

Definitions		
symbol	description	unit
z	vertical coordinate	cm
ϵ_k	abundance of element k with respect to H	–
n_i	particle density of molecule i	cm^{-3}
$n_{\langle\text{H}\rangle}$	total hydrogen nuclei density	cm^{-3}
$\rho = \sum_i n_i m_i = \mu_{\text{H}} n_{\langle\text{H}\rangle} = \mu \sum_i n_i$	gas mass density	g cm^{-3}
m_i	mass of particle i	g
m_k	mass of element k	g
$\mu = \sum_i n_i m_i / \sum_i n_i$	mean gas particle mass	g
$\mu_{\text{H}} = \sum_k \epsilon_k m_k$	proportionality constant	g
T	gas temperature	K
$p = \frac{\rho}{\mu} kT = \sum_i n_i kT$	gas pressure	dyn cm^{-2}
D	diffusion coefficient	$\text{cm}^2 \text{s}^{-1}$

1 The Atmospheric Diffusion Problem

To simultaneously model the evolution of the chemical composition of the atmosphere and the “upper crust” of a hot rocky planet, we use the implicit/explicit time-dependent second-order diffusion solver DIFFUSE developed by Peter Woitke. At the top of the atmosphere, a modified fixed-flux boundary condition will be applied to allow for the Jeans escape of H_2 and He (see Sect. 4.1), whereas at the bottom, a modified fixed-concentration boundary condition will be applied to treat the outgassing/deposition of elements to/from the crust (see Sect. 4.2). By book-keeping the element fluxed through the lower boundary, we can also predict how the crust composition slowly changes.

Between the boundaries, we solve the second-order diffusion problem

$$\frac{\partial(n_{\langle\text{H}\rangle} \epsilon_k)}{\partial t} + \nabla(\vec{v} n_{\langle\text{H}\rangle} \epsilon_k) = \nabla(n_{\langle\text{H}\rangle} D \nabla \epsilon_k) . \quad (1)$$

Assuming a 1-d plane-parallel and static ($\vec{v} = 0$) atmosphere, the equations to solve are

$$\frac{d}{dt}(n_{\langle\text{H}\rangle} \epsilon_k) = \frac{d}{dz} \left(n_{\langle\text{H}\rangle} D \frac{d\epsilon_k}{dz} \right) . \quad (2)$$

Assuming, in addition, a constant density structure $n_{\langle\text{H}\rangle} = n_{\langle\text{H}\rangle}(z)$ the equations simplify to

$$\frac{d\epsilon_k}{dt} = \frac{1}{n_{\langle\text{H}\rangle}} \frac{d}{dz} \left(n_{\langle\text{H}\rangle} D \frac{d\epsilon_k}{dz} \right) , \quad (3)$$

where $\epsilon_k(z, t)$ are the height-dependent and time-dependent element abundances we wish to determine. We note that $n_{\langle\text{H}\rangle} = n_{\langle\text{H}\rangle}(z)$ and $D = D(z)$ vary by many orders of magnitude throughout the atmosphere, often more than the element abundances we wish to determine, therefore we *cannot* neglect the $\frac{dn_{\langle\text{H}\rangle}}{dz}$ and $\frac{dD}{dz}$ terms as

$$\frac{d\epsilon_k}{dt} \neq D \frac{d^2 \epsilon_k}{dz^2} . \quad (4)$$

The diffusion constant $D(z)$ is determined by proper microscopic diffusion in the uppermost layers D_{micro} , but otherwise mainly by turbulent quasi-diffusion D_{mix} due to vertical mixing processes excited by convective instabilities.

$$D(z) = D_{\text{micro}}(z) + D_{\text{mix}}(z) \quad (5)$$

The turbulent diffusion takes place on large spatial scales, hence we can safely assume that D_{mix} does not depend on the molecule (or element) we want to diffuse. Concerning D_{micro} , however, there is a small dependence on the size and reduced mass of the molecule with respect to collisions with H_2 , see Voitke & Helling (2003, Eq. (26) therein), which causes deviations of D as function of molecule by a factor of about 2 to 3. We will simply neglect these deviations in the following.

This neglect allows us to consider the diffusion of elements instead of calculating the diffusion of all individual molecules separately. The total particle density [cm^{-3}] of element k is given by

$$n_{\langle\text{H}\rangle} \epsilon_k = \sum_i s_{i,k} n_i \quad (6)$$

where $s_{i,k}$ is the stoichiometric factor of element k in molecule i , for example $s_{\text{H}_2\text{O},\text{H}} = 2$. Using Eq. (2), the diffusion of that total nuclei particle density is given by

$$\frac{d}{dt} (n_{\langle\text{H}\rangle} \epsilon_k) = \frac{d}{dt} \sum_i s_{i,k} n_i = \sum_i s_{i,k} \frac{dn_i}{dt} \quad (7)$$

$$= \sum_i s_{i,k} \frac{d}{dz} \left(n_{\langle\text{H}\rangle} D \frac{d}{dz} \left(\frac{n_i}{n_{\langle\text{H}\rangle}} \right) \right) \quad (8)$$

$$= \frac{d}{dz} \left(n_{\langle\text{H}\rangle} D \frac{d}{dz} \left(\sum_i s_{i,k} \frac{n_i}{n_{\langle\text{H}\rangle}} \right) \right) \quad (9)$$

$$= \frac{d}{dz} \left(n_{\langle\text{H}\rangle} D \frac{d\epsilon_k}{dz} \right) \quad (10)$$

2 Spatial Grid and Initial Conditions

We use an equidistant grid in $\{z_i | i = 0, \dots, I\}$ to discretise the 1d atmosphere. At the lower boundary ($i = 0$), the atmospheric gas is in contact with the crust. The crust is represented by column densities $N_{\text{cond},j}$ for a number of $j = 1, \dots, J$ condensed species.

At initialisation ($t = 0$)

3 Element Conservation

4 Boundary Conditions

4.1 Upper boundary condition

We specify the escaping element fluxes through the upper boundary of the atmosphere by applying the formula for Jeans escape \rightarrow Tian (2015).

Where to apply that formula? \rightarrow exobase \rightarrow Volkov et al. (2011).

Here, we need to decompose the elements into molecules, apply the Jeans escape to each molecule, and lump the escaping fluxes together to get the escaping element fluxes.

4.2 Lower boundary condition

At the bottom of the atmosphere, the atmospheric gas is in contact with the crust. At pressures are the order of fractions of a bar to several bars, hence the collision rates of gas particles with the crust are huge, leading to a fast relaxation towards phase equilibrium. We will therefore assume that the gas at the bottom of the atmosphere is saturated (in phase equilibrium) with respect to the solid/liquid materials forming the crust.

$$S_j(\epsilon_k) \Big|_{z=0} = 1 \quad (11)$$

As shown in Appendix B of (Woitke et al., 2017), the number N of simultaneously present condensates in phase equilibrium is limited and cannot exceed the number of elements K contained in them,

$$\begin{aligned} J &= \text{number of condensates in crust } j = 1, \dots, J \\ K &= \text{number of condensed elements in crust } k = 1, \dots, K \\ J &\leq K \\ D &= K - J \geq 0 \text{ number of non-limiting elements,} \end{aligned}$$

i.e. the number of condensing elements can (and usually will) exceed the number of condensates in the crust. In order to formulate the phase equilibrium boundary condition for the diffusion experiment, we need to solve J Eq. 11 for K elements which is not possible if $N < K$. In order to solve this problem, we need to make a case differentiation

- case 1: element not affected by condensation $j_k = 0$
- case 2: limiting element solution of ϵ_k from Eq. (11)
- case 3: non-limiting element use stoichiometric constraints to compute j_k

5 Updating the crust thickness and composition

References

- TIAN, F. (2015, May). Atmospheric Escape from Solar System Terrestrial Planets and Exoplanets. *Annual Review of Earth and Planetary Sciences* **43**, 459–476.
- VOLKOV, A. N., JOHNSON, R. E., TUCKER, O. J., ERWIN, J. T. (2011, March). Thermally Driven Atmospheric Escape: Transition from Hydrodynamic to Jeans Escape. *ApJL* **729**, L24.
- WOITKE, P., HELLING, C. (2003, February). Dust in brown dwarfs. II. The coupled problem of dust formation and sedimentation. *A&A* **399**, 297–313.
- WOITKE, P., HELLING, C., HUNTER, G. H., MILLARD, J. D., TURNER, G. E., WORTERS, M., BLECIC, J., STOCK, J. W. (2017). Equilibrium chemistry down to 100 K - Impact of silicates and phyllosilicates on carbon/oxygen ratio. *ArXiv e-prints* **1712.01010**, December.