

Lecture #19

Fast Spin Echo, CPMG, and J coupling

- Spin echo vs Fast Spin Echo imaging
- Spin locking
- Decoupling
- References
 - Stables, et al, Analysis of J Coupling-Induced Fat Suppression in DIET Imaging, JMR, **136**, 143–151 (1999)
 - van de Ven, Chp 3.9, 4.9.

Journal of Magnetic Resonance **136**, 143–151 (1999)

Article ID jmre.1998.1628, available online at <http://www.idealibrary.com> on **IDEAL**®



Analysis of J Coupling-Induced Fat Suppression in DIET Imaging

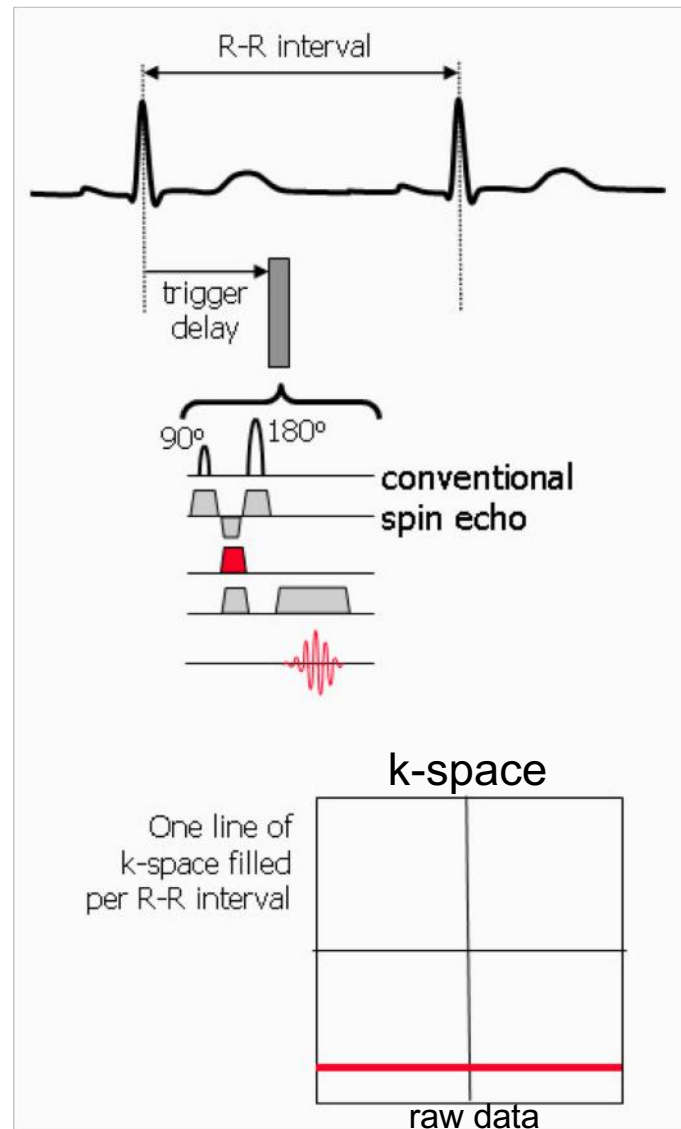
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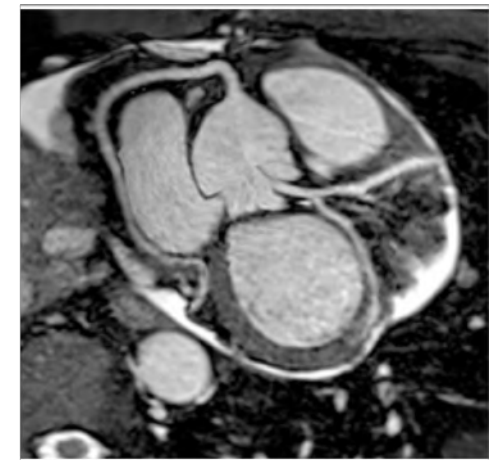
Spin Echo Imaging

- One k-space line collected each TR

Cardiac MRI
example:



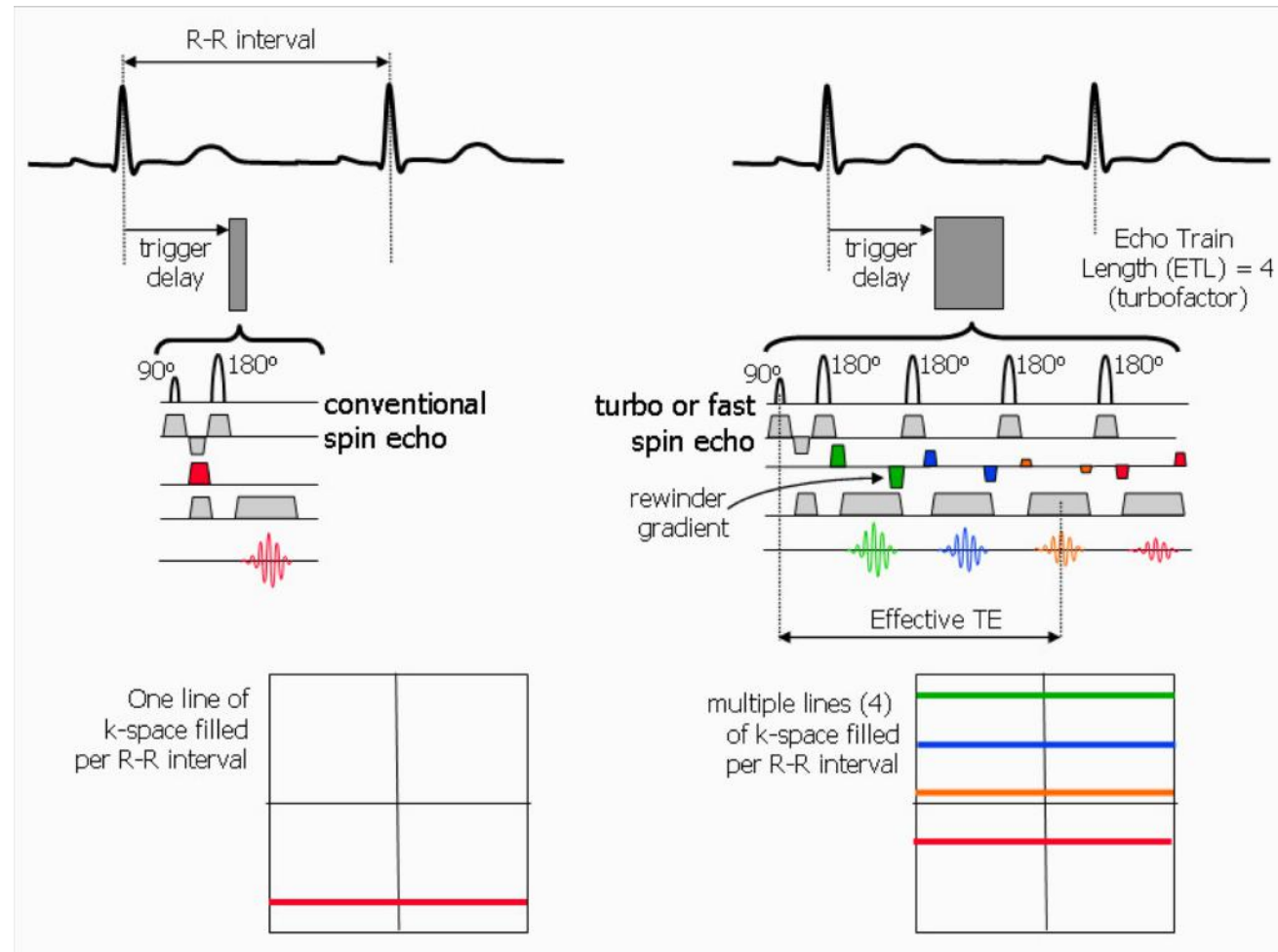
After n lines of data (e.g. 256) are acquired, a FT creates the image



Fast Spin Echo

- A train of 180s is used to acquire multiple k-space lines each TR

Cardiac MRI
example:

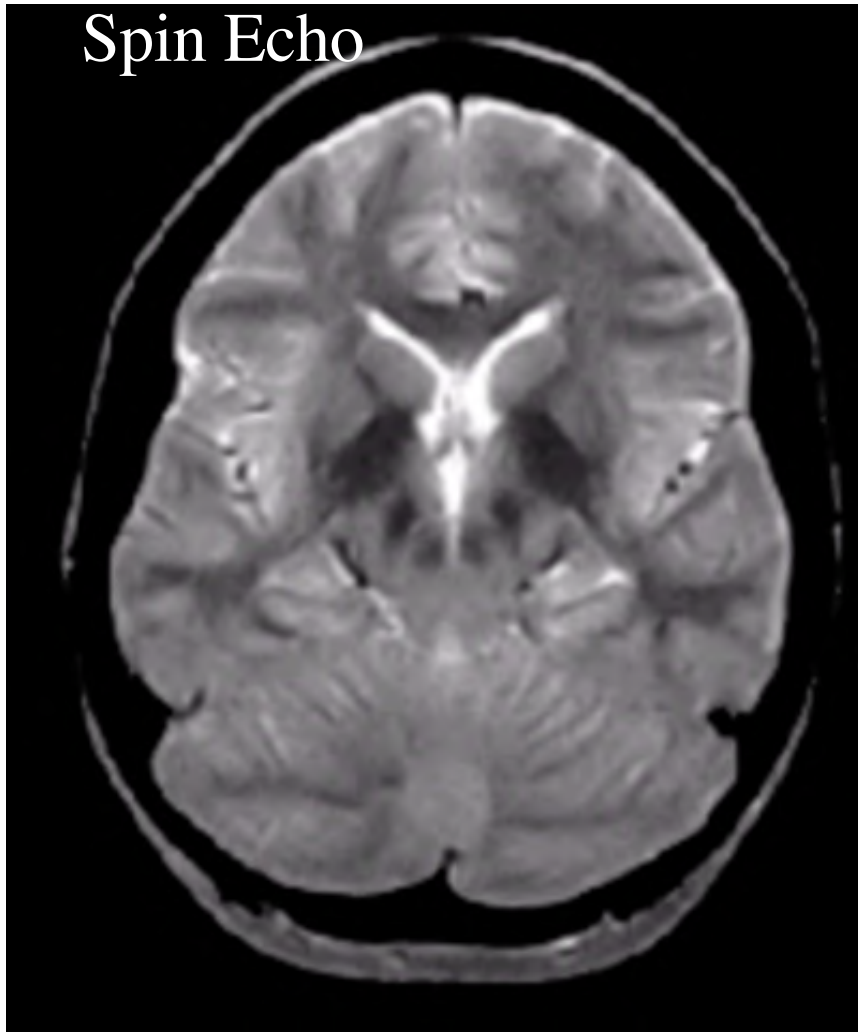


- Provides the ability to acquire images much faster, while retaining T_2 weighting if desired.

FSE Neuro Example

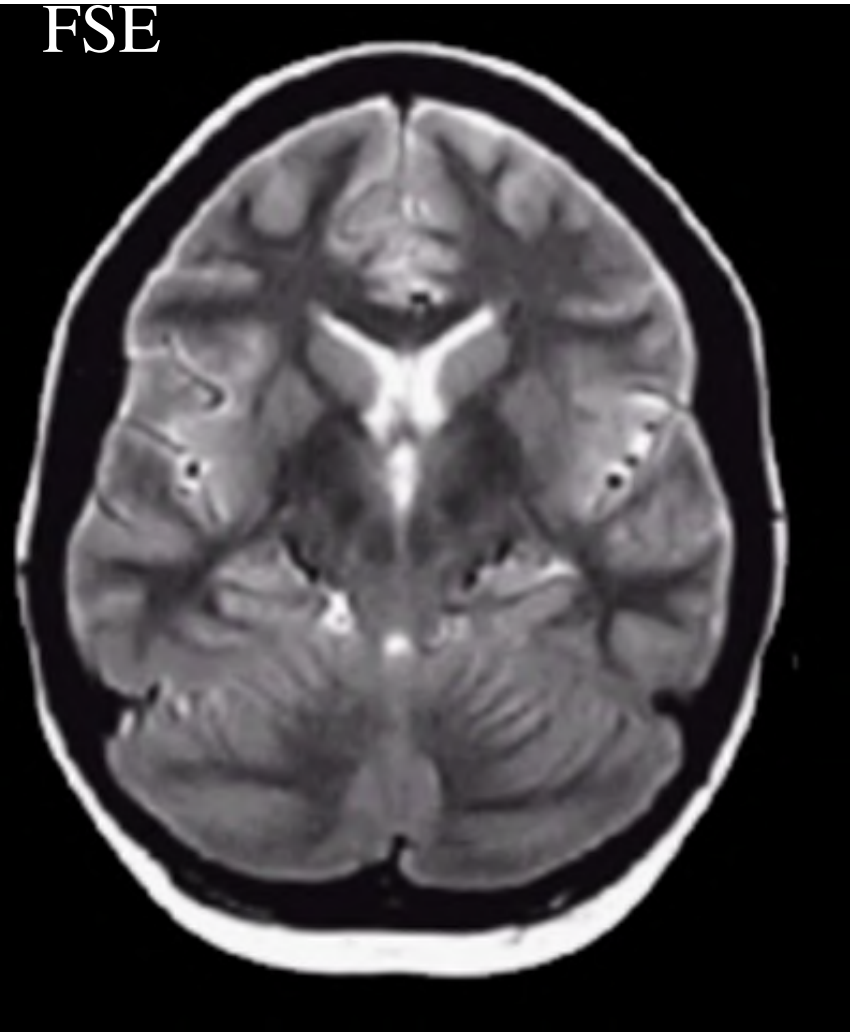
T₂-weighted images in much less time (3T, TE/TR = 80/2000 ms)

Spin Echo



Acq. time \approx 16 min

FSE



Acq. time \approx 1 min
echo train length (ETL) = 16

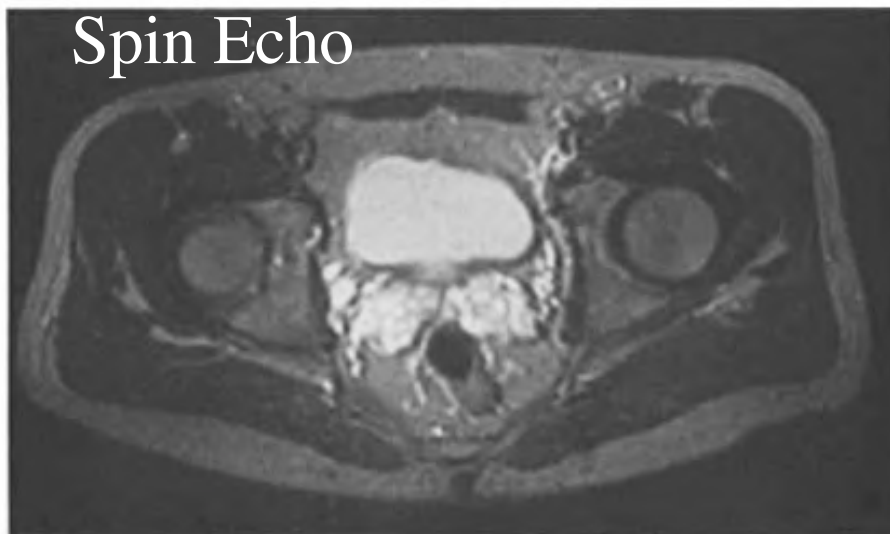
Why is fat bright in FSE images?



Lipid $T_2 = ?$

SE: lipid $T_2 \sim 35$ ms

FSE: lipid $T_2 \sim 135$ ms



Multi-Spin Systems

- Ignoring relaxation, the Hamiltonian has the following general form:

$$\hat{H} = \underset{\substack{\nearrow \\ \text{Static field + J couplings}}}{\hat{H}^0} + \hat{H}_{Rf}$$

- For a system with multiple spins, the total x, y, and z coherences are

$$\hat{F}_x = \sum_j \hat{I}_{xj} \quad \hat{F}_y = \sum_j \hat{I}_{yj} \quad \hat{F}_z = \sum_j \hat{I}_{zj}$$

and the first term of the Hamiltonian can be written as:

$$\hat{H}^0 = \sum_{j=1} \delta_j \hat{I}_{zj} + \sum_{j < k} 2\pi J_{jk} \vec{\hat{I}}_j \cdot \vec{\hat{I}}_k$$

- We wish to consider two cases.
 1. Magnetically equivalent spins, i.e. all δ_j s are equal.
 2. Non-equivalent spins, i.e. unequal δ_j s.

Equivalent Spins

- Multi-spin system for equivalent spins

$$\hat{H}^0 = \sum_j \hat{I}_{zj} + \sum_{j < k} 2\pi J_{jk} \vec{\hat{I}}_j \cdot \vec{\hat{I}}_k = \hat{H}_1 + \hat{H}_2 \quad \text{One can show } \left[\sum_j \hat{I}_{pj}, \vec{\hat{I}}_j \cdot \vec{\hat{I}}_k \right] = 0, \quad p = x, y, z$$

2-spin example

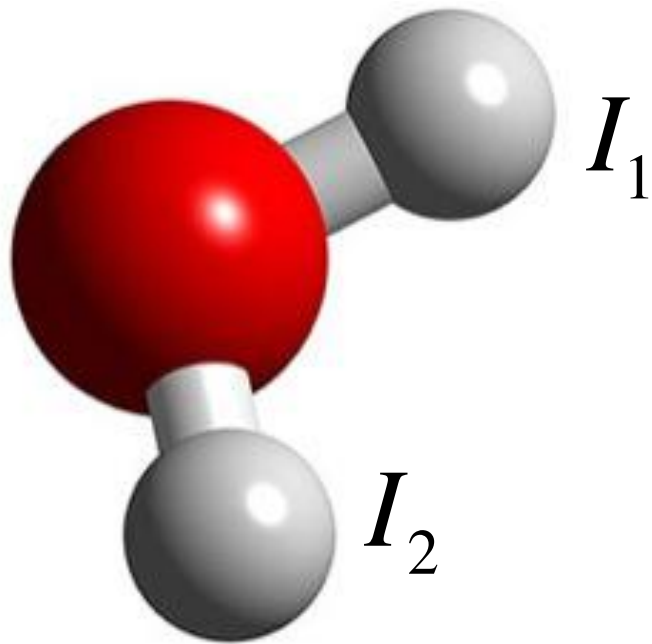
$$\begin{aligned} [\hat{I}_z + \hat{S}_z, \hat{I}_x \hat{S}_x + \hat{I}_y \hat{S}_y + \hat{I}_z \hat{S}_z] &= \hat{I}_y \hat{S}_x - \hat{I}_x \hat{S}_y + \hat{I}_x \hat{S}_y - \hat{I}_y \hat{S}_x = 0 \\ [\hat{I}_x + \hat{S}_x, \hat{I}_x \hat{S}_x + \hat{I}_y \hat{S}_y + \hat{I}_z \hat{S}_z] &= \hat{I}_z \hat{S}_y - \hat{I}_y \hat{S}_z + \hat{I}_y \hat{S}_z - \hat{I}_z \hat{S}_y = 0 \\ [\hat{I}_y + \hat{S}_y, \hat{I}_x \hat{S}_x + \hat{I}_y \hat{S}_y + \hat{I}_z \hat{S}_z] &= \hat{I}_z \hat{S}_x - \hat{I}_x \hat{S}_z + \hat{I}_x \hat{S}_z - \hat{I}_z \hat{S}_x = 0 \end{aligned}$$

- Theorem: Let $\hat{H} = \hat{H}_1 + \hat{H}_2$ and $[\hat{H}_1, \hat{H}_2] = [\hat{F}_x, \hat{H}_2] = [\hat{F}_y, \hat{H}_2] = [\hat{F}_z, \hat{H}_2] = 0$
then the observed signal is independent of \hat{H}_2

Proof: $\text{Tr}[\hat{F}_p e^{-i\hat{H}t} \sigma(0) e^{i\hat{H}t}] = \text{Tr}[\hat{F}_p e^{-i(\hat{H}_1 + \hat{H}_2)t} \sigma(0) e^{i(\hat{H}_1 + \hat{H}_2)t}]$

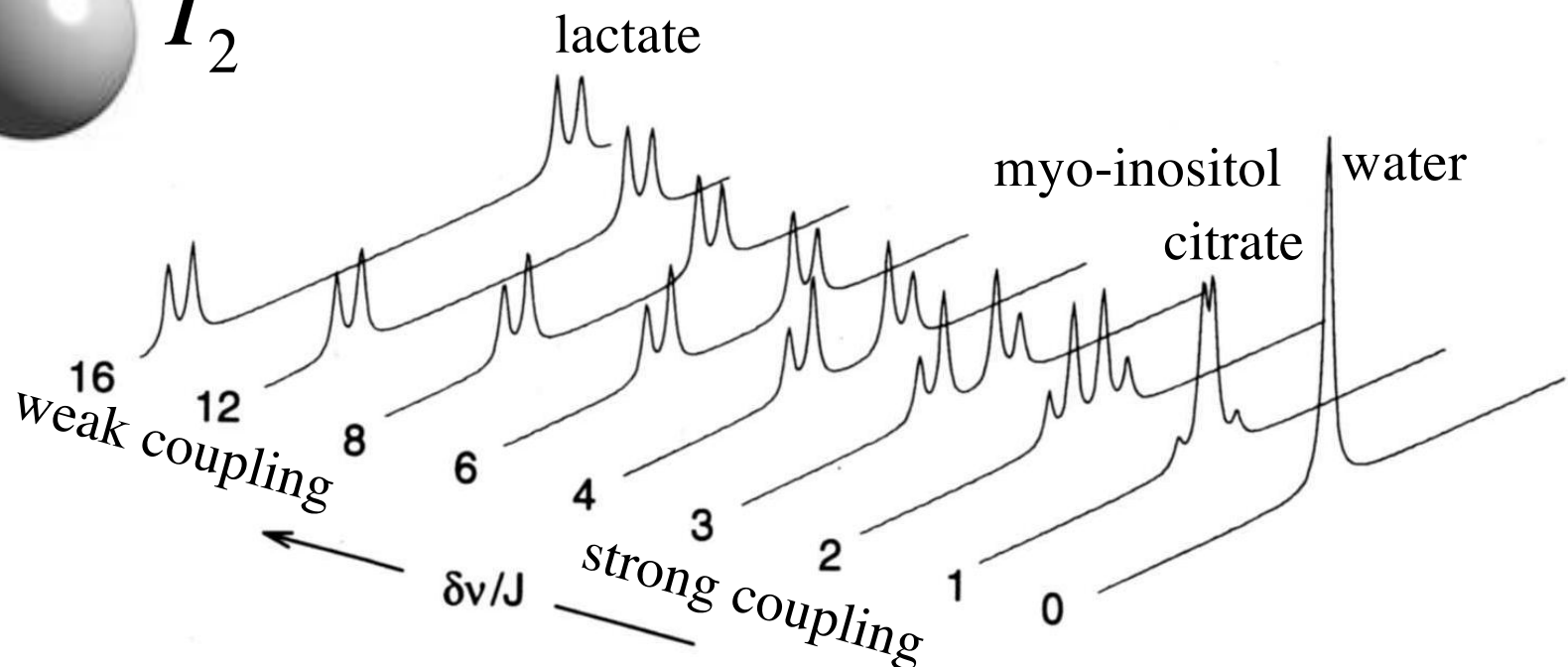
$$\begin{aligned} &= \text{Tr}[\hat{F}_p e^{-i\hat{H}_2t} e^{-i\hat{H}_1t} \sigma(0) e^{i\hat{H}_1t} e^{i\hat{H}_2t}] \\ &= \text{Tr}[e^{i\hat{H}_2t} \hat{F}_p e^{-i\hat{H}_2t} e^{-i\hat{H}_1t} \sigma(0) e^{i\hat{H}_1t}] \quad \leftarrow \text{Why is this step legitimate?} \\ &= \text{Tr}[\hat{F}_p e^{i\hat{H}_2t} e^{-i\hat{H}_2t} e^{-i\hat{H}_1t} \sigma(0) e^{i\hat{H}_1t}] \\ &= \text{Tr}[\hat{F}_p e^{-i\hat{H}_1t} \sigma(0) e^{i\hat{H}_1t}] \quad \text{Independent of } \hat{H}_2! \end{aligned}$$

Water and J-coupling



$$\hat{H}^0 = \Omega \hat{I}_{z1} + \Omega \hat{I}_{z2} + 2\pi J \vec{\hat{I}}_1 \cdot \vec{\hat{I}}_2$$

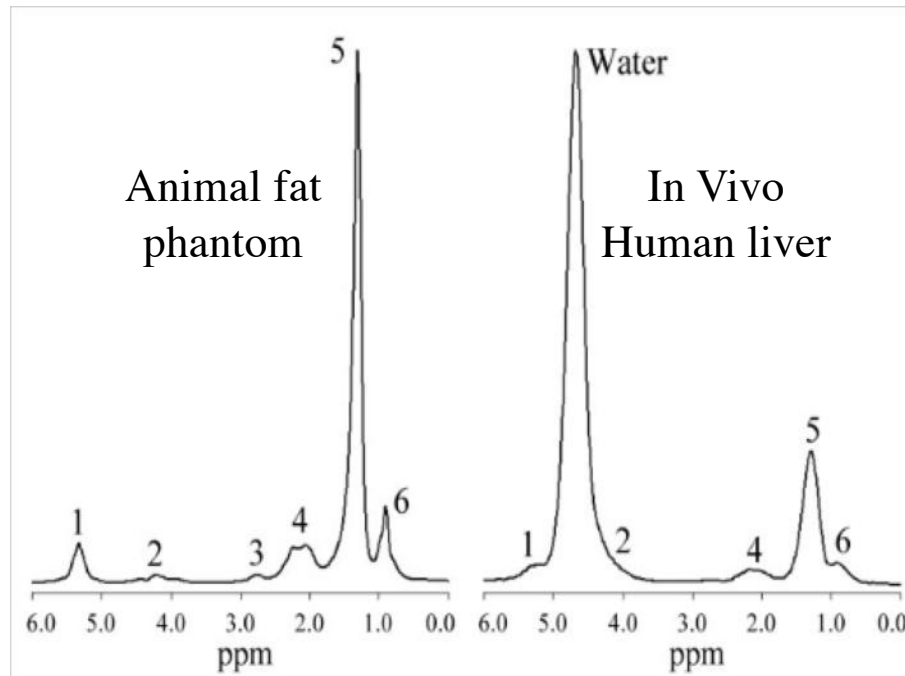
$$\left[\hat{I}_{z1} + \hat{I}_{z2}, \vec{\hat{I}}_1 \cdot \vec{\hat{I}}_2 \right] = 0$$



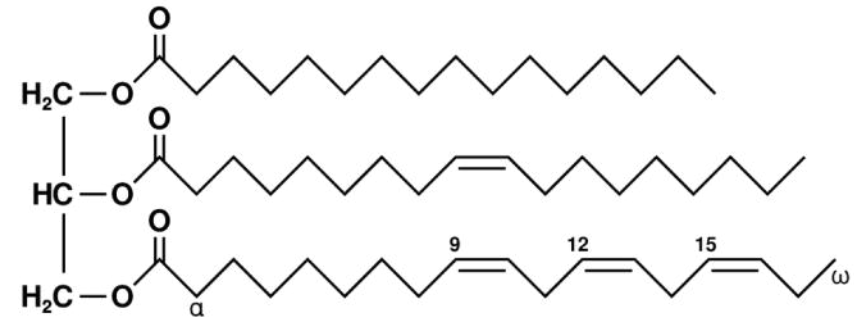
The two water ^1H spins are equivalent, hence show no effects due to J-coupling. 9

Lipids

- Lipids consist of multiple J-coupled resonances, and lipid ^1H s are not equivalent!



Peak	Location	Assignment	Peak type
1	5.30 ppm	$-\text{CH}=\text{CH}-$	Multiplet
	5.19 ppm	$-\text{CH}-\text{O}-\text{CO}-\text{R}$	Multiplet
2	4.20 ppm	$-\text{CH}_2-\text{O}-\text{CO}-\text{R}$	Multiplet
3	2.75 ppm	$-\text{CH}=\text{CH}-\text{CH}_2-\text{CH}=\text{CH}-$	Multiplet
4	2.20 ppm	$-\text{CO}-\text{CH}_2-\text{CH}_2-$	Multiplet
	2.02 ppm	$-\text{CH}_2-\text{CH}=\text{CH}-\text{CH}_2-$	Multiplet
5	1.60 ppm	$-\text{CO}-\text{CH}_2-\text{CH}_2-$	Multiplet
	1.30 ppm	$-(\text{CH}_2)_n-$	Multiplet
6	0.90 ppm	$-(\text{CH}_2)_n-\text{CH}_3$	Triplet



$$\hat{H}^0 = \sum_j \delta_j \hat{I}_{zj} + \sum_{j < k} 2\pi J_{jk} \vec{\hat{I}}_j \cdot \vec{\hat{I}}_k$$

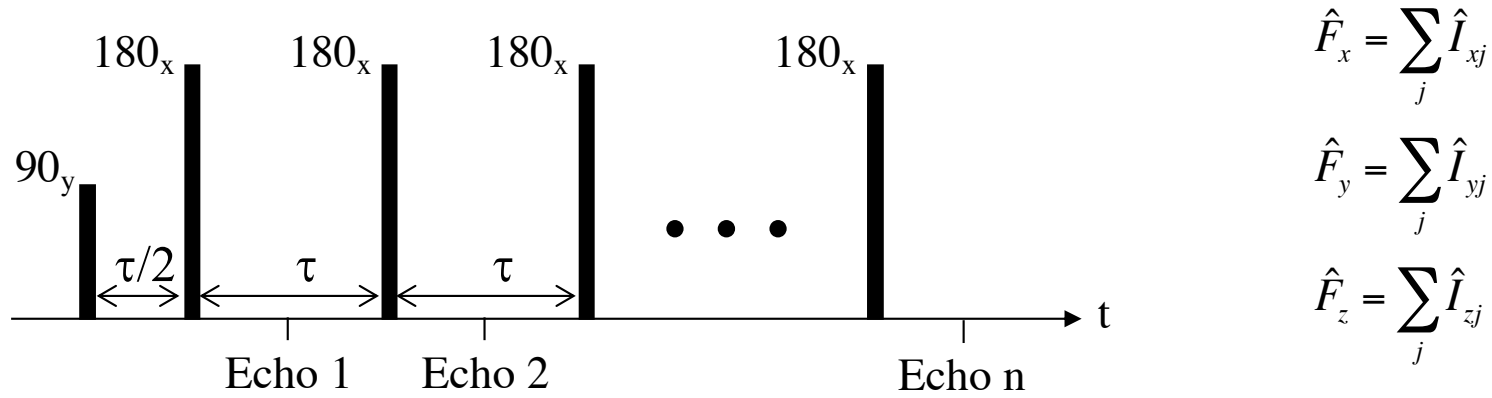
δ_j s not all equal, thus

$$\left[\sum_j \delta_j \hat{I}_{pj}, \vec{\hat{I}}_j \cdot \vec{\hat{I}}_k \right] \neq 0$$

Hence J-coupling can NOT be ignored.

Carr-Purcell-Meiboom-Gill (CPMG)

- Consider a multi-spin system with j spins, coupling constants J_{jk} , rotating frame resonance frequencies δ_j , and the following pulse sequence:



- Assuming hard RF pulses with $\omega_1 \gg \delta_j, J_{jk}$ for all j , find the Hamiltonian for each time interval.
 - 90_y Rf pulse: $\hat{H}_1 = \omega_1 \sum_j \hat{I}_{yj} = \omega_1 \hat{F}_y$ where $\omega_1 = -\gamma B_1 \gg \delta_j, J_{jk}$ for all j, k .
 - 180_x Rf pulses: $\hat{H}_2 = \omega_1 \hat{F}_x$
 - between Rf pulses: $\hat{H}_3 = \sum_j \delta_j \hat{I}_{zj} + \sum_{j < k} 2\pi J_{jk} \vec{\hat{I}}_j \cdot \vec{\hat{I}}_k$

CPMG Product Operator Analysis

- Ignoring relaxation...

At thermal equilibrium: $\hat{\sigma}_0 \propto \hat{F}_z$

After the 90_y : $\hat{\sigma} \propto \hat{F}_x$

Before the first 180_x : $\hat{\sigma} \propto e^{-\frac{i}{2}\hat{H}_3\tau} \hat{F}_x e^{\frac{i}{2}\hat{H}_3\tau}$

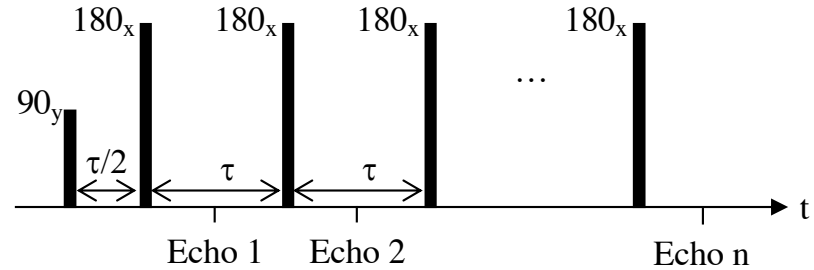
After the first 180_x : $\hat{\sigma} \propto e^{-i\pi\hat{F}_x} e^{-\frac{i}{2}\hat{H}_3\tau} \hat{F}_x e^{\frac{i}{2}\hat{H}_3\tau} e^{i\pi\hat{F}_x}$

$$\hat{\sigma} \propto \hat{B}\hat{A}\hat{F}_x\hat{A}^{-1}\hat{B}^{-1} \quad \text{where} \quad \hat{A} = e^{-\frac{i}{2}\hat{H}_3\tau} \quad \text{and} \quad \hat{B} = e^{-i\pi\hat{F}_x}$$

At Echo 1: $\hat{\sigma}_1 \propto (\hat{A}\hat{B}\hat{A})\hat{F}_x(\hat{A}\hat{B}\hat{A})^{-1}$

Continuing, the spin density at the n th echo will be:

$$\hat{\sigma}_n \propto (\hat{A}\hat{B}\hat{A})^n \hat{F}_x (\hat{A}\hat{B}\hat{A})^{-n}$$



CPMG Product Operator Analysis

- Let examine $\hat{\sigma}_n \propto (\hat{A}\hat{B}\hat{A})^n \hat{F}_x (\hat{A}\hat{B}\hat{A})^{-n}$ more closely for the case where τ is short, i.e.
 $|J_{jk}\tau|, |\delta_j\tau| \ll 1$ for all spin groups j, k .

- Expanding \hat{A} to first order in a Taylor series yields

$$\hat{A} = e^{-\frac{i}{2}\hat{H}_3\tau} \approx 1 - \frac{i}{2}\hat{H}_3\tau = 1 - \frac{i}{2} \sum_j \delta_j \hat{I}_{zj} \tau - \frac{i}{2} \sum_{j < k} 2\pi J_{jk} \vec{\hat{I}}_j \cdot \vec{\hat{I}}_k \tau$$

$$\text{Substituting... } \hat{A}\hat{B}\hat{A} = e^{-\frac{i}{2}\hat{H}_3\tau} e^{-i\pi\hat{F}_x} e^{-\frac{i}{2}\hat{H}_3\tau} \xrightarrow{\text{short } \tau} \left(1 + i \sum_{j < k} 2\pi J_{jk} \vec{\hat{I}}_j \cdot \vec{\hat{I}}_k \tau \right) e^{-i\pi\hat{F}_x}$$

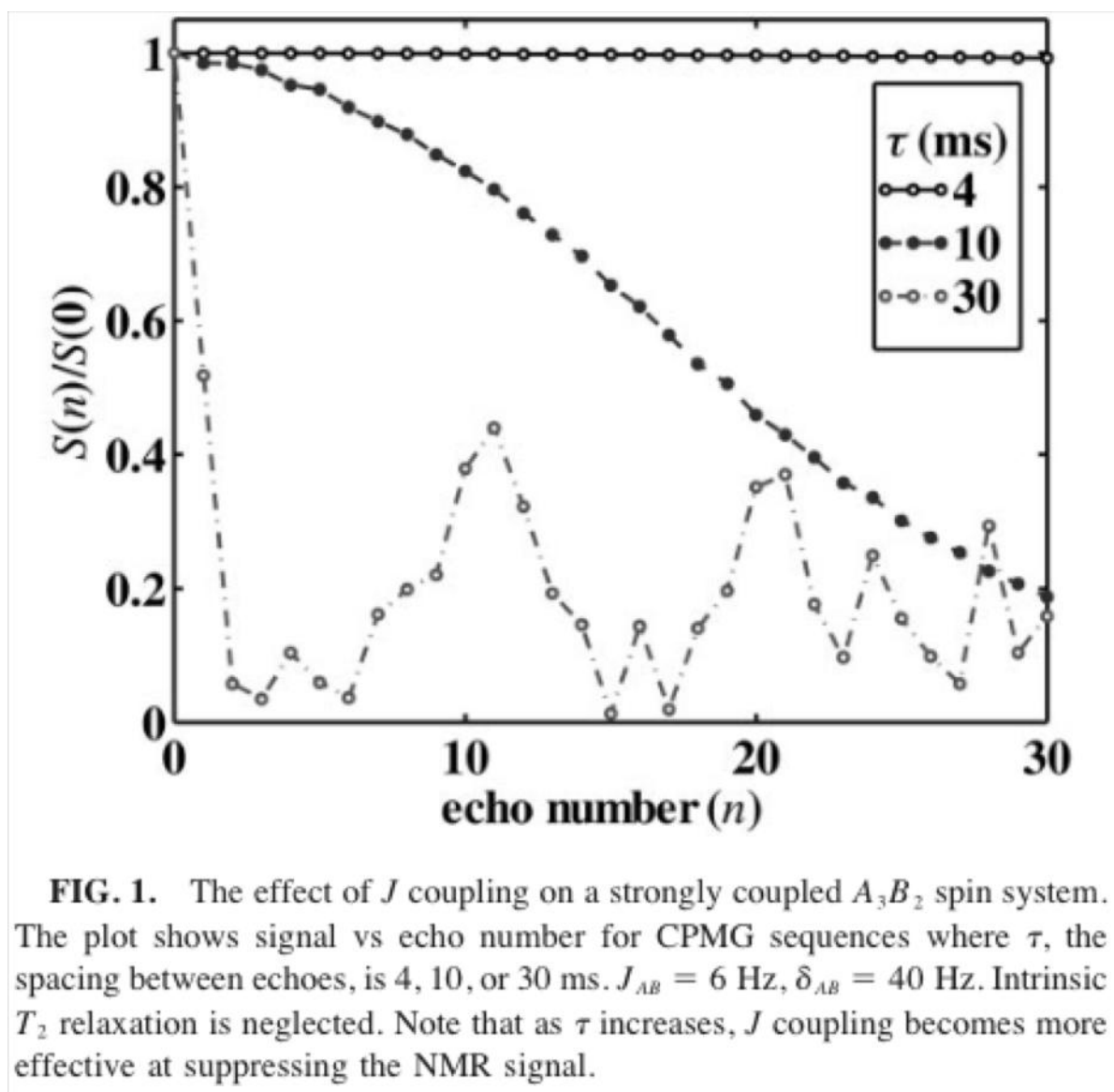
However, \hat{F}_x and $e^{-i\pi\hat{F}_x}$ commute...

$$\text{Therefore, in a CPMG sequence with } |J_{jk}\tau|, |\delta_j\tau| \ll 1, \quad \hat{\sigma}_n \xrightarrow{\text{short } \tau} \hat{F}_x$$

Independent of n ,
 δ_j , and J_{jk}

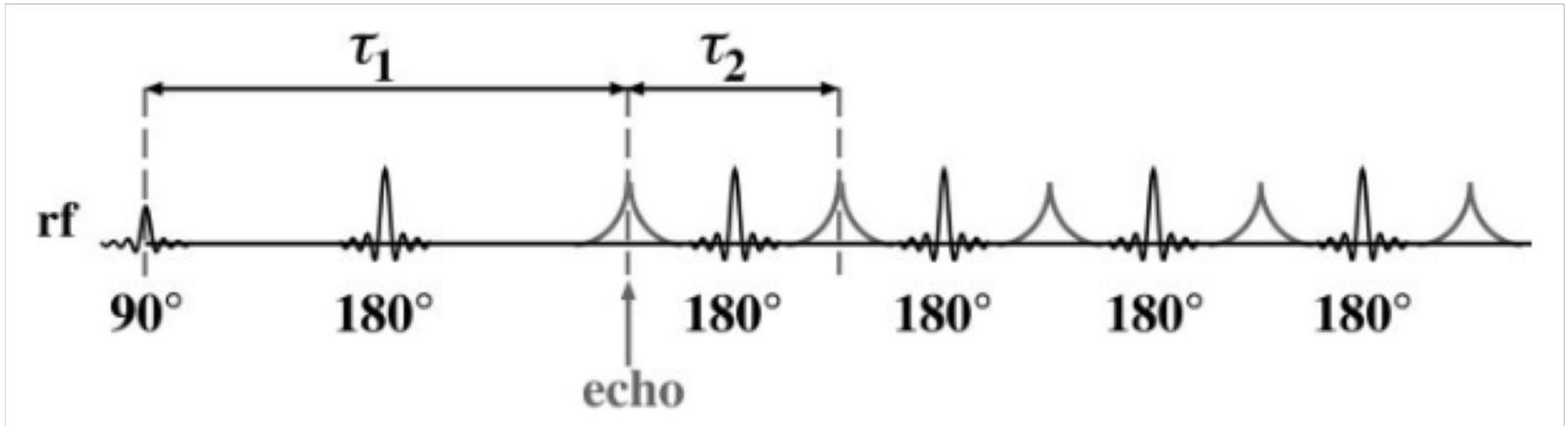
- Therefore, for this rapidly refocused CPMG sequence, lipids will decay with their true T_2 s free from the additional dephasing due to multiple J-couplings.

FSE Simulations



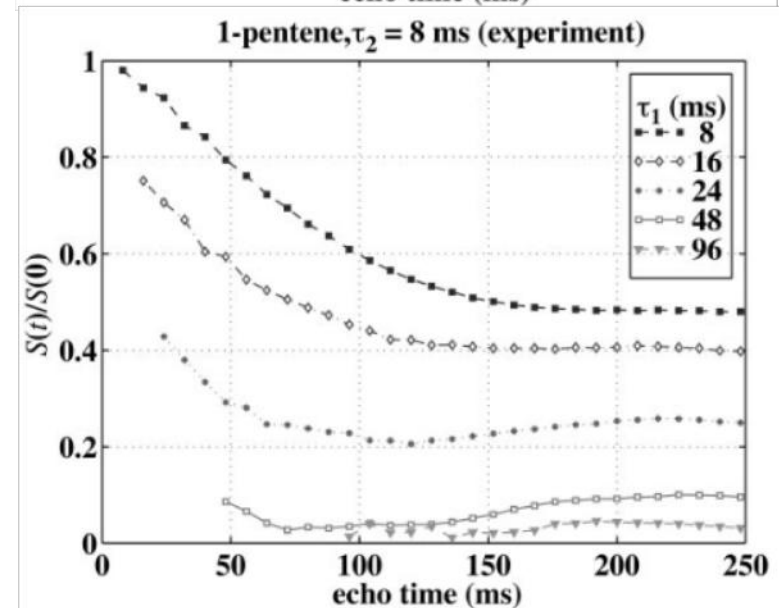
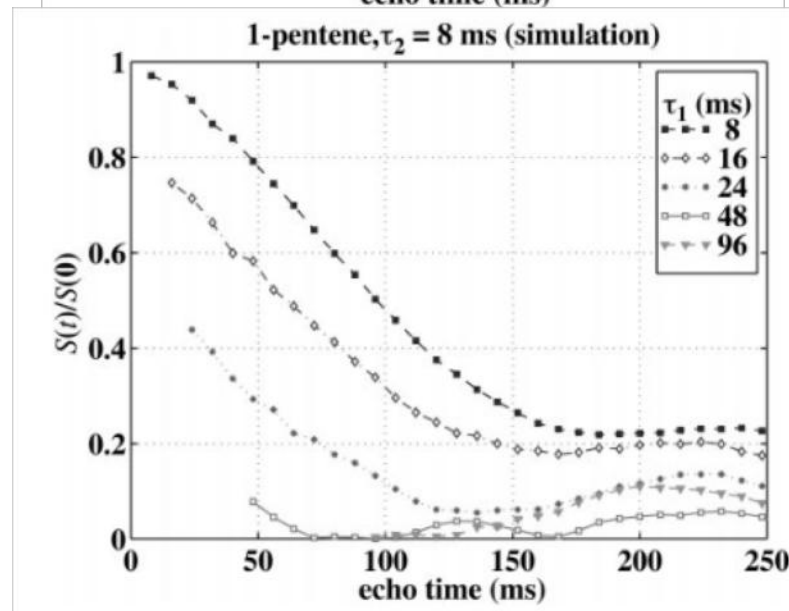
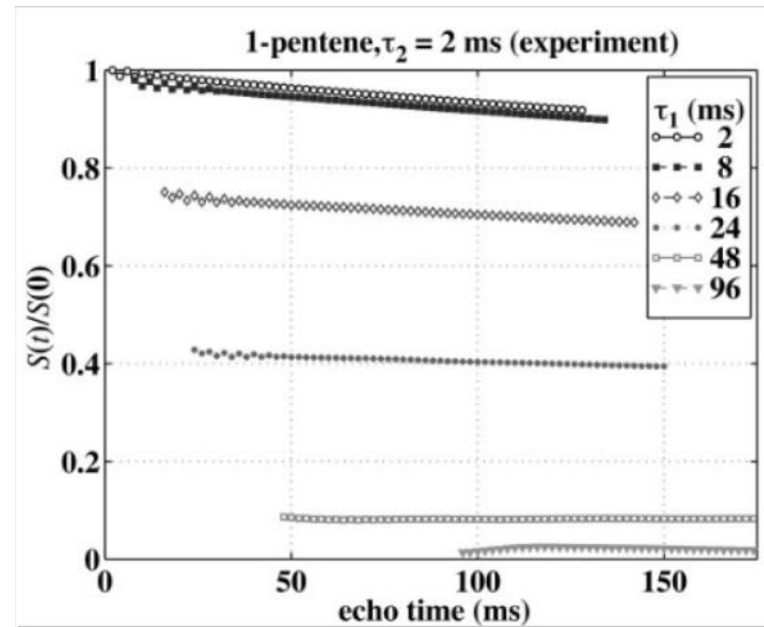
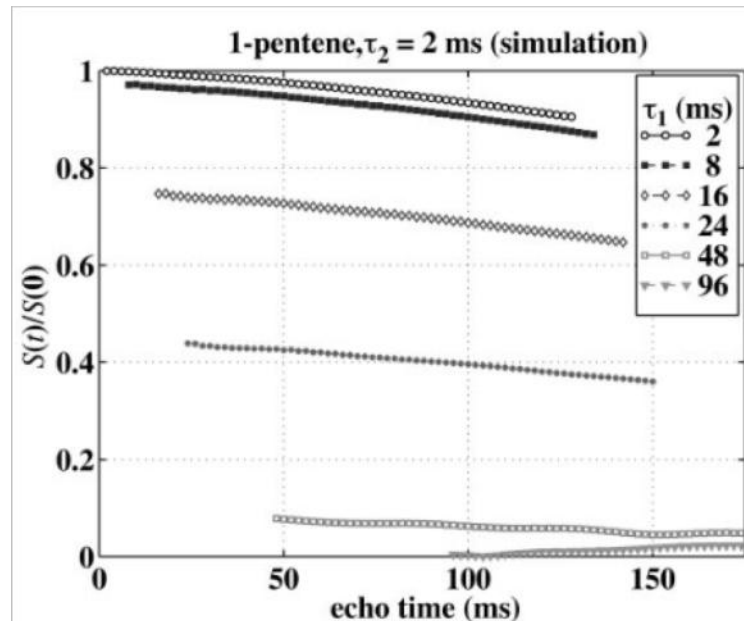
Dual Interval Echo Train (DIET) FSE

- Goal: lipid suppression



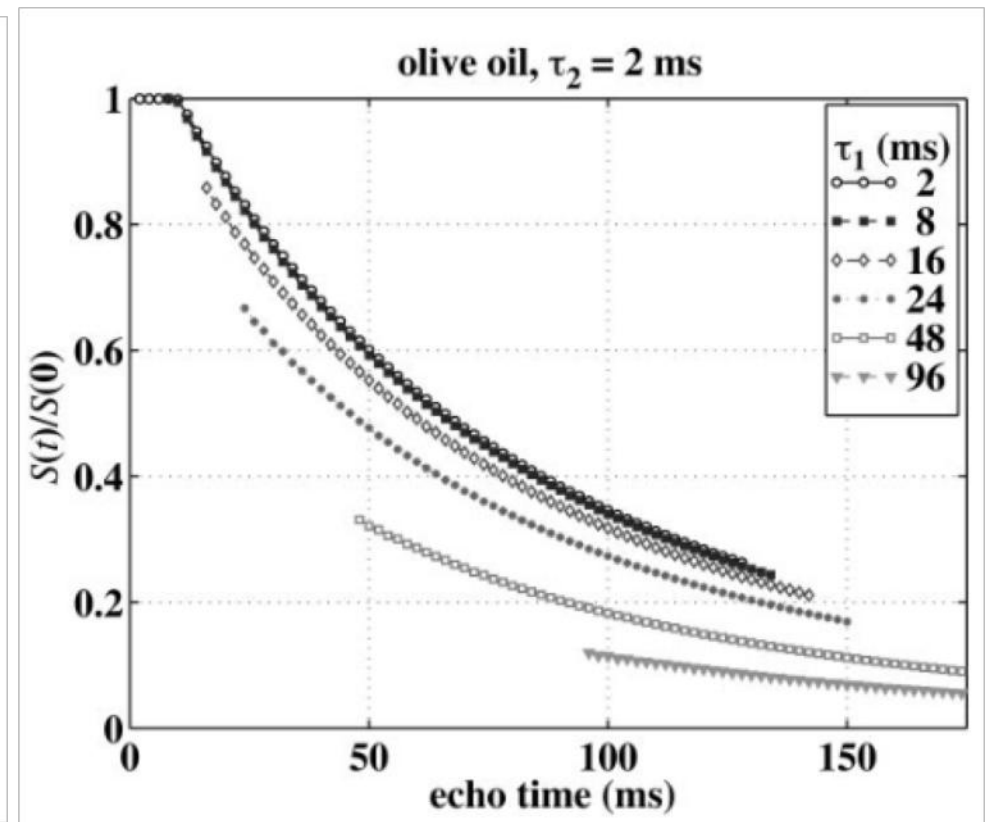
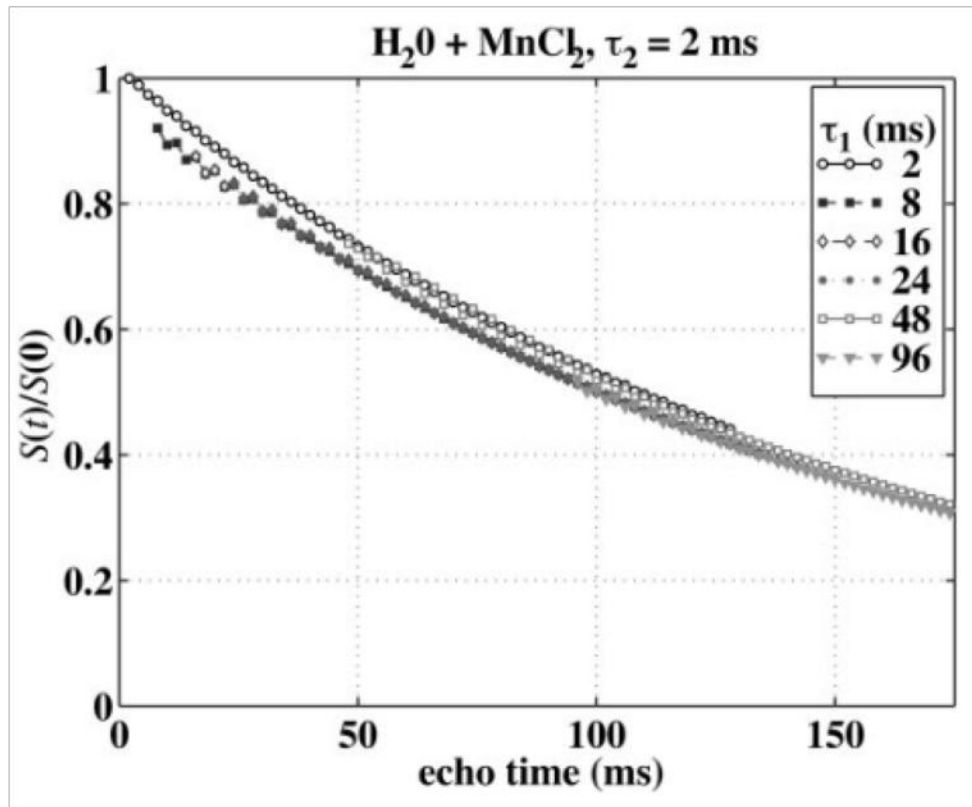
- The relatively long first echo allows for J-coupling-induced dephasing to occur.
- This signal loss is not recovered in the remaining echo train.

DIET Results

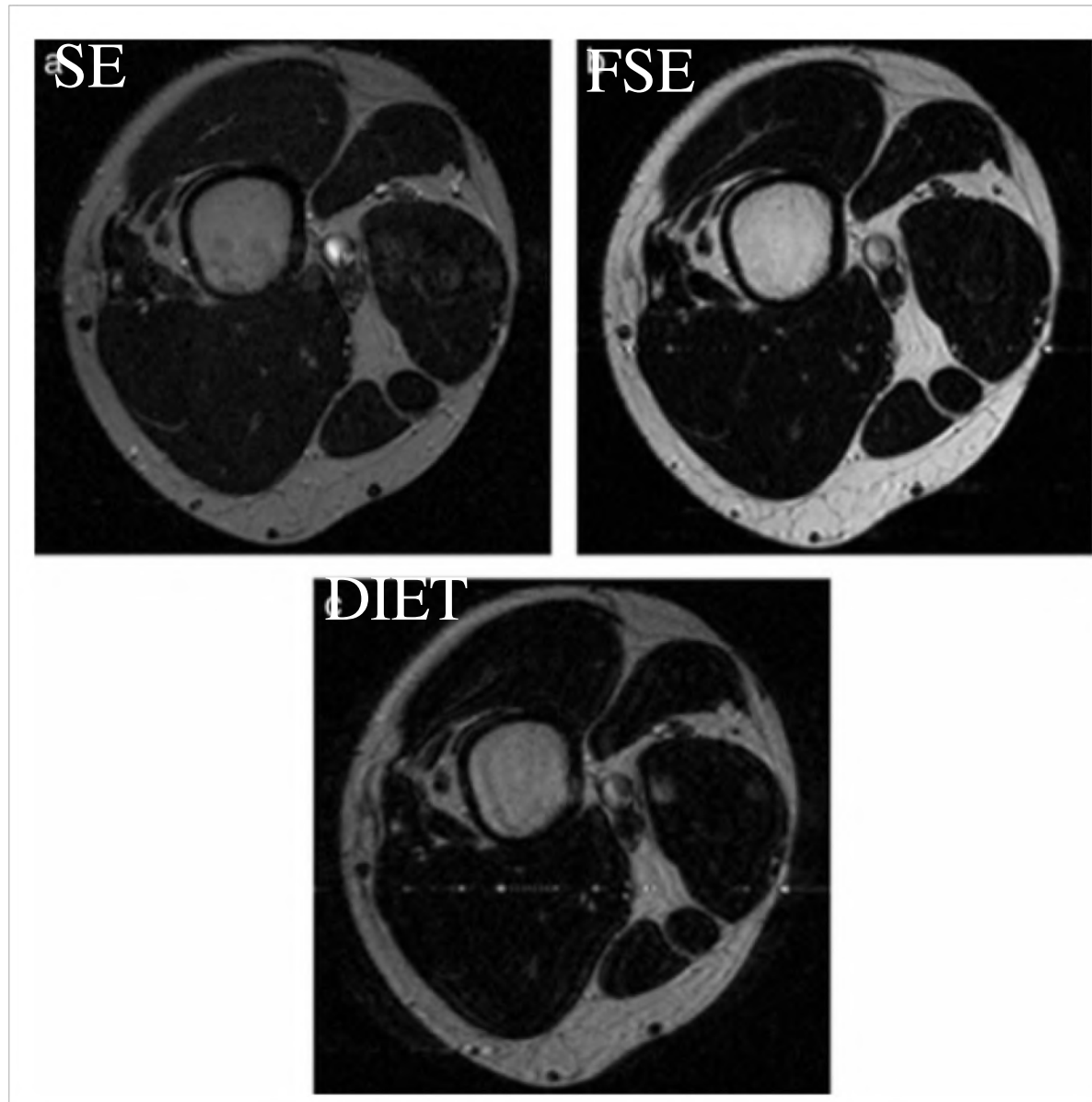


Stables, et al, JMR, **136**, 143–151 (1999)

DIET Results



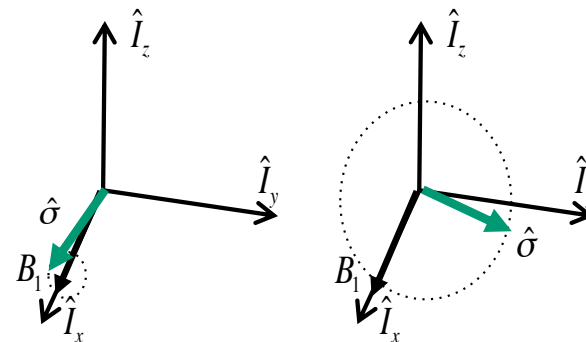
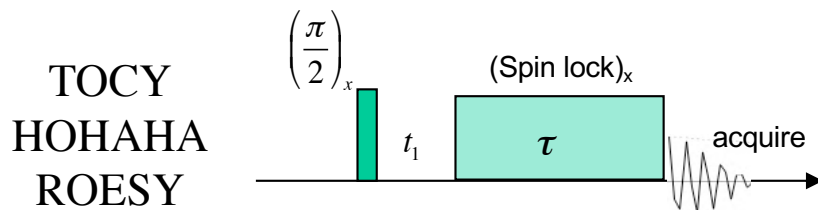
Results



Stables, et al, JMR, **136**, 143–151 (1999)

Spin Locking

- The suppression of J-coupling during a FSE sequence is considered a nuisance. However, the effect can also have advantages.
- Spin locking: the application of a long, strong continuous Rf pulse along a specified axis, e.g. x, in the rotating frame.
 - Chemical shift is suppressed, and spins are rendered effectively equivalent
 - Coherences along x are retained, those along y are dephased due to Rf inhomogeneity



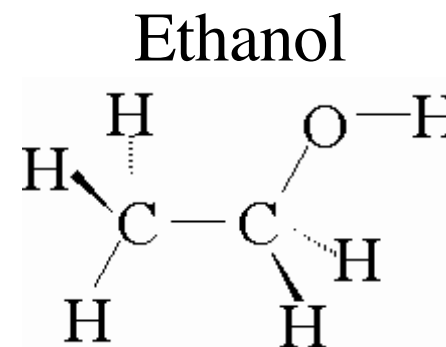
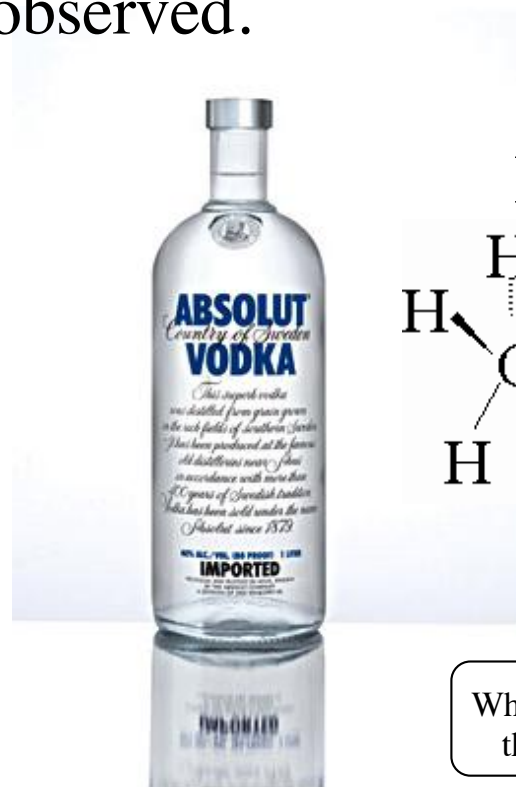
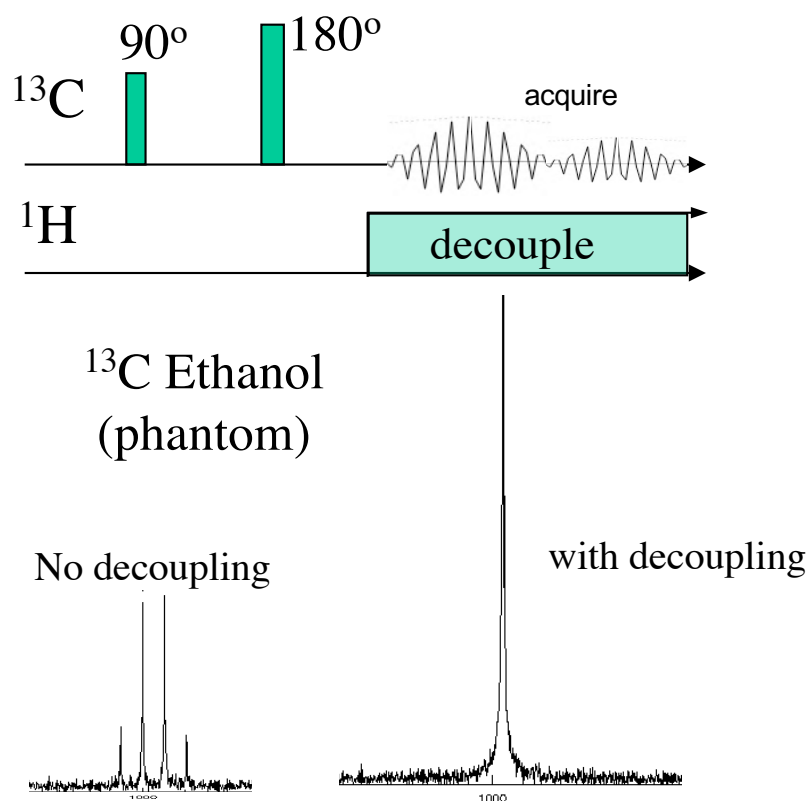
$$\hat{I}_x + \hat{S}_x \xrightarrow{\text{spin lock}} \hat{I}_x + \hat{S}_x \quad \hat{I}_x - \hat{S}_x \xrightarrow{\text{spin lock}} (\hat{I}_x - \hat{S}_x) \cos 2\pi J\tau + (2\hat{I}_y \hat{S}_z - 2\hat{I}_z \hat{S}_y) \sin 2\pi J\tau$$

- The observable magnetization for truly equivalent spins does not evolve under J coupling, however spins rendered temporarily equivalent can show much more complex behavior, as they can enter the spin-lock period in a variety of initial states.

Decoupling

- Line splitting reduces the ability to detect and quantify in vivo peaks.
- Decoupling involves the use of a long strong Rf pulse on the coupled partner of the spin being observed.

Example



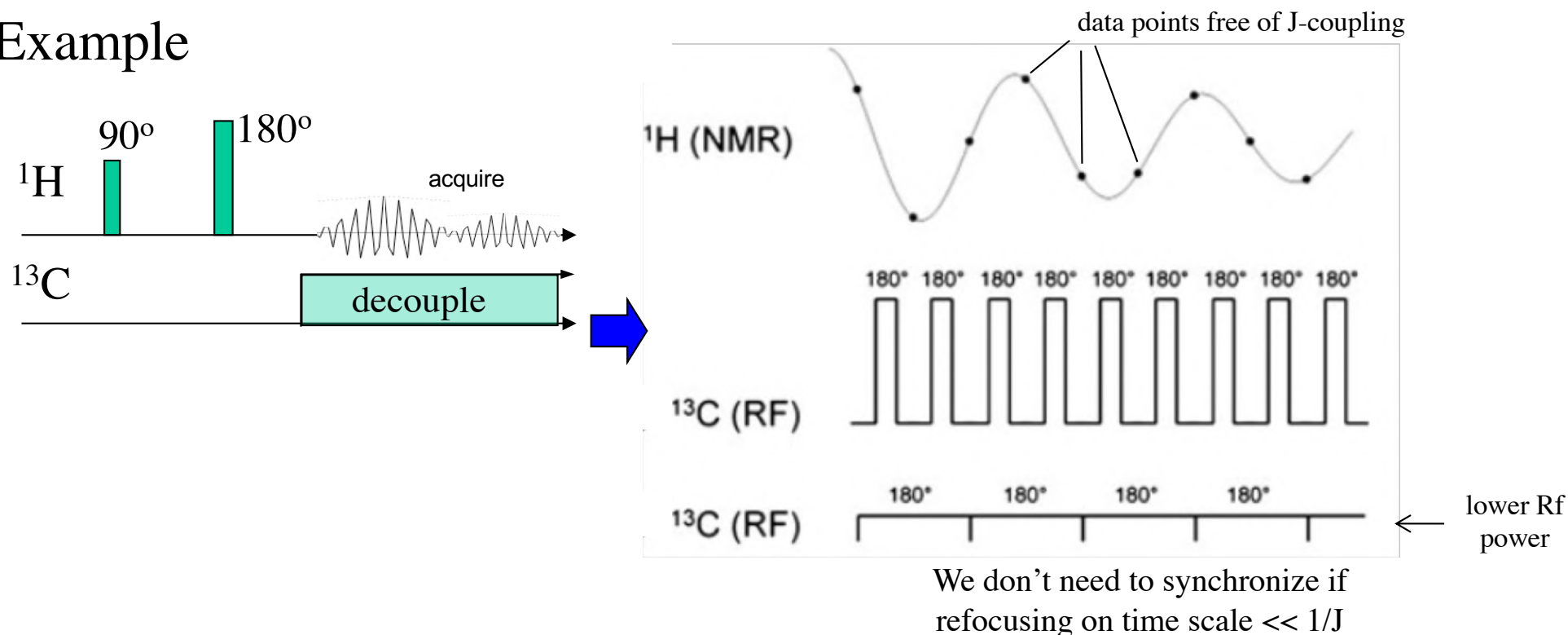
Which carbon has the ^{13}C label?



Decoupling

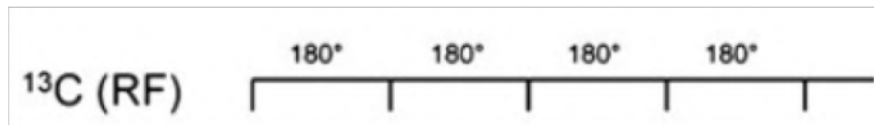
- Hamiltonian: $\hat{H} = \hat{H}_{Zeeman} + \hat{H}_{dipole} + 2\pi \underbrace{J\hat{I} \cdot \hat{S}}$
 time average=0 with “decouple” using RF pulses
 rapid molecular tumbling to rapidly flip the S spin

Example

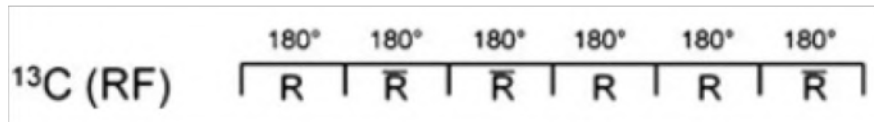


Broadband Decoupling

- Problem: long Rf pulses have narrow bandwidths.



- Phase cycling the 180° pulses improves off-resonance behavior.



- Composite 180 pulses are even better:

$$R = 90_x^\circ 180_{-x}^\circ 270_x^\circ = 1\bar{2}3$$

...and are typically used in supercycles.

$$1\bar{2}3 \ 1\bar{2}3 \ \bar{1}2\bar{3} \ \bar{1}2\bar{3}$$

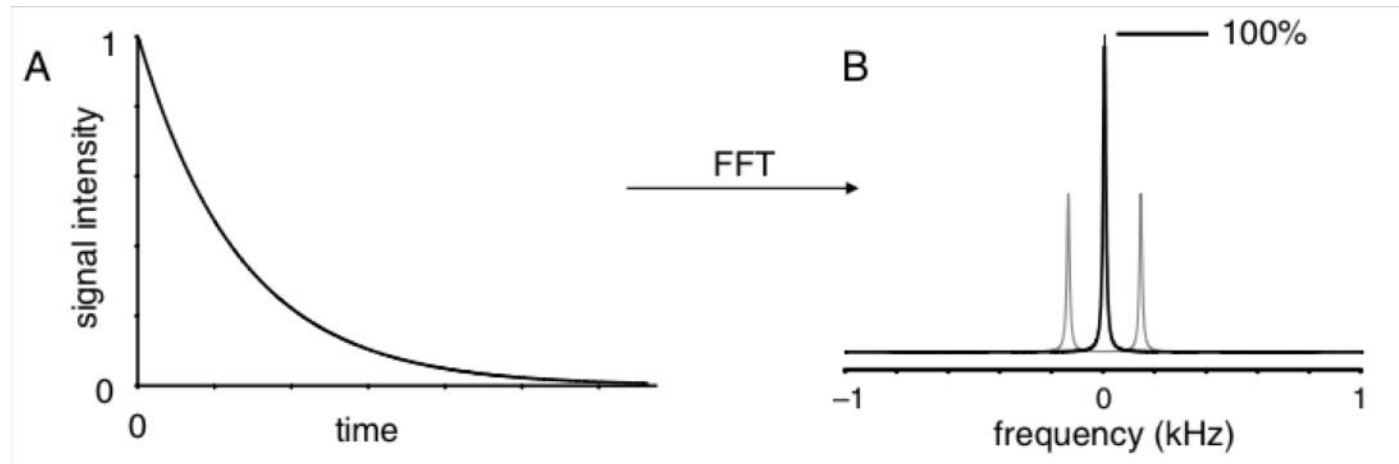
Wideband Alternating Phase Low-power
Technique for Zero-residue Splitting

WALTZ

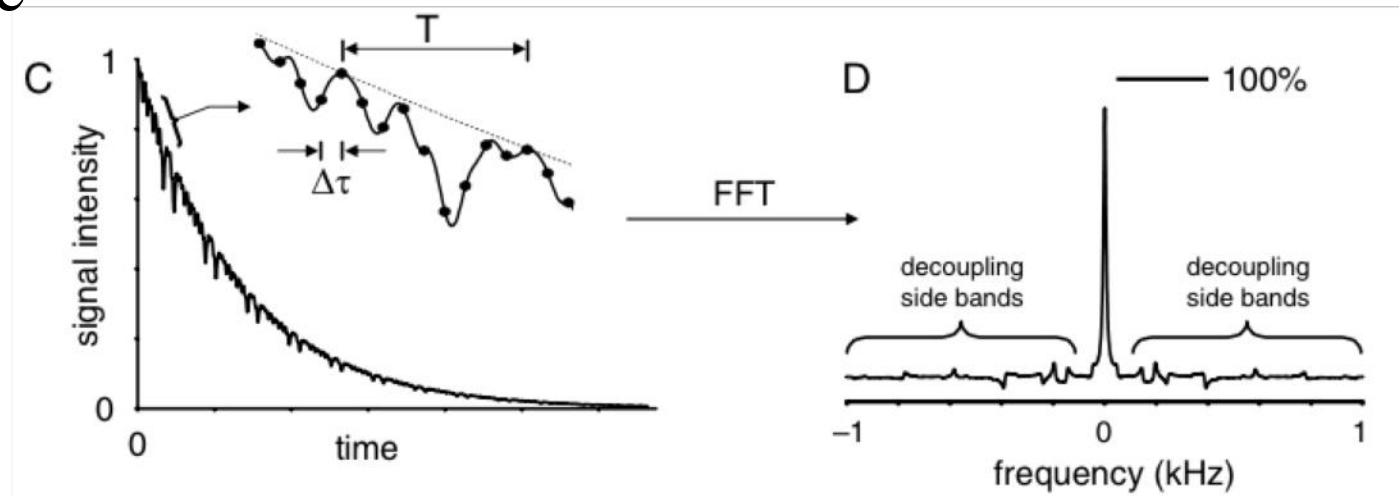


Decoupling in practice

- Theory

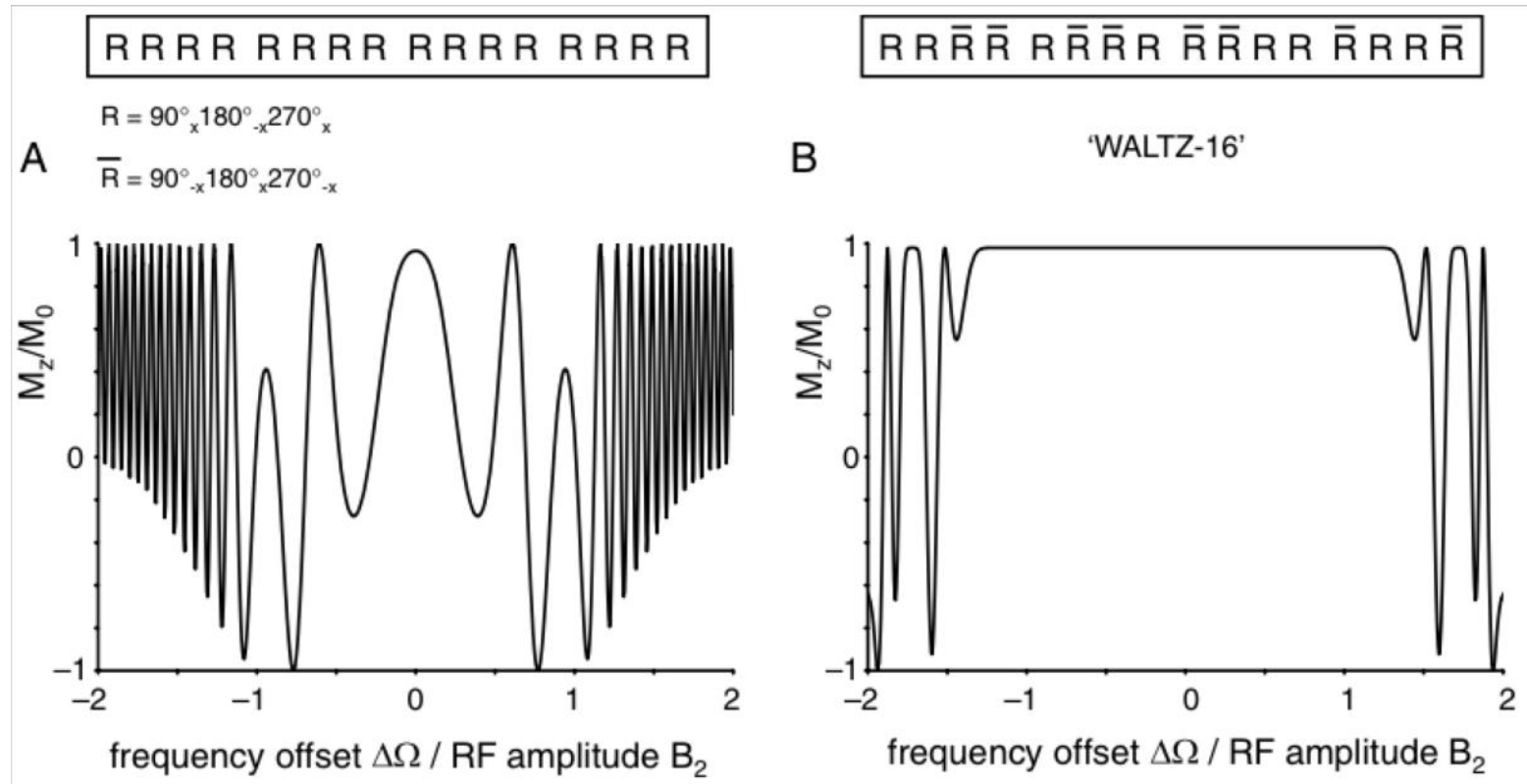


- In practice

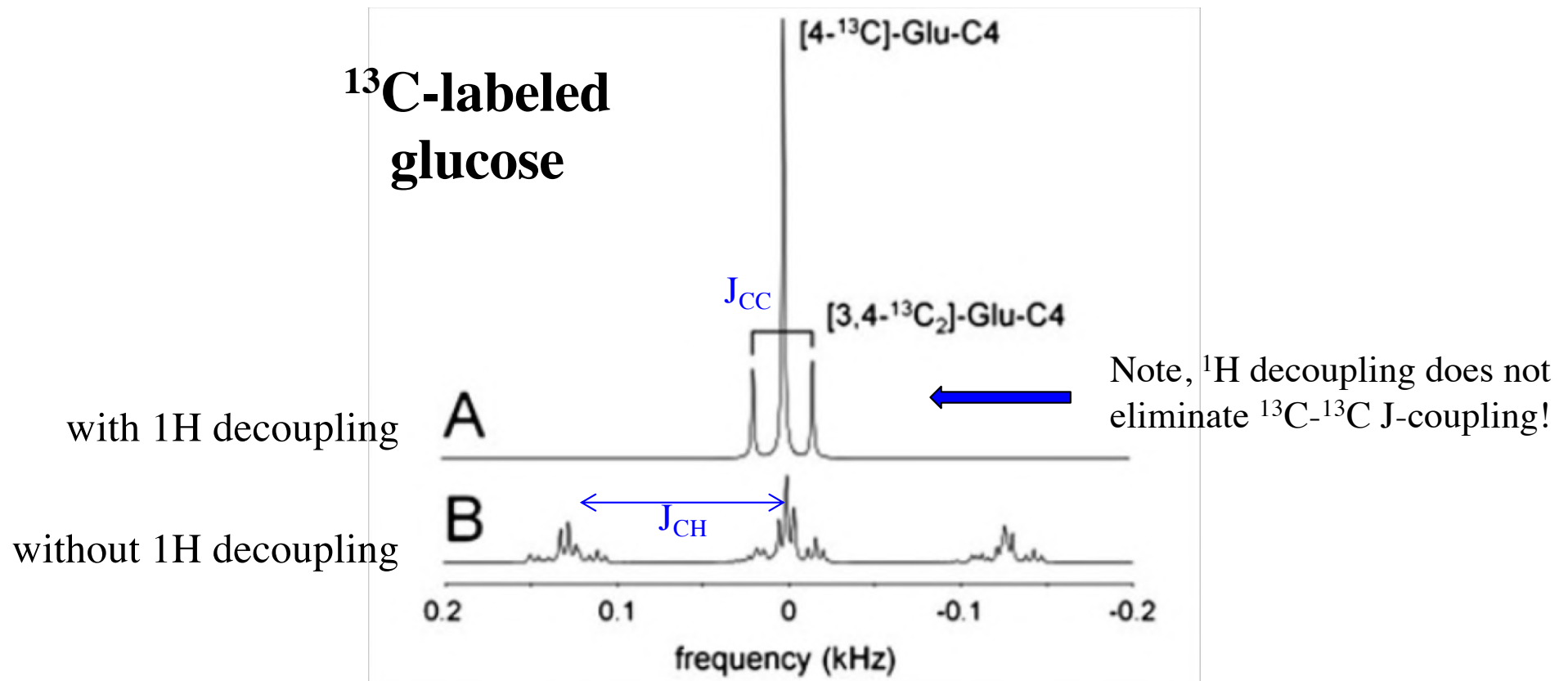


- Rf power deposition, typically measured as Specific Absorption Rate (SAR), is usually the limiting factor.

WALTZ Decoupling

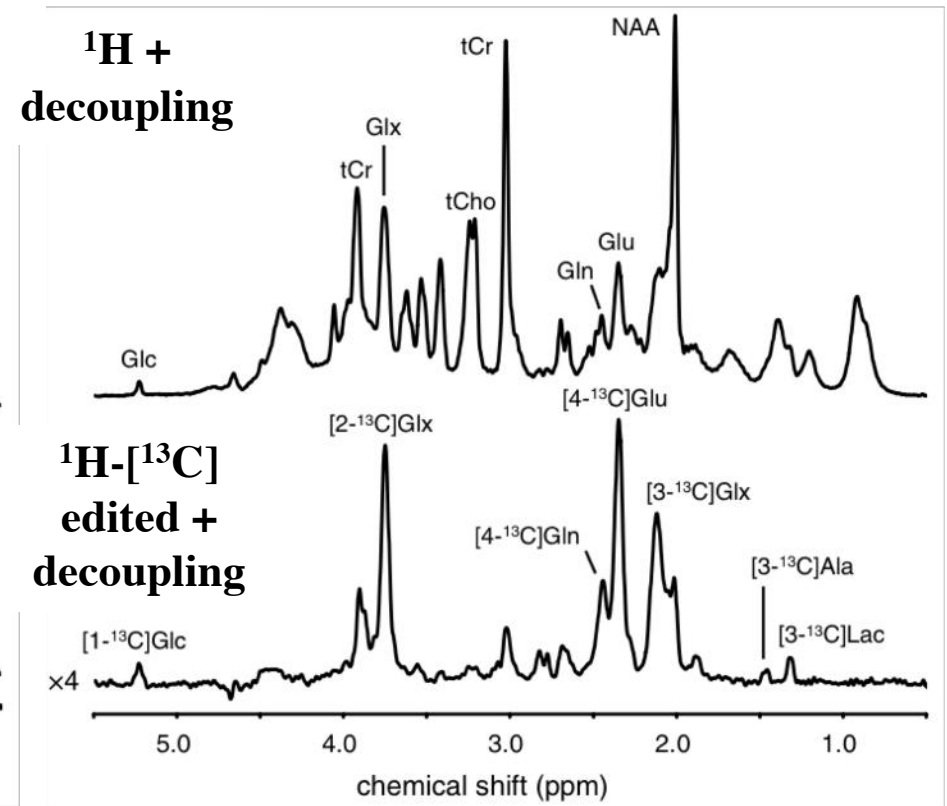
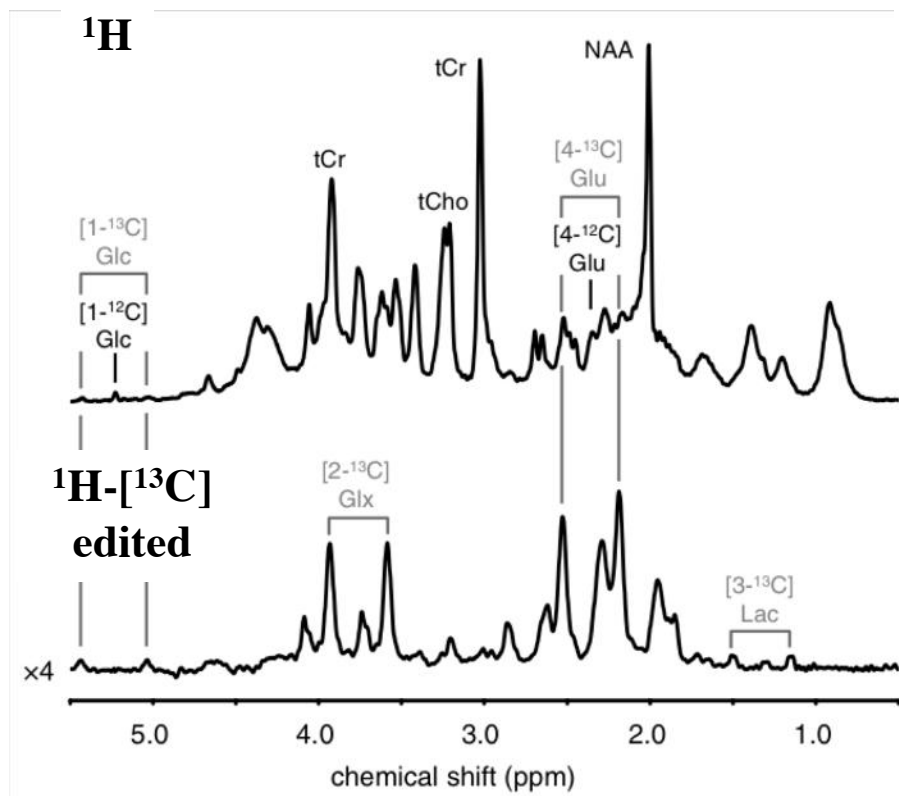


In Vitro Example: ^{13}C MRS



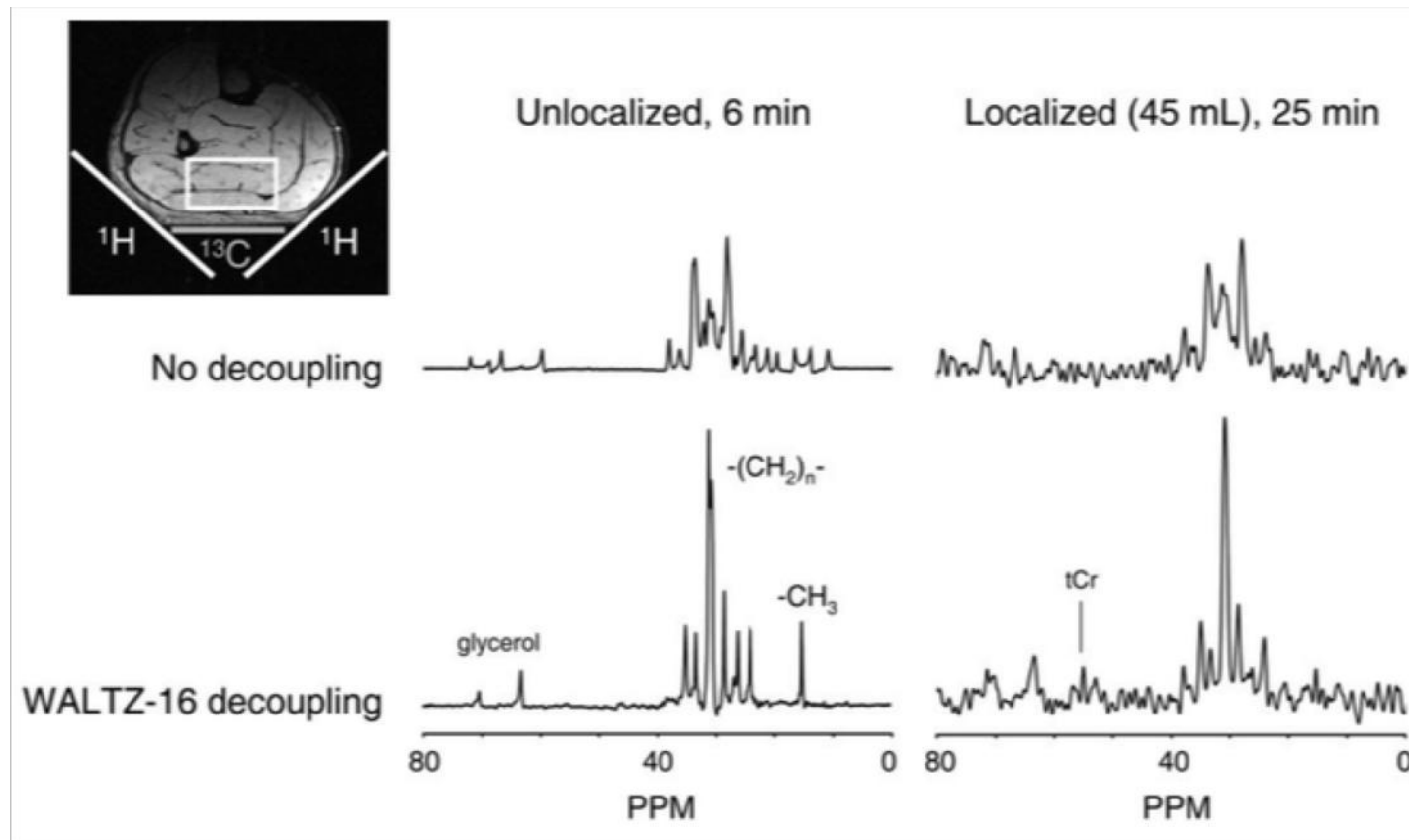
^{13}C -Glucose Infusion: ^1H MRS

Rat brain: 180 μl , TR/TE=4000/8.5 ms, 9.4 T
2 hrs post $[1,6-^{13}\text{C}_2]$ glucose infusion

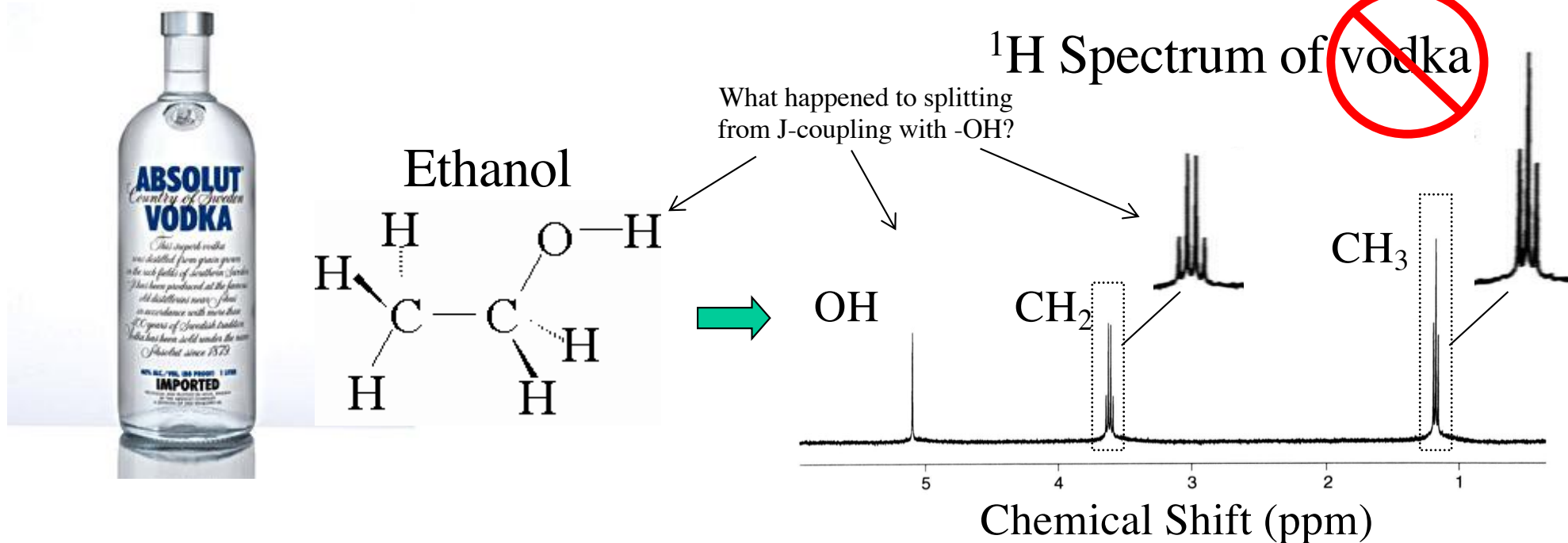


In Vivo Human: ^{13}C MRS

Calf muscle, 4 T, polarization transfer acquisition, natural abundance ^{13}C



Decoupling without Rf?



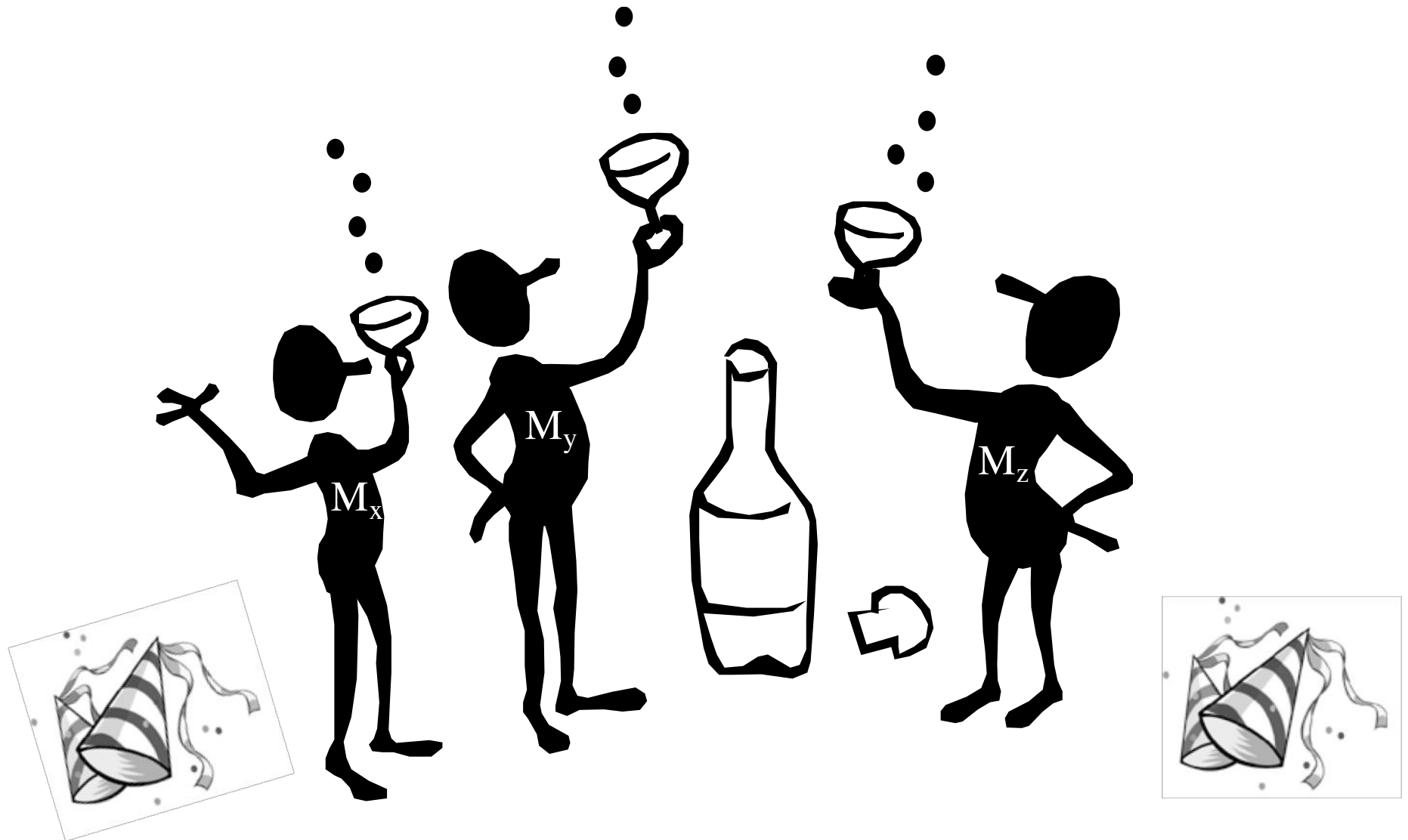
100% Ethanol \rightarrow ^1H Spectrum = ?

Ethanol + 5% H₂O \rightarrow ^1H Spectrum = ?

Hint: Consider the effects of chemical exchange

Ethanol + 5% D₂O \rightarrow ^1H Spectrum = ?

The End.



(Please fill out the course evaluations forms!)