

# Flat Cosmological Fluctuation Equations

We work to first order with Minkowski background:

$$g_{\mu\nu} = g_{\mu\nu}^{(0)} + g_{\mu\nu}^{(1)} = \eta_{\mu\nu} + h_{\mu\nu}.$$

In 3+1 splitting, we define vector tangent to 3-surfaces of constant (zero) curvature:

$$n_\mu = (-1, 0, 0, 0), \quad n_\mu n^\mu = -1.$$

Tensors are raised/lowered using the induced metric

$$\gamma_{\mu\nu} = \delta_{ij}, \quad \gamma_{0\mu} = 0.$$

The fluctuation equations in Minkowski standard gravity are:

$$\delta G_{\mu\nu} = -\frac{1}{2}\partial_\alpha\partial^\alpha\delta g_{\mu\nu} + \frac{1}{2}\partial_\alpha\partial_\mu\delta g_\nu{}^\alpha + \frac{1}{2}\partial_\alpha\partial_\nu\delta g_\mu{}^\alpha + \frac{1}{2}g_{\mu\nu}\partial^\alpha\partial_\alpha\delta g^\gamma{}_\gamma - \frac{1}{2}g_{\mu\nu}\partial_\gamma\partial^\alpha\delta g_\alpha{}^\gamma - \frac{1}{2}\partial_\mu\partial_\nu\delta g^\alpha{}_\alpha$$

The fluctuation equations in Minkowski conformal gravity are:

$$\begin{aligned} \delta W_{\mu\nu} = & \frac{1}{2}\partial_\beta\partial^\beta\partial_\alpha\partial^\alpha\delta g_{\mu\nu} + \frac{1}{6}g_{\mu\nu}\partial_\gamma\partial^\gamma\partial_\alpha\partial_\beta\delta g^{\alpha\beta} - \frac{1}{6}g_{\mu\nu}\partial_\gamma\partial^\gamma\partial_\beta\partial^\beta\delta g^\alpha{}_\alpha - \frac{1}{2}\partial_\beta\partial^\beta\partial_\mu\partial_\alpha\delta g_\nu{}^\alpha \\ & - \frac{1}{2}\partial_\beta\partial^\beta\partial_\nu\partial_\alpha\delta g_\mu{}^\alpha + \frac{1}{3}\partial_\mu\partial_\nu\partial_\alpha\partial_\beta\delta g^{\alpha\beta} + \frac{1}{6}\partial_\mu\partial_\nu\partial_\beta\partial^\beta\delta g^\alpha{}_\alpha. \end{aligned}$$

## Standard Gravity

### Newton Gauge

Gauge:

$$\begin{aligned} h_{\mu\nu} &= -2\phi n_\mu n_\nu - 2\psi\gamma_{\mu\nu} \\ ds^2 &= -(1+2\phi)dt^2 + \delta_{ij}(1-2\psi)dx^i dx^j \end{aligned}$$

Fluctuation Equations:

$$\begin{aligned} \delta G_{00} &= 2\nabla^2\psi \\ \delta G_{0i} &= 2\nabla_i\dot{\psi} \\ \delta G_{ij} &= \nabla_i\nabla_j(\psi - \phi) + \delta_{ij}(2\ddot{\psi} + \nabla^2\phi - \nabla^2\psi) \end{aligned}$$

## Conformal Gravity

### Newton Gauge

Gauge:

$$\begin{aligned} h_{\mu\nu} &= -2\phi n_\mu n_\nu - 2\psi\gamma_{\mu\nu} \\ ds^2 &= -(1+2\phi)dt^2 + \delta_{ij}(1-2\psi)dx^i dx^j \end{aligned}$$

Fluctuation Equations:

$$\begin{aligned}\delta W_{00} &= \frac{2}{3} \nabla^4 (\phi + \psi) \\ \delta W_{0i} &= \frac{2}{3} \nabla_i \nabla^2 (\dot{\phi} + \dot{\psi}) \\ \delta W_{ij} &= \nabla_i \nabla_j (\ddot{\phi} + \ddot{\psi} - \frac{1}{3} \nabla^2 \phi - \frac{1}{3} \nabla^2 \psi) - \frac{1}{3} \delta_{ij} \nabla^2 (\ddot{\phi} + \ddot{\psi} - \nabla^2 \phi - \nabla^2 \psi)\end{aligned}$$

## Standard Gravity

### No Gauge

We now look at full perturbations without a choice of gauge, separating them into their spin parts.

Gauge:

$$\begin{aligned}h_{\mu\nu} &= -2\phi n_\mu n_\nu - (B_\nu + \bar{\nabla}_\nu B) n_\mu - (B_\mu + \bar{\nabla}_\mu B) n_\nu - 2\gamma_{\mu\nu} \psi + \bar{\nabla}_\mu E_\nu + \bar{\nabla}_\nu E_\mu + 2\bar{\nabla}_\mu \bar{\nabla}_\nu E + 2E_{\mu\nu} \\ ds^2 &= -(1 + 2\phi) dt^2 + 2(B_i + \nabla_i B) dx^i dt + [-2\delta_{ij} \psi + (\nabla_i E_j + \nabla_j E_i) + 2\nabla_i \nabla_j E + 2E_{ij}] dx^i dx^j\end{aligned}$$

where

$$\bar{\nabla}_\mu = (0, \nabla_i), \quad n^\mu B_\mu = 0, \quad n^\mu n^\nu E_{\mu\nu} = 0, \quad \bar{\nabla}^\mu B_\mu = 0, \quad \bar{\nabla}^\mu E_\mu = 0, \quad \bar{\nabla}^\mu E_{\mu\nu} = 0.$$

Fluctuation Equations:

$$\begin{aligned}\delta G_{00} &= 2\nabla^2 \psi \\ \delta G_{0i} &= 2\nabla_i \dot{\psi} - \frac{1}{2} \nabla^2 B_i + \frac{1}{2} \nabla^2 \dot{E}_i \\ \delta G_{ij} &= 2\delta_{ij} \ddot{\psi} + (\delta_{ij} \nabla^2 - \nabla_i \nabla_j) (\phi - \psi + \dot{B} - \ddot{E}) - \frac{1}{2} (\nabla_i \dot{B}_j + \nabla_j \dot{B}_i) + \frac{1}{2} (\nabla_i \ddot{E} + \nabla_j \ddot{E}_i) + \ddot{E}_{ij} - \nabla^2 E_{ij}\end{aligned}$$

We see that the Einstein tensor itself separates into linear combinations of different spin tensors:

$$\begin{aligned}\delta G_{00}^{(S)} &= \delta G_{00} = 2\nabla^2 \psi \\ \delta G_{0i}^{(S)} &= 2\nabla_i \dot{\psi} \\ \delta G_{0i}^{(V)} &= -\frac{1}{2} \nabla^2 B_i + \frac{1}{2} \nabla^2 \dot{E}_i \\ \delta G_{ij}^{(S)} &= 2\delta_{ij} \ddot{\psi} + (\delta_{ij} \nabla^2 - \nabla_i \nabla_j) (\phi - \psi + \dot{B} - \ddot{E}) \\ \delta G_{ij}^{(V)} &= -\frac{1}{2} (\nabla_i \dot{B}_j + \nabla_j \dot{B}_i) + \frac{1}{2} (\nabla_i \ddot{E} + \nabla_j \ddot{E}_i) \\ \delta G_{ij}^{(T)} &= \ddot{E}_{ij} - \nabla^2 E_{ij}\end{aligned}$$

## Conformal Gravity

### No Gauge

Gauge:

$$\begin{aligned}h_{\mu\nu} &= -2\phi n_\mu n_\nu - (B_\nu + \bar{\nabla}_\nu B) n_\mu - (B_\mu + \bar{\nabla}_\mu B) n_\nu - 2\gamma_{\mu\nu} \psi + \bar{\nabla}_\mu E_\nu + \bar{\nabla}_\nu E_\mu + 2\bar{\nabla}_\mu \bar{\nabla}_\nu E + 2E_{\mu\nu} \\ ds^2 &= -(1 + 2\phi) dt^2 + 2(B_i + \nabla_i B) dx^i dt + [-2\delta_{ij} \psi + (\nabla_i E_j + \nabla_j E_i) + 2\nabla_i \nabla_j E + 2E_{ij}] dx^i dx^j\end{aligned}$$

where

$$\bar{\nabla}_\mu = (0, \nabla_i), \quad n^\mu B_\mu = 0, \quad n^\mu n^\nu E_{\mu\nu} = 0, \quad \bar{\nabla}^\mu B_\mu = 0, \quad \bar{\nabla}^\mu E_\mu = 0, \quad \bar{\nabla}^\mu E_{\mu\nu} = 0.$$

Fluctuation Equations:

$$\begin{aligned}
\delta W_{00} &= \frac{2}{3} \nabla^4 (\phi + \psi + \dot{B} - \ddot{E}) \\
\delta W_{0i} &= \frac{2}{3} \nabla_i \nabla^2 (\dot{\phi} + \dot{\psi} + \ddot{B} - \ddot{E}) + \frac{1}{2} \nabla^2 (\ddot{B}_i - \ddot{E}_i - \nabla^2 B_i + \nabla^2 \dot{E}_i) \\
\delta W_{ij} &= \frac{1}{3} \delta_{ij} \nabla^2 \left( -\ddot{\phi} - \ddot{\psi} - \ddot{B} + \ddot{E} + \nabla^2 (\phi + \psi + \dot{B} - \ddot{E}) \right) + \nabla_i \nabla_j \left( \ddot{\phi} + \ddot{\psi} + \ddot{B} - \ddot{E} - \frac{1}{3} \nabla^2 (\phi + \psi + \dot{B} - \ddot{E}) \right) \\
&\quad + \frac{1}{2} \nabla_i \left( \ddot{B}_j - \ddot{E}_j - \nabla^2 (\dot{B}_j - \dot{E}_j) \right) + \frac{1}{2} \nabla_j \left( \ddot{B}_i - \ddot{E}_i - \nabla^2 (\dot{B}_i - \dot{E}_i) \right) \\
&\quad - \ddot{E}_{ij} + 2 \nabla^2 \ddot{E}_{ij} - \nabla^4 E_{ij}
\end{aligned}$$

Again, the Weyl tensor itself separates into linear combinations of different spin tensors:

$$\begin{aligned}
\delta W_{00}^{(S)} &= \delta W_{00} = \frac{2}{3} \nabla^4 (\phi + \psi + \dot{B} - \ddot{E}) \\
\delta W_{0i}^{(S)} &= \frac{2}{3} \nabla_i \nabla^2 (\dot{\phi} + \dot{\psi} + \ddot{B} - \ddot{E}) \\
\delta W_{0i}^{(V)} &= \frac{1}{2} \nabla^2 (\ddot{B}_i - \ddot{E}_i - \nabla^2 B_i + \nabla^2 \dot{E}_i) \\
\delta W_{ij}^{(S)} &= \frac{1}{3} \delta_{ij} \nabla^2 \left( -\ddot{\phi} - \ddot{\psi} - \ddot{B} + \ddot{E} + \nabla^2 (\phi + \psi + \dot{B} - \ddot{E}) \right) + \nabla_i \nabla_j \left( \ddot{\phi} + \ddot{\psi} + \ddot{B} - \ddot{E} - \frac{1}{3} \nabla^2 (\phi + \psi + \dot{B} - \ddot{E}) \right) \\
\delta W_{ij}^{(V)} &= \frac{1}{2} \nabla_i \left( \ddot{B}_j - \ddot{E}_j - \nabla^2 (\dot{B}_j - \dot{E}_j) \right) + \frac{1}{2} \nabla_j \left( \ddot{B}_i - \ddot{E}_i - \nabla^2 (\dot{B}_i - \dot{E}_i) \right) \\
\delta W_{ij}^{(T)} &= -\ddot{E}_{ij} + 2 \nabla^2 \ddot{E}_{ij} - \nabla^4 E_{ij} = -\bar{\nabla}_\alpha \bar{\nabla}^\alpha \bar{\nabla}_\beta \bar{\nabla}^\beta E_{ij}.
\end{aligned}$$

## Supplementary

### Ricci Tensor

Gauge:

$$\begin{aligned}
h_{\mu\nu} &= -2\phi n_\mu n_\nu - (B_\nu + \bar{\nabla}_\nu B) n_\mu - (B_\mu + \bar{\nabla}_\mu B) n_\nu - 2\gamma_{\mu\nu} \psi + \bar{\nabla}_\mu E_\nu + \bar{\nabla}_\nu E_\mu + 2\bar{\nabla}_\mu \bar{\nabla}_\nu E + 2E_{\mu\nu} \\
ds^2 &= -(1 + 2\phi) dt^2 + 2(B_i + \nabla_i B) dx^i dt + [-2\delta_{ij} \psi + (\nabla_i E_j + \nabla_j E_i) + 2\nabla_i \nabla_j E + 2E_{ij}] dx^i dx^j
\end{aligned}$$

where

$$\bar{\nabla}_\mu = (0, \nabla_i), \quad n^\mu B_\mu = 0, \quad n^\mu n^\nu E_{\mu\nu} = 0, \quad \bar{\nabla}^\mu B_\mu = 0, \quad \bar{\nabla}^\mu E_\mu = 0, \quad \bar{\nabla}^\mu E_{\mu\nu} = 0.$$

Fluctuation Equations:

$$\begin{aligned}
\delta R_{00} &= 3\ddot{\psi} + \nabla^2 (\phi + \dot{B} - \ddot{E}) \\
\delta R_{0i} &= -\frac{1}{2} \nabla^2 (B_i + \dot{E}_i) + 2\nabla_i \dot{\psi} \\
\delta R_{ij} &= -\delta_{ij} (\ddot{\psi} - \nabla^2 \psi) + \nabla_i \nabla_j (-\phi + \psi - \dot{B} + \ddot{E}) - \frac{1}{2} \nabla_i (\dot{B}_j - \dot{E}_j) - \frac{1}{2} \nabla_j (\dot{B}_i - \dot{E}_i) + \ddot{E}_{ij} - \nabla^2 E_{ij}
\end{aligned}$$

Separation of the Riemann tensor into spin parts:

$$\begin{aligned}
\delta R_{00}^{(S)} &= \delta R_{00} = 3\ddot{\psi} + \nabla^2 (\phi + \dot{B} - \ddot{E}) \\
\delta R_{0i}^{(S)} &= 2\nabla_i \dot{\psi}
\end{aligned}$$

$$\begin{aligned}
\delta R_{0i}^{(V)} &= -\frac{1}{2}\nabla^2(B_i + \dot{E}_i) \\
\delta R_{ij}^{(S)} &= -\delta_{ij}(\ddot{\psi} - \nabla^2\psi) + \nabla_i\nabla_j(-\phi + \psi - \dot{B} + \ddot{E}) \\
\delta R_{ij}^{(V)} &= -\frac{1}{2}\nabla_i(\dot{B}_j - \ddot{E}_j) - \frac{1}{2}\nabla_j(\dot{B}_i - \ddot{E}_i) \\
\delta R_{ij}^{(T)} &= \ddot{E}_{ij} - \nabla^2 E_{ij}.
\end{aligned}$$