

## SVT Boundary Conditions

Taking  $h_{0i}$  as an example, if we decompose  $h_{0i}$  according to the identity

$$h_{0i} = \underbrace{h_{0i} - \nabla_i \int D\nabla^j h_{0j}}_{B_i} + \nabla_i \underbrace{\int D\nabla^j h_{0j}}_B, \quad (0.1)$$

then naturally it follows that

$$\nabla^k B_k = 0, \quad \nabla^k h_{0k} = \nabla^k \nabla_k B. \quad (0.2)$$

Thus a decomposition into transverse and longitudinal components appears to have been satisfied, and proper behavior of  $B$  only requires that  $\int D\nabla^j h_{ij}$  does not diverge.

By directly taking  $B = \int D\nabla^j h_{0j}$ , it would seem we do not have to specify any specific behavior of  $B$  on the boundary to achieve decomposition.

However, in using Green's identity

$$(\nabla_i \nabla^i D)B = D(\nabla_i \nabla^i B) + \nabla^i [(\nabla_i D)B - D(\nabla_i B)] \quad (0.3)$$

$$B = \int D(\nabla_i \nabla^i B) + \oint dS^i [(\nabla_i D)B - D(\nabla_i B)], \quad (0.4)$$

we see that since  $\nabla_i \nabla^i B = \nabla^i h_{0i}$  and since we have defined  $B = \int D\nabla^j h_{0j}$  it must follow that

$$B = B + \oint dS^i [(\nabla_i D)B - D(\nabla_i B)] \quad (0.5)$$

$$\implies 0 = \oint dS^i [(\nabla_i D)B - D(\nabla_i B)]. \quad (0.6)$$

Put in different terms, Green's identity allows us to decompose any scalar into harmonic and non-harmonic parts (harmonic according to math convention, meaning  $\nabla^2 f(x) = 0$ )

$$B = \underbrace{\int D(\nabla_i \nabla^i B)}_{B^{NH}} + \underbrace{\oint dS^i [(\nabla_i D)B - D(\nabla_i B)]}_{B^H}. \quad (0.7)$$

We see that our initial definition of  $B$  via  $B = \int D\nabla^j h_{0j}$  is in fact the non-harmonic projection  $B^{NH}$  by virtue of  $\nabla^2 B = \nabla^j h_{0j}$ . Thus such a definition automatically enforces that  $B$  vanish on the boundary.

Given the integral relation for  $\psi$  (according to a definition that never requires an integration by parts) and its derivative relation,

$$\psi = \frac{1}{4} \left[ \int D\nabla^l h_{kl} - g^{ab} h_{ab} \right] \quad (0.8)$$

$$\nabla^2 \psi = \frac{1}{4} \nabla^2 [\nabla^l h_{kl} - g^{ab} h_{ab}], \quad (0.9)$$

we see that  $\psi \neq \int D\nabla^2\psi$ , and thus is not required to vanish on the boundary. However, if we define the integral relation for  $\psi$  as

$$\psi = \frac{1}{4} \left[ \int D(\nabla^l h_{kl} - g^{ab} h_{ab}) \right], \quad (0.10)$$

then it does follow that  $\psi$  is nonharmonic, namely  $\psi = \int D\nabla^2\psi$  and thus must vanish on the boundary.

The tradeoff between having an  $E_{ij}$  that is automatically transverse + traceless but not itself gauge invariant, vs an  $E_{ij}$  that is automatically gauge invariant but requires integration by parts to be transverse and traceless is discussed below (2.9) in the decomposition paper.