Phenomenology of scotogenic models

LHC signals



Diego Restrepo

November 15, 2016

Instituto de Física Universidad de Antioquia Phenomenology Group http://gfif.udea.edu.co



Volcán de Fuego (Caroline Kish)

Focus on

arXiv: arXiv:1308.3655 (JHEP), arXiv:1504.07892 (PRD), arXiv:1509.06313 (PRD), arXiv:1511.01873 (JHEP), arXiv:1605.01129 (PRD)

In collaboration with

G. Palacio, F. von der Pahlen, D. Portillo, A. Rivera, M. Sánchez, O. Zapata (UdeA)

C. Arbeláez (USM), W. Tangarife (Tel Aviv U.), C. Yaguna (Heidelberg, Max Planck Inst.).

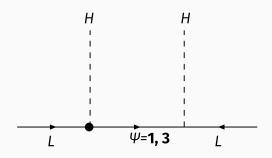
Table of Contents

- 1. General framework
- 2. Proposal: $pp \rightarrow l^+l^- + E_T^{\text{miss}}$
- 3. Specific examples
- 4. Lepton flavor dependence
- 5. Prospects for run-II
- 6. Vector-like fermion mediation

General framework

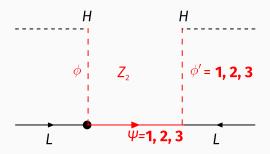
ν -DM models

small neutrino masses



ν -DM models

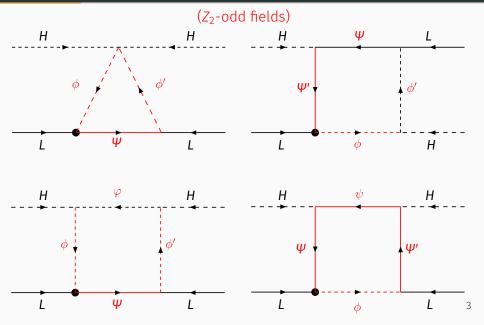
small neutrino masses $\Leftarrow Z_2 \Rightarrow$ dark matter



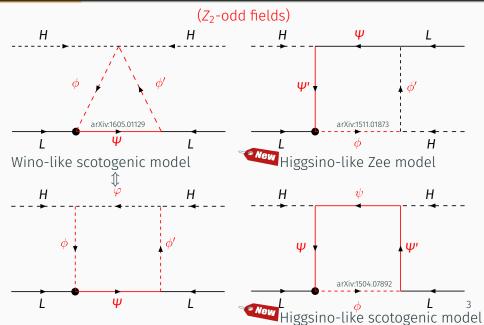
35 non-equivalent dark matter models classified in D.R., C. Yaguna, O. Zapata, arXiv:1308.3655 (JHEP)

2. Neutrinos talk to a different **Higgs boson**

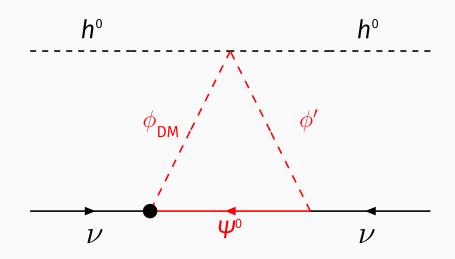
Weinberg operator at one-loop



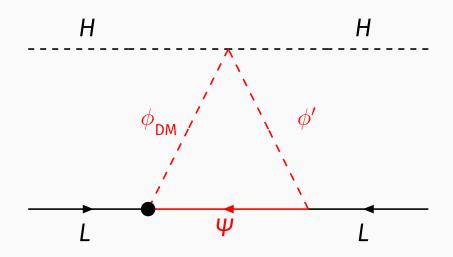
Weinberg operator at one-loop



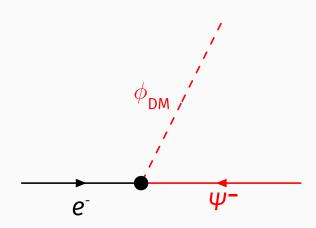
Typical radiative neutrino mass diagram.



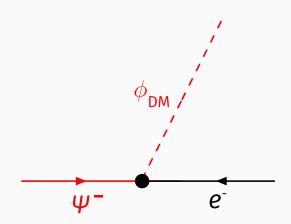
In term of general SU(2)_L multiplets,



may be also contain charged particles,

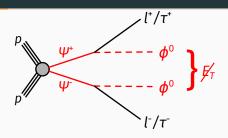


which may decay into the dark matter particle.

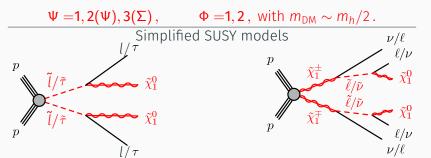


Proposal: $pp \rightarrow l^+l^- + E_T^{\text{miss}}$

Dilepton plus transverse missing energy signal



 $SU(2)_L$ assignments:



Smaller cross sections. Intermediate states and smaller lepton p_{\uparrow}

Specific examples

Specific examples

- Wino-like scotogenic models
 - Radiative type-III seesaw: 1605.01129, F. von der Pahlen, G. Palacio, DR, O. Zapata
- Higgsino-like scotogenic models
 - 1. SDFM with scalars: 1504.07892, DR, et. al..
 - 2. Inert Zee: 1511.01873, R. Longas, D. Portillo, DR, O. Zapata.
 - 3. Radiative type-II seesaw: 1511.06375, S. Fraser, C. Kownacki, E. Ma, O. Popov

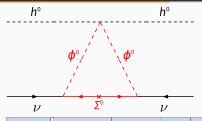
1609.01018, S. Guo, Z. Han, Y, Liao

- Bino-like scotogenic models
 - · In progress ...

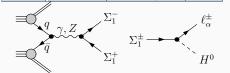
Wino-like scotogenic model

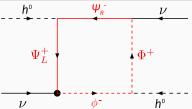
 \sum_{k}

Higgsino-like inert Zee model



	/	`\	_	_
ν	$\hat{\Sigma^0}$		$\overline{\nu}$	_
	SU(2) _L	U(1) _Y	Z_2	S
Н	2	1	+	0
Ф	2	1	_	0
L_{α}	2	-1	+	1/2





ν	ϕ^-		h⁰	
	SU(2) _L	U(1) _Y	Z_2	S
ϕ^-	1	-2	_	0
Ф	2	_	0	-
ψ^-	1	-2	_	0
$\Psi_{L,R}$	2	±1	_	1/2

$$\Sigma^+ o \psi^+$$

Wino-like scotogenic model

Higgsino-like scotogenic model



h ⁰	ψ^{0}	h ^o
	•	4
1110		111 0
$\Psi_{\scriptscriptstyle L}^{ \circ lack}$	4	Ψ_{R}^{0}
$\overline{\nu}$	φ ⁰	$\overline{\nu}$

	SU(2) _L	U(1) _Y	Z_2	S
Н	2	1	+	0
Ф	2	1	_	0
L_{lpha}	2	-1	+	1/2
Σ_k	3	0	_	1/2

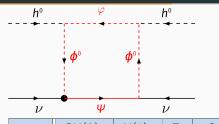
		SU(2) _L	U(1) _Y	Z_2	S
Ī	ϕ^-	1	-2	_	0
	Ф	2	_	0	-
	ψ^-	1	-2	_	0
	$\Psi_{L,R}$	2	±1	_	1/2
	ψ	1	0	-	1/2
	ϕ	1	0	-	0

$\equiv q_{\gamma,Z}$	Σ_1^-	ℓ_{α}^{\pm}
	Σ_1^{\pm} Σ_1^{+}	H^0

$$\Sigma^+ o \psi^+$$

Wino-like scotogenic model

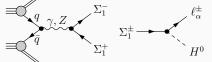
Higgsino-like model



h ^o	ψ^0	h°
$\psi_{\iota}{}^{\circ}\cdot$		Ψ _R °
$\overline{\nu}$	ϕ^{0}	$\overline{\nu}$

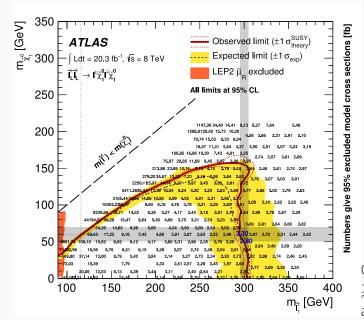
	SU(2) _L	U(1) _Y	Z_2	S
Н	2	1	+	0
Ф	2	1	_	0
L_{lpha}	2	-1	+	1/2
Σ_k	3	0	_	1/2

	SU(2) _L	U(1) _Y	Z_2	S
ϕ^-	1	-2	_	0
Ф	2	_	0	-
ψ^-	1	-2	_	0
$\Psi_{L,R}$	2	±1	_	1/2
ψ	1	0	-	1/2
ϕ	1	0	-	0
	ψ ⁻ Ψ _{L,R}		ϕ^{-} 1 -2 ϕ 2 - ψ^{-} 1 -2 $\psi_{L,R}$ 2 ±1	$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$

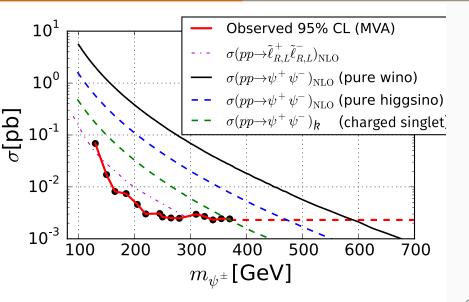


$$\Sigma^+ \rightarrow \psi^+$$

ATLAS arXiv:1403.5294 (JHEP)



CMS \gtrsim 260 GeV $_{8}$ arXiv:1405.7570



Lepton flavor dependence

Neutrino masses

$$(\mathcal{M}_{\nu})_{\alpha\beta} = \sum_{k=1}^{n_{\Sigma}} [\mathbf{Y}^{\mathsf{T}} \Lambda \mathbf{Y}]_{\alpha\beta} , \qquad \alpha, \beta = 1, 2, 3,$$

From neutrino oscillation data, we can get a set of Y choosing the angles for R, an arbitrary complex orthogonal matrix

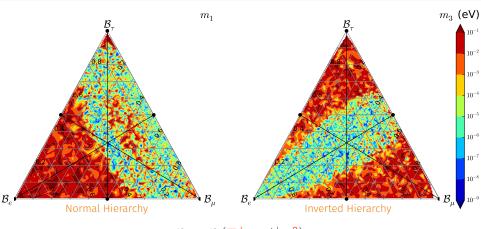
$$\mathbf{Y} = \sqrt{\Lambda}^{-1} \mathbf{R} \operatorname{diag}(\sqrt{m_{\nu_1}}, \sqrt{m_{\nu_2}}, \sqrt{m_{\nu_3}}) U_{\mathrm{PMNS}}^{\dagger}, \tag{1}$$

$$\hat{\mathbf{Y}}_{\alpha} \equiv \hat{\mathbf{Y}}_{1\alpha} = \mathbf{Y}_{1\alpha} / \sqrt{\sum_{\alpha=e,\mu,\tau} |\mathbf{Y}_{1\alpha}|^2}$$
 $\mathbf{\mathcal{B}}_{\alpha} \equiv \operatorname{Br}(\mathbf{\Sigma}_{1}^{\pm} \to \ell_{\alpha} \mathbf{H}^{0}) = |\hat{\mathbf{Y}}_{\alpha}|^2.$

Input parameters: 3 complex angles and 1 phase.

Casas-Ibarra parametrization

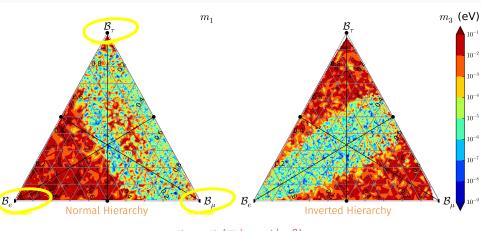
In wino-like scotogenic model (may be in general)



$$\mathcal{B}_l = \mathcal{B}\left(\mathbf{\Sigma^{\pm}}
ightarrow l^{\pm}H^0\right)$$

Casas-Ibarra parametrization

In wino-like scotogenic model (may be in general)



$$\mathcal{B}_l = \mathcal{B}\left(\mathbf{\Sigma^{\pm}} \rightarrow l^{\pm}\mathbf{H^0}\right)$$

Exploration of flavor space

Wino-like scotogenic model: Recast for $B_{\mu}+B_{e}\gtrsim0.1$ and

$$m_{H^0} < m_{\Sigma^{\pm}} = m_{\Sigma^0} < m_{A^0}, m_{H^{\pm}}$$

Start with Signal regions as in ATLAS-arXiv:1403.5294 for E_T with e^+e^- , $\mu^+\mu^-$, $e^\pm\mu^\mp$.

SARAH/FeynRules

 \Downarrow

micrOMEGAS (Experimental and theoretical constraints)

11

MadGraph

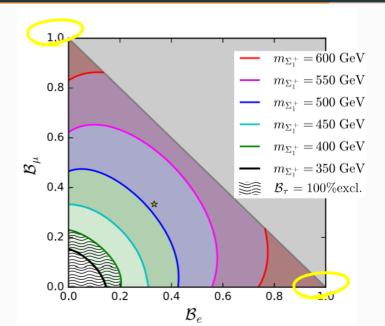
 \Downarrow

Pythia 6 (hep format)

 \Downarrow

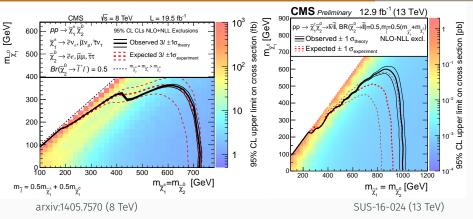
checkMATE (CL-calculation)

Combination



Prospects for run-II

Golden EW SUSY channel: trilepton and $\not E_T$



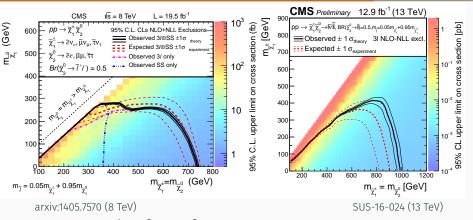
Improvement by a factor of 1.4

For a similar improvement we could expect exclusions at the level of 900 GeV in the wino-like scotogenic model,

700 GeV in Higgsino-like scotogenic models.

500 GeV in Bino-like scotogenic models.

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Vector-like fermion mediation

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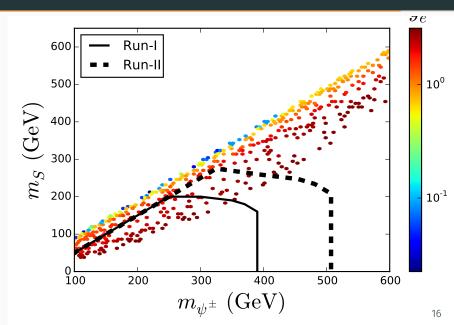
Straightforward way to avoid DD constraints in scalar dark matter:

Name	Symbol	$SU(3)_c$	$SU(2)_L$	U(1) _Y	Z_2
$\left(\nu_{L} e_{L}\right)^{T}$	$\begin{pmatrix} \xi_{1\alpha} & \xi_{2\alpha} \end{pmatrix}^T$	1	2	-1/2	+1
$(e_R)^{\dagger}$	η_1^{lpha}	1	1	+1	+1
$(\psi_{R})^\dagger$	η_2^{lpha}	1	1	+1	-1
ψ_{L}	ξ_{3lpha}	1	1	-1	-1
S		1	1	0	-1

$$\mathcal{L} \supset y_e S(e_R)^{\dagger} \frac{\psi_L}{\psi_L} + m_{\psi^{\pm}} (\psi_R)^{\dagger} \frac{\psi_L}{\psi_L} + \text{h.c} + \frac{1}{2} m_S S^2 + \underline{\lambda}_{HS} S^2 H^{\dagger} H$$

See: arXiv:1307.6181 and arXiv:1307.6480

LHC constraints: Preliminary



Conclusions

Opposite sign dilepton plus missing transverse energy signal at LHC

The use of scotogenic models to interpret dilepton plus missing transverse energy searches, allow for larger sensitivities and full lepton flavor exploration

Additional motivation for fermion vectorlike mediation with zero three-level direct detection cross section and challenging compressed spectra.

