Radiative seesaw and baryogenesis

with gauged Lepton number

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Focus on

arXiv:xxxx.xxxxx

In collaboration with

W. Novelo, L. Leite, O. Peres (UNICAMP)

Model building

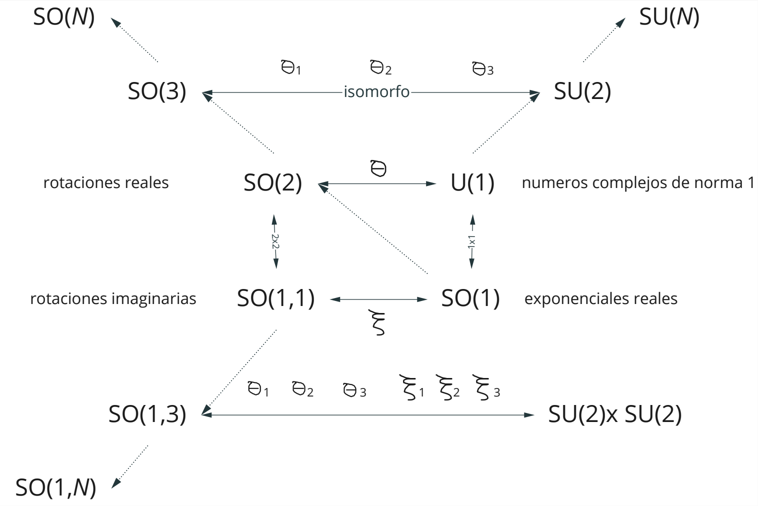


Figure 1: Grupos de Lie

Considere el generador 1×1

$$K = -i\,, \tag{1}$$

que genera el elemento del grupo dilaton, SO(1), [**?**] $R(\xi)$

$$\lambda(\xi) = \mathrm{e}^{\xi}\,, \tag{2}$$

que corresponde simplemente al grupo de las exponenciales reales. Un número real puede sufrir una transformación

$$x \to x' = \mathrm{e}^{\xi}\, x\,, \tag{3}$$

que corresponde a su vez a un boost por la cantidad e^{ξ}. Podemos definir un producto escalar invariante como la división de números reales tal que

$$x \cdot y \to x' \cdot y' \equiv \frac{x'}{y'} = \frac{\mathrm{e}^{\xi}\, x}{\mathrm{e}^{\xi}\, y} = \frac{x}{y} = x \cdot y\,. \tag{4}$$

Queremos obtener una representación 2×2 del álgebra

$$K^2 = -\mathbf{1}\,, \tag{5}$$

donde K es el único generador. Para hallar una representación de esta álgebra en términos de matrices 2×2 considere el generador

$$K = \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix}. \tag{6}$$

que genera un elemento del grupo $\mathrm{SO}(1,1)$ con parámetro ξ

$$\Lambda = \exp\left(i\xi K\right)\,. \tag{7}$$

Para realizar la expansión de Taylor, considere

$$
\begin{aligned}
K^0 &= \mathbf{1}_{2\times 2}\,, \qquad K = \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix}, \quad K^2 = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix}, \qquad K^3 = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}, \ldots \\
K^{2n} &= (-1)^n \mathbf{1}_{2\times 2}\,, \qquad K^{2n+1} = (-1)^n \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix}.
\end{aligned}
$$

Entonces,

$$
\begin{aligned}
\Lambda = \exp\left(i\xi K\right) &= \sum_{n=0}^{\infty} \frac{\left(i\xi K\right)^n}{n!} \\
&= \sum_{n=0}^{\infty} (i)^{2n} \frac{\left(\xi K\right)^{2n}}{2n!} + \sum_{n=0}^{\infty} (i)^{2n+1} \frac{\left(\xi K\right)^{2n+1}}{(2n+1)!} \\
&= \sum_{n=0}^{\infty} (-1)^n \frac{\xi^{2n}}{2n!} (-1)^n \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \sum_{n=0}^{\infty} i(-1)^n \frac{\xi^{2n+1}}{(2n+1)!} (-1)^n \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix} \\
&= \sum_{n=0}^{\infty} \frac{\xi^{2n}}{2n!} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \sum_{n=0}^{\infty} \frac{\xi^{2n+1}}{(2n+1)!} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\
&= \begin{pmatrix} \cosh\xi & 0 \\ 0 & \cosh\xi \end{pmatrix} + \begin{pmatrix} 0 & \sinh\xi \\ \sinh\xi & 0 \end{pmatrix} \\
&= \begin{pmatrix} \cosh\xi & \sinh\xi \\ \sinh\xi & \cosh\xi \end{pmatrix}, \qquad (8)
\end{aligned}
$$

Podemos entonces definir el grupo $SO(1,1)$ como el grupo de las matrices 2×2 que satisfacen la condición

$$\Lambda^T g \Lambda = g\,, \tag{9}$$

donde

$$g = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \tag{10}$$

Bajo una transformación de Lorentz.

$$x^{\mu} \to x'^{\mu} = \Lambda^{\mu}{}_{\nu} x^{\nu}. \tag{11}$$

Introducimos ahora un cuadrivector que lleva intrínsicamente el índice abajo

$$\partial_{\mu} = \frac{\partial}{\partial x^{\mu}} = \left(\frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) = (\partial_0, \boldsymbol{\nabla}). \tag{12}$$

Las propiedades de transformación para ∂_{μ}

$$\begin{aligned} \left(\Lambda^{-1}\right)^{\mu}{}_{\alpha} x'^{\alpha} =& \left(\Lambda^{-1}\right)^{\mu}{}_{\alpha} \Lambda^{\alpha}{}_{\nu} x^{\nu} \\ =& \delta^{\mu}_{\nu} x^{\nu} \\ =& x^{\mu}, \end{aligned} \tag{13}$$

$$\frac{1}{x'^{\nu}} = \left(\Lambda^{-1}\right)^{\mu}{}_{\nu} \frac{1}{x^{\mu}}\,, \tag{14}$$

o

$$\frac{1}{x'^{\mu}} = \left(\Lambda^{-1}\right)^{\nu}{}_{\mu} \frac{1}{x^{\nu}}\,, \tag{15}$$

de modo que la transformación de Lorentz para $\partial_\mu = \partial/\partial x^\mu$, es

$$\begin{aligned} \frac{\partial}{\partial x'^{\mu}} &= \left(\Lambda^{-1}\right)^{\nu}{}_{\mu} \frac{\partial}{\partial x^{\nu}} \\ \partial'_{\mu} &= \left(\Lambda^{-1}\right)^{\nu}{}_{\mu} \partial_{\nu}\,. \end{aligned} \tag{16}$$

De esta manera, un producto escalar invariante entre la cuadriderivada y un cuadrivector de Lorentz $A^\mu(x)$ es

$$\partial_\mu A^\mu \to \partial'_\mu A'^\mu \to \partial_\mu A^\mu\,. \tag{17}$$

Nombre	Símbolo	SU(N)
N-plete escalar	Ψ	$U\Psi$
anti-N-plete escalar	$\Psi^\dagger$	$\Psi^\dagger U^\dagger$

Nombre	Símbolo	Lorentz
fotón	A^μ	$\Lambda^\mu{}_\nu A^\nu$
derivada	∂_μ	$\partial_\nu \left(\Lambda^{-1}\right)^\nu{}_\mu$

Table 1: Productos escalares: $\Psi^\dagger\Psi$, $\partial_\mu A^\mu\,, \quad A^\nu A_\nu\,, \quad \partial_\mu\partial^\mu$

donde, $g_{\alpha\beta} = \Lambda^\mu{}_\alpha\, g_{\mu\nu}\Lambda^\nu{}_\beta$, $g^{\mu\nu} = \left(\Lambda^{-1}\right)^\mu{}_\alpha\, g^{\alpha\beta}\left(\Lambda^{-1}\right)^\nu{}_\beta$.

Nombre	Símbolo	Lorentz	$U(1)$
e_L: electrón izquierdo	ξ_α	$S_\alpha{}^\beta \xi_\beta$	$e^{i\theta}\xi_\alpha$
$(e_R)^\dagger$: positrón izquierdo	η^α	$\eta^\beta \left[S^{-1}\right]_\beta{}^\alpha$	$\eta^\alpha\, e^{-i\theta}$
$(e_L)^\dagger$: positrón derecho	$(\xi_\alpha)^\dagger = \xi^\dagger_{\dot\alpha}$	$\xi^\dagger_{\dot\beta}\left[S^\dagger\right]^{\dot\beta}{}_{\dot\alpha}$	$\xi^\dagger_{\dot\alpha} e^{-i\theta}$
e_R: electrón derecho	$(\eta^\alpha)^\dagger = \eta^{\dagger\,\dot\alpha}$	$\left[\left(S^{-1}\right)^\dagger\right]^{\dot\alpha}{}_{\dot\beta}\eta^{\dagger\,\dot\beta}$	$e^{i\theta}\eta^{\dagger\dot\alpha}$

Table 2: Definición de transformaciones de Lorentz

Productos escalares

- Escalares de Majorana: $\xi^\alpha\xi_\alpha + \xi^\dagger_{\dot\alpha}\xi^{\dagger\dot\alpha}$, $\eta^\alpha\eta_\alpha + \eta^\dagger_{\dot\alpha}\eta^{\dagger\dot\alpha}$.
- Escalar de Dirac: $\eta^\alpha\xi_\alpha + \xi^\dagger_{\dot\alpha}\eta^{\dagger\,\dot\alpha}$.
- Escalar subgrupo SL(2, C) pero vector bajo SO(1, 3): $S^\dagger\overline{\sigma}^\mu S = \Lambda^\mu{}_\nu\overline{\sigma}^\nu$

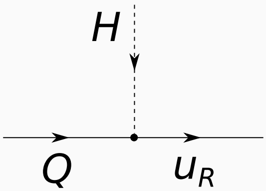
Campos	Lorentz	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$
Q	ξ_α^1	$\mathbf{3}$	$\mathbf{2}$	$1/6$
L	ξ_α^2	$\mathbf{1}$	$\mathbf{2}$	$-1/2$
$\left(u_R^-\right)^\dagger$	η_1^α	$\overline{\mathbf{3}}$	$\mathbf{1}$	$-2/3$
$\left(d_R^-\right)^\dagger$	η_2^α	$\overline{\mathbf{3}}$	$\mathbf{1}$	$1/3$
$\left(e_R^-\right)^\dagger$	η_3^α	$\mathbf{1}$	$\mathbf{1}$	1
H	-	$\mathbf{1}$	$\mathbf{2}$	$1/2$

Table 3: Campos fundamentales del modelo estándar

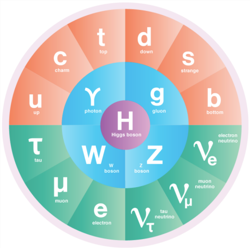
como por ejemplo,

$$(u_R)^\dagger \, Q \cdot H \, , \tag{18}$$

Que se puede representar con la Ley De Kircchoff:



Dark sectors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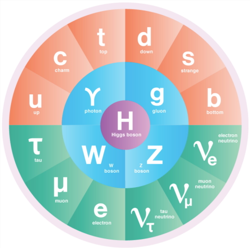
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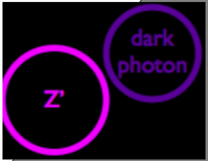
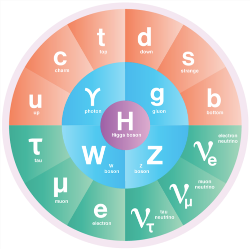

dark
photon
Z'



$F_{\mu\nu}$ $V^{\mu\nu}$

Local $U(1)_\mathcal{X}$

$$\mathcal{L} = -\frac{1}{4} V_{\mu\nu} V^{\mu\nu} + i \sum_i \chi_i^\dagger \not{\mathcal{D}} \chi_i$$

$$-h\left(\chi_1 \chi_2 \Phi + \mathrm{h.c}\right)$$

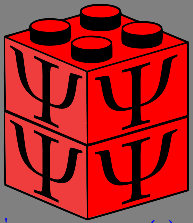
Anomalons: SM-singlet Dirac fermion

dark matter $m_\Psi = h\langle\Phi\rangle$

LHC production:

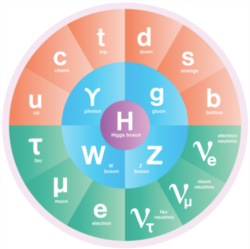
Gauged Symmetry: $\mathcal{X} \to B$: $q\bar{q} \to Z' \to \text{jets}$

Gauged Symmetry: $\mathcal{X} \to L$:



$$\overline{\Psi}\Psi = \chi_1\chi_2 + \chi_1^\dagger\chi_2^\dagger \to \chi_\alpha\chi_\beta\Phi^{(*)}\,, \qquad \alpha = 1, \ldots N' \to N' > 4$$

Local $U(1)_\mathcal{X}$



$$\mathcal{L} = -\frac{1}{4} V_{\mu\nu} V^{\mu\nu} + i \sum_i {\chi_i}^\dagger \not{\mathcal{D}} \chi_i$$

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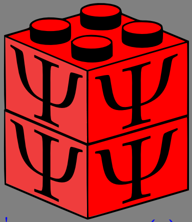

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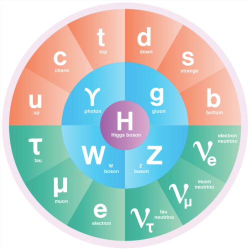

$F_{\mu\nu}$ $V^{\mu\nu}$

Local $U(1)_{\mathcal{X}}$

$$\mathcal{L} = -\frac{1}{4} V_{\mu\nu} V^{\mu\nu} + i \sum_i \chi_i^{\dagger} \not{D} \chi_i$$

$$-h\left(\chi_1 \chi_2 \Phi + \text{h.c}\right)$$

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dark matter $m_\Psi = h\langle\Phi\rangle$

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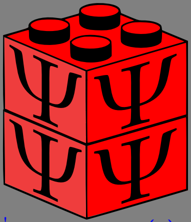
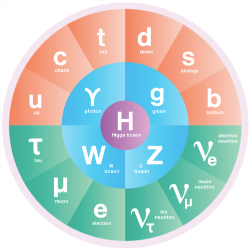


multi-component

dark matter

$\overline{\Psi}\Psi = \chi_1\chi_2 + \chi_1^{\dagger}\chi_2^{\dagger} \to \chi_\alpha \chi_\beta \Phi^{(*)}$, $\alpha = 1, \ldots N' \to N' > 4$



Local $U(1)_{\mathcal{X}}$

$$\mathcal{L} = -\frac{1}{4}V_{\mu\nu}V^{\mu\nu} + i\sum_i \chi_i^\dagger \not{D}\chi_i$$
$$-y(\chi_1\chi_2 S + \text{h.c})$$

Anomalons: SM-singlet Dirac fermion

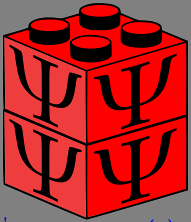
CP violation Yukawa y

LHC production:

$F_{\mu\nu}\quad V^{\mu\nu}$

Gauged Symmetry: $\mathcal{X} \to B$: $q\bar{q} \to Z' \to$ jets

Gauged Symmetry: $\mathcal{X} \to L$:



$$\overline{\Psi}\Psi = \chi_1\chi_2 + \chi_1^\dagger\chi_2^\dagger \to \chi_\alpha\chi_\beta\Phi^{(*)}\,, \qquad \alpha = 1,\dots N' \to N' > 4$$

Standard model extended with $\mathrm{U}(1)_{\mathcal{X}=L \text{ or } B}$ gauge symmetry

Fields	$\mathrm{SU}(2)_L$	$\mathrm{U}(1)_Y$	$\mathrm{U}(1)_{\mathcal{X}=B \text{ or } L}$
$Q_i^\dagger$	$\mathbf{2}$	$-1/6$	Q
d_{Ri}	$\mathbf{1}$	$-1/2$	d
u_{Ri}	$\mathbf{1}$	$+2/3$	u
$L_i^\dagger$	$\mathbf{2}$	$+1/2$	L
e_{Ri}	$\mathbf{1}$	-1	e
H	$\mathbf{2}$	$1/2$	$h=0$
χ_α	$\mathbf{1}$	0	z_α
$(L'_L)^\dagger$	$\mathbf{2}$	$1/2$	$-x'$
L''_R	$\mathbf{2}$	$-1/2$	x''
e'_R	$\mathbf{1}$	-1	x'
$(e''_L)^\dagger$	$\mathbf{1}$	1	$-x''$
Φ	$\mathbf{1}$	0	ϕ
S	$\mathbf{1}$	0	s

Table 4: A minimal set of new fermion content: $L=e=0$ for $\mathcal{X}=B$. Or $Q=u=d=0$ for $\mathcal{X}=L$. $i=1,2,3,\ \alpha=1,2,\ldots,N'$

Effective Dirac neutrino mass operator

$$\chi_1 \to \nu_{R1}, \cdots, \chi_{N_\nu} \to \nu_{R\,N_\nu}, \qquad 2 \leq N_\nu \leq 3\,, \tag{19}$$

$$\mathcal{L}_{\text{eff}} = h_\nu^{\alpha i}\,(\nu_{R\alpha})^\dagger\,\epsilon_{ab}\,L_i^a\,H^b\left(\frac{\Phi^*}{\Lambda}\right)^\delta + \text{H.c.}, \qquad \text{with } i = 1, 2, 3\,,$$

S is the complex singlet scalar responsible for the SSB of the anomaly-free gauge symmetry with D or X-charge

$$\phi = -(\nu + L)/\delta\,, \tag{20}$$

Anomaly cancellation I

The anomaly-cancellation conditions on $[\mathrm{SU}(3)_c]^2\,\mathrm{U}(1)_X$, $[\mathrm{SU}(2)_L]^2\,\mathrm{U}(1)_X$, $[\mathrm{U}(1)_Y]^2\,\mathrm{U}(1)_X$, allow us to express three of the X-charges in terms of the others

$$u = -e - \frac{2}{3}L - \frac{1}{9}\left(x' - x''\right), \quad d = e + \frac{4}{3}L - \frac{1}{9}\left(x' - x''\right), \quad Q = -\frac{1}{3}L + \frac{1}{9}\left(x' - x''\right), \tag{21}$$

while the $[\mathrm{U}(1)_X]^2\,\mathrm{U}(1)_Y$ anomaly condition reduces to

$$(e + L)(x' - x'') = 0\,. \tag{22}$$

- Previously: $x' = x''$
- We choose instead ($h = 0$):

$$e = -L\,, \tag{23}$$

so that (L is still a free parameter)

$$Q = -u = -d = -\frac{1}{3}L + \frac{1}{9}\left(x' - x''\right). \tag{24}$$

If $B = 0 \rightarrow \mathrm{U}(1)_L$

Anomaly cancellation II

The gravitational anomaly, $[\mathrm{SO}(1,3)]^2\,\mathrm{U}(1)_Y$, and the cubic anomaly, $[\mathrm{U}(1)_X]^3$, can be written as the following system of Diophantine equations, respectively,

$$\sum_{\alpha=1}^{N} z_\alpha = 0\,, \qquad \sum_{\alpha=1}^{N} z_\alpha^3 = 0\,, \tag{25}$$

where $N = N' + 5$ and

$$z_{N'+1} = -x'\,, \qquad z_{N'+2} = x''\,, \qquad z_{N'+2+i} = L\,, \quad i = 1, 2, 3 \tag{26}$$

$\rightarrow$

$$9Q = -\sum_{\alpha=N'+1}^{N'+5} z_\alpha = -x' + x'' + L + L + L\,, \tag{27}$$

$Q = 0 \not\to \mathrm{U}(1)_L$

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September 24, 2021

Dataset Open Access

Set of N integers between -30 and 30 with sum and cubic sum up to zero for 4<N<13

Diego Restrepo

Anomalies

Solutions obtained with the python package: anomalies based on the method to find anomaly free solutions of the standard model extended with an Abelian Dark Symmetry with *N* right-handed singlet chiral fields described in arXiv:1905.13729 [PRL]:

Data scheme

- 'l': integer lists → input to obtain the 'solution' by using the anomalies package
- 'k': integer lists → input to obtain the 'solution' by using hte anomalies package
- 'solution': list → of integers, z_i which satisfy $\sum_{i=1}^{N} z_i = 0$ and $\sum_{i=1}^{N} z_i^3 = 0$.
- 'n': integer → number of integers in 'solution', *N*.

USAGE

```
#Example of JSON file usage in Python with pandas (see also json module)
>>> import pandas as pd
>>> df=pd.read_json('solutions.json')
>>> df[:2]
          l         k              solution  gcd  n
0    [1, 2]  [0, -3]   [1, 5, -7, -8, 9]    1  5
1  [-2, -1]  [0, -1]  [2, 4, -7, -9, 10]    1  5
```

Data:

390074 solutions with $5 \leq N \leq 12$ integers until `|32|` [JSON]

Publication date:
September 24, 2021

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Keyword(s):
Anomaly free Diophantine equations Abelian symmetry Gauge Symmetry

Versions

Version 1 Sep 24, 2021
10.5281/zenodo.5526707

$\mathrm{U}(1)_L$ selection

- $B = 0 \to L = -3$ $(-x' + x'' + 3L = 0)$

$$(-3, -3, -3, -6, -6, -6, 2, 7, 4, 4, 5, 5)$$

$U(1)_L$ selection

- $B = 0 \to L = -3$ $(-x' + x'' + 3L = 0)$
- Effective neutrino mass $\nu = -6$:
  $\phi = -(\nu + L) = 9$

$$(-3, -3, -3, -6, -6, -6, 2, 7, 4, 4, 5, 5)$$

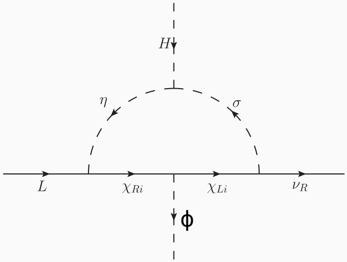
$\mathrm{U}(1)_L$ selection

- $B = 0 \to L = -3$ $(-x' + x'' + 3L = 0)$
- Effective neutrino mass $\nu = -6$:
  $\phi = -(\nu + L) = 9$
- Electroweak-scale vector-like fermions:
  $(L'_L)^\dagger L''_R \Phi^* \to x' = -2,\ x'' = 7$

$$(-3, -3, -3, -6, -6, -6, 2, 7, 4, 4, 5, 5)$$

$U(1)_L$ selection

- $B = 0 \to L = -3 \; (-x' + x'' + 3L = 0)$
- Effective neutrino mass $\nu = -6$:
  $\phi = -(\nu + L) = 9$
- Electroweak-scale vector-like fermions:
  $(L'_L)^\dagger L''_R \Phi^* \to x' = -2, \; x'' = 7$
- At least two generations of heavy Dirac-fermionic DM:
  $(\chi_{Li})^\dagger \chi_{Ri} \Phi^* \to z_3 = 4, \; z_4 = 5$

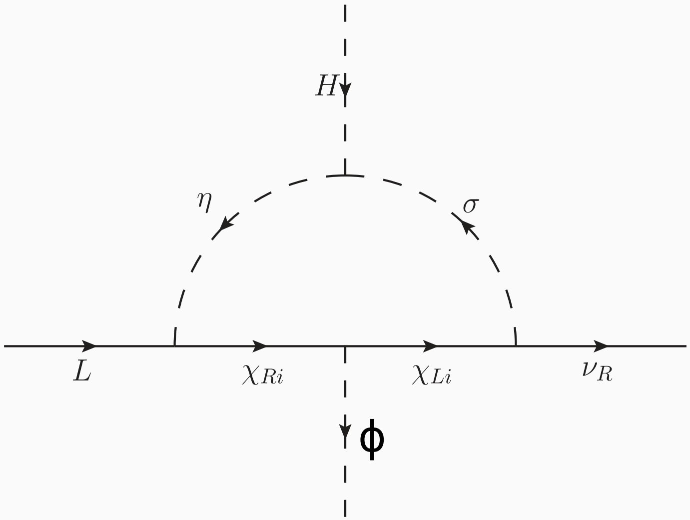


$$(-3, -3, -3, -6, -6, -6, 2, 7, 4, 4, 5, 5)$$

Unique solution from $\sim 400,000$!

Scotogenic realization

Any realization which does not affect anomaly cancellation is allowed



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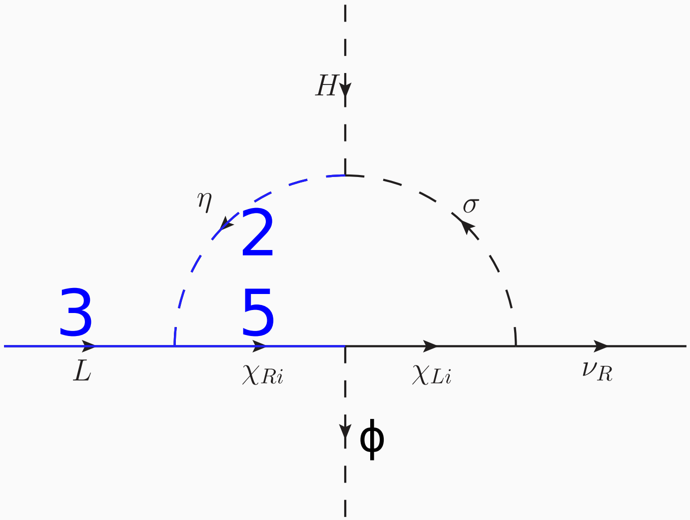
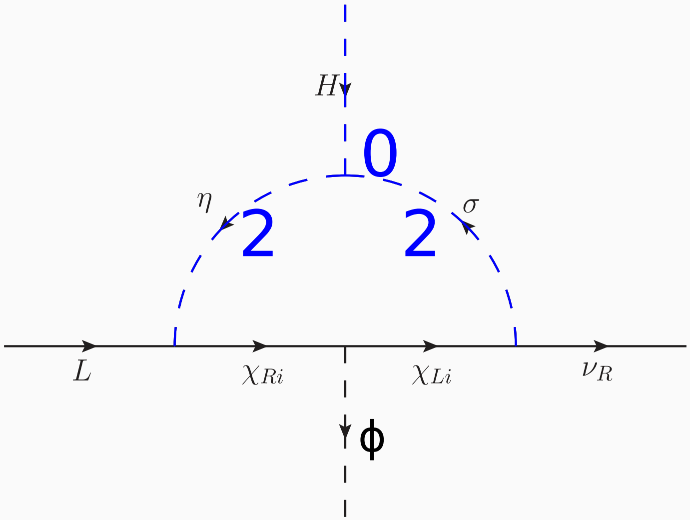


Any realization which does not affect anomaly cancellation is allowed



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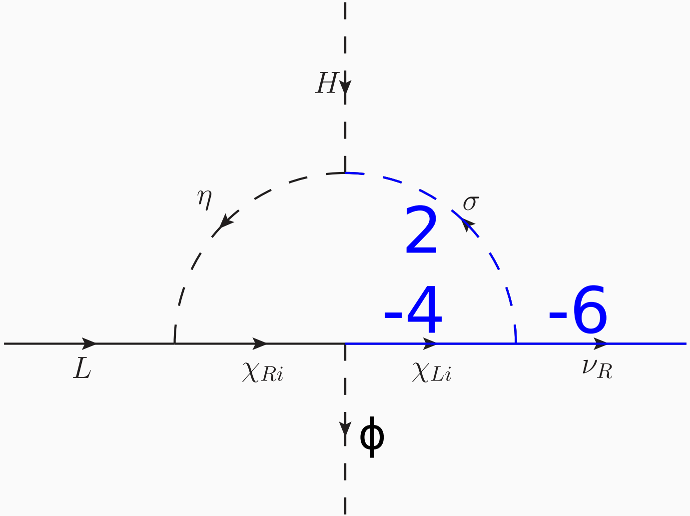

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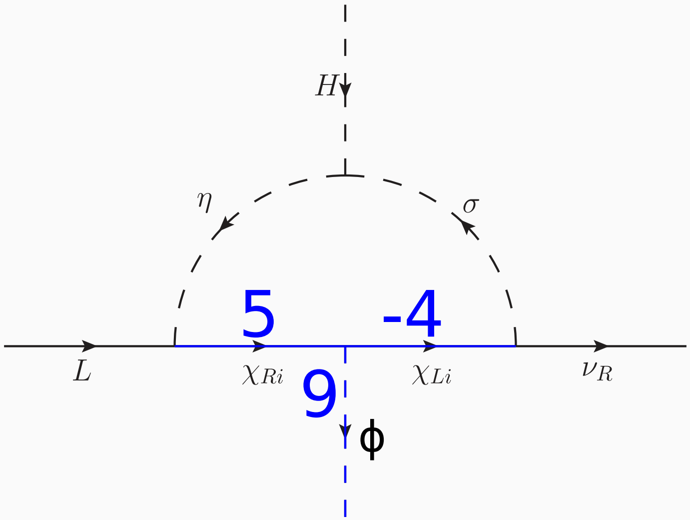


Full $\mathrm{U}(1)_L$ model

Field	$SU(2)_L$	$U(1)_Y$	$\mathrm{U}(1)_B$
u_{Ri}	$\mathbf{1}$	$2/3$	$u=0$
d_{Ri}	$\mathbf{1}$	$-1/3$	$d=0$
$(Q_i)^\dagger$	$\mathbf{2}$	$-1/6$	$Q=0$
$(L_i)^\dagger$	$\mathbf{2}$	$1/2$	$L=1$
e_R	$\mathbf{1}$	-1	$e=-1$
$\left(L'_L\right)^\dagger$	$\mathbf{2}$	$1/2$	$-x'=-2/3$
e'_R	$\mathbf{1}$	-1	$x'=2/3$
L''_R	$\mathbf{2}$	$-1/2$	$x''=-7/3$
$\left(e''_L\right)^\dagger$	$\mathbf{1}$	1	$-x''=7/3$
$\nu_{R,i}$	$\mathbf{1}$	0	2
χ_R	$\mathbf{1}$	0	$-5/3$
$(\chi_L)^\dagger$	$\mathbf{1}$	0	$-4/3$
H	$\mathbf{2}$	$1/2$	0
Φ	$\mathbf{1}$	0	-3
η	$\mathbf{2}$	$1/2$	$2/3$
σ	$\mathbf{1}$	0	$2/3$

Table 5: $i=1,2,3$ normalized Lepton number charges with a global factor $-1/3$.

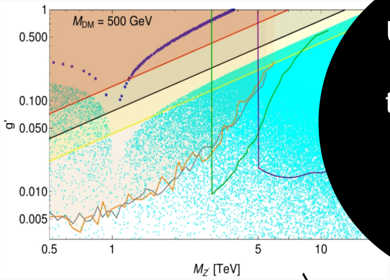
Electroweak baryogenesis

Problems

- Standard model (SM) $m_h \sim 40$ GeV. 😞
- Beyond the SM: Source of CP contains fields charged under SM $\rightarrow$ too large electric dipole moments 😩

Dark sectors

- Inert SM-singlet complex scalar field which acquires vev with temperature to have strong electroweak phase transition
- CP violation (CPV) triggered in dark sectors through SM gauge singlets
  $\rightarrow$ CPV Yukawa between SM-singlet complex scalar and SM-singlet quiral fermions



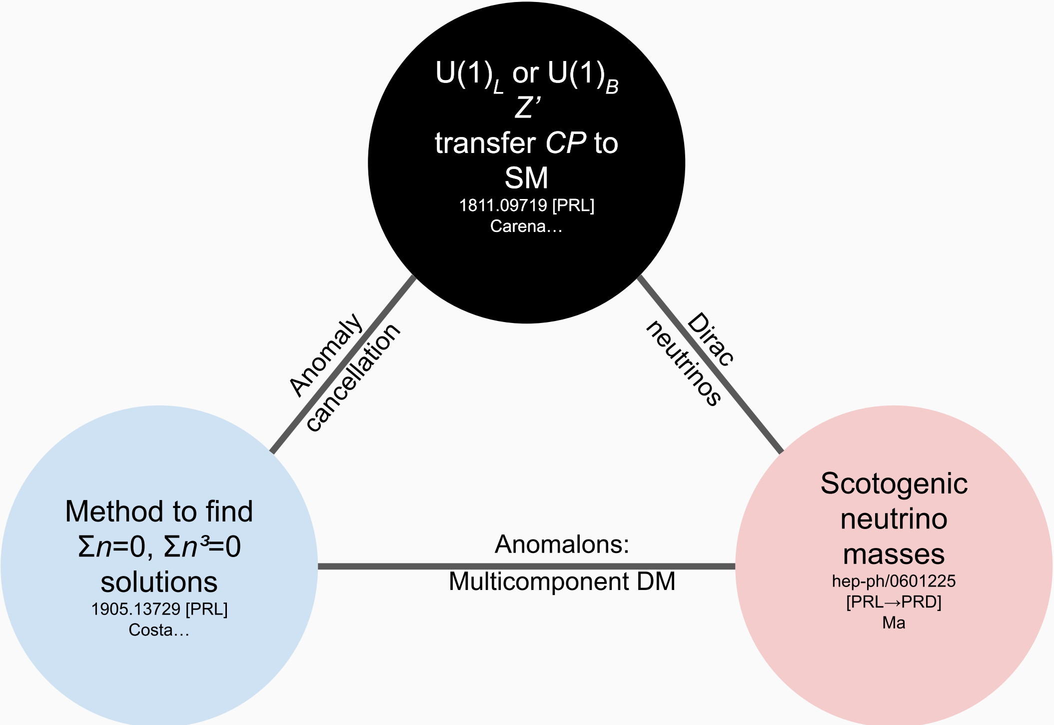
U(1)$_L$ or U(1)$_B$
Z'
transfer *CP* to SM
1811.09719 [PRL]
Carena…

Anomaly cancellation

Dirac neutrinos

Anomalons:

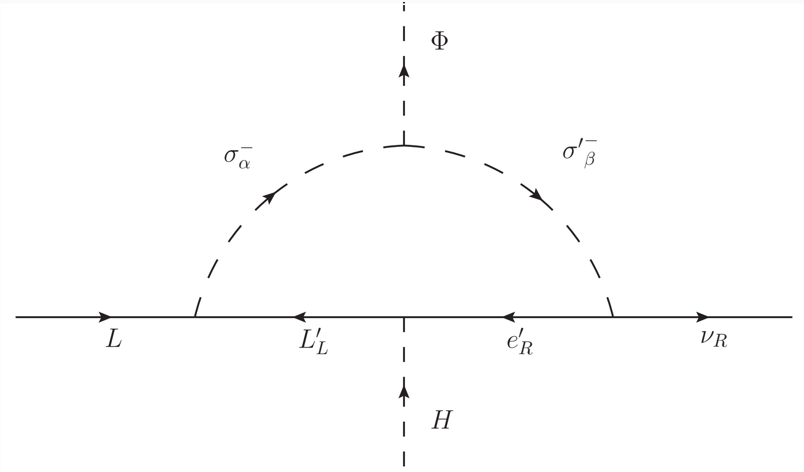
DM



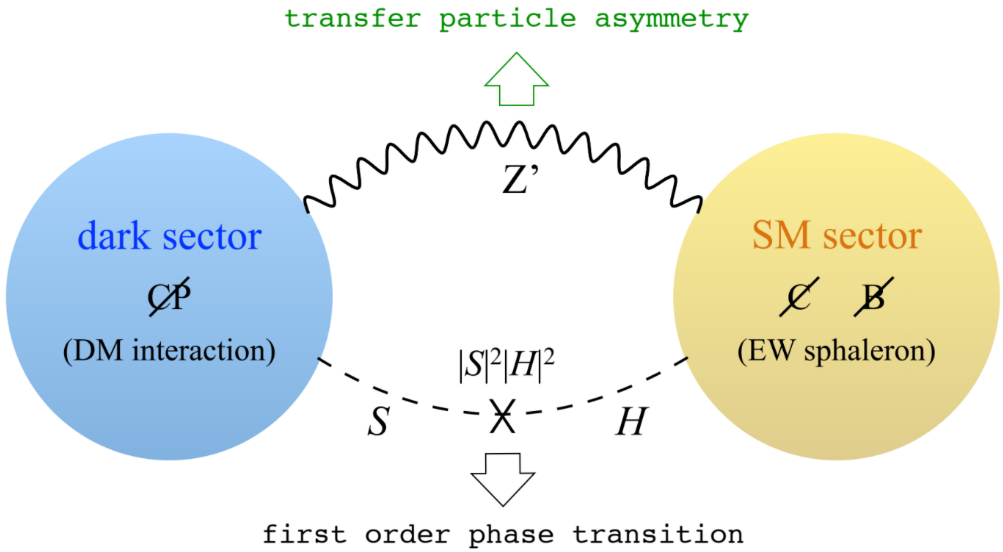
$U(1)_L$ or $U(1)_B$
Z'
transfer CP to SM
1811.09719 [PRL]
Carena…
Anomaly cancellation
Dirac neutrinos
Method to find Σn=0, Σn^3=0 solutions
1905.13729 [PRL]
Costa…
Anomalons:
Multicomponent DM
Scotogenic neutrino masses
hep-ph/0601225
[PRL→PRD]
Ma

Scotogenic realization

Any realization which does not affect anomaly cancellation is allowed



Dark sector baryogenesis



from: [1] Carena, *et al*, arXiv:1908.04818 [PRD] (see also: arXiv:1811.09719 [PRL])

Baryogenesis

CP violation occurs in the dark sector and is transmited to SM sector by the new Z' gauge boson.

- High scale fields: Φ , $(\langle\Phi\rangle \to L'_L, L''_R, e'_L, e''_R$: EW-scale vector-like anomalons)
- Electroweak scale (EW) fields: $\boxed{Z'_\mu, S, \chi_L, \chi_R}$
- *CP*-violation

$$\begin{aligned}\mathcal{L}_{\text{Dirac DM}} =& h\,(\chi_L)^\dagger\,\chi_R\Phi^* + y\,(\chi_L)^\dagger\,\chi_R S^* + \text{h.c}\,, \qquad\qquad y\in\mathbb{C}\\ \supset& \left(m_\chi + |y|\,\mathrm{e}^{i\theta}\,|S|\right)(\chi_L)^\dagger\,\chi_R + \text{h.c}\,.\end{aligned}$$

- *CP*-violation Portal

$$\mathcal{L}_{\text{anomalous}} \supset g'\,Z'_\mu\,\left[3\bar{\chi}_L\gamma^\mu\chi_L - 2\bar{\chi}_R\gamma^\mu\chi_R + \bar{Q}_i\gamma^\mu Q_i + \bar{q}_{Ri}\gamma^\mu q_{Ri}\right]$$

- Strong electroweak phase transition (EWPT) portal

$$\mathcal{L}_{\text{first order EWPT}} \supset -\lambda_{SH}H^\dagger H\,S^*\,S\,.$$

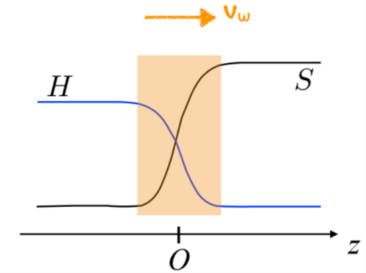
First-order phase transition: Effective potential ($T \neq 0$)

$\boxed{h = H/\sqrt{2},\ s = |S|}$ with vevs: $v(T)$ and $w(T)$ such that $v(T_c) = w(T_c)$

$$V_T(h,s) = \frac{\lambda_H v_c^4}{4}\left(\frac{h^2}{v_c^2} + \frac{s^2}{w_c^2} - 1\right)^2 + \frac{\lambda_H v_c^2}{m_{s,c}^2 w_{0,c}^4} h^2 s^2 + (T^2 - T_c^2)(c_h h^2 + c_s s^2)\,, \qquad (28)$$

where

$$c_h = \frac{1}{48}\left(9g_2^2 + 3g_1^2 + 12y_t^2 + 24\lambda_H + \lambda_{HS}\right)\,, \quad c_s = \frac{1}{12}\left(3\lambda_S + 2\lambda_{HS}\right)\,. \qquad (29)$$



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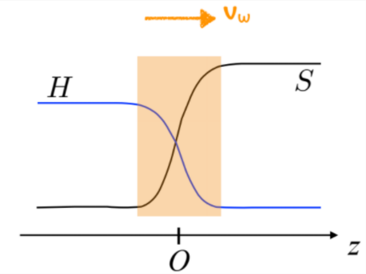
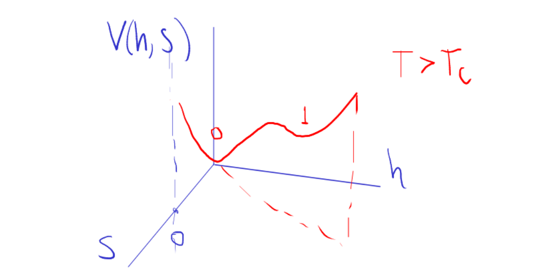
arXiv: Sec. 4.1 arXiv:1107.5451





0: (0, W₀(T)) ; 1: (v(T), w(T))

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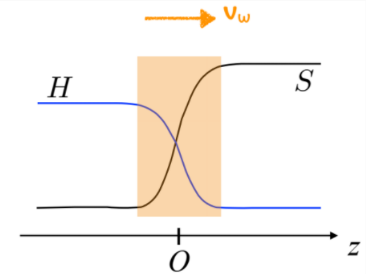

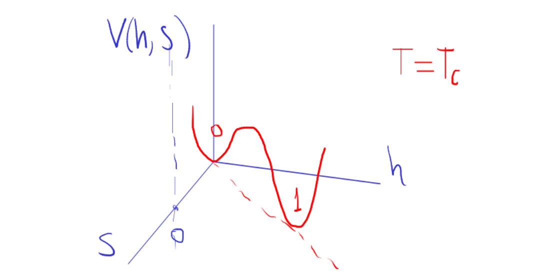

0: (0, w_0(T_c)) ; 1: (v(T_c), w(T_c))

First-order phase transition: Effective potential ($T \neq 0$)

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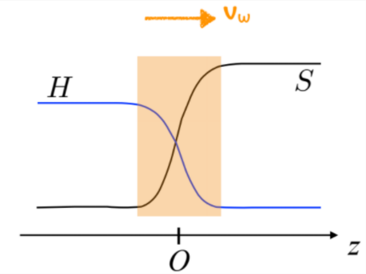
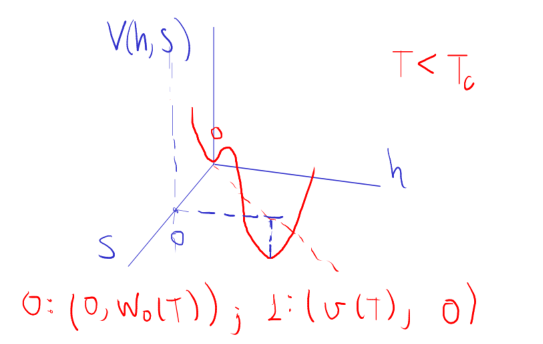




Using the thin wall approximantion for the nucleation bubbles, we use the ansatz in which the space dependence of the fields is given by

$$h(z) = \frac{1}{2}v(T_n)\left(1 - \tanh\left(z/L_w\right)\right), \qquad s(z) = \frac{1}{2}w_0(T_n)\left(1 + \tanh\left(z/L_w\right)\right),$$

where z is the direction normal to the wall and L_w is the wall width.

CP assymetry generation ii

The nucleation temperature, T_n, is defined by the condition

$$\exp\left(-S_3/T_n\right) = \frac{3}{4\pi}\left(\frac{H(T_n)}{T_n}\right)^4\left(\frac{2\pi\, T_n}{S_3}\right)^{\frac{3}{2}}\,,$$

where S_3 is the Euclidean action of the bubble and $H(T)$ is the Hubble rate.

The CP violating phase, θ, from

$$M_\chi(z) = m_\chi(z) + |y|\, \mathrm{e}^{i\theta}\, |S(z)|\,, \tag{30}$$

$\xi_i(z) \equiv \mu_i(z)/T = 6\left(n_i - \overline{n}_i\right)/T^3,$

$$-D_L\xi''_{\chi_L} - v_w\xi'_{\chi_L} + \Gamma_L(\xi_{\chi_L} - \xi_{\chi_R}) = S_{\not{CP}},$$

where D_L is the diffusion constant for χ_L, which is related to the scattering rate Γ_L by

$$D_L = \frac{3x+2}{x^2+3x+2}\frac{1}{3\Gamma_L}, \qquad x \equiv m_\chi/T \tag{31}$$

and

$$S_{\not{CP}} = -\frac{\lambda}{2}\frac{v_w D_L}{\frac{3x+2}{x^2+3x+2}T}\frac{(1-x)e^{-x}+x^2E_1(x)}{4m_\chi^2K_2(x)}\frac{m_\chi w_0(T_n)\lambda\left(-2+\cosh\left(\frac{2z}{L_w}\right)\right)\sin\theta}{L_w^3\cosh^4\left(\frac{z}{L_w}\right)}, \tag{32}$$

where v_w is the wall's velocity $E_1(x)$ is the error function and $K_2(x)$ is the modified Bessel function of the second kind.

Transfer DM assymetry to SM quarks

The chiral particle give rise to a non-zero $\mathrm{U}(1)_B$ charge density in the proximity of the wall. This results in a Z' background that couples to the SM fields with $\mathrm{U}(1)_B$ charge,

$$\langle Z_0'(z)\rangle = \frac{g_B\,(q_{\chi_L} - q_{\chi_R})\,T_n^3}{6M_{Z'}} \int_{-\infty}^{\infty} \mathrm{d}\,z_1\,\xi_{\chi_L}(z_1)\,e^{-M_{Z'}|z-z_1|}\,,$$

which generates a chemical potential for the SM quarks,

$$\mu_Q(z) = \mu_{d_R,u_R}(z) = 3 \times \frac{5}{9} \times g_B\langle Z_0'(z)\rangle.$$

This chemical potential sources a thermal-equilibrium asymmetry in the quarks, $\Delta n_Q^{\mathrm{EQ}}(z) \sim T_n^2\,\mu_Q(z)$.

From [1]

If the Z' is sufficiently light, it mediates a long range force that extends into the region outside the bubble wall with unbroken electroweak symmetry.

Finally, the baryon-number asymmetry is then given by

$$n_B = \frac{\Gamma_{\rm sph}}{v_w} \int_0^\infty \mathrm{d}\, z\, n_Q^{\rm EQ}(z)\, \exp\left(-\frac{\Gamma_{\rm sph}}{v_w}\, z\right)\,,$$

where $\Gamma_{\rm sph}$ is the sphaleron rate. The baryon-to-photon-number ratio is then obtained by

$$\eta_B = \frac{n_B}{s(T_n)}, \quad s(T) \equiv \frac{2\pi^2}{45} g_{*S}(T)\, T^3\,,$$

where $g_{*S}(T)$ is the effective number of relativistic degrees of freedom.

Our goal is to find what regions of the parameter space yield

$$0.82 \times 10^{-10} < \eta_B < 0.92 \times 10^{-10}\,. \tag{33}$$

https://github.com/anferivera/DarkBariogenesis

- SARAH→SPheno→MicroMegas
- η_B calculation code
- Python notebook with the scan

arXiv:1810.08055

Ten Simple Rules for Reproducible Research in Jupyter Notebook Fernando Pérez,*et al*

[...] In this paper, we address several questions about reproducibility [...] Combined with software repositories and open source licensing, notebooks are powerful tools for transparent, collaborative, reproducible, and reusable data analyses.

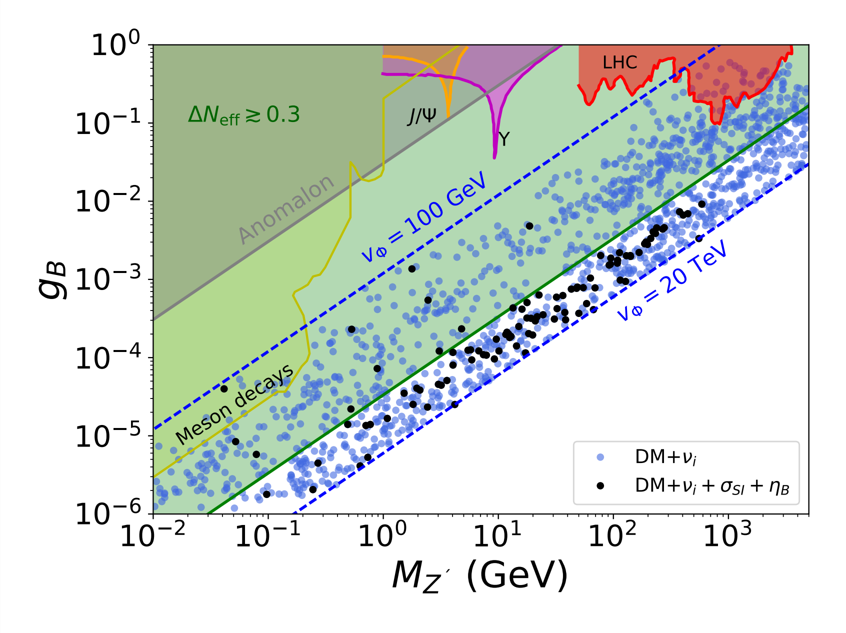
Results

We vary the typical Dirac-fermion DM parameter space and for each point that satisfy neutrino oscillation data, relic density and DM direct detection constraints. For each point we ...

Parameter	Range
θ	$(-\pi/2\,,\pi/2)$
$w_0(T_n)/\mathrm{GeV}$	$100 - 500$
T_n/GeV	$100 - 200$
L_w/GeV^{-1}	$1/T_n - 10/T_n$
v_w	$0.05 - 0.5$

Table 6: Scan ranges for the free parameters that are involved in the baryogenesis mechanism.

Black points: Dirac neutrinos with proper DM and baryon assymetry



Conclusions

A $U(1)_B$ is presented as an example of models where all new fermions required to cancel out the anomalies are used to solve phenomenological problems of the standard model (SM):

- EW-scale fermion vector-like doublets and iso-singlet charged singlets, in conjunction with right-handed neutrinos with repeated Abelian charges, participate in the generation of small neutrino masses through the Dirac-dark Zee mechanism
- The other SM-singlets are used to explain the dark matter in the universe, while their coupling to an inert singlet scalar is the source of the CP violation.

In the presence of a strong first-order electroweak phase transition, this “dark” CP violation allows for successful electroweak baryogenesis by using long range force mediated by a sufficiently light Z' which transfers the assymmetry from the Dark sector into the SM.