

Uses my GA30.m package to compute the values of the energy momentum tensor multivectors, and relate those to the conventional tensor description of the same.

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<< GA30` ;
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## Calculation of components of the energy momentum tensor $T(a) = F a F^\dagger \epsilon/2$

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ClearAll[e1, e2, e3, i, vecE, vecH,
  one, f, fDagger, vecJ, vecM, j, jDagger, cross]
e1 = Vector[1, 1];
e2 = Vector[1, 2];
e3 = Vector[1, 3];
i = e1 ** e2 ** e3;
one = Scalar[1];

vecE = "E"_1 e1 + "E"_2 e2 + "E"_3 e3;
vecH = "H"_1 e1 + "H"_2 e2 + "H"_3 e3;
vecJ = "J"_1 e1 + "J"_2 e2 + "J"_3 e3;
vecM = "M"_1 e1 + "M"_2 e2 + "M"_3 e3;
j =  $\eta$  (c  $\rho$  - vecJ) + i ** (c  $\rho_m$  - vecM);
jDagger =  $\eta$  (c  $\rho$  - vecJ) - i ** (c  $\rho_m$  - vecM);
cross[a_, b_] := -i ** (a  $\wedge$  b);

ClearAll[v, vv]
(*t[v_] :=  $\epsilon$  ( vecE ** v ** vecE + vecH ** v ** vecH +
  i ** ( vecH ** v ** vecE - vecE ** v ** vecH))/2*)
f := vecE + i ** vecH ( $\eta$ );
fDagger := vecE - i ** vecH ( $\eta$ );
t[v_] := ( $\epsilon/2$ ) ( f ** v ** fDagger)
vv = { t[one], t[e1], t[e2], t[e3] };

(*vv[[1]] // TexForm
  vv[[2]]
  vv[[3]]
  vv[[4]]
*)

Grid[ { Subscript["T", #], " = ", vv[[# + 1]]} & /@ (Range[4] - 1)]

$$T_0 = \frac{1}{2} \in (E_1^2 + E_2^2 + E_3^2 + \eta^2 H_1^2 + \eta^2 H_2^2 + \eta^2 H_3^2) +$$


$$\in \eta (-E_3 H_2 + E_2 H_3) \mathbf{e}_1 + \in \eta (E_3 H_1 - E_1 H_3) \mathbf{e}_2 + \in \eta (-E_2 H_1 + E_1 H_2) \mathbf{e}_3$$


$$T_1 = \in \eta (E_3 H_2 - E_2 H_3) + \frac{1}{2} \in (E_1^2 - E_2^2 - E_3^2 + \eta^2 H_1^2 - \eta^2 H_2^2 - \eta^2 H_3^2) \mathbf{e}_1 +$$


$$\in (E_1 E_2 + \eta^2 H_1 H_2) \mathbf{e}_2 + \in (E_1 E_3 + \eta^2 H_1 H_3) \mathbf{e}_3$$


$$T_2 = \in \eta (-E_3 H_1 + E_1 H_3) + \in (E_1 E_2 + \eta^2 H_1 H_2) \mathbf{e}_1 -$$


$$\frac{1}{2} \in (E_1^2 - E_2^2 + E_3^2 + \eta^2 H_1^2 - \eta^2 H_2^2 + \eta^2 H_3^2) \mathbf{e}_2 + \in (E_2 E_3 + \eta^2 H_2 H_3) \mathbf{e}_3$$


$$T_3 = \in \eta (E_2 H_1 - E_1 H_2) + \in (E_1 E_3 + \eta^2 H_1 H_3) \mathbf{e}_1 +$$


$$\in (E_2 E_3 + \eta^2 H_2 H_3) \mathbf{e}_2 - \frac{1}{2} \in (E_1^2 + E_2^2 - E_3^2 + \eta^2 H_1^2 + \eta^2 H_2^2 - \eta^2 H_3^2) \mathbf{e}_3$$


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Same calculation for the adjoint of the energy momentum tensor  $\overline{T}(a)$   
 $= F^\dagger a F \epsilon / 2$

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ClearAll[t2, vv2]
t2[v_] := (ϵ / 2) ( (vecE - i ** vecH (η)) ** v ** (vecE + i ** vecH (η)) )
vv2 = { t2[one], t2[e1], t2[e2], t2[e3] };

Grid[ { Subscript["T", #], " = ", vv2[[# + 1]] } & /@ (Range[4] - 1)]
Grid[ { Subscript["T", #], " = ", vv[[# + 1]] - vv2[[# + 1]] } & /@ (Range[4] - 1)]

T0 =  $\frac{1}{2} \in (E_1^2 + E_2^2 + E_3^2 + \eta^2 H_1^2 + \eta^2 H_2^2 + \eta^2 H_3^2) +$   

 $\in \eta (E_3 H_2 - E_2 H_3) \mathbf{e}_1 + \in \eta (-E_3 H_1 + E_1 H_3) \mathbf{e}_2 + \in \eta (E_2 H_1 - E_1 H_2) \mathbf{e}_3$ 
T1 =  $\in \eta (-E_3 H_2 + E_2 H_3) + \frac{1}{2} \in (E_1^2 - E_2^2 - E_3^2 + \eta^2 H_1^2 - \eta^2 H_2^2 - \eta^2 H_3^2) \mathbf{e}_1 +$   

 $\in (E_1 E_2 + \eta^2 H_1 H_2) \mathbf{e}_2 + \in (E_1 E_3 + \eta^2 H_1 H_3) \mathbf{e}_3$ 
T2 =  $\in \eta (E_3 H_1 - E_1 H_3) + \in (E_1 E_2 + \eta^2 H_1 H_2) \mathbf{e}_1 -$   

 $\frac{1}{2} \in (E_1^2 - E_2^2 + E_3^2 + \eta^2 H_1^2 - \eta^2 H_2^2 + \eta^2 H_3^2) \mathbf{e}_2 + \in (E_2 E_3 + \eta^2 H_2 H_3) \mathbf{e}_3$ 
T3 =  $\in \eta (-E_2 H_1 + E_1 H_2) + \in (E_1 E_3 + \eta^2 H_1 H_3) \mathbf{e}_1 +$   

 $\in (E_2 E_3 + \eta^2 H_2 H_3) \mathbf{e}_2 - \frac{1}{2} \in (E_1^2 + E_2^2 - E_3^2 + \eta^2 H_1^2 + \eta^2 H_2^2 - \eta^2 H_3^2) \mathbf{e}_3$ 

T0 =  $2 \in \eta (-E_3 H_2 + E_2 H_3) \mathbf{e}_1 + 2 \in \eta (E_3 H_1 - E_1 H_3) \mathbf{e}_2 + 2 \in \eta (-E_2 H_1 + E_1 H_2) \mathbf{e}_3$ 
T1 =  $2 \in \eta (E_3 H_2 - E_2 H_3)$ 
T2 =  $2 \in \eta (-E_3 H_1 + E_1 H_3)$ 
T3 =  $2 \in \eta (E_2 H_1 - E_1 H_2)$ 
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tt = (ε / 2) (fDagger ** j + jDagger ** f)
tt2 = - ε η (⟨vecE, vecJ⟩ + ⟨vecH, vecM⟩) +
      ε c η (ρ vecE + ρm vecH) + ε cross[vecE, vecM] + ε η2 cross[vecJ, vecH]
tt - tt2
Simplify[tt2 /. {η → Sqrt[μ / ε], c -> Sqrt[1 / (μ ε)]}, Assumptions → μ > 0 && ε > 0]
-ε η (E1 J1 + E2 J2 + E3 J3 + H1 M1 + H2 M2 + H3 M3) +
ε ⎛ η √(1 / (ε μ)) ρ E1 + η2 H3 J2 - η2 H2 J3 - E3 M2 + E2 M3 + η √(1 / (ε μ)) H1 ρm ⎞ e1 +
ε ⎛ η √(1 / (ε μ)) ρ E2 - η2 H3 J1 + η2 H1 J3 + E3 M1 - E1 M3 + η √(1 / (ε μ)) H2 ρm ⎞ e2 +
ε ⎛ η √(1 / (ε μ)) ρ E3 + η2 H2 J1 - η2 H1 J2 - E2 M1 + E1 M2 + η √(1 / (ε μ)) H3 ρm ⎞ e3
-ε η (E1 J1 + E2 J2 + E3 J3 + H1 M1 + H2 M2 + H3 M3) +
ε ⎛ η2 (H3 J2 - H2 J3) - E3 M2 + E2 M3 + η √(1 / (ε μ)) (ρ E1 + H1 ρm) ⎞ e1 +
ε ⎛ η √(1 / (ε μ)) ρ E2 - η2 H3 J1 + η2 H1 J3 + E3 M1 - E1 M3 + η √(1 / (ε μ)) H2 ρm ⎞ e2 +
ε ⎛ η √(1 / (ε μ)) ρ E3 + η2 H2 J1 - η2 H1 J2 - E2 M1 + E1 M2 + η √(1 / (ε μ)) H3 ρm ⎞ e3
0
-√(ε μ) (E1 J1 + E2 J2 + E3 J3 + H1 M1 + H2 M2 + H3 M3) +
(ρ E1 + μ H3 J2 - μ H2 J3 - ε E3 M2 + ε E2 M3 + H1 ρm) e1 +
(ρ E2 - μ H3 J1 + μ H1 J3 + ε E3 M1 - ε E1 M3 + H2 ρm) e2 +
(ρ E3 + μ H2 J1 - μ H1 J2 - ε E2 M1 + ε E1 M2 + H3 ρm) e3

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