TT Projection Curved Space v2

1 Maximally Symmetric Space TT

$$h_{\mu\nu} = h_{\mu\nu}^{T\theta} + \nabla_{\mu}W_{\nu} + \nabla_{\nu}W_{\mu} - \frac{g_{\mu\nu}}{D-1}(\nabla^{\sigma}W_{\sigma} - h) + \frac{2-D}{D-1}\left(\nabla_{\mu}\nabla_{\nu} - \frac{g_{\mu\nu}R}{D(D-1)}\right) \int D(x,x')\nabla^{\sigma}W_{\sigma} - \frac{1}{D-1}\left(\nabla_{\mu}\nabla_{\nu} - \frac{g_{\mu\nu}R}{D(D-1)}\right) \int D(x,x')h \quad (1.1)$$

$$\left(\nabla_{\alpha}\nabla^{\alpha} - \frac{R}{D-1}\right)D(x,x') = g^{-1/2}\delta^{(D)}(x-x') \tag{1.2}$$

$$\nabla^{\mu} h_{\mu\nu} = \left(\nabla_{\alpha} \nabla^{\alpha} - \frac{R}{D}\right) W_{\nu} \tag{1.3}$$

With the covariant operator $(\nabla^2 - R/D)$ mixing indices of W_{ν} , the particular integral solution for W_{ν} involves a bi-tensor Green's function $D_{\sigma\rho'}$ which obeys

$$\left(\nabla^{\alpha}\nabla_{\alpha} - \frac{R}{D}\right)D_{\sigma\rho'}(x, x') = g_{\sigma\rho'}g^{-1/2}\delta^{4}(x - x'). \tag{1.4}$$

Here $g_{\sigma\rho'}$ represents a parallel propagator, defined in terms of Vierbeins e^a_{μ} :

$$g^{\alpha'}{}_{\beta}(x,x') = e^{\alpha'}{}_{a}(x')e^{a}{}_{\beta}(x), \qquad g_{\mu\nu} = \eta_{ab}e^{a}_{\mu}e^{b}_{\nu}.$$
 (1.5)

In terms of (1.4), W_{ν} has particular solution

$$W_{\nu} = g^{1/2} \int D_{\nu}^{\rho'}(x, x') \nabla^{\sigma'} h_{\rho'\sigma'}. \tag{1.6}$$

2 Curved Space TT

To generalize to curved space, we assume $h_{\mu\nu}$ to be of the form:

$$h_{\mu\nu} = h_{\mu\nu}^{T\theta} + \underbrace{\left(\nabla_{\mu}W_{\nu} + \nabla_{\nu}W_{\mu} - \frac{2}{D}g_{\mu\nu}\nabla^{\alpha}W_{\alpha}\right)}_{W_{\mu\nu}} + \underbrace{\frac{1}{D-1}\left(g_{\mu\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\mu}\nabla_{\nu}\right)\Psi}_{S_{\mu\nu}}$$
(2.1)

Taking the trace of (2.1), we find the vector sector $W_{\mu\nu}$ is decoupled from the trace and Ψ can easily be inverted,

$$g^{\mu\nu}W_{\mu\nu} = 0 (2.2)$$

$$g^{\mu\nu}S_{\mu\nu} = \nabla_{\alpha}\nabla^{\alpha}\Psi = h \qquad \rightarrow \Psi = \int g^{1/2}D(x,x')h$$
 (2.3)

Taking the divergence of (2.1), we have

$$\nabla^{\mu}h_{\mu\nu} = \nabla^{\mu}W_{\mu\nu} + \nabla^{\mu}S_{\mu\nu}(h) \tag{2.4}$$

By substituting (2.3), the above serves to define an equation for W_{μ} in terms of h and $h_{\mu\nu}$, namely

$$\nabla_{\alpha}\nabla^{\alpha}W_{\nu} + \nabla^{\alpha}\nabla_{\nu}W_{\alpha} - \frac{2}{D}\nabla_{\nu}\nabla^{\alpha}W_{\alpha} = \nabla^{\alpha}h_{\alpha\nu} - \frac{1}{D-1}\left(\nabla_{\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\alpha}\nabla^{\alpha}\nabla_{\nu}\right)\int g^{1/2}D(x,x')h \quad (2.5)$$

Commuting derivatives, (2.5) can be expressed in the equivalent forms,

$$\left[g_{\nu\alpha}\nabla_{\beta}\nabla^{\beta} + \nabla_{\alpha}\nabla_{\nu} - \frac{2}{D}\nabla_{\nu}\nabla_{\alpha}\right]W^{\alpha} = \nabla^{\alpha}h_{\alpha\nu} - \frac{1}{D-1}\left(\nabla_{\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\alpha}\nabla^{\alpha}\nabla_{\nu}\right)\int g^{1/2}D(x,x')h, (2.6)$$

$$\left[g_{\nu\alpha}\nabla_{\beta}\nabla^{\beta} + \left(\frac{D-2}{D}\right)\nabla_{\nu}\nabla_{\alpha} - R_{\nu\alpha}\right]W^{\alpha} = \nabla^{\alpha}h_{\alpha\nu} - \frac{1}{D-1}R_{\nu\alpha}\nabla^{\alpha}\int g^{1/2}D(x,x')h. \tag{2.7}$$

Similar to (1.4), the requisite Green's function that solves W_{α} is a bi-tensor defined as

$$\left[g_{\nu\alpha}\nabla_{\beta}\nabla^{\beta} + \left(\frac{D-2}{D}\right)\nabla_{\nu}\nabla_{\alpha} - R_{\nu\alpha}\right]D^{\alpha\gamma'} = g^{\alpha\gamma'}g^{-1/2}\delta^{(D)}(x,x'). \tag{2.8}$$

Hence, W_{μ} takes the form

$$W_{\mu} = \int g^{1/2} D_{\mu}^{\sigma'} \left[\nabla^{\rho'} h_{\sigma'\rho'} - \frac{1}{D-1} R_{\sigma'\rho'} \nabla^{\rho'} \int g^{1/2} D(x', x'') h \right]. \tag{2.9}$$

To summarize, in curved space $h_{\mu\nu}$ may be decomposed according to

$$h_{\mu\nu} = h_{\mu\nu}^{T\theta} + \left(\nabla_{\mu}W_{\nu} + \nabla_{\nu}W_{\mu} - \frac{2}{D}g_{\mu\nu}\nabla^{\alpha}W_{\alpha}\right) + \frac{1}{D-1}\left(g_{\mu\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\mu}\nabla_{\nu}\right)\Psi$$
 (2.10)

with Ψ and W_{μ} obeying

$$\Psi = \int g^{1/2} D(x, x') h \tag{2.11}$$

$$W_{\mu} = \int g^{1/2} D_{\mu}^{\sigma'} \left[\nabla^{\rho'} h_{\sigma'\rho'} - \frac{1}{D-1} R_{\sigma'\rho'} \nabla^{\rho'} \int g^{1/2} D(x', x'') h \right]$$
(2.12)

$$\nabla_{\alpha} \nabla^{\alpha} D(x, x') = g^{-1/2} \delta^{(D)}(x - x')$$
(2.13)

$$\left[g_{\nu\alpha}\nabla_{\beta}\nabla^{\beta} + \left(\frac{D-2}{D}\right)\nabla_{\nu}\nabla_{\alpha} - R_{\nu\alpha}\right]D^{\alpha\gamma'} = g^{\alpha\gamma'}g^{-1/2}\delta^{(D)}(x,x'). \tag{2.14}$$

2.1 SVTD Decomposition

Starting with

$$h_{\mu\nu} = h_{\mu\nu}^{T\theta} + \left(\nabla_{\mu}W_{\nu} + \nabla_{\nu}W_{\mu} - \frac{2}{D}g_{\mu\nu}\nabla^{\alpha}W_{\alpha}\right) + \frac{1}{D-1}\left(g_{\mu\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\mu}\nabla_{\nu}\right)\Psi, \tag{2.15}$$

we decompose W_{μ} into transverse and longitudinal components viz.

$$W_{\mu} = \underbrace{W_{\mu} - \nabla_{\mu} \int g^{1/2} D(x, x') \nabla^{\sigma} W_{\sigma}}_{F_{\mu}} + \nabla_{\mu} \underbrace{\int g^{1/2} D(x, x') \nabla^{\sigma} W_{\sigma}}_{H}. \tag{2.16}$$

Setting $h_{\mu\nu}^{T\theta} = 2F_{\mu\nu}$, (2.28) becomes

$$h_{\mu\nu} = 2F_{\mu\nu} + \nabla_{\mu}F_{\nu} + \nabla_{\nu}F_{\mu} + 2\nabla_{\mu}\nabla_{\nu}H - \frac{2}{D}g_{\mu\nu}\nabla_{\alpha}\nabla^{\alpha}H + \frac{1}{D-1}\left(g_{\mu\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\mu}\nabla_{\nu}\right)\Psi. \tag{2.17}$$

Upon further defining

$$F = H - \frac{1}{2(D-1)}\Psi \tag{2.18}$$

$$\chi = \frac{1}{D} \nabla_{\alpha} \nabla^{\alpha} H - \frac{1}{2(D-1)} \nabla_{\alpha} \nabla^{\alpha} \Psi, \qquad (2.19)$$

we may express (2.28) as the desired SVTD form:

$$h_{\mu\nu} = -2g_{\mu\nu}\chi + 2\nabla_{\mu}\nabla_{\nu}F + \nabla_{\mu}F_{\nu} + \nabla_{\nu}F_{\mu} + 2F_{\mu\nu}. \tag{2.20}$$

$$\chi = \frac{1}{D} \nabla^{\sigma} W_{\sigma} - \frac{1}{2(D-1)} h \tag{2.21}$$

$$F = \int g^{1/2} D(x, x') \nabla^{\sigma} W_{\sigma} - \frac{1}{2(D-1)} \int g^{1/2} D(x, x') h$$
 (2.22)

$$F_{\mu} = W_{\mu} - \nabla_{\mu} \int g^{1/2} D(x, x') \nabla^{\sigma} W_{\sigma}$$

$$(2.23)$$

$$2F_{\mu\nu} = 2g_{\mu\nu}\chi - 2\nabla_{\mu}\nabla_{\nu}F - \nabla_{\mu}F_{\nu} - \nabla_{\nu}F_{\mu} - h_{\mu\nu} \tag{2.24}$$

$$\left[g_{\nu\alpha}\nabla_{\beta}\nabla^{\beta} + \left(\frac{D-2}{D}\right)\nabla_{\nu}\nabla_{\alpha} - R_{\nu\alpha}\right]W^{\alpha} = \nabla^{\alpha}h_{\alpha\nu} - \frac{1}{D-1}\left(\nabla_{\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\alpha}\nabla^{\alpha}\nabla_{\nu}\right)\Psi$$
 (2.25)

$$\frac{2(D-1)}{D}\nabla_{\alpha}\nabla^{\alpha}\nabla^{\sigma}W_{\sigma} - \nabla^{\alpha}RW_{\alpha} - 2R^{\alpha\beta}\nabla_{\alpha}W_{\beta} = \nabla^{\alpha}\nabla^{\beta}h_{\alpha\beta} - \frac{1}{(D-1)}\left[\frac{1}{2}\nabla^{\alpha}R\nabla_{\alpha} + R^{\alpha\beta}\nabla_{\alpha}\nabla_{\beta}\right]\Psi(2.26)$$

2.2 Curved TT in Max. Symmetric Space (Incomplete)

In a space of maximal symmetry defined by

$$R_{\lambda\mu\nu\kappa} = k(g_{\mu\nu}g_{\lambda\kappa} - g_{\lambda\nu}g_{\mu\kappa})$$

$$R_{\mu\kappa} = k(1-D)g_{\mu\kappa} = \frac{R}{D}g_{\mu\kappa}$$

$$R = kD(1-D), \qquad (2.27)$$

the defining equation for W_{μ} reduces to

$$\left[g_{\nu\alpha}\left(\nabla_{\beta}\nabla^{\beta} - \frac{R}{D}\right) + \left(\frac{D-2}{D}\right)\nabla_{\nu}\nabla_{\alpha}\right]W^{\alpha} = \nabla^{\alpha}h_{\alpha\nu} - \frac{R}{D-1}\nabla_{\nu}\int g^{1/2}D(x,x')h \tag{2.28}$$

2.3 Curved TT in Minkowski

$$h_{\mu\nu} = h_{\mu\nu}^{T\theta} + \left(\nabla_{\mu}W_{\nu} + \nabla_{\nu}W_{\mu} - \frac{2}{D}g_{\mu\nu}\nabla^{\alpha}W_{\alpha}\right) + \frac{1}{D-1}\left(g_{\mu\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\mu}\nabla_{\nu}\right)\Psi, \tag{2.29}$$

In a Minkowski geometry, the defining equation for W_{μ} reduces to

$$\left[g_{\nu\alpha}\nabla_{\beta}\nabla^{\beta} + \left(\frac{D-2}{D}\right)\nabla_{\nu}\nabla_{\alpha}\right]W^{\alpha} = \nabla^{\alpha}h_{\alpha\nu}$$
(2.30)

$$W_{\mu} = \int g^{1/2} D_{\mu}^{\sigma'} \nabla^{\rho'} h_{\sigma'\rho'} \tag{2.31}$$

$$\left[g_{\nu\alpha}\nabla_{\beta}\nabla^{\beta} + \left(\frac{D-2}{D}\right)\nabla_{\nu}\nabla_{\alpha}\right]D^{\alpha\gamma'} = g^{\alpha\gamma'}\delta^{(D)}(x,x')$$
(2.32)

$$\nabla_{\alpha}\nabla^{\alpha}\Psi = h \tag{2.33}$$

$$\Psi = \int g^{1/2} D(x, x') h \tag{2.34}$$

$$\nabla_{\alpha} \nabla^{\alpha} D(x, x') = g^{-1/2} \delta^{(D)}(x - x')$$
(2.35)

Decompose W_{μ} into transverse and longitudinal components viz.

$$W_{\mu} = \underbrace{W_{\mu} - \nabla_{\mu} \int g^{1/2} D(x, x') \nabla^{\sigma} W_{\sigma}}_{F_{\mu}} + \nabla_{\mu} \underbrace{\int g^{1/2} D(x, x') \nabla^{\sigma} W_{\sigma}}_{H}. \tag{2.36}$$

$$h_{\mu\nu} = h_{\mu\nu}^{T\theta} + \left(\nabla_{\mu}W_{\nu} + \nabla_{\nu}W_{\mu} - \frac{2}{D}g_{\mu\nu}\nabla^{\alpha}W_{\alpha}\right) + \frac{1}{D-1}\left(g_{\mu\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\mu}\nabla_{\nu}\right)\Psi$$
 (2.37)

$$h_{\mu\nu} = -2g_{\mu\nu}\chi + 2\nabla_{\mu}\nabla_{\nu}F + \nabla_{\mu}F_{\nu} + \nabla_{\nu}F_{\mu} + 2F_{\mu\nu}. \tag{2.38}$$

$$\chi = \frac{1}{D} \nabla^{\sigma} W_{\sigma} - \frac{1}{2(D-1)} h \tag{2.39}$$

$$F = \int g^{1/2} D(x, x') \nabla^{\sigma} W_{\sigma} - \frac{1}{2(D-1)} \int g^{1/2} D(x, x') h$$
 (2.40)

$$F_{\mu} = W_{\mu} - \nabla_{\mu} \int g^{1/2} D(x, x') \nabla^{\sigma} W_{\sigma}$$

$$(2.41)$$

$$\left[g_{\nu\alpha}\nabla_{\beta}\nabla^{\beta} + \left(\frac{D-2}{D}\right)\nabla_{\nu}\nabla_{\alpha} - R_{\nu\alpha}\right]W^{\alpha} = \nabla^{\alpha}h_{\alpha\nu} - \frac{1}{D-1}\left(\nabla_{\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\alpha}\nabla^{\alpha}\nabla_{\nu}\right)\Psi$$
 (2.42)

$$\nabla_{\alpha}\nabla^{\alpha}W_{\nu} + \left(\frac{D-2}{D}\right)\nabla_{\nu}\nabla^{\alpha}W_{\alpha} - R_{\nu\alpha}W^{\alpha} = \nabla^{\alpha}h_{\alpha\nu} - \frac{1}{D-1}\left(\nabla_{\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\alpha}\nabla^{\alpha}\nabla_{\nu}\right)\Psi$$
 (2.43)

$$\nabla_{\alpha}\nabla^{\alpha}W_{\nu} + \nabla^{\alpha}\nabla_{\nu}W_{\alpha} - \frac{2}{D}\nabla_{\nu}\nabla^{\alpha}W_{\alpha} = \nabla^{\alpha}h_{\alpha\nu} - \frac{1}{D-1}\left(\nabla_{\nu}\nabla_{\alpha}\nabla^{\alpha} - \nabla_{\alpha}\nabla^{\alpha}\nabla_{\nu}\right)\Psi$$
 (2.44)

$$\frac{2(D-1)}{D}\nabla_{\alpha}\nabla^{\alpha}\nabla^{\sigma}W_{\sigma} - \nabla^{\alpha}RW_{\alpha} - 2R^{\alpha\beta}\nabla_{\alpha}W_{\beta} \quad = \quad \nabla^{\alpha}\nabla^{\beta}h_{\alpha\beta} - \frac{1}{(D-1)}\left[\frac{1}{2}\nabla^{\alpha}R\nabla_{\alpha} + R^{\alpha\beta}\nabla_{\alpha}\nabla_{\beta}\right]\Psi\left(2.45\right)$$