### Effective Dirac neutrino masses and baryogenesis



#### with gauged Baryon number

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# Focus on arXiv:¿¿¿¿¿.????? In collaboration with

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# Electroweak baryogenesis

#### **Problems**

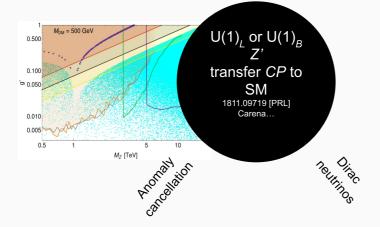
- Standard model (SM)  $m_h \sim$  40 GeV.  $\odot$
- Beyond the SM: Source of CP contains fields charged under SM
  - ightarrow too large electric dipole moments extstyle extsty



#### Dark sectors

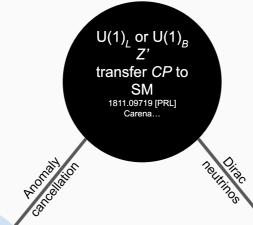
- Inert SM-singlet complex scalar field which acquires vev with temperature to have strong electroweak phase transition
- CP violation (CPV) triggered in dark sectors through SM gauge singlets
  - → CPV Yukawa between SM-singlet complex scalar and SM-singlet quiral fermions \(\to\)





Anomalons:

DM



Method to find  $\Sigma n=0$ ,  $\Sigma n^3=0$  solutions 1905.13729 [PRL] Costa...

Anomalons:

Multicomponent DM

Scotogenic neutrino masses

hep-ph/0601225 [PRL→PRD] Ma

# Dark sectors











 $\mathcal{L} = -\frac{1}{4}V_{\mu\nu}V^{\mu\nu} + i\sum_{i}\chi_{i}^{\dagger}\mathcal{D}\chi_{i}$ 

$$-h(\chi_1\chi_2\Phi + h.c)$$

Anomalons: SM-singlet Dirac fermion dark matter  $m_{\Psi} = h\langle \Phi \rangle$ 

LHC productio

Gauged Symmetry:  $\mathcal{X} \to B$ :  $q\overline{q} \to Z' \to \text{jets}$ 

Gauged Symmetry:  $\mathcal{X} \to L$ :



$$\overline{\Psi}\Psi = \chi_1\chi_2 + \chi_1^{\dagger}\chi_2^{\dagger} \rightarrow \chi_{\alpha}\chi_{\beta}\Phi^{(*)}, \qquad \alpha = 1, \dots N' \rightarrow N' > 4$$



 $\mathcal{L} = -rac{1}{4}V_{\mu
u}V^{\mu
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Gauged Symmetry:  $\mathcal{X} \rightarrow \mathcal{L}$ :



multi-component dark matter

 $\alpha=1,\ldots N' o N'>4$ 



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multi-component dark matter

 $\alpha=1,\ldots N' o N'>4$ 



Local  $U(1)\chi$   $\mathcal{L} = -\frac{1}{4}V_{\mu\nu}V^{\mu\nu} + i\sum_{i}\chi_{i}^{\dagger}\mathcal{D}\chi_{i}$ 

$$-y(\chi_1\chi_2S+h.c)$$

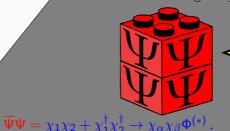
Anomalons: SM-singlet Dirac fermion

CP violation Yukawa y

LHC productio

Gauged Symmetry:  $\mathcal{X} \to B$ :  $q\overline{q} \to Z' \to \text{jets}$ 

Gauged Symmetry: 
$$\mathcal{X} \rightarrow \mathcal{L}$$
:



multi-component dark matter

 $\alpha = 1, \dots, N' \rightarrow N' > 4$ 

#### Standard model extended with $U(1)_{\mathcal{X}}$ gauge symmetry

Fields	$SU(2)_L$	$U(1)_Y$	$U(1)_{\mathcal{X}=B \text{ or } L}$
$Q_i^{\dagger}$	2	-1/6	Q
$d_{Ri}$	1	-1/2	d
$u_{Ri}$	1	+2/3	u
$L_i^{\dagger}$	2	+1/2	L
$e_{Ri}$	1	-1	e
Н	2	1/2	h = 0
$\chi_{\alpha}$	1	0	$z_{\alpha}$
$(L'_L)^{\dagger}$	2	1/2	-x'
$L_R^{\prime\prime}$	2	-1/2	x''
$e_R'$	1	-1	×′
$(e_L^{\prime\prime})^\dagger$	1	1	-x''
Ф	1	0	$\phi$
S	1	0	s

**Table 1:** L = e = 0 for  $\mathcal{X} = B$ . Or Q = u = d = 0 for  $\mathcal{X} = L$ .  $i = 1, 2, 3, \ \alpha = 1, 2, ..., N'$ 

#### Effective Dirac neutrino mass operator

$$\chi_1 \to \nu_{R1}, \cdots, \chi_{N_{\nu}} \to \nu_{RN_{\nu}}, \qquad 2 \le N_{\nu} \le 3,$$
(1)

$$\mathcal{L}_{\mathrm{eff}} = h_{
u}^{lpha i} \left(
u_{Rlpha}
ight)^{\dagger} \, \epsilon_{ab} \, \mathcal{L}_{i}^{a} \, \mathcal{H}^{b} \left(rac{\Phi^{*}}{\Lambda}
ight)^{\delta} + \mathrm{H.c.}, \qquad \mathrm{with} \, \, i=1,2,3 \, ,$$

S is the complex singlet scalar responsible for the SSB of the anomaly-free gauge symmetry with D or X-charge

$$\phi = -(\nu + \mathbf{L})/\delta \,, \tag{2}$$

#### Anomaly cancellation I

The anomaly-cancellation conditions on  $[SU(3)_c]^2 U(1)_X$ ,  $[SU(2)_L]^2 U(1)_X$ ,  $[U(1)_Y]^2 U(1)_X$ , allow us to express three of the X-charges in terms of the others

$$\mathbf{u} = -\mathbf{e} - \frac{2}{3}\mathbf{L} - \frac{1}{9}\left(x' - x''\right) , \quad \mathbf{d} = \mathbf{e} + \frac{4}{3}\mathbf{L} - \frac{1}{9}\left(x' - x''\right) , \quad \mathbf{Q} = -\frac{1}{3}\mathbf{L} + \frac{1}{9}\left(x' - x''\right) , \quad (3)$$

while the  $[U(1)_X]^2 U(1)_Y$  anomaly condition reduces to

$$(e+L)(x'-x'')=0.$$
 (4)

- Previously: x' = x''
- We choose instead (h = 0):

$$e = -L, (5)$$

so that (L is still a free parameter)

$$Q = -u = -d = -\frac{1}{3}L + \frac{1}{9}(x' - x'').$$
 (6)

If 
$$L=0 \rightarrow U(1)_B$$

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#### **Anomaly cancellation II**

The gravitational anomaly,  $[SO(1,3)]^2 U(1)_Y$ , and the cubic anomaly,  $[U(1)_X]^3$ , can be written as the following system of Diophantine equations, respectively,

$$\sum_{\alpha=1}^{N} z_{\alpha} = 0, \qquad \sum_{\alpha=1}^{N} z_{\alpha}^{3} = 0, \qquad (7)$$

where N = N' + 5 and

$$z_{N'+1} = -x',$$
  $z_{N'+2+i} = L, \quad i = 1, 2, 3$  (8)

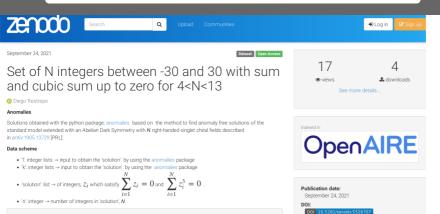
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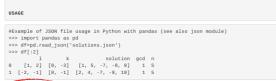
$$9Q = -\sum_{\alpha=N'+1}^{N'+5} z_{\alpha} = -x' + x'' + L + L + L,$$
 (9)

If 
$$Q = 0 \rightarrow U(1)_L$$









390074 solutions with  $5 \le N \le 12$  integers until '1321' [JSON]



Anomaly free | Diophantine equations | Abelian symmetry

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Keyword(s):

License (for files):

## $U(1)_{\mathbf{B}}$ selection

• 
$$L=0$$

$$(5,5,-3,-2,1,-6)$$

### $U(1)_{\mathbb{B}}$ selection

- L = 0
- Effective neutrino mass  $\phi = \nu = -5$

$$(5, 5, -3, -2, 1, -6)$$

### $U(1)_{\mathbb{B}}$ selection

- L = 0
- Effective neutrino mass  $\phi = \nu = -5$

• 
$$(L'_L)^{\dagger} L''_R \Phi^* \to x' = -1, \ x'' = 6$$

$$(5,5,-3,-2,1,-6)$$

### $U(1)_{\bf B}$ selection

- L=0
- Effective neutrino mass  $\phi=\nu=-5$

• 
$$(L'_L)^{\dagger} L''_R \Phi^* \to x' = -1, \ x'' = 6$$

• 
$$(\chi_L)^{\dagger} \chi_R'' \Phi^* \to z_3 = -3, \ z_4 = -2$$

959 solutions from  $\sim$  400,000

$$(5, 5, -3, -2, 1, -6)$$