

# Physics of electron rings

Zgoubi workshop

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# Overview

- Electron rings past and present
- How to store an electron
- Fast overview of symplectic transverse and longitudinal dynamics
- Addition of Synchrotron radiation in modelling and its impact

# Some electron storage rings past and present

Colliders:

- LEP (1989)
- HERA (1992-2007)
- PEP-II (1998-2008)
- KEK-B (2003-2010)

Synchrotron light sources:

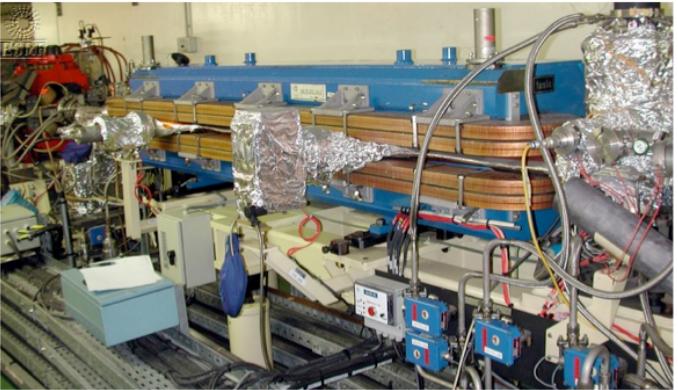
> 50 such facilities! (see <https://lightsources.org/>)

New round of upgrades decreasing emittance <200pm,  
largely based on hybrid MBA lattice from ESRF  
(MBA first done at Max-IV)  
Pushes towards diffraction limit

# How to store a high energy electron bunch?

- Electrons accelerated with linac+booster
- Inject electrons into a ring
- Dipole magnets bend the trajectory
- Quadrupole magnets focus the trajectory
- Sextupole magnets fix “chromaticity”
- RF cavities replenish synchrotron radiation energy lost and provide longitudinal focussing

# Storage ring components (pre-EBS ESRF)



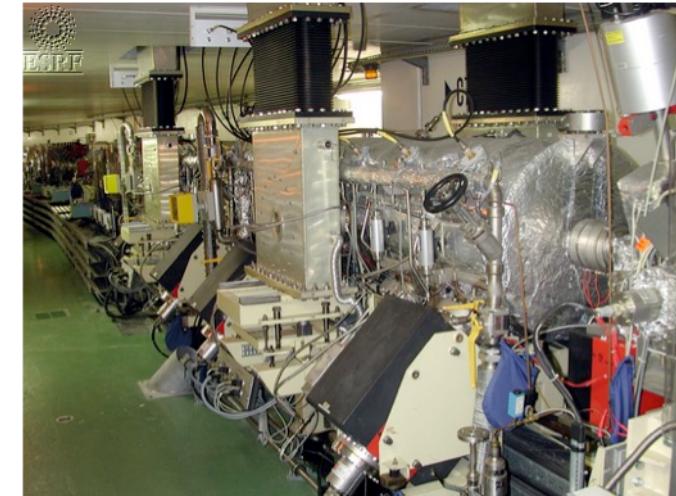
dipole



quadrupole



sextupole



RF cavity

# Transverse dynamics

## Equations of motion

Apply Newton's 2<sup>nd</sup> law for Lorentz force:

$$\dot{\vec{P}} = e(v\hat{z} \times \vec{B}) = ev(-B_1\hat{x} + B_1\hat{y})$$

divide by  $P_0$  and change time derivative to s- derivative, take  $v=c$ , and we have:

$$x'' + k_x(s)x = 0$$

$$y'' + k_y(s)y = 0$$

$$k_x = -\frac{B_1}{B\rho}$$

$$k_y = \frac{B_1}{B\rho}$$

$$\text{recall } B\rho = \frac{P}{e}$$

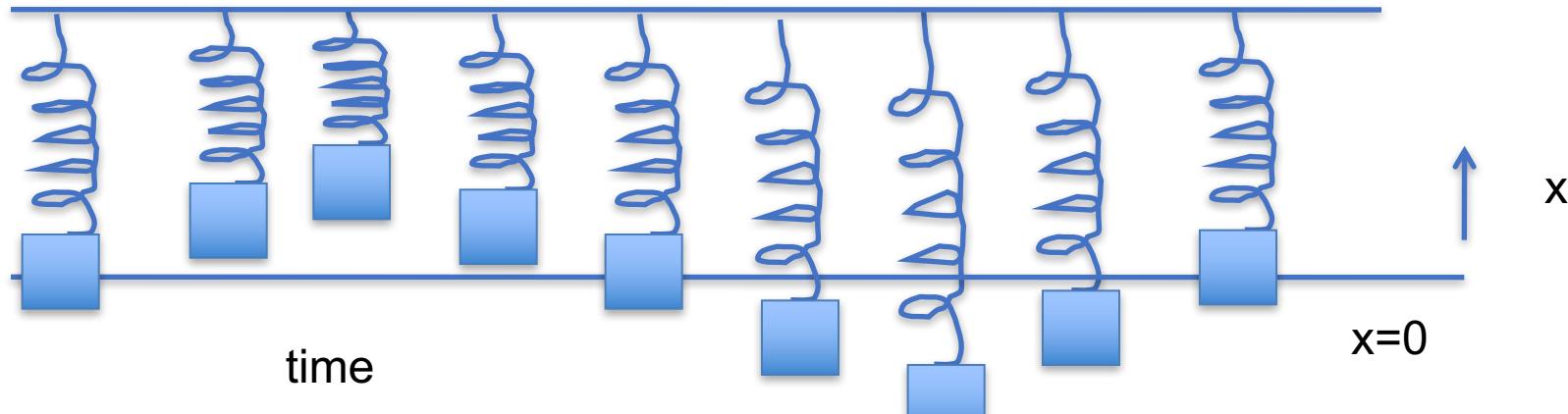
Harmonic oscillator with s-dependent, periodic spring constant. Known as Hill's equation.

$$k_{x,y}(s) = k_{x,y}(s + C)$$

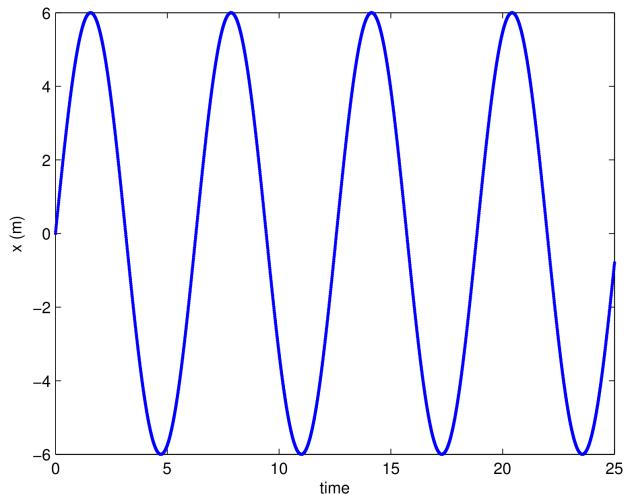
Can be derived from this Hamiltonian

$$H(x, x', y, y') = \frac{1}{2} \left( k_x x^2 + k_y y^2 + x'^2 + y'^2 \right)$$
$$x' = \frac{p_x}{P_0} = \frac{dx}{ds}$$
$$y' = \frac{p_y}{P_0} = \frac{dy}{ds}$$

# Phase space



configuration  
space  $x$  vs. time



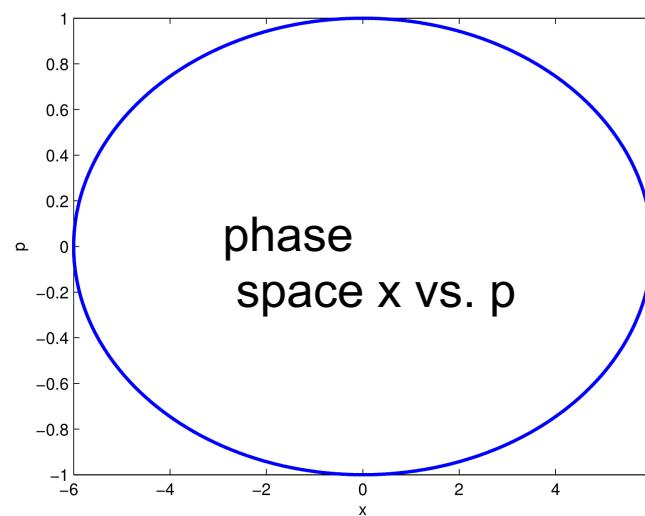
$$p_{x,y} = \gamma m v_{x,y}$$

for electron, we  
normalize with

$$P_0 = \gamma m v_s$$

and use

$$x' = \frac{p_x}{P_0} = \frac{dx}{ds}$$



# 2-D Matrix Analysis

## Courant-Snyder transformation

Given one turn map matrix, can we transform it into a rotation?

$$A^{-1}MA = R \quad R = \begin{pmatrix} \cos \mu & \sin \mu \\ -\sin \mu & \cos \mu \end{pmatrix}$$

again,  $\mu = 2\pi\nu$

One option for A:

together with

$$A = \begin{pmatrix} \frac{1}{\sqrt{\beta}} & 0 \\ \frac{\alpha}{\sqrt{\beta}} & \sqrt{\beta} \end{pmatrix} \quad \gamma = \frac{1 + \alpha^2}{\beta}$$

$$J_x = \gamma x^2 + 2\alpha x x' + \beta x'^2$$

is Courant-Snyder invariant

also, one can show that

$$\alpha = -\frac{1}{2} \frac{d\beta}{ds}$$

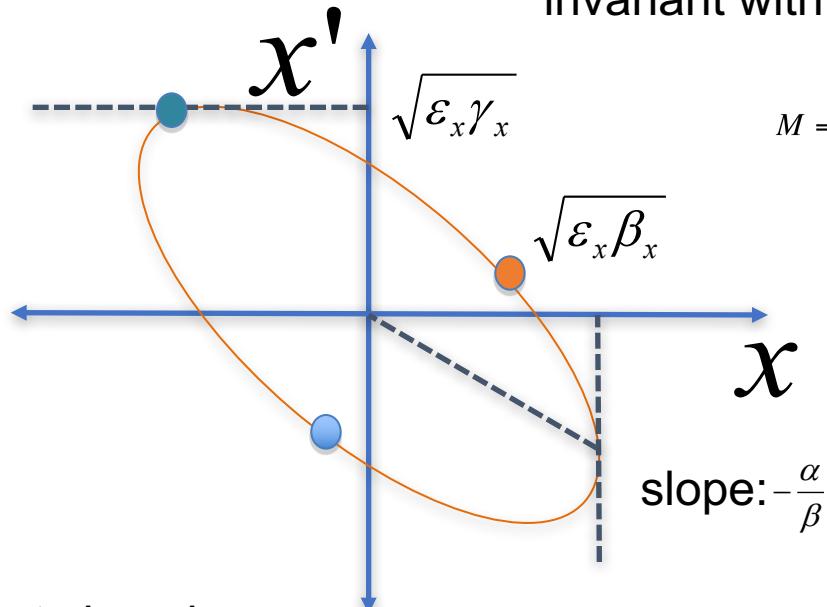
# Twiss Parameters

$$\begin{pmatrix} \alpha(s) \\ \beta(s) \\ \gamma(s) \end{pmatrix}$$

are known as 'Twiss Parameters'

measuring the position over time, it will oscillate

tune is defined by  
number of oscillations about closed orbit over 1 turn around ring.  
Note that the matrix only captures fractional part.



$$\gamma_x \eta_x^2 + 2\alpha_x \eta_x \eta_x' + \beta_x \eta_x'^2$$

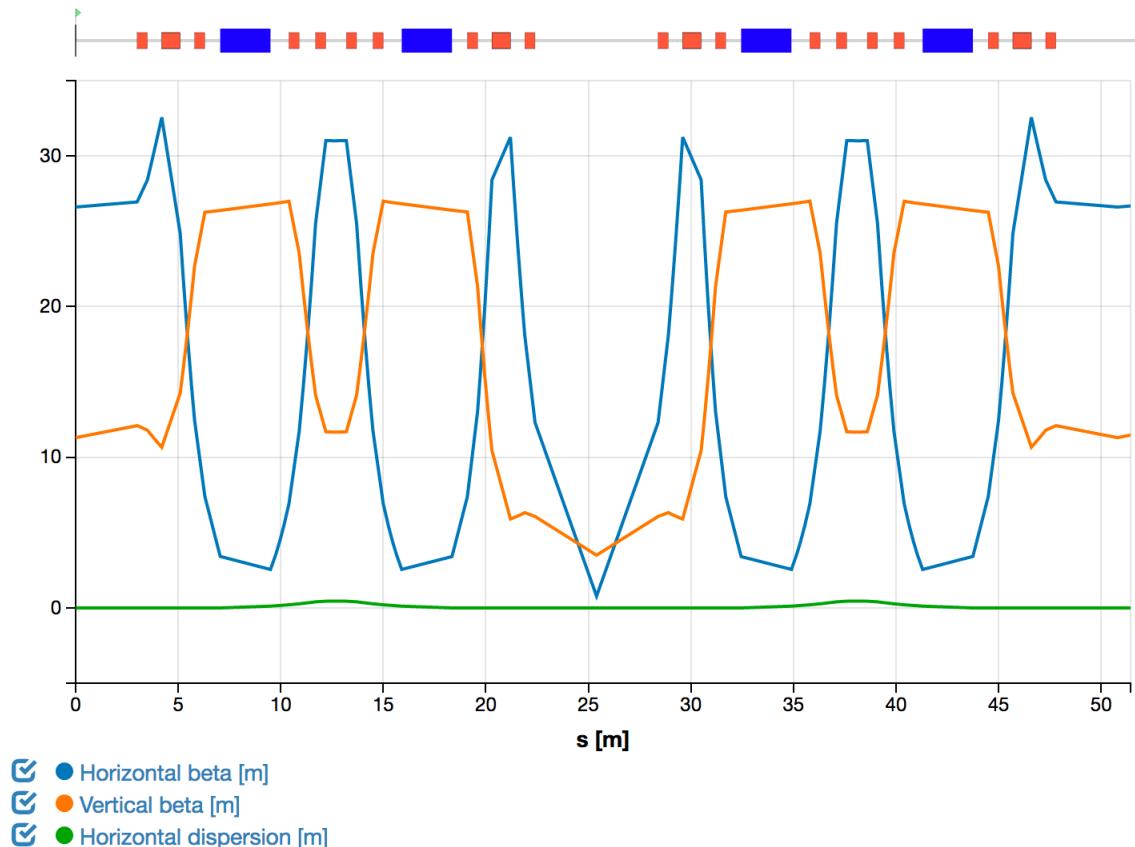
invariant with position around ring!

$$M = \begin{pmatrix} \cos \mu + \alpha \sin \mu & \beta \sin \mu \\ -\gamma \sin \mu & \cos \mu - \alpha \sin \mu \end{pmatrix}$$

- turn 1
- turn 2
- turn 3

This is at one position in the ring.

# Results for our model ESRF lattice



$\text{nu}_x=36.197428 \quad \text{nu}_y=11.202545$

# Longitudinal Dynamics

radiation with electrons, but not protons

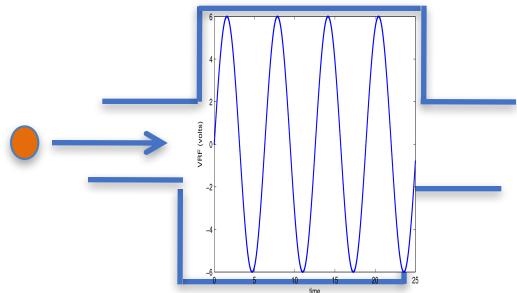
energy loss  
per turn is

$$U_0 = \frac{C_\gamma \beta^3 E_0^4}{\rho} \quad C_\gamma = 8.85 \times 10^{-5} \frac{m}{GeV^3}$$

$$U_0 = 4.88 MeV \quad \text{for present ESRF}$$

We need to provide this energy back, and also focus longitudinally.

# RF cavity dynamics



$$V(t) = \varepsilon \sin(\phi_{RF}(t) + \phi_0)$$

$$\phi_{RF}(t) = h\omega_0 t$$

To energy constant, we need

$$\Delta E = U_0$$

energy loss from  
synchrotron radiation

$$\Delta E = e\varepsilon g T \sin \phi_s = eV_{RF} \sin \phi_s$$

g=gap  
T=transit time

(here we assume just one  
cavity)

# Longitudinal dynamics (2)

Now consider small variations in energy and arrival time (phase)

matrix for cavity

$$\begin{pmatrix} \delta \\ ct \end{pmatrix}_2 = \begin{pmatrix} 1 & 0 \\ \frac{heV_{RF} \cos \phi_s}{C2\pi\beta^2 E_0} & 1 \end{pmatrix} \begin{pmatrix} \delta \\ ct \end{pmatrix}_1$$

$$\phi_s = \pi - \arcsin\left(\frac{U_0}{eV_{RF}}\right)$$

synchronous phase is phase needed  
to recover energy lost,  $U_0$

# Longitudinal dynamics (3)

## Momentum compaction

Now, we need to consider the rest of the ring.

Nominal energy  $E_0$

$$\frac{\Delta\tau}{\tau} = \frac{\Delta C}{C} - \frac{\Delta\beta}{\beta} = \eta_s \frac{dE}{E}$$

Change in orbit length vs. energy has two terms:

momentum slip factor

$$\eta_s = \alpha_c - \frac{1}{\gamma^2}$$

$$\alpha_c = \frac{1}{C} \frac{d\Delta C}{d\delta}$$

$$\alpha_c = \frac{1}{\eta_s C} \int \frac{\eta(s)}{\rho(s)} ds$$

$$\alpha_c = 1.78 \times 10^{-4}$$

$$\frac{1}{\gamma^2} = 7.2 \times 10^{-9}$$

for ESRF

momentum compaction factor

transfer matrix is:

$$\eta_s = 0$$

transition energy:

$$\gamma_T = \frac{1}{\sqrt{\alpha_c}} = 75$$

Above transition:  
negative mass!

$$T = \begin{pmatrix} 1 & -C\eta_s \\ 0 & 1 \end{pmatrix}$$

# Synchrotron tune

Multiplying cavity matrix with matrix for rest of ring, we find

$$\begin{pmatrix} \delta \\ ct \end{pmatrix}_2 = \begin{pmatrix} 1 & -C\eta_s \\ \frac{heV_{RF} \cos \phi_s}{C2\pi\beta^2 E_0} & 1 \end{pmatrix} \begin{pmatrix} \delta \\ ct \end{pmatrix}_1$$

One can now analyze this analogously as to the transverse

We derive the synchrotron tune:

$$v_s = \sqrt{\frac{h e V_{RF} \eta \cos \phi_s}{2\pi\beta^2 E_0}}$$

and a longitudinal ‘beta function’

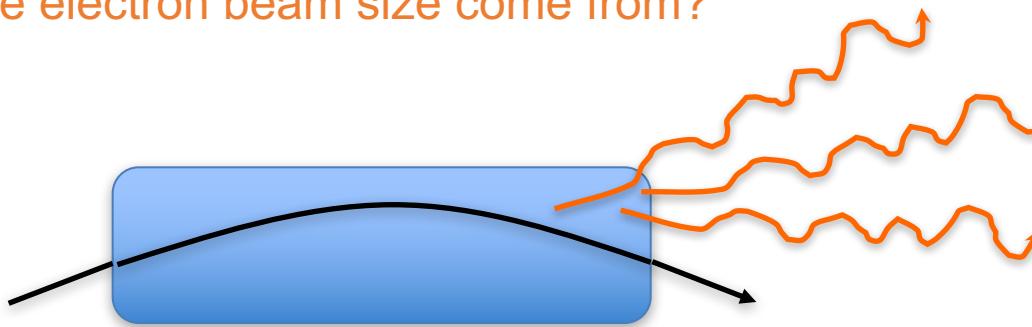
$$\beta_z = \frac{C\alpha_c}{\mu_z}$$

for  $V_{RF}=8\text{MV}$

$v_z = 6 \times 10^{-3}$  or 1 oscillation every 166 turns

Where does the electron beam size come from?

radiation damping (1)



The power radiated by a charged particle undergoing a bend:

$$P_\gamma = \frac{cC_\gamma}{2\pi} \frac{E^4}{\rho^2} \quad C_\gamma = \frac{4\pi}{3} \frac{r_e}{(mc^2)^3} = 8.85 * 10^{-5} \frac{m}{GeV^3}$$

recall:

$$B\rho = \frac{P}{e}$$

Why do electrons radiate so much more than protons?

mass of proton is 1836 times mass of electron

At ESRF, the undulator radiation is typically 15-20% of the radiation emitted by the dipoles. Max 30%. Here, we focus on dipole radiation.

# Radiation from electrons

Add this up for all the dipoles around the ring, and we get

$$U_0 = \int P_\gamma dt = \frac{C_\gamma E^4}{2\pi} \oint \frac{ds}{\rho^2}$$

$$U_0 = 4.88 \text{ MeV}$$

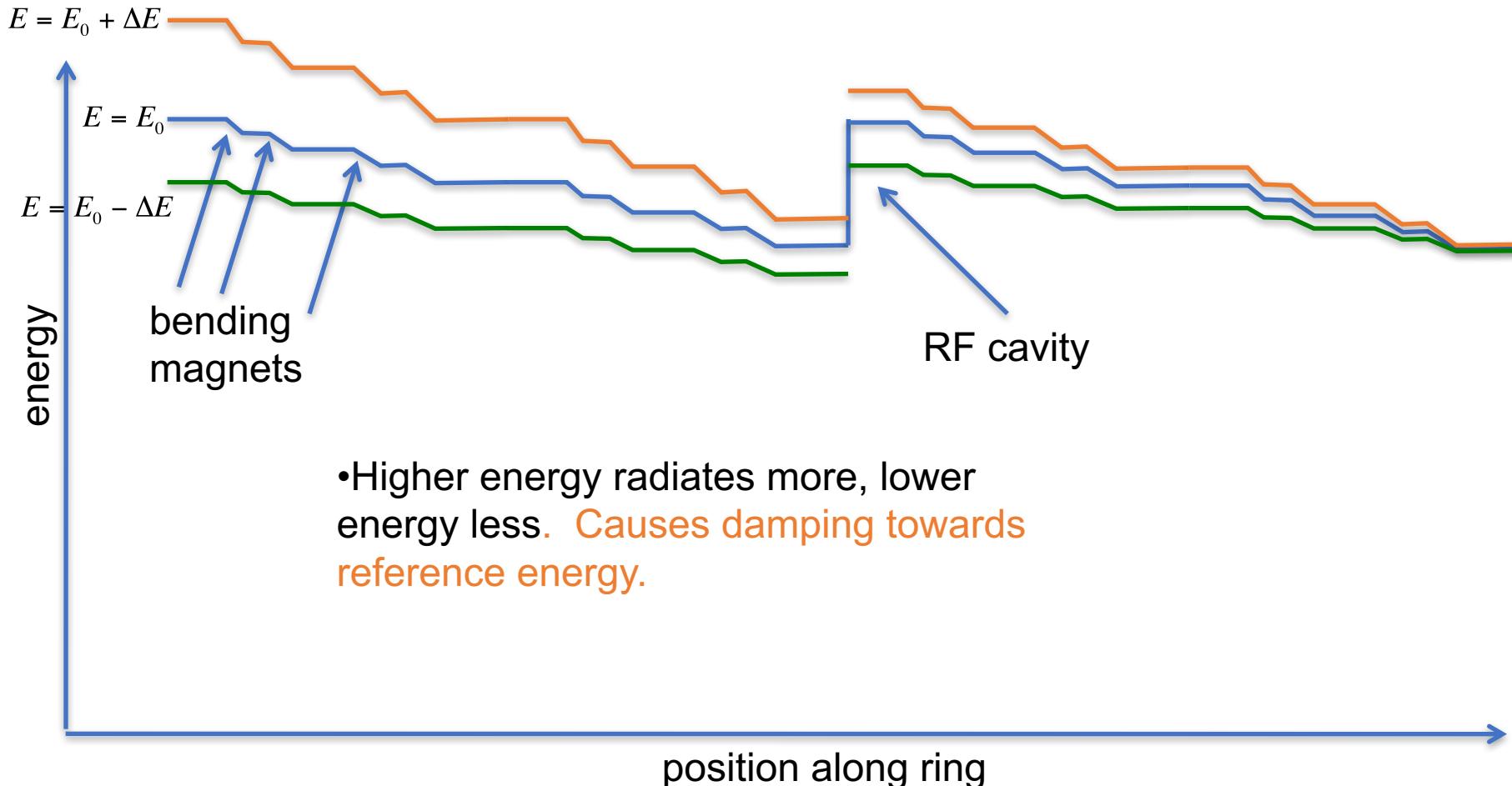
for current ESRF

New lattice will have decreased energy loss.  
More, weaker dipoles with larger bending radius.

$$U_0 = 2.6 \text{ MeV}$$

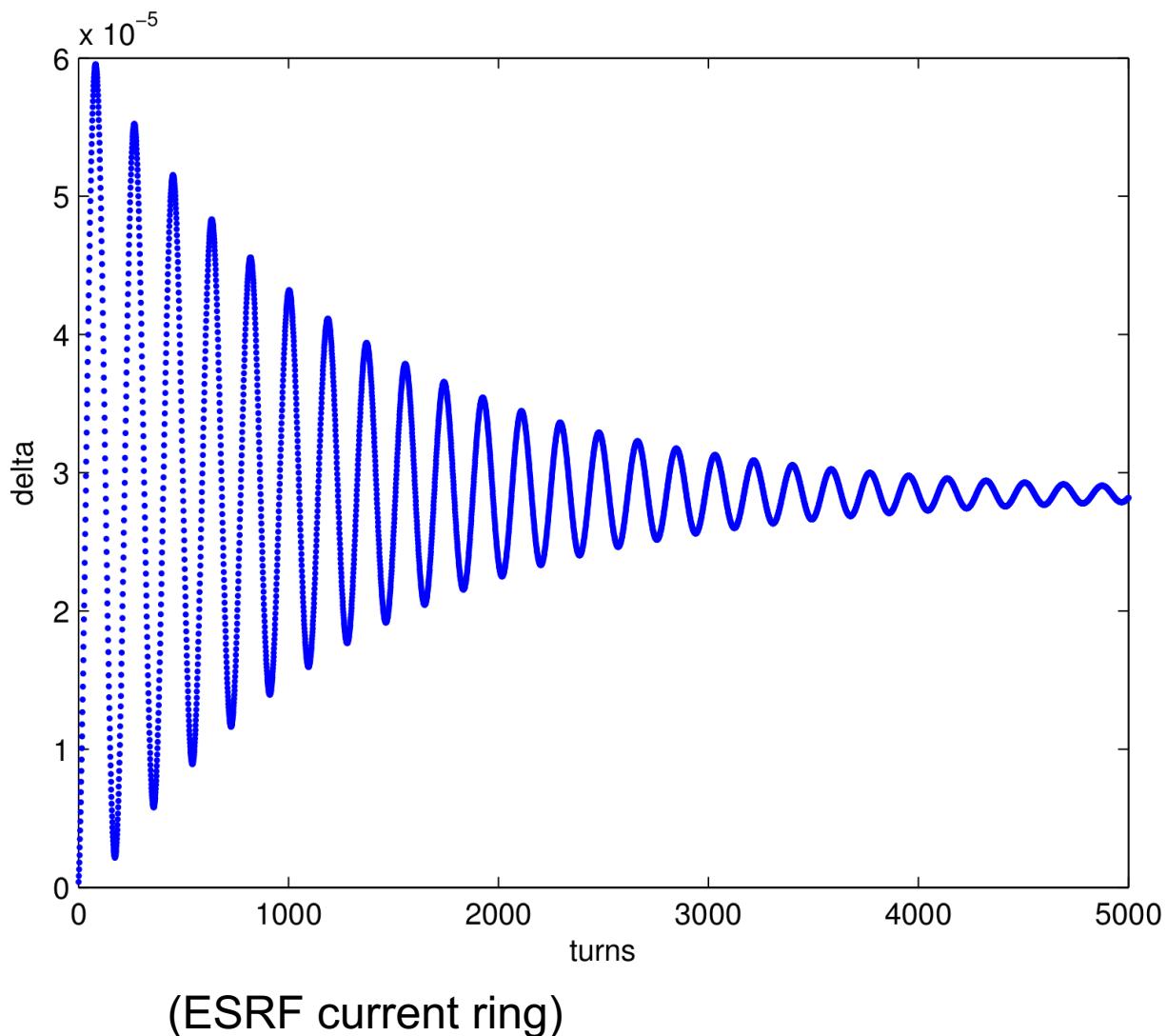
for upgrade ESRF

# Radiation effect on Longitudinal dynamics

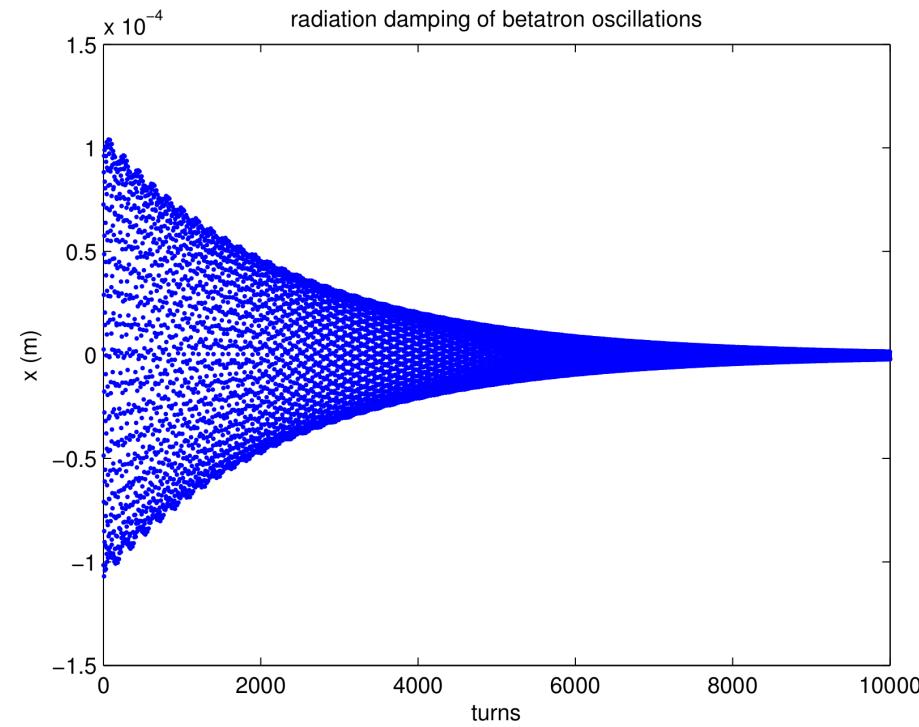
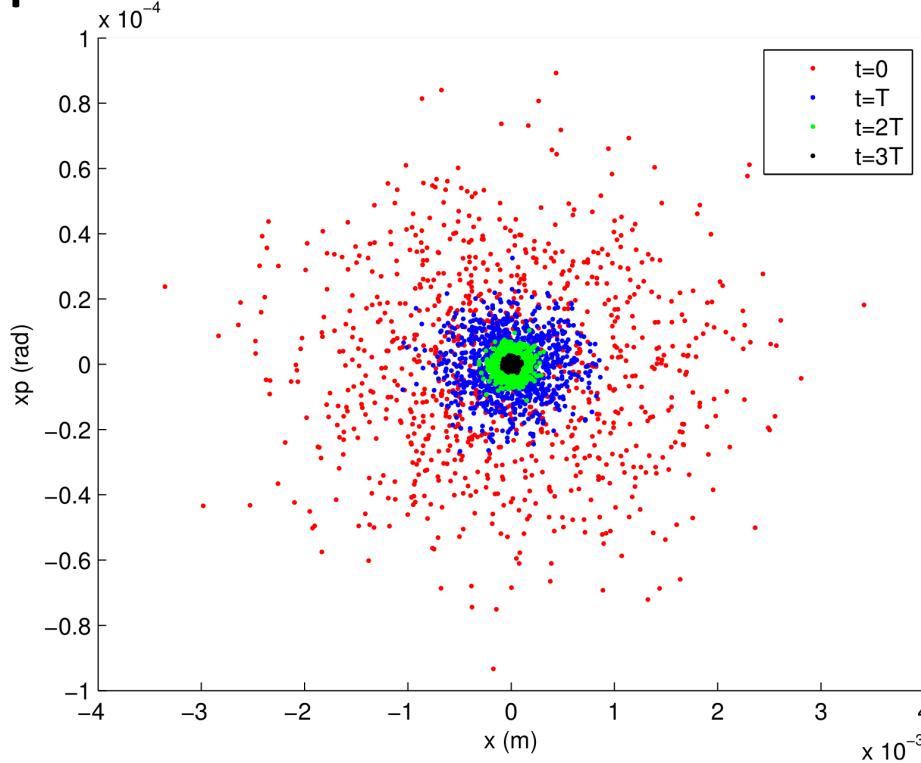


# Radiation damping effect on energy over many turns

$$\delta = \frac{E - E_0}{E_0}$$



# A question arises



What prevents the beam from collapsing to a point?

What is the physics that sets the size of the electron beam??

# Look more closely at the radiation process

(follow Sands, Slac report 121 (1970) )

The radiation power spectrum coming out of a dipole is given by:  
 (first computed by Schwinger (1948) )

$$\wp(\omega) = \frac{P_\gamma}{\omega_c} S\left(\frac{\omega}{\omega_c}\right) \quad \omega_c = \frac{3}{2} \frac{c\gamma^3}{\rho} \quad \begin{array}{l} \text{(critical frequency)} \\ \text{(18.8 KeV for current ESRF )} \end{array}$$

$$S(\xi) = \frac{9\sqrt{3}}{8\pi} \xi \int_{\xi}^{\infty} K_{5/3}(x) dx$$

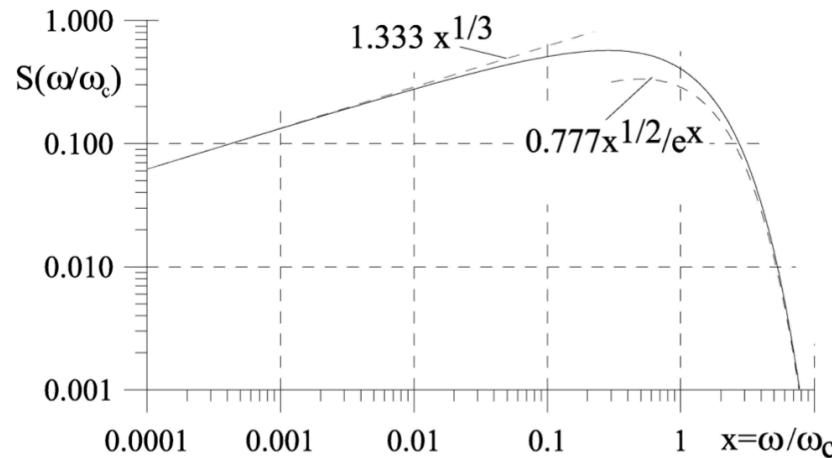


Fig. 22.11. Universal function:  $S(\xi) = \frac{9\sqrt{3}}{8\pi} \xi \int_{\xi}^{\infty} K_{5/3}(x) dx$ , with  $\xi = \omega/\omega_c$

In fact, this radiation will be emitted from the electron as photons

We relate the power spectrum to the distribution of the number of emitted photons per unit time as follows:

$$un(u)du = \wp(u/\hbar)du/\hbar$$

$$u = \hbar\omega$$

The total emission rate is given by:  $\dot{N} = \int_0^\infty n(u)du = \frac{15\sqrt{3}}{8} \frac{P_\gamma}{u_c}$

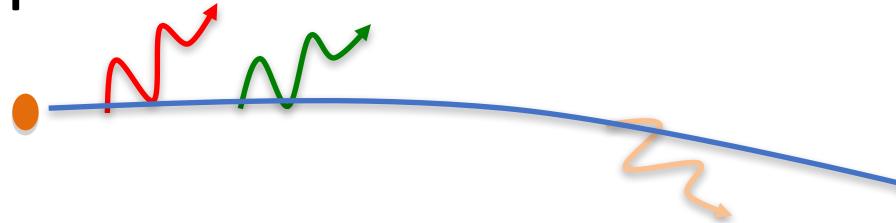
One may compute the average # of photons emitted per radian:

$$= \frac{5}{2\sqrt{3}} \frac{\gamma}{137}$$

For 6.04 GeV electrons, this results in 756 photons per turn.

Or, for ESRF, an average of about 1 photon emitted per meter!  
(approx. 12 photons per dipole)

# Graininess of photon emission



Two sources of randomness:  
emission time of photons are random: Poisson process

Energy emitted is also a random process, with the power spectrum as the probability distribution for each photon.

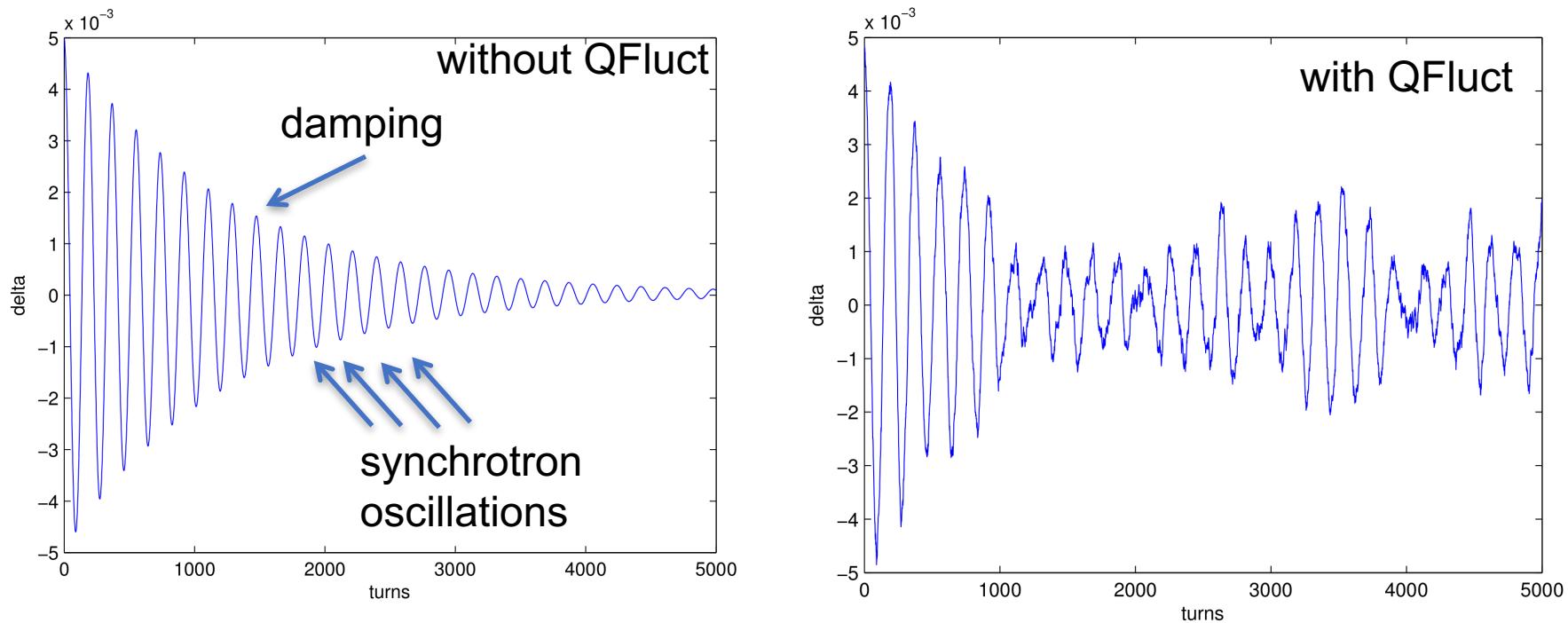
We need to compute the diffusion coefficient for this random walk.

$$d = \mathbb{E} \langle u^2 \rangle = \int_0^\infty u^2 n(u) du$$

The result of this calculation is:

$$d = \frac{55}{24\sqrt{3}} \alpha \left( \frac{\hbar}{mc} \right) \frac{\gamma^5}{|\rho|^3}$$

# Quantum fluctuation effect on longitudinal dynamics



Tracking with and without quantum fluctuations

# Implementation of Monte Carlo method in Zgoubi

- probability of emission of one or more photons, by a particle with rigidity  $B\rho$  (energy  $E$ ), over an integration step  $\Delta s$ , under the effect of  $1/\rho$  curvature,

$$p(k) = \frac{\Lambda^k}{k!} e^{-\Lambda} \quad \text{with} \quad \Lambda = \langle k \rangle = \langle k^2 \rangle \quad (2.1)$$

with  $\Lambda = \frac{5er_0}{2\hbar\sqrt{3}}B\rho \frac{\Delta s}{\rho}$  the average number of photons radiated over  $\Delta s$  ( $r_0 = e^2/4\pi\epsilon_0 m_0 c^2$  is the classical radius of the electron,  $e$  the elementary charge,  $m_0$  the electron rest mass,  $\epsilon_0 = 1/36\pi 10^9$ ,  $\hbar$  is the Plank constant),

- energy  $\epsilon$  of the emitted photon(s), following the probability

$$\mathcal{P}(\epsilon/\epsilon_c) = \frac{3}{5\pi} \int_0^{\epsilon/\epsilon_c} \frac{d\epsilon}{\epsilon_c} \int_{\epsilon/\epsilon_c}^{\infty} K_{5/3}(x) dx \quad (2.2)$$

with  $K_{5/3}$  the modified Bessel function,  $\epsilon_c = 3\hbar\gamma^3 c/2\rho$  the critical energy of the radiation ( $\gamma = E/E_0$  with  $E_0 = m_0 c^2$  the rest energy).

From F. Méot, JINST, 2015,  
“Simulation of radiation damping in rings, using stepwise  
ray-tracing methods”

# Equilibrium Energy spread

$$\frac{\partial f(z, \delta)}{\partial t} = [H, f] - b \frac{\partial f}{\partial \delta} + \frac{d}{2} \frac{\partial^2 f}{\partial \delta^2}$$

Equilibrium is a Gaussian

$$f(z, \delta) = \frac{1}{2\pi\sigma_z\sigma_\delta} e^{-\frac{z^2}{2\sigma_z^2} - \frac{\delta^2}{2\sigma_\delta^2}}$$

bunch length related  
to energy spread

$$\sigma_z = \frac{C\alpha_c}{2\pi\nu_s} \sigma_\delta$$

Calculation gives

$$\sigma_\delta^2 \approx \frac{C_q \gamma^2}{2\rho} \quad C_q = \frac{55}{32\sqrt{3}} \frac{h}{mc} = 3.84 \times 10^{-13} \text{ m}$$

$$(\gamma = 11,800 \quad \rho = 23 \text{ m}) \quad \Rightarrow \sigma_\delta = 1.1 \times 10^{-3} \quad \text{for ESRF}$$

# Energy scaling of ESRF SR dynamics quantities

	Energy loss $U_s$ (MeV/turn)	$\epsilon_{l,\text{eq}}$ (μeV.s)	$\sigma_l$ (mm)	$\tau_l$ (ms)	$\epsilon_{x,\text{eq}}$ (nm)	$\tau_x$ (ms)	$\tau_y$ (ms)
Scaling:	$\gamma^4$	$\gamma^{3/2}$	$1/\gamma^{1/2}$	$1/\gamma^3$	$\gamma^2$	$1/\gamma^3$	$1/\gamma^3$
6 GeV	4.5956 [4.5956]	196 [192]	9.37 [9.309]	1.769 [1.769]	6.90 [6.83]	3.547 [3.546]	3.501 [3.540]
9 GeV	23.263 [23.263]	358 [352]	7.67 [7.601]	0.548 [0.524]	$\frac{15.87}{15.60}$ [15.37]	1.020 [1.051]	1.040 [1.049]
12 GeV	73.518 [73.518]	554 [542]	6.67 [6.582]	0.225 [0.221]	$\frac{28.18}{28.04}$ [27.32]	0.447 [0.443]	0.439 [0.443]
18 GeV	372.16 [372.16]	1022 [996]	5.42 [5.375]	0.068 [0.066]	$\frac{65.77}{63.24}$ [61.46]	0.132 [0.131]	0.130 [0.131]
Theory	$\frac{\mathcal{C}_\gamma}{2\pi} E_s^4 I_2 =$ $\mathcal{C}_\gamma \frac{E_s^4}{\rho}$	$\frac{\alpha E_s}{2\Omega_s} \sigma_{\frac{\Delta E}{E}}^2 =$ $\frac{\alpha E_s}{\Omega_s} \frac{\mathcal{C}_q \gamma^2}{J_l \rho}$	$\frac{\alpha c}{\Omega_s} \sigma_{\frac{\Delta E}{E}}$	$\frac{3T_{\text{rev}}}{2r_0 \gamma^3 (2I_2 + I_4)}$ = $\frac{T_{\text{rev}} E_s}{U_s J_l}$	$\frac{\mathcal{C}_q \gamma^2}{J_x} \frac{I_5}{I_2}$ = $\frac{\mathcal{C}_q \gamma^2}{J_x \rho} \bar{\mathcal{H}}$	$\frac{3T_{\text{rev}}}{2r_0 \gamma^3 (I_2 - I_4)}$ = $\frac{T_{\text{rev}} E_s}{U_s J_x}$	$\frac{3T_{\text{rev}}}{2r_0 \gamma^3 I_2}$ = $\frac{T_{\text{rev}} E_s}{U_s J_y}$

(ref. F. Méot, 2015)

# Conclusions

- Electron storage rings are used for colliders, synchrotron light sources damping rings, and more!
- Transverse and longitudinal symplectic dynamics same as for hadron rings, but synchrotron radiation breaks this
- SR causes both damping and diffusion, resulting in equilibrium emittance
- Zgoubi (and other codes) implements SR with a Monte Carlo model. Both damping and diffusion result from single photon emission statistics!

Thanks for your attention!!