

# Constraints on local primordial non-Gaussianity from the large-scale clustering of DESI Luminous Red Galaxies

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## ABSTRACT

This paper uses the large-scale clustering of luminous red galaxies selected from the Dark Energy Spectroscopic Instrument Legacy Imaging Surveys Data Release 9 to constrain the local primordial non-Gaussianity (PNG) parameter  $f_{\text{NL}}$ . Using the angular power spectrum, we thoroughly investigate the impact of various photometric systematic effects, such as those caused by Galactic extinction and varying survey depth. Simulations are utilized to construct covariance matrices, evaluate the robustness of our pipeline, and perform statistical tests to assess whether spurious fluctuations are properly mitigated and calibrated. Using modes from  $\ell = 2$  to 300, we find  $X1 < f_{\text{NL}} < X2$  with our conservative and  $X1 < f_{\text{NL}} < X2$  with extreme treatment of imaging systematics, both at 68% confidence. While our results are consistent with zero PNG, but we show that the understanding of imaging systematics is of paramount importance to obtain unbiased constraints on  $f_{\text{NL}}$ .

**Key words:** cosmology: inflation - large-scale structure of the Universe

## 1 INTRODUCTION

Current observations of the cosmic microwave background, large-scale structure, and supernovae are explained by a cosmological model that consists of dark energy, dark matter, and ordinary luminous matter, which has gone through a phase of rapid expansion known as *inflation* (see, e.g., Weinberg et al. 2013). At the end of inflation, the universe was reheated and primordial fluctuations are created to seed the subsequent growth of structure. Statistical properties of primordial fluctuations still remain as one of the unsolved puzzles in modern observational cosmology. Analyses of cosmological data have revealed that initial conditions of the universe are consistent with Gaussian fluctuations; however, there are some classes of models that predict some levels of non-Gaussianities in the primordial gravitational field.

Primordial non-Gaussianity of the local type is parameterized by Komatsu & Spergel (2001),

$$\Phi = \phi + f_{\text{NL}}(\phi^2 - \langle \phi^2 \rangle), \quad (1)$$

where  $\phi$  represents the primordial gravitational field and  $f_{\text{NL}}$  is the coefficient of the nonlinear correction to the primordial field. The most stringent constraints on  $f_{\text{NL}}$  comes from the three-point statistics of the CMB temperature anisotropies,  $f_{\text{NL}} = 0.9 \pm 5.0$  (Akrami et al. 2019). However, it has been shown that PNG has a scale-dependent signature on the two-point clustering of biased tracers of dark matter gravitational field. The clustering of biased tracers is boosted on large scales by a scale-dependent shift proportional to  $k^{-2}$  (see, Dalal et al. 2008).

Previous studies of  $f_{\text{NL}}$  with galaxy and quasar clustering

have been hindered dramatically by spurious fluctuations in target density, which are due to the variation of imaging properties across the sky (Ho et al. 2015). For instance, Pullen & Hirata (2013) found that the level of systematic contamination in the quasar sample of SDSS DRX does not allow a robust  $f_{\text{NL}}$  measurement. These imaging systematic issues are expected to be severe for wide-area galaxy surveys that observe the night sky closer to the Galactic plane and select faint targets. Assuming imaging systematics are under control, the next generation of galaxy surveys such as DESI and the Rubin Observatory are forecast to yield unprecedented constraints on  $f_{\text{NL}}$ . Combined with CMB data, the limits can be enhanced to a precision level required to differentiate between single and multi-field inflationary models.

In this paper, we use photometric luminous red galaxies from the DESI Legacy Imaging Surveys Data Release 9 to constrain the primordial non-Gaussianity parameter  $f_{\text{NL}}$ , while marginalizing over bias and shotnoise parameters. We make use of spectroscopic data from DESI Survey Validation to determine the redshift distribution of galaxies. We cross correlate target density fields with the templates of imaging systematics to quantify systematic error and assess the effectiveness of systematics treatments and the significance of residual systematic error. The methodologies and statistical tools presented in this work will pave the path for future imaging surveys.

## 2 DATA

Luminous red galaxies are massive galaxies that lack active star formation and considered as one of the highly biased tracers of large scale structure. They are widely targeted in previous galaxy surveys CITE, and their clustering and redshift properties are well known CITE. We use the photometric DESI Luminous Red Galaxies (LRG), selected from the imaging surveys (Dey et al. 2018) using color-magnitude cuts described in the  $g$ ,  $r$ ,  $z$ , and  $W1$  bands (see, Zhou et al. 2021), and summarized in Tab. 1. The LRG sample are masked for bright stars, foreground bright galaxies as well as clusters of galaxies<sup>1</sup>, and then binned into HEALPix (Gorski et al. 2005) at  $N_{\text{SIDE}} = 256$  to construct the LRG density field, with an average density of  $800 \text{ deg}^{-2}$  with a coverage around 14000 square degrees of the sky. Fig. 1 shows observed density field of DR9 LRGs in  $\text{deg}^{-2}$ . There are some disconnected islands, hereafter referred to as *spurious islands*, in the DECaLS North region at Declination below  $-11$ , which are removed from the sample to minimize potential calibration issues. Additionally, parts of the DECaLS South with Declination below  $-30$  are cut from the sample, since similar calibration issues might infect our analysis. Section 4 presents how these cuts might alter constraints. DESI imaging is a multi-epoch dataset, and thus LRG density map is accounted for pixel incompleteness using a catalog of random points, uniformly scattered over the footprint with the same cuts and masks as DR9 LRGs. Fig. 2 shows the redshift distribution of DR9 LRGs inferred from DESI Survey Validation CITE and the evolution of galaxy bias for our LRG sample adapted from Zhou et al. (2021), consistent with the assumption of constant clustering amplitude.

We study the correlation coefficient between the LRG density map and potential sources of systematic error, mapped into HEALPix at the same  $N_{\text{SIDE}}$ . The maps used in this work are local stellar density constructed from point-like sources with a  $g$ -band

magnitude in the range  $12 \leq g < 17$  from Gaia Data Release 2 (see, Gaia Collaboration et al. 2018; Myers et al. 2022), Galactic extinction  $E[B-V]$  from Schlegel et al. (1998), and other imaging properties including survey depth (galaxy depth in  $grz$  and PSF depth in  $W1$ ) and seeing in  $grz$  from DESI imaging. Fig. 4 shows the Spearman correlation between galaxy density and imaging properties.

## 3 METHODOLOGY

### 3.1 Measuring Power Spectrum

Galaxy density contrast in pixel  $i$  is defined as,

$$\hat{\delta}_i = \frac{\rho_i}{\hat{\rho}} - 1, \quad (2)$$

where  $\rho$  is the density of galaxies accounted for pixel area  $f_{\text{pix},i}$  and  $\hat{\rho}$  is the mean galaxy density estimated by,

$$\hat{\rho} = \frac{\sum_i \rho_i f_{\text{pix},i}}{\sum_i f_{\text{pix},i}}. \quad (3)$$

By definition, Eqs. 2 and 3 ensure that the integral of the observed quantity over the footprint vanishes:

$$\sum_i \hat{\delta}_i f_{\text{pix},i} = 0, \quad (4)$$

To estimate power spectrum, we expand the galaxy overdensity field in terms of Legendre polynomials,

$$\hat{\delta}_i = \sum_{\ell=0}^{\ell_{\text{max}}} \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\theta_i, \phi_i), \quad (5)$$

where  $\theta, \phi$  represent the polar and azimuthal angular coordinates of pixel  $i$ , respectively. The cutoff at  $\ell = \ell_{\text{max}}$  assumes that modes with  $\ell > \ell_{\text{max}}$  do not contribute significantly to signal power. The coefficients  $a_{\ell m}$  are then obtained by integrating the density contrast field over the total number of non-empty pixels  $N_{\text{pix}}$  and using the orthogonality of Legendre polynomials:

$$\hat{a}_{\ell m} = \frac{4\pi}{N_{\text{pix}}} \sum_{i=1}^{N_{\text{pix}}} \hat{\delta}_i f_{\text{pix},i} Y_{\ell m}^*(\theta_i, \phi_i), \quad (6)$$

where  $*$  represents the complex conjugate. Then, the angular power spectrum estimator is defined as the variance of  $\hat{a}_{\ell m}$  coefficients:

$$\hat{C}_{\ell} = \frac{1}{2\ell+1} \sum_{m=-\ell}^{\ell} \hat{a}_{\ell m} \hat{a}_{\ell m}^*. \quad (7)$$

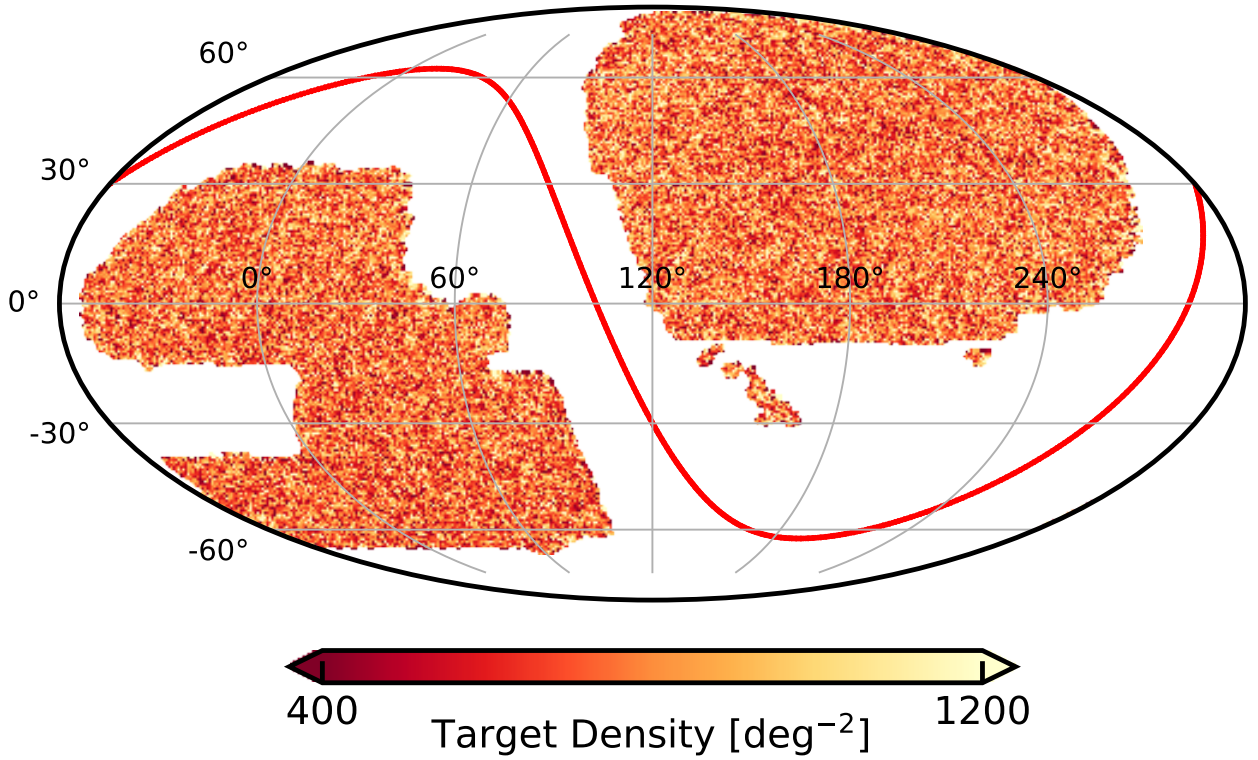
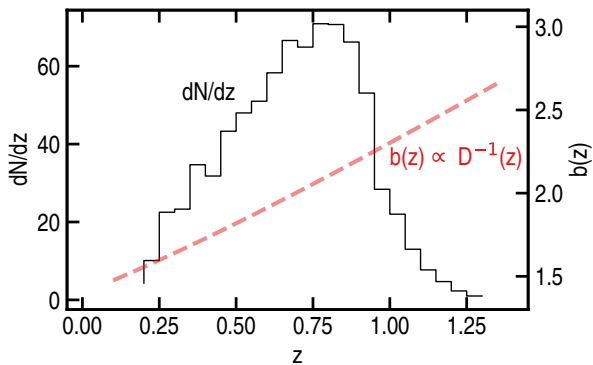
In order to extract  $\hat{a}_{\ell m}$  and compute the angular power spectrum,  $C_{\ell}$ , we make use of the ANAFast function from HEALPix (Gorski et al. 2005) with the third order iteration of the quadrature to increase the accuracy<sup>2</sup>. Due to the survey geometry implicit in the summation over the non-empty pixels and explicit in  $f_{\text{pix},i}$ , our estimator does not return an unbiased estimate of power spectrum, and thus the same effect must be accounted in the modeling of power spectrum. We ignore  $\ell = 0$  and 2 and bin spectra with  $\Delta\ell = 10$  up to  $\ell_{\text{max}} =$

<sup>1</sup> See <https://www.legacysurvey.org/dr9/bitmasks/>

<sup>2</sup> We refer the reader to <https://healpix.sourceforge.io/pdf/subroutines.pdf>, p. 104-105.

**Table 1.** Selection criteria for the LRG targets.

Footprint	Criterion	Description
DECaLS	$z_{\text{fiber}} < 21.7$	Faint limit
	$z - W1 > 0.8 \times (r - z) - 0.6$	Stellar rejection
	$[(g - r > 1.3) \text{ AND } ((g - r) > -1.55 \times (r - W1) + 3.13)] \text{ OR } (r - W1 > 1.8)$	Remove low- $z$ galaxies
	$[(r - W1 > (W1 - 17.26) \times 1.8) \text{ AND } (r - W1 > W1 - 16.36)] \text{ OR } (r - W1 > 3.29)$	Luminosity cut
BASS+MzLS	$z_{\text{fiber}} < 21.71$	Faint limit
	$z - W1 > 0.8 \times (r - z) - 0.6$	Stellar rejection
	$[(g - r > 1.34) \text{ AND } ((g - r) > -1.55 \times (r - W1) + 3.23)] \text{ OR } (r - W1 > 1.8)$	Remove low- $z$ galaxies
	$[(r - W1 > (W1 - 17.24) \times 1.83) \text{ AND } (r - W1 > W1 - 16.33)] \text{ OR } (r - W1 > 3.39)$	Luminosity cut


**Figure 1.** Observed density field of DESI Luminous Red Galaxies Data Release 9 in  $\text{deg}^{-2}$ . Spurious disconnected islands from the DECaLS North footprint at Declination below  $-11$  and parts of the DECaLS South with Declination below  $-30$  are dropped from the DR9 sample due to potential calibration issues.

**Figure 2.** Redshift distribution and bias evolution of DESI LRGs with the assumption of constant clustering amplitude.

### 3.2 Modeling Power Spectrum

The projected angular power spectrum of galaxies in the presence of redshift space distortions and local primordial non-Gaussianity is related to the 3D linear power spectrum  $P(k)$  and shotnoise  $N_{\text{shot}}$  by (see, e.g., Slosar et al. 2008),

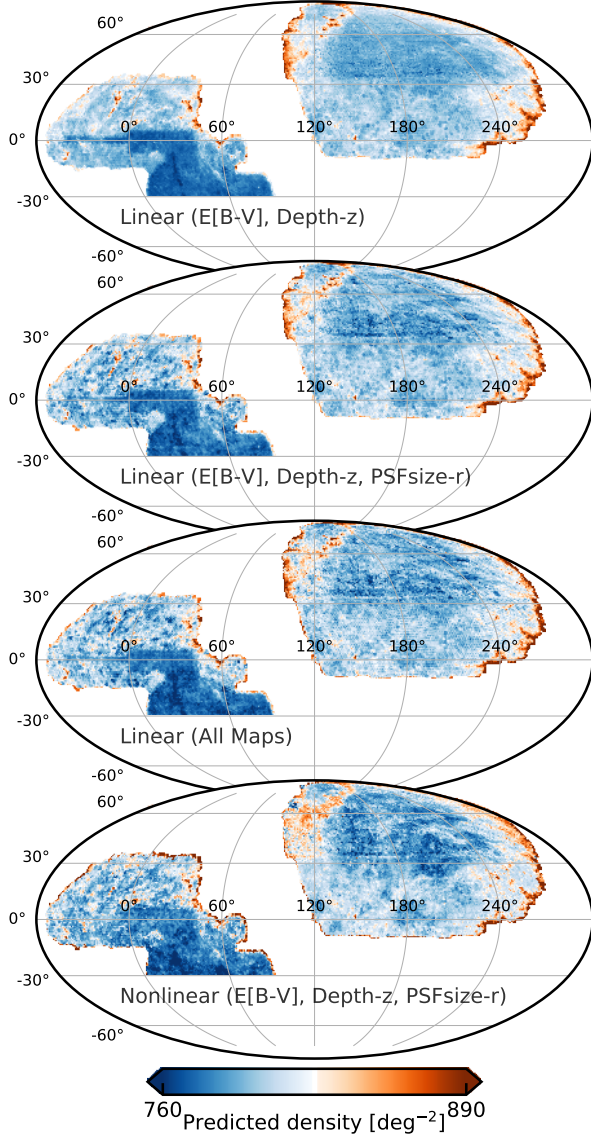
$$C_\ell = \frac{2}{\pi} \int_0^\infty \frac{dk}{k} k^3 P(k) |\Delta_\ell(k)|^2 + N_{\text{shot}}, \quad (8)$$

where  $\Delta_\ell(k) = \Delta_\ell^g(k) + \Delta_\ell^{\text{RSD}}(k) + \Delta_\ell^{\text{fNL}}(k)$  and,

$$\Delta_\ell^g(k) = \int \frac{dr}{r} r b(r) D(r) \frac{dN}{dr} j_\ell(kr), \quad (9)$$

$$\Delta_\ell^{\text{RSD}}(k) = - \int \frac{dr}{r} r f(r) D(r) \frac{dN}{dr} j_\ell''(kr), \quad (10)$$

$$\Delta_\ell^{\text{fNL}}(k) = f_{\text{NL}} \frac{\alpha}{k^2 T(k)} \int \frac{dr}{r} r [b(r) - p] \frac{dN}{dr} j_\ell(kr), \quad (11)$$

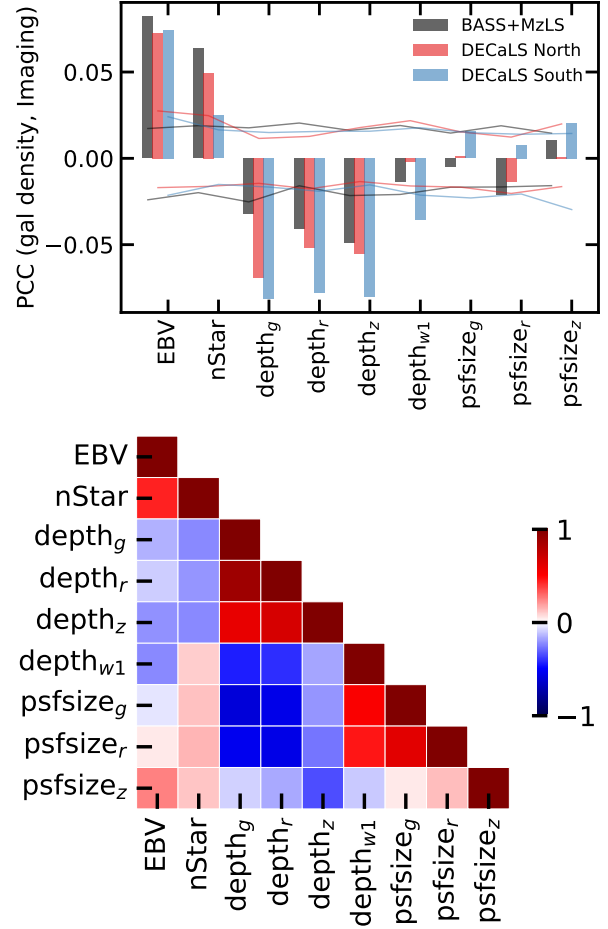


**Figure 3.** Predicted galaxy counts from template regression. Baseline approach uses imaging maps from Zhou et al. (2022): EBV, galaxy depth in  $rgz$ ,  $psfdepth$  in  $W1$ , and  $psfsize$  in  $grz$ . Conservative I uses EBV and galaxy depth in  $z$ , and Conservative II uses EBV, galaxy depth in  $z$ , and  $psfsize$  in  $r$ . In all approaches, the models are regressed on BASS+MzLS, DECaLS North, and DECaLS South separately.

where  $\alpha = 3\delta_c\Omega_M(H_0/c)^2$ ,  $b(r)$  is the linear bias,  $dN/dr$  is the normalized redshift distribution of galaxies<sup>3</sup>,  $D(r)$  is the normalized growth factor such that  $D(0) = 1$ ,  $f(r)$  is the growth rate, and  $r$  is the comoving distance. The parameter  $p$  is the response of the tracer to halo's gravitational field, e.g., 1 for luminous red galaxies and 1.6 for recent mergers. In order to overcome rapid oscillations in spherical Bessel functions, we employ the FFTLog<sup>4</sup> algorithm and its extension as implemented in ? to compute the inner integrations over  $d \ln r$ .

<sup>3</sup>  $dN/dr = (dN/dz) * (dz/dr) \propto (dN/dz) * H(z)$

<sup>4</sup> [github.com/xfangcosmo/FFTLog-and-beyond](https://github.com/xfangcosmo/FFTLog-and-beyond)



**Figure 4.** Top: Pearson-r correlation coefficient between galaxy density and imaging properties in the three imaging regions (top) and between imaging properties themselves for the full DESI footprint (bottom). Solid curves represent the range of correlations observed in 100 randomly selected mock realizations.

### 3.2.1 Survey Geometry

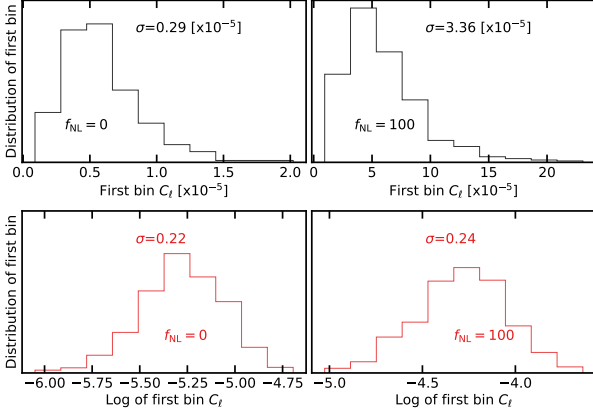
For a galaxy survey that observes the sky partially, the measured power spectrum is convolved with the survey geometry. This means that the pseudo-power spectrum  $\hat{C}_\ell$  obtained by the direct Spherical Harmonic Transforms of a partial sky map, differs from the full-sky angular spectrum  $C_\ell$ . However, their ensemble average is related by (?)

$$\langle \hat{C}_\ell \rangle = \sum_{\ell'} M_{\ell\ell'} \langle C_{\ell'} \rangle, \quad (12)$$

where  $M_{\ell\ell'}$  represents the mode-mode coupling from the partial sky coverage. This is known as the Window Function effect and a proper assessment of this effect is crucial for a robust measurement of the large-scale clustering of galaxies. This window effect is a source of observational systematic error and impacts the measured galaxy clustering, especially on scales comparable to survey size.

We follow a similar approach to that of (?) to model the window function effect on the theoretical power spectrum  $C_\ell$  rather than correcting the measured pseudo-power spectrum from data. First, we use HEALPIX to compute the pseudo-power spectrum of the window  $\hat{C}_\ell^{\text{window}}$ , which is defined by a mask file in ring order-





**Figure 5.** Distribution of the first bin power spectrum for  $f_{\text{NL}} = 0$  and 100 mocks.

ing format with `NSIDE= 256`. Then, we transform it to correlation function by,

$$\omega^{\text{window}}(\theta) = \frac{1}{4\pi} \sum_{\ell} (2\ell + 1) \hat{C}_{\ell}^{\text{window}} P_{\ell}(\cos \theta). \quad (13)$$

Next, we normalize  $\omega^{\text{window}}$  such that it is normalized to one at  $\theta = 0$ . Finally, we multiply the theory correlation function by  $\omega^{\text{window}}$  and transform the result back to  $\ell$ -space,

$$\hat{\omega}^{\text{model}} = \omega^{\text{model}} \omega^{\text{window}} \quad (14)$$

$$\hat{C}_{\ell}^{\text{model}} = 2\pi \int d\theta \hat{\omega}^{\text{model}}(\theta) P_{\ell}(\cos \theta). \quad (15)$$

### 3.2.2 Integral Constraint

The integral of the galaxy density contrast  $\delta$  on the footprint is bound to zero, which is often referred to as the *Integral Constraint*. We account for this effect in the modeling by,

$$\hat{C}_{\ell}^{\text{model,IC}} = \hat{C}_{\ell}^{\text{model}} - \hat{C}_{\ell=0}^{\text{model}} \left( \frac{\hat{C}_{\ell}^{\text{window}}}{\hat{C}_{\ell=0}^{\text{window}}} \right) \quad (16)$$

## 3.3 Characterization of systematic effects

We use the diagnostic tests presented in [Rezaie et al 2021](#) based on cross power spectrum between galaxy density field and imaging maps and mean density contrast as a function of imaging properties to quantify the significance of imaging systematic effects.

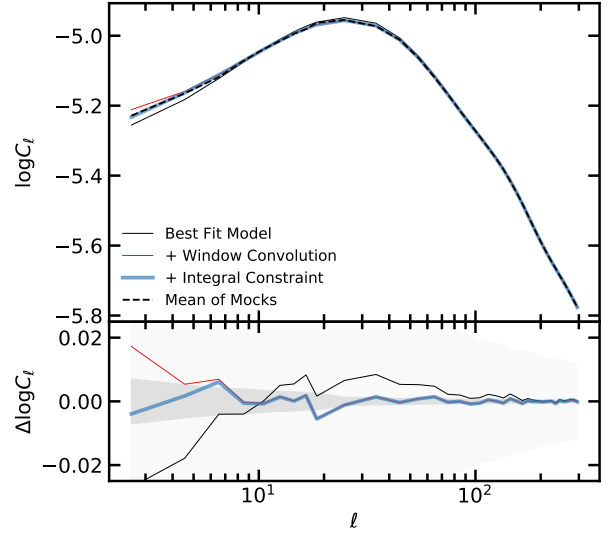
### 3.3.1 Cross Spectrum

Taking  $C_{\ell}^{g,x}$  as the cross power spectrum between galaxy density contrast field and imaging map, one can normalize this quantity by auto power spectrum of imaging map itself:

$$\hat{C}_{x,\ell} = \frac{(\hat{C}_{\ell}^{g,x})^2}{\hat{C}_{\ell}^{x,x}}, \quad (17)$$

and then construct a vector from cross spectra against all other imaging maps:

$$\hat{C}_{X,\ell} = [\hat{C}_{x_1,\ell}, \hat{C}_{x_2,\ell}, \hat{C}_{x_3,\ell}, \dots, \hat{C}_{x_9,\ell}]. \quad (18)$$



**Figure 6.** Mean power spectrum of 1000 mocks with  $f_{\text{NL}} = 0$  and best fit theoretical prediction after accounting for various theoretical systematic effects.

We bin the  $C_X$  measurements with  $\ell$  edges defined at 2, 10, 18, 26, 40, 60, 80, and 100. The mean and standard deviation of  $\hat{C}_{X,\ell}$  for 1000 mocks with and without  $f_{\text{NL}}$  are shown in [Fig. 7](#). Finally, cross power spectrum  $\chi^2$  can be defined as,

$$\chi^2 = C_{X,\ell}^T C^{-1} C_{X,\ell}, \quad (19)$$

where covariance matrix  $C = \langle C_{X,\ell} C_{X,\ell'} \rangle$  is constructed from mocks without systematic effects. This statistics is measured for every mock realization with the leave-one-out technique to construct a histogram, which is then compared to the  $\chi^2$  value observed from the DR9.

### 3.3.2 Mean Density

As an alternative test, we calculate the histogram of the density contrast field relative to each imaging map.

$$\delta_x = (\hat{\rho})^{-1} \frac{\sum_i \rho_i f_{\text{pix},i}}{\sum_i f_{\text{pix},i}}, \quad (20)$$

where the summations are over pixels in each bin of imaging map  $x$ . Similarly, we construct the mean density contrast vector against all imaging maps,

$$\delta_X = [\delta_{x_1}, \delta_{x_2}, \delta_{x_3}, \dots, \delta_{x_9}], \quad (21)$$

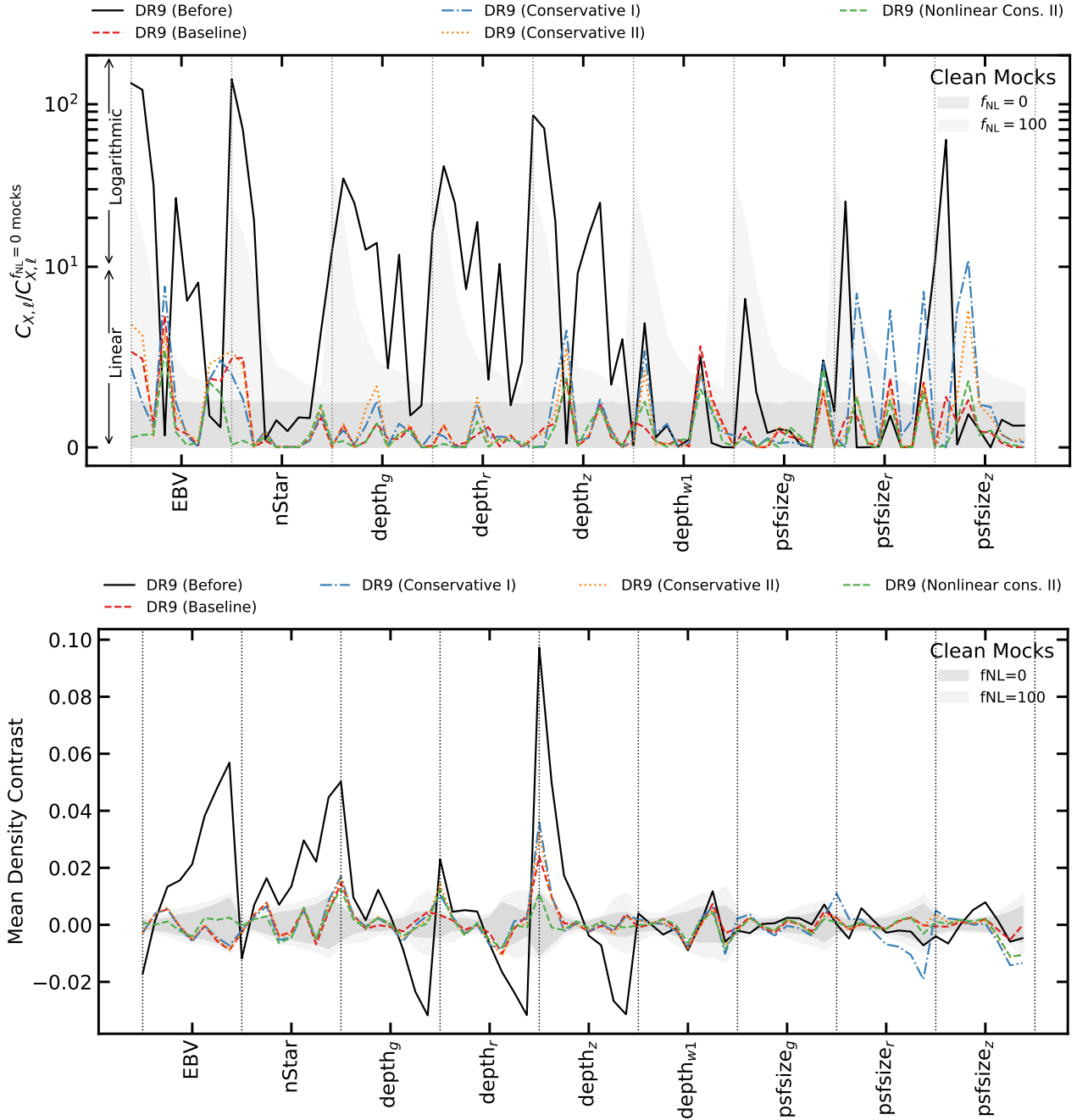
and the total residual error as,

$$\chi^2 = \delta_X^T C^{-1} \delta_X, \quad (22)$$

where the covariance matrix  $C = \langle \delta_X \delta_X \rangle$  is constructed from mocks without systematic effects. [Fig. 7](#) shows the cross power spectrum and mean density contrast for the DR9 LRG sample. The shades represent the  $1\sigma$  level fluctuations observed in 1000 clean mocks with  $f_{\text{NL}} = 0$  and 100.

## 4 RESULTS

**TODO: The mean density and cross power spectrum  $\chi^2$  diagnostics do not reveal any signature of remaining systematic effects – given**

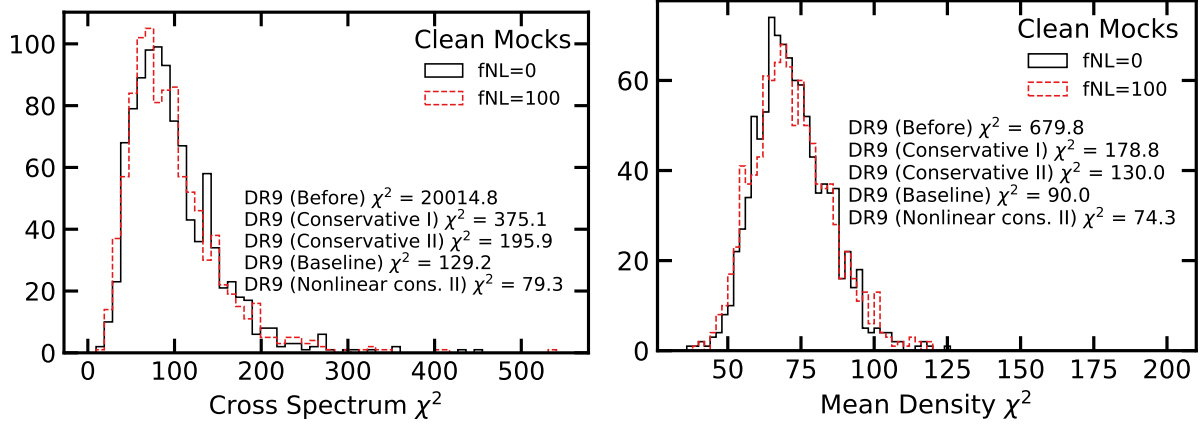


**Figure 7.** Residual systematic fluctuations of DR9 LRGs against imaging maps. Left: cross spectrum. Right: mean density contrast. Shades represent  $1\sigma$  dispersion of 1000 clean mocks with and without  $f_{NL}$ .

all the available templates. But we could be missing some unknown maps. There are some calibration maps made by Aaron which I need to look into. I did not use the Gaia stellar map for training—Zhou et al has a discussion on why it should not be used. But I plan to apply some cuts based on imaging (e.g., poor depth, high stellar density and extinction) to see if those impact the best fit estimates. Next, I will study how changing the lowest ell used could affect our results.

#### 4.1 Lognormal Mocks

Corner plots of the PNG parameter  $f_{NL}$  and bias coefficient are shown in Fig. 9 for fitting the mean power spectrum of the mocks, with and without  $f_{NL}$ . Maximum-A-Posteriori estimates and marginalized mean, median and  $1\sigma$  quantiles are summarized in Tab. 2. Comparing DECaLS North with sky coverage 0.14 to full DESI with 0.40, we find the constraint improve by a factor of 1.9 which is slightly more than  $\sim f_{SKY}^{-1/2}$ , 1.7. while As a robustness test, we also fit the mean power spectrum of the  $f_{NL} = 100$  mocks using the covariance matrix estimated from the  $f_{NL} = 0$  mocks. We



**Figure 8.** Left: Cross power spectrum  $\chi^2$  diagnostic. Right: Mean density contrast diagnostic. The values observed in DR9 before and after linear and nonlinear treatments are quoted and the histograms are constructed from 1000 realizations of clean mocks with  $f_{NL} = 0$  and 100.

find that the constraints improve by a factor of 4.2, due to a higher signal to noise ratio.

## 4.2 DR9 LRGs

## 5 CONCLUSION

## ACKNOWLEDGEMENTS

## DATA AVAILABILITY

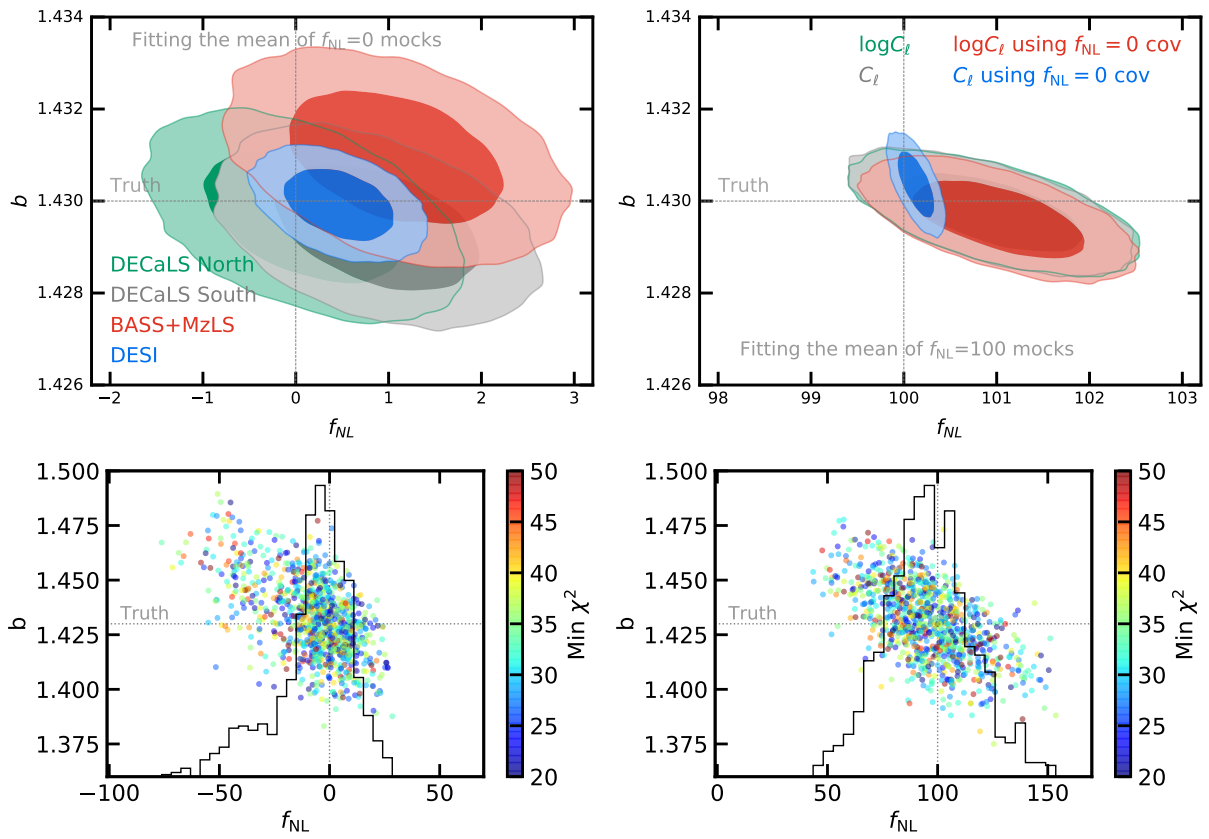
## REFERENCES

- Akrami Y., et al., 2019, arXiv preprint arXiv:1905.05697
- Dalal N., Dore O., Huterer D., Shirokov A., 2008, Physical Review D, 77, 123514
- Dey A., et al., 2018, arXiv preprint arXiv:1804.08657
- Gaia Collaboration et al., 2018, *A&A*, 616, A1
- Gorski K. M., Hivon E., Banday A. J., Wandelt B. D., Hansen F. K., Reinecke M., Bartelmann M., 2005, The Astrophysical Journal, 622, 759
- Ho S., et al., 2015, *J. Cosmology Astropart. Phys.*, 2015, 040
- Komatsu E., Spergel D. N., 2001, Physical Review D, 63, 063002
- Myers A. D., et al., 2022, arXiv e-prints, p. arXiv:2208.08518
- Pullen A. R., Hirata C. M., 2013, Publications of the Astronomical Society of the Pacific, 125, 705
- Schlegel D. J., Finkbeiner D. P., Davis M., 1998, The Astrophysical Journal, 500, 525
- Slosar A., Hirata C., Seljak U., Ho S., Padmanabhan N., 2008, Journal of Cosmology and Astroparticle Physics, 2008, 031
- Weinberg D. H., Mortonson M. J., Eisenstein D. J., Hirata C., Riess A. G., Rozo E., 2013, Physics reports, 530, 87
- Zhou R., et al., 2021, Monthly Notices of the Royal Astronomical Society, 501, 3309

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**Table 2.** Maximum-A-Posteriori (MAP) and marginalized mean estimates for  $f_{\text{NL}}$  from fitting the mean power spectrum of the mocks. Degree of freedom is 33 (37 data points - 3 parameters - 1).

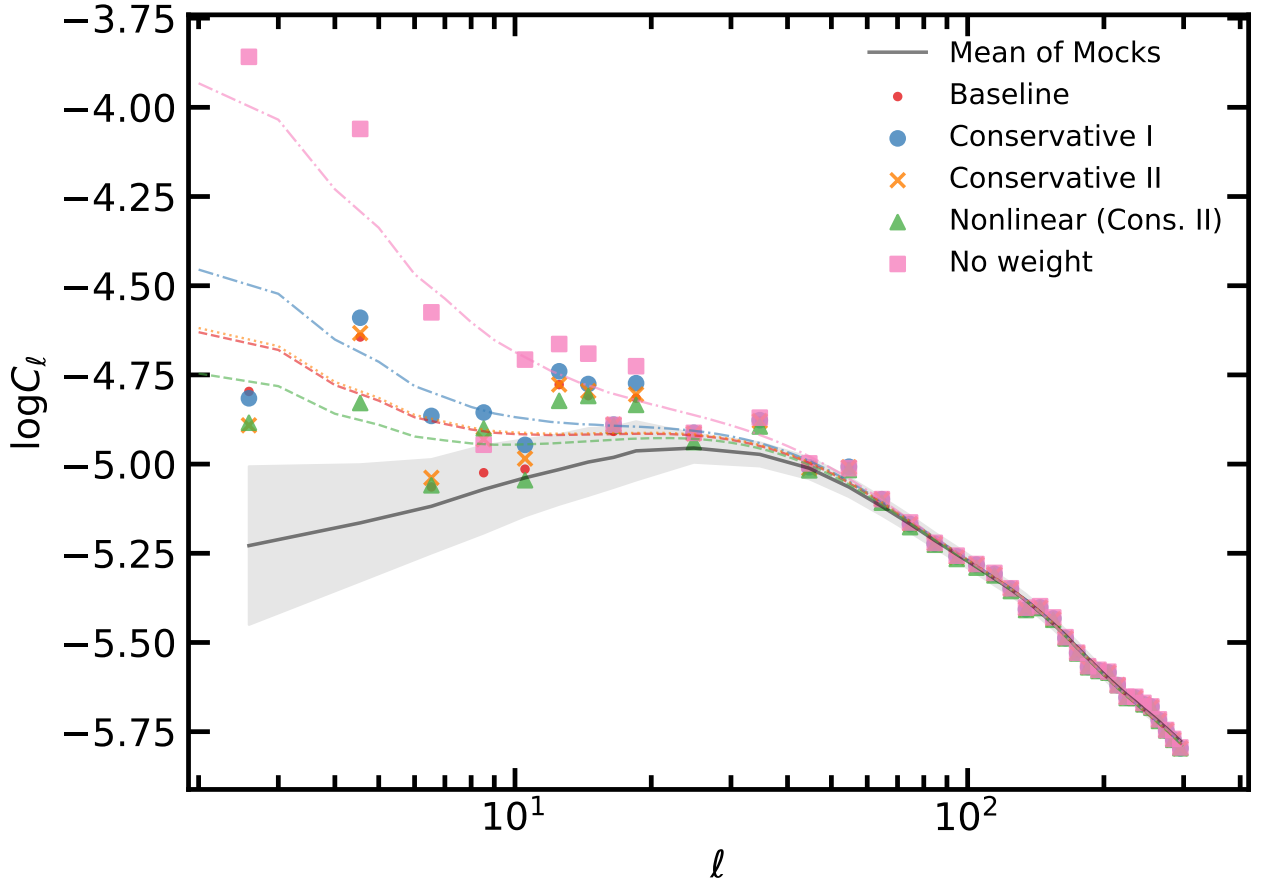
Mock $f_{\text{NL}}$	Footprint	Observable	Best fit	Mean	68% CL	95% CL	$\chi^2$
$f_{\text{NL}} = 100$	DESI	$\log C_\ell$	100.97	100.97	$100.33 < f_{\text{NL}} < 101.61$	$99.72 < f_{\text{NL}} < 102.23$	38.8
$f_{\text{NL}} = 100$	DESI	$C_\ell$	100.97	100.95	$100.32 < f_{\text{NL}} < 101.58$	$99.71 < f_{\text{NL}} < 102.18$	39.0
$f_{\text{NL}} = 100$	DESI	$\log C_\ell$ using $f_{\text{NL}} = 0$ cov	101.01	101.03	$100.42 < f_{\text{NL}} < 101.63$	$99.85 < f_{\text{NL}} < 102.21$	39.9
$f_{\text{NL}} = 100$	DESI	$C_\ell$ using $f_{\text{NL}} = 0$ cov	100.13	100.13	$100.00 < f_{\text{NL}} < 100.26$	$99.88 < f_{\text{NL}} < 100.38$	207.6
$f_{\text{NL}} = 0$	DESI	$\log C_\ell$	0.47	0.47	$0.08 < f_{\text{NL}} < 0.85$	$-0.30 < f_{\text{NL}} < 1.22$	35.7
$f_{\text{NL}} = 0$	DECaLS North	$\log C_\ell$	0.09	0.08	$-0.61 < f_{\text{NL}} < 0.78$	$-1.29 < f_{\text{NL}} < 1.45$	26.7
$f_{\text{NL}} = 0$	DECaLS South	$\log C_\ell$	0.87	0.88	$0.16 < f_{\text{NL}} < 1.58$	$-0.52 < f_{\text{NL}} < 2.27$	34.3
$f_{\text{NL}} = 0$	BASS+MzLS	$\log C_\ell$	1.07	1.07	$0.32 < f_{\text{NL}} < 1.82$	$-0.40 < f_{\text{NL}} < 2.55$	39.4

**Figure 9.** Top: 68% and 95% confidence contours for  $f_{\text{NL}} = 0$  (left) and 100 (right) mocks. Using the  $\log C_\ell$  fitting yield constraints that are insensitive to the covariance used. Bottom: best fit estimates from fitting 1000 lognormal mocks with  $f_{\text{NL}} = 0$  (left) and 100 (right) in the DESI footprint. The truth values are represented by vertical and horizontal lines.



**Table 3.** Maximum-A-Posteriori (MAP) and marginalized mean estimates for  $f_{\text{NL}}$  from fitting power spectrum of DR9 LRGs before and after correcting for systematics. Degree of freedom is 33 (37 data points - 3 parameters - 1).

Footprint	Method	Best fit	Mean	68% CL	95% CL	$\chi^2$
DESI	No Weight	147.13	150.13	$127.58 < f_{\text{NL}} < 172.76$	$108.56 < f_{\text{NL}} < 197.07$	44.4
DESI	Linear (all maps)	46.87	49.04	$33.97 < f_{\text{NL}} < 63.98$	$21.21 < f_{\text{NL}} < 81.00$	41.1
DESI	Linear (Conservative I)	64.46	66.69	$49.67 < f_{\text{NL}} < 83.63$	$35.64 < f_{\text{NL}} < 102.59$	38.8
DESI	Linear (Conservative II)	47.62	49.54	$34.21 < f_{\text{NL}} < 64.81$	$21.27 < f_{\text{NL}} < 82.06$	39.6
DESI	Nonlinear (Conservative II)	37.15	38.73	$24.58 < f_{\text{NL}} < 52.77$	$12.32 < f_{\text{NL}} < 68.55$	34.6
BASS+MzLS	Nonlin. (Cons. II)	20.06	24.71	$-1.52 < f_{\text{NL}} < 51.26$	$-24.95 < f_{\text{NL}} < 82.62$	35.6
DECaLS North	Nonlin. (Cons. II)	53.33	58.36	$30.33 < f_{\text{NL}} < 86.81$	$6.45 < f_{\text{NL}} < 120.93$	41.1
DECaLS North	Nonlin.(Cons. II+CALIBZ+HI)	72.10	78.57	$47.81 < f_{\text{NL}} < 109.27$	$23.22 < f_{\text{NL}} < 146.65$	38.4
DECaLS North (no cut)	Nonlin. (Cons. II)	53.36	58.27	$30.65 < f_{\text{NL}} < 85.91$	$8.33 < f_{\text{NL}} < 118.84$	40.7
DECaLS South	Nonlin.(Cons. II)	40.61	43.17	$19.36 < f_{\text{NL}} < 68.11$	$-6.64 < f_{\text{NL}} < 96.65$	30.2
DECaLS South	Nonlin. (Cons. II+CALIBZ+HI)	43.93	48.74	$23.02 < f_{\text{NL}} < 74.64$	$-0.40 < f_{\text{NL}} < 105.22$	30.8
DECaLS South (no cut)	Nonlin. (Cons. II)	56.93	60.83	$39.20 < f_{\text{NL}} < 82.44$	$21.30 < f_{\text{NL}} < 107.53$	23.8



**Figure 10.** Measured power spectrum of the DR9 LRG sample before and after correcting for systematics with their corresponding best fit theory predictions. The shade represents  $1\sigma$  error constructed from the  $f_{\text{NL}} = 0$  mocks.

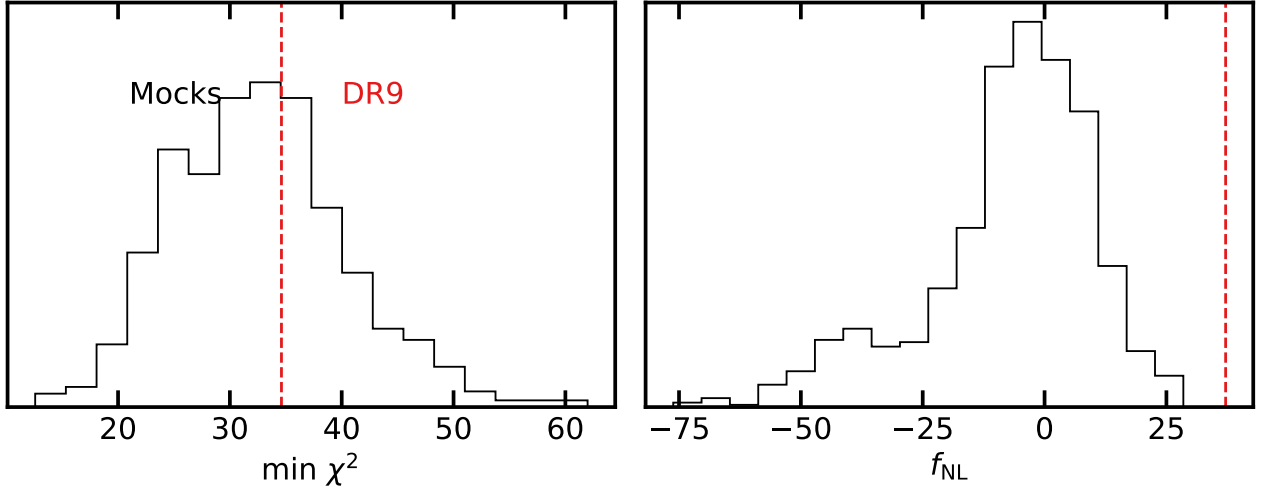


Figure 11.

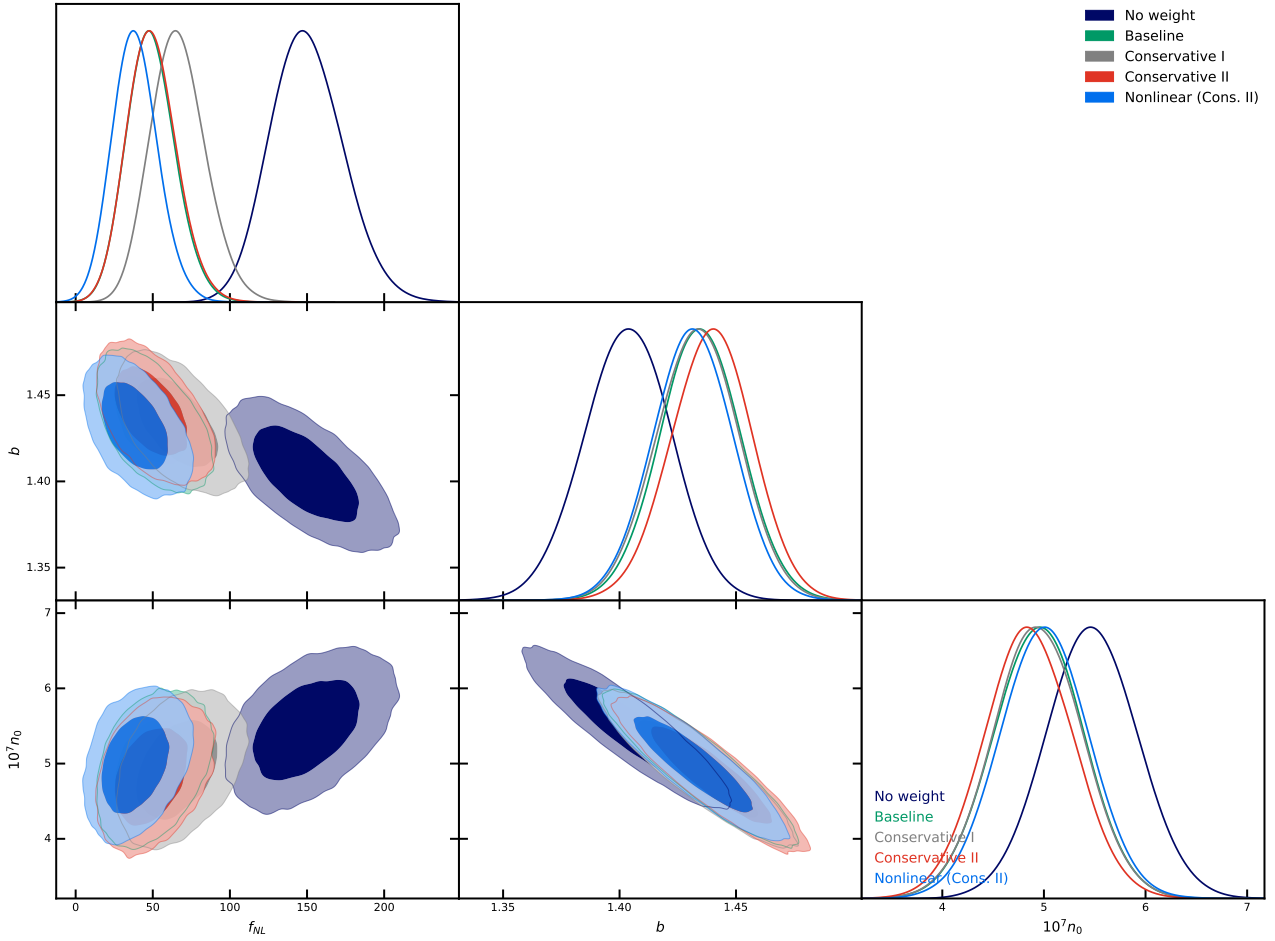


Figure 12. DR9 LRGs for DESI footprint.

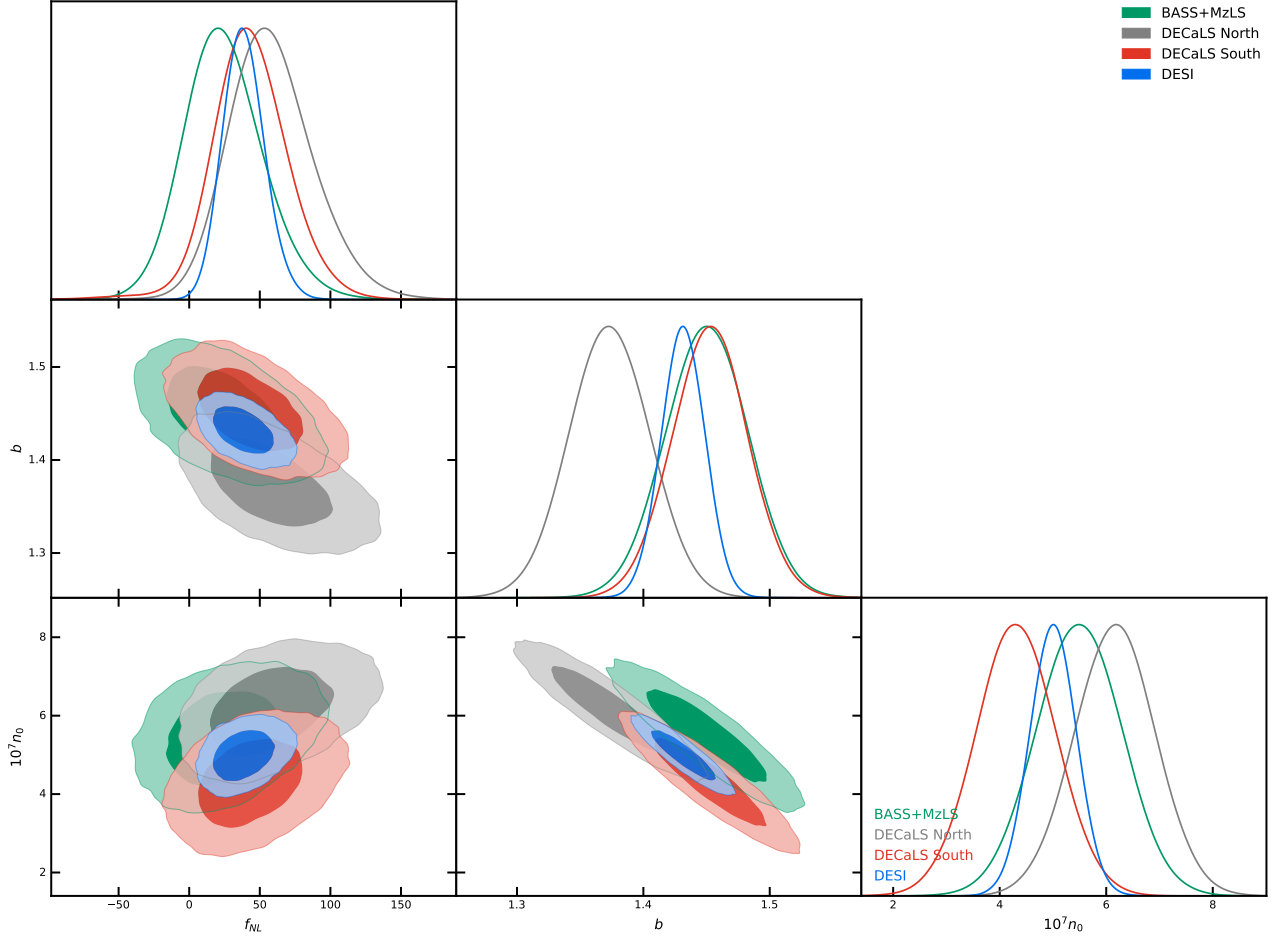
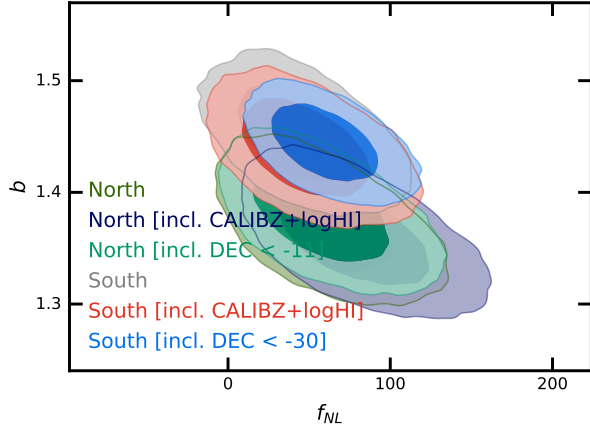


Figure 13. DR9 LRGs for each individual imaging survey.



**Figure 14.** Impact of cutting out spurious islands in DECaLS North and DEC < -30 in DECaLS South.