

3Space Conformal Transformations v1

1 Conformal Transformation of Curvature Tensors

For $D = 3$ with $\mu, \nu = 1, 2, 3$ the Ricci tensor and scalar transform under conformal transformation $g_{\mu\nu} \rightarrow \Omega^2 g_{\mu\nu}$ as

$$R_{\mu\nu} \rightarrow R_{\mu\nu} + g_{\mu\nu} \Omega^{-1} \nabla_\alpha \nabla^\alpha \Omega + \Omega^{-1} \nabla_\mu \nabla_\nu \Omega - 2\Omega^{-2} \nabla_\mu \Omega \nabla_\nu \Omega$$

$$R \rightarrow \Omega^{-2} R + 4\Omega^{-3} \nabla_\alpha \nabla^\alpha \Omega - 2\Omega^{-4} \nabla_\alpha \Omega \nabla^\alpha \Omega \quad (1.1)$$

and thus the Einstein tensor transforms as

$$G_{\mu\nu} \rightarrow G_{\mu\nu} + g_{\mu\nu} (\Omega^{-2} \nabla_\alpha \Omega \nabla^\alpha \Omega - \Omega^{-1} \nabla_\alpha \nabla^\alpha \Omega) + \Omega^{-1} \nabla_\mu \nabla_\nu \Omega - 2\Omega^{-2} \nabla_\mu \Omega \nabla_\nu \Omega \quad (1.2)$$

Perturbing the above we find

$$\delta G_{\mu\nu} \rightarrow \delta G_{\mu\nu} + \delta S_{\mu\nu} \quad (1.3)$$

$$\begin{aligned} \delta S_{\mu\nu} = & -h_{\mu\nu} \Omega^{-1} \nabla_\alpha \nabla^\alpha \Omega + \frac{1}{2} \Omega^{-1} \nabla_\alpha h_{\mu\nu} \nabla^\alpha \Omega - \frac{1}{2} g_{\mu\nu} \Omega^{-1} \nabla_\alpha h \nabla^\alpha \Omega + h_{\mu\nu} \Omega^{-2} \nabla_\alpha \Omega \nabla^\alpha \Omega \\ & + g_{\mu\nu} \Omega^{-1} \nabla^\alpha \Omega \nabla_\beta h_{\alpha}{}^\beta - g_{\mu\nu} h_{\alpha\beta} \Omega^{-2} \nabla^\alpha \Omega \nabla^\beta \Omega + g_{\mu\nu} h_{\alpha\beta} \Omega^{-1} \nabla^\beta \nabla^\alpha \Omega \\ & - \frac{1}{2} \Omega^{-1} \nabla^\alpha \Omega \nabla_\mu h_{\nu\alpha} - \frac{1}{2} \Omega^{-1} \nabla^\alpha \Omega \nabla_\nu h_{\mu\alpha}. \end{aligned} \quad (1.4)$$

2 Background $G_{ij}^{(0)} = -\kappa_3^2 T_{ij}^{(0)}$

Since $G_{\mu\nu}$ vanishes in a flat geometry, the background equation is given as

$$g_{ij} (\Omega^{-2} \nabla_a \Omega \nabla^a \Omega - \Omega^{-1} \nabla_a \nabla^a \Omega) + \Omega^{-1} \nabla_i \nabla_j \Omega - 2\Omega^{-2} \nabla_i \Omega \nabla_j \Omega = -\kappa_3^2 \Lambda \Omega^2 g_{ij}. \quad (2.1)$$

In the covariant formulation, the background equation fixed $k = -\kappa_3^2 \Lambda$. Taking the trace of (2.1)

$$-2\Omega^{-1} \nabla_a \nabla^a \Omega + \Omega^{-2} \nabla_a \Omega \nabla^a \Omega = \frac{12k}{(1+k\rho^2)^2} = 3\Omega^2 k = -3\kappa_3^2 \Lambda \Omega^2. \quad (2.2)$$

We also find

$$\begin{aligned} G_{\rho\rho}^{(0)} &= k\Omega^2 g_{\rho\rho} \\ \frac{4k}{(1+k\rho^2)^2} &= k\Omega^2 g_{\rho\rho} \end{aligned} \quad (2.3)$$

However, for $\theta\theta$ we find

$$\begin{aligned} G_{\theta\theta}^{(0)} &= k\Omega^2 g_{\theta\theta} \\ \frac{2k\rho^2(3+k\rho^2)}{(1+k\rho^2)^2} &\neq \frac{k\rho^2}{(1+k\rho^2)^2} \end{aligned} \quad (2.4)$$

Similarly for $\phi\phi$

$$\begin{aligned} G_{\phi\phi}^{(0)} &= k\Omega^2 g_{\phi\phi} \\ \frac{2k\rho^2(3+k\rho^2)\sin^2\theta}{(1+k\rho^2)^2} &\neq \frac{k\rho^2\sin^2\theta}{(1+k\rho^2)^2} \end{aligned} \quad (2.5)$$

3 Perturbation $\delta G_{ij} = -\kappa_3^2 \delta T_{ij}$

The perturbed energy momentum tensor takes the form

$$\begin{aligned}
-\kappa_3^2 \delta T_{ij} &= -\kappa_3^2 \Lambda \Omega^2 h_{ij} \\
&= k\Omega^2 (-2g_{ij}\psi + 2\nabla_i \nabla_j E + \nabla_i E_j + \nabla_j E_i + 2E_{ij}) \\
-\kappa_3^2 g^{ij} \delta T_{ij} &= k\Omega^2 (-6\psi + 2\nabla_a \nabla^a E)
\end{aligned} \tag{3.1}$$

If we evaluate (1.3) and (1.4) we find

$$\begin{aligned}
\delta G_{ij} &= g_{ij} \nabla_a \nabla^a \psi + g_{ij} \Omega^{-1} \nabla^a \Omega \nabla_b \nabla^b \nabla_a E - 2g_{ij} \Omega^{-2} \nabla^a \Omega \nabla_b \nabla_a E \nabla^b \Omega \\
&\quad + 2g_{ij} \Omega^{-1} \nabla_b \nabla_a \Omega \nabla^b \nabla^a E + \Omega^{-1} \nabla_i \Omega \nabla_j \psi + \Omega^{-1} \nabla_i \psi \nabla_j \Omega - 2\Omega^{-1} \nabla_a \nabla^a \Omega \nabla_j \nabla_i E \\
&\quad + 2\Omega^{-2} \nabla_a \Omega \nabla^a \Omega \nabla_j \nabla_i E - \nabla_j \nabla_i \psi - \Omega^{-1} \nabla^a \Omega \nabla_j \nabla_i \nabla_a E \\
&\quad + g_{ij} \Omega^{-1} \nabla^a \Omega \nabla_b \nabla^b E_a - 2g_{ij} \Omega^{-2} \nabla_a \Omega \nabla_b \Omega \nabla^b E^a + 2g_{ij} \Omega^{-1} \nabla_b \nabla_a \Omega \nabla^b E^a \\
&\quad - \Omega^{-1} \nabla_a \nabla^a \Omega \nabla_i E_j + \Omega^{-2} \nabla_a \Omega \nabla^a \Omega \nabla_i E_j - \Omega^{-1} \nabla_a \nabla^a \Omega \nabla_j E_i + \Omega^{-2} \nabla_a \Omega \nabla^a \Omega \nabla_j E_i \\
&\quad - \Omega^{-1} \nabla^a \Omega \nabla_j \nabla_i E_a \\
&\quad + \nabla_a \nabla^a E_{ij} - 2E_{ij} \Omega^{-1} \nabla_a \nabla^a \Omega + \Omega^{-1} \nabla_a E_{ij} \nabla^a \Omega + 2E_{ij} \Omega^{-2} \nabla_a \Omega \nabla^a \Omega \\
&\quad + 2E^{ab} g_{ij} \Omega^{-1} \nabla_b \nabla_a \Omega - 2E_{ab} g_{ij} \Omega^{-2} \nabla^a \Omega \nabla^b \Omega - \Omega^{-1} \nabla^a \Omega \nabla_i E_{ja} - \Omega^{-1} \nabla^a \Omega \nabla_j E_{ia}.
\end{aligned} \tag{3.2}$$

The only two gauge invariant quantities are

$$\begin{aligned}
\bar{\psi} + \Omega^{-1} (\tilde{\nabla}_k \bar{E} + \bar{E}_k) \tilde{\nabla}^k \Omega &= \psi + \Omega^{-1} (\tilde{\nabla}_k E + E_k) \tilde{\nabla}^k \Omega \\
\bar{E}_{ij} &= E_{ij}.
\end{aligned} \tag{3.3}$$

Hence any E_i or E_j (specifically with index i or j) term must vanish identically in the full $\delta G_{ij} = -\kappa_3^2 \delta T_{ij}$.

Looking only at the relevant vector pieces, we see

$$\begin{aligned}
&= \delta G_{ij}^{(V)} = -\kappa_3^2 \delta T_{ij}^{(V)} \\
-\Omega^{-1} \nabla_a \nabla^a \Omega \nabla_i E_j + \Omega^{-2} \nabla_a \Omega \nabla^a \Omega \nabla_i E_j + (i \leftrightarrow j) &= k\Omega^2 \nabla_i E_j + (i \leftrightarrow j)
\end{aligned} \tag{3.4}$$

which implies

$$(-\Omega^{-1} \nabla_a \nabla^a \Omega + \Omega^{-2} \nabla_a \Omega \nabla^a \Omega) \nabla_i E_j = (-\frac{2}{3} \Omega^{-1} \nabla_a \nabla^a \Omega + \frac{1}{3} \Omega^{-2} \nabla_a \Omega \nabla^a \Omega) \nabla_i E_j. \tag{3.5}$$

The above equation along with (2.4) and (2.5) hints at the necessary form for δS_{ij} and thus $G_{ij}^{(0)}$ and serves as another check upon the conformal flat form of G_{ij} .

$$\delta T_{\mu\nu} = \delta T_{\mu\nu}^{T\theta} + \frac{g_{\mu\nu}}{D-1}\delta T - \frac{1}{D-1}\nabla_\mu\nabla_\nu\int D\delta T. \quad (3.6)$$

$$\begin{aligned} \delta G_{\mu\nu}^{T\theta} &= \delta T_{\mu\nu}^{T\theta} \\ \delta G &= \delta T \end{aligned} \quad (3.7)$$

$$\begin{aligned} \delta G_{\mu\nu} &= \nabla^2 E_{\mu\nu} + (D-2)(g_{\mu\nu}\nabla^2\psi - \nabla_\mu\nabla_\nu\psi) \\ \delta G &= (D-2)(D-1)\nabla^2\psi \end{aligned} \quad (3.8)$$

$$\delta G_{\mu\nu}^{T\theta} = \nabla^2 E_{\mu\nu} - (D-2)\nabla_\mu\nabla_\nu\left[\psi - \int D\nabla^2\psi\right] \quad (3.9)$$

$$\begin{aligned} h_{\mu\nu}^{T\theta} &= h_{\mu\nu} - \nabla_\mu W_\nu - \nabla_\nu W_\mu + \frac{g_{\mu\nu}}{D-1}(\nabla^\alpha W_\alpha - h) \\ &\quad + \frac{D-2}{D-1}\nabla_\mu\nabla_\nu\int D\nabla^\alpha W_\alpha + \frac{\nabla_\mu\nabla_\nu}{D-1}\int Dh \end{aligned} \quad (3.10)$$

$$\mathcal{N}(h_{\mu\nu}^{T\theta}) = \int D\nabla^2 h_{\mu\nu}^{T\theta} \quad (3.11)$$

$$(\mathcal{N}h_{\mu\nu})^{T\theta} = \int D\nabla^2 h_{\mu\nu}^{T\theta} \quad (3.12)$$