

Estimation and Control Library

Pan

Abstract—A library for optimal estimation and control, as well as classical/modern control theories. Algorithms are applied on simulated robots using Ignition Gazebo.

I. INTRODUCTION

A library for optimal estimation and control, as well as classical/modern control theories. Algorithms are applied on simulated robots using ROS [1] and Ignition Gazebo [2].

II. CARPOLE DYNAMICS

This section derives the equation of motion of the cartpole system in the following two ways: (1) standard vector analysis; (2) using Lagrangian.

A. Standard vector analysis

The position of the cart is given as

$$\mathbf{r}_1 = x\mathbf{E}_x \quad (1)$$

Computing the 1st and 2nd order rate of change of \mathbf{r}_1 in reference frame \mathcal{F}_i , we obtain the velocity and acceleration as

$$\mathbf{v}_1 = {}^{\mathcal{F}_i} \frac{d}{dt}(\mathbf{r}_1) = \dot{x}\mathbf{E}_x \quad (2)$$

$$\mathbf{a}_1 = {}^{\mathcal{F}_i} \frac{d}{dt}(\mathbf{v}_1) = \ddot{x}\mathbf{E}_x \quad (3)$$

The position of the pole is given as

$$\mathbf{r}_2 = x\mathbf{E}_x + l\mathbf{e}_r \quad (4)$$

Similarly, the velocity and acceleration of the pole can be computed as

$$\mathbf{v}_2 = {}^{\mathcal{F}_i} \frac{d}{dt}(\mathbf{r}_2) = \dot{x}\mathbf{E}_x + l\dot{\theta}\mathbf{e}_\theta \times \mathbf{e}_r = \dot{x}\mathbf{E}_x - l\dot{\theta}\mathbf{e}_z \quad (5)$$

$$\mathbf{a}_2 = {}^{\mathcal{F}_i} \frac{d}{dt}(\mathbf{v}_2) = \ddot{x}\mathbf{E}_x - l\ddot{\theta}\mathbf{e}_z - l\dot{\theta}^2\mathbf{e}_r \quad (6)$$

Rewritten \mathbf{a}_2 in \mathbf{E}_x and \mathbf{E}_z as

$$\mathbf{a}_2 = (\ddot{x} + l\ddot{\theta}\cos\theta - l\dot{\theta}^2\sin\theta)\mathbf{E}_x - (l\ddot{\theta}\sin\theta + l\dot{\theta}^2\cos\theta)\mathbf{E}_z \quad (7)$$

Applying Newton 2nd law on the \mathbf{E}_x -direction of the cart and pole, we have

$$F - T_x = m_1\ddot{x} \quad (8)$$

$$T_x = m_2(\ddot{x} + l\ddot{\theta}\cos\theta - l\dot{\theta}^2\sin\theta) \quad (9)$$

Adding the above two equations we get the first equation,

$$F = (m_1 + m_2)\ddot{x} + m_2(l\ddot{\theta}\cos\theta - l\dot{\theta}^2\sin\theta) \quad (10)$$

The cartpole system is consisted of two systems, i.e. the cart and the pole. The internal force can be canceled out by analyzing the two systems as a whole. A simpler way is to apply the balance of angular momentum relative to the cart.

$$\frac{d}{dt}\mathbf{H}_c = \mathbf{M}_c - (\mathbf{r}_2 - \mathbf{r}_1) \times m_2\mathbf{a}_1 \quad (11)$$

where \mathbf{H}_c is the angular momentum of the cartpole system relative to the cart.

$$\mathbf{H}_c = (\mathbf{r}_2 - \mathbf{r}_1) \times m_2(\mathbf{v}_2 - \mathbf{v}_1) = m_2l^2\dot{\theta}\mathbf{E}_y \quad (12)$$

And \mathbf{M}_c is the moment relative to the cart.

$$\mathbf{M}_c = (\mathbf{r}_2 - \mathbf{r}_1) \times (-m_2g)\mathbf{E}_z = m_2gl\sin\theta\mathbf{E}_y \quad (13)$$

$$(\mathbf{r}_2 - \mathbf{r}_1) \times m_2\mathbf{a}_1 = l\mathbf{e}_r \times m_2\ddot{x}\mathbf{E}_x = m_2l\ddot{x}\cos\theta\mathbf{E}_y \quad (14)$$

Thus, we get the second equation

$$\cos\theta\ddot{x} + l\ddot{\theta} = g\sin\theta \quad (15)$$

In summary, the equation of motion is

$$(m_1 + m_2)\ddot{x} + m_2(l\ddot{\theta}\cos\theta - l\dot{\theta}^2\sin\theta) = F \quad (16)$$

$$\cos\theta\ddot{x} + l\ddot{\theta} = g\sin\theta \quad (17)$$

B. Linearization

Assuming $\mathbf{q} = [x, \theta]^T$, the equations of motion of the cartpole system can be written in the standard form as

$$\mathbf{H}(\mathbf{q})\ddot{\mathbf{q}} + \mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})\dot{\mathbf{q}} + \mathbf{G}(\mathbf{q}) = \mathbf{B}\mathbf{u} \quad (18)$$

where $\mathbf{q} = [x, \theta]^T$ is a n -vector called the *generalized coordinates vector*, $\mathbf{H}(\mathbf{q})$ is a $n \times n$ nonsingular symmetric positive-definite matrix called the *mass matrix*, $\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})$ is a $n \times n$ matrix called the *centrifugal/Coriolis/friction matrix*, $\mathbf{G}(\mathbf{q})$ is a n -vector sometimes called the *conservative forces vector*. **This equation of motion is valid for systems that follow classical Newton-Euler mechanics or Lagrangian mechanics** with a kinetic energy that is quadratic in the derivative of the generalized coordinates and a potential energy that may depend on the generalized coordinates, but not on its derivative. Such systems include robot arms, mobile robots, airplanes, helicopters, underwater vehicles, hovercraft, etc.

$$\mathbf{H}(\mathbf{q}) = \begin{bmatrix} m_1 + m_2 & m_2l\cos\theta \\ \cos\theta & l \end{bmatrix}, \mathbf{C}(\mathbf{q}, \dot{\mathbf{q}}) = \begin{bmatrix} 0 & -m_2l\dot{\theta}\sin\theta \\ 0 & 0 \end{bmatrix},$$

$$\mathbf{G}(\mathbf{q}) = \begin{bmatrix} 0 \\ -g\sin\theta \end{bmatrix}, \mathbf{B} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}, \quad (19)$$

¹ The Pennsylvania State University, University Park, PA 16802, USA. pan.liu@psu.edu

Let $x_1 = \mathbf{q}$, $x_2 = \dot{\mathbf{q}}$, and $x = [x_1, x_2]^T$. Then we have the **standard form for linearization**, i.e. $\dot{x} = f(x, u)$

$$\dot{x} = \begin{bmatrix} x_2 \\ \mathbf{H}(\mathbf{q})^{-1}(-\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})\dot{\mathbf{q}} - \mathbf{G}(\mathbf{q}) + \mathbf{B}u) \end{bmatrix} \quad (20)$$

Linearize around the equilibrium point $x^* = 0$,

$$\dot{x} = Ax + Bu \quad (21)$$

where

$$A = \begin{bmatrix} \mathbf{0} & \mathbf{I} \\ -\mathbf{H}(\mathbf{q})^{-1} \frac{\partial}{\partial \mathbf{q}} \mathbf{G}(\mathbf{q}) & -\mathbf{H}(\mathbf{q})^{-1} \mathbf{C}(\mathbf{q}, \dot{\mathbf{q}}) \end{bmatrix}, \quad (22)$$

$$B = \begin{bmatrix} 0 \\ \mathbf{H}^{-1} \mathbf{B} \end{bmatrix} \quad (23)$$

Note that the term involving $\frac{\partial}{\partial \mathbf{q}} \mathbf{H}(\mathbf{q})^{-1}$ disappears because $\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})\dot{\mathbf{q}} - \mathbf{G}(\mathbf{q}) + \mathbf{B}u$ must be zero at the fixed point. Many of the $\mathbf{C}(\mathbf{q}, \dot{\mathbf{q}})\dot{\mathbf{q}}$ derivatives drop out, too, because $\dot{\mathbf{q}}^* = 0$. The $-\mathbf{H}(\mathbf{q})^{-1}$ can be calculated by Matlab symbolic inversion,

$$\mathbf{H}(\mathbf{q})^{-1} = \begin{bmatrix} \frac{1}{m_1 + m_2 - m_2 \cos^2 \theta} & \frac{-m_2 \cos \theta}{m_1 + m_2 - m_2 \cos^2 \theta} \\ \frac{-\cos \theta}{m_1 l + m_2 l - m_2 l \cos^2 \theta} & \frac{m_1 + m_2}{m_1 + m_2 - m_2 \cos^2 \theta} \end{bmatrix} \quad (24)$$

and

$$\frac{\partial}{\partial \mathbf{q}} \mathbf{G}(\mathbf{q}) = \begin{bmatrix} 0 & 0 \\ 0 & -g \cos \theta \end{bmatrix} \quad (25)$$

Substitute $x^* = 0$ and rearrange x to be $x = [x, \dot{x}, \theta, \dot{\theta}]$, we have

$$A = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & -\frac{m_2}{m_1} g & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & \frac{m_1 + m_2}{m_1 l} g & 0 \end{bmatrix}, B = \begin{bmatrix} 0 \\ 1/m \\ 0 \\ -1/(m_1 l) \end{bmatrix} \quad (26)$$

REFERENCES

- [1] M. Quigley, K. Conley, B. Gerkey, J. Faust, T. Foote, J. Leibs, R. Wheeler, and A. Y. Ng, "Ros: an open-source robot operating system," in *ICRA workshop on open source software*, vol. 3, no. 3.2. Kobe, Japan, 2009, p. 5.
- [2] I. Gazebo, "https://ignitionrobotics.org."