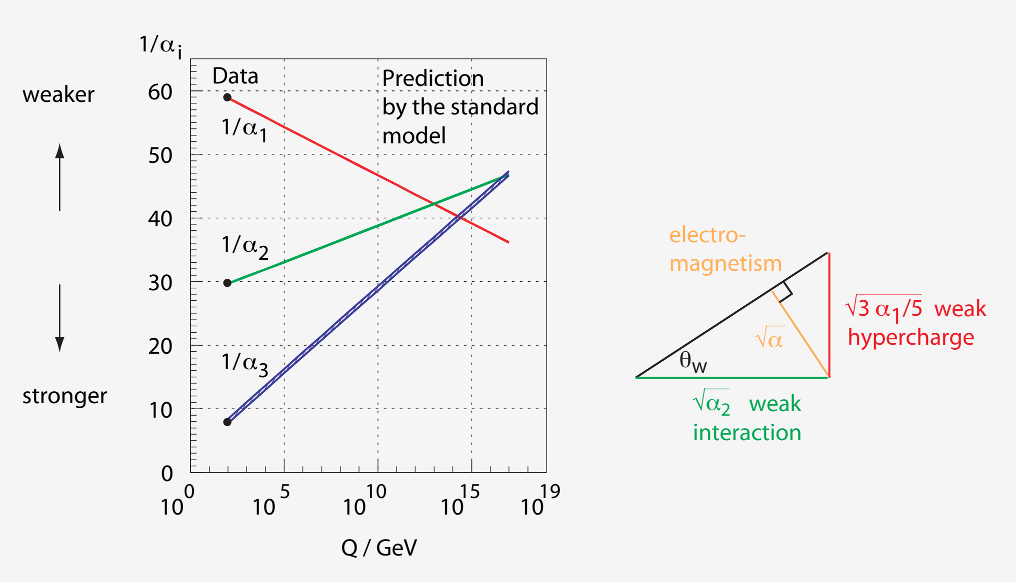

FIGURE 120 Left: How the three coupling constants (squared) change with energy, as predicted by the standard model of particle physics; the graph shows the constant $\alpha_1 = \frac{5}{3}\alpha/\cos^2\theta_W$ for the weak hypercharge coupling (related to the electromagnetic fine structure constant α through the weak mixing angle θ_W and a historical factor 5/3 that is useful in grand unification), the coupling constant $\alpha_2 = \alpha_w = \alpha/\sin^2\theta_W$ for the weak interaction, and the coupling constant $\alpha_3 = \alpha_s$ for the strong interaction. The three black points are measurement points; at lower and slightly higher energies, data and calculation match within experimental errors. (Courtesy Wim de Boer) Right: The relation between the coupling constants α for the electromagnetic $U(1)_{EM}$, $\alpha_2 = \alpha_w$ for the weak SU(2), α_1 for the weak hypercharge $U(1)_Y$ gauge groups and the weak mixing angle θ_W.

COUPLING CONSTANTS AND UNIFICATION

In nature, electric, weak and strong charge are *quantized*. No experiment has ever found even the smallest deviation from charge quantization. *All charges in nature are integer multiples of a smallest charge unit.* Specifically, the electric charge of every free particle is observed to be an integer multiple of the positron electric charge. We call the integer the *electric charge quantum number*.

In nature, the *strength* of a gauge interaction for a unit charge is described by its coupling constant. The coupling constant gives the probability with which a unit charge emits a virtual gauge boson, or, equivalently, the average phase change produced by the absorption of a gauge boson. There are three charge types and three coupling constants: for the electromagnetic, for the weak and for the strong interaction. All particles with a given charge type and value share the same coupling constant, even if their masses differ. The three coupling constants depend on energy. The known data and the change with energy predicted by the standard model of particle physics are shown in Figure 120.

In nature, the *fine structure constant* α, i.e., the electromagnetic coupling constant, at the lowest possible energy, 0.511 MeV, has the well-known measured value Ref. 5

$$\alpha = 1/137.035\,999\,139(31)\ . \qquad (206)$$

Equivalently, the electromagnetic coupling of the positron can also be described by the equivalent number

$$\sqrt{\alpha} = 1/11.706\,237\,6167(13) = 0.085424543114(10)\,, \qquad (207)$$

which is also called the *electric charge unit* (at low energy). Quantum electrodynamics predicts the precise change with energy of this charge unit; the experiments performed so far, up to over 100 GeV, agree with this prediction. Quantum electrodynamics also predicts that the charge unit, when extrapolated right up to the Planck energy, would have a value of 1/10.2(1). These predictions are shown, in a common, but somewhat scrambled way, in Figure 120.

Explaining the value of α, which determines all colours and all material properties in nature, is the most famous millennium issue. If the strand model cannot reproduce every observation about α and the other coupling constants, it is wrong. In particular, we thus need to understand, using the strand model, the quantization of charges on the one hand, and the origin of the mysterious value of the charge unit – either at low energy or at Planck energy – on the other hand.

Interaction strengths and strands

In the strand model, *all three gauge interactions are due to shape changes of tangle cores.* We first classify the possible shape changes. Given a tangle core, the following shape changes can occur:

— *Small* changes of core shape do not produce any crossing switch. Small shape changes thus have no physical significance: for a given observer, they leave all observables unchanged.
— *Twist* shape changes of a strand segment in the core produce an *electric field*, if the particle is charged. More precisely, the electric field around a particle is the difference between the average number p_{tr} of right twists and the average number p_{tl} of inverse, left twists that a particle tangle produces per unit time.
— *Poke* shape changes of a strand segment in the core produce a *weak interaction field*. More precisely, the weak field is the asymmetry among the probabilities $p_{\mathrm{p}x}$, $p_{\mathrm{p}y}$ and $p_{\mathrm{p}z}$ for the three fundamental poke types and their inverses.
— *Slide* shape changes of a strand segment in the core produce a *colour field*, if the particle has colour. More precisely, the colour field is the asymmetry among the probabilities $p_{\mathrm{s}1}$ to $p_{\mathrm{s}8}$ for the eight fundamental slide types and their inverses.
— A combination of these moves can also appear.

In the strand model, the fluctuation probabilities for each Reidemeister move – twist, poke or slide – determine the coupling constants. We thus need to determine these probability values. We can directly deduce a number of conclusions, without any detailed calculation:

— The coupling constants are not free parameters, but are specified by the geometric, *three-dimensional shape* of the particle tangles.
— By relating coupling constants to shape fluctuation probabilities, the strand model

predicts that coupling constants are *positive* numbers and *smaller than 1* for all energies.

▷ $\alpha < 1$,

▷ $\alpha_w < 1$,

▷ $\alpha_s < 1$

This is indeed observed.

— A still stricter bound for coupling constants can also be deduced. The sum of all possible fluctuations for a particular tangle has unit probability. We thus have

$$1 = p_{\text{small}} + p_{\text{tr}} + p_{\text{tl}} + \sum_{w=x,y,z} (p_{\text{p}w} + p_{\text{p}-w}) + \sum_{g=1}^{8} (p_{\text{s}g} + p_{\text{s}-g}) + p_{\text{combination}} . \quad (208)$$

The strand model thus predicts that the *sum* of the three charge units must be *strictly smaller than 1,* for every energy value. This is easily checked, both with the data and with the prediction of quantum field theory. In quantum field theory, the three (modified) coupling constants are given, as a function of energy, in the popular graph shown in Figure 120. The values are a combination of experimental data – for low energies – and theoretical extrapolations – for high energies. In this popular graph, the electromagnetic coupling is traditionally multiplied by $5/(3\cos^2\theta_W)$. (This is done in order to test grand unification; we keep the traditional factor, even though grand unification is shown by experiment and predicted by the strand model not to apply to nature.) The graph allows us to confirm that the sum of the three unmodified charge units is indeed smaller than 1 for all energy values, as predicted by the strand model.

— The strand model also predicts that the three coupling constants are related by small numbers, as the corresponding fluctuations differ only in the number of involved strands. This is also observed, as Figure 120 shows – especially if we remember that the couplings are the square roots of the values shown in the graph, corrected for the traditional factor.

— The strand model further predicts that the coupling constants are independent of time and space, and that in particular, they do not depend on the age of the universe. This is also observed, despite occasional claims to the contrary. Ref. 257

— Finally, strand model predicts that the coupling constants are the same for particles and antiparticles, as is observed.

In summary, the strand model implies, like quantum field theory, that *coupling constants are probabilities.* In addition, strands imply

▷ $\alpha + \alpha_w + \alpha_s < 1$ and $\sqrt{\alpha} + \sqrt{\alpha_w} + \sqrt{\alpha_s} < 1$.

Despite the agreement with experiment, we have not deduced any deep result yet – except one.

Strands imply unification

In fact, one new point is made by the strand model. Each gauge interaction is due to a different Reidemeister move. However, given a specific tangle core deformation, different observers will classify the deformation as a different Reidemeister move. Indeed, *every* Reidemeister move can be realized by the *same* deformation of a single strand: for each Reidemeister move, it is sufficient to add a curved section to a straight strand segment. Such a deformation can look like a type I Reidemeister move for one observer, like a type II move for another, and like a type III move for a third one.
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Because all interactions follow from the same kind of strand deformation of tangle cores, the strand model thus provides *unification* of the interactions. This result is new: in fact, this kind of strand unification of the interactions differs completely from any other approach ever proposed. And in contrast to several other approaches, strand unification does *not* require that the three coupling constants have the same value at high energy.

A given shape deformation thus has five probabilities associated to it: the probabilities describe what percentage of observers sees this deformation as a type I move, as a type II move, as a type III move, as a combination of such moves, or as no move at all, i.e., as a small move without any crossing switch. On the other hand, at energies measurable in the laboratory, the moves can *almost always* be distinguished, because for a given reaction, usually all probabilities but one practically vanish, due to the time averaging and spatial scales involved.* In short, at energies measurable in the laboratory, the three gauge interactions almost always differ.
Challenge 215 e

Calculating coupling constants

The strand model predicts that the calculation of the three coupling constants is a problem of tangle geometry and fluctuation statistics. Thus it can be approached, at each energy scale, either *analytically* or with *computer calculations*. The calculations need to determine the probabilities of the corresponding Reidemeister moves. If the results do not agree with the experimental values, the strand model is false. We note that there is no freedom to tweak the calculations towards the known experimental results.

In particular, in the strand model, one way to proceed is the following. The (square root of the) fine structure constant is the probability for the emission of twists by a fluctuating chiral tangle.

▷ The strand model predicts that the fine structure constant can be calculated by determining the probability of twists, i.e., Reidemeister I moves, in the *fluctuating tangle shapes* of a given particle with nonzero electrical charge.

In other words, the strand model must show that the probability of the first Reidemeister move in chiral particle tangles is *quantized*. This probability must be an integer multiple of a unit that is common to all tangles; and this coupling unit must be the fine structure

* The strand model thus appears to predict that at extremely high energy, meaning near the Planck energy, for each gauge interaction, also particles with zero charge can interact. At Planck energy, when horizons form, the time averaging is not perfect, and interactions become possible even with zero charge. But then the particle concept make little sense at those energies.

constant. Any check for the existence of a coupling unit requires the calculation of twist appearance probabilities for *each* chiral particle tangle. The strand model is only correct if all particles with the *same* electric charge yield the *same* twist emission probability.

Instead of emission, also absorption can be used to calculate the fine structure constant:

▷ The strand model predicts that the fine structure constant can be calculated from the average angle that a tangle core rotates when absorbing a photon.

We will pursue this alternative shortly.

So far, there do not seem to exist any analytical tool that permits the calculation of shape deformation probabilities. Thus, at present, computer calculations seem to be the only possible choice. Of all existing software programs, the most adapted to calculating fluctuation probabilities are the programs that simulate the dynamics of tangled polymers; but also the programs that simulate the dynamics of cosmic strings or the dynamics of helium vortices are candidates. The main issue, apart from a large computer time, is
Challenge 216 r
the correct and self-consistent specification of the shape fluctuation distribution at each energy scale.

In summary, using the strand model we expect to be able to calculate the electromagnetic coupling constant and to understand its validity across all elementary particles. The same expectation obviously also holds for the two nuclear interactions. If any of the expectations on tangle interactions are found to be incorrect, the strand model is false. *The strand model must yield quantized tangle equivalence classes for the electromagnetic, weak and colour charge.* Even though the calculation issues are still subject of research, there are encouraging hints that these expectations will be validated.

First hint: the energy dependence of physical quantities

In nature, all effective charges, i.e., the coupling constants, change with energy. One also
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says that they *run* with energy. Figure 120 shows the details. Running also occurs for masses and mixing angles. All other intrinsic particle properties, such as spin, parities and all other quantum numbers, are found *not* to change with energy. For the coupling constants, the measured changes between everyday energy and about 100 GeV agree with the prediction from quantum field theory.*

The strand model predicts

▷ Coupling constants, masses and mixing angles change with energy because they are quantities that *depend* on the average *geometrical details,* and in particular, on the *scale* of the underlying particle tangles.

* In the standard model of particle physics, the running of the electromagnetic and weak coupling constants
Ref. 229
– the slope in Figure 120 – depends on the number of existing Higgs boson types. The (corrected) strand
Page 330
model predicts that this number is one. Measuring the running of the constants thus allows checking the number of Higgs bosons. Unfortunately, the difference is small; for the electromagnetic coupling, the slope changes by around 2 % if the Higgs number changes by one. But in future, such a measurement accuracy might be possible.

More precisely, the running quantities depend on the fluctuations of the geometric tangle shapes, and these fluctuations depend somewhat on the spatial and thus the energy scale under consideration. We note that the strand model predicts a running *only* for these three types of observables; all the other observables – spin, parities or other quantum numbers – are predicted to depend on the *topology* of the particle tangles, and thus to be *independent* of energy. This prediction agrees with observation. Therefore, we can now explore the details of the running.

Second hint: the running of the coupling constants at low energy

The strand model proposes a new view on the screening and antiscreening effects that are part of quantum field theory. In the strand model, screening effects are consequences of the statistics of shape deformations for loose tangle cores that are embedded into the strands that form the vacuum. Since these statistical effects can in principle be calculated, it is expected that such calculations can be compared with the predictions of quantum field theory shown in Figure 120. This check is in progress. A few results, however, can be deduced without any calculations at all.

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In the strand model, the electromagnetic interaction is due to the first Reidemeister move, the twist. For a charged particle – thus one with a chiral tangle core – the average difference in the occurrence of right and left twists determines the effective charge. It is expected that this difference *decreases* when the strand core is loose, because the loose strands are more similar to those of the surrounding vacuum, so that the differences due to the chirality of the tangle will be washed out. In the language of quantum field theory, the virtual particle-antiparticle pairs – created by the fluctuations of the vacuum strands – screen the central, naked charge. The screening is reduced when the energy is increased, and thus when the scales are reduced. In other words, the strand model predicts that the electromagnetic coupling *increases with energy*, as is observed:

$$\triangleright \quad \frac{\mathrm{d}\alpha}{\mathrm{d}E} > 0 \, .$$

Also for the two nuclear interactions, the washing out effect for loose tangle cores by the vacuum does occur as predicted by quantum field theory. In the weak interaction, the *antiscreening* of the weak charge appears in this way. In the strong interaction, both virtual quark–antiquark pairs and virtual gluon pairs can appear from the strands that make up the vacuum. Virtual quark–antiquark pairs lead to screening, as virtual electron–antielectron pairs do for the electromagnetic interaction. In addition, however, we have seen that the strand model of mesons implies that virtual gluon pairs lead to antiscreening. (In contrast, virtual photon pairs do not lead to such an effect.) Because the strand model fixes the number of quark and gluons, the strand model is consistent with the result that the screening of the colour charge by quark pairs is overcompensated by the antiscreening of the virtual gluon pairs.

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In other words, the strand model reproduces the observed signs for the slopes of the coupling constants in Figure 120, for the same reason that it reproduces the quantum field theoretic description of the three gauge interactions. The predicted running could also be checked quantitatively, by taking statistical averages of tangle fluctuations of varying dimension. This is a challenge for future research.

Third hint: further predictions at low energy

As we just saw, the complete explanation of the running of the couplings depends on the explicit boson and fermion content of nature and on the fact that the strand model reproduces quantum field theory. Interestingly, the strand model also proposes a simpler, though less precise explanation of the running.

At energies much smaller than the Planck energy, such as everyday energies, the strand model implies that the average size of the tangle core is of the order of the position uncertainty of a particle. In other words, any thickness of the strands – real or effective – can be neglected at low energies. Therefore, at low energies, the average strand length within a particle tangle core is also of the order of the de Broglie wavelength. Low, everyday energy thus implies *large*, *loose* and *spherical/ellipsoidal* tangle cores.

At low energies, shape fluctuations can lead to any of the three Reidemeister moves. The probabilities of such shape deformations will scale with some power of the average strand length within the tangle core. In other words, coupling constants depend on energy. But how exactly?

We note directly that *higher* Reidemeister moves, which involve larger numbers of strand segments, will scale with *larger* power values. In particular, the longer the strand in the core – i.e., the lower the energy – the more the relative probability for the higher Reidemeister moves will increase.

In summary, the strand model predicts that when a tangle is loose and long, i.e., when energies are low, the strong nuclear interaction, due to the third Reidemeister move, is the strongest gauge interaction, followed by the weak nuclear interaction, due to the second Reidemeister move, in turn followed by the electromagnetic interaction:

$$\triangleright \quad \alpha_{\mathrm{em}} < \alpha_{\mathrm{w}} < \alpha_{\mathrm{s}} \; .$$

The prediction matches observations. Unfortunately, this argument is not reliable. If the strand number were the *only* cause of the running, the argument would imply that the three slopes for the running of the three coupling constants should behave like 3:2:1. However, the graph of Figure 120 shows otherwise, even if the difference between the electromagnetic coupling and the weak hypercharge coupling is taken into account. Indeed, the running of the coupling constants is not due to strand number only, but also to the explicit boson and fermion content of nature, as we just saw.

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The running of the coupling constants up to Planck energy

At energies near the Planck energy, quantum field theory is modified: effects due to the strand diameter start to play a role. Near Planck energy, tangles get tighter and tighter and fluctuations get weaker, because there is less room for them. In other words, near Planck energy tangles tend to approach the structure of horizons. Therefore, near the Planck energy, the strand model predicts deviations from the energy dependence of the coupling constants that is predicted by quantum field theory. So far, estimating such deviations has not been possible.

Another calculation might seem more promising: to calculate the coupling constants near Planck energy. It could be argued that the approach to calculate the low-energy coupling constants from Planck-energy values seems unsatisfactory, due to the approximations and extrapolations involved. But it is possible if we are convinced that quantum field theory is correct up to Planck energy. And this is just what the strand model predicts. Such a Planck-scale calculation might then allow us to estimate the low-energy coupling constants from their Planck energy values. However, so far, also this approach has not led to success, despite a number of attempts. The challenge seems to be to understand core deformation for case of tight tangle cores. We keep this option in mind.

Limits for the fine structure constant do not provide explanations

When searching for ways to determine the fine structure constant, we need to be careful. Here is an example that explains why.

Numerous observations of nature imply a limit on the fine structure constant. A pretty
Ref. 258 one appeared in a post on the internet in 2017. The *electrostatic repulsion* between two electrons at a given distance must be larger than the *radiation force* between to small neutral black holes at that same distance. In other words,

$$\frac{e^2}{4\pi\varepsilon_0}\,\frac{1}{r^2} > \frac{L_{\rm bh}}{c}\,\frac{\pi R_{\rm bh}^2}{4\pi r^2}\ . \tag{209}$$

Here it is assumed that thermal radiation from one black hole acts on the cross section of the other black hole by pushing it away. Multiplying both sides by $r^2/\hbar c$ and inserting
Vol. V, page 151 the expressions for the black hole luminosity $L_{\rm bh}$ and the black hole radius $R_{\rm bh}$ gives

$$\alpha > \frac{1}{15\,320\,\pi}\ . \tag{210}$$

The bound is not tight, but is obviously correct.

Various researchers are looking for observations that give the best possible bound for α. Such a search can indeed yield much better bounds. However, such a search cannot *explain* the value of α. We can indeed use thermodynamics, gravity or other observed
Ref. 259 properties to deduce *observational limits* on α. Many formulae of physics contain α in a more or less obvious way. Maybe, one day, known physics will be able to yield very tight upper and lower bounds for α. Still, *the explanation of the value of α would still lack.*

To explain the fine structure constant α, we need an approach based on the complete theory, not one based on known, millennium physics, such as expression (209). Millennium physics can *measure* α, but *cannot explain* it. To explain the fine structure constant, a unified theory is needed. In our case, we need to check whether we can calculate α with strands. Therefore, we now explore tangle topology, tangle shapes and tangle motion with this aim in mind.

Charge quantization and topological writhe

In the strand model, electric charge is related to the *chirality* of a tangle. Only chiral tangles are electrically charged. The strand model thus implies that a topological quantity for tangles – defined for each tangle in the tangle family corresponding to a specific elementary particle – must represent electric charge. Which quantity could this be?

The first candidate for charge in the strand model is provided by knot theory:

▷ The usual topological quantity to determine chirality of knots and tangles is the *topological writhe*.

To determine its value, we draw a *minimal projection*, i.e., a two-dimensional knot or tangle diagram with the *smallest* number of crossings possible. We then count the right-handed crossings and subtract the number of left-handed crossings. This difference, an integer, is the topological writhe. Topological writhe is thus a two-dimensional concept and does not depend on the shape of a knot or tangle. We note:

— The topological writhe of the W boson tangles is +3 or −3, depending on which mirror image we look at; the topological writhe of the Z boson and Higgs boson tangles vanishes. The topological writhe of any unknotted strand also vanishes. In this way, if we define the electric charge quantum number as *one third* of the topological writhe, we recover the correct electric charge quantum number of the weak and all other gauge bosons. We note that the Higgs boson does not change this result, so that all family members of a particle share the same topological writhe.

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— The tangles of the quarks show that if we define the electric charge quantum number as *one third* of the topological writhe, we recover the correct electric charge quantum number of all quarks. Adding Higgs bosons has no effect on this definition.

Page 326
— The tangles of the leptons show that if we define the electric charge quantum number as the topological writhe of the *centre region* only, we recover the correct electric charge quantum number of all leptons. Again, adding Higgs bosons does not change this result.

In other terms, the electric charge quantum number can be reproduced with the help of topological writhe. And indeed, the electric charge of massless bosons, i.e., photons and gravitons, vanishes.

Let us sum up. In nature, electric charge is quantized. The strand model describes charged particles with the help of fluctuating alternating tangles, and charge quantization is a topological effect that results because all particles are made of strands. In particular,

▷ The *electric charge quantum number* behaves similarly to topological writhe (times one third or times one): it is quantized, has two possible signs, vanishes for achiral tangles, is a topological invariant – and thus is conserved.

In short, a topological quantity, namely topological writhe, reproduces the electric charge quantum number in the strand model. Three issues remain. First, given that every particle is described by a tangle family with an infinite number of members, how is the electric charge, i.e., the topological writhe of the other tangle family members ac-

Challenge 217 e

counted for? It is not hard to see that family members do not change topological
Challenge 218 ny writhe. The second issue is more thorny: why is the charge definition different for leptons? We skip this problem for the time being. The third issue is the central one: What is the origin of the peculiar value of the charge unit, whose square has the value
Ref. 5 $\alpha = 1/137.035\,999\,139(31)$ at low energy?

Charge quantization and linking number

An alternative conjecture for charge quantization is the following:

> ▷ Electric charge, i.e., twist emission probability, might be proportional to the *linking number* of ribbons formed by strand pairs.

The following arguments speak in favour of this conjecture.

— In knot theory, a ribbon is the strip associated to and limited by two strands.
— The *linking number* of a ribbon is the number of times that the two edges of a ribbon wind around each other. The linking number is a topological invariant and an integer.
— In particle tangles, only wound up, i.e., linked ribbons should lead to (net) boson emission. For tangles made of three strands, we define a total linking number as the sum of all three possible linking numbers.
— The linking number of the Higgs boson strand pairs is zero; that of the Z boson strand pairs is the sum of 1, 0 and -1, thus also zero. The linking number for the W boson is 3 or -3, that of the quarks is 1, -1, 2 or -2. We thus conjecture that the charge quantum number is one third of the total linking number.
— Massless bosons, i.e., photons, gluons and gravitons, have no electric charge.

In short, linking number, an integer, might be a better topological quantity to explain electric charge quantization than topological writhe. On the other hand, it might well be that linking number, being a quantity that depends on *two* strands, is related to the *weak charge* rather than to the electric charge.

If the conjectured relation between linking number and electric or weak charge is correct, it might lead to a calculation of the corresponding coupling constant, once the tangle shape or, better, once the tangle dynamics is included in the proper way. For example, the photon emission probability could depend on the *writhe* or on the *twist* of the (averaged) ribbons. Both these properties might lead to virtual photon emission. (The sum of writhe and twist of a ribbon is given by the linking number, as explained by Calugareanu's theorem.)

In this and any topological definition of electric charge, we face two slight hurdles: First, we have to watch out for the graviton: it is uncharged. Secondly, we have to explain why the strand model for the simplest family member of the d quark is not chiral. Both
Challenge 219 e hurdles can be overcome.

If the linking of *two* strands is connected to weak charge, it might well be that a similar quantity defined for *three* strands is related to colour charge. All these possibilities are topic of research.

How to calculate coupling constants

The strand model suggests that crossing number and linking number somehow define electric and weak charge. In simple words, the model suggests that quantization of all charge types is a topological effect; quantization is due to the multiple ways in which strands cross inside tangles.

Coupling constants describe the probability of interaction with gauge bosons. Experiments show that these quantities are slightly *scale*-dependent, since they run with energy. But in the strand model, coupling constants are not really *shape*-dependent: electrons, muons and antiprotons have the same electric charge and fine structure constant values despite being described by different tangles. Coupling constants do not depend on the *kind* of tangle. Experiments show that they just depend somewhat on its size. In short,

> ▷ We need a definition of each coupling constant that is *tangle-independent* and *shape-independent*, and only depends on a topological invariant of tangles.

In fact, this conclusion eliminates many speculations, including a number of calculation approaches that were included in this chapter in previous editions. We are left with just a few options. To explore them, we start with an overview.

Coupling constants in the strand model

In experiments, there are the following gauge interactions with their charges:
1. The electromagnetic interaction with electric charge and U(1) symmetry.
2. The weak interaction with weak isospin and SU(2) symmetry.
3. The strong interaction with colour and SU(3) symmetry.

In the strand model, the *gauge interactions* are modelled as transfers of Reidemeister moves:
1. The electromagnetic interaction is twist transfer and the electric charge is preferred twist transfer to or from a massive particle. Twists can be added abd form a circle: they form a U(1) Lie group. They change the tangle phase by exchanging one observable crossing.
2. The weak interaction is poke transfer and the weak isospin is preferred poke transfer to or from a massive particle. Pokes exist in three linearly independent directions and their generators behave like the belt trick: they generate an SU(2) Lie group. They change the tangle phase by exchanging two observable crossings.
3. The strong interaction is slide transfer and the colour charge is preferred slide transfer to or from a massive particle. Slides can be added, its generators have a Z_3 symmetry and they form an SU(3) Lie group. They change the tangle phase by exchanging two or three? crossings.

In the strand model, *neutral particles* are those that cannot receive Reidemeister moves or that receive them all in equal way:
1. Electromagnetism: Neutral 'tangles' are made of one strand (e.g., the photon) or are topologically achiral (the Z and the neutrinos).

2. Weak interaction: Neutral tangles are made of one strand (e.g., the photon) or of two straight or unpokeable strand pairs (the Z, the right-handed leptons and quarks).
3. Strong interaction: Neutral tangles are made of one strand or of three strands.

In the strand model, *charged particles* are specific tangles:
1. Electric charge is due to the observability of crossings during photon emission or absorption, i.e., when twists are applied. Particles with electric charge, i.e., with preferred twist transfer, have a global asymmetry, global twistedness, namely topological chirality. Locally, electrically charged particles have crossings; electric charge is positive or negative. Charge is 1/3 of the signed crossing number. Examples are the charged leptons, the quarks and the W boson.
2. Weak charge is thus due to the observability of crossings during W or Z emission or absorption, i.e., when pokes are applied. Particles with weak isospin, i.e., with preferred poke transfer, have a global asymmetry that prevents all pokes to act equally effectively: For fermions, such an asymmetry arises when tangle twistedness and the belt trick have the same sign; thus all left-handed fermions and right-handed antifermions have weak isospin. Locally, weakly charged fermions behave like a belt buckle that rotates in the appropriate direction. Due to their tangle topology, some fermions have positive, others negative weak isospin. For the W boson, the asymmetry is built into the tangle; due to the tangle structure, the W and its antiparticle have plus or minus twice the weak isospin of fermions.
3. Colour, strong charge, is due to the observability of crossings during gluon emission or absorption, i.e., when slides are applied. Particles with colour charge, i.e., with preferred slide transfer, have a global asymmetry that prevents all slides to act equally effectively: Coloured particles are made of exactly two strands with tails in tetrahedron skeleton directions. Only two-stranded tangles allow certain slides and prevent others. Therefore only quarks have colour charge. Locally, red, blue and green colours correspond to three directions in one plane that differ by an angle of $2\pi/3$.

Coupling strength is the ease of crossing rotation, of poke creation, and of slide induction. These connections allow calculating the coupling strength values.

Deducing α from precession

In nature, magnetic fields rotate charged particles. In the strand model, as shown in
Page 234 Figure 53, magnetic fields are made of moving twists. In fact, from the strand definition of the electromagnetic interaction and the electric charge and from the drawing in
Page 229 Figure 50, we deduce:

▷ Moving twists rotate crossings.

We note that this description differs slightly from a pure twist transfer. But this formulation is the key to calculating α.

We assume that the typical, average crossing is lying in the paper plane, as in the drawing of the fundamental principle. For an average crossing, the two strands lie along the x and y axes. When a photon, i.e., a twist, arrives along the diagonal in the first quadrant, it rotates the crossing completely, by one turn. If the twist arrives at a different angle, its effect is lowered. We approximate this angle effect with simple trigonometry: we assume that the angular projection describes the reduction of the effect with the incoming angle of the twist.

For the incident photon, we call γ the angle from the y-axis and β the angle out of the paper plane. The average rotation angle induced by an absorbed photon or twist on a charged particle with *three* crossings, corresponding to *one* elementary charge, can be calculated. We include $\sin\gamma$ for the volume element in spherical coordinates and average over the possible angle values δ between the strands at the crossing. Further terms arise from the trigonometric approximation. In particular, a second power arises from the two tails, and a further squaring is required to get probabilities. Nevertheless, the expression remains open to dispute:

$$\sqrt{\alpha_{\text{calc}}} = \frac{3}{2\pi^2} \int_{\delta=0}^{\pi} \int_{\gamma=-\delta/2}^{\delta/2} \int_{\beta=-\pi/2}^{\pi/2} \cos\beta\,(\sin\delta)^2\,(\cos(\gamma\pi/\delta)\ \cos\beta)^4\,\mathrm{d}\beta\,\mathrm{d}\gamma\,\mathrm{d}\delta = 0.15\ . \tag{211}$$

The resulting value of 0.15 is not an acceptable approximation to reality, in which $\sqrt{\alpha} = 0.08542454311(1)$ at low energy and $\sqrt{\alpha} = 0.10(1)$ at Planck energy. Neither is the value a good approximation to the hypercharge coupling, which changes from $\sqrt{\alpha_1} = 0.10(1)$ at 100 GeV to $\sqrt{\alpha_1} = 0.13(1)$ at Planck energy. We need a better approximation for the value of the electromagnetic coupling strength.

Deducing the weak coupling

Weak fields deform strand (crossing) pairs by adding or transferring generalized pokes. Weak fields are collections of pokes; pokes represent virtual weak bosons. The weak isospin, the weak charge, is related to the orientation of the strand pairs. The weak interaction occurs through an incoming poke that deforms a strand pair:

▷ A moving poke rotates a pair of strands.

This process is the key to calculating α_{w}. We note that there is a certain similarity to the setting used for calculating the electromagnetic coupling: in both cases, the incoming boson acts on a target consisting of two strands. This similarity is the reason for electroweak mixing.

We calculate the coupling constant for a single belt buckle, assuming parallel strands. The average rotation angle induced by one incoming weak (unbroken) boson (out of the three possible cases) is one full turn when the impact is perpendicular to the two strands and to the plane defined by them. For a general incidence angle the induced rotation angle is lower. We again use trigonometrical projection to approximate the induced crossing rotation angle in the general case, with the same issues as in the previous case. We call γ the angle from ideal incidence, and β the longitude. The average angle is then given by

$$\sqrt{\alpha_{\text{w calc}}} = \int_{\gamma=0}^{\pi/2} \int_{\beta=0}^{2\pi} \sin\gamma\,(\cos^2\gamma\ \cos^2\beta)^4\,\mathrm{d}\gamma\,\mathrm{d}\beta \approx 0.19\ . \tag{212}$$

If we need to average over the different angles between the strands that make up the pair experiencing the poke, we get a different value.

The calculated value of the weak coupling is not an acceptable approximation to reality, in which $\sqrt{\alpha_w} = 0.18$ at the (low) energy of 100 GeV and $\sqrt{\alpha_w} = 0.14$ at Planck energy. We need a better approximation.

Deducing the strong coupling

Strong fields deform specific three-strand configurations by adding generalized slides. The generalized slides are due to gluons. Strong colour is related to the order and orientation of the strands in these specific three-strand configurations. In short:

▷ Incoming, moving slides deform three-strand configurations.

This is the key to calculating α_s.

We assume that one of eight possible gluons is incident. In an average triple strand configuration, the three strands are oriented in a way that in the paper plane they look like three symmetrically arranged rays. One ray lies along the y axis. When a gluon arrives, it performs a slide. For an incident gluon, we call γ the angle from the y-axis to the next strand and β the angle out of the paper. In the trigonometric approximation, the average slide angle induced on a coloured particle is given by

$$\sqrt{\alpha_{\text{s calc}}} = \int_{\gamma=0}^{\pi/2} \int_{\beta=0}^{2\pi} \sin\gamma\, (\cos^3\gamma\ \cos^3\beta)^4\, d\beta\, d\gamma \approx 0.11\ . \qquad (213)$$

This is not an acceptable approximation to reality, in which $\sqrt{\alpha_s} = 0.7(1)$ at the (low) energy of 1 GeV and $\sqrt{\alpha_s} = 0.13(1)$ at Planck energy. We need a better approximation for the strong coupling.

Open challenge: calculate coupling constants with precision

The approximations used above for estimating the coupling constants can be dismissed as mere educated guesses. Despite this objection, these guesses show that a determination of the coupling constants from the strand model is within reach, and that it can be realized with limited effort. It is sufficient to improve the three approximations; this is can be realized by using computer simulations for the transfer of Reidemeister moves or by finding an improved analytical model.

Calculating all three coupling constants ab initio with high precision will allow checking the statements of this section in an independent manner and, above all, will allow
Challenge 220 r
testing the strand model. The calculations should be performed at different energies, to confirm the energy dependence of the couplings. Also the influence of the effective strand diameter on the fine structure constant should be explored.

In order to reach highest precision, the effects of the various tangle family members might have to be taken into account, because in the strand model, each particle is de-
Page 352
scribed by a family of tangles. On the other hand, the strand model predicts that family members have a small effect on the coupling constant, so that the family issue can be

neglected in the beginning.

In the case of the nuclear coupling constants, Arnold's results on plane curves may help in the estimations and calculations. Ref. 260

Electric dipole moments

Experimental physicists are searching for electric dipole moments of elementary particles. No non-zero value has been detected yet. The idea of electric dipole moment is based on a non-spherical distribution of electric charge in space.

In the strand model, particles are tangles. As a consequence, the electric charge distribution – the distribution of the crossings in a tangle – is intrinsically a slightly non-spherical quantity, thus a quantity unequally distributed in space. However, it is only non-local on a scale of the order of a Planck length. In other terms, the electric dipole moment d of elementary particles is predicted to be

$$\triangleright \quad d = f\, e\, l_{\mathrm{Pl}} \,,$$

where the factor f arises from averaging the tangle and is of order one. Similar values are predicted by the standard model in the absence of supersymmetry and grand unification. Ref. 261 However, the sensitivity of measurements has not reached these values yet, by several orders of magnitude.

We note that the strand model predicts that the dipole moment changes, or 'runs', with energy. This follows from the shape-dependence of the dipole moment. Such a dependence is also predicted by quantum field theory.

In summary, we expect that up to a region close to a Planck length, the *strand* model should not yield dipole moments that differ in order of magnitude from those predicted by the *standard* model of particle physics. In the future, more precise calculations and measurements could allow testing the strand model using dipole moments.*

Five key challenges about coupling strengths

There are many ways to evaluate candidates for unified models. A concrete evaluation focuses on four key challenges about coupling constants. These challenges must be resolved by any candidate model in order to be of interest.

1. So far, we explained particle charges with topological properties of the tangle models of the particles, and we explained coupling strengths with the transfer of crossings, pokes and slides. This allowed deducing a rough approximation of coupling constants. By doing so, we have settled a first key challenge:

▷ The strand model explains why the fine structure constant, or equivalently, the electric charge, is the *same* for electrons and protons.

Deducing this equality is a key challenge for any unified model. In fact, all coupling constants must be independent of particle type. This is indeed the case in the strand

* It might even be that software packages allowing to do this already exist. The package of Christian Beck that he used in arxiv.org/abs/hep-th/0702082 might be an example. Others might also exist.

model.

2. The second key challenge was the energy-dependence of the coupling constants. The strand model predicts that coupling constants run with energy in exactly the way that is predicted by QED, QCD and electroweak theory. We could also argue that this is not a real challenge for any unified model that reproduces these theories. In the strand model, the running of the electromagnetic coupling constant can be seen as a consequence of the gradual tightening of tangles with energy. For a typical electrically charged particle at low energy, the tangle is very loose; therefore:

▷ The Planck scale number of crossings is shielded by an additional cloud of crossings created by the loose strands of the tangle.

In this way, the strand model explains the running of the fine structure constant in exactly the same way as QED.

3. The third key challenge has only been touched upon very briefly:

▷ Any unified model needs to clarify the relation between the hypercharge, the electric charge and the weak isospin (the 'weak charge').

The strand model explains electromagnetism as acting on crossings and the weak interaction as acting on parallel strands. This general statement contains the required explanation; but the details still need to be worked out. It is expected that in electromagnetism, a *single* crossing is rotated, mainly by rotating *one* strand around the other. In contrast, in the weak interaction, *two* strands are rotated together, producing ar switching *two* crossings. The number of crossings differs between electromagnetism and the weak interaction, but the total number of involved strands is tow in both cases. As a result of this similarity, the two interactions mix. The final explanation of electroweak mixing might even allow to deduce a intuitive geometric meaning of θ_{w} , the weak mixing angle or Weinberg angle.

4. The fourth key challenge, related to the previous one, still needs to be explored in more detail:

▷ Any unified model must explain why the mass ratio of the intermediate weak vector bosons is related to the coupling ratio of the weak and the electromagnetic interaction as

$$\left(\frac{m_W}{m_Z}\right)^2 + \frac{\alpha}{\alpha_{\mathrm{w}}} = 1 \ . \qquad (214)$$

The strand model strongly suggests that it can explain the relation, but the detailed argument must yet be provided. Using more drastic language, we can repeat what many have said already in the past: explaining the electroweak mixing expression (214) is the key challenge for any unified model.

In the strand model, the two electroweak coupling constants are measures for interaction probabilities of crossings with twists and with pokes. In contrast, masses are interaction probabilities of crossings with spatial curvature. Why are they related by expression (214)? Here is a short *brainstorm* on the issue.

In the strand model, mass appears by tail braiding. Tail braiding adds crossings, and in this way adds mass. Added crossings also imply added weak and sometimes electric charges. The Z boson arises from vacuum by different tail braidings than the W. The W arises by the braiding of two tail pairs at 90 degrees; the Z arises by braiding one tail pair at 90 degrees.

In case of the W and the Z bosons, the Z tangle produces a larger disturbance of the vacuum than the W; therefore it is more massive than the W.

At which angle does a clasp start to form a ‘‘enclosed space in between’’? Surprisingly, this happens for any angled strand pair, even if the angle is near π. How does this enclosed space change with scale, given that scale might change the clasp angle? This question might be related to the running of masses, mixing angles or coupling constants. In particular, we should answer the following question: Which physical observable does this enclosed space influence? Mass, couplings, or mixings? Is mass more related to ropelength or more related to the enclosed space?

5. The fifth key challenge is, of course, the precise calculation of the coupling constants.

Summary on coupling constants

The strand model implies that coupling constants are geometric properties of tangle families that correspond to charged particles. As a consequence, strands explain why the coupling constants are not free parameters in nature, but fixed constants. Strands predict that coupling constants are the same for particles with the same charge, and that coupling constants are constant during the macroscopic evolution of the universe. Strands predict small electric dipole moments for elementary particles, compatible with and lower than present measurement limits. Strands also predict the correct sequence of the coupling constants at low energy and the correct sign of their running with energy. Strands thus reproduce all observed qualitative properties of coupling constants. No other unified model achieves this yet.

Using tangle shapes, the strand model proposes several ways to calculate coupling constants ab initio. First estimates based on the new *tangled* particle models yield results of the order of the observed ones; nevertheless, the errors due to the approximations are still larger than the measurement errors. In 2019, a first publication has appeared . Improved calculations are ongoing and will allow either to confirm or to refute the strand model.

Ref. 262

Chapter 13

A PICTORIAL SUMMARY OF THE STRAND MODEL

> La forma universal di questo nodo
> credo ch'i' vidi, ... **
>
> Dante, *La Divina Commedia*, Paradiso, XXXIII, 91-92.

Deducing all of modern physics from simple pictures is possible! Indeed, describing observations with the help of fluctuating strands agrees with all data and all experiments – without exception. More precisely, in the *strand model* or *strand conjecture*, all properties of nature follow by assigning the Planck units to fundamental events. This fundamental principle is illustrated in Figure 121: every fundamental event is a strand crossing switch occurring at the Planck scale. With this description, the strand model visualizes and explains quantum theory, the standard model of particle physics, general relativity and cosmology.

Quantum theory, inclusive wave functions and spin 1/2, is illustrated in Figure 122, Figure 123 and Figure 124. Gauge interactions and the origin of gauge couplings are illustrated in Figure 125 and Figure 126. The strand model proposes a tangle structure for each elementary particle; they are given in Figure 127 and Figure 128. The tangles determine masses and mixing angles. No modification or extension of the particle spectrum arises. With all the assigned particle tangles, Figure 129 and Figure 130 show that the strand model reproduces all Feynman diagrams of the standard model. No additional Feynman diagram is possible. In short: *No modification or extension of the standard model arises.*

The strand model also yields a tangle structure for flat vacuum, curved vacuum, gravitons, everyday gravity, and black hole horizons. They are shown in Figure 131, Figure 132, Figure 133, Figure 134 and Figure 135. These figures imply that *strands reproduce gravity and general relativity, without extension or modification at sub-galactic scales.*

Finally, the strand model proposes a tangle structure for the universe, from the origin to the present, as illustrated in Figure 136 and Figure 137. Strands predict a single universe with an expanding horizon, trivial topology and a matter density close to the critical value. The cosmological implications are still subject of research.

Strands fluctuating at the Planck scale thus explain the colours of flowers, the motion of butterflies, and the sight of the sky at night.

** 'The universal form of that knot, I think I saw, ...'

The **origin** of the strand conjecture

1. Planck units are limit values
$W \geqslant \hbar$
$F \leqslant c^4/4G$
$v \leqslant c$
$S \geqslant k$

2. Tethers explain spin 1/2

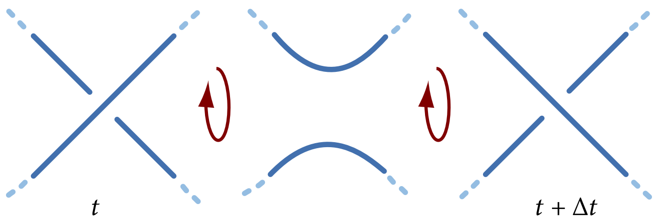
3. Extension explains black hole entropy and space microstructure

4. *Ur* or *qubit* as fundamental constituent

The **fundamental Planck-scale principle** of the strand conjecture

Strand conjecture

Observation



$W = \hbar/2$
$\Delta l \geqslant 2\sqrt{\hbar G/c^3}$
$\Delta t \geqslant 2\sqrt{\hbar G/c^5}$
$S = k/2$

A fundamental event localized in space

Predictions:

Events, observations and measurements are due to crossing switches.

All measurements are electromagnetic.

FIGURE 121 Many arguments lead to a description of nature with extended constituents at the Planck scale. The fundamental principle of the strand conjecture defines the simplest observation in nature, the almost point-like fundamental event. Every event results from a *skew strand crossing switch*, at a given position in three-dimensional space. The strands themselves are not observable; they are impenetrable and best imagined with Planck radius. The crossing switch all fundamental constants. The double Planck length limit and the double Planck time limit arise, respectively, from the smallest and from the fastest crossing switch possible.

Strand crossings have the same properties as **wave functions**

FIGURE 122 The geometry of a (skew) strand crossing suggests a relation to wave functions. In both cases, absolute phase around the orientation axis can be chosen freely. In contrast, phase differences due to rotations around that axis are always uniquely defined.

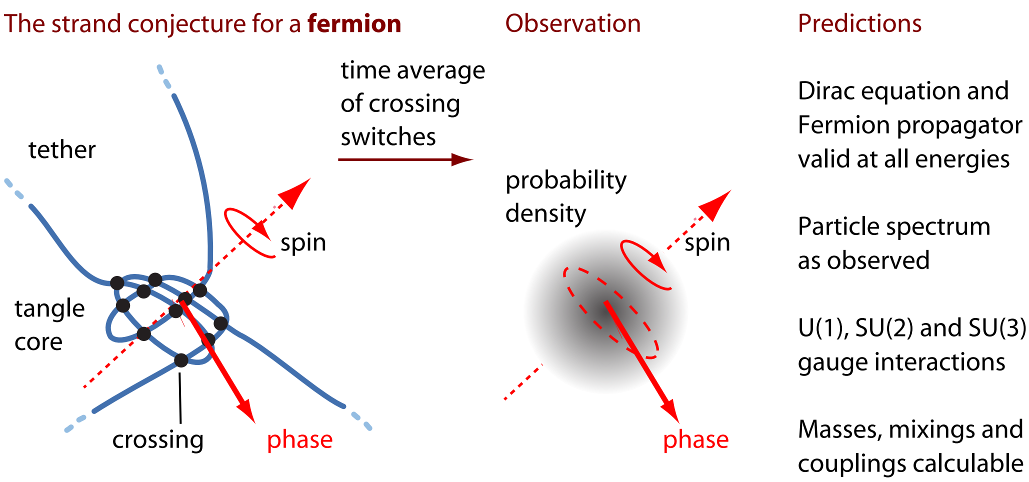


FIGURE 123 In the strand conjecture, the wave function and the probability density are due, respectively, to crossings and to crossing switches at the Planck scale. The wave function arises as time average of crossings in fluctuating tangled strands; a Hilbert space also arises. The probability density arises as time average of the crossing switches in a tangle. The tethers – connections that continue up to large spatial distances – generate spin 1/2 behaviour under rotations and fermion behaviour under particle exchange. The tangle model ensures that fermions are massive and move slower than light.

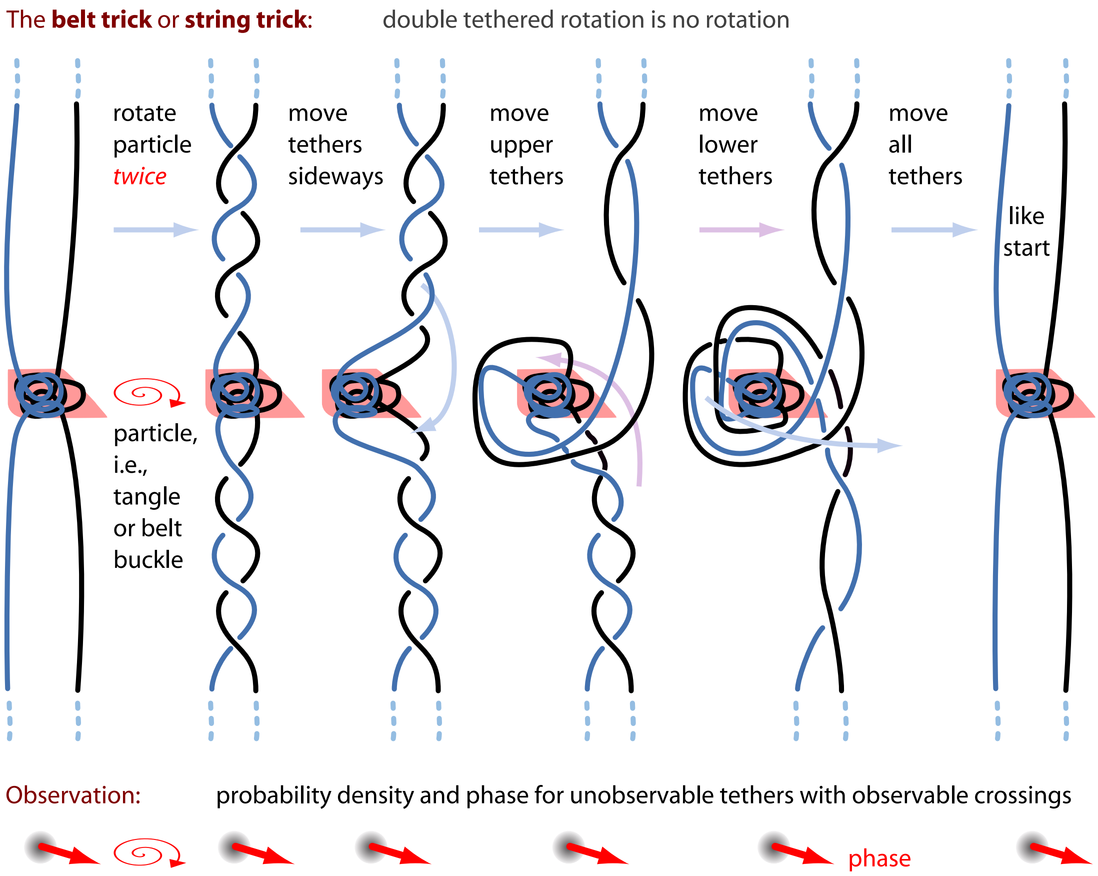


FIGURE 124 The *belt trick* or *string trick*: a rotation by 4π of a tethered particle, such as a belt buckle or a tangle, is equivalent to no rotation – when the tethers are allowed to fluctuate and untangle as shown. This equivalence allows the tethered particle to rotate forever. Untangling is impossible after a rotation by 2π only. This illustrates *spin 1/2*.
In addition (not shown), when two tethered particles are interchanged twice, all tethers can be untangled. Untangling is impossible after a single interchange only. This illustrates *fermion statistics*.
Both equivalences work for any number of tethers and assume that tethers are not observable, but crossing switches are.
The probability for the spontaneous occurrence of the belt trick depends on the complexity and details of the tangle core. This spontaneous process, together with Higgs emission and absorption, leads to the *mass* of elementary particles.

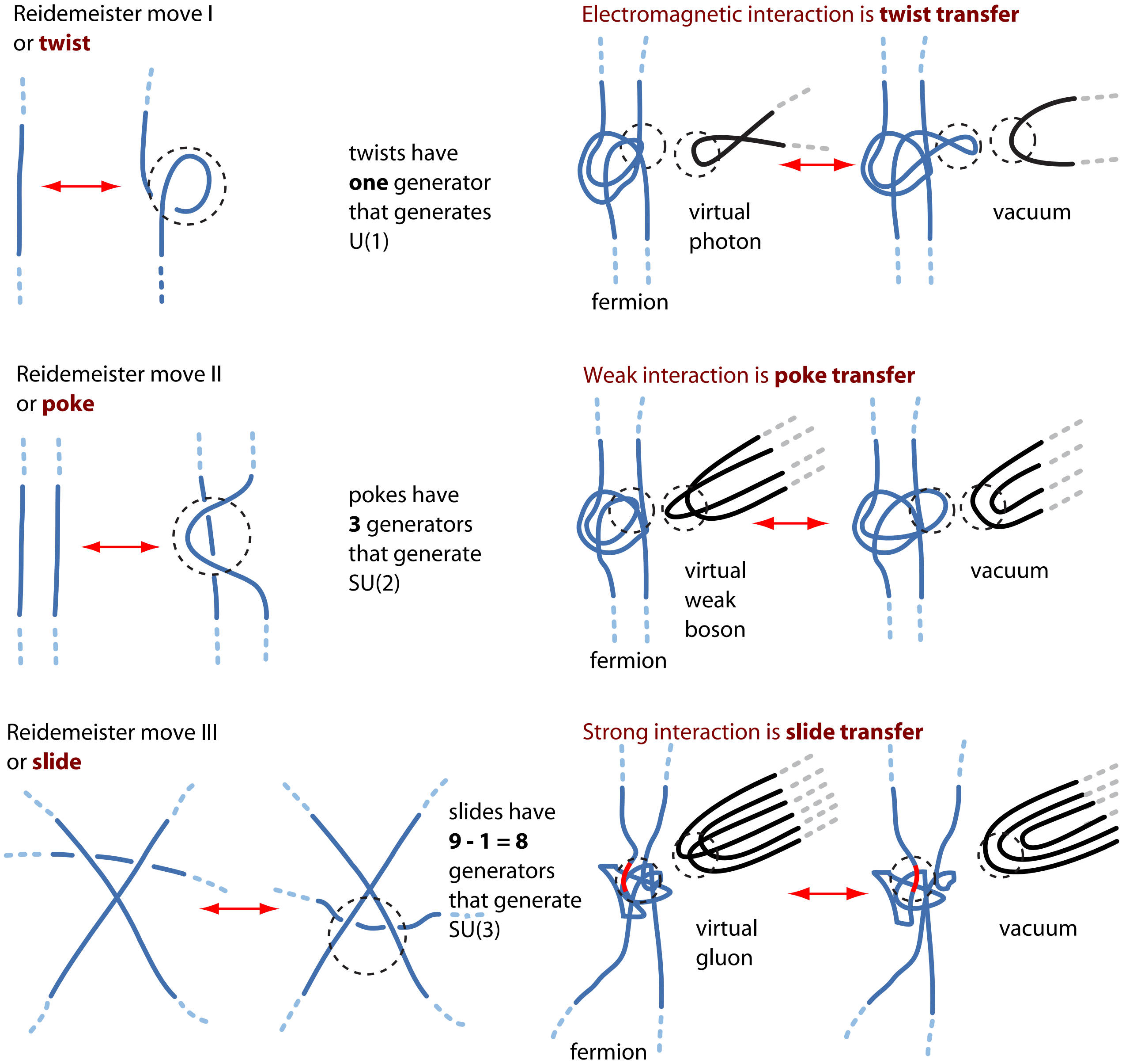


FIGURE 125 The three Reidemeister moves, i.e., the three possible deformations of tangle cores, determine the generators of the observed gauge interactions and thus determine their generator algebra. The generators rotate the regions enclosed by dotted circles by π. The full gauge groups arise by generalizing these rotations to continuous angles. Also the gauge coupling strengths arise.

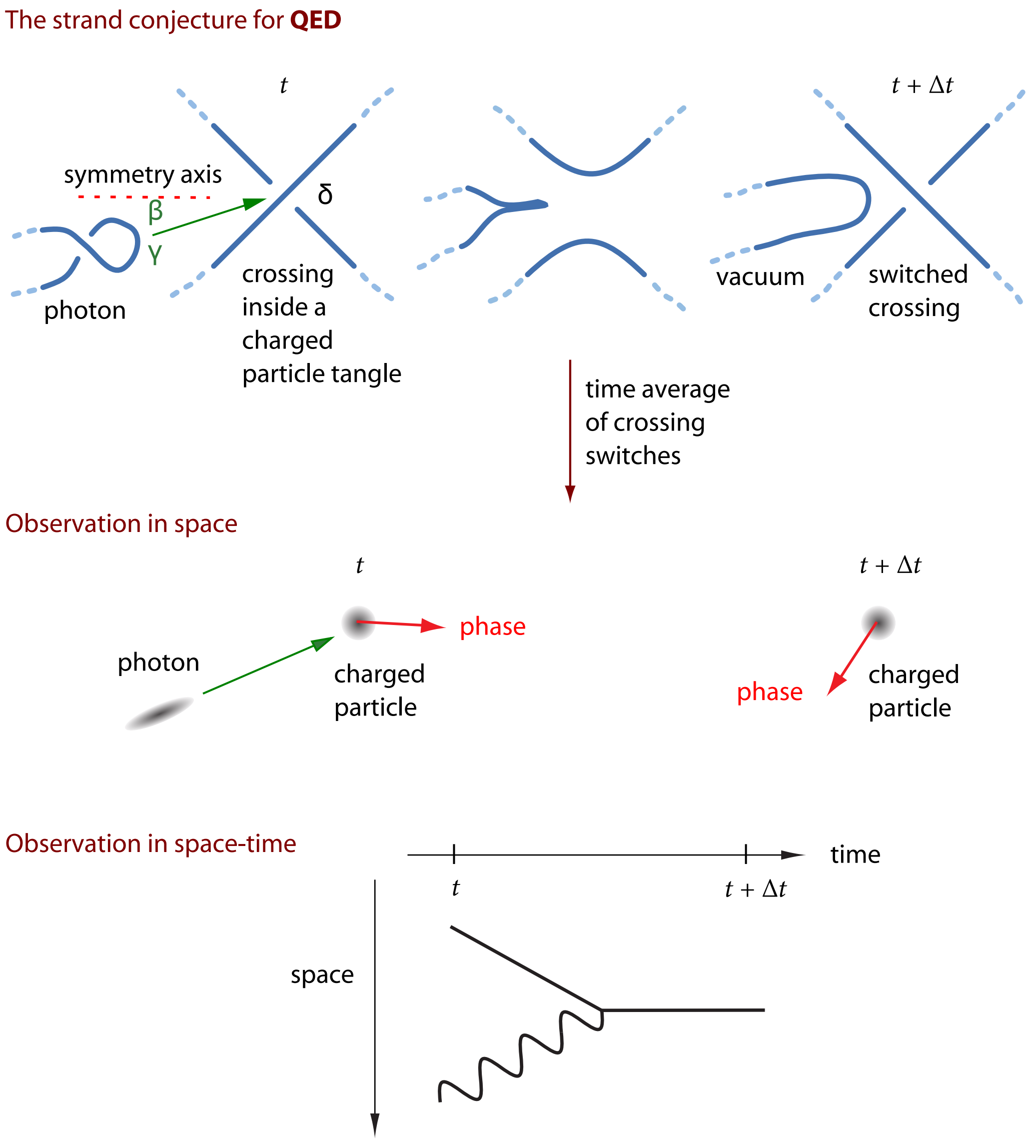


FIGURE 126 The geometric details of the basic process of quantum electrodynamics (QED). Top: the absorption of a photon by a tangle region carrying the charge $e/3$, viewed along the shortest distance of the crossing. Centre: the corresponding observation at usual experimental scales. Bottom: the basic Feynman diagram of QED. In short: twist transfer generates minimal coupling and determines the electromagnetic gauge coupling strength, thus the fine structure constant.

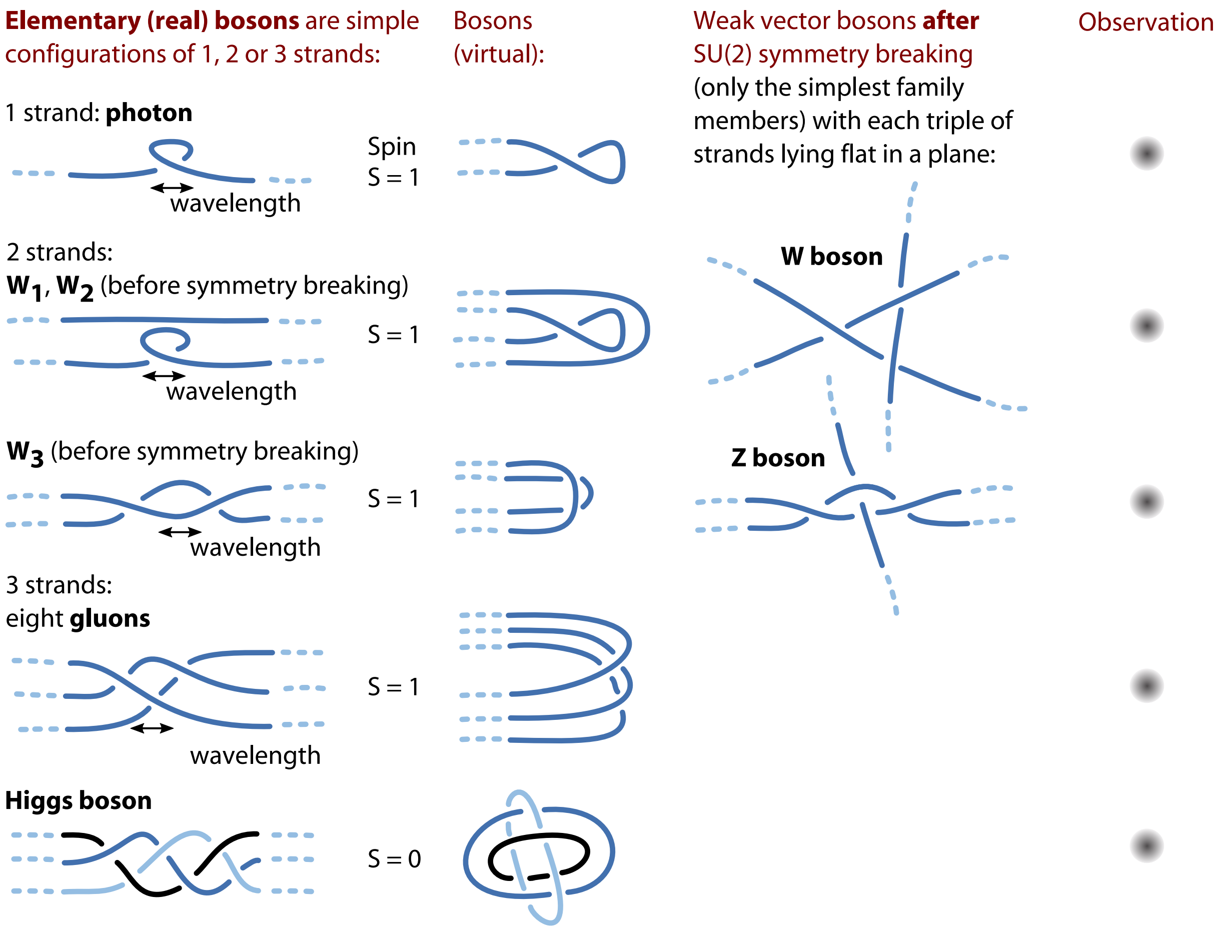


FIGURE 127 The tangle models for the elementary bosons. These tangles determine the spin values, the corresponding propagators, and ensure that the massless photons and gluons move with the speed of light. Apart from the graviton, no additional elementary boson appears to be possible. The tangle structure determines masses and the weak mixing angle.

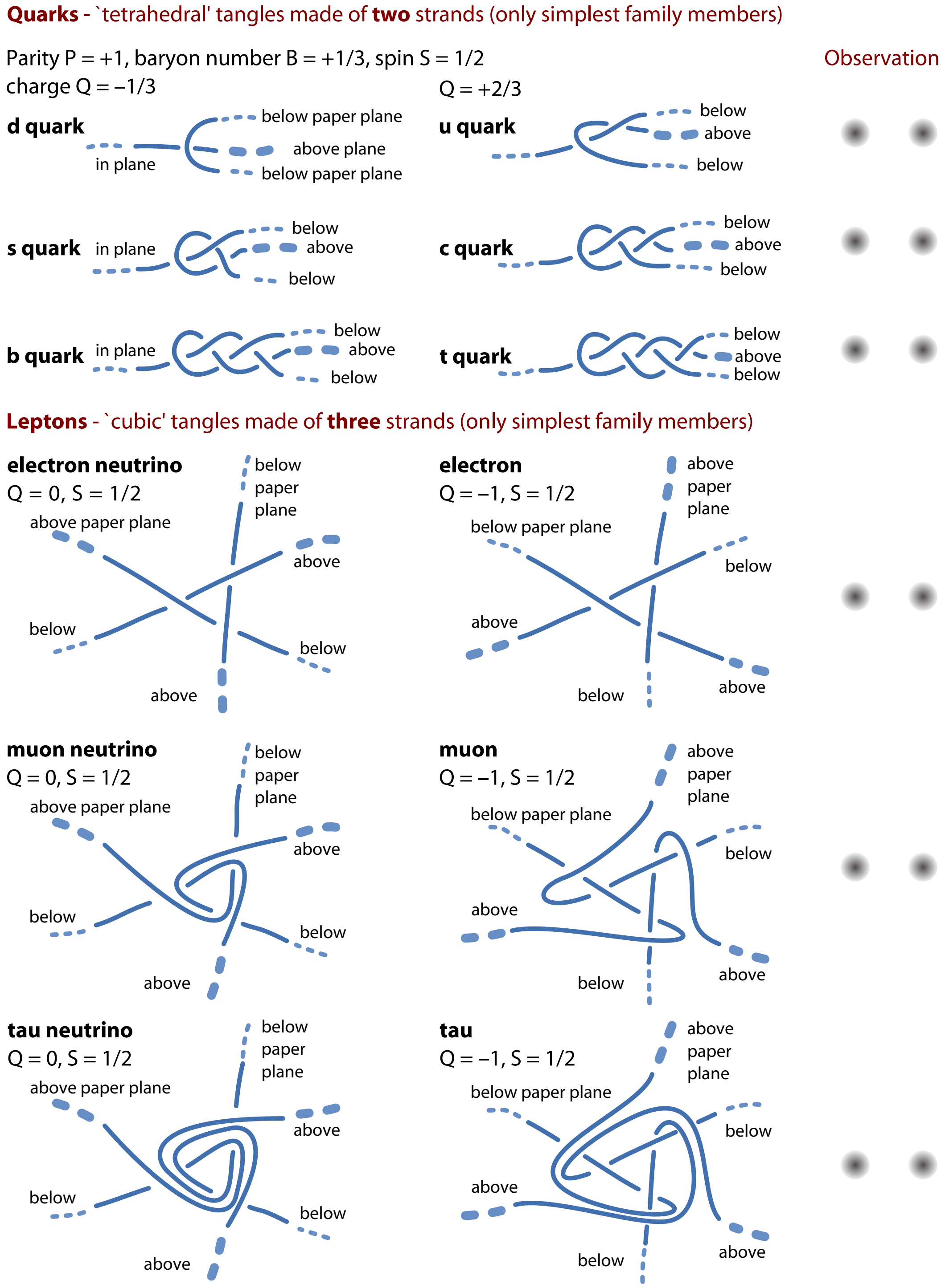


FIGURE 128 The *simplest* tangle models for the elementary fermions. Elementary fermions are described by rational, i.e., unknotted tangles. Their structures lead to coupling to the Higgs, as illustrated in Figure 130, produce positive mass values, and limit the number of generations to 3. Each fermion is represented by a tangle *family*, consisting of the simplest member and of tangles with added Higgs braids. The tangles determine the specific fermion propagators, including the values of masses and mixing angles. The tethers of the quark tangles follow the axes of a tetrahedron. The neutrino cores are simpler when seen in three dimensions: they are twisted triples of strands. The tethers of all lepton tangles approach the three coordinate axes at large distances from the core. No additional elementary fermions appear to be possible.

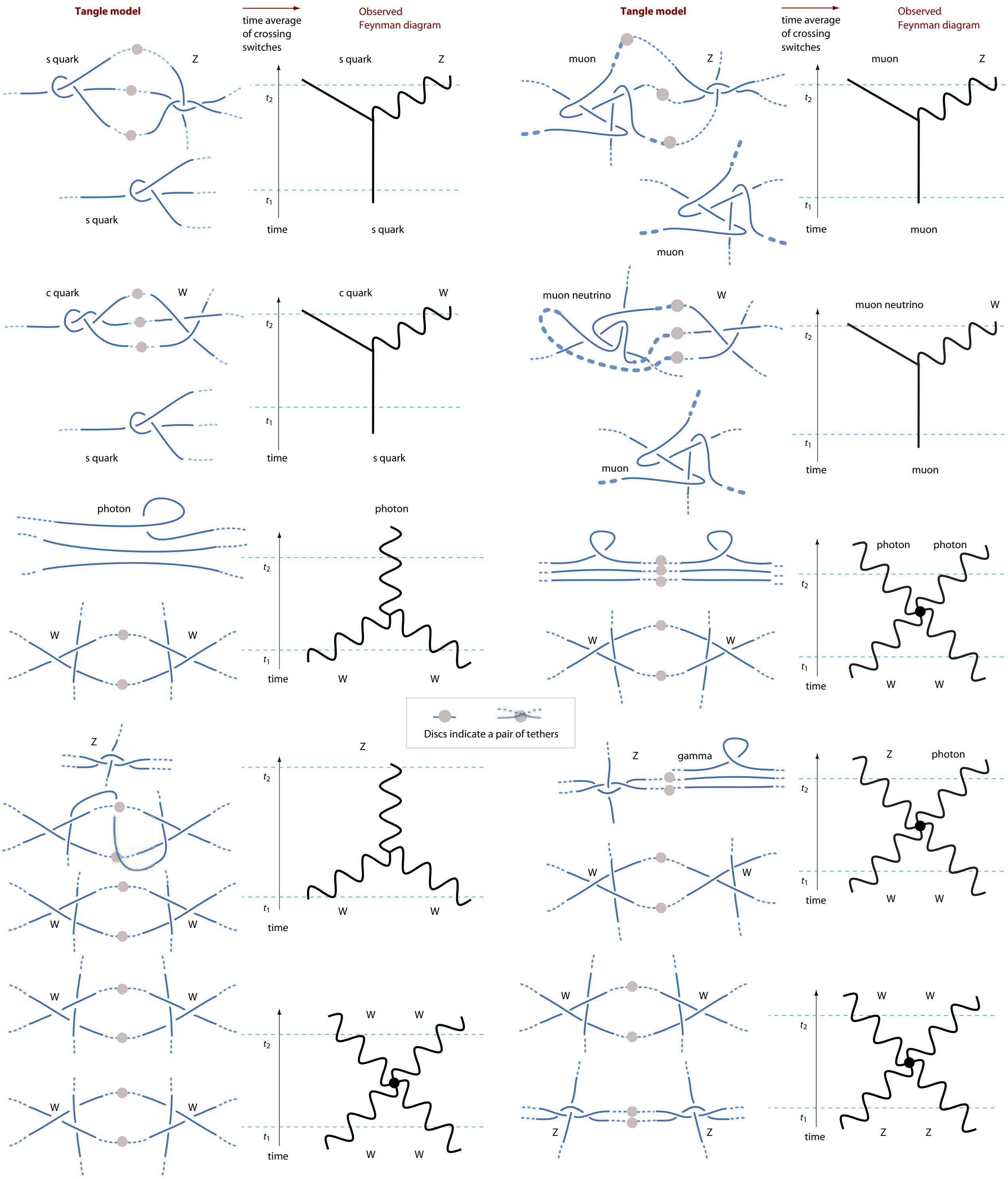


FIGURE 129 The vertices in Feynman diagrams allowed by the topology of the fermion and boson tangles (part one).

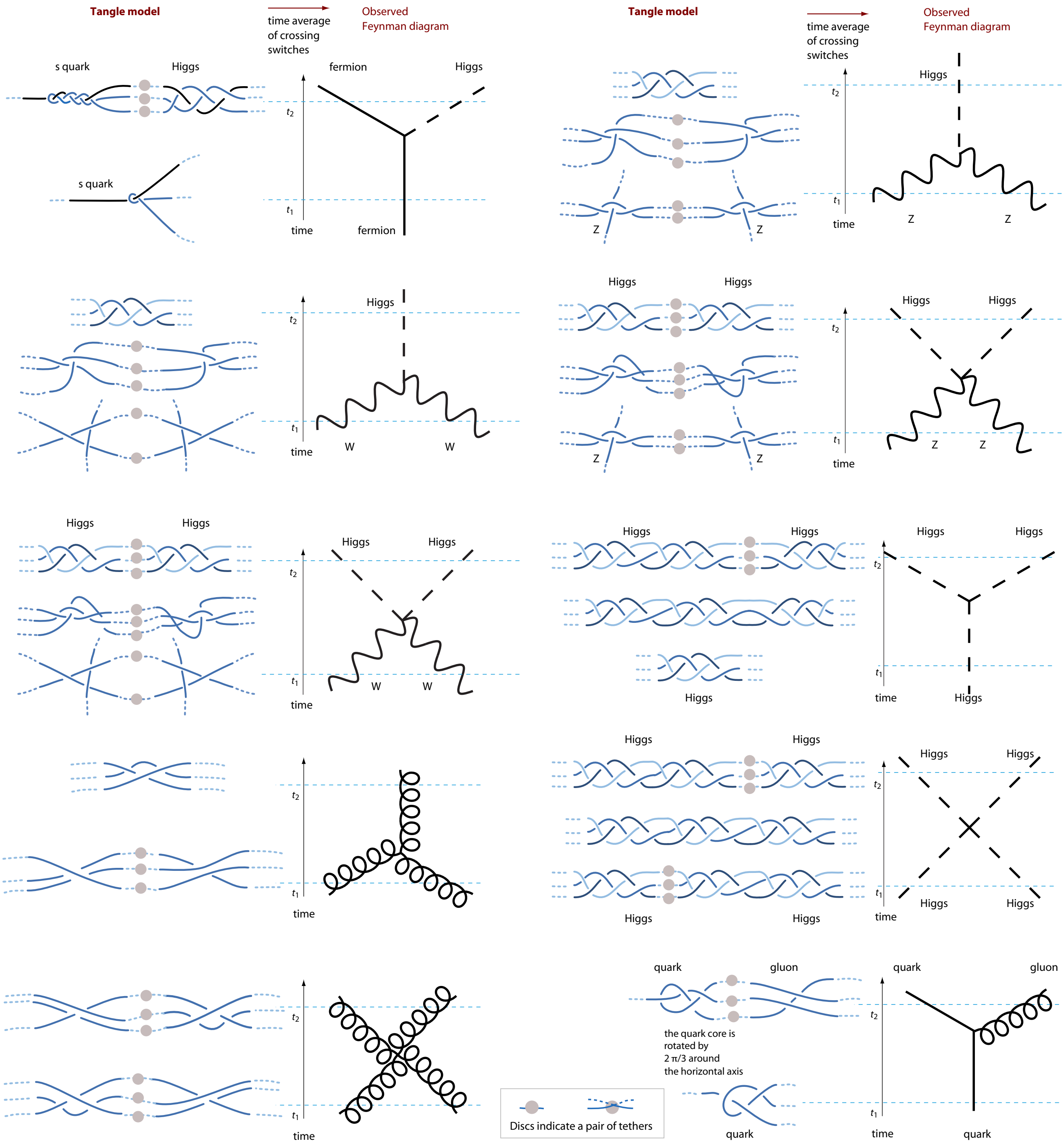
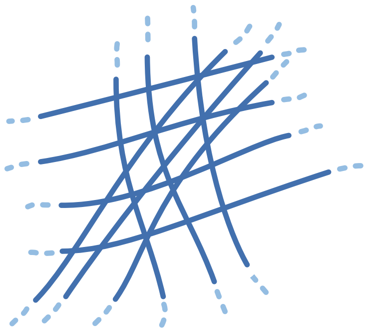


FIGURE 130 The vertices in Feynman diagrams allowed by the topology of the fermion and boson tangles (part two).

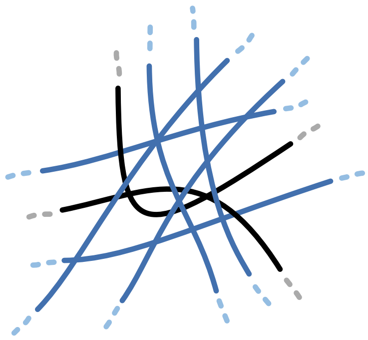
The strand conjecture for the **vacuum**



Observation	Predictions
Nothing (for long observation times)	Vanishing energy
Virtual pairs (for short observation times)	Emergent, Lorentz-invariant, and unique vacuum

FIGURE 131 An illustration of the strand conjecture for a flat vacuum: for sufficiently long time scales, the lack of crossing switches leads to a vanishing energy density; for short time scales, particle–antiparticle pairs, i.e., rational tangle–antitangle pairs, arise.

The strand conjecture for **curved space**



Observation	Predictions
Curved space	Black hole entropy
Non-trivial metric	Pure general relativity
Black holes	$P < c^5/4G$, $F < c^4/4G$
	Gravitons hard to detect

FIGURE 132 An illustration of the strand conjecture for a curved vacuum. The strand configuration is half way between that of a horizon and that of a flat vacuum. The black strands differ in their configuration from those in a flat vacuum.

The strand conjecture for the **graviton**

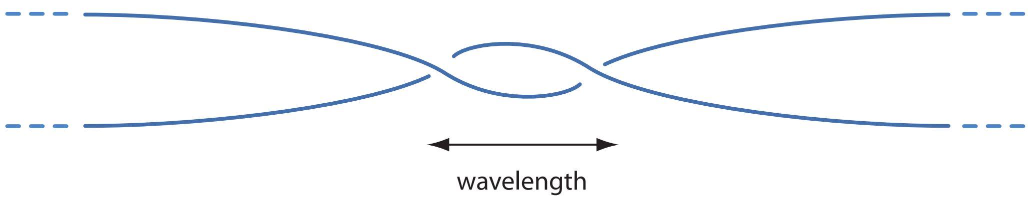


FIGURE 133 The strand conjecture for the graviton: a twisted pair of strands has spin 2, boson behaviour and vanishing mass.

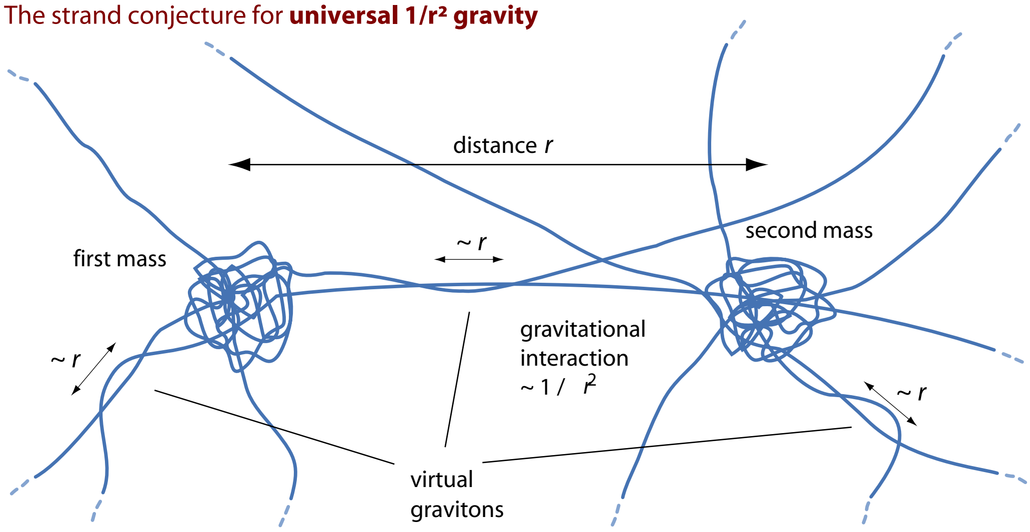


FIGURE 134 Gravitational attraction results from strands. For low speeds and negligible curvature, twisted tether pairs (virtual gravitons) from a mass lead to a $1/r^2$ attraction. The average length of twisted tether pairs scales with r and yields a $1/r^3$ decay of curvature. These conclusions are valid for infinite space without cosmological horizon.

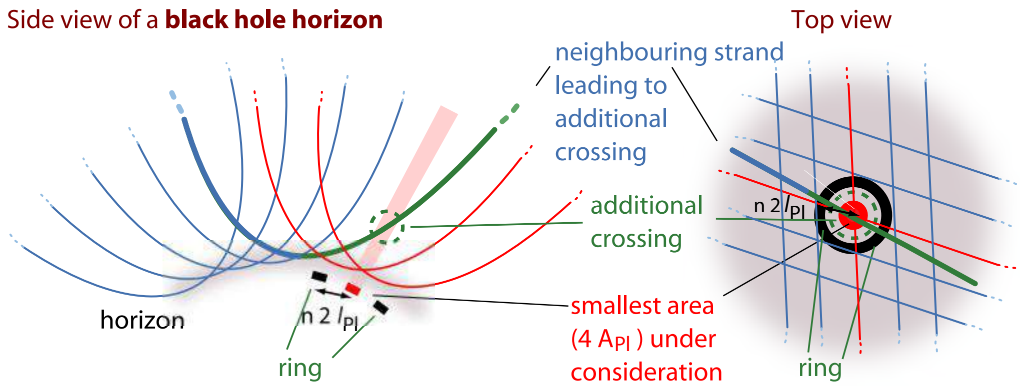


FIGURE 135 The strand conjecture for a Schwarzschild black hole: the black hole horizon is a cloudy or fuzzy surface produced by the crossing switches of the strands woven into it. Due to the additional crossings on the side of the observer, the number of micro-states per smallest area is larger than 2.

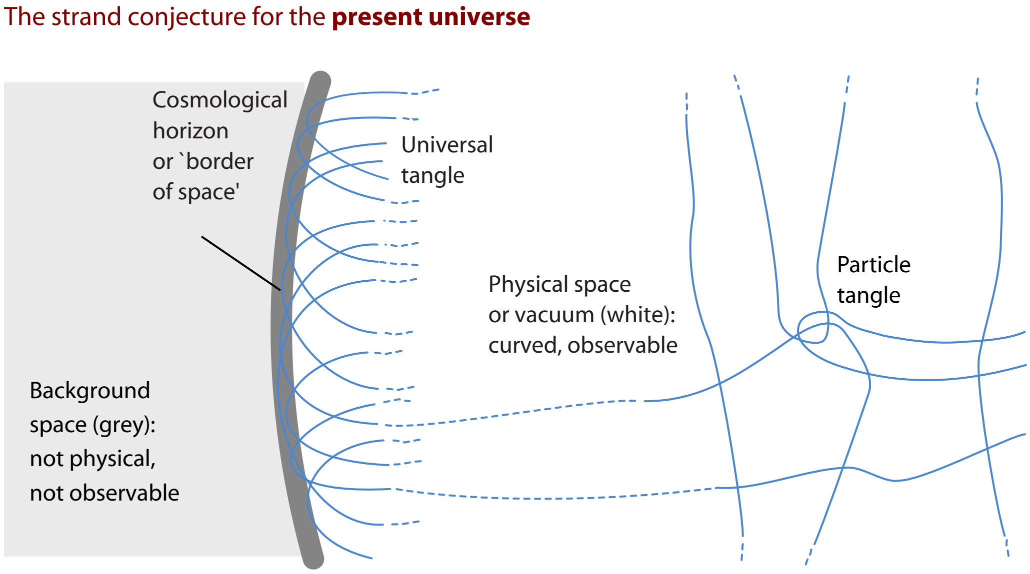


FIGURE 136 In the strand conjecture, the universe is limited by a cosmological (particle) horizon, as schematically illustrated here. Physical space (white) matches background space (grey) only inside the horizon. Physical space thus only exists *inside* the cosmic horizon.

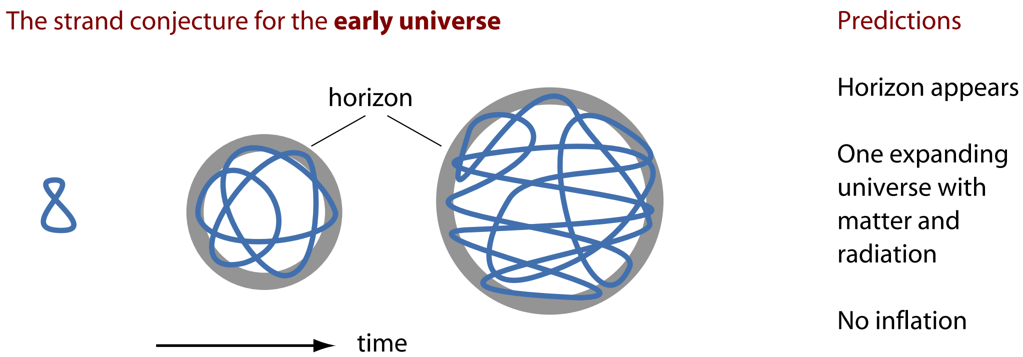
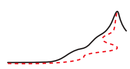


FIGURE 137 In the strand conjecture for the early universe, the universe increases in complexity over time and thereby forms a boundary: the cosmological horizon.



CHAPTER 14

EXPERIMENTAL PREDICTIONS OF THE STRAND MODEL

> Es gibt viele Theorien,
> die sich jedem Test entziehen.
> Diese aber kann man checken,
> elend wird sie dann verrecken.**
>
> Anonymous

Around the world, numerous researchers are involved in experiments that are searching for new effects. They are searching for new observations that are unexplained by the standard model of particle physics or by the conventional view of cosmology. At the same time, all these experiments are testing the strand model presented here. In fact, most people working on these experiments have not heard about the strand model, so that there is not even the danger of unconscious bias.

To simplify the check with experiments, the most important predictions of the strand model that we deduced in our adventure are listed in Table 17.

TABLE 17 The main predictions of the strand model that follow from the fundamental principle. The typeface distinguishes predictions that are unsurprising, that are *unconfirmed or unique* to the strand model, and those that are both **unconfirmed and unique**.

	Experiment	Prediction (most from 2008/2009)	Status (2019)
Page 37	Planck units (c, $\hbar$, k, $c^4/4G$)	are limit values.	None has been exceeded, but more checks are possible.
Page 330	Higgs boson	2012: does exist.	Verified.
Page 386	*Running of the coupling constants*	*2012: implies one Higgs.*	*Correct so far.*
Page 328 Page 330	*Longitudinal W and Z boson scattering*	*2012: show no non-local effects at the Large Hadron Collider.*	*None found yet.*
Page 328	*Longitudinal W and Z boson scattering*	*is unitary at the LHC.*	*Obvious.*
	W boson g-factor	is near to 2.	Is observed.
Page 352	*Unknown fermions (supersymmetric particles, magnetic monopoles, dyons, heavy neutrinos etc.)*	*do not exist.*	*None found yet.*

** No adequate translation is possible of this rhyme, inspired by Wilhelm Busch, claiming that any theory that can be tested is bound to die miserably.

TABLE 17 (Continued) The main predictions of the strand model that follow from the fundamental principle. The typeface distinguishes predictions that are unsurprising, that are *unconfirmed or unique* to the strand model, and those that are both **unconfirmed and unique**.

	Experiment	Prediction (most from 2008/2009)	Status (2019)
Page 314	*Unknown bosons (other gauge bosons, supersymmetric particles, axions etc.)*	*do not exist.*	*None found yet.*
Page 278, page 318	*Unknown interactions, energy scales and symmetries (grand unification, supersymmetry, quantum groups, technicolour etc.)*	*do not exist.*	*None found yet.*
Page 313	**Particle masses, mixing angles and coupling constants**	**are calculable by modifying existing software packages.**	**Most not yet calculated; approximations very encouraging.**
Page 313	Particle masses, mixing angles and coupling constants	are constant in time and space.	Is observed.
	Particle masses, mixing angles, coupling constants and g-factors	are identical for antimatter.	Is observed.
Page 376	Mixing matrix for quarks	is unitary.	Is observed.
Page 379	*Mixing matrix for neutrinos*	*is unitary.*	*No data yet.*
Page 339	*Neutrinos*	*are Dirac particles.*	*No data yet.*
Page 380	*Neutrinos*	*violate CP symmetry.*	*No data yet.*
Page 327	Neutrino-less double beta decay	does not exist.	Not yet found.
Page 278, page 327	**Electric dipole moments of elementary particles, magnetic dipole moment of neutrinos**	**have extremely small, calculable values.**	**No interesting data yet.**
Page 342	Tetraquarks	exist.	Likely.
Page 344, page 339	*Glueballs*	*probably do not exist; if they do, the spectrum can be compared to the strand model.*	*Not yet observed.*
Page 278, page 327	Proton decay and other rare decays, neutron-antineutron oscillations	occur at extremely small, standard model rates.	Not yet observed.
Page 340	Neutron decay	follows the standard model.	No deviations found.
Page 340	Neutron charge	vanishes.	None observed.
Page 336	*Hadron masses and form factors*	*can be calculated ab initio.*	*Not yet calculated; value sequences and signs correct.*
Page 355	Dark matter	is conventional matter plus black holes.	Partly confirmed by black hole mergers and lack of other results.
	Standard model of particle physics	*2012: is correct for all measurable energies.*	*All data agrees.*

TABLE 17 (Continued) The main predictions of the strand model that follow from the fundamental principle. The typeface distinguishes predictions that are unsurprising, that are *unconfirmed or unique* to the strand model, and those that are both **unconfirmed and unique**.

	Experiment	Prediction (most from 2008/2009)	Status (2019)
Page 148	Additional dimensions	do not exist.	Not observed.
Page 148	Non-commutative space-time	does not exist.	Not observed.
Page 295	General relativity	is correct at all accessible energies.	No deviation found.
Page 295	Short-distance deviations from universal gravitation and modified gravity	do not exist.	All data agrees.
Page 293	Space-time singularities, cosmic strings, wormholes, time-like loops, negative energy regions, domain walls	do not exist.	None observed.
Page 299, page 279	Quantum gravity effects	will not be found.	None observed yet.
Page 306	Behind a horizon	nothing exists.	Nothing observed.
Page 309	Cosmic inflation	did not occur.	Data not in contrast.
Page 380	Leptogenesis	did not occur.	Data are inconclusive.
Page 310	Cosmic topology	is trivial.	As observed.
Page 356	Vacuum	is stable and unique.	As observed.
	In summary: all motion	**results from strands.**	**Not yet falsified.**

In this list, the most interesting predictions of the strand model are the *numerical* predictions on the various mass ratios and mass sequences – including the Z/W and Higgs/W mass ratios – and the relative strength of the three gauge interactions. There is the clear option to calculate all fundamental constants in the foreseeable future.

In addition, the strand model reproduces the quark model, gauge theory, wave functions and general relativity; at the same time, the model predicts the lack of measurable deviations. The strand model solves conceptual problems such as the dark matter problem, inflation, confinement, the strong CP problem and the anomaly issue; by doing so, the strand model predicts the lack of unknown effects in these domains.

The strand model deduces all its experimental predictions from a single and simple fundamental principle: *events and Planck units are due to crossing switches of strands.* Provided there are no errors of reasoning, there is no way to change the predictions summarized here. The strand model is both simple and unmodifiable.

Naturally, errors of reasoning in the preceding chapters are well possible. A few have occurred in the past. The exploration was performed at high speed – possibly too high. If any experiment ever contradicts a prediction of the strand model, the model is doomed. When the above experimental predictions were first deduced in 2008 and 2009, they were quite unpopular. Practically all other attempts at unification predicted the existence of yet undiscovered particles and effects. However, so far, experiment does not confirm these other attempts; in fact, no prediction of the strand model has been falsified yet.

These predictions are not intellectual speculations. The author is prepared to take bets on each prediction of the above table – and has taken a few already.

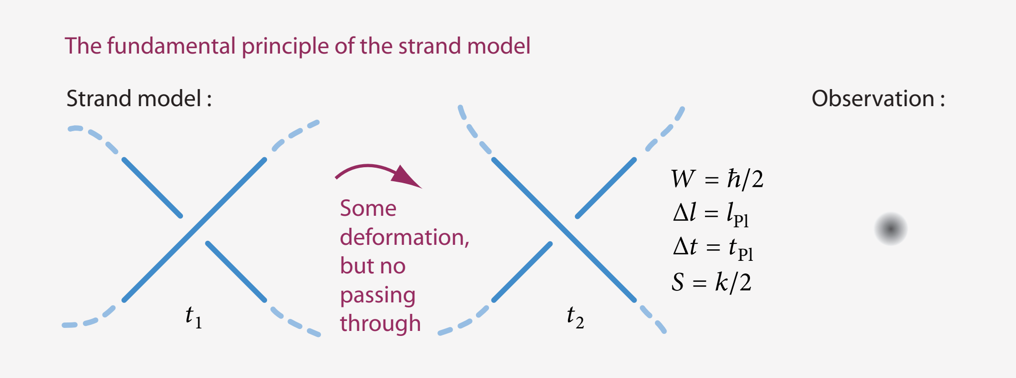


FIGURE 138 The fundamental principle of the strand model: *Planck units are defined by a crossing switch in three spatial dimensions.* With this principle, as shown in the previous chapters, the fundamental principle implies general relativity and the standard model of particle physics.

FINAL SUMMARY ABOUT THE MILLENNIUM ISSUES

In our adventure, we have argued that Planck's natural units should be modelled with the fundamental principle for strands, which is shown again in Figure 138. As we discovered, the fundamental principle explains the following measured properties of nature:

— Strands explain the principle of least action and the invariance of c, $\hbar$, G and k.
— Strands explain the three dimensions of space, the existence of gravitation, curvature and horizons, the equations of general relativity, the value of black hole entropy and the observations of modern cosmology.
— Strands explain all the concepts used in the Lagrangian of the standard model of particle physics, including wave functions, the Dirac equation and the finite, discrete and small mass of elementary particles.
— Strands explain the existence of electromagnetism and of the two nuclear interactions, with their gauge groups and all their other observed properties.
— Strands describe the observed gauge and Higgs bosons, their charges, their quantum numbers and their mass mass ranges.
— Strands explain the three generations of quarks and leptons, their charges and quantum numbers, their mixing, their mass sequences, as well as their confinement properties.
— Strands explain the quark model of hadrons, including CP violation, mass sequences, signs of quadrupole moments, the lack of unobserved hadrons, common Regge slopes and the existence of tetraquarks.
— Strands *do not allow* arbitrary values for masses, coupling constants, mixing angles and CP violating phases.
— Strands *enable* calculations of particle masses, their coupling constants, their mixing angles and the CP violating phases. First rough estimates of these values agree with the (much more precise) experimental data. Computer calculations will allow us to improve these checks in the near future.
— Strands *predict* the lack of unknown dark matter and of unknown inflation mechan-

isms.

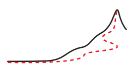
— Finally, strands predict that nature *does not hide* any unknown elementary particle, fundamental interaction, fundamental symmetry or additional dimension. In particular, strands predict that no additional mathematical or physical concepts are required for a unified theory.

All these results translate to specific statements on experimental observations. So far, there is no contradiction between the strand model and experiments. These results allow us to sum up our adventure in three statements:

1. *Strands solve all open issues.* With one simple fundamental principle, the strand model solves or at least proposes a way to solve *all* issues from the millennium list of open issues in fundamental physics. All fundamental constants can be calculated. Page 19
2. *Strands agree with all observations.* In particular, the strand model implies that general relativity, quantum theory and the standard model of elementary particles are a *precise* description of motion for all practical purposes.
3. *Nothing new will be discovered in fundamental physics.* Unexpectedly but convincingly, strands predict that general relativity, quantum theory and the standard model of elementary particles are a *complete* description of motion for all practical purposes.

We have not yet literally reached the top of Motion Mountain – because certain numerical predictions of the fundamental constants are not yet precise enough – but if no cloud has played a trick on us, we have seen the top from nearby. In particular, we finally know the origin of colours.

The last leg, the accurate calculation of the constants of the standard model of particle physics, is still under way. The drive for simplicity and the spirit of playfulness that we invoked at the start have been good guides. Page 22



Chapter 15

THE TOP OF MOTION MOUNTAIN

> All things are full of gods.
>
> Thales**

Who am I? Where do I come from? What shall I do next? Where does the world come from? Can the whole world really come to a sudden end? What will happen in the future? What is beauty? All these questions have a common aspect: they are questions about motion. But what is motion? Our search for an answer led us to study motion in all its details. In this quest, every increase in the precision of our description of motion was a step towards the peak of Motion Mountain. Now that we arrived there, we can savour what we have achieved and recall the emotions that we have experienced.

In our ascent, we have learned how we move, how we experience our environment, how we grow, what parts we are made of, and how our actions and our convictions about them can be understood. We have learned a lot about the history and a bit about the future of matter, of radiation and of space. We have experienced and understood the many ways in which beauty appears in nature: as colours, as shapes, as rhythms and most of all: as simplicity.

Savouring our achievement means that first of all, we now can look back to where we came from. Then we enjoy the view we are offered and look out for what we could not see before. After that, we search for what is still hidden from our sight. And finally, we take a different path back down to where we live.

OUR PATH TO THE TOP

> The labour we delight in physics pain.
>
> William Shakespeare, *Macbeth*.

Our walk had a simple aim: to talk accurately about all motion. This 2500 year old quest drove us to the top of this mountain. We can summarize our path in three legs: everyday life, general relativity plus quantum theory, and unification.

** Thales of Miletus (*c.* 624 – *c.* 546 BCE) was the first known philosopher, mathematician and scientist.

Everyday life: the rule of infinity

Galilean physics is the description of everyday life. We all learned Galilean physics before secondary school. Galilean physics is the exploration and description of the motion of stones, water, trees, heat, the weather, electricity and light. To achieve this description of our environment, our first and main act in life is to partition experience into experience*s*. In other words, our first intellectual act is the invention of *parts*; we invented the *plural*.

Ref. 1, Ref. 3

The act of partitioning allows us to define sequences among our experiences, and thus to define the concept of *time*. The concept of *space* arises similarly by our possibility to distinguish observations that occur at the same time. By comparing parts with other parts, we define *measurement*. Using all of this, we become able to define *velocity*, *mass* and *electric charge*, among others. These allow us to introduce *action*, the quantity that quantifies change.

For a simple description of observations, we assume that division is possible without end: thus we introduce the infinitely small. We also assume that widening our scope of observation is possible without end. Thus we introduce the infinitely large. Defining parts thus leads us to introduce infinity.

Using parts and, with them, the infinitely small and the infinitely large, we found, in volumes I and III, that everyday motion has six main properties: it is continuous, conserved, relative, reversible, mirror-invariant and lazy. Motion is lazy – or efficient – because it produces as little change as possible.

Nature minimizes change. This is Galilean physics, the description of everyday motion, in one statement. It allows us to describe all our everyday experiences with stones, fluids, stars, electric current, heat and light. The idea of change-minimizing motion is based on a concept of motion that is continuous and predictable, and a concept of nature that contains the infinitely small and the infinitely large in every observable.

Relativity and quantum theory: the absence of infinity

> “Vorhin haben wir gesehen, daß in der Wirklichkeit das Unendliche nirgends zu finden ist, was für Erfahrungen und Beobachtungen und welcherlei Wissenschaft wir auch heranziehen.*”
>
> David Hilbert

The idea that nature offers an infinite range of possibilities is often voiced with deep personal conviction. However, the results of relativity and quantum theory show the opposite. In nature, speeds, forces, sizes, ages and actions are limited. No quantity in nature is infinitely large or infinitely small. No quantity in nature is defined with infinite precision. There never are infinitely many examples of a situation; the number of possibilities is always finite. The world around us is not infinite; neither its size, nor its age, nor its content. *Nature is not infinite.* This is general relativity and quantum theory in one statement.

Ref. 2, Ref. 4

Relativity and quantum theory show that the idea of infinity appears only in *approximate* descriptions of nature; it disappears when talking with precision. Nothing in nature

* ‘Above we have seen that in the real world, the infinite is nowhere to be found, whatever experiences and observations and whatever knowledge we appeal to.’

Ref. 263

is infinite. For example, we found in volume II that the sky is dark at night (also) because space is not infinite. And we found, in volumes IV and V, that quantum theory contains probabilities because there is a smallest action value in nature. In fact, the statement that a quantity is infinitely large or infinitely small cannot be confirmed or reproduced by any experiment. Worse, such a statement is falsified by every measurement. In short, we found that infinity is a fantasy of the human mind. In nature, it does not appear. *Infinity about nature is always a lie.*

The number of particles, their possible positions, the states they can have, our brain, our creativity, our possible thoughts: all this is not infinite. Nevertheless, quantum theory
Ref. 4 and relativity changed the world: they allowed building ultrasound imaging, magnetic resonance imaging, lasers, satellite navigation systems, music players and the internet.

Despite the vast progress due to modern physics and the related technologies, one result remains: nothing in our environment is infinite – neither our life, nor our experiences, nor our memories, not even our dreams or our fantasies. Neither the information necessary to describe the universe, nor the paper to write down the formulae, nor the necessary ink, nor the time necessary to understand the formulae is infinite. Nature is not infinite. On the other hand, we also know that the illusion of the existence of infinity in nature is one the most persistent prejudices and myths ever conceived. Why did we use it in the first place?

The habit to use infinity to describe the world has many emotional reasons. For some, it reflects the deep-rooted experience of smallness that we carry within us as a remnant our personal history, when the world seemed so large and powerful. For others, the idea of our smallness allows us to deny somehow the responsibility for our actions or the existence of death. For others again, the idea of a finite universe often, at a first glance, produces deception, disbelief and discouragement. The absence of infinity means that we cannot achieve everything we want, and that our dreams and our possibilities are limited. Clinging to the idea of infinity is a way to avoid confronting this reality.

However, once we face and accept the absence of infinity, we make a powerful experience. We gain in strength. We are freed from the power of those who use this myth
Challenge 221 e to put themselves above others. It is an illuminating experience to reread all those sentences on nature, on the world and on the universe containing the term 'infinite', knowing that they are incorrect, and then clearly experience the manipulations behind them. The desire to make others bow to what is called the infinite is a common type of human violence.

At first, the demise of infinity might also bring panic fear, because it can appear as a lack of guidance. But at closer inspection, the absence of infinity brings strength. Indeed, the elimination of infinity takes from people one of the deepest fears: the fear of being weak and insignificant.

Moreover, once we face the limits of nature, we react like in all those situations in which we encounter a boundary: the limit becomes a challenge. For example, the experience that all bodies unavoidably fall makes parachuting so thrilling. The recognition that our life is finite produces the fire to live it to the full. The knowledge of death gives meaning to our actions. In an infinite life, every act could be postponed without any consequence. The disappearance of infinity generates creativity. A world without limits is discouraging and depressing. Infinity is empty; limits are a source of strength and pour passion into our life. Only the limits of the world ensure that every additional step in life brings us forward. Only in a limited universe is progress possible and sensible. Who is wiser, the one who denies limits, or the one who accepts them? And who lives more intensely?

Unification: the absence of finitude

> Pray be always in motion. Early in the morning go and see things; and the rest of the day go and see people. If you stay but a week at a place, and that an insignificant one, see, however, all that is to be seen there; know as many people, and get into as many houses as ever you can.
> Philip Stanhope,* *Letters to his Son on the Fine Art of Becoming a Man of the World and a Gentleman.*

The last part of our adventure, described in this volume, produced an unexpected result. Not only is nature not infinite; nature is not finite either. None of the quantities which were supposed to be finite turn out to be so. Finitude turns out to be an approximation, or better, an illusion, though a subtle one. *Nature is not finite.* This is the unification of physics in one statement.

Precise observation shows that nothing in nature can be counted. If nature were finite it would have to be (described by) a set. However, the exploration of Planck scales shows that such a description is intrinsically incomplete and inaccurate. Indeed, a description of nature by a set can never explain the number of its elements, and thus cannot explain finitude itself. In other words, any approach that tries to describe nature as finite is a belief, and is never correct. *Finitude is a lie.*

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We thus lost our security of thought a second time. Nature is neither infinite nor finite. We explored the possibilities left over and found that only one option is left: *Nature is indivisible.* In other words, all parts that we experience are approximations. Both finitude and infinity are approximation of nature. All distinctions are approximate. This central conclusion solved the remaining open issues about motion. *Nature has no parts.*

The impossibility to count and the lack of parts imply that nature is not a computer, not an automaton, nor a physical system. *Nature is not discrete.*

Recognizing all distinctions as being approximate abolishes the distinction between the permanent aspects of nature ('objects', described by mass, charge, spin, etc.) and the changing aspects ('states', described by position, momentum, energy). Taking all distinctions as approximate introduces extended constituents: fluctuating strands. Looking even closer, these extended constituents are all the same one. Space, formally only used to describe states, also acquires changing aspects: it is made from fluctuating strands. Also properties like mass or charge, which formally were seen as static, become aspects of the ever changing interplay between these fundamental constituents. Describing nature as one fluctuating strand allows us to avoid finitude and to answer all questions left open by quantum theory and general relativity.

In a sense, the merging of objects and states is a resolution of the contrasting views on motion of the Greek thinkers Parmenides – 'there is no motion', i.e., in physical language, 'there are no states, there is only permanence' – and Heraclitus – 'everything

* Philip D. Stanhope (b. 1694 London, d. 1773 London) was a statesman and writer.

moves', i.e., in physical language 'there is no permanence, there are only states'. Both turn out to be right.

We can thus sum up the progress during our adventure of physics in the following table:

TABLE 18 The progress of physics.

Step 1	Galilean Physics	Nature is continuous.	We live in Galilean space.
Step 2	Relativity	Nature has no infinitely large.	We live in Riemannian space.
Step 3	Quantum field theory	Nature has no infinitely small.	We live in a Hilbert/Fock space.
Step 4	Unification	Nature is not finite. Nature has no parts.	We do not live in any space; we are space.

In summary, we are made of space. More precisely, we are made of the same constituents as space. In fact, the fascination of this result goes further than that.

NEW SIGHTS

"Nel suo profondo vidi che s'interna,
legato con amore in un volume,
ciò che per l'universo si squaderna:

sustanze e accidenti e lor costume
quasi conflati insieme, per tal modo
che ciò ch'i' dico è un semplice lume.

La forma universal di questo nodo
credo ch'i' vidi, perché più di largo,
dicendo questo, mi sento ch'i' godo.*"

Dante, *La Divina Commedia*, Paradiso, XXXIII, 85-93.

Modelling nature as a complicated web of fluctuating strands allowed us to describe at the same time empty space, matter, radiation, horizons, kefir, stars, children and all our other observations. All everyday experiences are consequence of everything in nature being made of one connected strand. This result literally widens our horizon.

* 'In its depth I saw gathered, bound with love into one volume, that which unfolds throughout the universe: substances and accidents and their relations almost joined together, in such a manner that what I say is only a simple image. The universal form of that knot, I think I saw, because, while I am telling about it, I feel deep joy.' This is, in nine lines, Dante's poetic description of his deepest mystical experience: the vision of god. For Dante, god, at the depth of the light it emanates, is a knot. That knot spreads throughout the universe, and substances and accidents – physicists would say: particles and states – are aspects of that knot. Dante Alighieri (b. 1265 Florence, d. 1321 Ravenna) was one of the founders and the most important poet of the Italian language. Most of the Divine Comedy, his magnum opus, was written in exile, after 1302, the year when he had been condemned to death in Florence.

The beauty of strands

> Someday, surely, we will see the principle underlying existence itself as so simple, so beautiful, so obvious, that we will all say to each other, "Oh, how could we all have been so blind, so long."
> John Wheeler, *A Journey Into Gravity And Spacetime.*

Describing everything as connected does not come natural to us humans. After all, in our life, we perform only one act: to partition. We define pluralities. There is no way we can avoid doing this. To observe, to think, to talk, to take a decision, to move, to suffer, to love or to enjoy life is impossible without partitioning.

Our walk showed us that there are limits to the ability to distinguish. Any kind of partitioning is always approximate. In fact, most people can summarize their personal experience by saying that they learned to make finer and finer distinctions. However, talking with highest precision about a part of the world inevitably leads to talk about the whole universe. The situation resembles a person who gets a piece of rope in his hand, and by following it, discovers a large net. He continues to pull and finally discovers that everything, including himself, is part of the net.

For the strand model, the term 'theory of everything' is therefore not acceptable. Nature cannot be divided into 'things'. In nature, things are never separable. There is no way to speak of 'every' thing; there are no sets, no elements and no parts in nature. A theory describing all of nature cannot be one of 'everything', as 'things' are only approximate entities: properly speaking, they do not exist. The strand model is not a theory of everything; it is a *complete theory*.

The strand model shows that nature is not made of related parts. Nature is made of relations only. Parts only exist approximately. The strand model also shows: being in motion is intrinsic to being a part. Parts, being approximate, are always in motion. As soon as we divide, we observe motion. The act of dividing, of partitioning, of defining parts is the very one which produces order out of chaos. Strands force us to rethink this habit.

Despite being so tough to grasp, strands yield a precise description of motion that unifies quantum field theory and general relativity. The strand model for the unification of motion is both simple and powerful. There are no free parameters. There are no questions left. Our view from the top of the mountain is thus complete. No uncertainty, no darkness, no fear and no insecurity are left over. Only wonder remains.

Can the strand model be generalized?

> Die Natur kann besser Physik als der beste Physiker.*
> Carl Ramsauer

Page 166 As mentioned above, mathematical physicists are fond of *generalizing* models. Despite this fondness, we required that any complete, unified description must be unique: any complete, unified description must be impossible to reduce, to modify or to generalize.

* 'Nature knows physics better than the best physicist.' Carl Ramsauer (b. 1879 Oldenburg, d. 1955 Berlin), influential physicist, discovered that electrons behave as waves.

In particular, a unified theory must neither be a generalization of particle physics nor of general relativity. Let us check this.

The strand model is not a generalization of general relativity: the definitions of curvature, of gravitons and of horizons differ radically from general relativity's approach. The strand model is also not a generalization of particle physics: the definitions of particle and of interactions differ radically from the concepts of quantum field theory. Indeed, we have shown that quantum field theory and general relativity are *approximations* to the strand model; they are neither special cases nor reductions of the strand model.

But what about the other requirements for a unified theory? Can the strand model be modified or generalized? We have seen that the model does not work in more spatial dimensions, does not work with more families of quarks, does not work with more interactions, and does not work with other evolution equations in general relativity or particle physics. The strand model does not work with other fundamental constituents, such as bifurcating entities, membranes, bands, or networks. (Though it does work with Page 172 the equivalent *funnels*, as explained earlier on, but that description is equivalent to the one with strands.) The strand model does not work with any modified fundamental principle. Obviously, exploring all possible variations and modifications remains a challenge Challenge 222 r for the years to come. If an actual modification of the strand model can be found, the strand model instantly loses its value: in that case, it would need to be shelved as a failure. Only a *unique* unified model can be correct.

In summary, one of the beautiful aspects of the strand model is its radical departure from twentieth-century physics in its basic concepts, combined with its almost incredible uniqueness. No generalization, no specialization and no modification of the strand Ref. 156 model seems possible. In short, the strand model qualifies as a unified, complete theory.

What is a requirement to one person, is a criticism to another. A number of researchers deeply dislike the strand model precisely because it doesn't generalize previous theories and because it cannot be generalized. This attitude deserves respect, as it is born from the admiration for several ancient masters of physics. However, the strand model points into a different direction.

What is nature?

> “Nature is what is whole in each of its parts.”
> Hermes Trismegistos, *Book of Twenty-four Philosophers*.

Ref. 264

At the end of our long adventure, we discovered that nature is not a set: everything is connected. Nature is only *approximately* a set. The universe has no topology, because space-time is not a manifold. Nevertheless, the approximate topology of the universe is that of an open Riemannian space. The universe has no definite particle number, because the universe is not a container; the universe is made of the same stuff of which particles are made. Nevertheless, the approximate particle density in the universe can be deduced.

In nature, everything is connected. This observation is reflected in the conjecture that all of nature is described by a single strand.

We thus arrive at the (slightly edited) summary given around the year 1200 by the author who wrote under the pen name Hermes Trismegistos: *Nature is what is whole in* *each of its parts.* But in contrast to that author, we now also know how to draw testable conclusions from the statement.

Quantum theory and the nature of matter and vacuum

> In everything there is something of everything.
>
> Anaxagoras of Clazimenes (500 –428 bce Lampsacus)

The strand model shows that as soon as we separate the universe into space-time and the rest, i.e., as soon as we introduce the coordinates x and t, quantum mechanics appears automatically. More precisely, *quantum effects are effects of extension*. Quantum theory appears when we realize that observations are composed of smallest events due to crossing switches, each with a change given by the quantum of action. All events and observations appear through the fluctuations of the strand that composes nature.

We found that *matter is made of tangled strands.* In fact, the correct way would be to say: matter is made of tangled strand *segments.* This connection leads to Schrödinger's equation and to Dirac's equation.

Insofar as matter is of the same fabric as the vacuum, we can rightly say that *everything is made of vacuum* and that *matter is made of nothing.* But the most appropriate definition arises when we realize that matter is not made from something, but that matter is a certain aspect of the *whole* of nature. Unification showed that every single elementary particle results from an arrangement of strands that involves the whole of nature, or, if we prefer, the entire universe. In other words, we can equally say: *matter is made of everything.*

We can also turn the equivalence of matter and vacuum around. Doing so, we arrive at the almost absurd statement: *vacuum is made of everything.*

> Der heutigen Physik liegt die Frage nicht mehr ferne, ob nicht etwa alles, was ist, aus dem Äther geschaffen sei. Diese Dinge sind die äußersten Ziele unserer Wissenschaft, der Physik.*
>
> Heinrich Hertz

Cosmology

The strand model also showed us how to deduce general relativity. The strand model clarified the fabric of horizons and explained the three dimensions of space. Most fascinating is the idea of a universe as the product of a single strand. A single strand implies that there was nothing before the big bang, and that there is nothing outside the night sky. For example, the strand model implies that there is no 'multiverse' and that there are no hidden worlds of any kind. And the fluctuating strand explains all observations of our universe.

The cosmological constant is not constant; it only measures the present age and size of
Page 8 the universe. Therefore, the constant does not need to appear in Figure 1. In other words, the cosmological constant simply measures the time from the big bang to the present.

* 'Modern physics is not far from the question whether everything that exists could possibly be made from aether. These things are the extreme goals of our science, physics.' Hertz said this in a well-known speech he gave in 1889. If we recall that 'aether' was the term of the time for 'vacuum', the citation is particularly striking.

The ‘big bang’ is the name for what we observe if we try to make observations approaching the limits of nature. The ‘big bang’ appears automatically from the strand model whenever we observe nature at the most distant times, the largest distances or at the largest energies: ‘big bang’ is the name for Planck scale physics.

The universe consists of a single strand. There are many particles in nature, because the strand is tangled up in complicated ways. What we call the ‘horizon’ of the universe is the place where new tangles appear.

The belief that the big bang or the horizon are examples of creation is incorrect. What happened at the big bang still happens at the horizon today. Both the black sky at night and the big bang are nature’s way to tell us: ‘Galilean physics is approximate! Quantum theory is approximate! General relativity is approximate!’

Musings about unification and strands

> “Continuing motion masters coldness.
> Continuing rest masters heat.
> Motion based on rest:
> Measure of the all-happening for the single one.”
> Lao Tse,* *Tao Te King*, XXXXV.

All is made from one sort of thing: all is one substance. This idea, *monism*, sounds a lot like what the influential philosopher Baruch Spinoza (b. 1632 Amsterdam, d. 1677 The Hague) held as conviction. Monism, though mixed up with the idea of god, is also the basis of the philosophical ideas that Gottfried Wilhelm Leibniz (b. 1646 Leipzig, d. 1716 Hannover) presents in his text *La Monadologie*.

* *

Any complete theory of motion, also the strand model, is built on a single statement about nature: The *many* exists only approximately. Nature is approximately multiple. The etymological meaning of the term ‘multiple’ is ‘it has many folds’; in a very specific sense, nature thus has many folds. Ref. 265

* *

Any precise description of nature is free of arbitrary choices, because the divisions that we have to make in order to think are all common to everybody, and logically inescapable. Because physics is a consequence of this division, it is also ‘theory-free’ and ‘interpretation-free’. This consequence of the complete theory will drive most philosophers up the wall.

* *

For over a century, physics students have been bombarded with the statement: ‘Symmetries are beautiful.’ Every expert on beauty, be it a painter, an architect, a sculptor, a musician, a photographer or a designer, fully and completely disagrees, and rightly so. Beauty has no relation to symmetry. Whoever says the contrary is blocking out his experiences of a beautiful landscape, of a beautiful human figure or of a beautiful work of art.

* Lao Tse (sixth century BCE) was an influential philosopher and sage.

The correct statement is: 'Symmetries simplify descriptions.' Symmetries simplify physical theories. That is the background for the statement of Werner Heisenberg: 'In the beginning there was symmetry.' On the other hand, the strand model shows that even this statement is incorrect. In fact, neither the search for beauty nor the search for symmetry were the right paths to advance towards unification. Such statements have always been empty marketing phrases. In reality, the progress of fundamental theoretical physics was always driven by the search for *simplicity*.

Ref. 266

* *

Strands unify physics. In particular, strands extend our views on quantum theory and mathematical physics, on particle physics and field theory, on axiomatic physics and algebraic physics, on polymer physics and gauge theory, on general relativity and cosmology. It will take several years before all these extensions will have been explored.

* *

The description of nature with strands is surprisingly simple, mainly because it uses so few basic concepts. Is this result astonishing? In our daily life, we describe our experiences with the help of a few thousand words, e.g. taking them from the roughly 350 000 words which make up the English language, or from a similar number from another language. This set is sufficient to talk about everything, from love to suffering, from beauty to happiness. And these terms are constructed from no more than about 35 basic concepts, as we have seen already. We should not be too surprised that we can in fact talk about the whole universe using only a few basic concepts: the act and the results of (approximate) distinction, or more specifically, a basic event – the crossing switch – and its observation.

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* *

Almost all discoveries in physics were made at least 30 years too late. The same is true for the strand model. If we compare the strand model with what many physicists believed in the twentieth century, we can see why: researchers had too many wrong ideas about unification. All these wrong ideas can be summarized in the following statement:

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— 'Unification requires generalization of existing theories.'

This statement is subtle: it was rarely expressed explicitly but widely believed. But the statement is wrong, and it led many astray. On the other hand, the development of the strand model also followed a specific guiding idea, namely:

— 'Unification requires simplification.'

Hopefully this guiding idea will not become a dogma itself: in many domains of life, simplification means not to pay attention to the details. This attitude does a lot of harm.

* *

The strand model shows that achieving unification is not a feat requiring difficult abstraction. Unification was not hidden in some almost inaccessible place that can reached only by a few select, well-trained research scientists. No, unification is accessible to everyone who has a basic knowledge of nature and of physics. No Ph.D. in theoretical physics is

FIGURE 139 Motion Mountain does not resemble Cerro Torre, but a gentle hill (© Davide Brighenti, Myriam70)

needed to understand or to enjoy it. The knowledge presented in the previous volumes of this series is sufficient.

When Andrew Wiles first proved Fermat's last theorem after three centuries of attempts by the brightest and the best mathematicians, he explained that his search for a proof was like the exploration of a dark mansion. And seen the conceptual difficulties he had to overcome, the analogy was fitting. Recalling how many more people have already searched for unification without success, the first reaction is to compare the search for unification to the exploration of something even bigger, such as a complex dark cave system. But that analogy was not helpful. In contrast to the proof of Fermat's theorem, the goal of the quest for unification turned out to be simple and lying out in the open. Researchers had simply overlooked it, because they were convinced that the goal was complex, hidden in the dark and hard to reach. It was not.

The adventure of climbing Motion Mountain is thus not comparable to climbing Cerro Torre, which might be the toughest and most spectacular challenge that nature offers to mountain climbers. Figure 139 gives an impression of the peak. Motion Mountain does not resemble this peak at all. Neither does Motion Mountain resemble the Langtang Lirung peak in the Nepalese Himalayas shown on the cover of this volume. Climbing Motion Mountain is more like walking up a gentle green hill, alone, with a serene mind, on a sunny day, while enjoying the surrounding beauty of nature.

* *

Page 85 The strand model settles all questions about *determinism*. Quantum theory and general relativity are deterministic. Nevertheless, when both descriptions are combined, time turns out to be an approximate, low-energy concept. The same applies to determinism. Even though nature is deterministic for all practical purposes and shows no surprises, determinism shares the fate of all its conceivable opposites, such as fundamental randomness, indeterminism of all kinds, existence of wonders, creation out of nothing, or divine intervention: determinism is an *incorrect* description of nature at the Planck scale – like all its alternatives.

* *

The strand model also settles most so-called *really big questions* that John Wheeler used to ask: Why the quantum? How come existence? It from bit? A "participatory universe"? What makes "meaning"? Enjoy the exploration. Challenge 223 e

* *

Any unified model of nature encompasses a lot of ideas, issues and knowledge. Due to the sheer amount of material, publishing it in a journal will be challenging.

* *

Ref. 267 The strand model is so simple that it fits on a tombstone - or on a T-shirt. This would surely be god's favourite T-shirt. It is available at www.motionmountain.net/gfts.html.

* *

Historically, the strand model evolved from an exploration, started in the 1990s, of the maximum force in nature, the belt trick and the entropy of black holes. After the first six chapters of the present volume were completed in 2002, meditating on their implications led to the strand model and its fundamental principle.

Above all, it was the description of general relativity with the help of the maximum force that triggered the search for a unified description that was purely based on Planck units. Another essential point was the drive to search for a complete theory directly, from its requirements ('top down' in Figure 1), and *not* from the unification of quantum theory and general relativity ('bottom up'). Page 8 In the years from 2002 to 2007, most of the ideas of the strand model took shape, mainly in Munich's underground trains, while commuting between home and work. In those years, it appeared that strands could explain the Dirac equation, the entropy of black holes, general relativity and the particle spectrum with the three particle generations. While walking in the woods and fields around Munich, on 13 October 2008, it appeared that interactions are core deformations; in subsequent walks during 2008 and 2009, it appeared that strands explain the three gauge interactions, predict the lack of a Higgs boson – a bad mistake due to faulty reasoning, as turned out in 2012 – and of any new physical effects beyond the standard model, and Page 330 allow calculating the unexplained constants of particle physics. The model thus yielded almost all its main predictions before the accelerator experiments at the Large Hadron Collider at CERN in Geneva were switched on in autumn 2010. Thus much of the work was done in a haste – future will show what is of lasting value.

In 2012, the discovery of the Higgs boson, and in 2014, the comments by Sergei Fadeev led to an improvement and simplification of the strand model, eliminating knotted strands. From 2016 onwards, the experimental results of the LHC groups, of dark matter searches, and of the LIGO observatory confirmed the lack of deviations from the standard model of particle physics and from general relativity, as predicted by the strand model.

* *

Many researchers believed during all their life that the complete theory is something useful, important and valuable. This common belief about the importance and seriousness of the quest has led, over the past decades, to an increasingly aggressive atmosphere among these researchers. This unprofessional atmosphere, combined with the dependence of researchers on funding, has delayed the discovery of the complete theory by several decades.

In fact, the complete theory is *not useful:* it adds nothing of practical relevance to the combination of the standard model and general relativity. The unified theory is also *not important:* it has no application in everyday life or in industry and does not substantially change our view of the world; it just influences teaching – somewhat. Finally, the unified theory is *not valuable:* it does not help people in their life or make them happier. In short, the complete theory is what all fundamental theoretical research is: entertaining ideas.

Even if the strand model were to be replaced by another model, the conclusion remains: the unified theory is not useful, not important and not valuable. Knowing about unification does not confer any special powers. But it is enjoyable, comparable to a walk through a beautiful garden.

* *

The strand model will take a long time to get accepted. The first reason is obvious: *The strand model contradicts thinking habits* in many research fields. Researchers working on the foundations of quantum theory, on general relativity, on cosmic strings, on mathematical physics, on classical and quantum field theory, on polymer physics, on shape deformations, on quantum gravity, on strings, on the visualization of quantum mechanics, on knot theory, on higher dimensions, on supersymmetry, on the axiomatization of physics, on group theory, on the foundation of physics, on quantum optics and on particle physics have to give up many life-long thinking habits. So do all other physicists. *Strands supersede particles and points.*

There is also a second reason for the slow acceptance of the model presented here: *The strand model, in its simplicity, is only a small step away from present research.* Many researchers are finding out how close they have been to the ideas of the strand model, and for how long they were overlooking or ignoring such a simple option. The simplicity of the fundamental principle contrasts with the expectation of most researchers, namely that the unified theory is complicated, difficult and hard to discover. In fact, the opposite
Page 87
is true. *Strands are based on Planck units and provide a simple, almost algebraic description of nature.*

In summary, for many researchers and for many physicists, there is a mixture of confusion, anger and disappointment. It will take time before these feelings subside and are replaced by the clarity and fascination provided by the strand model.

“Only boring people get bored.

Anonymous”

The elimination of induction

> Cum iam profeceris tantum, ut sit tibi etiam tui reverentia, licebit dimittas pedagogum.*
>
> Seneca

The complete theory of motion has a consequence worth mentioning in detail: its lack of infinity and its lack of finitude eliminate the necessity of induction. This conclusion is of importance for general discussions on man's grasp of nature.

In physics, as in the other natural sciences, there is a tradition to state that a certain description of nature – once confusingly called a 'law' – is valid in *all* cases. In these statements, 'all' means 'for all values of the quantities appearing'. As a concrete example, the 'law' of universal gravitation is always claimed to be the same here and today, as well as at *all* other places and times, such as on the other end of the universe and in a few thousand years. The full list of such all-claims is part of the millennium list of open issues in twentieth-century physics. For many decades, the habit of claiming general validity from a limited and finite number of experiences, also called *induction*, has been seen, and rightly so, as a logically dubious manoeuvre, tolerated only because it works. But the developments described in this text show that this method is indeed justified.

Page 164

First of all, a claim of generality is not that enormous as it may seem, because the number of events that can be distinguished in nature is finite, not infinite. The preceding sections showed that the maximal number N of events that can be distinguished in the universe is of the order of $N = (T_0/t_{\text{Pl}})^4 = 10^{244\pm2}$, T_0 being the age of the universe and t_{Pl} the Planck time. This is a big, but certainly finite number.

The unified description of nature has thus first reduced the various all-claims from an apparently infinite to a finite number of cases, though still involving astronomically large numbers. This reduction results from the recognition that infinities do not appear in the description of nature. We now know that when talking about nature, 'all' cases never means an infinite number.

A second, important result is achieved by the description of nature with strands. In any all-claim about fundamental motion, the checking of each of the large number of possibilities is not necessary any more, because all events result from a single entity, in which we introduce distinctions with our senses and our brain. And the distinctions we introduce imply automatically that the symmetries of nature – the 'all-claims' or 'inductions' – that are used in the description of motion are correct. Nature does not contain separate parts. Therefore, there is no way that separate parts can behave differently. Induction is a consequence of the unity of nature.

Ultimately, the possibility to *verify* statements of nature is due to the fact that all the aspects of our experience are *related*. Complete separation is impossible in nature. The verification of all-claims is possible because the strand model achieves the full description of how all 'parts' of nature are related.

The strand model shows that we can talk and think about nature because we are a part of it. The strand model also shows that induction works because everything in nature is related to everything else: nature is one.

* 'When you have profited so much that you respect yourself you may let go your tutor.' Seneca, the influential Roman poet and philosopher, writes this in his *Epistulae morales ad Lucilium*, XXV, 6.

What is still hidden?

> "That which eludes curiosity can be grasped in action."
> Traditional saying.

Where do we come from? Where does the world come from? What will future bring? What is death? All these questions are questions about motion – and its meaning. To all such questions, the strand model does not provide answers. We are a collection of tangled strands. We are everything and nothing. The strand(s) we are made of will continue to fluctuate. Birth, life and death are aspects of tangled strands. The universe is a folded strand that grows in complexity.

Obviously, abstract statements about tangles do not help in any human quest. Indeed, we aimed at achieving a precise description of moving particles and bending space. Studying them was a sequence of riddles; but solving these riddles does not provide meaning, not even at the top of Motion Mountain. From the top we cannot see the evolution of complicated systems; in particular, we cannot see or describe the evolution of life, the biological evolution of species, or the growth of a human beings. Nor can we understand why we are climbing at all. Challenge 224 s

In short, from the top of Motion Mountain we cannot see the details down in the valleys of human relations or experience; strands do not provide advice or meaning. Remaining too long on the top is of no use. To find meaning, we have to descend back down to real life. Vol. I, page 16

A RETURN PATH: JE RÊVE, DONC JE SUIS

> "I hate reality. But it is the only place where one can get a good steak."
> Woody Allen

Enjoying life and giving it meaning requires to descend from the top of Motion Mountain. The return path can take various different directions. From a mountain, the most beautiful and direct descent might be the use of a paraglider. After our adventure, we take an equally beautiful way: we leave reality.

The usual trail to study motion, also the one of this text, starts from our ability to talk about nature to somebody else. From this ability we deduced our description of nature, starting from Galilean physics and ending with the strand model. The same results can be found by requiring to be able to talk about nature to ourselves. Talking to oneself is an example of thinking. We should therefore be able to derive all physics from René Descartes' sentence 'je pense, donc je suis' – which he translated into Latin as 'cogito ergo sum'. Descartes stressed that this is the only statement of which he is completely sure, in opposition to his observations, of which he is not. He had collected numerous examples in which the senses provide unreliable information. Ref. 268

However, when talking to ourselves, we can make more mistakes than when asking for checks from others. Let us approach this issue in a radically different way. We directly proceed to that situation in which the highest freedom is available and the largest number of mistakes are possible: the world of dreams. If nature would only be a dream, could we deduce from it the complete set of physical knowledge? Let us explore the issue.

— Dreaming implies the use of distinctions, of memory and of sight. Dreams contain *parts* and *motion*.
— Independently on whether dreams are due to previous observations or to fantasies, through memory we can define a sequence among them. The order relation is called *time*. The dream aspects being ordered are called *events*. The set of all (dream) events forms the (dream) *world*.

Ref. 269

— In a dream we can have several independent experiences at the same time, e.g. about thirst and about hunger. Sequences thus do not provide a complete classification of experiences. We call the necessary additional distinction *space*. Dream space has three dimensions.* Dreaming thus means to use space and time.
— We can distinguish between dream contents. Distinguishing means that we can count items in dreams. Counting means that we have a way to define measurements. Dreams are thus characterized by something which we can call 'observables'. Dream experiences at a given instant of time are characterized by a *state*.
— Because we can describe dreams, the dream contents exist independently of dream time. We can also imagine the same dream contents at different places and different times in the dream space. There is thus an invariance of dream concepts in space and time. There are thus symmetries in dream space.
— Dream contents can interact. Dreams appear to vary without end. Dreams seem to be infinite.

In other words, a large part of the world of dreams is described by a modified form of *Galilean physics*. We note that the biggest difference between dreams and nature is the lack of conservation. In dreams, observations can appear, disappear, start and stop. We also note that instead of dreams, we could equally explore cinema *films*. Films, like dreams, are described by a modified form of Galilean physics. And films, like dreams, do not follow conservation laws. But dreams teach us much more.

— Dreams show that space can warp.
— Dream motion, as you may want to check, shows a maximum speed.
— Dreams show a strange limit in distance. There is a boundary to our field of vision, even though we do not manage to see it.

Challenge 225 s

Pondering these issues shows that there are *limits* to dreams. In summary, the world of dreams has a maximum size, a maximum speed and three dimensions that can warp. The world of dreams and of films is described by a simple form of *general relativity*.

— Both the number of items we can dream of at the same time and the memory of previous dreams is finite.
— Dreams have colours.
— There are pixels in dreams, though we do not experience them directly. But we can do so indirectly: The existence of a highest number of things we can dream of at the same time implies that dream space has a smallest scale.

* Though a few mathematicians state that they can *think* in more than three spatial dimensions, all of them *dream* in three dimensions.

In summary, the world of dreams has something similar to a minimum change. The world of dreams and that of films is described by a simple form of *quantum theory*. The difference with nature is that in dreams and films, space is discrete from the outset. But there is still more to say about dreams.

— There is no way to say that dream images are made of mathematical points, as there is nothing smaller than pixels.
— In dreams, we cannot clearly distinguish objects ('matter') and environment ('space'); they often mix.
— In dreams, fluctuations appear both for images as well as for the background.
— In dreams, sharp distinctions are impossible. Dream space-time cannot be a set.
— Dream motion appears when approximate conservation (over time) is observed.
— In dreams, dimensionality at small distances is not clear; two and three dimensions are mixed up there.

In summary, the world of dreams seems to behave as if points and point particles do not exist; and since quantum theory and general relativity hold, the world of dreams seems to be described by extended constituents! We thus conclude this short exploration of the physics of dreams with a fascinating *conjecture*: even if nature would be a dream, an illusion or a fantasy, we might still get most of the results that we discovered in our ascent Challenge 226 s of Motion Mountain. (What differences with modern physics would be left?) Speaking with tongue in cheek, the fear of our own faults of judgement, so rightly underlined by Descartes and many others after him, might not apply to fundamental physics.

WHAT IS THE ORIGIN OF COLOURS?

All colours around us are determined by the fine structure constant α – the coupling constant for the electromagnetic interaction at low energy – with its measured value Ref. 5 1/137.035 999 139(31). The fine structure constant is also essential to describe most everyday devices and machines, as well as all human thoughts and movements. The constant is an aspect of every electric charge in nature.

The strand model showed us that electrical charge is a property of tangles of strands. In particular, the strand model showed:

▷ The fine structure constant describes the probability that a fluctuation adds a twist to the chiral tangles of electrically charged particles.

We have not yet deduced an accurate value for the fine structure constant, but we seem Ref. 262 to have found out how to do so. In short, we seem to glimpse the origin of all colours – and thus of all beauty around us. Strands provide a beautiful explanation for beauty.

SUMMARY: WHAT IS MOTION?

> Deep rest is motion in itself. Its motion rests in itself.
> Lao Tse, *Tao Te King*, VI, as translated by Walter Jerven.

We can now answer the question that drove us through our adventure:

▷ ***Motion*** is the observation of crossing switches of the one, unobservable, tangled and fluctuating strand that describes all of nature.

Nature's strand forms particles, horizons and space-time: these are the parts of nature. Particles are tangles of strands; horizons and space-time are weaves of strands. The parts of nature move. The parts move because their strands fluctuate.

Motion appears because all parts in nature are approximate. Indeed, the observation of crossing switches and the description of strand segments fluctuating in a background space result and are possible because we approximate from the one strand that makes up nature to the many parts inside nature. The one strand (approximately) forms the many elementary particles inside us. Strand segments and particles (approximately) lead us to introduce background space, matter and radiation. Introducing background space implies observing motion. Motion thus appears automatically when approximate parts of nature, such as humans, animals or machines, describe other approximate parts of nature, such as other bodies or systems.

The observation of motion is due to our introduction of the plural. Motion results from of our forced use of *many (approximate) parts* to describe the *unity* of nature. The observation of motion results from approximations. All these approximate distinctions are unavoidable and are due to the limitations of our experience.

Motion appears as soon as we divide the world into parts and then follow these parts. Dividing nature into parts is not a conscious act; our human nature – our senses and our brain – *force* us to perform it. And whenever we experience or talk about parts of the universe, we find motion. Our senses and our brain are made to distinguish and to divide – and cannot do otherwise. We need to distinguish in order to survive, to think and to enjoy life. In a sense, we can say that motion appears as a logical consequence of our limitations; the fundamental limitation is the one that makes us distinguish and introduce parts, including points and sets.

Motion is an 'artefact' of locality. Locality is an approximation and is due to our human nature. Distinction, localization and motion are inextricably linked.

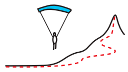
Motion is low energy concept. Motion does not exist at Planck scales, i.e., at the limits of nature.

Motion is an artefact due to our limitations. This conclusion resembles what Zeno of Elea stated 2500 years ago, that motion is an illusion. But in contrast to Zeno's pessimistic view, we now have a fascinating spectrum of results and tools at our disposition: they allow us to describe motion and nature with high precision. Most of all, these tools allow us to change ourselves and our environment for the better.

Vol. I, page 15

Ref. 270

> All the great things that have happened in the world first took place in a person's imagination, and how tomorrow's world will look like will largely depend on the power of imagination of those who are just learning to read right now.
>
> Astrid Lindgren*



* Astrid Lindgren (b. 1907 Näs, d. 2002 Stockholm) was a beloved writer of children books.

POSTFACE

Perhaps once you will read Plato's *Phaedrus*, one of the beautiful philosophical Greek texts. In it, Socrates is made to say that he almost never left the city walls because to him, as a 'lover of learning, trees and the open country do not teach anything, whereas men in the town do.' This is a veiled critique of Democritus, the most important and famous philosopher in Greece during Plato's time. Democritus was the natural philosopher par excellence, and arguably had learned from nature – with its trees and open country – more than anybody else after him.

After this adventure you can decide for yourself which of these two approaches is more congenial to you. It might be useful to know that Aristotle, Plato's pupil and the most influential Greek thinker, refused to choose and cultivated them both. There is no alternative in life to following one's own mind, and to enjoy doing so. If you enjoyed this particular trip, show it to your friends. For yourself, after this walk, sense intensively the pleasure of having accomplished something important. Many before you did not have the occasion. Enjoy the beauty of the view offered. Enjoy the vastness of horizon it provides. Enjoy the impressions that it creates inside you. Collect them and rest. You will have a treasure that will be useful in many occasions. Then, when you feel the desire of going further, get ready for another of the adventures life has to offer.

Plato's *Phaedrus*, written around 380 BCE, is available in many pocket editions. Do not waste your time learning ancient Greek to read it; the translated versions are as beautiful as the original.

Plato's lifelong avoidance of the natural sciences had two reasons. First of all, he was jealous of Democritus. Plato never even cites Democritus in his texts. Democritus was the most prolific, daring, admired and successful philosopher of his time (and maybe of all times). Democritus was a keen student of nature. His written works did not survive, because his studies were not congenial to the followers of christianity, and thus they were not copied by the monks in the Middle Ages. The loss of these texts is related to the second reason that kept Plato away from the natural sciences: he wanted to save his life. Plato had learned one thing from men in the town: talking about nature is dangerous. Starting around his lifetime, for over 2000 years people practising the natural sciences were regularly condemned to exile or to death for impiety. Fortunately, this is only rarely the case today. But such violence still occurs, and we can honour the dangers that those preceding us had to overcome in order to allow us enjoying this adventure.

Appendix A

KNOT AND TANGLE GEOMETRY

The following table provides a terse summary of the mathematics of knot shapes.

TABLE 19 Important properties of knot, links and tangles.

Concept	Defining property	Other properties
Knot / link / tangle	one closed / several closed / one or several open curves, all in 3d and without intersections	ropelength is integral of arclength; ropelength is shape-dependent.
Ideal knot, link, tangle (shape)	tightest possible knot, link or tangle (shape) assuming a rope of constant diameter that is infinitely flexible and infinitely slippery	at present, all non-trivial ideal shapes are only known approximately; most ideal knots (almost surely) have kinks.
Ribbon or *framing*	short perpendicular (or non-tangent) vector attached at each point of a curve	
Curvature of a curve	inverse curvature radius of ‘touching’ circle	measures departure from straightness, i.e., local bending of a curve.
Normal vector or curvature vector	local vector normal to the curve, in direction of the centre of the ‘touching’ circle, with length given by the curvature	is given by the second and first derivatives of the curve.
Binormal vector	local unit vector normal to the tangent and to the normal/curvature vector	
Torsion	local speed of rotation of the binormal vector; positive (negative) for right-handed (left-handed) helix	measures departure from flatness, i.e., local twisting or local handedness of a curve; essentially a third derivative of the curve.
Frenet frame at a curve point	‘natural’ local orthogonal frame of reference defined by *unit* tangent, *unit* normal/curvature and binormal vector	the Frenet frame differs at each curve point, the Frenet frame is *not* uniquely defined if the curve is locally straight.

TABLE 19 (Continued) Important properties of knot, links and tangles.

Concept	Defining property	Other properties
'Natural' framing or *Frenet ribbon*	defined by the local normal, i.e., local curvature vector	for a closed curve, it is always closed and two-sided, and thus never a Moebius band.
Linking number between two closed curves	sloppily, number of times that two curves wind around each other, or, equivalently, half the number of times that the curves 'swap' position	topological invariant, i.e., shape-independent; $\mathrm{Lk}(K1, K2) = \frac{1}{4\pi} \oint_{K2} \oint_{K1} \frac{\boldsymbol{r}_{12}\,(\mathrm{d}\boldsymbol{r}_1 \times \mathrm{d}\boldsymbol{r}_2)}{\boldsymbol{r}_{12}^3}$.
Linking number for a closed two-sided ribbon	number of times that the edges wind around each other	topological invariant, i.e., shape-independent; always an integer.
Self-linking number or 'natural' linking number for a knot	number of times that the edges of the natural/Frenet ribbon wind around each other	not a topological invariant, because of existence of inflection points.
Link integral for an open curve	generalization of the linking number for knots to open curves	usually not an integer.
Twist of a ribbon, open or closed	$\mathrm{Tw}(R)$ is the total angle, in units of 2π, by which the ribbon rotates around the central axis of the ribbon; sloppily said, it measures the *local* helicity; this type of twist has no relation to the first Reidemeister move	vanishes for ribbons that are everywhere flat.
Twist of a curve or knot	$\mathrm{Tw}(K)$ is the total angle, in units of 2π, by which the Frenet frame rotates around the tangent direction, or equivalently, (total) twist of the Frenet ribbon, also called the *total torsion* of the curve; this type of twist has no relation to the first Reidemeister move	not an integer even in case of knots; depends on curve/knot shape; is different from zero for chiral curves/knots; is zero for achiral curves/knots that have a rigid reflective symmetry; twist and torsion are only equal if the twist is defined with the Frenet ribbon – with other framings they differ.
Signed crossing number	sum of positive minus sum of negative crossings in a given oriented 2d projection of a curve or knot (sometimes called '2d-writhe')	always an integer; depends on shape.
2d-writhe of a knot, or *topological writhe*, or *Tait number*	signed crossing number for a *minimal* crossing number diagram/projection (sometimes the term '2d-writhe' is used for the signed crossing number of *any* configuration)	is shape-invariant; is always an integer; differs from 0 for all chiral knots; has the value 3 for the trefoil, 0 for the figure-eight knot, 5 for the 5_1 and 5_2 knots, 2 for the 6_1 knot, 7 for the 7_1 and 7_2 knots, 4 for the 8_1 knot, and 9 for the 9_2 knot.

TABLE 19 (Continued) Important properties of knot, links and tangles.

Concept	Defining property	Other properties
Writhing number or 3d-*writhe* of a knot	$\mathrm{Wr}(K)$ is the average, over all projection directions, of the signed crossing number; sloppily said, it measures how wrapped, coiled and chiral a knot is, i.e., it measures the *global* helicity	depends on knot shape; usually is not an integer; is different from zero for chiral knots; is zero for achiral knots that have a rigid reflective symmetry; $\mathrm{Wr}(K) = \frac{1}{4\pi}\oint_K\oint_K \frac{\boldsymbol{r}_{12}\,(\mathrm{d}\boldsymbol{r}_1\times\mathrm{d}\boldsymbol{r}_2)}{\boldsymbol{r}_{12}^3}$; uses no ribbon and thus is independent of the ribbon shape attached to the knot.
Writhe of ideal, alternating knots and of odd-component links	the value is quasi-quantized for alternating knots with small crossing numbers (<11) in values that differ from $m4/7$ by only a few per cent	is additive under knot addition for knots with small crossing numbers (<11) within less than 1 %.
Writhe of ideal, alternating even-component links	the value is quasi-quantized for alternating links with small crossing numbers (<11) in values that differ from $2/7 + m4/7$ by only a few per cent	
Writhe of a ribbon	sloppily said, measures how wrapped, coiled and chiral a ribbon is, i.e., measures its *global* helicity	
Writhe of an open curve		vanishes for plane curves.
Calugareanu's theorem	for any knot K and any ribbon G attached to it, $\mathrm{Lk}(K,G) = \mathrm{Tw}(K,G) + \mathrm{Wr}(K)$	for applying the theorem to *open* curves, a (standardized) closing of curves is required.

CHALLENGE HINTS AND SOLUTIONS

Challenge 2, page 29: Take $\Delta f \; \Delta t \geqslant 1$ and substitute $\Delta l = c/\Delta f$ and $\Delta a = c/\Delta t$.

Challenge 16, page 44: Yes. But we can also argue its opposite, namely that matter appears when space is compressed too much. Both viewpoints are correct.

Challenge 22, page 47: The strictest upper limits are those with the smallest exponent for length, and the strictest lower limits are those with the largest exponent of length.

Challenge 24, page 49: To my knowledge, no such limits have been published. Do it yourself!

Challenge 25, page 49: The system limits cannot be chosen in other ways; after the limits have been corrected, the limits given here should still apply.

Challenge 28, page 50: Just insert numbers to check this.

Challenge 30, page 51: No.

Challenge 32, page 53: If you ever write such a table, publish it and send me a copy. I will include it in the text.

Challenge 35, page 66: Sloppily speaking, such a clock is not able to move its hands in a way that guarantees precise time reading.

Challenge 39, page 83: The final energy E produced by a proton accelerator increases with its radius R roughly as $E \sim R^{1.2}$; as an example, CERN's LHC achieved about 13 TeV for a radius of 4.3 km. Thus we would get a radius of more than 100 light years for a Planck energy accelerator. Building an accelerator achieving Planck energy is impossible.

Nature has no accelerator of this power, but gets near it. The maximum measured value of cosmic rays, 10^{22} eV, is about one millionth of the Planck energy. The mechanism of acceleration is still obscure. Neither black holes nor the cosmic horizon seem to be sources, for some yet unclear reasons. This issue is still a topic of research.

Challenge 40, page 84: The Planck energy is $E_{\mathrm{Pl}} = \sqrt{\hbar c^5/G} = 2.0\,\mathrm{GJ}$. Car fuel delivers about 43 MJ/kg. Thus the Planck energy corresponds to the energy of 47 kg of car fuel, about a tankful.

Challenge 41, page 84: Not really, as the mass error is equal to the mass only in the Planck case.

Challenge 42, page 84: It is improbable that such deviations can be found, as they are masked by
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the appearance of quantum gravity effects. However, if you do think that you have a prediction for a deviation, publish it, and send the author an email.

Challenge 43, page 84: The minimum measurable distance is the same for single particles and systems of particles.

Challenge 44, page 85: There is no gravitation at those energies and there are no particles. There is thus no paradox.

Challenge 45, page 85: The issue is still being debated; a good candidate for a minimum momentum of a single particle is given by $\hbar/R$, where R is the radius of the universe. Is this answer satisfying?

Challenge 46, page 86: All mentioned options could be valid at the same time. The issue is not closed and clear thinking about it is not easy.

Challenge 47, page 86: The precise energy scale is not clear. The scale is either the Planck energy or within a few orders of magnitude from it; the lowest possible energy is thus around a thousandth of the Planck energy.

Challenge 49, page 88: If you can think of an experiment, publish the proposal, and send the author an email.

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Challenge 50, page 91: The table of aggregates shows this clearly.

Challenge 51, page 92: The cosmic background radiation is a clock in the widest sense of the term.

Challenge 52, page 93: The way to deduce cosmological limits is presented in detail in the section starting on page 46.

Challenge 63, page 101: Also measurement errors at Planck scales prevent the determination of topology at those scales.

Challenge 65, page 103: The measurement error is as large as the measurement result.

Challenge 69, page 105: You will not find one.

Challenge 71, page 106: If you find one, publish it, and send the author an email.

Challenge 73, page 107: For the description of nature this is a contradiction. Nevertheless, the term 'universe', 'set of all sets' and other mathematical terms, as well as many religious concepts are of this type.

Challenge 74, page 109: No, for the reasons mentioned earlier on: fundamental measurement errors for horizon measurements, as well as many other effects, prevent this. The speculation is another example of misguided fantasy about extremal identity.

Challenge 75, page 109: The physical concepts most related to 'monad' are 'strand' and 'universe', as shown in the second half of this text.

Challenge 76, page 109: The macroscopic content of the universe may be observer-dependent. But to speak about many universes (Many 'everythings'?) or a 'multiverse' (What is more than everything? Why only one multiverse?) is pure nonsense.

Vol. II, page 258

Challenge 79, page 109: True only if it were possible to do this. Because particles and space are indistinguishable, removing particles means to remove everything. (The strand model visualizes this connection most clearly.)

Vol. III, page 324

Challenge 81, page 109: True. Existence is the ability to interact. If the ability disappears, existence disappears. In other words, 'existence' is a low-energy concept.

Challenge 82, page 111: If you find a sensible statement about the universe, publish it! And send it to the author as well. The next challenge shows one reason why this issue is interesting. In addition, such a statement would contradict the conclusions on the combined effects of general relativity and quantum theory.

Challenge 83, page 111: Plotinus in the *Enneads* has defined 'god' in exactly this way. Later, Augustine in *De Trinitate* and in several other texts, and many subsequent theologians have taken up this view. (See also Thomas Aquinas, *Summa contra gentiles*, 1, 30.) The idea they propose is simple: it is possible to clearly say what 'god' is *not*, but it is impossible to say what 'god' *is*. This statement is also part of the official *Roman Catholic Catechism*: see part one, section one, chapter one, IV, 43, found at www.vatican.va/archive/ENG0015/__PC.HTM. Similar statements are found in Judaism, Hinduism and Buddhism.

In other terms, theologians admit that 'god' cannot be defined, that the term has no properties or content, and that therefore the term cannot be used in any positive sentence. The aspects common to ‘universe’ and to ‘god’ suggest the conclusion that both are the same. Indeed, the analogy between the two concepts can be expanded to a proof: both concepts have the same content, the same boundary, and the same domain of application. (This is an intriguing and fascinating exercise.) In fact, this might be the most interesting of all proofs of the existence of ‘god’, as it lacks all the problems that the more common ‘proofs’ have. Despite its interest, this proof of equivalence is not found in any book on the topic yet. The reason is twofold. First, the results of modern physics – showing that the concept of universe has all these strange properties – are not common knowledge yet. Secondly, the result of the proof, the identity of ‘god’ and ‘universe’ – also called *pantheism* – is a heresy for most religions. It is an irony that the catholic catechism, together with modern physics, can be used to show that pantheism is correct, because any catholic who defends pantheism (or other heresies following from modern physics) incurs automatic excommunication, latae sententiae, without any need for a formal procedure.

Challenge 227 e

If one is ready to explore the identity of universe and ‘god’, one finds that a statement like ‘god created the universe’ translates as ‘the universe implies the universe’. The original statement is thus not a lie any more, but is promoted to a tautology. Similar changes appear for many other – but not all – statements using the term ‘god’. (The problems with the expression ‘in the beginning’ remain, though.) In fact, one can argue that statements about ‘god’ are only sensible and true if they remain sensible and true after the term has been exchanged with ‘universe’. Enjoy the exploration of such statements.

Challenge 228 e

Challenge 85, page 114: If you find one, publish it and send it also to me. The conjecture is that no such effects exist.

Challenge 87, page 114: In fact, no length below the Planck length itself plays any role in nature.

Challenge 89, page 115: You need quantum humour, because the result obviously contradicts a previous one given on page 94 that includes general relativity.

Challenge 92, page 124: The number of spatial dimensions must be given first, in order to talk about spheres.

Challenge 93, page 128: This is a challenge to you to find out. It is fun, it may yield a result in contradiction with the arguments given so far (publish it in this case), or it may yield an independent check of the results of the section.

Challenge 95, page 132: This issue is open and still a subject of research. The conjecture of the author is that the answer is negative. If you find an alternative, publish it, and send the author an email.

Challenge 97, page 137: The lid of a box must obey the indeterminacy relation. It cannot be at perfect rest with respect to the rest of the box.

Challenge 99, page 138: No, because the cosmic background is not a Planck scale effect, but an effect of much lower energy.

Challenge 100, page 138: Yes, at Planck scales all interactions are strand deformations; therefore collisions and gravity are indistinguishable there.

Challenge 101, page 138: No. Time is continuous only if *either* quantum theory and point particles *or* general relativity and point masses are assumed. The argument shows that only the combination of *both* theories with continuity is impossible.

Challenge 102, page 138: You should, because at Planck scales nature’s inherent measurement errors cannot clearly distinguish between different measurement results.

Challenge 103, page 138: We still have the chance to find the best approximate concepts possible. There is no reason to give up.

Challenge 104, page 138: Here are a few thoughts. A beginning of the big bang does not exist; something similar is given by that piece of continuous entity which is encountered when going backwards in time as much as possible. This has several implications.

— Going backwards in time as far as possible – towards the 'beginning' of time – is the same as zooming to smallest distances: we find a single strand of the amoeba.
— In other words, we speculate that the whole world is one single piece, fluctuating, and possibly tangled, knotted or branched.
— Going far away into space – to the border of the universe – is like taking a snapshot with a short shutter time: strands everywhere.
— Whenever we sloppily say that extended entities are 'infinite' in size, we only mean that they reach the horizon of the universe.

In summary, no starting point of the big bang exists, because time does not exist there. For the same reason, no initial conditions for particles or space-time exist. In addition, this shows that the big bang involved no creation, because without time and without possibility of choice, the term 'creation' makes no sense.

Challenge 105, page 138: The equivalence follows from the fact that all these processes require Planck energy, Planck measurement precision, Planck curvature, and Planck shutter time.

Page 369 **Challenge 106**, page 138: No, as explained later on in the text.

Challenge 107, page 139: Probably there is nothing wrong with the argument. For example, in the strand model, all observables are composed of fundamental events, and so, in some way, all observables are fundamentally indistinguishable.

Challenge 108, page 139: If not, force yourself. Brainstorming is important in life, as is the subsequent step: the checking of the speculations.

Challenge 113, page 151: The author would like to receive a mail on your reasons for disagreement.

Challenge 114, page 153: Let the author know if you succeed. And publish the results.

Challenge 115, page 153: Energy is action per time. Now, the Planck constant is the unit of action, and is defined by a crossing switch. A system that continuously produces a crossing switch for every Planck time running by thus has Planck energy. An example would be a tangle that is rotating extremely rapidly, once per Planck time, producing a crossing switch for every turn.

Momentum is action per length. A system that continuously produces a crossing switch whenever it advances by a Planck length has Planck momentum. An example would be a tangle configuration that lets a switch hop from one strand to the next under tight strand packing.

Force is action per length and time. A system that continuously produces a crossing switch for every Planck time that passes by and for every Planck length it advances through exerts a Planck force. A tangle with the structure of a screw that rotates and advances with sufficient speed would be an example.

Challenge 119, page 163: Yes; the appearance of a crossing does not depend on distance or on the number of strands in between.

Challenge 120, page 163: No; more than three dimensions do not allow us to define a crossing switch.

Challenge 121, page 163: If so, let the author know. If the generalization is genuine, the strand model is not correct.

Challenge 133, page 190: The magnitude at a point should be related to the vectorial sum of all inverse shortest crossing distances at that point.

Challenge 139, page 201: This algebraic transformation is shown in all textbooks that treat the Pauli equation. It can also be checked by writing the two equations out component by component.

Challenge 143, page 222: Yes, as can easily be checked by rereading the definitions with the spinor tangle description in mind.

Challenge 146, page 222: No contradiction is known.

Challenge 147, page 222: In the relativistic case, local space curvature is also taken into account.

Challenge 149, page 222: Find out, publish the result, and let the author know.

Challenge 150, page 223: If the strand interpenetration is allowed *generally*, quantum theory is impossible to derive, as the spinor behaviour would not be possible. If strand interpenetration were allowed only *under certain conditions* (such as only for a strand with itself, but not among two different strands), quantum theory might still possible. A similar process lies at the basis of mass generation, as shown in the section on the weak interaction.

Challenge 152, page 223: The belt trick would imply that a wheel rolls over its own blood supply at every second rotation.

Challenge 161, page 245: The author bets that you cannot find a deviation of the strand model from QED. If you find one, publish it!

Challenge 172, page 274: No slide is possible, thus no crossing change appears; thus the situation has no observable effects. If we deform one slide before the slide – which is possible – we get back the situation already discussed above.

Challenge 177, page 275: For the Wightman axioms, this seems to be the case; however, a formal proof is still missing. The same is expected for the Haag–Kastler axioms.

Challenge 187, page 303: A black hole has at least one crossing, thus at least a Planck mass.

Challenge 193, page 324: These tangles are not rational. In the renewed strand model of 2015, they cannot form; they are not allowed and do not represent any particle.

Challenge 195, page 339: Such a tangle is composed of several gravitons.

Challenge 196, page 346: Tail braiding leads to tangledness, which in turn is the basis for core rotation. And core rotation is kinetic energy, not rest mass.

Challenge 198, page 346: The issue is topic of research; for symmetry reasons it seems that a state in which each of the six quarks has the same bound to the other five quarks cannot exist.

Challenge 205, page 372: If you find such an estimate, publish it and send it to the author. A really good estimate also answers the following question: why does particle mass increase with core complexity? A tangle with a complex core, i.e., with a core of large ropelength, has a large mass value. Any correct estimate of the mass must yield this property. But a more complex knot will have a smaller probability for the belt trick. We seem to be forced to conclude that particle mass is not due to the belt trick alone.

Challenge 206, page 373: If you find such an estimate, publish it and send it to the author.

Challenge 208, page 373: Probably not.

Challenge 209, page 373: Probably not.

Challenge 210, page 373: Probably not.

Challenge 211, page 374: Find out – and let the author know.

Challenge 213, page 378: This would be an interesting result worth a publication.

Challenge 216, page 386: If you plan such a calculation, the author would be delighted to help.

Challenge 220, page 395: Take up the challenge!

Challenge 222, page 423: There is a good chance, however, that such alternatives can be eliminated rather quickly. If you cannot do so, do publish the argument, and let the author know.

Challenge 224, page 431: Nobody can really answer 'why'-questions about human actions. Climbing, like every other passion, is also a symbolic activity. Climbing can be a search for adventure, for meaning, for our mother or father, for ourselves, for happiness, or for peace.

Challenge 225, page 432: Also in dreams, speeds can be compared; and also in dreams, a kind of causality holds (though not a trivial one). Thus there is an invariant and therefore a maximum speed.

Challenge 226, page 433: Probably none. The answer depends on whether the existence of strands can be deduced from dreams. If strands can be deduced from dreams, all of physics follows. The conjecture is that this deduction is possible. If you find an argument against or in favour of this conjecture, let the author know.

BIBLIOGRAPHY

> The only end of writing is to enable the readers better to enjoy life, or better to endure it.
>
> Samuel Johnson*

1 See the first volume of the Motion Mountain series, *Fall, Flow and Heat*, available as free download at www.motionmountain.net. Cited on pages 17 and 418.

2 See the second volume of the Motion Mountain series, *Relativity*, available as free download at www.motionmountain.net. Cited on pages 17, 18, 418, and 448.

3 See the third volume of the Motion Mountain series, *Light, Charges and Brains*, available as free download at www.motionmountain.net, as well as the mentioned fourth and fifth volumes. Cited on pages 17 and 418.

4 See the fourth and fifth volumes of the Motion Mountain series, *The Quantum of Change* and *Pleasure, Technology and the Stars*, available as free download at www.motionmountain.net. Cited on pages 18, 418, 419, and 447.

5 The most precise value of the fine structure constant is determined from a weighted world average of high-precision measurements by a special international scientific committee called CODATA. Its website is www.codata.org/committees-and-groups/fundamental-physical-constants. The site also provides the latest official publication with the values of the fundamental constants. The most recent value of the fine structure constant is published at physics.nist.gov/cgi-bin/cuu/Value?alphinv and physics.nist.gov/cgi-bin/cuu/Value?alph. Cited on pages 18, 229, 382, 391, and 433.

6 See for example, the book by Robert Laughlin, *A Different Universe: Reinventing Physics from the Botton Down* Basic Books, 2005. Of the numerous books that discuss the idea of a final theory, this is the only one worth reading, and the only one cited in this bibliography. The opinions of Laughlin are worth pondering. Cited on page 21.

7 Many physicists, including Steven Weinberg, regularly – and incorrectly – claim in interviews that the measurement problem is not solved yet. Cited on page 21.

8 Undocumented sentences to this effect are regularly attributed to Albert Einstein. Because Einstein was a pantheist, as he often explained, his statements on the 'mind of god' are not really to be taken seriously. They were all made – if at all – in a humorous tone. Cited on page 21.

9 For an example for the inappropriate fear of unification, see the theatre play *Die Physiker* by the Swiss author Friedrich Dürrenmatt. Several other plays and novels took over this type of disinformation. Cited on page 22.

* This is a statement from the brilliant essay by the influential writer Samuel Johnson, *Review of Soame Jenyns' "A Free Enquiry Into the Nature and Origin of Evil"*, 1757. See www.samueljohnson.com.

10 Exploring the spirit of play is the subject of research of the famous National Institute for Play, founded by Stuart Brown, and found at www.nifplay.org. Cited on page 22.

11 See e.g. the 1922 lectures by Lorentz at Caltech, published as H. A. Lorentz, *Problems of Modern Physics*, edited by H. Bateman, Ginn and Company, 1927, page 99. Cited on page 28.

12 Bohr explained the indivisibilty of the quantum of action in his famous Como lecture, printed in N. Bohr, *Atomtheorie und Naturbeschreibung*, Springer, 1931. It was translated into English language as N. Bohr, *Atomic Theory and the Description of Nature*, Cambridge University Press, 1934. More statements about the indivisibility of the quantum of action can be found in N. Bohr, *Atomic Physics and Human Knowledge*, Science Editions, New York, 1961. For summaries of Bohr's ideas by others see Max Jammer, *The Philosophy of Quantum Mechanics*, Wiley, first edition, 1974, pp. 90–91, and John Honner, *The Description of Nature – Niels Bohr and the Philosophy of Quantum Physics*, Clarendon Press, 1987, p. 104. Cited on page 29.

13 For an overview of the quantum of action as a basis of quantum theory, see the first chapter of the fourth volume of the Motion Mountain series, Ref. 4. Cited on page 30.

Vol. IV, page 15

Vol. IV, page 182

14 An overview of EBK quantization can be found in the volume on quantum theory. Cited on page 30.

15 Minimal entropy is discussed by L. Szilard, *Über die Entropieverminderung in einem thermodynamischen System bei Eingriffen intelligenter Wesen*, Zeitschrift für Physik **53**, pp. 840–856, 1929. This classic paper can also be found in English translation in his collected works. Cited on page 31.

16 See for example A. E. Shalyt-Margolin & A. Ya. Tregubovich, *Generalized uncertainty relation in thermodynamics*, preprint at arxiv.org/abs/gr-qc/0307018, or J. Uffink & J. van Lith-van Dis, *Thermodynamic uncertainty relations*, Foundations of Physics **29**, pp. 655–692, 1999. Cited on page 31.

17 See also the fundamental paper by A. DiSessa, *Momentum flow as an alternative perspective in elementary mechanics*, **48**, p. 365, 1980, and A. DiSessa, *Erratum: "Momentum flow as an alternative perspective in elementary mechanics" [Am. J. Phys. 48, 365 (1980)]*, **48**, p. 784, 1980. Cited on page 32.

18 The observations of black holes at the centre of galaxies and elsewhere are summarised by R. Blandford & N. Gehrels, *Revisiting the black hole*, Physics Today **52**, June 1999. Their existence is now well established. Cited on page 32.

19 It seems that the first published statement of the maximum force as a *fundamental principle* was around the year 2000, in this text, in the chapter on gravitation and relativity. The author discovered the maximum force principle, not knowing the work of others, when searching for a way to derive the results of the last part of this adventure that would be so simple that it would convince even a secondary-school student. In the year 2000, the author told his friends in Berlin about his didactic approach for general relativity.

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The *concept* of a maximum force was first proposed, most probably, by H.-J. Treder in 1985, followed by by Venzo de Sabbata and C. Sivaram in 1993. Also this physics discovery was thus made much too late. In 1995, Corrado Massa took up the idea. Independently, Ludwik Kostro in 1999, Christoph Schiller just before 2000 and Gary Gibbons in the years before 2002 arrived at the same concept. Gary Gibbons was inspired by a book by Oliver Lodge; he explains that the maximum force value follows from general relativity; he does not make a statement about the converse, nor do the other authors. The statement of maximum force as a *fundamental principle* seems original to Christoph Schiller.

The temporal order of the first papers on maximum force seems to be H.-J. Treder, *The Planckions as Largest Elementary Particles and as Smallest Test Bodies*, Foundations of Physics **15**, pp. 161–166, 1985, followed by V. de Sabbata & C. Sivaram, *On limiting field strengths in gravitation*, Foundations of Physics Letters **6**, pp. 561–570, 1993, then by C. Massa, *Does the gravitational constant increase?*, Astrophysics and Space Science **232**, pp. 143–148, 1995, and by L. Kostro & B. Lange, *Is c^4/G the greatest possible force in nature?*, Physics Essays **12**, pp. 182–189, 1999. The next references are the paper by G. W. Gibbons, *The maximum tension principle in general relativity*, Foundations of Physics **32**, pp. 1891–1901, 2002, preprint at arxiv.org/abs/hep-th/0210109 – though he developed the ideas before that date – and the older versions of the present text, i.e., Christoph Schiller, *Motion Mountain – The Adventure of Physics*, a free pdf available at www.motionmountain.net. Then came C. Schiller, *Maximum force and minimum distance: physics in limit statements*, preprint at arxiv.org/abs/physics/0309118, and C. Schiller, *General relativity and cosmology derived from principle of maximum power or force*, International Journal of Theoretical Physics **44**, pp. 1629–1647, 2005, preprint at arxiv.org/abs/physics/0607090. See also R. Beig, G. W. Gibbons & R. M. Schoen, *Gravitating opposites attract*, Classical and Quantum Gravity **26**, p. 225013, 2009. preprint at arxiv.org/abs/0907.1193.

A detailed discussion of maximum force and power is given in the volume on general relativity, Ref. 2. Cited on pages 33, 43, 296, and 458.

Vol. II, page 107

20 Maximal luminosity is often mentioned in connection with gravitational wave detection; nevertheless, the general power maximum has never been mentioned before. See for example L. Ju, D. G. Blair & C. Zhao, *Detection of gravitational waves*, Reports on Progress in Physics **63**, pp. 1317–1427, 2000. See also C. W. Misner, K. S. Thorne & J. A. Wheeler, *Gravitation*, Freeman, 1973, page 980. Cited on page 33.

21 See for example Wolfgang Rindler, *Relativity – Special, General and Cosmological*, Oxford University Press, 2001, p. 70 ss, or Ray d'Inverno, *Introducing Einstein's Relativity*, Clarendon Press, 1992, p. 36 ss. Cited on page 34.

22 T. Jacobson, *Thermodynamics of spacetime: the Einstein equation of state*, Physical Review Letters **75**, pp. 1260–1263, 1995, preprint at arxiv.org/abs/gr-qc/9504004; this deep article remains fascinating to this day. Even the author was scared to draw all the possible conclusions. The general concepts are explained, almost without formulae, in L. Smolin, *On the nature of quantum fluctuations and their relation to gravitation and the principle of inertia*, Classical and Quantum Gravity **3**, pp. 347–359, 1986. Cited on pages 34 and 295.

23 This relation was pointed out by Achim Kempf. The story is told in A. D. Sakharov, General Relativity and Gravitation **32**, pp. 365–367, 2000, a reprint of his paper Doklady Akademii Nauk SSSR **177**, pp. 70–71, 1967. Cited on pages 35 and 44.

24 Indeterminacy relations in general relativity are discussed in C. A. Mead, *Possible connection between gravitation and fundamental length*, Physical Review B **135**, pp. 849–862, 1964. The generalized indeterminacy relation is implicit on page 852, but the issue is explained rather unclearly. Probably the author considered the result too simple to be mentioned explicitly. (That paper took 5 years to get published; comments on the story, written 37 years later, are found at C. A. Mead, *Walking the Planck length through history*, Physics Today **54**, p. 15 and p. 81, 2001, with a reply by Frank Wilczek.) See also P. K. Townsend, *Small-scale structure of space-time as the origin of the gravitational constant*, Physical Review D **15**, pp. 2795–2801, 1977, or the paper by M.-T. Jaekel & S. Renaud, *Gravitational quantum limit for length measurement*, Physics Letters A **185**, pp. 143–148, 1994. Cited on pages 36, 67, 68, 69, 72, and 120.

25 M. Kramer & al., *Tests of general relativity from timing the double pulsar*, preprint at arxiv.org/abs/astro-ph/060941. Cited on page 36.

26 Minimal length and minimal time intervals are discussed, for example, by G. Amelino-Camelia, *Limits on the measurability of space-time distances in (the semiclassical approximation of) quantum gravity*, Modern Physics Letters A **9**, pp. 3415–3422, 1994, preprint at arxiv.org/abs/gr-qc/9603014, and by Y. J. Ng & H. van Dam, *Limit to space-time measurement*, Modern Physics Letters A **9**, pp. 335–340, 1994. Many other authors have explored the topic. Cited on pages 38 and 68.

27 Maximal curvature, as well as area and volume quantization, are discussed in A. Ashtekar, *Quantum geometry and gravity: recent advances*, preprint at arxiv.org/abs/gr-qc/0112038 and in A. Ashtekar, *Quantum geometry in action: big bang and black holes*, preprint at arxiv.org/abs/math-ph/0202008. Cited on pages 38, 76, and 457.

28 Maximons, elementary particles of Planck mass, are discussed by A. D. Sakharov, *Vacuum quantum fluctuations in curved space and the theory of gravitation*, Soviet Physics – Doklady **12**, pp. 1040–1041, 1968. Cited on pages 40, 79, and 123.

29 Wolfgang Rindler, *Relativity – Special, General and Cosmological*, Oxford University Press, 2001, p. 230. Cited on page 42.

30 Several incorrect counterclaims to the entropy limit were made in R. Bousso, *The holographic principle*, Review of Modern Physics **74**, pp. 825–874, 2002, preprint at arxiv.org/abs/hep-th/0203101. However, this otherwise good review has some errors in its arguments, as explained on page 147 in volume V. Bousso has changed his position in the meantime; he now accepts the entropy limit. Cited on pages 44, 48, 455, and 457.

31 Gamma-ray bursts are discussed by G. Preparata, R. Ruffini & S.-S. Xue, *The dyadosphere of black holes and gamma-ray bursts*, Astronomy and Astrophysics **338**, pp. L87–L90, 1998, and C. L. Bianco, R. Ruffini & S.-S. Xue, *The elementary spike produced by a pure e^+e^- pair-electromagnetic pulse from a black hole: the PEM pulse*, Astronomy and Astrophysics **368**, pp. 377–390, 2001. Cited on page 45.

32 See for example the review in C.W.J. Beenakker & al., *Quantum transport in semiconductor nanostructures*, pp. 1–228, in H. Ehrenreich & D. Turnbull editors, *Solid State Physics*, volume 44, Academic Press, 1991. Cited on page 45.

33 A discussion of a different electrical indeterminacy relation, between current and charge, can be found in Y-Q. Li & B. Chen, *Quantum theory for mesoscopic electronic circuits and its applications*, preprint at arxiv.org/abs/cond-mat/9907171. Cited on page 45.

34 Hans C. Ohanian & Remo Ruffini, *Gravitation and Spacetime*, W.W. Norton & Co., 1994. Cited on pages 46 and 462.

35 The entropy limit for black holes is discussed by J. D. Bekenstein, *Entropy bounds and black hole remnants*, Physical Review D **49**, pp. 1912–1921, 1994. See also J. D. Bekenstein, *Universal upper bound on the entropy-to-energy ratio for bounded systems*, Physical Review D **23**, pp. 287–298, 1981. Cited on pages 48 and 133.

36 The statement is also called the *Kovtun-Son-Starinets conjecture*. It was published as P. Kovtun, D. T. Son & A. O. Starinets, *A viscosity bound conjecture*, preprint at arxiv.org/abs/hep-th/0405231. See also P. Kovtun, D. T. Son & A. O. Starinets, *Viscosity in strongly interacting quantum field theories from black hole physics*, Physical Review Letters **44**, p. 111601, 2005. For an experimental verification, see U. Hohm, *On the ratio of the shear viscosity to the density of entropy of the rare gases and H2, N2, CH4, and CF4*, Chemical Physics **444**, pp. 39–42, 2014. Cited on page 49.

37 Brian Greene, *The Elegant Universe – Superstrings, Hidden Dimensions, and the Quest for the Ultimate Theory*, Jonathan Cape 1999. Cited on page 53.

38 S. Weinberg, *The cosmological constant problem*, Reviews of Modern Physics **61**, pp. 1–23, 1989. Cited on page 58.

39 Steven Weinberg, *The Quantum Theory of Fields*, Cambridge University Press, volumes I, 1995, and II, 1996. Cited on page 58.

40 See the excellent presentation on the cosmological constant in general relativity by E. Bianchi & C. Rovelli, *Why all these prejudices against a constant?*, preprint at arxiv.org/abs/1002.3966 Cited on page 58.

41 The difficulties are summarised by B. S. DeWitt, *Quantum field theory in curved space-time*, Physics Reports **19**, pp. 295–357, 1975. Cited on page 58.

42 C. W. Misner, K. S. Thorne & J. A. Wheeler, *Gravitation*, Freeman, 1973. Cited on pages 59, 60, and 68.

43 J. A. Wheeler, in *Relativity, Groups and Topology*, edited by C. DeWitt & B. S. DeWitt, Gordon and Breach, 1994. See also J. A. Wheeler, *Physics at the Planck length*, International Journal of Modern Physics A **8**, pp. 4013–4018, 1993. However, his claim that spin 1/2 *requires* topology change is *refuted* by the strand model of the vacuum. Cited on page 59.

44 J. L. Friedman & R. D. Sorkin, *Spin 1/2 from gravity*, Physical Review Letters **44**, pp. 1100–1103, 1980. Cited on page 59.

45 A. P. Balachandran, G. Bimonte, G. Marmo & A. Simoni, *Topology change and quantum physics*, Nuclear Physics B **446**, pp. 299–314, 1995, preprint at arxiv.org/abs/hep-th/9503046. Cited on page 59.

46 J. Ehlers, *Introduction – Survey of Problems*, pp. 1–10, in J. Ehlers, editor, *Sistemi gravitazionali isolati in relatività generale*, Rendiconti della scuola internazionale di fisica "Enrico Fermi", LXVII° corso, Società Italiana di Fisica/North Holland, 1979. Cited on page 59.

47 See C. Schiller, *Le vide diffère-t-il de la matière?* in E. Gunzig & S. Diner editors, *Le Vide – Univers du tout et du rien – Des physiciens et des philosophes s'interrogent*, Les Éditions de l'Université de Bruxelles, 1998. An older, English-language version is available as C. Schiller, *Does matter differ from vacuum?* preprint at arxiv.org/abs/gr-qc/9610066. Cited on pages 59, 120, 122, 123, 124, 135, and 136.

48 See for example Richard P. Feynman, Robert B. Leighton & Matthew Sands, *The Feynman Lectures on Physics*, Addison Wesley, 1977. Cited on page 60.

49 Steven Weinberg, *Gravitation and Cosmology*, Wiley, 1972. Cited on pages 60, 66, and 68.

50 The argument is given e.g. in E. P. Wigner, *Relativistic invariance and quantum phenomena*, Reviews of Modern Physics **29**, pp. 255–258, 1957. Cited on page 65.

51 The starting point for the following arguments is taken from M. Schön, *Operative time definition and principal indeterminacy*, preprint at arxiv.org/abs/gr-qc/9304024, and from T. Padmanabhan, *Limitations on the operational definition of space-time events and quantum gravity*, Classical and Quantum Gravity **4**, pp. L107–L113, 1987; see also Padmanabhan's earlier papers referenced there. Cited on page 65.

52 W. Heisenberg, *Über den anschaulichen Inhalt der quantentheoretischen Kinematik und Mechanik*, Zeitschrift für Physik **43**, pp. 172–198, 1927. Cited on page 65.

53 E. H. Kennard, *Zur Quantenmechanik einfacher Bewegungstypen*, Zeitschrift für Physik **44**, pp. 326–352, 1927. Cited on page 65.

54 M. G. Raymer, *Uncertainty principle for joint measurement of noncommuting variables*, American Journal of Physics **62**, pp. 986–993, 1994. Cited on page 65.

55 H. Salecker, & E. P. Wigner, *Quantum limitations of the measurement of space-time distances*, Physical Review **109**, pp. 571–577, 1958. Cited on pages 66, 94, and 115.

56 E. J. Zimmerman, *The macroscopic nature of space-time*, American Journal of Physics **30**, pp. 97–105, 1962. Cited on pages 66, 94, and 115.

57 J. D. Bekenstein, *Black holes and entropy*, Physical Review D 7, pp. 2333–2346, 1973. Cited on pages 66, 133, and 292.

58 S. W. Hawking, *Particle creation by black holes*, Communications in Mathematical Physics **43**, pp. 199–220, 1975; see also S. W. Hawking, *Black hole thermodynamics*, Physical Review D **13**, pp. 191–197, 1976. Cited on pages 66, 133, and 292.

59 P. Gibbs, *The small scale structure of space-time: a bibliographical review*, preprint at arxiv.org/abs/hep-th/9506171. Cited on pages 66 and 85.

60 The impossibility of determining temporal ordering in quantum theory is discussed by J. Oppenheimer, B. Reznik & W. G. Unruh, *Temporal ordering in quantum mechanics*, Journal of Physics A **35**, pp. 7641–7652, 2001, preprint at arxiv.org/abs/quant-ph/0003130. Cited on page 67.

61 M. -T. Jaekel & S. Renaud, *Gravitational quantum limit for length measurement*, Physics Letters A **185**, pp. 143–148, 1994. Cited on page 68.

62 D. V. Ahluwalia, *Quantum measurement, gravitation and locality*, Physics Letters B **339**, pp. 301–303, 1994, preprint at arxiv.org/abs/gr-qc/9308007. Cited on page 68.

63 L. Garay, *Quantum gravity and minimum length*, International Journal of Modern Physics A **10**, pp. 145–165, 1995, preprint at arxiv.org/abs/gr-qc/9403008. This paper also includes an extensive bibliography. See also R. J. Adler & D. I. Santiago, *On gravity and the uncertainty principle*, Modern Physics Letters A **14**, pp. 1371–1381, 1999, preprint at arxiv.org/abs/gr-qc/9904026. Cited on page 68.

64 C. Rovelli & L. Smolin, *Discreteness of area and volume in quantum gravity*, Nuclear Physics B **442**, pp. 593–619, 1995. R. Loll, *The volume operator in discretized quantum gravity*, preprint at arxiv.org/abs/gr-qc/9506014. See also C. Rovelli, *Notes for a brief history of quantum gravity*, preprint at arxiv.org/abs/gr-qc/0006061. Cited on page 69.

65 D. Amati, M. Ciafaloni & G. Veneziano, *Superstring collisions at Planckian energies*, Physics Letters B **197**, pp. 81–88, 1987. D. J. Gross & P. F. Mende, *The high energy behavior of string scattering amplitudes*, Physics Letters B **197**, pp. 129–134, 1987. K. Konishi, G. Paffuti & P. Provero, *Minimum physical length and the generalized uncertainty principle*, Physics Letters B **234**, pp. 276–284, 1990. P. Aspinwall, *Minimum distances in non-trivial string target spaces*, Nuclear Physics B **431**, pp. 78–96, 1994, preprint at arxiv.org/abs/hep-th/9404060. Cited on page 69.

66 M. Maggiore, *A generalised uncertainty principle in quantum mechanics*, Physics Letters B **304**, pp. 65–69, 1993. Cited on page 69.

67 A simple approach is S. Doplicher, K. Fredenhagen & J. E. Roberts, *Space-time quantization induced by classical gravity*, Physics Letters B **331**, pp. 39–44, 1994. Cited on pages 69 and 84.

68 A. Kempf, *Uncertainty relation in quantum mechanics with quantum group symmetry*, Journal of Mathematical Physics **35**, pp. 4483–4496, 1994. A. Kempf, *Quantum groups*

and quantum field theory with nonzero minimal uncertainties in positions and momenta, Czechoslovak Journal of Physics **44**, pp. 1041–1048, 1994. Cited on page 69.

69 E. J. Hellund & K. Tanaka, *Quantized space-time*, Physical Review **94**, pp. 192–195, 1954. Cited on page 70.

70 This intriguing extract from a letter by Einstein was made widely known by John J. Stachel, in his paper *The other Einstein: Einstein contra field theory*, that is best found in his book *Einstein from 'B' to 'Z'*, Birkhäuser, 2002. The German original of the letter is found in Robert Schulmann, A. J. Knox, Michel Janssen & József Illy, *The Collected Papers of Albert Einstein, Volume 8A – The Berlin Years: Correspondence, 1914–1917*, letter 299, Princeton University Press, 1998. Barbara Wolff helped in clarifying several details in the German original. The letter is now available online, at einsteinpapers.press.princeton.edu/vol8a-doc/463. Cited on page 70.

71 A. Peres & N. Rosen, *Quantum limitations on the measurement of gravitational fields*, Physical Review **118**, pp. 335–336, 1960. Cited on page 72.

72 It is the first definition in Euclid's *Elements*, *c.* 300 bce. For an English translation see T. Heath, *The Thirteen Books of the Elements*, Dover, 1969. Cited on page 73.

73 A beautiful description of the Banach–Tarski paradox is the one by Ian Stewart, *Paradox of the spheres*, New Scientist, 14 January 1995, pp. 28–31. Cited on page 73.

74 H. S. Snyder, *Quantized space-time*, Physical Review **71**, pp. 38–41, 1947. H. S. Snyder, *The electromagnetic field in quantized space-time*, Physical Review **72**, pp. 68–74, 1947. A. Schild, *Discrete space-time and integral Lorentz transformations*, Physical Review **73**, pp. 414–415, 1948. E. L. Hill, *Relativistic theory of discrete momentum space and discrete space-time*, Physical Review **100**, pp. 1780–1783, 1950. H. T. Flint, *The quantization of space-time*, Physical Review **74**, pp. 209–210, 1948. A. Das, *Cellular space-time and quantum field theory*, Il Nuovo Cimento **18**, pp. 482–504, 1960. Cited on page 75.

75 D. Finkelstein, *'Superconducting' causal nets*, International Journal of Theoretical Physics **27**, pp. 473–519, 1985. Cited on page 75.

76 N. H. Christ, R. Friedberg & T. D. Lee, *Random lattice field theory: general formulation*, Nuclear Physics B **202**, pp. 89–125, 1982. G. 't Hooft, *Quantum field theory for elementary particles – is quantum field theory a theory?*, Physics Reports **104**, pp. 129–142, 1984. Cited on page 75.

77 For a discussion, see R. Sorabji, *Time, Creation and the Continuum: Theories in Antiquity and the Early Middle Ages*, Duckworth, 1983. Cited on page 75.

78 See, for example, L. Bombelli, J. Lee, D. Meyer & R. D. Sorkin, *Space-time as a causal set*, Physical Review Letters **59**, pp. 521–524, 1987. G. Brightwell & R. Gregory, *Structure of random space-time*, Physical Review Letters **66**, pp. 260–263, 1991. Cited on page 75.

79 The false belief that particles like quarks or electrons are composite is slow to die out. See for example: S. Fredriksson, *Preon prophecies by the standard model*, preprint at arxiv.org/abs/hep-ph/0309213. Preon models gained popularity in the 1970s and 1980s, in particular through the papers by J. C. Pati & A. Salam, *Lepton number as the fourth "color"*, Physical Review D **10**, pp. 275–289, 1974, H. Harari, *A schematic model of quarks and leptons*, Physics Letters B **86**, pp. 83–86, 1979, M. A. Shupe, *A composite model of leptons and quarks*, Physics Letters B **86**, pp. 87–92, 1979, and H. Fritzsch & G. Mandelbaum, *Weak interactions as manifestations of the substructure of leptons and quarks*, Physics Letters B **102**, pp. 319–322, 1981. Cited on page 77.

80 N. F. Ramsey & A. Weis, *Suche nach permanenten elektrischen Dipolmomenten: ein Test der Zeitumkehrinvarianz*, Physikalische Blätter **52**, pp. 859–863, 1996. See also W. Bernreuther & M. Suzuki, *The electric dipole moment of the electron*, Reviews of Modern Physics **63**, pp. 313–340, 1991, and the musings in Hans Dehmelt, *Is the electron a composite particle?*, Hyperfine Interactions **81**, pp. 1–3, 1993. Cited on page 78.

81 K. Akama, T. Hattori & K. Katsuura, *Naturalness bounds on dipole moments from new physics*, preprint at arxiv.org/abs/hep-ph/0111238. Cited on page 78.

82 The paper by J. Baron & al., *Order of magnitude smaller limit on the electric dipole moment of the electron*, preprint at arxiv.org/abs/1310.7534 gives an upper experimental limit to the dipole moment of the electron of $8.7 \cdot 10^{-31}\,e$ m. Cited on page 78.

83 C. Wolf, *Upper limit for the mass of an elementary particle due to discrete time quantum mechanics*, Il Nuovo Cimento B **109**, pp. 213–218, 1994. Cited on page 80.

84 W. G. Unruh, *Notes on black hole evaporation*, Physical Review D **14**, pp. 870–875, 1976. W. G. Unruh & R. M. Wald, *What happens when an accelerating observer detects a Rindler particle*, Physical Review D **29**, pp. 1047–1056, 1984. Cited on page 82.

85 The first example was J. Magueijo & L. Smolin, *Lorentz invariance with an invariant energy scale*, Physical Review Letters **88**, p. 190403, 2002, preprint at arxiv.org/abs/hep-th/0112090. They propose a modification of the mass energy relation of the kind

$$E = \frac{c^2 \gamma m}{1 + \frac{c^2 \gamma m}{E_{\rm Pl}}} \quad \text{and} \quad p = \frac{\gamma m v}{1 + \frac{c^2 \gamma m}{E_{\rm Pl}}} \,. \qquad (215)$$

Another, similar approach of recent years, with a different proposal, is called 'doubly special relativity'. A recent summary is G. Amelino-Camelia, *Doubly-special relativity: first results and key open problems*, International Journal of Modern Physics **11**, pp. 1643–1669, 2002, preprint at arxiv.org/abs/gr-qc/0210063. The paper shows how conceptual problems hinder the advance of the field. Another such discussion R. Aloisio, A. Galante, A. F. Grillo, E. Luzio & F. Méndez, *Approaching space-time through velocity in doubly special relativity*, preprint at arxiv.org/abs/gr-qc/0410020. The lesson from these attempts is simple: special relativity *cannot* be modified to include a limit energy without also including general relativity and quantum theory. Cited on pages 84 and 282.

86 W. Jauch, *Heisenberg's uncertainty relation and thermal vibrations in crystals*, American Journal of Physics **61**, pp. 929–932, 1993. Cited on page 84.

87 H. D. Zeh, *On the interpretation of measurement in quantum theory*, Foundations of Physics **1**, pp. 69–76, 1970. Cited on page 85.

88 See Y. J. Ng, W. A. Christiansen & H. van Dam, *Probing Planck-scale physics with extragalactic sources?*, Astrophysical Journal **591**, pp. L87–L90, 2003, preprint at arxiv.org/abs/astro-ph/0302372; D. H. Coule, *Planck scale still safe from stellar images*, Classical and Quantum Gravity **20**, pp. 3107–3112, 2003, preprint at arxiv.org/abs/astro-ph/0302333. Negative experimental results (and not always correct calculations) are found in R. Lieu & L. Hillman, *The phase coherence of light from extragalactic sources – direct evidence against first order Planck scale fluctuations in time and space*, Astrophysical Journal **585**, pp. L77–L80, 2003, and R. Ragazzoni, M. Turatto & W. Gaessler, *The lack of observational evidence for the quantum structure of spacetime at Planck scales*, Astrophysical Journal **587**, pp. L1–L4, 2003. Cited on page 88.

89 B. E. Schaefer, *Severe limits on variations of the speed of light with frequency*, Physical Review Letters **82**, pp. 4964–4966, 21 June 1999. Cited on page 88.

90 A.A. Abdo & al., (Fermi GBM/LAT collaborations) *Testing Einstein's special relativity with Fermi's short hard gamma-ray burst GRB090510*, preprint at arxiv.org/0908.1832. Cited on page 88.

91 G. Amelino-Camelia, J. Ellis, N.E. Mavromatos, D.V. Nanopoulos & S. Sakar, *Potential sensitivity of gamma-ray-burster observations to wave dispersion in vacuo*, Nature **393**, pp. 763–765, 1998, preprint at arxiv.org/abs/astro-ph/9712103. Cited on page 88.

92 G. Amelino-Camelia, *Phenomenological description of space-time foam*, preprint at arxiv.org/abs/gr-qc/0104005. The paper includes a clearly written overview of present experimental approaches to detecting quantum gravity effects. See also his update G. Amelino-Camelia, *Quantum-gravity phenomenology: status and prospects*, preprint at arxiv.org/abs/gr-qc/0204051. Cited on pages 88 and 89.

93 G. Amelino-Camelia, *An interferometric gravitational wave detector as a quantum gravity apparatus*, Nature **398**, pp. 216–218, 1999, preprint at arxiv.org/abs/gr-qc/9808029. Cited on page 88.

94 F. Karolyhazy, *Gravitation and quantum mechanics of macroscopic objects*, Il Nuovo Cimento **A42**, pp. 390–402, 1966. Y. J. Ng & H. van Dam, *Limit to space-time measurement*, Modern Physics Letters A **9**, pp. 335–340, 1994. Y. J. Ng & H. van Dam, Modern Physics Letters A *Remarks on gravitational sources*, **10**, pp. 2801–2808, 1995. The discussion is neatly summarised in Y. J. Ng & H. van Dam, *Comment on 'Uncertainty in measurements of distance'*, preprint at arxiv.org/abs/gr-qc/0209021. See also Y. J. Ng, *Spacetime foam*, preprint at arxiv.org/abs/gr-qc/0201022. Cited on pages 88 and 94.

95 L. J. Garay, *Spacetime foam as a quantum thermal bath*, Physics Review Letters **80**, pp. 2508–2511, 1998, preprint at arxiv.org/abs/gr-qc/9801024. Cited on page 89.

96 G. Amelino-Camelia & T. Piran, *Planck-scale deformation of Lorentz symmetry as a solution to the UHECR and the TeV-γ paradoxes*, preprint at arxiv.org/astro-ph/0008107, 2000. Cited on page 89.

97 R. P. Woodard, *How far are we from the quantum theory of gravity?*, preprint at arxiv.org/abs/0907.4238. For a different point of view, see L. Smolin, *Generic predictions of quantum theories of gravity*, preprint at arxiv.org/abs/hep-th/0605052. Cited on pages 89 and 306.

98 A similar point of view, often called monism, was proposed by Baruch Spinoza, *Ethics Demonstrated in Geometrical Order*, 1677, originally in Latin; an affordable French edition is Baruch Spinoza, *L'Ethique*, Folio-Gallimard, 1954. For a discussion of his ideas, especially his monism, see Don Garret editor, *The Cambridge Companion to Spinoza*, Cambridge University Press, 1996, or any general text on the history of philosophy. Cited on page 89.

99 See the lucid discussion by G. F. R. Ellis & T. Rothman, *Lost horizons*, American Journal of Physics **61**, pp. 883–893, 1993. Cited on pages 94, 98, and 99.

100 See, for example, the Hollywood film *Contact* by Robert Zemeckis, based on the book by Carl Sagan, *Contact*, Simon & Schuster, 1985. Cited on page 100.

101 See, for example, the international bestseller by Stephen Hawking, *A Brief History of Time – From the Big Bang to Black Holes*, 1988. Cited on page 102.

102 L. Rosenfeld, *Quantentheorie und Gravitation*, in H.-J. Treder, editor, *Entstehung, Entwicklung und Perspektiven der Einsteinschen Gravitationstheorie*, Springer Verlag, 1966. Cited on page 105.

103 Holography in high-energy physics is connected with the work of 't Hooft and Susskind. See for example G. 't Hooft, *Dimensional reduction in quantum gravity*, pp. 284–296, in A. Ali, J. Ellis & S. Randjbar-Daemi, *Salaamfeest*, 1993, or the much-cited paper by L. Susskind, *The world as a hologram*, Journal of Mathematical Physics **36**, pp. 6377–6396, 1995, preprint at arxiv.org/abs/hep-th/9409089. A good modern overview is Ref. 30. Cited on pages 106 and 114.

104 D. Bohm & B. J. Hiley, *On the intuitive understanding of nonlocality as implied by quantum theory*, Foundations of Physics 5, pp. 93–109, 1975. Cited on page 107.

105 S. Lloyd, *Computational capacity of the universe*, Physical Review Letters **88**, p. 237901, 2002. Cited on page 109.

106 Gottfried Wilhelm Leibniz, *La Monadologie*, 1714. Written in French, it is available freely at www.uqac.uquebec.ca/zone30/Classiques_des_sciences_sociales and in various other languages on other websites. Cited on page 109.

107 See, for example, H. Wussing & P. S. Alexandrov editors, *Die Hilbertschen Probleme*, Akademische Verlagsgesellschaft Geest & Portig, 1983, or Ben H. Yandell, *The Honours Class: Hilbert's Problems and their Solvers*, A.K. Peters, 2002. Cited on page 110.

108 A large part of the study of dualities in string and M theory can be seen as investigations into the detailed consequences of extremal identity. For a review of dualities, see P. C. Argyres, *Dualities in supersymmetric field theories*, Nuclear Physics Proceedings Supplement **61**, pp. 149–157, 1998, preprint at arxiv.org/abs/hep-th/9705076. A classical version of duality is discussed by M. C. B. Abdalla, A. L. Gadelka & I. V. Vancea, *Duality between coordinates and the Dirac field*, preprint at arxiv.org/abs/hep-th/0002217. Cited on page 114.

109 See L. Susskind & J. Uglum, *Black holes, interactions, and strings*, preprint at arxiv.org/abs/hep-th/9410074, or L. Susskind, *String theory and the principle of black hole complementarity*, Physical Review Letters **71**, pp. 2367–2368, 1993, and M. Karliner, I. Klebanov & L. Susskind, *Size and shape of strings*, International Journal of Modern Physics A **3**, pp. 1981–1996, 1988, as well as L. Susskind, *Structure of hadrons implied by duality*, Physical Review D **1**, pp. 1182–1186, 1970. Cited on pages 119 and 134.

110 M. Planck, *Über irreversible Strahlungsvorgänge*, Sitzungsberichte der königlich-preußischen Akademie der Wissenschaften zu Berlin pp. 440–480, 1899. Today it is commonplace to use Dirac's $\hbar = h/2\pi$ instead of Planck's h, which Planck originally called b. Cited on page 120.

111 P. Facchi & S. Pascazio, *Quantum Zeno and inverse quantum Zeno effects*, pp. 147–217, in E. Wolf editor, *Progress in Optics*, **42**, 2001. Cited on page 123.

112 Aristotle, *Of Generation and Corruption*, book I, part 2. See Jean-Paul Dumont, *Les écoles présocratiques*, Folio Essais, Gallimard, p. 427, 1991. Cited on page 123.

113 See for example the speculative model of vacuum as composed of Planck-size spheres proposed by F. Winterberg, Zeitschrift für Naturforschung **52a**, p. 183, 1997. Cited on page 124.

114 The Greek salt-and-water argument and the fish argument are given by Lucrece, in full Titus Lucretius Carus, *De natura rerum*, *c.* 60 bce. Cited on pages 125 and 140.

115 J. H. Schwarz, *The second superstring revolution*, Colloquium-level lecture presented at the Sakharov Conference in Moscow, May 1996, preprint at arxiv.org/abs/hep-th/9607067. Cited on pages 126 and 128.

116 Simplicius, *Commentary on the Physics of Aristotle*, 140, 34. This text is cited in Jean-Paul Dumont, *Les écoles présocratiques*, Folio Essais, Gallimard, p. 379, 1991. Cited on page 127.

117 D. Olive & C. Montonen, *Magnetic monopoles as gauge particles*, Physics Letters **72B**, pp. 117–120, 1977. Cited on page 128.

118 A famous fragment from Diogenes Laërtius (IX 72) quotes Democritus as follows: 'By convention hot, by convention cold, but in reality, atoms and void; and also in reality we know nothing, since truth is at the bottom.' Cited on page 129.

119 This famous statement is found at the beginning of chapter XI, 'The Physical Universe', in Arthur Eddington, *The Philosophy of Physical Science*, Cambridge, 1939. Cited on page 130.

120 Plato, *Parmenides*, *c.* 370 bce. It has been translated into most languages. Reading it aloud, like a song, is a beautiful experience. A pale reflection of these ideas is Bohm's concept of 'unbroken wholeness'. Cited on page 130.

121 P. Gibbs, *Event-symmetric physics*, preprint at arxiv.org/abs/hep-th/9505089; see also his website www.weburbia.com/pg/contents.htm. Cited on page 130.

122 J. B. Hartle, & S. W. Hawking, *Path integral derivation of black hole radiance*, Physical Review D **13**, pp. 2188–2203, 1976. See also A. Strominger & C. Vafa, *Microscopic origin of Bekenstein–Hawking entropy*, Physics Letters B **379**, pp. 99–104, 1996, preprint at arxiv.org/abs/hep-th/9601029. For another derivation of black hole entropy, see G. T. Horowitz & J. Polchinski, *A correspondence principle for black holes and strings*, Physical Review D 55, pp. 6189–6197, 1997, preprint at arxiv.org/abs/hep-th/9612146. Cited on pages 133 and 142.

123 J. Maddox, *When entropy does not seem extensive*, Nature **365**, p. 103, 1993. The issue is now explored in all textbooks discussing black holes. John Maddox (b. 1925 Penllergaer, d. 1999 Abergavenny) was famous for being one of the few people who was knowledgeable in most natural sciences. Cited on page 133.

124 L. Bombelli, R. K. Koul, J. Lee & R. D. Sorkin, *Quantum source of entropy of black holes*, Physical Review D **34**, pp. 373–383, 1986. Cited on page 133.

125 The analogy between polymers and black holes is due to G. Weber, *Thermodynamics at boundaries*, Nature **365**, p. 792, 1993. Cited on page 133.

126 See the classic text by Pierre-Gilles de Gennes, *Scaling Concepts in Polymer Physics*, Cornell University Press, 1979. Cited on page 134.

127 See for example S. Majid, *Introduction to braided geometry and q-Minkowski space*, preprint at arxiv.org/abs/hep-th/9410241, or S. Majid, *Duality principle and braided geometry*, preprint at arxiv.org/abs/hep-th/9409057. Cited on pages 135 and 136.

128 The relation between spin and statistics has been studied recently by M. V. Berry & J. M. Robbins, *Quantum indistinguishability: spin–statistics without relativity or field theory?*, in R. C. Hilborn & G. M. Tino editors, *Spin–Statistics Connection and Commutation Relations*, American Institute of Physics, 2000. Cited on page 137.

129 A. Gregori, *Entropy, string theory, and our world*, preprint at arxiv.org/abs/hep-th/0207195. Cited on pages 138 and 139.

130 String cosmology is a pastime for many. Examples include N. E. Mavromatos, *String cosmology*, preprint at arxiv.org/abs/hep-th/0111275, and N. G. Sanchez, *New developments in string gravity and string cosmology – a summary report*, preprint at arxiv.org/abs/hep-th/0209016. Cited on page 139.

131 On the present record, see en.wkipedia.org/wiki/Ultra-high-energy_cosmic_ray and fr.wkipedia.org/wiki/Zetta-particule. Cited on page 140.

132 P. F. Mende, *String theory at short distance and the principle of equivalence*, preprint at arxiv.org/abs/hep-th/9210001. Cited on page 140.

133 An example is given by A. A. Slavnov, *Fermi–Bose duality via extra dimension*, preprint at arxiv.org/abs/hep-th/9512101. See also the standard work by Michael Stone editor, *Bosonization*, World Scientific, 1994. Cited on page 140.

134 A weave model of space-time appears in certain approaches to quantum gravity, such as Ref. 27. On a slightly different topic, see also S. A. Major, *A spin network primer*, preprint at arxiv.org/abs/gr-qc/9905020. Cited on page 140.

135 L. Smolin & Y. Wan, *Propagation and interaction of chiral states in quantum gravity*, preprint at arxiv.org/abs/0710.1548, and references therein. Cited on page 140.

136 A good introduction into his work is the paper D. Kreimer, *New mathematical structures in renormalisable quantum field theories*, Annals of Physics **303**, pp. 179–202, 2003, erratum ibid. 305, p. 79, 2003, preprint at arxiv.org/abs/hep-th/0211136. Cited on page 141.

137 Introductions to holography include E. Alvarez, J. Conde & L. Hernandez, *Rudiments of holography*, preprint at arxiv.org/abs/hep-th/0205075, and Ref. 30. The importance of holography in theoretical high-energy physics was underlined by the discovery of J. Maldacena, *The large N limit of superconformal field theories and supergravity*, preprint at arxiv.org/abs/hep-th/9711200. Cited on page 141.

138 X.-G. Wen, *From new states of matter to a unification of light and electrons*, preprint at arxiv.org/abs/0508020. Cited on page 141.

139 J. S. Avrin, *A visualizable representation of the elementary particles*, Journal of Knot Theory and Its Ramifications **14**, pp. 131–176, 2005. Cited on pages 141 and 347.

140 The well-known ribbon model is presented in S. Bilson-Thompson, *A topological model of composite preons*, preprint at arxiv.org/hep-ph/0503213; S. Bilson-Thompson, F. Markopoulou & L. Smolin, *Quantum gravity and the standard model*, preprint at arxiv.org/hep-th/0603022; S. Bilson-Thompson, J. Hackett, L. Kauffman & L. Smolin, *Particle identifications from symmetries of braided ribbon network invariants*, preprint at arxiv.org/abs/0804.0037; S. Bilson-Thompson, J. Hackett & L. Kauffman, *Particle topology, braids, and braided belts*, preprint at arxiv.org/abs/0903.1376. Cited on pages 141, 168, and 347.

141 R. J. Finkelstein, *A field theory of knotted solitons*, preprint at arxiv.org/abs/hep-th/0701124. See also R. J. Finkelstein, *Trefoil solitons, elementary fermions, and* $SU_q(2)$, preprint at arxiv.org/abs/hep-th/0602098, R. J. Finkelstein & A. C. Cadavid, *Masses and interactions of q-fermionic knots*, preprint at arxiv.org/abs/hep-th/0507022, and R. J. Finkelstein, *A knot model suggested by the standard electroweak theory*, preprint at arxiv.org/abs/hep-th/0408218. Cited on pages 141 and 347.

142 Louis H. Kauffman, *Knots and Physics*, World Scientific, 1991. A wonderful book. Cited on pages 141 and 277.

143 S. K. Ng, *On a knot model of the* π^+ *meson*, preprint at arxiv.org/abs/hep-th/0210024, and S. K. Ng, *On a classification of mesons*, preprint at arxiv.org/abs/hep-ph/0212334. Cited on pages 141 and 347.

144 For a good introduction to superstrings, see the lectures by B. Zwiebach, *String theory for pedestrians*, agenda.cern.ch/fullAgenda.php?ida=a063319. For an old introduction to superstrings, see the famous text by M. B. Green, J. H. Schwarz & E. Witten, *Superstring Theory*, Cambridge University Press, volumes 1 and 2, 1987. Like all the other books on superstrings, they contain no statement that is applicable to or agrees with the strand model. Cited on pages 141 and 348.

145 See A. Sen, *An introduction to duality symmetries in string theory*, in *Les Houches Summer School: Unity from Duality: Gravity, Gauge Theory and Strings (Les Houches, France, 2001)*, Springer Verlag, **76**, pp. 241–322, 2002. Cited on page 141.

146 Brian Greene regularly uses the name *string conjecture*. For example, he did so in a podium discussion at TED in 2009; the video of the podium discussion can be downloaded at www.ted.org. Cited on page 142.

147 L. Susskind, *Some speculations about black hole entropy in string theory*, preprint at arxiv.org/abs/hep-th/9309145. G. T. Horowitz & J. Polchinski, *A correspondence principle for black holes and strings*, Physical Review D 55, pp. 6189–6197, 1997, preprint at arxiv.org/abs/hep-th/9612146. Cited on pages 142 and 463.

148 F. Wilczek, *Getting its from bits*, Nature **397**, pp. 303–306, 1999. Cited on page 143.

149 M. R. Douglas, *Understanding the landscape*, preprint at arxiv.org/abs/hep-th/0602266; his earlier papers also make the point. For the larger estimate, see W. Taylor & Y.-N. Wang, *The F-theory geometry with most flux vacua*, preprint at arxiv.org/abs/1511.03209. Cited on page 144.

150 The difficulties of the string conjecture are discussed in the well-known internet blog by Peter Woit, *Not even wrong*, at www.math.columbia.edu/~woit/blog. Several Nobel Prize winners for particle physics dismiss the string conjecture: Martin Veltman, Sheldon Glashow, Burton Richter, Richard Feynman and since 2009 also Steven Weinberg are among those who did so publicly. Cited on pages 144 and 168.

151 The present volume was originally started with the aim to clarify the basic principles of string theory and to simplify it as much as possible. In particular, the first six chapters and the last chapter were conceived, structured and written with that aim. They are older than the strand model. Later on, the project took an unexpected direction, as explained in Ref. 19. Cited on page 144.

152 Searches for background-free approaches are described by E. Witten, *Quantum background independence in string theory*, preprint at arxiv.org/abs/hep-th/9306122 and E. Witten, *On background-independent open string theory*, preprint at arxiv.org/abs/hep-th/9208027. Cited on page 145.

153 In fact, no other candidate model that fulfils all requirements for the unified theory is available in the literature so far. This might change in the future, though. Cited on page 151.

154 In December 2011 already I did not recall when I conceived and first explored the fundamental principle. It was between 2001 and 2006. It was a gradual process started by the exploration of the belt trick and driven by the aim to describe space in a way that produces Lorentz invariance based on extended entities (as I used to call them back then). The idea and the graphs of extended constituents for spin 1/2 is part of my volume 4 (on spin and particle exchange) since the 1990s. For many years I searched for a way to extend that model to include vacuum and the rest of physics. This was my pastime when taking the subway to work during those years. Cited on page 158.

155 S. Carlip, *The small scale structure of spacetime*, preprint at arxiv.org/abs/1009.1136. This paper deduces the existence of fluctuating lines in vacuum from a number of arguments that are completely independent of the strand model. Steven Carlip has dedicated much of his research to the exploration of this topic. One summary is S. Carlip, *Spontaneous dimensional reduction in quantum gravity*, preprint at arxiv.org/abs/1605.05694; its is also instructive to read his review S. Carlip, *Dimension and dimensional reduction in quantum gravity*, Classical and Quantum Gravity **34**, p. 193001, 2017, preprint at arxiv.org/abs/1705.05417. With the strand model in the back of one's mind, these results are even more fascinating. Cited on pages 163 and 302.

156 David Deutsch states that any good explanation must be 'hard to vary'. This must also apply to a unified model, as it claims to explain everything that is observed. See D. Deutsch, *A new way to explain explanation*, video talk at www.ted.org. Cited on pages 167 and 423.

157 L. Bombelli, J. Lee, D. Meyer & R. D. Sorkin, *Space-time as a causal set*, Physical Review Letters **59**, pp. 521–524, 1987. See also the review by J. Henson, *The causal set approach to quantum gravity*, preprint at arxiv.org/abs/gr-qc/0601121. Cited on pages 168 and 301.

158 D. Finkelstein, *Homotopy approach to quantum gravity*, International Journal of Theoretical Physics 47, pp. 534–552, 2008. Cited on pages 168 and 301.

159 L. H. Kauffman & S. J. Lomonaco, *Quantum knots*, preprint at arxiv.org/abs/quant-ph/0403228. See also S. J. Lomonaco & L. H. Kauffman, *Quantum knots and mosaics*, preprint at arxiv.org/abs/0805.0339. Cited on page 168.

160 Immanuel Kant, *Critik der reinen Vernunft*, 1781, is a famous but long book that every philospher pretends to have read. In his book, Kant introduced the 'a priori' existence of space and time. Cited on page 170.

161 The literature on circularity is rare. For two interesting exceptions, see L. H. Kauffman, *Knot logic*, downloadable from www2.math.uic.edu/~kauffman, and L. H. Kauffman, *Reflexivity and eigenform*, Constructivist Foundations **4**, pp. 121–137, 2009. Cited on page 171.

162 Information on the belt trick is scattered across many books and few papers. The best source of information on this topic are websites. For belt trick visualizations see www.evl.uic.edu/hypercomplex/html/dirac.html, www.evl.uic.edu/hypercomplex/html/handshake.html, or www.gregegan.net/APPLETS/21/21.html. For an excellent literature summary and more movies, see www.math.utah.edu/~palais/links.html. None of these sites or the cited references seem to mention that there are *many* ways to perform the belt trick; this seems to be hidden knowledge. In September 2009, Greg Egan took up my suggestion and changed his applet to show an additional version of the belt trick. Cited on pages 178 and 180.

163 There is an interesting exploration behind this analogy between a non-dissipative system – a free quantum particle moving in vacuum – and a dissipative system – a macroscopic body drawn through a viscous liquid, say honey. The first question is to discover why this analogy is possible at all. (A careful distinction between the cases with spin 0, spin 1 and spin 1/2 are necessary.) The second question is the exploration of the motion of bodies of general shape in viscous fluids at low Reynolds numbers and under constant force. For the best overview of this question, see the beautiful article by O. Gonzalez, A. B. A. Graf & J. H. Maddocks, *Dynamics of a rigid body in a Stokes fluid*, Journal of Fluid Mechanics **519**, pp. 133–160, 2004. Cited on pages 197 and 362.

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164 D. Bohm, R. Schiller & J. Tiomno, *A causal interpretation of the Pauli equation (A)*, Supplementi al Nuovo Cimento **1**, pp. 48 – 66, 1955, and D. Bohm & R. Schiller, *A causal interpretation of the Pauli equation (B)*, Supplementi al Nuovo Cimento **1**, pp. 67–91, 1955. The authors explore an unusual way to interpret the wavefunction, which is of little interest here; but doing so, they give and explore the description of Pauli spinors in terms of Euler angles. Cited on page 200.

165 Richard P. Feynman, *QED – The Strange Theory of Light and Matter*, Princeton University Press 1988. This is one of the best summaries of quantum theory ever written. Every physicist should read it. Cited on pages 201, 207, 219, 225, and 460.

166 S. Kochen & E. P. Specker, *The problem of hidden variables in quantum mechanics*, 17, pp. 59–87, 1967. This is a classic paper. Cited on page 205.

167 A. Aspect, J. Dalibard & G. Roger, *Experimental tests of Bell's inequalities using time-varying analyzers*, Physical Review Letters **49**, pp. 1804–1807, 1982, Cited on page 209.

168 L. Kauffman, *New invariants of knot theory*, American Mathematical Monthly **95**, pp. 195–242, 1987. See also the image at the start of chapter 6 of Louis H. Kauffman, *On Knots*, Princeton University Press, 1987. Cited on page 211.

169 The details on the speed of photons are explained in any textbook on quantum electrodynamics. The issue is also explained by Feynman in Ref. 165 on page 89. Cited on page 214.

170 J.-M. Lévy-Leblond, *Nonrelativitic particles and wave equations*, Communications in Mathematical Physics **6**, pp. 286–311, 1967. See also A. Galindo & C. Sánchez del Río, *Intrinsic magnetic moment as a nonrelativistic phenomenon*, American Journal of Physics **29**, pp. 582–584, 1961, and V. I. Fushchich, A. G. Nikitin & V. A. Salogub, *On the non-relativistic motion equations in the Hamiltonian form*, Reports on Mathematical Physics **13**, pp. 175–185, 1978. Cited on page 216.

171 L. Lerner, *Derivation of the Dirac equation from a relativistic representation of spin*, European Journal of Phyics **17**, pp. 172–175, 1996. Cited on page 216.

172 E. P. Battey-Pratt & T. J. Racey, *Geometric model for fundamental particles*, International Journal of Theoretical Physics **19**, pp. 437–475, 1980. Without knowing this work, C. Schiller had deduced the same results in 2008. Cited on pages 217, 218, and 470.

173 A. Abraham, *Prinzipien der Dynamik des Elektrons*, Annalen der Physik **10**, pp. 105–179, 1903, J. Frenkel, *Die Elektrodynamik des rotierenden Elektrons*, Zeitschrift für Physik **37**, pp. 243–262, 1926, L. H. Thomas, *The motion of a spinning electron*, Nature April 10, p. 514, 1926, and L. H. Thomas, *The kinematics of an electron with an axis*, Philosophical Magazine **3**, pp. 1–22, 1927. See also W. E. Baylis, *Surprising symmetries in relativistic charge dynamics*, preprint at arxiv.org/abs/physics/0410197. See also W. E. Baylis, *Quantum/classical interface: a geometric approach from the classical side*, pp. 127–154 and W. E. Baylis, *Geometry of paravector space with applications to relativistic physics*, pp. 363–387 in *Computational Noncommutative Algebra and Applications*, Proceedings of the NATO Advanced Study Institute, NATO Science Series II, vol. 136, ed. J. Byrnes, Kluwer Academic 2004. W. E. Baylis, R. Cabrera & D. Keselica, *Quantum/classical interface: fermion spin*, preprint at arxiv.org/abs/0710.3144. D. Hestenes, *Zitterbewegung Modelling*, Foundations of Physics **23**, pp. 365–386, 1993. D. Hestenes, *Zitterbewegung in quantum mechanics – a research program*, preprint at arxiv.org/abs/0802.2728. See also D. Hestenes, *Reading the electron clock*, preprint at arxiv.org/abs/0802.3227 and his webpage modelingnts.la.asu.edu/html/GAinQM.html. A. Loinger & A. Sparzani, *Dirac equation without Dirac matrices*, Il Nuovo Cimento **39**, pp. 1140–1145, 1965. D. Bohm, P. Hillion, T. Takabayasi & J.-P. Vigier, *Relativistic rotators and bilocal theory*, Progress of Theoretical Physics **23**, pp. 496–511, 1960. A. Challinor, A. Lasenby, S. Gill & C. Doran, *A relativistic, causal account of a spin measurement*, Physics Letters A **218**, pp. 128–138, 1996. E. Santamato, *The role of Dirac equation in the classical mechanics of the relativistic top*, preprint at arxiv.org/abs/0808.3237. Cited on page 219.

174 The concept of Zitterbewegung was formulated in E. Schrödinger, *Über die kräftefreie Bewegung in der relativistischen Quantenmechanik*, Berliner Berichte pp. 418–428, 1930, and *Zur Quantendynamik des Elektrons*, Berliner Berichte pp. 63–72, 1931. Numerous subsequent papers discuss these publications. Cited on page 219.

175 See for example the book by Martin Rivas, *Kinematic Theory of Spinning Particles*, Springer, 2001. Cited on page 219.

176 The basic papers in the field of stochastic quantization are W. Weizel, *Ableitung der Quantentheorie aus einem klassischen, kausal determinierten Modell*, Zeitschrift für Physik A **134**, pp. 264–285, 1953, W. Weizel, *Ableitung der Quantentheorie aus einem klassischen Modell – II*, Zeitschrift für Physik A **135**, pp. 270–273, 1954, W. Weizel, *Ableitung der quantenmechanischen Wellengleichung des Mehrteilchensystems aus einem klassischen Modell*, Zeitschrift für Physik A **136**, pp. 582–604, 1954. This work was taken up by E. Nelson, *Derivation of the Schrödinger equation from Newtonian mechanics*, Physical Review **150**, pp. 1079–1085, 1969, and in Edward Nelson, *Quantum Fluctuations*, Princeton University Press 1985, also downloadable at www.math.princeton.edu/~nelson/books.html, and the book Edward Nelson, *Stochastic Quantization*, Princeton University Press 1985. See also L. Fritsche & M. Haugk, *A new look at the derivation of the Schrödinger equation from Newtonian mechanics*, Annalen der Physik **12**, pp. 371–402, 2003. A summary of Nelson's approach is also given in F. Markopoulou & L. Smolin, *Quantum theory from quantum gravity*, Physical Review D **70**, p. 124029, 2004, preprint at arxiv.org/abs/gr-qc/0311059. See also the important criticism by T. C. Wallstrom, *Inequivalence between the Schrödinger equation and the Madelung hydrodynamic equation*, Physical Review A **49**, pp. 1613–1617, 1994, and T. C. Wallstrom, *The stochastic mechanics of the Pauli equation*, Transactions of the American Mathematical Society **318**, pp. 749–762, 1990. A proposed answer is L. Smolin, *Could quantum mechanics be an approximation to another theory?*, preprint at arxiv.org/quant-ph/abs/0609109. See also S. K. Srinivasan & E. C. G. Sudarshan, *A direct derivation of the Dirac equation via quaternion measures*, Journal of Physics A **29**, pp. 5181–5186, 1996. Cited on page 220.

177 Julian Schwinger, *Quantum Mechanics – Symbolism of Atomic Measurements*, Springer, 2001. Cited on page 222.

178 H. Nikolić, *How (not) to teach Lorentz covariance of the Dirac equation*, European Journal of Physics **35**, p. 035003, 2014, preprint at arxiv.org/abs/1309.7070. Cited on page 222.

179 For such an attempt, see the proposal by M. Rainer, *Resolution of simple singularities yielding particle symmetries in space-time*, Journal of Mathematical Physics **35**, pp. 646–655, 1994. Cited on page 225.

180 C. Schiller, *Deducing the three gauge interactions from the three Reidemeister moves*, preprint at arxiv.org/abs/0905.3905. Cited on pages 225 and 228.

181 G. T. Horowitz & J. Polchinski, *Gauge/gravity duality*, preprint at arxiv.org/abs/gr-qc/0602037. Note also the statement in the introduction that a graviton might be a composite of two spin-1 bosons, which is somewhat reproduced by the strand model of the graviton. A more concrete approach to gauge–gravity duality is made by M. van Raamsdonk, *Building up spacetime with quantum entanglement*, preprint at arxiv.org/1005.3035. This approach to gauge–gravity duality is close to that of the strand model. No citations.

182 K. Reidemeister, *Elementare Begründung der Knotentheorie*, Abhandlungen aus dem Mathematischen Seminar der Universität Hamburg **5**, pp. 24–32, 1926. Cited on pages 228 and 276.

183 L. Burns, *Maxwell's Equations are Universal for* Locally *Conserved Quantities*, Advances in Applied Clifford Algebras **29**, p. 62, 2019. See also J. A. Heras, *Can Maxwell's equations be obtained from the continuity equation?*, American Journal of Physics 75, pp. 652–657, 2007. Cited on pages 229, 235, and 243.

184 For an attempt to reconcile braided particle models and SU(5) GUT, see D. Cartin, *Braids as a representation space of SU(5)*, preprint at arxiv.org/pdf/1506.08067. Cited on page 246.

185 Sheldon Glashow, confirmed this to the author in an email; Richard Feynman, makes the point in James Gleick, *Genius: The Life and Science of Richard Feynman*, Vintage Books, 1991, page 288 and also in Robert Crease & Charles Mann, *The Second Creation: Makers of the Revolution in Twentieth-Century Physics*, Macmillan Publishing, page 418; Martin Veltman, writes this in his Nobel Prize Lecture, available on www.nobel.org. Cited on page 257.

186 For some of the background on this topic, see F. Wilczek & A. Zee *Appearance of gauge structures in simple dynamical systems*, Physical Review Letters **52**, pp. 2111–2114, 1984, A. Shapere & F. Wilczek, *Self-propulsion at low Reynolds number*, Physical Review Letters **58**, pp. 2051–2054, 1987, and A. Shapere & F. Wilczek, *Gauge kinematics of deformable bodies*, American Journal of Physics **57**, pp. 514–518, 1989. Cited on page 274.

187 R. Britto, F. Cachazo, B. Feng & E. Witten, *Direct proof of tree-level recursion relation in Yang–Mills theory*, preprint at arxiv.org/abs/hep-th/0501052. Cited on page 275.

188 D. V. Ahluwalia-Khalilova, *Operational indistinguishability of double special relativity from special relativity*, Classical and Quantum Gravity **22**, pp. 1433–1450, 2005, preprint at arxiv.org/abs/gr-qc/0212128; see also N. Jafari & A. Shariati, *Doubly special relativity: a new relativity or not?*, preprint at arxiv.org/abs/gr-qc/0602075. Cited on page 282.

189 E. Verlinde, *On the origin of gravity and the laws of Newton*, preprint at arxiv.org/abs/1001.0785. Cited on page 283.

190 G.-L. Lesage, *Lucrèce Newtonien*, Nouveaux mémoires de l'Académie Royale des Sciences et Belles Lettres pp. 404–431, 1747, or www3.bbaw.de/bibliothek/digital/struktur/03-nouv/1782/jpg-0600/00000495.htm. See also en.wikipedia.org/wiki/Le_Sage's_theory_of_gravitation. In fact, the first to propose the idea of gravitation as a result of small particles pushing masses was Nicolas Fatio de Duillier in 1688. Cited on page 285.

191 G. 't Hooft, *Dimensional reduction in quantum gravity*, preprint at arxiv.org/abs/gr-qc/9310026. Many of the ideas of this paper become easier to understand and to argue when the strand model is used. Cited on page 290.

192 S. Carlip, *Logarithmic corrections to black hole entropy from the Cardy formula*, Classical and Quantum Gravity **17**, pp. 4175–4186, 2000, preprint at arxiv.org/abs/gr-qc/0005017. Cited on page 291.

193 D. N. Page, *The Bekenstein Bound*, preprint at arxiv.org/abs/1804.10623. Cited on page 293.

194 On the limit for angular momentum of black holes, see Ref. 34. Cited on page 293.

195 F. Tamburini, C. Cuofano, M. Della Valle & R. Gilmozzi, *No quantum gravity signature from the farthest quasars*, preprint at arxiv.org/abs/1108.6005. Cited on page 298.

196 B.P. Abbott & al., (LIGO Scientific Collaboration and Virgo Collaboration) *Observation of gravitational waves from a binary black hole merger*, Physical Review Letters **116**, p. 061102, 2016, also available for free download at journals.aps.org/prl/pdf/10.1103/PhysRevLett.116.061102. See also the website www.ligo.caltech.edu. More about this discovery and its implications is told in volume II of the Motion Mountain series. Cited on page 299.

197 On torsion, see the excellent review by R. T. Hammond, *New fields in general relativity*, Contemporary Physics **36**, pp. 103–114, 1995. Cited on page 301.

198 H. Kleinert, & J. Zaanen, *World nematic crystal model of gravity explaining the absence of torsion*, Physics Letters A **324**, pp. 361–365, 2004. Cited on page 301.

199 The analogy between the situation around line defects and general relativity is explained in Ekkehart Kröner, *Kontinuumstheorie der Versetzungen und Eigenspannungen*, Springer, 1958, These ideas have been taken up and pursued by J. D. Eshelby, B. A. Bilby, and many others after them. Cited on page 301.

200 Loop quantum gravity is a vast research field. The complete literature is available at arxiv.org/archive/gr-qc. Cited on page 301.

201 G. 't Hooft, *Crystalline Gravity*, International Journal of Modern Physics A **24**, pp. 3243–3255, 2009, and also G. 't Hooft, *A locally finite model of gravity*, preprint at arxiv.org/abs/0804.0328. Cited on page 301.

202 L. Susskind, *New concepts for old black holes*, preprint at arxiv.org/abs/1311.3335, and also reference Ref. 147. Cited on page 302.

203 M. Botta Cantcheff, *Spacetime geometry as statistic ensemble of strings*, preprint at arxiv.org/abs/1105.3658. Cited on page 302.

204 N. Arkani-Hamed, L. Motl, A. Nicolis & C. Vafa, *The string landscape, black holes and gravity as the weakest force*, preprint at arxiv.org/abs/hep-th/0601001. The paper contradicts the strand model in multiple ways. Cited on page 303.

205 M. van Raamsdonk, *Comments on quantum gravity and entanglement*, preprint at arxiv.org/abs/0907.2939. Cited on page 303.

206 W. H. Zurek & K. S. Thorne, *Statistical mechanical origin of the entropy of a rotating, charged black hole*, Physical Review Letters **54**, pp. 2171–2175, 1985. Cited on page 303.

207 M. Shaposhnikov & C. Wetterich, *Asymptotic safety of gravity and the Higgs boson mass*, preprint at arxiv.org/abs/0912.0208. Cited on page 304.

208 M. M. Anber & J. F. Donoghue, *On the running of the gravitational constant*, preprint at arxiv.org/abs/1111.2875. Cited on page 304.

209 The 2016 data about modified Newtonian dynamics is found in S. McGaugh, F. Lelli & J. Schombert, *The radial acceleration relation in rotationally supported galaxies*, preprint at arxiv.org/abs/1609.05917, and in F. Lelli, S. S. McGaugh, J. M. Schombert & M. S. Pawlowski, *One law to rule them all: the radial acceleration relation of galaxies*, preprint at arxiv.org/abs/1610.08981. Cited on page 305.

210 C. H. Lineweaver & T. M. Davis, *Misconceptions about the big bang*, Scientific American pp. 36–45, March 2005. Cited on page 307.

211 Supernova Search Team Collaboration, A.G. Riess & al., *Observational evidence from supernovae for an accelerating universe and a cosmological constant*, Astronomical Journal **116**, pp. 1009–1038, 1998, preprint at arxiv.org/abs/astro-ph/9805201. Cited on page 307.

212 Stephen Hawking & Roger Penrose, *The Nature of Space and Time*, Princeton University Press, 1996. Cited on page 308.

213 C. Balázs & I. Szapudi, *Naturalness of the vacuum energy in holographic theories*, preprint at arxiv.org/abs/hep-th/0603133. See also C. Bambi & F. R. Urban, *Natural extension of the generalised uncertainty principle*, preprint at arxiv.org/abs/0709.1965. The same point is made by D. A. Easson, P. H. Frampton & G. F. Smoot, *Entropic accelerating universe*, preprint at arxiv.org/abs/1002.4278. No citations.

214 W. Fischler & L. Susskind, *Holography and Cosmology*, preprint at arxiv.org/abs/hep-th/9806039. No citations.

215 For a review of recent cosmological data, see D. N. Spergel, R. Bean, O. Doré, M. R. Nolta, C. L. Bennett, G. Hinshaw, N. Jarosik, E. Komatsu,

L. Page, H. V. Peiris, L. Verde, C. Barnes, M. Halpern, R. S. Hill, A. Kogut, M. Limon, S. S. Meyer, N. Odegard, G. S. Tucker, J. L. Weiland, E. Wollack & E. L. Wright, *Wilkinson Microwave Anisotropy Probe (WMAP) three year results: implications for cosmology*, preprint at arxiv.org/abs/astro-ph/0603449. Cited on page 309.

216 There is a large body of literature that has explored a time-varying cosmological constant, especially in relation to holography. An example with many references is L. Xu, J. Lu & W. Li, *Time variable cosmological constants from the age of the universe*, preprint at arxiv.org/abs/0905.4773. No citations.

217 D. Wiltshire, *Gravitational energy and cosmic acceleration*, preprint at arxiv.org/abs/0712.3982 and D. Wiltshire, *Dark energy without dark energy*, preprint at arxiv.org/abs/0712.3984. A newer and well-argued paper is D. L. Wiltshire, *Cosmic structure, averaging and dark energy*, preprint at arxiv.org/abs/1311.3787. See also T. Buchert, A. A. Coley, H. Kleinert, B. F. Roukema & D. L. Wiltshire, *Observational challenges for the standard FLRW model*, preprint at arxiv.org/abs/1512.03313. No citations.

218 The attribution to Voltaire could not be confirmed. Cited on page 313.

219 V. Crede & C. A. Meyer, *The experimental status of glueballs*, Progress in Particle and Nuclear Physics **63**, pp. 74–116, 2009. Cited on page 344.

220 E. Klempt & A. Zaitsev, *Glueballs, hybrids, multiquarks. Experimental facts versus QCD inspired concepts*, Physics Reports **454**, 2007, preprint at arxiv.org/abs/0708.4016. Cited on pages 342 and 344.

221 R. V. Buniy & T. W. Kephart, *A model of glueballs*, preprint at arxiv.org/pdf/hep-ph/0209339; R. V. Buniy & T. W. Kephart, *Universal energy spectrum of tight knots and links in physics*, preprint at arxiv.org/pdf/hep-ph/0408025; R. V. Buniy & T. W. Kephart, *Glueballs and the universal energy spectrum of tight knots and links*, preprint at arxiv.org/pdf/hep-ph/0408027. See also J. P. Ralston, *The Bohr atom of glueballs*, preprint at arxiv.org/pdf/hep-ph/0301089. Cited on page 344.

222 A. J. Niemi, *Are glueballs knotted closed strings?*, pp. 127–129, in H. Suganuma, N. Ishii, M. Oka, H. Enyo, T. Hatsuda, T. Kunihiro & K. Yazaki editors, *Color confinement and hadrons in quantum chromodynamics*, World Scientific, 2003, preprint at arxiv.org/pdf/hep-th/0312133. See also Y. M. Cho, B. S. Park & P. M. Zhang, *New interpretation of Skyrme theory*, preprint at arxiv.org/pdf/hep-th/0404181; K. Kondo, A. Ono, A. Shibata, T. Shinohara & T. Murakami, *Glueball mass from quantized knot solitons and gauge-invariant gluon mass*, Journal of Physics A **39**, pp. 13767–13782, 2006, preprint at arxiv.org/abs/hep-th/0604006. Cited on page 344.

223 See the one million dollar prize described at www.claymath.org/millennium/Yang-Mills_Theory. Cited on page 345.

224 For a clear review on the topic and the planned experiments, see E. Fiorini, *Measurement of neutrino mass in double beta decay*, Europhysics News **38**, pp. 30–34, 2007, downloadable at www.europhysicsnews.org. Cited on page 327.

225 For example, see the detailed discussion of neutrino properties at pdg.web.cern.ch or, in print, in Ref. 233. Cited on page 327.

226 For a possible third approach, see A. F. Nicholson & D. C. Kennedy, *Electroweak theory without Higgs bosons*, International Journal of Modern Physics A **15**, pp. 1497–1519, 2000, preprint at arxiv.org/abs/hep-ph/9706471. Cited on page 329.

227 M. Veltman, *The Higgs system*, lecture slides at www.nikhef.nl/pub/theory/academiclectures/Higgs.pdf. See also his CERN Yellow Report 97-05, *Reflections on the*

Higgs system, 1997, and the paper H. Veltman & M. Veltman, *On the possibility of resonances in longitudinally polarized vector boson scattering*, Acta Physics Polonica B **22**, pp. 669–695, 1991. Cited on page 329.

228 J. W. Moffat & V. T. Thot, *A finite electroweak model without a Higgs particle*, preprint at arxiv.org/abs/0812.1991. The ideas go back to D. Evens, J. W. Moffat, G. Kleppe & R. P. Woodard, *Nonlocal regularizations of gauge theories*, Physical Review D **43**, pp. 499–519, 1991. For more details on how to introduce non-locality while maintaining current conservation and unitarity, see G. Kleppe & R. P. Woodard, *Non-local Yang-Mills*, Nuclear Physics B **388**, pp. 81–112, 1992, preprint at arxiv.org/abs/hep-th/9203016. For a different approach that postulates no specific origin for the W and Z masses, see J. W. Moffat, *Ultraviolet complete electroweak model without a Higgs particle*, preprint at arxiv.org/abs/1006.1859. Cited on page 329.

229 H. B. Nielsen & P. Olesen, *A vortex line model for dual strings*, Nuclear Physics B **61**, pp. 45–61, 1973. Cited on pages 333 and 386.

230 B. Andersson, G. Gustafson, G. Ingelman & T. Sjöstrand, *Parton fragmentation and string dynamics*, Physics Reports **97**, pp. 31–145, 1983. Cited on page 333.

231 C. B. Thorn, *Subcritical string and large N QCD*, preprint at arxiv.org/abs/0809.1085. Cited on page 333.

232 A. J. Buchmann & E. M. Henley, *Intrinsic quadrupole moment of the nucleon*, Physical Review C **63**, p. 015202, 2000. Alfons Buchmann also predicts that the quadrupole moment of the other, strange $J = 1/2$ octet baryons is positive, and predicts a prolate structure for all of them (private communication). For the decuplet baryons, with $J = 3/2$, the quadrupole moment can often be measured spectroscopically, and is always negative. The four Δ baryons are thus predicted to have a negative intrinsic quadrupole moment and thus an oblate shape. This explained in A. J. Buchmann & E. M. Henley, *Quadrupole moments of baryons*, Physical Review D **65**, p. 073017, 2002. For recent updates, see A. J. Buchmann, *Charge form factors and nucleon shape*, pp. 110–125, in C. N. Papanicolas & Aron Bernstein editors, *Shape of Hadrons Workshop Conference*, Athens, Greece, 27-29 April 2006, AIP Conference Proceedings 904. Cited on pages 336 and 340.

233 C. Patrignani & al., (Particle Data Group), Chinese Physics C **40**, p. 100001, 2016, or pdg.web.cern.ch. Cited on pages 337, 338, 339, 361, 363, 376, 377, 378, 379, 464, and 466.

234 A review on Regge trajectories and Chew-Frautschi plots is W. Drechsler, *Das Regge-Pol-Modell*, Naturwissenschaften **59**, pp. 325–336, 1972. See also the short lecture on courses.washington.edu/phys55x/Physics557_lec11.htm. Cited on page 337.

235 Kurt Gottfried & Victor F. Weisskopf, *Concepts of Particle Physics*, Clarendon Press, Oxford, 1984. Cited on page 338.

236 G. 't Hooft, G. Isidori, L. Maiani, A. D. Polosa & V. Riquer, *A theory of scalar mesons*, Physics Letters B **662**, pp. 424–430, 2008, preprint at arxiv.org/abs/0801.2288. However, other researchers, such as arxiv.org/abs/1404.5673, argue against the tetraquark interpretation. The issue is not closed. Cited on page 342.

237 M. Karliner, *Doubly heavy tetraquarks and baryons*, preprint at arxiv.org/abs/1401.4058. Cited on page 346.

238 J. Viro & O. Viro, *Configurations of skew lines*, Leningrad Mathematical Journal **1**, pp. 1027–1050, 1990, and updated preprint at arxiv.org/abs/math.GT/0611374. Cited on page 347.

239 W. Thomson, *On vortex motion*, Transactions of the Royal Society in Edinburgh pp. 217–260, 1868. This famous paper stimulated much work on knot theory. Cited on page 347.

240 H. Jehle, *Flux quantization and particle physics*, Physical Review D **6**, pp. 441–457, 1972, and H. Jehle, *Flux quantization and fractional charge of quarks*, Physical Review D **6**, pp. 2147–2177, 1975. Cited on page 347.

241 T. R. Mongan, *A holographic charged preon model*, preprint at arxiv.org/abs/0801.3670. Cited on page 347.

242 The arguments can be found in A. H. Chamseddine, A. Connes & V. Mukhanov, *Geometry and the quantum: basics*, preprint at arxiv.org/abs/1411.0977 and in A. H. Chamseddine & A. Connes, *Why the standard model*, Journal of Geometry and Physics **58**, pp. 38–47, 2008, preprint at arxiv.org/abs/0706.3688. Cited on page 348.

243 Jacob's rings are shown, for example, in the animation on www.prestidigitascience.fr/index.php?page=anneaux-de-jacob. They are already published in the book by Tom Tit, *La science amusante*, 1870, and the images were reprinted the popular science books by Edi Lammers, and, almost a century later on, even in the mathematics column and in one of the books by Martin Gardner. See also www.lhup.edu/~dsimanek/scenario/toytrick.htm. Cited on page 351.

244 R. Boughezal, J. B. Tausk & J. J. van der Bij, *Three-loop electroweak corrections to the W-boson mass and* $\sin^2\theta_{\text{eff}}$ *in the large Higgs mass limit*, Nuclear Physics B **725**, pp. 3–14, 2005, preprint at arxiv.org/abs/hep-ph/0504092. Cited on page 361.

245 The topic of the g-factor of the W boson and of charged fermions is covered in the delightful paper by Barry R. Holstein, *How large is the "natural" magnetic moment?*, American Journal of Physics **74**, pp. 1104–1111, 2006, preprint at arxiv.org/abs/hep-ph/0607187. Cited on page 363.

246 The calculations have been performed in August 2016 by Eric Rawdon. Cited on pages 361 and 363.

247 The calculations have been performed by Eric Rawdon and Maria Fisher. Cited on page 365.

248 The quark masses at Planck energy are due to a private communication by Xing Zhizhong and Zhou Shun. They are calculated following the method presented in *Quark mass hierarchy and flavor mixing puzzles*, preprint at arxiv.org/abs/1411.2713 and Zhi-zhong Xing, He Zhang & Shun Zhou, *Updated values of running quark and lepton masses*, preprint at arxiv.org/abs/0712.1419. Cited on page 365.

249 See H. Fritzsch, A.D. Özer, *A scaling law for quark masses*, preprint at arxiv.org/abs/hep-ph/0407308. Cited on page 366.

250 K. A. Meissner & H. Nicolai, *Neutrinos, axions and conformal symmetry*, preprint at arxiv.org/abs/0803.2814. Cited on pages 369 and 370.

251 M. Shaposhnikov, *Is there a new physics between electroweak and Planck scale?*, preprint at arxiv.org/abs/0708.3550. Cited on page 370.

252 Y. Diao, C. Ernst, A. Por & U. Ziegler, *The roplength of knots are almost linear in terms of their crossing numbers*, preprint at arxiv.org/abs/0912.3282. Cited on page 373.

253 H. Fritzsch & Z.-Z. Xing, *Lepton mass hierarchy and neutrino mixing*, preprint at arxiv.org/abs/hep-ph/0601104 Cited on page 378.

254 The effects of neutrino mixing, i.e., neutrino oscillations, were measured in numerous experiments from the 1960s onwards; most important were the experiments at Super-Kamiokande in Japan and at the Sudbury Neutrino Observatory in Canada. See Ref. 233. Cited on page 379.

255 M. Fukugita & T. Yanagida, *Baryogenesis without grand unification*, Physics Letters B **174**, pp. 45–47, 1986. Cited on page 380.

256 J. M. Cline, *Baryogenesis*, preprint at arxiv.org/abs/hep-ph/0609145 or the review by L. Canetti, M. Drewes & M. Shaposhnikov, *Matter and Antimatter in the Universe*, preprint at arxiv.org/abs/1204.4186. They explain the arguments that the standard model with its CKM-CP violation is not sufficient to explain baryogenesis. The opposite view, by the same authors, is found in L. Canetti, M. Drewes, T. Frossard & M. Shaposhnikov, *Dark matter, baryogenesis and neutrino oscillations from right handed neutrinos*, preprint at arxiv.org/abs/1208.4607; another opposing view is found in T. Brauner, *CP violation and electroweak baryogenesis in the Standard Model*, EPJ Web of Conferences **70**, p. 00078, 2014. This issue is not settled yet. Cited on page 380.

257 Several claims that the coupling constants changed with the age of the universe have appeared in the literature. The first claim was by J. K. Webb, V. V. Flambaum, C. W. Churchill, M. J. Drinkwater & J. D. Barrow, *A search for time variation of the fine structure constant*, Physical Review Letters **82**, pp. 884–887, 1999, preprint at arxiv.org/abs/astro-ph/9803165. Several subsequent claims have appeared. However, non of these claims has been confirmed by subsequent measurements. Nowadays, there is no data showing that the fine structure constant changes in time. Cited on page 384.

258 The poster on www.physicsoverflow.org referred to J. P. Lestone, *Physics based calculation of the fine structure constant*, preprint at arxiv.org/abs/physics/0703151. The preprint has never been published. Cited on page 389.

259 For a highly questionable, but still intriguing argument based on black hole thermodynamics that claims to deduce the limit $\alpha > \ln 3/48\pi \approx 1/137.26$, see S. Hod, *Gravitation, thermodynamics, and the fine-structure constant*, International Journal of Modern Physics D **19**, pp. 2319–2323, 2010. It might well be that similar or other arguments based on textbook physics will yield more convincing or even better limits in the future. Cited on page 389.

260 V. Arnold, *Topological Invariants of Plane Curves and Caustics*, American Mathematical Society, 1994. Cited on page 396.

261 See M. Pospelov & A. Ritz, *Electric dipole moments as probes of new physics*, preprint at arxiv.org/abs/hep-ph/0504231. Cited on page 396.

262 C. Schiller, *A conjecture on deducing general relativity and the standard model with its fundamental constants from rational tangles of strands*, Physics of Particles and Nuclei **50**, pp. 259–-299, 2019. Cited on pages 398 and 433.

263 D. Hilbert, *Über das Unendliche*, Mathematische Annalen **95**, pp. 161–190, 1925. Cited on page 418.

264 The *Book of Twenty-four Philosophers*, *c.* 1200, is attributed to the god Hermes Trismegistos, but was actually written in the middle ages. The text can be found in F. Hudry, ed., *Liber viginti quattuor philosophorum*, Turnholt, 1997, in the series *Corpus Christianorum, Continuatio Mediaevalis*, CXLIII a, tome III, part 1, of the Hermes Latinus edition project headed by P. Lucentini. There is a Spinozian cheat in the quote: instead of 'nature', the original says 'god'. The reason why this substitution is applicable is given above. Cited on page 423.

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265 As a disappointing example, see Gilles Deleuze, *Le Pli – Leibniz et le baroque*, Les Editions de Minuit, 1988. In this unintelligible, completely crazy book, the author pretends to investigate the implications of the idea that the *fold* (in French 'le pli') is the basic entity of matter and 'soul'. Cited on page 425.

266 Werner Heisenberg, *Der Teil und das Ganze*, Piper, 1969. The text shows well how boring the personal philosophy of an important physicist can be. Cited on page 426.

267 John Barrow wrote to the author saying that he might indeed have been the first to have used the T-shirt image, in his 1988 Gifford Lectures at Glasgow that were a precursor to his book John D. Barrow, *Theories of Everything: The Quest for Ultimate Explanation*, 1991. He added that one can never be sure, though. Cited on page 428.

268 René Descartes, *Discours de la méthode*, 1637. He used and discussed the sentence again in his *Méditations métaphysiques* 1641, and in his *Les principes de la philosophie* 1644. These books influenced many thinkers in the subsequent centuries. Cited on page 431.

269 D. D. Kelly, *Sleep and dreaming*, in *Principles of Neural Science*, Elsevier, New York, 1991. The paper summarises experiments made on numerous humans and shows that even during dreams, people's estimate of time duration corresponds to that measured by clocks. Cited on page 432.

270 Astrid Lindgren said this in 1977, in her speech at the fiftieth anniversary of Oetinger Verlag, her German publisher. The German original is: 'Alles was an Großem in der Welt geschah, vollzog sich zuerst in der Phantasie eines Menschen, und wie die Welt von morgen aussehen wird, hängt in großem Maß von der Einbildungskraft jener ab, die gerade jetzt lesen lernen.' The statement is found in Astrid Lindgren, *Deshalb brauchen Kinder Bücher*, Oetinger Almanach Nr. 15, p. 14, 1977. Cited on page 435.

CREDITS

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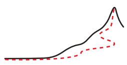

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Christoph Schiller, PhD Université Libre de Bruxelles, is a physicist and physics popularizer. This entertaining book is for students, teachers and anybody interested in modern research about fundamental physics.

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