#### I. SIMULATION OF IONS IN A PENNING TRAP

Ions are confined radially by a magnetic field  $\mathbf{B}$ , with an ideal quadrupolar trapping potential of  $V_T$ . The potential and field near the centre of the trap is approximated as

$$\Phi_T(x,y,z) = V_T(\gamma' - \frac{\alpha'}{2l^2}(x^2 + y^2 - 2z^2)) \qquad \mathbf{E} = -\nabla\Phi_T(x,y,z) = \frac{\alpha}{l^2}(-x\mathbf{i} - y\mathbf{j} + 2z\mathbf{k})$$
(1)

where  $\gamma' = 1/3$  and  $\alpha' = 2.77373$  are geometric factors for the cubic trap, and l is the edge length[1]. The force experienced by an ion in the trap is

$$\mathbf{F} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B}) \tag{2}$$

Charges and mass are simulated in atomic units, lengths in nm and time in ns.

### A. Integration scheme

The code uses the Boris integrator[2] as formulated in Birdsall and Langdon[3]. This is a modified leapfrog scheme in which positions are calculated at times ..., n - 1, n, n + 1, ... and velocities at times ...,  $n^{-1/2}$ ,  $n^{+1/2}$ , ....

$$\mathbf{v}^{-} = \mathbf{v}^{n-1/2} + \frac{q\mathbf{E}}{m} \frac{\Delta t}{2} \tag{3}$$

$$\mathbf{v}' = \mathbf{v}^- + \mathbf{v}^- \times \mathbf{t} \qquad \qquad \mathbf{t} = \frac{q\mathbf{B}}{m} \frac{\Delta t}{2}$$
 (4)

$$\mathbf{v}^{+} = \mathbf{v}^{-} + \mathbf{v}' \times \mathbf{s} \qquad \qquad \mathbf{s} = \frac{2\mathbf{t}}{1 + t^{2}}$$
 (5)

$$\mathbf{v}^{n+1/2} = \mathbf{v}^+ + \frac{q\mathbf{E}}{m} \frac{\Delta t}{2} \tag{6}$$

## B. Ion motion

The cyclotron frequency  $\omega_c$  is predicted by  $\omega_c=\frac{qB}{m}[1]$  (S.I. units - see Table III). To achieve 1% phase error in simulated cyclotron frequency requires  $\Omega\Delta t\lesssim 0.3[3]$ , [4]. The cyclotron radius r is predicted by  $r=\frac{mv}{|q|B}$  and is measured in simulation for a single particle by the distance between maxima in the x dimension.

The electric field gives rise to a magnetron motion of a lower frequency  $\omega_z$ . In a quadrupolar trapping potential of  $V_T$ , the modified cyclotron frequency  $\omega_+$  is predicted by

$$\omega_{+} = \frac{\omega_{c}}{2} + \sqrt{\frac{\omega_{c}^{2}}{4} - \frac{\omega_{z}^{2}}{2}} \qquad \qquad \omega_{z} = \left(\frac{2\alpha q V_{T}}{m l^{2}}\right)^{1/2} \tag{7}$$

The modified cyclotron frequency is measured by peaks in the Fourier transform of the induced current time signal.

# C. Induced current

FTICR-MS measures the current induced between detector plates on opposite walls of the cube parallel to the magnetic field. In the simulation, current is induced by the movement of the 'image'  $\mathbf{E}_{image}(\mathbf{r})$  associated with each ion, that is, the difference in the electric field generated by the ion at each of the two detector plates[1].

$$I = \sum_{i=1}^{N} q_i \mathbf{v}_i \cdot \mathbf{E}_{image}(\mathbf{r}_i) \qquad \qquad \mathbf{E}_{image}(\mathbf{r}) = -\frac{\beta'}{l} r_j \qquad \qquad \beta' = 0.72167$$
 (8)

## D. Evaluation of electrostatic potential and forces

In addition to the influence of the trapping field and the ion-image interaction, each ion experiences a repulsive Coulomb force from every other ion in the packet. The calculation of these forces is nominally  $O(N^2)$ , which makes it infeasible for simulation of large ion packets unless some approximation is used to reduce the computational complexity. Previous simulations have either used the particle-in-cell approximation[5] or virtual particles to reduce the number of interactions to be simulated[6]. Particle-in-cell simulations are best suited to uniform particle distributions with a low required accuracy[7]. In contrast, the adaptive fast multipole algorithm[8] provides guaranteed error bounds for highly non-uniform distributions.

TABLE I REPLICATING: HAN AND SHIN[9]. INPUT, PREDICTED AND MEASURED SIMULATION PARAMETERS

$B = 0.7646T, V_T = 1.0V, l = 0.047$											
			predicted		measured		error: timestep				
species	charge	mass	$v_0$ (m/s)	$\omega_+/2\Pi$ (Hz)	r (mm)	$\omega'_+/2\Pi$ (Hz)	r' (mm)	$\Delta t$ (ns)	$\epsilon : \Delta t$	$\epsilon$ : $\Delta t$ / 10	$\epsilon$ : $\Delta t$ * 10
HCO+	1	29.0182	$10^{4}$	394,022	3.933	404,220	3.933	118			
$CH_3CO^+$	1	43.04462	$10^{4}$	262,028	5.835	272,460	5.835	175			

TABLE II REPLICATING: LEACH ET AL.[5]. INPUT, PREDICTED AND MEASURED SIMULATION PARAMETERS

$B = 7.0T, V_T = 1.0V, l = 0.0508$											
				predicted		measured		error: timestep			
species	charge	mass	$v_0$ (m/s)	$\omega_+/2\Pi$ (Hz)	r (mm)	$\omega'_+/2\Pi$ (Hz)	r' (mm)	$\Delta t$ (ns)	$\epsilon : \Delta t$	$\epsilon$ : $\Delta t$ / 10	$\epsilon : \Delta t * 10$
Cs <sup>+</sup>	1	132.9054	$2.7 \times 10^4$	807,826	5.313	807,680	5.319	59			
$Xx^+(\frac{m}{q}=150)$	1	150.0	$2.7 \times 10^4$	715,652	5.996	715,840	5.997	67			

### E. Simulation results

### REFERENCES

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TABLE III CONVERSION FACTORS

property	unit	SI unit
mass	amu	$1.660538921 \times 10^{-27} kg$
charge	e	$1.60217653 \times 10^{-19}C$