

Electrons in a weak periodic potential

We solve

$$\left(-\frac{\nabla^2}{2m} + U(\vec{r}) \right) \psi(\vec{r}) = \epsilon \psi(\vec{r})$$

$$U(\vec{r} + \vec{R}) = U(\vec{r})$$

with the boundary conditions $\psi(\vec{r} + N_i \vec{a}_i) = \psi(\vec{r})$

We can expand

$$\psi(\vec{r}) = \sum_{\vec{q}} C_{\vec{q}} e^{i\vec{q}\vec{r}} \quad \text{with} \quad \vec{q} = \sum_i \frac{m_i}{N_i} \vec{b}_i$$

From the periodicity of $U(\vec{r})$

$$U(\vec{r}) = \sum_{\vec{K}} U_{\vec{K}} e^{i\vec{K}\vec{r}}$$

\vec{K} - vector of the reciprocal lattice

$$U_{\vec{K}} = \frac{1}{V} \int_{\text{Prim cell}} d\vec{r} e^{-i\vec{K}\vec{r}} U(\vec{r})$$

We choose $U_0 = 0$ and we have $U_{-\vec{K}} = U_{\vec{K}}^*$

The kinetic energy

$$\frac{\vec{p}^2}{2m} \psi = -\frac{\nabla^2}{2m} \psi = \sum_{\vec{q}} \frac{q^2}{2m} C_{\vec{q}} e^{i\vec{q}\vec{r}}$$

The term in the potential energy

$$\begin{aligned}
 U\psi &= \left(\sum_{\mathbf{k}} U_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{r}} \right) \left(\sum_{\mathbf{q}} c_{\mathbf{q}} e^{i\mathbf{q}\cdot\mathbf{r}} \right) = \\
 &= \sum_{\mathbf{k}, \mathbf{q}} U_{\mathbf{k}} c_{\mathbf{q}} e^{i(\mathbf{k}+\mathbf{q})\cdot\mathbf{r}} = \sum_{\mathbf{k}, \mathbf{q}'} U_{\mathbf{k}} c_{\mathbf{q}'-\mathbf{k}} e^{i\mathbf{q}'\cdot\mathbf{r}}
 \end{aligned}$$

$$\sum_{\mathbf{q}} e^{i\mathbf{q}\cdot\mathbf{r}} \left\{ \left(\frac{1}{2m} q^2 - \epsilon \right) c_{\mathbf{q}} + \sum_{\mathbf{k}'} U_{\mathbf{k}'} c_{\mathbf{q}-\mathbf{k}'} \right\} = 0$$

The plane waves in (1) are orthogonal, so

$$\left(\frac{1}{2m} q^2 - \epsilon \right) c_{\mathbf{q}} + \sum_{\mathbf{k}'} U_{\mathbf{k}'} c_{\mathbf{q}-\mathbf{k}'} = 0$$

$$\mathbf{q} \stackrel{\text{def}}{=} \mathbf{k} - \mathbf{K}, \quad \mathbf{K}' \rightarrow \mathbf{K}' - \mathbf{K}$$

$$\mathbf{q} - \mathbf{K}' \rightarrow \mathbf{k} - \mathbf{K} - (\mathbf{K}' - \mathbf{K}) = \mathbf{k} - \mathbf{K}'$$

$$\left[\frac{1}{2m} (\mathbf{k} - \mathbf{K})^2 - \epsilon \right] c_{\mathbf{k}-\mathbf{K}} + \sum_{\mathbf{K}'} U_{\mathbf{K}'-\mathbf{K}} c_{\mathbf{k}-\mathbf{K}'} = 0$$

Summation over recip. lat. vectors \mathbf{K}' . For fixed \mathbf{k} there is an eq. for any \mathbf{K} .

The infinitely many solutions for each \mathbf{k} are labeled with the band index n . \mathbf{k} is in the ^{first} Brillouin zone.

In the free electron case

$$(\epsilon_{k-K}^0 - \epsilon) C_{k-K} = 0$$

$$\epsilon_q^0 = \frac{1}{2m} q^2$$

So either $\epsilon = \epsilon_{k-K}^0$ or $C_{k-K} = 0$, so

$$\epsilon = \epsilon_{k-K}^0 \quad \psi_k \propto e^{i(k-K)r}$$

(When there is no degeneracies).

Now consider a small U Fix k, K_1

and assume that for all other $K \neq K_1$

$$|\epsilon_{k-K_1}^0 - \epsilon_{k-K}^0| \gg U$$

Then

$$[\epsilon - \epsilon_{k-K_1}^0] C_{k-K_1} = \sum_{K \neq K_1} U_{K-K_1} C_{k-K}$$

$$C_{k-K} = \frac{U_{K_1-K} C_{k-K_1}}{\epsilon - \epsilon_{k-K}^0} + \sum_{K' \neq K_1} \frac{U_{K'-K} C_{k-K'}}{\epsilon - \epsilon_{k-K}^0}$$

$$= \frac{U_{K_1-K} C_{k-K_1}}{\epsilon - \epsilon_{k-K}^0} + O(U^2)$$

(13)

$$(\epsilon - \epsilon_{k-k_1}^0) C_{k-k_1} = \sum_k \frac{U_{k-k_1} U_{k_1-k}}{\epsilon - \epsilon_{k-k}^0} C_{k-k_1} + O(U^3)$$

$$\epsilon = \epsilon_{k-k_1}^0 + \sum_k \frac{|U_{k-k_1}|^2}{\epsilon_{k-k_1}^0 - \epsilon_{k-k}^0} + O(U^3)$$

Weakly perturbed non-degenerate levels weakly repel each other (lower levels push ϵ_{k-k_1} up and higher levels push it down)

Suppose the value of \vec{k} is such that there are recip. lat. vec k_1, \dots, k_m such that

$\epsilon_{k-k_1}, \dots, \epsilon_{k-k_m}$ are all within U of each other but far apart from ^{the} other ϵ_{k-k}^0 on the scale of U . In this case we must use a degenerate perturb. theory. Consider the case when only two levels are within U of each other

$$(\epsilon - \epsilon_{k-k_1}^0) C_{k-k_1} = U_{k_2-k_1} C_{k-k_2}$$

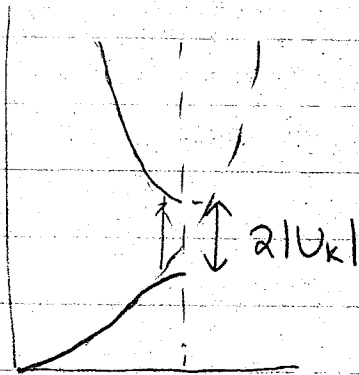
$$(\epsilon - \epsilon_{k-k_2}^0) C_{k-k_2} = U_{k_1-k_2} C_{k-k_1}$$

$$q = k - k_1 \quad K = k_2 - k_1$$

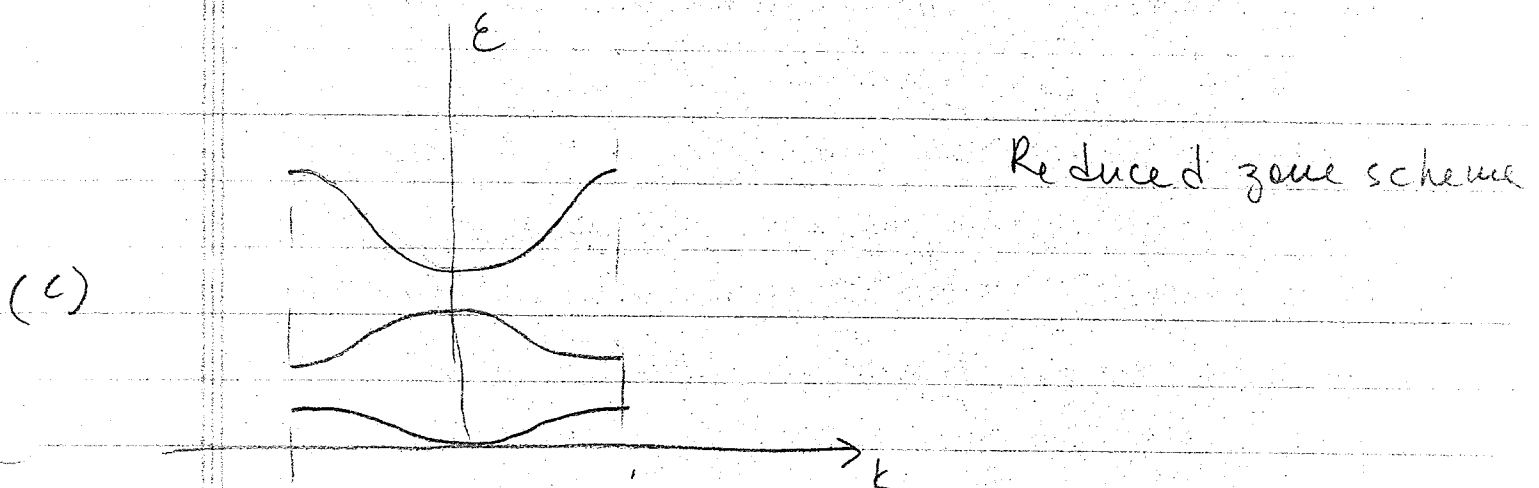
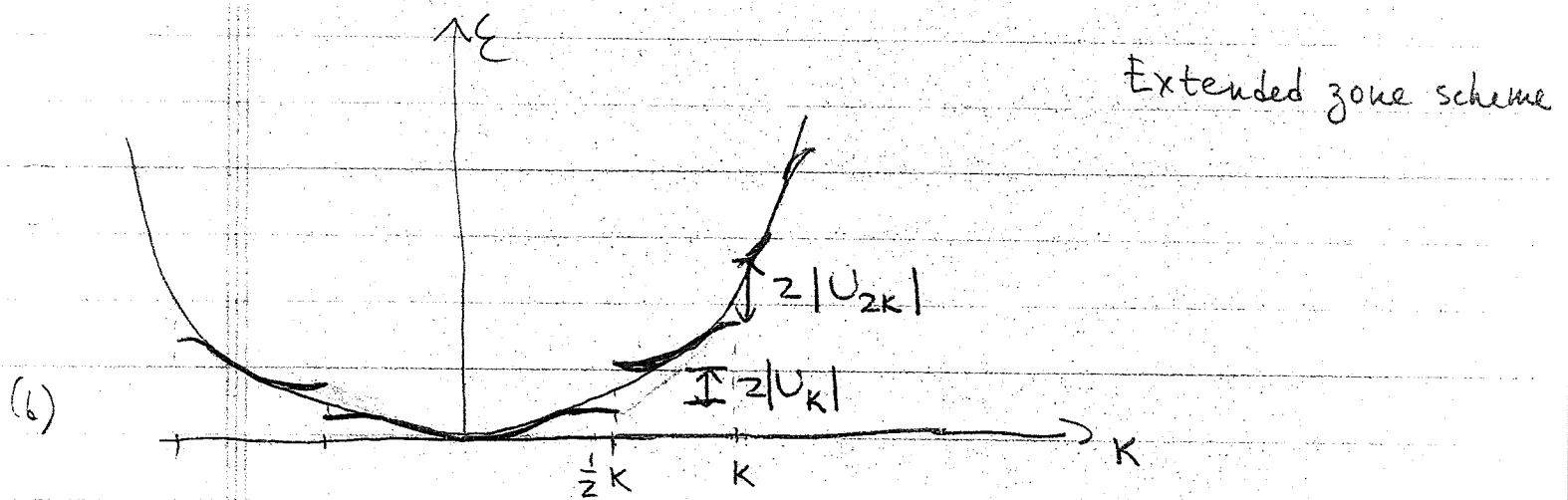
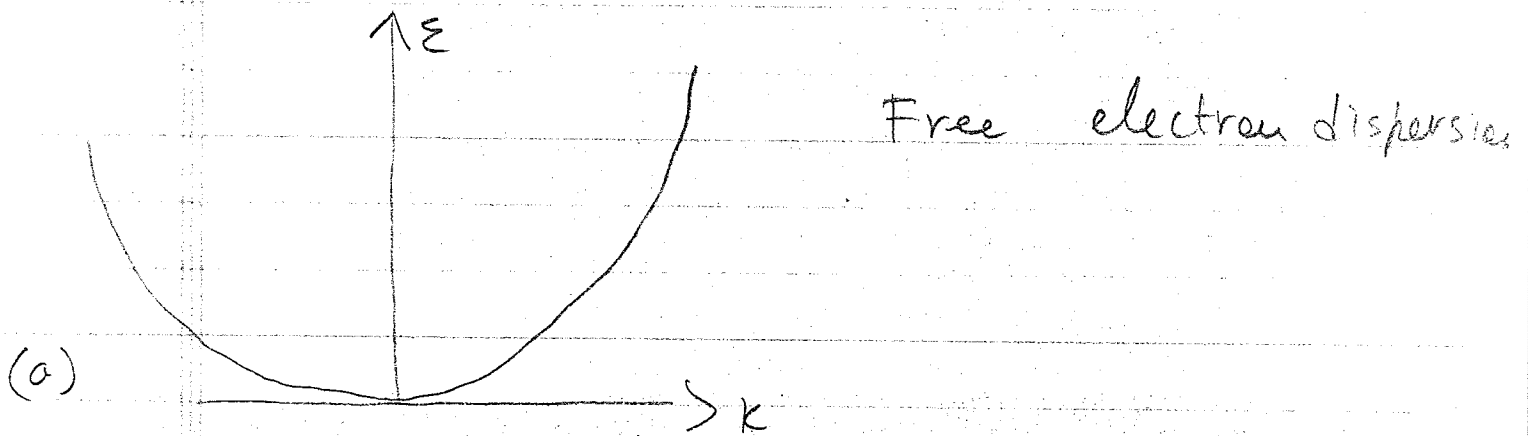
$$\begin{cases} (\varepsilon - \varepsilon_q^0) c_q = U_k c_{q-k} \\ (\varepsilon - \varepsilon_{q-k}^0) c_{q-k} = U_k^* c_q \end{cases}$$

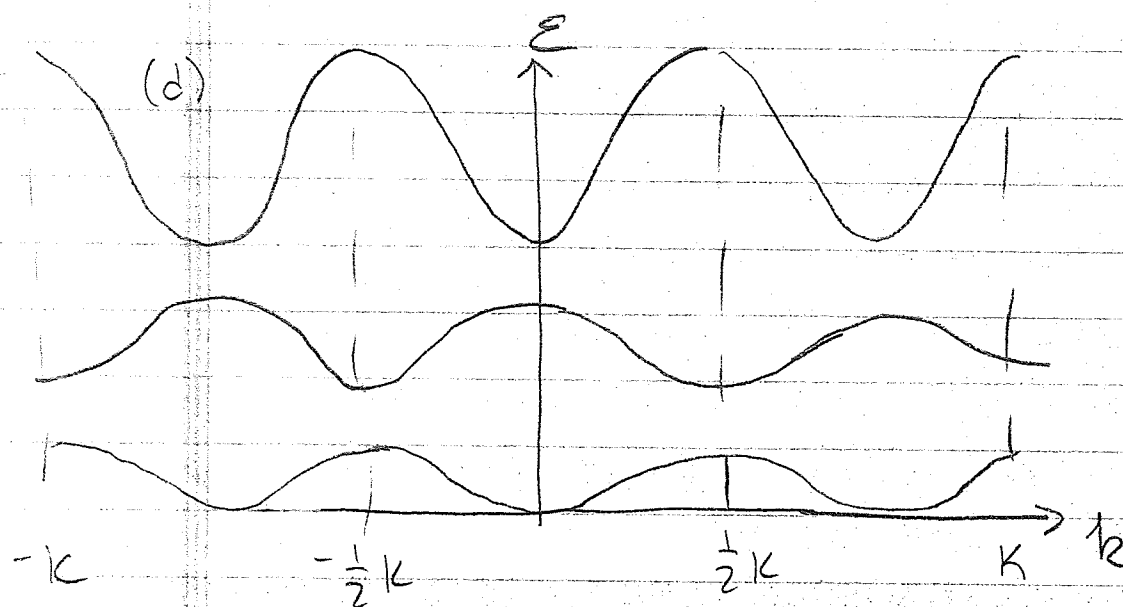
$$\begin{vmatrix} \varepsilon - \varepsilon_q^0 & -U_k \\ -U_k^* & \varepsilon - \varepsilon_{q-k}^0 \end{vmatrix} = 0$$

$$\varepsilon = \frac{1}{2} (\varepsilon_q^0 + \varepsilon_{q-k}^0) \pm \left[\left(\frac{\varepsilon_q^0 - \varepsilon_{q-k}^0}{2} \right)^2 + |U_k|^2 \right]^{1/2}$$



Construction of the Band structure in 1d for a weak potential





Repeated zone
scheme