

$$\delta W_{\mu\nu}(h_{\mu\nu}^{T\theta}) \text{ in } g_{\mu\nu}^{(0)} = \Omega^2(x)\eta_{\mu\nu} \text{ v1}$$

## 1 Transverse Tracless Projection $h_{\mu\nu}^{T\theta}$

A general rank 2 scalar may be decomposed into transverse and longitudinal pieces, and further into transverse traceless, longitudinal traceless, and trace terms. The decomposition takes the form

$$\begin{aligned} h_{\mu\nu} &= h_{\mu\nu}^T + h_{\mu\nu}^L \\ &= h_{\mu\nu}^{T\theta} + h_{\mu\nu}^{L\theta} + \frac{1}{D-1} g_{\mu\nu} g^{\sigma\tau} h_{\sigma\tau} - \frac{1}{D-1} \nabla_\mu \nabla_\nu \int d^D y D^{(D)}(x-y) g^{\sigma\tau} h_{\sigma\tau} \\ &\equiv h_{\mu\nu}^{T\theta} + h_{\mu\nu}^{L\theta} + h_{\mu\nu}^{tr}, \end{aligned} \quad (1)$$

with the following properties:

$$\begin{aligned} \nabla^\mu h_{\mu\nu}^T &= 0, \quad \nabla^\mu h_{\mu\nu}^L = \nabla^\mu h_{\mu\nu}, \quad \nabla^\mu \nabla^\nu h_{\mu\nu}^L = \nabla^\mu \nabla^\nu h_{\mu\nu} \\ g^{\mu\nu} h_{\mu\nu}^T &\equiv h^T, \quad g^{\mu\nu} h_{\mu\nu}^L \equiv h^L, \quad h = h^T + h^L \\ \nabla^\mu h_{\mu\nu}^{T\theta} &= 0, \quad \nabla^\mu h_{\mu\nu}^{L\theta} = \nabla^\mu h_{\mu\nu}, \quad \nabla^\mu \nabla^\nu h_{\mu\nu}^{L\theta} = \nabla^\mu \nabla^\nu h_{\mu\nu} \\ g^{\mu\nu} h_{\mu\nu}^{T\theta} &= 0, \quad g^{\mu\nu} h_{\mu\nu}^{L\theta} = 0 \\ \nabla^\mu h_{\mu\nu}^{tr} &= 0, \quad g^{\mu\nu} h_{\mu\nu}^{tr} = g^{\mu\nu} h_{\mu\nu} = h. \end{aligned} \quad (2)$$

The decomposition can be carried about through projectors. The transverse projector (E24) is

$$\begin{aligned} T_{\mu\nu\sigma\tau} &= \eta_{\mu\sigma} \eta_{\nu\tau} - \nabla_\mu \int d^D y D^{(D)}(x-y) \eta_{\nu\tau} - \nabla_\nu \int d^D y D^{(D)}(x-y) \eta_{\mu\sigma} \nabla_\tau \\ &\quad + \nabla_\mu \nabla_\nu \int d^D y D^{(D)}(x-y) \nabla_\sigma \int d^D z D^{(D)}(y-z) \nabla_\tau, \end{aligned} \quad (3)$$

where  $\nabla_i$  denotes the flat space derivative throughout this section. Together with a complementary trace projector (E29)

$$Q_{\mu\nu\sigma\tau} = \frac{1}{D-1} \left[ \eta_{\mu\nu} - \nabla_\mu \nabla_\nu \int d^D y D^{(D)}(x-y) \right] \left[ \eta_{\sigma\tau} - \nabla_\sigma \int d^D z D^{(D)}(y-z) \nabla_\tau \right], \quad (4)$$

a transverse traceless projector may be composed as

$$P_{\mu\nu\sigma\tau} = T_{\mu\nu\sigma\tau} - Q_{\mu\nu\sigma\tau}, \quad (5)$$

which may be used as

$$P_{\mu\nu\sigma\tau} h^{\sigma\tau} = h_{\sigma\tau}^{T\theta}. \quad (6)$$

As given in (C.72 *Brane Gravity*), the fluctuation of the flat  $\delta W_{\mu\nu}$  may be expressed entirely in terms of the gauge invariant  $h_{\mu\nu}^{T\theta}$ .

To see this, apply  $\square^2$  to both  $T_{\mu\nu\sigma\tau}$  and  $Q_{\mu\nu\sigma\tau}$ :

$$\nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta T_{\mu\nu\sigma\tau} h^{\sigma\tau} = \nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta h_{\mu\nu} - \nabla_\alpha \nabla^\alpha \nabla_\mu \nabla_\sigma h^\sigma{}_\nu - \nabla_\alpha \nabla^\alpha \nabla_\nu \nabla_\sigma h^\sigma{}_\mu + \nabla_\mu \nabla_\nu \nabla_\sigma \nabla_\tau h^{\sigma\tau} \quad (7)$$

$$\nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta Q_{\mu\nu\sigma\tau} h^{\sigma\tau} = \frac{1}{3} (\eta_{\mu\nu} \nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta h - \eta_{\mu\nu} \nabla_\alpha \nabla^\alpha \nabla_\sigma \nabla_\tau h^{\sigma\tau} + \nabla_\alpha \nabla^\alpha \nabla_\mu \nabla_\nu h + \nabla_\mu \nabla_\nu \nabla_\sigma \nabla_\tau h^{\sigma\tau}) \quad (8)$$

Summing the result gives

$$\begin{aligned} \nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta (T_{\mu\nu\sigma\tau} - Q_{\mu\nu\sigma\tau}) h^{\sigma\tau} &= \nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta P_{\mu\nu\sigma\tau} h^{\sigma\tau} \\ &= \nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta h_{\mu\nu} - \nabla_\alpha \nabla^\alpha \nabla_\mu \nabla_\sigma h^\sigma{}_\nu - \nabla_\alpha \nabla^\alpha \nabla_\nu \nabla_\sigma h^\sigma{}_\mu \\ &\quad + \frac{2}{3} \nabla_\mu \nabla_\nu \nabla_\sigma \nabla_\tau h^{\sigma\tau} - \frac{1}{3} \eta_{\mu\nu} \nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta h - \frac{1}{3} \nabla_\alpha \nabla^\alpha \nabla_\mu \nabla_\nu h \\ &\quad + \frac{1}{3} \eta_{\mu\nu} \nabla_\alpha \nabla^\alpha \nabla_\sigma \nabla_\tau h^{\sigma\tau}. \end{aligned} \quad (9)$$

Comparison to  $\delta W_{\mu\nu}$  evaluated in a flat background

$$\begin{aligned} \delta W_{\mu\nu} &= \frac{1}{2} \left( \nabla_\beta \nabla^\beta \nabla_\alpha \nabla^\alpha h_{\mu\nu} - \nabla_\beta \nabla^\beta \nabla_\mu \nabla_\alpha h^\alpha{}_\nu - \nabla_\beta \nabla^\beta \nabla_\nu \nabla_\alpha h^\alpha{}_\mu + \frac{2}{3} \nabla_\nu \nabla_\mu \nabla_\beta \nabla_\alpha h^{\alpha\beta} \right. \\ &\quad \left. - \frac{1}{3} \eta_{\mu\nu} \nabla_\beta \nabla^\beta \nabla_\alpha \nabla^\alpha h + \frac{1}{3} \nabla_\nu \nabla_\mu \nabla_\alpha \nabla^\alpha h + \frac{1}{3} \eta_{\mu\nu} \nabla_\gamma \nabla^\gamma \nabla_\beta \nabla_\alpha h^{\alpha\beta} \right), \end{aligned} \quad (10)$$

shows that  $\delta W_{\mu\nu}$  may be expressed in terms of  $\square^2$  onto the transverse traceless projector,

$$\begin{aligned} \delta W_{\mu\nu} &= \frac{1}{2} \nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta (T_{\mu\nu\sigma\tau} - Q_{\mu\nu\sigma\tau}) h^{\sigma\tau} \\ &= \frac{1}{2} \nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta P_{\mu\nu\sigma\tau} h^{\sigma\tau} \\ &= \frac{1}{2} \nabla_\alpha \nabla^\alpha \nabla_\beta \nabla^\beta h_{\mu\nu}^{T\theta}. \end{aligned} \quad (11)$$

Based on the properties of  $\delta W_{\mu\nu}$  under conformal transformation, namely  $\delta \bar{W}_{\mu\nu} = \Omega^{-2} \delta W_{\mu\nu}$ , it follows that within a conformal to flat background the fluctuation equations can still be expressed in terms of  $\square^2$  onto a transverse traceless projection, without an imposition of gauge.

## 2 $\delta W_{\mu\nu}(h_{\mu\nu})$

Starting with (APM 43), the direct perturbation of  $W_{\mu\nu}$ , we evaluate in a background  $g_{\mu\nu}^{(0)} = \Omega^2(x) \eta_{\mu\nu}$  which yields the 153 term

$$\begin{aligned} \delta W_{\mu\nu}(h_{\mu\nu}) &= -6\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-6} \partial_\alpha \Omega \partial_\beta h_{\nu\eta} \partial_\gamma \partial_\mu \Omega - 6\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-6} \partial_\alpha \Omega \partial_\beta h_{\mu\eta} \partial_\gamma \partial_\nu \Omega \\ &\quad - 6\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-6} \partial_\alpha \Omega \partial_\gamma \partial_\nu \Omega \partial_\eta h_{\mu\beta} - 6\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-6} \partial_\alpha \Omega \partial_\beta \partial_\nu \Omega \partial_\eta h_{\mu\gamma} \\ &\quad - 48\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-7} \partial_\alpha \Omega \partial_\beta \Omega \partial_\gamma \Omega \partial_\eta h_{\mu\nu} + 24\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-6} \partial_\alpha \Omega \partial_\gamma \partial_\beta \Omega \partial_\eta h_{\mu\nu} \\ &\quad - 6\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-6} \partial_\alpha \Omega \partial_\gamma \partial_\mu \Omega \partial_\eta h_{\nu\beta} - 6\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-6} \partial_\alpha \Omega \partial_\beta \partial_\mu \Omega \partial_\eta h_{\nu\gamma} \\ &\quad + 20\eta^{\alpha\beta} \eta^{\gamma\kappa} \eta^{\eta\lambda} \eta_{\mu\nu} h_{\kappa\lambda} \Omega^{-8} \partial_\alpha \Omega \partial_\beta \Omega \partial_\gamma \Omega \partial_\eta \Omega + 60\eta^{\alpha\beta} \eta^{\gamma\eta} h_{\mu\nu} \Omega^{-8} \partial_\alpha \Omega \partial_\beta \Omega \partial_\gamma \Omega \partial_\eta \Omega \\ &\quad + 2\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-5} \partial_\alpha \partial_\nu \Omega \partial_\eta \partial_\beta h_{\mu\gamma} - 4\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-5} \partial_\gamma \partial_\alpha \Omega \partial_\eta \partial_\beta h_{\mu\nu} \\ &\quad + 2\eta^{\alpha\beta} \eta^{\gamma\eta} \Omega^{-5} \partial_\alpha \partial_\mu \Omega \partial_\eta \partial_\beta h_{\nu\gamma} + 6\eta^{\alpha\beta} \eta^{\gamma\eta} h_{\mu\nu} \Omega^{-6} \partial_\gamma \partial_\alpha \Omega \partial_\eta \partial_\beta \Omega \\ &\quad - 6\eta^{\alpha\beta} \eta^{\gamma\eta} h_{\nu\alpha} \Omega^{-6} \partial_\gamma \partial_\mu \Omega \partial_\eta \partial_\beta \Omega - 6\eta^{\alpha\beta} \eta^{\gamma\eta} h_{\mu\alpha} \Omega^{-6} \partial_\gamma \partial_\nu \Omega \partial_\eta \partial_\beta \Omega \end{aligned}$$





$$+ \frac{1}{3} \eta^{\alpha\beta} \eta^{\gamma\eta} \eta^{\kappa\lambda} \eta_{\mu\nu} \partial_\lambda \partial_\kappa \partial_\eta \partial_\beta f_{\alpha\gamma} \Big). \quad (15)$$

### $\delta W_{\mu\nu}(K_{\mu\nu})$

In terms of  $K_{\mu\nu}$ , we may start with (??) and substitute  $h_{\mu\nu} = K_{\mu\nu} + \frac{1}{4}\Omega^2\eta_{\mu\nu}h$ , which evaluates to the 151 term

[illegible]



$$\begin{aligned}
& -16\eta^{\alpha\beta}\eta^{\gamma\eta}\Omega^{-7}\partial_\alpha\Omega\partial_\eta K_{\beta\gamma}\partial_\mu\Omega\partial_\nu\Omega + 2\eta^{\alpha\beta}\eta^{\gamma\eta}\Omega^{-6}\partial_\eta\partial_\beta K_{\alpha\gamma}\partial_\mu\Omega\partial_\nu\Omega \\
& -8\eta^{\alpha\beta}\eta^{\gamma\eta}K_{\alpha\gamma}\Omega^{-7}\partial_\eta\partial_\beta\Omega\partial_\mu\Omega\partial_\nu\Omega - \frac{2}{3}\eta^{\alpha\beta}\eta^{\gamma\eta}\Omega^{-5}\partial_\gamma\partial_\alpha\Omega\partial_\nu\partial_\mu K_{\beta\eta} \\
& + 2\eta^{\alpha\gamma}\eta^{\beta\eta}\Omega^{-6}\partial_\alpha\Omega\partial_\beta\Omega\partial_\nu\partial_\mu K_{\gamma\eta} - 8\eta^{\alpha\gamma}\eta^{\beta\eta}K_{\gamma\eta}\Omega^{-7}\partial_\alpha\Omega\partial_\beta\Omega\partial_\nu\partial_\mu\Omega \\
& + 4\eta^{\alpha\beta}\eta^{\gamma\eta}\Omega^{-6}\partial_\alpha\Omega\partial_\eta K_{\beta\gamma}\partial_\nu\partial_\mu\Omega - \frac{2}{3}\eta^{\alpha\beta}\eta^{\gamma\eta}\Omega^{-5}\partial_\eta\partial_\beta K_{\alpha\gamma}\partial_\nu\partial_\mu\Omega \\
& + 2\eta^{\alpha\beta}\eta^{\gamma\eta}K_{\alpha\gamma}\Omega^{-6}\partial_\eta\partial_\beta\Omega\partial_\nu\partial_\mu\Omega.
\end{aligned} \tag{16}$$

The dependence upon  $h$  drops out as expected since  $W_{\mu\nu}^{(0)} = 0$ .

Now taking  $K_{\mu\nu} = \Omega^2 k_{\mu\nu}$  with trace

$$k = \eta^{\sigma\tau} k_{\sigma\tau} = \Omega^{-2} \eta^{\sigma\tau} K_{\sigma\tau} = 0, \tag{17}$$

(??) evaluates to

$$\begin{aligned}
\delta W_{\mu\nu}(k_{\mu\nu}) = & \frac{1}{2}\Omega^{-2} \left( \frac{2}{3}\eta^{\alpha\beta}\eta^{\gamma\eta}\partial_\eta\partial_\beta\partial_\nu\partial_\mu k_{\alpha\gamma} + \eta^{\alpha\beta}\eta^{\gamma\eta}\partial_\eta\partial_\gamma\partial_\beta\partial_\alpha k_{\mu\nu} - \eta^{\alpha\beta}\eta^{\gamma\eta}\partial_\eta\partial_\gamma\partial_\beta\partial_\mu k_{\nu\alpha} \right. \\
& \left. - \eta^{\alpha\beta}\eta^{\gamma\eta}\partial_\eta\partial_\gamma\partial_\beta\partial_\nu k_{\mu\alpha} + \frac{1}{3}\eta^{\alpha\beta}\eta^{\gamma\eta}\eta^{\kappa\lambda}\eta_{\mu\nu}\partial_\lambda\partial_\kappa\partial_\eta\partial_\beta k_{\alpha\gamma} \right).
\end{aligned} \tag{18}$$

## Summary

Evaluated within a conformal to Minkowski background

$$ds^2 = -(g_{\mu\nu}^{(0)} + h_{\mu\nu})dx^\mu dx^\nu = -\Omega^2(x)(\eta_{\mu\nu} + f_{\mu\nu})dx^\mu dx^\nu, \tag{19}$$

the first order  $\delta W_{\mu\nu}$  may be expressed entirely in terms of the transverse traceless  $f_{\mu\nu}^{T\theta}$  as

$$\delta W_{\mu\nu} = \frac{1}{2}\Omega^{-2}\partial_\alpha\partial^\alpha\partial_\beta\partial^\beta f_{\mu\nu}^{T\theta}. \tag{20}$$

I have not checked, but such an  $f_{\mu\nu}^{T\theta}$  should be gauge invariant.

Alternatively,  $\delta W_{\mu\nu}$  may be expressed within the same geometry of (??) in terms of the traceless  $k_{\mu\nu} = \Omega^{-2}K_{\mu\nu}$  as

$$\begin{aligned}
\delta W_{\mu\nu}(k_{\mu\nu}) = & \frac{1}{2}\Omega^{-2} \left( \frac{2}{3}\eta^{\alpha\beta}\eta^{\gamma\eta}\partial_\eta\partial_\beta\partial_\nu\partial_\mu k_{\alpha\gamma} + \eta^{\alpha\beta}\eta^{\gamma\eta}\partial_\eta\partial_\gamma\partial_\beta\partial_\alpha k_{\mu\nu} - \eta^{\alpha\beta}\eta^{\gamma\eta}\partial_\eta\partial_\gamma\partial_\beta\partial_\mu k_{\nu\alpha} \right. \\
& \left. - \eta^{\alpha\beta}\eta^{\gamma\eta}\partial_\eta\partial_\gamma\partial_\beta\partial_\nu k_{\mu\alpha} + \frac{1}{3}\eta^{\alpha\beta}\eta^{\gamma\eta}\eta^{\kappa\lambda}\eta_{\mu\nu}\partial_\lambda\partial_\kappa\partial_\eta\partial_\beta k_{\alpha\gamma} \right).
\end{aligned} \tag{21}$$

Imposing the conformal gauge such that  $\partial^\nu k_{\mu\nu} = 0$  reduces the fluctuation to

$$\delta W_{\mu\nu}(k_{\mu\nu}) = \frac{1}{2}\Omega^{-2}\partial_\alpha\partial^\alpha\partial_\beta\partial^\beta k_{\mu\nu}. \tag{22}$$

These results coincide with the analogous (E37) and (E38).