Scalar Gauge Invariant RW SVT4 v1

1 Background

$$ds^2 = \Omega^2(\tau) \left(-d\tau^2 + \tilde{g}_{ij} dx^i dx^j \right), \qquad R_{ij} = -2k \tilde{g}_{ij}$$
(1.1)

$$R_{\lambda\mu\nu\kappa} = -\frac{1}{6}g_{\lambda\nu}g_{\mu\kappa}R + \frac{1}{6}g_{\lambda\kappa}g_{\mu\nu}R - \frac{1}{2}g_{\mu\nu}R_{\lambda\kappa} + \frac{1}{2}g_{\mu\kappa}R_{\lambda\nu} + \frac{1}{2}g_{\lambda\nu}R_{\mu\kappa} - \frac{1}{2}g_{\lambda\kappa}R_{\mu\nu}$$
(1.2)

$$R_{\mu\nu} = (A+B)U_{\mu}U_{\nu} + g_{\mu\nu}B, \qquad R = 3B - A \tag{1.3}$$

$$G_{\mu\nu} = \frac{1}{2}Ag_{\mu\nu} - \frac{1}{2}Bg_{\mu\nu} + AU_{\mu}U_{\nu} + BU_{\mu}U_{\nu}$$
 (1.4)

$$g^{\mu\nu}G_{\mu\nu} = A - 3B \tag{1.5}$$

$$T_{\mu\nu} = (\rho + p)U_{\mu}U_{\nu} + pg_{\mu\nu} \tag{1.6}$$

$$g^{\mu\nu}T_{\mu\nu} = 3p - \rho \tag{1.7}$$

$$\Delta_{\mu\nu}^{(0)} = \frac{1}{2}Ag_{\mu\nu} - \frac{1}{2}Bg_{\mu\nu} + g_{\mu\nu}p + AU_{\mu}U_{\nu} + BU_{\mu}U_{\nu} + pU_{\mu}U_{\nu} + U_{\mu}U_{\nu}\rho \tag{1.8}$$

$$g^{\mu\nu}\Delta^{(0)}_{\mu\nu} = A - 3B + 3p - \rho \tag{1.9}$$

$$A = -\frac{1}{2}(3p + \rho)$$

= $-3\dot{\Omega}^2\Omega^{-4} + 3\ddot{\Omega}\Omega^{-3}$ (1.10)

$$B = \frac{1}{2}(p - \rho)$$

= $-\dot{\Omega}^2 \Omega^{-4} - \ddot{\Omega}\Omega^{-3} - 2k\Omega^{-2}$ (1.11)

$$\rho = \frac{1}{2}(-A - 3B)
= 3\dot{\Omega}^2 \Omega^{-4} + 3k\Omega^{-2}$$
(1.12)

$$p = \frac{1}{2}(-A+B)$$

= $\dot{\Omega}^2 \Omega^{-4} - 2\ddot{\Omega}\Omega^{-3} - k\Omega^{-2}$ (1.13)

1.1 Identities

$$\nabla_{\alpha}R = 3\nabla_{\alpha}B - \nabla_{\alpha}A \tag{1.14}$$

$$= -2\nabla_{\alpha}A - (A+B)(U^{\beta}\nabla_{\alpha}U_{\beta} - 2U_{\alpha}\nabla_{\beta}U^{\beta}) \tag{1.15}$$

$$= 2\nabla_{\alpha}B + 2U_{\alpha}U^{\beta}(\nabla_{\beta}A + \nabla_{\beta}B) + 2(A+B)U_{\alpha}\nabla_{\beta}U^{\beta}$$
(1.16)

2 Trace $g^{\mu\nu}\Delta_{\mu\nu}$

$$ds^2 = (g_{\mu\nu} + h_{\mu\nu})dx^{\mu}dx^{\nu} \tag{2.1}$$

$$h_{\mu\nu} = -2g_{\mu\nu}\chi + 2\nabla_{\mu}\nabla_{\nu}F + \nabla_{\mu}F_{\nu} + \nabla_{\nu}F_{\mu} + 2F_{\mu\nu}$$
 (2.2)

$$g^{\mu\nu}F_{\mu\nu} = 0, \quad \nabla^{\mu}F_{\mu\nu} = 0, \quad \nabla^{\mu}F_{\mu} = 0$$
 (2.3)

$$U^{\mu}\delta U_{\mu} = \frac{1}{2}U^{\mu}U^{\nu}h_{\mu\nu}, \qquad U^{\mu}U_{\mu} = -1 \tag{2.4}$$

$$Q_{\mu} \equiv F_{\mu} + \nabla_{\mu}F \tag{2.5}$$

$$g^{\mu\nu}\delta G_{\mu\nu} = 6\nabla_{\alpha}\nabla^{\alpha}\chi - R\nabla_{\alpha}Q^{\alpha} - Q^{\alpha}\nabla_{\alpha}R + 2R_{\alpha\beta}\nabla^{\beta}Q^{\alpha} + 4F^{\alpha\beta}R_{\alpha\beta}$$
 (2.6)

$$g^{\mu\nu}\delta T_{\mu\nu} = 3\delta p - \delta\rho - 2R\chi + R\nabla_{\alpha}Q^{\alpha} - 2R_{\alpha\beta}\nabla^{\beta}Q^{\alpha} - 2F^{\alpha\beta}R_{\alpha\beta}$$
 (2.7)

$$g^{\mu\nu}\Delta_{\mu\nu} = 3\delta p - \delta\rho - 2R\chi + 6\nabla_{\alpha}\nabla^{\alpha}\chi - Q^{\alpha}\nabla_{\alpha}R + 2F^{\alpha\beta}R_{\alpha\beta}$$
 (2.8)

2.1 L^{-2} Scalars

Let S_2 be a linear superposition of all scalars with dimension L^{-2} ,

$$S_2 = A_1 \nabla^{\alpha} \nabla^{\beta} h_{\alpha\beta} + A_2 \nabla_{\alpha} \nabla^{\alpha} h + A_3 R h + A_4 R^{\alpha\beta} h_{\alpha\beta}$$
 (2.9)

$$= -2(4A_3 + A_4)R\chi - 2(A_1 + 4A_2)\nabla_{\alpha}\nabla^{\alpha}\chi + 2A_3R\nabla_{\alpha}Q^{\alpha} - A_1Q^{\alpha}\nabla_{\alpha}R +2(A_1 + A_2)\nabla_{\beta}\nabla^{\beta}\nabla_{\alpha}Q^{\alpha} + 2(-A_1 + A_4)R_{\alpha\beta}\nabla^{\beta}Q^{\alpha} + 2A_4F^{\alpha\beta}R_{\alpha\beta}.$$
(2.10)

Setting $\{A_1, A_2, A_3, A_4\} = \{1, -1, 0, 1\}$ we have

$$S_2^{\{1,-1,0,1\}} = h_{\alpha\beta}R^{\alpha\beta} - \nabla_{\alpha}\nabla^{\alpha}h + \nabla^{\alpha}\nabla^{\beta}h_{\alpha\beta}$$

$$= 6\nabla_{\alpha}\nabla^{\alpha}\chi - 2R\chi - Q^{\alpha}\nabla_{\alpha}R + 2F^{\alpha\beta}R_{\alpha\beta}.$$
(2.11)

This combination was chosen because in a maximally symmetric space, it follows that

$$S_2^{\{1,-1,0,1\}} = 6\nabla_\alpha \nabla^\alpha \chi - 2R\chi,$$
 (2.12)

which aligns with the gauge invariant scalar found in $TT_Projection_Curved_v4$. However, if our background does not possess maximal symmetry, the combination $h_{\alpha\beta}R^{\alpha\beta} - \nabla_{\alpha}\nabla^{\alpha}h + \nabla^{\alpha}\nabla^{\beta}h_{\alpha\beta}$ is not gauge invariant. Specifically, it transforms as

$$(h_{\alpha\beta}R^{\alpha\beta} - \nabla_{\alpha}\nabla^{\alpha}h + \nabla^{\alpha}\nabla^{\beta}h_{\alpha\beta}) \to (h_{\alpha\beta}R^{\alpha\beta} - \nabla_{\alpha}\nabla^{\alpha}h + \nabla^{\alpha}\nabla^{\beta}h_{\alpha\beta}) - \epsilon^{\alpha}\nabla_{\alpha}R. \tag{2.13}$$

Now, since we also note that

$$g^{\mu\nu}\Delta_{\mu\nu} - (3\delta p - \delta\rho) = -2R\chi + 6\nabla_{\alpha}\nabla^{\alpha}\chi - Q^{\alpha}\nabla_{\alpha}R + 2F^{\alpha\beta}R_{\alpha\beta}$$
 (2.14)

$$= h_{\alpha\beta}R^{\alpha\beta} - \nabla_{\alpha}\nabla^{\alpha}h + \nabla^{\alpha}\nabla^{\beta}h_{\alpha\beta}, \tag{2.15}$$

we conclude from the gauge invariance of $\Delta_{\mu\nu}$ that

$$(3\delta p - \delta \rho) \to (3\delta p - \delta \rho) + \epsilon^{\alpha} \nabla_{\alpha} R. \tag{2.16}$$

By analyzing the gauge transformation of the general S_2 , we conclude that is there does not exist a scalar of order L^{-2} that is gauge invariant.

2.2 L^{-4} Scalars

From Asanka's $SVT_{in}RW_{full}$, we see that we must go to fourth order derivatives to find the SVT3 scalar gauge invariants. If such a gauge invariant is fourth order, i.e. L^{-4} , then it must reduce to a factorized expression of $6\nabla_{\alpha}\nabla^{\alpha}\chi - 2R\chi$ within the maximally symmetric limit. Analogous to S_2 , we find that there are 18 possible scalars with dimension L^{-4} .

$$S_{4} = (A_{15} + 2A_{6})\nabla_{\alpha}\nabla_{\beta}\epsilon^{\beta}\nabla^{\alpha}R + (2A_{13} + A_{17})R^{\alpha\beta}\nabla_{\beta}\nabla_{\alpha}\nabla_{\gamma}\epsilon^{\gamma} + (A_{15} - A_{8})\nabla^{\alpha}R\nabla_{\beta}\nabla^{\beta}\epsilon_{\alpha} + \nabla_{\alpha}\epsilon^{\alpha}(\frac{2}{3}(-A_{4} + A_{8} + 3A_{9})R^{2} + 2(A_{12} + A_{4} - A_{8})R_{\beta\gamma}R^{\beta\gamma} + 2A_{1}\nabla_{\beta}\nabla^{\beta}R) + 2(A_{11} + A_{2})R\nabla_{\beta}\nabla^{\beta}\nabla_{\alpha}\epsilon^{\alpha} + \frac{1}{6}\epsilon^{\alpha}\left((-6A_{11} + A_{4} - A_{8})R\nabla_{\alpha}R - 6A_{8}\nabla_{\alpha}\nabla_{\beta}\nabla^{\beta}R - (6A_{15} + A_{4} - 7A_{8})R_{\alpha\beta}\nabla^{\beta}R + 6R^{\beta\gamma}\left((A_{4} - A_{8})\nabla_{\alpha}R_{\beta\gamma} - (A_{17} + A_{4} - A_{8})\nabla_{\gamma}R_{\alpha\beta}\right)\right) + \nabla^{\beta}\epsilon^{\alpha}\left(\frac{2}{3}(3A_{10} - 3A_{11} + 4A_{4} - 4A_{8})RR_{\alpha\beta} + (2A_{14} - A_{17} - 6A_{4} + 6A_{8})R_{\alpha}{}^{\gamma}R_{\beta\gamma} + 2(A_{16} - A_{8})\nabla_{\beta}\nabla_{\alpha}R + 2(A_{3} - A_{8})\nabla_{\gamma}\nabla^{\gamma}R_{\alpha\beta}\right) + 2(A_{7} + A_{8})\nabla_{\gamma}\nabla^{\gamma}\nabla_{\beta}\nabla^{\beta}\nabla_{\alpha}\epsilon^{\alpha} + (A_{17} + 2A_{4} - 2A_{8})R_{\alpha\gamma}\nabla^{\gamma}\nabla_{\beta}\nabla^{\beta}\epsilon^{\alpha} + (A_{18}(\nabla_{\alpha}R_{\beta\gamma} + \nabla_{\beta}R_{\alpha\gamma}) + 2(A_{5} - 2A_{8})\nabla_{\gamma}R_{\alpha\beta})\nabla^{\gamma}\nabla^{\beta}\epsilon^{\alpha}$$

$$(2.18)$$