

MLMC: Machine Learning Monte Carlo for Lattice Gauge Theory



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[saforem2/lattice23](#)

Overview

1. Background: {MCMC, HMC}

- Leapfrog Integrator
- Issues with HMC
- Can we do better?

2. L2HMC: Generalizing MD

- 4D $SU(3)$ Model
- Results

3. References

4. Extras

Background: MCMC

Markov Chain Monte Carlo (MCMC)

Goal

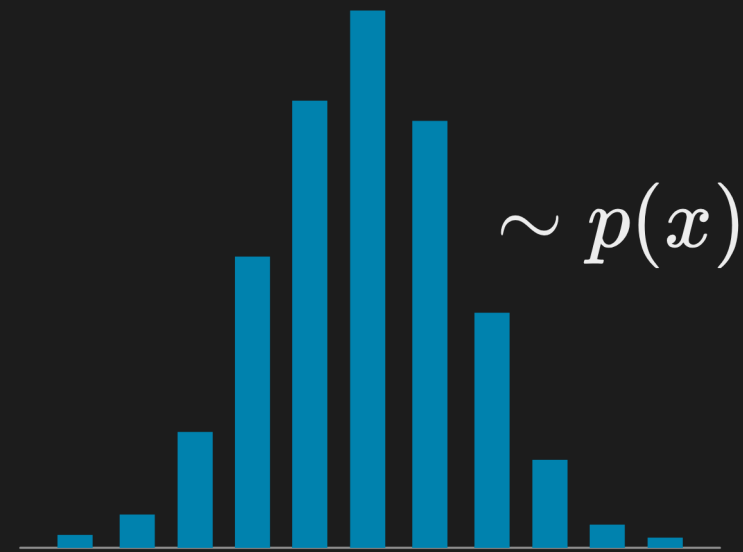
Generate **independent** samples $\{x_i\}$, such that¹

$$\{x_i\} \sim p(x) \propto e^{-S(x)}$$

where $S(x)$ is the *action* (or potential energy)

- Want to calculate observables \mathcal{O} :

$$\langle \mathcal{O} \rangle \propto \int [\mathcal{D}x] \mathcal{O}(x) p(x)$$



If these were *independent*, we could approximate: $\langle \mathcal{O} \rangle \simeq \frac{1}{N} \sum_{n=1}^N \mathcal{O}(x_n)$

$$\sigma_{\mathcal{O}}^2 = \frac{1}{N} \text{Var}[\mathcal{O}(x)] \implies \sigma_{\mathcal{O}} \propto \frac{1}{\sqrt{N}}$$

1. Here, \sim means “is distributed according to”

Markov Chain Monte Carlo (MCMC)

Goal

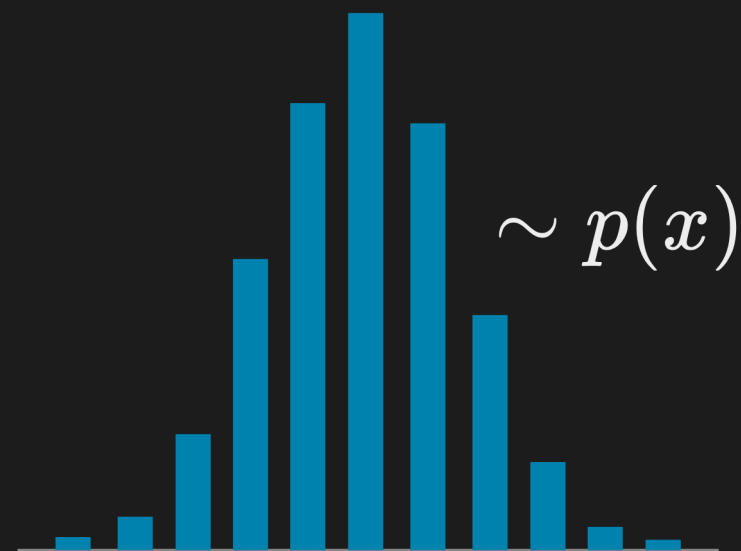
Generate **independent** samples $\{x_i\}$, such that¹

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where $S(x)$ is the *action* (or potential energy)

- Want to calculate observables \mathcal{O} :

$$\langle \mathcal{O} \rangle \propto \int [\mathcal{D}x] \mathcal{O}(x) p(x)$$



Instead, nearby configs are **correlated**, and we incur a factor of $\tau_{\text{int}}^{\mathcal{O}}$:

$$\sigma_{\mathcal{O}}^2 = \frac{\tau_{\text{int}}^{\mathcal{O}}}{N} \text{Var}[\mathcal{O}(x)]$$

1. Here, \sim means “is distributed according to”

Background: HMC

Hamiltonian Monte Carlo (HMC)

- Want to (sequentially) construct a chain of states:

$$x_0 \rightarrow x_1 \rightarrow x_i \rightarrow \cdots \rightarrow x_N$$

such that, as $N \rightarrow \infty$:

$$\{x_i, x_{i+1}, x_{i+2}, \cdots, x_N\} \xrightarrow{N \rightarrow \infty} p(x) \propto e^{-S(x)}$$

</> Trick

- Introduce **fictitious** momentum $v \sim \mathcal{N}(0, 1)$
 - Normally distributed **independent** of x , i.e.

$$p(x, v) = p(x) p(v) \propto e^{-S(x)} e^{-\frac{1}{2}v^T v} = e^{-[S(x) + \frac{1}{2}v^T v]} = e^{-H(x, v)}$$

Hamiltonian Monte Carlo (HMC)

- **Idea:** Evolve the (\dot{x}, \dot{v}) system to get new states $\{x_i\}$!
- Write the **joint distribution** $p(x, v)$:

$$p(x, v) \propto e^{-S[x]} e^{-\frac{1}{2}v^T v} = e^{-H(x, v)}$$

Hamiltonian Dynamics

$$H = S[x] + \frac{1}{2}v^T v \implies$$

$$\dot{x} = +\partial_v H, \quad \dot{v} = -\partial_x H$$

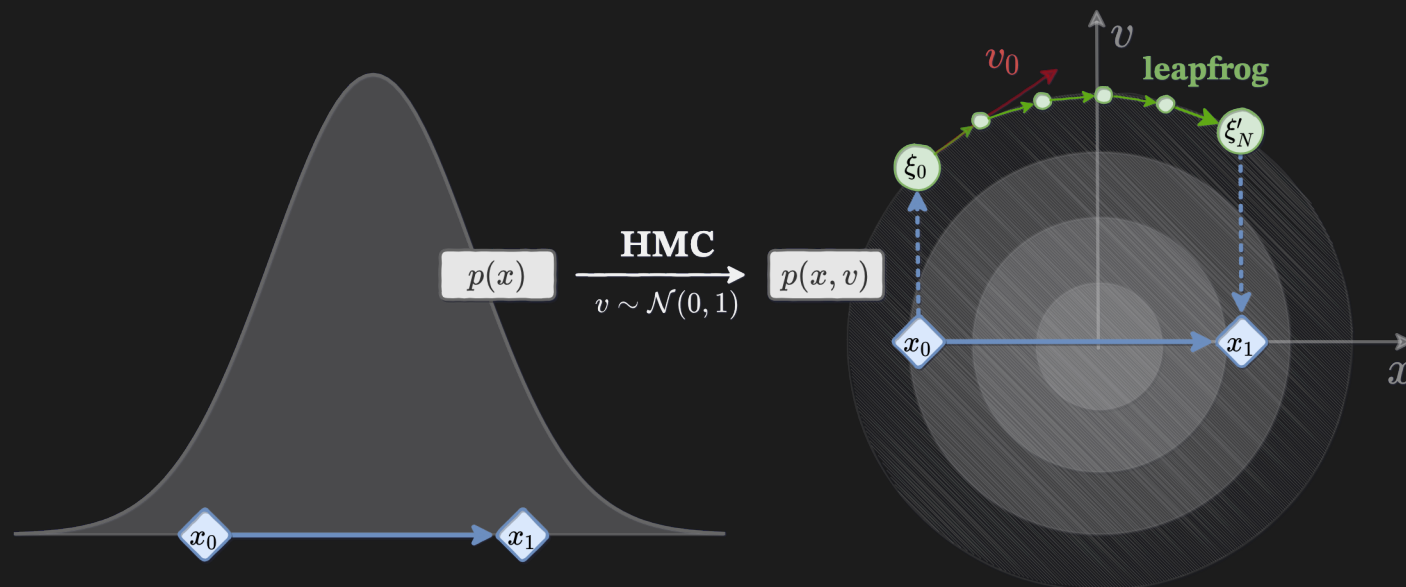


Figure 1: Overview of HMC algorithm

Leapfrog Integrator (HMC)

</> Hamiltonian Dynamics

$$(\dot{x}, \dot{v}) = (\partial_v H, -\partial_x H)$$

⌚ Leapfrog Step

input $(x, v) \rightarrow (x', v')$ output

$$\tilde{v} := \Gamma(x, v) = v - \frac{\varepsilon}{2} \partial_x S(x)$$

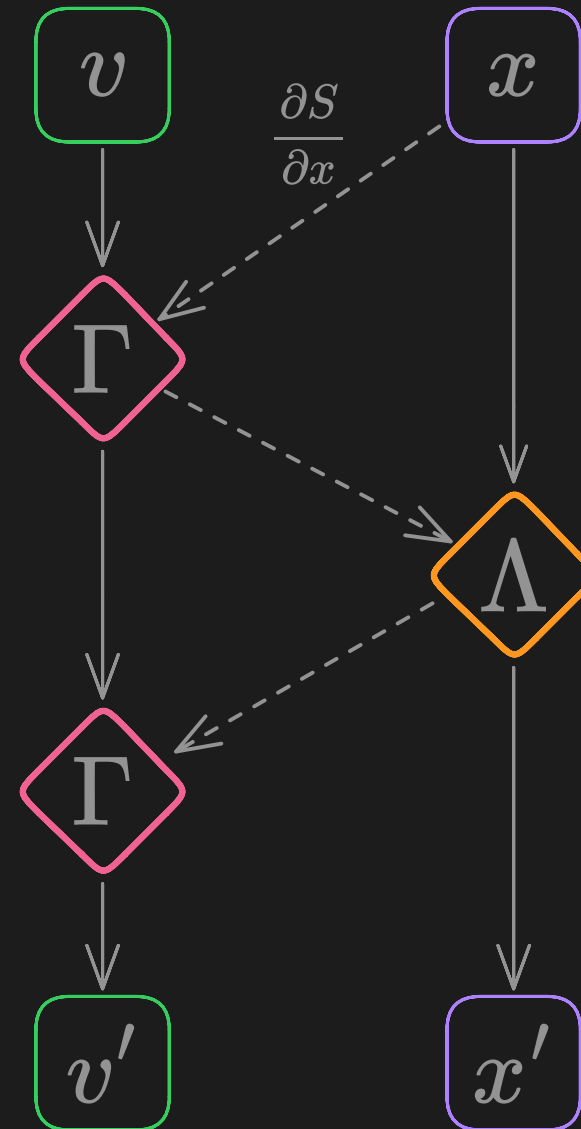
$$x' := \Lambda(x, \tilde{v}) = x + \varepsilon \tilde{v}$$

$$v' := \Gamma(x', \tilde{v}) = \tilde{v} - \frac{\varepsilon}{2} \partial_x S(x')$$

🚨 Warning!

Resample $v_0 \sim \mathcal{N}(0, 1)$
at the **beginning** of each trajectory

Note: $\partial_x S(x)$ is the *force*



HMC Update

- We build a trajectory of N_{LF} **leapfrog steps**¹

$$(x_0, v_0) \rightarrow (x_1, v_1) \rightarrow \cdots \rightarrow (x', v')$$

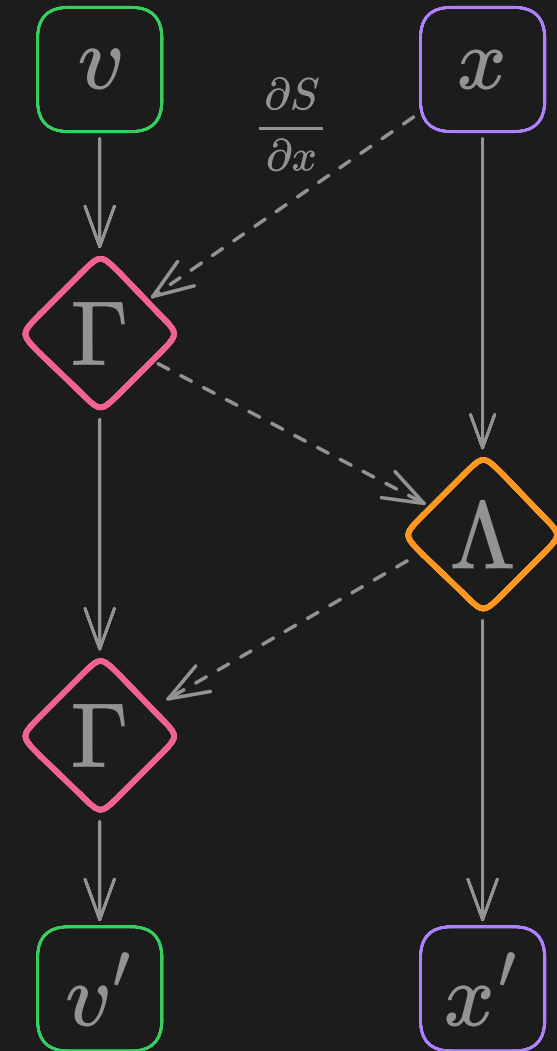
- And propose x' as the next state in our chain

$$\Gamma : (x, v) \rightarrow v' := v - \frac{\varepsilon}{2} \partial_x S(x)$$

$$\Lambda : (x, v) \rightarrow x' := x + \varepsilon v$$

- We then accept / reject x' using Metropolis-Hastings criteria,

$$A(x'|x) = \min \left\{ 1, \frac{p(x')}{p(x)} \left| \frac{\partial x'}{\partial x} \right| \right\}$$

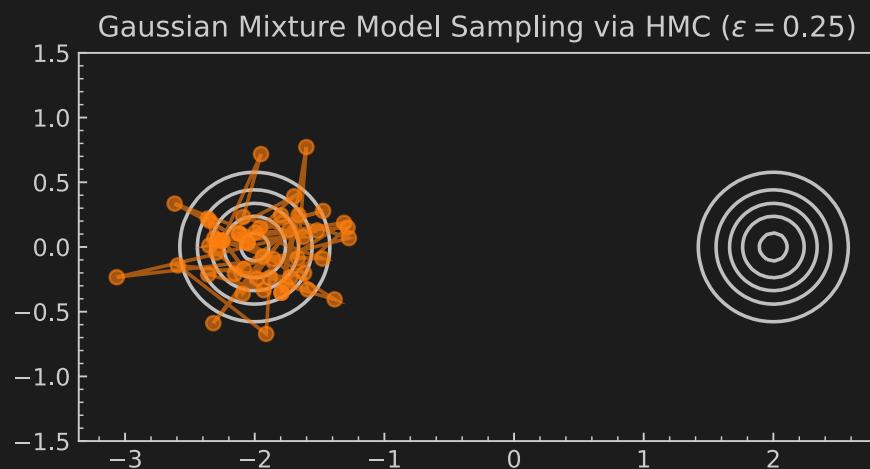


1. We **always** start by resampling the momentum, $v_0 \sim \mathcal{N}(0, 1)$

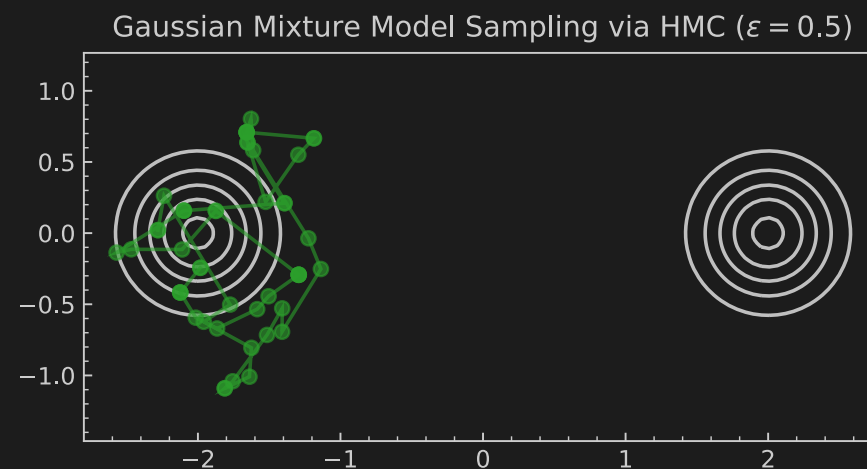
HMC Demo

Figure 2: HMC Demo

- What do we want in a good sampler?
 - **Fast mixing** (small autocorrelations)
 - **Fast burn-in** (quick convergence)
- Problems with HMC:
 - Energy levels selected randomly → **slow mixing**
 - Cannot easily traverse low-density zones → **slow convergence**



HMC Samples with $\epsilon = 0.25$



HMC Samples with $\epsilon = 0.5$

Figure 3: HMC Samples generated with varying step sizes ϵ

Topological Freezing

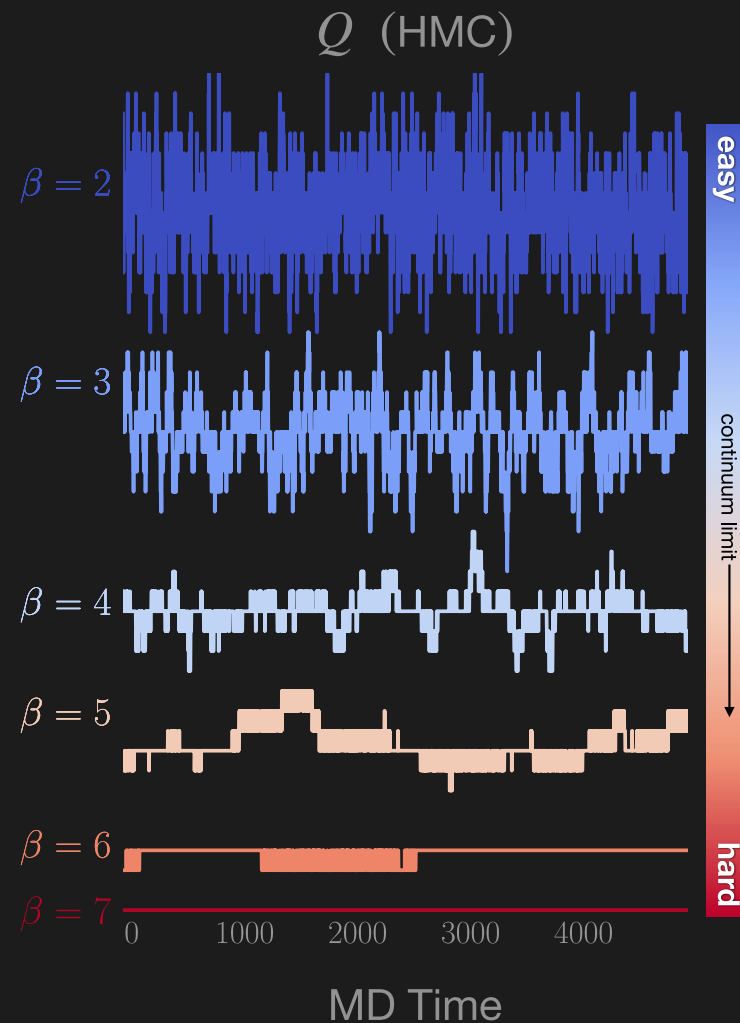
Topological Charge:

$$Q = \frac{1}{2\pi} \sum_P [x_P] \in \mathbb{Z}$$

note: $[x_P] = x_P - 2\pi \left\lfloor \frac{x_P + \pi}{2\pi} \right\rfloor$

🔥 Critical Slowing Down

- Q gets stuck!
 - as $\beta \rightarrow \infty$:
 - $Q \rightarrow \text{const.}$
 - $\delta Q = (Q^* - Q) \rightarrow 0 \implies$
 - # configs required to estimate errors **grows exponentially:** $\tau_{\text{int}}^Q \rightarrow \infty$



Note $\delta Q \rightarrow 0$ at increasing β

Can we do better?

- Introduce two (**invertible NNs**) **vNet** and **xNet**¹:
 - **vNet**: $(x, F) \longrightarrow (s_v, t_v, q_v)$
 - **xNet**: $(x, v) \longrightarrow (s_x, t_x, q_x)$
- Use these (s, t, q) in the *generalized* MD update:
 - $\Gamma_{\theta}^{\pm} : (x, v) \xrightarrow{s_v, t_v, q_v} (x, v')$
 - $\Lambda_{\theta}^{\pm} : (x, v) \xrightarrow{s_x, t_x, q_x} (x', v)$

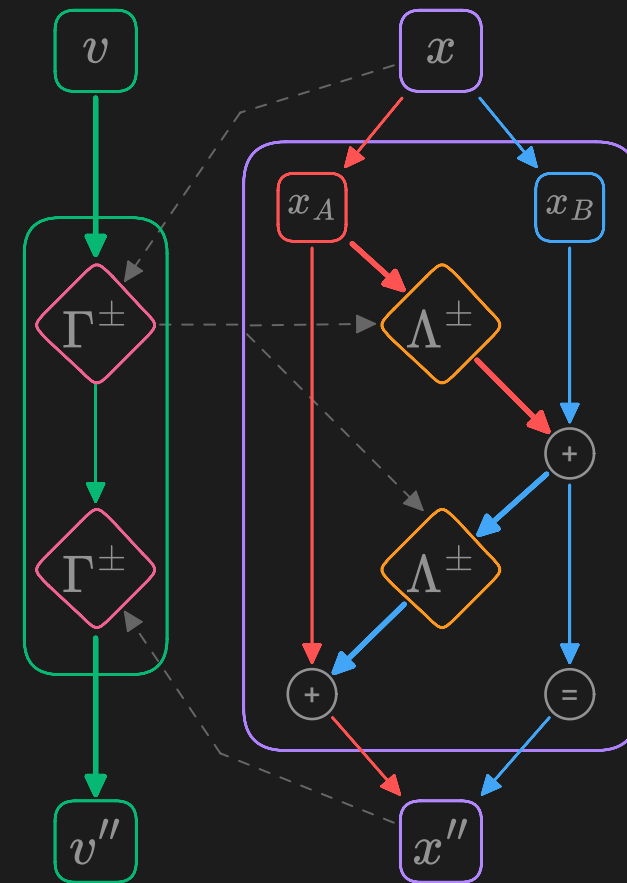



Figure 4: Generalized MD update where $\Lambda_{\theta}^{\pm}, \Gamma_{\theta}^{\pm}$ are **invertible NNs**

1. L2HMC:  (Foreman, Jin, and Osborn 2021, 2022)

L2HMC: Generalizing the MD Update

L2HMC Update

- Introduce $d \sim \mathcal{U}(\pm)$ to determine the direction¹ of our update

- $v' = \Gamma^\pm(x, v)$ update v
- $x' = x_B + \Lambda^\pm(x_A, v')$ update first **half**: x_A
- $x'' = x'_A + \Lambda^\pm(x'_B, v')$ update other half: x_B
- $v'' = \Gamma^\pm(x'', v')$ update v

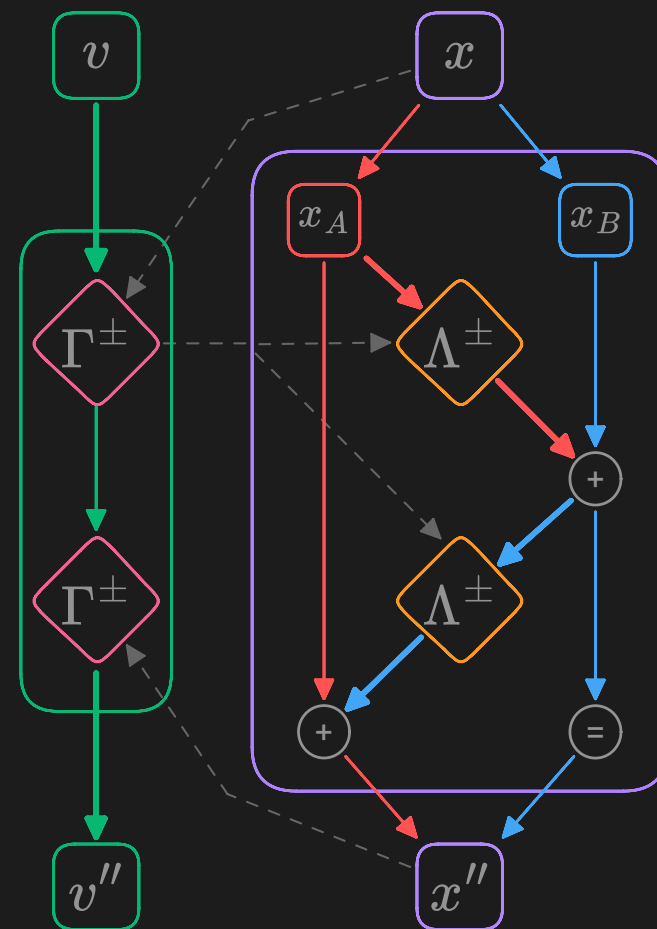


Figure 5: Generalized MD update with $\Lambda_\theta^\pm, \Gamma_\theta^\pm$
invertible NNs

- d resampled at start of each trajectory, to ensure reversibility

L2HMC: Leapfrog Layer

1. Update \mathbf{v} :

$$\mathbf{v}' = \Gamma^\pm[\mathbf{v}; \zeta_{\mathbf{v}}]$$

2. Update half of \mathbf{x} via $\bar{\mathbf{m}}_k \odot \mathbf{x}_k$:

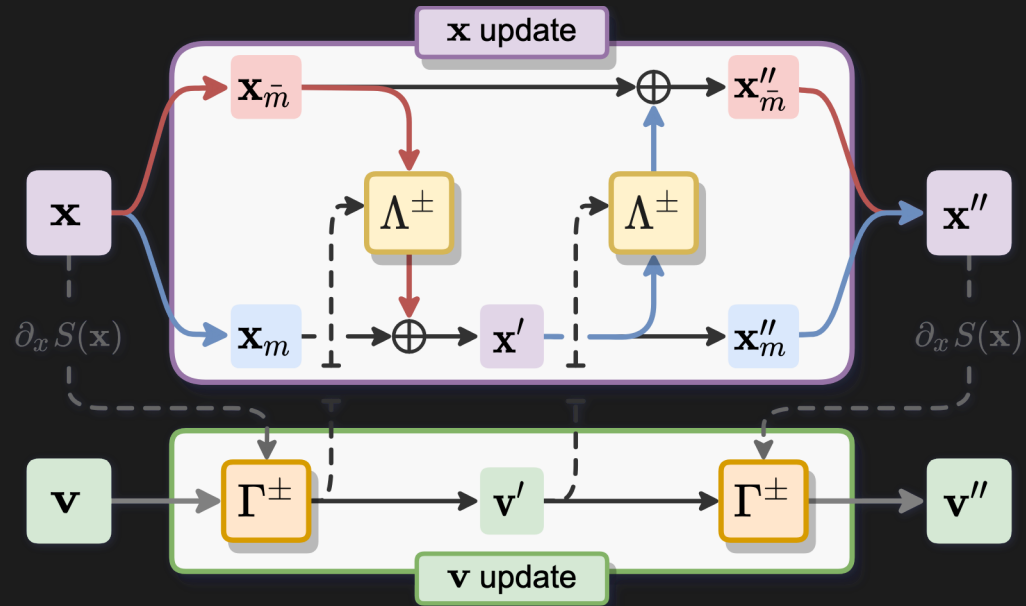
$$\mathbf{x}' = \mathbf{x}_m + \bar{\mathbf{m}} \odot \Lambda^\pm[\mathbf{x}_{\bar{\mathbf{m}}}; \zeta_{\bar{\mathbf{x}}_k}]$$

3. Update (other) half via $\mathbf{m}^k \odot \mathbf{x}'_k$:

$$\mathbf{x}'' = \mathbf{x}'_m + \bar{\mathbf{m}} \odot \Lambda^\pm[\mathbf{x}'_m; \zeta_{\mathbf{x}'}]$$

4. Half-step full \mathbf{v} update:

$$\mathbf{v}'' = \Gamma^\pm[\mathbf{v}'; \zeta_{\mathbf{v}'}]$$



$$\Gamma^+[\mathbf{v}_k; \zeta_{\mathbf{v}}] \equiv \mathbf{v}_k \odot \exp\left(\frac{\varepsilon_{\mathbf{v}}^k}{2} s_{\mathbf{v}}^k(\zeta_{\mathbf{v}_k})\right) - \frac{\varepsilon_{\mathbf{v}}^k}{2} \left[\partial_x S(x_k) \odot \exp(\varepsilon_{\mathbf{v}}^k q_{\mathbf{v}}^k(\zeta_{\mathbf{v}_k})) + t_{\mathbf{v}}^k(\zeta_{\mathbf{v}_k}) \right]$$

trainable step sizes

$$\Lambda^+[\bar{\mathbf{x}}_k; \zeta_{\bar{\mathbf{x}}_k}] \equiv \bar{\mathbf{x}}_k \odot \exp(\varepsilon_{\bar{\mathbf{x}}}^k s_{\bar{\mathbf{x}}}^k(\zeta_{\bar{\mathbf{x}}_k})) + \varepsilon_{\bar{\mathbf{x}}}^k \left[v'_k \odot \exp(\varepsilon_{\bar{\mathbf{x}}}^k q_{\bar{\mathbf{x}}}^k(\zeta_{\bar{\mathbf{x}}_k})) + t_{\bar{\mathbf{x}}}^k(\zeta_{\bar{\mathbf{x}}_k}) \right]$$

Algorithm

1. **input:** x

- Resample: $v \sim \mathcal{N}(0, 1)$; $d \sim \mathcal{U}(\pm)$
- Construct initial state: $\xi = (x, v, \pm)$

2. **forward:** Generate proposal ξ' by passing initial ξ through N_{LF} leapfrog layers

$$\xi \xrightarrow{\text{LF layer}} \xi_1 \longrightarrow \cdots \longrightarrow \xi_{N_{\text{LF}}} = \xi' := (x'', v'')$$

- Accept / Reject:

$$A(\xi'|\xi) = \min \left\{ 1, \frac{\pi(\xi')}{\pi(\xi)} |\mathcal{J}(\xi', \xi)| \right\}$$

3. **backward** (if training):

- Evaluate the **loss function**¹ $\mathcal{L} \leftarrow \mathcal{L}_\theta(\xi', \xi)$ and backprop

4. **return:** x_{i+1}

Evaluate MH criteria (1) and return accepted config,

$$x_{i+1} \leftarrow \begin{cases} x'' & \text{w/ prob } A(\xi''|\xi) \quad \checkmark \\ x & \text{w/ prob } 1 - A(\xi''|\xi) \quad \otimes \end{cases}$$

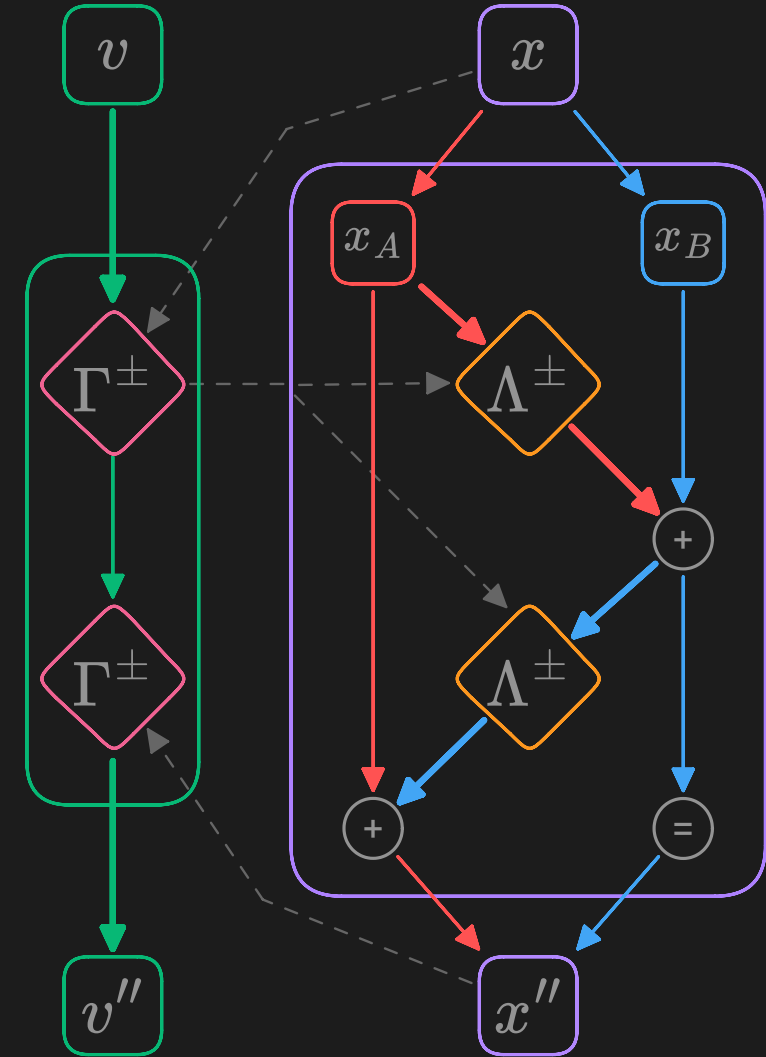


Figure 6: **Leapfrog Layer** used in generalized MD update

1. For simple $\mathbf{x} \in \mathbb{R}^2$ example, $\mathcal{L}_\theta = A(\xi^*|\xi) \cdot (\mathbf{x}^* - \mathbf{x})^2$

4D $SU(3)$ Model

🌀 Link Variables

- Write link variables $U_\mu(x) \in SU(3)$:

$$\begin{aligned} U_\mu(x) &= \exp \left[i \omega_\mu^k(x) \lambda^k \right] \\ &= e^{iQ}, \quad \text{with } Q \in \mathfrak{su}(3) \end{aligned}$$

where $\omega_\mu^k(x) \in \mathbb{R}$, and λ^k are the generators of $SU(3)$

</> Conjugate Momenta

- Introduce $P_\mu(x) = P_\mu^k(x) \lambda^k$ conjugate to $\omega_\mu^k(x)$

🔥 Wilson Action

$$S_G = -\frac{\beta}{6} \sum \text{Tr} [U_{\mu\nu}(x) + U_{\mu\nu}^\dagger(x)]$$

where $U_{\mu\nu}(x) = U_\mu(x) U_\nu(x + \hat{\mu}) U_\mu^\dagger(x + \hat{\nu}) U_\nu^\dagger(x)$

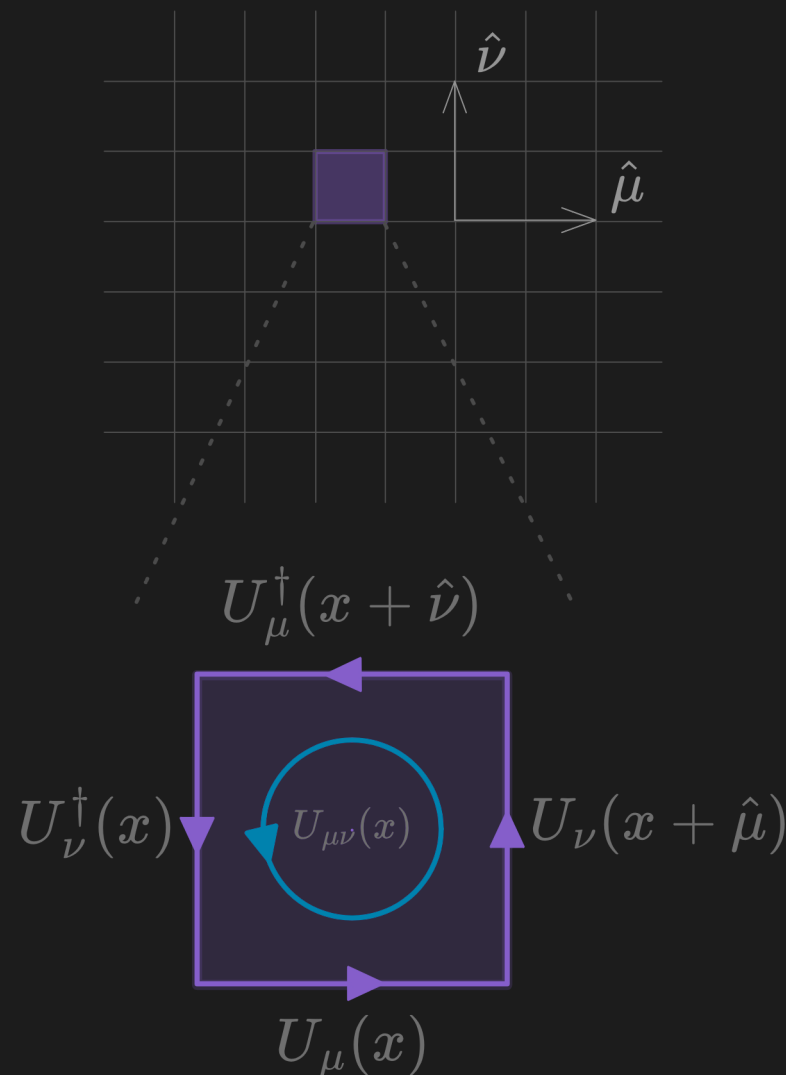


Figure 7: Illustration of the lattice

HMC: 4D $SU(3)$

Hamiltonian: $H[P, U] = \frac{1}{2}P^2 + S[U] \implies$

- $\underline{U \text{ update: } \frac{d\omega^k}{dt} = \frac{\partial H}{\partial P^k}}$

$$\frac{d\omega^k}{dt} \lambda^k = P^k \lambda^k \implies \frac{dQ}{dt} = P$$

$$Q(\epsilon) = Q(0) + \epsilon P(0) \implies$$

$$-i \log U(\epsilon) = -i \log U(0) + \epsilon P(0)$$

$$U(\epsilon) = e^{i\epsilon P(0)} U(0) \implies$$

$$\Lambda : U \longrightarrow U' := e^{i\epsilon P'} U$$

ϵ is the step size

- $\underline{P \text{ update: } \frac{dP^k}{dt} = -\frac{\partial H}{\partial \omega^k}}$

$$\frac{dP^k}{dt} = -\frac{\partial H}{\partial \omega^k} = -\frac{\partial H}{\partial Q} = -\frac{dS}{dQ} \implies$$

$$P(\epsilon) = P(0) - \epsilon \left. \frac{dS}{dQ} \right|_{t=0}$$

$$= P(0) - \epsilon F[U]$$

$$\Gamma : P \longrightarrow P' := P - \frac{\epsilon}{2} F[U]$$

$F[U]$ is the force term

HMC: 4D $SU(3)$

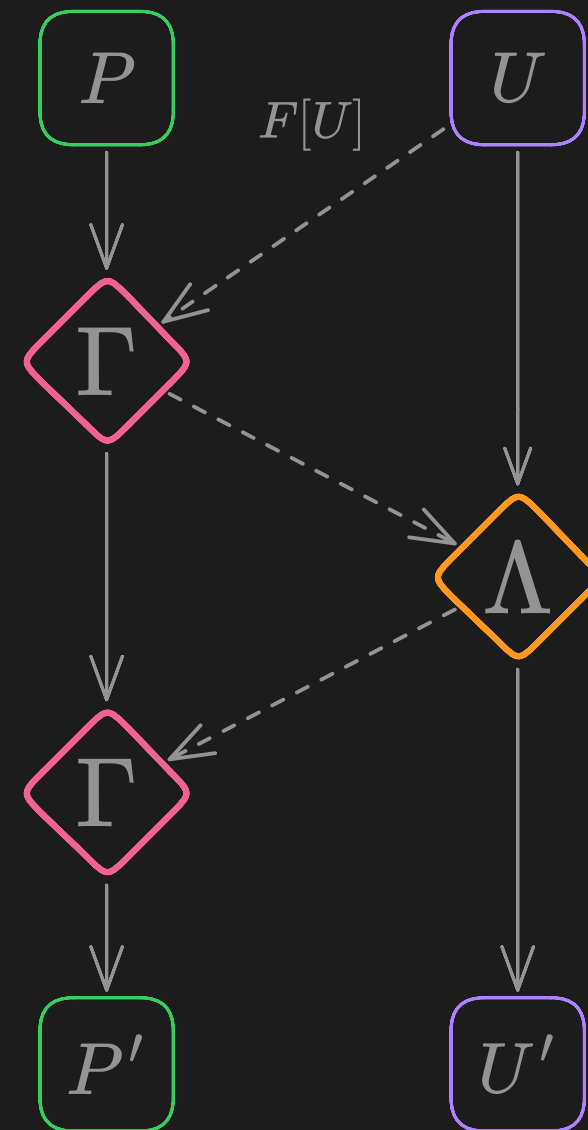
- Momentum Update:

$$\Gamma : P \longrightarrow P' := P - \frac{\varepsilon}{2} F[U]$$

- Link Update:

$$\Lambda : U \longrightarrow U' := e^{i\varepsilon P'} U$$

- We maintain a batch of N_b lattices, all updated in parallel
 - $U.dtype = complex128$
 - $U.shape = [N_b, 4, N_t, N_x, N_y, N_z, 3, 3]$



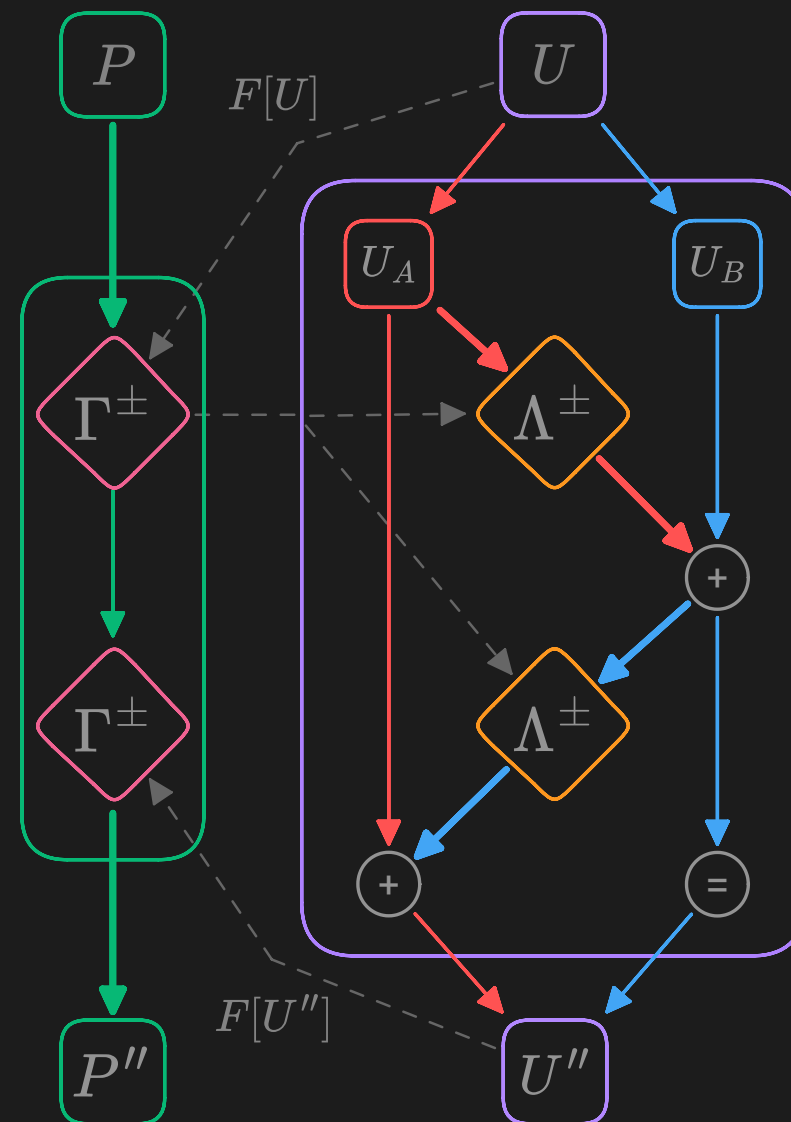
Networks 4D $SU(3)$

U -Network:

$$\text{UNet: } (U, P) \longrightarrow (s_U, t_U, q_U)$$

P -Network:

$$\text{PNet: } (U, P) \longrightarrow (s_P, t_P, q_P)$$



Networks 4D $SU(3)$

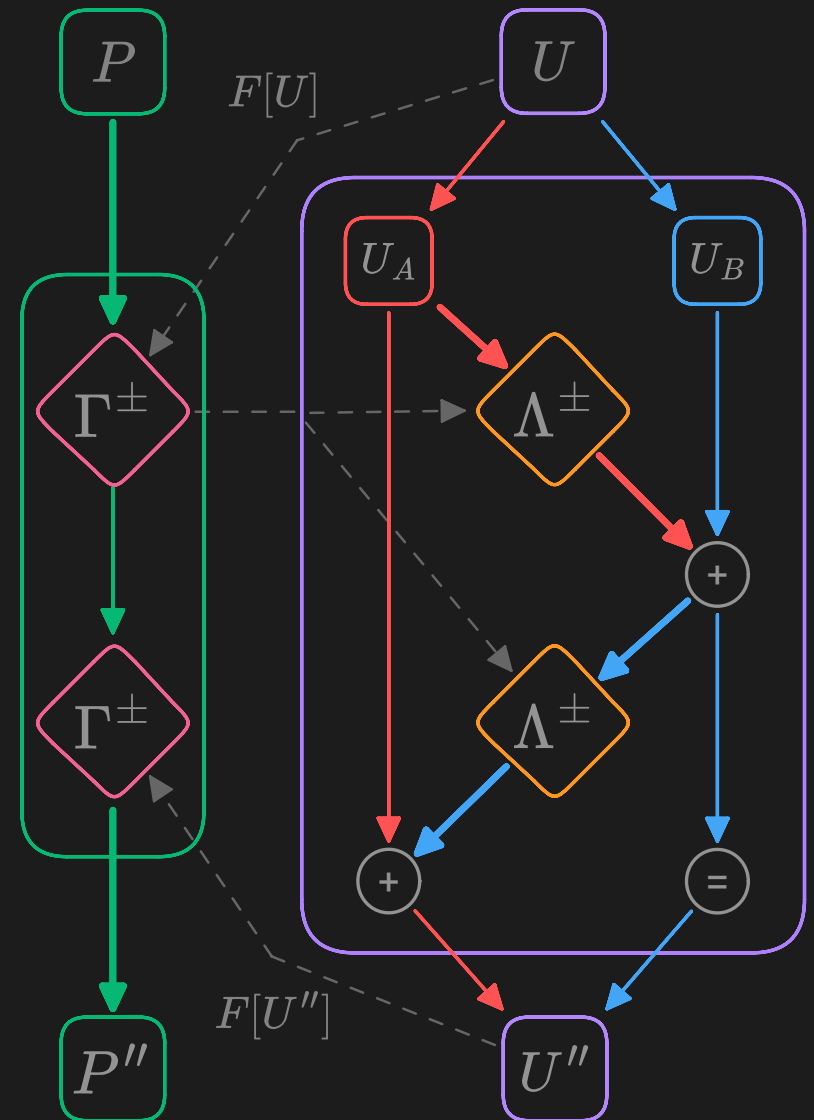
U -Network:

$$\text{UNet: } (U, P) \longrightarrow (s_U, t_U, q_U)$$

P -Network:

$$\text{PNet: } (U, P) \longrightarrow (s_P, t_P, q_P)$$

↑
let's look at this



P -Network (pt. 1)



- input¹: $(U, F) := (e^{iQ}, F)$

$$h_0 = \sigma(w_Q Q + w_F F + b)$$

$$h_1 = \sigma(w_1 h_0 + b_1)$$

$$\vdots$$

$$h_n = \sigma(w_{n-1} h_{n-2} + b_n)$$

$$z := \sigma(w_n h_{n-1} + b_n) \longrightarrow$$

- output²: (s_P, t_P, q_P)

$$\blacksquare s_P = \lambda_s \tanh(w_s z + b_s)$$

$$\blacksquare t_P = w_t z + b_t$$

$$\blacksquare q_P = \lambda_q \tanh(w_q z + b_q)$$

1. $\sigma(\cdot)$ denotes an activation function

2. $\lambda_s, \lambda_q \in \mathbb{R}$ are trainable parameters

P -Network (pt. 2)



- Use (s_P, t_P, q_P) to update $\Gamma^\pm : (U, P) \rightarrow (U, P_\pm)^1$:
 - **forward** ($d = +$):

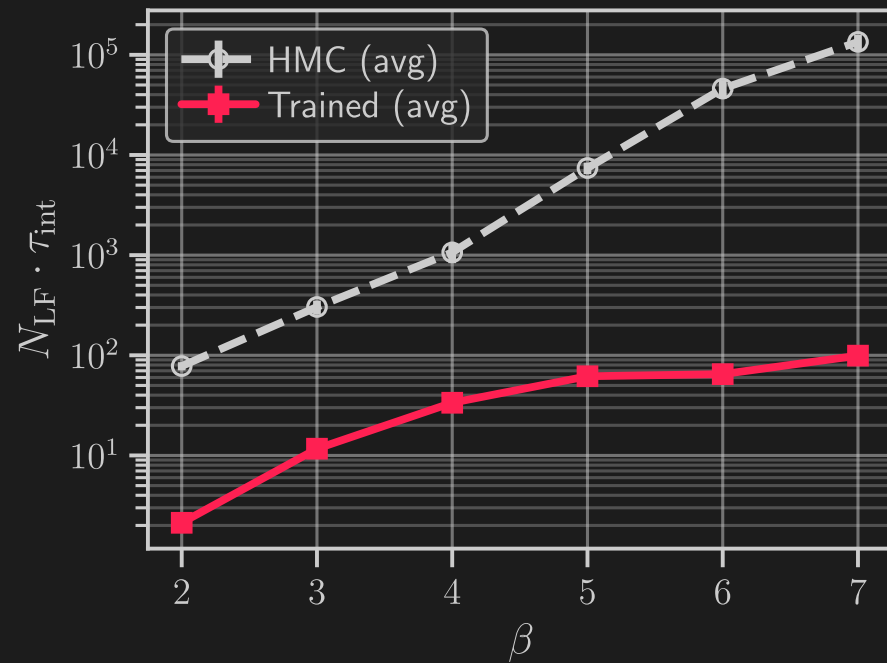
$$\Gamma^+(U, P) := P_+ = P \cdot e^{\frac{\varepsilon}{2}s_P} - \frac{\varepsilon}{2} [F \cdot e^{\varepsilon q_P} + t_P]$$

- **backward** ($d = -$):

$$\Gamma^-(U, P) := P_- = e^{-\frac{\varepsilon}{2}s_P} \left\{ P + \frac{\varepsilon}{2} [F \cdot e^{\varepsilon q_P} + t_P] \right\}$$

1. Note that $(\Gamma^+)^{-1} = \Gamma^-$, i.e. $\Gamma^+ [\Gamma^-(U, P)] = \Gamma^- [\Gamma^+(U, P)] = (U, P)$

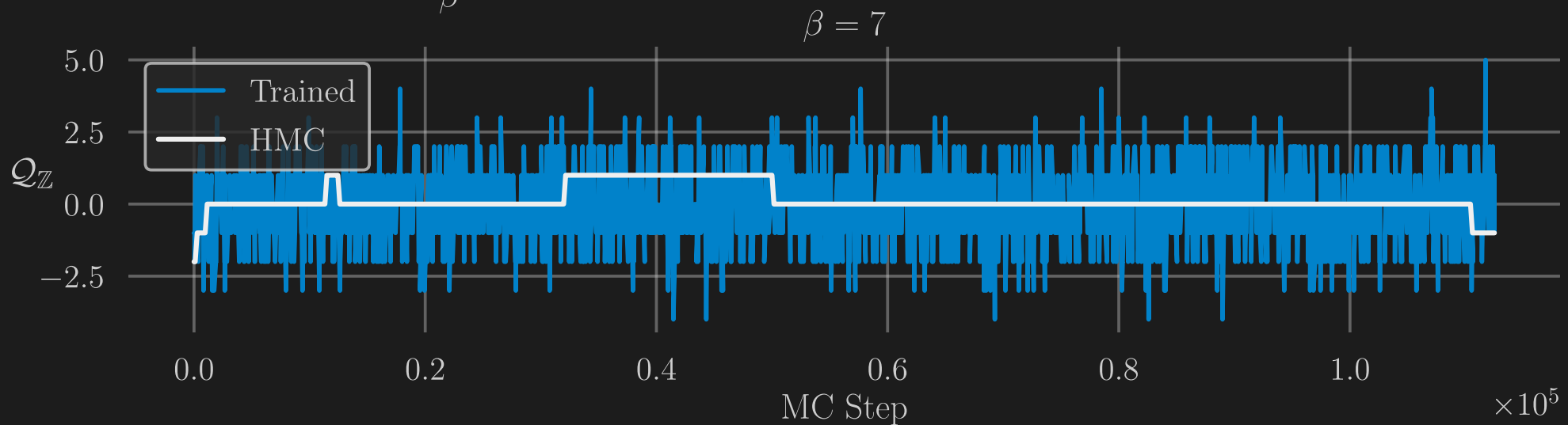
Results: 2D $U(1)$



Improvement

We can measure the performance by comparing τ_{int} for the **trained model** vs. **HMC**.

Note: lower is better



Interpretation

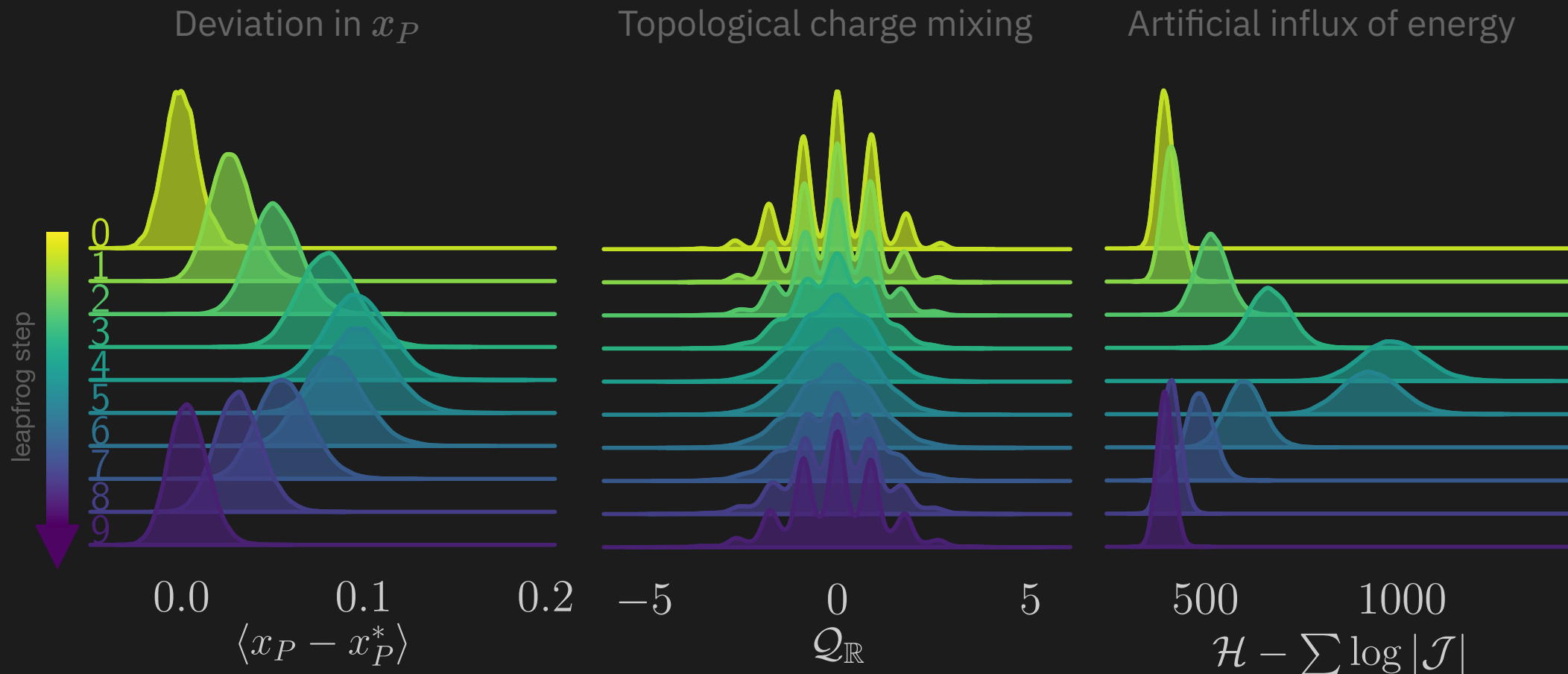
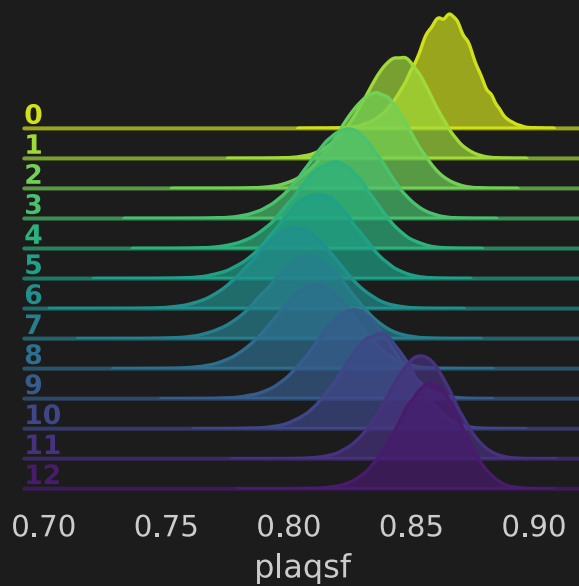
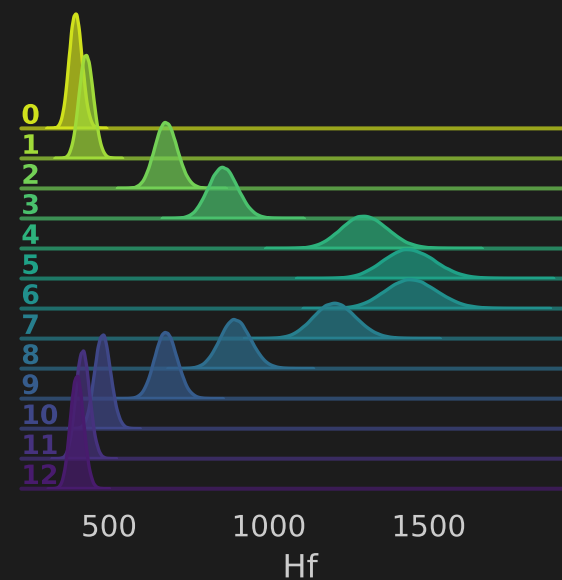


Figure 8: Illustration of how different observables evolve over a single L2HMC trajectory.

Interpretation



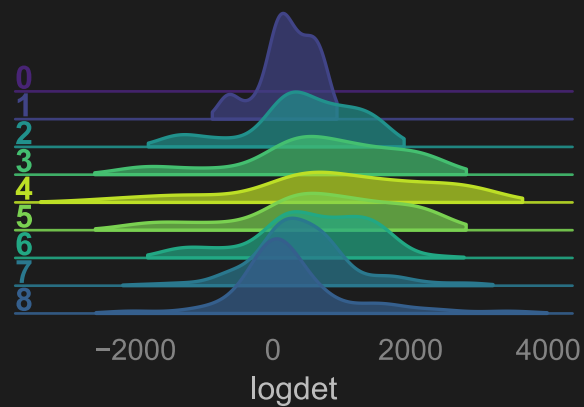
Average plaquette: $\langle x_P \rangle$ vs LF step



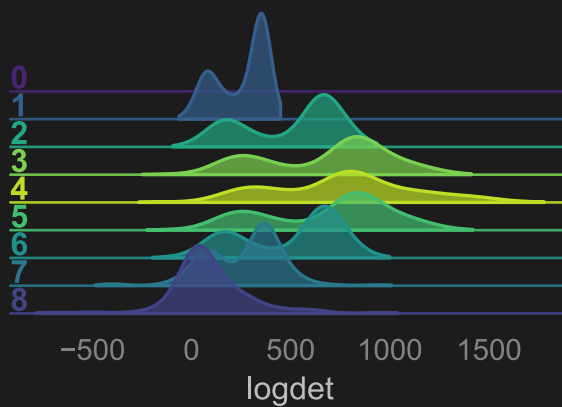
Average energy: $H - \sum \log |\mathcal{J}|$

Figure 9: The trained model artificially increases the energy towards the middle of the trajectory, allowing the sampler to tunnel between isolated sectors.

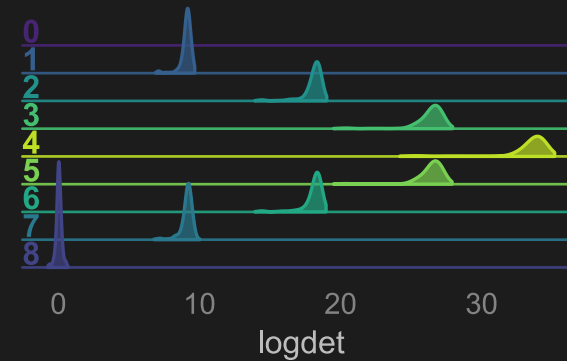
4D $SU(3)$ Results



(a) 100 train iters



(b) 500 train iters



(c) 1000 train iters

Figure 10: $\log|\mathcal{J}|$ vs. N_{LF} during training

4D $SU(3)$ Results: $\delta U_{\mu\nu}$

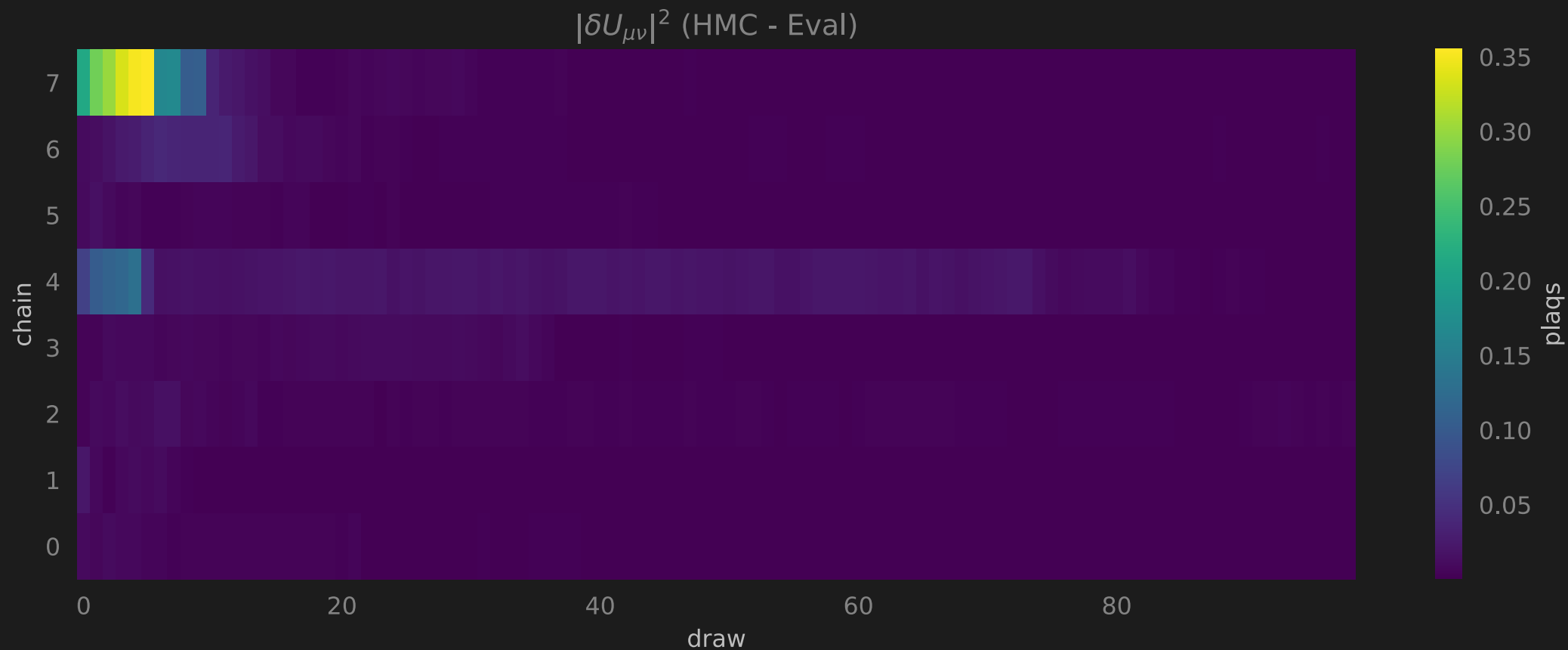


Figure 11: The difference in the average plaquette $|\delta U_{\mu\nu}|^2$ between the trained model and HMC

4D $SU(3)$ Results: $\delta U_{\mu\nu}$

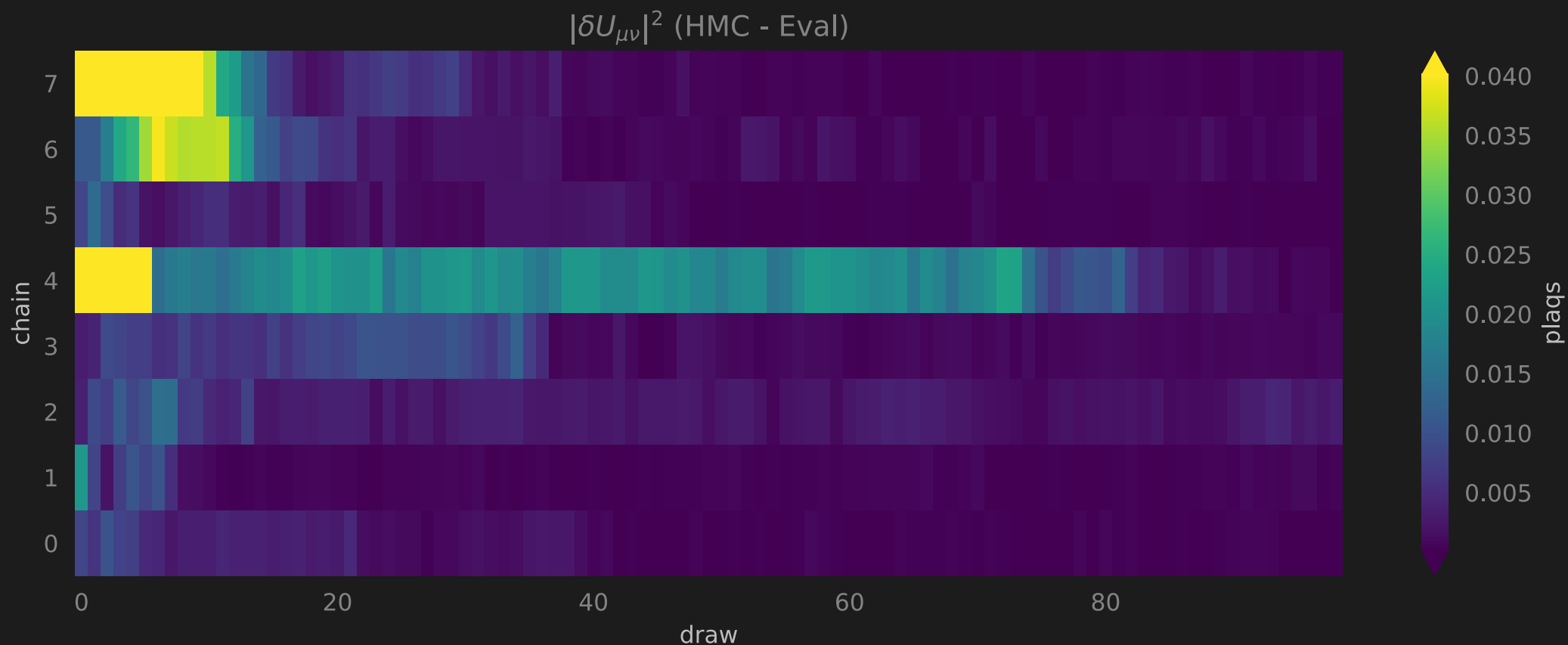




Figure 12: The difference in the average plaquette $|\delta U_{\mu\nu}|^2$ between the trained model and HMC

Next Steps

- Further code development
 -  [saforem2/12hmc-qcd](https://github.com/saforem2/12hmc-qcd)
- Continue to use / test different network architectures
 - Gauge equivariant NNs for $U_\mu(x)$ update
- Continue to test different loss functions for training
- Scaling:
 - Lattice volume
 - Network size
 - Batch size
 - # of GPUs

Thank you!

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 341 / 21233  l2hmc-qcd  codefactor  A

arXiv 2112.01582 arXiv 2105.03418




Config  Hydra  PyTorch  TensorFlow  Visualize in W&B

Acknowledgements

- **Links:**

-  Link to github
-  reach out!

- **References:**

- Link to HMC demo
- Link to slides
 -  link to github with slides
-  (Foreman et al. 2022; Foreman, Jin, and Osborn 2022, 2021)
-  (Boyda et al. 2022; Shanahan et al. 2022)



- Huge thank you to:

- Yannick Meurice
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- **ALCF Staff + Datascience Group**

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Links + References

- This talk:  [saforem2/lattice23](https://github.com/saforem2/lattice23)
- Code repo  [saforem2/12hmc-qcd](https://github.com/saforem2/12hmc-qcd)
- Slides saforem2.github.io/lattice23
- Title Slide Background (worms) animation

References

- Boyda, Denis et al. 2022. “Applications of Machine Learning to Lattice Quantum Field Theory.” In *Snowmass 2021*. <https://arxiv.org/abs/2202.05838>.
- Foreman, Sam, Taku Izubuchi, Luchang Jin, Xiao-Yong Jin, James C. Osborn, and Akio Tomiya. 2022. “HMC with Normalizing Flows.” *PoS LATTICE2021*: 073. <https://doi.org/10.22323/1.396.0073>.
- Foreman, Sam, Xiao-Yong Jin, and James C. Osborn. 2021. “Deep Learning Hamiltonian Monte Carlo.” In *9th International Conference on Learning Representations*. <https://arxiv.org/abs/2105.03418>.
- . 2022. “LeapfrogLayers: A Trainable Framework for Effective Topological Sampling.” *PoS LATTICE2021* (May): 508. <https://doi.org/10.22323/1.396.0508>.
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Extras

Integrated Autocorrelation Time

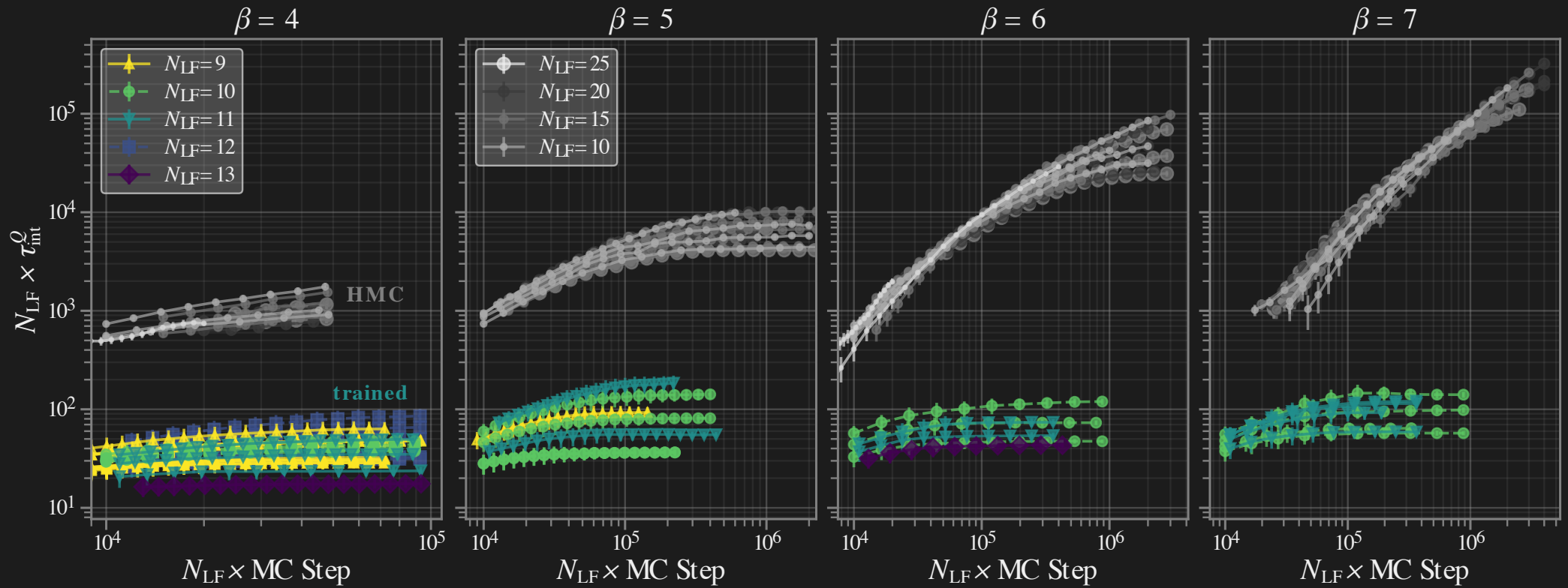
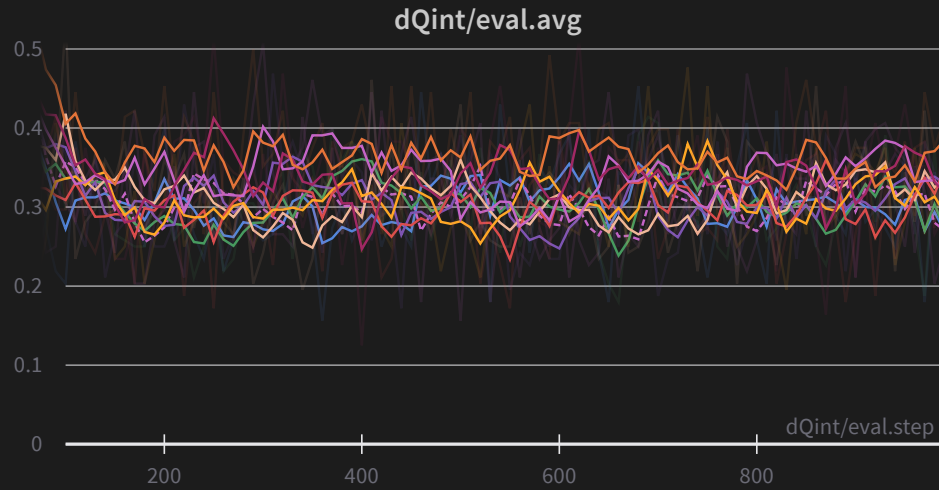
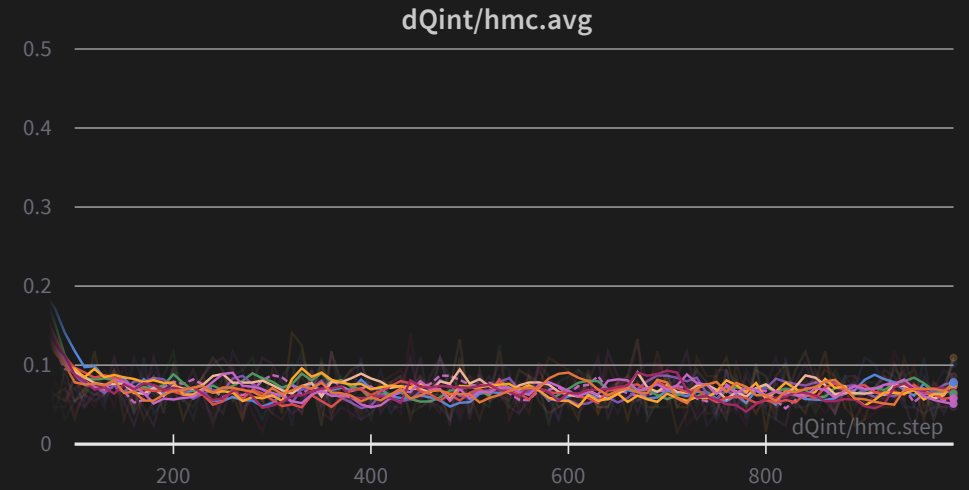


Figure 13: Plot of the integrated autocorrelation time for both the trained model (colored) and HMC (greyscale).

Comparison



(a) *Trained model*



(b) *Generic HMC*

Figure 14: Comparison of $\langle \delta Q \rangle = \frac{1}{N} \sum_{i=k}^N \delta Q_i$ for the trained model Figure 14 (a) vs. HMC Figure 14 (b)

Plaquette analysis: x_P

Deviation from $V \rightarrow \infty$ limit, x_P^*

Average $\langle x_P \rangle$, with x_P^* (dotted-lines)

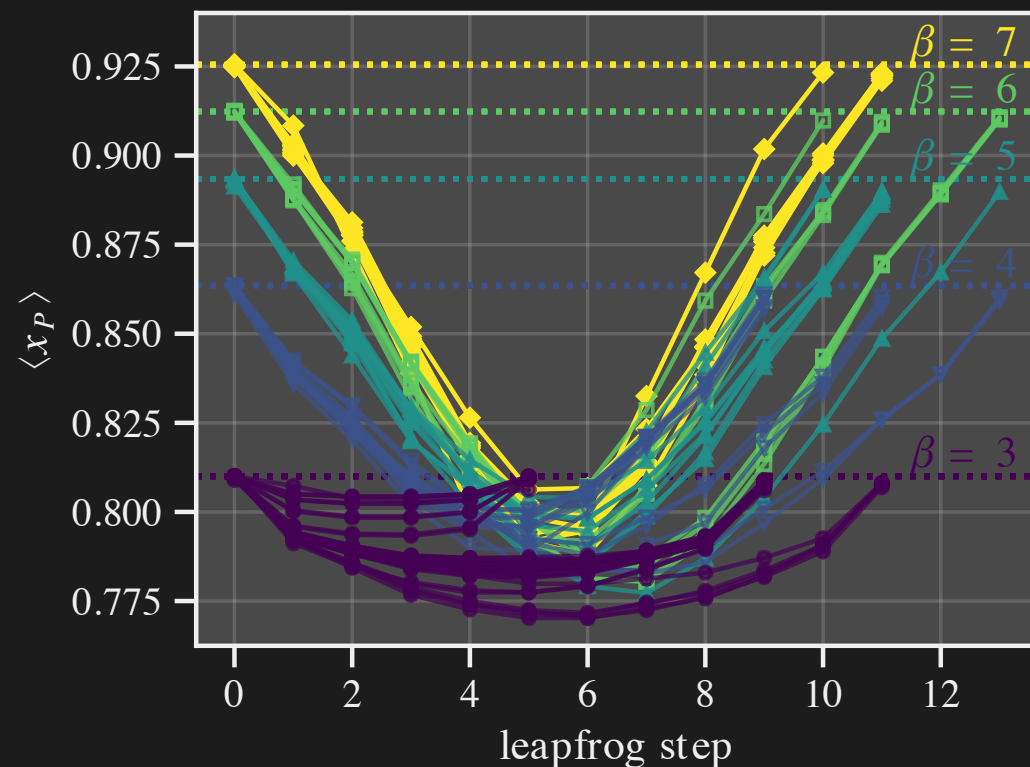
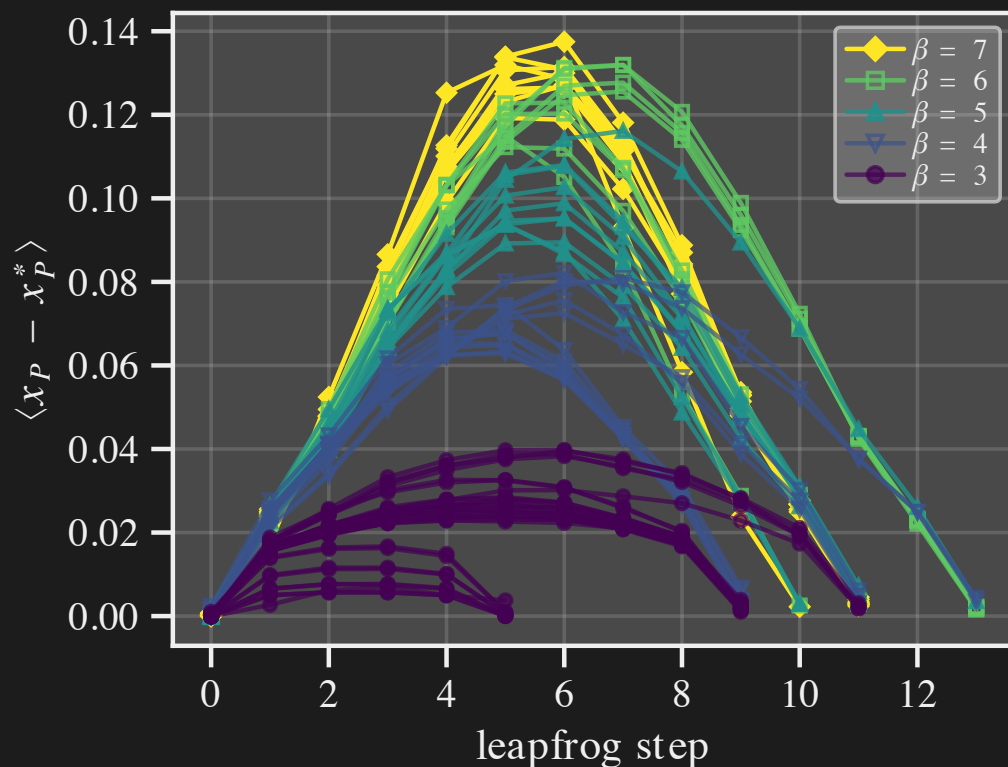


Figure 15: Plot showing how **average plaquette**, $\langle x_P \rangle$ varies over a single trajectory for models trained at different β , with varying trajectory lengths N_{LF}

Loss Function

- Want to maximize the *expected* squared charge difference¹:

$$\mathcal{L}_\theta (\xi^*, \xi) = \mathbb{E}_{p(\xi)} \left[-\delta Q^2 (\xi^*, \xi) \cdot A(\xi^*|\xi) \right]$$

- Where:

- δQ is the *tunneling rate*:

$$\delta Q(\xi^*, \xi) = |Q^* - Q|$$

- $A(\xi^*|\xi)$ is the probability² of accepting the proposal ξ^* :

$$A(\xi^*|\xi) = \min \left(1, \frac{p(\xi^*)}{p(\xi)} \left| \frac{\partial \xi^*}{\partial \xi^T} \right| \right)$$

1. Where ξ^* is the *proposed* configuration (prior to Accept / Reject)

2. And $\left| \frac{\partial \xi^*}{\partial \xi^T} \right|$ is the Jacobian of the transformation from $\xi \rightarrow \xi^*$

v -Update¹

- forward ($d = +$):

$$\Gamma^+ : (x, v) \rightarrow v' := v \cdot e^{\frac{\varepsilon}{2} s_v} - \frac{\varepsilon}{2} [F \cdot e^{\varepsilon q_v} + t_v]$$

- backward ($d = -$):

$$\Gamma^- : (x, v) \rightarrow v' := e^{-\frac{\varepsilon}{2} s_v} \left\{ v + \frac{\varepsilon}{2} [F \cdot e^{\varepsilon q_v} + t_v] \right\}$$

1. Note that $(\Gamma^+)^{-1} = \Gamma^-$, i.e. $\Gamma^+ [\Gamma^-(x, v)] = \Gamma^- [\Gamma^+(x, v)] = (x, v)$

x -Update

- forward ($d = +$):

$$\Lambda^{+}(x, v) = x \cdot e^{\frac{\varepsilon}{2} s_x} - \frac{\varepsilon}{2} [v \cdot e^{\varepsilon q_x} + t_x]$$

- backward ($d = -$):

$$\Lambda^{-}(x, v) = e^{-\frac{\varepsilon}{2} s_x} \left\{ x + \frac{\varepsilon}{2} [v \cdot e^{\varepsilon q_x} + t_x] \right\}$$

Lattice Gauge Theory (2D $U(1)$)

🌀 Link Variables

$$U_\mu(n) = e^{ix_\mu(n)} \in \mathbb{C}, \quad \text{where}$$

$$x_\mu(n) \in [-\pi, \pi)$$

🔥 Wilson Action

$$S_\beta(x) = \beta \sum_P \cos x_P,$$

$$x_P = [x_\mu(n) + x_\nu(n + \hat{\mu}) - x_\mu(n + \hat{\nu}) - x_\nu(n)]$$

Note: x_P is the product of links around 1×1 square, called a “plaquette”

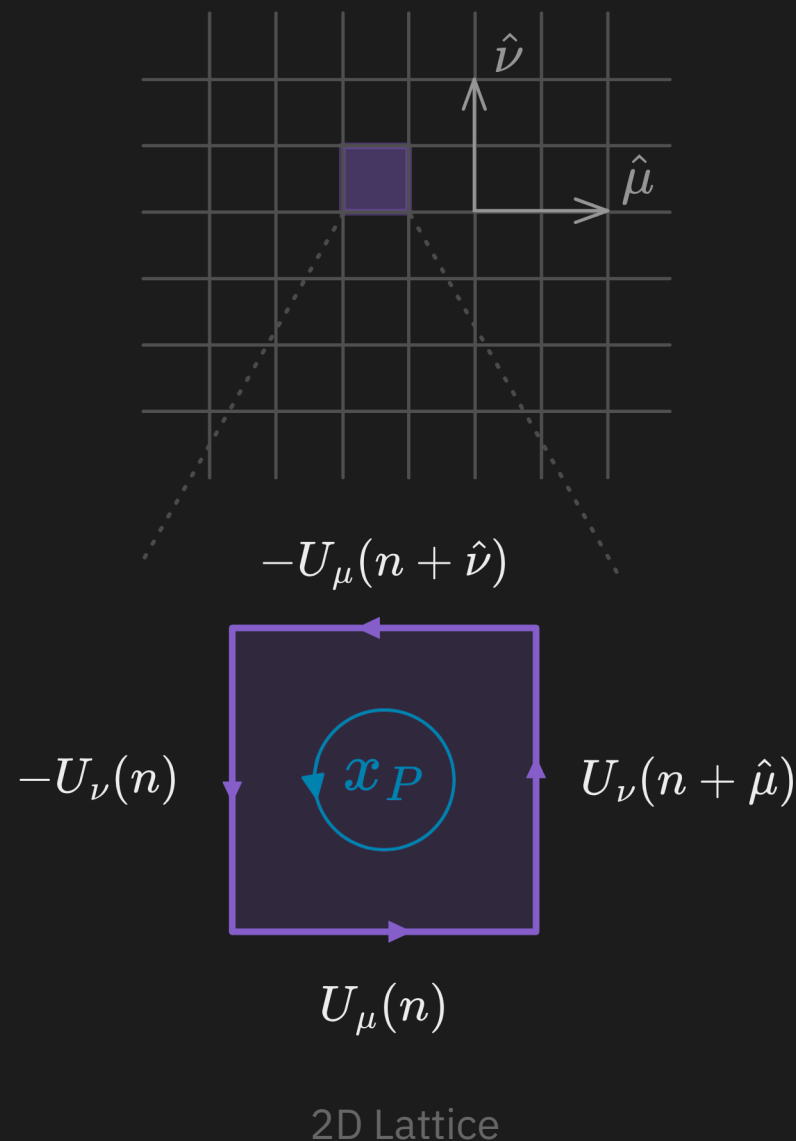


Figure 16: Jupyter Notebook



Annealing Schedule

- Introduce an *annealing schedule* during the training phase:

$$\{\gamma_t\}_{t=0}^N = \{\gamma_0, \gamma_1, \dots, \gamma_{N-1}, \gamma_N\}$$

where $\gamma_0 < \gamma_1 < \dots < \gamma_N \equiv 1$, and $|\gamma_{t+1} - \gamma_t| \ll 1$

- **Note:**
 - for $|\gamma_t| < 1$, this rescaling helps to reduce the height of the energy barriers \implies
 - easier for our sampler to explore previously inaccessible regions of the phase space

Networks 2D $U(1)$

- Stack gauge links as $\text{shape}(U_\mu) = [\text{Nb}, 2, \text{Nt}, \text{Nx}] \in \mathbb{C}$

$$x_\mu(n) := [\cos(x), \sin(x)]$$

with $\text{shape}(x_\mu) = [\text{Nb}, 2, \text{Nt}, \text{Nx}, 2] \in \mathbb{R}$

- x -Network:
 - $\psi_\theta : (x, v) \longrightarrow (s_x, t_x, q_x)$
- v -Network:
 - $\varphi_\theta : (x, v) \longrightarrow (s_v, t_v, q_v)$

Networks 2D $U(1)$

- Stack gauge links as $\text{shape}(U_\mu) = [\text{Nb}, 2, \text{Nt}, \text{Nx}] \in \mathbb{C}$

$$x_\mu(n) := [\cos(x), \sin(x)]$$

with $\text{shape}(x_\mu) = [\text{Nb}, 2, \text{Nt}, \text{Nx}, 2] \in \mathbb{R}$

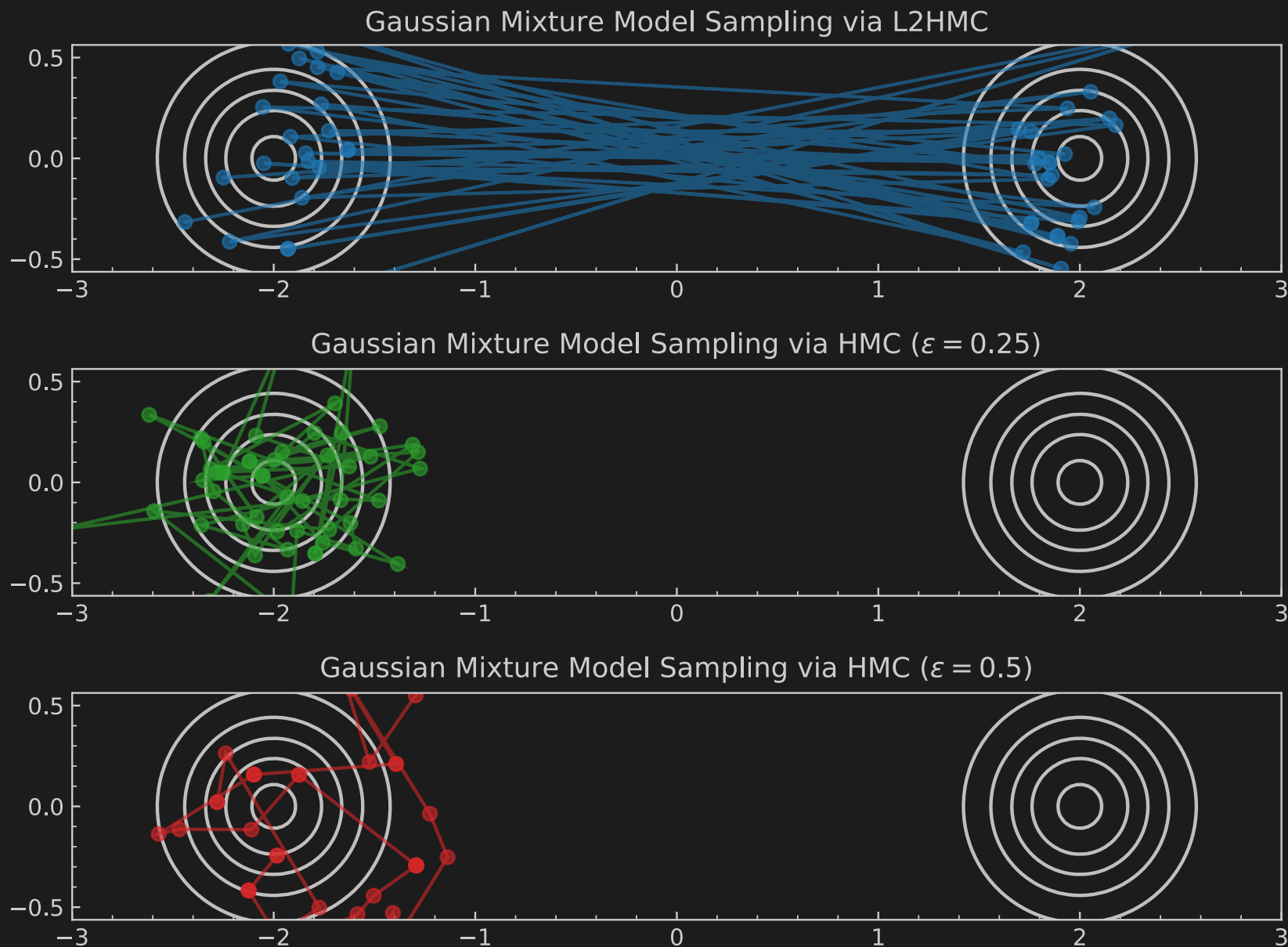
- x -Network:

- $\psi_\theta : (x, v) \longrightarrow (s_x, t_x, q_x)$

- v -Network:

- $\varphi_\theta : (x, v) \longrightarrow (s_v, t_v, q_v) \longleftarrow$ lets look at this

Toy Example: $\text{GMM} \in \mathbb{R}^2$



Physical Quantities

- To estimate physical quantities, we:
 - calculate physical observables at **increasing** spatial resolution
 - perform extrapolation to continuum limit

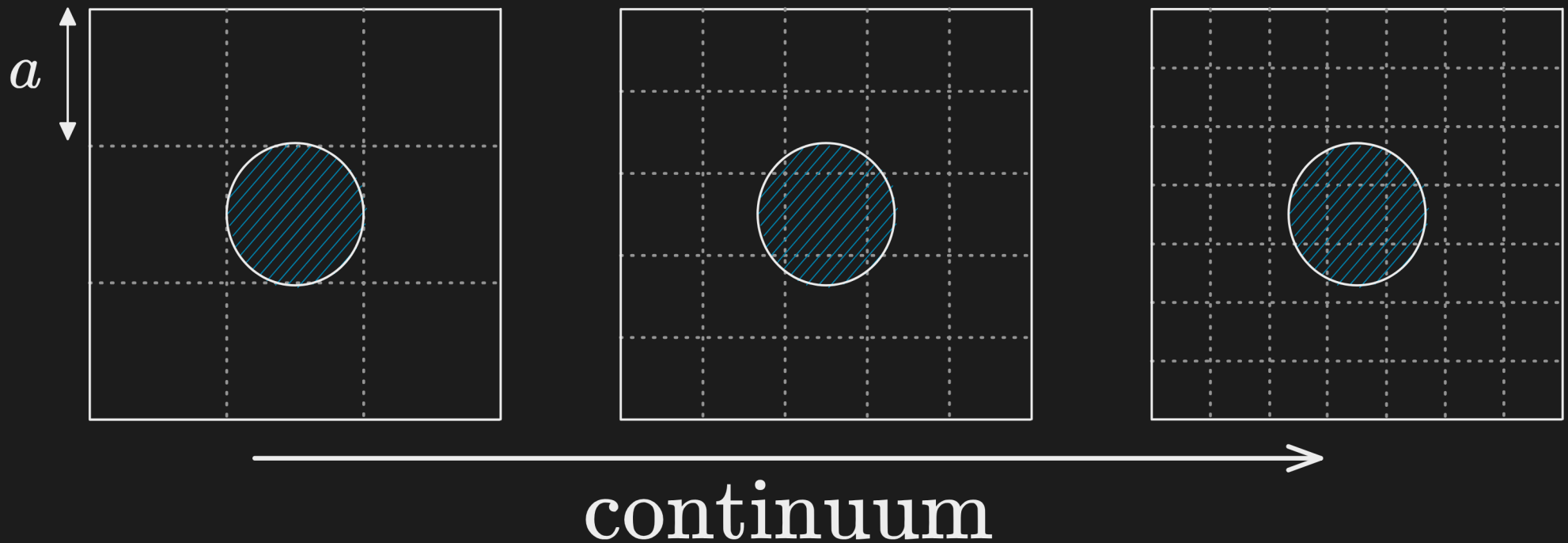


Figure 17: Increasing the physical resolution ($a \rightarrow 0$) allows us to make predictions about numerical values of physical quantities in the continuum limit.

