

Stern-Gerlach Experiment

Dr. Suchetana Chatterjee

November 2, 2015

The Stern-Gerlach experiment is of great historical and philosophical importance (and interest) in physics. It is historically interesting, because it is this experiment that led to the discovery of spin, and it has since played a very important role in projecting a simple visualization of quantum mechanics. In this tutorial, we shall briefly explore these aspects of the “Stern-Gerlach experiment”.

1. An electron is

- I. an electric dipole.
- II. an electric monopole.
- III. a magnetic dipole.
- IV. a magnetic quadrupole.

Ans. The answers are **II** and **III**. The electron has unit quantum charge, that makes it an electric monopole. But due to the presence of electron spin, it acts as a magnetic dipole with dipole moment $\vec{\mu} = \gamma \vec{s}$, where γ is called the ‘gyromagnetic ratio’ of the electron, and \vec{s} is the spin angular momentum.

2. What is the energy associated with a magnetic dipole in an external magnetic field?

- I. $H = -\vec{\mu} \cdot \vec{B}$.
- II. $H = -\vec{\mu} \times \vec{B}$.

III. $H = \vec{\nabla} (\vec{\mu} \cdot \vec{B}) \times \vec{\mu}.$

IV. $H = \vec{\nabla} (\vec{\mu} \times \vec{B}).$

Ans. The energy associated with a magnetic dipole moment $\vec{\mu}$ in an external magnetic field \vec{B} is given by $H = -\vec{\mu} \cdot \vec{B}.$

3. What is the force on a magnetic dipole due to an external magnetic field?

I. $\vec{\mu} \times \vec{B}.$

II. $|\vec{\mu}| \vec{B}.$

III. $\vec{\nabla} \times (\vec{\mu} \times \vec{B}).$

IV. $\vec{\nabla} (\vec{\mu} \cdot \vec{B}).$

Ans. The force originates due to a gradient in the potential energy. Since the energy of a magnetic dipole of magnetic moment $\vec{\mu}$ in an external magnetic field \vec{B} is given by $H = -\vec{\mu} \cdot \vec{B}$, the force is given by $\vec{F} = -\vec{\nabla} H = -\vec{\nabla} (-\vec{\mu} \cdot \vec{B}) = \vec{\nabla} (\vec{\mu} \cdot \vec{B}).$

4. Which of the following are non-uniform magnetic fields?

I. $\vec{B} = B_0 \hat{i} + B'_0 \hat{j}.$

II. $\vec{B} = B_0 (\hat{i} + \hat{j} + \hat{k}).$

III. $\vec{B} = B_0 \sin \frac{x}{a} \hat{k}.$

IV. $\vec{B} = -\alpha x \hat{i} + (B_0 + \alpha z) \hat{k}.$

Ans. Magnetic fields in options **III** and **IV** have spatial dependence, whereas the magnetic fields in options **I** and **II** are constant throughout space.

5. Calculate the force on a magnetic dipole $\vec{\mu} = \gamma \vec{s}$ under the influence of the magnetic field $\vec{B} = \alpha z \hat{k}$.

I. 0.

II. $\gamma \hat{S}_z \alpha z \hat{i}$.

III. $\gamma \hat{S}_z \alpha z \hat{k}$.

IV. $\alpha \gamma (\hat{S}_x + \hat{S}_y + \hat{S}_z)$.

Ans. The choice is **III**. Force is given as $\vec{\nabla} (\vec{\mu} \cdot \vec{B})$. Now, $\vec{\mu} = \gamma \vec{s}$. The associated magnetic energy is then,

$$\begin{aligned} -\vec{\mu} \cdot \vec{B} &= -\gamma (\hat{S}_x \hat{i} + \hat{S}_y \hat{j} + \hat{S}_z \hat{k}) \cdot (\alpha z \hat{k}) \\ &= -\gamma [0 + 0 + \hat{S}_z \alpha z] \\ &= -\gamma \hat{S}_z \alpha z. \end{aligned}$$

Thus, the force is,

$$\begin{aligned} -\vec{\nabla} H &= \vec{\nabla} (\vec{\mu} \cdot \vec{B}) \\ &= \alpha \gamma \vec{\nabla} (\hat{S}_z z) \\ &= \alpha \gamma \left[\frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k} \right] (\hat{S}_z z) \\ &= \alpha \gamma [0 + 0 + \hat{S}_z \hat{k}] = \alpha \gamma \hat{S}_z \hat{k}. \end{aligned}$$

In a Stern-Gerlach experiment, a non-uniform magnetic field is used. The setup is as follows:

In Stern-Gerlach experiment, if we send a beam of charged particles through the apparatus, then the Lorentz force arising from the interaction of the charged particles with the non-uniform magnetic field will be quite complicated, and we do not want to add additional complications to the experiment. So, typically neutral atoms with single unpaired electrons (for example, silver atoms) are used as the constituents of the beam.

6. For Ag atoms, total angular momentum $L = 0$ because the outermost shell is a 's' shell. It has $s = \frac{1}{2}$, because there is only one electron in the outermost shell. Can you calculate the force on such an atom due to an external magnetic field $\vec{B} = \alpha z \hat{k}$? [Hint: Think about the previous question, and γ is the gyromagnetic ratio of the electron.]

- I. $\pm \frac{1}{2} \alpha \gamma \hat{k}$.
- II. $\alpha \gamma \hat{k}$.
- III. $\frac{\alpha + \gamma}{2} \hat{k}$.
- IV. $\alpha \gamma z \hat{k}$.

Ans. In the previous question, we calculated the force for any \vec{s} as $\vec{F} = \alpha \gamma s_z \hat{k}$. Here,

$s = \frac{1}{2}$, which implies $s_z = \pm \frac{1}{2}$. So, the force is $\pm \frac{1}{2} \alpha \gamma \hat{k}$.

7. Calculate the magnetic energy for electron in the last shell of the Ag atom in an external magnetic field $\vec{B} = \alpha z \hat{k}$.

[Hint: Consider the energy of the only unpaired electron in the outermost shell, whose spin generates a magnetic dipole moment that contributes to the magnetic interaction energy of the Ag atom.]

I. $\mp \gamma \alpha \hbar z / 2$.

II. $\mp \alpha \gamma \hbar z$.

III. $\mp \alpha \gamma \hbar / 2$.

IV. None of the above.

Ans. The answer is **I**. The magnetic energy is given by $-\vec{\mu} \cdot \vec{B}$, where $\vec{\mu}$ is the magnetic moment of the electron given by:

$$\vec{\mu} = \gamma \left[s_x \hat{i} + s_y \hat{j} + s_z \hat{k} \right].$$

Hence, $-\vec{\mu} \cdot \vec{B} = -\gamma \alpha s_z z$. Since $s_z = \pm \frac{\hbar}{2}$, $H = \mp \gamma \alpha \hbar z / 2$, which is the energy of the electron.

Here will be a picture!!!

8. Consider a Stern-Gerlach set-up where neutral Ag atoms are sent across a non-uniform magnetic field $\vec{B} = \alpha z \hat{k}$. The magnetic field is turned off at $t = T$. Then which of the following is the correct expression for the Hamiltonian of the Ag atom due to the interaction with the magnetic field?

- I. $H(t) = \mp \gamma \alpha \hbar z / 2$ all along its path.
- II. $H(t) = \mp \gamma \alpha \hbar z / 2$ when $0 \leq t \leq T$ and 0 otherwise.
- III. $H(t) = 0$ throughout.
- IV. None of the above.

Ans. The answer is **II**. In the previous problem, we have worked out the case of the hamiltonian of magnetic interaction due to external magnetic field, which turns out to be $H(t) = \mp \gamma \alpha \hbar z / 2$. But in this setup, the magnetic field turns off after a time T . So, after this time has elapsed, the Ag atoms behave like free particles with no magnetic interaction due to the absence of the magnetic field. Hence $H(t) = 0$ for $t > T$.

9. If $\chi_+ = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ and $\chi_- = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ are the up and down spin states in the z direction, then in this basis, how would you express a general spin state?

- I. $\chi(t) = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \end{pmatrix}$, where such that $a^2 + b^2 = 1$.
- II. $\chi(t) = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \end{pmatrix}$, where a and b are arbitrary numbers.
- III. $\chi(t) = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \end{pmatrix}$, where such that $a + b = 1$.
- IV. $\chi(t) = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} \times b \begin{pmatrix} 0 \\ 1 \end{pmatrix}$, where such that $a^2 + b^2 = 1$.

Ans. Any general state $\chi(t)$ can be represented as a linear superposition of the basis vectors. But in quantum mechanics, all states are normalized. So, $|\chi(t)|^2$ must come out as 1. Since $\chi(t)$ is a linear superposition of χ_+ and χ_- , we can write it as $\chi(t) = a\chi_+ + b\chi_-$. To normalize this to 1, we must have $a^2 + b^2 = 1$. Hence, choice **I** is correct.

10. A neutral particle sets out in the state $\chi(t = 0) = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \end{pmatrix}$, under the influence of a magnetic field $\vec{B} = \alpha z \hat{k}$. What will be the state of the particle after time t ?

- I. $\chi(t) = (a\chi_+ + b\chi_-) e^{i\frac{\gamma\alpha zt}{2}}.$
- II. $\chi(t) = \left(\frac{a+b}{2}\right) (\chi_+ + \chi_-) e^{i\frac{\gamma\alpha zt}{2}}.$
- III. $\chi(t) = a\chi_+ e^{+i\frac{\gamma\alpha zt}{2}} + b\chi_- e^{-i\frac{\gamma\alpha zt}{2}}.$
- IV. $\chi(t) = a\chi_+ e^{+i\frac{\gamma\alpha zt}{2}} \times b\chi_- e^{-i\frac{\gamma\alpha zt}{2}} = ab\chi_+\chi_-.$

Ans. A quantum state evolves over time according to the *time evolution operator*, which is $e^{-i\frac{\hat{H}}{\hbar}t}$. Now, we have already determined the hamiltonian for such particles in previous problems:

$$\hat{H} = -\gamma\alpha\hat{S}_z z.$$

Hence,

$$\begin{aligned}
\chi(t) &= \mathcal{O}(t)\chi(0), \\
&= e^{-i\frac{\hat{H}}{\hbar}t}\chi(0), \\
&= e^{i\frac{\gamma\alpha}{\hbar}\hat{S}_z zt}\chi(0), \\
&= e^{i\frac{\gamma\alpha}{\hbar}\hat{S}_z zt} (a\chi_+ + b\chi_-), \\
&= a \left(e^{i\frac{\gamma\alpha}{\hbar}zt\hat{S}_z} \chi_+ \right) + b \left(e^{i\frac{\gamma\alpha}{\hbar}zt\hat{S}_z} \chi_- \right), \\
&= a \left(e^{i\frac{\gamma\alpha}{\hbar}zt(+\frac{\hbar}{2})} \chi_+ \right) + b \left(e^{i\frac{\gamma\alpha}{\hbar}zt(-\frac{\hbar}{2})} \chi_- \right), \\
&= a \left(e^{+i\frac{1}{2}\gamma\alpha zt} \chi_+ \right) + b \left(e^{-i\frac{1}{2}\gamma\alpha zt} \chi_- \right), \\
&= a\chi_+ e^{+i\frac{\gamma\alpha zt}{2}} + b\chi_- e^{-i\frac{\gamma\alpha zt}{2}},
\end{aligned}$$

which is choice **III**.

11. In problem 8, we considered the case of a time-varying Hamiltonian, where the particle escapes the magnetic field at $t = T$. What will be $\chi(t)$ when the particle escapes?

- I. $\chi(T) = a\chi_+ e^{+i\frac{\gamma\alpha zT}{2}} + b\chi_- e^{-i\frac{\gamma\alpha zT}{2}}.$
- II. $\chi(T) = 0.$
- III. $\chi(T) = a\chi_+ + b\chi_-.$
- IV. None of the above.

Ans. The choice is **I**, because from the problem where we discussed the time evolution of a general spin state under the influence of a magnetic field (problem 10), we can see that,

$$\chi(t) = a\chi_+e^{+i\frac{\gamma\alpha z t}{2}} + b\chi_-e^{-i\frac{\gamma\alpha z t}{2}}.$$

Now, putting $t = T$ in this expression, we obtain

$$\chi(T) = a\chi_+e^{+i\frac{\gamma\alpha z T}{2}} + b\chi_-e^{-i\frac{\gamma\alpha z T}{2}},$$

which is **I**.

Consider the following conversation between Andy and Caroline :

Andy : The two terms in $\chi(t)$ will be separated spatially because now they carry momentum in \hat{z} direction; the spin up component having momentum $p_+ = +\frac{\gamma\alpha T}{2}\hbar$ and the spin down component carrying momentum $p_- = -\frac{\gamma\alpha T}{2}\hbar$.

Caroline : No, as soon as the particle leaves the magnetic field, the magnetic interaction will stop and the particle will revert back to its initial state $\chi = a\chi_+ + b\chi_-$.

Who is correct?

- I.** Andy is correct.
- II.** Caroline is correct.
- III.** Both are correct.
- IV.** Both are mistaken.

Ans. Andy is correct. The Hamiltonian evolves the wave function while the particle interacts with the magnetic field, whereby the particles gain momentum in the \hat{z} direction, which results in spatial separation of the particles, as described in the picture.

To determine the momentum in \hat{z} direction, we can simply operate on the state with the momentum operator, which is $\frac{\hbar}{i} \frac{\partial}{\partial z}$. Initially, we have,

$$\chi(t=0) = a\chi_+ + b\chi_-.$$

So,

$$\begin{aligned}\hat{p}_z\chi(0) &= \hat{p}_z(a\chi_+ + b\chi_-) \\ &= \frac{\hbar}{i} \frac{\partial}{\partial z} (a\chi_+ + b\chi_-) \\ &= 0,\end{aligned}$$

since χ_{\pm} are not dependent on z . Hence, it can be seen that initially the particles were not carrying any momentum in the \hat{z} direction. But at time $t = T$, let us determine the momentum:

$$\begin{aligned}\hat{p}_z\chi(T) &= \hat{p}_za\chi_+e^{+i\frac{\gamma\alpha zT}{2}} + b\chi_-e^{-i\frac{\gamma\alpha zT}{2}} \\ &= \frac{\hbar}{i} \frac{\partial}{\partial z} \left(a\chi_+e^{+i\frac{\gamma\alpha zT}{2}} + b\chi_-e^{-i\frac{\gamma\alpha zT}{2}} \right) \\ &= \frac{\hbar}{i} \cdot +i\frac{\gamma\alpha T}{2} a\chi_+e^{+i\frac{\gamma\alpha zT}{2}} \\ &\quad + \frac{\hbar}{i} \cdot -i\frac{\gamma\alpha T}{2} b\chi_-e^{-i\frac{\gamma\alpha zT}{2}} \\ &= \left(+\frac{\gamma\alpha\hbar T}{2} \right) a\chi_+e^{+i\frac{\gamma\alpha zT}{2}} + \left(-\frac{\gamma\alpha\hbar T}{2} \right) b\chi_-e^{-i\frac{\gamma\alpha zT}{2}}.\end{aligned}$$

Hence, the up component has momentum $\left(+\frac{\gamma\alpha\hbar T}{2}\right)$, and the down component has momentum $\left(-\frac{\gamma\alpha\hbar T}{2}\right)$.

12. For $s = \frac{1}{2}$ the beam splits in two parts. If we use particles with $s = \frac{3}{2}$ instead, in how many parts will the beam split?

- I. 3.
- II. 2.
- III. 4.
- IV. 5.

Ans. For spin s , there can be $2s + 1$ values for s_z . So, for $s = \frac{1}{2}$, we have $2 \times \frac{1}{2} + 1 = 2$ values. Hence, for $s = \frac{3}{2}$, we have $2 \times \frac{3}{2} + 1 = 4$ values for s_z . Now, each value of spin corresponds to a different energy value, and thus a different momentum. So we shall obtain 4 spots for spin $\frac{3}{2}$ (choice **III**).

13. In many multi-electron atoms, it is possible to get $s = 0$, where the total spin adds up to 0. In such a case, how many spots do you expect to get, for the same system we have been using?

- I. 1.
- II. 4.
- III. 2.
- IV. 3.

Ans. The answer is **I**. This is apparent from the logic of the previous problem: $2 \times 0 + 1 = 1$. But, there is also a simple physical interpretation of the answer. As the spin is zero, the magnetic moment due to spin is also zero. Hence there is no spin-magnetic field interaction that will separate the particles.

14. For $s = 1$, how many spots do you expect?

- I. 1.
- II. 2.
- III. 3.

IV. 4.

Ans. The answer is $2 \times 1 + 1 = 3$. So choice **III** is correct.

You can see that we are getting odd number of spots only for integral values of s , and we are getting even number of spots for half-integral values of s . This can be explained in the following way: Since angular momentum (spin or orbital) can only change in integral steps, there will be odd number of possible values of s_z for integral s , and even number of possible values for half integral s . Now, since orbital angular momentum is always integral, we can expect even number of spots in a Stern-Gerlach experiment due to angular momentum. Whereas in the Stern-Gerlach experiment with Ag atom (which first led to the discovery of spin), we obtain 2 (even) spots. This mystery was incompatible with the orbital angular momentum model. Then, Goudsmit and Uhlenbeck's conjecture of half-integral spin angular momentum properly explained this mystical phenomena of even splitting.

But apart from this, there is another important aspect of the Stern-Gerlach experiment. This experiment has played a pivotal role in explaining the philosophy of quantum mechanics regarding preparation of quantum states, as well as the importance of the process of measurement. We shall address the following issues related to preparation of quantum states and measurement:

- Preparation of a quantum system in a certain quantum state.
- Difference between state preparation and measurement.
- Compatible and incompatible observables.
- Effect of measurement on future evolution of the system.

Insert image: Notations used and descriptions.

Consider *silver* atoms which have $|\chi\rangle = |\uparrow\rangle_z = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ ie. eigenstates of \hat{S}_z with eigenvalue $\frac{\hbar}{2}$.

15. Consider the following situation:

- I.** The up detector never clicks.
- II.** The up detector clicks only 50% of the time.
- III.** the up detector clicks all the time.
- IV.** Can not say.

Ans. III. Since $|\uparrow\rangle_z$ is an eigenstate of \hat{S}_z and the Stern-Gerlach device has also a field gradient in \hat{z} direction, the state is still an eigenstate of \hat{S}_z after passing through the Stern-Gerlach device, which evolves the state of the particle according to $e^{-i\frac{\hat{H}}{\hbar}t}$, where \hat{H} is simply $-\gamma\alpha\hat{S}_zz$. Therefore, in this case, the *up* detector will click all the time, and the *down* detector will never click at all.

16. Consider the following setup and determine which of the following propositions is correct:

- I.** The up detector never clicks.
- II.** the up detector clicks half of the time.
- III.** the up detector clicks all the time.
- IV.** Can not say.

Ans. Choice **II** is correct. $|\uparrow\rangle_z$ is not an eigenstate of \hat{S}_x operator, whereas the Stern-Gerlach apparatus used in this case evolves the state with a function of the \hat{S}_x operator (see problem 15). But, $|\uparrow\rangle_z$ can be written as a linear superposition of the eigenstates of \hat{S}_x in the following way:

$$|\uparrow\rangle_z = \frac{1}{\sqrt{2}} (|\uparrow\rangle_x + |\downarrow\rangle_x).$$

Therefore, the probability that the *up* detector clicks is just $\left(\frac{1}{\sqrt{2}}\right)^2 = \frac{1}{2}$. Hence, the *up* detector will click half of the time.

- 17.** Consider the following setup and determine which of the following propositions is correct:

- I.** The up detector never clicks.
- II.** the up detector clicks half of the time.
- III.** the up detector clicks all the time.
- IV.** Can not say.

Ans. Again choice **II** is correct, as $|\uparrow\rangle_z$ not being an eigenstate of the \hat{S}_y operator, it can be written as $|\uparrow\rangle_z = \frac{1}{\sqrt{2}} (|\uparrow\rangle_y + |\downarrow\rangle_y)$. Then, there is 50% chance that the *up* detector will click.

Since the probabilities of getting *up* or *down* states are equal in the two cases considered above, it is impossible to know a-priori which of the detectors will click, i.e. to which state the initial state will *collapse* after going through the Stern-Gerlach apparatus.

18. Consider the following case and determine which of the following is correct:

- I.** The up detector never clicks.
- II.** the up detector clicks half of the time.
- III.** the up detector clicks all the time.
- IV.** Can not say.

Ans. The first SG_X makes $|\uparrow\rangle_z = \frac{1}{\sqrt{2}} (|\uparrow\rangle_y + |\downarrow\rangle_y)$. Since the second SG_X has opposite gradient, it recombines the states to $|\uparrow\rangle_z$. The third Stern-Gerlach instrument imposes $\mathcal{O}(\hat{S}_z)$ on the state, so the state is unaltered and we obtain 100% clicks in the *up* z detector, making choice **III** correct.

Andy : I disagree with the answer to the previous question. Since the first SG_X splits $|\uparrow\rangle_z$ into a spatially separated superposition of $|\uparrow\rangle_x$ and $|\downarrow\rangle_x$, the second SG_X , which is

simply imposing $\mathcal{O}(\hat{S}_x)$ on the states, can not recombine them to $|\uparrow\rangle_z$ again, as those states are eigenstates of the \hat{S}_x operator, and the initial interaction is irreversible.

Caroline : No, when the first Stern-Gerlach instrument splits $|\uparrow\rangle_z$ into $|\uparrow\rangle_x$ and $|\downarrow\rangle_x$, the process is only a time evolution of the state, not a measurement. I think only measurements alter the states irreversibly. So we CAN do the inverse transform with an opposite SG_X machine and recombine the states back to $|\uparrow\rangle_z$.

19. Which of the following statements are correct?

- I.** Andy is correct.
- II.** Caroline is correct.
- III.** Both are correct.
- IV.** Both are mistaken.

Ans. Caroline is correct. In the previous case, nowhere a measurement was made after passing the beam through the SG_X apparatuses to determine the state in which the particle going through the devices were. Hence, the initial states were not altered, and it was a time evolution, but not a measurement.

Preparation of states:

Consider the above setup, and choose which of the following states you shall obtain in the particle beam:

- I.** $|\downarrow\rangle_z$.

II. $|\uparrow\rangle_x$.

III. $|\downarrow\rangle_x$.

IV. $\frac{1}{\sqrt{2}}(|\downarrow\rangle_x + |\uparrow\rangle_x)$.

V. $\frac{1}{2\sqrt{2}}(|\uparrow\rangle_x + |\downarrow\rangle_x)$.

Ans. Since the detector measures for $|\downarrow\rangle_x$, the moment the measurement is done, $|\uparrow\rangle_z$ collapses to either $|\uparrow\rangle_x$ or $|\downarrow\rangle_x$ with equal probability, and whenever $|\downarrow\rangle_x$ is found, it is trapped in the detector. So the obtained beam contains only $|\uparrow\rangle_x$ states. So the answer is choice **II**.

Thus, we can produce pure states with the aid of a Stern-Gerlach apparatus.

20. Consider the following setup:

What do we obtain as the final state?

I. $|\uparrow\rangle_z$.

II. $|\uparrow\rangle_x$.

III. $|\downarrow\rangle_x$.

IV. $|\downarrow\rangle_y$.

Ans. The answer is **IV**, because of the reasons similar to the ones discussed in the previous problem.

Consider the following conversation between Andy and Caroline :

Andy : : The $|\uparrow\rangle_z$ state is a linear superposition of $|\uparrow\rangle_x$ and $|\downarrow\rangle_x$, and it is not an eigenstate of $|\uparrow\rangle_x$.

Caroline : : Exactly, since \hat{S}_x and \hat{S}_z do not commute with each other, they are incompatible operators and they do not have simultaneous eigenstates.

21. Which of the following is true?

- I.** Andy is correct but Caroline is wrong.
- II.** Andy is wrong but Caroline is correct.
- III.** Both are wrong.
- IV.** Both are correct.

Ans. Both are correct.

22. Consider the Stern-Gerlach set-up:

What is the possible output of the SG_X ?

I.

II.

III. None of the above.

Ans. This Stern-Gerlach instrument will produce $|\uparrow\rangle_x$ and $|\downarrow\rangle_x$ states, with equal probability. Hence, the answer is **I**.

23. Consider the following setup:

What is the output in this case?

- I.** $|\uparrow\rangle_x$ and $|\downarrow\rangle_x$.
- II.** $|\uparrow\rangle_x$ only.
- III.** $|\downarrow\rangle_x$ only.
- IV.** $|\uparrow\rangle_z$ only.

Ans. The answer is **II**. The detector detects the $|\downarrow\rangle_x$ states, so only $|\uparrow\rangle_x$ states enter the second Stern-Gerlach instrument. Since $|\uparrow\rangle_x$ is an eigenstate of the evolution operator imposed due to the second Stern-Gerlach instrument, the state is unaltered after passing through the instrument, and we obtain $|\uparrow\rangle_x$ as the final state.

24. Consider the following setup and select the right output:

- I.** $|\uparrow\rangle_x$.
- II.** $|\uparrow\rangle_z$ and $|\downarrow\rangle_z$.
- III.** $|\uparrow\rangle_x$ and $|\downarrow\rangle_x$. None of the above.

Ans. The correct answer is **III**. The first SG_X instrument splits the incident $|\uparrow\rangle_z$ beam into $|\uparrow\rangle_x$ and $|\downarrow\rangle_x$ beams with half intensity of the initial beam for each. The $|\downarrow\rangle_x$ state

is blocked in the detector, so certainly the $|\uparrow\rangle_x$ state enters the second Stern-Gerlach instrument, which imposes an evolution operator to which $|\uparrow\rangle_x$ is an eigenstate, so the beam leaves this instrument unaltered in state, and the final SG_Z apparatus splits the incident $|\uparrow\rangle_x$ beam equally into $|\uparrow\rangle_z$ and $|\downarrow\rangle_z$, with 25% of the intensity of the actual $|\uparrow\rangle_z$ beam each.

25. Now consider a silver atom beam in a spin state $\langle\chi| = a|\uparrow\rangle_z + b|\downarrow\rangle_z$, with $a^2 + b^2 = 1$ incident on the $-SG_Z$ apparatus. Then which of the following is true?

- I. The up detector clicks 100% of the time.
- II. The up detector clicks 50% of the time, and the down detector clicks 50% of the time.
- III. The down detector clicks 100% of the time.
- IV. The up detector clicks $100a^2\%$ of the time and the down detector clicks $100b^2\%$ of the time.

Ans. The correct answer is **IV**. The probability of getting $|\uparrow\rangle_z$ from $\langle\chi|$ is a^2 , so percentage wise it gives $100a^2$, similarly it turns out as $100b^2$ for $|\downarrow\rangle_z$.

26. Consider the following setup:

The output is:

- I. $|\uparrow\rangle_z$.
- II. $|\downarrow\rangle_z$.
- III. $|\uparrow\rangle_z$ and $|\downarrow\rangle_z$.
- IV. None of the above.

Ans. The correct choice is **I**. The detector blocks the $|\downarrow\rangle_z$, so a pure $|\uparrow\rangle_z$ state is produced. This can be extended to more general cases as well.

Do it yourself:

Consider the Stern-Gerlach setup with silver atoms in the $|\uparrow\rangle_y$ state, and determine the outcome. Also give adequate explanation: