

Dark matter in left-right symmetric standard model

triplet scalar dark matter portal



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May 21, 2017

5th Uniandes Particle Physics School

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Focus on

arXiv:1703.08148 (PRD)

In collaboration with

C. Arbeláez (UFSM)& M. Hirsch (IFIC)

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Review of SM

Lorentz transformation

Boost in x

$$\{x^\mu\} = \begin{pmatrix} t \\ x \\ y \\ z \end{pmatrix} \rightarrow \begin{pmatrix} t' \\ x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \frac{t+vx}{\sqrt{1-v^2}} \\ \frac{x+vt}{\sqrt{1-v^2}} \\ y \\ z \end{pmatrix} = \begin{pmatrix} \cosh \xi & \sinh \xi & 0 & 0 \\ \sinh \xi & \cosh \xi & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} t \\ x \\ y \\ z \end{pmatrix}$$
$$x'^\mu \rightarrow x^\mu = \Lambda_x{}^\mu{}_\nu x^\nu,$$

where ($c = 1$)

$$\cosh \xi = \gamma \quad \sinh \xi = v\gamma, \quad \text{and} \quad \gamma = \frac{1}{\sqrt{1-v^2}}.$$

In general

$$\Lambda^\mu{}_\nu = [\exp(\boldsymbol{\xi} \cdot \mathbf{K} + i\boldsymbol{\theta} \cdot \mathbf{L})]^\mu{}_\nu$$

$$x' \cdot y' = x'_\mu y'^\mu = x_\mu y^\mu, \quad \text{if:} \quad g_{\alpha\beta} = \Lambda^\mu{}_\alpha g_{\mu\nu} \Lambda^\nu{}_\beta.$$

Theorem of Noether

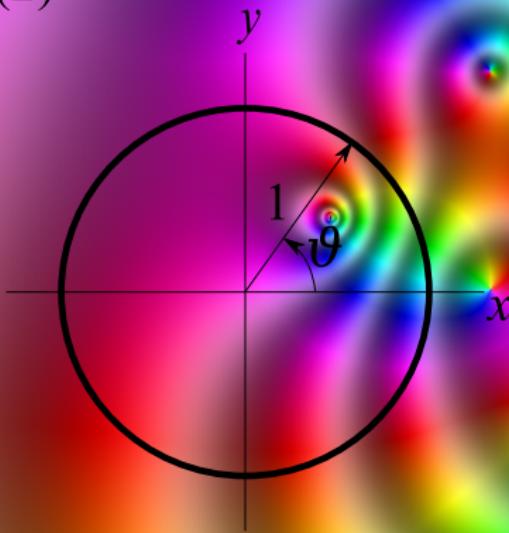
For each continuous symmetry there is at least one conserved charge.



Emmy Noether (1882-1935)

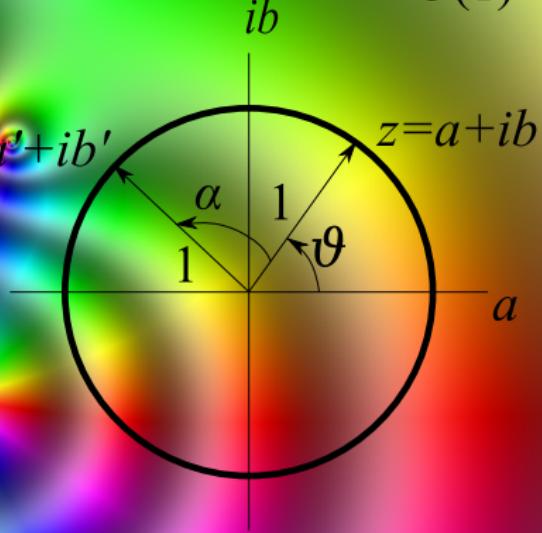
círculo de
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$R(2)$



$$z' = a' + ib'$$

$U(1)$



$$\sqrt{x^2 + y^2} = 1$$

$$|z| = \sqrt{a^2 + b^2} = 1$$

Only fields

Lorentz :

$$\mathcal{L} = -\frac{1}{4} (\partial^\mu A^\nu - \partial^\nu A^\mu) (\partial_\mu A_\nu - \partial_\nu A_\mu) . \quad (1)$$

$$A^\mu(x) \rightarrow A'^\mu(x) = \Lambda^\mu{}_\nu A^\nu(\Lambda^{-1}x) \quad \text{Vector field}$$

$$\phi(x) \rightarrow \phi'(x) = \phi(\Lambda^{-1}x) \quad \text{Scalar field}$$

$$\psi_\alpha(x) \rightarrow \psi'_\alpha(x) = [S(\Lambda)]_\alpha{}^\beta \psi_\beta(\Lambda^{-1}x), \quad \text{Left Weyl spinor}$$

$$(\psi_\alpha(x))^\dagger = \psi_{\dot{\alpha}}^\dagger(x) \rightarrow \psi'_{\dot{\alpha}}^\dagger(x) = [S^*(\Lambda)]_{\dot{\alpha}}{}^{\dot{\beta}} \psi_{\dot{\beta}}^\dagger(\Lambda^{-1}x), \quad \text{Right anti-Weyl spinor}$$

With

$$S(\Lambda) = \exp \left(\boldsymbol{\xi} \cdot \frac{\boldsymbol{\sigma}}{2} + i \boldsymbol{\theta} \cdot \frac{\boldsymbol{\sigma}}{2} \right) ,$$

where $\boldsymbol{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$ are the Pauli matrices. $\overline{\boldsymbol{\sigma}} \equiv -\boldsymbol{\sigma}$.

Only fields

Lorentz+ $U(1)$:

$$\mathcal{L} = \partial_\mu \phi^* \partial^\mu \phi - \textcolor{red}{m}^2 \phi^* \phi - \lambda (\phi^* \phi)^2 . \quad (1)$$

$$A^\mu(x) \rightarrow A'^\mu(x) = \Lambda^\mu{}_\nu A^\nu(\Lambda^{-1}x) \quad \text{Vector field}$$

$$\phi(x) \rightarrow \phi'(x) = \phi(\Lambda^{-1}x) \quad \text{Scalar field}$$

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$$(\psi_\alpha(x))^\dagger = \psi_{\dot{\alpha}}^\dagger(x) \rightarrow \psi'_{\dot{\alpha}}^\dagger(x) = [S^*(\Lambda)]_{\dot{\alpha}}{}^{\dot{\beta}} \psi_{\dot{\beta}}^\dagger(\Lambda^{-1}x), \quad \text{Right anti-Weyl spinor}$$

With

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where $\boldsymbol{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$ are the Pauli matrices. $\overline{\boldsymbol{\sigma}} \equiv -\boldsymbol{\sigma}$.

Weyl spinor

$$\begin{aligned}\mathcal{L} &= i\psi_{\dot{\alpha}}^{\dagger} \bar{\sigma}^{\mu\dot{\alpha}\alpha} \partial_{\mu} \psi_{\alpha} - m \left(\psi^{\alpha} \psi_{\alpha} + \psi_{\dot{\alpha}}^{\dagger} \psi^{\dagger\dot{\alpha}} \right) \\ &= i\psi^{\dagger} \bar{\sigma}^{\mu} \partial_{\mu} \psi - m \left(\psi \psi + \psi^{\dagger} \psi^{\dagger} \right).\end{aligned}\tag{2}$$

Only if

$$S^{\dagger} a^{\nu} S = (\Lambda)^{\nu}_{\rho} a^{\rho},$$

Scalar product: $\alpha_\alpha \alpha^\dot{\alpha}$ and $\dot{\alpha}^\dot{\alpha}$.

Name	Symbol Name	Lorentz N
e_L : left electron	ξ_α	$[S]_\alpha^\beta$
$(e_L)^\dagger = e_R^\dagger$: right positron	$(\xi_\alpha)^\dagger = \xi_{\dot{\alpha}}^\dagger$	$[S^*]_{\dot{\alpha}}^\beta$
e_R : right electron	$(\eta^\alpha)^\dagger = \eta^\dagger{}^\dot{\alpha}$	$[(S^{-1})^\dagger]^{\dot{\alpha}}$
$(e_R)^\dagger = e_L^\dagger$: left positron	η^α	$[(S^{-1})^T]^\alpha$

$$\begin{aligned}\mathcal{L} &= i\xi_{\dot{\alpha}}^\dagger \bar{\sigma}^\mu {}^{\dot{\alpha}\alpha} \partial_\mu \xi_\alpha + i\eta^\alpha \sigma^\mu_{\alpha\dot{\alpha}} \partial_\mu \eta^\dagger{}^{\dot{\alpha}} - m \left(\eta^\alpha \xi_\alpha + \xi_{\dot{\alpha}}^\dagger \eta^\dagger{}^{\dot{\alpha}} \right) \\ &= i\xi^\dagger \bar{\sigma}^\mu \partial_\mu \xi + i\eta \sigma^\mu \partial_\mu \eta^\dagger - m \left(\eta \xi + \xi^\dagger \eta^\dagger \right) + .\end{aligned}$$



Scalar product: $\alpha_\alpha \alpha_\alpha$ and $\dot{\alpha} \dot{\alpha}$.

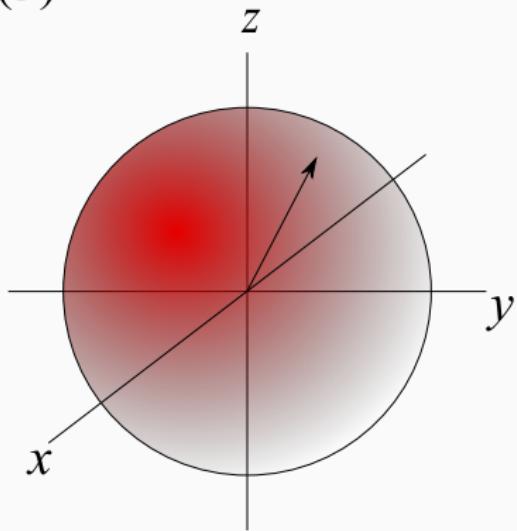
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$(e_R)^\dagger = e_L^\dagger$: left positron	η^α	$[(S^{-1})^T]^\alpha$

$$\mathcal{L} = i \begin{pmatrix} \xi_{\dot{\alpha}}^\dagger & \eta^\alpha \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma_{\alpha\dot{\alpha}}^\mu \\ \bar{\sigma}^{\mu\dot{\alpha}\alpha} & 0 \end{pmatrix} \begin{pmatrix} \partial_\mu \xi_\alpha \\ \partial_\mu \eta^\dagger \dot{\alpha} \end{pmatrix} - m \begin{pmatrix} \xi_{\dot{\alpha}}^\dagger & \eta^\alpha \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

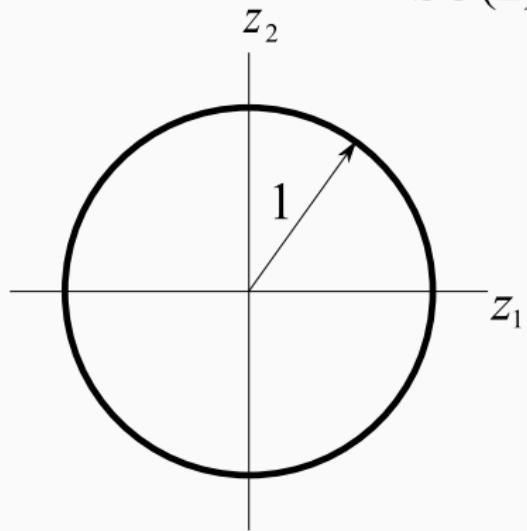


3 generadores

$R(3)$

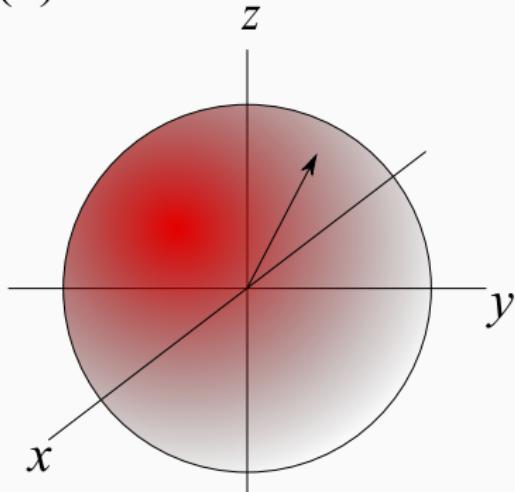


$SU(2)$

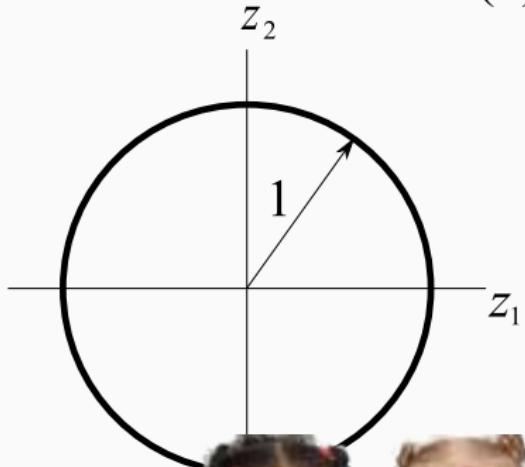


3 generadores

$R(3)$



$SU(2)$

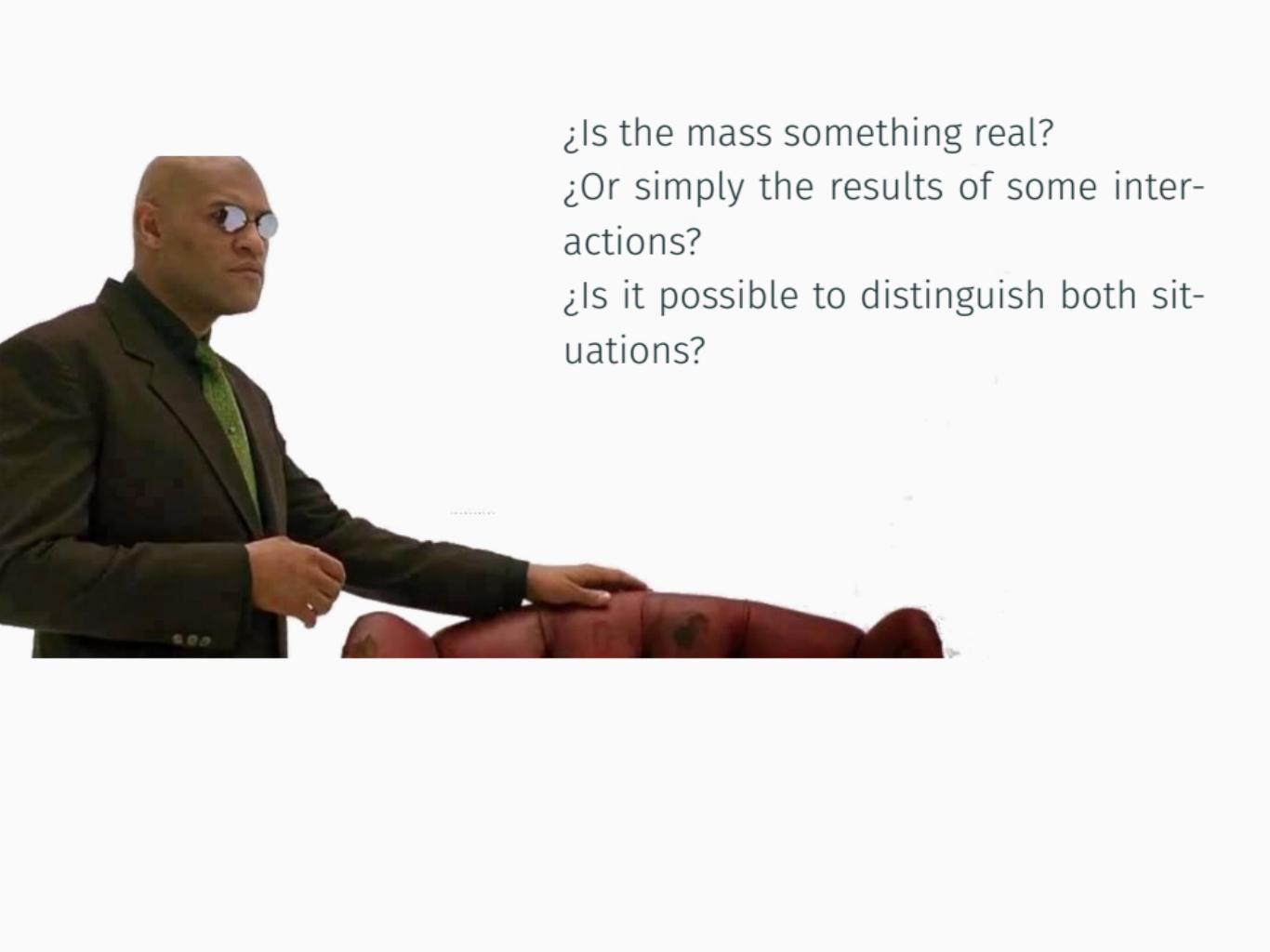


$SU(2)_L$ in an ideal Universe

$$L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix} \quad Q = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$$

Massless states with same (hyper)charge!





¿Is the mass something real?
¿Or simply the results of some interactions?
¿Is it possible to distinguish both situations?



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¿Is it possible to distinguish both situations?

Answer

- Prediction of charm quark (Bjorken-Glashow 1964)
- W and Z mass and SM Higgs boson



¿Is the mass something real?
¿Or simply the results of some interactions?
¿Is it possible to distinguish both situations?

Answer

- Prediction of charm quark (Bjorken-Glashow 1964)
- W and Z mass and **SM Higgs boson**
- $m_{\text{proton}} \gg 2m_u + m_d$

Consider the left-quark up: u

$$\Psi_L^u = \begin{pmatrix} u_L \\ u_L \\ u_L \end{pmatrix}$$

From global to local

Let $U(x) \in SU(N)$ and Ψ a multiplet of left spinors under $SU(N)$. We replace ∂_μ by the covariant derivative \mathcal{D}_μ :

$$\Psi \rightarrow \Psi' = U\Psi, \quad \Psi^\dagger \rightarrow \Psi'^\dagger = \Psi^\dagger U^\dagger,$$

$$\mathcal{D}_\mu \Psi \rightarrow (\mathcal{D}_\mu \Psi)' = U(\mathcal{D}_\mu \Psi), \quad (\mathcal{D}_\mu \Psi)^\dagger \rightarrow (\mathcal{D}_\mu \Psi)'^\dagger = (\mathcal{D}_\mu \Psi)^\dagger U^\dagger,$$

since the mass term is forbidden

$$\mathcal{L}_{SU(N)} = i\Psi \bar{\sigma}^\mu \mathcal{D}_\mu \Psi - \frac{1}{2} \text{Tr}(G_{\mu\nu} G^{\mu\nu})$$

with the following $N \times N$ matrix equations ($\alpha_N \equiv g^2/4\pi$ is the $SU(N)$ coupling)

$$\mathcal{D}_\mu = \not{1}\partial_\mu - igG_\mu, \quad \widehat{G}_{\mu\nu} = \frac{i}{g} [\mathcal{D}_\mu, \mathcal{D}_\nu]$$

Details: <http://fisica.udea.edu.co> → cursos → Física Subatómica

Standard Model Lagrangian (simplified version)

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi + h.c. \\ & + \bar{\chi}_i Y_{ij} \chi_j \phi + h.c. \\ & + |\partial_\mu \phi|^2 - V(\phi)\end{aligned}$$



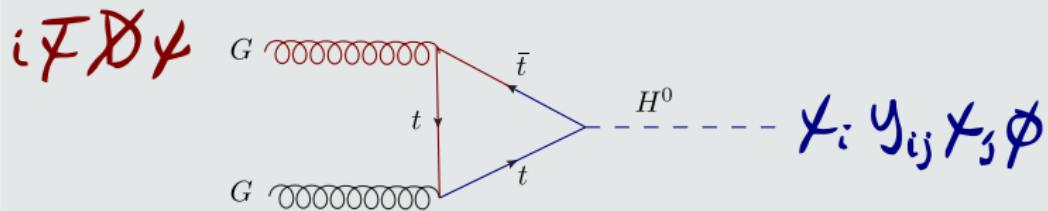
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$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi + h.c. \\ & + \bar{\chi}_i Y_{ij} \chi_j \phi + h.c. \\ & + |\partial_\mu \phi|^2 - V(\phi)\end{aligned}$$

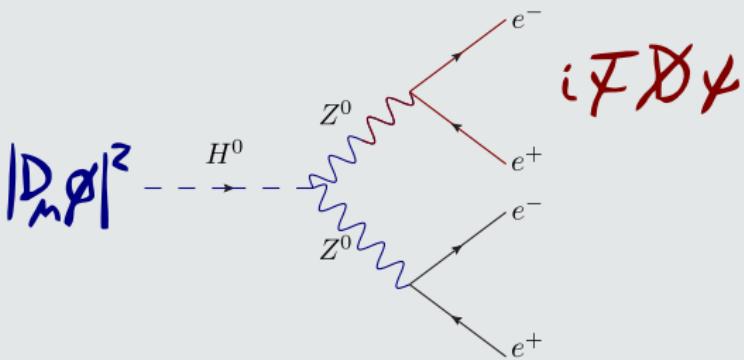


Higgs production and decay

Gluon fusion



Golden channel



Standard Model: particles

$$SU(3)_c$$

$$\textcolor{red}{d_L^1} \quad d_L^2 \quad d_L^3$$

$$\textcolor{red}{d_R^1} \quad d_R^2 \quad d_R^3$$

Standard Model: particles

$$SU(3)_c$$

$$\begin{array}{ccc} \color{red}d_L^1 & d_L^2 & d_L^3 \\ u_L^1 & u_L^2 & u_L^3 \\ \color{red}d_R^1 & d_R^2 & d_R^3 \\ u_R^1 & u_R^2 & u_R^3 \end{array}$$

Standard Model: particles

$SU(3)_c$

$$\begin{array}{cccccccccc} d_L^1 & d_L^2 & d_L^3 & s_L^1 & s_L^2 & s_L^3 & b_L^1 & b_L^2 & b_L^3 \\ u_L^1 & u_L^2 & u_L^3 & c_L^1 & c_L^2 & c_L^3 & t_L^1 & t_L^2 & t_L^3 \\ d_R^1 & d_R^2 & d_R^3 & s_R^1 & s_R^2 & s_R^3 & b_R^1 & b_R^2 & b_R^3 \\ u_R^1 & u_R^2 & u_R^3 & c_R^1 & c_R^2 & c_R^3 & t_R^1 & t_R^2 & t_R^3 \end{array}$$

Standard Model: particles

$$SU(3)_c$$

$$\begin{array}{l} G_1 \\ G_2 \\ \vdots \\ G_8 \end{array} \left\{ \begin{array}{ccccccccccccc} d_L^1 & d_L^2 & d_L^3 & s_L^1 & s_L^2 & s_L^3 & b_L^1 & b_L^2 & b_L^3 & \rightarrow & d_L^\alpha & s_L^\alpha & b_L^\alpha \\ u_L^1 & u_L^2 & u_L^3 & c_L^1 & c_L^2 & c_L^3 & t_L^1 & t_L^2 & t_L^3 & \rightarrow & u_L^\alpha & c_L^\alpha & t_L^\alpha \\ d_R^1 & d_R^2 & d_R^3 & s_R^1 & s_R^2 & s_R^3 & b_R^1 & b_R^2 & b_R^3 & \rightarrow & d_R^\alpha & s_R^\alpha & b_R^\alpha \\ u_R^1 & u_R^2 & u_R^3 & c_R^1 & c_R^2 & c_R^3 & t_R^1 & t_R^2 & t_R^3 & \rightarrow & u_R^\alpha & c_R^\alpha & t_R^\alpha \end{array} \right.$$

Standard Model: particles

$SU(3)_c$

$$\begin{array}{ccccccccc}
 & & & & & & & e_R & \mu_R & \tau_R \\
 & & & & & & & e_L & \mu_L & \tau_L \\
 & & & & & & & \nu_L^e & \nu_L^\mu & \nu_L^\tau \\
 G_1 & \left\{ \begin{array}{ccccccccc} d_L^1 & d_L^2 & d_L^3 & s_L^1 & s_L^2 & s_L^3 & b_L^1 & b_L^2 & b_L^3 \\ u_L^1 & u_L^2 & u_L^3 & c_L^1 & c_L^2 & c_L^3 & t_L^1 & t_L^2 & t_L^3 \\ d_R^1 & d_R^2 & d_R^3 & s_R^1 & s_R^2 & s_R^3 & b_R^1 & b_R^2 & b_R^3 \\ u_R^1 & u_R^2 & u_R^3 & c_R^1 & c_R^2 & c_R^3 & t_R^1 & t_R^2 & t_R^3 \end{array} \right. & \rightarrow & d_L^\alpha & s_L^\alpha & b_L^\alpha \\
 G_2 & & & & & & & u_L^\alpha & c_L^\alpha & t_L^\alpha \\
 \vdots & & & & & & & d_R^\alpha & s_R^\alpha & b_R^\alpha \\
 G_8 & & & & & & & u_R^\alpha & c_R^\alpha & t_R^\alpha
 \end{array}$$

Standard Model: particles

$$SU(3)_c \times SU(2)_L$$

$$\begin{array}{c}
W^\pm \quad \left\{ \begin{array}{l} \\
G_1 \quad \left\{ W^\pm \quad \left\{ \begin{array}{cccccccccc}
d_L^1 & d_L^2 & d_L^3 & s_L^1 & s_L^2 & s_L^3 & b_L^1 & b_L^2 & b_L^3 & \rightarrow & d_L^\alpha & s_L^\alpha & b_L^\alpha \\
u_L^1 & u_L^2 & u_L^3 & c_L^1 & c_L^2 & c_L^3 & t_L^1 & t_L^2 & t_L^3 & \rightarrow & u_L^\alpha & c_L^\alpha & t_L^\alpha \\
d_R^1 & d_R^2 & d_R^3 & s_R^1 & s_R^2 & s_R^3 & b_R^1 & b_R^2 & b_R^3 & \rightarrow & d_R^\alpha & s_R^\alpha & b_R^\alpha \\
u_R^1 & u_R^2 & u_R^3 & c_R^1 & c_R^2 & c_R^3 & t_R^1 & t_R^2 & t_R^3 & \rightarrow & u_R^\alpha & c_R^\alpha & t_R^\alpha
\end{array} \right\} \quad Q_1^\alpha \quad Q_2^\alpha \quad Q_3^\alpha
\end{array} \right\} \quad L_1 \quad L_2 \quad L_3
\end{array}$$

$$m_i = 0$$

Standard Model: particles

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

$$\begin{aligned}
& Z^0 \left\{ \begin{array}{l} W^\pm \\ G_1 \\ G_2 \\ \vdots \\ G_8 \end{array} \right\} \left\{ \begin{array}{l} W^\pm \\ \begin{array}{ccccccccccccc} d_L^1 & d_L^2 & d_L^3 & s_L^1 & s_L^2 & s_L^3 & b_L^1 & b_L^2 & b_L^3 \\ u_L^1 & u_L^2 & u_L^3 & c_L^1 & c_L^2 & c_L^3 & t_L^1 & t_L^2 & t_L^3 \\ d_R^1 & d_R^2 & d_R^3 & s_R^1 & s_R^2 & s_R^3 & b_R^1 & b_R^2 & b_R^3 \\ u_R^1 & u_R^2 & u_R^3 & c_R^1 & c_R^2 & c_R^3 & t_R^1 & t_R^2 & t_R^3 \end{array} \end{array} \right\} \rightarrow \begin{array}{l} e_R \\ e_L \\ \nu_L^e \end{array} \begin{array}{l} \mu_R \\ \mu_L \\ \nu_L^\mu \end{array} \begin{array}{l} \tau_R \\ \tau_L \\ \nu_L^\tau \end{array} \left\{ \begin{array}{l} L_1 \\ L_2 \\ L_3 \end{array} \right\} \begin{array}{l} Q_1^\alpha \\ Q_2^\alpha \\ Q_3^\alpha \end{array}
\end{aligned}$$

$$m_i = 0$$

$$Y_{u_L} = Y_{d_L}$$

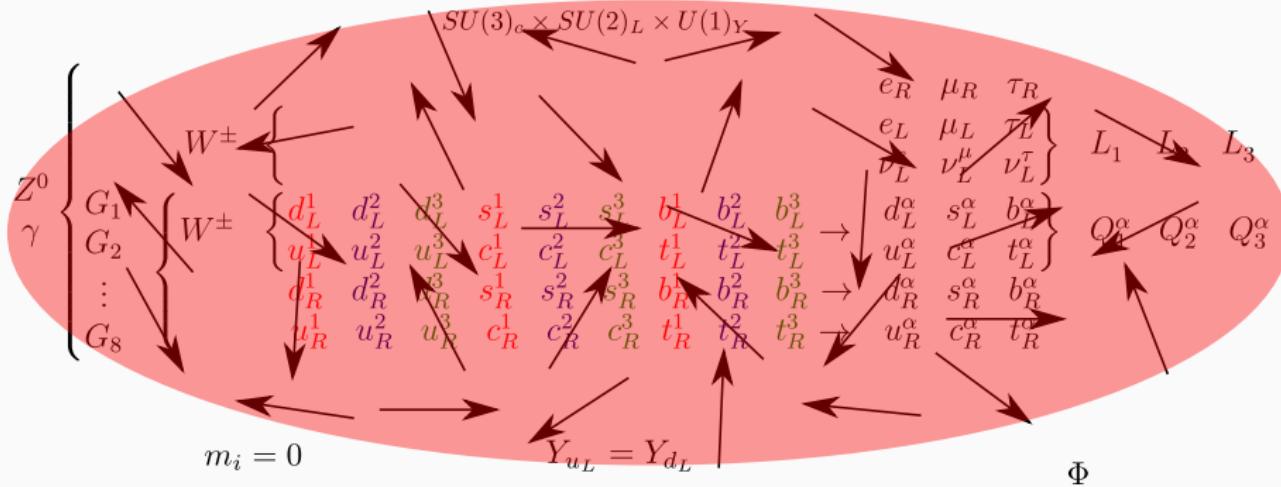
Standard Model: particles

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

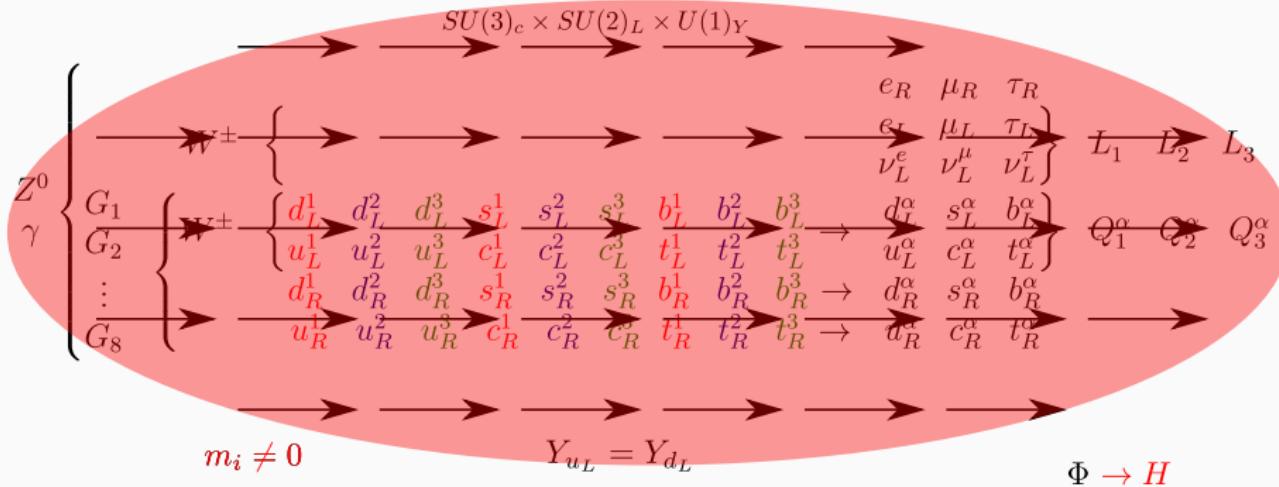
$$m_i = 0$$

$$Y_{u_L} = Y_{d_L}$$

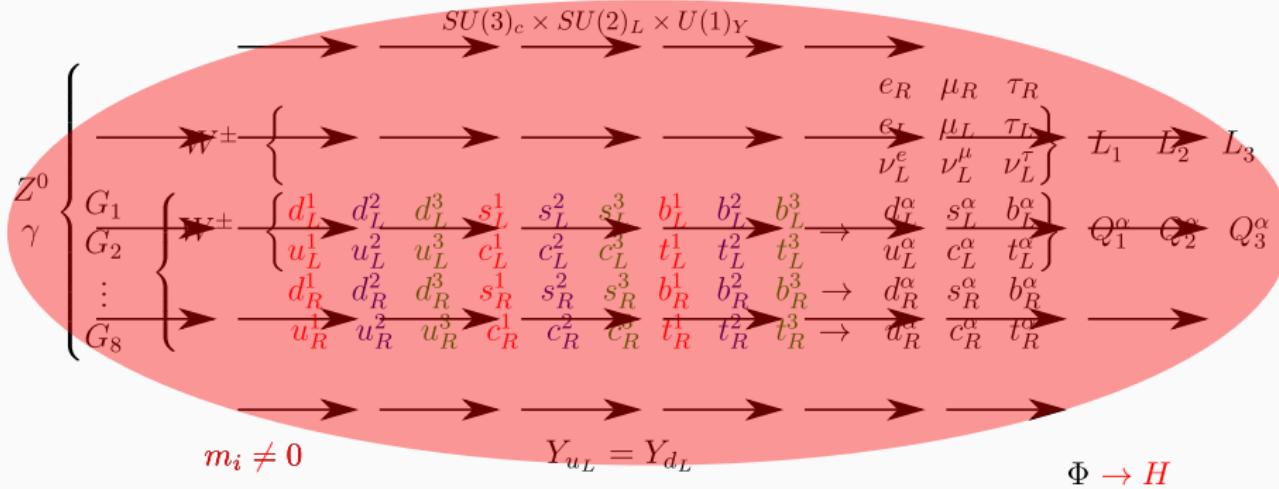
Standard Model: particles



Standard Model: particles



Standard Model: particles



SM fields + Gauge + Lorentz Invariance \rightarrow Lagrangian

Electroweak superconductor (*Higgs condensate*) with emergent masses which allows

- confined baryons (nuclear physics)
- atoms (chemistry)

Links

<http://fisica.udea.edu.co> → cursos

- CNF-665: Física subatómica:
- CNF-666: Teoría Cuántica de Campos



cursos fisica udea

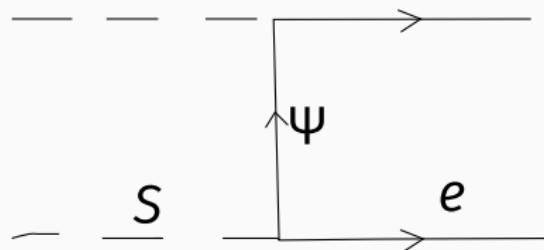


Dark matter and unification

Dark matter

Name	Symbol	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	Z_2
$(\nu_L \ e_L)^\dagger$	$(\xi_{1\alpha} \ \xi_{2\alpha})^\dagger$	1	2	-1/2	+1
$(e_R)^\dagger$	η_1^α	1	1	+1	+1
$(\psi_R)^\dagger$	η_2^α	1	1	+1	-1
ψ_L	$\xi_{3\alpha}$	1	1	-1	-1
S		1	1	0	-1

$$\mathcal{L} = M_\psi \left[(\psi_R)^\dagger \psi_L (\psi_L)^\dagger \psi_R \right] + h_S \left[S (e_R)^\dagger \psi_L + S (\psi_L)^\dagger e_R \right] \quad (3)$$



SARAH



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Latest commit 63e9c7

- [SARAH](#) Yse intialized to 0.01
- [SPHENO](#) Yse intialized to 0.01
- [SSP](#) toolbox 1.2.8
- [autom4te.cache](#) Update to 1.2.10
- [calchep](#) Fix compilation
- [m4](#) Initial relase after untar
- [madgraph](#) nf
- [micromegas](#) Misising files



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..

Scotogenic.ipynb	Scotogenic beta version	
SimplifiedDM-DFDM.ipynb	Feynrules model by Andres	11
SimplifiedDM-IDM.ipynb	Bug in micromegas for degenerate masses fixed by hand	11
SimplifiedDM-SSSFDM.ipynb	nf	18
cmdlike.py	merge this changes	11
hep.py	Fix for single channel with direct detection	7
pdg.py	merge this changes	11
pdg_series.py	Including pdg object in hep model class	10
test_toolbox.py	Updated SSDM model	
tripletLR.ipynb	Bug in micromegas for degenerate masses fixed by hand	11

Reproducibility

Singlet Scalar Singlet (charged) Fermion dark matter model

SSSFDM

- [arXiv:1307.6181](#)
- [arXiv:1307.6480](#)

Particle content

In [61]:

```
%%latex
\begin{array}{lllllll}
\text{Name} & \text{Symbol} & \text{SU}(3)_c & \text{SU}(2)_L & \text{U}(1)_Y \\
\begin{pmatrix} \nu_L & e_L \end{pmatrix}^T & \begin{pmatrix} \xi_{1\alpha} & \xi_{2\alpha} \end{pmatrix}^T & \mathbf{1} & \mathbf{2} & -1/2 & +1 \\
\begin{pmatrix} \psi_R^\dagger & \eta_1^\alpha & \psi_L & \xi_{3\alpha} \end{pmatrix}^T & \begin{pmatrix} \psi_R^\dagger & \eta_1^\alpha & \psi_L & \xi_{3\alpha} \end{pmatrix}^T & \mathbf{1} & \mathbf{1} & +1 & +1 \\
& & \mathbf{1} & \mathbf{1} & +1 & -1
\end{array}
```

Name	Symbol	SU(3) _c	SU(2) _L	U(1) _Y	Z ₂
$(\nu_L \quad e_L)^T$	$(\xi_{1\alpha} \quad \xi_{2\alpha})^T$	1	2	-1/2	+1
$(e_R)^\dagger$	η_1^α	1	1	+1	+1
$(\psi_R)^\dagger$	η_2^α	1	1	+1	-1
ψ_L	$\xi_{3\alpha}$	1	1	-1	-1

After the spontaneous symmetry breaking, the relevant Yukawa terms are

$$\begin{aligned} \mathcal{L} &= \frac{h_e v}{\sqrt{2}} (\eta_1^\alpha \xi_{2\alpha} + \xi_{2\dot{\alpha}}^\dagger \eta_1^{\dot{\alpha}}) + M_\psi (\eta_2^\alpha \xi_{3\alpha} + \xi_{3\dot{\alpha}}^\dagger \eta_2^{\dot{\alpha}}) + h_S (S \eta_1^\alpha \xi_{3\alpha} + S \xi_{3\dot{\alpha}}^\dagger \eta_1^{\dot{\alpha}}) \\ &= \frac{h_e v}{\sqrt{2}} \begin{pmatrix} \eta_1^\alpha & \xi_{2\dot{\alpha}}^\dagger \end{pmatrix} \begin{pmatrix} \xi_{2\alpha} \\ \eta_1^{\dot{\alpha}} \end{pmatrix} + M_\psi \begin{pmatrix} \eta_2^\alpha & \xi_{3\dot{\alpha}}^\dagger \end{pmatrix} \begin{pmatrix} \xi_{3\alpha} \\ \eta_2^{\dot{\alpha}} \end{pmatrix} \\ &\quad + \left[h_S S \begin{pmatrix} \eta_1^\alpha & \xi_{2\dot{\alpha}}^\dagger \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} \xi_{3\alpha} \\ \eta_2^{\dot{\alpha}} \end{pmatrix} + \text{h.c.} \right] \\ &= \frac{h_e v}{\sqrt{2}} \begin{pmatrix} \xi_{2\dot{\alpha}}^\dagger & \eta_1^\alpha \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \xi_{2\alpha} \\ \eta_1^{\dot{\alpha}} \end{pmatrix} + M_\psi \begin{pmatrix} \xi_{3\dot{\alpha}}^\dagger & \eta_2^\alpha \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \xi_{3\alpha} \\ \eta_2^{\dot{\alpha}} \end{pmatrix} \\ &\quad + \left[h_S S \begin{pmatrix} \xi_{2\dot{\alpha}}^\dagger & \eta_1^\alpha \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} \xi_{3\alpha} \\ \eta_2^{\dot{\alpha}} \end{pmatrix} + \text{h.c.} \right] \end{aligned}$$

Defining

$$e = \begin{pmatrix} \xi_{2\alpha} \\ \eta_1^{\dot{\alpha}} \end{pmatrix} = \begin{pmatrix} e_L \\ e_R \end{pmatrix} \quad \Psi = \begin{pmatrix} \xi_{3\alpha} \\ \eta_2^{\dot{\alpha}} \end{pmatrix} = \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix}$$

the relevant Yukawa Lagrangian in terms of Dirac fermions is

$$\begin{aligned} \mathcal{L} &= \frac{h_e v}{\sqrt{2}} e^\dagger \gamma^0 e + M_\psi \Psi^\dagger \gamma^0 \Psi + (h_S S e^\dagger \gamma^0 P_L \Psi + \text{h.c.}) \\ &= \frac{h_e v}{\sqrt{2}} \bar{e} e + M_\psi \bar{\Psi} \Psi + h_S (S \bar{e} \Psi_L + \text{h.c.}) , \end{aligned}$$

where

$$\Psi_L = P_L \Psi = \begin{pmatrix} \psi_L \\ 0 \end{pmatrix} .$$

```
In [1]: %pylab inline
```

Populating the interactive namespace from numpy and matplotlib

```
In [2]: import pandas as pd
import numpy as np
from matplotlib.colors import LogNorm
import os, sys, inspect
import commands
from hep import *
```

Check one point

```
In [3]: a=hep(MODEL='SimplifiedDMSSSFDM')
```

a-object is an object with many attributes and methods. Use the tab to explore them. Some of them are

- a.Series: [pandas](#) Series object with the "relevant" variables
- a.LHA: Input LesHouces file as [pyslha](#) object
- a.runSPheno() -> a.LHA_out: return LHA output files as [pyslha](#) object
- a.runcinemegas() -> a.runSPheno() -> Updated the a-object with micrOMEGAS "relevant" output

```
In [4]: pd.Series(a.LHA.blocks['MINPAR'].entries)
```

```
Out[4]: 1    2.5500000E-01 # LambdaIN
2    0.0000000E+00 # LamSHIN
3    0.0000000E+00 # LamSIN
4    2.0000000E+02 # MS2Input
5    2.0000000E+02 # MSFIN
```

File Edit View Insert Cell Kernel Help

Python 2

dtype: object

```
In [5]: v=a.vev
lambda_1=0.26
lambda_SH=0.
MS=150**2
MF=200
Yse=1.9
Ymu=0
Ytau=0
devnull=commands.getoutput('rm -f SPheno.spc.%s' %a.MODEL)
a.LHA.blocks['SPHENOINPUT'].entries[55]='0' # Calculate one loop
a.LHA.blocks['SPHENOINPUT'].entries[520]='0.' # Write effective F
a.LHA.blocks['MINPAR'][1]='%0.8E' # lambda1' %lambda_1
a.LHA.blocks['MINPAR'][2]='%0.8E' # lambdaSH' %lambda_SH
a.LHA.blocks['MINPAR'][4]='%0.8E' # MS' %MS
a.LHA.blocks['MINPAR'][5]='%0.8E' # MF' %MF
a.LHA.blocks['YSIN'][1]='%0.8E' # Ys(1)' %Yse
a.LHA.blocks['YSIN'][2]='%0.8E' # Ys(1)' %Ymu
a.LHA.blocks['YSIN'][3]='%0.8E' # Ys(3)' %Ytau
moc=a.runmicromegas(Direct_Detection=True)
print 'Omega h^2, SI proton, neutron =',a.Series.Omega_h2,a.Series.proton_SI
```

Omega h^2, SI proton, neutron = 0.111 0.0 0.0

Scan the parameter space

```
In [ ]: import time
st=time.time()
a=hep(MODEL='SimplifiedDMSSSFDM')
v=a.vev
Omega_h2_delta=0.0022
```

Scan the parameter space

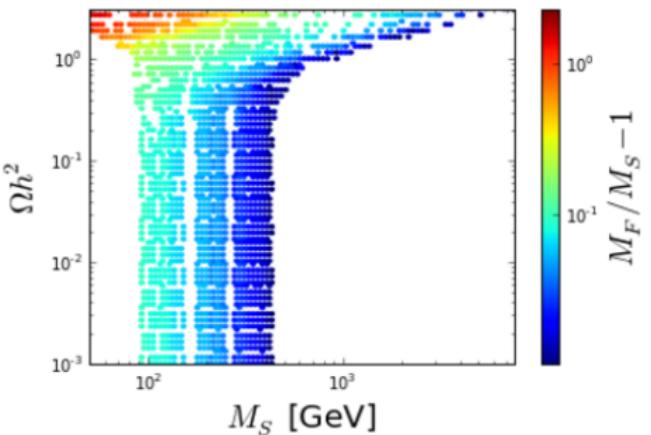
```
In [ ]: import time
st=time.time()
a=hep(MODEL='SimplifiedDMSSSFDM')
v=a.vev
Omega_h2_delta=0.0022
CL=3
Omega_h2=0.1197
Omega_h2_exp=[Omega_h2-CL*Omega_h2_delta,Omega_h2,Omega_h2+CL*Omega_h2_delta]
lambda_1=0.26
lambda_SH=0
a.LHA.blocks['SPHENOINPUT'].entries[55]='0'                                # Calculate one loop
a.LHA.blocks['SPHENOINPUT'].entries[520]='0.'                               # Write effective F
df=pd.DataFrame()
a.LHA.blocks['SPHENOINPUT'].entries[55]='0'                                # Calculate one loop
dfmin=100 #40
dfmax=600 #1E4
npoints=1000
df_masses=np.logspace(np.log10(dfmin),np.log10(dfmax),npoints) #np.array([200,400,600])
DEBUG=False
for i in range(1):
    for MF in df_masses:
        rml1=10**np.random.uniform(np.log10(1E-2),np.log10(3))
        r=rml1+1.
        M_S=MF/r
        MS2=M_S**2-a.vev**2*lambda_SH
        Yse_range=np.logspace(np.log10(np.pi),np.log10(1E-3),200)
        Omega_h2_old=1E32
        for Yse in Yse_range:
            devnull=commands.getoutput('rm -f SPheno.spc.%s' %a.MODEL)
            Yml1=A #10**np.random.uniform(-log10(1E-31),log10(np.pi))
```

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Python 2

```
In [34]: df=df[df.MF>=100]
plt.hexbin(df.ss,df.Ys1,df.MF/df.ss-1,xscale='log',yscale='log',norm=LogNorm
cl=plt.colorbar()
cl.set_label(r'$M_F/M_S-1$',size=20)
plt.xlim(50,df.MF.max())
plt.xlabel(r'$M_S$ [GeV]',size=20)
plt.ylabel(r'$\Omega h^2$',size=20)
```

```
Out[34]: <matplotlib.text.Text at 0x7fc18defe5d0>
```

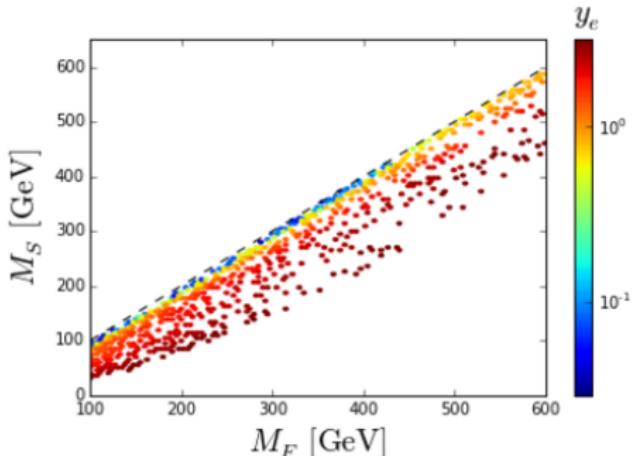


```
In [20]: MF_max=600
plt.hexbin(df[df.MF<MF_max].MF,df[df.MF<MF_max].ss,df[df.MF<MF_max].Ys1,norm
cb=plt.colorbar()
x=np.linspace(100,1000,10)
plt.plot(x,x,'k--')
```

```
In [20]: MF_max=600
plt.hexbin(df[df.MF<MF_max].MF,df[df.MF<MF_max].ss,df[df.MF<MF_max].Ys1,norm=norm)
cb=plt.colorbar()
x=np.linspace(100,1000,10)
plt.plot(x,x,'k--')

plt.xlabel(r'$M_F$ [GeV]',size=20)
plt.ylabel(r'$M_S$ [GeV]',size=20)
plt.text(630,680,r'$y_e$',size=20)
plt.xlim(100,600)
plt.ylim(0,650)

plt.tight_layout()
#plt.savefig('singlet_exc.pdf')
```



File Edit View Insert Cell Kernel Help | Python 2

In [7]: `from madgraph import *`

In [10]: `generate_cross_section(a.MODEL,processes=['generate p p > fre frebar'],sqrtS=200)`

In [11]: `launch_cross_section(a.MODEL)`

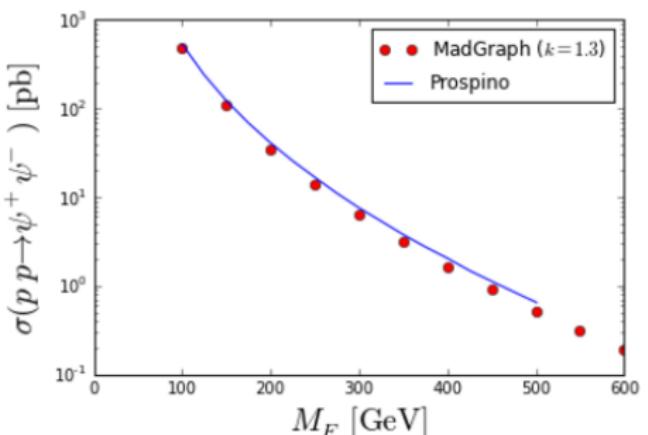
Out[11]: `(0.02713, 5e-05)`

In [12]: `from madgraph import *`
`a=hep(MODEL='SimplifiedDMSSSFDM')`

`generate_cross_section(a.MODEL,processes=['generate p p > fre frebar'],sqrtS=200)`
`v=a.vvv`
`lambda_1=0.26`
`lambda_SH=0`
`a.LHA.blocks['SPHENOINPUT'].entries[55]='0' # Calculate one loop`
`a.LHA.blocks['SPHENOINPUT'].entries[520]='0' # Write effective F`
`df=pd.DataFrame()`
`dfmin=100 #40`
`dfmax=600 #1E4`
`npoints=11`
`df_masses=np.linspace(dfmin,dfmax,npoints) #np.array([200]) 1E-4`
`for MF in df_masses:`
 `M_S=20`
 `MS2=M_S**2-a.vvv**2*lambda_SH`
 `Yse=1.`
 `devnull=commands.getoutput('rm -f SPheno.spc.%s' %a.MODEL)`
 `Ymu=0. #10**np.random.uniform(log10(1E-3),np.log10(np.pi))`
 `Ytau=0. #10`
 `a.LHA.blocks['MINPAR'][1]='%0.8E # lambda1' %lambda_1`
 `a.LHA.blocks['MINPAR'][2]='%0.8E # lambdaSH' %lambda_SH`
 `a.LHA.blocks['MINPAR'][4]='%0.8E # MS2' %MS2`
 `a.LHA.blocks['MINPAR'][51]='%0.8E # MF' %MF`

```
In [19]: plt.semilogy(df.MF,df.cs*1000*1.3,'ro',label='MadGraph ($k=1.3)')  
plt.semilogy(pr.mcl,pr.cs,label='Prospino')  
plt.xlabel(r'$M_F$ [GeV]',size=20)  
plt.ylabel(r'$\sigma(p p \rightarrow \psi^+ \psi^-)$ [pb]',size=20)  
plt.legend(loc='best')
```

Out[19]: <matplotlib.legend.Legend at 0x7f4243c7b290>



Results in Poster Session:

Marta Liliana Sánchez Pélaez

Unification: $SO(10)$

$$16_{F_i} = \begin{pmatrix} u_R^\dagger \\ u_R^\dagger \\ u_R^\dagger \\ u_L \\ u_L \\ u_L \\ d_L \\ d_L \\ d_L \\ d_L^\dagger \\ d_R^\dagger \\ d_R^\dagger \\ \nu_L \\ e_L \\ e_R^\dagger \\ N \end{pmatrix}_i \Rightarrow \mathcal{L}_{SM} \supset h \, 16_F \times 16_F \times 10_S + \text{h.c}$$



Not-susy $SO(10) \rightarrow SU(5) \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y \times Z_2$

Standard Model: Z_2 -even

Fermions: 16_F

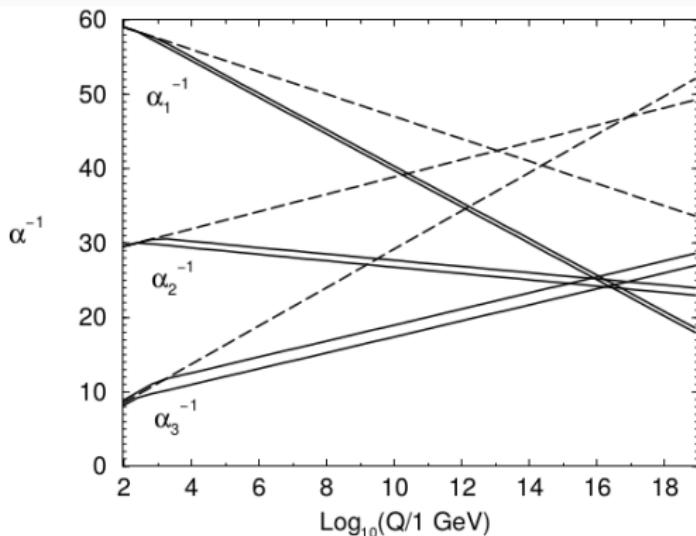
Scalars: $10_H, 45_H \dots$

New Z_2 -odd particles

$10_F, 45_F, \dots$

$16_H, \dots$

Lightest Odd Particle (LOP) may be a suitable dark matter candidate, and can improve gauge coupling unification



$$\text{Not-susy } SO(10) \rightarrow SU(5) \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y \times Z_2$$

Standard Model: Z_2 -even

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Scalars: $10_H, 45_H \dots$

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Lightest Odd Particle (LOP) may be a suitable dark matter candidate. and can improve gauge coupling unification

$SU(2)_L \times U(1)_Y$ representation	fermions even $SO(10)$ representations	scalars odd $SO(10)$ representations
1_0	$45, 54, 126, 210$	$16, 144$
$2_{\pm 1/2}$	$10, 120, 126, 210, 210'$	$16, 144$
3_0	$45, 54, 210$	144

$SU(3)_C : 3 (T), 6, 8 (\Lambda)$

$$m_{3_0} = 2.7 \text{ TeV}, \quad m_{\Lambda} \sim 10^{10} \text{ TeV}, \quad m_{\text{GUT}} \sim 10^{16} \text{ GeV}.$$

$$\text{Not-susy } SO(10) \rightarrow SU(5) \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y \times Z_2$$

Standard Model: Z_2 -even

Fermions: 16_F

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3_0	$45, 54, 210$	144

$SU(3)_C : 3 (T), 6, 8 (\Lambda)$

Split-SUSY like

arXiv:1509.06313 (C. Arbelaez, R. Longas, D.R, O. Zapata)

Not-susy $SO(10) \rightarrow SU(5) \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y \times Z_2$

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Fermions: 16_F

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$SU(3)_C : 3 (T), 6, 8 (\Lambda)$

Radiative hybrid seesaw (Parida
1106.4137) or 1509.06313

Partial Split-SUSY-like spectrum: bino-higgsino-wino

+

↓

$10'_H$ with fermion DM or,

$16_H, \dots$ with scalar DM

Not-susy $SO(10) \rightarrow SU(5) \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y \times Z_2$

Standard Model: Z_2 -even

Fermions: 16_F

Scalars: $10_H, 45_H \dots$

New Z_2 -odd particles

$10_F, 45_F, \dots$

$16_H, \dots$

Lightest Odd Particle (LOP) may be a suitable dark matter candidate. and can improve gauge coupling unification

$SU(2)_L \times U(1)_Y$ representation	fermions even $SO(10)$ representations	scalars odd $SO(10)$ representations
1_0	$45, 54, 126, 210$	$16, 144$
$2_{\pm 1/2}$	$10, 120, 126, 210, 210'$	$16, 144$
3_0	$45, 54, 210$	144
$SU(3)_C : [3 (T), 6, 8 (\Lambda)]$	1509.06313	

SUSY-like spectrum: bino-higgsino-wino

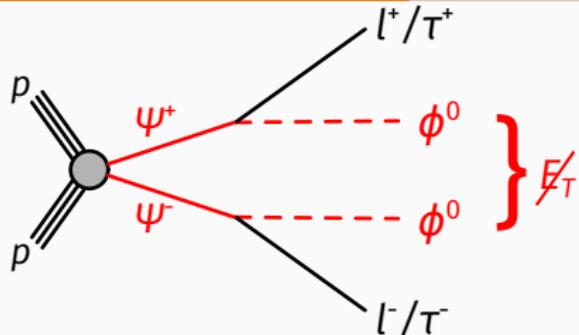
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\downarrow

$10'_H$ with fermion DM or,

$16_H, \dots$ with scalar DM

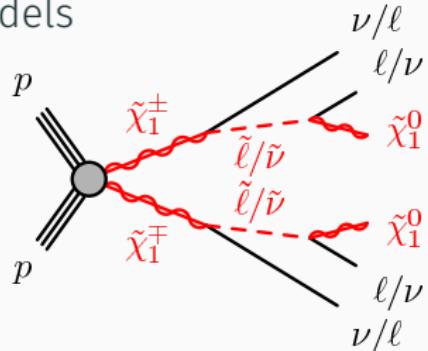
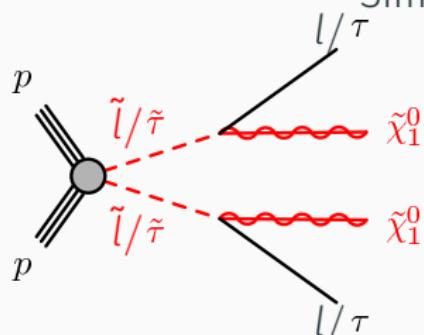
Dilepton plus transverse missing energy signal



SU(2)_L assignments:

$$\Psi = 1, 2(\Psi), 3(\Sigma), \quad \Phi = 1, 2, \text{ with } m_{\text{DM}} \sim m_h/2.$$

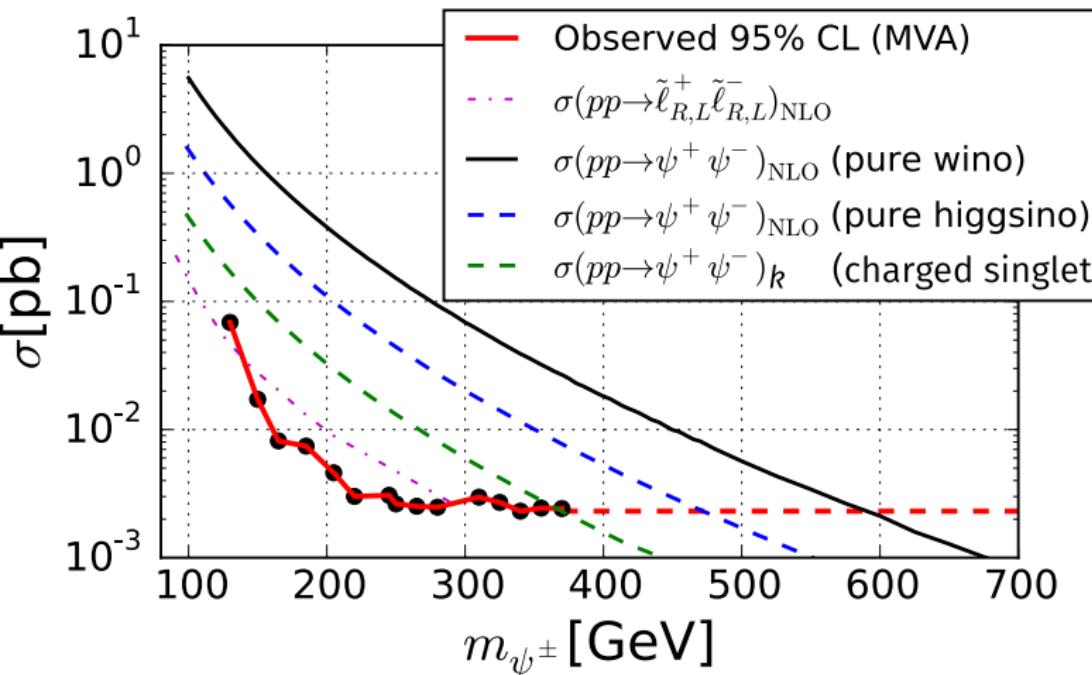
Simplified SUSY models



Smaller cross sections.

Intermediate states and smaller lepton p_T

$$m_{\phi^0} = 60 \text{ GeV}$$



Full analysis on flavor space: F. von der Pahlen, G. Palacio, DR,
O. Zapata arXiv:1605.01129 [PRD]

Singlet-Doublet-Triplet fermion dark-matter

The most general SO(10) invariant Lagrangian contains the following Yukawa terms

$$-\mathcal{L} \supset Y \mathbf{10}_F \mathbf{45}_F \mathbf{10}_H + M_{\mathbf{45}_F} \mathbf{45}_F \mathbf{45}_F + M_{\mathbf{10}_F} \mathbf{10}_F \mathbf{10}_F$$

Basis

$$\psi^0 = (N, \Sigma^0, \psi_L^0, (\psi_R^0)^\dagger)^T$$

$\mathcal{M}_{\psi^0} =$

$$\begin{pmatrix} M_N & 0 & -yc_\beta v/\sqrt{2} & ys_\beta v/\sqrt{2} \\ 0 & M_\Sigma & fc_\beta' v/\sqrt{2} & -fs_\beta' v/\sqrt{2} \\ -yc_\beta v/\sqrt{2} & fc_\beta' v/\sqrt{2} & 0 & -M_D \\ ys_\beta v/\sqrt{2} & -fs_\beta' v/\sqrt{2} & -M_D & 0 \end{pmatrix},$$

$$\mathbf{10}_F \rightarrow \psi_L, (\psi_R)^\dagger$$

$$\mathbf{45}_F \rightarrow \Sigma, \Lambda$$

$$\mathbf{45}'_F \rightarrow N$$

Singlet-Doublet-Triplet fermion dark-matter

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$$-\mathcal{L} \supset Y \mathbf{10}_F \mathbf{45}_F \mathbf{10}_H + M_{\mathbf{45}_F} \mathbf{45}_F \mathbf{45}_F + M_{\mathbf{10}_F} \mathbf{10}_F \mathbf{10}_F$$

Basis

$$\psi^0 = \left(N, \Sigma^0, \psi_L^0, (\psi_R^0)^\dagger \right)^T$$

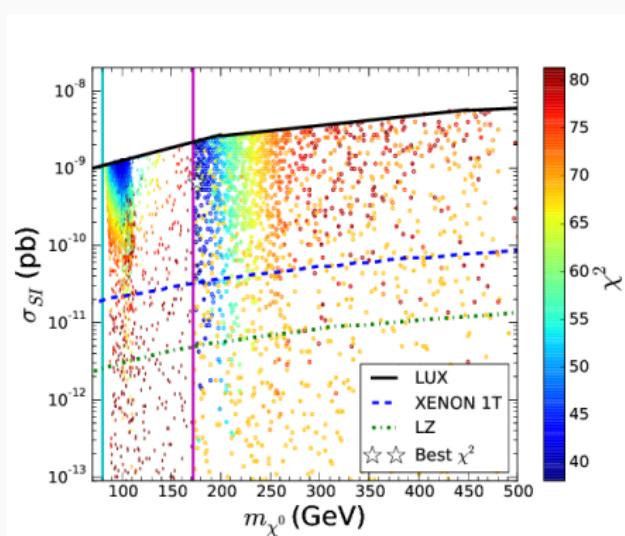
\mathcal{M}_{ψ^0} =

$$\begin{pmatrix} M_N & 0 & -yc_\beta v/\sqrt{2} & ys_\beta v/\sqrt{2} \\ 0 & M_\Sigma & fc_\beta v/\sqrt{2} & -fs_\beta v/\sqrt{2} \\ -yc_\beta v/\sqrt{2} & fc_\beta v/\sqrt{2} & 0 & -M_D \\ ys_\beta v/\sqrt{2} & -fs_\beta v/\sqrt{2} & -M_D & 0 \end{pmatrix},$$

$$\mathbf{10}_F \rightarrow \psi_L, (\psi_R)^\dagger$$

$$\mathbf{45}_F \rightarrow \Sigma, \Lambda$$

$$\mathbf{45}'_F \rightarrow N$$



S. Horiuchi, O. Macias, DR, A. Rivera, O. Zapata, 1602.04788
(JCAP)

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The most general SO(10) invariant Lagrangian contains the following Yukawa terms

$$-\mathcal{L} \supset Y \mathbf{10}_F \mathbf{45}_F \mathbf{10}_H + M_{\mathbf{45}_F} \mathbf{45}_F \mathbf{45}_F + M_{\mathbf{10}_F} \mathbf{10}_F \mathbf{10}_F$$

Basis

$$\psi^0 = \left(N, \Sigma^0, \psi_L^0, (\psi_R^0)^\dagger \right)^T$$

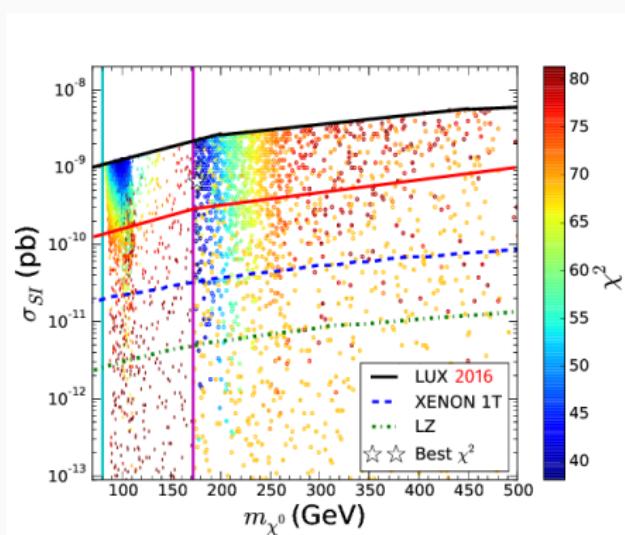
\mathcal{M}_{ψ^0} =

$$\begin{pmatrix} M_N & 0 & -yc_\beta v/\sqrt{2} & ys_\beta v/\sqrt{2} \\ 0 & M_\Sigma & fc_\beta v/\sqrt{2} & -fs_\beta v/\sqrt{2} \\ -yc_\beta v/\sqrt{2} & fc_\beta v/\sqrt{2} & 0 & -M_D \\ ys_\beta v/\sqrt{2} & -fs_\beta v/\sqrt{2} & -M_D & 0 \end{pmatrix},$$

$\mathbf{10}_F \rightarrow \psi_L, (\psi_R)^\dagger$

$$\mathbf{45}_F \rightarrow \Sigma, \Lambda$$

$$\mathbf{45}'_F \rightarrow N$$



S. Horiuchi, O. Macias, DR, A. Rivera, O. Zapata, 1602.04788
(JCAP)

Singlet-Doublet-Triplet fermion dark-matter

The most general SO(10) invariant Lagrangian contains the following Yukawa terms

$$-\mathcal{L} \supset Y_{10_F} 45_F 10_H + M_{45_F} 45_F 45_F + M_{10_F} 10_F 10_F + \mathcal{L}(10_\Phi) .$$

Basis

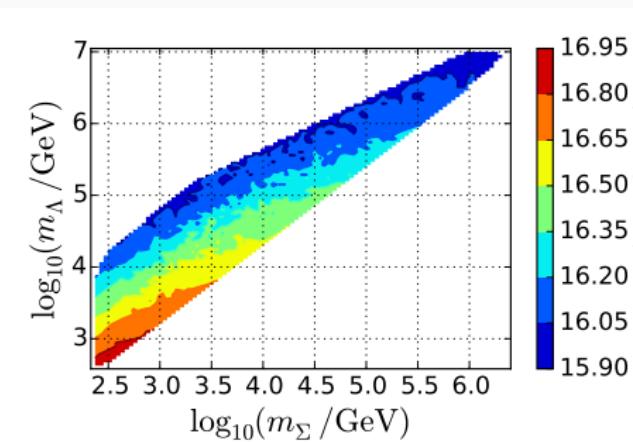
$$\psi^0 = \left(N, \Sigma^0, \psi_L^0, (\psi_R^0)^\dagger \right)^T$$

$\mathcal{M}_{\psi^0} =$

$$\begin{pmatrix} M_N & 0 & -yc_\beta v/\sqrt{2} & ys_\beta v/\sqrt{2} \\ 0 & M_\Sigma & fc_\beta' v/\sqrt{2} & -fs_\beta' v/\sqrt{2} \\ -yc_\beta v/\sqrt{2} & fc_\beta' v/\sqrt{2} & 0 & -M_D \\ ys_\beta v/\sqrt{2} & -fs_\beta' v/\sqrt{2} & -M_D & 0 \end{pmatrix},$$

$$45_F \rightarrow \Sigma, \Lambda$$

$$45'_F \rightarrow N$$

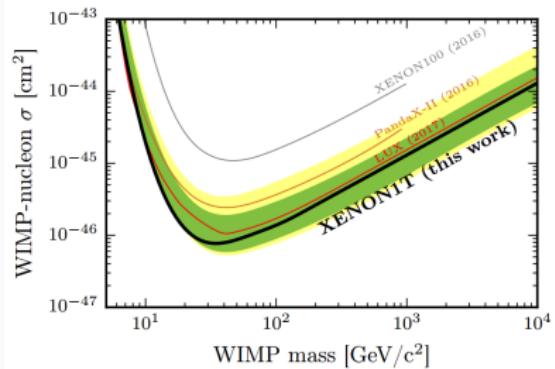


Split-SUSY: like $M_\Phi = 2$ TeV

Is the glass half empty or half full?

Tree-level SM-portal could be fully excluded in the near future

- Singlet scalar dark matter
- Inert doublet model
- Tree-level SM-portal dark matter ...

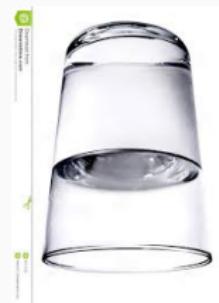


In this talk we explore

Is the glass half empty or half full?

Tree-level SM-portal could be fully excluded in the near future

- Singlet scalar dark matter
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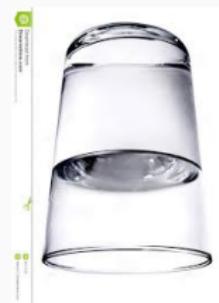
In this talk we explore

- Recover SM-portals in LR models

Is the glass half empty or half full?

Tree-level SM-portal could be fully excluded in the near future

- Singlet scalar dark matter
- Inert doublet model
- Tree-level SM-portal dark matter ...



In this talk we explore

- Recover SM-portals in LR models
- New portals in LR models

Left-Right symmetric realization

Singlet-doublet fermion dark matter

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Φ	1	(1, 2, 2, 0)	0	10
χ, χ^c	1	(1, 2, 2, 0)	1/2	10
N	1	(1, 1, 1, 0)	1/2	45

Table 1: The relevant part of the field content. Note that, the two fermion doublets χ and χ^c come from an only fermionic LR bidoublet. In the third column the relevant fields are characterized by their $SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$ quantum numbers while their SO(10) origin is specified in the fourth column.

Unification

m_{LR} (GeV)	$3_c 2_L 2_R 1_{B-L}$	m_G (GeV)
2×10^3	$\Phi_{1,2,2,0} + 2\Phi_{1,1,3,-2} + \Psi_{1,1,3,0} + \Phi_{1,1,3,0} + \Phi_{8,1,1,0}$	1.65×10^{16}
\vdots	\vdots	\vdots

Table 2: $\Delta_{L,R} = 2\Phi_{1,1,3,-2}$. m_{LR} and m_G are given in GeV.

Triplets

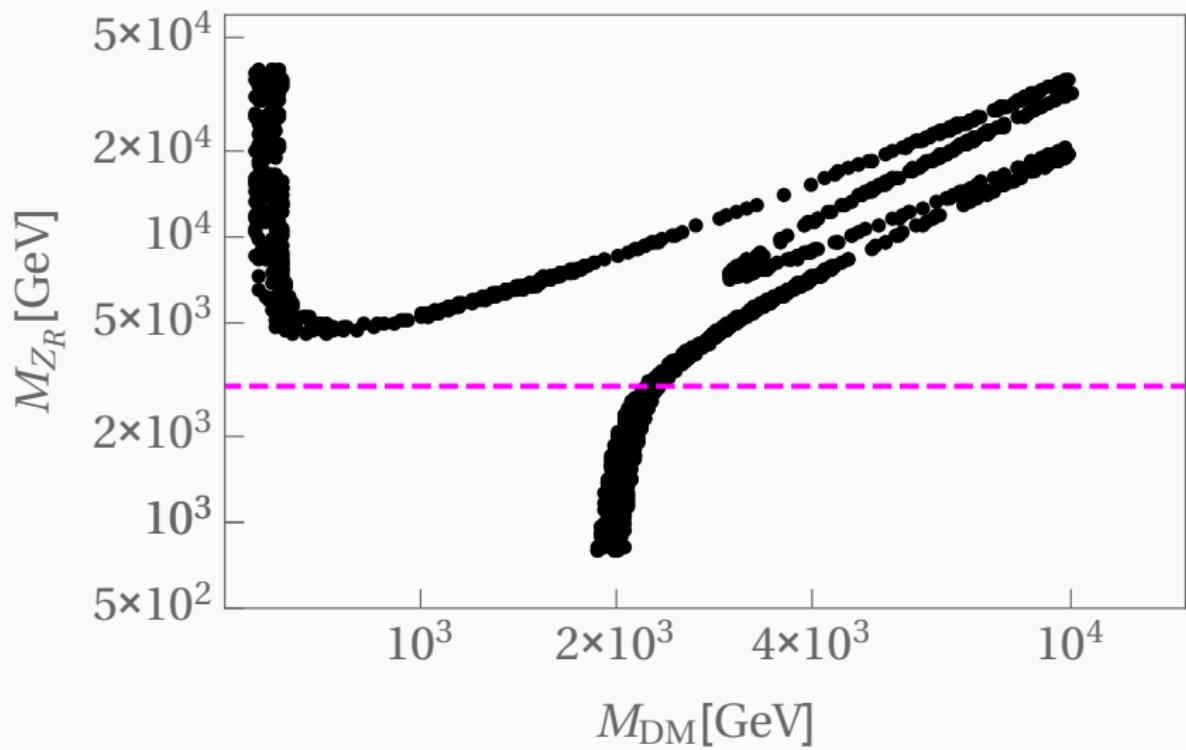
Minimal Left-Right Symmetric Standard Model

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	1/2	16
L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126

Left-singlet right-triplet DM

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	1/2	16
L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1130}	1	$(1, 1, 3, 0)$	1/2	45

$$\Omega h^2 = 0.1199 \pm 0.0027 \text{ at } 3\sigma$$



Mixed Left-singlet right-triplet DM

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	1/2	16
L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1130}	1	$(1, 1, 3, 0)$	1/2	45

Mixed Left-singlet right-triplet DM

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	1/2	16
L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1130}	1	$(1, 1, 3, 0)$	1/2	45
Ψ_{1132}	1	$(1, 1, 3, 2)$	1/2	126
Ψ_{113-2}	1	$(1, 1, 3, -2)$	1/2	$\overline{126}$

Setup

$$\Psi_{1132} = \begin{pmatrix} \Psi^+/\sqrt{2} & \Psi^{++} \\ \Psi^0 & -\Psi^+/\sqrt{2} \end{pmatrix}, \quad \bar{\Psi}_{113-2} = \begin{pmatrix} \Psi^-/\sqrt{2} & \bar{\Psi}^0 \\ \Psi^{--} & -\Psi^-/\sqrt{2} \end{pmatrix}. \quad (4)$$

$$L \supset M_{11} \operatorname{Tr}(\Psi_{1130} \Psi_{1130}) + M_{23} \operatorname{Tr}(\Psi_{1132} \bar{\Psi}_{113-2}) \\ + \lambda_{13} \operatorname{Tr}(\Delta_R \bar{\Psi}_{113-2} \Psi_{1130}) + \lambda_{12} \operatorname{Tr}(\Delta_R^\dagger \Psi_{1132} \Psi_{1130}), \quad (5)$$

$$\tan \gamma = \frac{\lambda_{13}}{\lambda_{12}}, \quad \lambda = \sqrt{\lambda_{12}^2 + \lambda_{13}^2}. \quad (6)$$

Blind spot at

$$M_{23} \sin 2\gamma - M_{\text{DM}} = 0 \quad (7) \quad 30$$

Proper relic density scan

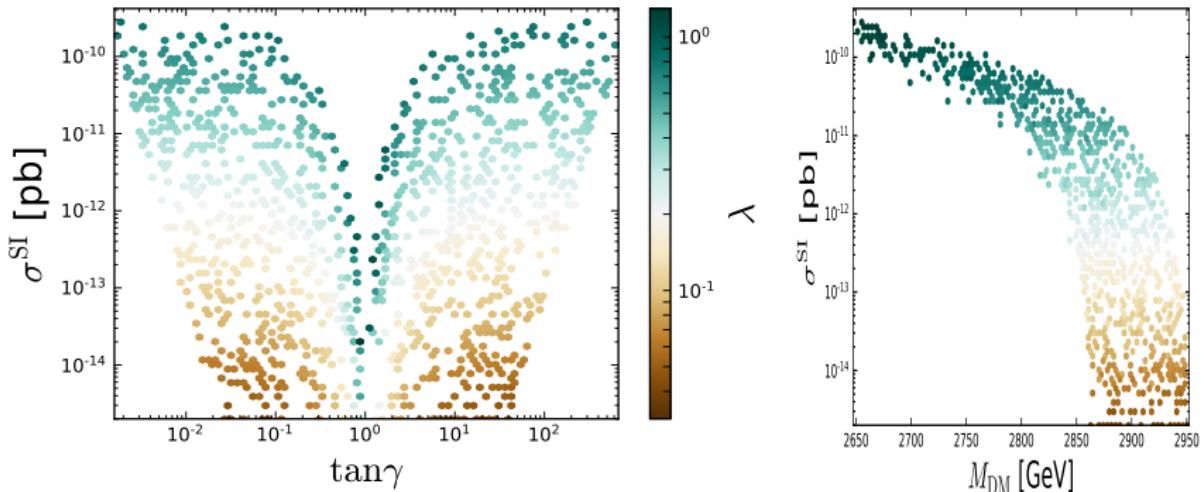


Figure 2: $M_{11} = 50$ TeV $2.7 < M_{23}/\text{TeV} < 3.1$

(Right: $\tan \gamma > 5$)

$$\Omega h^2 = 0.1199 \pm 0.0027 \text{ at } 3\sigma$$

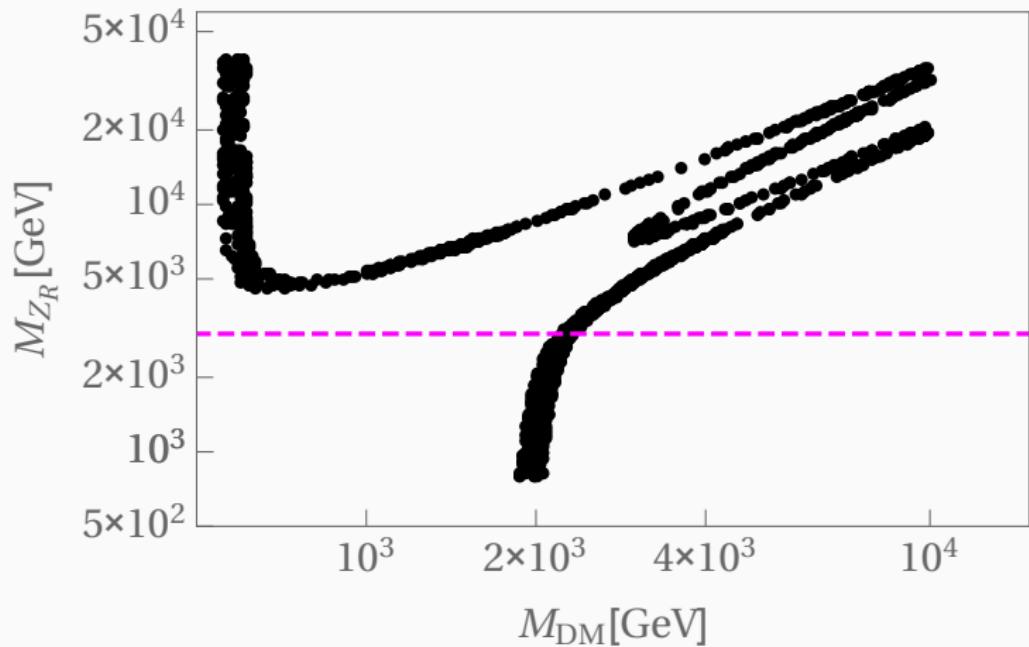


Figure 3:

$$\Omega h^2 = 0.1199 \pm 0.0027 \text{ at } 3\sigma$$

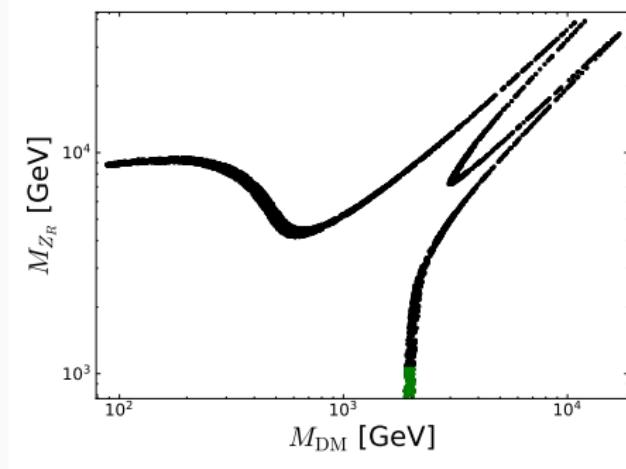


Figure 3: Proper relic density scan: $v_R : [2, 50]$ TeV, $M_{23} : [0.2, 50]$ TeV, $M_{11} : 50$ TeV, $\tan \gamma = -1$ and $\lambda = 0.14$.

Direct detection cross section

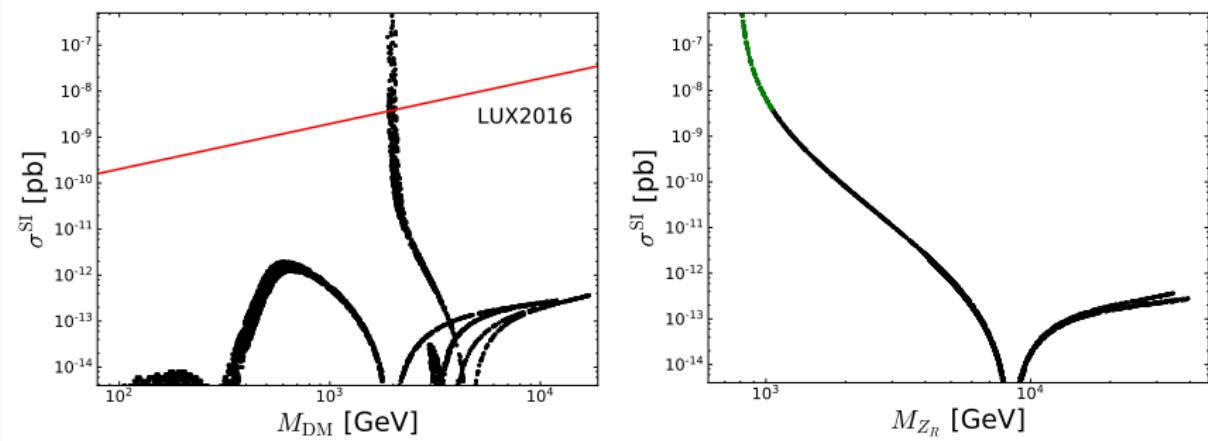


Figure 4: $v_R : [2, 50]$ TeV, $M_{23} : [0.2, 50]$ TeV, $M_{11} : 50$ TeV, $\tan \gamma = -1$ and $\lambda = 0.14$.

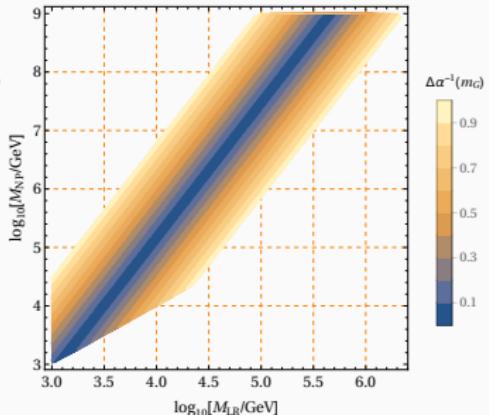
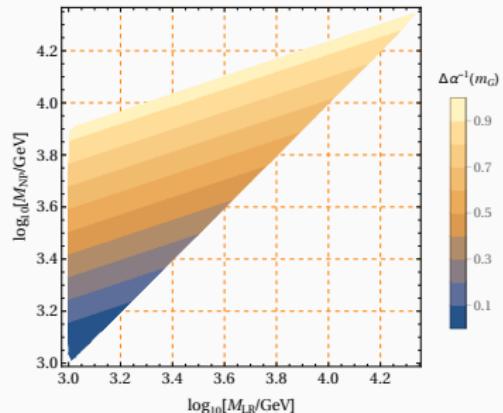
Unification

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	1/2	16
L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1130}	1	$(1, 1, 3, 0)$	1/2	45
Ψ_{1132}	1	$(1, 1, 3, 2)$	1/2	126
Ψ_{113-2}	1	$(1, 1, 3, -2)$	1/2	$\overline{126}$

Unification

Field	Multiplicity	$3c2_L2_R1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	1/2	16
L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1130}	1	$(1, 1, 3, 0)$	1/2	45
Ψ_{1132}	1	$(1, 1, 3, 2)$	1/2	126
Ψ_{113-2}	1	$(1, 1, 3, -2)$	1/2	$\overline{126}$
Ψ_{1310}	1	$(1, 3, 1, 0)$	1/2	45
Ψ_{8110}	1	$(1, 1, 8, 0)$	1/2	45
$\Psi_{321\frac{1}{3}}$	1	$(3, 2, 1, 1/3)$	1/2	16
$\Psi_{321-\frac{1}{3}}$	1	$(1, 2, 3, -1/3)$	1/2	$\overline{16}$

Unification quality



Conclusions

In addition to accommodate usual simplified dark matter models, Left-right symmetric standard models have additional DM portals:

New Δ_R portal for direct detection of left-singlet right-triplet mixed dark matter, in companion with left-singlets charged and doubly charged fermions.

Next: Search for them in compressed spectra scenarios at the LHC

Thanks!