



UNIVERSIDAD DE ANTIOQUIA
1805

Dark matter

Before and after reheating

Diego Restrepo

July 30, 2018

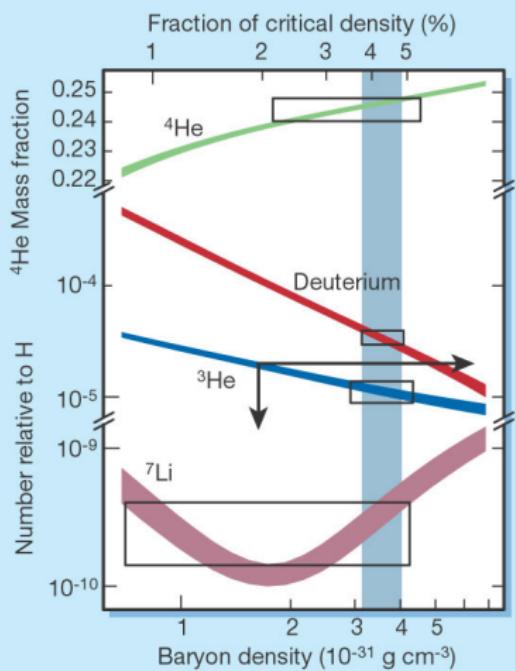
Instituto de Física
Universidad de Antioquia
Phenomenology Group
[http:
/gfif.udea.edu.co](http://gfif.udea.edu.co)

In collaboration with
Valentina Montoya

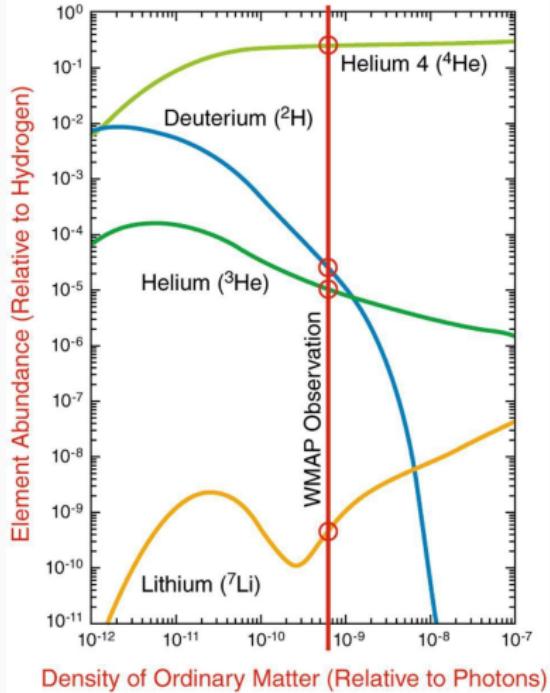


Gravitational evidence of dark matter

Big Bang Nucleosynthesis



Credits: Chris Mihos ASTR/PHYS 328/428



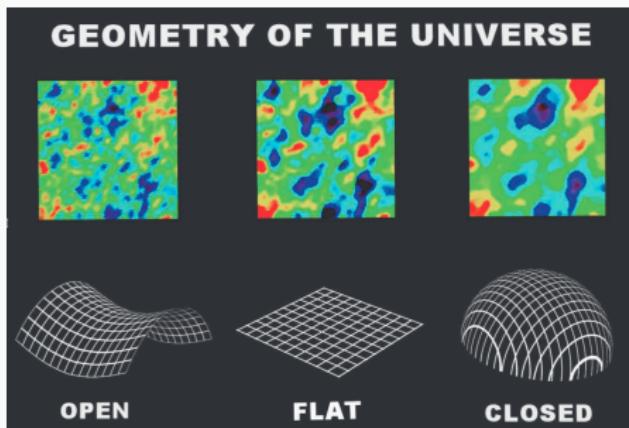
NASA WMAP Science Team
WMAP101087

Element Abundance graphic: Steigman, Encyclopedia of Astronomy and Astrophysics [Institute of Physics] December, 2000

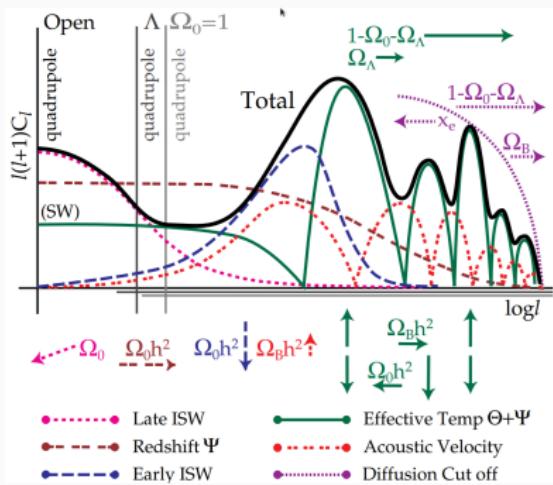
Credits: NASA/WMAP Science TEAM

Cosmic microwave background

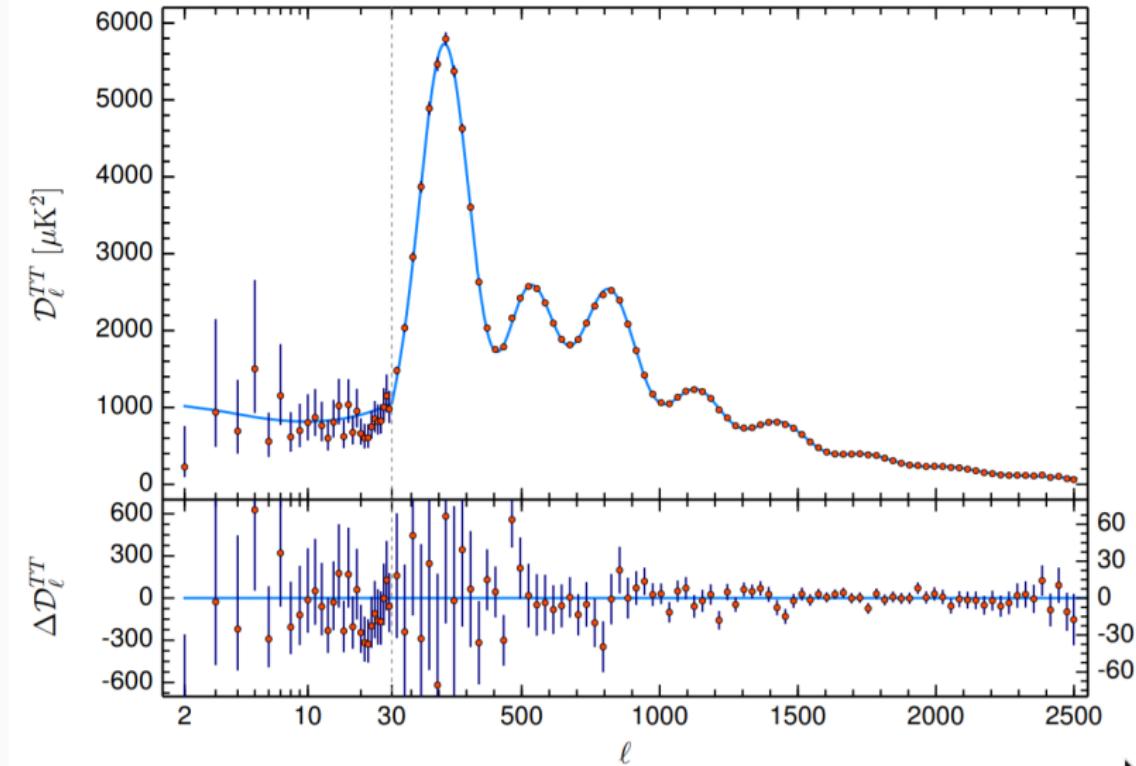
You are seeing the very photons that were emitted when the Universe was just 380,000 years old, and they display a spectrum indicative of a very precise temperature.



Credit: E. SIEGEL / BEYOND THE GALAXY



Credit: Hu, Sugiyama, Silk, astro-ph/9504057



Credit: Planck 2018

Why was the temperature of the CMB the same in all directions?

What was the origin of the small temperature fluctuations?

Cosmic Miso Soup

- When matter and radiation were hotter than 3000 K, matter was completely ionised. The Universe was filled with plasma, which behaves just like a soup
- Think about a Miso soup (if you know what it is). Imagine throwing Tofus into a Miso soup, while changing the density of Miso
- And imagine watching how ripples are created and propagate throughout the soup

Credit: Komatsu, ICTP Summer School on Cosmology 2018¹

¹Video available

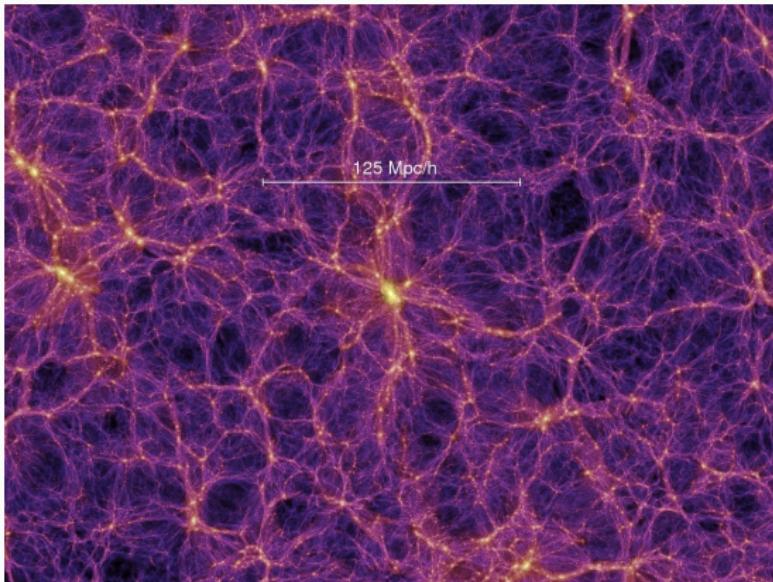


Credit: Komatsu, ICTP Summer School on Cosmology 2018¹

¹Video available

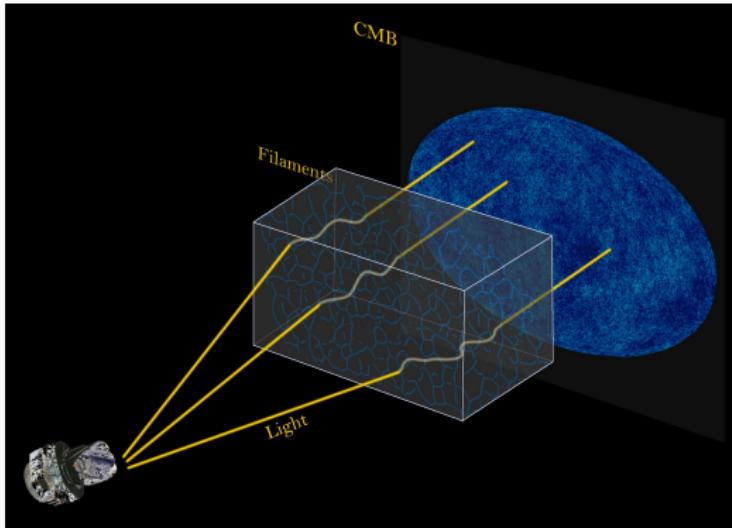
Cooking the soup: Cosmic web

Dark matter in the universe evolves through gravity to form a complex network of halos, filaments, sheets and voids, that is known as the cosmic web [arXiv:1801.09070]



Milennium simulation: <https://wwwmpa.mpa-garching.mpg.de/galform/virgo/millennium/>

CMB filaments



Credits: Planck/ESA

A.C. Rodríguez, *et al* in arXiv:1801.09070 used data from the Baryon Oscillation Spectroscopic Survey, or BOSS, an Earth-based sky survey that captured light from about 1.5 million galaxies

Cosmic Anatomy

Baryons

Missing Baryons

Dark Matter



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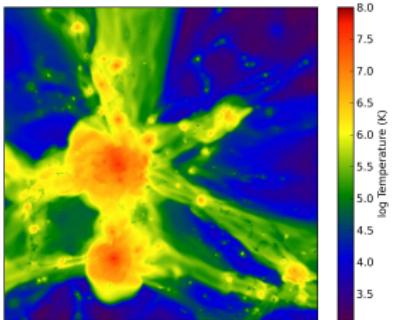
The muscles



Direct observations of filaments

Where are the Baryons? (Cen, Ostriker, astro-ph/9806281 [AJ])

Thus, not only is the universe dominated by dark matter, but more than one half of the normal matter is yet to be detected. (the muscles)



Warm-hot intergalactic medium (WHIM)
Density-weighted temperature projection
of a portion of the refinement box of the
C run of size $(18 h^{-1}\text{Mpc})^3$.
Low temperature WHIM confirmed by O VI
line that peak at $T \sim 3 \times 10^5 \text{ K}$

Credit: Cen, arXiv:1112.4527 [AJ]

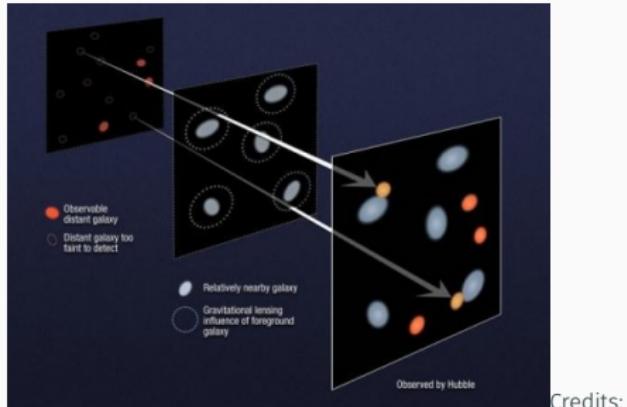
Hotter phases of the WHIM: Observations of the missing baryons in the warm-hot intergalactic medium (Nicastro, et al. arXiv:1806.08395 [Nature]).

The skeleton



Credit: <https://www.disnola.com>

Weak gravitational lensing



Credits:

NASA/WMAP Science TEAM

Weak gravitational lensing.
This technique works somewhat like an X-ray: The light passes through the dark matter until it reaches us.

Upcoming surveys will lead to the discovery of thousands of new gravitational lens systems

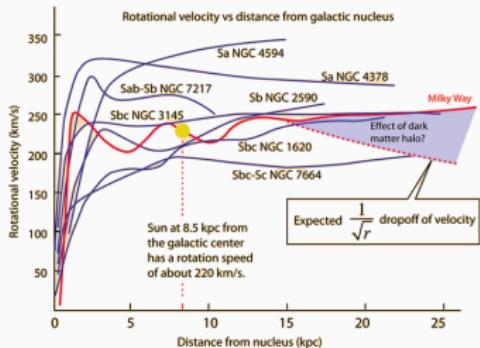
DM allows for a (free) factor of 10 increase in the discovery potential of very large telescopes that try to look at First Light objects

Bullet cluster



Further evidences

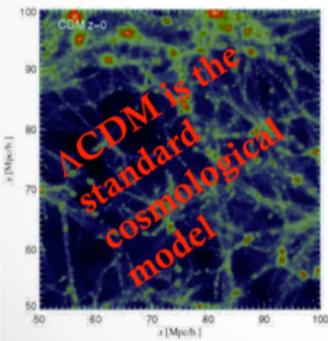
- For cluster of Galaxies
- For stars inside Galaxies
- Inside our Galaxy



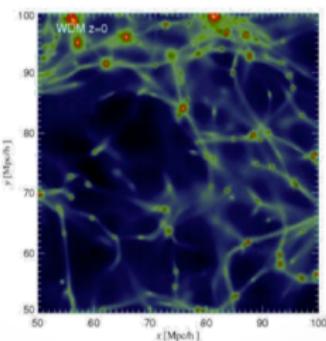
Λ CDM paradigm

Dark matter simulations

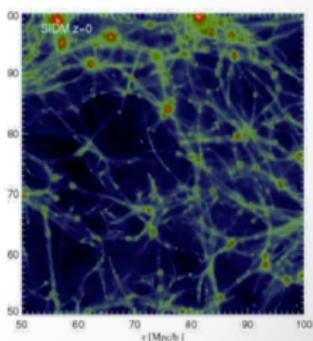
Cold Dark Matter
(Slow moving)
 $m \sim \text{GeV-TeV}$
Small structures form
first, then merge



Warm Dark Matter
(Fast moving)
 $m \sim \text{keV}$
Small structures are
erased



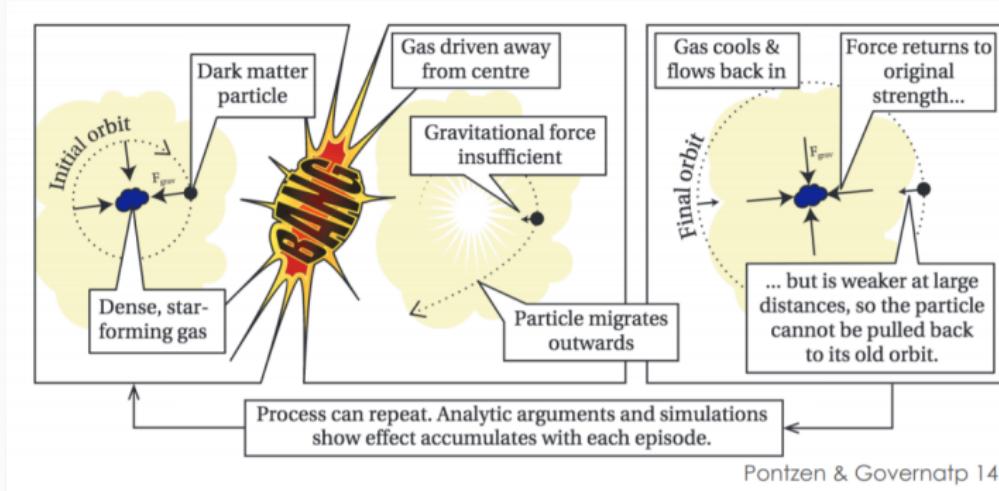
Self-Interacting Dark Matter
Strongly interact with itself
Large scale similar to CDM,
Small galaxies are different



Credit: Arianna Di Cintio (Conference on Shedding Light on the Dark Universe with Extremely Large Telescopes, ICTP

- 2018)

Baryonic effects



Once the effect of baryonic physics is included, it is hard to distinguish between WDM/SIDM/CDM

Dark matter interpretations

| | | | | | | | | |
|-------------|---------------|-----------|---------------|-----------------|-----------|----------------|------------------|-----------|
| u | u | u | c | c | c | t | t | t |
| d | d | d | s | s | s | b | b | b |
| e^- | ν_e | | μ^- | ν_μ | | τ^- | ν_τ | |
| \bar{u} | \bar{u} | \bar{u} | \bar{c} | \bar{c} | \bar{c} | \bar{t} | \bar{t} | \bar{t} |
| \bar{d} | \bar{d} | \bar{d} | \bar{s} | \bar{s} | \bar{s} | \bar{b} | \bar{b} | \bar{b} |
| \bar{e}^+ | $\bar{\nu}_e$ | | $\bar{\mu}^+$ | $\bar{\nu}_\mu$ | | $\bar{\tau}^+$ | $\bar{\nu}_\tau$ | |
| g | g | g | g | g | g | Y | W^- | W^+ |
| | | | | | | | Z^0 | H |

Z_2 Singlet scalar dark matter

Minimal number of fields:

SM (even) + Real Scalar singlet (odd)

$$\mathcal{L} = \mathcal{L}_{\text{SM}} \frac{1}{2} \partial_\mu S \partial^\mu S - \lambda_{HS} H^\dagger H S^2 - M_S^2 S^2 - \lambda_S S^4.$$

S is stable and a dark matter particle

At some high temperature after reheating the scalar dark matter particle is in thermal equilibrium with the standard model plasma. We rather high yield. We need the Higgs portal interaction to exponential decrease the abundance until the correct relic density value. Freeze-out at $M_S/T \sim 20$

See <http://bit.ly/singletscalar>

requires $\lambda_{HS} \sim 0.1$ (weak-like annihilation cross sections)

Dark Matter Boltzmann equation

This program is made to reproduce the behavior of dark matter yield in WIMP and FIMP frameworks based on Chapter 5th, Kolb Turner (Early Universe)

```
[ ] %pylab inline
import numpy as np
from numpy import arange
from scipy.integrate import odeint

[ ] # parameters
Ms = 100 #GeV Singlet Mass
Mp = 1.22e19 #GeV Planck Mass
g = 100 # Degrees of freedom
gs = 106.75 # Entropy degrees of freedom
H0 = 2.133*(0.7)*1e-42 # GeV Hubble parameter (unused)
```

Boltzmann equation

The general expression for the thermal evolution of DM is as follows (see eq (5.26) Kolb and Turner):

$$\frac{x}{Y_{EQ}(x)} \frac{dY}{dx} = -\frac{n_{EQ}(x)\langle\sigma v\rangle}{H(x)} \left[\left(\frac{Y}{Y_{EQ}(x)} \right)^2 - 1 \right],$$

donde

$$n_{EQ}(x) = 2 \left(\frac{M^2}{2\pi x} \right)^{3/2} e^{-x}$$

and [see (eq 5.16) Kolb & Turner]

$$H(x) = 1.67x^{-2}g_*^{1/2} \frac{M^2}{Mp}$$

The equilibrium distribution of this particles is consider for the non-relativistic case, as follows (see eq 5.25):

$$Y_{EQ}(x) = \frac{45}{2\pi^4} \frac{g}{g_{*s}} x^{3/2} e^{-x} = 0.145 \frac{g}{g_{*s}} x^{3/2} e^{-x},$$

where $x = M/T$ and $M = 100$ GeV is the singlet mass taken as constant.

WIMP

The initial condition to solve the evolution equation is $Y(x_i) = Y_{EQ}$, where $x_i = 0.01$, such that $T_i = M/x_i = 10^4$ GeV.

```
[7] def Yeq(x):
    return 0.145*(g/gs)*(x)**(3/2)*np.exp(-x).

xi=1E-4
xe=1000
npts=3000
# For several order of magnitude:
x = np.linspace(0.01, 1000, 1000)

sigmav=[1.747556819623999e-09,1.747556819623999e-06]
def eqd(yl,x,Ms = 100,ov = sigmav[0]):
    ...
    Ms [GeV]      : Singlet Mass
    ov: [1/GeV^2] : (ov)
    ...

    Mp = 1.22e19
    g = 100          # Degrees of freedom
    gs = 106.75      # Entropy degrees of freedom

    H = 1.67*g**((1/2)*Ms**2/Mp)

    dyl = -2*((Ms**2/(2*np.pi*x))**(3/2)*np.exp(-x))*ov/(x**(-2)*H*x))*(yl**2) - (0.145*(g/gs)*(x)**(3/2)
    ...
    return dyl
```

CODE TEXT ↑ CELL ↓ CELL

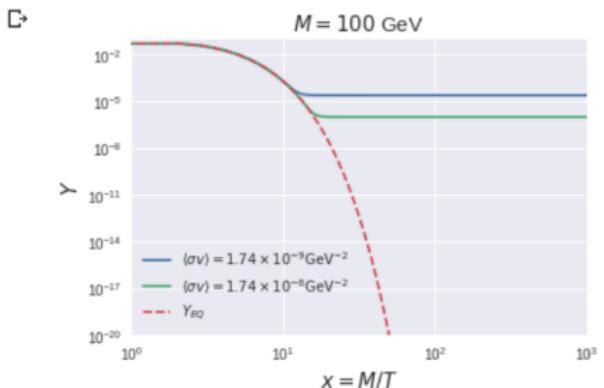
CONNECTED

EDITING

```
[9] #Initial conditions
y0 = Yeq(x[0])
yl = odeint( eqd, y0, x,args=(Ms,sigmap[0]) )
yll = odeint( eqd, y0, x,args=(Ms,sigmap[1]) )
```

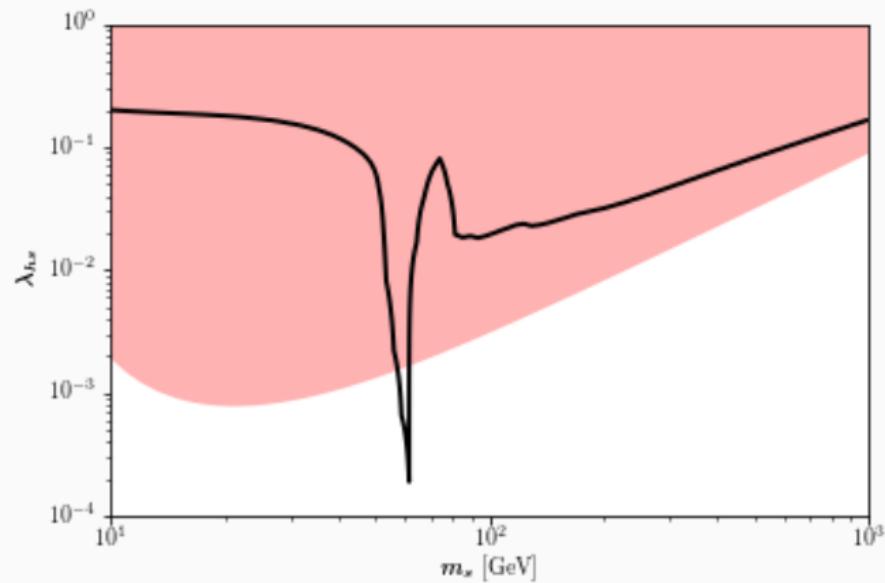
The following plot can be find in the reference book (Figure 5.1)

```
[10] plt.loglog(x,yl, label = r'$\langle \sigma v \rangle = 1.74 \times 10^{-9} \text{ GeV}^{-2}$')
plt.loglog(x,yll, label = r'$\langle \sigma v \rangle = 1.74 \times 10^{-6} \text{ GeV}^{-2}$')
plt.loglog(x,Yeq(x), '--', label = 'Y_EQ')
plt.ylim(ymax=0.1, ymin=1e-20)
plt.xlim(xmax=1e3, xmin=1)
plt.xlabel('x = M/T', size= 15)
plt.ylabel('Y', size= 15)
plt.title('M = 100 GeV', size= 15)
plt.legend(loc='best', fontsize=10)
plt.grid(True)
```

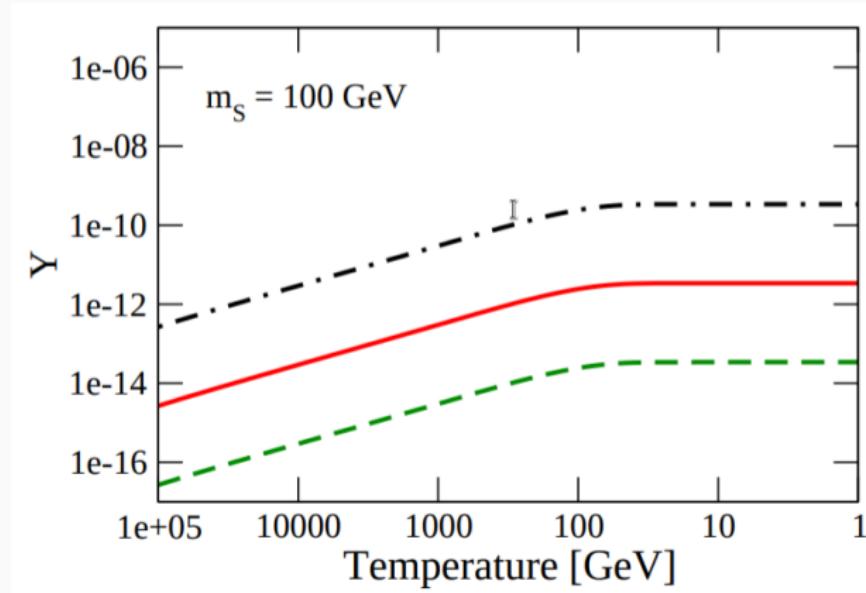


Full wimp parameter space

Restricted by direct detection cross section Xenon1T
(arXiv:1805.12562)



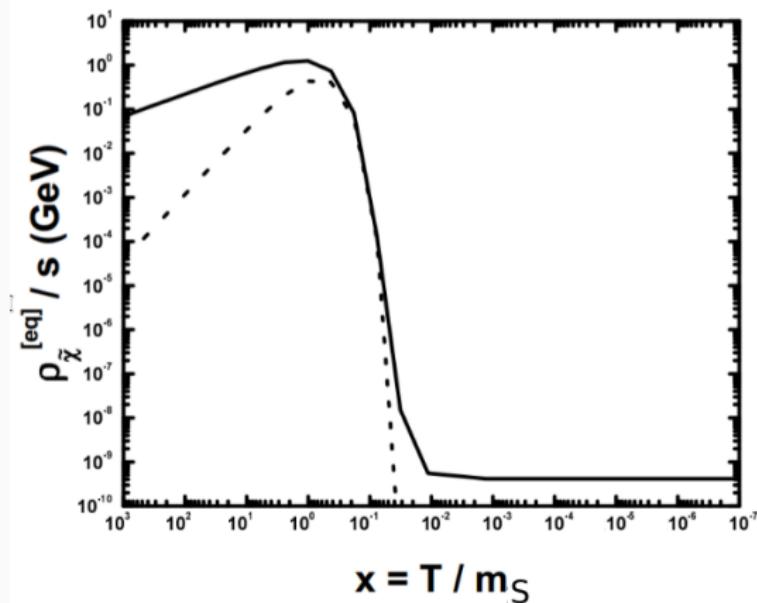
At some high temperature after reheating the abundance of scalar dark matter particle is zero. A feebly interaction allows the Higgs to produce dark matter particles until $M_S/T \sim 1$ where the freeze-in is reached. See C. Yaguna arXiv:1105.1654 [JHEP]



$\lambda_{HS} = 10^{-10}, 10^{-11}, 10^{-12}$. Not signals at all!

WIMP during reheating

The freeze-out occurs at $M/T_{\text{RH}} \sim 10^3$. Very slow reheating for $M \sim 1\text{TeV}$ - After reheating the dark matter particle not longer termalize and the freeze out is kept. See C. Pallis hep-ph/0402033 [APP]



$M_\phi = 10^6 \text{ GeV}$, $M_S = 350 \text{ GeV}$, $T_{\text{RH}} = 5 \text{ GeV}$, $N_S^i = 1.4 \times 10^{-7}$. λ_{HS} fixed to be compatible with direct detection constraints

FIMP during reheating

Produced abundance can be an initial condition for the FIMP after reheating.

See poster session (Valentina Montoya)

Conclusions

Even a simple model like the Singlet Scalar Dark matter model can span many orders of magnitude in mass and in the strength of the interactions (Higgs portal λ_{HS} , self-interaction λ_S) after non-thermal and non-standard cosmology scenarios are considered

Mass: few MeV to PeV λ_{SH} : 10^{-16} to 10

May be the window to search for dark matter is too wide