

# Symplectic Tracking in Zgoubi

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# Hamiltonian Equations

The default integrator in Zgoubi is based on directly solving the Lorentz equation:

$$\frac{d(m\mathbf{v})}{dt} = e(\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

The relativistic Hamiltonian:

$$H = \sqrt{c^2(\mathbf{P} - e\mathbf{A})^2 + m^2c^4} + e\Phi$$

Change the independent variable from time  $t$  to displacement  $s$ :

$$H = -\sqrt{P_T^2 - \left(\gamma_0\beta_0\right)^{-2} - P_X^2 - P_Y^2} - \frac{e}{P_0}A_z(X, Y)$$

For the case of  $\Phi = 0$ ,  $A_x = 0$ ,  $A_y = 0$ , and  $X, Y, P_X, P_Y, T, P_T$  are scaled.

For a general magnetic multipole:  $A_z(x, y) = -\text{Re} \sum_{m=1}^{\infty} \frac{1}{m} (b_m + ia_m)(x + iy)^m$

# Hamiltonian Equations

$$H = -\sqrt{P_T^2 - (\gamma_0 \beta_0)^{-2} - P_X^2 - P_Y^2} - \frac{e}{P_0} A_z(X, Y)$$

The Hamiltonian equations:

$$X' = \frac{\partial H}{\partial P_X}, \quad P'_X = -\frac{\partial H}{\partial X}$$

$$Y' = \frac{\partial H}{\partial P_Y}, \quad P'_Y = -\frac{\partial H}{\partial Y}$$

$$T' = \frac{\partial H}{\partial P_T}, \quad P'_T = -\frac{\partial H}{\partial T}$$

Here  $X' = \frac{dX}{ds}$

Multipoles:

$$A_{\text{di}} = -b_1 x$$

$$A_{\text{quad}} = -\frac{1}{2} b_2 (x^2 - y^2)$$

$$A_{\text{sext}} = -\frac{1}{3} b_3 (x^3 - 3xy^2)$$

$$A_{\text{oct}} = -\frac{1}{4} b_4 (x^4 - 6x^2 y^2 + y^4)$$

# Drift

Case (1):  $A_z = 0$

solution:

$$X' = \frac{P_X}{P_S}, \quad P'_X = 0$$

$$Y' = \frac{P_Y}{P_S}, \quad P'_Y = 0$$

$$T' = -\frac{P_T}{P_S}, \quad P'_T = 0$$

$$\begin{pmatrix} X \\ P_X \\ Y \\ P_Y \\ T \\ P_T \end{pmatrix} \rightarrow \begin{pmatrix} X + LP_X / P_S \\ P_X \\ Y + LP_Y / P_S \\ P_Y \\ T - LP_T / P_S \\ P_T \end{pmatrix}$$

$$\text{where } P_S = \sqrt{P_T^2 - (\gamma_0 \beta_0)^{-2} - P_X^2 - P_Y^2}$$

# Quadrupole

Case (2):  $H_{\text{quad}} = -\sqrt{P_T^2 - (\gamma_0 \beta_0)^{-2}} - P_X^2 - P_Y^2 + \frac{eb_2}{2P_0}(X^2 - Y^2)$

$$H_{\text{quad}} = H_1 + H_2$$

$$H_1 = -\sqrt{P_T^2 - (\gamma_0 \beta_0)^{-2}} - P_X^2 - P_Y^2$$

solution: **Drift (D)**

$$H_2 = \frac{eb_2}{2P_0}(X^2 - Y^2) \quad \text{solution: **kick (K)**}$$

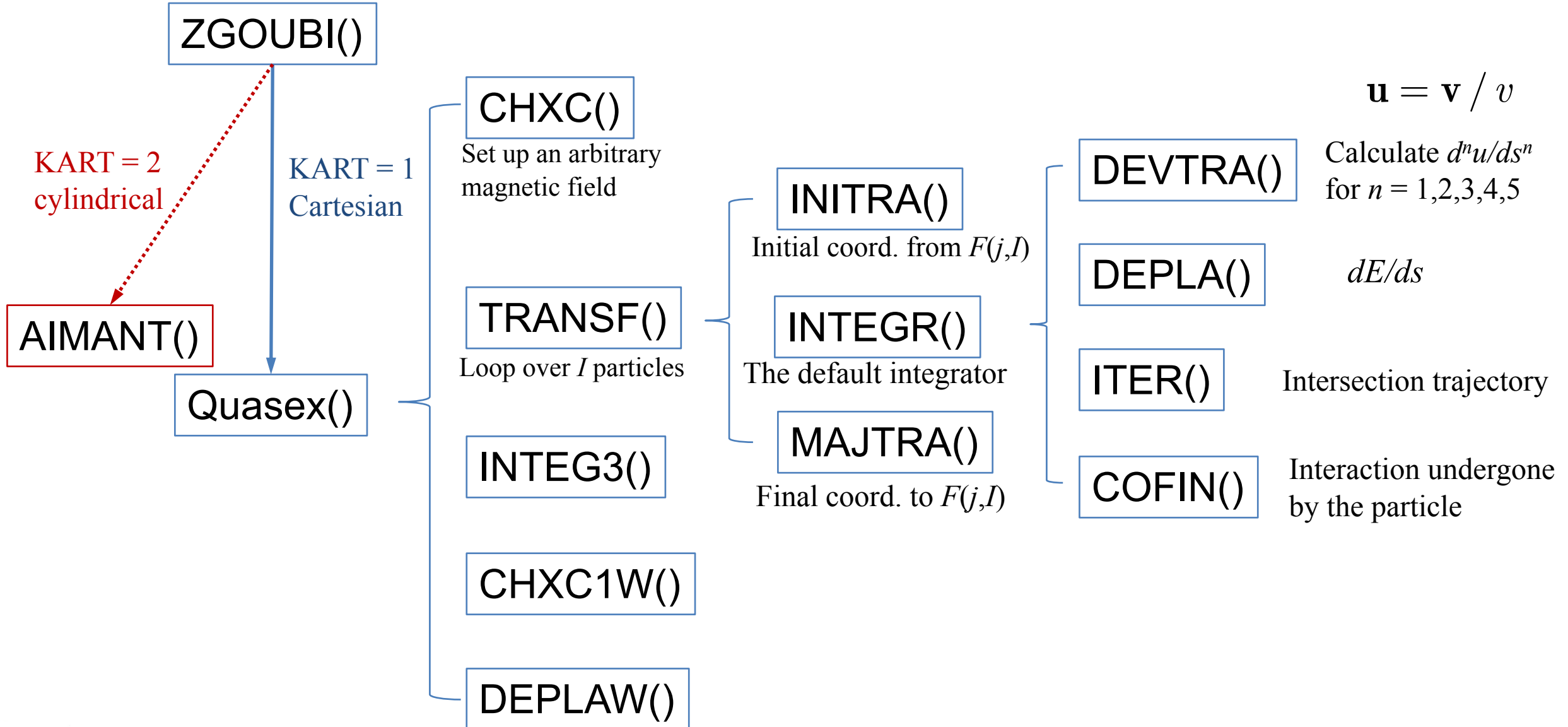
The symplectic drift-kick-drift integrator:

$$Z \rightarrow D(L/2)K_2(L)D(L/2)Z$$

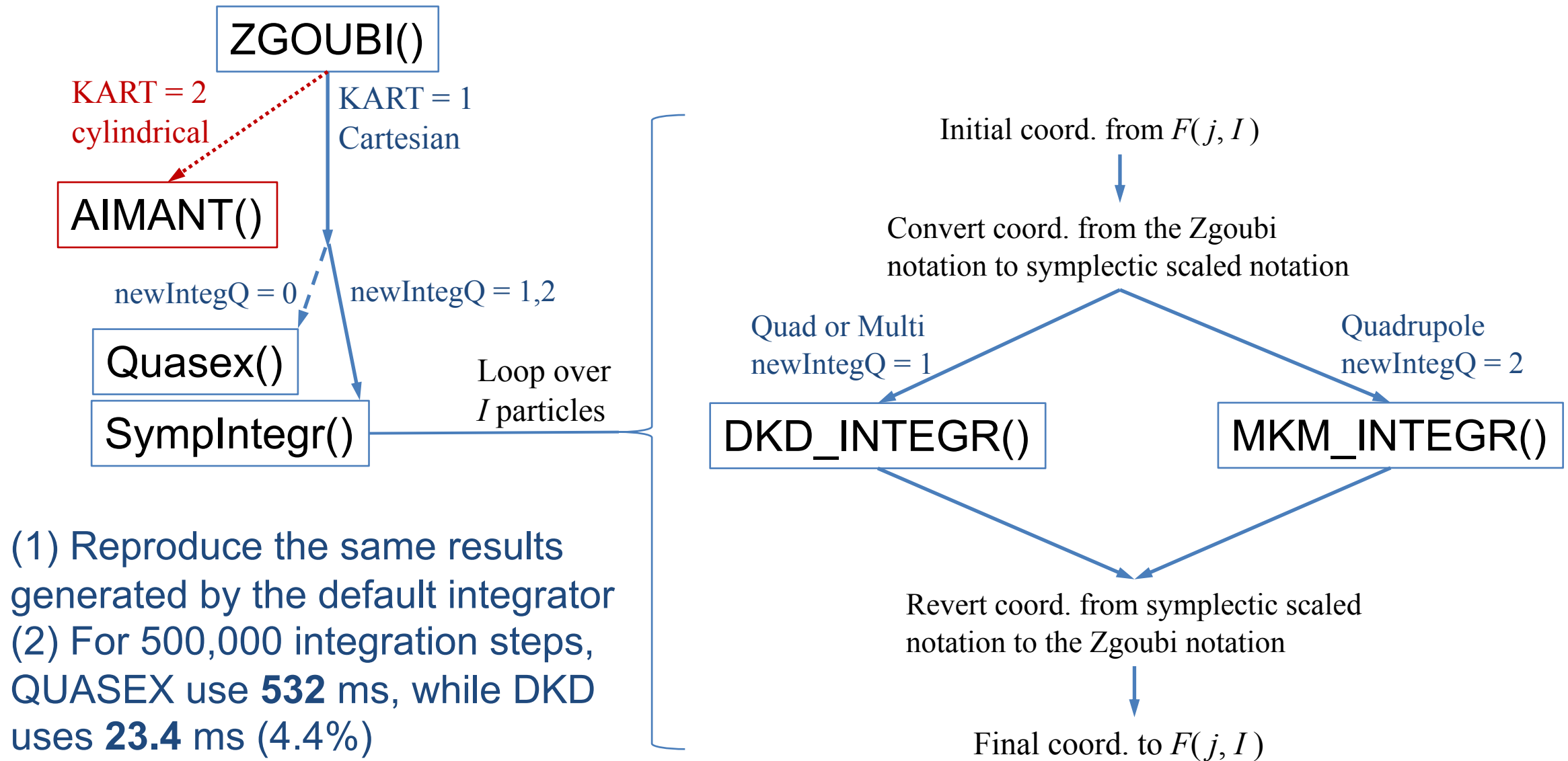
$$\begin{pmatrix} X \\ P_X \\ Y \\ P_Y \\ T \\ P_T \end{pmatrix} \rightarrow \begin{pmatrix} X \\ P_X - \kappa_2 L X \\ Y \\ P_Y + \kappa_2 L Y \\ T \\ P_T \end{pmatrix}$$

where  $\kappa_2 = eb_2 / P_0$

# Default Particle Motion Integrator in Zgoubi



# Symplectic Integrator for Quadrupole and Multipole



# Symplectic Integrator for Quadrupole and Multipole

- 1: the symplectic dkd integrator (any multipole)
- 2: the symplectic mkm integrator (quadrupole only)
- 0 or blank: the default integrator
- Otherwise: end the job

```
'QUADRUPO'                QP  1
0 1
50.0  10.  .763695                .763695
0.  0.
6  .1122  6.2671 -1.4982  3.5882 -2.1209  1.723
0.  0.
6  .1122  6.2671 -1.4982  3.5882 -2.1209  1.723
1.0  ! Quad
1 0.  0.  0.
```



# The Matrix-Kick-Matrix (MKM) Integrator

$$H_{\text{quad}} = H_{q2} + H_{qk}$$

$$H_{qk} = -\sqrt{P_T^2 - (\gamma_0 \beta_0)^{-2} - P_X^2 - P_Y^2} - \frac{1}{2P} (P_X^2 + P_Y^2)$$

$$H_{q2} = \frac{1}{2P} (P_X^2 - P_Y^2) + \frac{eb_2}{2P_0} (X^2 - Y^2)$$

Here  $K = \sqrt{\frac{eb_2}{P_0 P}}$

$$K_Q(L) : \begin{pmatrix} X \\ P_X \\ Y \\ P_Y \\ T \\ P_T \end{pmatrix}$$

$$P_s = \sqrt{P_T^2 - (\beta_0 \gamma_0)^{-2} - P_X^2 - P_Y^2}$$

$$= \sqrt{P^2 - P_X^2 - P_Y^2}$$

$$\rightarrow \begin{pmatrix} X + \left( \frac{1}{P_s} - \frac{1}{P} \right) L P_X \\ P_X \\ Y + \left( \frac{1}{P_s} - \frac{1}{P} \right) L P_Y \\ P_Y \\ T - \left( \frac{1}{P_s} - \frac{P_X^2 + P_Y^2}{2P^3} \right) L P_T \\ P_T \end{pmatrix}$$

# The Matrix-Kick-Matrix (MKM) Integrator

$$H_{q2} = \frac{1}{2P} (P_X^2 - P_Y^2) + \frac{eb_2}{2P_0} (X^2 - Y^2)$$

$$M_Q(L) : \begin{pmatrix} X \\ P_X \\ Y \\ P_Y \\ T \\ P_T \end{pmatrix} \rightarrow$$

$$\left[ \begin{array}{l} X \cos(KL) + P_X \sin(KL) / KP \\ P_X \cos(KL) - KP \cdot X \sin(KL) \\ Y \cosh(KL) + P_Y \sinh(KL) / KP \\ P_Y \cosh(KL) + KP \cdot Y \sinh(KL) \\ T - \frac{P_T}{2P} \left[ \left( \frac{P_X}{P} \right)^2 \frac{2KL + \sin(2KL)}{4K} + \left( \frac{P_Y}{P} \right)^2 \frac{\sinh(2KL) + 2KL}{4K} \right. \\ \left. + (KX)^2 \frac{2KL - \sin(2KL)}{4K} + (KY)^2 \frac{\sinh(2KL) - 2KL}{4K} \right. \\ \left. + KX \frac{P_X}{P} \frac{\cos(2KL) - 1}{2K} + KY \frac{P_Y}{P} \frac{\cosh(2KL) - 1}{2K} \right] \\ P_T \end{array} \right]$$

$$Z \rightarrow M_Q(L/2) K_Q(L) M_Q(L/2) Z$$

# The DKD Integrator for a multipole

$$Z \rightarrow D(L/2)K(L)D(L/2)Z$$

$$D(L) : \begin{pmatrix} X \\ P_X \\ Y \\ P_Y \\ T \\ P_T \end{pmatrix} \rightarrow \begin{pmatrix} X + LP_X / P_S \\ P_X \\ Y + LP_Y / P_S \\ P_Y \\ T - LP_T / P_S \\ P_T \end{pmatrix} \quad K(L) : \begin{pmatrix} X \\ P_X \\ Y \\ P_Y \\ T \\ P_T \end{pmatrix} \rightarrow \begin{pmatrix} X \\ P_X - L \cdot \frac{\partial}{\partial_X} H_{\text{kick}}(X, Y) \\ Y \\ P_Y - L \cdot \frac{\partial}{\partial_Y} H_{\text{kick}}(X, Y) \\ T \\ P_T \end{pmatrix}$$

$$H_{\text{kick}} = \frac{e}{P_0} \text{Re} \sum_{m=1}^{\infty} \frac{1}{m} (b_m + ia_m) (X + iY)^m$$

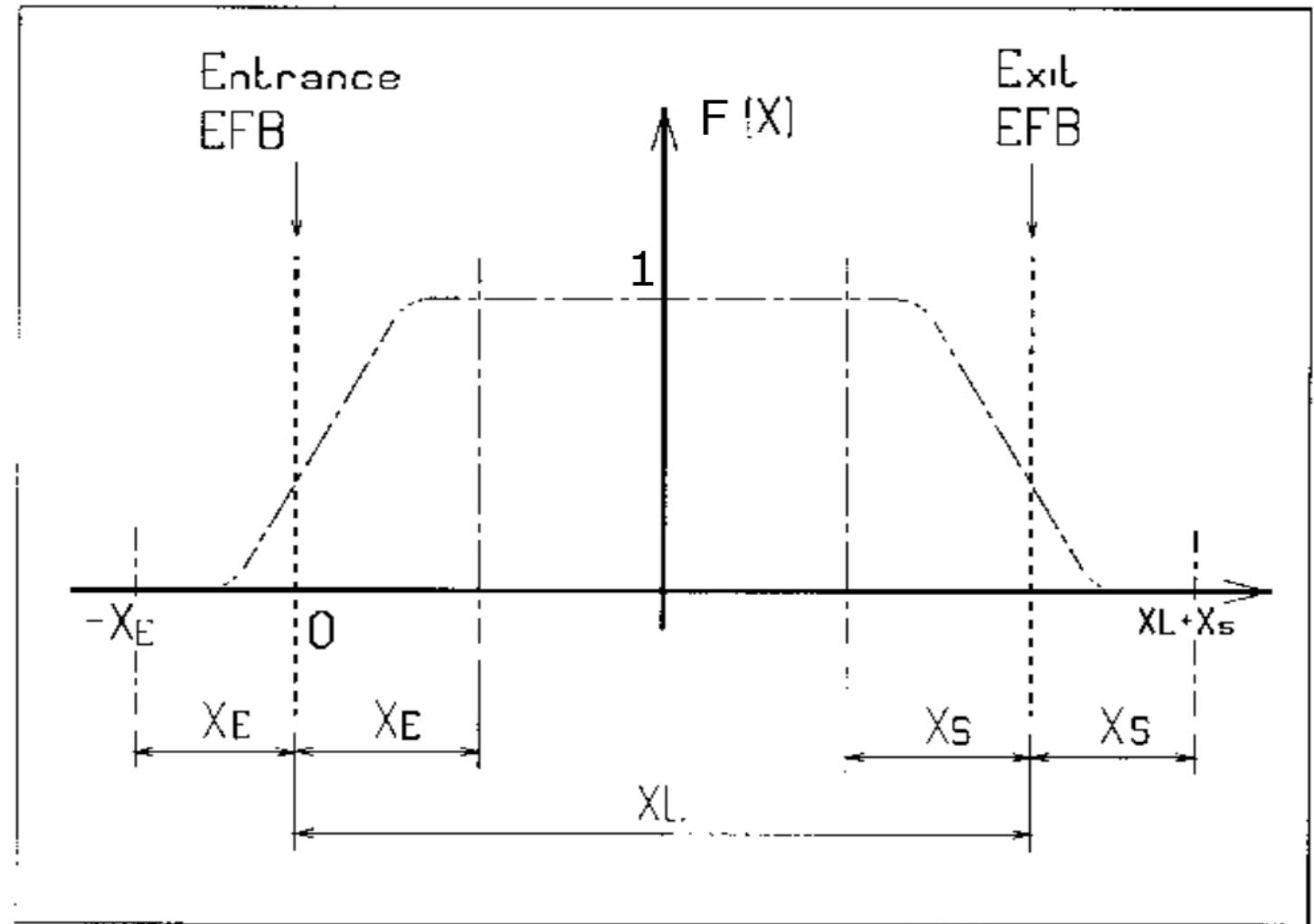
# Fringe Fields

For the case of quadrupole:

$$b_2(Z) = F(Z)b_2$$

$$F(Z) = \frac{1}{1 + e^{P[d(Z)]}}$$

$$P(d) = \sum_{i=0}^5 C_i \left( d / \lambda \right)^i$$



# Fringe Fields for Quadrupole

$$A_z = -\frac{b_2}{2P_0}(X^2 - Y^2)F(Z)$$

However, this violates the Hamiltonian equation

Solutions satisfying the Hamiltonian equation:

$$A_z = -\frac{b_2}{2P_0}(X^2 - Y^2) \left[ F(Z) - \frac{1}{12} \frac{d^2 F(Z)}{dZ^2} (X^2 + Y^2) + \dots \right]$$

$$A_\rho = \frac{b_2}{4P_0} \left[ \frac{dF(Z)}{dZ} \rho^3 + \dots \right] e^{i2\varphi} \quad \text{Here} \quad \begin{aligned} x &= \rho \cos \varphi, & y &= \rho \sin \varphi \\ A_x &= A_\rho \cos \varphi, & A_y &= A_\rho \sin \varphi \end{aligned}$$

We have implemented the symplectic integrator for quadrupole with fringe fields

# Summary and Future Work

- Implemented the symplectic integrators for particle tracking
  - drift-kick-drift for an arbitrary multipole
  - drift-kick-drift and matrix-kick-matrix for quadrupole
  - fringe fields are treated for quadrupole (working on a general multipole)
- More efficient than the default integrator
  - Using the same step size, only 4.4% of CPU time
  - Reaching convergence with larger step size
- Future work
  - Motion in bend (rectangular and sector), etc.
  - Electric field ( $\Phi \neq 0$ ), spin, etc.
  - Discrete Euler-Lagrange equations