




# PeriDEM – High-fidelity modeling of granular media consisting of deformable complex-shaped particles

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## Summary

Granular materials plays a crucial role in wide range of sectors including geotechnical, manufacturing, and mining. Predictive modeling of these materials under large loading becomes a challenging task due to deformation and breakage of particles and complex contact mechanism between complex-shaped particles undergoing considerable deformation. Focusing on the scenarios when particle deformation and breakage are crucial, PeriDEM model introduced in (Jha et al., 2021) is implemented in the PeriDEM library. The underlying idea is that individual particles are modeled as deformable solid using peridynamics theory, and the contact between two deforming particles are applied at locally at the contact region allowing modeling of complex-shaped particles. The integration of peridynamics within discrete element method (DEM) provides a flexible, hybrid framework that handles the contact mechanics at the particle boundary while accounting for the internal material response, including deformation and fracture. This opens up new avenues for exploring the interactions in granular systems, including developing constitutive laws for phenomenological continuum models, understanding effective behavior when subjected to large loading, and impact of particle shape on particle dynamics.

## Statement of Need

Granular materials are prevalent in numerous industrial sectors, including geotechnical, manufacturing, and mining. Current modeling techniques such as DEM struggle with accurately capturing the behavior of granular materials under extreme conditions, especially when dealing with complex geometries and deformable particles. PeriDEM overcomes the challenges and implements a high-fidelity framework combining DEM and peridynamics to allow for accurate simulations of granular systems under extreme loading conditions. PeriDEM library makes the implementation of high-fidelity approach transparent. The library depends on very limited external libraries and is easier to build on ubuntu and mac systems allowing quick testing and extension to user specific needs.

## Background

PeriDEM model was introduced in (Jha et al., 2021), where it demonstrated the ability to model both inter-particle contact and intra-particle fracture for arbitrarily shaped particles. The model is briefly described next.

## 35 Brief Introduction to PeriDEM Model

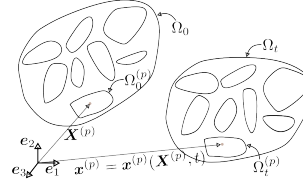


Figure 1: Motion of particle system.

36 Suppose a fixed frame of reference and  $\{e_i\}_{i=1}^d$  are orthonormal bases. Consider a collection  
 37 of  $N_P$  particles  $\Omega_0^{(p)}$ ,  $1 \leq p \leq N_P$ , where  $\Omega_0^{(p)} \subset \mathbb{R}^d$  with  $d = 2, 3$  represents the initial  
 38 configuration of particle  $p$ . Suppose  $\Omega_0 \supset \bigcup_{p=1}^{N_P} \Omega_0^{(p)}$  is the domain containing all particles; see  
 39 Figure 1. The particles in  $\Omega_0$  are dynamically evolving due to external boundary conditions and  
 40 internal interactions; let  $\Omega_t^{(p)}$  denote the configuration of particle  $p$  at time  $t \in (0, t_F]$ , and  
 41  $\Omega_t \supset \bigcup_{p=1}^{N_P} \Omega_t^{(p)}$  domain containing all particles at that time. The motion  $x^{(p)} = x^{(p)}(X^{(p)}, t)$   
 42 takes point  $X^{(p)} \in \Omega_0^{(p)}$  to  $x^{(p)} \in \Omega_t^{(p)}$ , and collectively, the motion is given by  $x = x(X, t) \in$   
 43  $\Omega_t$  for  $X \in \Omega_0$ . We assume the media is dry and not influenced by factors other than  
 44 mechanical loading (e.g., moisture and temperature are not considered). The configuration of  
 45 particles in  $\Omega_t$  at time  $t$  depends on various factors, such as material and geometrical properties,  
 46 contact mechanism, and external loading. Essentially, there are two types of interactions  
 47 present in the media:

- 48 (1.) *Intra-particle interaction* that models the deformation and internal forces in the particle  
 49 and
- 50 (2.) *Inter-particle interaction* that accounts for the contact between particles and the boundary  
 51 of the domain the particles are contained in.

52 In DEM, the first interaction is ignored, assuming particle deformation is insignificant compared  
 53 to the inter-particle interaction. On the other hand, PeriDEM, accounts for both interactions.

54 The balance of linear momentum for particle  $p$ ,  $1 \leq p \leq N_P$ , takes the form:

$$\rho^{(p)} \ddot{u}^{(p)}(X, t) = f_{int}^{(p)}(X, t) + f_{ext}^{(p)}(X, t), \quad \forall (X, t) \in \Omega_0^{(p)} \times (0, t_F), \quad (1)$$

55 where  $\rho^{(p)}$ ,  $f_{int}^{(p)}$ , and  $f_{ext}^{(p)}$  are density, and internal and external force densities. The above equa-  
 56 tion is complemented with initial conditions,  $u^{(p)}(X, 0) = u_0^{(p)}(X)$ ,  $\dot{u}^{(p)}(X, 0) = \dot{u}_0^{(p)}(X)$ ,  $X \in$   
 57  $\Omega_0^{(p)}$ .

### 58 Internal force - State-based peridynamics

59 Since all expressions in this paragraph are for a fixed particle  $p$ , we drop the superscript  $p$ ,  
 60 noting that material properties and other quantities can depend on the particle  $p$ . Following  
 61 (Silling et al., 2007) and simplified expression of state-based peridynamics force in (Jha et al.,  
 62 2021), the internal force takes the form, for  $X \in \Omega_0^{(p)}$ ,

$$f_{int}^{(p)}(X, t) = \int_{B_\epsilon(X) \cap \Omega_0^{(p)}} (T_X(Y) - T_Y(X)) \, dY, \quad (2)$$

63 where  $T_X(Y) - T_Y(X)$  is the force on  $X$  due to nonlocal interaction with  $Y$ . Let  $R = |Y - X|$   
 64 be the reference bond length,  $r = |x(Y) - x(X)|$  current bond length,  $s(Y, X) = (r - R)/R$   
 65 bond strain, then  $T_X(Y)$  is given by (Jha et al., 2021; Silling et al., 2007)

$$T_X(Y) = h(s)J(R/\epsilon) \left[ R\theta_X \left( \frac{3\kappa}{m_X} - \frac{15G}{3m_X} \right) + (r - R) \frac{15G}{m_X} \right] \frac{x(Y) - x(X)}{|x(Y) - x(X)|}, \quad (3)$$

66 where

$$\begin{aligned} m_X &= \int_{B_\epsilon(X) \cap \Omega_0^{(p)}} R^2 J(R/\epsilon) dY, \\ \theta_X &= h(s) \frac{3}{m_X} \int_{B_\epsilon(X) \cap \Omega_0^{(p)}} (r - R) R J(R/\epsilon) dY, \\ h(s) &= \begin{cases} 1, & \text{if } s < s_0 := \sqrt{\frac{\mathcal{G}_c}{(3G + (3/4)^4 [\kappa - 5G/3])\epsilon}}, \\ 0, & \text{otherwise.} \end{cases} \end{aligned} \quad (4)$$

67 In the above,  $J : [0, \infty) \rightarrow \mathbb{R}$  is the influence function,  $\kappa, G, \mathcal{G}_c$  are bulk and shear moduli  
68 and critical energy release rate, respectively. These parameters, including nonlocal length scale  
69  $\epsilon$ , could depend on the particle  $p$ .

#### 70 DEM-inspired contact forces

71 The external force density  $f_{ext}^{(p)}$  is generally expressed as

$$f_{ext}^{(p)} = \rho^{(p)} b + f^{\Omega_0, (p)} + \sum_{q \neq p} f^{(q), (p)}, \quad (5)$$

72 where  $b$  is body force per unit mass,  $f^{\Omega_0, (p)}$  and  $f^{(q), (p)}$  are contact forces due to interaction  
73 between particle  $p$  and container  $\Omega_0$  and neighboring particles  $q$ , respectively. In (Jha et al.,  
74 2021; ?), the contact between two particles is applied locally where the contact takes place;  
75 this is exemplified in ?? where contact between points  $y$  and  $x$  of two distinct particles  $p$   
76 and  $q$  is activated when they get sufficiently close. The contact forces are shown using a  
77 spring-dashpot-slider system. To fix the contact forces, consider a point  $X \in \Omega_0^{(p)}$  and let  
78  $R_c^{(q), (p)}$  be the critical contact radius (points in particles  $p$  and  $q$  interact if the distance is  
79 below this critical distance). Further, define the relative distance between two points  $Y \in \Omega_0^{(q)}$   
80 and  $X \in \Omega^{(p)}$  and normal and tangential directions as follows:

$$\begin{aligned} \Delta^{(q), (p)}(Y, X) &= |x^{(q)}(Y) - x^{(p)}(X)| - R_c^{(q), (p)}, \\ e_N^{(q), (p)}(Y, X) &= \frac{x^{(q)}(Y) - x^{(p)}(X)}{|x^{(q)}(Y) - x^{(p)}(X)|}, \\ e_T^{(q), (p)}(Y, X) &= [I - e_N^{(q), (p)}(Y, X) \otimes e_N^{(q), (p)}(Y, X)] \frac{\dot{x}^{(q)}(Y) - \dot{x}^{(p)}(X)}{|\dot{x}^{(q)}(Y) - \dot{x}^{(p)}(X)|}. \end{aligned} \quad (6)$$

81 Then the force on particle  $p$  due to contact with particle  $q$  can be written as (Jha et al.,  
82 2021}):

$$f^{(q), (p)}(X, t) = \int_{Y \in \Omega_0^{(q)} \cap B_{R^{(q), (p)}}(X)} (f_N^{(q), (p)}(Y, X) + f_T^{(q), (p)}(Y, X)) dY, \quad (7)$$

83 with normal and tangential forces following (Desai et al., 2019; Jha et al., 2021) given by, if  
84  $\Delta^{(q), (p)}(Y, X) < 0$ ,

$$f_N^{(q), (p)}(Y, X) = [\kappa_N^{(q), (p)} \Delta^{(q), (p)}(Y, X) - \beta_N^{(q), (p)} \dot{\Delta}^{(q), (p)}(Y, X)], \quad (8)$$

85 else  $f_N^{(q), (p)}(Y, X) = 0$ , and

$$f_T^{(q), (p)}(Y, X) = -\mu_T^{(q), (p)} |f_N^{(q), (p)}(Y, X)| e_T^{(q), (p)}. \quad (9)$$

## Implementation

PeriDEM is implemented as an open-source library in GitHub; see [PeriDEM](#). It is based on C++, and uses only handful external libraries which are included in the library in the external folder, allowing the code to be built and tested in ubuntu and mac systems relatively easily. Taskflow ([Huang et al., 2021](#)) is used for asynchronous multithreaded computation and Nanoflann ([Blanco & Rai, 2014](#)) for tree search to calculate neighbors for contact forces. MPI and metis ([Karypis & Kumar, 1997](#)) have recently been integrated to implement the distributed parallelism in near future. VTK is used to output the simulation files which can be visualized using paraview.

## Features

- Hybrid modeling using peridynamics and DEM for intra-particle and inter-particle interactions.
- Support for arbitrarily shaped particles, allowing for realistic simulation scenarios.
- Future work includes developing an adaptive modeling approach to enhance efficiency without compromising accuracy.
- Open-source implementation with support for HPC environments, leveraging modern multi-threading techniques for scalability.

## Brief implementation details

The main implementation of the model is carried out in the model directory [dem](#). The model is implemented in class [DEMModel](#). Function `DEMModel::run()` performs the simulation. We next look at some key methods in `DEMModel` in more details:

### `DEMModel::run()`

This function does three tasks:

```
void model::DEMModel::run(inp::Input *deck) {
    // initialize data
    init();

    // check for restart
    if (d_modelDeck_p->d_isRestartActive)
        restart(deck);

    // integrate in time
    integrate();
}
```

In `DEMModel::init()`, the simulation is prepared by reading the input files (such as `.yaml`, `.msh`, `particle_locations.csv` files).

### `DEMModel::integrate()`

Key steps in `DEMModel::integrate()` are

```
void model::DEMModel::run(inp::Input *deck) {
    // apply initial condition
    if (d_n == 0)
        applyInitialCondition();

    // apply loading
    computeExternalDisplacementBC();
}
```

```
computeForces();
```

```
// time step
```

```
for (size_t i = d_n; i < d_modelDeck_p->d_Nt; i++) {
```

```
    // advance simulation to next step
```

```
    integrateStep();
```

```
    // perform output if needed
```

```
    output();
```

```
}
```

```
}
```

113 In DEMModel::integrateStep(), we either utilize the central-difference scheme, implemented in DEMModel::integrateCD(), or the velocity-verlet scheme, implemented in  
114 DEMModel::integrateVerlet(). As an example, we look at DEMModel::integrateCD()  
115 method below:  
116

```
void model::DEMModel::integrateVerlet() {
```

```
    // update current position, displacement, and velocity of nodes
```

```
{
```

```
    tf::Executor executor(util::parallel::getNThreads());
```

```
    tf::Taskflow taskflow;
```

```
    taskflow.for_each_index(
```

```
        (std::size_t) 0, d_fPdCompNodes.size(), (std::size_t) 1,
```

```
        [this, dt, dim](std::size_t II) {
```

```
            auto i = this->d_fPdCompNodes[II];
```

```
            const auto rho = this->getDensity(i);
```

```
            const auto &fix = this->d_fix[i];
```

```
            for (int dof = 0; dof < dim; dof++) {
```

```
                if (util::methods::isFree(fix, dof)) {
```

```
                    this->d_v[i][dof] += 0.5 * (dt / rho) * this->d_f[i][dof];
```

```
                    this->d_u[i][dof] += dt * this->d_v[i][dof];
```

```
                    this->d_x[i][dof] += dt * this->d_u[i][dof];
```

```
                }
```

```
            }
```

```
        } // loop over nodes
```

```
    ); // for_each
```

```
    executor.run(taskflow).get();
```

```
}
```

```
// advance time
```

```
d_n++;
```

```
d_time += dt;
```

```
// update displacement bc
```

```
computeExternalDisplacementBC();
```

```
// compute force
```

```
computeForces();
```

```
// update velocity of nodes (similar to the above)
```

```
}
```

## 117 DEMModel::computeForces()

118 The key method in time integration is DEMModel::computeForces() In this function, we  
119 compute internal and external forces at each node of a particle and also account for the  
120 external boundary conditions. This function looks like

```
void model::DEMModel::computeForces() {
    // update the point cloud (make sure that d_x is updated along with displacement)
    auto pt_cloud_update_time = d_nsearch_p->updatePointCloud(d_x, true);
    pt_cloud_update_time += d_nsearch_p->setInputCloud();

    // reset forces to zero ...

    // compute peridynamic forces
    computePeridynamicForces();

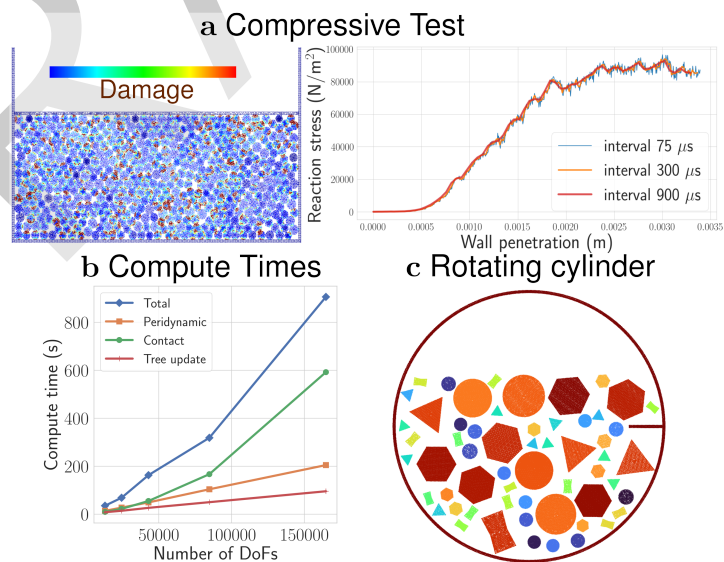
    // compute contact forces between particles
    computeContactForces();

    // Compute external forces
    computeExternalForces();
}
```

## 121 Further reading

122 Above gives the basic idea of simulation steps. For more thorough understanding of the  
123 implementation, interested readers can look at [demModel.cpp](#).

## 124 Examples



**Figure 2:** Nonlinear response under compression, {b} exponential growth of compute time due to nonlocality of internal and contact forces, and {c} rotating cylinder with nonspherical particles.

125 Examples are described in [examples/README.md](#) of the library. One of the key result is  
126 the compression of 502 circular and hexagon particles in a rectangular container by moving  
127 the top wall. The stress on the moving wall as a function of wall penetration becomes  
128 increasingly nonlinear, and media shows signs of yielding as the damage becomes extensive; see

129 [Figure 2a](#). Preliminary compute time analysis with an increasing number of particles shows an  
 130 exponential increase in compute time of contact and peridynamics forces, which is unsurprising  
 131 as both computations are nonlocal. Demonstration examples also include attrition of various  
 132 non-circular particles in a rotating cylinder [Figure 2c](#).

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