Lecture 5

02/03/2016

eap:

harmonic solutions
$$A(\vec{r},t) = e^{-i\omega t} A(\vec{r})$$
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 $A(\vec{r},t) = e^{-i\omega t}$

 $r << \lambda$ or kr << 1 $\begin{array}{c|c}
\hline
K | \vec{r} - \vec{r}' | \sim \kappa r << 1 \\
\hline
A & (\vec{r}) = \frac{\kappa}{4\pi} \int \frac{f'(\vec{r}') d^3r'}{|\vec{r} - \vec{r}'|} & \text{Static field} \\
\hline
A & (\vec{r}) = \frac{\kappa}{4\pi} \int \frac{f'(\vec{r}') d^3r'}{|\vec{r} - \vec{r}'|} & \text{(Biot-Sowart)}
\end{array}$

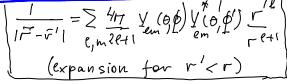


multipole expansion)

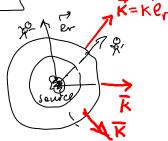
$$\overline{\mathcal{A}}_{\omega}(\vec{r}) = \frac{\mu_{o}}{4\pi} \sum_{\ell,m} \frac{4\pi}{2\ell+1} \frac{Y_{em}(\theta_{l}\phi)}{r^{e+1}} \int_{0}^{3} (\vec{r}' \cdot \vec{r}') \underbrace{Y_{em}(\theta'\phi')}_{\ell,m} d^{3}r'$$

The radiation 20 Me.

r>>> => K/F-F'/>>1 (d<<s)



 $r > d \sim r' = \sum_{r' \in r'} |\vec{r} - \vec{r}'| = r - \vec{r} \cdot \vec{\nabla}(r) + \dots = r - \vec{r}' \cdot \vec{r}' = r - \vec{r} \cdot \vec{e}'$ $f(\vec{r} - \vec{r}') = e^{-\vec{r}'} \vec{\nabla}(\vec{r}') \qquad Unit vector \vec{e} = \vec{r}'$ $r' \ll r$ $|\vec{A}(\vec{r}') = \frac{\mu_0}{4\pi} \underbrace{e^{i\kappa r}}_{r'} \int \vec{J}(\vec{r}') e^{-i\kappa \vec{e}_r \cdot \vec{r}'}_{d^3r'}$ source



The function F(R) depends on the K-vector

Spherical wowe i (kr-wt)

The expansion with the spherical Bessel function Math. -> | eikr = 4n \ il Y (k) Y (r) Je(kr) The spherical Bessel Function fector):

equation for
the fe(kr) $\left\{\frac{1}{r^2}\frac{d}{dr}\left(r^2\frac{d}{dr}\right) + k^2 - \frac{\ell(\ell+1)}{r^2}\right\} de(kr) = 0$ $fe(x) = (-1)^{\ell} \left(\frac{1}{x} \frac{d}{dx}\right)^{\ell} \frac{mx}{x}$; $f_0(x) = \sin x/x$ f1(x)= \frac{\sin X}{\sqrt{2}} - \frac{\cos X}{\times} General solution of the wave equation in the radiation zone: $\vec{A}(\vec{r},t) = e^{-i\omega t} \vec{A}_{\omega}(\vec{r}) = \frac{e^{i(\kappa r - \omega t)}}{r} \int d^{3}r' \vec{J}(\vec{r}') e^{-i\kappa r'}$ = the ei(kr-wt) =4π il Yem') (d3' yr') Yem') Je(kr') $Y_{em}^{*}(-\hat{k}) = (-1)^{\ell} Y_{em}^{*}(\hat{k})$ $\vec{A}_{\omega}(\vec{r}) = \frac{\mu_{0}}{4\pi} \frac{e^{i\kappa r}}{r} \sum_{\ell,m}^{\ell} (-i)^{\ell} Y_{em}^{*}(\hat{k}) \int Y_{em}^{(A)} (\kappa r) \vec{y} \vec{r} dr dr$ $(\vec{k} \rightarrow -\vec{k}; \theta \rightarrow \pi - A)$ $K r' = 2\pi \frac{r'}{\lambda} \ll \Delta$ radiation zone: $f_{\ell}(\kappa r') \approx \frac{(\kappa r')^{\ell}}{2 \cdot 3 \cdot 5 \cdot 120 + \Lambda} = \frac{(\kappa r')^{\ell}}{(2\ell + \Lambda)!} \quad (\kappa r' < 1)$ For small values of Kr': $\overrightarrow{A}(\overrightarrow{r},t) \simeq \underbrace{\frac{e^{i\overrightarrow{kr}-i\omega t}}{4\pi}}_{e,m} \underbrace{\sum_{e,m} (-i)^{\ell} \underbrace{Y^{*}_{em}(\overrightarrow{k})}_{(2\ell+4)!!}}_{\underbrace{Y^{*}_{em}(\overrightarrow{r}')}} \underbrace{Y^{*}_{em}(\overrightarrow{r}')}_{(kr')} \underbrace{\overrightarrow{J}_{\omega}(\overrightarrow{r}')}_{d^{2}r'}$ The leading term $\ell=0$: $\overrightarrow{A}(\overrightarrow{r},t) = \frac{e^{i(kr-\omega t)}}{4\pi} \int \overrightarrow{J}_{\omega}(\overrightarrow{r}') d^{3}r'$ same result can be obtained, if we neglect an exponential factor: $\vec{A}(\vec{r},t) = \frac{\mu_0}{\mu_{\Pi}} \underbrace{e^{i(\mu r - \omega t)}}_{r} \underbrace{\vec{J}_{\omega}(\vec{r}')e^{-i\vec{k}\vec{r}'}_{s,r'}}_{r} \simeq \underbrace{\frac{\mu_0}{4\pi}}_{r} \underbrace{e^{i(kr - \omega t)}}_{r} \underbrace{\vec{J}_{\omega}(\vec{r}')d\vec{r}}_{r}^{1}$