

# $\nu$ -Dark matter models

## Collider phenomenology



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### Focus on

arXiv: arXiv:1308.3655 (JHEP), arXiv:1504.07892 (PRD), arXiv:1509.06313, arXiv:1511.01873, arXiv:1511.nnnnn. In collaboration with

G. Palacio, F. von der Pahlen, D. Portillo, A. Rivera, M. Sánchez, O. Zapata (UdeA)

C. Arbeláez (USM), W. Tangarife (Tel Aviv U.), C. Yaguna (Heidelberg, Max Planck Inst.).

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2. Proposal
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4. Recast at the LHC Run-I
5. Lepton flavor dependence
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7. Compressed spectra
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# General framework

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If **neutrino masses** arise radiatively it may originate from new physics at the TeV scale in join with **dark matter** (DM)

It may be, though, that they are related to each other.

In this direction, models with one-loop radiative neutrino masses and viable dark matter candidates have now a complete classification given in

R.D., Yaguna, C, Zapata, O, arXiv:1308.3655 (JHEP)

There, the **new fields are odd under a  $Z_2$**  symmetry which ensures the stability of the DM particle, while the SM particles are even.

### FeynRules/SARAH



```
graph TD; A[FeynRules/SARAH] --> B[micrOMEGAS (CalcHEP)]; A --> C[MadGraph(UFO)]; B --> D[DM]; C --> E[LHC]
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**micrOMEGAS (CalcHEP)**

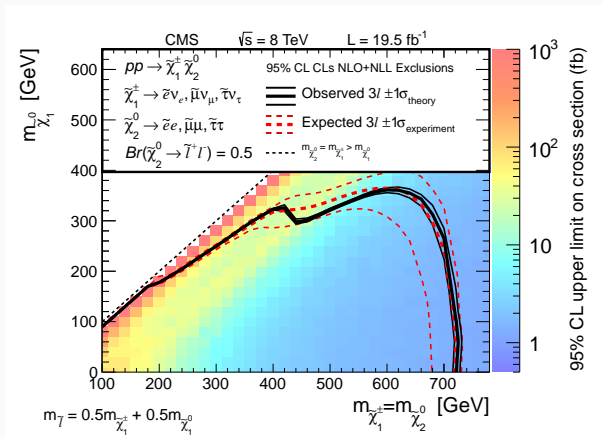
**DM**

**MadGraph(UFO)**

**LHC**

- We have only qualitatively described the particle content and the dark matter candidates of each model. A more specific analysis of some of these models is certainly desirable.
- The collider and dark matter phenomenology of many of these viable models have yet to be studied in detail.
- Some strategies to systematically search for this kind of models at colliders would be designed.
- New particles allowed to be even under  $Z_2$  could give rise to new possibilities.

## Largest Wino-production to three leptons and MET



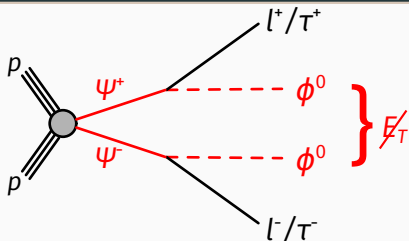
Interpretation of the results of the three-lepton search in the flavor-democratic signal model with slepton mass parameter  $\xi_{\tilde{l}} = 0.5$

Simplified SUSY to dileptons and MET: slepton pair production  
(see below)

# Proposal

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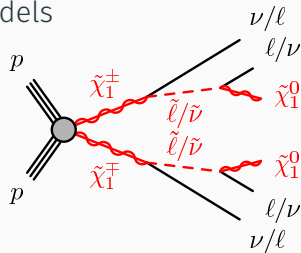
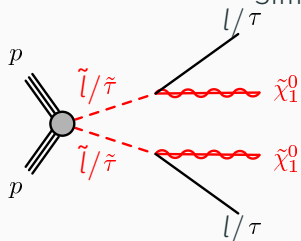
# Dilepton plus transverse missing energy signal



SU(2)<sub>L</sub> assignments:

$\Psi = 1, 2, 3$ ,  $\Phi = 1, 2$ , with  $m_{DM} \sim m_h/2$ .

Simplified SUSY models



Smaller cross sections.

Intermediate states and smaller lepton  $p_T$



## New simplified models

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# Not-susy $SO(10) \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y \times Z_2$

Standard Model:  $Z_2$ -even

Fermions:  $16_F$

Scalars:  $10_H, 45_H \dots$

New  $Z_2$ -odd particles

$10_F, 45_F, \dots$

$16_H, \dots$

Lightest Odd Particle (LOP) may be a suitable dark matter candidate.

Example (arXiv:1509.06313)

Split-SUSY-like spectrum: bino-higgsino-wino

+

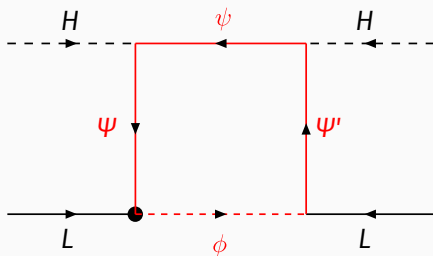
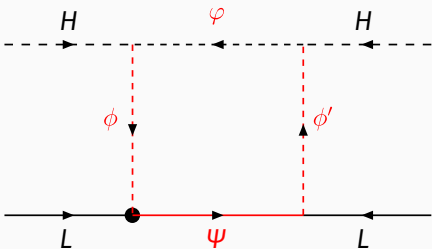
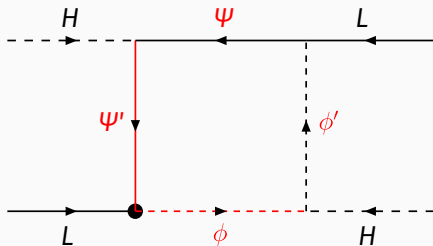
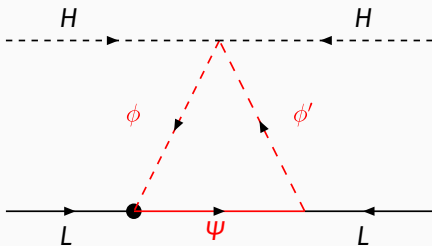
↓

$10'_H$  with fermion DM or,

$16_H, \dots$  with scalar DM

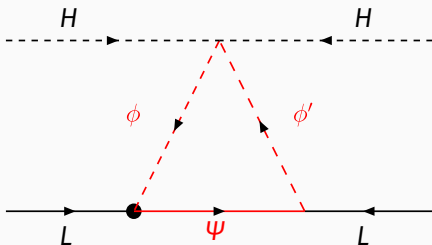
# Weinberg operator at one-loop

( $Z_2$ -odd fields)

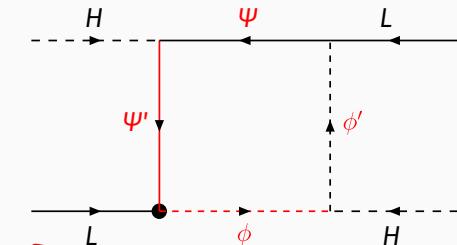
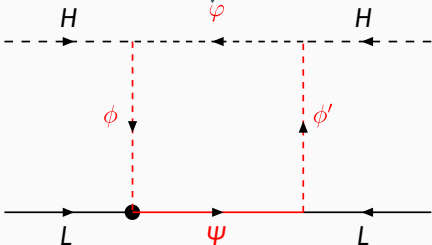


# Weinberg operator at one-loop

( $Z_2$ -odd fields)

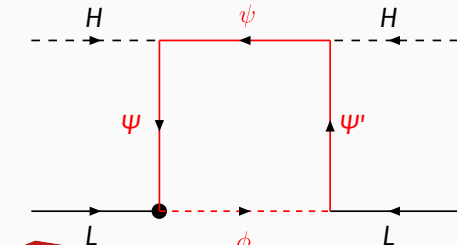


Wino-like scotogenic model



**New**

Higgsino-like Zee model



**New**

Higgsino-like scotogenic model

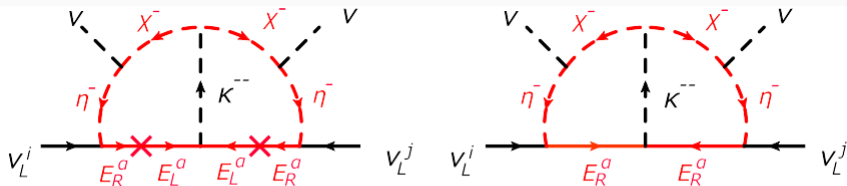
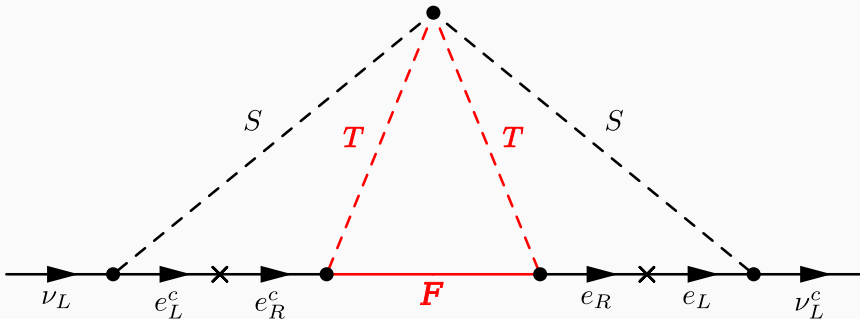


FIG. 2: Feynman diagrams for the neutrino mass generation at the two-loop level. Particles indicated by the red color have the  $Z_2$  odd parity.



Three-loop diagram for neutrino mass. Here,  $S \sim (1, 1, 2)$  and  $T \sim (1, 2n + 1, 2)$  are beyond-SM scalars while  $F \sim (1, 2n + 1, 0)$  is a beyond-SM fermion.

## Recast at the LHC Run-I

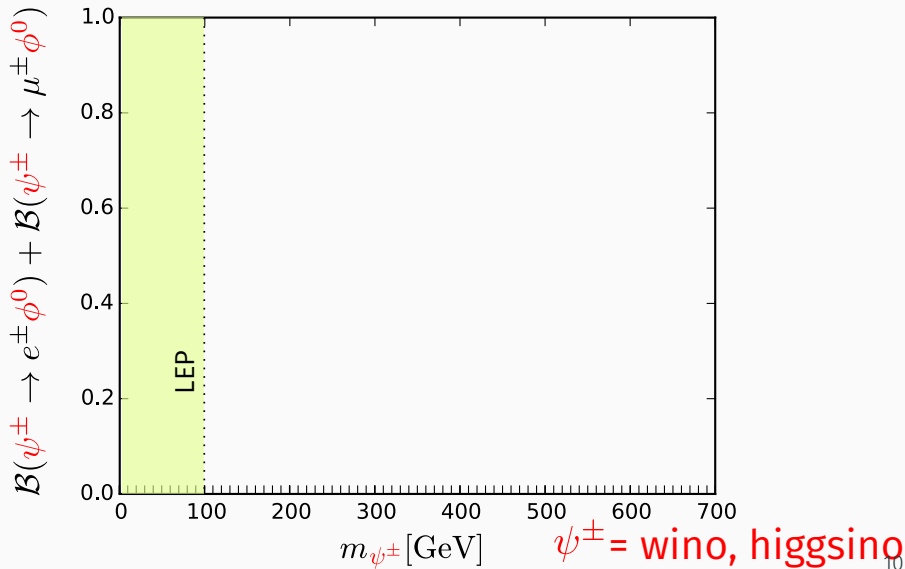
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$$pp \rightarrow \psi^\pm \psi^\mp \rightarrow l_i^\pm l_j^\mp \phi^0 \phi^0 \quad l_i = e, \mu, \tau, \quad \phi^0 = H^0, S^0$$

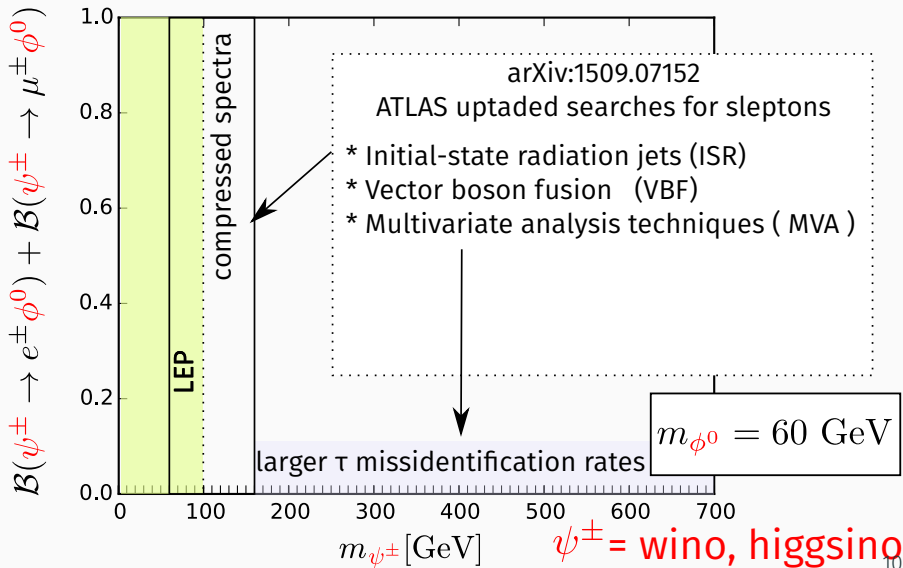
$$\psi^\pm = \text{wino, higgsino}$$

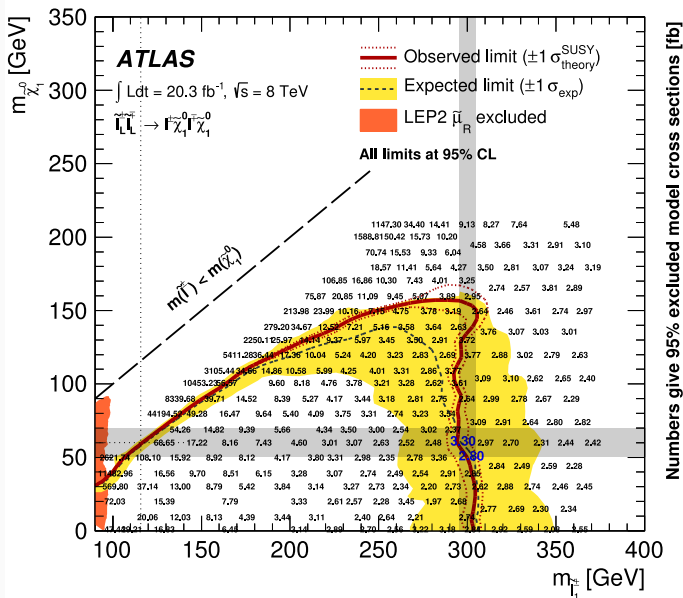


$$pp \rightarrow \psi^\pm \psi^\mp \rightarrow l_i^\pm l_j^\mp \phi^0 \phi^0 \quad l_i = e, \mu, \tau, \quad \phi^0 = H^0, S^0$$



$$pp \rightarrow \psi^\pm \psi^\mp \rightarrow l_i^\pm l_j^\mp \phi^0 \phi^0 \quad l_i = e, \mu, \tau, \quad \phi^0 = H^0, S^0$$

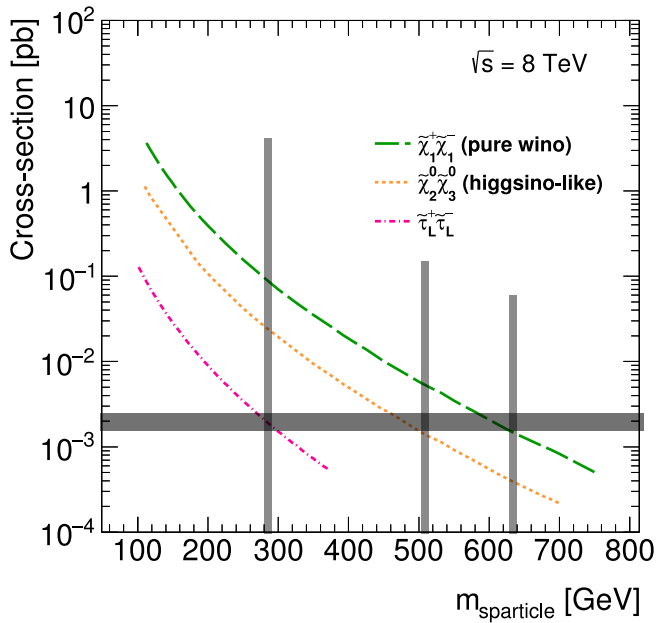




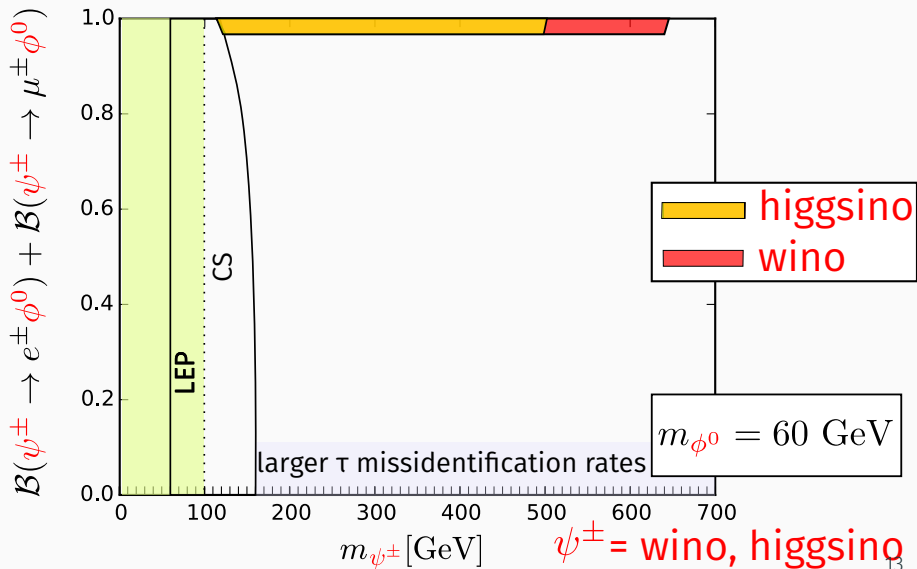
CMS

 $\gtrsim 260 \text{ GeV}$ arXiv:1405.7570<sup>11</sup>

# Excluded cross section for higgsino and wino fermions



$$pp \rightarrow \psi^\pm \psi^\mp \rightarrow l_i^\pm l_j^\mp \phi^0 \phi^0 \quad l_i = e, \mu, \tau, \quad \phi^0 = H^0, S^0$$

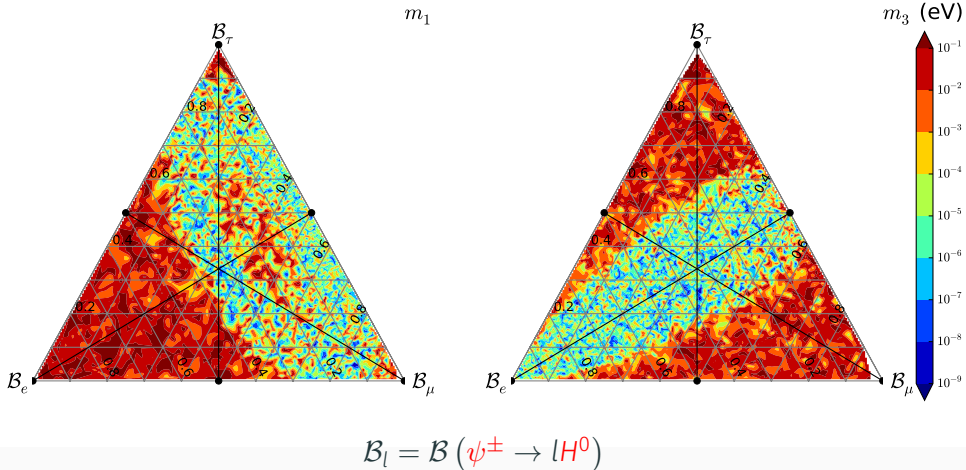


# Lepton flavor dependence

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# Casas-Ibarra parametrization

In wino-like scotogenic model (may be in general)



# Exploration of flavor space

Wino-like escotogenic model: Recast for  $B_\mu + B_e \gtrsim 0.1$  and

$$m_{H^0} < m_{\psi^\pm} = m_{\psi^0} < m_{A^0}, m_{H^\pm}$$

Start with Signal regions as in ATLAS-arXiv:1403.5294 for  $\cancel{E}_T$  with  $e^+e^-$ ,  $\mu^+\mu^-$ ,  $e^\pm\mu^\mp$ .

FeynRules



micrOMEGAS (Experimental and theoretical constraints)



MadGraph



Pythia 6 (hep format)

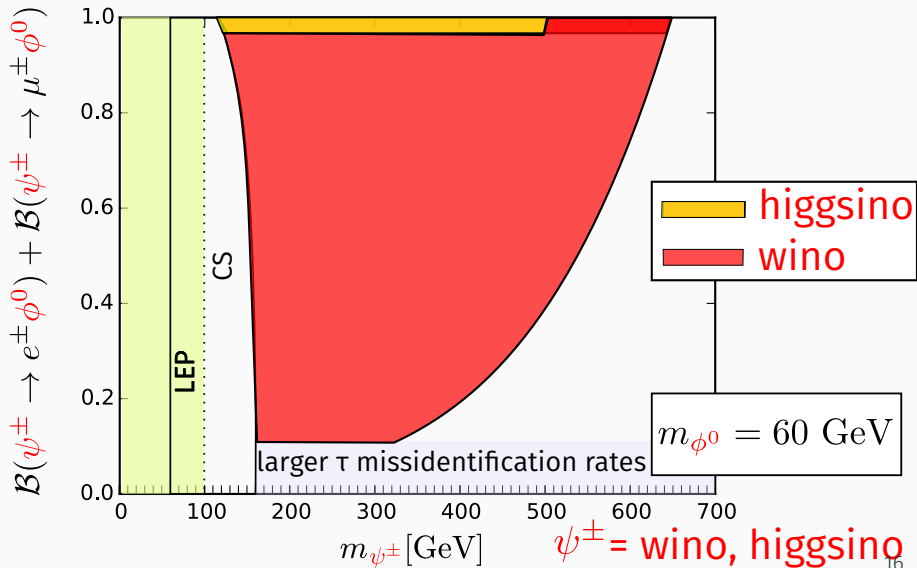


checkMATE (CL-calculation)

$$CL = CL^{ee} CL^{\mu\mu} CL^{\mu e}$$

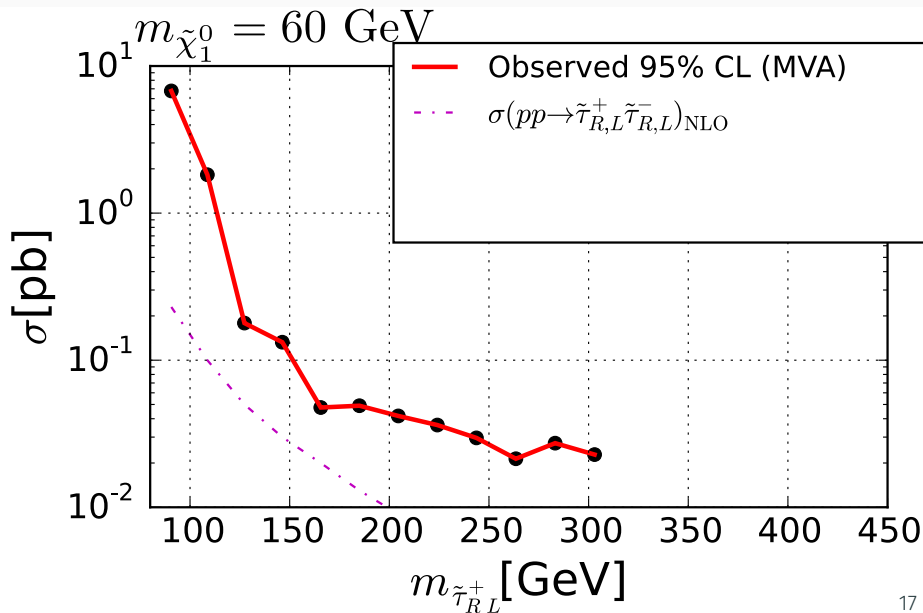


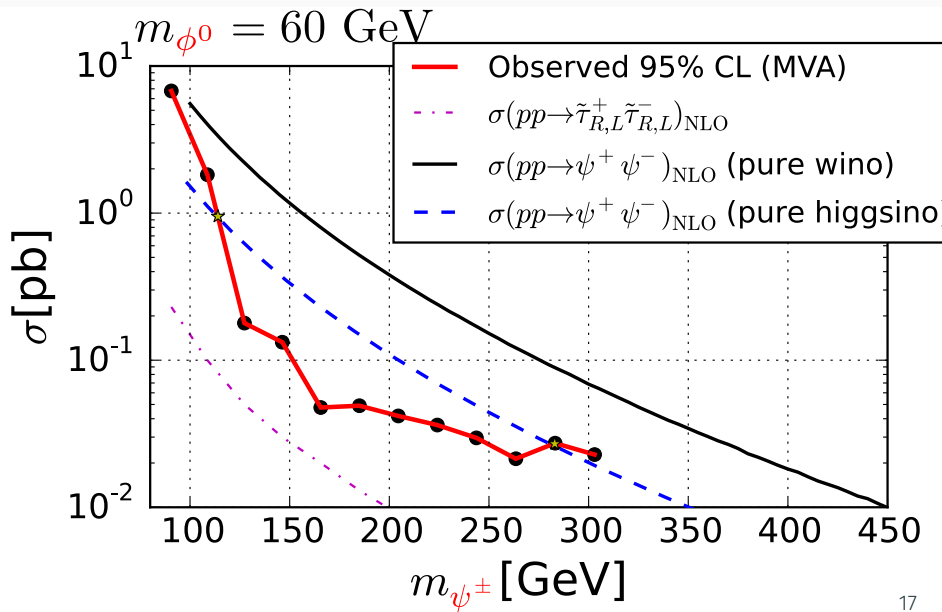
$$pp \rightarrow \psi^\pm \psi^\mp \rightarrow l_i^\pm l_j^\mp \phi^0 \phi^0 \quad l_i = e, \mu, \tau, \quad \phi^0 = H^0, S^0$$



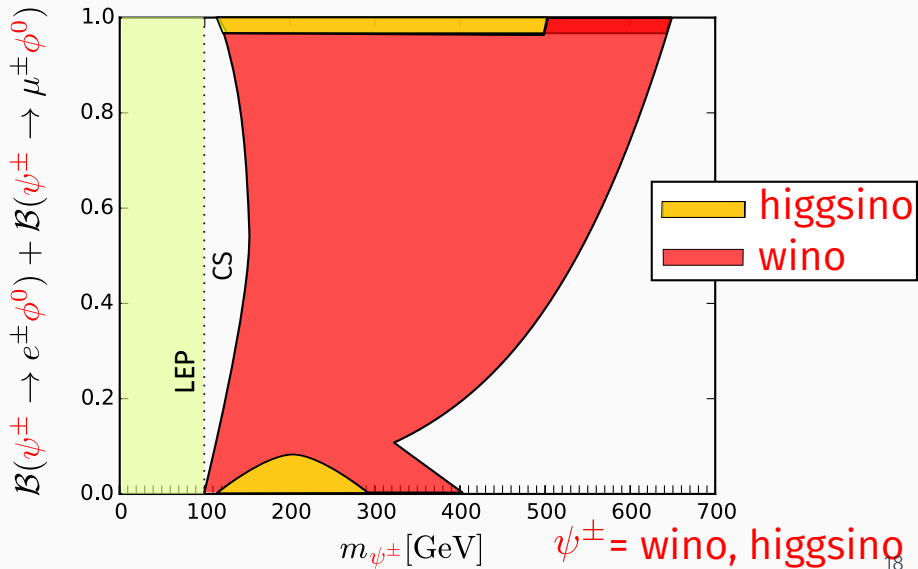
Pure  $\tau$  case

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$$pp \rightarrow \psi^\pm \psi^\mp \rightarrow l_i^\pm l_j^\mp \phi^0 \phi^0 \quad l_i = e, \mu, \tau, \quad \phi^0 = H^0, S^0$$

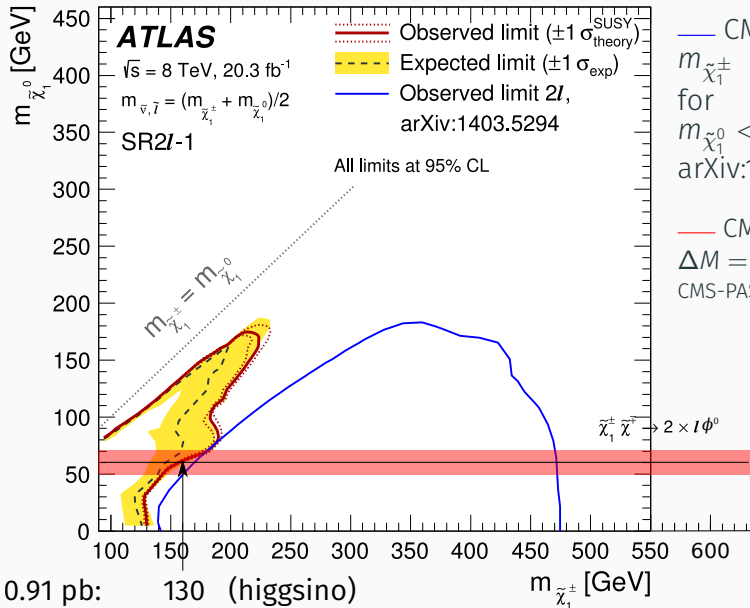


# Compressed spectra

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# wino production to dilepton plus missing energy (ISR)

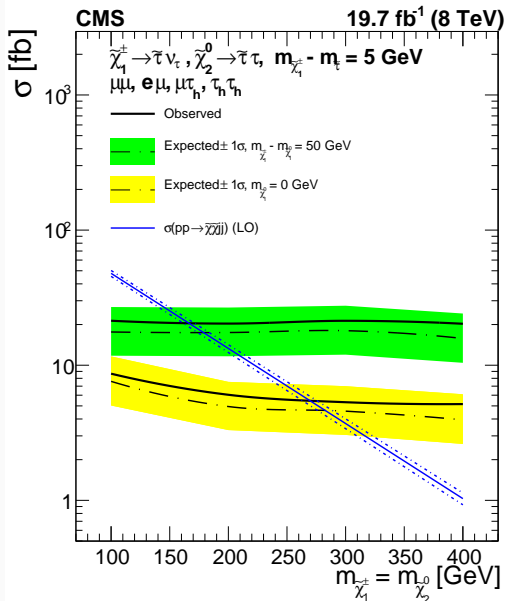
Simplified Model:  $\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp \rightarrow 2 \times \tilde{l} \nu (\tilde{\nu} l) \rightarrow 2 \times l \nu \tilde{\chi}_1^0$



CMS  
 $m_{\tilde{\chi}_1^\pm} \gtrsim 510 \text{ GeV}$   
 for  
 $m_{\tilde{\chi}_1^0} < 100 \text{ GeV}$   
 arXiv:1405.7570

CMS  $\tilde{\chi}_1^\pm \tilde{\chi}_1^0$   
 $\Delta M = 20 \text{ GeV (ISR)}$   
 CMS-PAS-SUS-14-021

# CMS 1508.07628: di-tau plus missing energy (VBF)



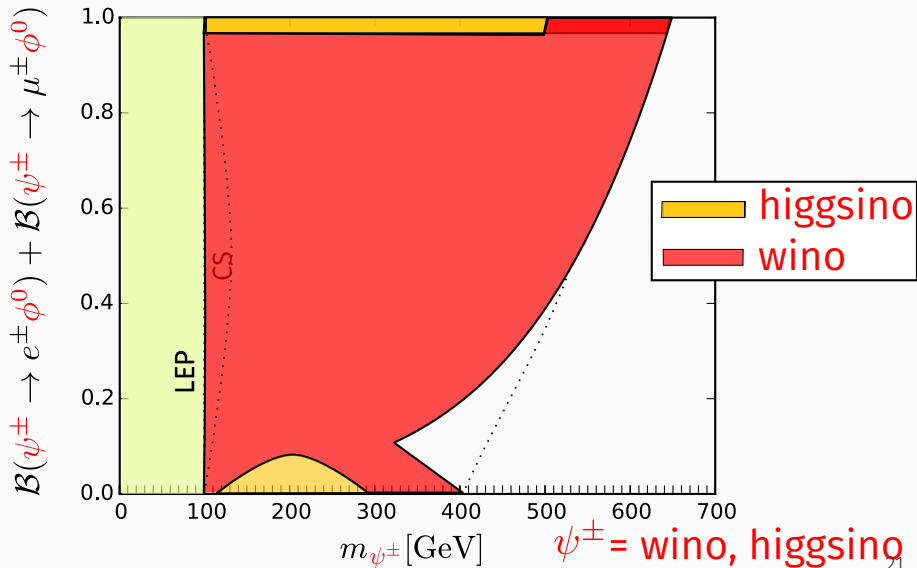
~~$\psi^\pm \psi^0$~~

$\psi^\pm \psi^\mp$

$\psi^\pm \psi^\pm$



$$pp \rightarrow \psi^\pm \psi^\mp \rightarrow l_i^\pm l_j^\mp \phi^0 \phi^0 \quad l_i = e, \mu, \tau, \quad \phi^0 = H^0, S^0$$



# Conclusions

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The wino-like scotogenic model (*radiative type-III seesaw*) with  $\mathcal{B}(\psi^+ \rightarrow l_i \phi^0) = 1$  is a good simplified model for the interpretation of experimental results at the LHC.

At Run-I, the charged fermion have been **excluded** until 630 GeV (400 GeV) for  $\mathcal{B}(\psi^+ \rightarrow e^+ \phi^0) = 1$  ( $\mathcal{B}(\psi^+ \rightarrow \tau^+ \phi^0) = 1$ ),

Including the compressed-spectra region

*TODO: Higgsino-like case*

Thanks!