Dark matter from SM gauge extensions

with neutrino masses

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July 23, 2019 - PPC2019 - Cartagena

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Focus on

In collaboration with
M. Hirsch (IFIC), C. Álvarez (UTFSM), A. Flórez (UniAndes), B. Dutta(Texas A& M), C. Yaguna (UPTC), J. Calle, O. Zapata, A. Rivera (UdeA), W. Tangarife (Loyola University Chicago)

Hidden sectors

SM

35 models

SM

+

$$m^{\nu}_{\text{Majorana}} = \frac{1}{\Lambda} L \cdot H L \cdot H \text{ (1-loop)}$$

arXiv:1308.3655 [JHEP] with C. Yaguna and Ó. Zapata

SM

+

$$m^{\nu}_{\text{Majorana}} = \frac{1}{\Lambda} L \cdot H L \cdot H$$

$$m^{\nu}_{\text{Dirac}} = \frac{1}{\Lambda} (\nu_R)^{\dagger} L \cdot H S$$

Local $U(1)_X \to Z_N$

SM

+

$$m^{\nu}_{\text{Majorana}} = \frac{1}{\Lambda} L \cdot H L \cdot H$$

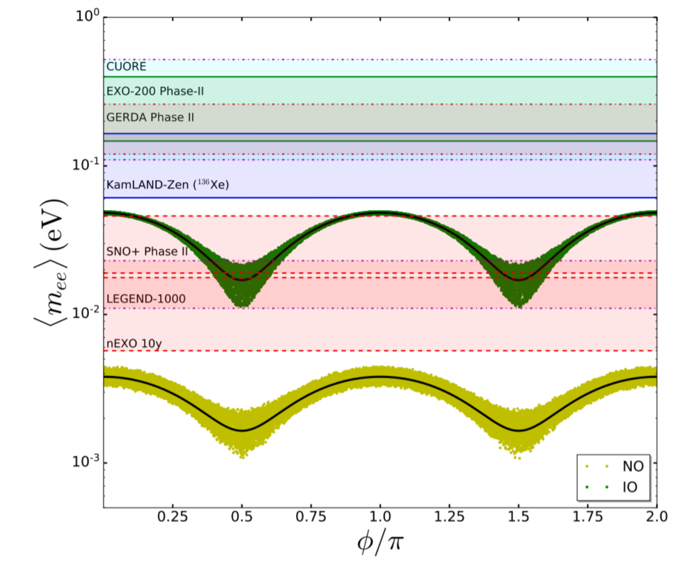
$$m^{\nu}_{\text{Dirac}} = \frac{1}{\Lambda} (\nu_R)^{\dagger} L \cdot H S$$

Neutrino masses

Lepton number

- Lepton number (L) is an accidental discret or Abelian symmetry of the standard model (SM).
- Without neutrino masses L_e, L_μ, L_τ are also conserved.
- The processes which violates individual L are called Lepton flavor violation (LFV) processes.
- All the neutrino mass models predict, to some extent, LFV processes
- Only models with Majorana neutrinos predict processes with total $L = L_e + L_\mu + L_\tau$ violation, like neutrino less doublet beta decay (NLDBD).
- NLDBD is experimentally challenging, specially if there is a massless neutrino in the spectrum.

NLDBD prospects for a model with a massless neutrino (arXiv:1806.09977 [PLB] with Reig, Valle and Zapata)



with M.Reig, J.W.F Valle, O. Zapata, arXiv:1806.09977

Total lepton number: $L = L_e + L_\mu + L_\tau$

Majorana ~~U(1)~~$_L$

Field	Z_2 ($\omega^2 = 1$)
SM	1
L	ω
$(e_R)^\dagger$	ω
$(\nu_R)^\dagger$	ω

$$\mathcal{L}_\nu = h_D\,(\nu_R)^\dagger\, L \cdot H + M_R\,\nu_R\nu_R + \text{h.c}\,.$$

$$h_D \sim \mathcal{O}(1)$$

Dirac U(1)$_L$

Field	Z_3 ($\omega^3 = 1$)
SM	1
L	ω
$(e_R)^\dagger$	ω^2
$(\nu_R)^\dagger$	ω^2

$$\mathcal{L}_\nu = h_D\,(\nu_R)^\dagger\, L \cdot H + \text{h.c}\,.$$

$$h_D \sim \mathcal{O}(10^{-11})$$

Total lepton number: $L = L_e + L_\mu + L_\tau$

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Dirac $\mathrm{U}(1)_{B-L}$

Field	Z_3 ($\omega^3 = 1$)
SM	1
L	ω
$(e_R)^\dagger$	ω^2
$(\nu_R)^\dagger$	ω^2

$$\mathcal{L}_\nu = h_D\,(\nu_R)^\dagger\, L\cdot H + \text{h.c}\,.$$

$$h_D \sim \mathcal{O}(10^{-11})$$

Explain smallness ala Peccei-Quinn:

$$\mathrm{U}(1)_{B-L} \underbrace{\longrightarrow}_{\langle S\rangle} Z_N\,, \quad N \geq 3\,.$$

Small Dirac neutrino masses

To explain the smallness of Dirac neutrino masses choose $\mathrm{U}(1)_{B-L}$ which:

- Forbids tree-level mass (TL) term ($Y(H) = +1/2$)

$$\begin{aligned}\mathcal{L}_{\mathrm{T.L}} =& h_D \epsilon_{ab} \left(\nu_R\right)^\dagger L^a H^b + \mathrm{h.c} \\ =& h_D \left(\nu_R\right)^\dagger L \cdot H + \mathrm{h.c}\end{aligned}$$

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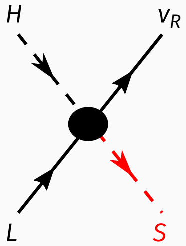
$$= h_D (\nu_R)^\dagger L \cdot H + \mathrm{h.c}$$

- Forbids Majorana term: $\nu_R \nu_R$
- Realizes of the 5-dimension operator which conserves lepton number in $\mathrm{SU}(3)_c \times \mathrm{SU}(2)_L \times \mathrm{U}(1)_Y \times \mathrm{U}(1)_{B-L}$:

$$\mathcal{L}_{5-D} = \frac{h_\nu}{\Lambda} (\nu_R)^\dagger L \cdot H S + \mathrm{h.c}$$

$$\mathrm{U}(1)_{B-L} \xrightarrow{\langle S \rangle} Z_N$$



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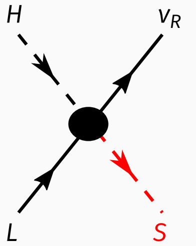

- Forbids Majorana term: $\nu_R \nu_R$
- Realizes of the 5-dimension operator which conserves lepton number in $\mathrm{SU}(3)_c \times \mathrm{SU}(2)_L \times \mathrm{U}(1)_Y \times \mathrm{U}(1)_{B-L}$:

$$\mathcal{L}_{5-D} = \frac{h_\nu}{\Lambda} \left(\nu_R\right)^\dagger L \cdot H S + \text{h.c}$$

- Enhancement to the *effective number of degrees of freedom in the early Universe* $\Delta N_{\text{eff}} = N_{\text{eff}} - N_{\text{eff}}^{\text{SM}}$ (see arXiv:1211.0186)

See E. Ma, Rahul Srivastava: arXiv:1411.5042 [PLB] for tree-level realization

SM

+

$$m^{\nu}_{\text{Majorana}} = \frac{1}{\Lambda} L \cdot H L \cdot H$$

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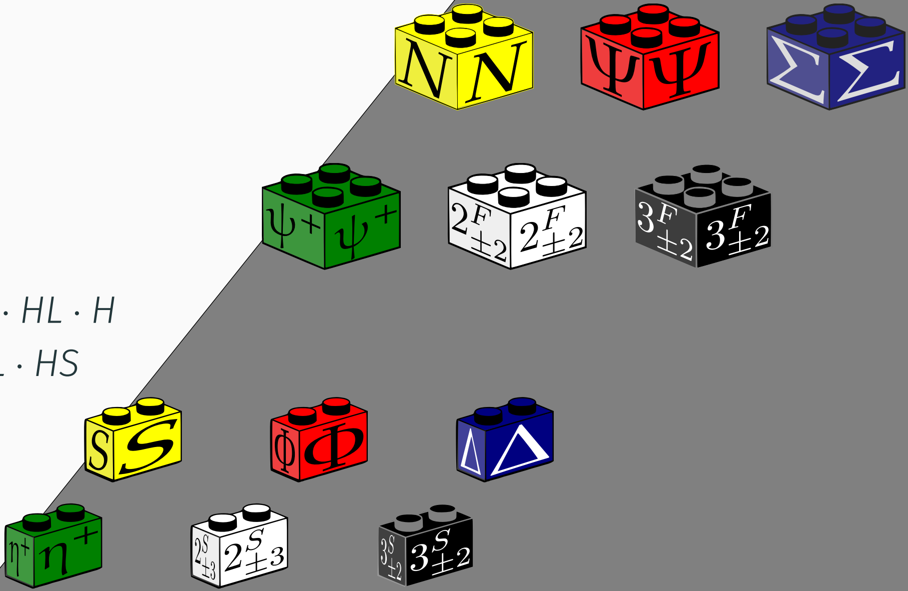

Dark matter and unification

Unification: SO(10)

$$= \mathbf{16}_{F_i} = \begin{pmatrix} (u_R)^\dagger \\ (u_R)^\dagger \\ (u_R)^\dagger \\ u_L \\ u_L \\ u_L \\ (d_R)^\dagger \\ (d_R)^\dagger \\ (d_R)^\dagger \\ d_L \\ d_L \\ d_L \\ \nu_L \\ e_L \\ (e_R)^\dagger \\ N \end{pmatrix}_i \qquad \Rightarrow \mathcal{L}_{SM} \supset h\, \mathbf{16}_F \times \mathbf{16}_F \times \mathbf{10}_S + \text{h.c}$$

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SO(10) breakings

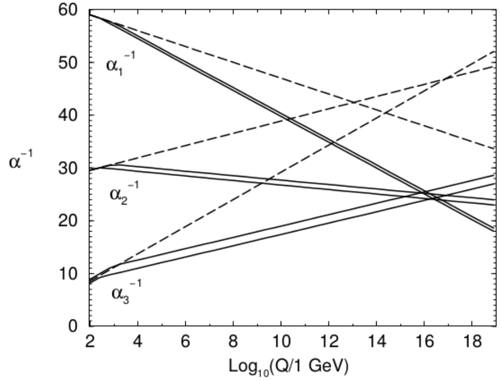
$$\mathrm{SO}(10) \to \begin{cases} \mathrm{SU}(5) \times \mathrm{U}(1)_X & \text{with } \begin{cases} Z_X \text{ at GUT: } \mathrm{U}(1)_X \to Z_N \\ Z_X \text{ at EW} \end{cases} \\ \mathrm{SU}(3)_c \times \mathrm{SU}(2)_R \times \mathrm{SU}(2)_L \times \mathrm{U}(1)_Y \times \mathrm{U}(1)_{B-L} \end{cases}$$

Majorana neutrinos case

Not-susy SO(10) $\rightarrow SU(5) \rightarrow$ SU(3)$_C \times SU(2)_L \times$ U(1)$_Y \times Z_2$

Standard Model: Z_2-even	New Z_2-odd particles
Fermions: $\mathbf{16}_F$	$\mathbf{10}_F, \mathbf{45}_F, \cdots$
Scalars: $\mathbf{10}_H, \mathbf{45}_H \cdots$	$\mathbf{16}_H, \cdots$

Lightest Odd Particle (LOP) may be a suitable dark matter candidate, and can improve gauge coupling unification



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	fermions	scalars
$SU(2)_L \times U(1)_Y$ representation	even $SO(10)$ representations	odd $SO(10)$ representations
1_0	45, 54, 126, 210	16, 144
$2_{\pm 1/2}$	10, 120, 126, 210, 210′	16, 144
3_0	45, 54, 210	144

$SU(3)_C$: 3 (T), 6, 8 (Λ)

$$m_{3_0} = 2.7 \text{ TeV}, \qquad m_\Lambda \sim 10^{10} \text{ TeV}, \qquad m_{\text{GUT}} \sim 10^{16} \text{ GeV}.$$

arXiv:0912.1545 [PRD] (Frigerio-Hambye)

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$\boxed{2_{\pm 1/2}}$	$\boxed{10,}$ 120, 126, 210, 210′	16, 144
$\boxed{3_0}$	$\boxed{45,}$ 54, 210	144

$SU(3)_C$: 3 (T), 6, $\boxed{8\ (\Lambda)}$

Split-SUSY like

arXiv:1509.06313 [PRD] with C. Arbelaez, R. Longas, and O. Zapata.

Not-susy SO(10) $\to SU(5) \to \mathrm{SU}(3)_C \times SU(2)_L \times \mathrm{U}(1)_Y \times Z_2$

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Radiative hybrid seesaw (Parida 1106.4137) or 1509.06313

Partial Split-SUSY-like spectrum: bino-higgsino-wino

+ $\downarrow$

$\mathbf{10}'_H$ with fermion DM or, $\mathbf{16}_H, \cdots$ with scalar DM

Not-susy SO(10) $\to SU(5) \to \mathrm{SU}(3)_C \times SU(2)_L \times \mathrm{U}(1)_Y \times Z_2$

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Fermions: $\mathbf{16}_F$	$\mathbf{10}_F, \mathbf{45}_F, \cdots$
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1509.06313

SUSY-like spectrum: bino-higgsino-wino

\+

$\downarrow$

$\mathbf{10}'_H$ with fermion DM or, $\mathbf{16}_H, \cdots$ with scalar DM

Singlet-Doublet-Triplet fermion dark-matter

The most general SO(10) invariant Lagrangian contains the following Yukawa terms

$$-\mathcal{L} \supset Y \mathbf{10}_F \mathbf{45}_F \mathbf{10}_H + M_{\mathbf{45}_F} \mathbf{45}_F \mathbf{45}_F + M_{\mathbf{10}_F} \mathbf{10}_F \mathbf{10}_F$$

Basis $\boldsymbol{\psi}^0 = \left(N, \Sigma^0, \psi_L^0, (\psi_R^0)^\dagger \right)^T$

$$\mathcal{M}_{\boldsymbol{\psi}^0} =$$

$$\begin{pmatrix} M_N & 0 & -y\, c_\beta v/\sqrt{2} & y\, s_\beta v/\sqrt{2} \\ 0 & M_\Sigma & f\, c_{\beta'} v/\sqrt{2} & -f\, s_{\beta'} v/\sqrt{2} \\ -y\, c_\beta v/\sqrt{2} & f\, c_{\beta'} v/\sqrt{2} & 0 & -M_D \\ y\, s_\beta v/\sqrt{2} & -f\, s_{\beta'} v/\sqrt{2} & -M_D & 0 \end{pmatrix},$$

$\mathbf{10}_F \to \psi_L, (\psi_R)^\dagger$

$\mathbf{45}_F \to \Sigma, \Lambda$

$\mathbf{45}'_F \to N$

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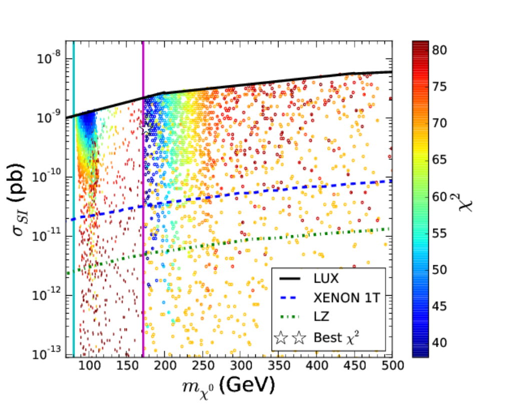
$$10_F \to \psi_L, (\psi_R)^\dagger$$

$$45_F \to \Sigma, \Lambda$$

$$45'_F \to N$$



S. Horiuchi, O. Macias, DR, A. Rivera, O. Zapata, 1602.04788 (JCAP)

Singlet-Doublet-Triplet fermion dark-matter

The most general SO(10) invariant Lagrangian contains the following Yukawa terms

$$-\mathcal{L} \supset Y 10_F 45_F 10_H + M_{45_F} 45_F 45_F + M_{10_F} 10_F 10_F + \mathcal{L}(10_\Phi \qquad).$$

Basis $\psi^0 = \left(N, \Sigma^0, \psi_L^0, (\psi_R^0)^\dagger\right)^T$

$$\mathcal{M}_{\psi^0} = \begin{pmatrix} M_N & 0 & -y\,c_\beta v/\sqrt{2} & y\,s_\beta v/\sqrt{2} \\ 0 & M_\Sigma & f\,c_{\beta'} v/\sqrt{2} & -f\,s_{\beta'} v/\sqrt{2} \\ -y\,c_\beta v/\sqrt{2} & f\,c_{\beta'} v/\sqrt{2} & 0 & -M_D \\ y\,s_\beta v/\sqrt{2} & -f\,s_{\beta'} v/\sqrt{2} & -M_D & 0 \end{pmatrix},$$

$10_F \to \psi_L, (\psi_R)^\dagger$

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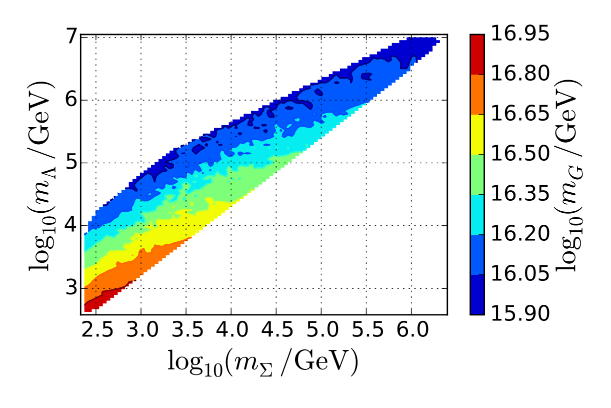

(See previous arXiv:1509.06313 [PRD]): Split-SUSY: like $M_\Phi = 2$ TeV

Not-susy SO(10) $\rightarrow$ SU(3)$_c$ $\times$ SU(2)$_R$ $\times$ SU(2)$_L$ $\times$ U(1)$_Y$ $\times$ U(1)$_{B-L}$

Fermionic triplet dark matter in an SO(10)-inspired left right model, with C. Arbeláez and M. Hirsch, arXiv:1703.08148 [PRD]

Minimal Left-Right Symmetric Standard Model

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	1/2	16
L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126

Left-singlet right-triplet DM

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	1/2	16
L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1132}	1	$(1, 1, 3, 2)$	1/2	126
Ψ_{113-2}	1	$(1, 1, 3, -2)$	1/2	$\overline{126}$

$\Omega h^2 = 0.1199 \pm 0.0027$ at 3σ

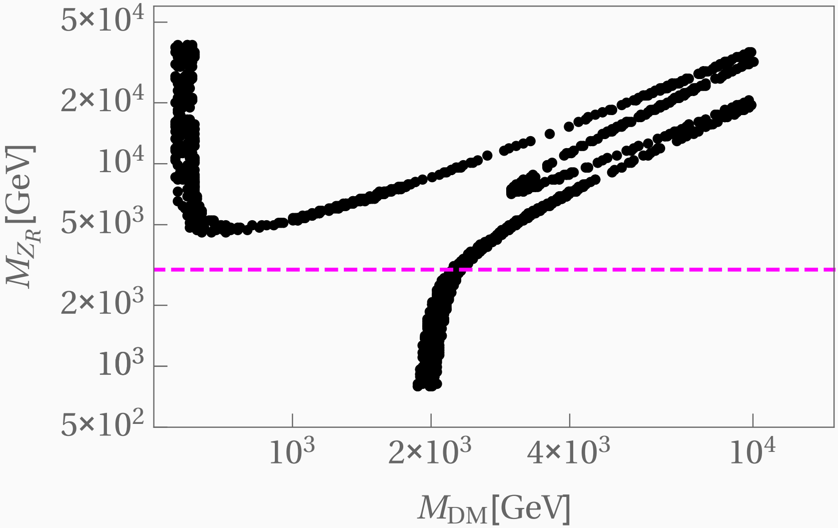

Figure 1: Proper relic density scan: $0.5 < v_R/\text{TeV} < 50$

Mixed Left-singlet right-triplet DM

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
Q^c	3	$(\overline{3}, 1, 2, -\frac{1}{3})$	1/2	16
L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1130}	1	$(1, 1, 3, 0)$	1/2	45
Ψ_{1132}	1	$(1, 1, 3, 2)$	1/2	126
Ψ_{113-2}	1	$(1, 1, 3, -2)$	1/2	$\overline{126}$

Mixed Left-singlet right-triplet DM

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	1/2	16
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L	3	$(1, 2, 1, -1)$	1/2	16
L^c	3	$(1, 1, 2, +1)$	1/2	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1130}	1	$(1, 1, 3, 0)$	1/2	45
Ψ_{1132}	1	$(1, 1, 3, 2)$	1/2	126
Ψ_{113-2}	1	$(1, 1, 3, -2)$	1/2	$\overline{126}$

Setup

$$\Psi_{1132} = \begin{pmatrix} \Psi^{+}/\sqrt{2} & \Psi^{++} \\ \Psi^{0} & -\Psi^{+}/\sqrt{2} \end{pmatrix}, \qquad \bar{\Psi}_{113-2} = \begin{pmatrix} \Psi^{-}/\sqrt{2} & \overline{\Psi}^{0} \\ \Psi^{--} & -\Psi^{-}/\sqrt{2} \end{pmatrix}. \tag{1}$$

$$\begin{aligned} L \supset & M_{11}\,\mathrm{Tr}(\Psi_{1130}\Psi_{1130}) + M_{23}\,\mathrm{Tr}(\Psi_{1132}\bar{\Psi}_{113-2}) \\ & + \lambda_{13}\,\mathrm{Tr}(\Delta_R\bar{\Psi}_{113-2}\Psi_{1130}) + \lambda_{12}\,\mathrm{Tr}(\Delta_R^{\dagger}\Psi_{1132}\Psi_{1130}), \end{aligned} \tag{2}$$

$$\tan\gamma = \frac{\lambda_{13}}{\lambda_{12}}, \qquad \lambda = \sqrt{\lambda_{12}^2 + \lambda_{13}^2}. \tag{3}$$

Blind spot at

$$M_{23}\sin 2\gamma - M_{\mathrm{DM}} = 0 \tag{4}$$

Proper relic density scan

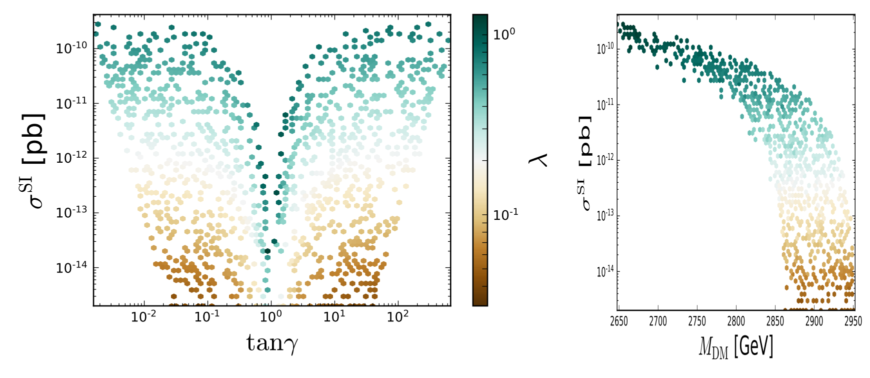


Figure 2: $M_{11} = 50$ TeV $2.7 < M_{23}/\mathrm{TeV} < 3.1$ (Right: $\tan\gamma > 5$)

$\Omega h^2 = 0.1199 \pm 0.0027$ at 3σ

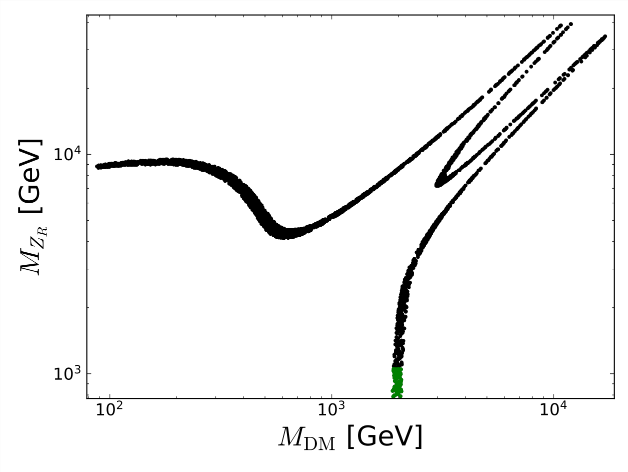


Figure 3:

Direct detection cross section

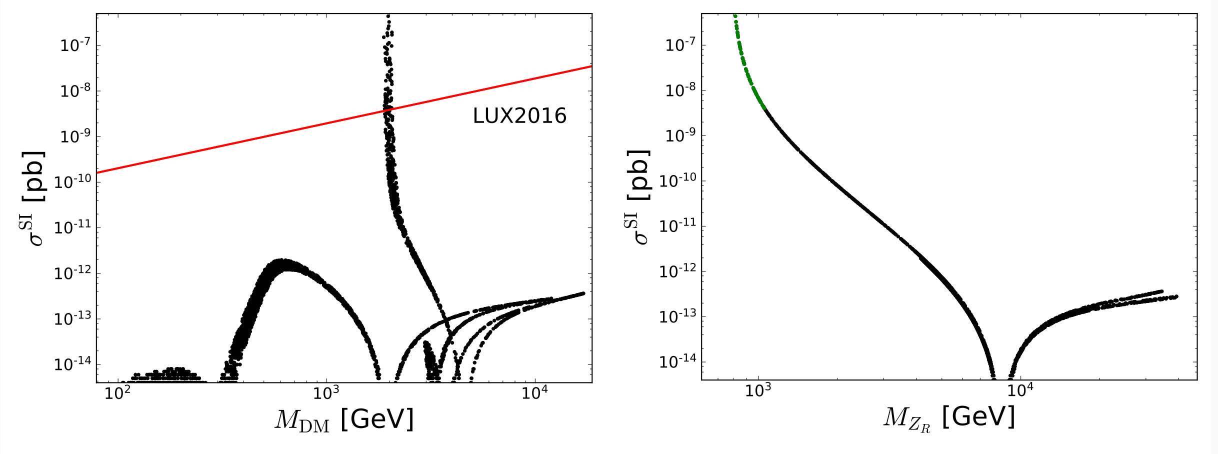


Figure 4: v_R : $[2, 50]$ TeV, M_{23} : $[0.2, 50]$ TeV, M_{11} : 50 TeV, $\tan\gamma = -1$ and $\lambda = 0.14$.

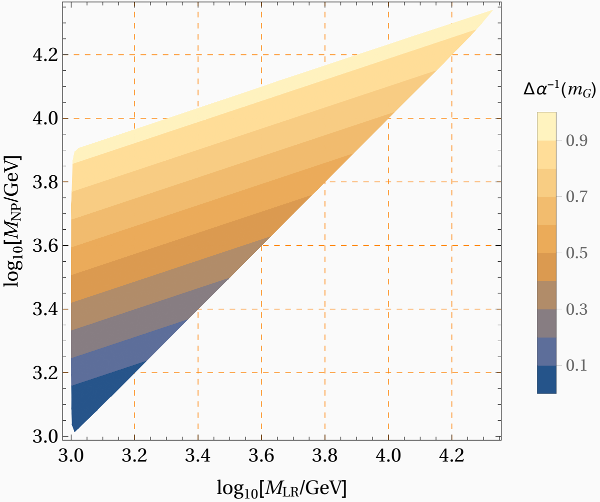
Unification

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	$1/2$	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	$1/2$	16
L	3	$(1, 2, 1, -1)$	$1/2$	16
L^c	3	$(1, 1, 2, +1)$	$1/2$	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1130}	1	$(1, 1, 3, 0)$	$1/2$	45
Ψ_{1132}	1	$(1, 1, 3, 2)$	$1/2$	126
Ψ_{113-2}	1	$(1, 1, 3, -2)$	$1/2$	$\overline{126}$

Unification

Field	Multiplicity	$3_c 2_L 2_R 1_{B-L}$	Spin	SO(10) origin
Q	3	$(3, 2, 1, +\frac{1}{3})$	$1/2$	16
Q^c	3	$(\bar{3}, 1, 2, -\frac{1}{3})$	$1/2$	16
L	3	$(1, 2, 1, -1)$	$1/2$	16
L^c	3	$(1, 1, 2, +1)$	$1/2$	16
Φ	1	$(1, 2, 2, 0)$	0	10
Δ_R	1	$(1, 1, 3, -2)$	0	126
Ψ_{1130}	1	$(1, 1, 3, 0)$	$1/2$	45
Ψ_{1132}	1	$(1, 1, 3, 2)$	$1/2$	126
Ψ_{113-2}	1	$(1, 1, 3, -2)$	$1/2$	$\overline{126}$
Ψ_{1310}	1	$(1, 3, 1, 0)$	$1/2$	45
Ψ_{8110}	1	$(1, 1, 8, 0)$	$1/2$	45
$\Psi_{321\frac{1}{3}}$	1	$(3, 2, 1, 1/3)$	$1/2$	16
$\Psi_{321-\frac{1}{3}}$	1	$(1, 2, 3, -1/3)$	$1/2$	$\overline{16}$

Unification quality



NP: Ψ_{1310}

Remarks

In addition to accommodate usual simplified dark matter models, Left-right symmetric standard models have additional DM portals:

New Δ_R portal for direct detection of left-singlet right-triplet mixed dark matter, in companion with left-singlets charged and doubly charged fermions.

Next: Search for them in compressed spectra scenarios at the LHC

Dirac neutrinos case:

$\mathrm{SO}(10) \to \mathrm{SU}(5) \times U(1)_X$

From 1210.6350 and 1805.02025: $\Delta N_{\text{eff}} = 3\left(T_{\nu_R}/T_{\nu_L}\right)^4$

$$\Gamma_{\nu_R}(T) = n_{\nu_R}(T)\sum_f \langle \sigma_f(\nu_R\bar{\nu}_R \to f\bar{f})v\rangle$$
$$= \sum_f \frac{g^2_{\nu_R}}{n_{\nu_R}} \int \frac{d^3p}{(2\pi)^3}\frac{d^3q}{(2\pi)^3} f_{\nu_R}(p) f_{\nu_R}(q)\sigma_f(s)(1-\cos\theta),$$

$$s = 2pq(1-\cos\theta), \qquad f_{\nu_R}(k) = 1/(e^{k/T}+1)$$

$$n_{\nu_R}(T) = g_{\nu_R}\int \frac{d^3k}{(2\pi)^3} f_{\nu_R}(k), \qquad \text{with } g_{\nu_R} = 2$$

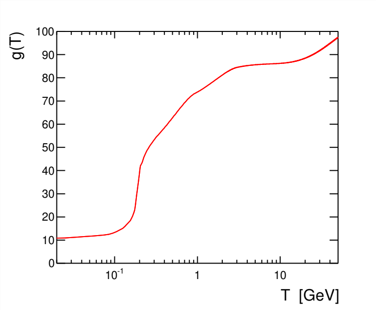
$$\sigma_f(s) \simeq \frac{N_C^f (Q^f_{BL})^2 Q^2 s}{12\pi}\left(\frac{g'}{M_{Z'}}\right)^4, \quad \text{In the limit } M^2_{Z'} \gg s\,.$$

with three right-handed neutrinos, the Hubble parameter is

$$H(T) = \sqrt{\frac{4\pi^3 G_N\,[g(T) + 21/4]}{45}}\,T^2.$$

The right-handed neutrinos decouple when

$$\Gamma_{\nu_R}(T^{\nu_R}_{\text{dec}}) = H(T^{\nu_R}_{\text{dec}}).$$



A. Solaguren-Beascoa, M. C. Gonzalez-Garcia: arXiv:1210.6350 [PLB]

From 1210.6350 and 1805.02025: $\Delta N_{\text{eff}} = 3\left(T_{\nu_R}/T_{\nu_L}\right)^4$

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$$s = 2pq(1-\cos\theta), \qquad f_{\nu_R}(k) = 1/(e^{k/T}+1)$$

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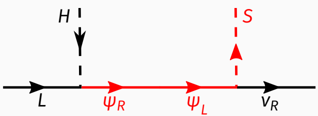
The right-handed neutrinos decouple when

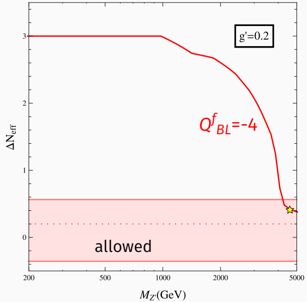
$$\Gamma_{\nu_R}(T^{\nu_R}_{\text{dec}}) = H(T^{\nu_R}_{\text{dec}}).$$

$\mathrm{U}(1)_{B-L} \xrightarrow{\langle S\rangle} Z_3$



ν_{R3}	ν_{R2}	ν_{R1}	ψ_L	ψ_R	S
-4	-4	$+5$	-1	-1	$+3$

E. Ma, R. Srivastava: arXiv:1411.5042 [PLB]



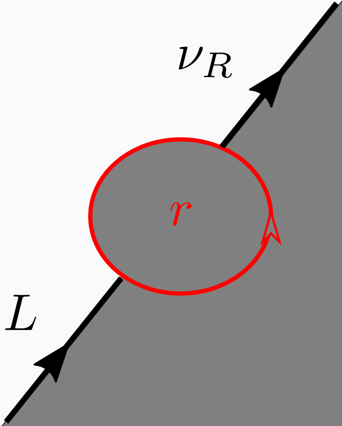
Z.-L. Han, W. Wang: arXiv:1805.02025 [EJPC]

(also: Planck 1807.06209, Riess *et al* 1903.07603)

One-loop realization of $\mathcal{L}_{5-D}$ with total L

$\mathrm{U}(1)_X$

Dirac neutrino masses + Dirac fermion dark matter



$$\mathrm{U}(1)_X$$

Dirac neutrino masses + Dirac fermion dark matter

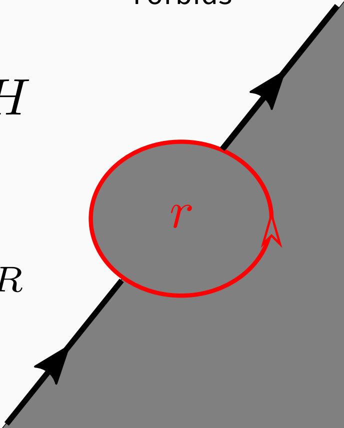
Forbids

$$\nu_R \nu_R$$

$$(\nu_R)^\dagger L H$$

$$\nu_R \psi_R$$

$$(\psi_L)^\dagger \nu_R$$



$$\psi_R \psi_R$$

$$\psi_L \psi_L$$

$\mathrm{U}(1)_X$

Dirac neutrino masses + Dirac fermion dark matter

Forbids

$\nu_R \nu_R$

$(\nu_R)^\dagger L H$

$\nu_R \psi_R$

$(\psi_L)^\dagger \nu_R$

$\psi_R \psi_R$

$\psi_L \psi_L$

r

Allows

$(\nu_R)^\dagger L H S$

$(\psi_R)^\dagger \psi_L$

$(\psi_R)^\dagger \psi_L S$

$$\mathrm{U}(1)_{B-L} \xrightarrow[\langle S \rangle]{} Z_N \qquad N \neq 2$$

Dirac neutrino masses + Dirac fermion dark matter

Forbids

$\nu_R \nu_R$

$(\nu_R)^\dagger L H$

$\nu_R \psi_R$

$(\psi_L)^\dagger \nu_R$

$\psi_R \psi_R$

$\psi_L \psi_L$

r

Allows

$X(L) \neq 0$

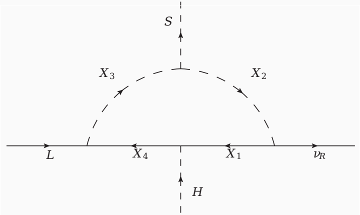
$(\nu_R)^\dagger L H S$

$(\psi_R)^\dagger \psi_L$

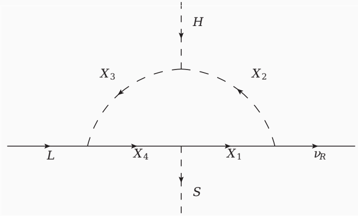
$(\psi_R)^\dagger \psi_L S$

normalized to$X(L) = -1$

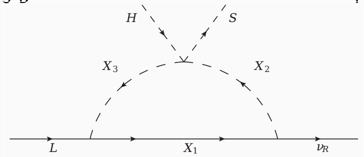
One loop topologies $U(1)_{B-L} \oplus Z_2 \oplus Z_2$



T1-3-D

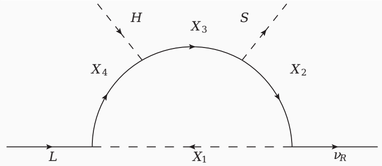


T1-3-E

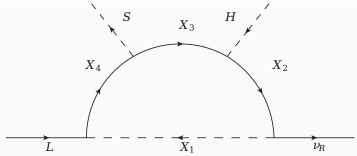


T3-1-A

Chang-Yuan Yao and Gui-Jun Ding, arXiv:1802.05231 [PRD]

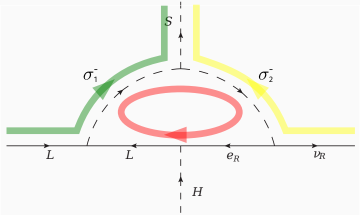


T1-2-A

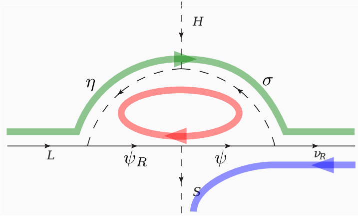


T1-2-B

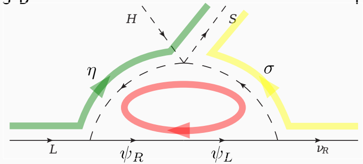
One loop topologies U(1)$_{B-L}$ only!



T1-3-D



T1-3-E



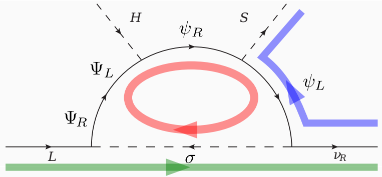
T3-1-A

$\psi_{L,R}$ $\rightarrow$Singlet fermions

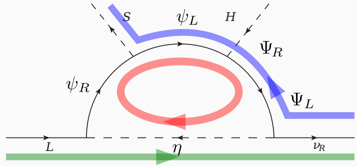
$\Psi_{L,R}$ $\rightarrow$Vector-like doublet fermions

σ $\rightarrow$Singlet scalar

with J. Calle, C. Yaguna, and O. Zapata, arXiv:1812.05523 [PRD]

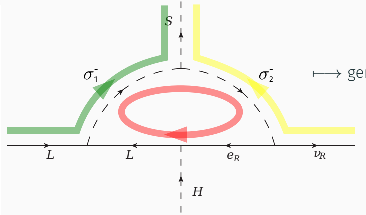


T1-2-A



T1-2-B

One loop topologies $U(1)_{B-L}$ only! with J. Calle, C. Yaguna, and O. Zapata, arXiv:1812.05523 [PRD]

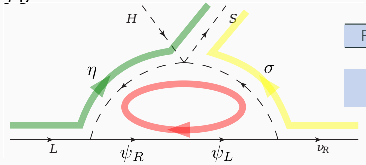


T1-3-D

$\longmapsto$ generalization to two and three loops: Shaikh Saad arXiv:1902.07259 [NPB]

$\psi_{L,R}$ $\to$Singlet fermions (vector-like)

σ $\to$Singlet scalar



T3-1-A

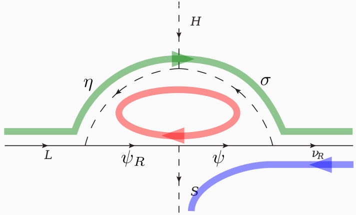
Fields: f_i	$(\nu_{R3})^\dagger$	$(\nu_{R2})^\dagger$	$(\nu_{R1})^\dagger$	ψ_L	$(\psi_R)^\dagger$	S
(A)	$+4$	$+4$	-5	$-r$	r	$+3$

Anomaly cancellation conditions

$$\sum_i f_i = 3$$

$$\sum_i f_i^3 = 3$$

One loop topologies $U(1)_{B-L}$ only! with J. Calle, C. Yaguna, and O. Zapata, arXiv:1812.05523 [PRD]



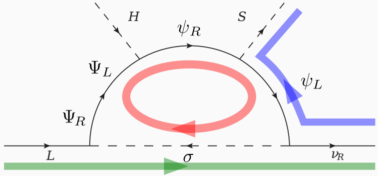
T1-3-E

$\psi_{L,R}$ →Singlet fermions (quiral)

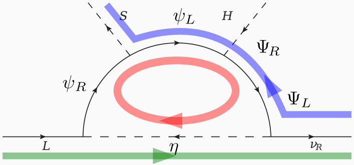
$\Psi_{L,R}$ →Vector-like doublet fermions

σ →Singlet scalar

Fields: f_i	$(\nu_{R3})^\dagger$	$(\nu_{R2})^\dagger$	$(\nu_{R1})^\dagger$	ψ_L	$(\psi_R)^\dagger$	S
(A)	$+4$	$+4$	-5	$-r$	r	$+3$
(B)	$+\frac{8}{5}$	$+\frac{8}{5}$	$+\frac{2}{5}$	$\frac{7}{5}$	$-\frac{10}{5}$	$+\frac{3}{5}$



T1-2-A



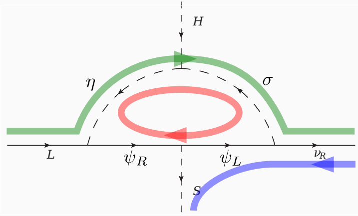
T1-2-B

Anomaly cancellation conditions

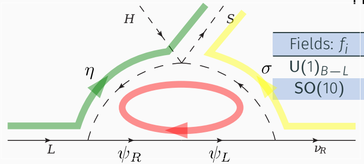
$$\sum_i f_i = 3$$

$$\sum_i f_i^3 = 3$$

Scotogenic Dirac from SO(10) E. Ma arXiv:1901.09091 [PLB]



T1-3-E



T3-1-A

$\psi_{L,R}$ $\longrightarrow$Singlet fermions (quiral)

σ $\longrightarrow$Singlet scalar

η $\longrightarrow$Doublet scalar

Fields: f_i	$(\nu_{R3})^\dagger$	$(\nu_{R2})^\dagger$	$(\nu_{R1})^\dagger$	ψ_L	$(\psi_R)^\dagger$	σ	η	S
$U(1)_{B-L}$	$+4$	$+4$	-5	$-r$	r	$4-r$	$1-r$	$+3$
SO(10)	16	16	16	126	126^*	16	144	672

Anomaly cancellation conditions

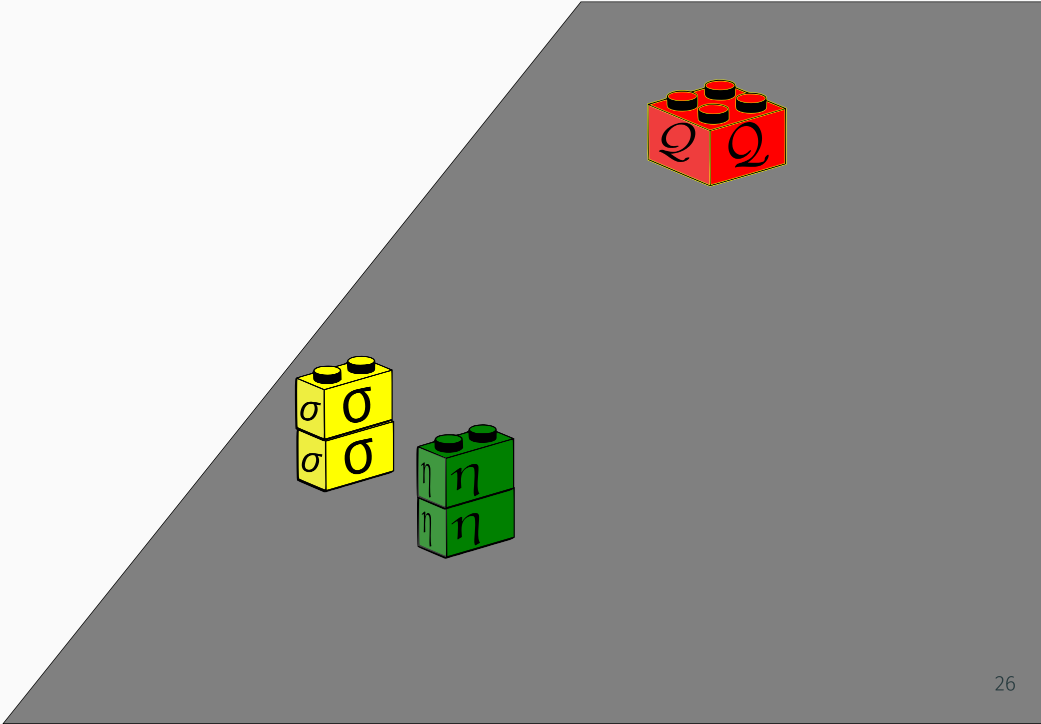
$$\sum_i f_i = 3$$

$$\sum_i f_i^3 = 3$$

SM

+

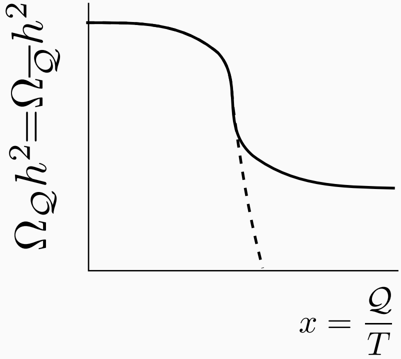
$(\nu_R)^\dagger_i$



(Switch to Dirac fermions)

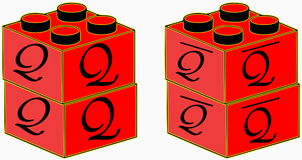
Because $\mathcal{Q}$ is a Dirac fermion, $\mathcal{Q}\mathcal{Q}$ is also stable

$$\mathcal{Q}\mathcal{Q} \not\to g\,, \qquad \overline{\mathcal{Q}\mathcal{Q}} \not\to g\,.$$



Step one

$\mathcal{Q}$-onlyum

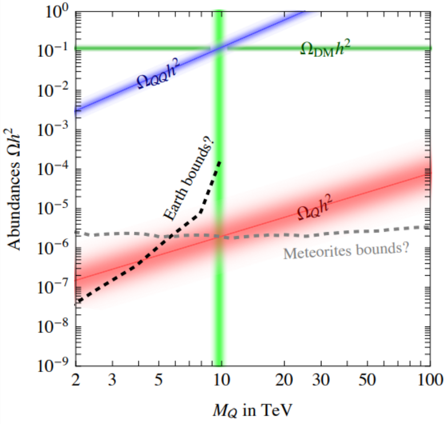


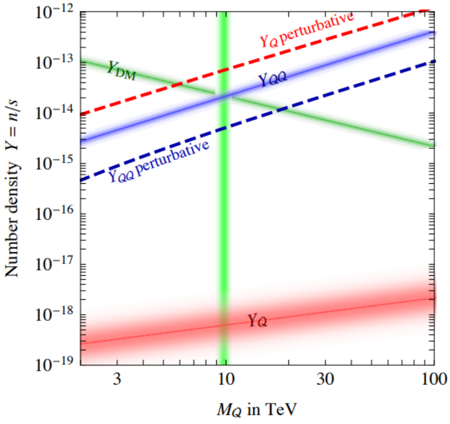
Step two

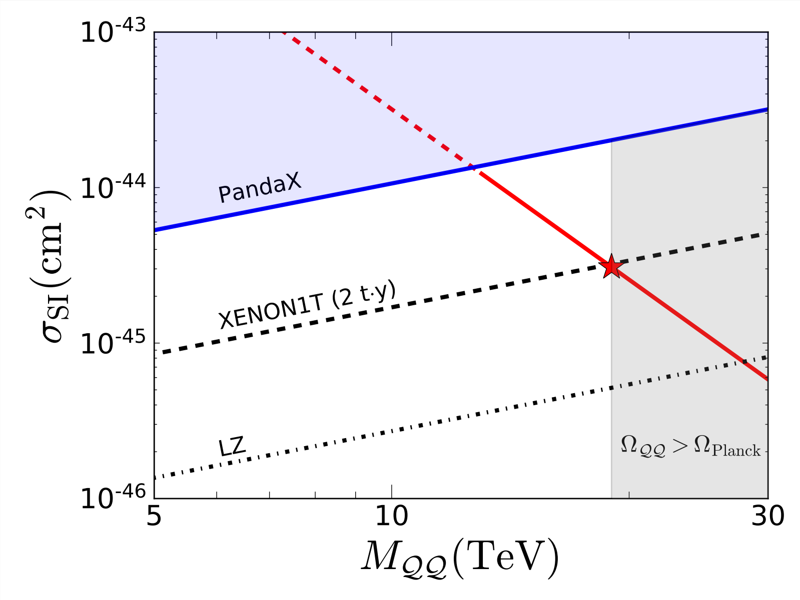
Relic abundance. From: http://bit.ly/Mitridate_53d_53rd_Rencontres_de_Moriond

$$M_Q \simeq 9.5\,\text{TeV}$$

$$\Omega_{\rm hyb} \sim 10^{-5}\Omega_{\rm DM}$$







10^{-43}
10^{-44}
10^{-45}
10^{-46}
$\sigma_{\rm SI}\,(\rm cm^2)$
PandaX
XENON1T (2 t·y)
LZ
$\Omega_{QQ} > \Omega_{\rm Planck}$
5
10
30
$M_{QQ}\,(\rm TeV)$

SD^{3}M+σ_i $(i = 1, 2)$ with J. Calle, C. Yaguna, and O. Zapata, arXiv:1812.05523 [PRD]

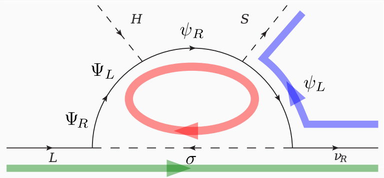
$\psi_{L,R}$ $\rightarrow$Singlet fermions (quiral)

$\Psi_{L,R}$ $\rightarrow$Vector-like doublet fermions : 10/5

σ $\rightarrow$Singlet scalar : 15/5

Fields: f_i	$(\nu_{R3})^\dagger$	$(\nu_{R2})^\dagger$	$(\nu_{R1})^\dagger$	ψ_L	$(\psi_R)^\dagger$	S	Ψ_L	$\widetilde{(\Psi_R)}$	σ
U$(1)_{B-L}$	$+\frac{8}{5}$	$+\frac{8}{5}$	$+\frac{2}{5}$	$\frac{7}{5}$	$-\frac{10}{5}$	$+\frac{3}{5}$	$\frac{10}{5}$	$-\frac{10}{5}$	$\frac{7}{5}$
SO(10)	16	16	16	45	45	126	10	10	16

Second Scotogenic Dirac from SO(10) E. Ma
arXiv:1901.09091 [PLB]



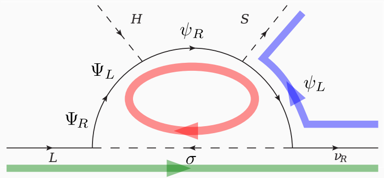
T1-2-A

SD^{3}M+σ_i $(i = 1, 2)$

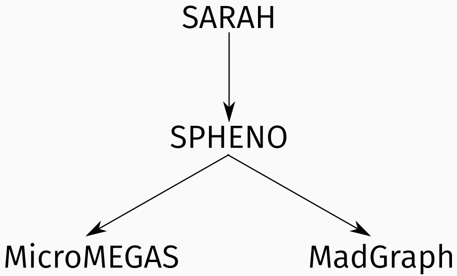
$M_\psi = h_1 \langle S \rangle$, $y_2 = 0$:

$$\mathcal{L} = \mathcal{L}_{\text{SD}^3\text{M}} + h_3^{ia}\widetilde{(\Psi_R)} \cdot L_i\, \sigma_a + h_2^{\beta a}\, (\nu_{R\beta})^\dagger\, \psi_L\, \sigma_a^* - V(\sigma_a, S, H)\,.$$

with A.F Rivera, W. Tangarife, arXiv:1906.09685



T1-2-A



Singlet-Doublet Dirac Dark Matter (SD^{3}M) By Carlos E. Yaguna. arXiv:1510.06151 [PRD].

The model extends the standard model (SM) particle content with Dirac Fermions: from SU(2) doublets of Weyl fermions: $\Psi_L = \left(\Psi_L^0,\, \Psi_L^-\right)^\mathsf{T}$, $\widetilde{(\Psi_R)} = \left((\Psi_R^-)^\dagger,\, -(\Psi_R^0)^\dagger\right)^\mathsf{T}$ and singlet Weyl fermions ψ_{LR} that interact among themselves and with the SM fields

$$\mathcal{L} \supset M_\psi\,(\psi_R)^\dagger\,\psi_L + M_\Psi\widetilde{(\Psi_R)}\cdot\Psi_L + y_1\,(\psi_R)^\dagger\,\Psi_L\cdot H + y_2\widetilde{(\Psi_R)}\cdot\widetilde{H}\psi_L + \text{h.c} \tag{5}$$

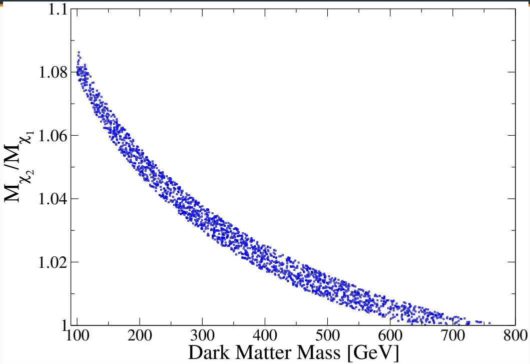
Four free parameters:

$$M_\psi, M_\Psi \;<2\text{ GeV}\,, \qquad\qquad y_1,\; y_2 > 10^{-6} \tag{6}$$

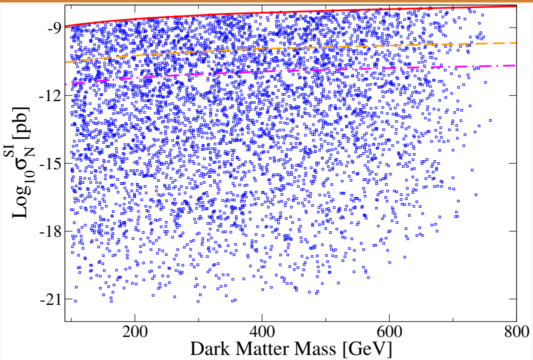
Two neutral Dirac fermion eigenstates:

$$M = \begin{pmatrix} M_\psi & y_2 v/\sqrt{2} \\ y_1 v/\sqrt{2} & M_D \end{pmatrix}, \qquad\qquad M_{\text{diag}} = \begin{pmatrix} M_{\chi_1} & 0 \\ 0 & M_{\chi_2} \end{pmatrix} = U_L^\dagger M U_R \tag{7}$$

SD^{3}M By Carlos E. Yaguna. arXiv:1510.06151 [PRD].

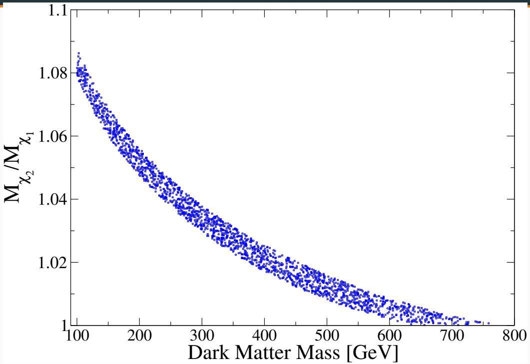


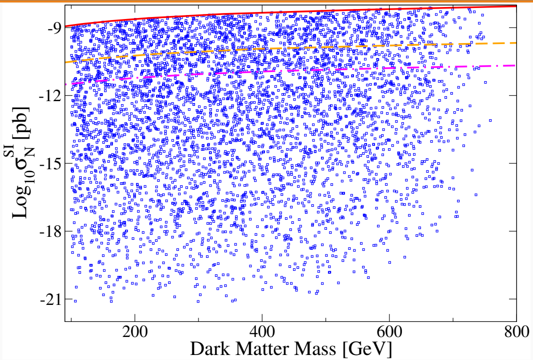
Compressed spectra region



LUX - XENON1T - LZ

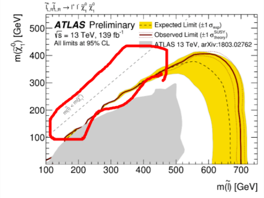
SD^{3}M By Carlos E. Yaguna. arXiv:1510.06151 [PRD].



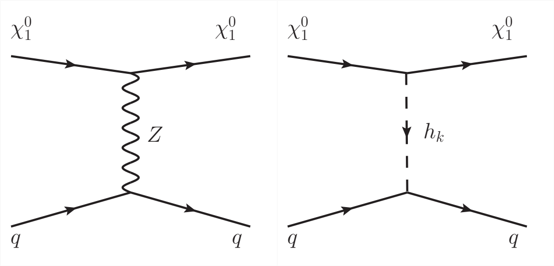


Compressed spectra region

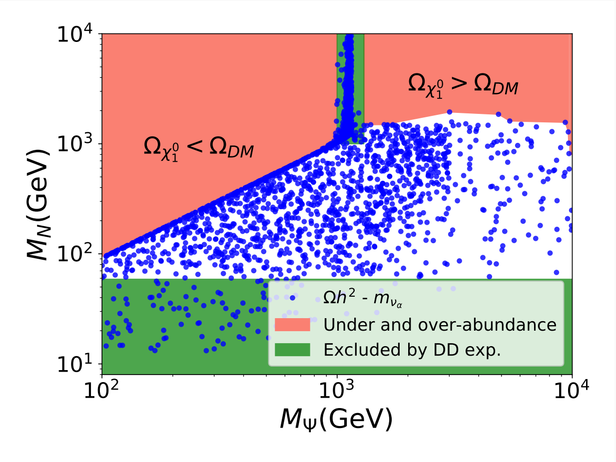
LUX - XENON1T - LZ



Spin independent (SI) direct detection cross section



Decoupled Z' limit



Vector SI (blue points) and scalar SI (green points)

Conclusions

A single $U(1)$ symmetry to explain both the smallnes of neutrino masses and the stability of fermion dark matter

Conclusions

A single $U(1)$ symmetry to explain both the smallnes of neutrino masses and the stability of fermion dark matter

Neutrino masses and DM

- Spontaneously broken $U(1)_X$ generates a radiative neutrino masses
- A remnant symmetry makes the lightest field circulating the loop stable and good dark matter candidate.
- For T1-2-A: Either Singet Doublet (Dirac) Dark Matter or Singlet Scalar Dark Matter with extra scalar and vector portal
- With relaxed direct detection constraints

Thanks!