

SE(3) WNOJ Interp Jacobians

$$\begin{aligned}\hat{\gamma}_k(t_k) &= \begin{bmatrix} \mathbf{0} \\ \boldsymbol{\varpi}_k \\ \dot{\boldsymbol{\varpi}}_k \end{bmatrix} \\ \hat{\gamma}_k(t_{k+1}) &= \begin{bmatrix} \ln(\mathbf{T}_{k+1,k})^\vee \\ \mathcal{J}_{k+1,k}^{-1} \boldsymbol{\varpi}_{k+1} \\ -\frac{1}{2}(\mathcal{J}_{k+1,k}^{-1} \boldsymbol{\varpi}_{k+1})^\wedge \boldsymbol{\varpi}_{k+1} + \mathcal{J}_{k+1,k}^{-1} \dot{\boldsymbol{\varpi}}_{k+1} \end{bmatrix} \\ \hat{\gamma}_k(t_{k+1}) &\approx \begin{bmatrix} \ln(\hat{\mathbf{T}}_{\text{op},k+1,k})^\vee + \mathcal{J}_{\text{op},k+1,k}^{-1}(\delta\boldsymbol{\xi}_{k+1} - \mathcal{T}_{\text{op},k+1,k}\delta\boldsymbol{\xi}_k) \\ \left(\mathcal{J}_{\text{op},k+1,k}^{-1} - \frac{1}{2} \left(\mathcal{J}_{\text{op},k+1,k}^{-1}(\delta\boldsymbol{\xi}_{k+1} - \mathcal{T}_{\text{op},k+1,k}\delta\boldsymbol{\xi}_k) \right)^\wedge \right) (\boldsymbol{\varpi}_{\text{op},k+1} + \delta\boldsymbol{\varpi}_{k+1}) \\ \mathbf{e}' + \mathcal{J}_{\text{op},k+1,k}^{-1} \dot{\boldsymbol{\varpi}}_{\text{op},k+1} + \frac{1}{2} \dot{\boldsymbol{\varpi}}_{\text{op},k+1} \mathcal{J}_{k+1,k}^{-1} \delta\boldsymbol{\xi}_{k+1,k} + \mathcal{J}_{\text{op},k+1,k}^{-1} \delta\dot{\boldsymbol{\varpi}}_{k+1} \end{bmatrix}\end{aligned}$$

From Tim Tang's thesis:

$$\begin{aligned}\mathbf{e}' &= -\frac{1}{2}(\mathcal{J}_{k+1,k}^{-1} \boldsymbol{\varpi}_{k+1})^\wedge \boldsymbol{\varpi}_{k+1} \\ &\approx \bar{\mathbf{e}} - \frac{1}{2} \left((\mathcal{J}_{\text{op},k+1,k}^{-1} \boldsymbol{\varpi}_{\text{op},k+1})^\wedge \delta\boldsymbol{\varpi}_{k+1} - \boldsymbol{\varpi}_{\text{op},k+1}^\wedge \mathcal{J}_{\text{op},k+1,k}^{-1} \delta\boldsymbol{\varpi}_{k+1} - \frac{1}{2} \boldsymbol{\varpi}_{\text{op},k+1}^\wedge \boldsymbol{\varpi}_{\text{op},k+1}^\wedge \mathcal{J}_{\text{op},k+1,k}^{-1} \delta\boldsymbol{\xi}_{k+1,k} \right)\end{aligned}$$

where $\delta\boldsymbol{\xi}_{k+1,k} = \mathcal{J}_{\text{op},k+1,k}^{-1}(\delta\boldsymbol{\xi}_{k+1} - \mathcal{T}_{\text{op},k+1,k}\delta\boldsymbol{\xi}_k)$.

The Jacobians for the prior cost terms are included in Tim Tang's thesis, these can be used to do covariance interpolation as well.

Interpolation Jacobians:

$$\begin{aligned}\mathbf{T}(\tau) &= \exp \left(\left(\Lambda_{12} \boldsymbol{\varpi}_k + \Lambda_{13} \dot{\boldsymbol{\varpi}}_k + \Omega_{11} \ln(\mathbf{T}_{k+1,k})^\vee \right. \right. \\ &\quad \left. \left. + \Omega_{12} \mathcal{J}_{k+1,k}^{-1} \boldsymbol{\varpi}_{k+1} + \Omega_{13} \left(-\frac{1}{2}(\mathcal{J}_{k+1,k}^{-1} \boldsymbol{\varpi}_{k+1})^\wedge \boldsymbol{\varpi}_{k+1} + \mathcal{J}_{k+1,k}^{-1} \dot{\boldsymbol{\varpi}}_{k+1} \right) \right)^\wedge \right) \mathbf{T}_k \\ \frac{\partial \mathbf{T}(\tau)}{\partial \mathbf{x}} &= \mathcal{J}_{\text{op},\tau} \frac{\partial \boldsymbol{\xi}_\tau}{\partial \mathbf{x}} + \mathcal{T}_{\text{op},\tau} \frac{\partial \mathbf{T}_k}{\partial \mathbf{x}} \\ \boldsymbol{\varpi}(\tau) &= \mathcal{J}(\boldsymbol{\xi}_\tau) \dot{\boldsymbol{\xi}}_\tau \\ \frac{\partial \boldsymbol{\varpi}(\tau)}{\partial \mathbf{x}} &= \mathcal{J}_{\text{op},\tau} \frac{\partial \dot{\boldsymbol{\xi}}_\tau}{\partial \mathbf{x}} - \frac{1}{2} \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge \frac{\partial \boldsymbol{\xi}_\tau}{\partial \mathbf{x}} \\ \dot{\boldsymbol{\varpi}}(\tau) &\approx \mathcal{J}(\boldsymbol{\xi}_\tau) (\ddot{\boldsymbol{\xi}}_\tau + \frac{1}{2} \dot{\boldsymbol{\xi}}_\tau^\wedge \boldsymbol{\varpi}(\tau)) \\ \frac{\partial \dot{\boldsymbol{\varpi}}(\tau)}{\partial \mathbf{x}} &= \mathcal{J}_{\text{op},\tau} \frac{\partial}{\partial \mathbf{x}} \left(\ddot{\boldsymbol{\xi}}_\tau + \frac{1}{2} \dot{\boldsymbol{\xi}}_\tau^\wedge \boldsymbol{\varpi}(\tau) \right) - \frac{1}{2} \left(\ddot{\boldsymbol{\xi}}_{\text{op},\tau} + \frac{1}{2} \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge \boldsymbol{\varpi}_{\text{op}}(\tau) \right)^\wedge \frac{\partial \boldsymbol{\xi}_\tau}{\partial \mathbf{x}} \\ \frac{\partial}{\partial \mathbf{x}} \left(\ddot{\boldsymbol{\xi}}_\tau + \frac{1}{2} \dot{\boldsymbol{\xi}}_\tau^\wedge \boldsymbol{\varpi}(\tau) \right) &= \frac{\partial \ddot{\boldsymbol{\xi}}_\tau}{\partial \mathbf{x}} + \frac{1}{2} \frac{\partial}{\partial \mathbf{x}} \dot{\boldsymbol{\xi}}_\tau^\wedge \boldsymbol{\varpi}(\tau) \\ &= \frac{\partial \ddot{\boldsymbol{\xi}}_\tau}{\partial \mathbf{x}} - \frac{1}{2} \boldsymbol{\varpi}_{\text{op}}(\tau)^\wedge \frac{\partial \dot{\boldsymbol{\xi}}_\tau}{\partial \mathbf{x}} + \frac{1}{2} \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge \frac{\partial \boldsymbol{\varpi}(\tau)}{\partial \mathbf{x}}\end{aligned}$$

$$\begin{aligned}\frac{\partial \dot{\boldsymbol{\varpi}}(\tau)}{\partial \mathbf{x}} &= \mathcal{J}_{\text{op},\tau} \left(\frac{\partial \ddot{\boldsymbol{\xi}}_\tau}{\partial \mathbf{x}} - \frac{1}{2} \boldsymbol{\varpi}_{\text{op}}(\tau)^\wedge \frac{\partial \dot{\boldsymbol{\xi}}_\tau}{\partial \mathbf{x}} + \frac{1}{2} \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge \left(\mathcal{J}_{\text{op},\tau} \frac{\partial \dot{\boldsymbol{\xi}}_\tau}{\partial \mathbf{x}} - \frac{1}{2} \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge \frac{\partial \boldsymbol{\xi}_\tau}{\partial \mathbf{x}} \right) \right) \\ &\quad - \frac{1}{2} \left(\ddot{\boldsymbol{\xi}}_{\text{op},\tau} + \frac{1}{2} \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge \boldsymbol{\varpi}_{\text{op}}(\tau) \right)^\wedge \frac{\partial \boldsymbol{\xi}_\tau}{\partial \mathbf{x}}\end{aligned}$$

$$\begin{aligned}\frac{\partial \dot{\boldsymbol{\varpi}}(\tau)}{\partial \mathbf{x}} &= \mathcal{J}_{\text{op},\tau} \frac{\partial \ddot{\boldsymbol{\xi}}_\tau}{\partial \mathbf{x}} \\ &\quad + \mathcal{J}_{\text{op},\tau} \left(-\frac{1}{2} \boldsymbol{\varpi}_{\text{op}}(\tau)^\wedge + \frac{1}{2} \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge \mathcal{J}_{\text{op},\tau} \right) \frac{\partial \dot{\boldsymbol{\xi}}_\tau}{\partial \mathbf{x}} \\ &\quad + \left(-\frac{1}{4} \mathcal{J}_{\text{op},\tau} \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge - \frac{1}{2} \left(\ddot{\boldsymbol{\xi}}_{\text{op},\tau} + \frac{1}{2} \dot{\boldsymbol{\xi}}_{\text{op},\tau}^\wedge \boldsymbol{\varpi}_{\text{op}}(\tau) \right)^\wedge \right) \frac{\partial \boldsymbol{\xi}_\tau}{\partial \mathbf{x}}\end{aligned}$$

$$\begin{aligned}\boldsymbol{\xi}_\tau = \boldsymbol{\xi}(\tau) &= \boldsymbol{\Lambda}_{12}(\tau) \boldsymbol{\varpi}_k + \boldsymbol{\Lambda}_{13}(\tau) \dot{\boldsymbol{\varpi}}_k + \boldsymbol{\Omega}_{11}(\tau) \ln(\mathbf{T}_{k+1,k})^\vee + \boldsymbol{\Omega}_{12}(\tau) \mathcal{J}_{k+1,k}^{-1} \boldsymbol{\varpi}_{k+1} \\ &\quad + \boldsymbol{\Omega}_{13} \left(-\frac{1}{2} \left(\mathcal{J}_{k+1,k}^{-1} \boldsymbol{\varpi}_{k+1} \right)^\wedge \boldsymbol{\varpi}_{k+1} + \mathcal{J}_{k+1,k}^{-1} \dot{\boldsymbol{\varpi}}_{k+1} \right)\end{aligned}$$

$$\begin{aligned}\frac{\partial \boldsymbol{\xi}_\tau}{\partial \delta \boldsymbol{\xi}_{k+1}} &= \boldsymbol{\Omega}_{11} \mathcal{J}_{\text{op},k+1,k}^{-1} + \frac{1}{2} \boldsymbol{\Omega}_{12} \boldsymbol{\varpi}_{\text{op},k+1}^\wedge \mathcal{J}_{\text{op},k+1,k}^{-1} \\ &\quad + \frac{1}{4} \boldsymbol{\Omega}_{13} \boldsymbol{\varpi}_{\text{op},k+1}^\wedge \boldsymbol{\varpi}_{\text{op},k+1}^\wedge \mathcal{J}_{\text{op},k+1,k}^{-1} + \frac{1}{2} \boldsymbol{\Omega}_{13} \dot{\boldsymbol{\varpi}}_{\text{op},k+1}^\wedge \mathcal{J}_{\text{op},k+1,k}^{-1}\end{aligned}$$

$$\frac{\partial \boldsymbol{\xi}_\tau}{\partial \delta \boldsymbol{\xi}_k} = -\frac{\partial \boldsymbol{\xi}_\tau}{\partial \delta \boldsymbol{\xi}_{k+1}} \boldsymbol{\tau}_{\text{op},k+1,k}$$

$$\frac{\partial \boldsymbol{\xi}_\tau}{\partial \delta \boldsymbol{\varpi}_k} = \boldsymbol{\Lambda}_{12} \mathbf{1}$$

$$\frac{\partial \boldsymbol{\xi}_\tau}{\partial \delta \boldsymbol{\varpi}_{k+1}} = \boldsymbol{\Omega}_{12} \mathcal{J}_{\text{op},k+1,k}^{-1} - \frac{1}{2} \boldsymbol{\Omega}_{13} \left(\left(\mathcal{J}_{\text{op},k+1,k}^{-1} \boldsymbol{\varpi}_{\text{op},k+1} \right)^\wedge - \boldsymbol{\varpi}_{\text{op},k+1}^\wedge \mathcal{J}_{\text{op},k+1,k}^{-1} \right)$$

$$\frac{\partial \boldsymbol{\xi}_\tau}{\partial \delta \dot{\boldsymbol{\varpi}}_k} = \boldsymbol{\Lambda}_{13} \mathbf{1}$$

$$\frac{\partial \boldsymbol{\xi}_\tau}{\partial \delta \dot{\boldsymbol{\varpi}}_{k+1}} = \boldsymbol{\Omega}_{13} \mathcal{J}_{\text{op},k+1,k}^{-1}$$

Note that the Jacobians of $\dot{\boldsymbol{\xi}}_\tau$ and $\ddot{\boldsymbol{\xi}}_\tau$ have the same form as above, except that we use the second and third row of the interpolation matrices respectively. For example, $\frac{\partial \dot{\boldsymbol{\xi}}_\tau}{\partial \delta \boldsymbol{\xi}_{k+1}} = \boldsymbol{\Omega}_{21} \mathcal{J}_{\text{op},k+1,k}^{-1} + \frac{1}{2} \boldsymbol{\Omega}_{22} \boldsymbol{\varpi}_{\text{op},k+1}^\wedge \mathcal{J}_{\text{op},k+1,k}^{-1} + \dots$

Extrapolation

WNOA:

$$\begin{aligned}\boldsymbol{\xi}_k(\tau) &= \ln(\mathbf{T}(t) \mathbf{T}_k^{-1})^\vee \\ \dot{\boldsymbol{\xi}}_k(\tau) &= \mathcal{J}(\boldsymbol{\xi}_k(\tau))^{-1} \boldsymbol{\varpi}(t)\end{aligned}$$

Past the end knot: $\boldsymbol{\xi}_k = \mathbf{0}^T$ and so $\mathcal{J}(\boldsymbol{\xi}_k(\tau))^{-1} = \mathbf{1}$, $\dot{\boldsymbol{\xi}}_k = \boldsymbol{\varpi}_k$.

$$\boldsymbol{\gamma}_\tau = \begin{bmatrix} \mathbf{1} & \Delta t \\ \mathbf{0} & \mathbf{1} \end{bmatrix} \begin{bmatrix} \boldsymbol{\xi}_k \\ \dot{\boldsymbol{\xi}}_k \end{bmatrix} = \begin{bmatrix} \Delta t \boldsymbol{\varpi}_k \\ \boldsymbol{\varpi}_k \end{bmatrix}$$

WNOJ:

Past the end knot: $\boldsymbol{\xi}_k = \mathbf{0}^T$ and so $\mathcal{J}(\boldsymbol{\xi}_k(\tau))^{-1} = \mathbf{1}$, $\dot{\boldsymbol{\xi}}_k = \boldsymbol{\varpi}_k$.

$$\begin{aligned}\ddot{\xi}_k &= -\frac{1}{2}\dot{\xi}_k(t)^\wedge \varpi(t) + \mathcal{J}(\xi_k(t))^{-1}\dot{\varpi}(t) \\ &= \dot{\varpi}_k\end{aligned}$$

So the extrapolation for WNOJ looks like this:

$$\gamma_\tau = \begin{bmatrix} \xi_\tau \\ \dot{\xi}_\tau \\ \ddot{\xi}_\tau \end{bmatrix} = \underbrace{\begin{bmatrix} 1 & \Delta t 1 & \frac{1}{2}\Delta t^2 1 \\ 0 & 1 & \Delta t 1 \\ 0 & 0 & 1 \end{bmatrix}}_{\Phi(\tau, t_k)} \underbrace{\begin{bmatrix} \xi_k \\ \dot{\xi}_k \\ \ddot{\xi}_k \end{bmatrix}}_{\gamma_k(t_k)} = \begin{bmatrix} \Delta t \varpi_k + \frac{1}{2}\Delta t^2 \dot{\varpi}_k \\ \varpi_k + \Delta t \dot{\varpi}_k \\ \dot{\varpi}_k \end{bmatrix}$$

We can use this formula for doing the extrapolation of any GP prior including WNOA, WNOJ, Singer, etc:

$$\gamma(\tau) = \Phi(\tau, t_k)\gamma_k(t_k)$$

In order to convert from local to global variables, we do the following:

$$\begin{aligned}\mathbf{T}(\tau) &= \exp(\xi_\tau^\wedge)\hat{\mathbf{T}}_k \\ \varpi(\tau) &= \mathcal{J}(\xi_\tau)\dot{\xi}_\tau \\ &\approx \dot{\xi}_\tau \\ \dot{\varpi}(\tau) &\approx \mathcal{J}(\xi_\tau)(\ddot{\xi}_\tau + \frac{1}{2}\dot{\xi}_\tau^\wedge \varpi(\tau)) \\ &\approx \ddot{\xi}_\tau\end{aligned}$$

where we have made the approximation that ξ_τ is small, and so $\mathcal{J}(\xi_\tau)$ is close to identity. Hence, our extrapolation formulas are approximate and this approximation only holds so long as ξ_τ is small.

$$\Phi = \begin{bmatrix} \mathbf{C}_{11} & \mathbf{C}_{12} & \mathbf{C}_{13} \\ \mathbf{0} & \mathbf{C}_{22} & \mathbf{C}_{23} \\ \mathbf{0} & \mathbf{0} & \mathbf{C}_{33} \end{bmatrix}$$

Φ is defined differently for the WNOJ and Singer priors.

Next, we can easily define the Jacobians with respect to the perturbations:

$$\begin{aligned}\frac{\partial \xi_\tau}{\partial \delta \varpi_k} &= \mathbf{C}_{12} \\ \frac{\partial \xi_\tau}{\partial \delta \dot{\varpi}_k} &= \mathbf{C}_{13} \\ \frac{\partial \dot{\xi}_\tau}{\partial \delta \varpi_k} &= \mathbf{C}_{22} \\ \frac{\partial \dot{\xi}_\tau}{\partial \delta \dot{\varpi}_k} &= \mathbf{C}_{23} \\ \frac{\partial \ddot{\xi}_\tau}{\partial \delta \dot{\varpi}_k} &= \mathbf{C}_{33}\end{aligned}$$