

Electrodynamics II

HW 4

Matthew Phelps

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1. A particle moving with velocity V dissociates “in flight” into two particles. Determine the relation between the angles of emergence of these particles and their energies.

Define θ as the angle of emergence in the Lab Frame, and \mathcal{E}_0 energy in c.o.m. frame and \mathcal{E} as lab frame energy. Then ($c = 1$ units)

$$\mathcal{E}_0 = \frac{\mathcal{E} - Vp \cos \theta}{\sqrt{1 - V^2}}$$

$$\Rightarrow \cos \theta = \frac{\mathcal{E} - \mathcal{E}_0 \sqrt{1 - V^2}}{V \sqrt{\mathcal{E}^2 - m^2}} = \frac{\mathcal{E} - \mathcal{E}_0 \sqrt{1 - V^2}}{Vp}$$

For $m_1 = m_2$ then $\theta_1 = \theta_2$ in the L.F. In c.o.m. frame, particles are separated by π (for $m_1 = m_2$). Taking our original equation

$$\mathcal{E}^2(1 - V^2 \cos^2 \theta) - 2\mathcal{E}\mathcal{E}_0\sqrt{1 - v^2} + \mathcal{E}_0^2(1 - V^2) + V^2m^2 \cos^2 \theta = 0$$

solving for \mathcal{E}

$$\mathcal{E} = \frac{2\mathcal{E}_0\sqrt{1 - V^2} \pm 2V \cos \theta \sqrt{\mathcal{E}_0^2(1 - V^2) - m^2(1 - V^2 \cos^2 \theta)}}{2(1 - V^2 \cos^2 \theta)}.$$

2. For the collision of two particles of equal mass m , express \mathcal{E}'_1 , \mathcal{E}' , χ in terms of the angle θ_1 of scattering in the L-system.

$$p_1 p'_1 \cos \theta_1 = \mathcal{E}'_1(\mathcal{E}_1 + m) - \mathcal{E}_1 m - m^2.$$

Use $c = 1$ units. Now use the energy momentum relation

$$p_i^2 = \mathcal{E}_i^2 - m_i^2$$

$$\Rightarrow (\mathcal{E}_1^2 - m^2)(\mathcal{E}'_1^2 - m^2) \cos^2 \theta_1 = (\mathcal{E}_1 + m)^2(\mathcal{E}'_1 - m)^2$$

$$\Rightarrow \mathcal{E}'_1 \cos^2 \theta_1 (\mathcal{E}_1 - m) + m_1 \cos^2 \theta_1 (\mathcal{E}_1 - m) = \mathcal{E}'_1(\mathcal{E}_1 + m) - m(\mathcal{E}_1 + m)$$

$$\Rightarrow \mathcal{E}'_1 = \frac{m_1 ((\mathcal{E}_1 + m) + (\mathcal{E}_1 - m) \cos^2 \theta)}{(\mathcal{E}_1 + m) - (\mathcal{E}_1 - m) \cos^2 \theta_1}.$$

Using energy conservation

$$\mathcal{E}'_2 = \mathcal{E}_1 + m - \mathcal{E}'_1 = \mathcal{E}_1 + m - m \frac{((\mathcal{E}_1 + m) + (\mathcal{E}_1 - m) \cos^2 \theta_1)}{(\mathcal{E}_1 + m) - (\mathcal{E}_1 - m) \cos^2 \theta_1}$$

$$\begin{aligned}
&= m + \left[\frac{\mathcal{E}_1(2m) + \mathcal{E}_1(\mathcal{E}_1 - m) \sin^2 \theta_1 - 2m\mathcal{E}_1 + m(\mathcal{E}_1 - m) \sin^2 \theta_1}{2m + (\mathcal{E}_1 - m) \sin^2 \theta_1} \right] \\
\mathcal{E}_2 &= m + \left[\frac{(\mathcal{E}_1^2 - m^2) \sin^2 \theta_1}{2m + (\mathcal{E}_1 - m) \sin^2 \theta_1} \right]
\end{aligned}$$

Lastly, for χ

$$\begin{aligned}
\mathcal{E}'_1 &= \mathcal{E}_1 - \frac{(\mathcal{E}_1 - m)}{2}(1 - \cos \chi) \\
\Rightarrow \mathcal{E}'_1 - \mathcal{E}_1 - m &= -m - \frac{(\mathcal{E}_1 - m)}{2}(1 - \cos \chi).
\end{aligned}$$

Now use

$$\mathcal{E}'_1 - \mathcal{E}_1 - m = -\mathcal{E}'_0$$

to arrive at

$$\begin{aligned}
\mathcal{E}'_2 &= m + \frac{(\mathcal{E}_1 - m)}{2}(1 - \cos \chi). \\
\Rightarrow 2(\mathcal{E}'_2 - m) &= \frac{2(\mathcal{E}_1^2 - m^2) \sin^2 \theta_1}{2m + (\mathcal{E}_1 - m) \sin^2 \theta_1} = (\mathcal{E}_1 - m)(1 - \cos \chi)
\end{aligned}$$

Solve for χ

$$\begin{aligned}
\cos \chi &= 1 - \left[\frac{2(\mathcal{E}_1 + m) \sin^2 \theta_1}{2m + (\mathcal{E}_1 - m) \sin^2 \theta_1} \right] \\
\Rightarrow \cos \chi &= \frac{2m - (\mathcal{E}_1 + 3m) \sin^2 \theta_1}{2m + (\mathcal{E}_1 - m) \sin^2 \theta_1}
\end{aligned}$$

3. Express the acceleration of a particle in terms of its velocity and the electric and magnetic field intensities.

Given the Lorentz force

$$\frac{d\mathbf{p}}{dt} = e\mathbf{E} + \frac{e}{c}(\mathbf{v} \times \mathbf{H})$$

substitute relation for \mathbf{p}

$$\mathbf{p} = \frac{\mathbf{v}\mathcal{E}_k}{c^2}, \quad \mathcal{E}_k = \frac{mc^2}{\sqrt{1 - v^2/c^2}}$$

then

$$\begin{aligned}
\Rightarrow \frac{d\mathbf{p}}{dt} &= \dot{v} \frac{\mathcal{E}_k}{c^2} + \frac{v\dot{\mathcal{E}}_k}{c^2} \\
&= \dot{v} \left(\frac{\mathcal{E}_k}{c^2} \right) + \frac{v}{c^2} e\mathbf{E} \cdot \mathbf{v} \\
&= \dot{v} \left(\frac{\mathcal{E}_k}{c^2} \right) + \frac{e}{c^2} \mathbf{v}(\mathbf{v} \cdot \mathbf{E})
\end{aligned}$$

where $\dot{\mathcal{E}}_k = e\mathbf{E} \cdot \mathbf{v}$. Thus

we have

$$\begin{aligned}
\dot{v} &= \frac{e}{m} \sqrt{1 - \left(\frac{v}{c} \right)^2} \left(\mathbf{E} + \frac{1}{c}(\mathbf{v} \times \mathbf{H}) - \frac{1}{c^2} \mathbf{v}(\mathbf{v} \cdot \mathbf{E}) \right). \\
\dot{v} \left(\frac{\mathcal{E}_k}{c^2} \right) &= c\mathbf{E} + \frac{e}{c}(\mathbf{v} \times \mathbf{H}) - \frac{e}{c^2} \mathbf{v}(\mathbf{v} \cdot \mathbf{E}) \\
&= \frac{ec^2}{\mathcal{E}_k} \left(\mathbf{E} + \frac{1}{c}(\mathbf{v} \times \mathbf{H}) - \frac{1}{c^2} \mathbf{v}(\mathbf{v} \cdot \mathbf{E}) \right)
\end{aligned}$$

Now using the energy relation

$$\mathcal{E}_k = \frac{mc^2}{\sqrt{1 - (v/c)^2}}$$

we have

$$\dot{v} = \frac{e}{m} \sqrt{1 - (v/c)^2} \left(\mathbf{E} + \frac{1}{c} (\mathbf{v} \times \mathbf{H}) - \frac{1}{c^2} \mathbf{v} (\mathbf{v} \cdot \mathbf{E}) \right)$$

4. Determine the relativistic motion of a charge in electric and magnetic fields which are mutually perpendicular and equal in magnitude.

Using the lorentz force

$$\frac{dp}{dt} = e\mathbf{E} + \frac{e}{c} (\mathbf{v} \times \mathbf{H})$$

we have

$$\dot{p}_x = \frac{e}{c} E v_y; \quad \dot{p}_y = eE \left(1 - \frac{v_x}{c}\right); \quad \dot{p}_z = 0.$$

Now we use

$$\dot{\mathcal{E}}_k = e\mathbf{E} \cdot \mathbf{v}, \quad \mathcal{E}_k = \frac{mc^2}{\sqrt{1 - (v/c)^2}}$$

so

$$\dot{\mathcal{E}}_k = eE v_y, \quad \dot{p}_x = \frac{\dot{\mathcal{E}}_k}{c} \Rightarrow \mathcal{E}_k - c p_x = \text{const} = \gamma$$

From this

$$\begin{aligned} \mathcal{E}_k^2 &= m^2 c^4 + c^2 (p_x^2 + p_y^2 + p_z^2) \\ &= c^2 p_x^2 + c^2 p_y^2 + \mathcal{E}^2 \end{aligned}$$

where we have used

$$\mathcal{E}^2 = m^2 c^4 + c^2 p_z^2 = \text{const.}$$

We may rewrite this as

$$\begin{aligned} \mathcal{E}_k^2 - c^2 p_x^2 &= c^2 p_y^2 + \mathcal{E}^2 \\ (\mathcal{E}_k - c p_x)(\mathcal{E}_k + c p_x) &= c^2 p_y^2 + \mathcal{E}^2 \\ \gamma(\mathcal{E}_k + c p_x) &= c^2 p_y^2 + \mathcal{E}^2 \end{aligned}$$

thus

$$\mathcal{E}_k + c p_x = \frac{1}{\gamma} (c^2 p_y^2 + \mathcal{E}^2)$$

Forging onward, we use the relations

$$\begin{aligned} \mathcal{E}_k &= \frac{\gamma}{2} + \frac{c^2 p_y^2 + \mathcal{E}^2}{2\gamma} \\ p_x &= -\frac{\gamma}{2c} + \frac{c^2 p_y^2 + \mathcal{E}^2}{2\alpha c} \\ \dot{p}_y &= eE(1 - v_x/c) \Rightarrow eE(\mathcal{E}_k - c p_x) = eE\gamma \end{aligned}$$

and we integrate over time

$$eE\gamma \int dt = \int dp_y \left(\frac{\gamma}{2} + \frac{\mathcal{E}^2}{2\gamma} + \frac{c^2 p_y^2}{2\gamma} \right)$$

$$\begin{aligned}\Rightarrow eE\gamma t &= p_y \left(\frac{\gamma}{2} + \frac{\mathcal{E}^2}{2\gamma} \right) + \frac{c^2 p_y^3}{6\gamma} \\ \Rightarrow 2eE\gamma t &= p_y \left(1 + \frac{\mathcal{E}^2}{\gamma^2} \right) + \frac{c^2 p_y^3}{3\alpha^2}.\end{aligned}$$

Transforming our variables

$$\begin{aligned}\frac{dx}{dt} &= \frac{c^2 p_x}{\mathcal{E}_k} \Rightarrow dt = \frac{\mathcal{E}_k}{eE\gamma} dp_y \\ \Rightarrow dx &= \frac{c^2 p_x}{\mathcal{E}_k} dt \\ &= \frac{c^2}{eE\gamma} \left(-\frac{\gamma}{2c} + \frac{\mathcal{E}^2}{2\gamma c} + \frac{c^2 p_y^2}{2\gamma c} \right) dp_y\end{aligned}$$

Now integrate

$$\begin{aligned}x &= \frac{c^2}{eE\gamma} \left(\frac{E^2}{2\gamma c} - \frac{\gamma}{2c} \right) p_y + \frac{c^2 p_y^3 c^2}{6\gamma c E e \gamma} \\ x &= \frac{c}{2eE} \left(\frac{\mathcal{E}^2}{\gamma^2} - 1 \right) p_y + \frac{c^3 p_y^3}{6\alpha^2 e E}.\end{aligned}$$

Now we repeat the same procedure to find both y and z . Starting with

$$\frac{dy}{dt} = \frac{c^2 p_y}{\mathcal{E}_k} \Rightarrow dy = \frac{c^2}{\mathcal{E}_k} p_y \frac{\mathcal{E}_k}{eE\gamma} dp_y \Rightarrow dt = \frac{\mathcal{E}_k}{eE\gamma} dp_y$$

this leads to

$$y = \frac{c^2 p_y^2}{2\gamma e E}.$$

As for the z component

$$\frac{dz}{dt} = \frac{c^2 p_z}{\mathcal{E}_k} \Rightarrow dz = \frac{p_z c^2}{\mathcal{E}_k} \frac{\mathcal{E}_k}{eE\gamma} dp_y$$

which leads to

$$z = \frac{p_z c^2 p_y}{eE\gamma}.$$

5. Show explicitly that two successive Lorentz transformations in the same direction are equivalent to a single Lorentz transformation with a velocity

$$v = \frac{v_1 + v_2}{1 + (v_1 v_2 / c)}$$

This is an alternative way to derive the parallel-velocity addition law.

First boost to S'

$$x'_0 = \gamma_1(x_0 - \beta_1 x_1), \quad x'_1 = \gamma_1(x_1 - \beta_1 x_0)$$

where we have

$$\gamma_1 = \frac{1}{\sqrt{1 - (v_1/c)^2}}, \quad \beta_1 = v_1/c$$

and likewise for a boost to S''

$$x''_0 = \gamma_2(x'_0 - \beta_2 x'_1), \quad x''_1 = \gamma_2(x'_1 - \beta_2 x'_0)$$

where again

$$\gamma_2 = \frac{1}{\sqrt{1 - (v_2/c)^2}}, \quad \beta_2 = v_2/c.$$

Substitute these into each other

$$x_0'' = \gamma_2 \gamma_1 ((1 + \beta_2 \beta_1) x_0 - (\beta_1 + \beta_2) x_1)$$

$$x_1'' = \gamma_2 \gamma_1 ((1 + \beta_2 \beta_1) x_1 - (\beta_1 + \beta_2) x_0).$$

The transformation to original frame S goes as

$$x_0'' = \gamma(x_0 - \beta x_1)$$

$$x_1'' = \gamma(x_1 - \beta x_0)$$

where

$$\gamma = \frac{1}{\sqrt{1 - (v/c)^2}}, \quad \beta = v/c.$$

Matching the coefficients from the equations above

$$\gamma_2 \gamma_1 (1 + \beta_2 \beta_1) = \gamma$$

$$\Rightarrow \frac{1}{\sqrt{1 - (v/c)^2}} = \frac{1}{\sqrt{1 - (v_2/c)^2}} \frac{1}{\sqrt{1 - (v_1/c)^2}} \left(1 + \frac{v_2 v_1}{c^2}\right).$$

Now solve for v

$$v = \sqrt{c^2 - \frac{(1 - (v_2/c)^2)(1 - (v_1/c)^2)}{(1 + v_2 v_1/c^2)^2}}$$

$$\Rightarrow v = \frac{v_1 + v_2}{1 + (v_1 v_2/c^2)}$$

6. A coordinate system K' moves with a velocity \mathbf{v} relative to another system K . In K' a particle has a velocity \mathbf{u}' and an acceleration \mathbf{a}' . Find the Lorentz transformation law for accelerations, and show that in the system K the components of acceleration parallel and perpendicular to \mathbf{v} are

$$\mathbf{a}_{||} = \frac{\left(1 - \frac{v^2}{c^2}\right)^{3/2}}{\left(1 + \frac{\mathbf{v} \cdot \mathbf{u}'}{c^2}\right)^3} \mathbf{a}'_{||}$$

$$\mathbf{a}_{\perp} = \frac{\left(1 - \frac{v^2}{c^2}\right)}{\left(1 + \frac{\mathbf{v} \cdot \mathbf{u}'}{c^2}\right)^3} \left(\mathbf{a}'_{\perp} + \frac{\mathbf{v}}{c^2} \times (\mathbf{a}' \times \mathbf{u}')\right)$$

Let us perform a boost in the x direction

$$x^0 = \gamma(x'^0 + \beta x'), \quad x = \gamma(x' + \beta x'^0).$$

Then for the K frame

$$\mathbf{u} = c \frac{\partial \mathbf{x}}{\partial x^0}, \quad \mathbf{a} = c \frac{\partial \mathbf{u}}{\partial x^0}$$

and for the K' frame

$$\mathbf{u}' = c \frac{\partial \mathbf{x}'}{\partial x'^0}, \quad \mathbf{a}' = c \frac{\partial \mathbf{u}'}{\partial x'^0}.$$

It follows that

$$\frac{dx^0}{dx'^0} = \gamma(1 + \beta u'_x/c)$$

and

$$\frac{dx'^0}{dx} = \frac{1}{\gamma(1 + \beta u'_x/c)}$$

and so

$$\begin{aligned} u_x &= c \frac{\partial x}{\partial x'^0} \\ &= c \frac{dx'^0}{dx^0} \frac{dx}{dx'^0} \\ &= \frac{c}{\gamma(1 + \beta u'_x/c)} \frac{d}{dx'^0} \gamma(x' + \beta x'^0) \\ &= \frac{u'_x + c\beta}{1 + \beta u'_x/c}. \end{aligned}$$

As for the y component

$$\begin{aligned} u_y &= c \frac{dy}{dx'^0} \\ &= c \frac{dx'^0}{dx^0} \frac{dy}{dx'^0} \\ &= \frac{c}{\gamma(1 + \beta u'_x/c)} \left(\frac{u'_y}{c} \right) \\ &= \frac{u'_y}{\gamma(1 + \beta u'_x/c)}. \end{aligned}$$

Now from using

$$\beta u'_x = \boldsymbol{\beta} \cdot \mathbf{u}'$$

and that \mathbf{x} and \mathbf{v} are parallel (and perpendicular to \mathbf{y}) we have

$$\begin{aligned} \mathbf{u}_{||} &= \frac{\mathbf{u}'_{||} + c\boldsymbol{\beta}}{1 + \boldsymbol{\beta} \cdot \mathbf{u}'/c} \\ \mathbf{u}_{\perp} &= \frac{\mathbf{u}'_{\perp}}{\gamma(1 + \boldsymbol{\beta} \cdot \mathbf{u}'/c)}. \end{aligned}$$

To find the accelerations

$$\begin{aligned} a_x &= c \frac{du_x}{dx'^0} \\ &= \frac{c}{\gamma(1 + \beta u'_x/c)} \frac{d}{dx'^0} \frac{u'_x + c\beta}{1 + \beta u'_x/c} \\ &= \frac{(1 - \beta^2) a'_x}{\gamma(1 + \beta u'_x/c)^3} \\ &= \frac{a'_x}{\gamma^3(1 + \beta u'_x/c)^3} \end{aligned}$$

and similarly for a_y we have

$$a_y = \frac{a'_y + \beta(u'_x a'_y - u'_y a'_x)/c}{\gamma^2(1 + \beta u'_x/c)^3}$$

Finally

$$a_{||} = a_x$$

so

$$a_{\parallel} = \frac{\mathbf{a}'_{\parallel}}{\gamma^3(1 + \beta \cdot \mathbf{u}'/c)^3}.$$

For the perpendicular component

$$a_{\perp} = \frac{a'_{\perp} + \mathbf{a}'(\beta \cdot \mathbf{u}') - \mathbf{u}'(\beta \cdot \mathbf{a}')/c}{\gamma^2(1 + \beta \cdot \mathbf{u}'/c)^3}$$

and so

$$a_{\perp} = \frac{a'_{\perp} + \beta \times (\mathbf{a}' \times \mathbf{u}')/c}{\gamma^2(1 + \beta \cdot \mathbf{u}'/c)^3}$$

7. A particle of mass M and 4-momentum P decays into two particles of masses m_1 and m_2 .

- (a) Use the conservation of energy and momentum in the form, $p_2 = P - p_1$, and the invariance of scalar products of 4-vectors to show that the total energy of the first particle in the rest frame of the decaying particle is

$$E_1 = \frac{M^2 + m_1^2 - m_2^2}{2M}$$

and that E_2 is obtained by interchanging m_1 and m_2 .

- (b) Show that the *kinetic energy* T_i of the i th particle in the same frame is

$$T_i = \Delta M \left(1 - \frac{m_i}{M} - \frac{\Delta M}{2M} \right)$$

where $\Delta M = M - m_1 - m_2$ is the mass excess or Q value of the process.

- (c) The charged pi-meson ($M = 139.6 \text{ MeV}$) decays into a mu-meson ($m_1 = 105.7 \text{ MeV}$) and a neutrino ($m_2 = 0$). Calculate the kinetic energies of the mu-meson and the neutrino in the pi-meson's rest frame. The unique kinetic energy of the muon is the signature of a two-body decay. It entered importantly in the discovery of the pi-meson in photographic emulsions by Powell and his coworkers in 1947.

- (a) In the center of mass frame

$$E = E_1 + E_2 = Mc^2$$

and

$$\mathbf{p} - \mathbf{p}_1 = \mathbf{p}_2 \Rightarrow \mathbf{p}_2 = -\mathbf{p}_1.$$

Thus

$$(E_1, \mathbf{p}_1) + (E_2, \mathbf{p}_2) = (Mc^2, -\mathbf{p}_1)$$

and using the energy momentum relation

$$\begin{aligned} E_1^2 - E_2^2 &= c^4(m_1^2 - m_2^2) \\ (E_1 + E_2)(E_1 - E_2) &= c^4(m_1^2 - m_2^2) \\ \Rightarrow E_1 - E_2 &= c^4 \frac{(m_1^2 - m_2^2)}{Mc^2} \\ \Rightarrow E_1 + E_2 &= Mc^2 \end{aligned}$$

Add these last two equations together

$$E_1 = \frac{c^2(m_1^2 - m_2^2 + M^2)}{2M}$$

or subtract the two

$$E_2 = \frac{c^2(M^2 + m_2^2 - m_1^2)}{2M}.$$

(b) Employing the $c = 1$ convenient units

$$\begin{aligned} T_1 &= E_1 - m_1 \\ &= \frac{m_1^2 - m_2^2 - M^2 - 2m_1M}{2M} \\ &= \frac{M^2 - m_1^2 - m_2^2 - 2m_1M}{2m} \\ &= \frac{\Delta M}{2m}(M - m_1 - m_2) \\ &= \frac{\Delta M}{2M}(M - m_1 + m_2 - \Delta M + M - m_1 - m_2) \\ &= \frac{\Delta M}{2M}(2M - 2m_1 - \Delta M) \\ &= \Delta M \left(1 - \frac{m_1}{M} - \frac{\Delta M}{2M} \right) \end{aligned}$$

thus

$$T_i = \Delta M \left(1 - \frac{m_i}{M} - \frac{\Delta M}{2M} \right).$$

(c) From part (a)

$$E_1 = \frac{M^2 + m_1^2 - m_2^2}{2M} = 109.8 \text{ (MeV)}$$

and so for the mu-meson

$$T_1 = E_1 - m_1 = 4.1 \text{ (MeV)}.$$

As for the pi-meson

$$E_1 = 29.8 \text{ (MeV)}$$

and so again

$$T_2 = E_2 - m_2 = 29.8 \text{ (MeV)}.$$